

# Light dark sector and core-collapse supernovae

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Dark Matter in Compact Objects, Stars, and Low-Energy Experiments  
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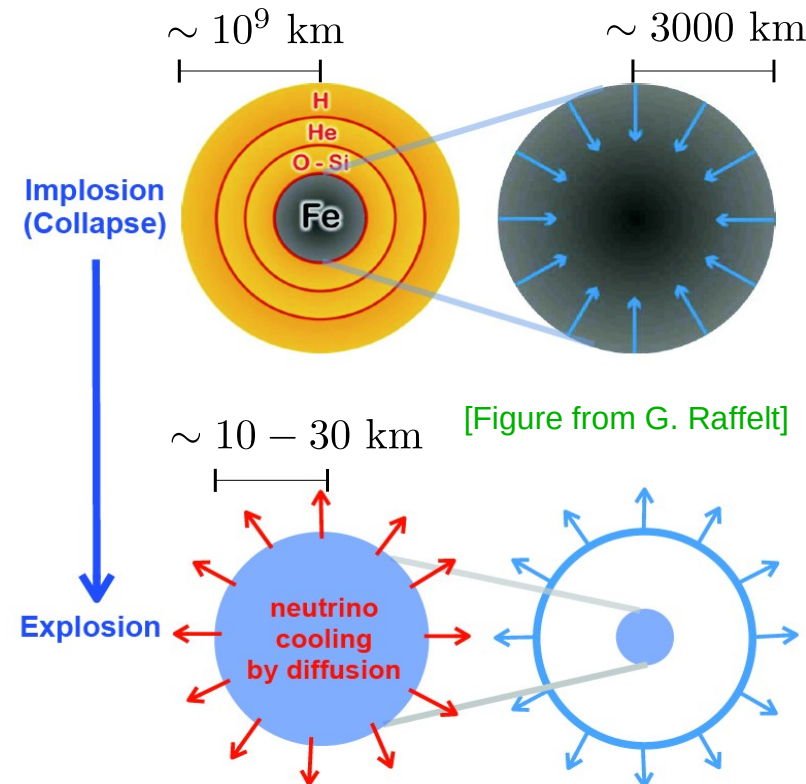
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# Core-collapse supernovae

- the death of massive stars  $\gtrsim 8 M_{\odot}$
- luminosity  $\simeq 10^9 L_{\odot}$  for  $\sim \mathcal{O}(100)$  days  
( $E_{\gamma} \sim 10^{49}$  erg)
- explosion energy  $\sim 10^{51}$  erg  $\equiv 1$  B(ethe)
- strong MeV neutrino emission  $\sim 10^{53}$  erg within  $\sim 10$  s ( $\sim 10^{58}$  neutrinos)

The high density ( $\rho_c \gtrsim 10^{14} \text{g cm}^{-3}$ ) and temperature  $T_c \gtrsim 30$  MeV of the proto-neutron stars make them interesting astrophysical “laboratory” complementary to terrestrial experiments

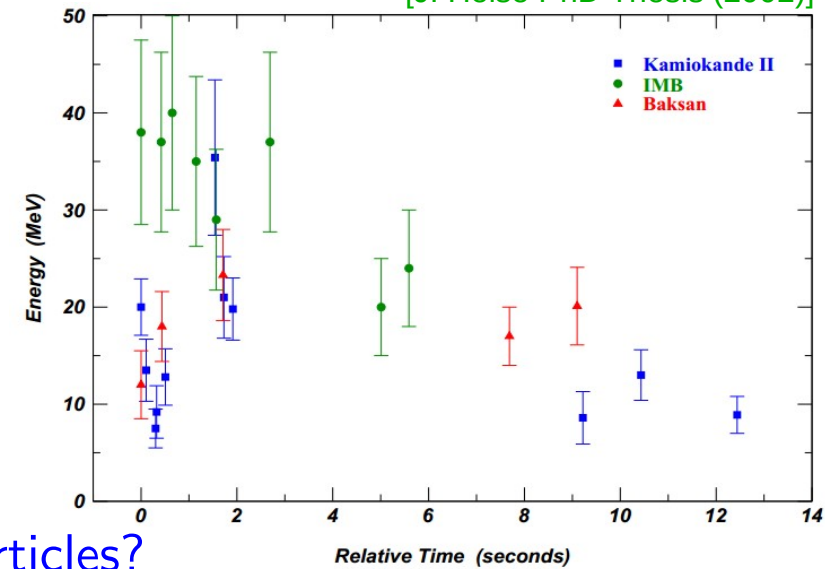
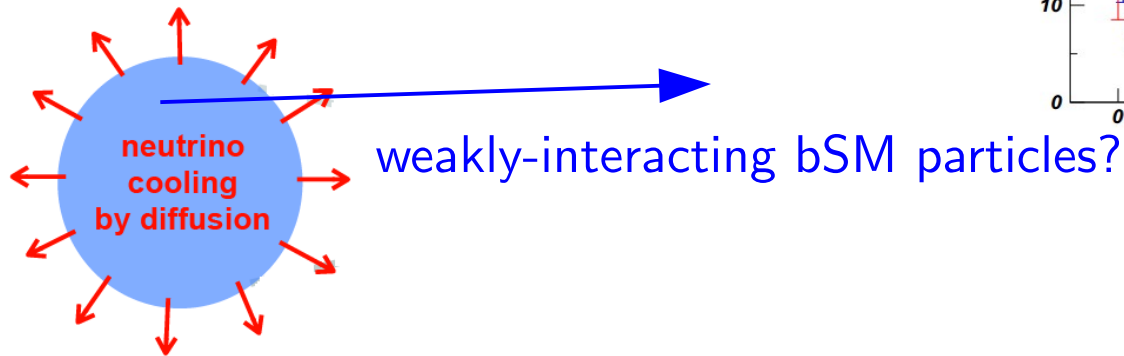
SN1987a (From AAO website)



## Supernova cooling bound

$\sim 20$  SN  $\bar{\nu}_e$  detected from SN1987a

$$L_{\bar{\nu}_e} \sim 5 \times 10^{52} \text{ erg}, \langle E_{\bar{\nu}_e} \rangle \sim 15 \text{ MeV}$$



If the energy carried away by any bSM particles is larger than that by neutrinos predicted within the Standard Model, the cooling timescale of the proto-neutron star becomes shorter than  $\sim 10$  s

Raffelt's criteria:  $L_{\text{new particle}} < L_{\nu} \sim 3 \times 10^{52} \text{ erg/s}$

axions, keV sterile neutrinos, dark photons,...

