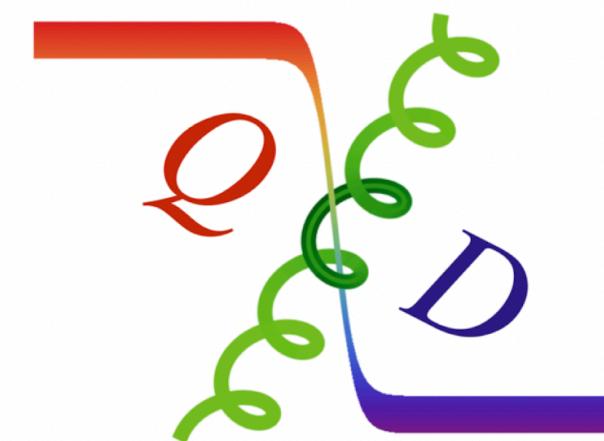


Analysis of nucleon excited states from lattice QCD with the Bayesian Reconstruction



**QUARK-GLUON
TOMOGRAPHY
COLLABORATION**

Bigeng Wang
Department of Physics and Astronomy, University of Kentucky

Collaborators: Raza Sabbir Sufian, Jian Liang, Terrence Draper, Keh-Fei Liu
(χ QCD collaboration)

Inverse Problems and UQ in Nuclear Physics
INT Workshop 24-88W
July 10th, 2024

Outline

- Introduction
 - Finite-volume spectrum from lattice QCD
- How to extract finite-volume spectrum from lattice QCD
 - Multi-exponential fit
 - Variational methods (GEVP)
 - An example: Roper state from lattice QCD
- Application of the Bayesian Reconstruction
- Conclusion and outlook

The path integral in the Euclidean space - lattice QCD

- Observables from the path integral (Minkowski):

$$\langle O \rangle = \frac{\int \mathcal{D}\phi O e^{iS_M(\phi)}}{\int \mathcal{D}\phi e^{iS_M(\phi)}}$$
 ← $Z_M(J) = \int \mathcal{D}\phi e^{i \int d^4x \mathcal{L}_M}$

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Perturbative expansions

Wick rotation

$$t \rightarrow -i\tau \\ k^0 \rightarrow ik_0^E$$

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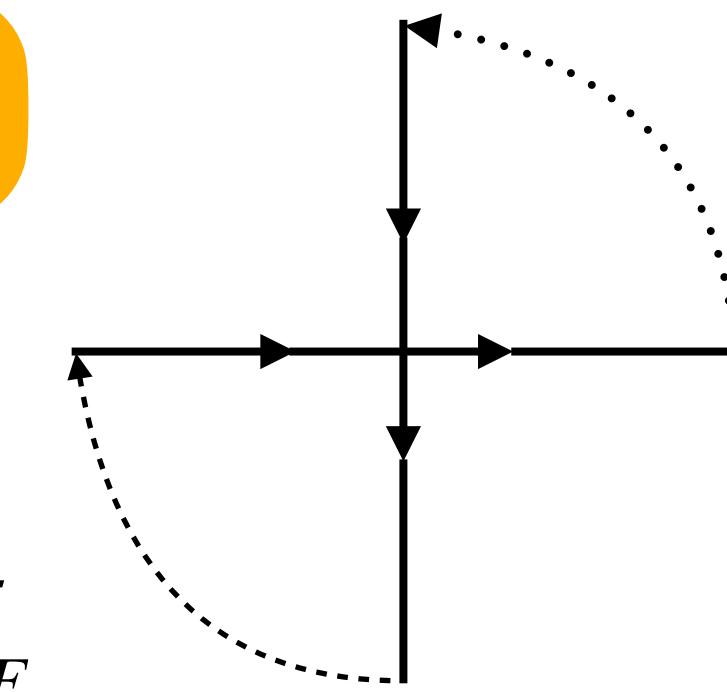
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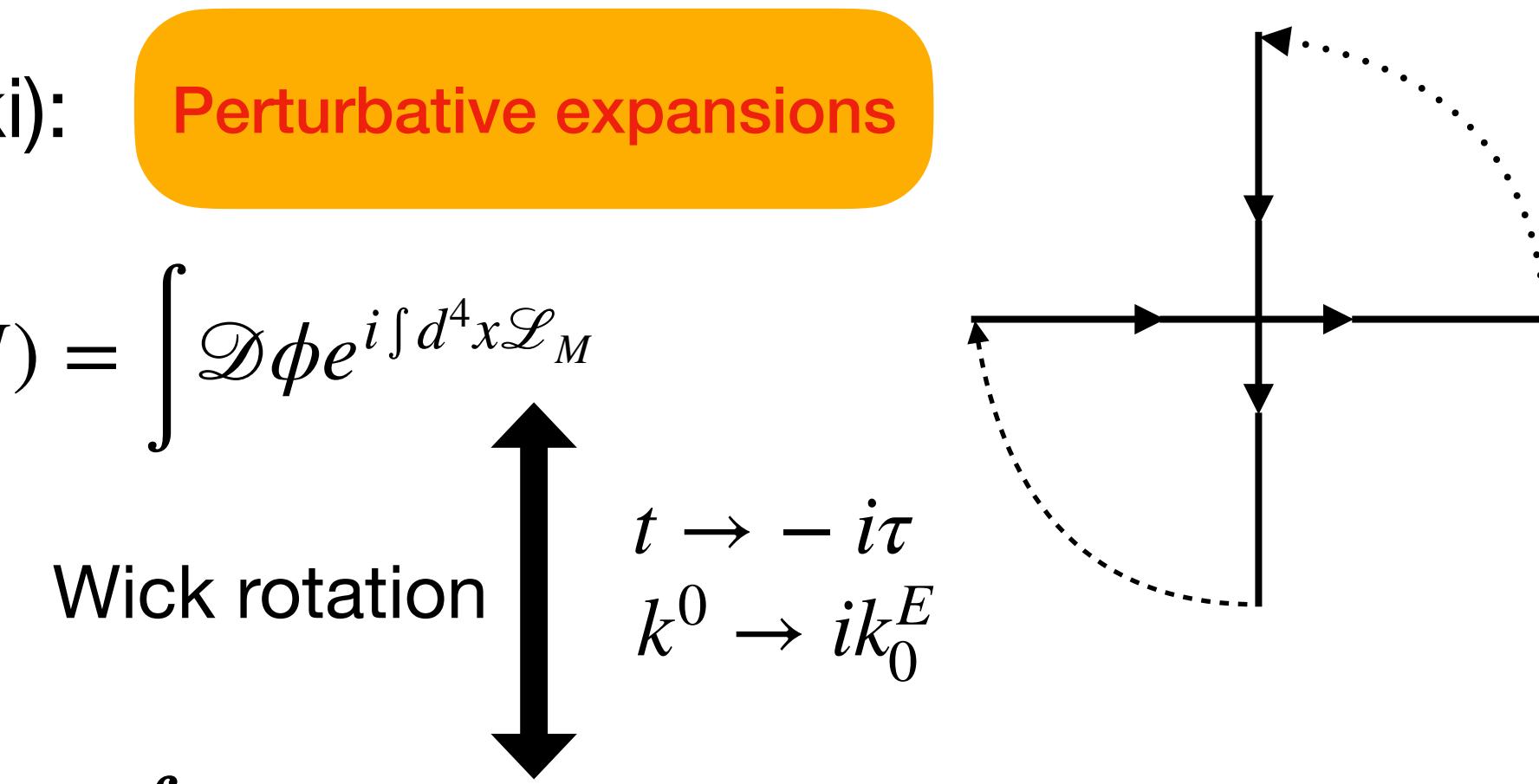
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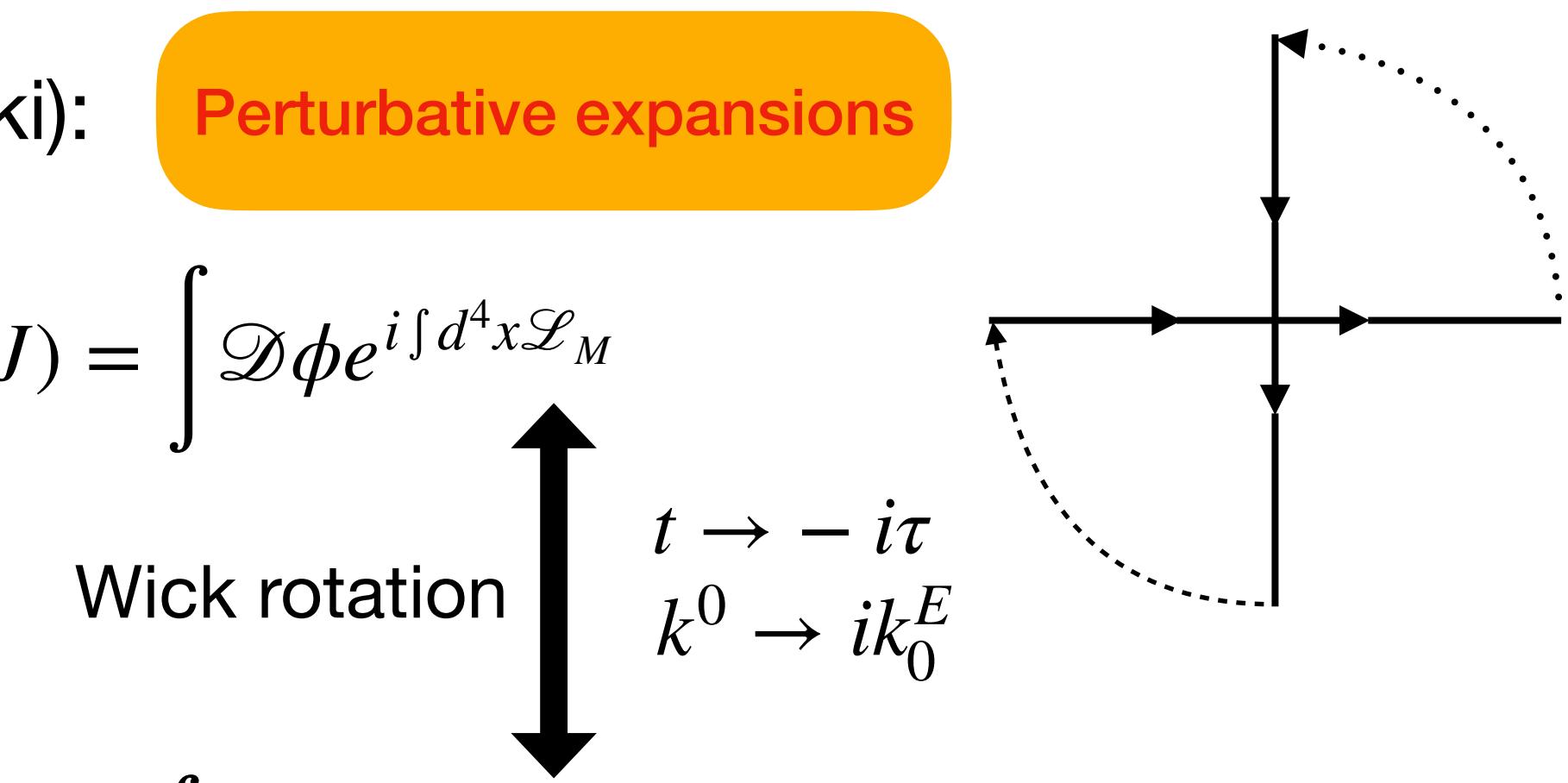


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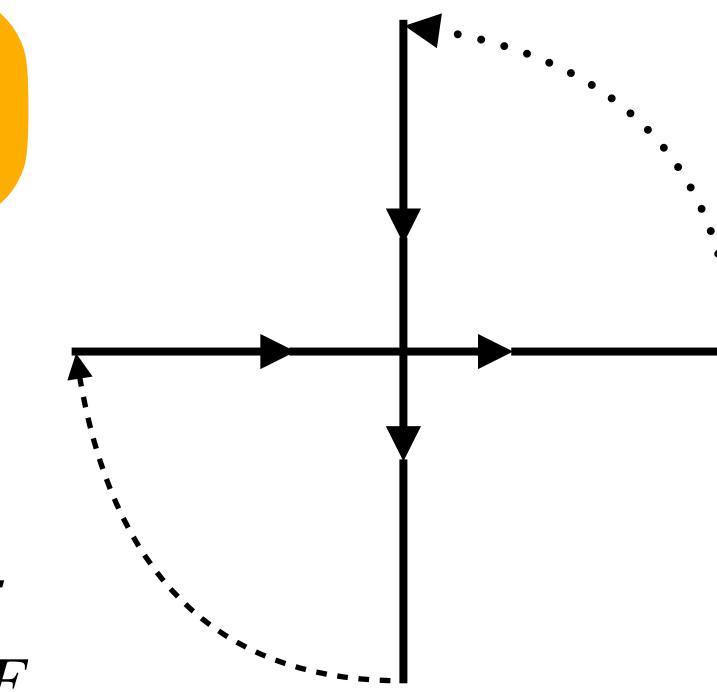
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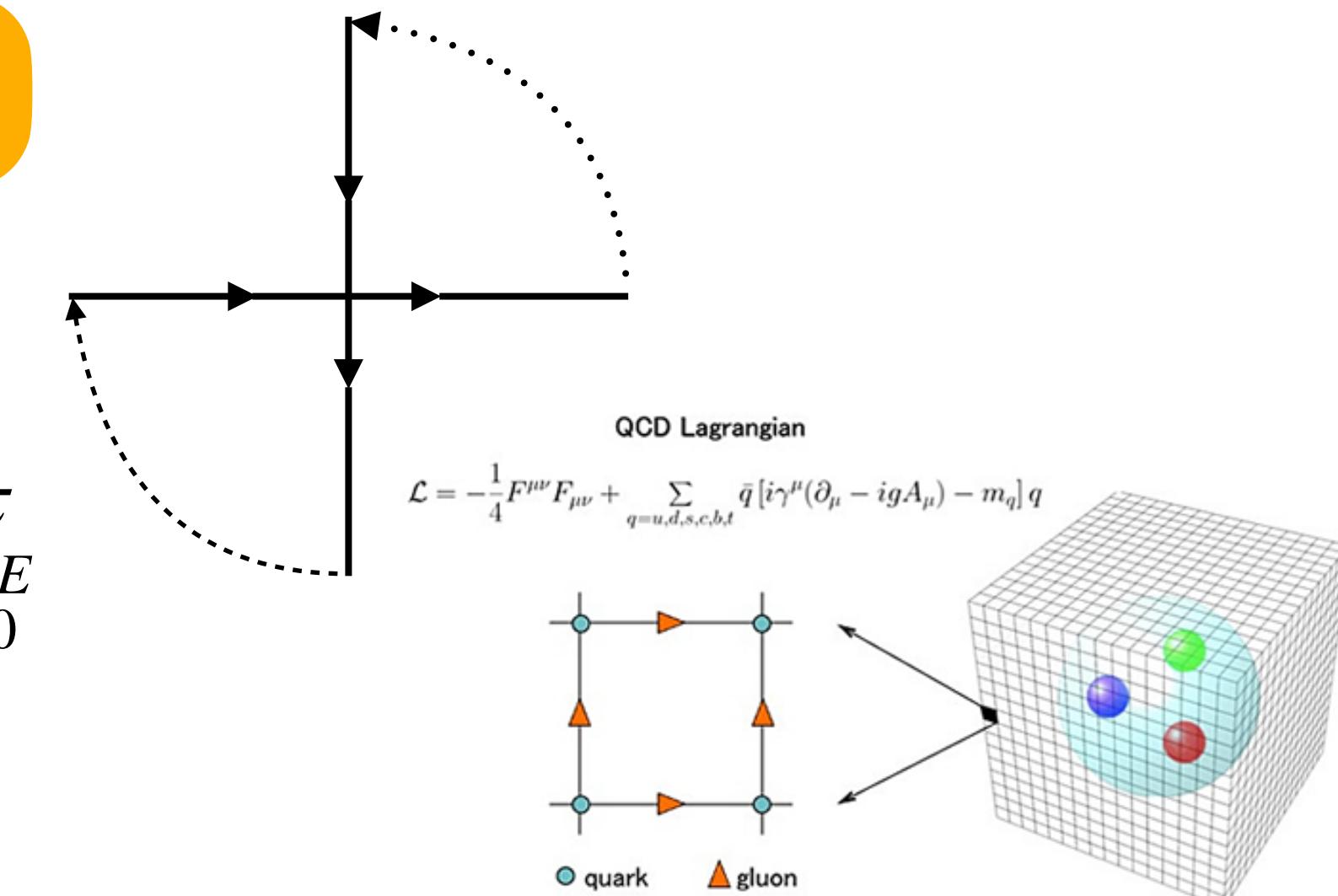
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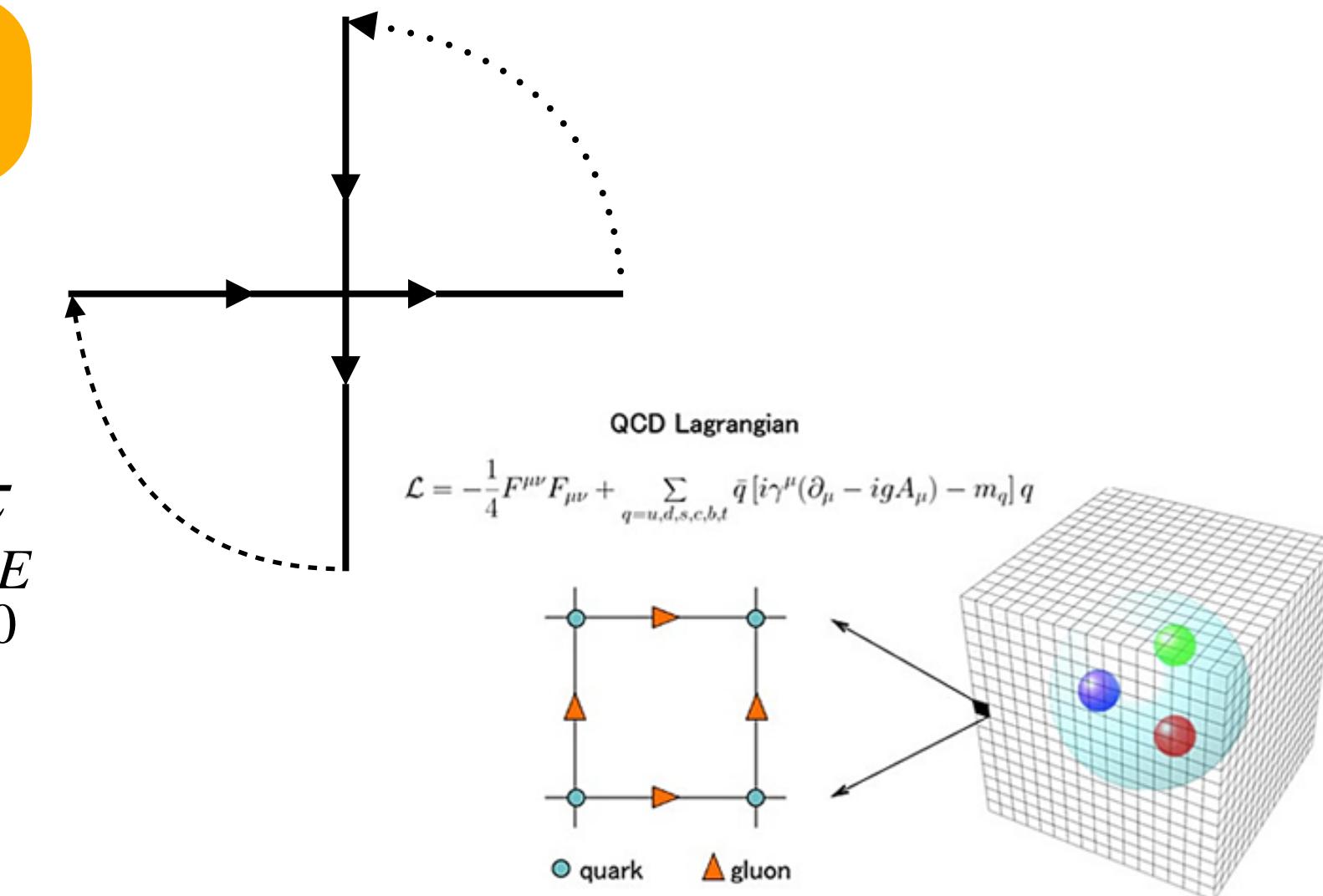
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Non-perturbative

Input parameters
 $a, m_q, \dots, \alpha_s, \dots$



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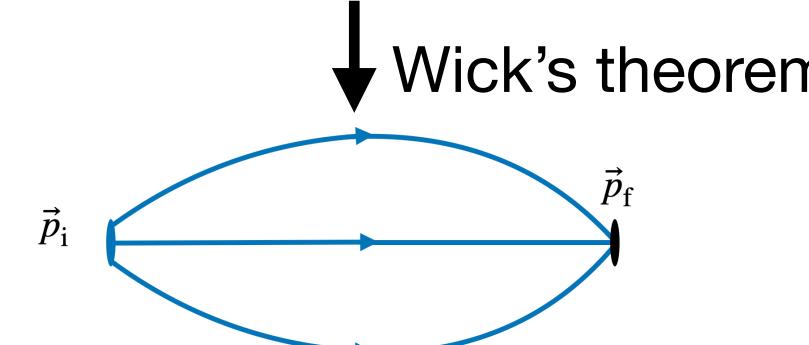
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E.g. n -point correlation functions

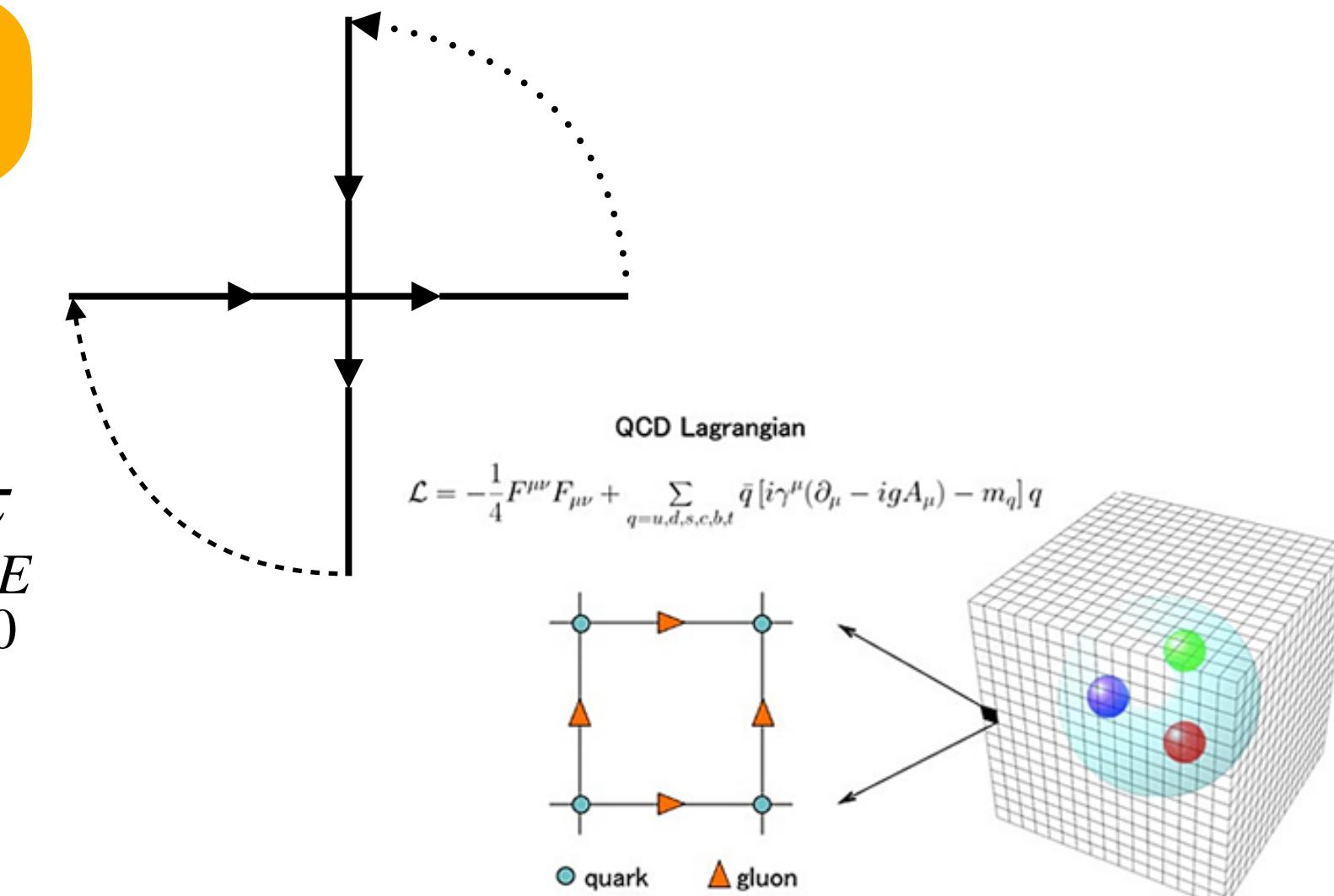
$$\langle \Omega | T\phi(x_1) \dots \phi(x_n) | \Omega \rangle$$



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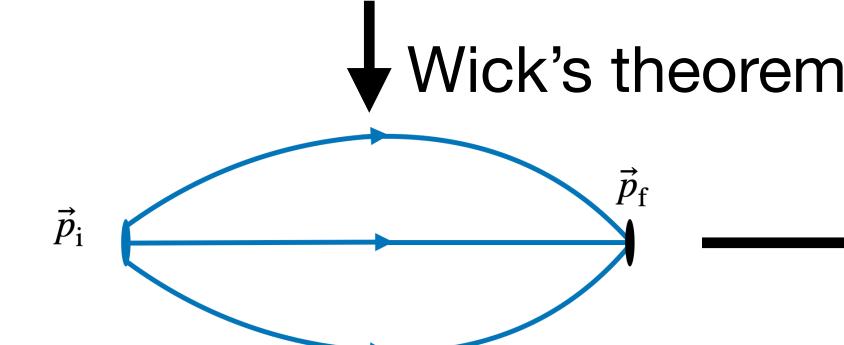
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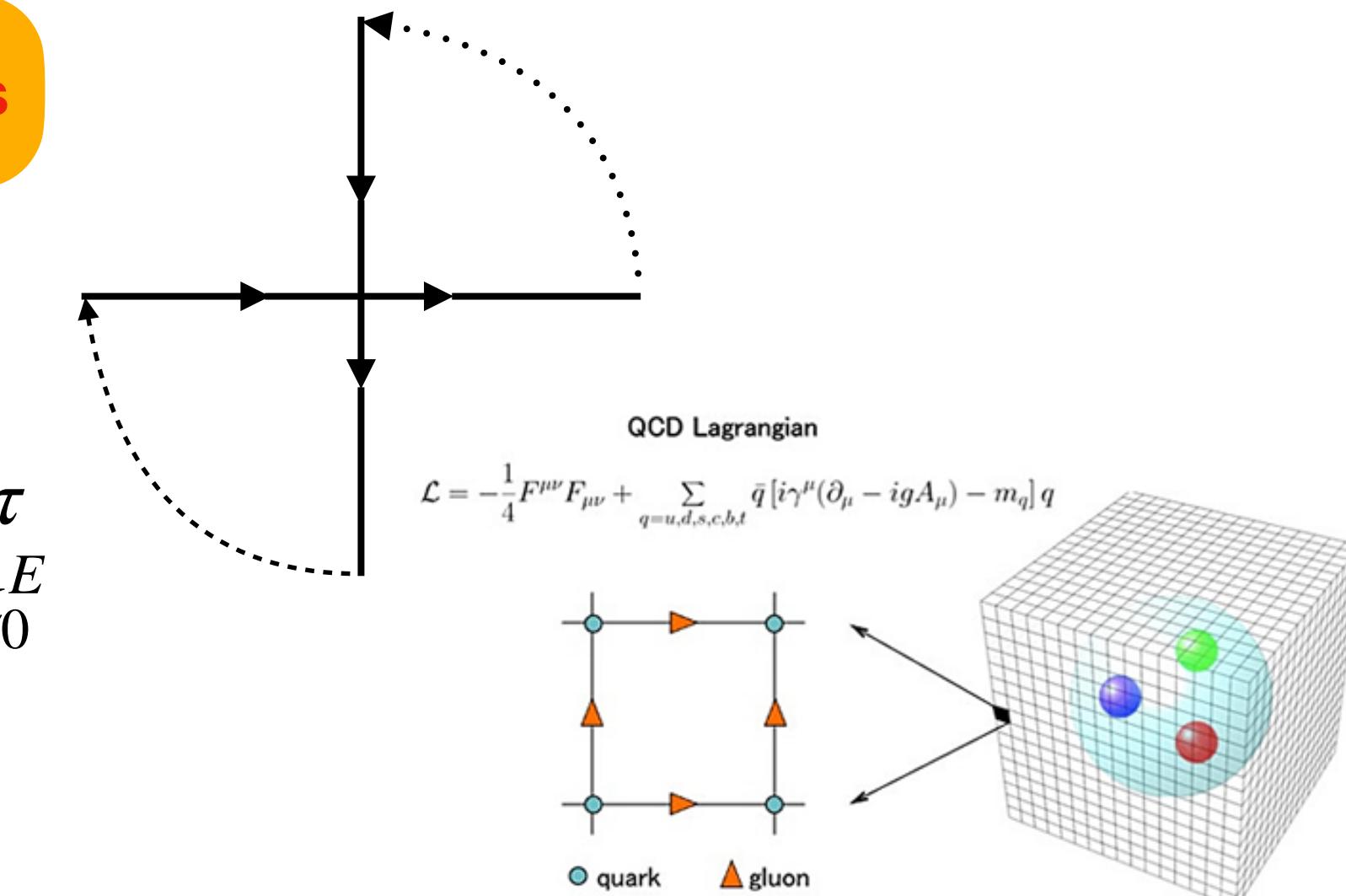


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Finite-volume spectrum, matrix elements ...



Extracting finite-volume spectrum

- What kind of observables do we measure on the lattice?
- To extract the finite-volume spectrum
 - Two-point correlation functions

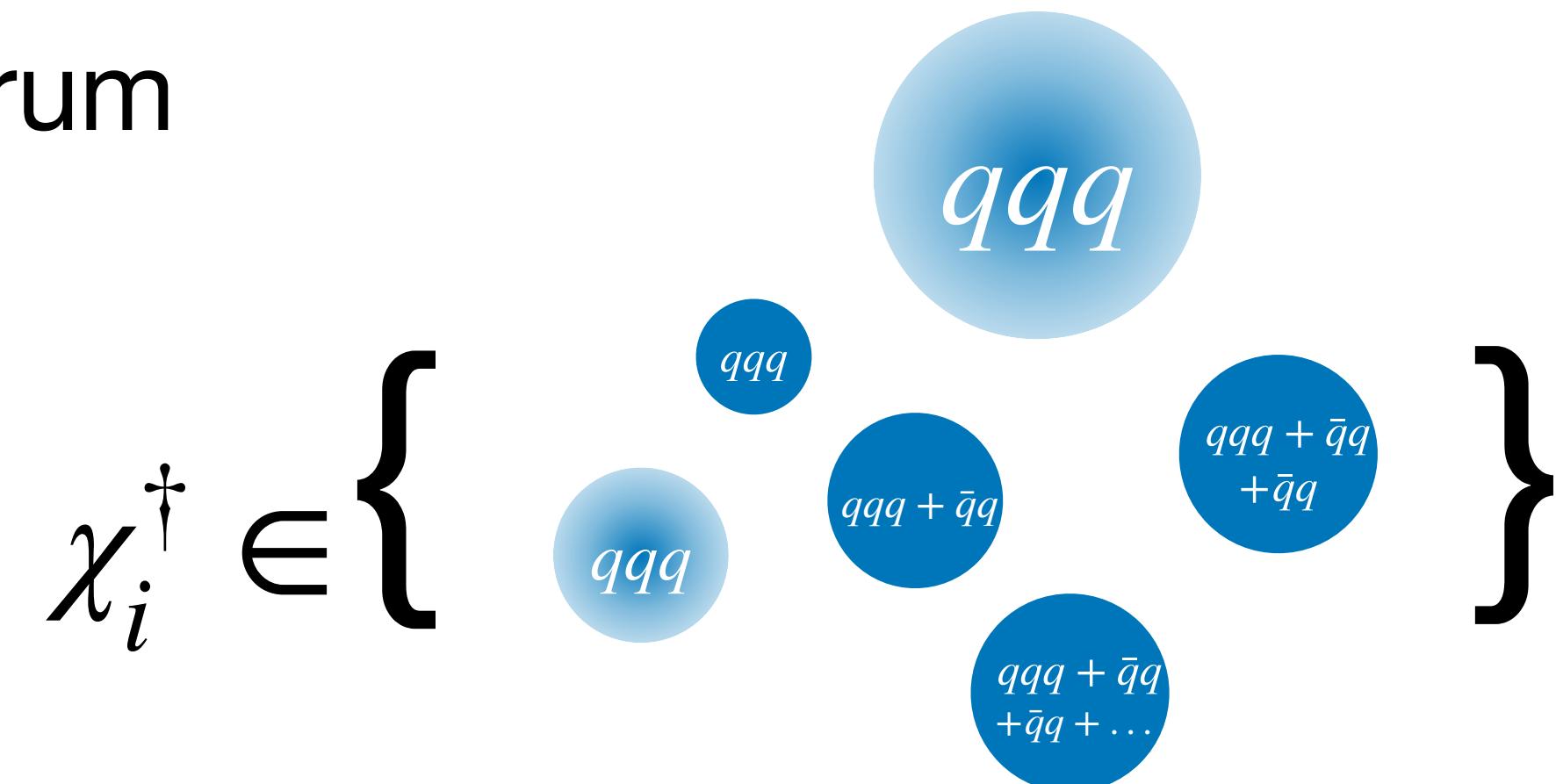
$$C_{ij}(t) = \langle \chi_i(t) \chi_j^\dagger(0) \rangle$$

The diagram shows a blue oval loop representing a two-point correlation function. The loop has arrows indicating direction. It starts at a point labeled \vec{p} on the left, goes up to a peak, then down to another peak, and finally back to \vec{p} on the right. Two black arrows point from the text $C_{ij}(t) = \langle \chi_i(t) \chi_j^\dagger(0) \rangle$ to the top and bottom peaks of the loop. Below the loop, there are two points labeled t_0 and t_f , with a horizontal arrow pointing from t_0 to t_f .

Extracting finite-volume spectrum

- What kind of observables do we measure on the lattice?
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interpolating operators with the same quantum numbers as the hadron state



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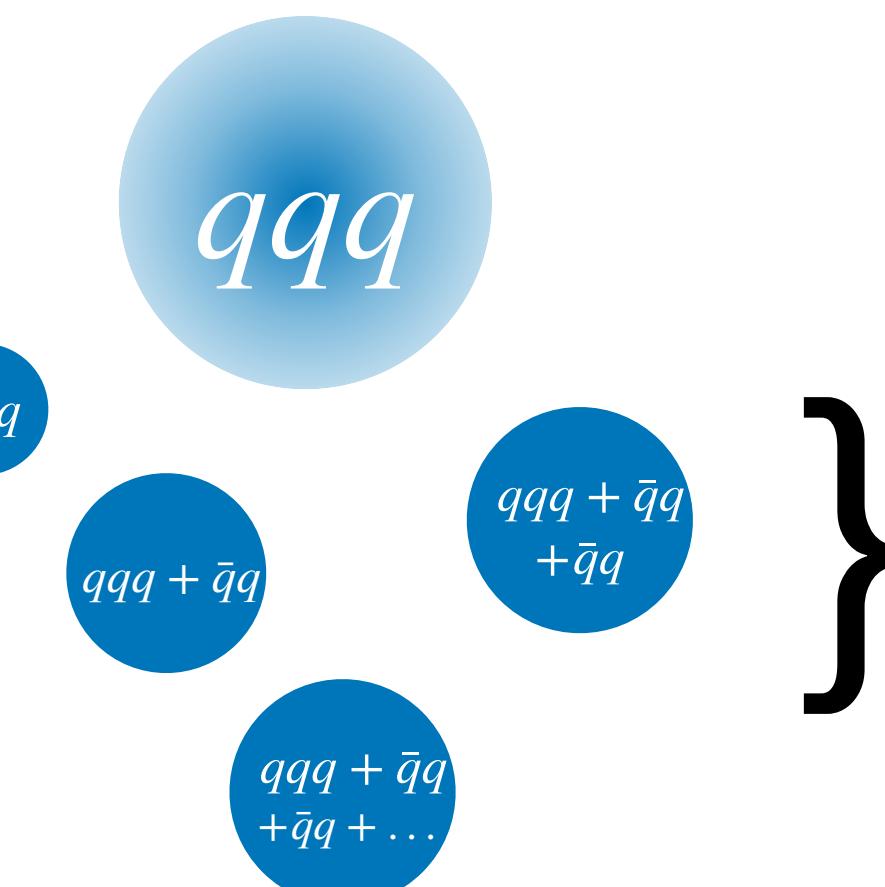
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$C_{ij}(t) = \langle \chi_i(t) \chi_j^\dagger(0) \rangle$
 $1 = \sum_n |n\rangle\langle n|$
 $\rightarrow \sum_{n=0}^{\infty} \langle \Omega | \chi_N | n \rangle \langle n | \chi_N^\dagger | \Omega \rangle e^{-E_n t}$

$qqq + \bar{q}q$
 $+ \bar{q}q + \dots$

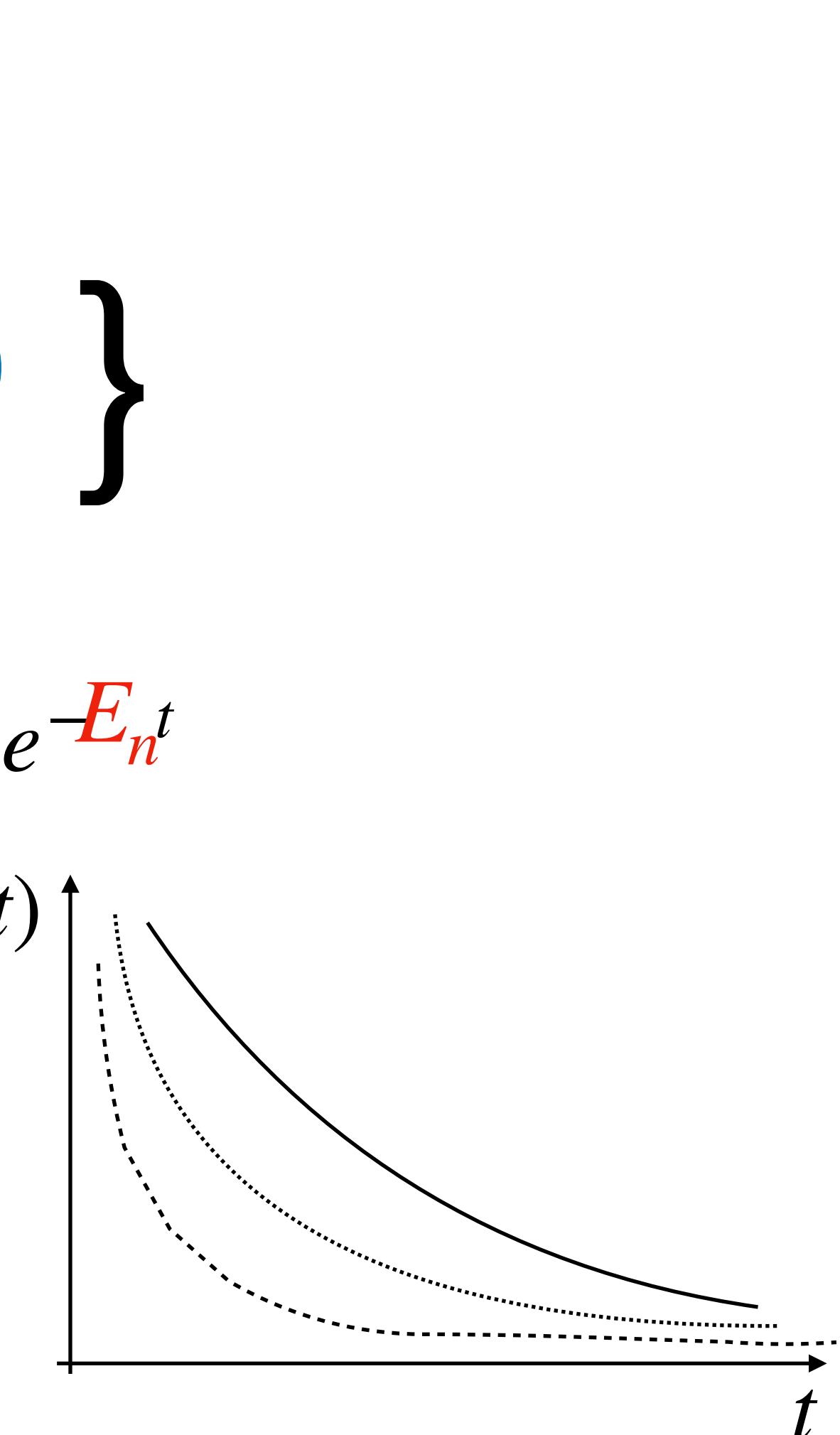


Extracting finite-volume spectrum

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The diagram illustrates the time evolution of a system from time t_0 to t_f . A blue oval represents the trajectory path. At t_0 , a state vector \vec{p} is shown at the left vertex. Two arrows point from the expression $C_{ij}(t) = \langle \chi_i(t) \chi_j^\dagger(0) \rangle$ towards the right vertex at t_f , indicating the final state of the system.

