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July 10th, 2026

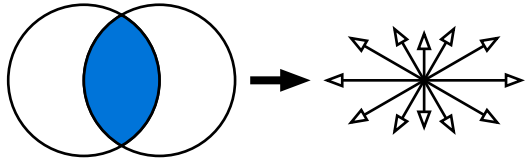
Spin and Magnetic fields in Heavy-Ion Collisions

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Hydrodynamics and Geometry

- Hydrodynamics translates initial-state geometry to momentum-space



- Ingredients for this mechanism:
 - EoM is a conservation law (e.g. $\nabla_\mu T^{\mu\nu} = 0$), hydrodynamic quantities are **transported**
 - Mean free path is smaller than average (pseudo)particle separation

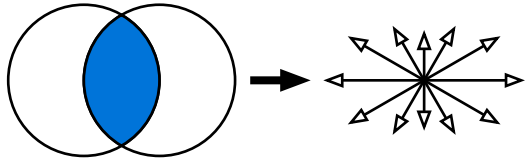
Takeaway: hydro is a geometry machine; what you put in, comes out.

- Fitting hydro models is a two-way street:
 - well-known hydro $\Rightarrow$ constrains initial state geometry
 - well-known initial state $\Rightarrow$ constrains hydro parameters
- Don't know either? Do a Global (Bayesian) fit!



Hydrodynamics and Geometry

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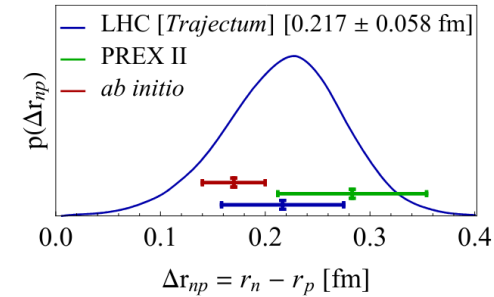
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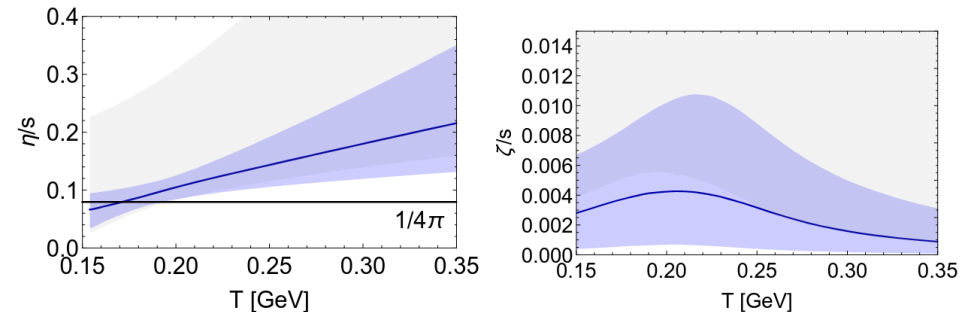
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Examples of fits using the hydro model *Trajectum*:

- Neutron skin of ^{208}Pb from ALICE data [1]



- Bulk and shear viscosities from global p_T -analysis [2]



Overview of the “Trajectum” nuclear collision framework

Trajectum[3] is...

