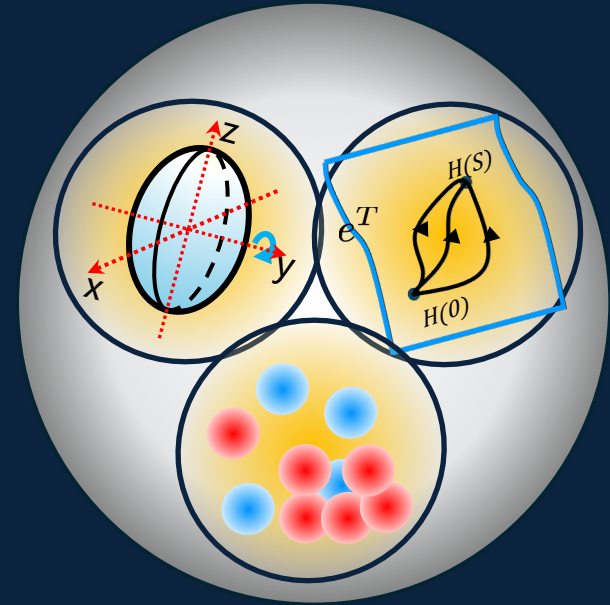


# multi-reference equation of motion

Zhonghao Sun

University of Notre Dame



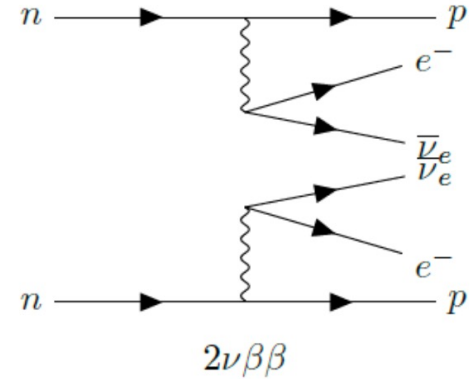
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# $2\nu\beta\beta$ decay nuclear matrix element

The half-life of  $2\nu\beta\beta$  depends on the nuclear matrix elements

$$(T_{1/2}^{2\nu})^{-1} = G^{2\nu} g_A^4 (M^{2\nu})^2,$$

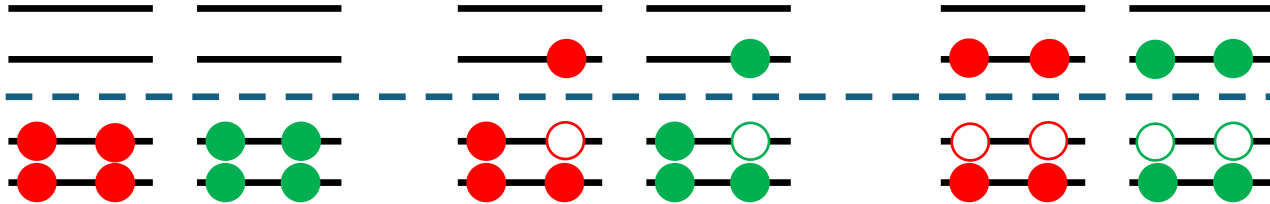
$$M^{2\nu} = \sum_k \frac{\langle 0_f^+ || \sum_a \tau_a^- \sigma_a || 1_k^+ \rangle \langle 1_k^+ || \sum_b \tau_b^- \sigma_b || 0_i^+ \rangle}{(E_k - (E_i + E_f)/2)/m_e}$$



Need to compute the ground state of initial and final nuclei, as well as the **all-excited  $1^+(0^+)$**  state of intermediate nuclei for the Gamow-Teller (Fermi) operator

**Challenges to get accurate high-lying states of the intermediate nuclei.**

# How do we compute the states of nuclei?



- all possible configurations
- Huge, sparse Matrix
- Diagonalization to get energy and wavefunction

$$H \begin{pmatrix} |\Psi_{\text{ex}}\rangle \\ \dots \\ |\Psi_0\rangle \end{pmatrix} = E \begin{pmatrix} |\Psi_{\text{ex}}\rangle \\ \dots \\ |\Psi_0\rangle \end{pmatrix}$$

$10^8 \approx 500$  MB in memory for one vector  
 $10^{12}$  is almost the limit of classic computer  
 $\text{dim} \approx N^A$  (exponential scale)  
 $\approx 2000^A$