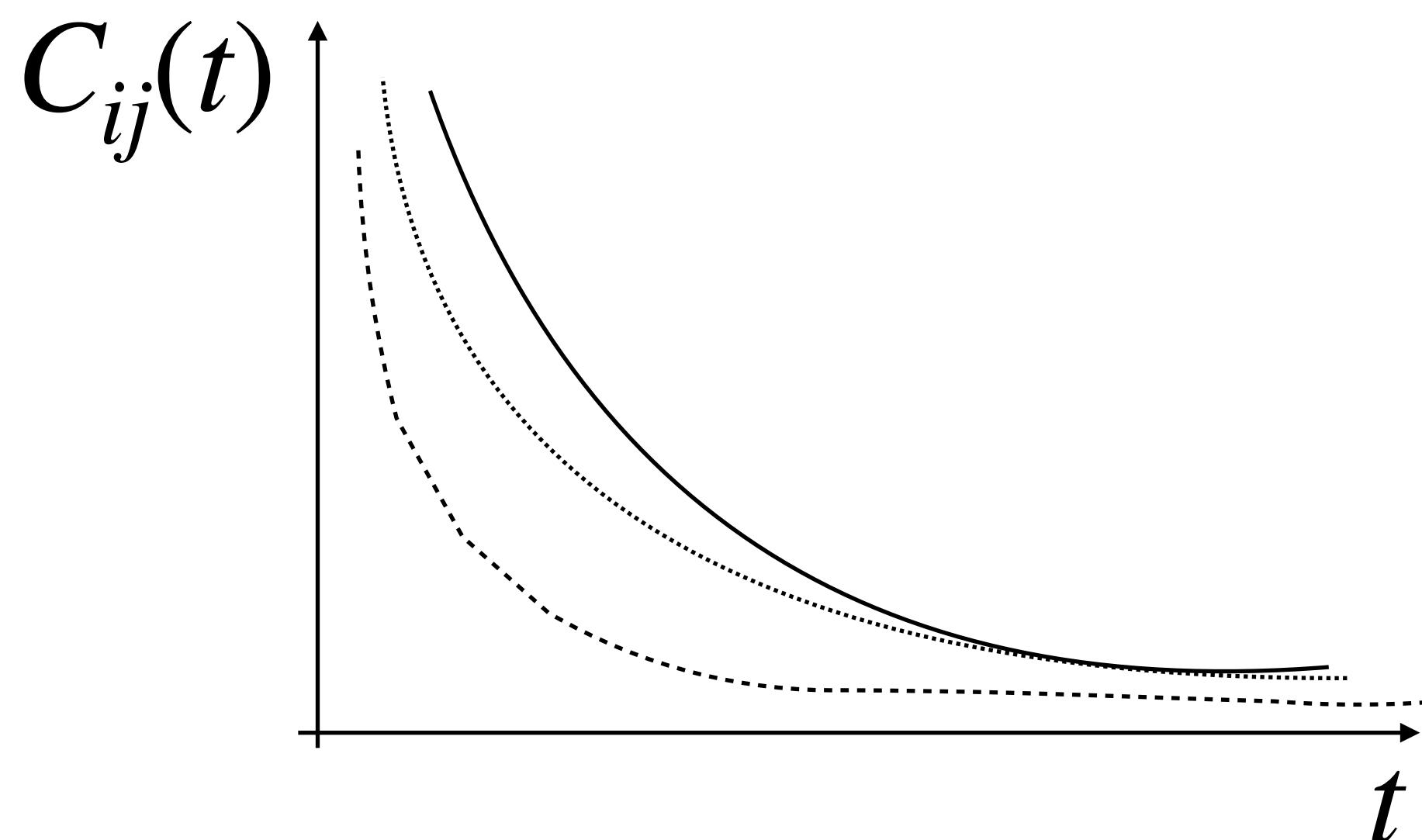


Multi-exponential fit

- To extract the finite-volume spectrum

- Two-point correlation functions $\chi_i^\dagger = \chi_j^\dagger = \text{qqq}$

$$C_{ij}(t) = \langle \chi_i(t) \chi_j^\dagger(0) \rangle = \sum_{n=0}^{\infty} W_n e^{-E_n t}$$



Constraint 1

$$C_{N,2\text{pt}}(t) = W_0 e^{-E_0 t}$$

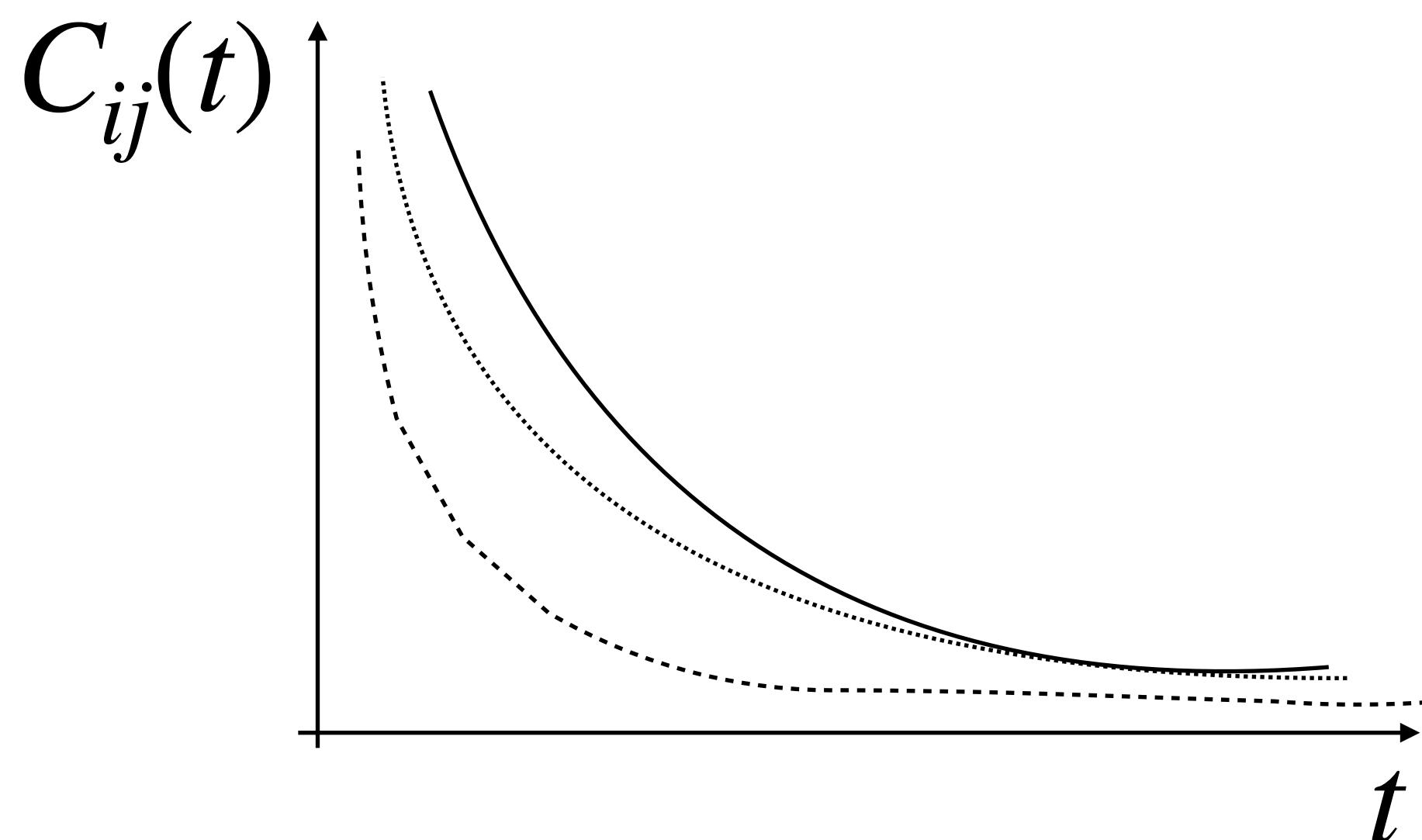
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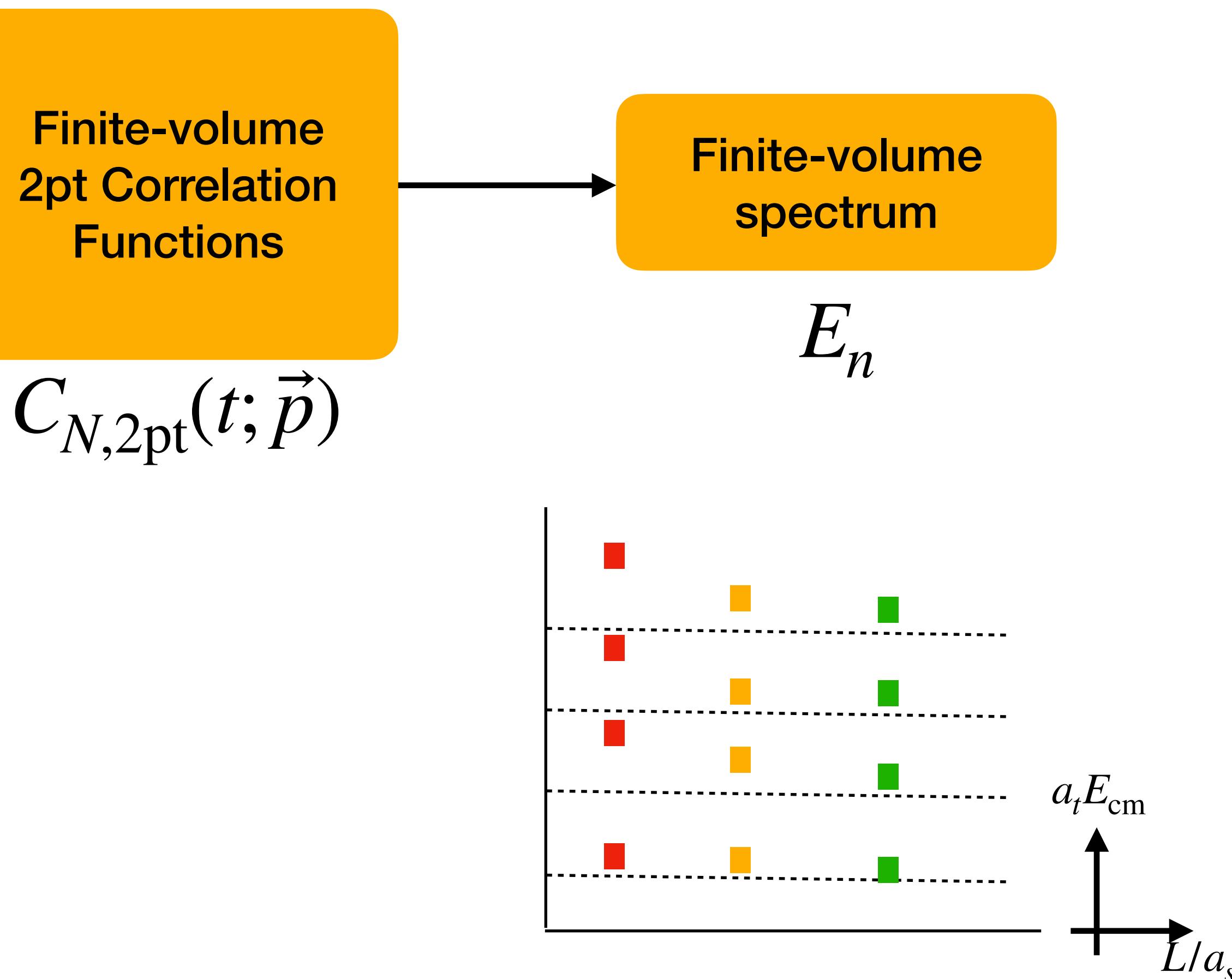
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$$\begin{array}{c} n=0 \\ t \gg 0 \end{array}$$

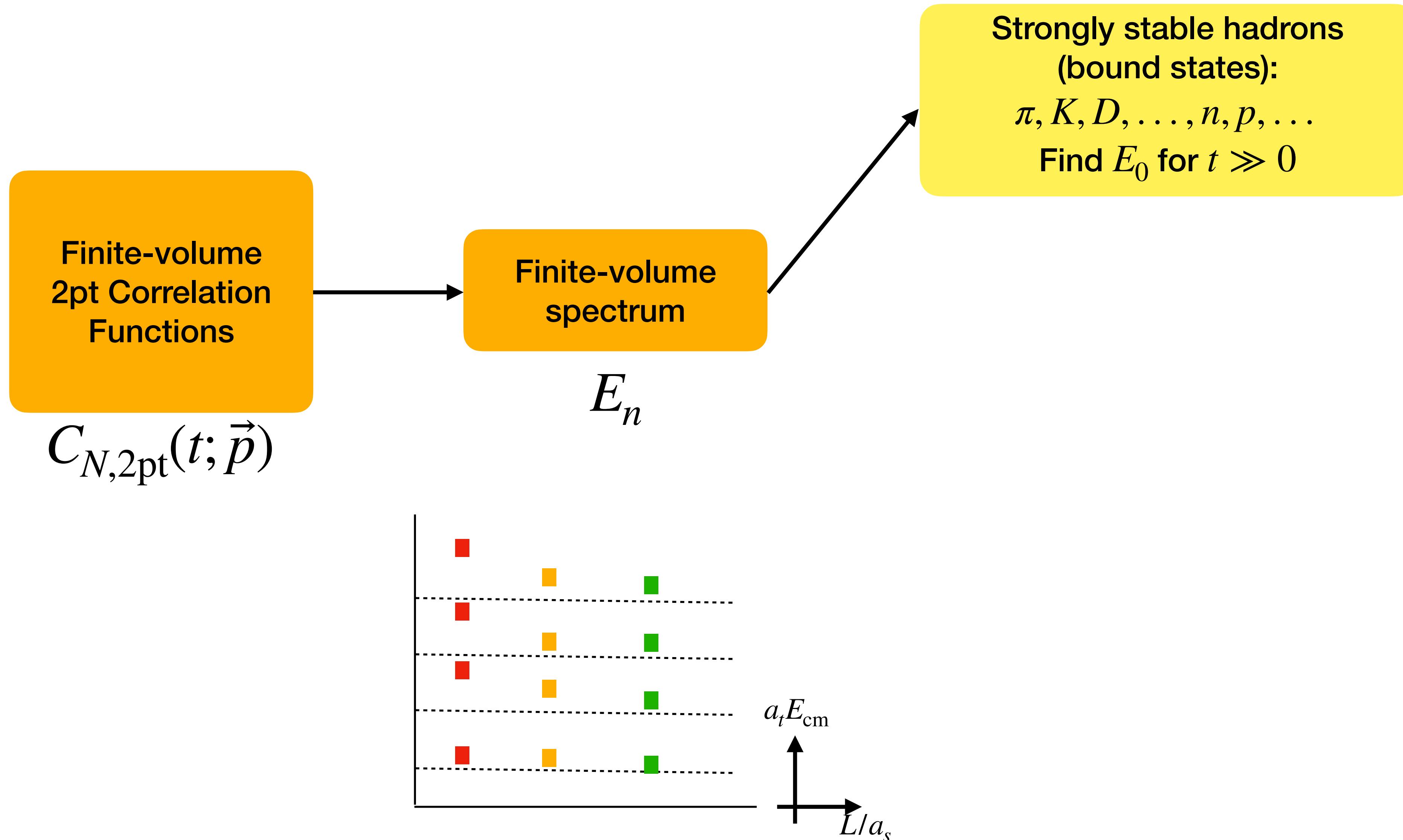
Strongly stable hadrons
(bound states):
 $\pi, K, D, \dots, n, p, \dots$
Find E_0

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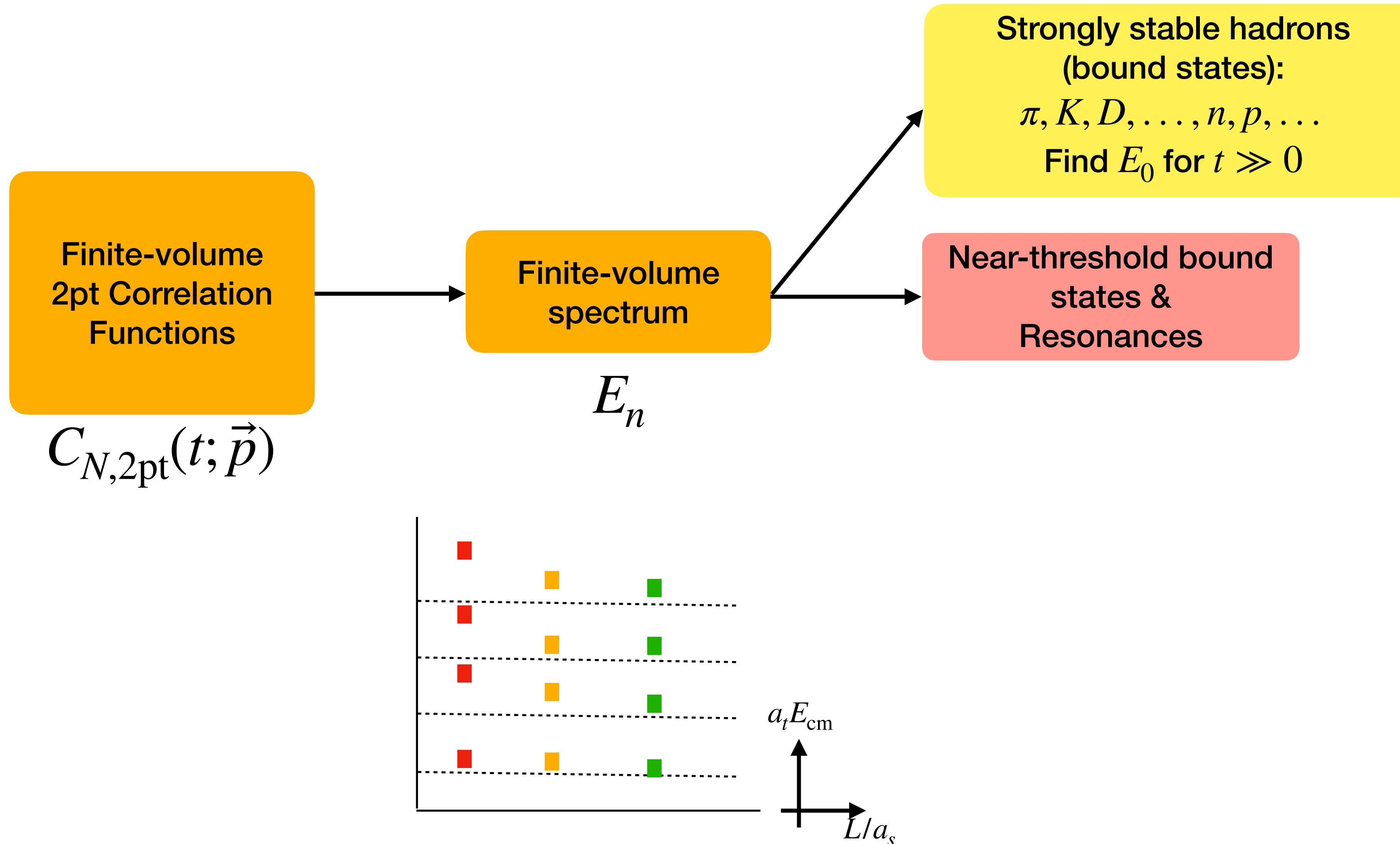
Finite-volume spectrum from lattice QCD



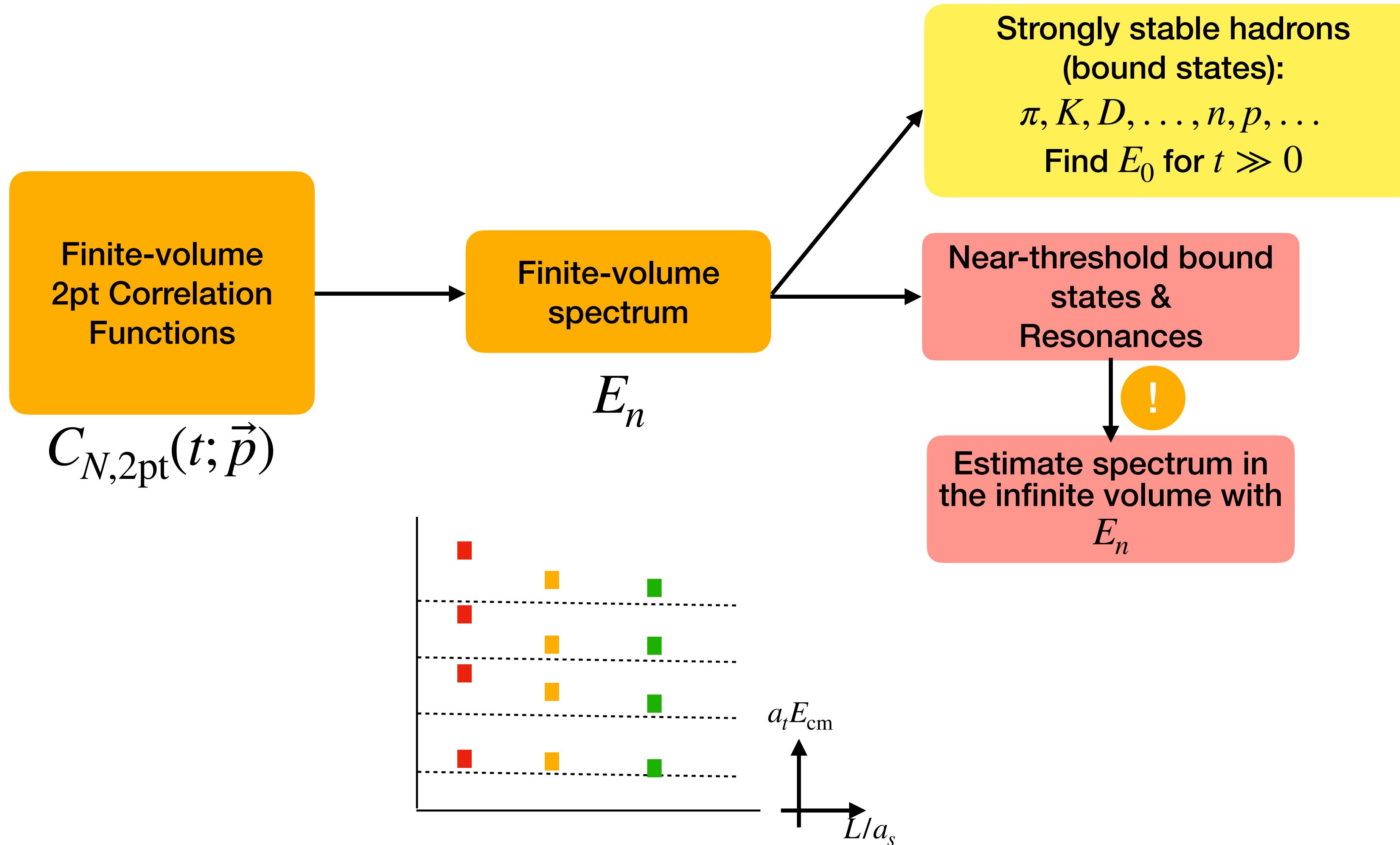
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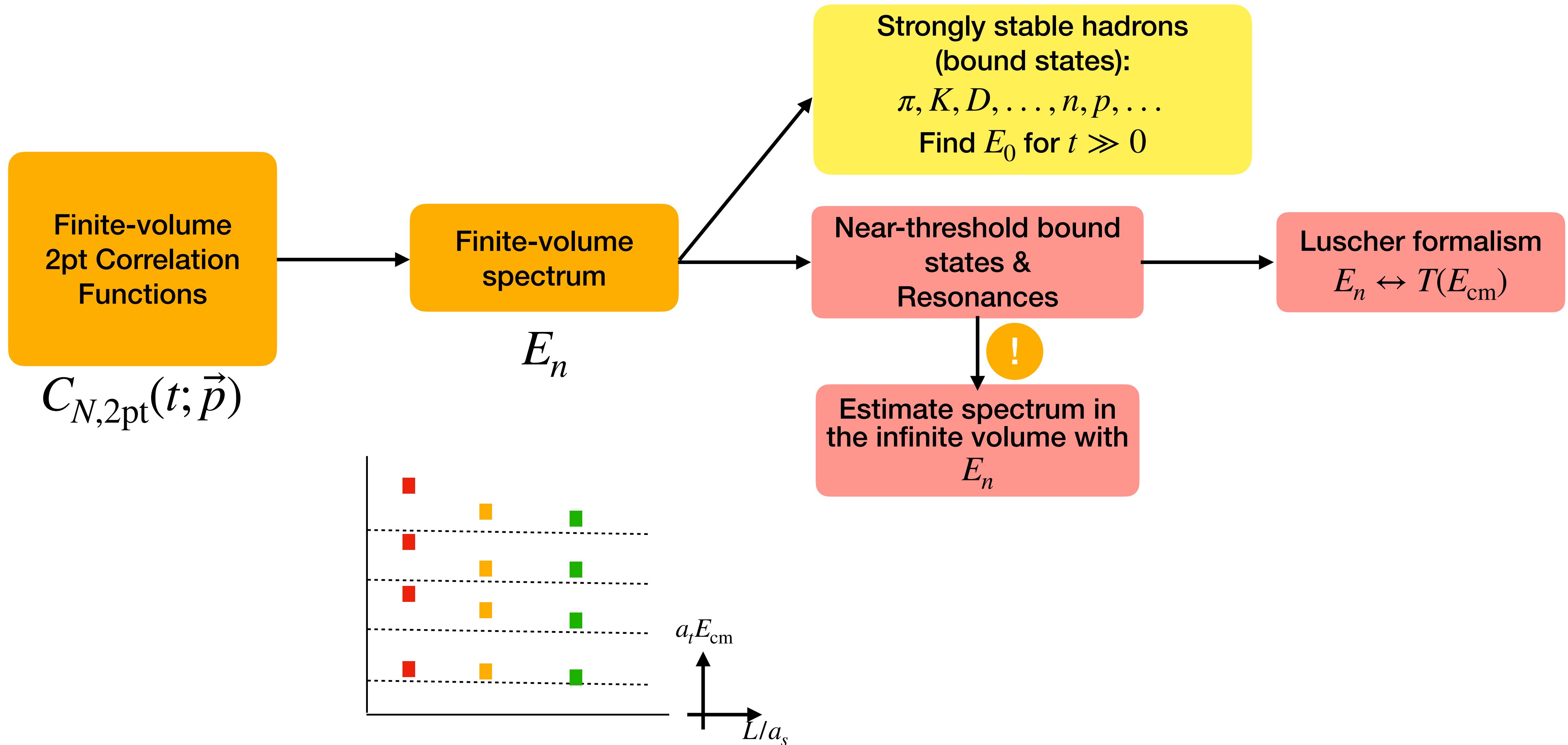
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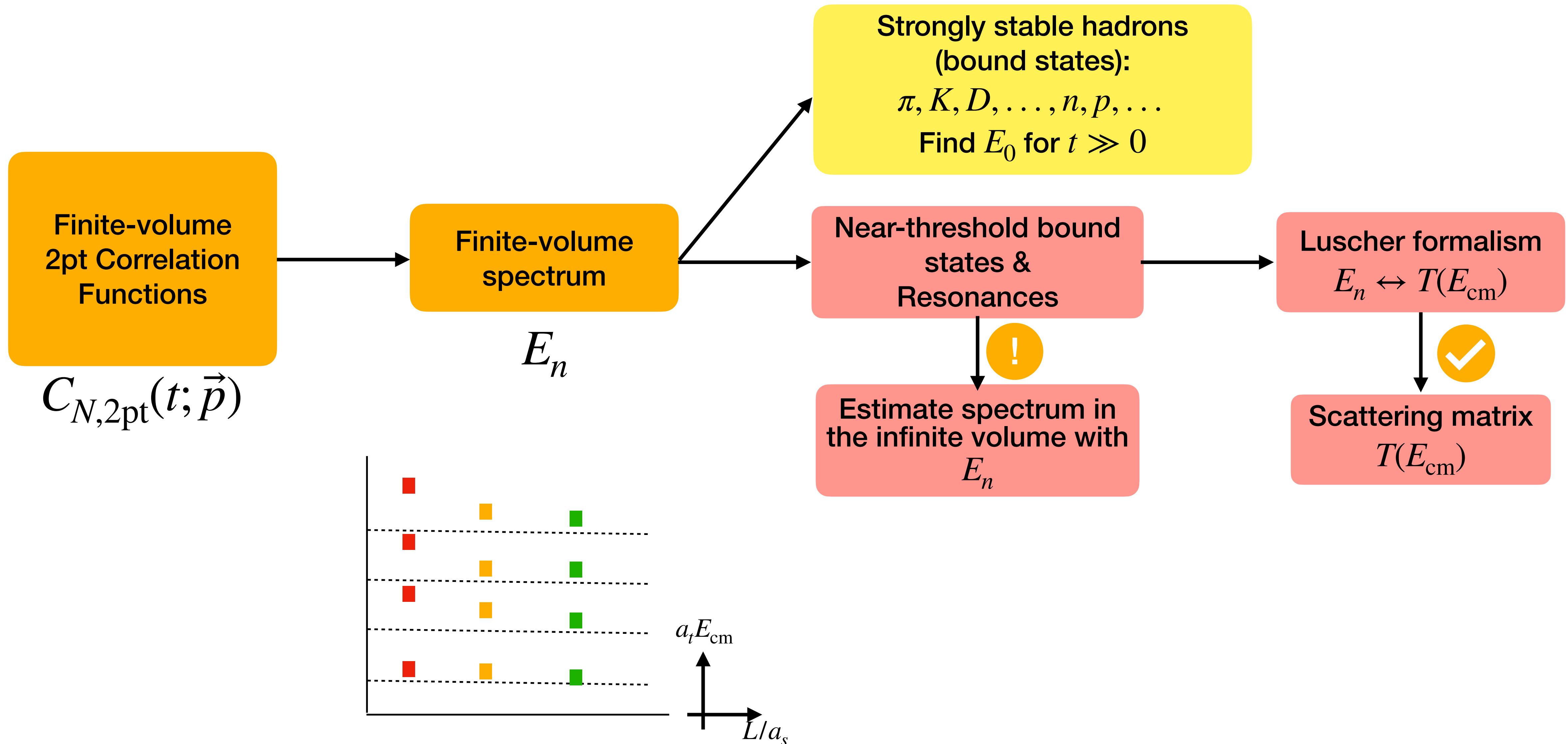
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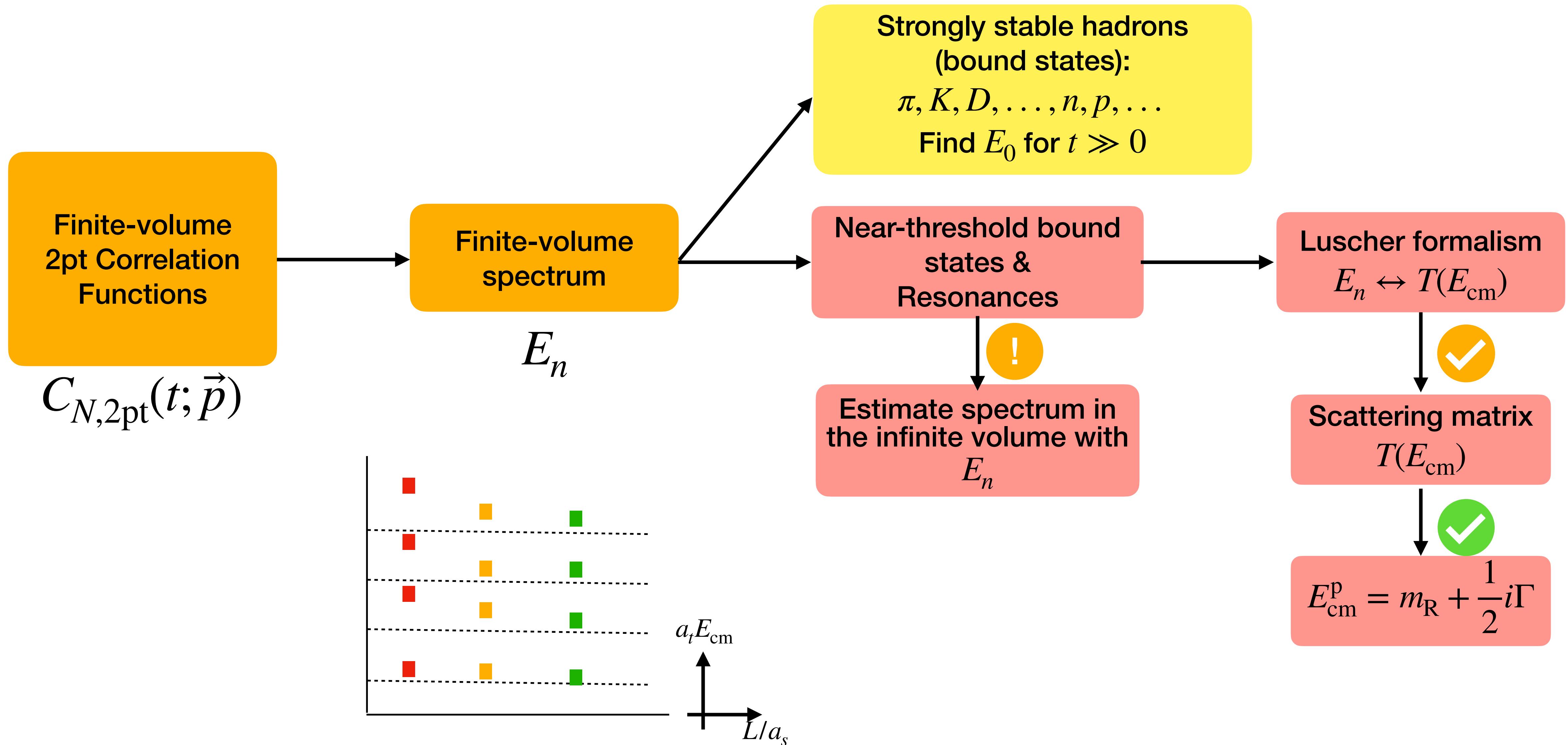
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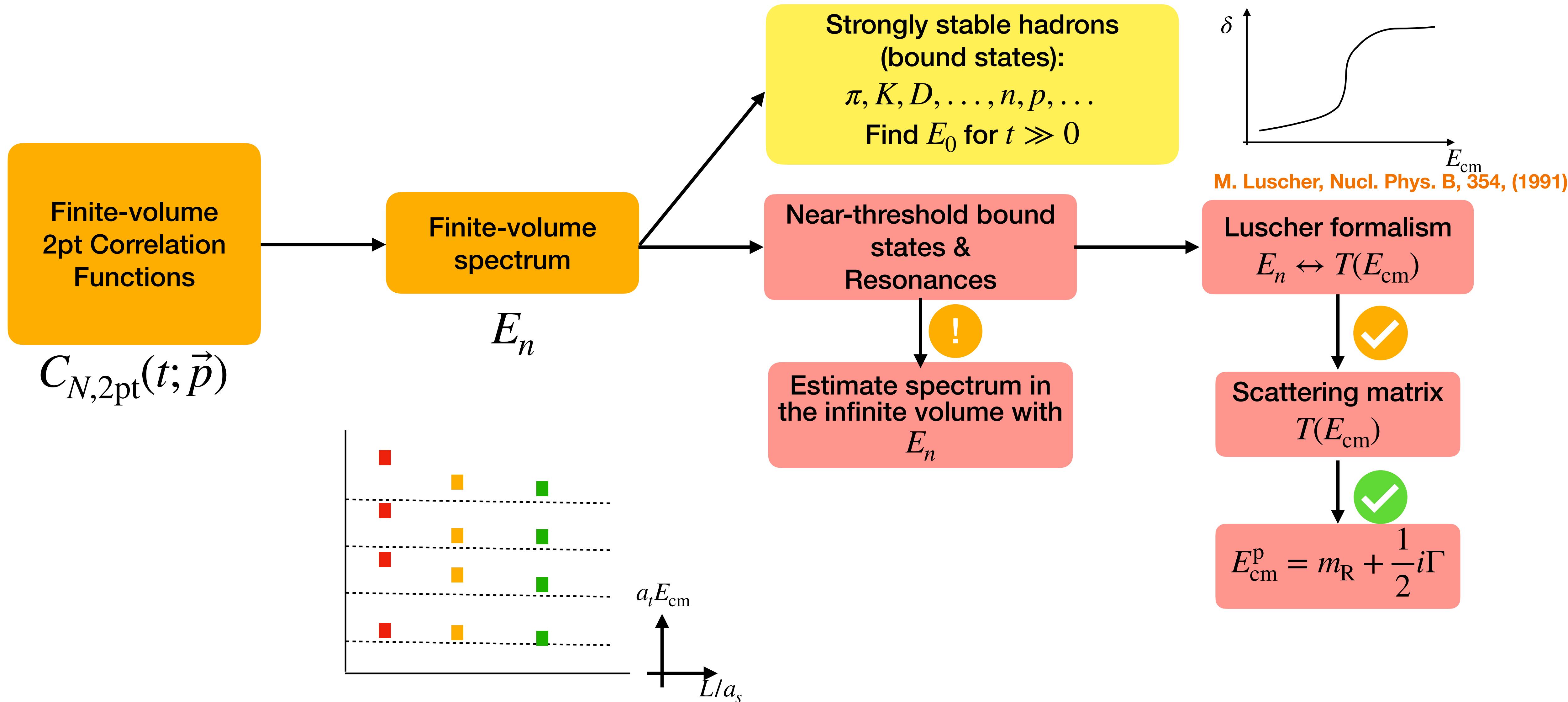
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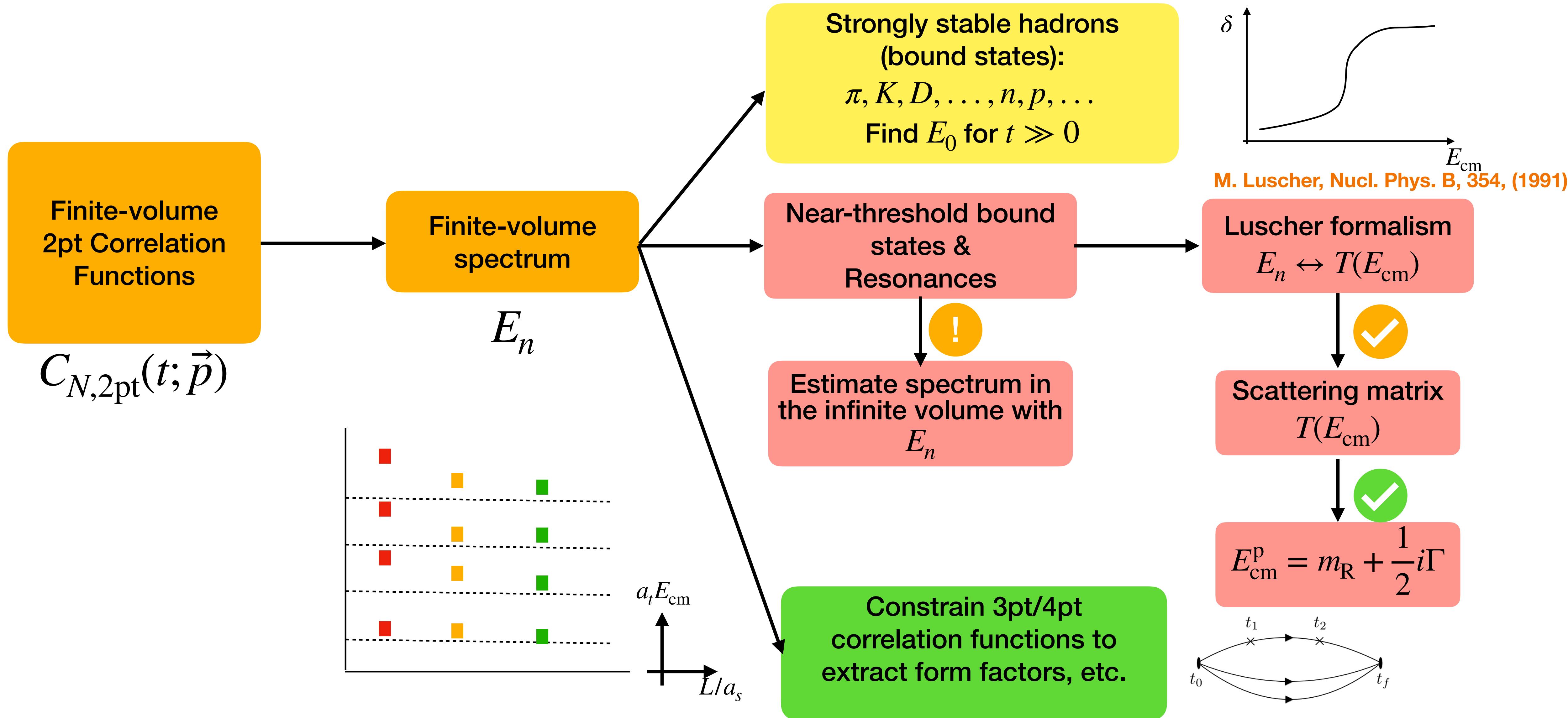
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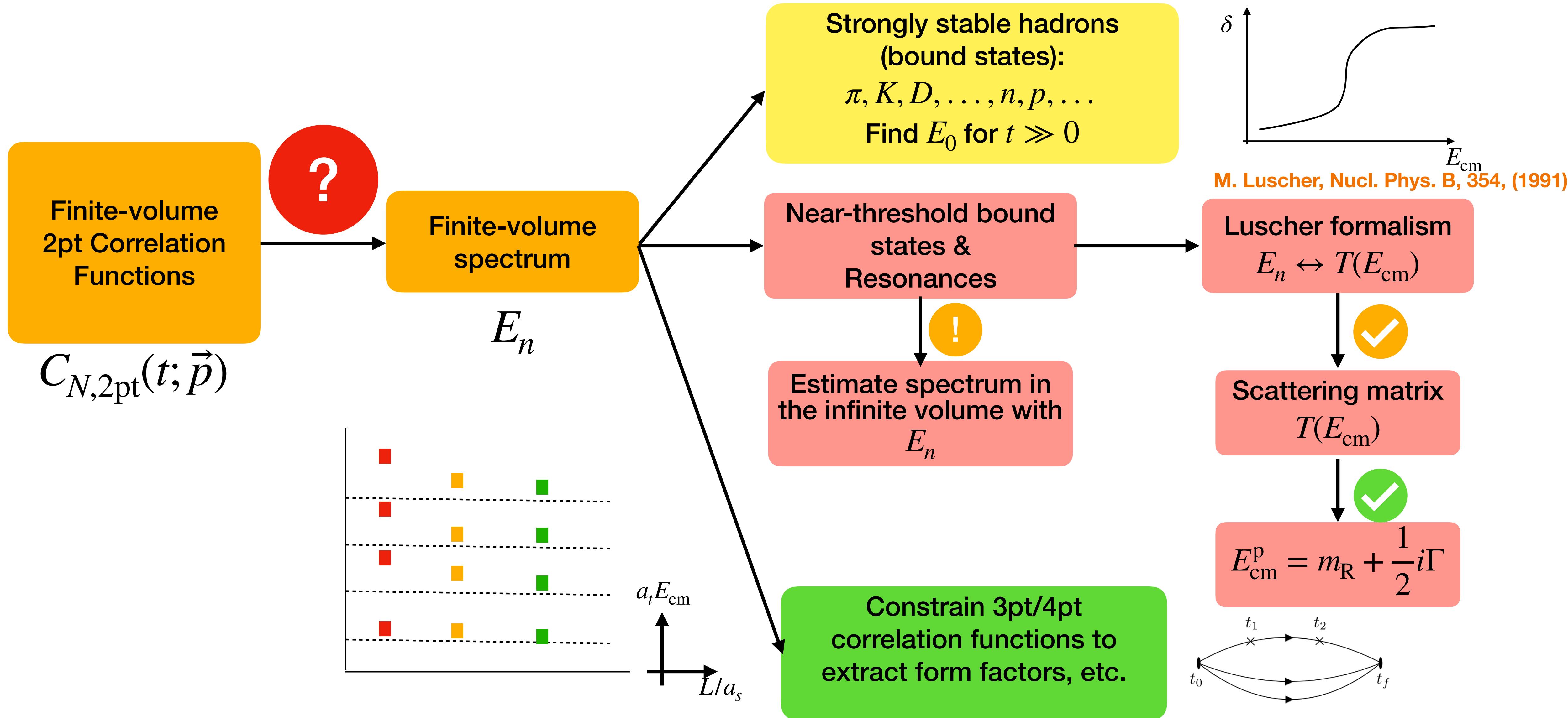
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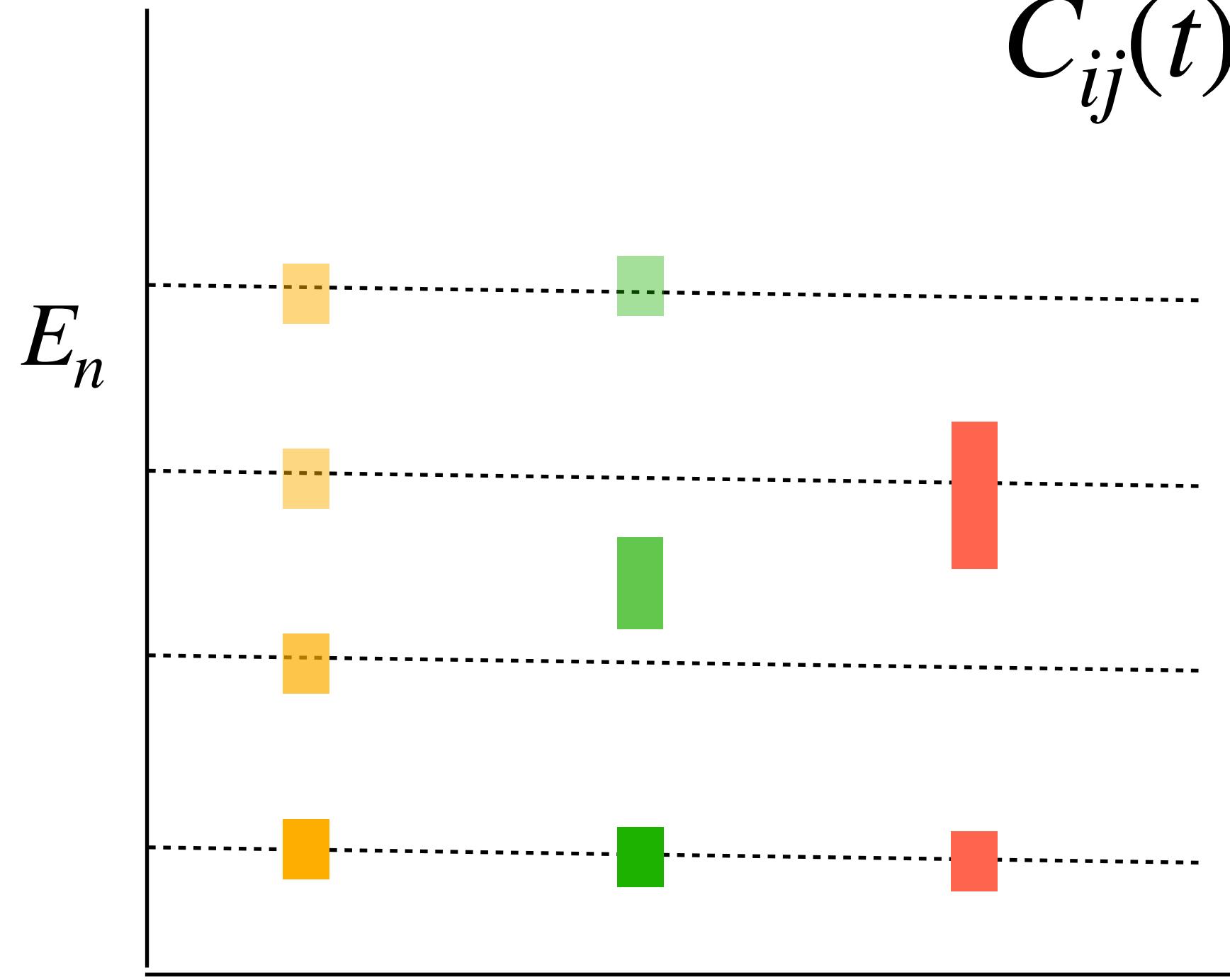
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Multi-exponential fit

