

Characterizing the nuclear models informed by PREX and CREX

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arXiv.2406.05267



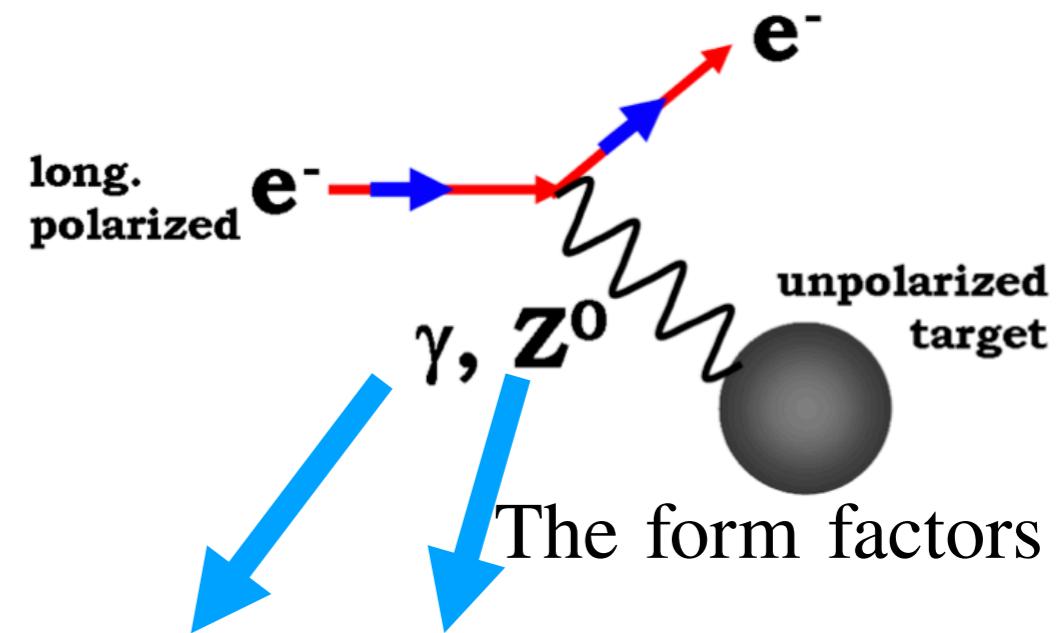
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Parity violating electron scattering

| | CREX | PREX |
|--------------|-------------------------------------|----------------------------------|
| (N,Z) | (28,20) Ca | (126,82) Pb |
| q (fm-1) | 0.8733 | 0.3977 |
| Fch, Rch(fm) | 0.1581, 3.481 | 0.409, 5.503 |
| Apv | 2668±106(stat) ±40(syst) | 550±16(stat) ±8(syst) |
| Fw | 0.1304±0.0052(stat) ±0.002(syst) | 0.368±0.013(exp) ±0.001(theo) |
| Fch-Fw | 0.0277±0.0052(stat) ±0.002(syst) | 0.041±0.013(exp) ±0.001(theo) |
| Rw | 3.64±0.026(exp) ±0.023(theo) | 5.8±0.075(tot) |
| Rw-Rch | 0.159±0.026(exp) ±0.023(theo) | 0.297±0.075(tot) |
| Rn-Rp | 0.121±0.026(exp) ±0.024(theo) | 0.283±0.071(tot) |

CREX 2022 PREX I 2012 PREX II 2021

MREX: see Tuesday's talk by Concettina



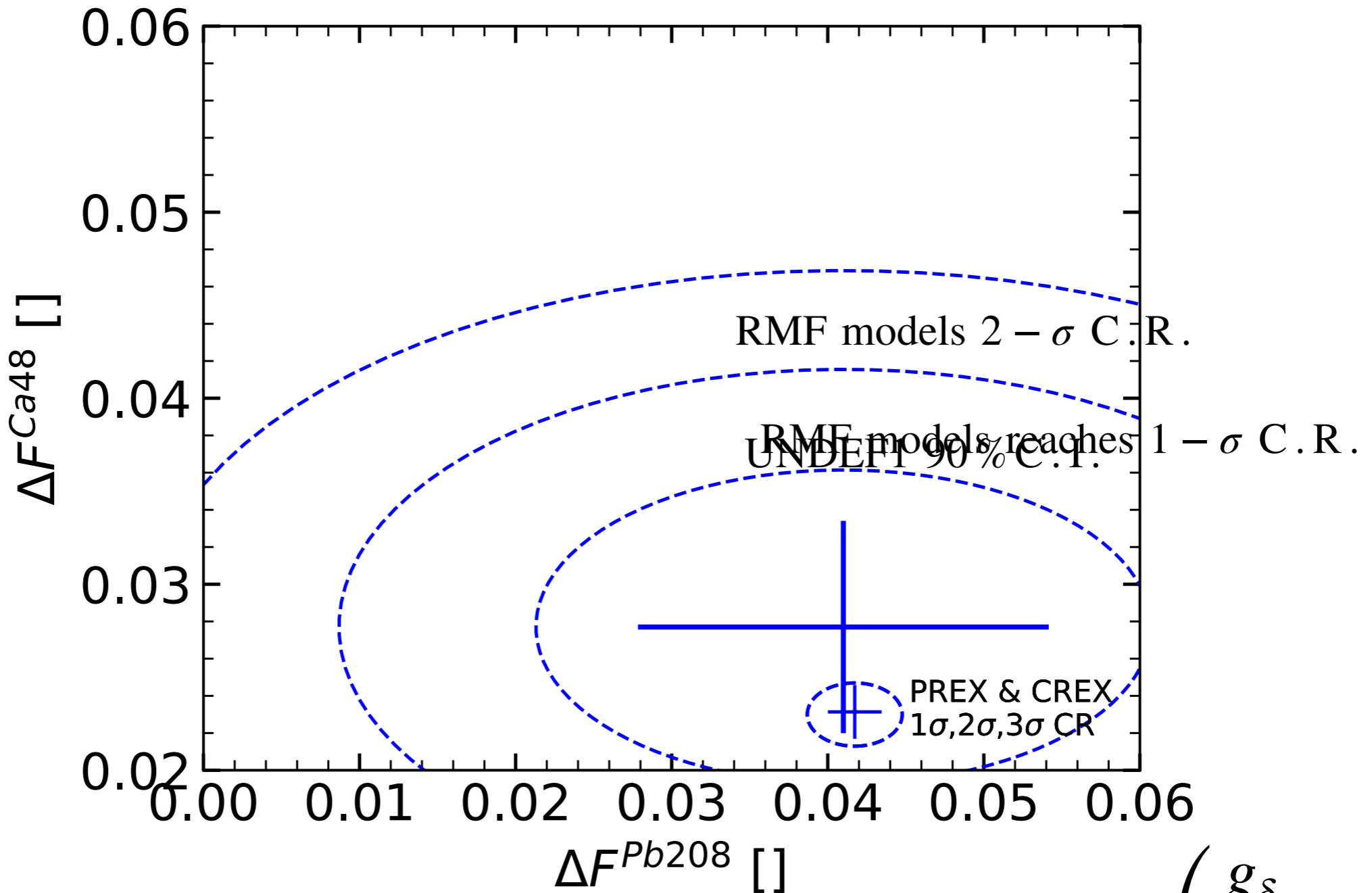
The parity violating asymmetry :

$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}$$

The weak interaction violates parity :

$$J_Z^\mu = -\frac{1}{2}\bar{\psi}_L \gamma^\mu \psi_L - \sin^2(\Theta)\bar{\psi} \gamma^\mu \psi$$