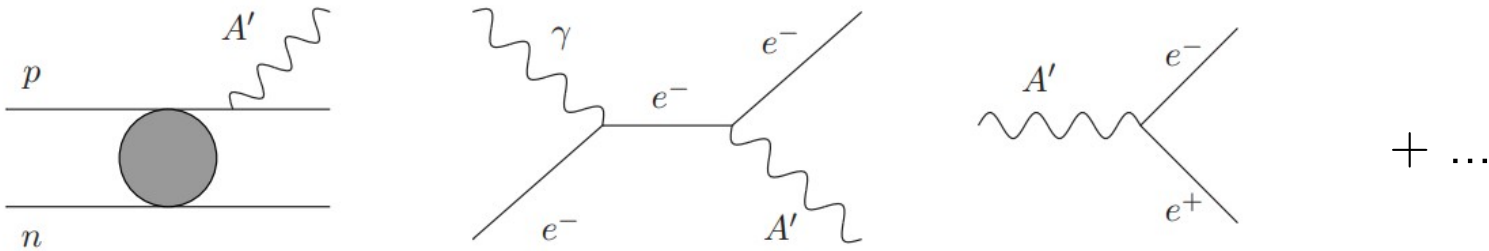
# Cooling constraint on dark photon

For dark photon portal dark sector:

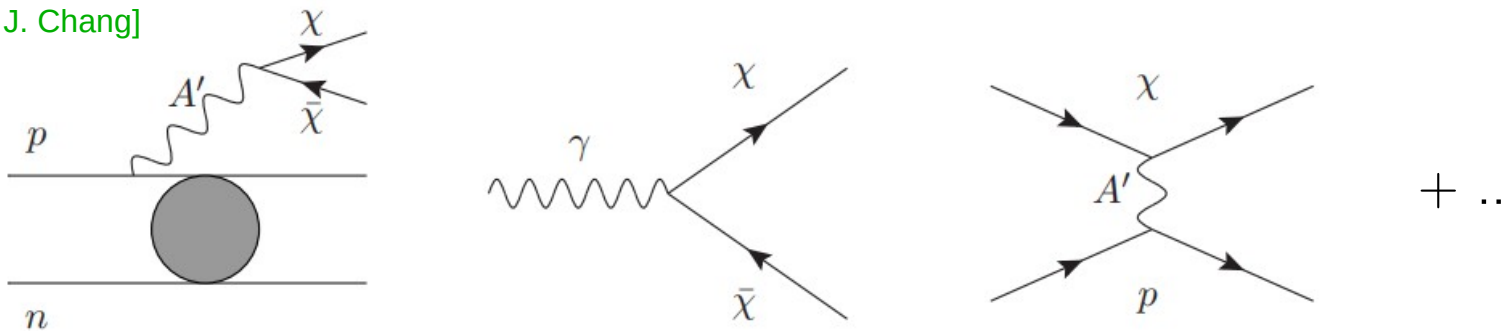
$$\mathcal{L} \supset -\frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} - \frac{\epsilon}{2}F'_{\mu\nu}F^{\mu\nu} + \frac{1}{2}m_{A'}^2 A'_\mu A'^\mu + \bar{\chi}(i\gamma^\mu \partial_\mu - m_\chi)\chi + g_D \bar{\chi}\gamma^\mu A'_\mu \chi$$

[Holdom 1986, Okun 1982, and many others]

They can be produced in PNS via a number of processes, e.g.,



[Adapted from J. Chang]



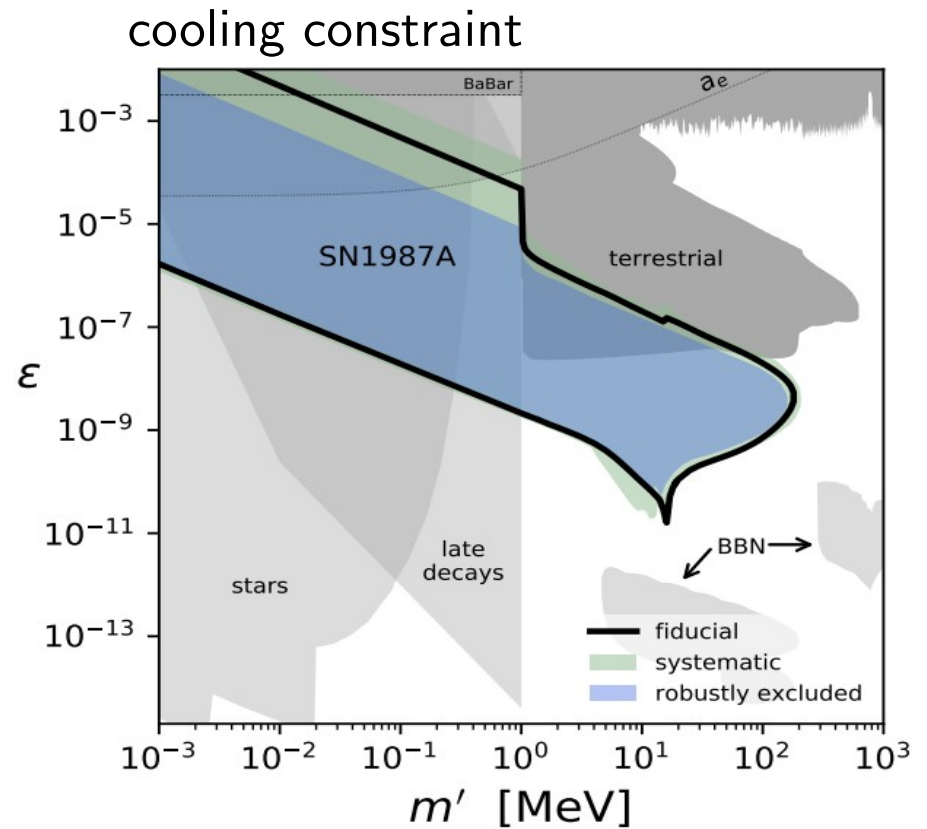
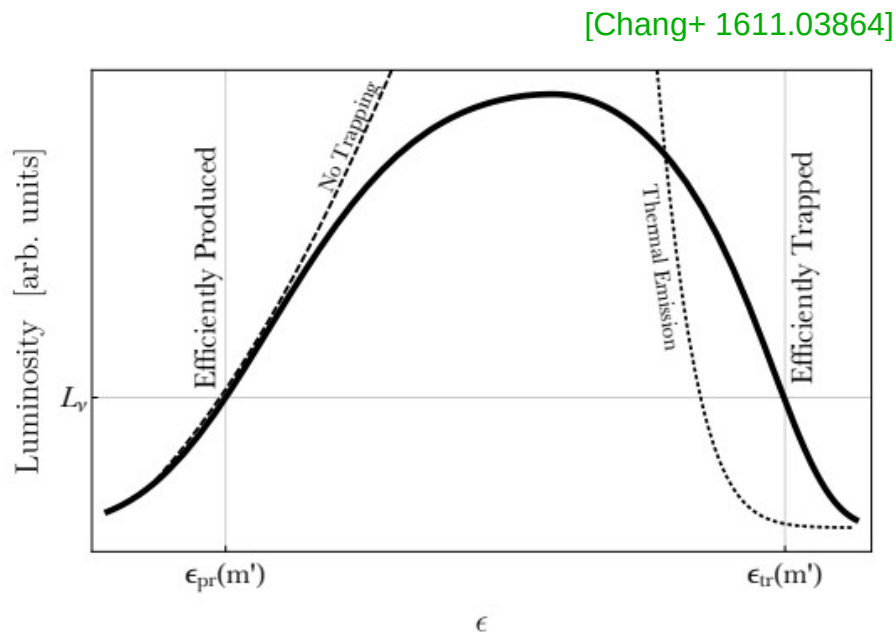
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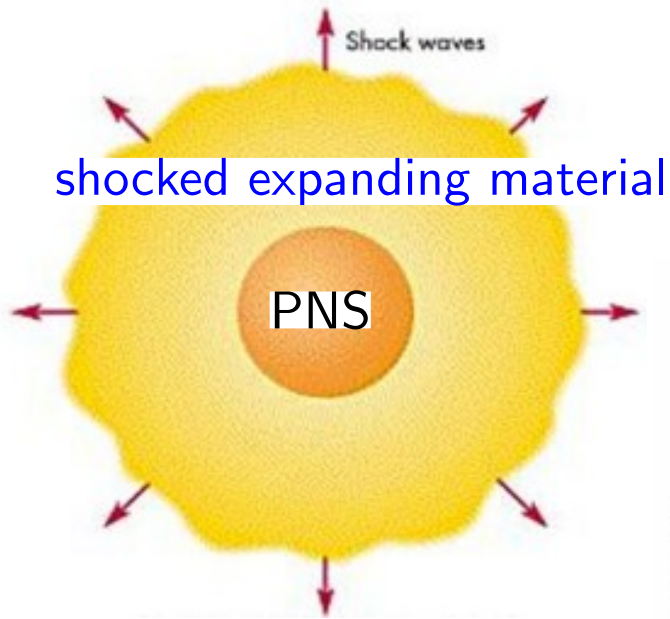
[Holdom 1986, Okun 1982, and many others]

Considering the degree of freedom of dark photon only:



[See also Dent+, Rrapaj+, Hardy+,...]