- To extract the finite-volume spectrum

- Two-point correlation functions $\chi_i^\dagger = \chi_j^\dagger =$



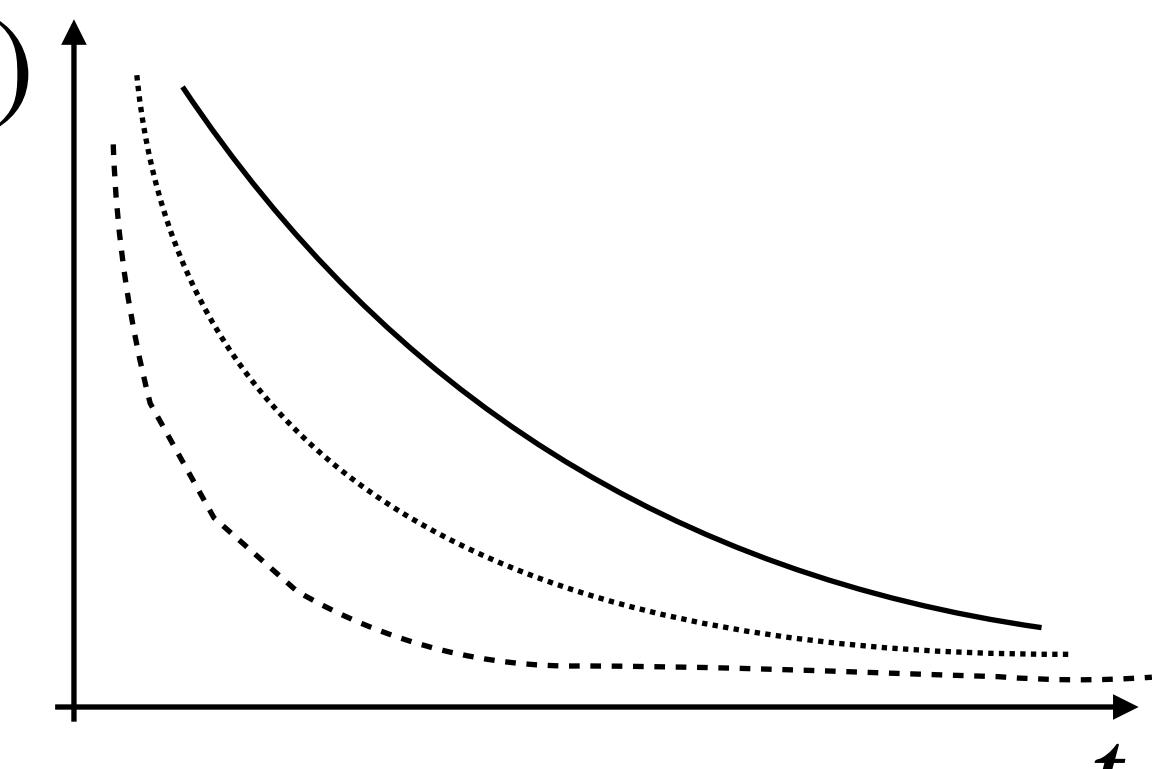
$$N = 4 \quad N = 3 \quad N = 2$$

$$C_{ij}(t) = \langle \chi_i(t) \chi_j^\dagger(0) \rangle = \sum_{n=0}^{\infty} W_n e^{-E_n t}$$

Constraint 1

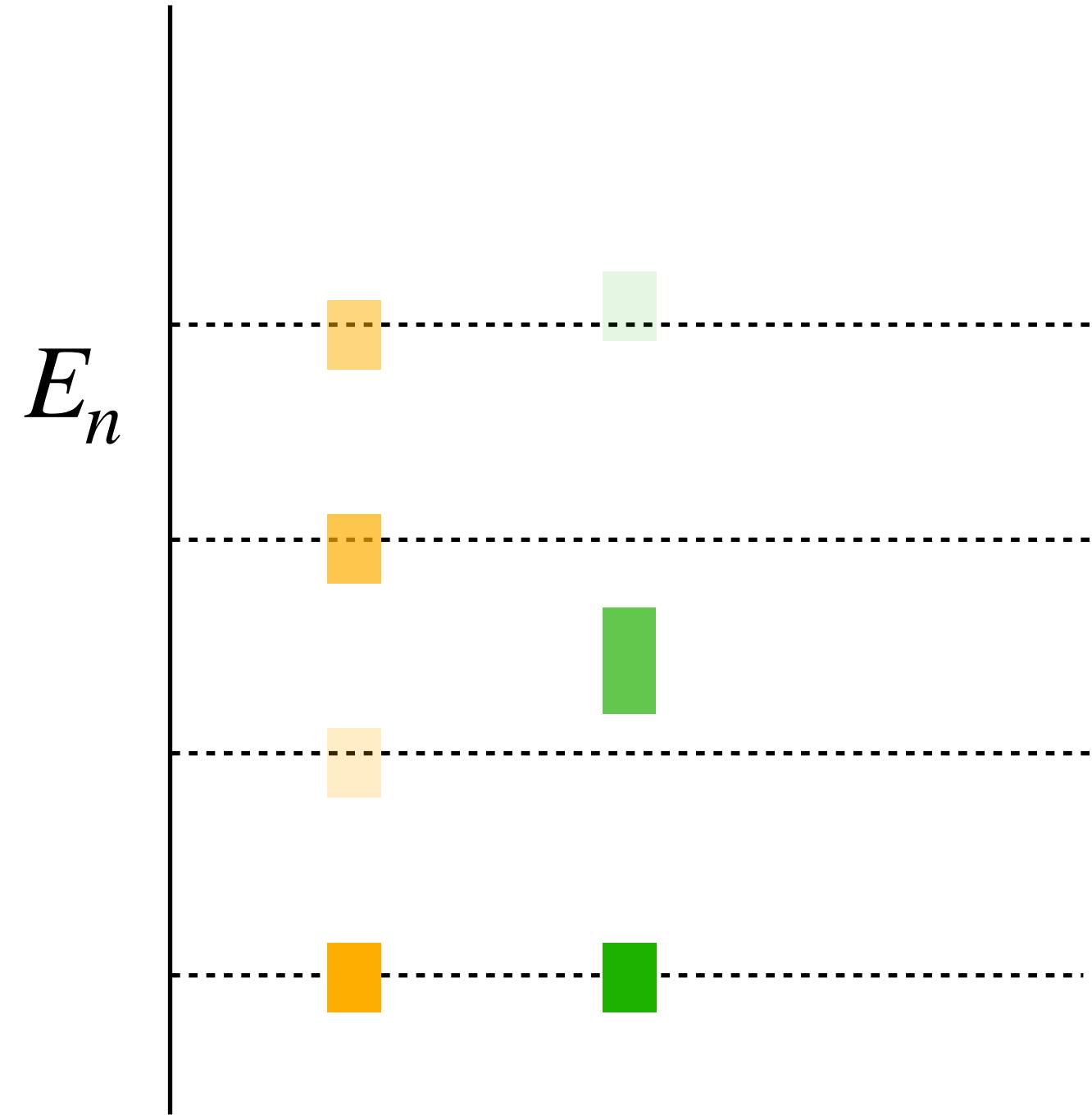
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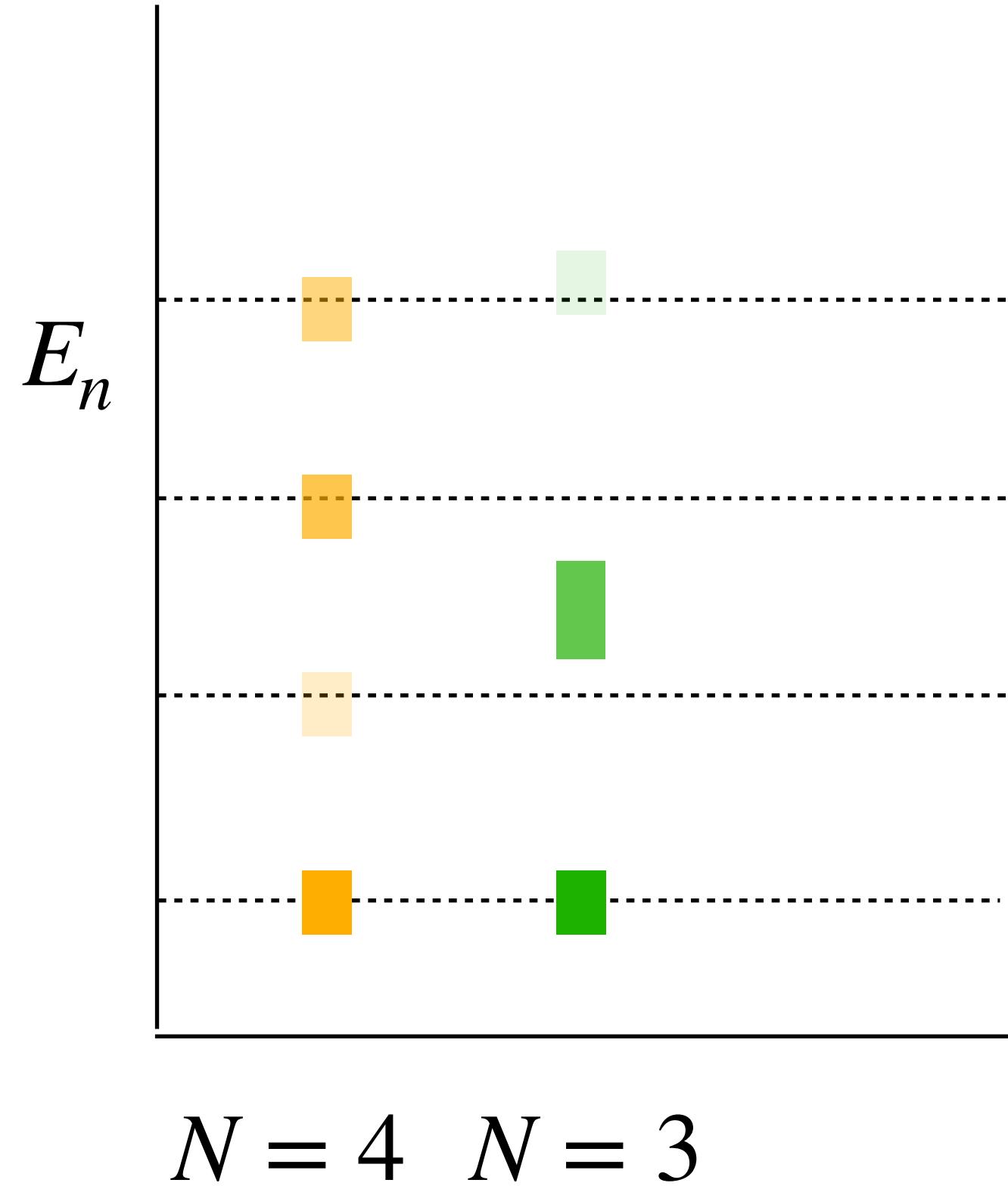
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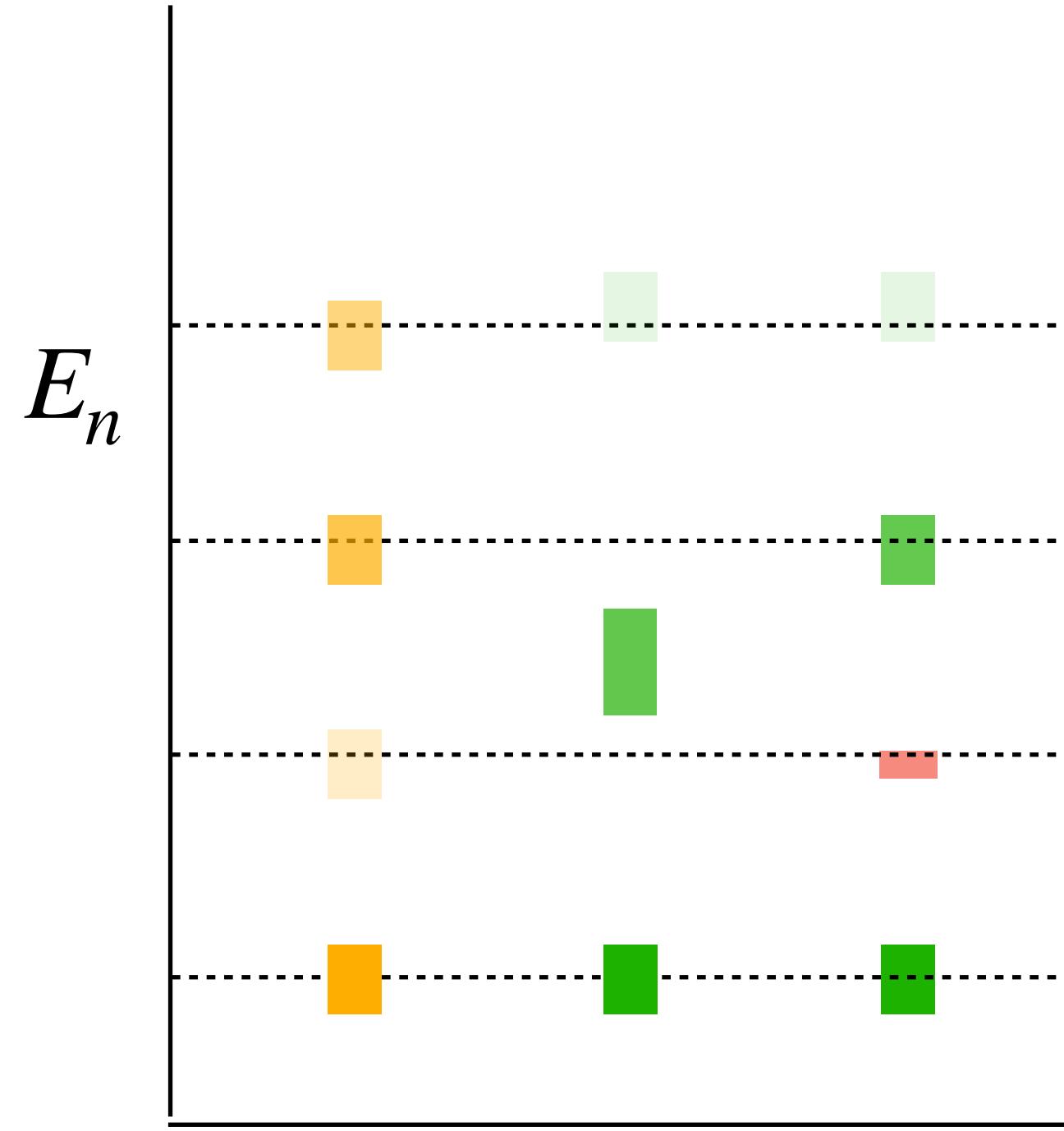
G. P. Lepage et al., arXiv: hep-lat/0110175



$$\chi_{\text{prior}}^2 \equiv \sum_n \frac{(W_n - \tilde{W}_n)^2}{\tilde{\sigma}_{W_n}^2} + \sum_n \frac{(E_n - \tilde{E}_n)^2}{\tilde{\sigma}_{E_n}^2}$$

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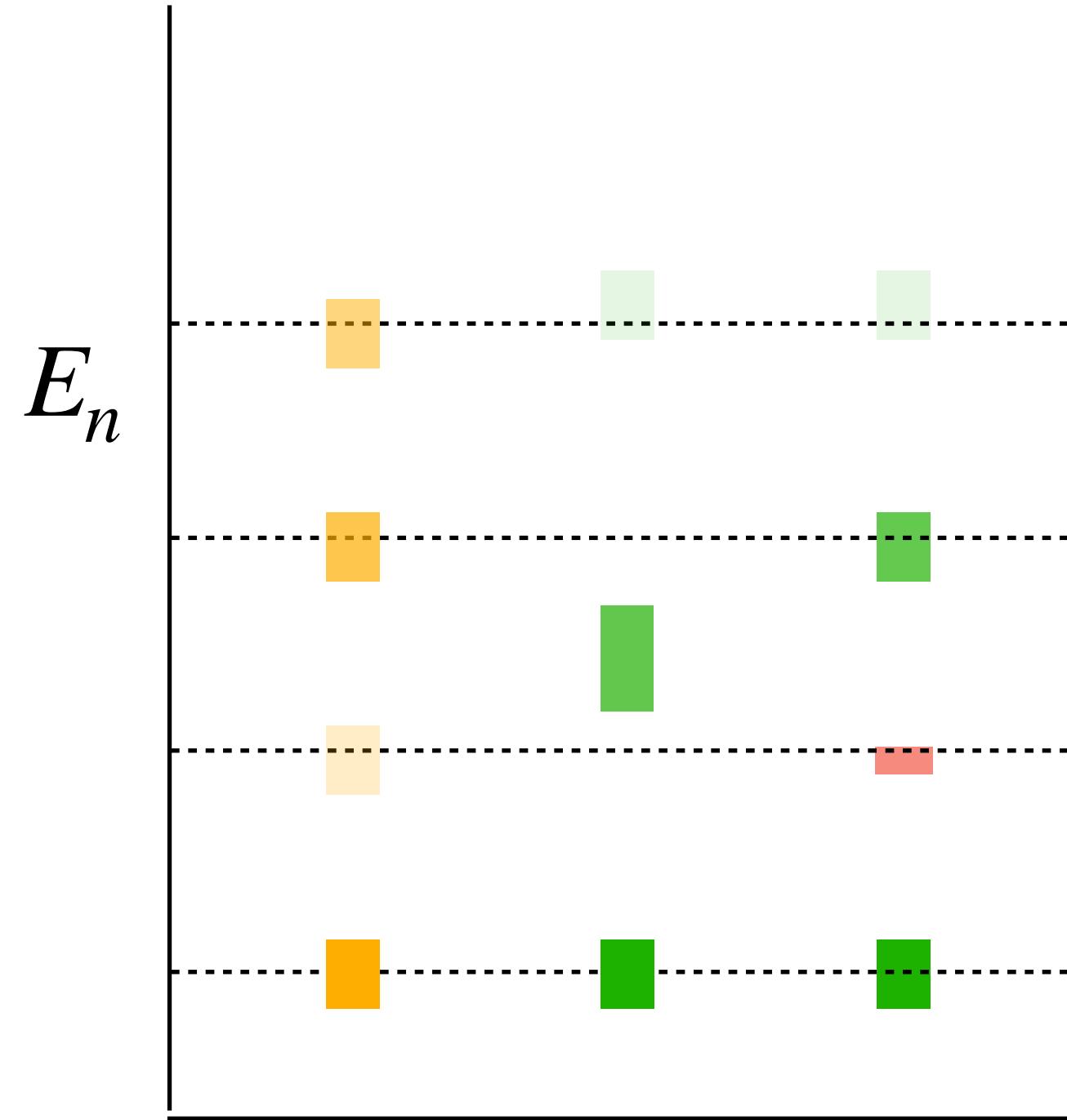
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Q1: How to choose
 N ,
 \tilde{W}_n , $\tilde{\sigma}_{W_n}$,
 \tilde{E}_n , $\tilde{\sigma}_{E_n}$? ?

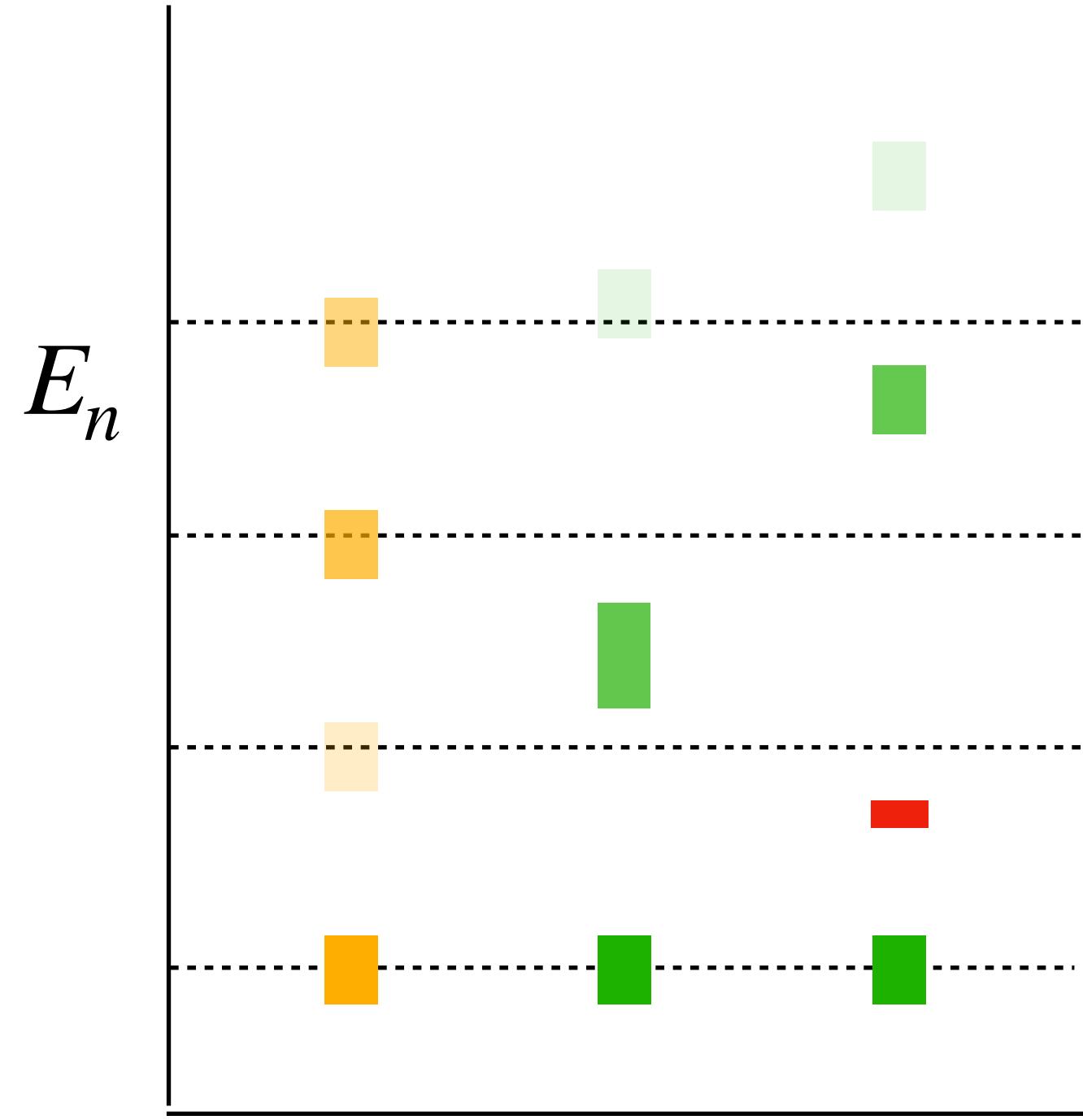
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- To extract the finite-volume spectrum
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$N = 4 \quad N = 3 \quad +Add prior$

$$C_{ij}(t) = \langle \chi_i(t) \chi_j^\dagger(0) \rangle = \sum_{n=0}^{\infty} W_n e^{-E_n t}$$

Constraint 1

$$C_{N,2\text{pt}}(t) = \sum_{n=0}^N W_n e^{-E_n t}$$

Q1: How to choose N , \tilde{W}_n , $\tilde{\sigma}_{W_n}$, \tilde{E}_n , $\tilde{\sigma}_{E_n}$?

$$\chi^2 \equiv \sum_{ij} (D_i - \tilde{D}_i) C_{ij}^{-1} (D_j - \tilde{D}_j)$$

G. P. Lepage et al., arXiv: hep-lat/0110175

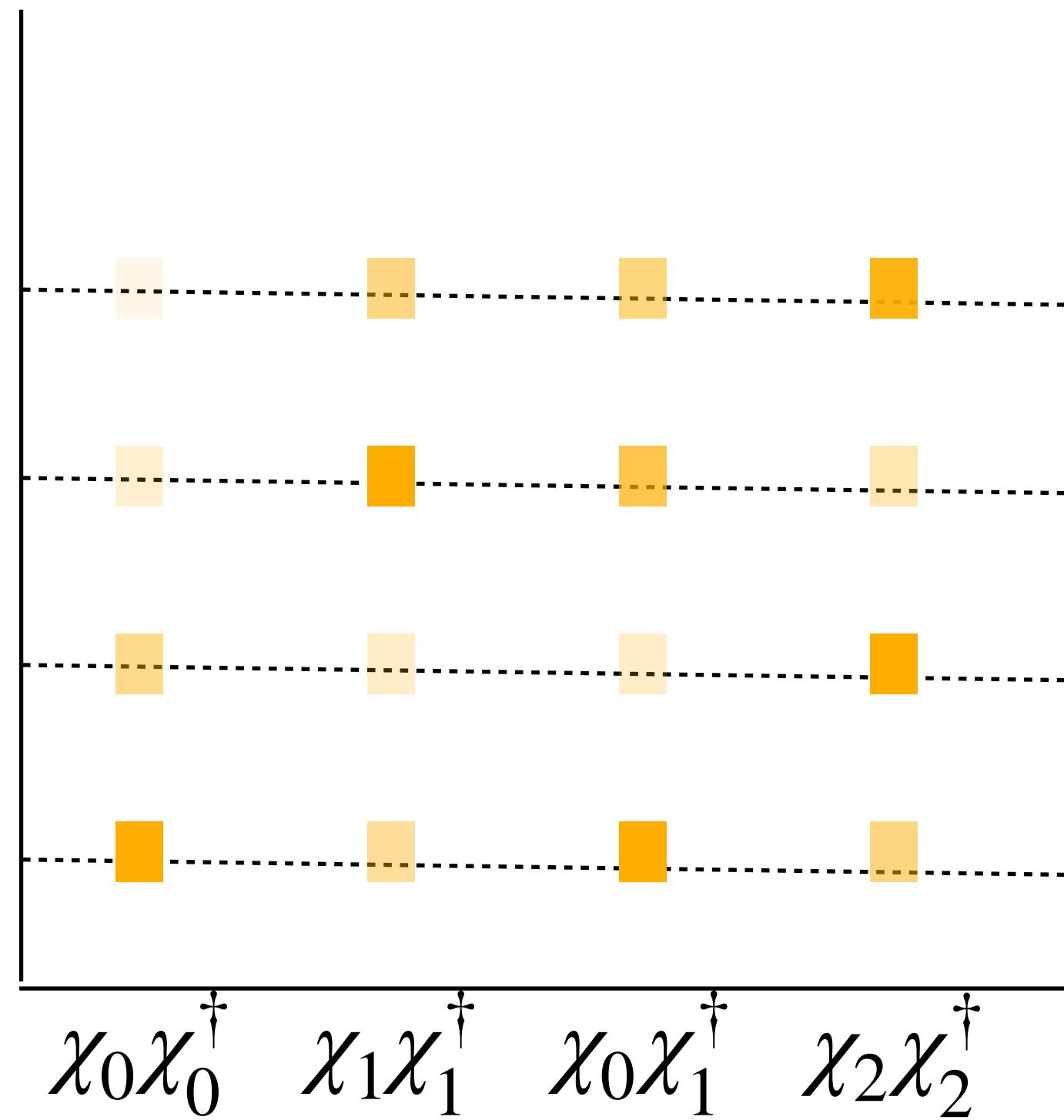


$$\chi_{\text{prior}}^2 \equiv \sum_n \frac{(W_n - \tilde{W}_n)^2}{\tilde{\sigma}_{W_n}^2} + \sum_n \frac{(E_n - \tilde{E}_n)^2}{\tilde{\sigma}_{E_n}^2}$$

Multi-exponential fit

- To extract the finite-volume spectrum
 - Two-point correlation functions

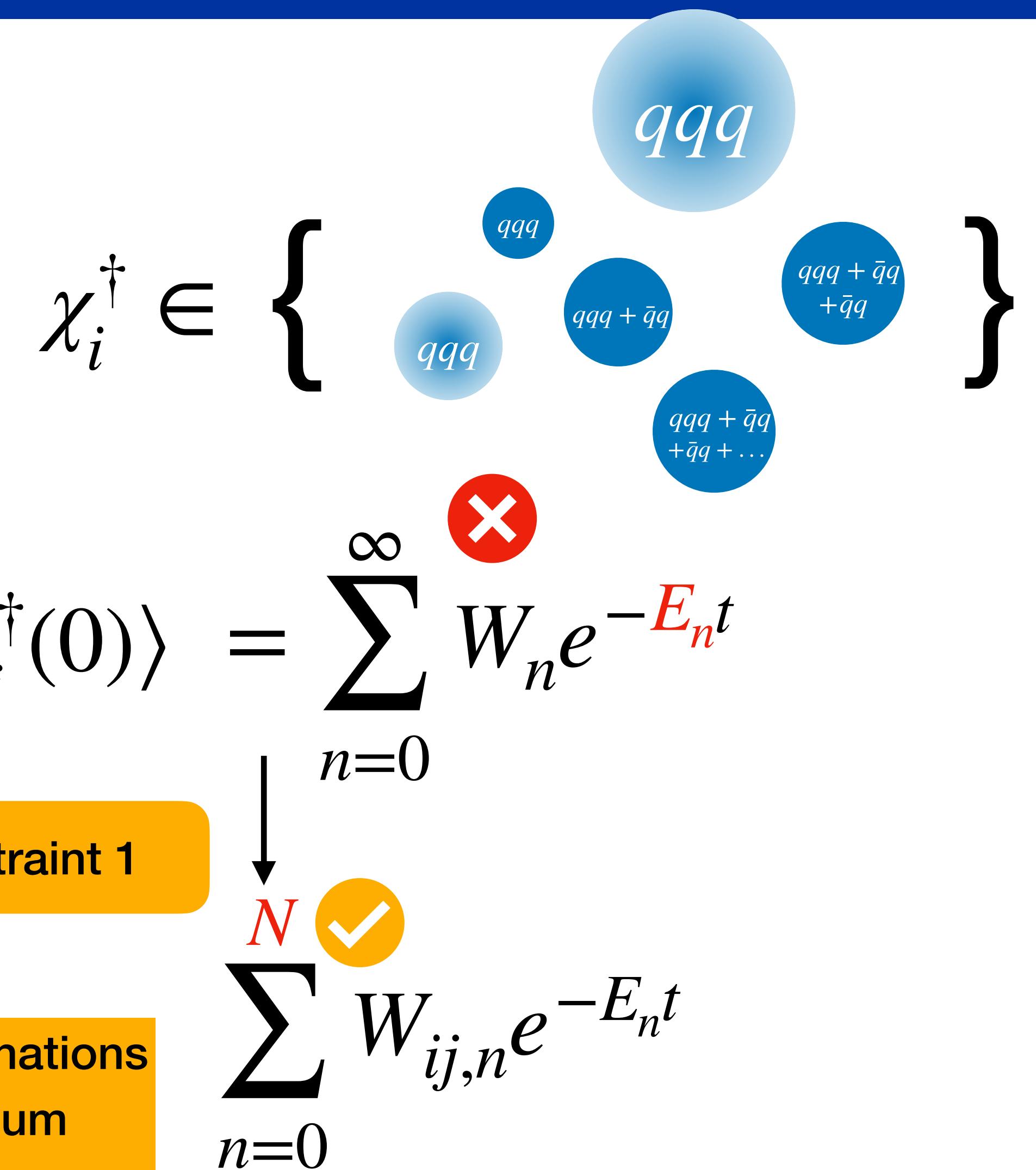
interpolating operators with the same quantum numbers as the hadron state



$$C_{ij}(t) = \langle \chi_i(t) \chi_j^\dagger(0) \rangle = \sum_{n=0}^{\infty} W_n e^{-E_n t}$$

Constraint 1

Why not find linear combinations
of χ_i^\dagger 's with the maximum
overlap to the states n 's?