Pre PREX-CREX era

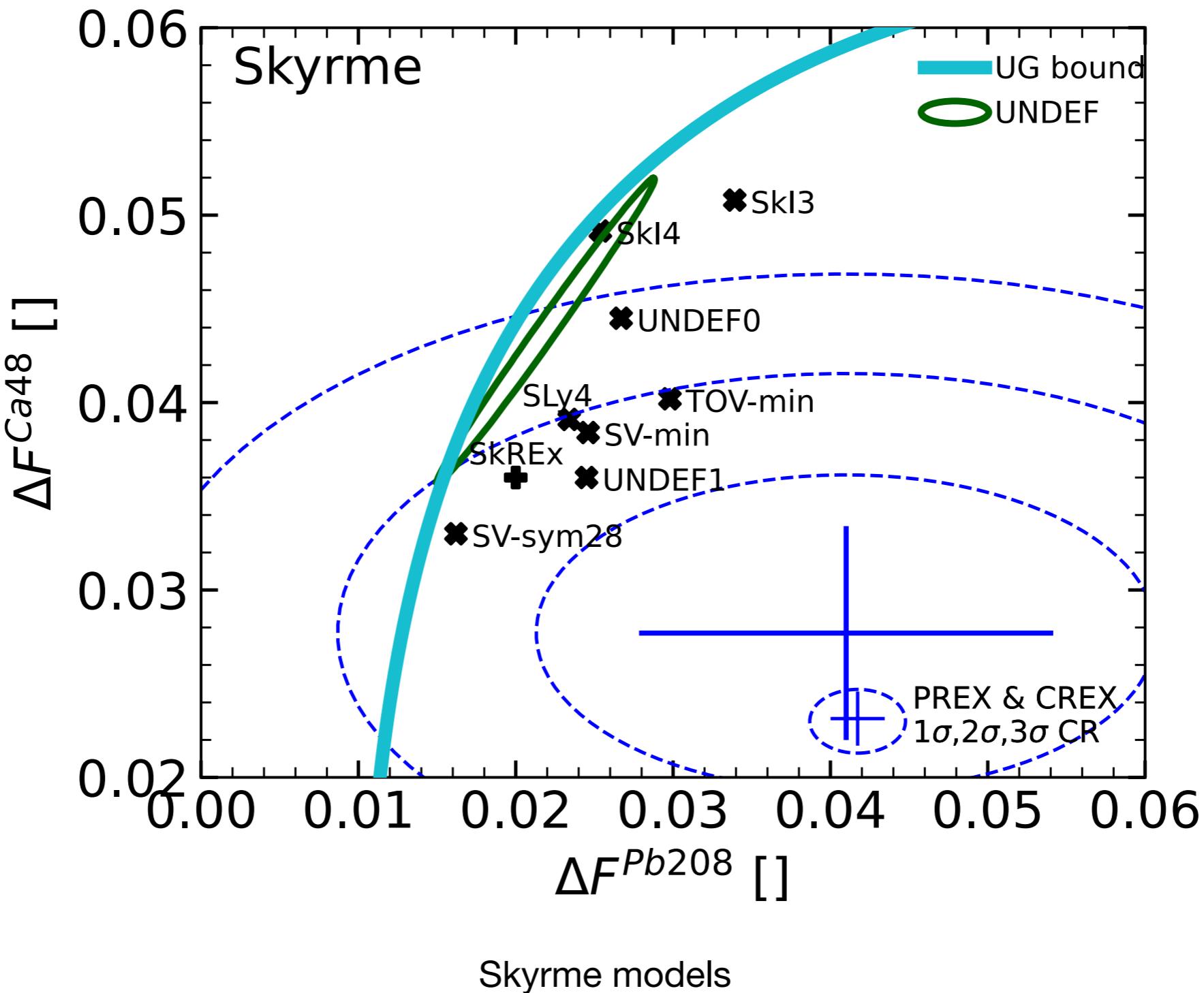


$$\psi \left(\frac{g_\delta}{2} \tau \cdot \delta \right) \psi$$

Post PREX-CREX era

- What nuclear properties can we learn from the experiment?
- Why are Skyrme models more compatible than RMF models?
- How may the mean-field model improve in the future?

Skyrme and RMF samples



Symmetry energy S_V

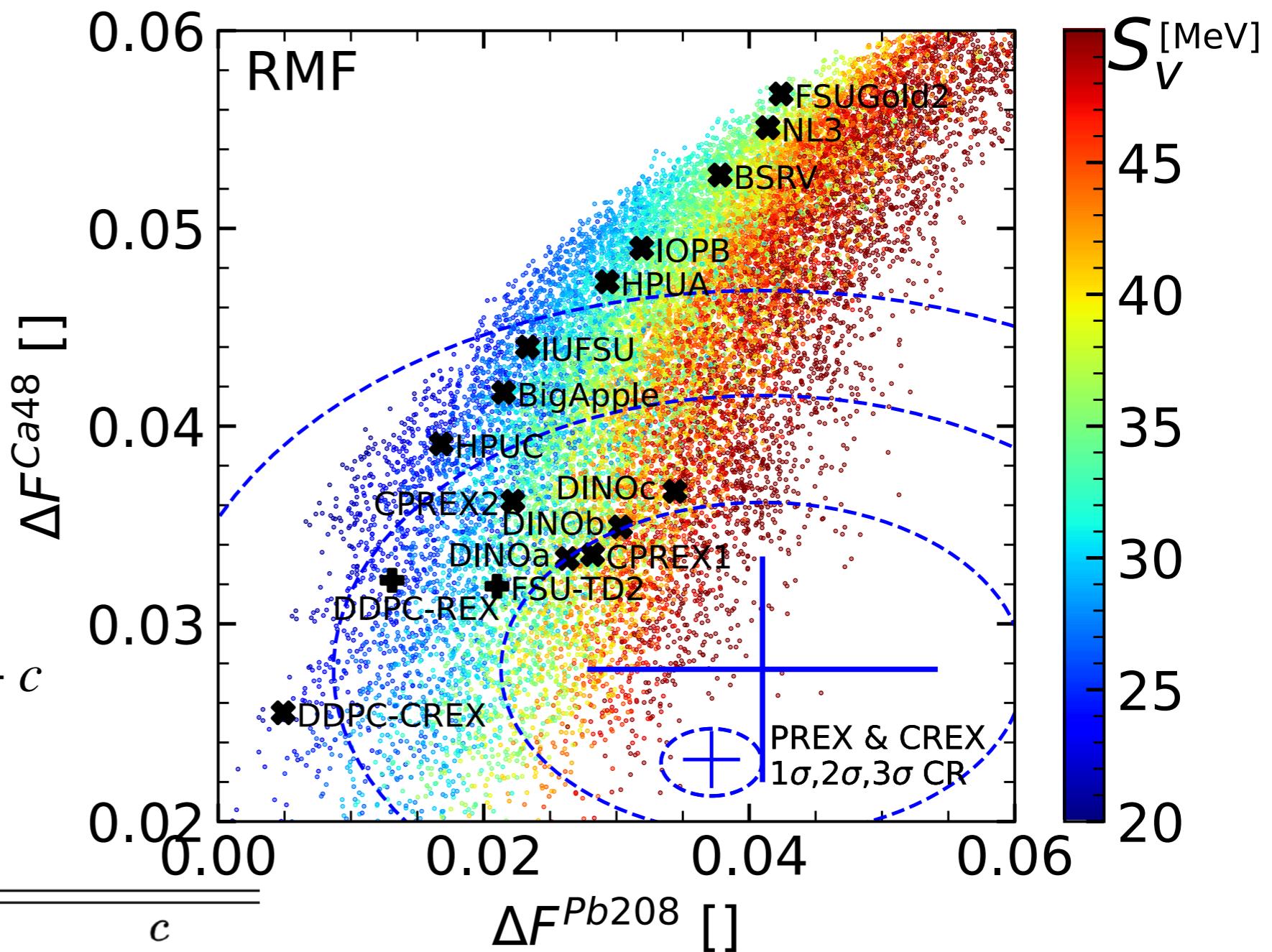
$$S(n) = S_V + \frac{L}{3} \left(\frac{n}{n_S} - 1 \right) + \dots$$

ΔF^{Ca48} and ΔF^{Pb208} are positively correlated for nuclear models with fixed S_V

The correlation is linear:

$$S_V = a\Delta F^{Ca48} + b\Delta F^{Pb208} + c$$

Fitting parameter for RMF (Skyrme) models:



| | a | b | c |
|--------|------------------|-----------------|----------------|
| RMF | -575.2 ± 5.1 | 916.3 ± 4.6 | 32.2 ± 3.7 |
| Skyrme | -503.2 ± 7.8 | 945.2 ± 5.5 | 31.9 ± 2.9 |

Symmetry energy slope L

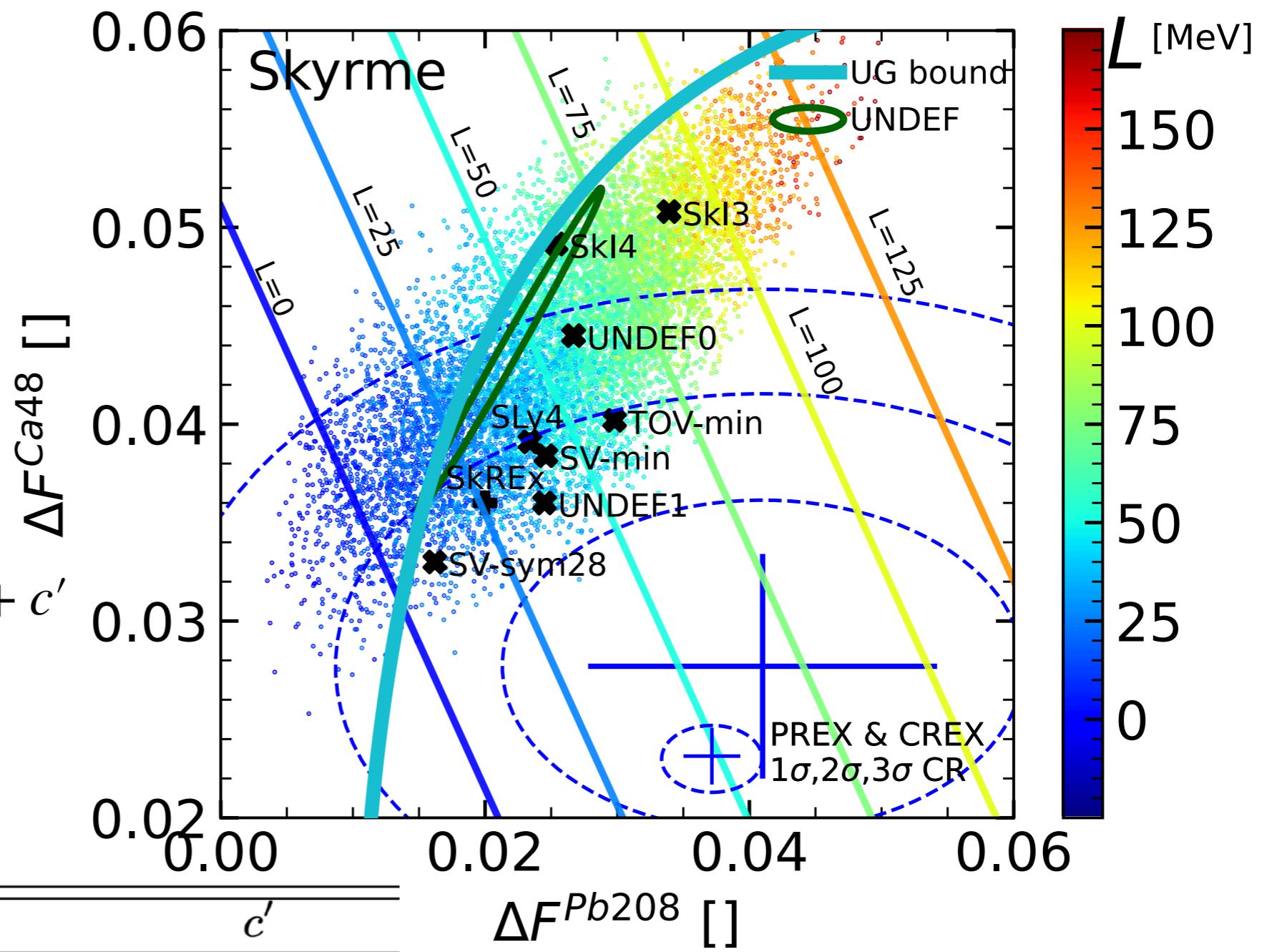
$$S(n) = S_V + \frac{L}{3} \left(\frac{n}{n_S} - 1 \right) + \dots$$