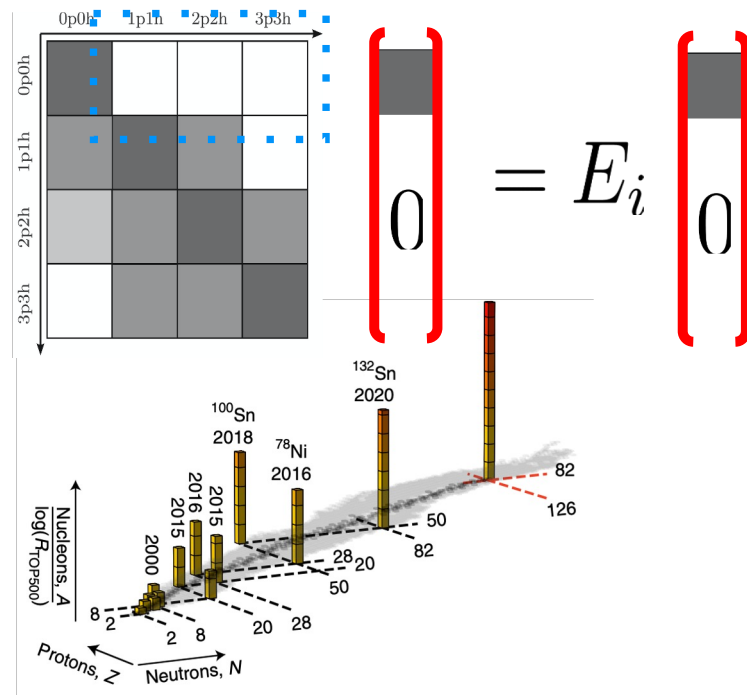
# Ab-initio nuclear structure calculation with polynomial scaling (Coupled Cluster, IM-SRG,...)

$$\langle \Psi_0 | e^{-T} H e^T | \Psi_0 \rangle = E$$

$$\langle \Psi_{ij\dots}^{ab\dots} | e^{-T} H e^T | \Psi_0 \rangle = 0$$

$$\hat{T} = \sum_i T_i \quad \text{Excitation operators}$$

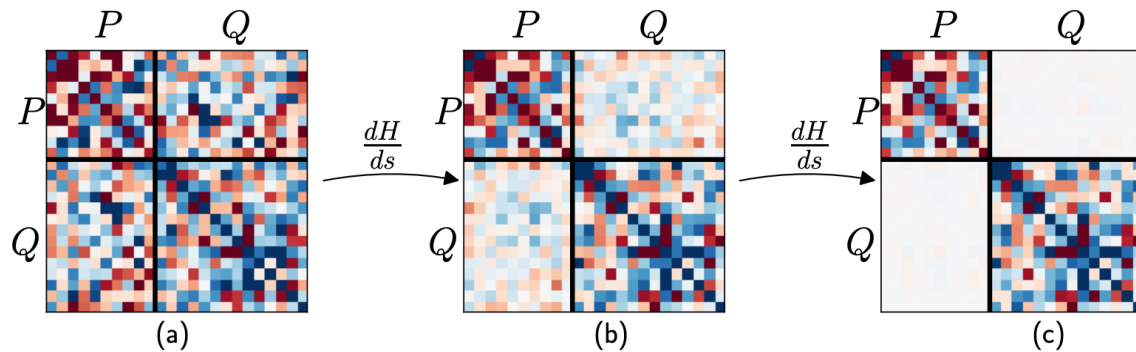
$$\hat{T}_1 = \sum_{ia} t_i \alpha_a^\dagger \alpha_i, \quad \hat{T}_2 = \frac{1}{2} \sum_{i < j, a < b} t_{ij}^{ab} \alpha_a^\dagger \alpha_b^\dagger \alpha_i \alpha_j$$



Hu, *et al*, Nature physics, **18**, 1196. (2022)

Similarity transformation, decoupled the ground state, **Polynomial scale**

# How we compute open shell nuclei (VS-IMSRG, CC)



Annu. Rev. Nucl. Part. Sci. 69:307

$$\eta^{\text{Wh}}(s) \equiv \frac{H^{\text{od}}(s)}{\Delta(s)}$$

$$|\Psi\rangle = C_0 \left( 1 + \frac{C_i^a}{C_0} a^\dagger i + \frac{C_{ij}^{ab}}{C_0} a^\dagger b^\dagger ij + \dots \right) |\Psi_0\rangle = P|\Psi_0\rangle + Q|\Psi_0\rangle \propto e^S |0\rangle$$

- The true state should be dominated by the reference, otherwise ( $C_0 = 0$ )  $T$ ,  $S$  or  $\Omega$  may diverge for operators with any rank.
- Low-lying states, the quality and accuracy depends on the **shell gap**
- **High energy states: multi-shell effective interaction, vanished energy denominator, intruder states, etc.**

# A long-term effort to tackle intruder states

## A shell-model description of $0^+$ intruder states in even-even nuclei

K. Heyde, J. Jolie<sup>1</sup>, J. Moreau<sup>2</sup>, J. Ryckebusch<sup>1</sup>, M. Waroquier<sup>3</sup>, P. Van Duppen, M. Huyse, J.L. Wood [Show more](#) ▾