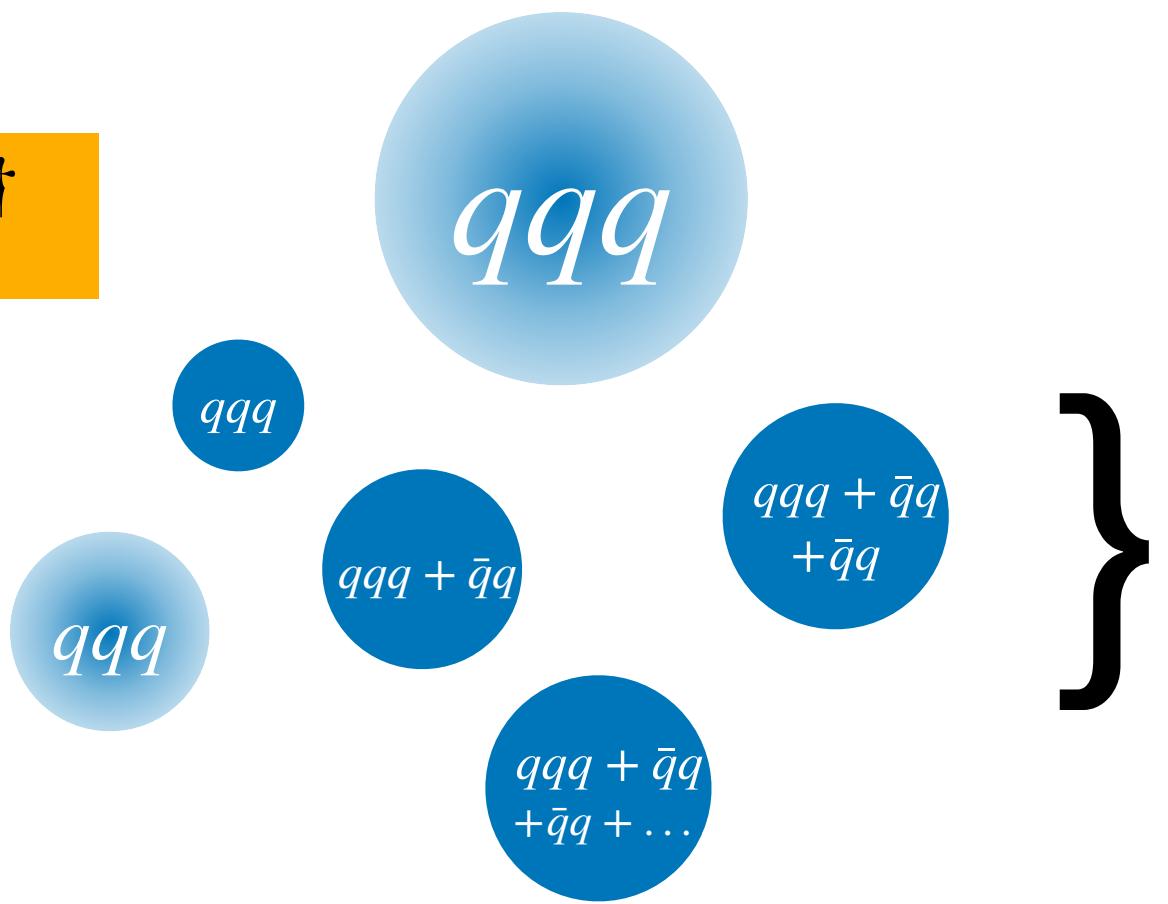


Multi-exponential fit

- To extract the finite-volume spectrum
- Two-point correlation functions N ✓

$$C_{ij}(t) = \langle \chi_i(t) \chi_j^\dagger(0) \rangle \rightarrow \sum_{n=0} W_{ij,n} e^{-E_n t} \quad \chi_i^\dagger \in \{$$

Choose N basis operators: χ_i^\dagger



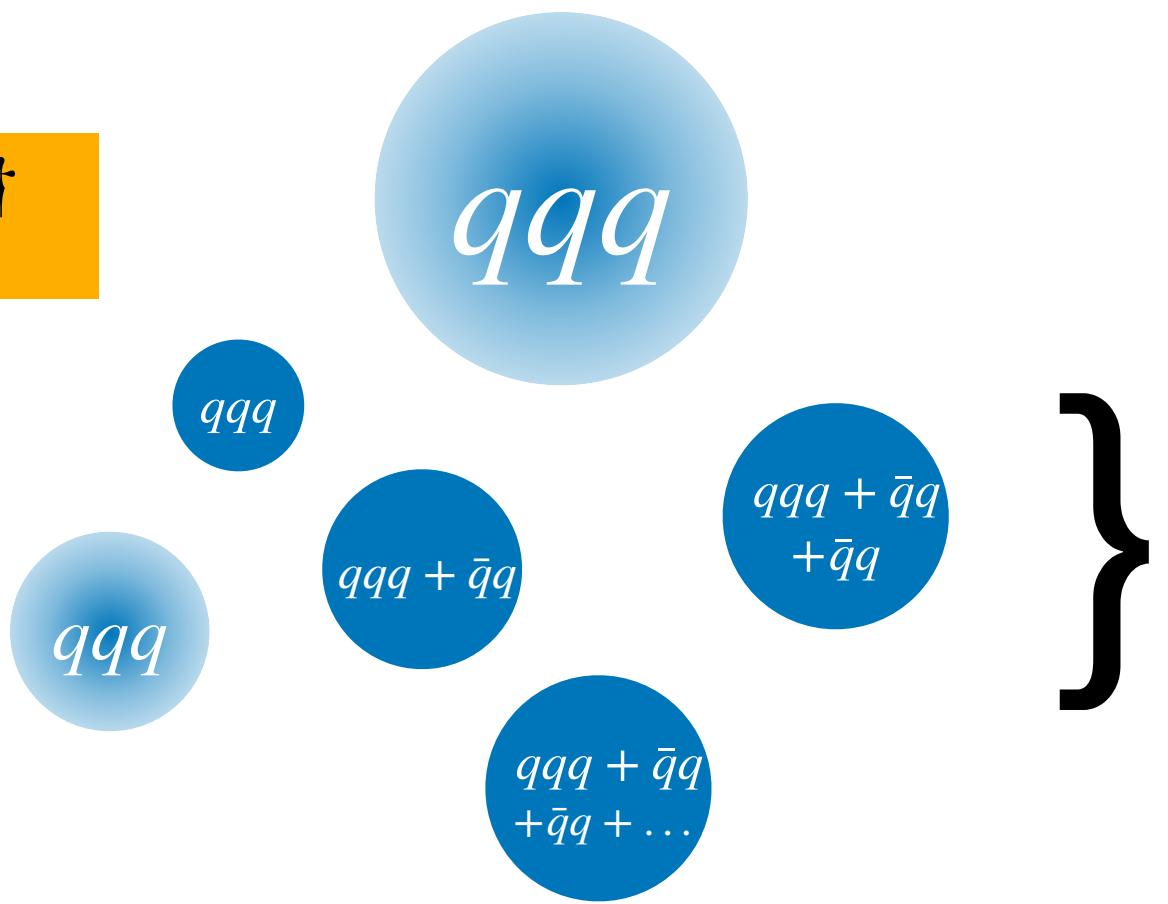
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Find linear combinations of χ_i^\dagger 's

$$X_n^\dagger = \sum_i v_i^{(n)} \chi_i^\dagger$$

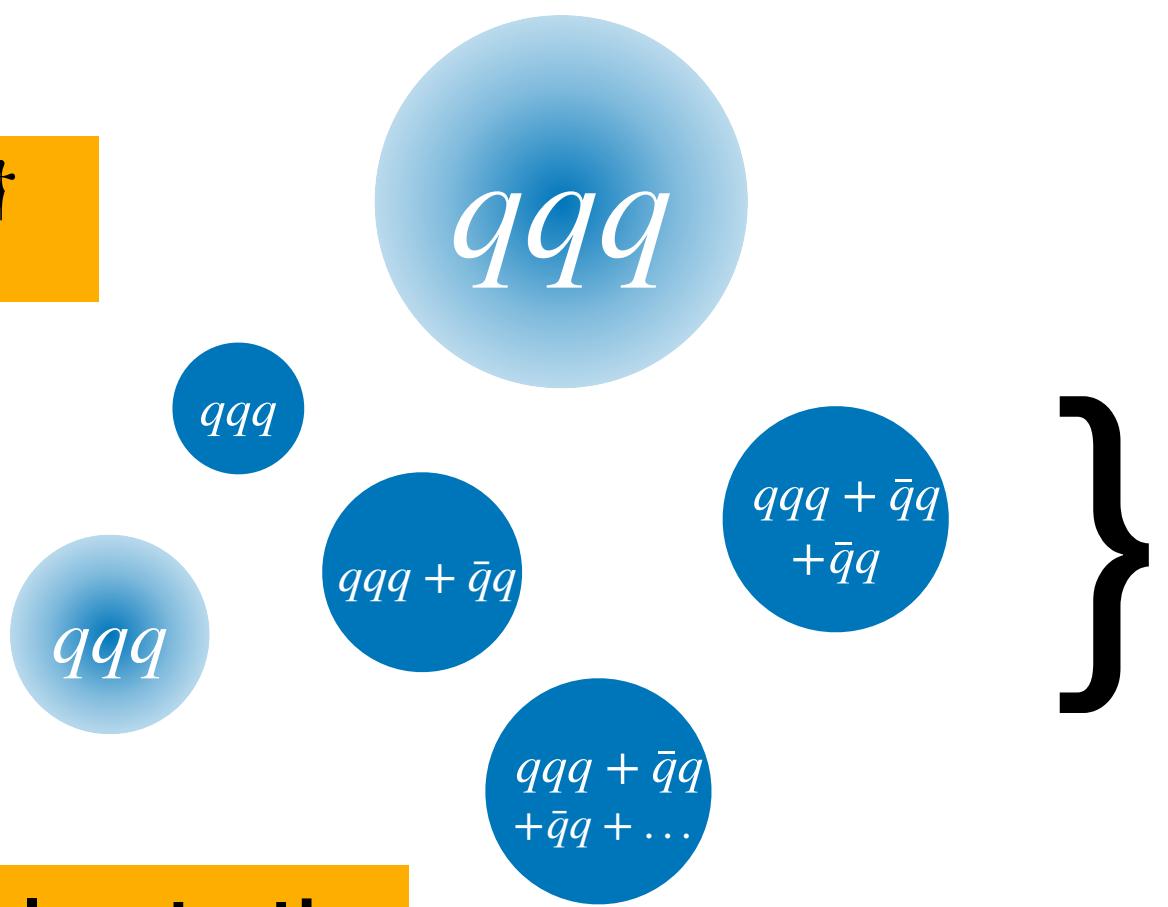
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Choose N basis operators: χ_i^\dagger



Find linear combinations of χ_i^\dagger 's

$$X_n^\dagger = \sum_i v_i^{(n)} \chi_i^\dagger$$

with the maximum overlap to the states n's

$$\tilde{C}_{\text{2pt}}^{(n)}(t) = \langle X^{(n)}(t) X^{(n)\dagger}(0) \rangle = 1 \times e^{-E_n t} + \sum_{k \neq n} |w_k|^2 e^{-E_k t}$$

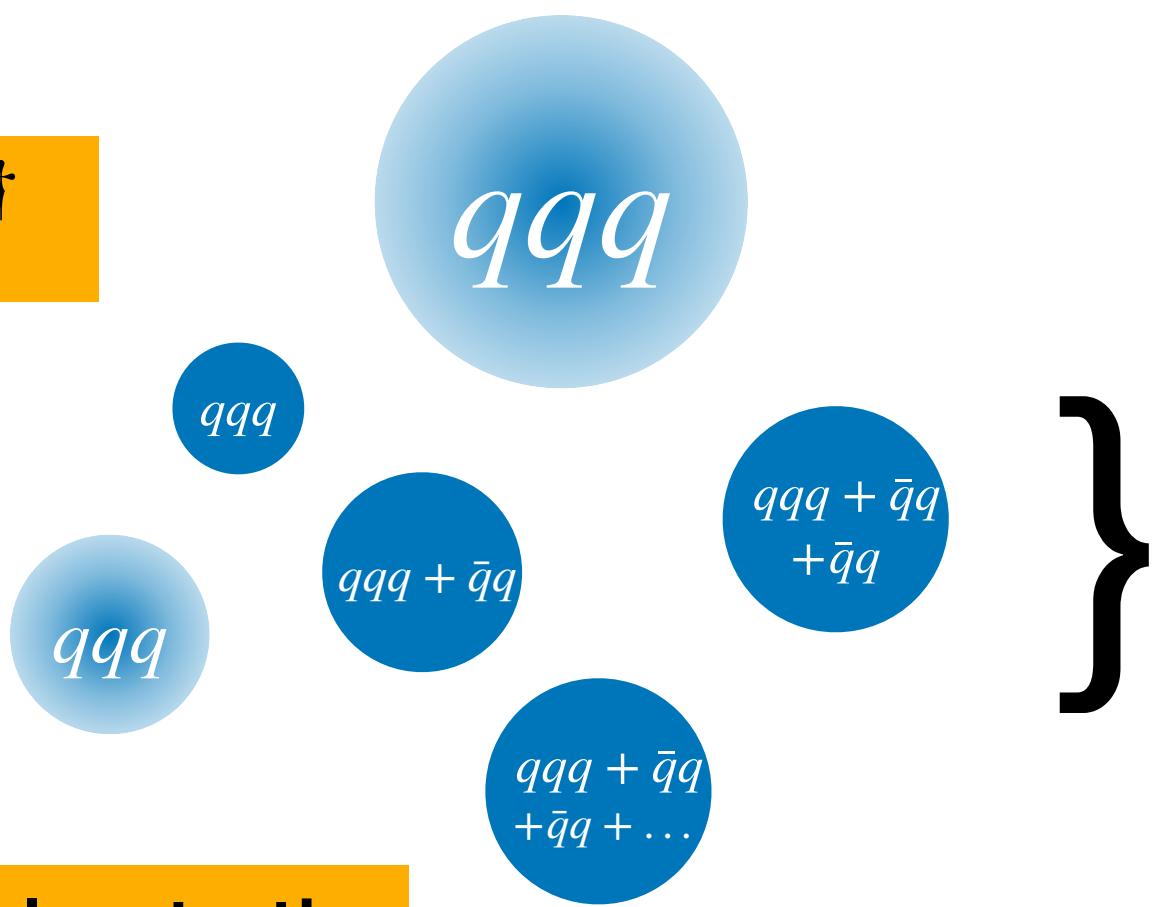
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$$\mathbf{C}(t) = \begin{pmatrix} C_{00}(t) & C_{01}(t) & \dots & C_{0N}(t) \\ C_{10}(t) & C_{11}(t) & \dots & C_{1N}(t) \\ \dots & \dots & \dots & \dots \\ C_{N0}(t) & C_{n1}(t) & \dots & C_{NN}(t) \end{pmatrix}$$

$$v_i^{(n)*} C_{ij}(t_0) v_j^{(n')} = \delta_{n,n'}$$

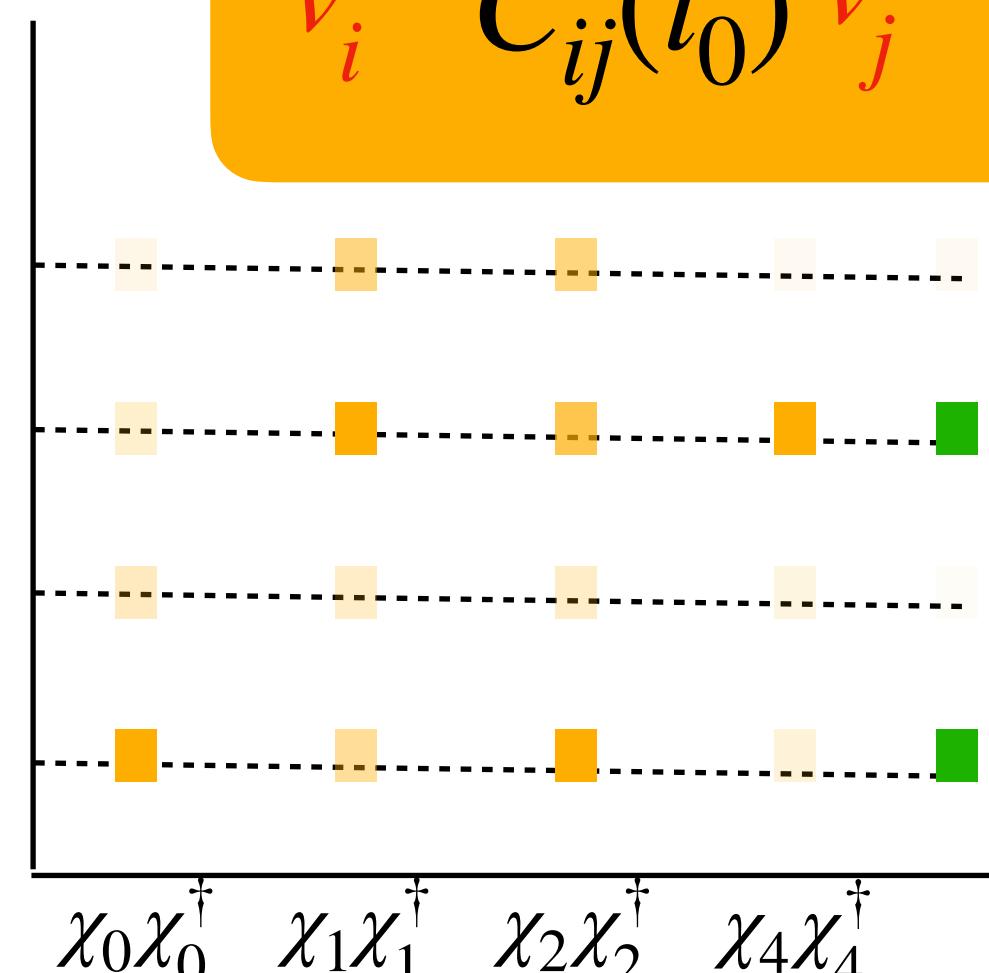
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Choose N basis operators: χ_i^\dagger

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with the maximum overlap to the states n 's

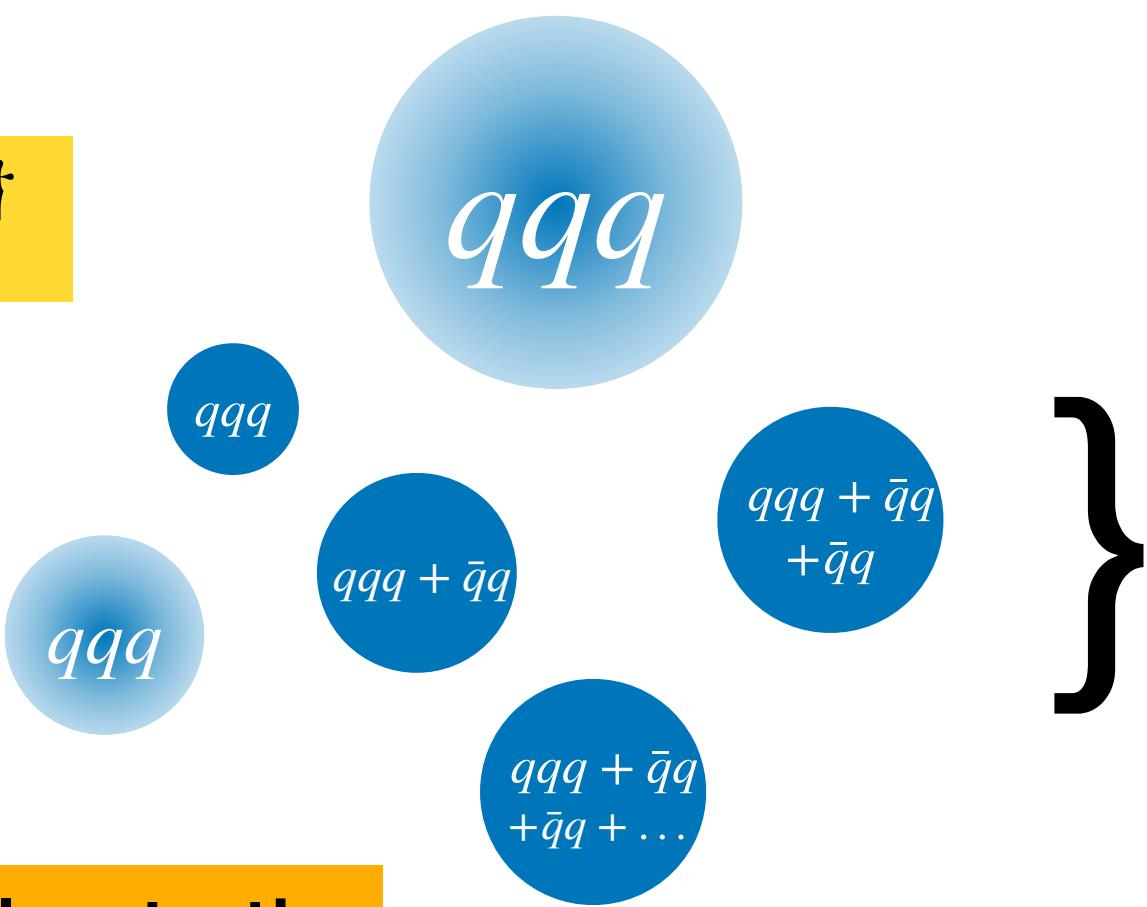
$$\mathbf{C}(t) v^{(n)} = \lambda^{(n)}(t) \mathbf{C}(t_0) v^{(n)}$$

\uparrow

$$= e^{-E_n t}$$

generalized eigenvalue problem (GEVP)

Q2: How to make sure we are not missing some of the E'_i 's ?



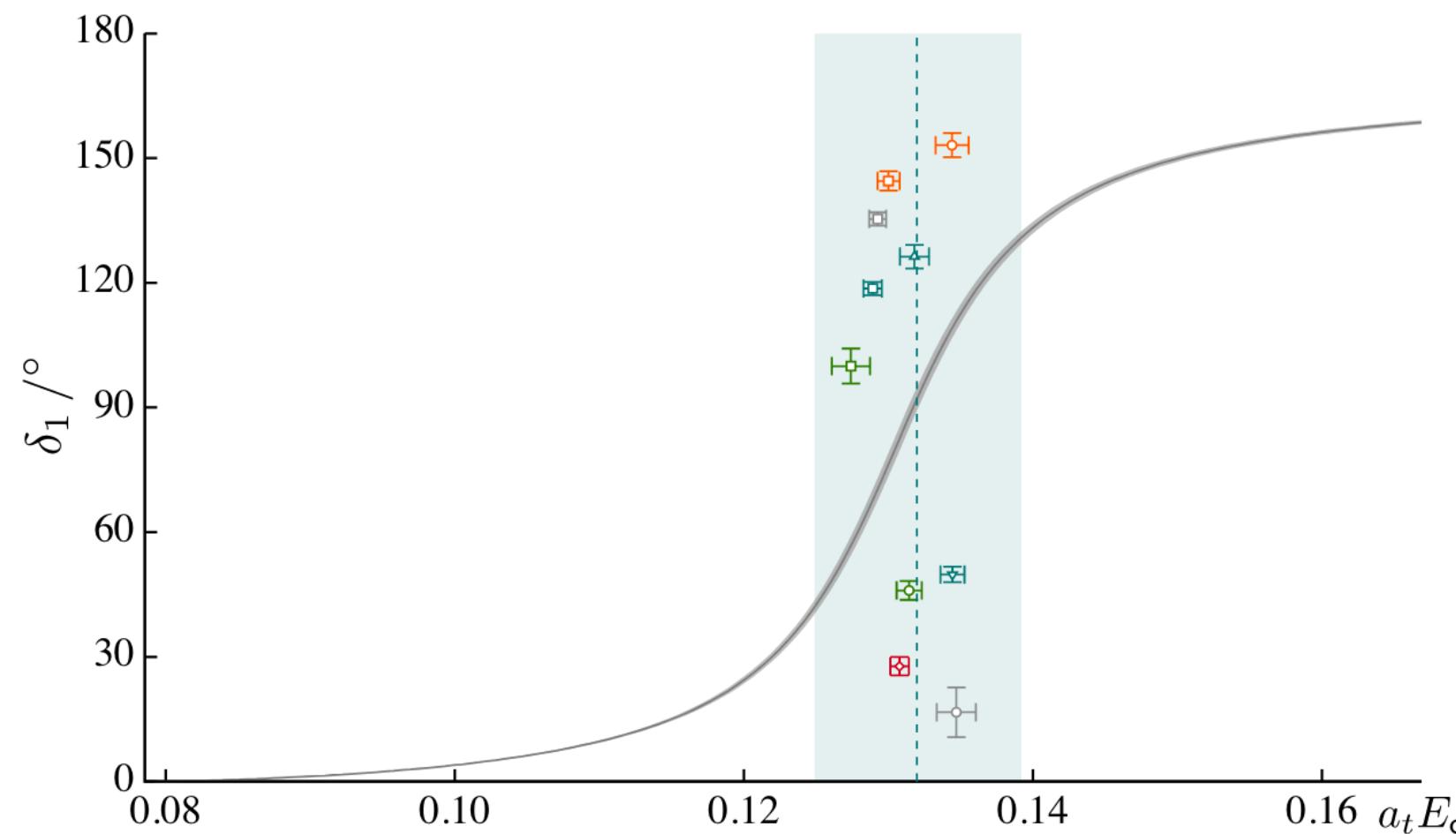
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$$\sum_{n=0}^N W_{ij,n} e^{-E_n t} \quad \chi_i^\dagger \in \left\{ \begin{array}{l} X_n^\dagger = \sum_i v_i^{(n)} \chi_i^\dagger \\ \mathbf{C}(t) v^{(n)} = \lambda^{(n)}(t) \mathbf{C}(t_0) v^{(n)} \end{array} \right.$$

with the maximum overlap to the states n 's

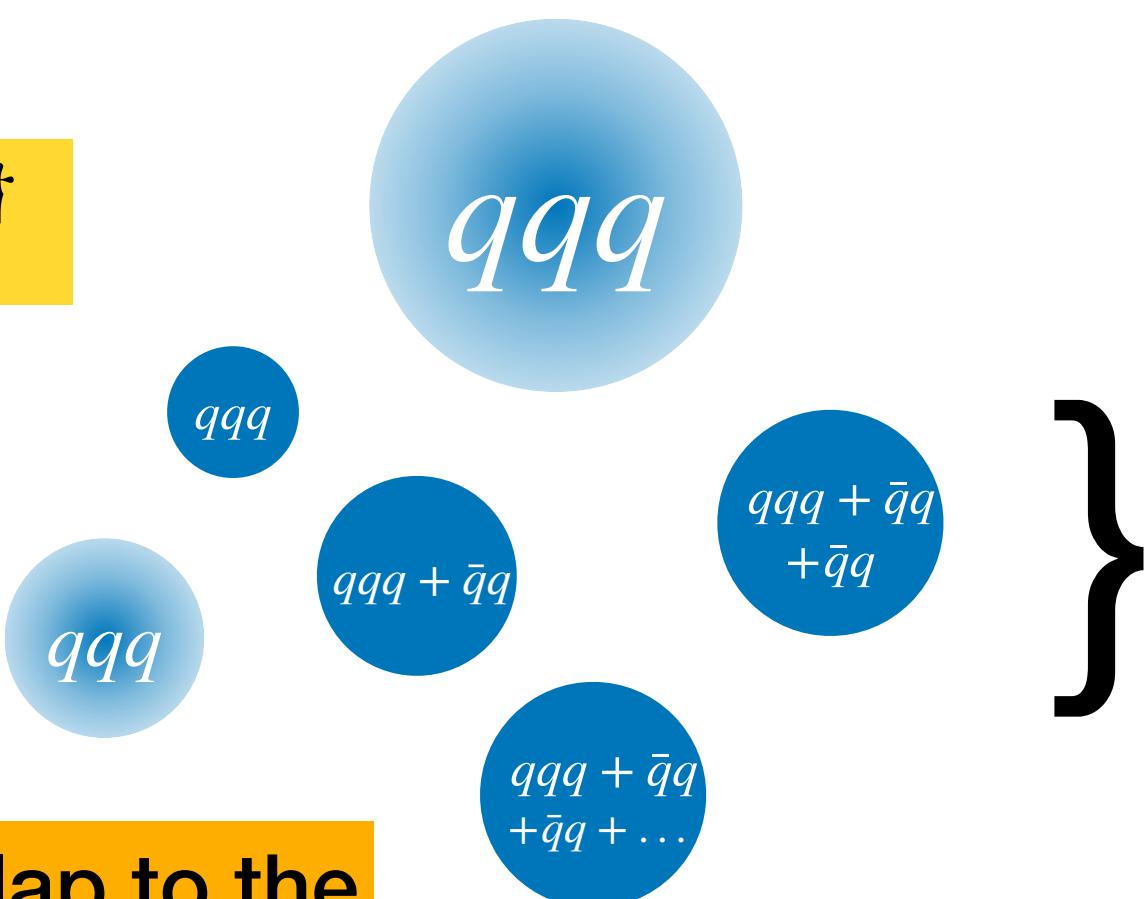
generalized eigenvalue problem (GEVP)

$$= e^{-E_n t}$$

M. Luscher, Nucl. Phys. B, 354, (1991)

Q2: How to make sure we are not missing some of the E_i 's ?