A similar correlation for L has an opposite slope!

The correlation is linear:

$$L = a' \Delta F^{Ca48} + b' \Delta F^{Pb208} + c'$$

Fitting parameter for RMF (Skyrme) models:



| | a' | b' | c' |
|--------|-------------------|-------------------|-------------------|
| RMF | 2938.7 ± 43.5 | 2420.6 ± 33.9 | -149.8 ± 25.6 |
| Skyrme | 1791.2 ± 27.2 | 2652.0 ± 19.0 | -91.5 ± 10.1 |

Linear correlation of form factor difference

Constraints on (S_V, L) from $(\Delta F^{Ca48}, \Delta F^{Pb208})$

S_V and L can be fixed by

ΔF^{Ca48} and(or) ΔF^{Pb208} :

$$S_V = a\Delta F^{Ca48} + b\Delta F^{Pb208} + c$$

$$L = a'\Delta F^{Ca48} + b'\Delta F^{Pb208} + c'$$

PREX:

$$\Delta F^{Pb208} = 0.041$$

± 0.013 (exp) ± 0.001 (theo)

CREX:

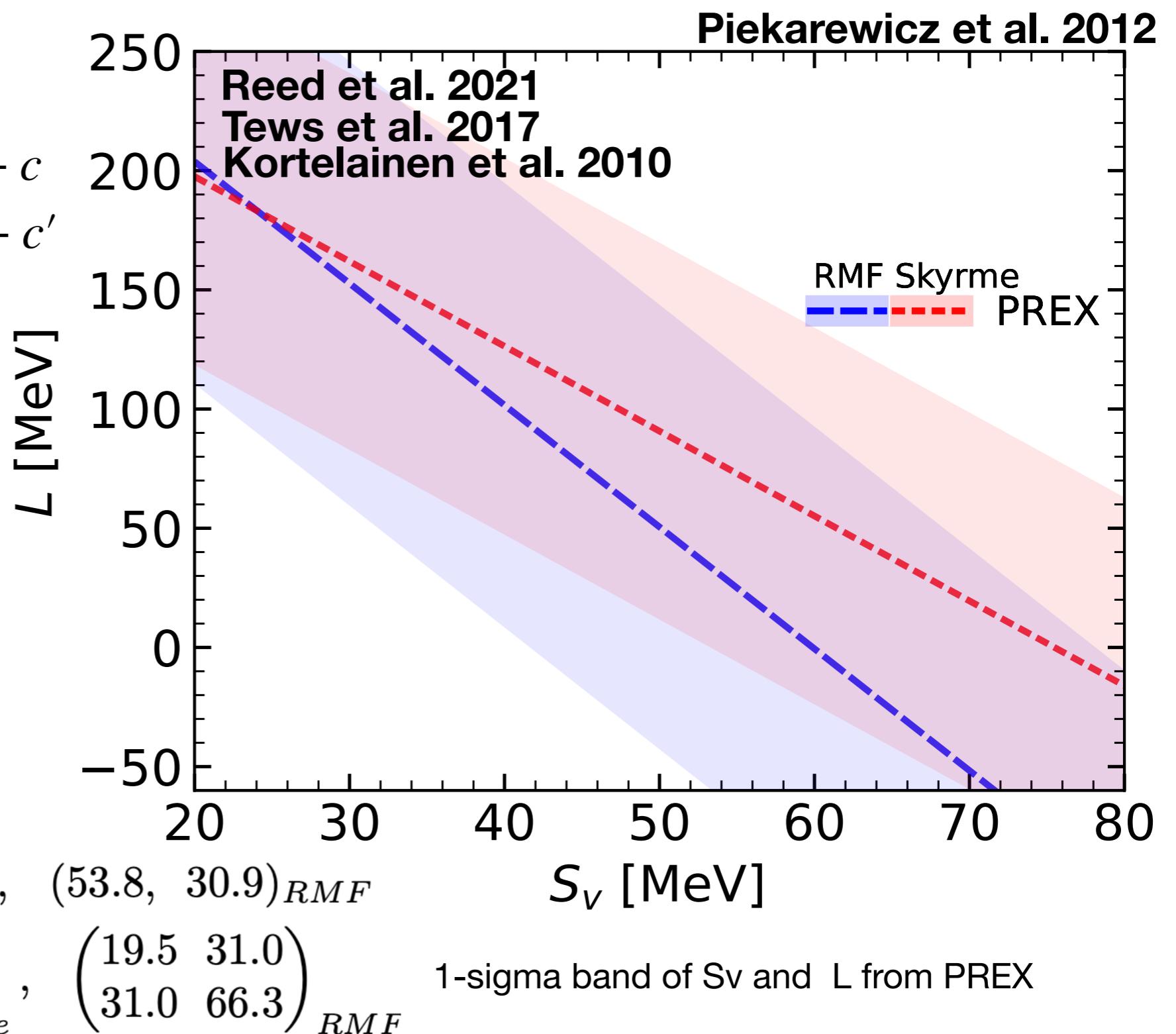
$$\Delta F^{Ca48} = 0.0277$$

± 0.0052 (stat) ± 0.002 (syst)

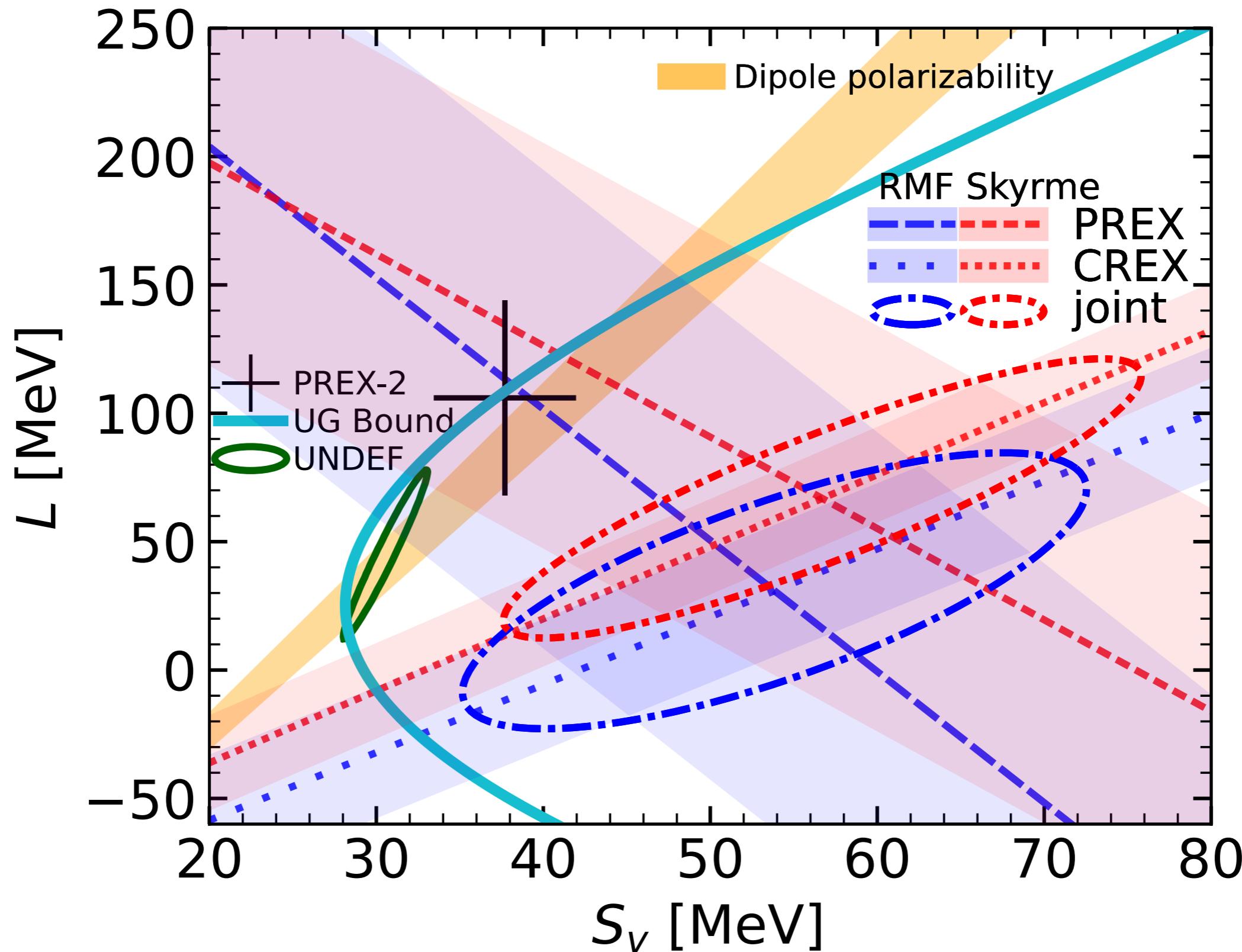
PREX+CREX:

$$(\bar{S}_V, \bar{L}) = (56.7, 66.8)_{Skyrme}, (53.8, 30.9)_{RMF}$$

$$\sqrt{\text{cov}} = \begin{pmatrix} 19.6 & 31.2 \\ 31.2 & 56.5 \end{pmatrix}_{Skyrme}, \begin{pmatrix} 19.5 & 31.0 \\ 31.0 & 66.3 \end{pmatrix}_{RMF}$$



Bayesian posterior



Isovector spin-orbit force

Isovector spin-orbit force is independent of S_V and L in Skyrme (not in RMF) model.

Spin-orbit force in Skyrme model:

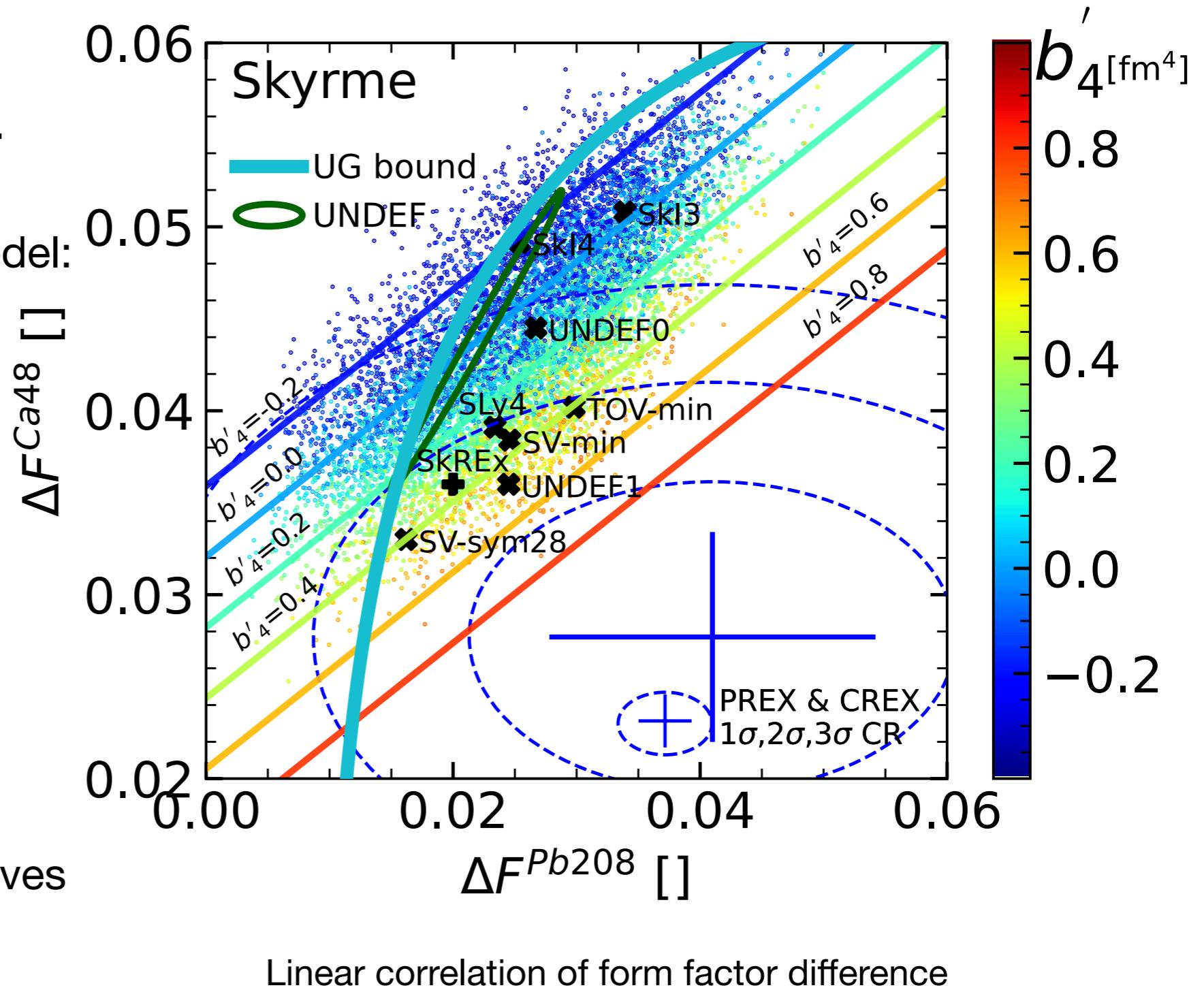
$$H_{SO} = b_4 \mathbf{J} \cdot \nabla \mathbf{n} + b'_4 (\mathbf{J}_n \cdot \nabla n_n + \mathbf{J}_p \cdot \nabla n_p)$$

The freedom b'_4 improves the Skyrme model performance.

$v \ll c$ limit of RMF model:

$$b'_4 \approx \frac{1}{8m^2} \left(\frac{g_\delta^2}{m_\delta^2} + \frac{g_\rho^2}{m_\rho^2} \right)$$

large δ -meson coupling improves the RMF models.



Impact of b'_4 on neutron skin ΔR_{np}

ΔR_{np} of ^{208}Pb increases with b'_4

ΔR_{np} of ^{48}Ca decreases with b'_4

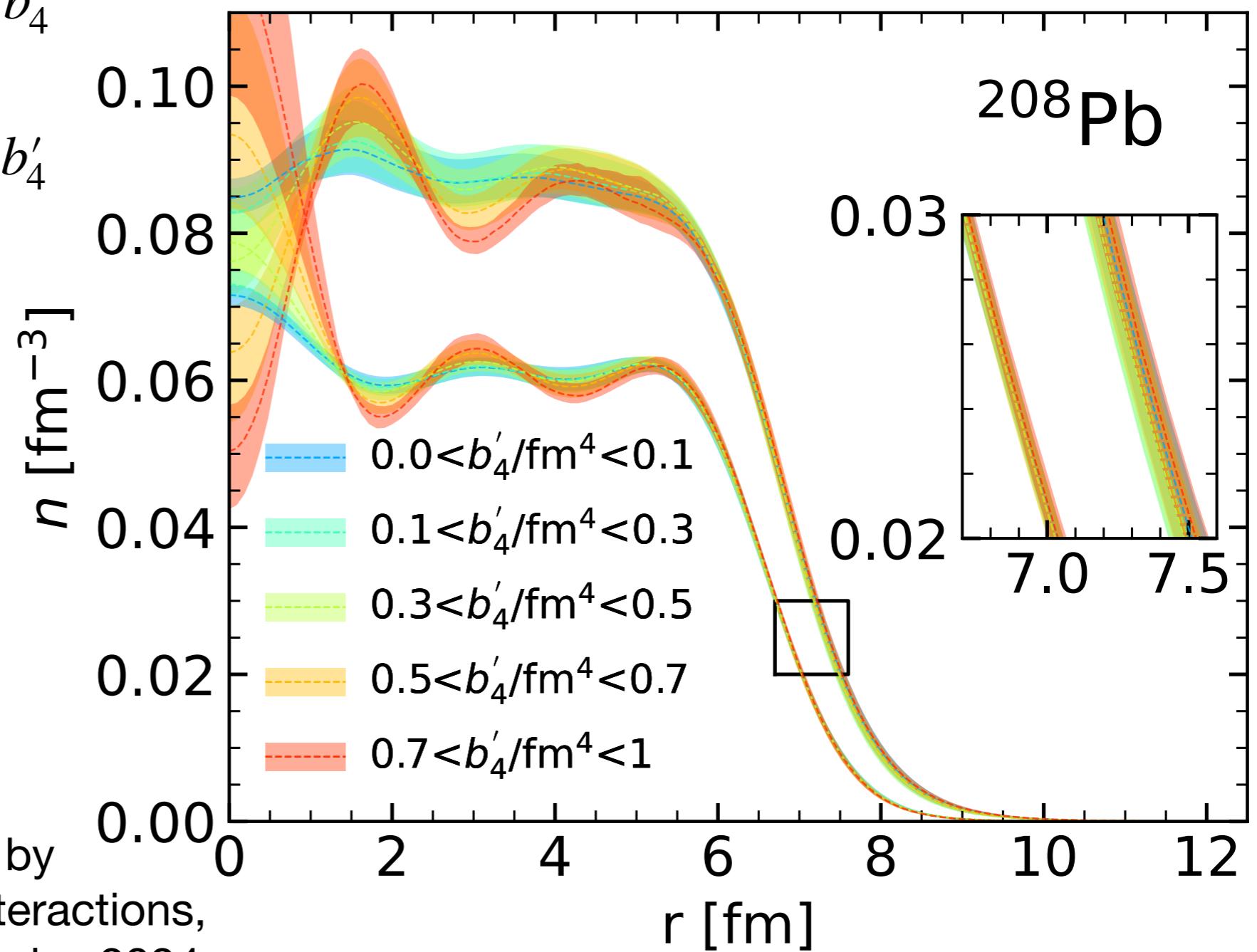
Large b'_4 reduces the tension between PREX and CREX.