# Supernova ejecta and explosion energy



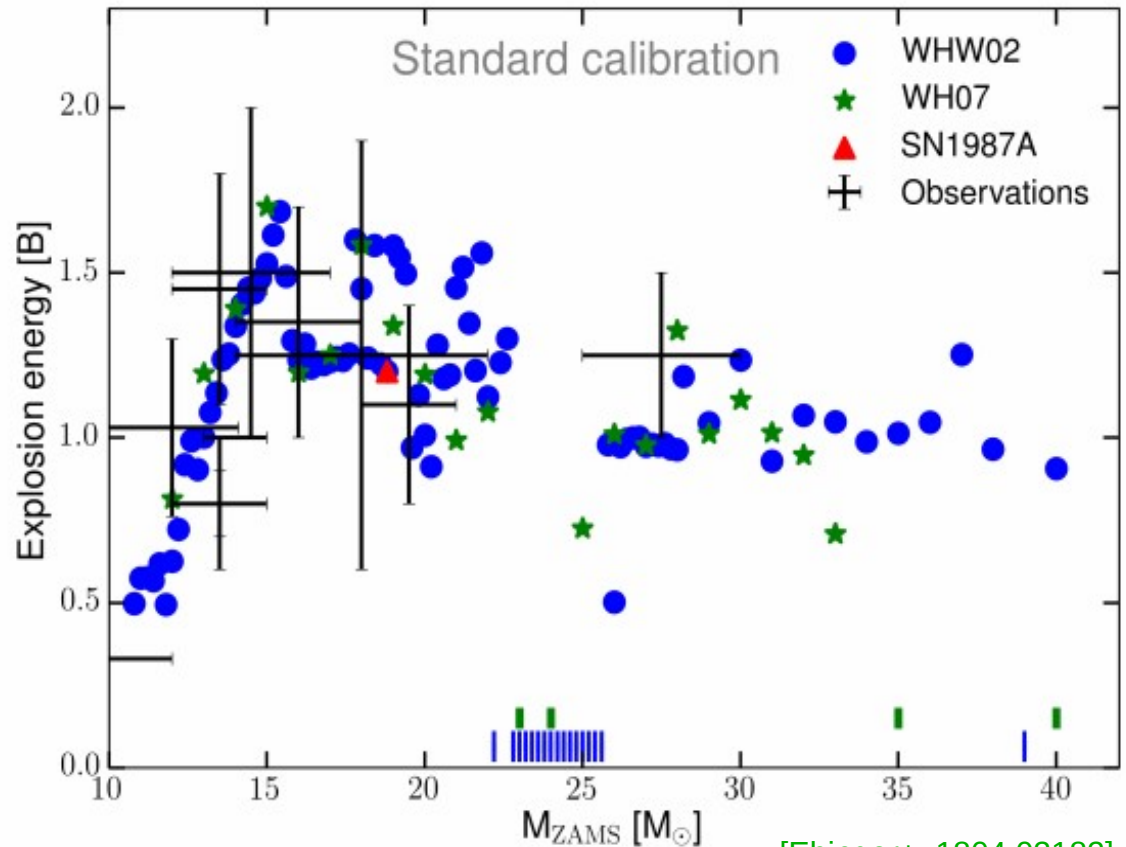
$$M_{\text{ej}} \sim 10 M_{\odot}, \quad v_{\text{ej}} \sim 0.01 c$$

$$E_{\text{expl}} \sim \text{K.E. of ejecta}$$

$$\sim \frac{1}{2} M_{\text{ej}} v_{\text{ej}}^2$$

$$\sim 10^{51} \text{ erg} \equiv 1 \text{ B(ethe)}$$

more observations than just SN1987a!



[Ebinger+, 1804.03182]

# Explosion energy constraints on radiative dark particle

[A. Sung, H. Tu, MRW, PRD 99 (2019) 121305, arXiv: 1903.07923]

The gravitational binding energy of the stellar envelope outside the proto-neutron star radius is

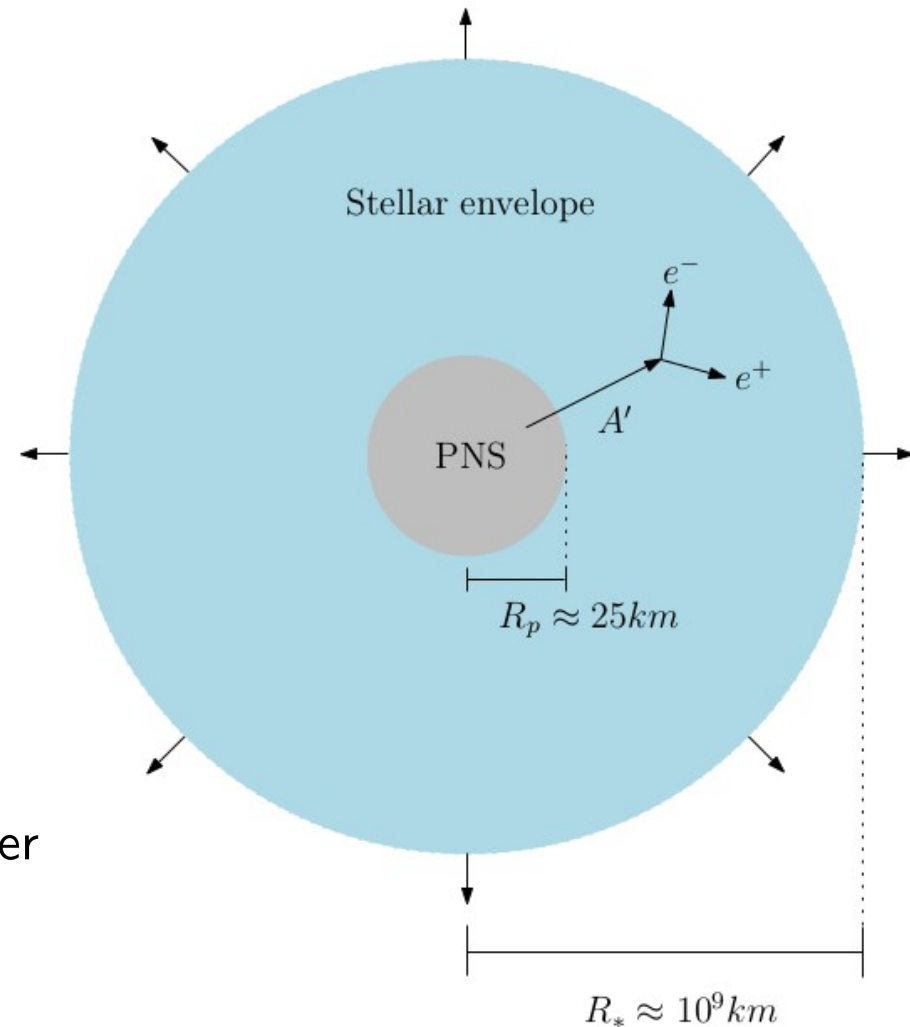
$$E_b \lesssim 10^{51} \text{ erg}$$

The observed explosion energy is

$$E_{\text{expl}} \lesssim 2 \times 10^{51} \text{ erg}$$

→ new particles that can deposit energy more than  $\simeq 3 \times 10^{51}$  erg into the stellar envelope & expanding material is NOT consistent with observation