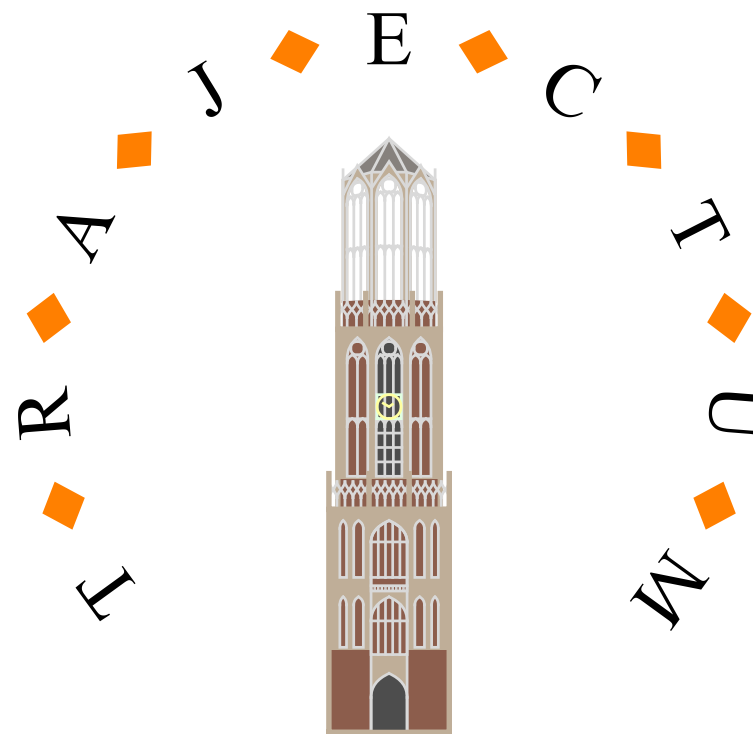
- ... a hydrodynamical nuclear collision framework...
- ... for modeling AA and pA collisions at LHC energies...
- ... purpose-built for doing Bayesian Analyses (BA).

Being purpose-built for BA means that

- We put a lot of effort into replicating event-by-event, particle based experimental observables.
- Computational efficiency is a priority, $\mathcal{O}(10^3)$ faster in comparison to models with similar physics content.

Physics content

- Projectile geometry input from nuclear community (NLEFT, PCGM, VMC, no-core shell model etc.)
- Various initial conditions (TRENTO, MCGlauber, MCDipper)
- Hybrid (holography + free streaming) prehydro stage
- 3+1D DNMR 2nd order viscous hydro evolution
- Cooper-Frye hadronization



Overview of the “Trajectory” nuclear collision framework

Trajectory Event Generator

Initial Condition $\lim_{\tau \downarrow 0} \tau \varepsilon(\vec{x}_\perp, \eta)$



Prehydrodynamic stage $\tau = \tau_{\text{hydro}}$



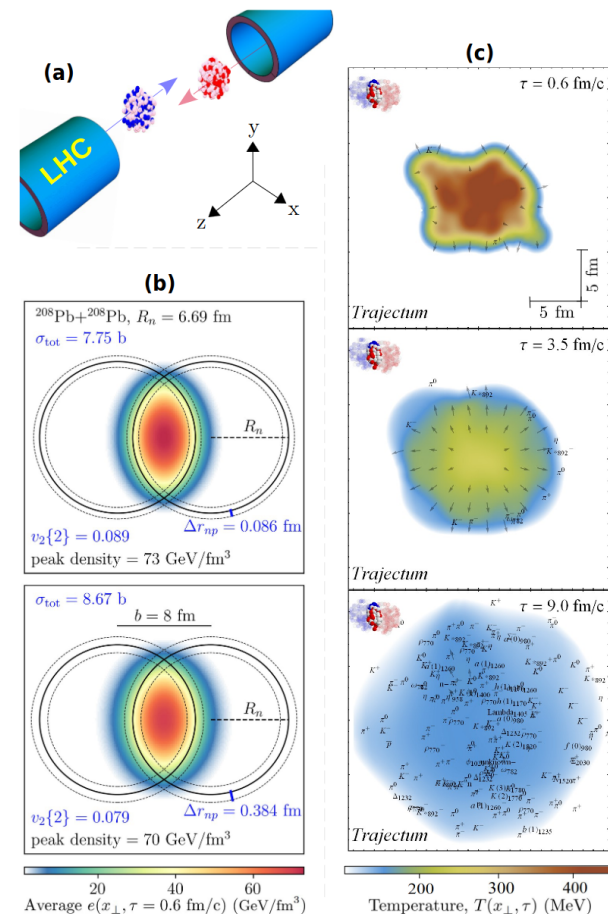
Hydro evolution $\tau_{\text{hydro}} < \tau < \tau_{\text{fs}}(x)$