90% lower bound of b'_4 :

$b'_4 \gtrsim 0.74 \text{ fm}^4$ (Skyrme)

$b'_4 \gtrsim 0.54 \text{ fm}^4$ (RMF)

The large density fluctuation inside nuclei may be reduced by introducing addition tensor interactions, see M. Salinas and J. Piekarewicz 2024
[arXiv:2312.13474](https://arxiv.org/abs/2312.13474)



Radial density profile of proton and neutron for ^{208}Pb

Free Tensor Interaction

Spin-orbit force in Skyrme model:

$$H_{SO} = b_4 \mathbf{J} \cdot \nabla n + b'_4 (\mathbf{J}_n \cdot \nabla n_n + \mathbf{J}_p \cdot \nabla n_p)$$

Tensor force in Skyrme model:

$$H_T = b_J \mathbf{J}^2 + b'_J (\mathbf{J}_n^2 + \mathbf{J}_p^2)$$

The freedom b'_4 , b'_4 and b'_4 improve the Skyrme model performance, see [arXiv.2406.03844](https://arxiv.org/abs/2406.03844):

$$S240 \text{ and } eS240: b'_4 = 0.6 \text{ fm}^{-4}$$

$$S500 \text{ and } eS500: b'_4 = 1.3 \text{ fm}^{-4}$$

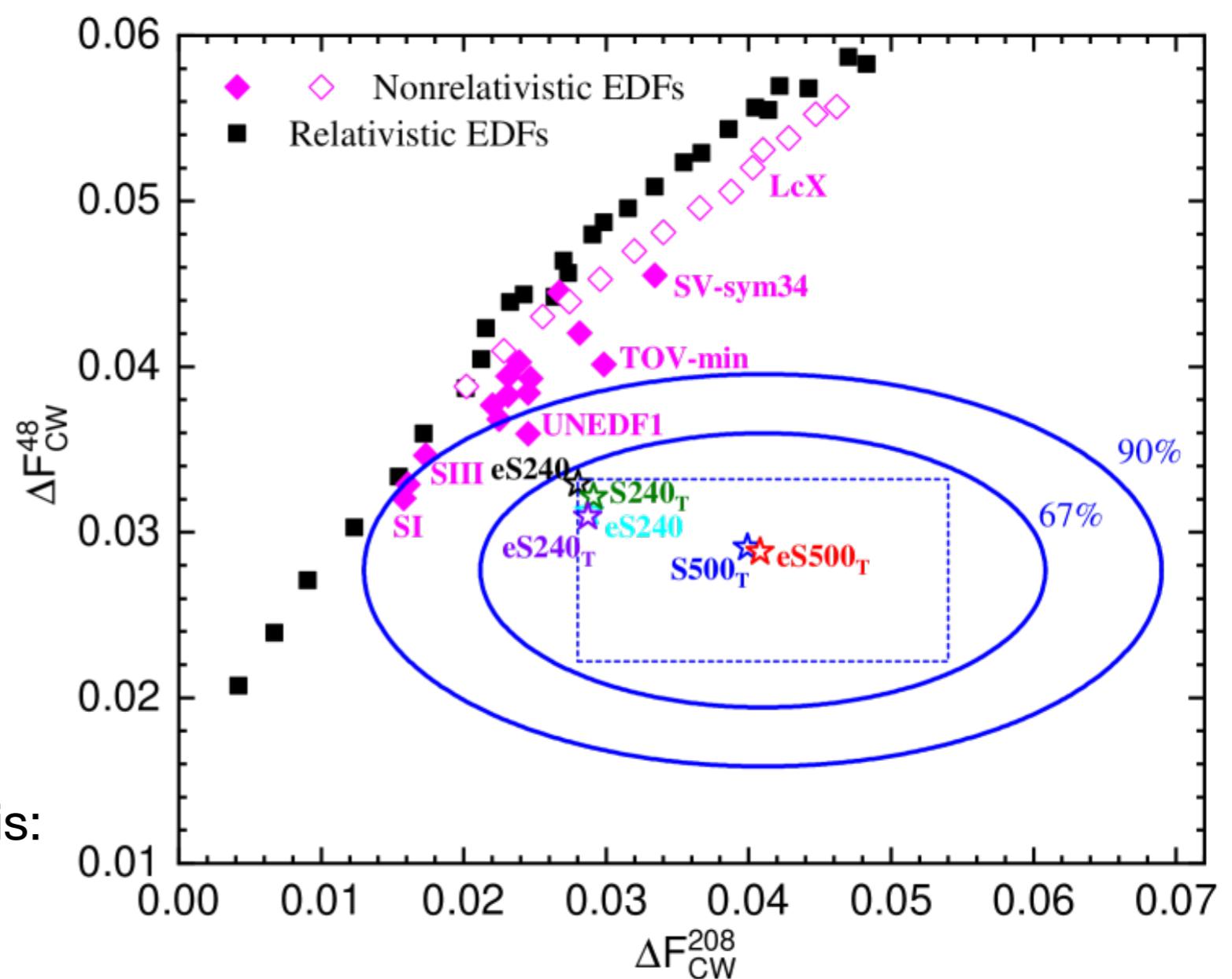
which is consistent with our analysis:

$$b'_4 = 1.37 \pm 0.49 \text{ fm}^{-4}$$

and 90% lower bound:

$$b'_4 \gtrsim 0.74 \text{ fm}^4 \text{ (Skyrme)}$$

$$b'_4 \gtrsim 0.54 \text{ fm}^4 \text{ (RMF)}$$

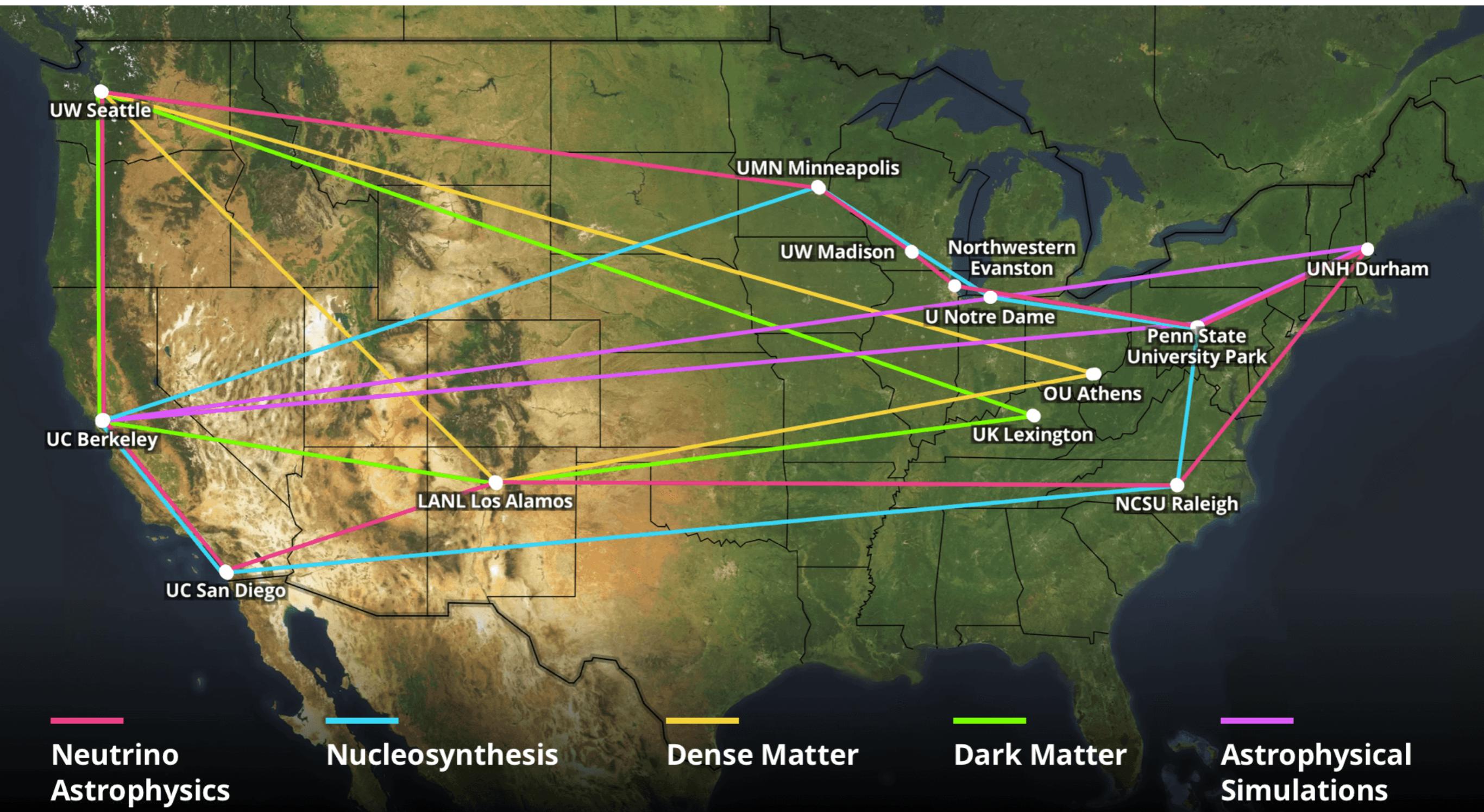


T.G. Yue, Z. Zhang, L.W. Chen arXiv.2406.03844

Post PREX-CREX era

- What nuclear properties can we learn from the experiment?
PREX+CREX prefers much Larger S_V than expected.
- Why are Skyrme models more compatible than RMF models?
The freedom in isovector spin-orbit interaction b'_4 .
- How may the mean-field model improve in the future?
Increase the degree of freedom on surface-related isovector interactions, e.g. isovector spin-orbit interaction, isovector tensor interaction.