Wilson et al., Phys. Rev. D, vol. 92, no. 9, p. 094502, Nov. 2015



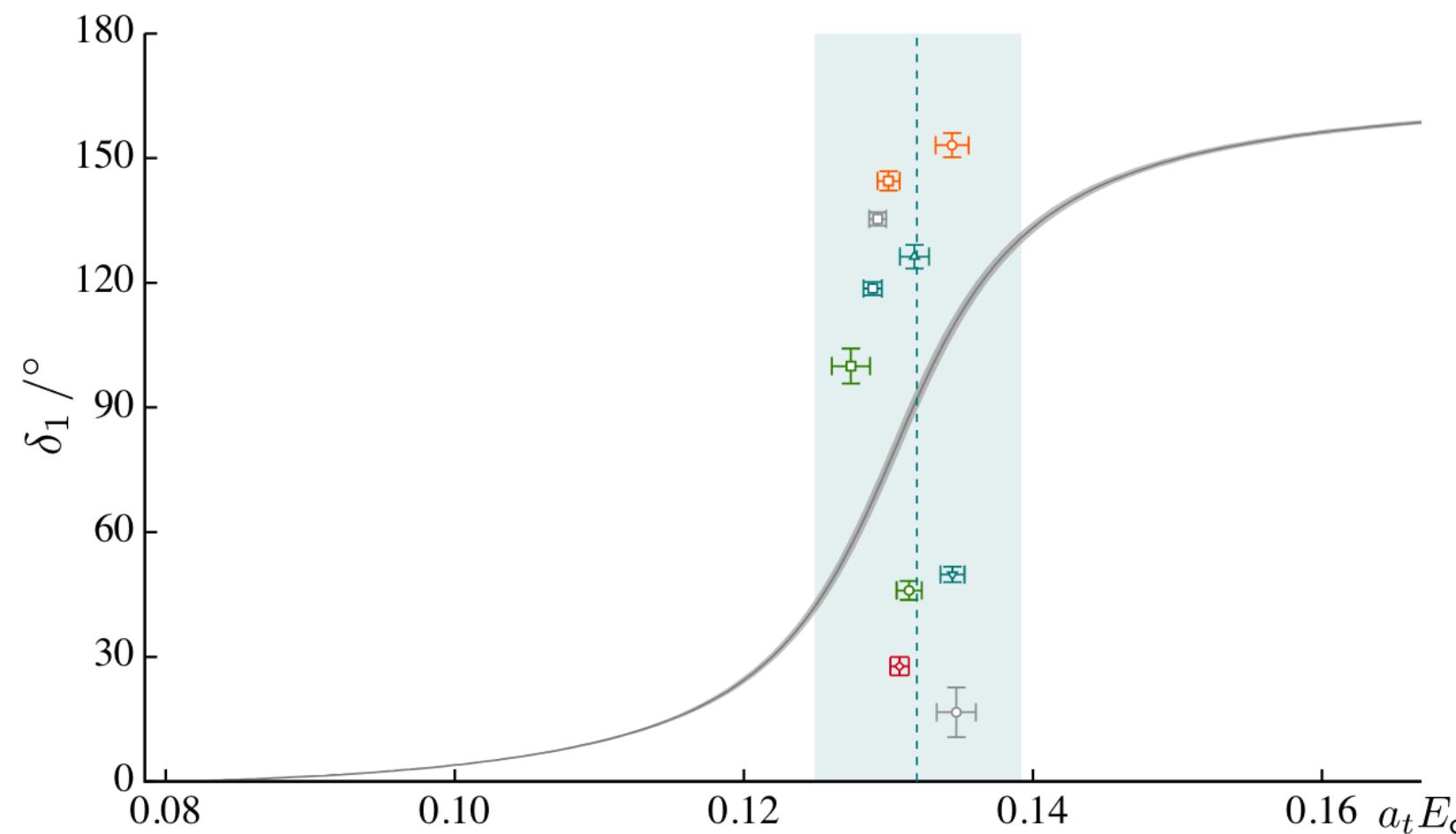
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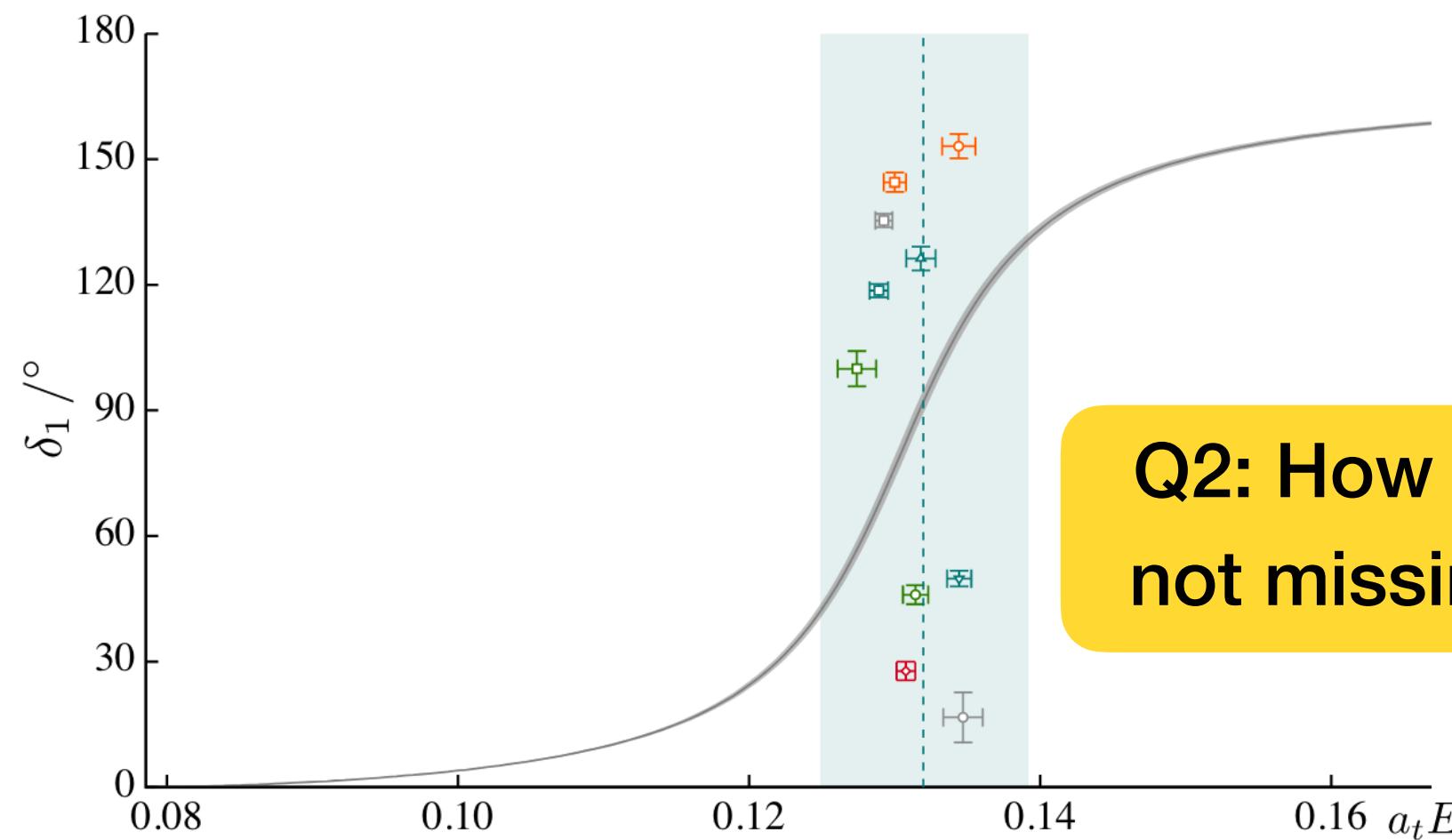
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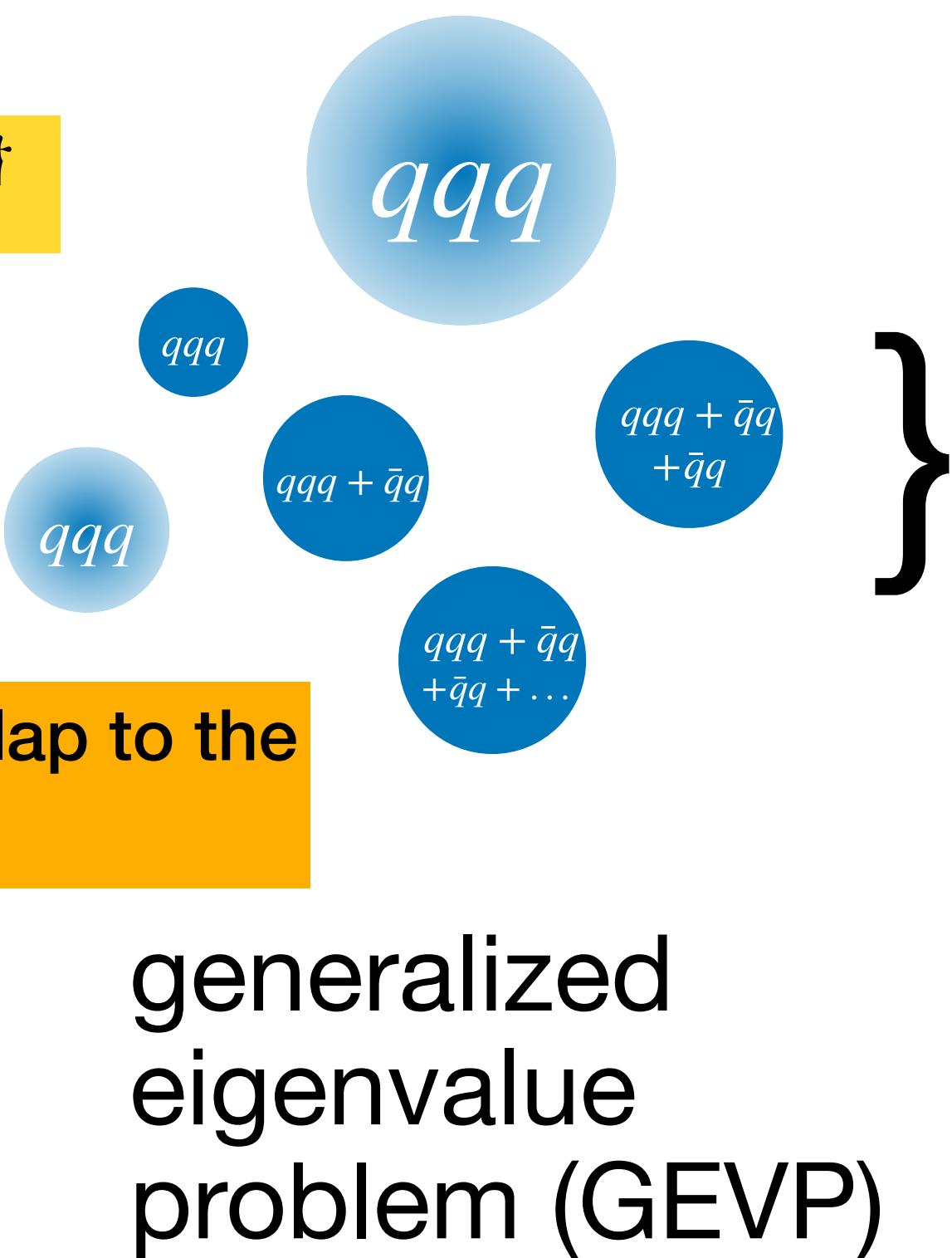
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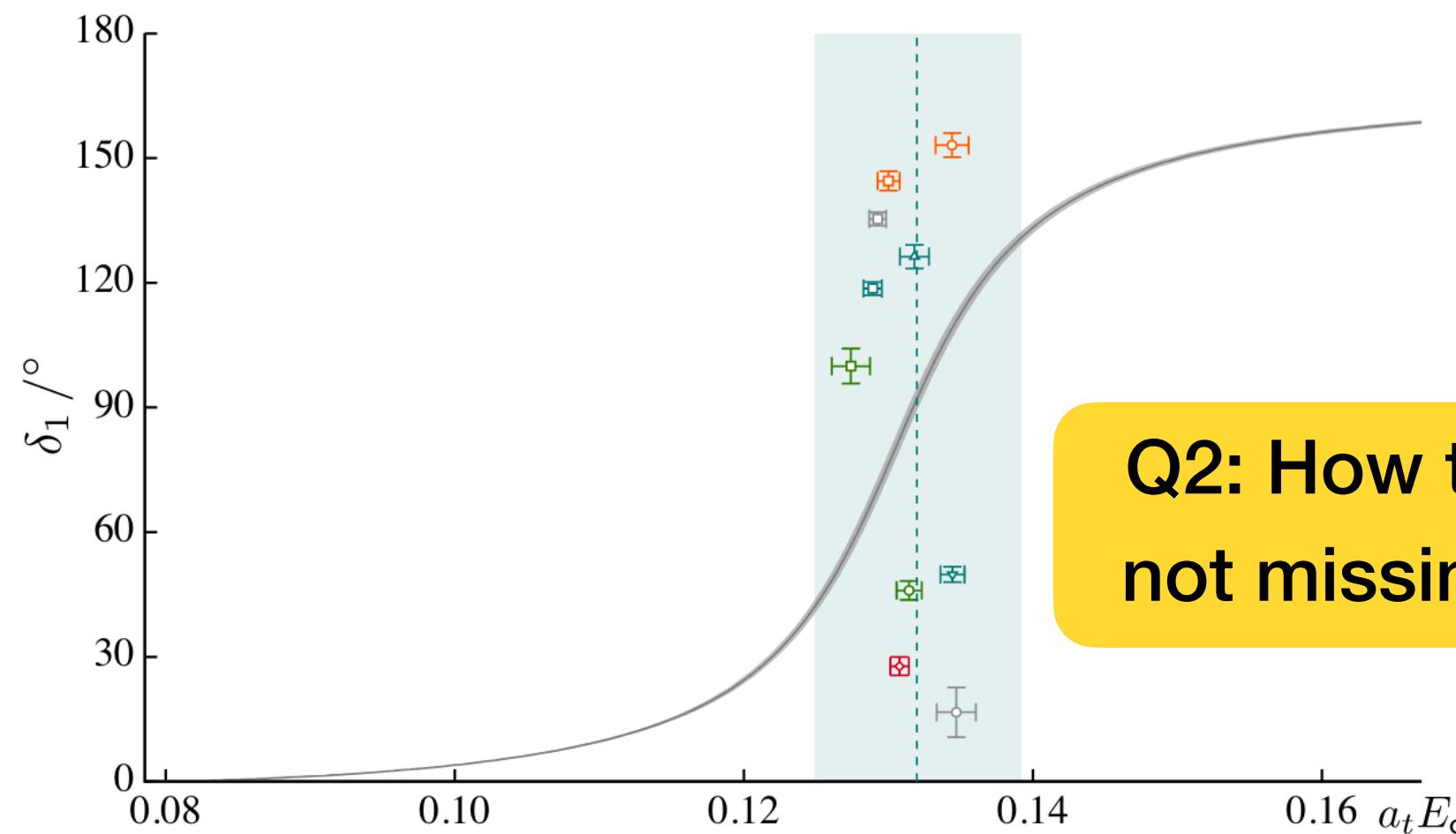
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Wilson et al., Phys. Rev. D, vol. 92, no. 9, p. 094502, Nov. 2015

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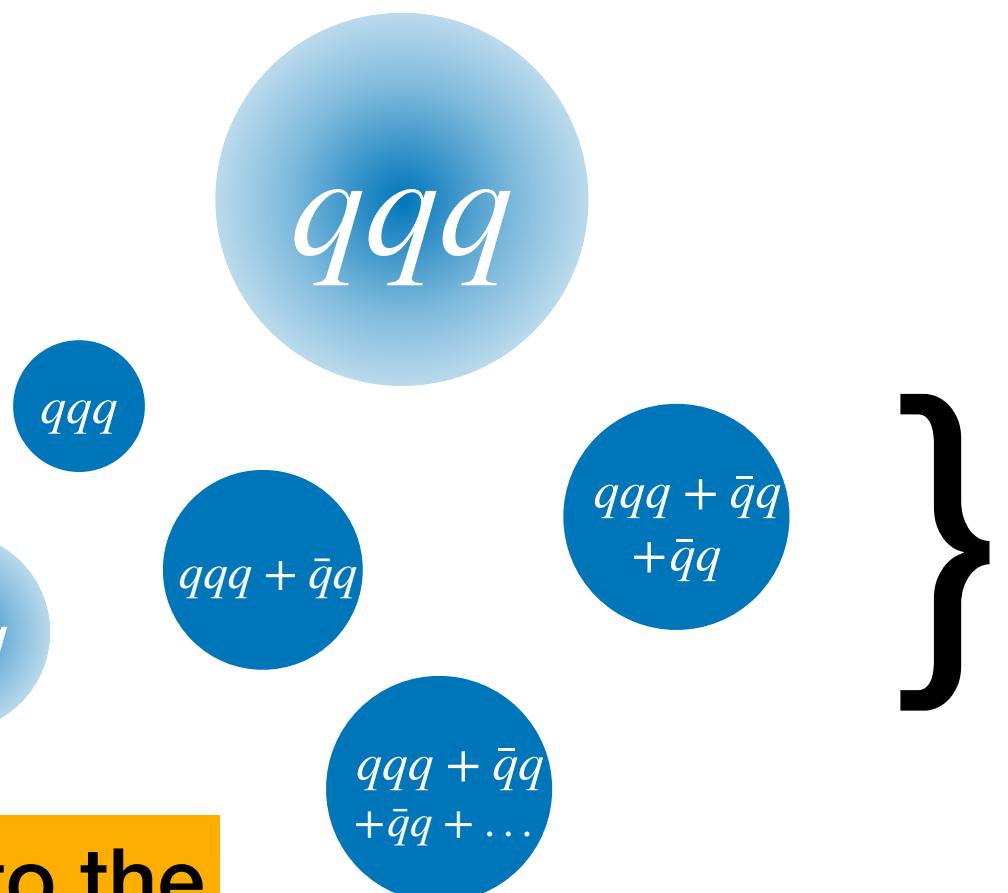
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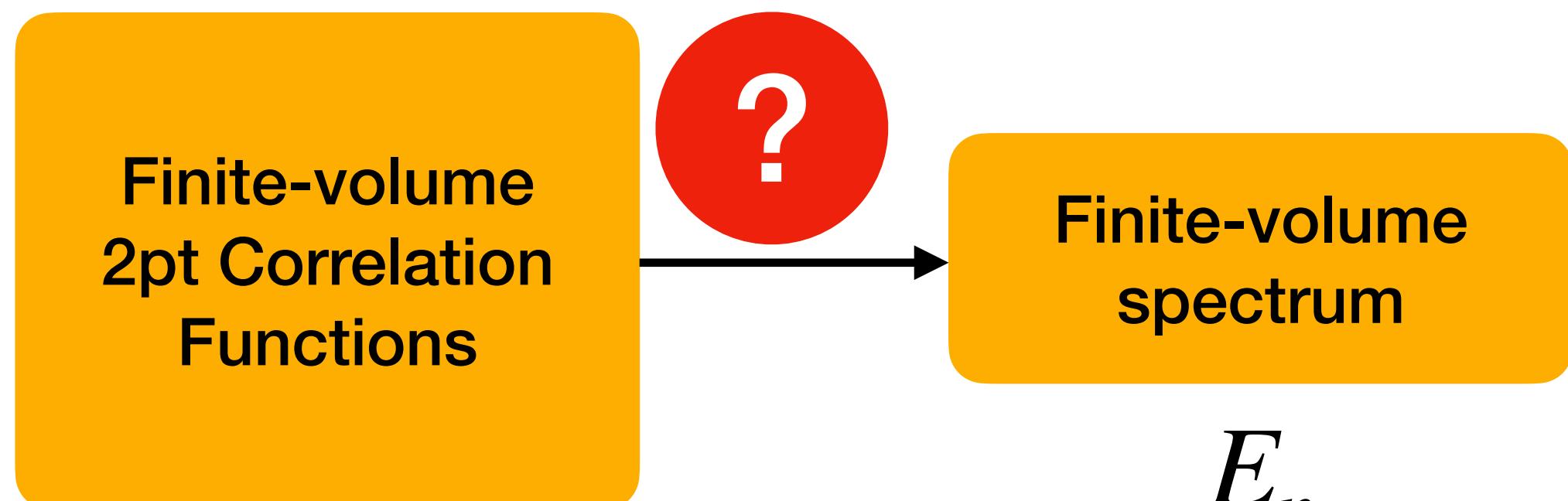
More operators!
 "distillation"

M. Peardon et al. (Hadron Spectrum),
 Phys. Rev. D80, 054506 (2009)

Cost for GEVP is high!!!

Finite-volume spectrum from lattice QCD

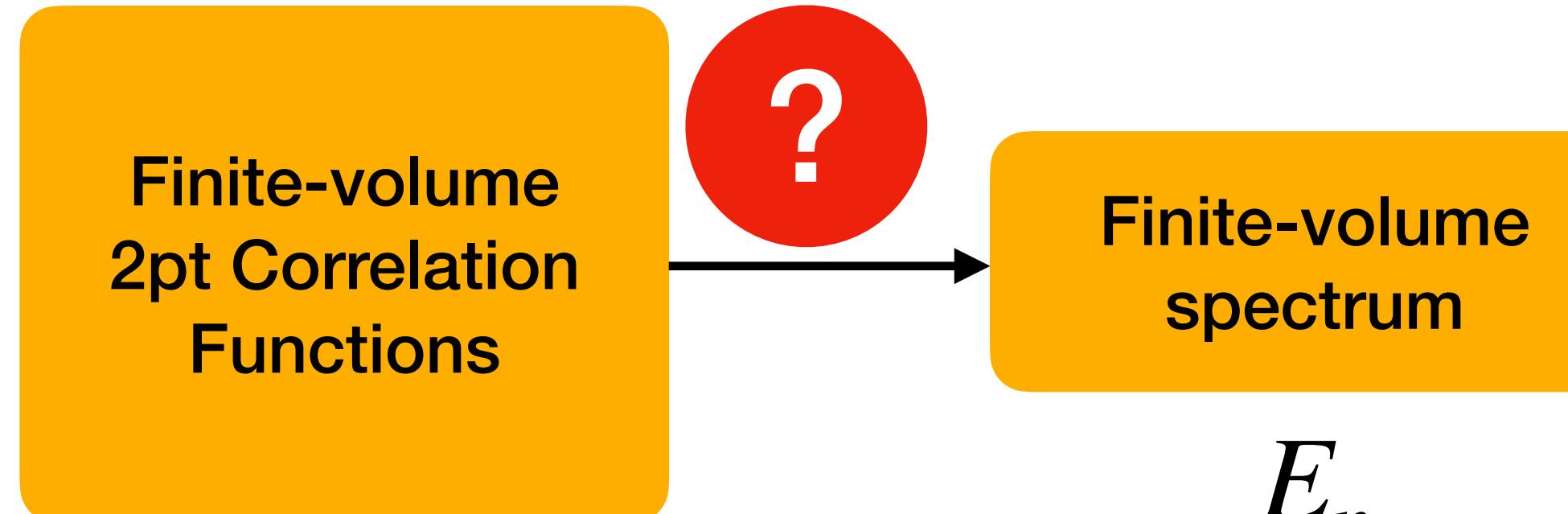
Q2*: Do we really need to make sure we are not missing (resolve) some of the E'_i 's ?



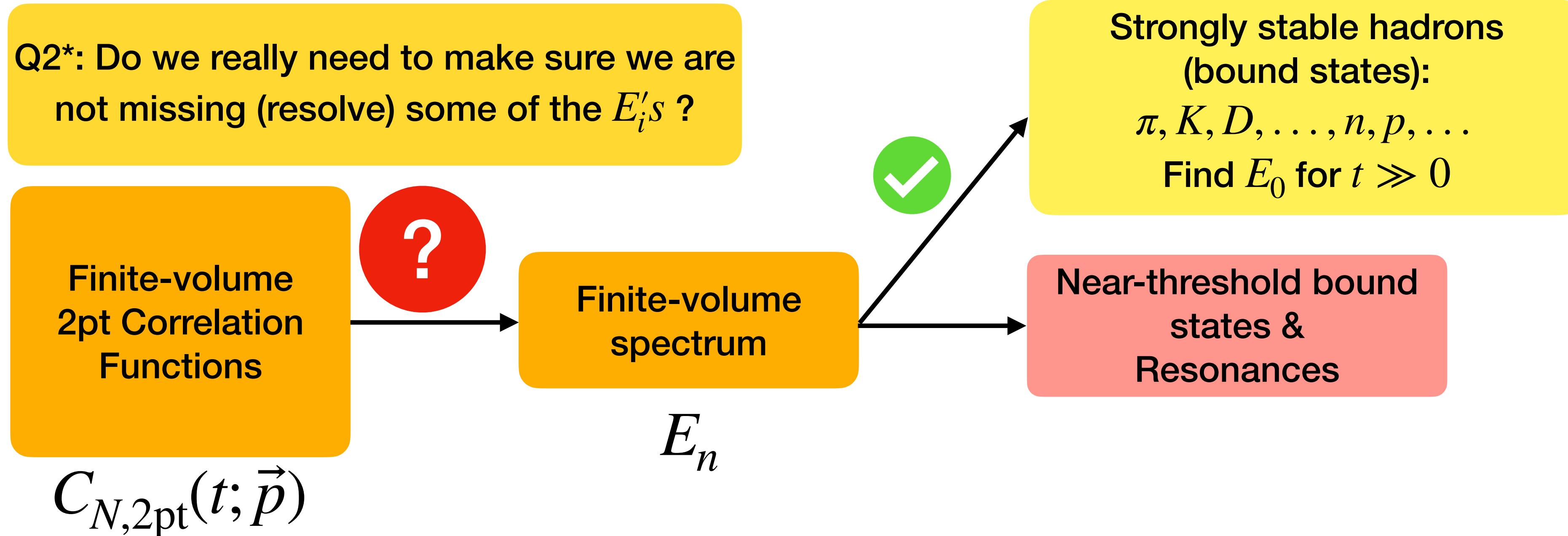
$$C_{N,2\text{pt}}(t; \vec{p})$$
$$E_n$$

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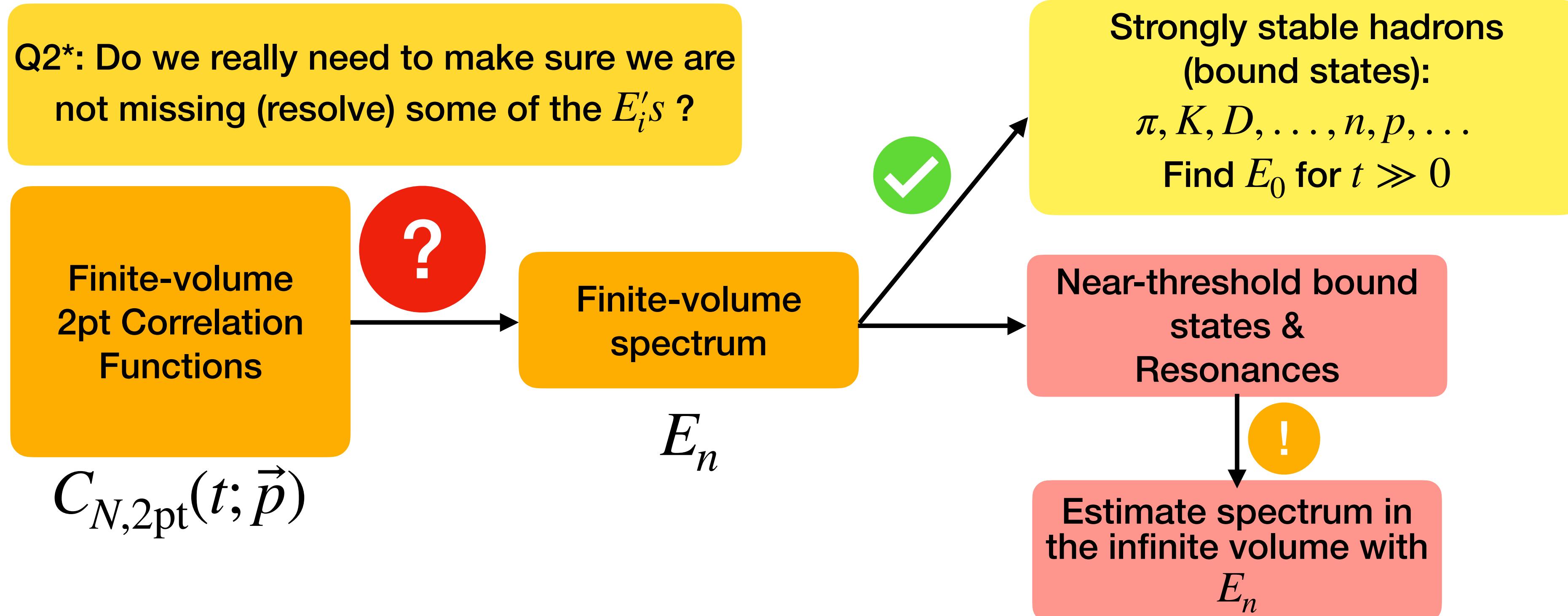
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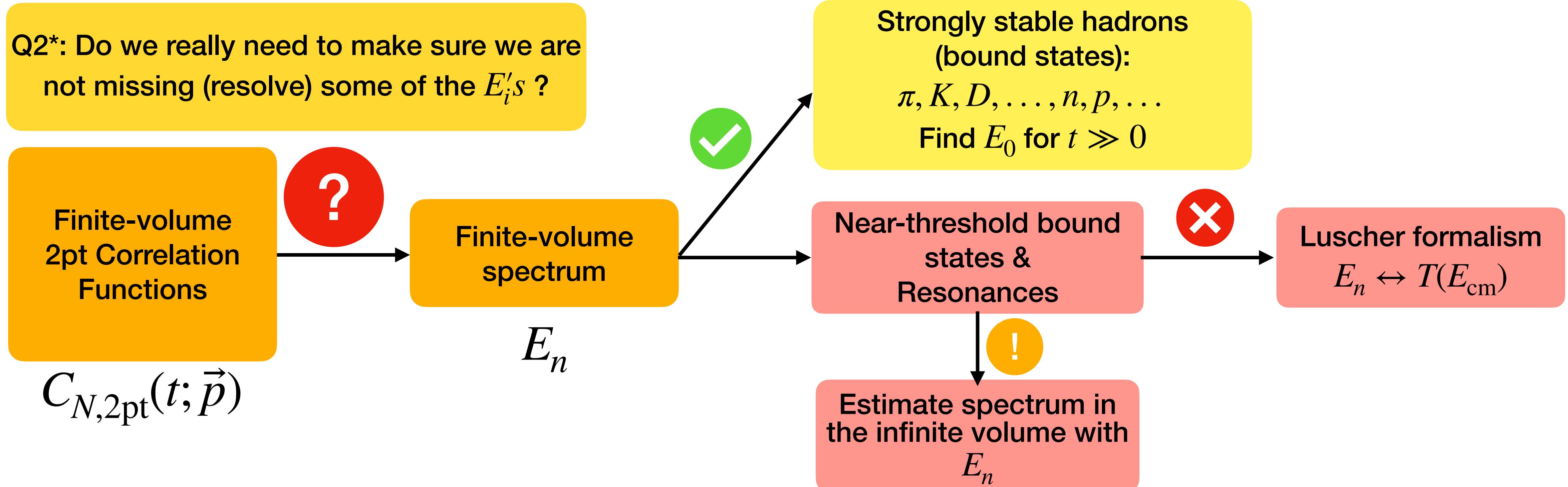
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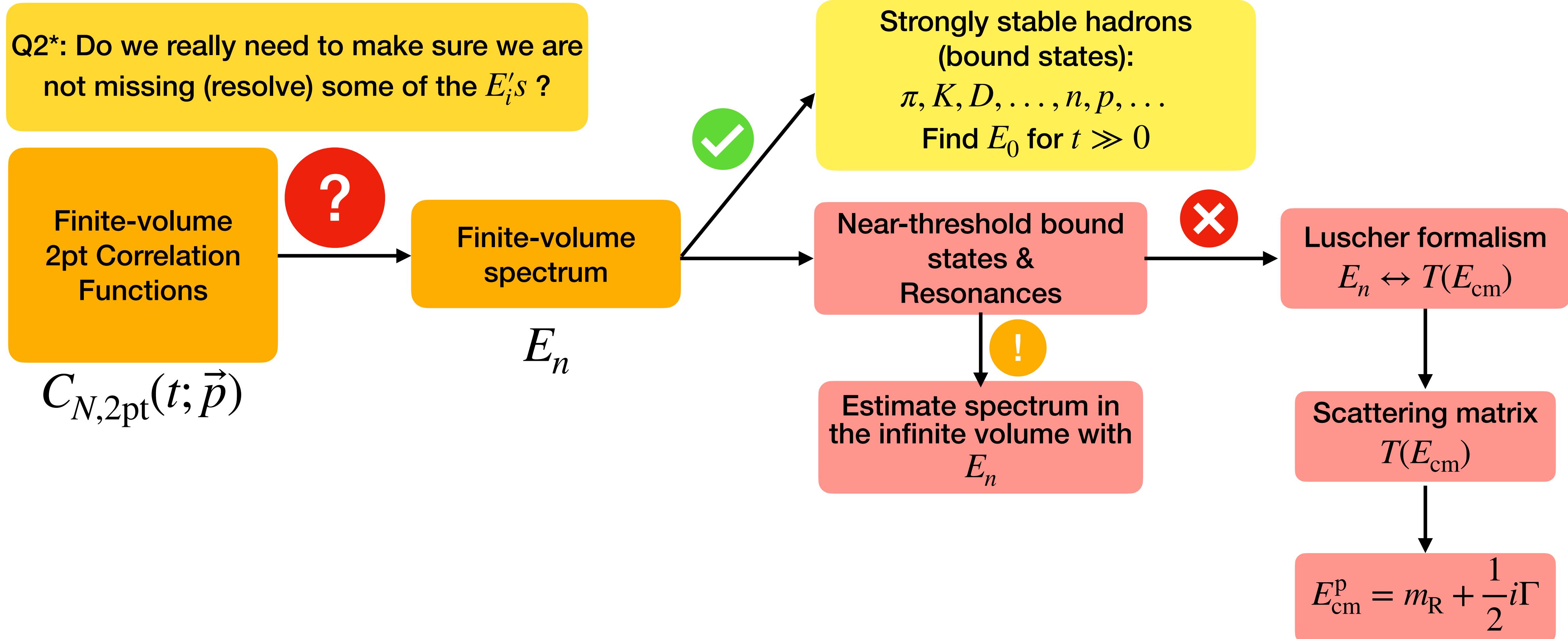
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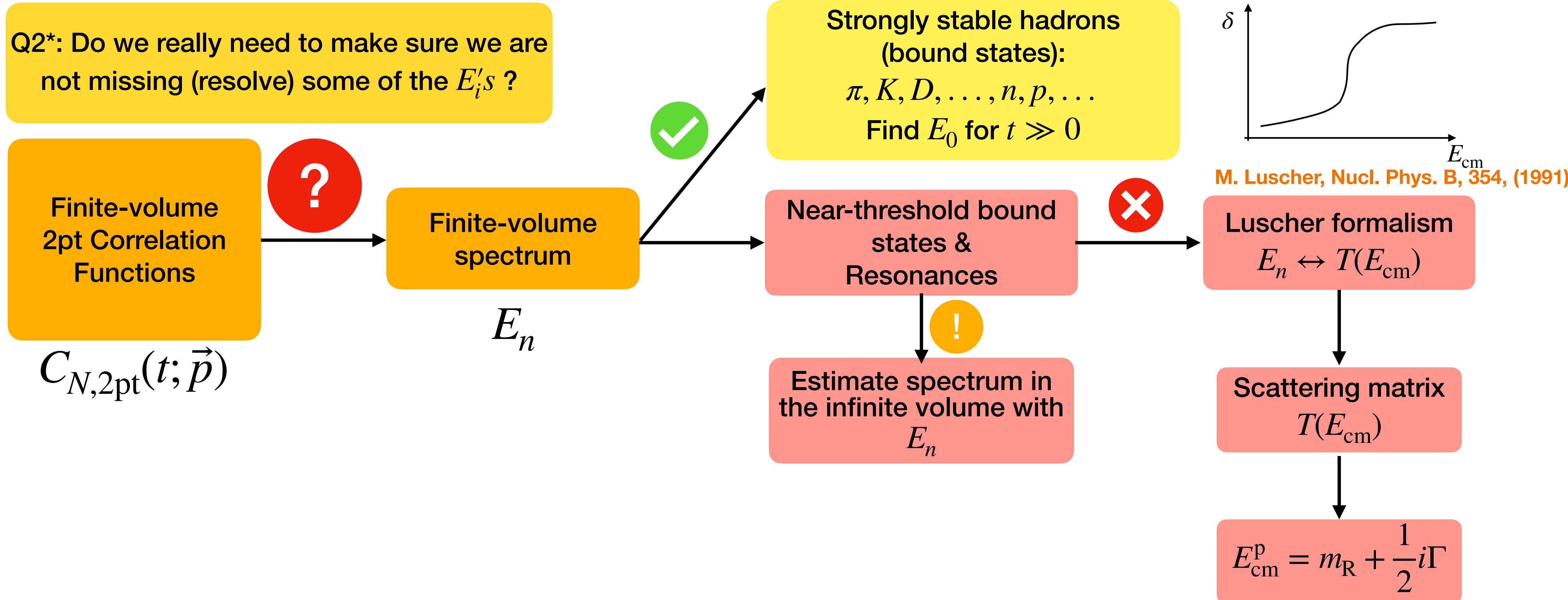
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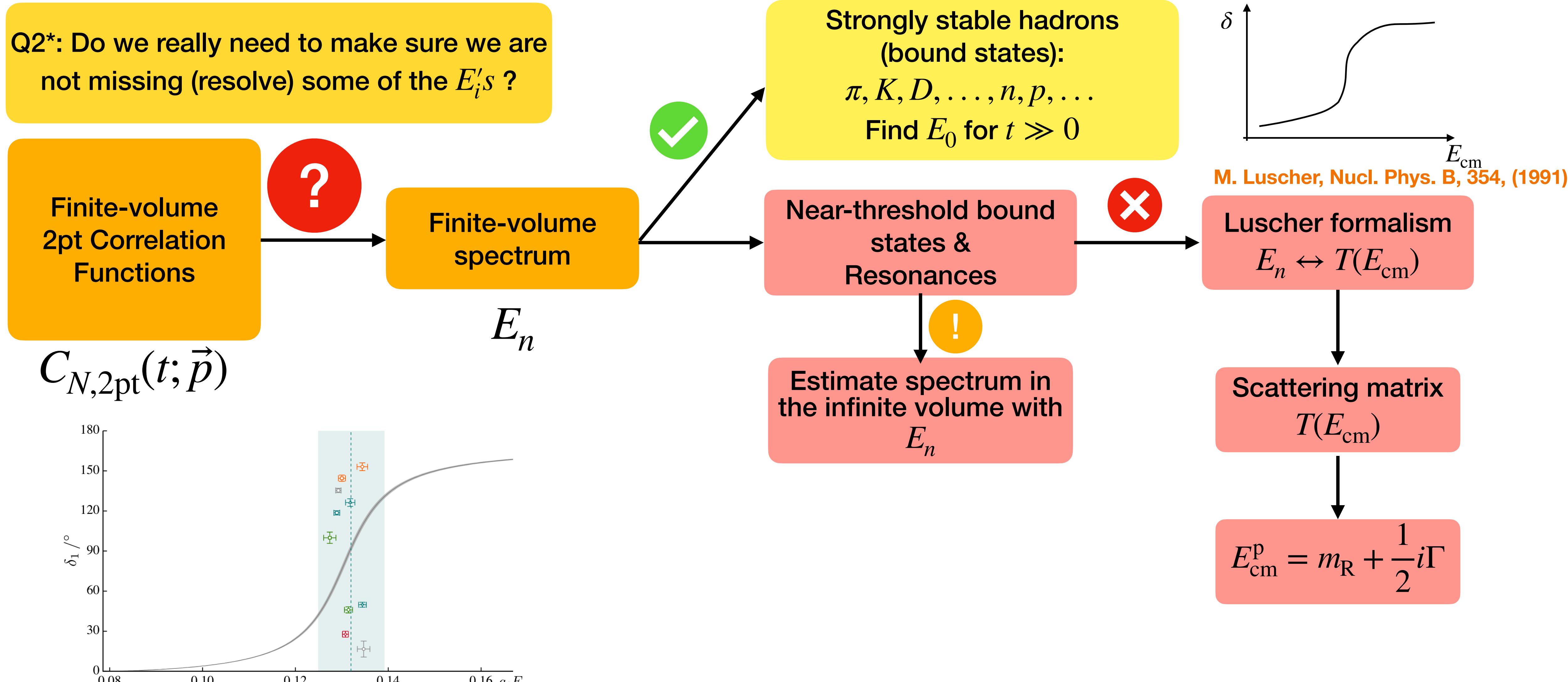
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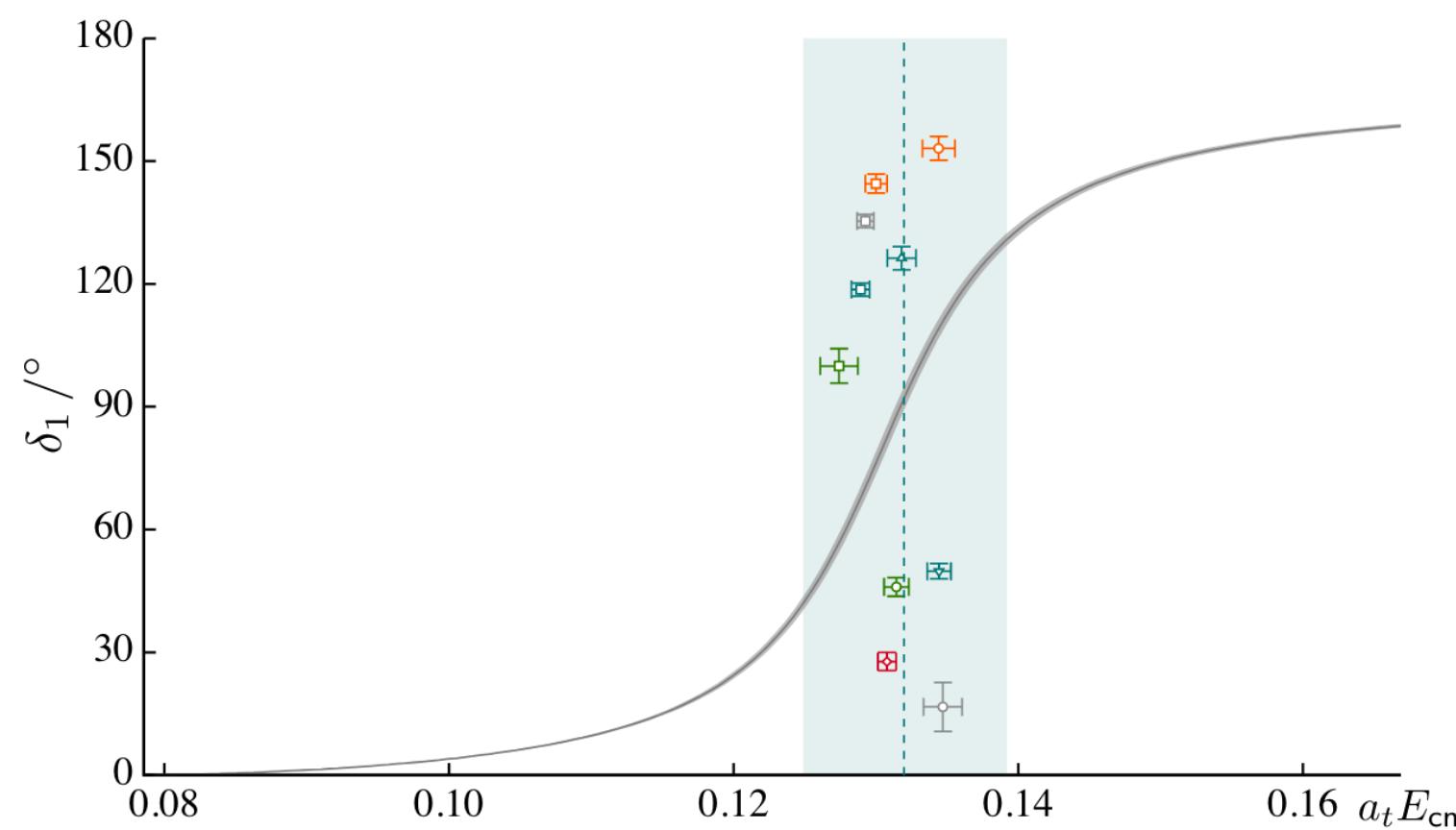
Wilson et al., Phys. Rev. D, vol. 92, no. 9, p. 094502, Nov. 2015