Isothermal hadronisation $\tau = \tau_{\text{fs}}(x)$



Afterburner (external, SMASH or UrQMD) $\tau > \tau_{\text{fs}}(x)$



Spin and angular momentum in heavy-ion collisions

Mantra: hydro is a geometry machine

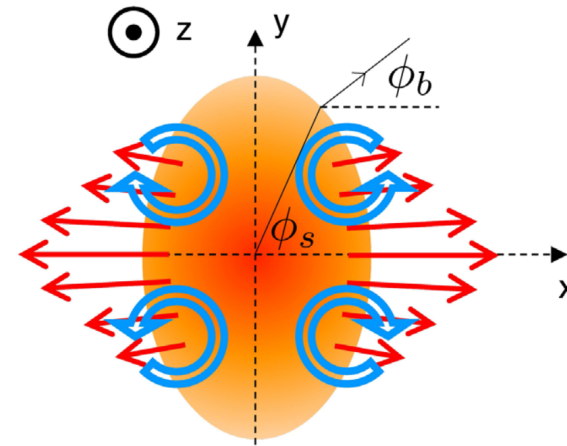
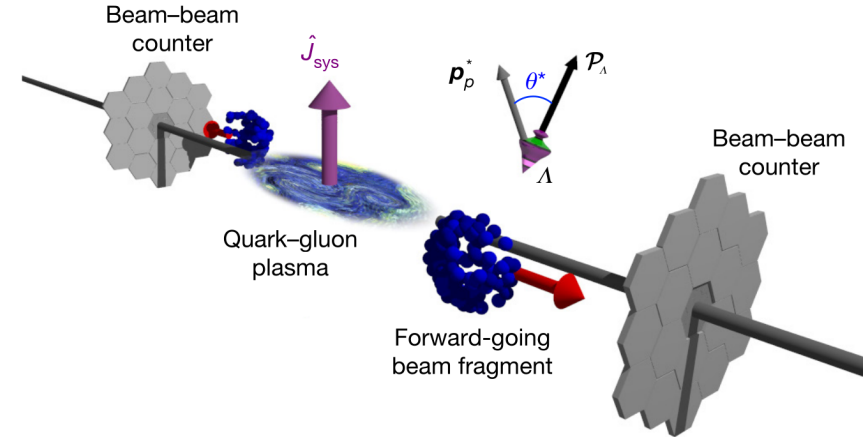
⇒ *geometric explanation for spin observables?*

- It is possible to measure the spin of weakly-decaying baryons by the angle between the decay products
- Spin of final-state hadrons is related to the vorticity and shear in the fluid through spin-orbit coupling
- We use a computation by [4], [5] of the expectation value of the Pauli-Lubański pseudo-vector $\hat{W}^\mu$ of a Dirac field in a thermal rotating state, normalised as $P^\mu = W^\mu / W$:

$$P^\mu(k) = \frac{1}{8m} \epsilon^{\mu\alpha\beta\gamma} (n_F(k) - 1) \left[2 \frac{k_\beta \Sigma_\gamma}{\Sigma \cdot k} \xi_{\alpha\delta} k^\delta + \varpi_{\alpha\beta} k_\gamma \right]$$