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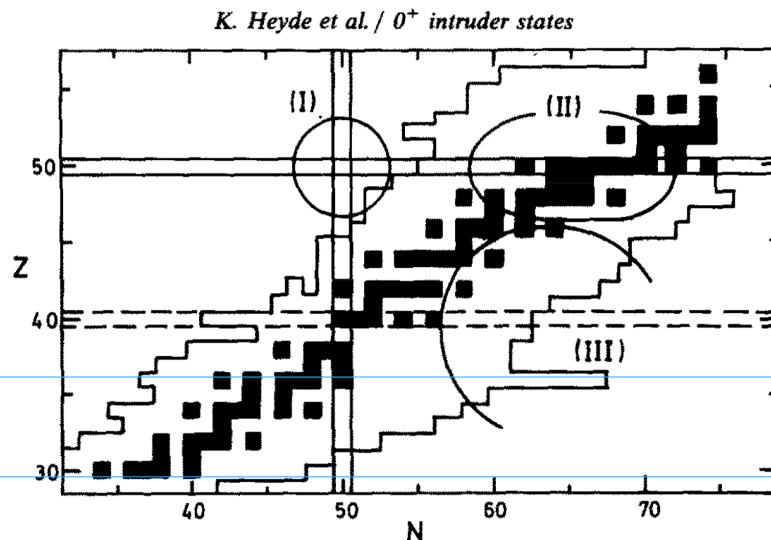


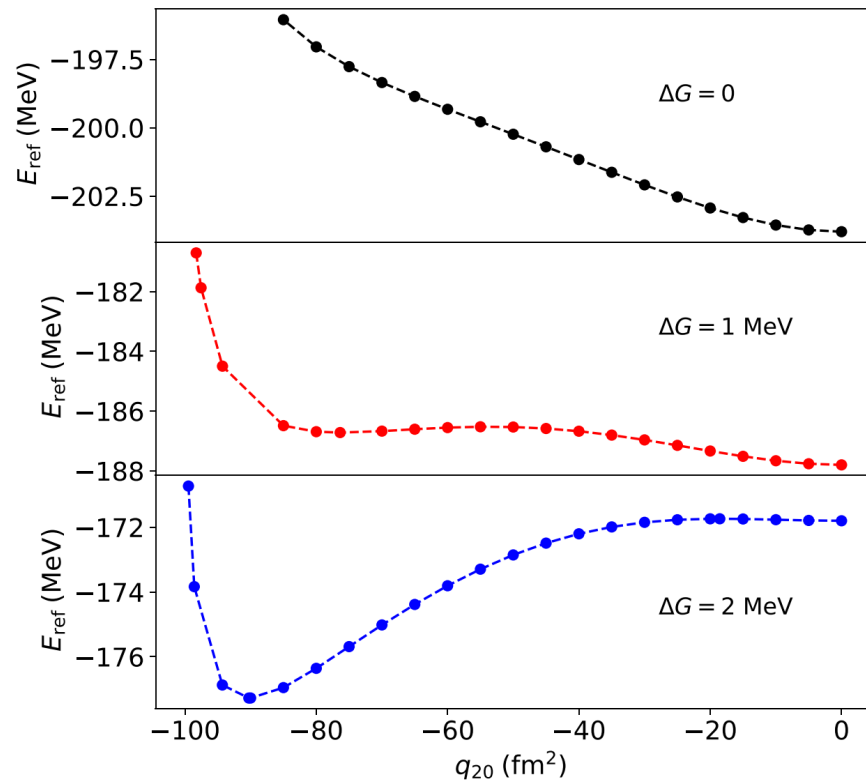
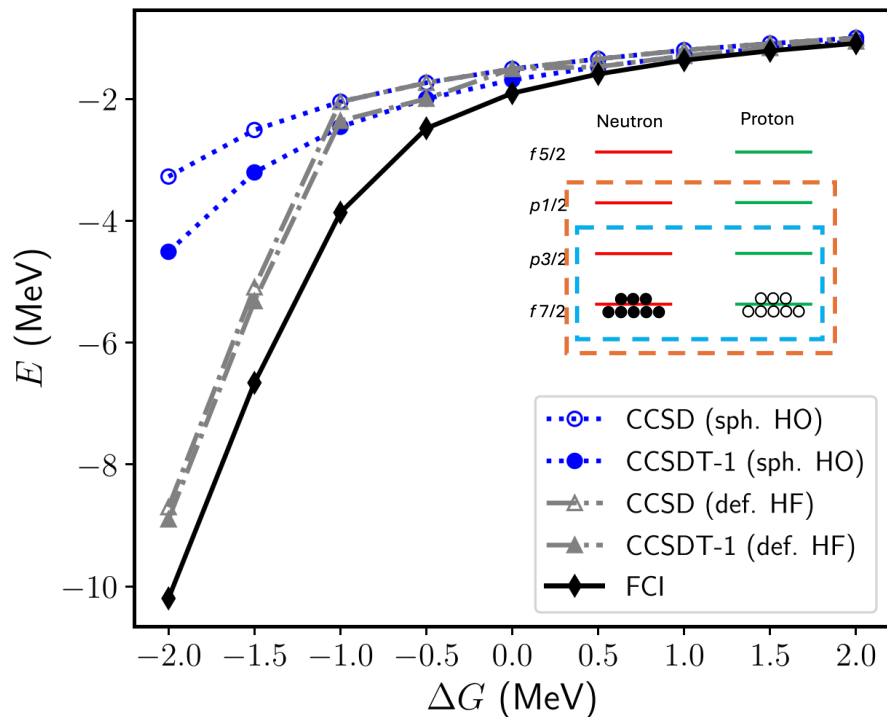
Fig. 1. Schematic division of nuclei in the  $Z \approx 40, 50$  and  $50 \leq N \leq 82$  mass table into three major regions, i.e. (I) region near doubly-closed shells; (II) region of possible intruder states; (III) region of strongly deformed nuclei.