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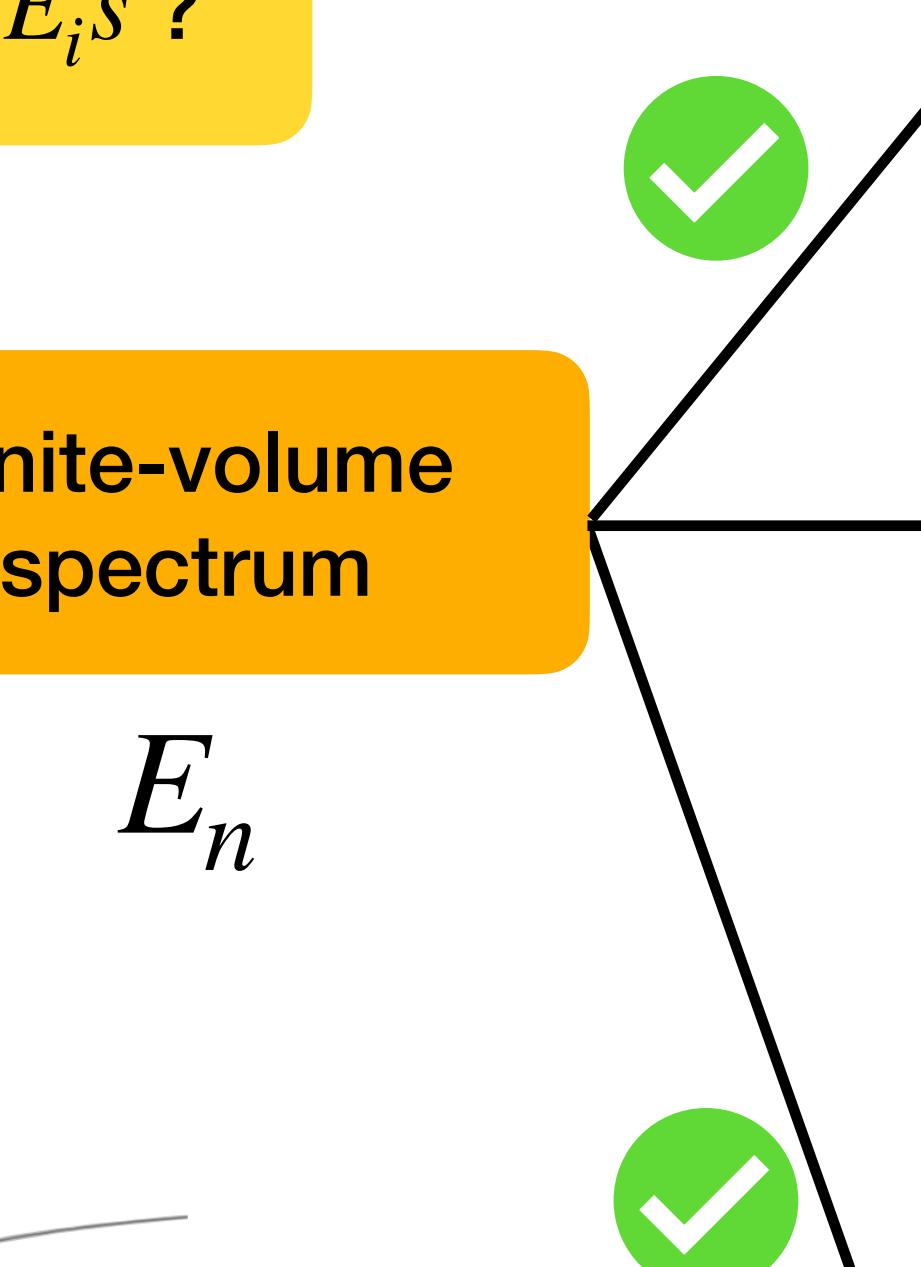
Finite-volume
2pt Correlation
Functions

$$C_{N,2\text{pt}}(t; \vec{p})$$



Wilson et al., Phys. Rev. D, vol. 92, no. 9, p. 094502, Nov. 2015

$$E_n$$

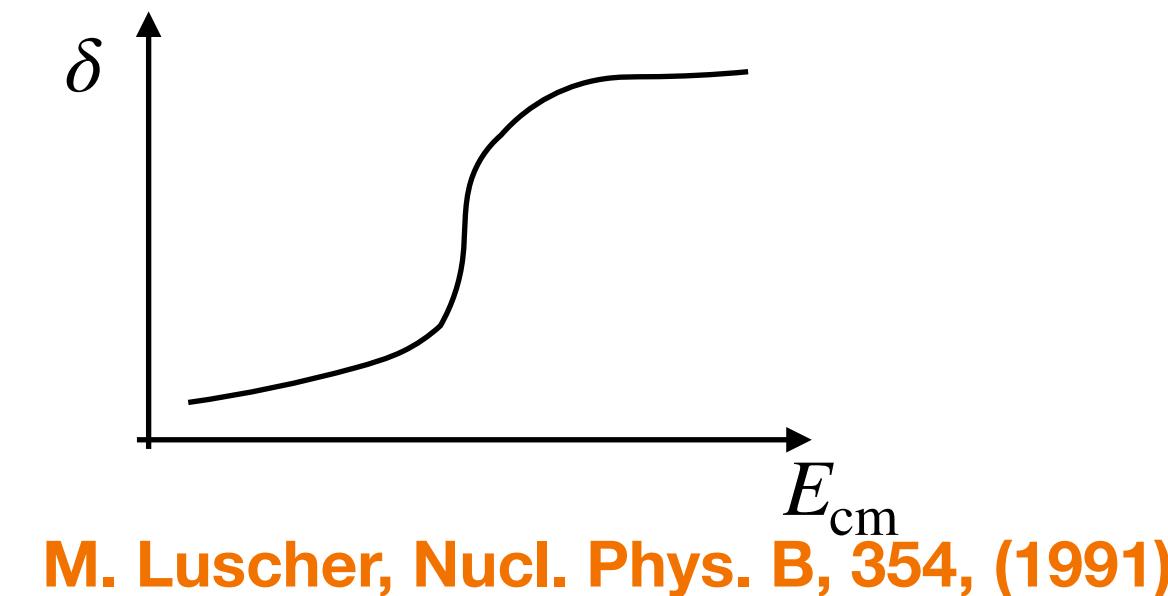


Strongly stable hadrons
(bound states):
 $\pi, K, D, \dots, n, p, \dots$
Find E_0 for $t \gg 0$

Near-threshold bound
states &
Resonances

Estimate spectrum in
the infinite volume with
 E_n

Constrain 3pt/4pt
correlation functions to
extract form factors, etc.

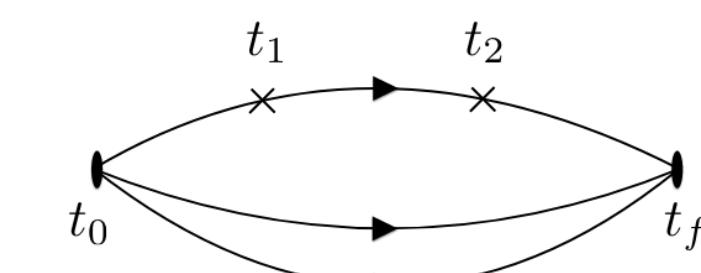


M. Luscher, Nucl. Phys. B, 354, (1991)

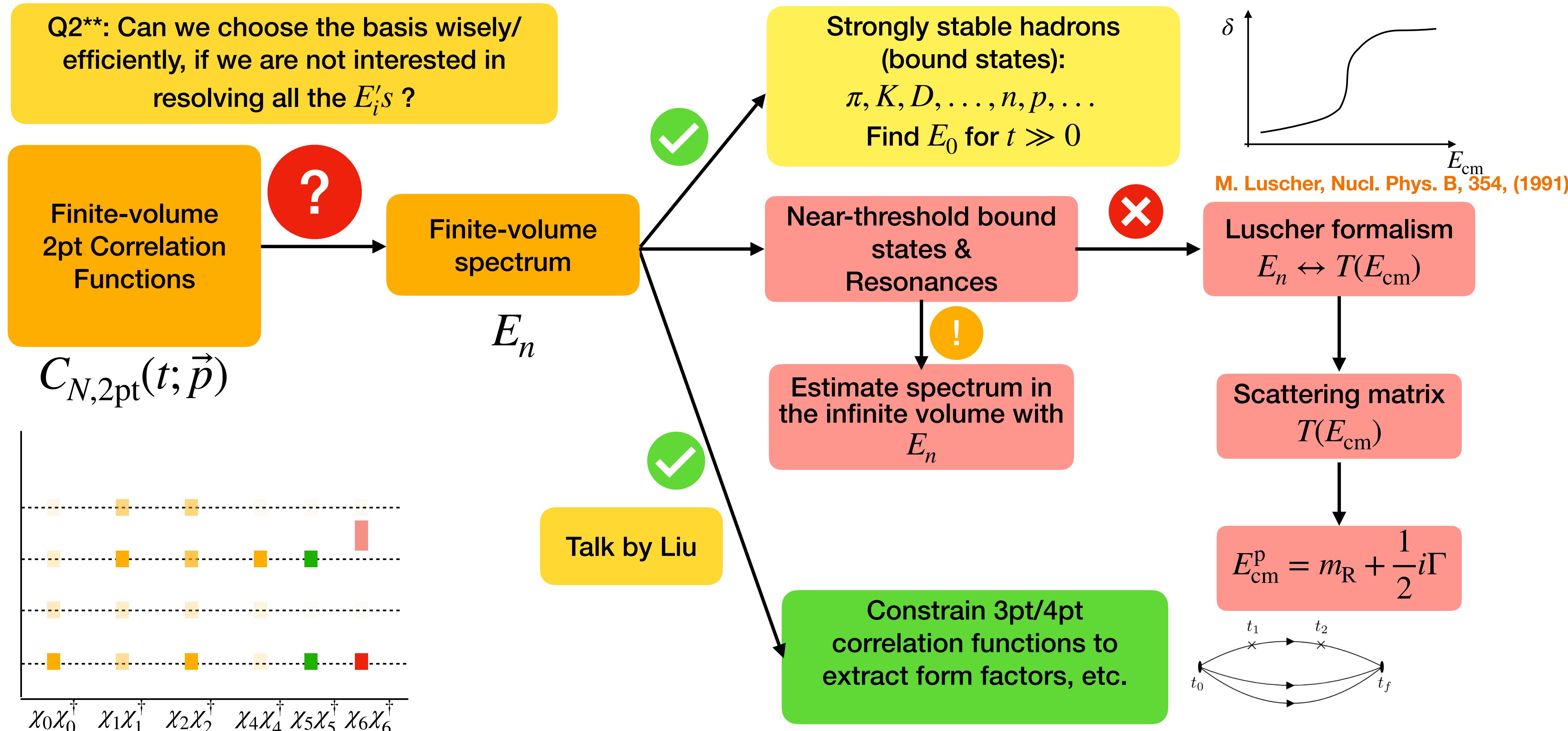
Luscher formalism
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Scattering matrix
 $T(E_{\text{cm}})$

$$E_{\text{cm}}^{\text{p}} = m_R + \frac{1}{2}i\Gamma$$



Finite-volume spectrum from lattice QCD



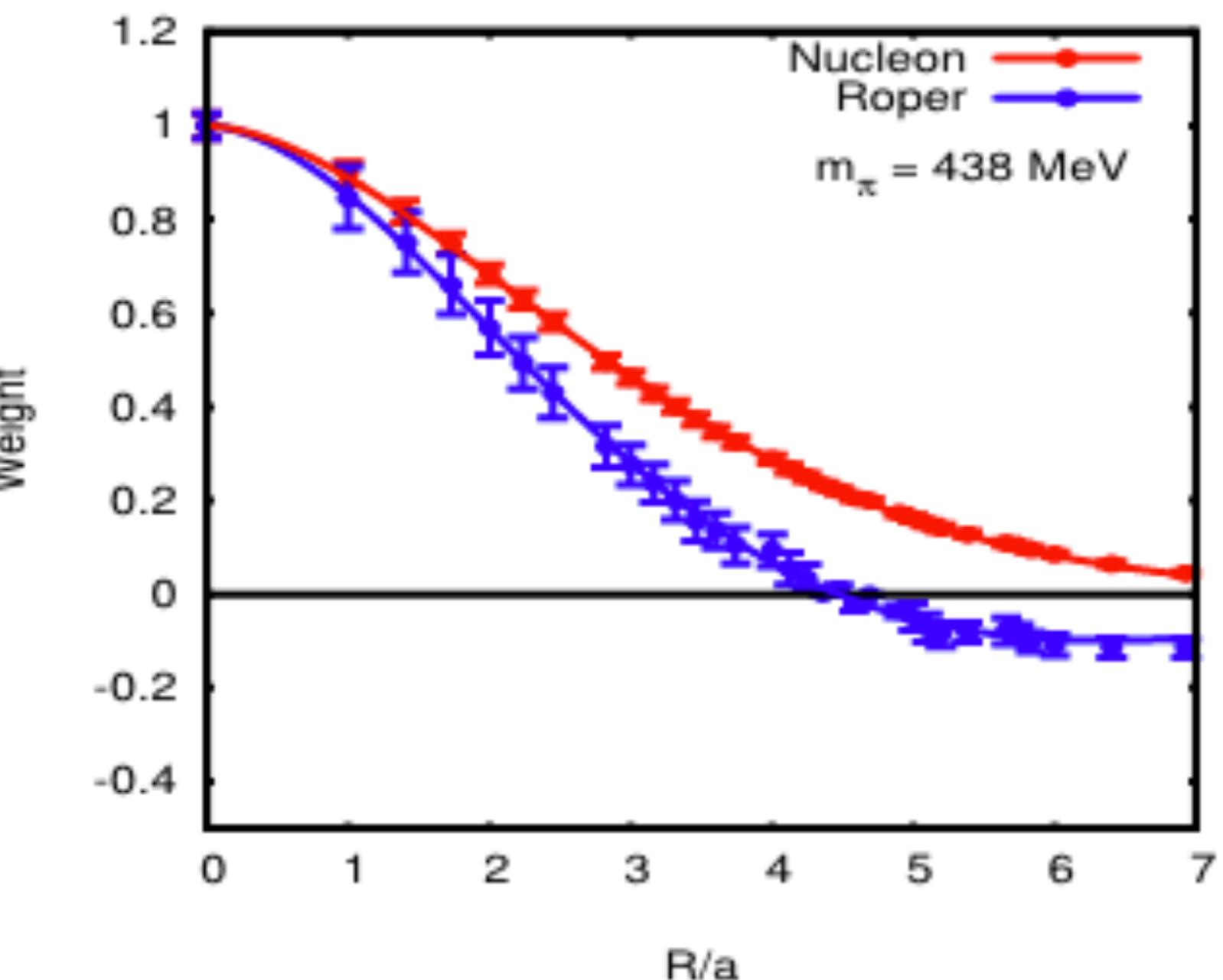
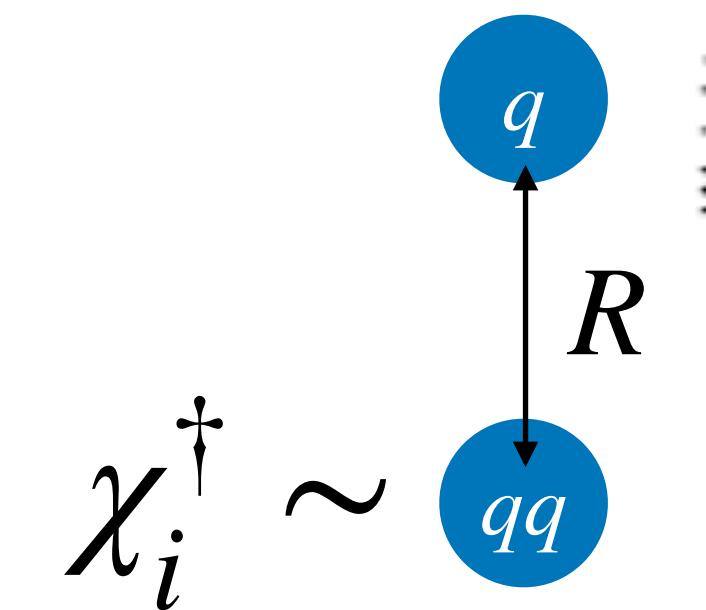
Example: Roper State from Overlap Fermions

$N(1440)$ $1/2^+$

$$I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$$

$\text{Re}(\text{pole position}) = 1360$ to 1380 (≈ 1370) MeV
 $-2\text{Im}(\text{pole position}) = 180$ to 205 (≈ 190) MeV
 Breit-Wigner mass = 1410 to 1470 (≈ 1440) MeV
 Breit-Wigner full width = 250 to 450 (≈ 350) MeV

$N(1440)$ DECAY MODES	Fraction (Γ_i/Γ)	p (MeV/c)
$N\pi$	55–75 %	398
$N\eta$	<1 %	†
$N\pi\pi$	17–50 %	347
$\Delta(1232)\pi$, P -wave	6–27 %	147
$N\sigma$	11–23 %	—
$p\gamma$, helicity=1/2	0.035–0.048 %	414
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Sensitive to the **size**
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Nucleon and Roper wavefunctions
in the Coulomb gauge

M. Sun et al., “Roper State from Overlap Fermions,” PhysRevD.101.054511

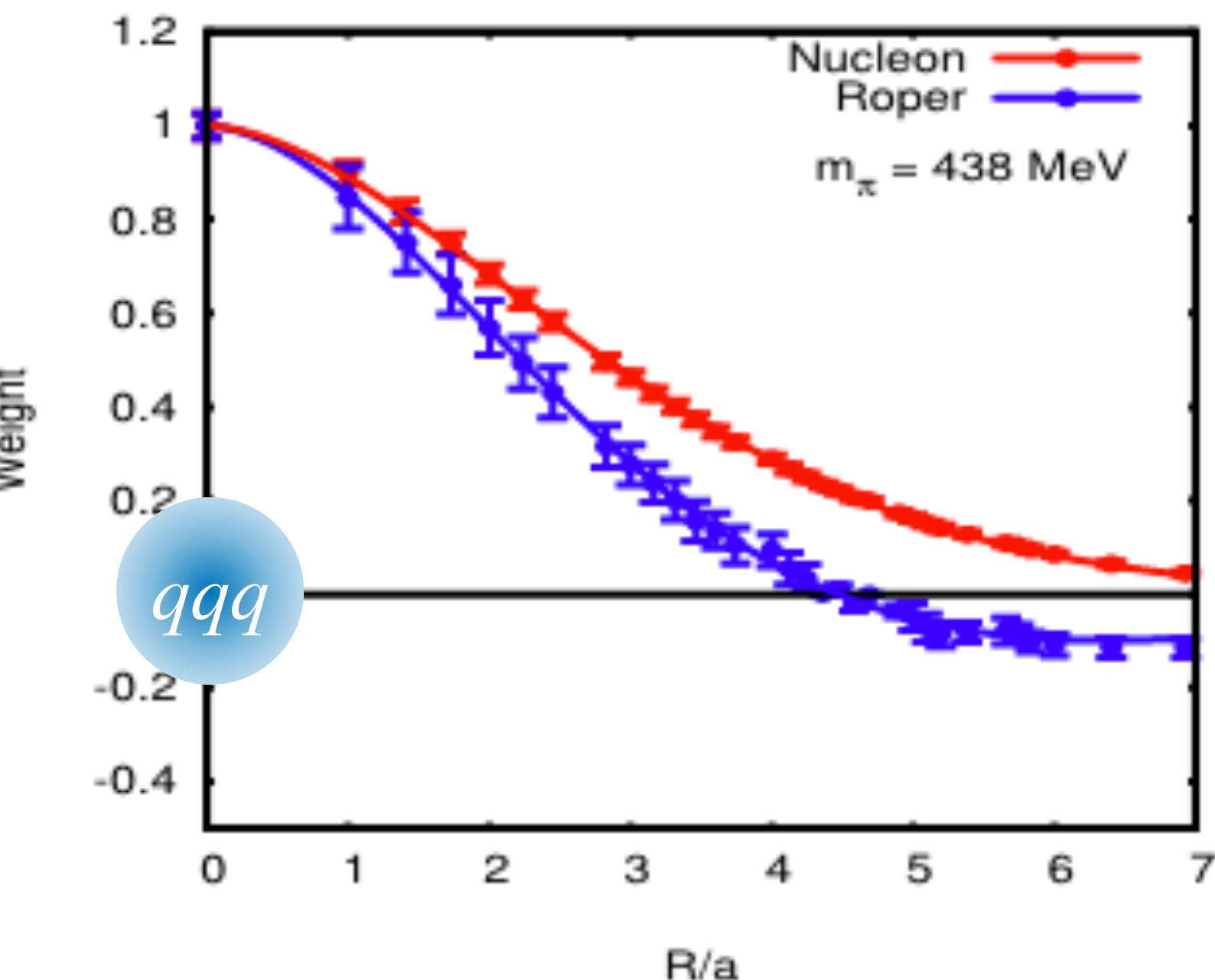
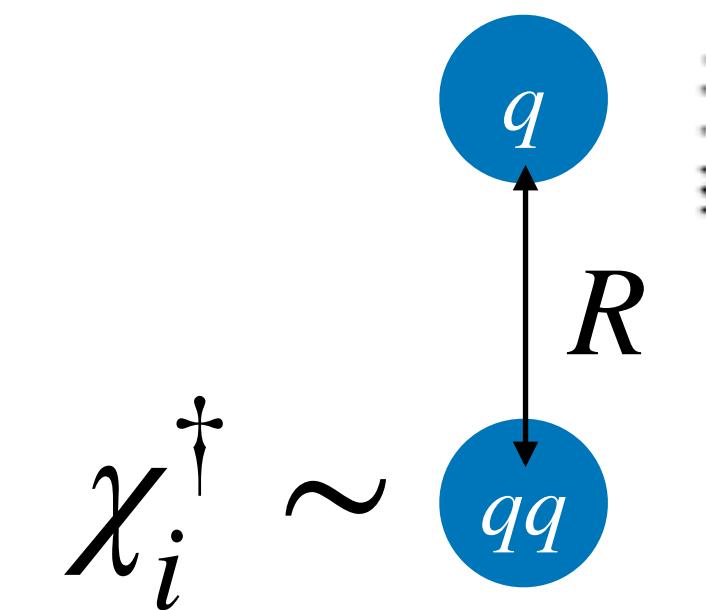
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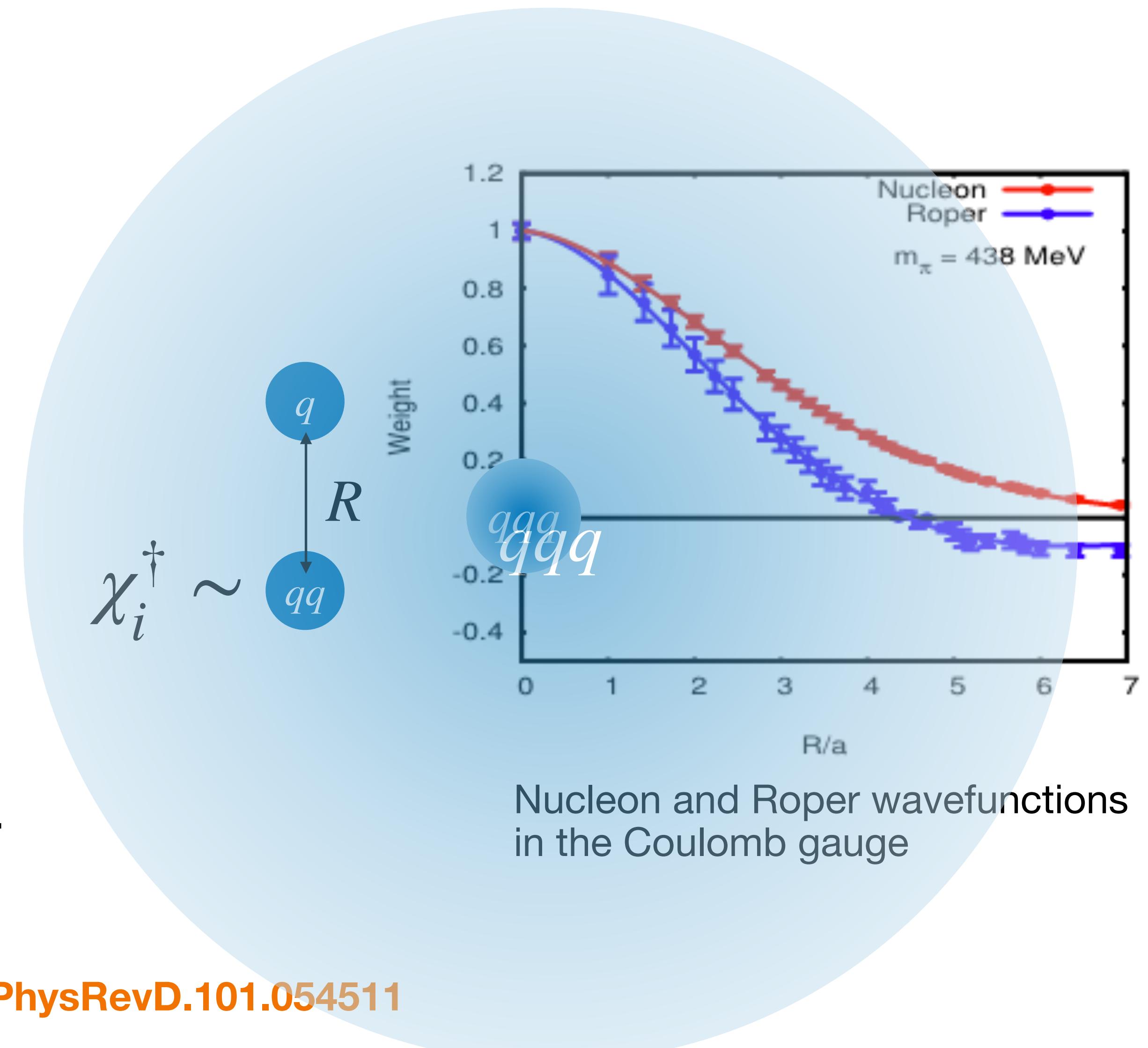
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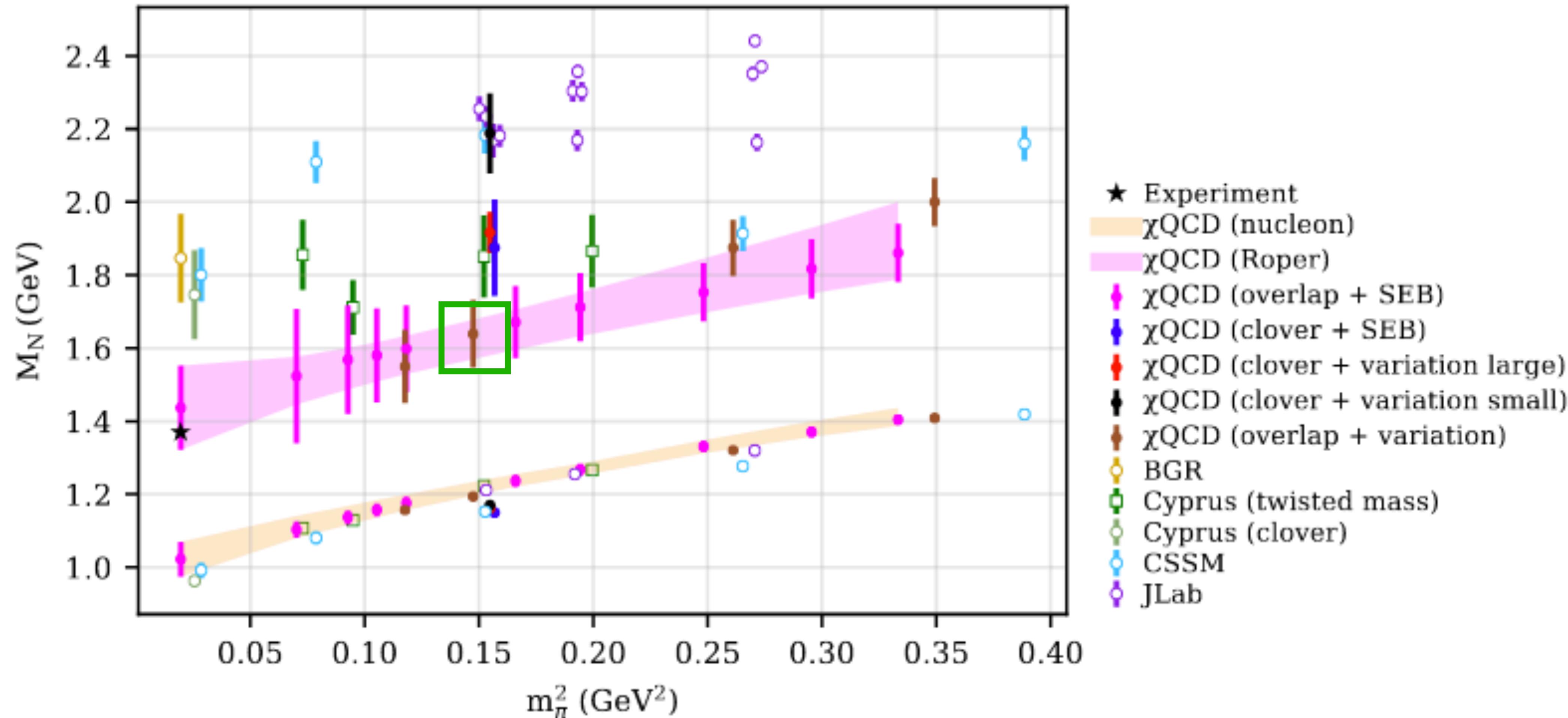
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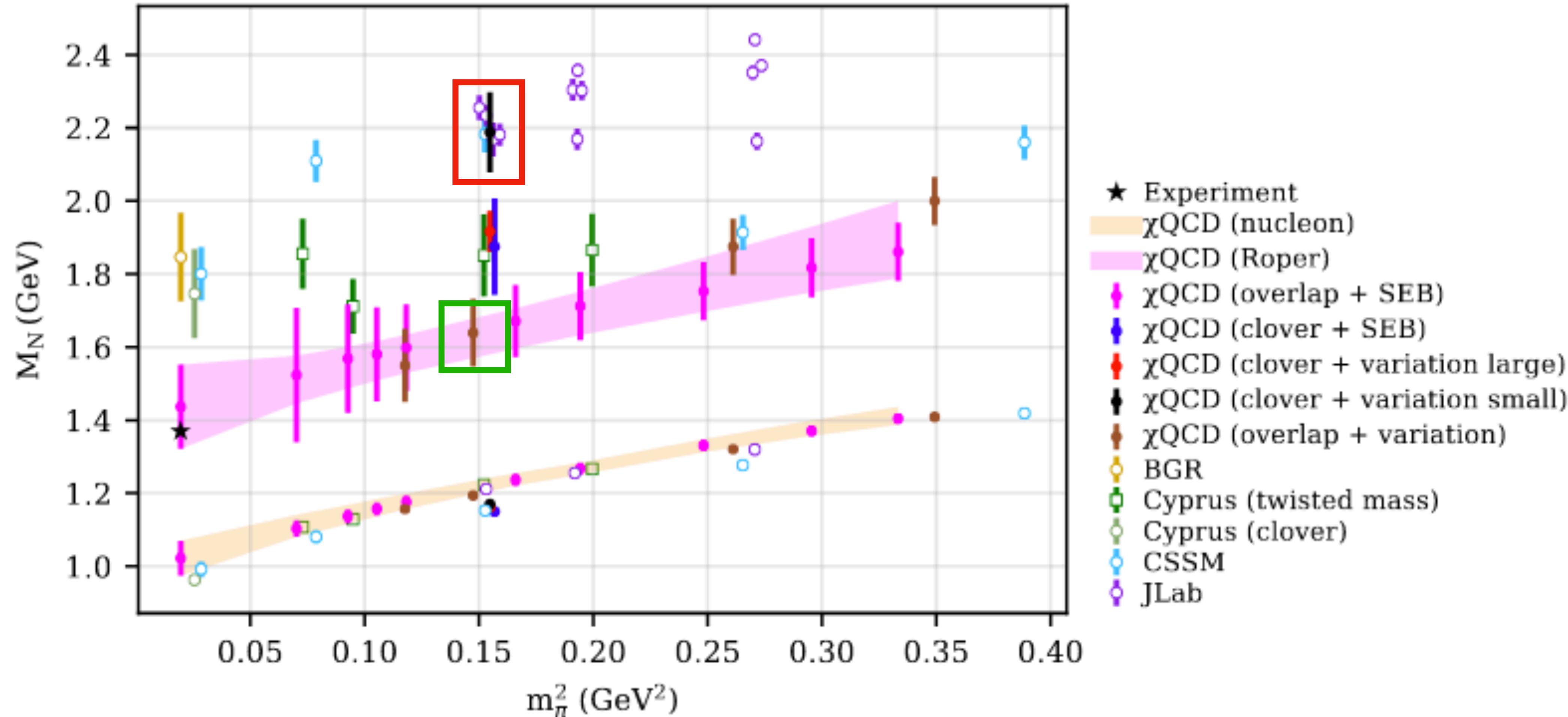
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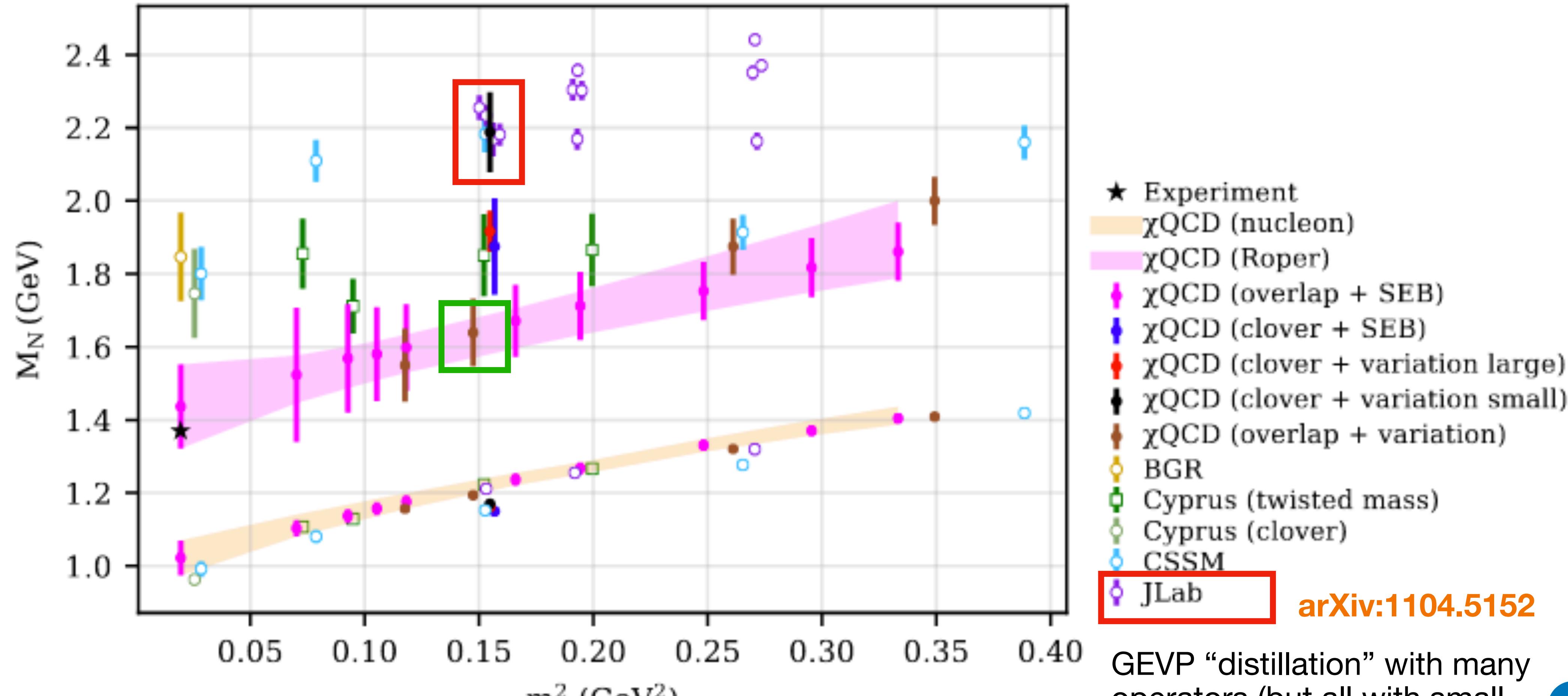
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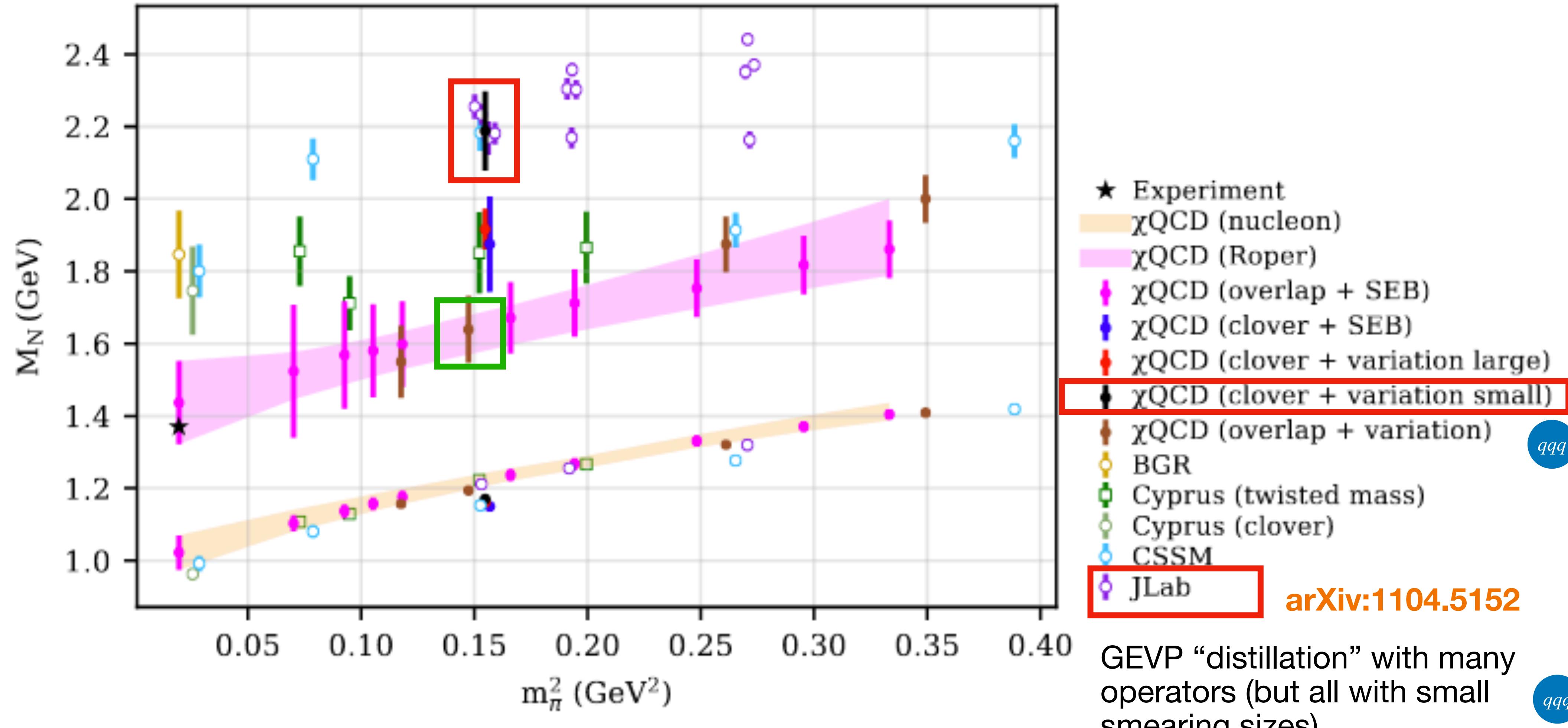