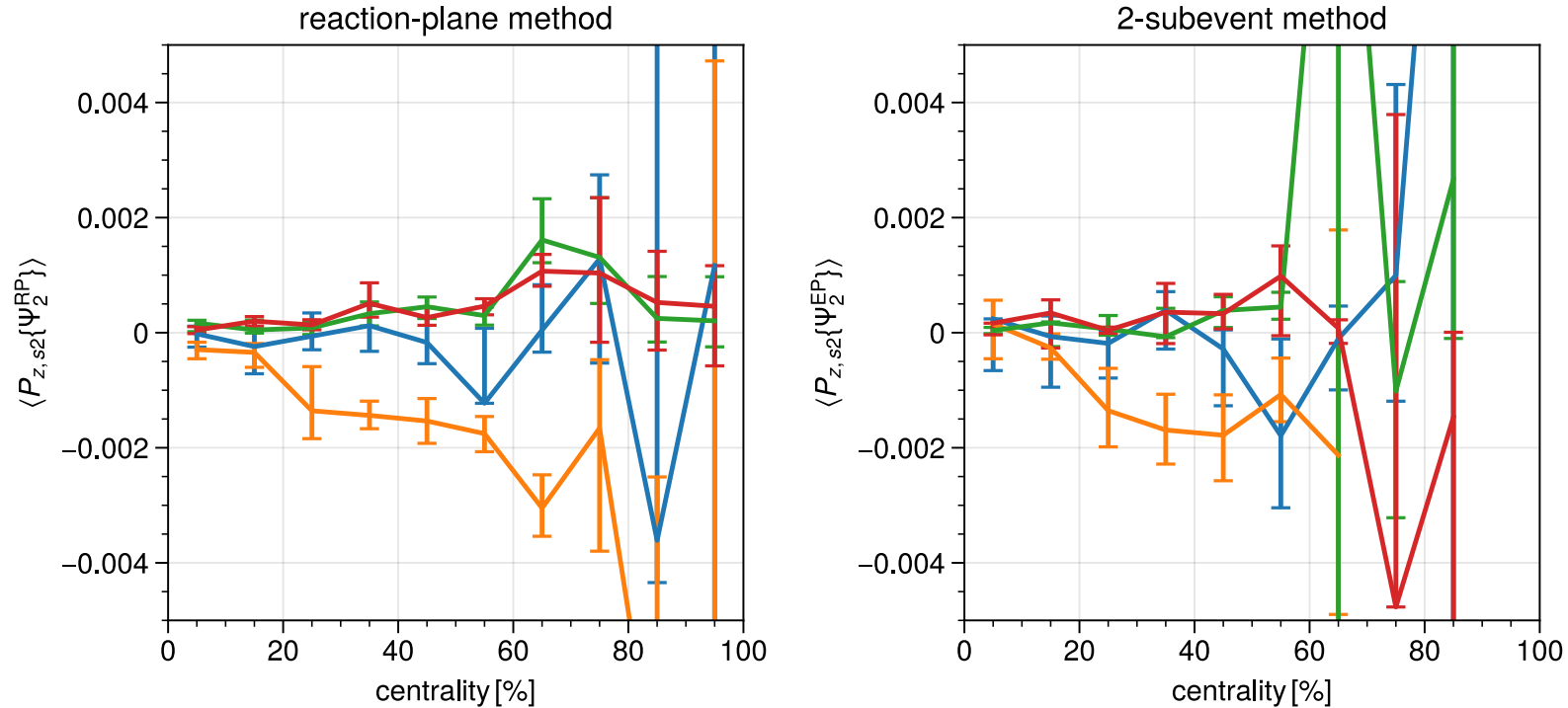
Two terms:

1. Thermal vorticity $\varpi_{\mu\nu} = -\partial_{[\mu}\beta_{\nu]}$ (yes really with a minus)
2. Thermal shear $\xi_{\mu\nu} = \partial_{(\mu}\beta_{\nu)}$