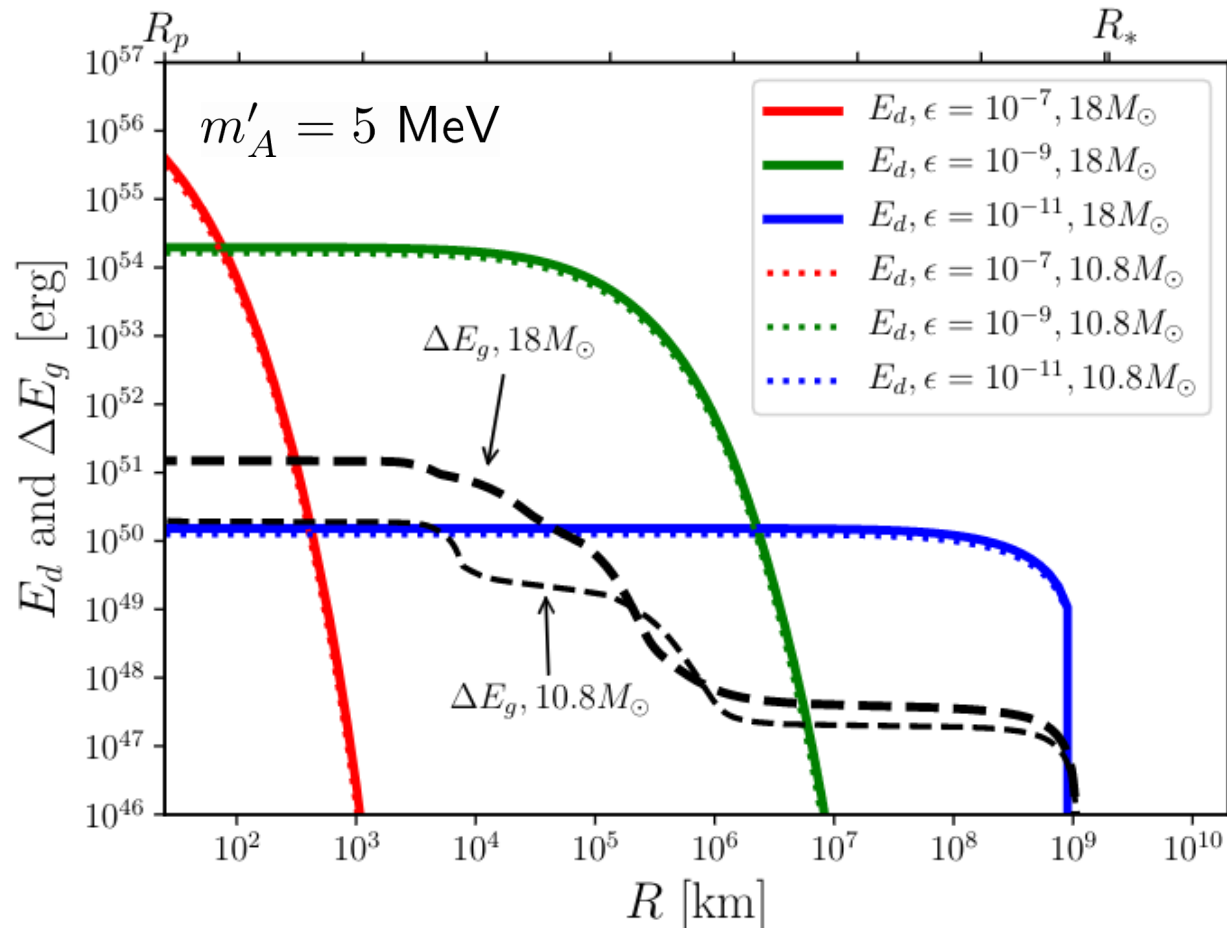
→ can improve the constraint on exotic particle to smaller coupling by  $\sim$  one order of magnitude



# Decay of dark photons outside PNS

$E_d$ : cumulative energy deposition above a stellar radius  $R$

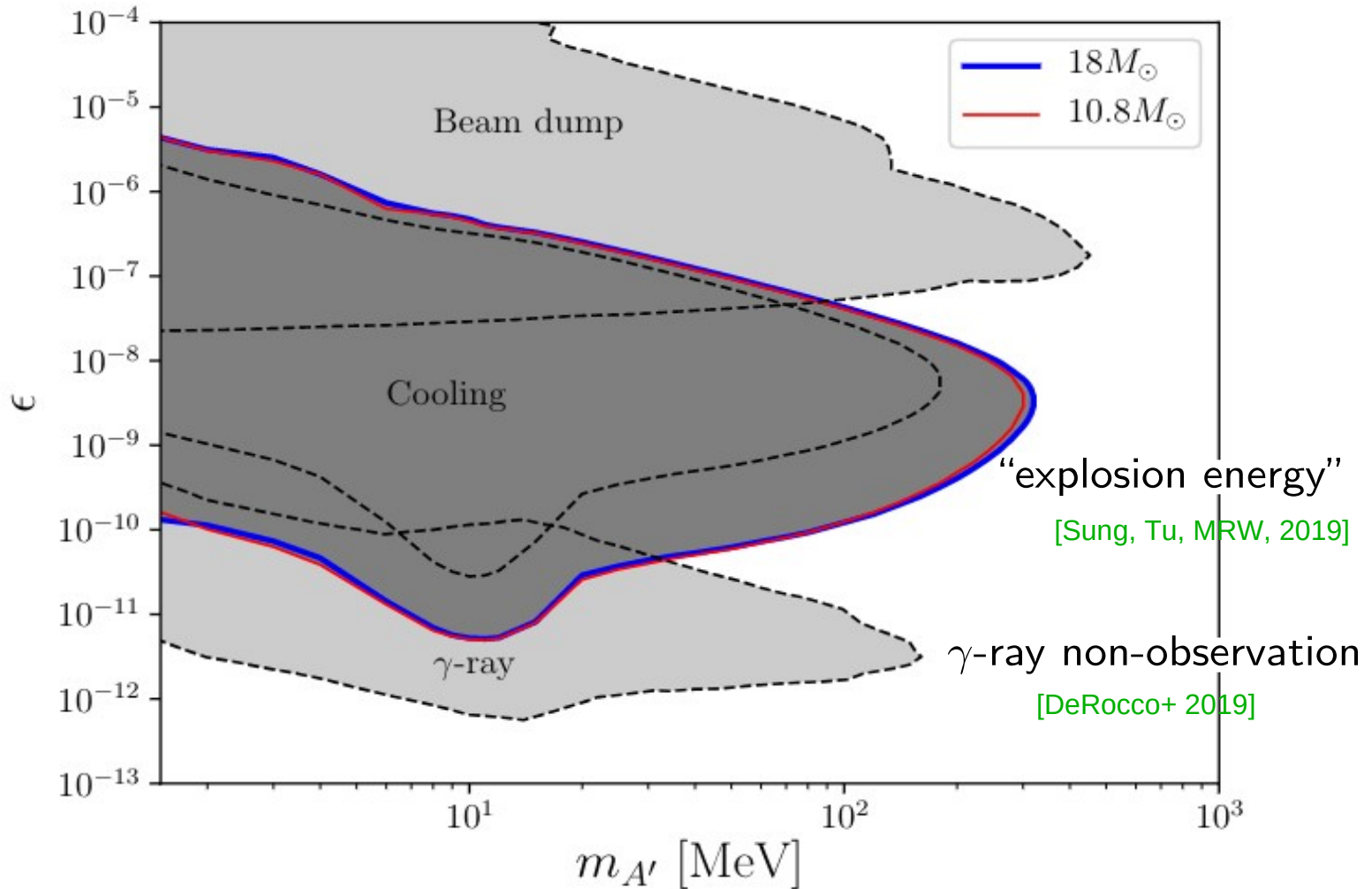
$\Delta E_g$ : gravitational binding energy above  $R$