Thanks to N3AS's support



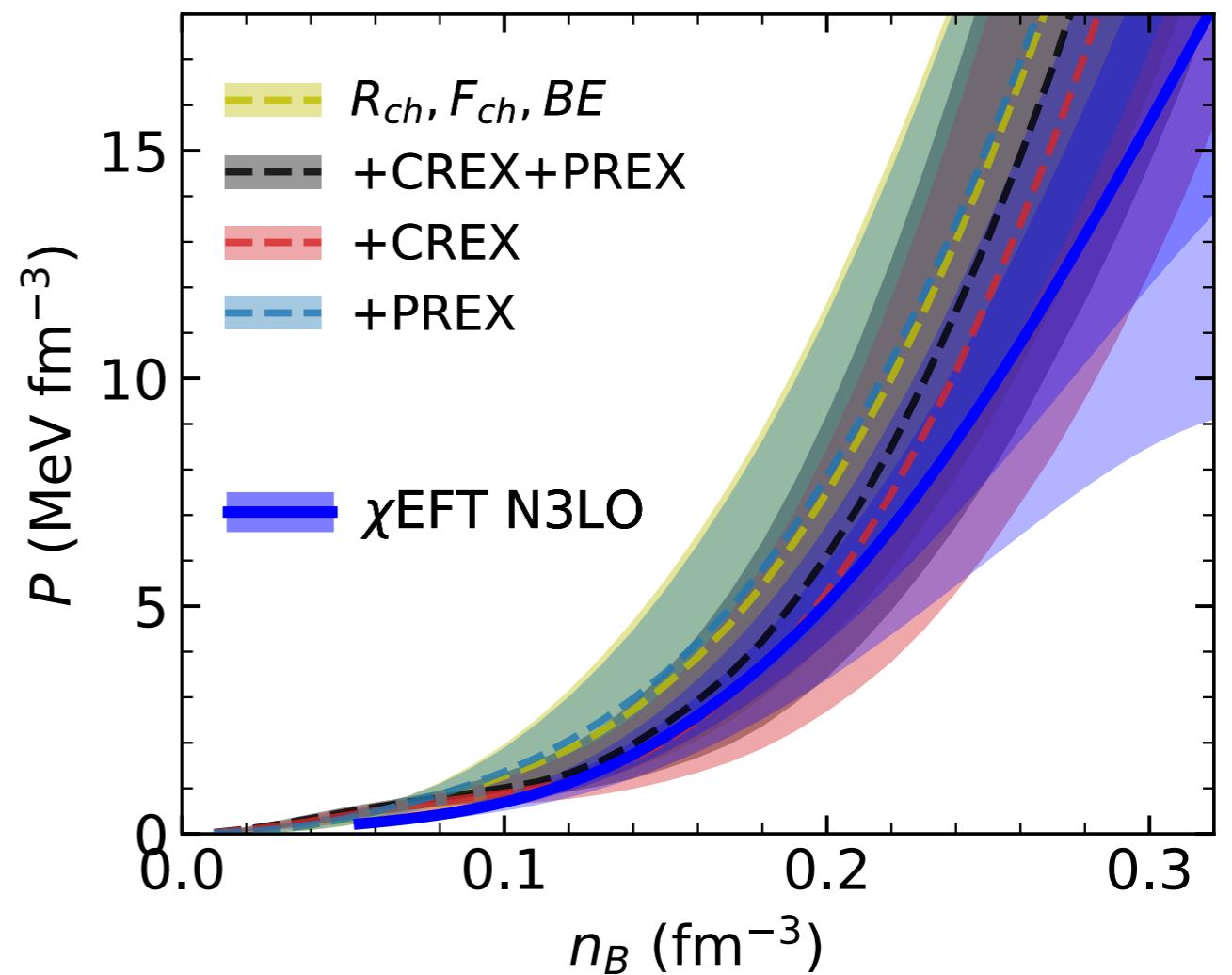
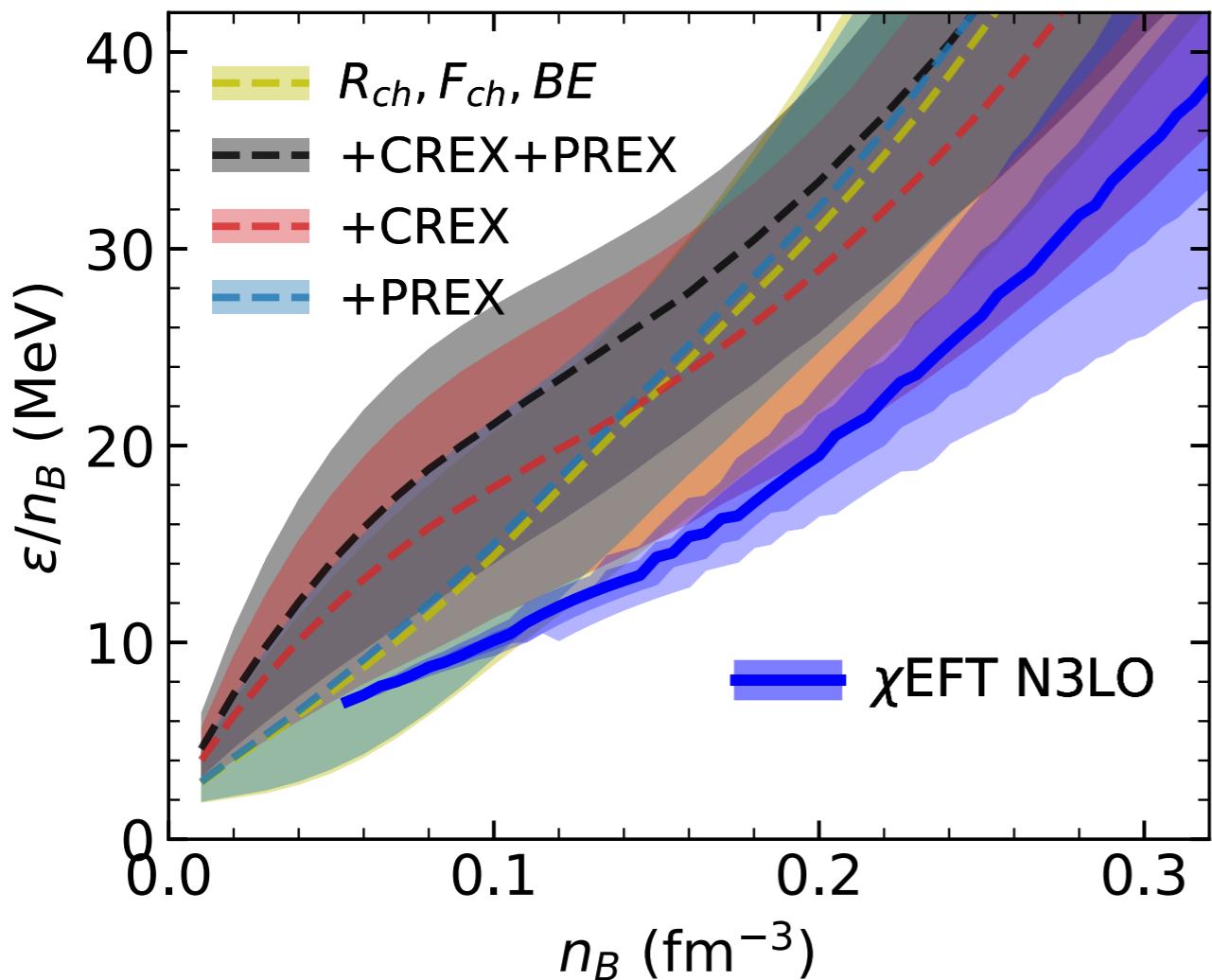
| | Experiment | NL3 | FSU2 | IOPB-I | IUFSU | BigApple | HPUC | BSRV | DINOa | DINOb | DINOc | CPREX1 | CPREX2 |
|-------------------|----------------------|---------------------|--------|--------|--------|----------|--------|--------|--------|--------|--------|--------|--------|
| ²⁰⁸ Pb | B/A [MeV] | 7.87 | 7.88 | 7.87 | 7.86 | 7.88 | 7.85 | 7.85 | 7.84 | 7.87 | 7.87 | 7.84 | 7.86 |
| | R_{ch} [fm] | 5.50 | 5.51 | 5.49 | 5.52 | 5.49 | 5.50 | 5.56 | 5.53 | 5.51 | 5.51 | 5.49 | 5.49 |
| | ΔR_{np} [fm] | 0.159 ± 0.017 | 0.2797 | 0.2862 | 0.2195 | 0.1618 | 0.1508 | 0.1196 | 0.2595 | 0.1746 | 0.1993 | 0.2235 | 0.1905 |
| | F_{ch} [] | 0.409 | 0.4067 | 0.4094 | 0.4052 | 0.4106 | 0.4080 | 0.3992 | 0.4043 | 0.4074 | 0.4075 | 0.4073 | 0.4100 |
| | ΔF [] | 0.041 ± 0.013 | 0.0414 | 0.0423 | 0.0319 | 0.0233 | 0.0214 | 0.0168 | 0.0378 | 0.0262 | 0.0303 | 0.0342 | 0.0282 |
| ⁴⁸ Ca | B/A [MeV] | 8.67 | 8.65 | 8.62 | 8.64 | 8.53 | 8.52 | 8.65 | 8.66 | 8.67 | 8.67 | 8.64 | 8.66 |
| | R_{ch} [fm] | 3.48 | 3.45 | 3.43 | 3.45 | 3.44 | 3.46 | 3.46 | 3.44 | 3.47 | 3.47 | 3.48 | 3.46 |
| | ΔR_{np} [fm] | 0.137 ± 0.015 | 0.2255 | 0.2318 | 0.1995 | 0.1736 | 0.1690 | 0.1479 | 0.2196 | 0.0994 | 0.1054 | 0.1141 | 0.1252 |
| | F_{ch} [] | 0.158 | 0.1604 | 0.1665 | 0.1616 | 0.1647 | 0.1582 | 0.1577 | 0.1621 | 0.1591 | 0.1589 | 0.1585 | 0.1537 |
| | ΔF [] | 0.0277 ± 0.0055 | 0.0551 | 0.0564 | 0.0490 | 0.0435 | 0.0413 | 0.0391 | 0.0527 | 0.0330 | 0.0345 | 0.0364 | 0.0335 |