Polarization turns out to be very sensitive to viscosity

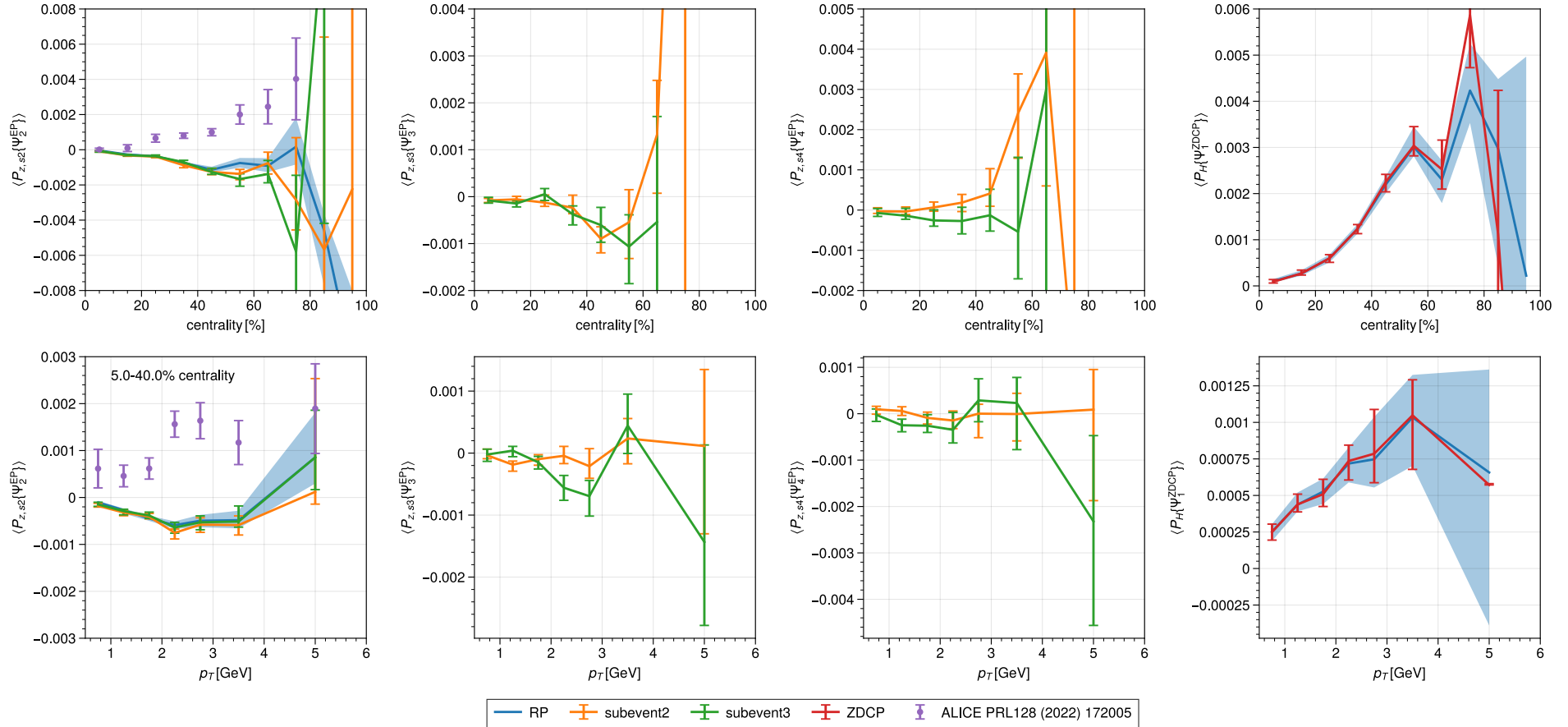
Trajectum (PRELIMINARY) - longitudinal polarisation of $\Lambda + \bar{\Lambda}$, second sine harmonic $P_{z,s2}$



- | | |
|---|---|
| 10k 2+1D Pb+Pb @ ~5020GeV ideal hydro | 10k 2+1D Pb+Pb @ ~5020GeV no bulk, high shear |
| 10k 2+1D Pb+Pb @ ~5020GeV high bulk, no shear | 10k 2+1D Pb+Pb @ ~5020GeV high bulk broad, no shear |

Tension between 3+1D fit* MAP value and ALICE

Trajectum (PRELIMINARY) $\Lambda + \bar{\Lambda}$ polarization 200k 3+1D Pb+Pb @ 5020GeV MAP, $0.5\text{GeV} \leq p_T \leq 10\text{GeV}$, $|y| \leq 0.8$



Magnetic fields in Heavy-Ion collisions

- Anisotropic electric and magnetic fields induce a difference in flow observables for charged particle/antiparticle pairs
 $\Rightarrow$ **charge-identified flow** Δv_n
- Model EM effect as perturbation to flow at freeze-out [6]:
 - EM fields are sourced by point charges with four-current:

$$j_{\text{free}}^\mu(y) = \sum_{\text{protons}, i} q \int d\tau V_i^\mu \tau \delta^4(y - b_i(\tau))$$

$$= q \delta^3(\vec{y}_\perp - \vec{b}_{\perp, i} - V_i^0 y^0 \mathbf{e}_3) [1, 0, 0, V_i]^T$$