The radiative decay of dark particle can deposit much more energy than the binding energy of stellar envelope + the explosion energy



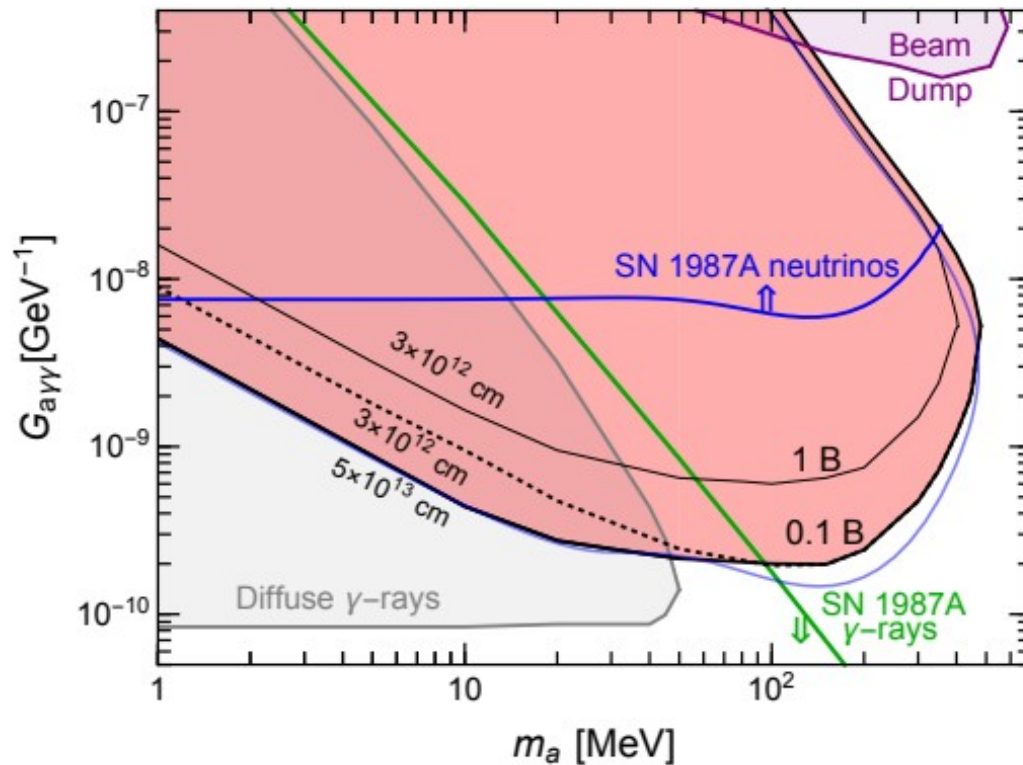
# Supernova bounds on dark photon



## Some discussions...

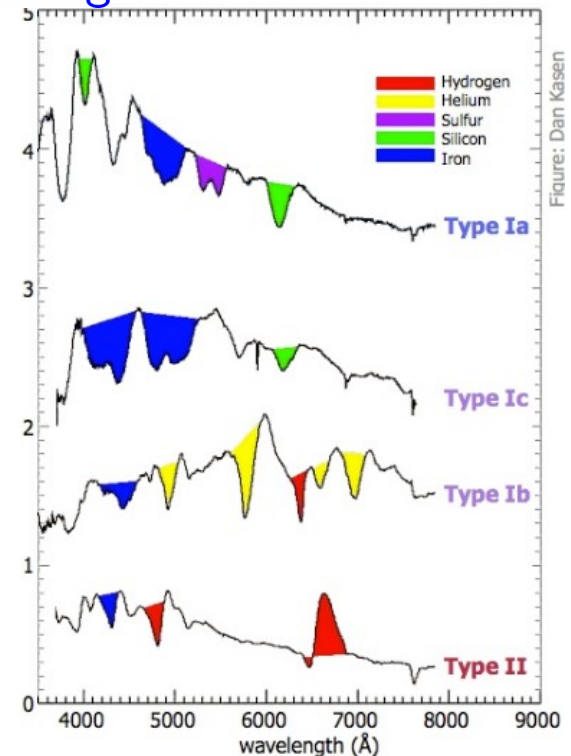
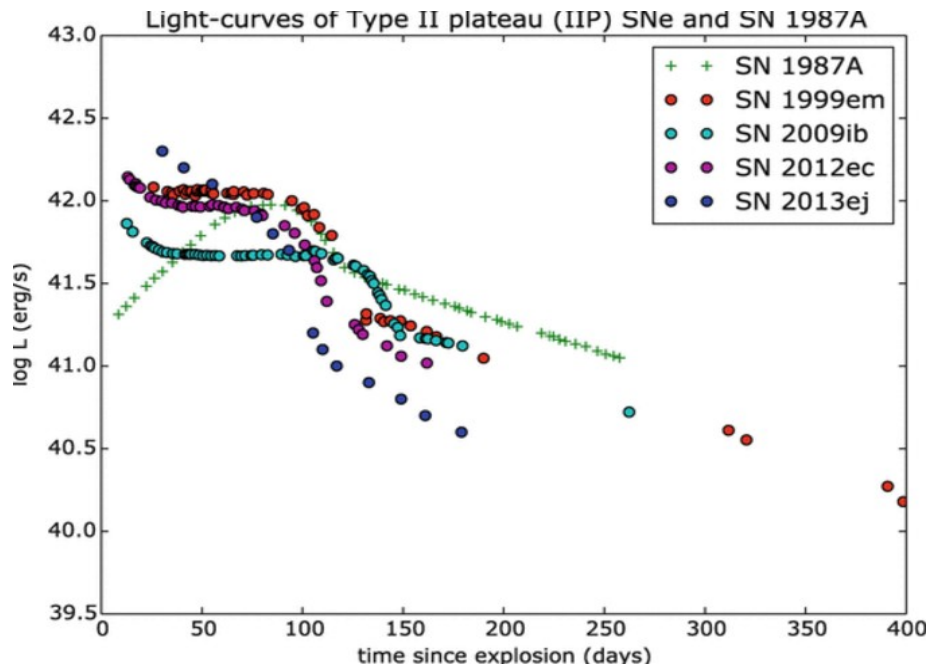
- low-energy supernovae  
(observed  $E_{\text{expl}} \simeq 0.1B$ )

[Caputo+ 2201.09890]