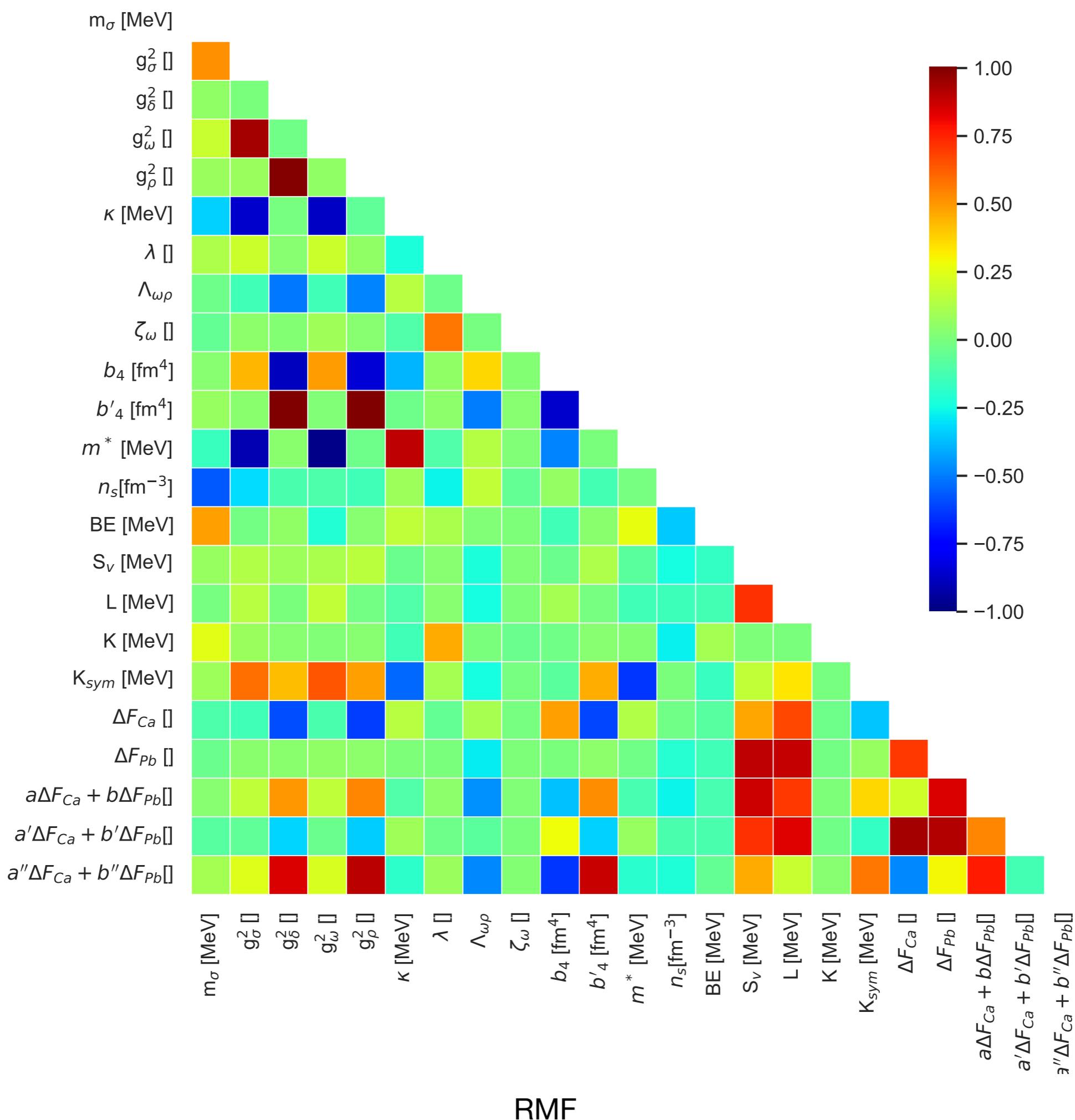
TABLE I. Experimental data for the binding energy per nucleon[1], charge radii[2], neutron skins (excluding PREX and CREX)[3], charge from factor and form factor difference from PREX[4] for ²⁰⁸Pb and CREX[5] for ⁴⁸Ca. Also displayed are the theoretical results obtained with NL3[6], FSUGold2[7], IOPB-I[8], IUFSU[9], BigApple[10], HPUC[11], BSRV[12], DINOa-c[13] and the two new parameterizations, CPREX1 and CPREX2.