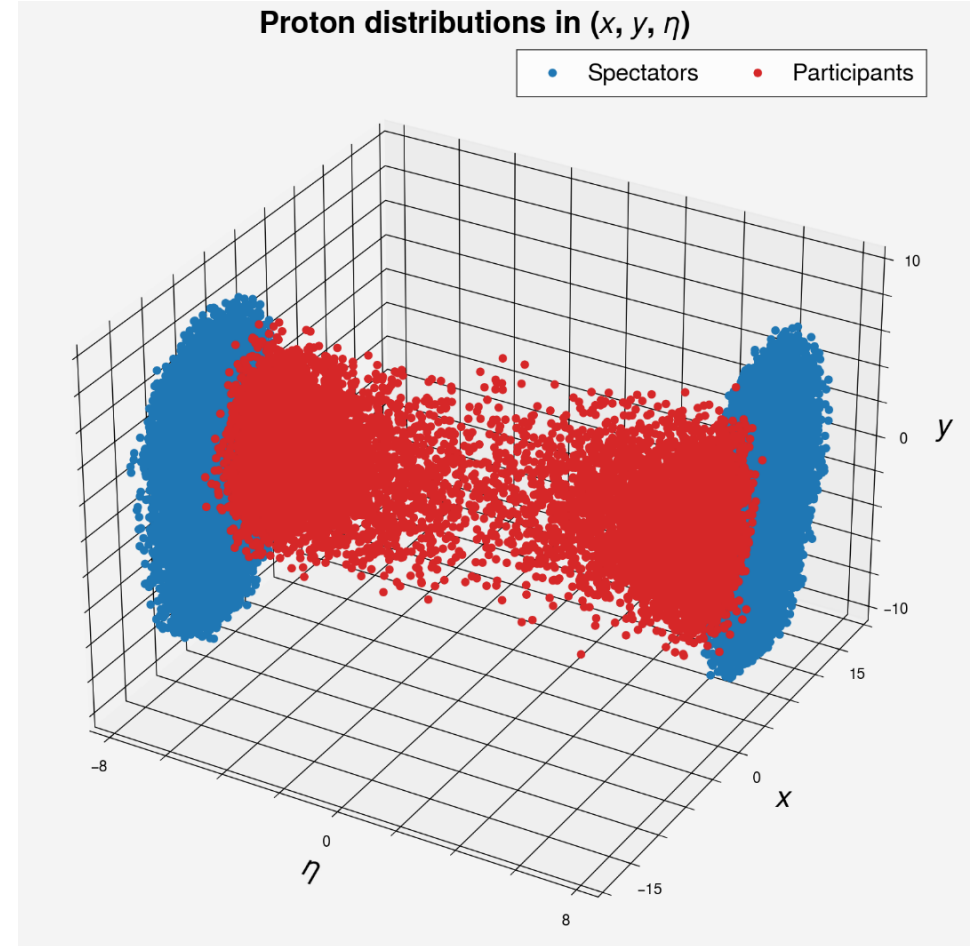
- Solve Maxwell equations in presence of conducting (conductivity σ) stationary ($\Upsilon^\mu = (1, \vec{0})$) medium:

$$[\square - \sigma \Upsilon^\mu \partial_\mu] \vec{B}(x) = -\nabla \times \vec{j}_{\text{free}}$$