$$\frac{\langle \Psi_\mu | Q | \Psi_\mu \rangle}{\langle \Psi_\mu | \Psi_\mu \rangle} = \frac{\rho_Q}{1 + \rho_Q} \lesssim \frac{1}{2}$$

K. Suzuki, Shyh Yuan Lee.  
*PTP.* 64, 2091(1980)

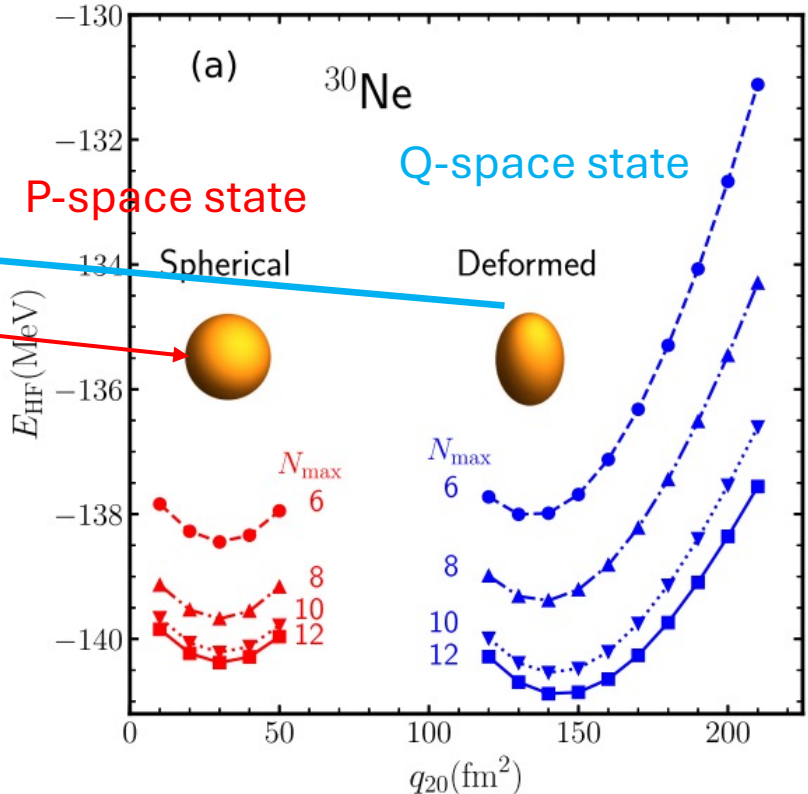
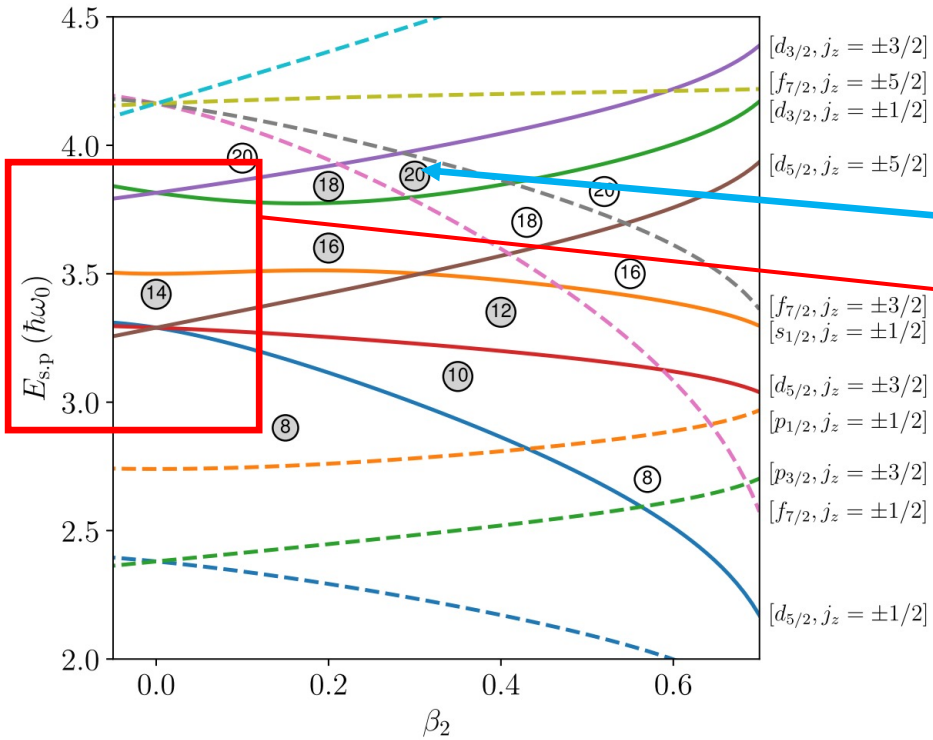
$P$  space components  
should be dominant in  
the wavefunction

# Example: vanished shell gap(quasi degenerate)



Strong entanglement, spontaneously rotational symmetry breaking, deformation.