## Some discussions...

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(observed  $E_{\text{expl}} \simeq 0.1B$ )
- modification of lightcurve and spectrum?  
e.g., altered ejecta structure, collision between bSM-driven & shock driven components, change of absorption feature, ...  
→ requires further hydro+radiative transfer modeling



## Some discussions...

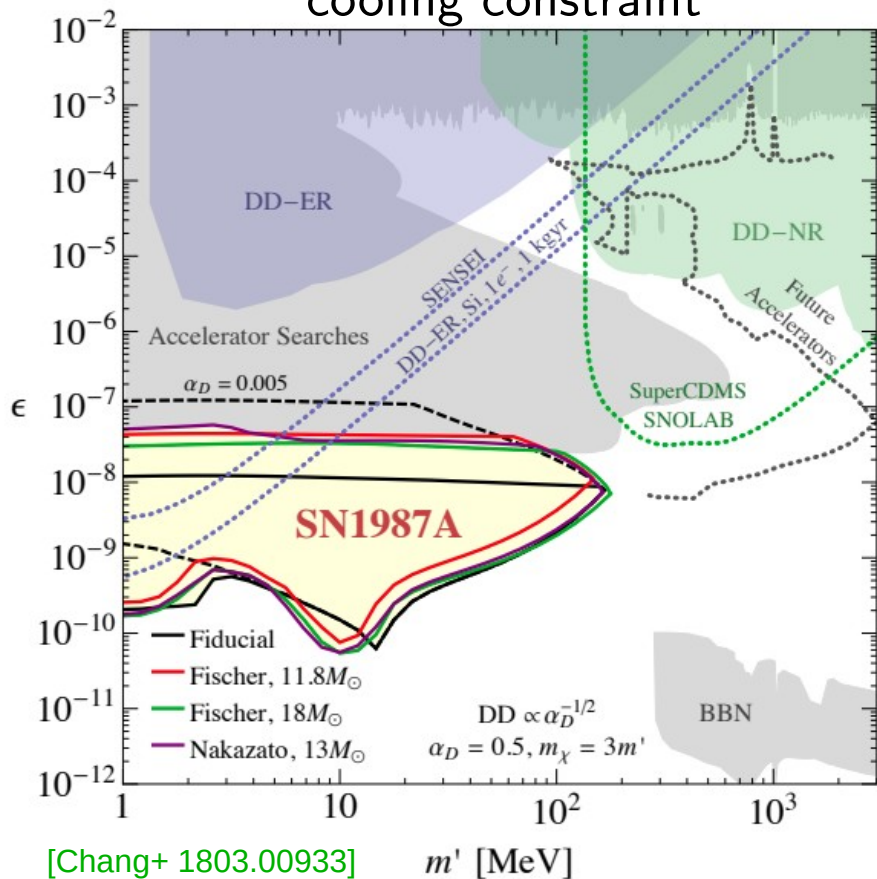
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- modification of lightcurve and spectrum?  
e.g., altered ejecta structure, collision between bSM-driven & shock driven components, change of absorption feature, ...  
→ requires further hydro+radiative transfer modeling
- disfavor SN explosions helped by bSM energy deposition  
Such scenario typically requires a sustained energy deposition rate of  $\gtrsim 10^{51}$  erg/s during the accretion phase

# SN constraints on light dark sector?

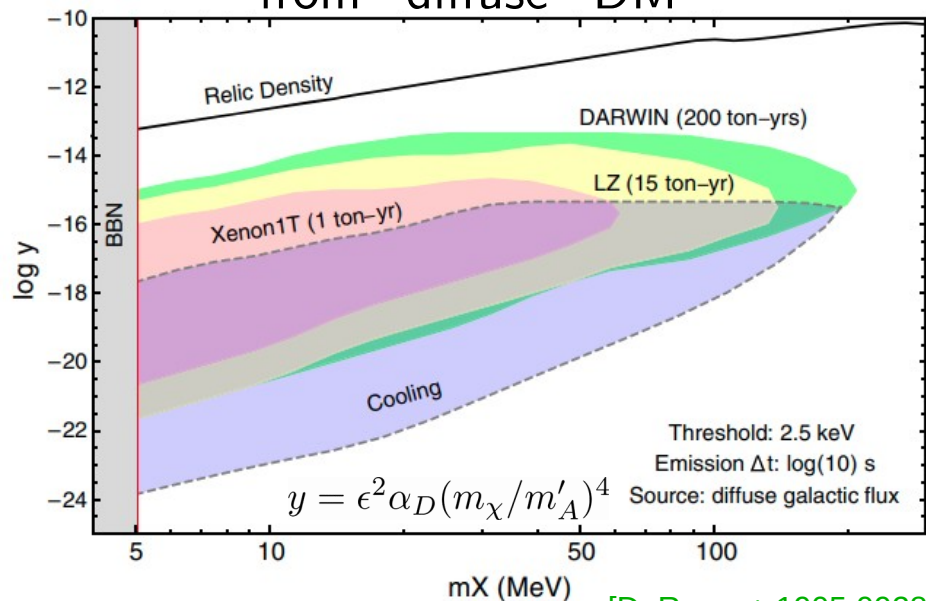
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Including both light dark fermion & dark photon

cooling constraint



from "diffuse" DM



Effect of dark sector self-interaction?

## Diffuse luminosity

For light & neutral particles, they can diffuse out from stars even if they are in thermal contact with matter (e.g., photons, neutrinos)

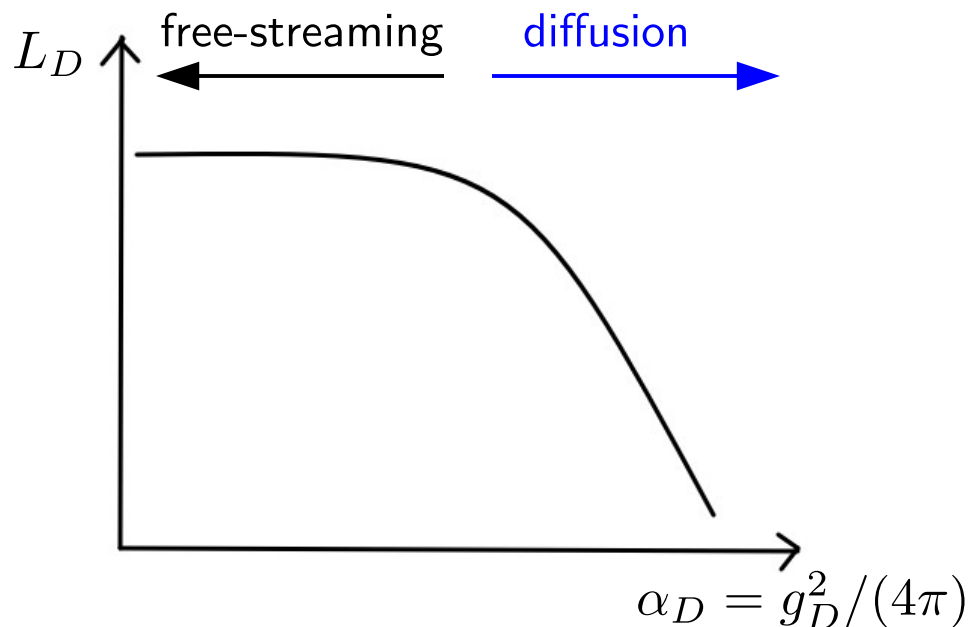
Diffusion luminosity across a surface (let's take the neutrinosphere):

$$L_i = -\frac{2g_i R_\nu^2 T_\nu^3}{3\pi} \left. \frac{dT}{dr} \right|_{R_\nu} \frac{1}{\langle \lambda_i^{-1}(R_\nu) \rangle} \times \int_{m_i/T_\nu}^{\infty} \xi^3 \sqrt{\xi^2 - \left(\frac{m_i}{T_\nu}\right)^2} \frac{e^\xi}{(e^\xi \pm 1)^2} d\xi$$

$\propto$  mean-free-path

larger dark sector coupling  
→ smaller mean-free-path  
→ smaller diffuse luminosity

diffusion due to self-trapping  
affects supernova (or in general,  
stellar) cooling bound!