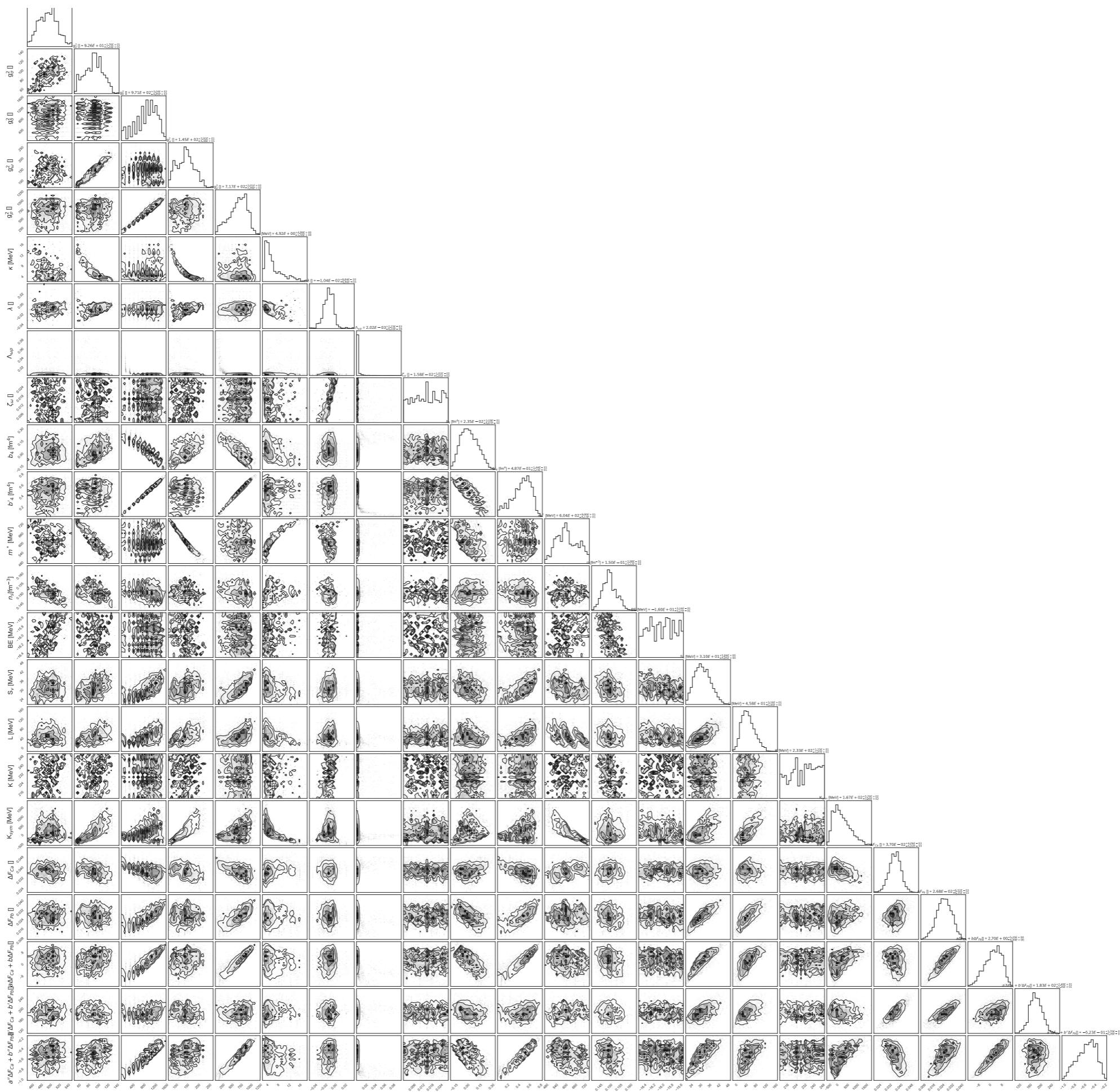
| | NL3 | FSU2 | IOPB-I | IUFSU | BigApple | HPUC | BSRV | DINOa | DINOb | DINOc | CPREX1 | CPREX2 | |
|--------------------|---------------------------|--------|--------|--------|----------|--------|--------|--------|--------|--------|--------|--------|--------|
| SNM | n_s [fm ⁻³] | 0.1483 | 0.1504 | 0.1495 | 0.1546 | 0.1546 | 0.1490 | 0.1480 | 0.1522 | 0.1525 | 0.1519 | 0.1516 | 0.1518 |
| | M^* [MeV] | 558.7 | 557.0 | 557.2 | 572.1 | 572.8 | 572.9 | 565.3 | 587.4 | 593.0 | 593.9 | 692.8 | 648.1 |
| | B [MeV] | 16.24 | 16.26 | 16.10 | 16.40 | 16.34 | 15.98 | 16.10 | 16.16 | 16.21 | 16.21 | 16.29 | 16.14 |
| | K [MeV] | 271.6 | 237.7 | 222.6 | 231.3 | 227.0 | 220.2 | 227.2 | 210.0 | 207.0 | 206.0 | 223.8 | 223.5 |
| | S_V [MeV] | 37.3 | 37.6 | 33.3 | 31.3 | 31.3 | 28.4 | 36.1 | 31.4 | 33.1 | 34.6 | 32.9 | 29.8 |
| PNM | L [MeV] | 118.2 | 112.7 | 63.6 | 47.2 | 39.8 | 41.6 | 84.6 | 50.0 | 70.0 | 90.0 | -3.5 | 0.4 |
| | K_{sym} [MeV] | 101.0 | 25.4 | -37.0 | 28.5 | 87.5 | 81.1 | -73.2 | 506.0 | 609.1 | 714.8 | -418.4 | -239.8 |
| | M_n^* [MeV] | 569.2 | 566.0 | 566.7 | 580.5 | 582.8 | 581.4 | 573.3 | 352.1 | 333.0 | 320.5 | 377.4 | 465.6 |
| | M_p^* [MeV] | 569.2 | 566.0 | 566.7 | 580.5 | 582.8 | 581.4 | 574.8 | 908.8 | 948.2 | 969.1 | 1062.5 | 870.1 |
| | S_V [MeV] | 38.3 | 38.6 | 34.7 | 32.9 | 33.1 | 29.9 | 37.2 | 46.5 | 50.6 | 53.4 | 54.3 | 38.4 |
| NS | L [MeV] | 121.2 | 115.9 | 67.7 | 49.5 | 40.6 | 42.7 | 88.7 | 172.1 | 216.4 | 247.8 | 211.2 | 75.9 |
| | K_{sym} [MeV] | 100.3 | 27.2 | -45.5 | 23.1 | 74.3 | 89.2 | -70.6 | 726.7 | 907.2 | 1021.2 | 801.8 | 76.4 |
| | M_{max} [M_\odot] | 2.77 | 2.07 | 2.15 | 1.94 | 2.60 | 2.05 | 2.04 | 2.17 | 2.15 | 2.15 | 2.04 | 2.12 |
| | $R_{1.0}$ [km] | 14.4 | 14.1 | 13.2 | 12.6 | 12.8 | 12.6 | 13.6 | 14.4 | 14.8 | 15.1 | 13.9 | 12.9 |
| | $R_{1.4}$ [km] | 14.5 | 13.9 | 13.2 | 12.6 | 13.1 | 12.8 | 13.4 | 14.4 | 14.6 | 14.9 | 13.4 | 12.9 |
| $\Lambda_{1.0}$ [] | 7797 | 6473 | 4347 | 3384 | 3918 | 3752 | 4903 | 6623 | 7572 | 8579 | 4543 | 3544 | |
| | 1275 | 876 | 687 | 500 | 719 | 593 | 689 | 1065 | 1150 | 1256 | 584 | 570 | |

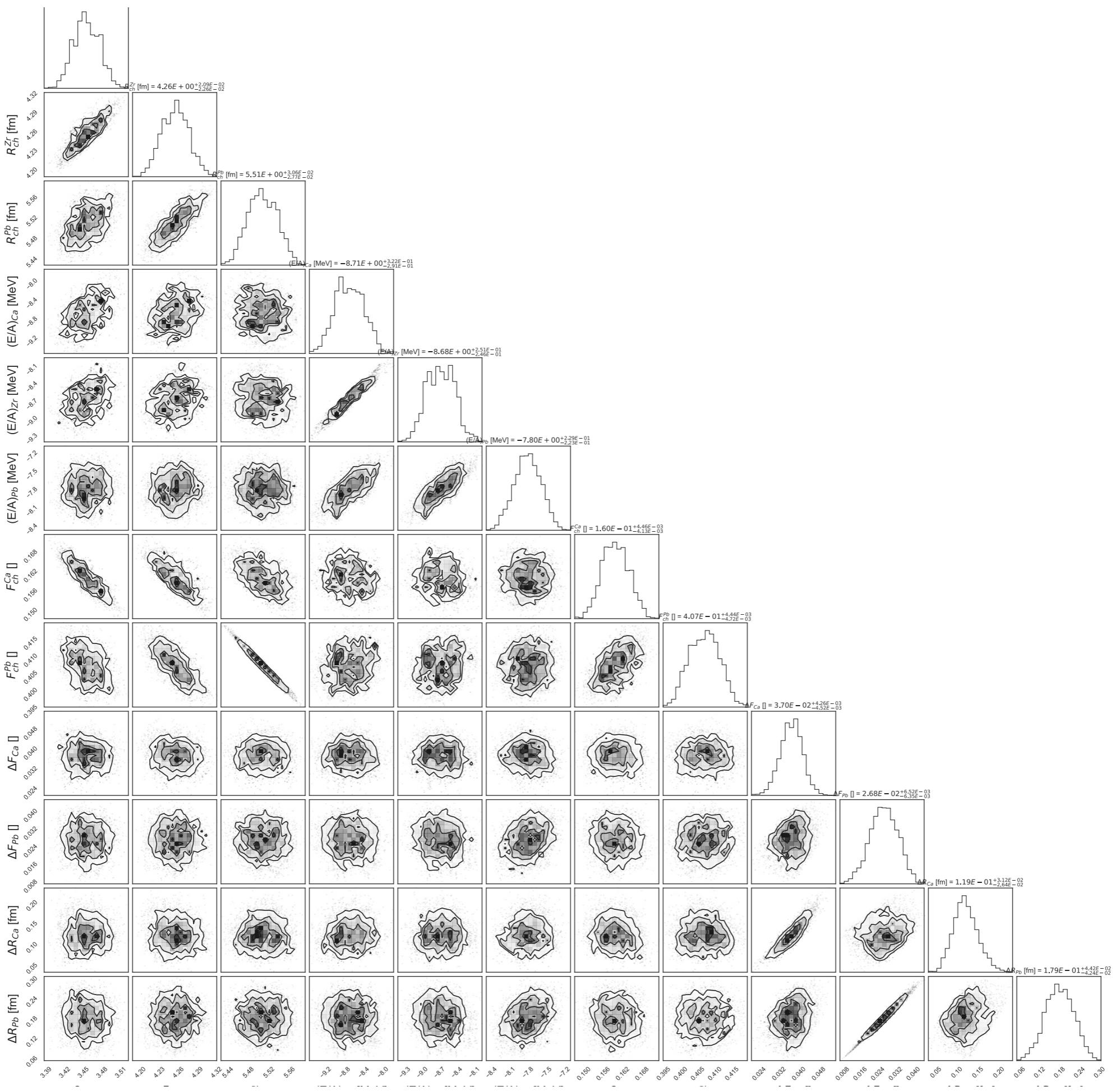
TABLE II. Saturation properties and neutron star properties of RMF models listed in Table I. Saturation properties for SNM and PNM are defined in the letter. Neutron star properties are calculated with the crust EOSs constructed with the compressible liquid droplet model repectively for various RMF models with fixed surface tension parameters $\sigma_s = 1.2$ MeV fm⁻², $S_S = 48$ MeV[14].

Neutron star EOS









RMF