# Deformation emerge as vanished shell gap



intruder state, level crossing, deformation, cluster structure...

# Equation of Motion method,

The standard problem:

$$H|\mu\rangle = E_\mu|\mu\rangle$$

We focus on specified excitation mode the excited state on top of a reference state

$$|\mu\rangle = Q_\mu^\dagger|0\rangle \quad Q_\mu|0\rangle = 0 \quad Q_\mu^\dagger = \sum_\alpha X_\alpha^\mu \eta_\alpha^\dagger$$

EOM Secular equation:

$$\sum_\beta \langle 0 | [\eta_\alpha, [H, \eta_\beta^\dagger]] | 0 \rangle X_\beta^\mu = \delta E_\mu \sum_\beta \langle 0 | [\eta_\alpha, \eta_\beta^\dagger] | 0 \rangle X_\beta^\mu$$

Systematically truncated to fit available computational power and the physics of interest

D.J. Rowe. Rev. Mod. Phys. **40**, 153

# Equation of motion coupled cluster

Equation of motion method to get excited states, particle attached, charge exchange

$$(\bar{H} \hat{R}_\mu^{(A\pm 2)})_C |\Phi_0\rangle = \omega_\mu \hat{R}_\mu^{(A\pm 2)} |\Phi_0\rangle$$

G. R. Jansen, et al, Phys. Rev. C 83, 054306

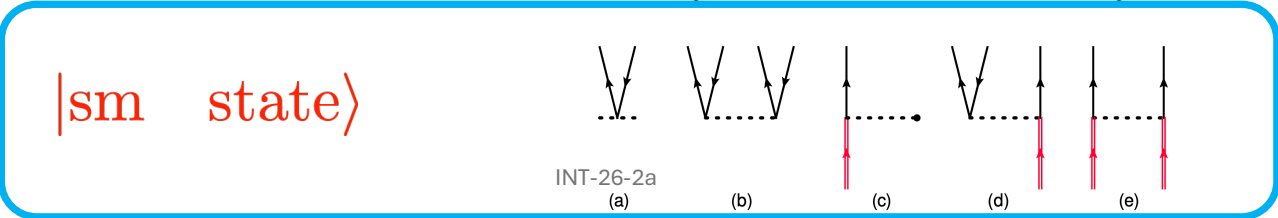
Operator	Expression
$\hat{R}_{2p-0h}^{A+2}$	$\frac{1}{2} \sum_{ab} r^{ab} a_a^\dagger a_b^\dagger$
$\hat{R}_{3p-1h}^{A+2}$	$\frac{1}{2} \sum_{ab} r^{ab} a_a^\dagger a_b^\dagger + \frac{1}{6} \sum_{abc} r_i^{abc} a_a^\dagger a_b^\dagger a_c^\dagger a_i$
$\hat{R}_{0p-2h}^{A-2}$	$\frac{1}{2} \sum_{ij} r_{ij} a_j a_i$
$\hat{R}_{1p-3h}^{A-2}$	$\frac{1}{2} \sum_{ij} r_{ij} a_j a_i + \frac{1}{6} \sum_{aijk} r_{ijk}^a a_a^\dagger a_k a_j a_i$

Method	Reference State $ 0\rangle$
HF	$ HF\rangle$
TDA	$ HF\rangle$
RPA	$ RPA\rangle$
HFB	$ HFB\rangle$
QRPA	$ QRPA\rangle$

Generator  $O_\nu^\dagger$

$$\begin{aligned} & a_m^\dagger a_i \\ & \sum_{mi} X_{mi}^\nu a_m^\dagger a_i \\ & \sum_{mi} \left( X_{mi}^\nu a_m^\dagger a_i - Y_{mi}^\nu a_i^\dagger a_m \right) \\ & \alpha_k^\dagger \alpha_{k'}^\dagger \\ & \sum_{k < k'} \left( X_{kk'}^\nu \alpha_k^\dagger \alpha_{k'}^\dagger - Y_{kk'}^\nu \alpha_{k'} \alpha_k \right) \end{aligned}$$