## When to switch to diffusion limit?

Computing precisely the amount of dark sector particle emission from PNS is difficult and requires solving non-equilibrium transport equations

[Sung, Guo, MRW, 2102.04601]

We estimate this by:

i) estimate whether dark particles can be abundantly produced in PNS within characteristic time scale

$$\Delta t_i = \begin{cases} t_{\text{free}}, & \text{if } t_{i,\text{diff}} \leq t_{\text{free}}, \\ t_{i,\text{diff}}, & \text{if } t_{\text{free}} < t_{i,\text{diff}} < t_{\text{cool}}, \\ t_{\text{cool}}, & \text{if } t_{\text{cool}} \leq t_{i,\text{diff}}. \end{cases}$$

ii) if so and if particle mean-free-path is smaller than the radius of neutrinosphere

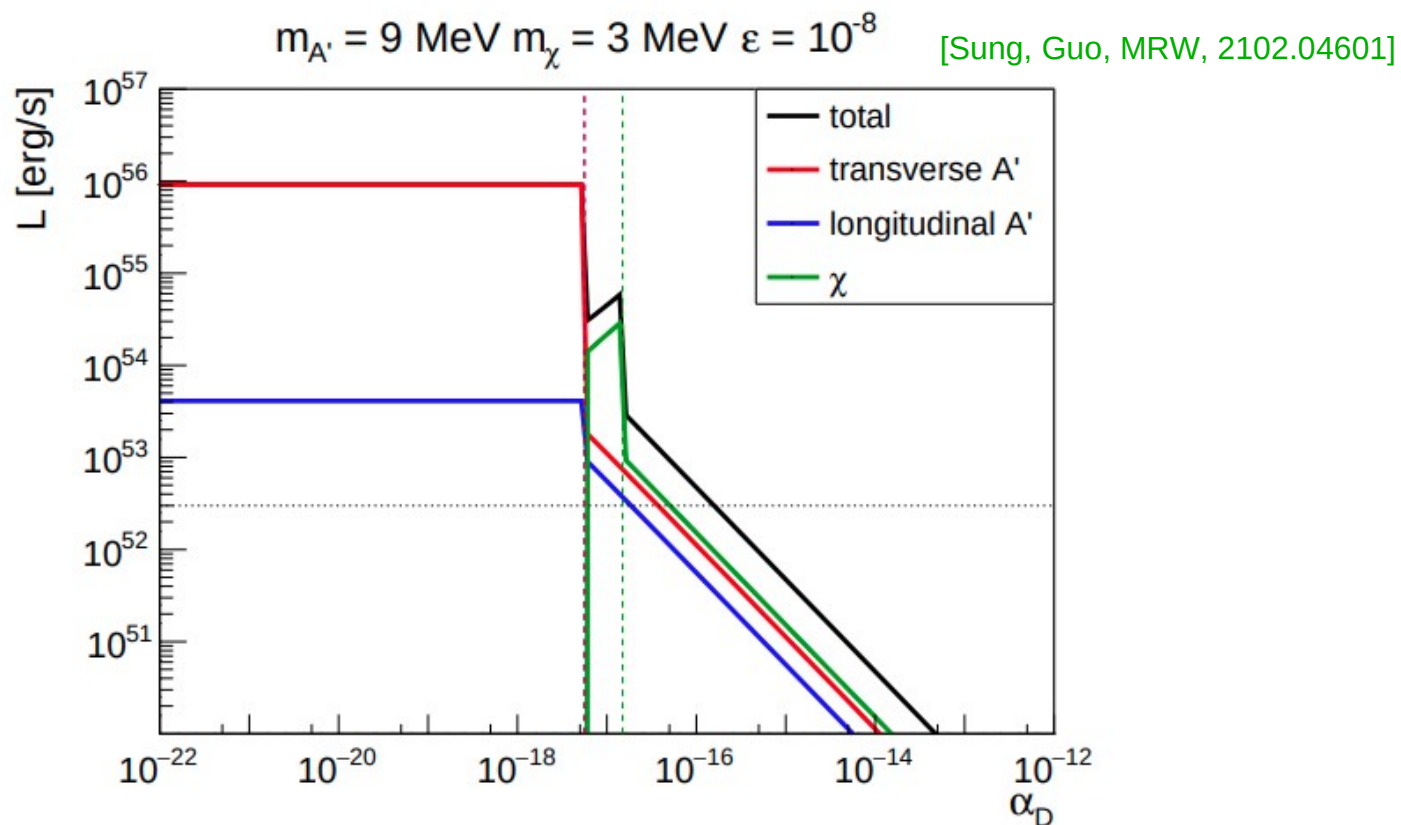
$$N_i > N_i^{\text{eq}}, \text{ and } \langle \lambda_i^{-1}(R_\nu) \rangle > R_\nu^{-1}$$

→ diffusion limit

Mass	Interaction	Type	Particle	Coupling
$m_{A'} < 2m_\chi$	$A'np \rightarrow np$	Abs.	SM	$\epsilon^2$
	$A'e^- \rightarrow e^-\gamma$	Abs.	SM	$\epsilon^2$
	$A' \rightarrow e^-e^+$	Abs.	SM	$\epsilon^2$
	$A'A' \rightarrow \chi\bar{\chi}$	Abs.	DS	$\alpha_D^2$
	$A'\chi \rightarrow \chi A'$	Sca.	DS	$\alpha_D^2$
$m_{A'} > 2m_\chi$	$A'np \rightarrow np$	Abs.	SM	$\epsilon^2$
	$A'e^- \rightarrow e^-\gamma$	Abs.	SM	$\epsilon^2$
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	$A' \rightarrow \chi\bar{\chi}$	Abs.	DS	$\alpha_D$