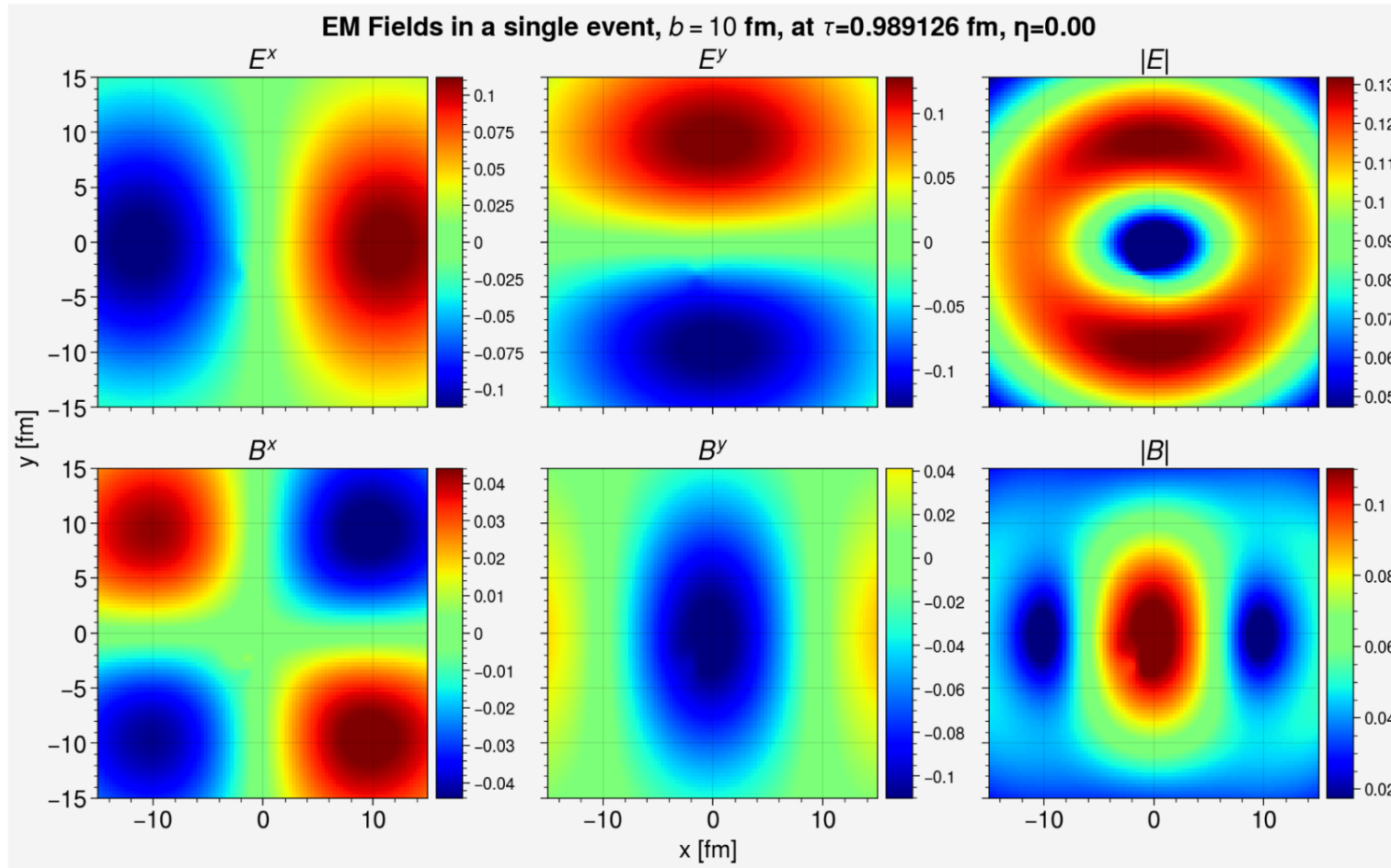
$$[\square - \sigma \Upsilon^\mu \partial_\mu] \vec{E}(x) = \nabla j_{\text{free}}^0 + \partial_0 \vec{j}_{\text{free}}$$

- Balance the force due to the EMF by adding a small drift velocity $u' = u + v$, that feels a drag force:

$$q_{\text{part.}} (\vec{E} - \vec{B} \times \vec{v}) = \mu m \vec{v}$$

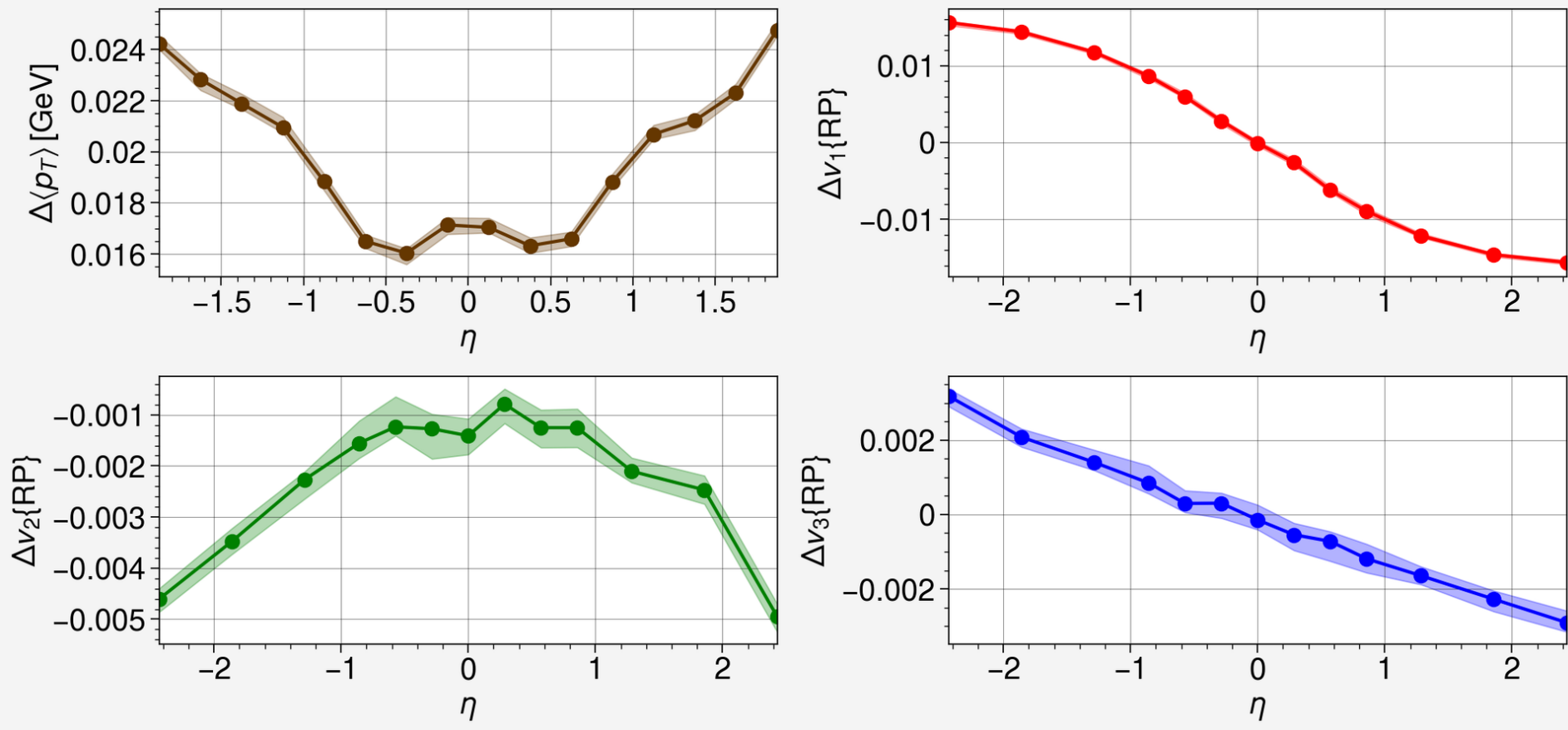


External EMF from initial state charge distribution



Charge identified observables @ $\sim 5020\text{GeV}$

Charge-dependent flow coefficients





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- [1] G. Giacalone, G. Nijs, and W. van der Schee, “Determination of the Neutron Skin of ^{208}Pb from Ultrarelativistic Nuclear Collisions”, *Phys. Rev. Lett.*, vol. 131, no. 20, p. 202302, Nov. 2023, doi: 10.1103/PhysRevLett.131.202302.
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- [5] F. Becattini, “Polarization in Relativistic Fluids: A Quantum Field Theoretical Derivation,” in *Strongly Interacting Matter under Rotation*, Springer International Publishing, 2021, pp. 15–52. doi: 10.1007/978-3-030-71427-7_2.
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