MR-EOM



# Equation of Motion method, excitations on top of a “good” reference

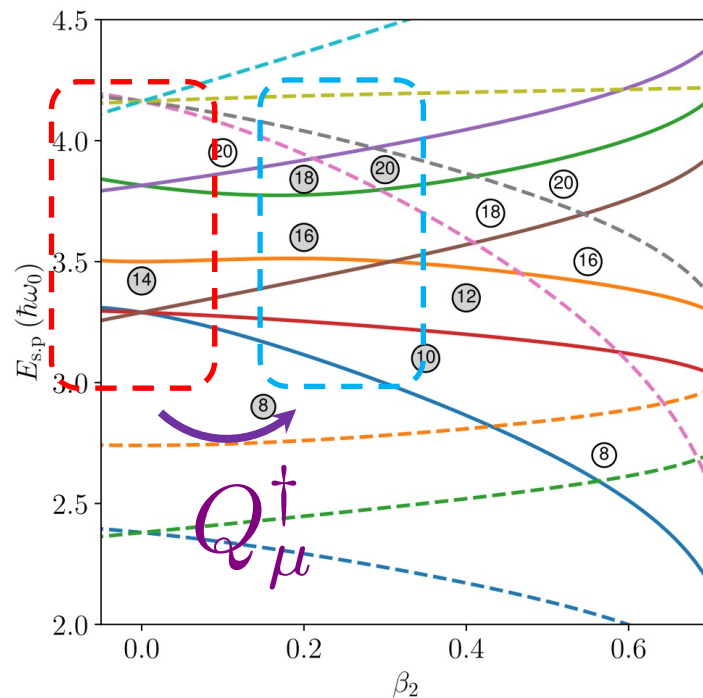
The shell model state as a reference state

$$|\Psi_0\rangle = e^{\Omega} |\nu^n\rangle$$

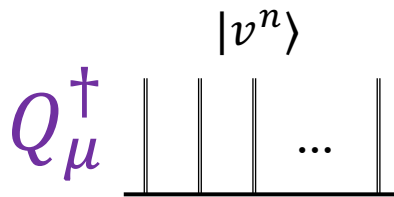
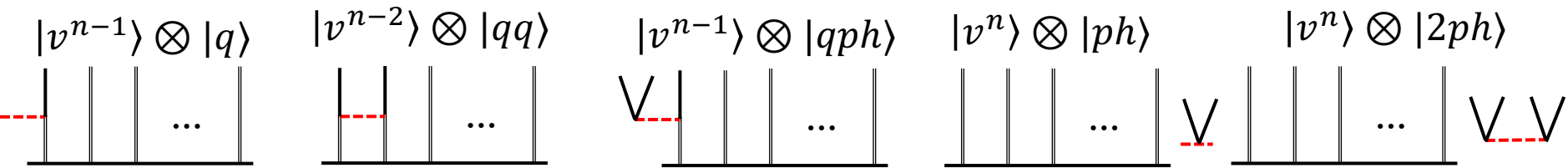
An intruder state could be a collective 1-particle-1-hole, 2-particle-2-hole excitation on top of the shell model state

$$|\Psi^\mu\rangle = Q_\mu^\dagger |\Psi_0\rangle$$

$Q_\mu^\dagger$  has a systematically hierarchy of order-by-order truncations

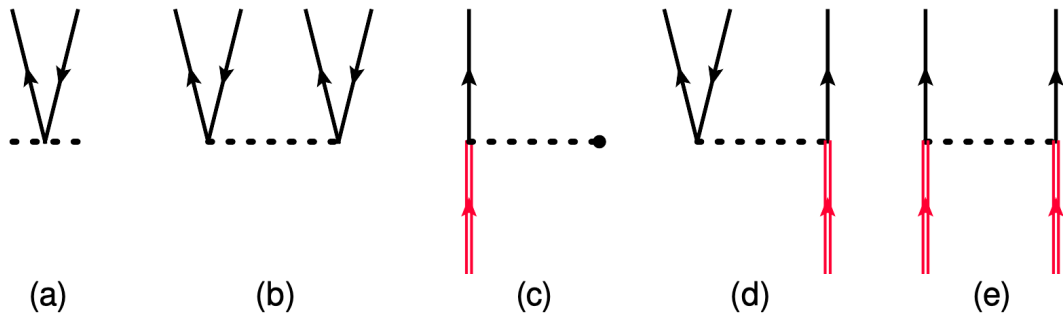


# Equation of motion on top of a shell model state



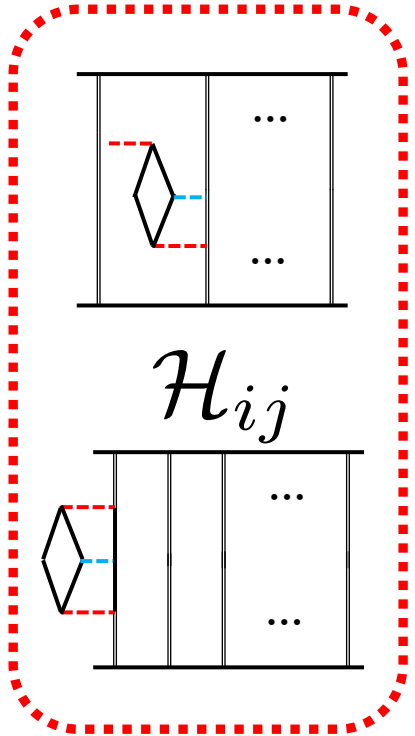
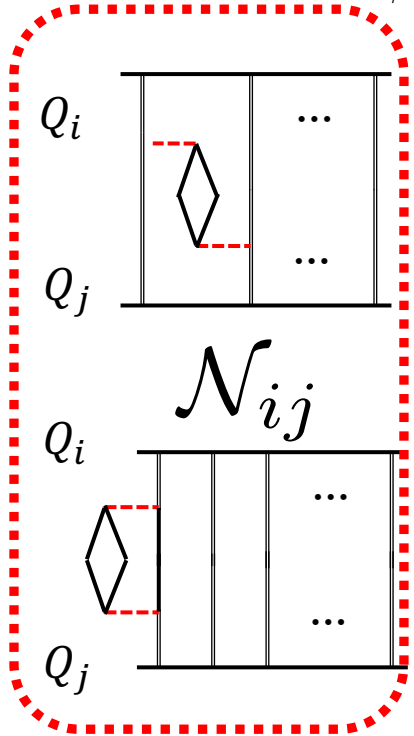
$$Q_\mu^\dagger |\Phi_0\rangle = 0$$