Mass	Interaction	Type	Particle	Coupling
$m_{A'} < 2m_\chi$	$\chi\bar{\chi}np \rightarrow np$	Abs.	SM	$\epsilon^2\alpha_D$
	$\chi\bar{\chi} \rightarrow e^-e^+$	Abs.	SM	$\epsilon^2\alpha_D$
	$\chi\bar{\chi} \rightarrow \gamma^*$	Abs.	SM	$\epsilon^2\alpha_D$
	$\chi p \rightarrow \chi p$	Sca.	SM	$\epsilon^2\alpha_D$
	$\chi e^- \rightarrow \chi e^-$	Sca.	SM	$\epsilon^2\alpha_D$
	$\chi\bar{\chi} \rightarrow A'A'$	Abs.	DS	$\alpha_D^2$
	$\chi\chi \rightarrow \chi\chi$	Sca.	DS	$\alpha_D^2$
	$\chi\bar{\chi} \rightarrow \chi\bar{\chi}$	Sca.	DS	$\alpha_D^2$
	$\chi A' \rightarrow A'\chi$	Sca.	DS	$\alpha_D^2$
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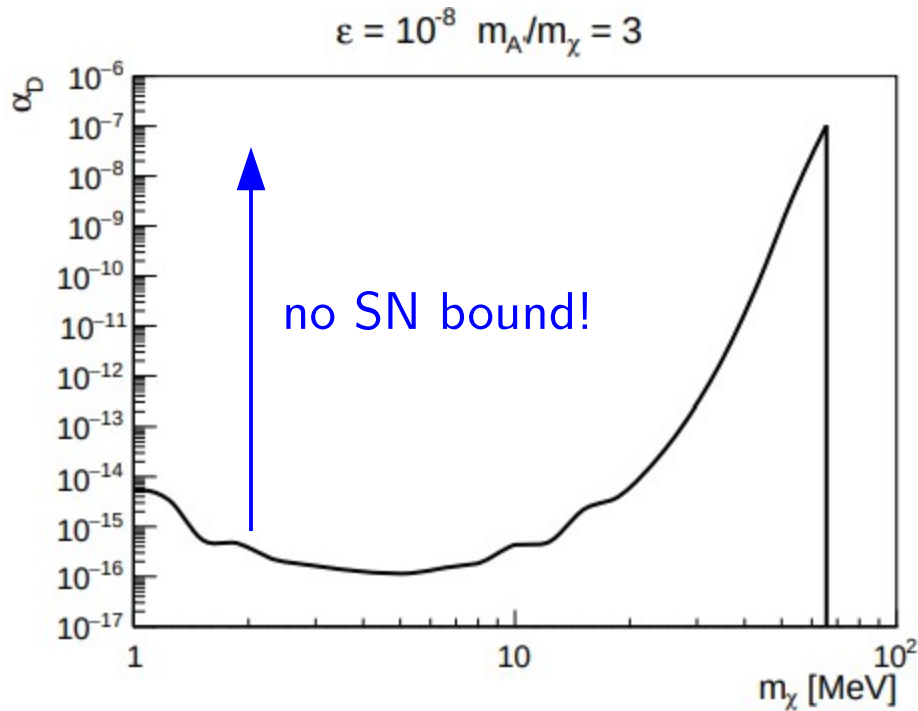
# Self-trapping and dark luminosity



- self-trapping effect can evade the SN bound for  $\alpha_D \gtrsim 10^{-16}$
- the corresponding  $\sigma_{\chi\bar{\chi}}$  (or equivalent) is  $\sim \mathcal{O}(10^{-40}) \text{ cm}^2$



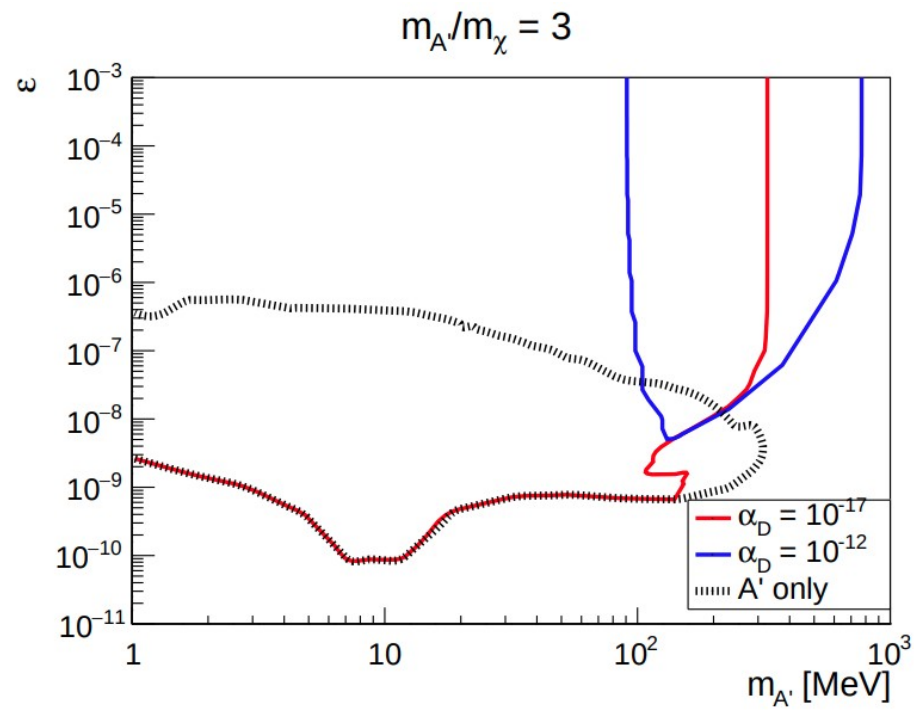
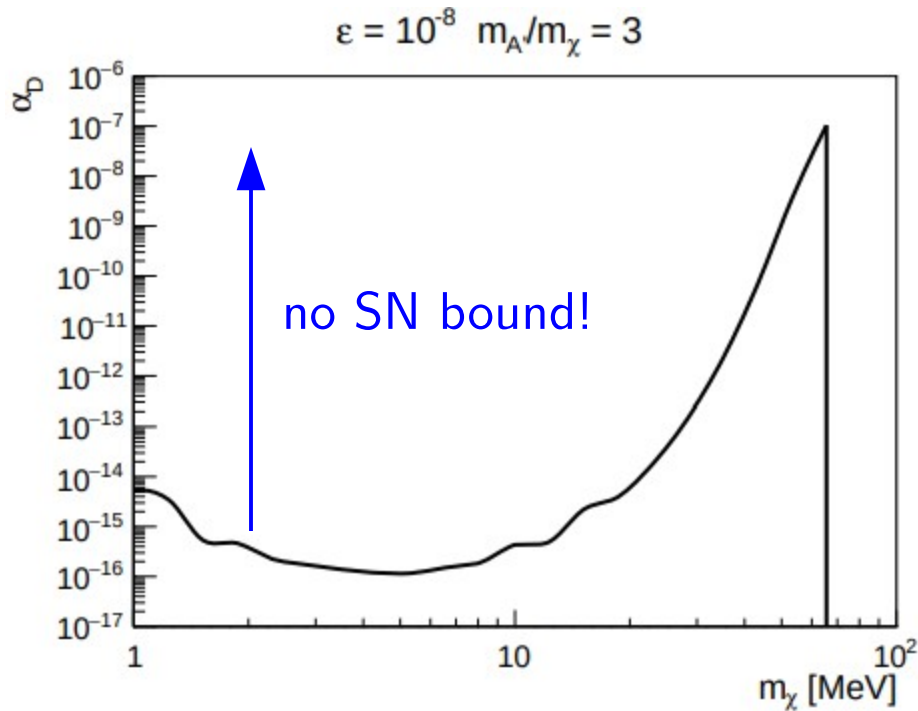
# Self-trapping effects on supernova bound



[Sung, Guo, MRW, 2102.04601]

- erase bounds for  $m_{A'} \lesssim 30$  MeV if  $\alpha_D \gtrsim 10^{-16}$

# Self-trapping effects on supernova bound



[Sung, Guo, MRW, 2102.04601]

- erase bounds for  $m_{A'} \lesssim 30$  MeV if  $\alpha_D \gtrsim 10^{-16}$
- may extend SN bounds to large  $\epsilon$  for small  $\alpha_D$ , due to the decay of  $A' \rightarrow \chi\bar{\chi}$

## Some discussions...

If these light dark sector particles got trapped...

- How do they modify the subsequent evolution of (proto)neutron star?
- Any potential asymmetry to seed a compact dark component inside NS to be probed by GW?
- Can the same mechanism work (in a meaningful way) to evade other stellar cooling bounds?
- A transient phase of "dark matter" burst (which may not be actually trapped)?