GEVP “distillation” with many operators (but all with small smearing sizes)



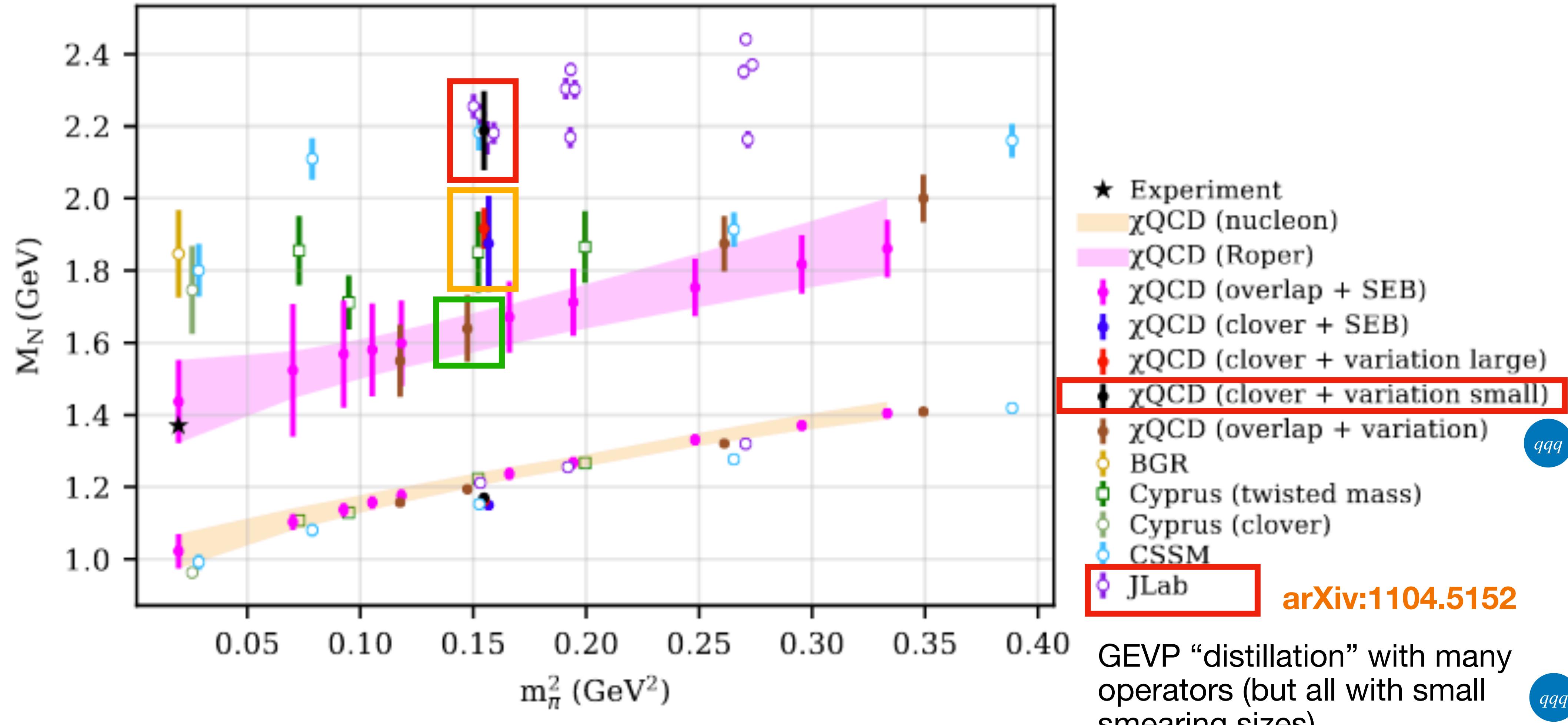
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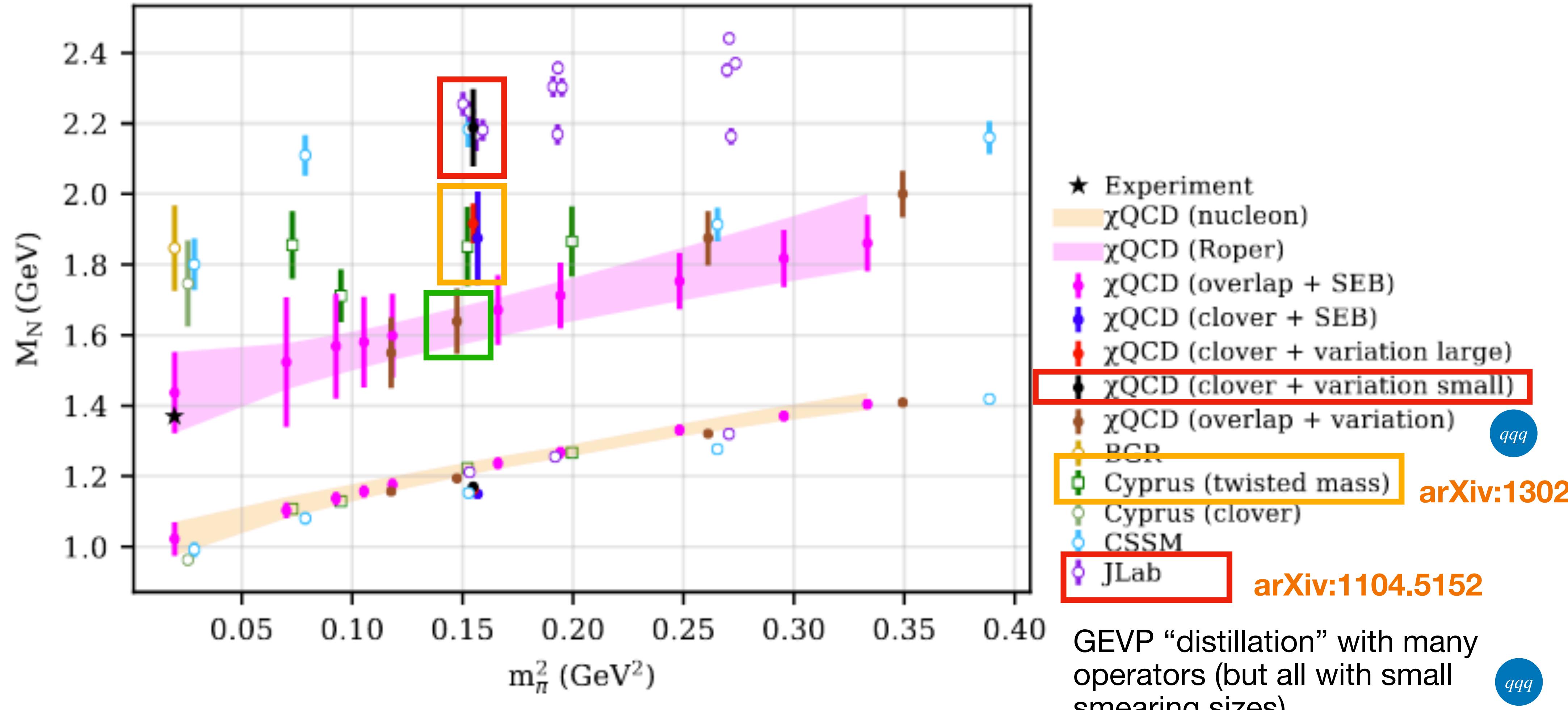
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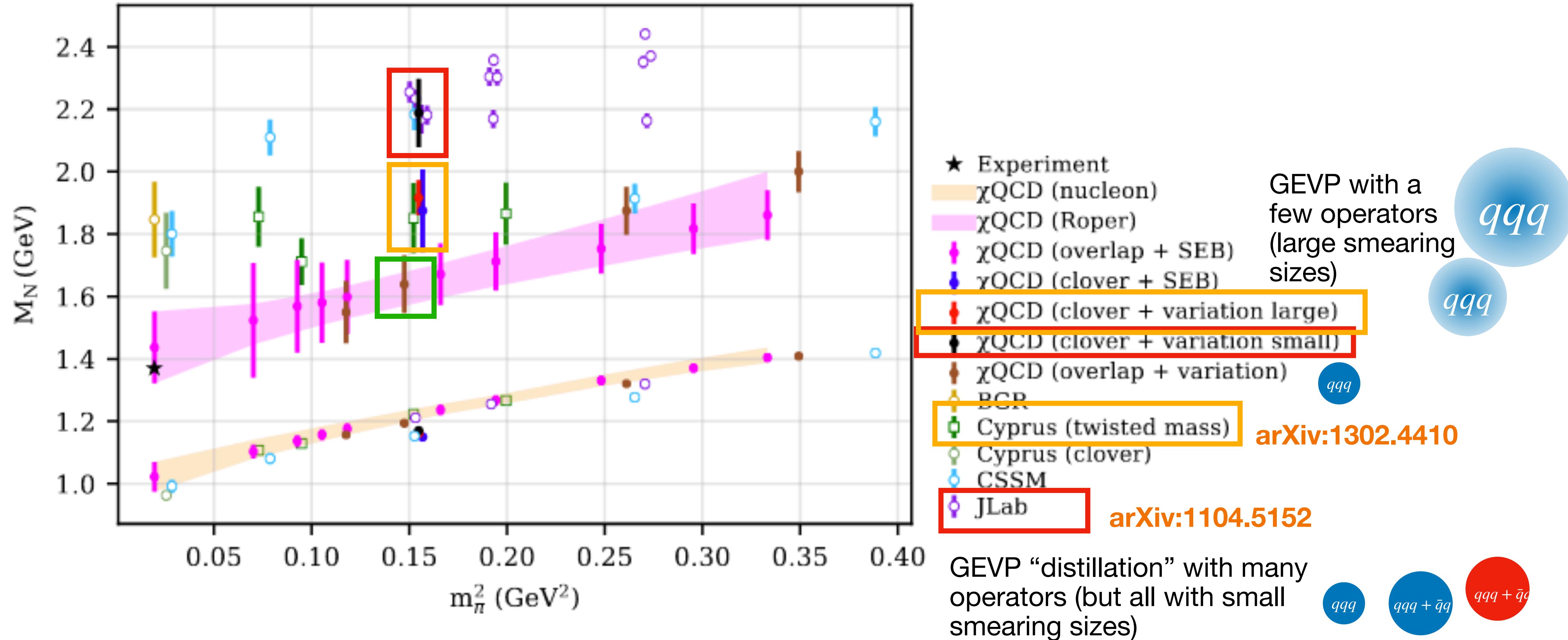
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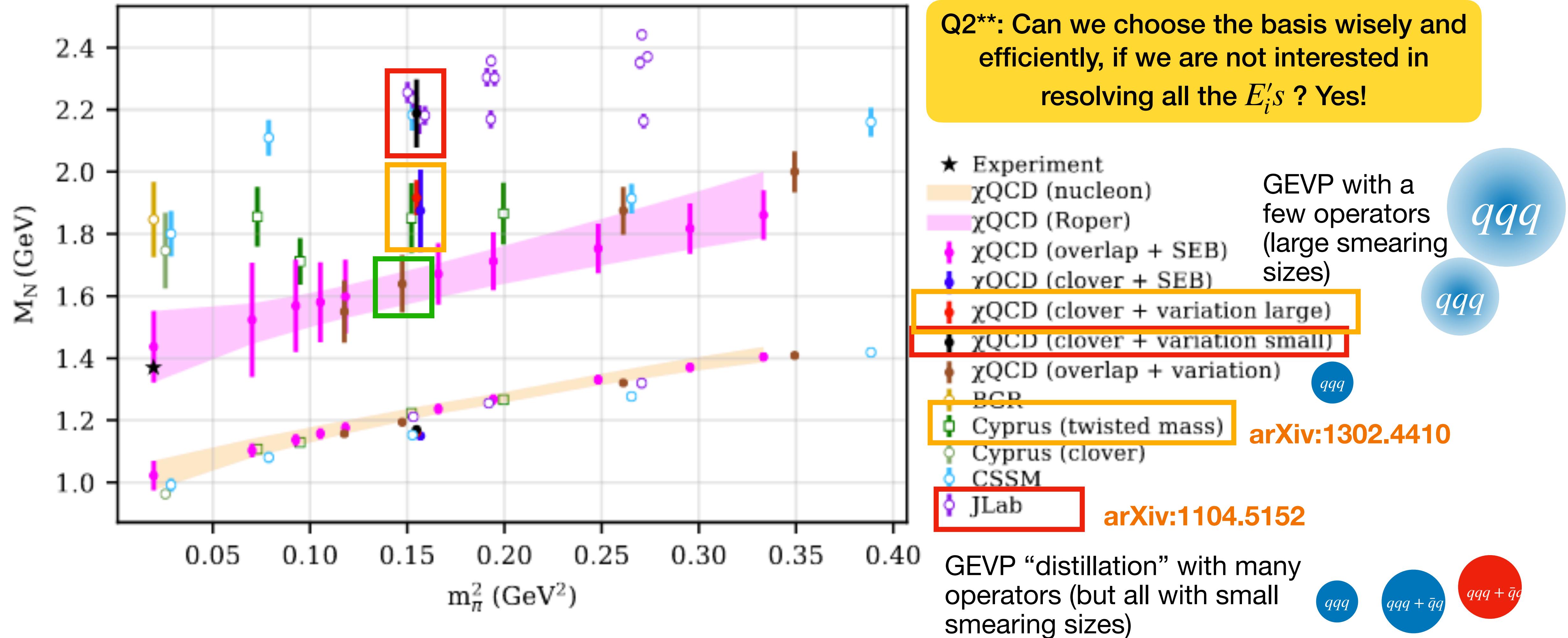
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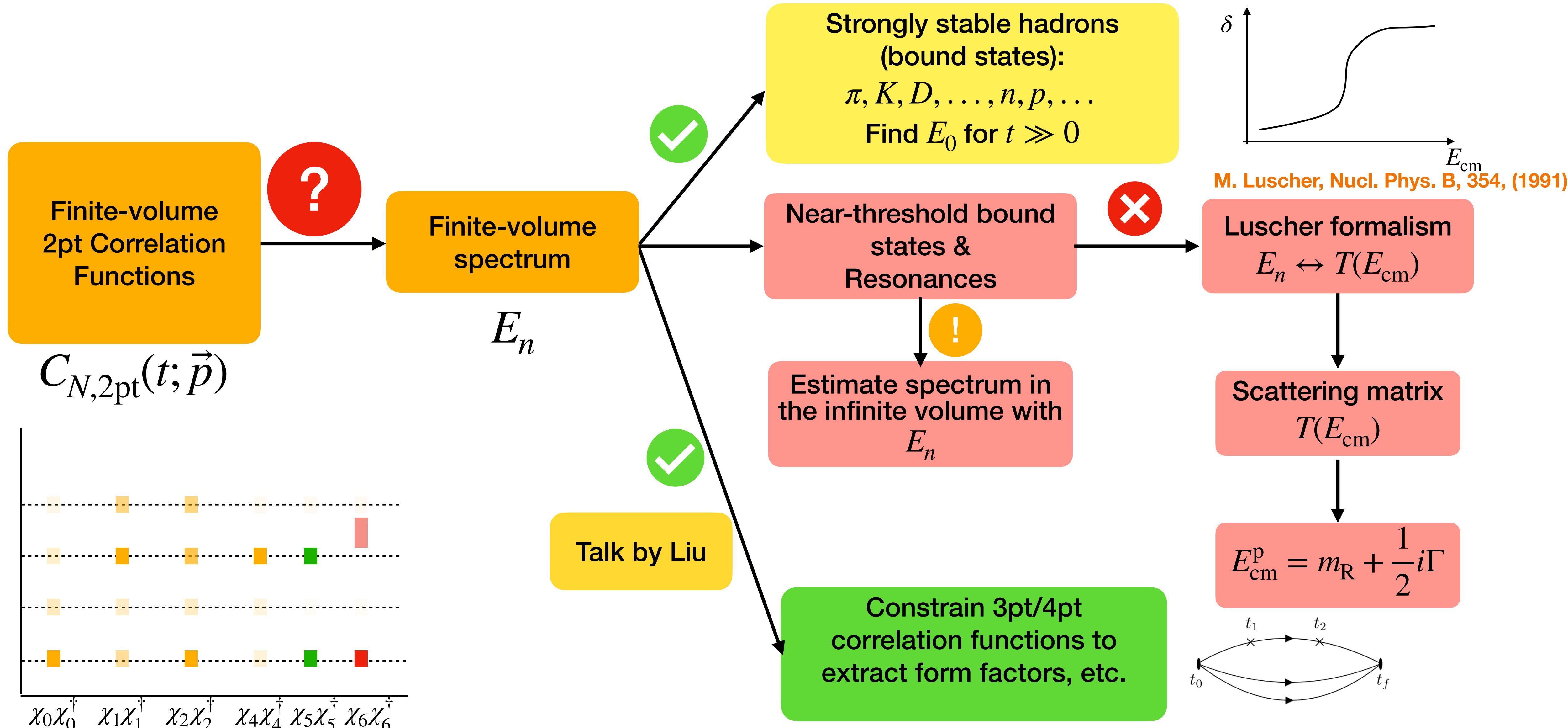


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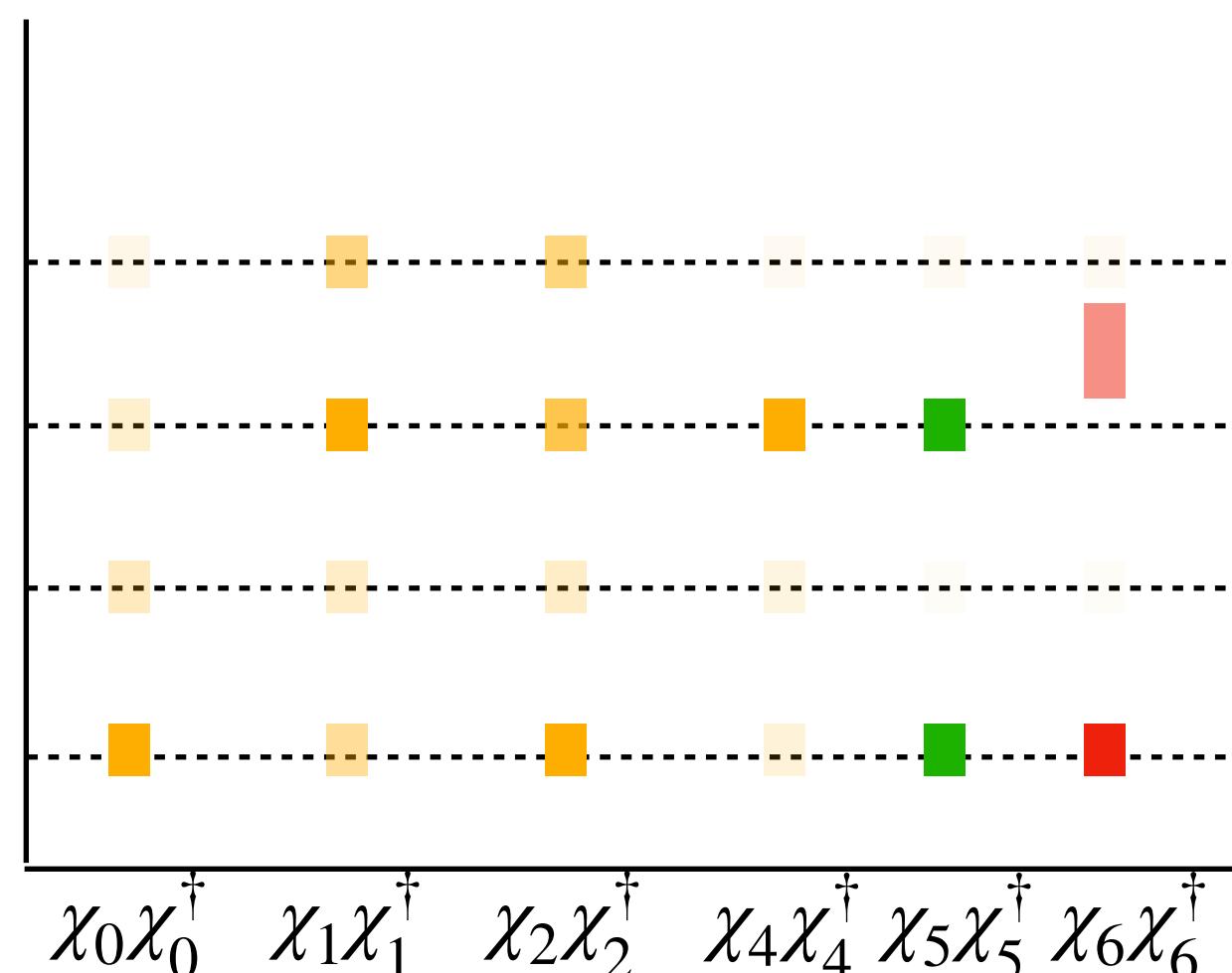
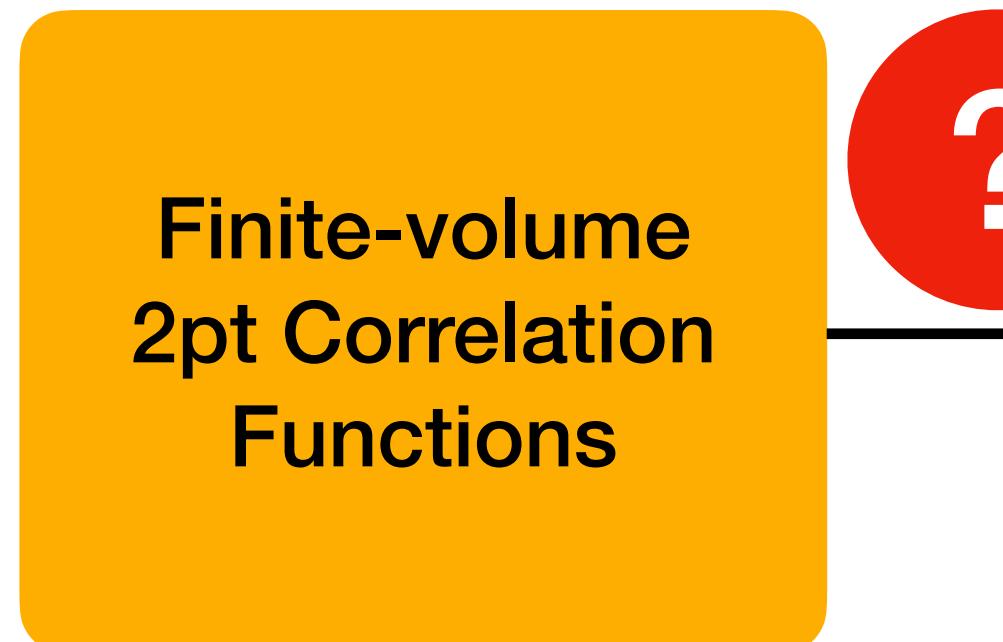


Finite-volume spectrum from lattice QCD



Finite-volume spectrum from lattice QCD

Q2**: Can we choose the basis wisely and efficiently, if we are not interested in resolving all the E'_i 's ? Yes!



Talk by Liu

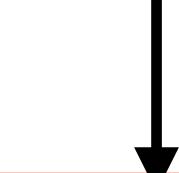


Strongly stable hadrons (bound states):
 $\pi, K, D, \dots, n, p, \dots$
 Find E_0 for $t \gg 0$

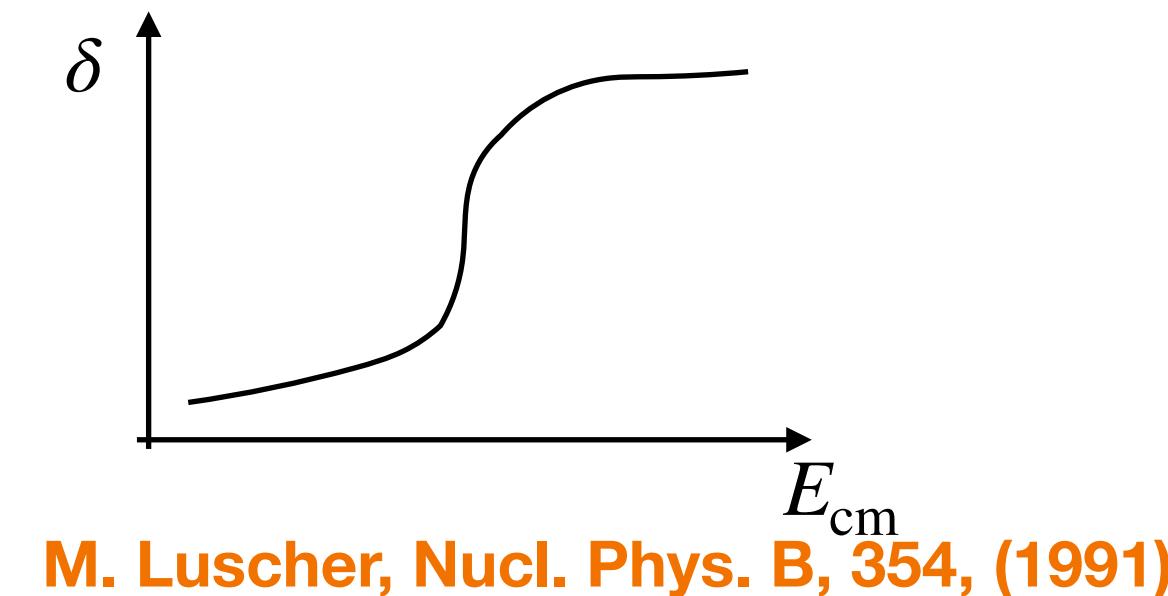
Near-threshold bound states & Resonances



Estimate spectrum in the infinite volume with E_n



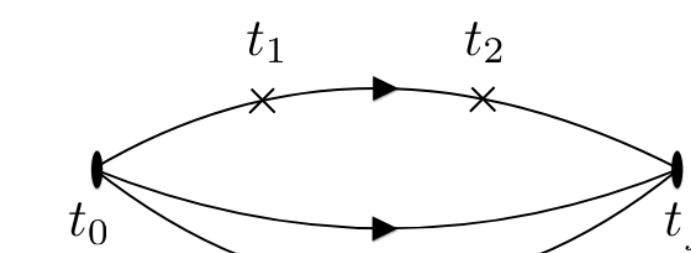
Constrain 3pt/4pt correlation functions to extract form factors, etc.



Luscher formalism
 $E_n \leftrightarrow T(E_{\text{cm}})$

Scattering matrix
 $T(E_{\text{cm}})$

$$E_{\text{cm}}^{\text{p}} = m_R + \frac{1}{2}i\Gamma$$



Finite-volume spectrum from lattice QCD

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Finite-volume
2pt Correlation
Functions

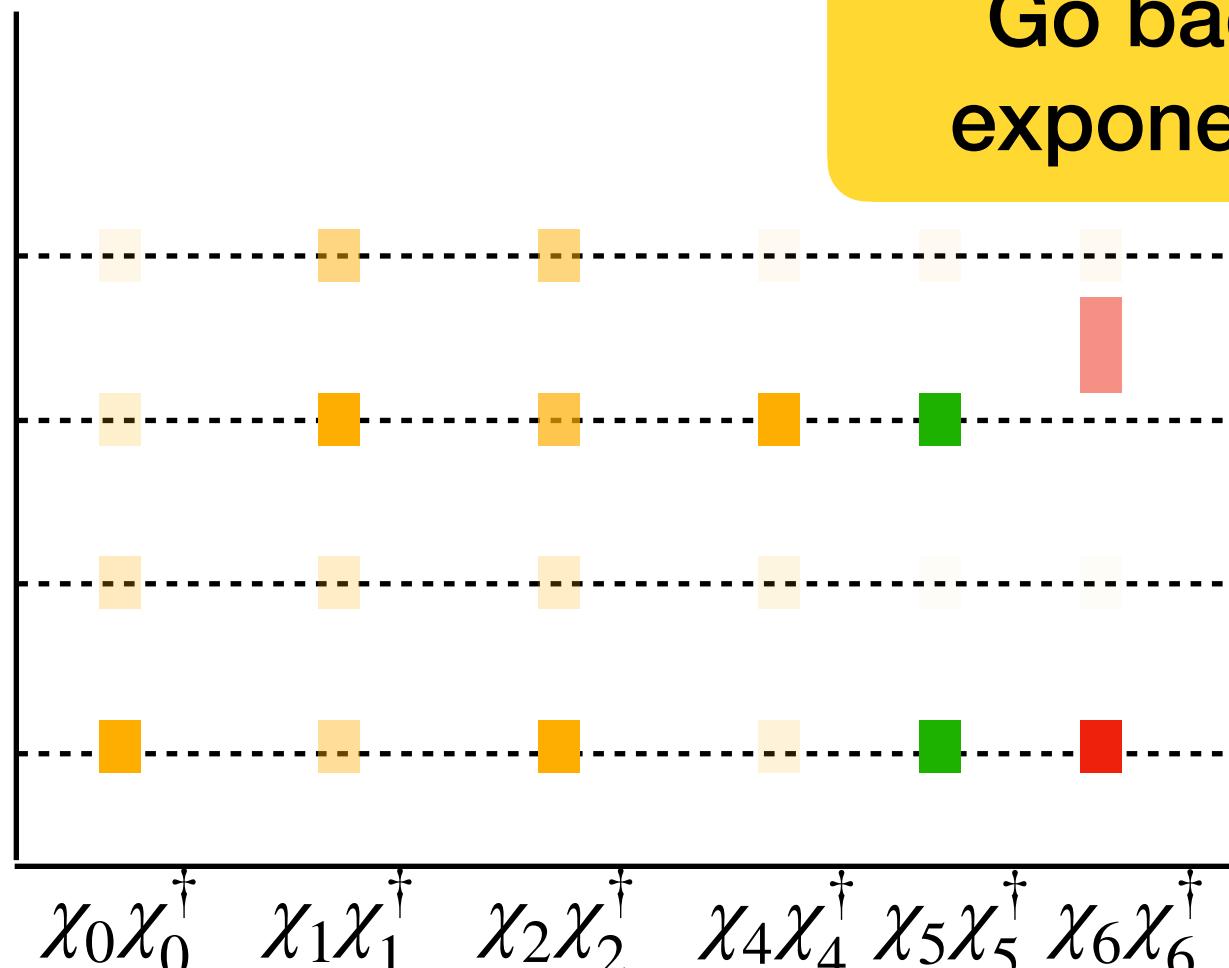
$$C_{N,2\text{pt}}(t; \vec{p})$$



Finite-volume
spectrum

$$E_n$$

Only one operator?
Go back to multi-exponential priors?



Talk by Liu



Strongly stable hadrons
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Find E_0 for $t \gg 0$

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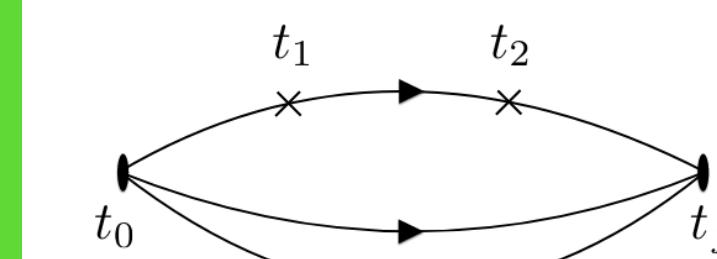
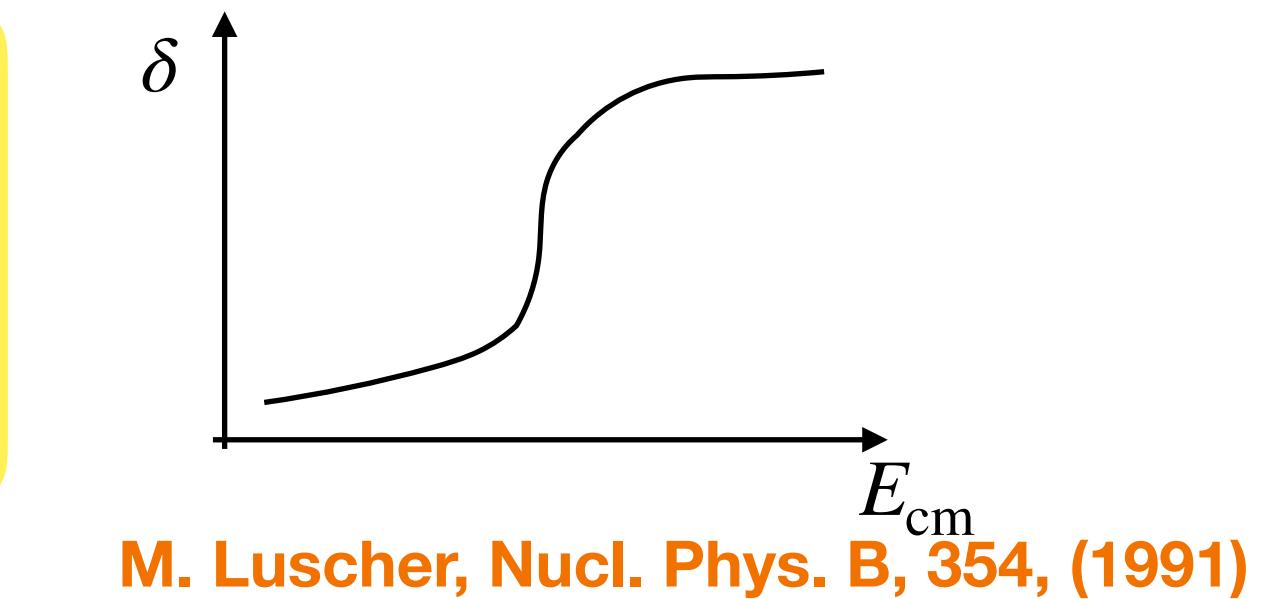


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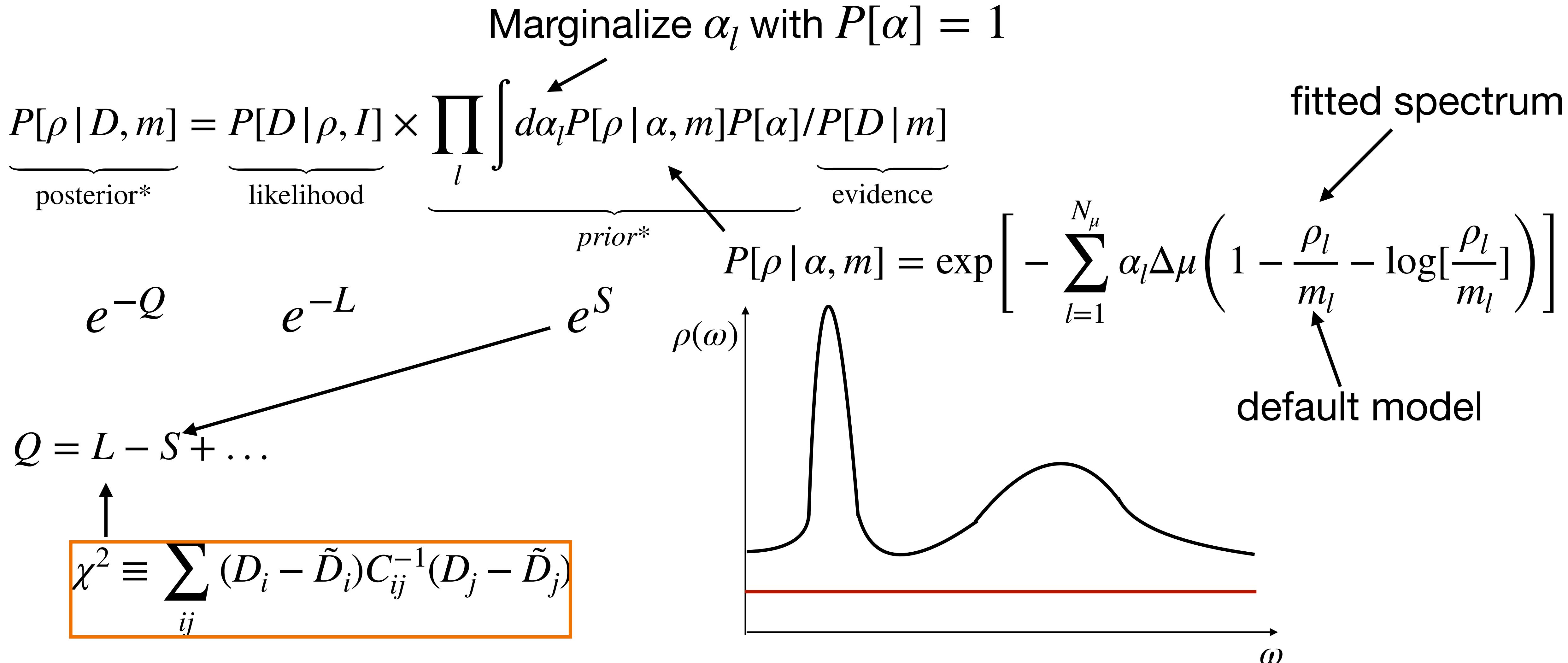
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The Bayesian Reconstruction for inverse problems