- Pick a reference state from the model space
- Generate excited states (automatically orthogonalized to all state inside the model space)
- Solve the equation of motion in the Fock space (two- and three-body operators)



# Equation of motion on top of a shell model state

$$\sum_{\beta} \langle 0 | [\eta_{\alpha}, [H, \eta_{\beta}^{\dagger}]] | 0 \rangle X_{\beta}^{\mu} = \delta E_{\mu} \sum_{\beta} \langle 0 | [\eta_{\alpha}, \eta_{\beta}^{\dagger}] | 0 \rangle X_{\beta}^{\mu}$$



- The Norm and Hamiltonian matrix (kernels) come as one-, two and three-body operators inside the model space. Its matrix elements are computed as expectation of these operators respect to the reference (shell model state).

$$\mathcal{H}_{ij} = \text{Tr}[(Q_i^{\dagger} Q_j)_c^{(k)} \rho^{(k)}]$$

- Order-by-order approximation on the truncation of  $Q_{\mu}$ .
- The EOM is solved in the **Fock space**, suppose to work for any nuclei for a given model space

# How to solve the EOM

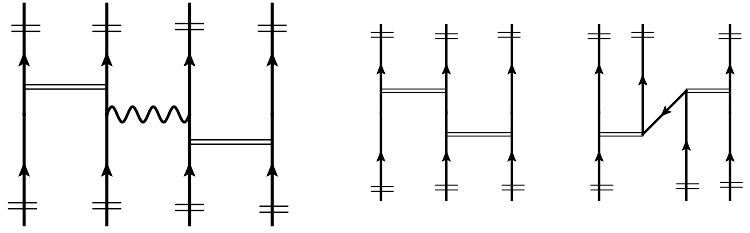
- Lanczos or Arnoldi algorithm

$$(HQ_i)_c|0\rangle = \beta_{i+1}Q_{i+1}|0\rangle + \alpha_iQ_i|0\rangle + \beta_iQ_{i-1}|0\rangle$$

$$\alpha_i = [Q_i, [H, Q_j]] \quad \text{Up to 4-body}$$

$$\beta_i = [\tilde{Q}_i, \tilde{Q}_i^\dagger] \quad \text{Up to 3-body}$$

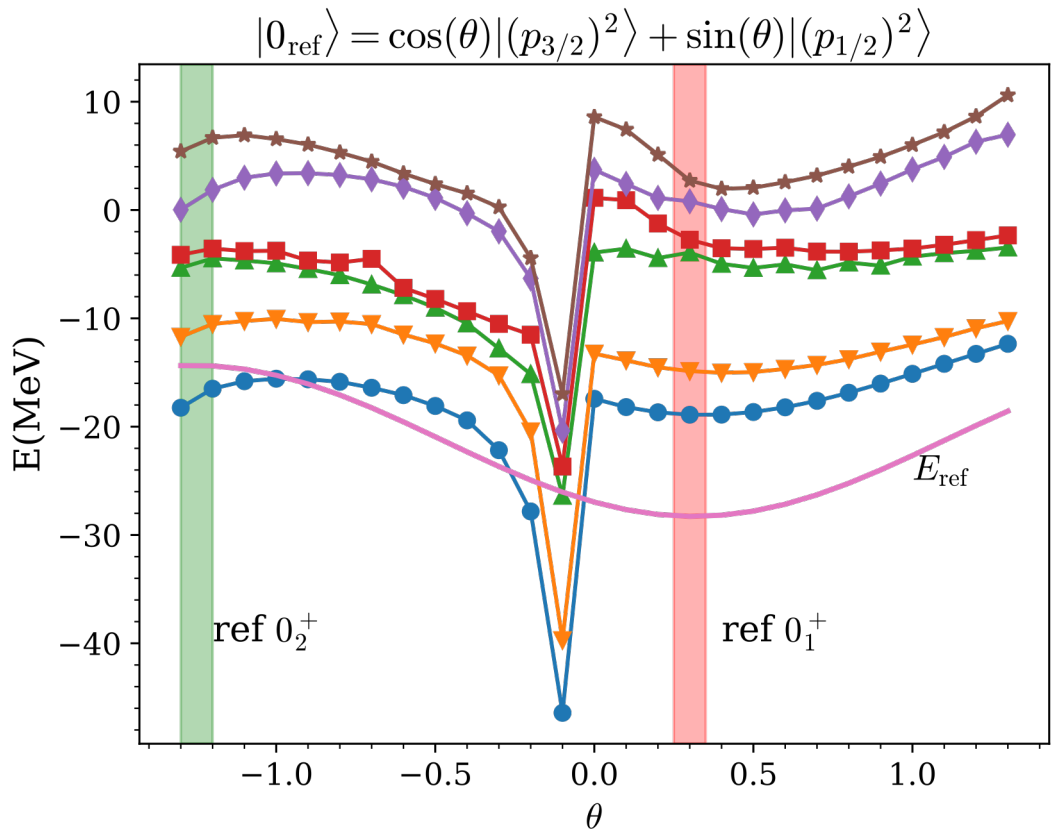
Problem:



1. truncation error leads to failure of Lanczos algorithm, include 3-body intermediate in Arnoldi.

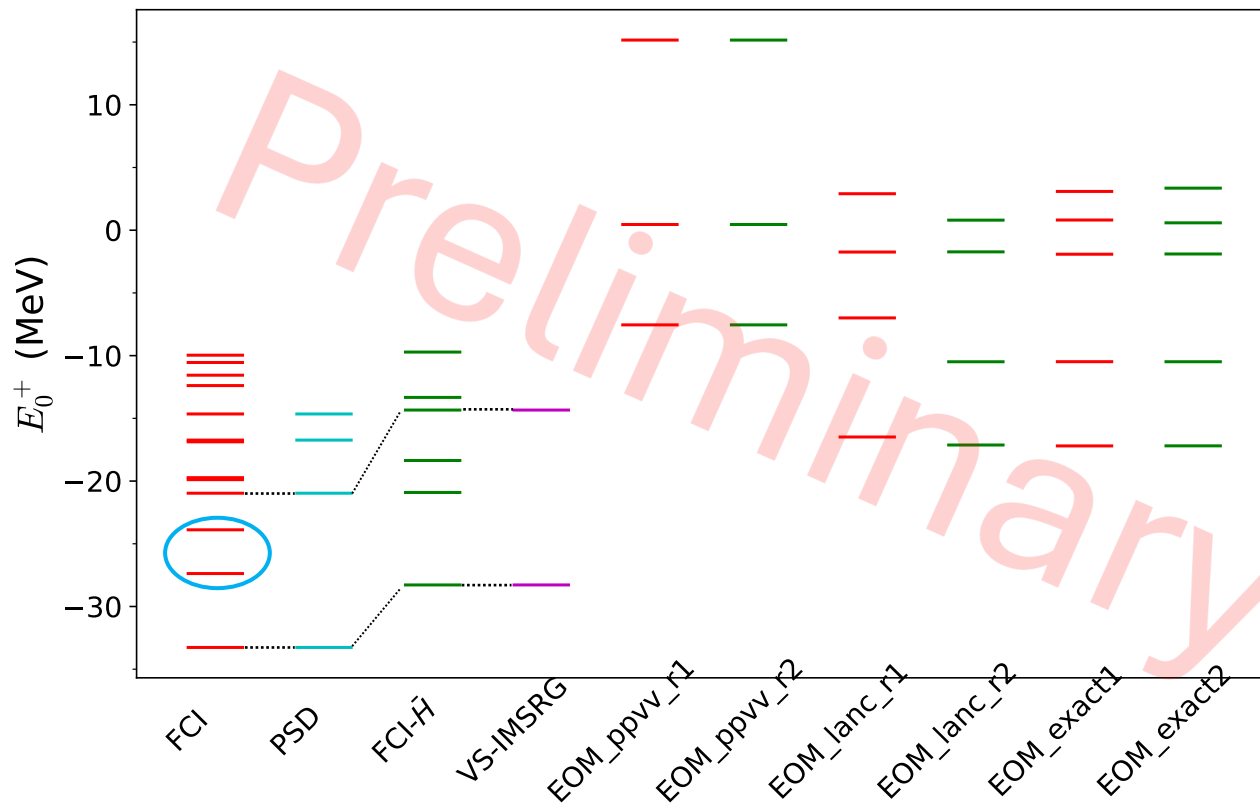
2. Overcomplete basis,  $\sum \|Q_i^k\|^2$  become numerical huge though is one  $\langle 0|[\tilde{Q}_i, \tilde{Q}_i^\dagger]|0\rangle$ , spurious states. **SVD on the Norm matrix** to remove redundant degree of freedom

# Picking accurate reference state



For Simple system: Based on two different valence states of  ${}^6\text{He}$ , we roughly get the same spectrum from EOM calculation.

# $0^+$ states of ${}^6\text{He}$ , $N_{max}=2$



The shell-model-EOM captured the intruder states, as well as other high lying states.



# Summary and outlook

- An equation of motion method on top of a shell model state (multiple reference).
- Excitations defined in Fock space as two body operators.
- The norm (overlap) and Hamiltonian kernel is computed as expectation of operators on shell model wavefunction,
- Applied to two valence particle system (He, Li), scalar excitation ( $0^+ \rightarrow 0^+$ ,  $1^+ \rightarrow 1^+$ )
- Working on generate tensor excitations ( $0^+ \rightarrow 1^+$ )

Thank you!