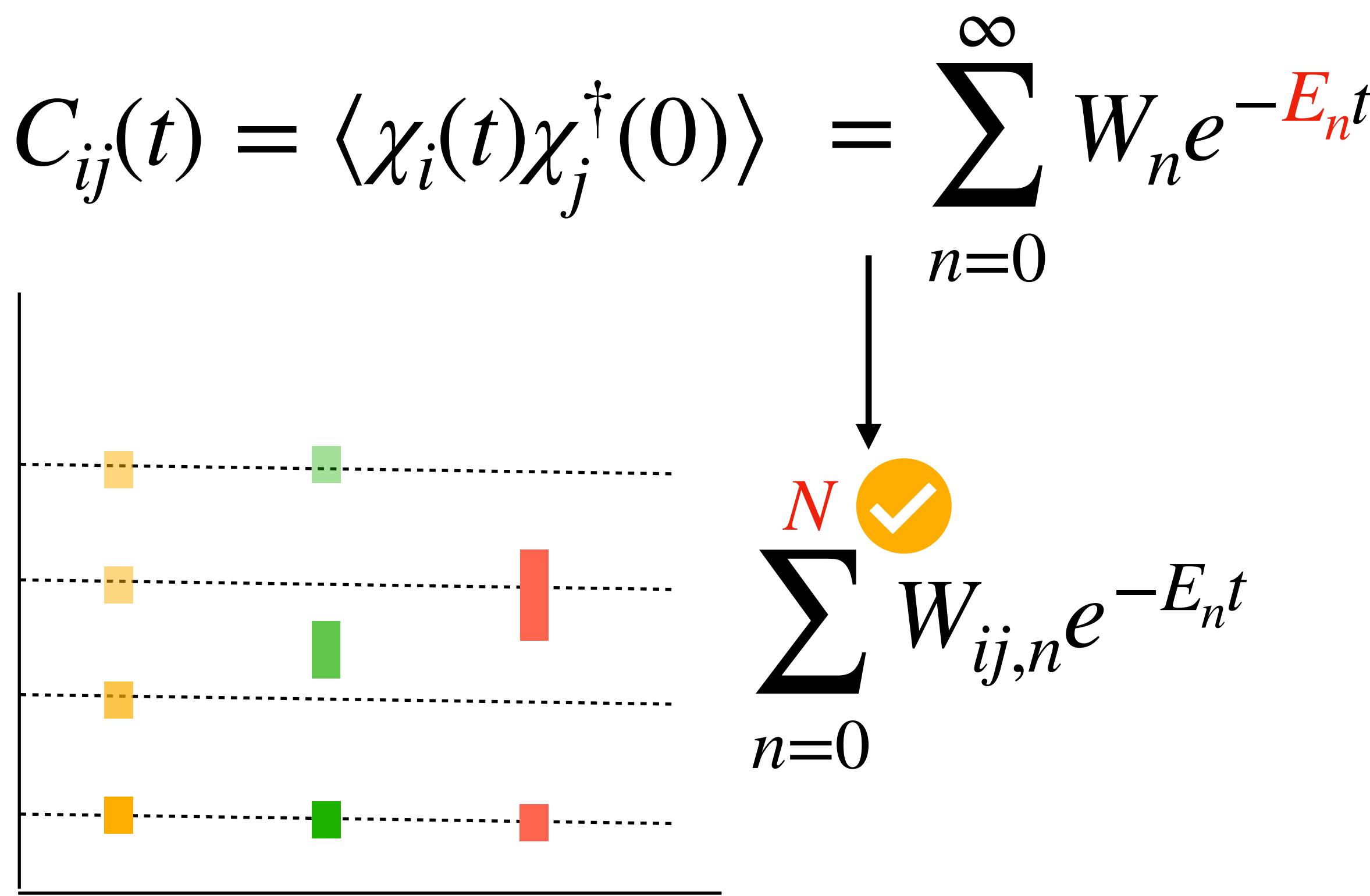
A. Rothkopf, “Bayesian inference of real-time dynamics from lattice QCD,” Front. Phys., arXiv:2208.13590

Yannis Burnier and Alexander Rothkopf, Phys. Rev. Lett. 111, 182003



The Bayesian Reconstruction on nucleon correlators

- To extract the finite-volume spectrum
 - Two-point correlation functions



$N = 4 \quad N = 3 \quad N = 2$

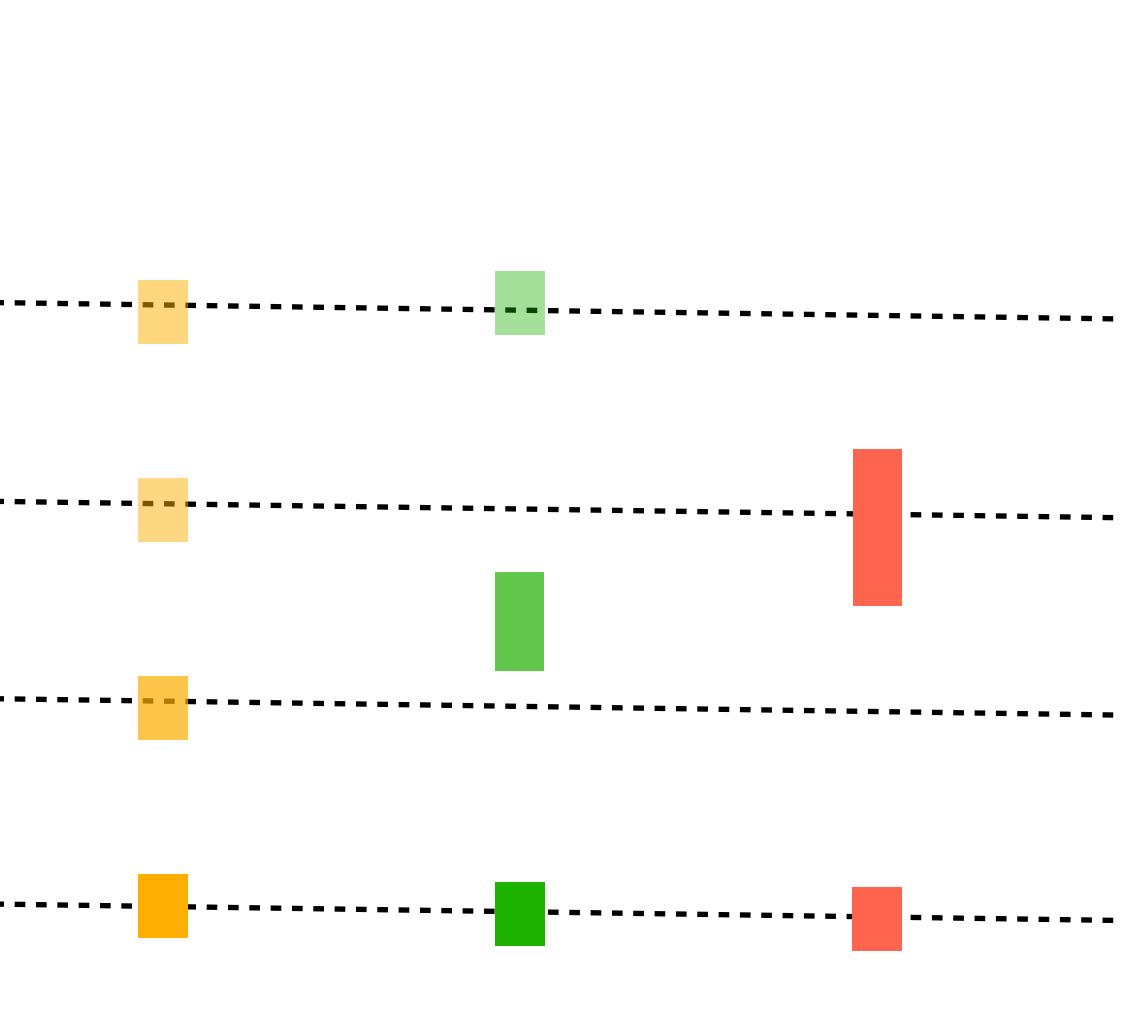
- Finite-volume spectrum found in terms of spectral function $\rho(\omega)$

The Bayesian Reconstruction on nucleon correlators

- To extract the finite-volume spectrum
 - Two-point correlation functions

$$C_{ij}(t) = \langle \chi_i(t) \chi_j^\dagger(0) \rangle = \sum_{n=0}^{\infty} W_n e^{-E_n t} \rightarrow \int d\omega \rho(\omega) e^{-\omega t}$$

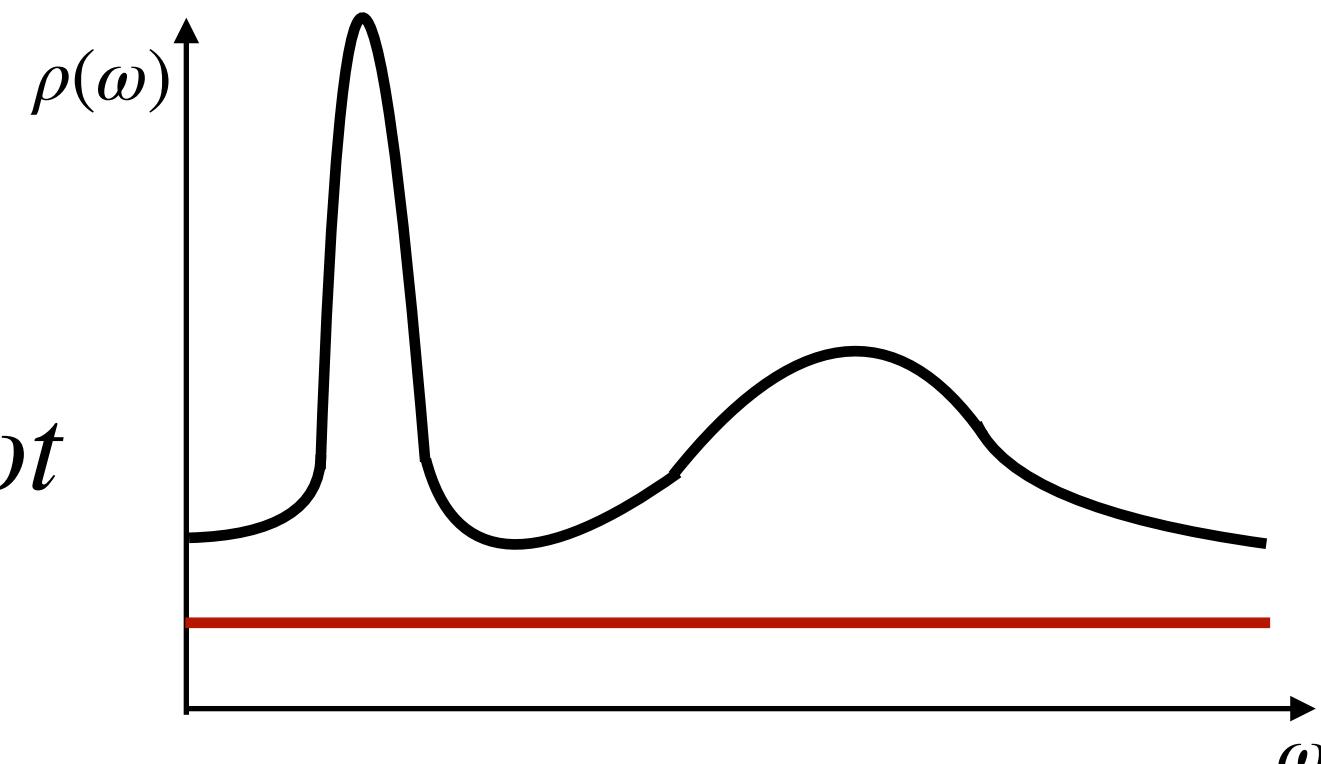
\downarrow

$$\sum_{n=0}^N W_{ij,n} e^{-E_n t}$$


$N = 4$ $N = 3$ $N = 2$

M. T. Hansen et al.,
Phys. Rev. D 96, 094513 (2017)

Talks by Liang,
Sufian

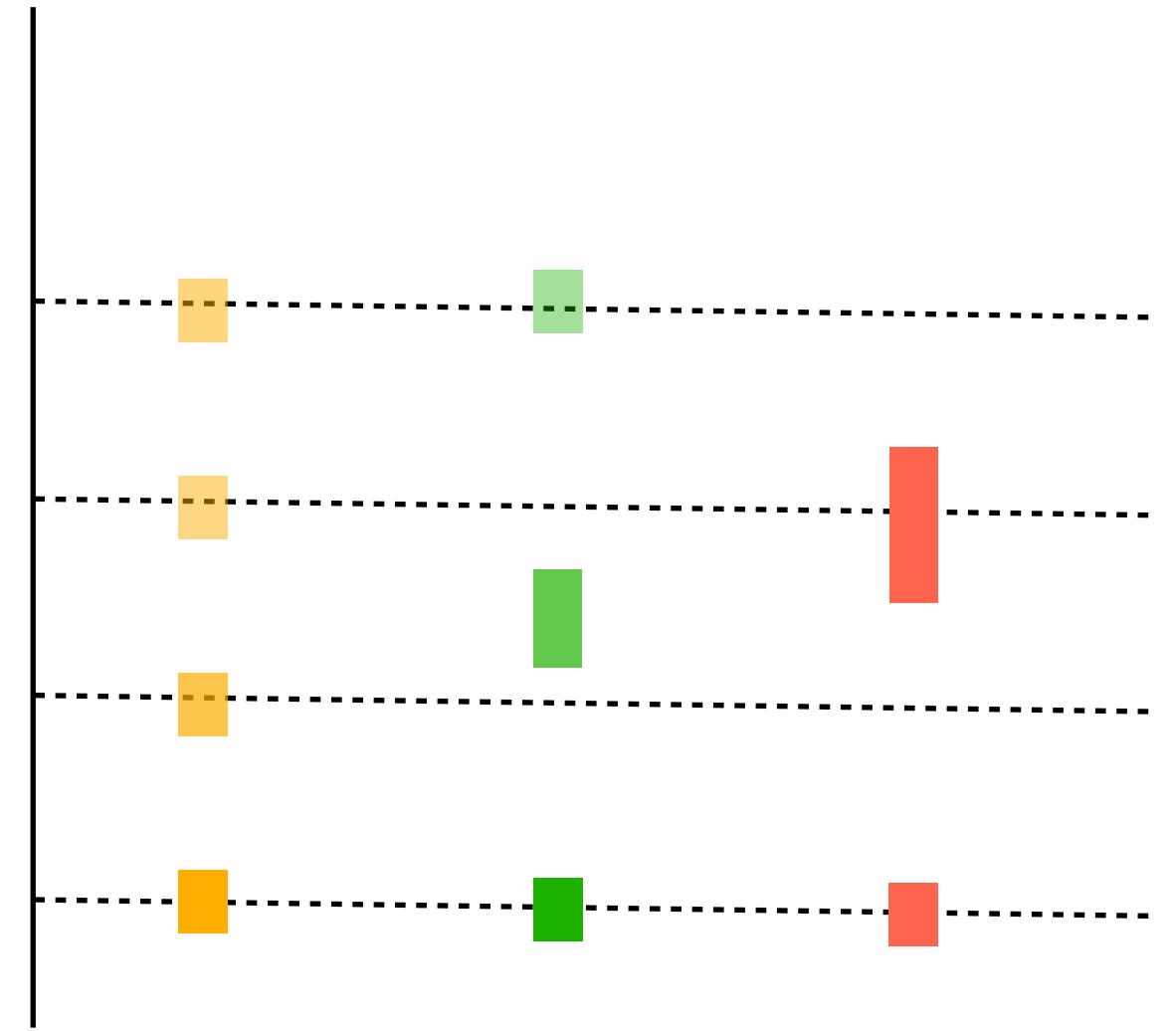


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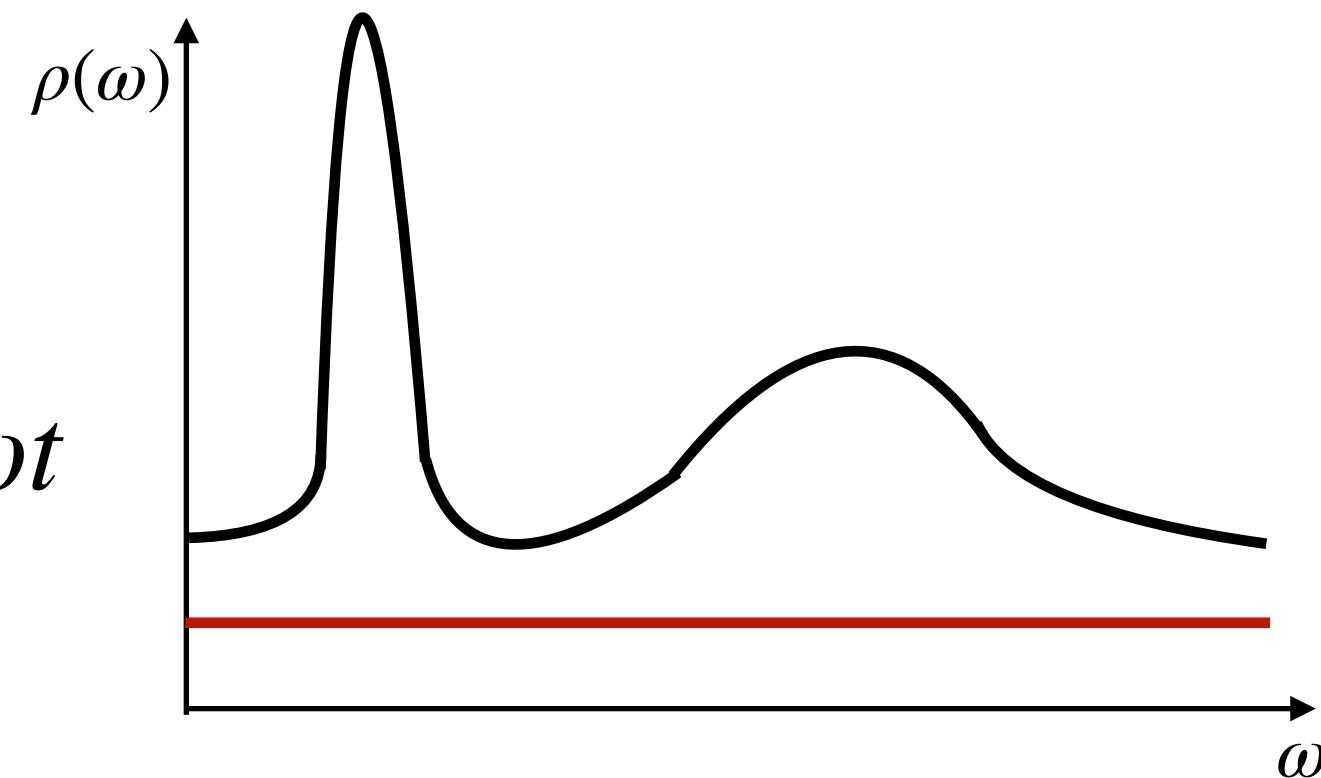


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M. T. Hansen et al.,
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$$\rightarrow \int d\omega \rho(\omega) e^{-\omega t}$$



$$\sum_{n=0}^N W_{ij,n} e^{-E_n t}$$

$$\rho(\omega) = \sum_{n=0}^N W_n \delta(\omega - E_n)$$

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\downarrow

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\checkmark

G. P. Lepage et al., arXiv: hep-lat/0110175

$N = 4 \quad N = 3 \quad N = 2$



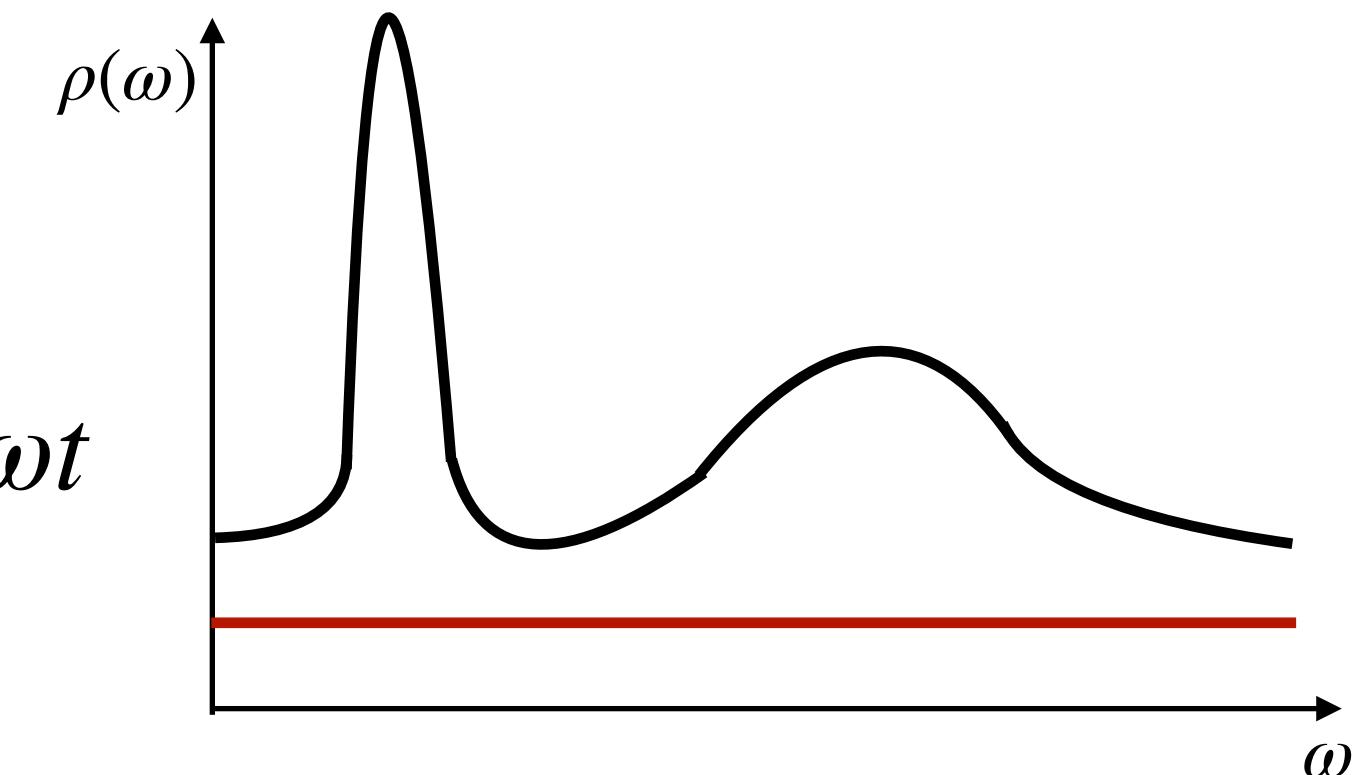
$$\chi^2 \equiv \sum_{ij} (D_i - \tilde{D}_i) C_{ij}^{-1} (D_j - \tilde{D}_j)$$

$$\chi_{\text{prior}}^2 \equiv \sum_n \frac{(W_n - \tilde{W}_n)^2}{\tilde{\sigma}_{W_n}^2} + \sum_n \frac{(E_n - \tilde{E}_n)^2}{\tilde{\sigma}_{E_n}^2}$$

M. T. Hansen et al.,
Phys. Rev. D 96, 094513 (2017)

Talks by Liang,
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$$C_{ij}(t) = \langle \chi_i(t)\chi_j^\dagger(0) \rangle = \sum_{n=0}^{\infty} W_n e^{-E_n t}$$

\downarrow

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N ✓

G. P. Lepage et al., arXiv: hep-lat/0110175

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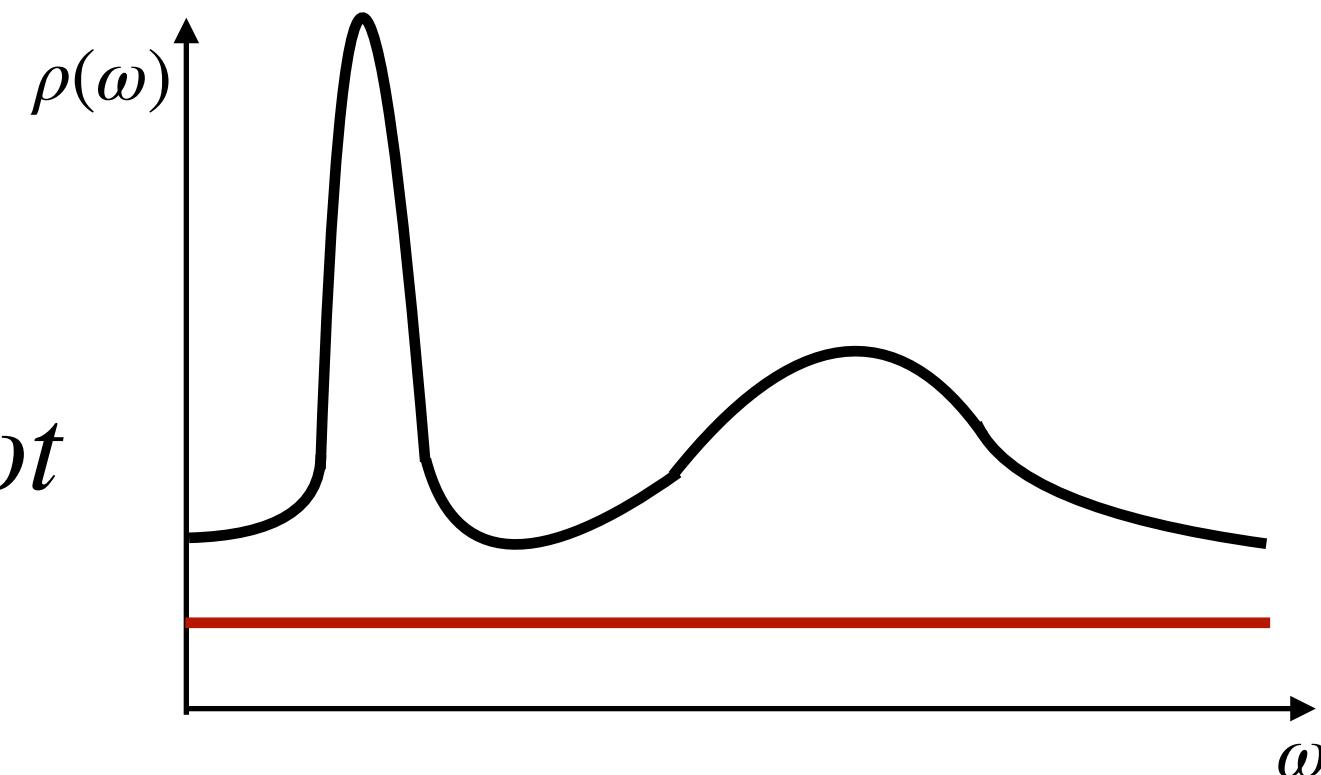
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M. T. Hansen et al.,
Phys. Rev. D 96, 094513 (2017)

Talks by Liang,
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$$\rightarrow \int d\omega \rho(\omega) e^{-\omega t}$$



$$\rho(\omega) = \sum_{n=0}^N W_n \delta(\omega - E_n)$$

$$Q = L - S + \dots$$

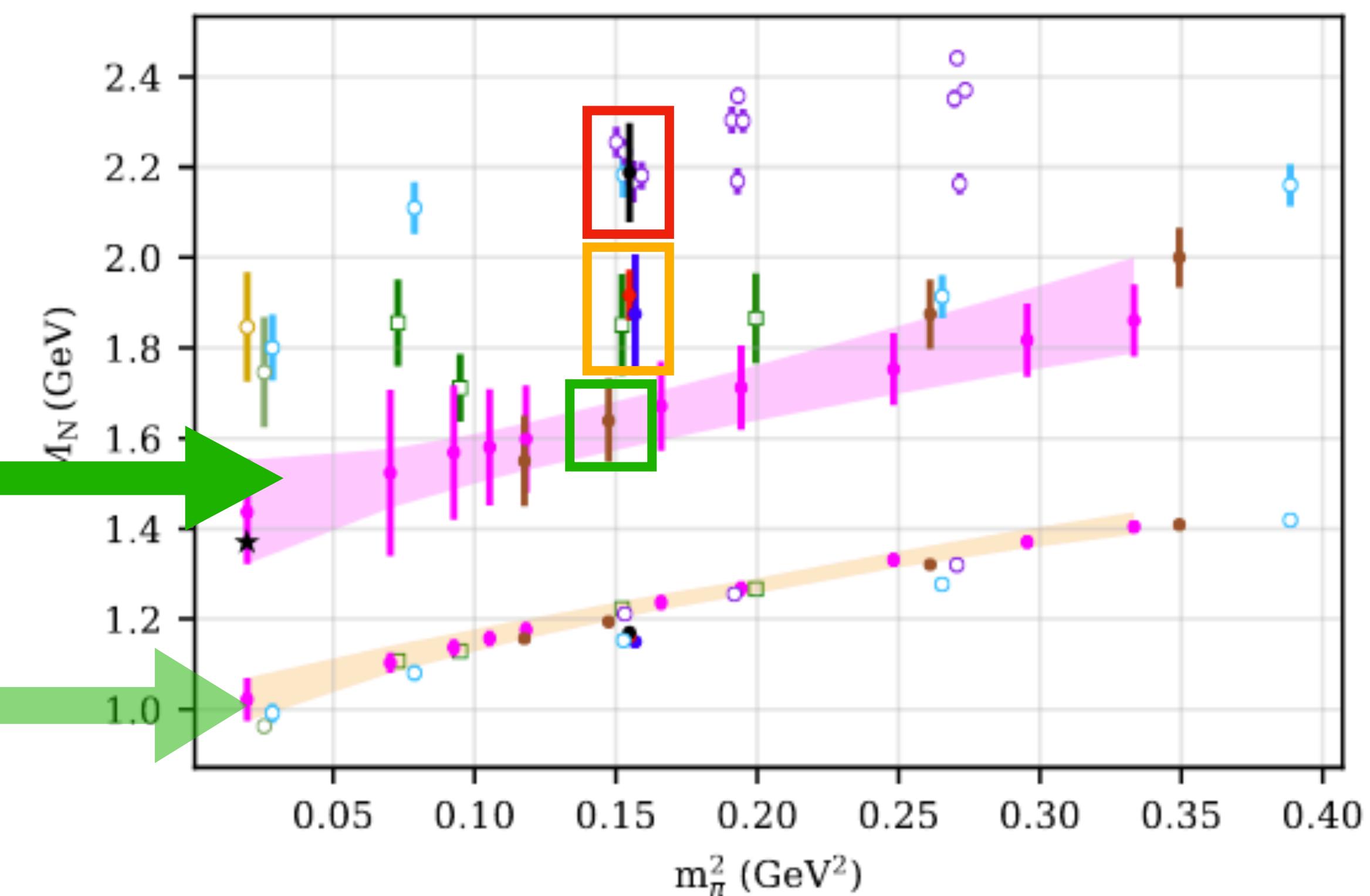
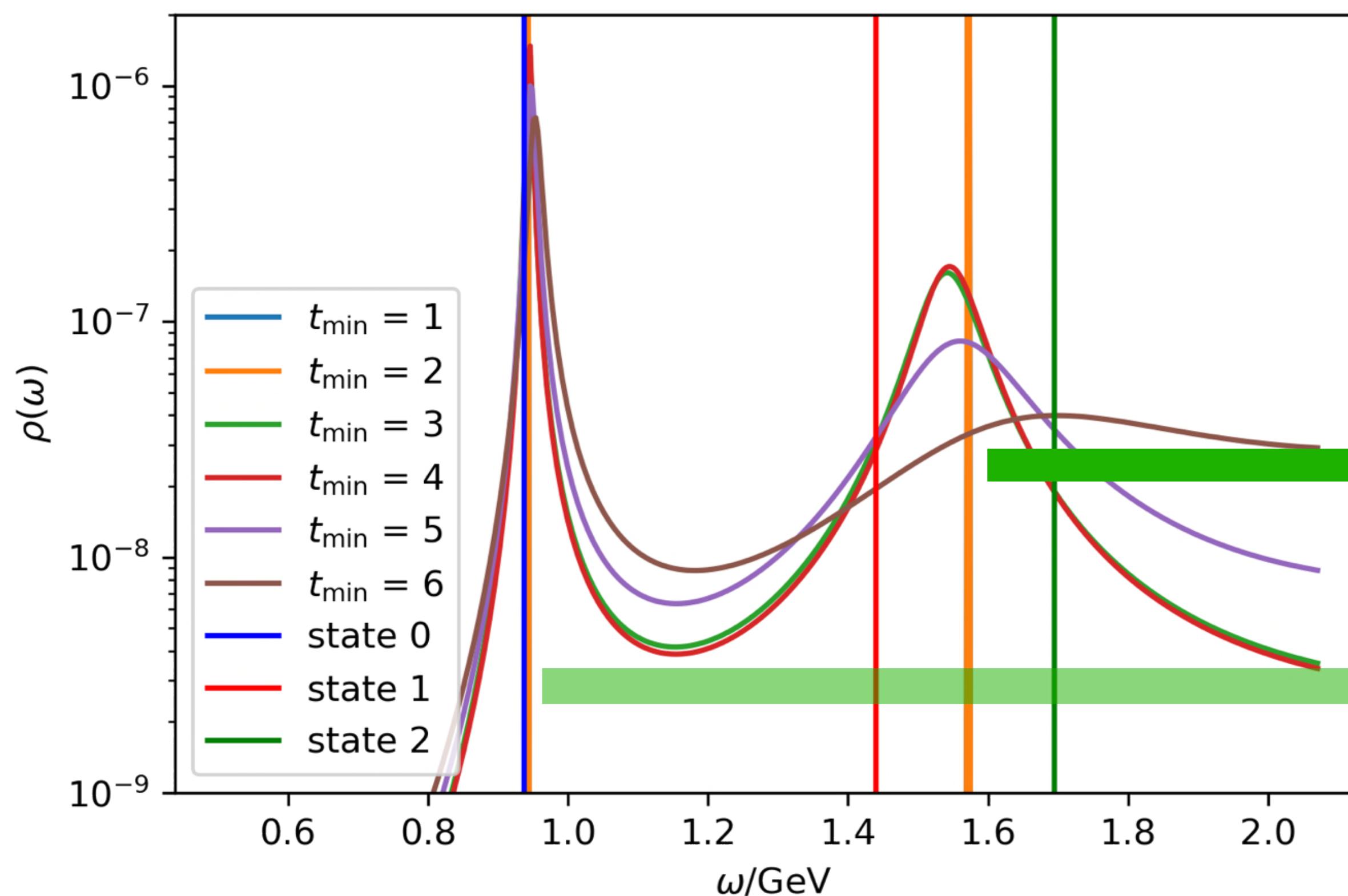
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Finite-volume spectrum with BR preliminary

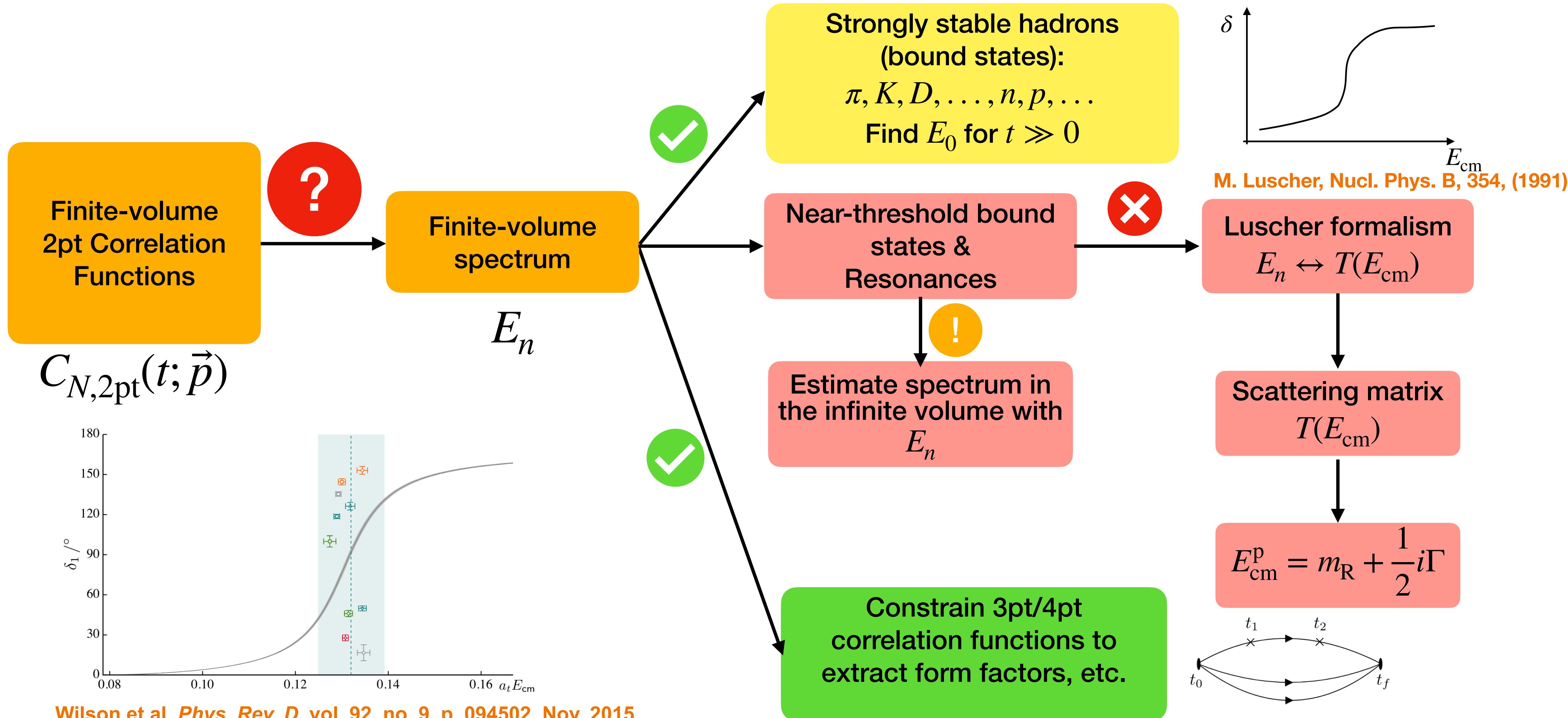
Ens.	Action (F+G)	$1/a$ (GeV)	Lattice volume	m_l	m_s	m_{res}	m_π (MeV)	Size (fm)
				(in lattice units)				
10	MDWF+I	1.730(4)	$48^3 \times 96 \times 24$	0.00078	0.0362	0.000614	139	5.5

T. Blum et al., Phys. Rev. D 93, 074505

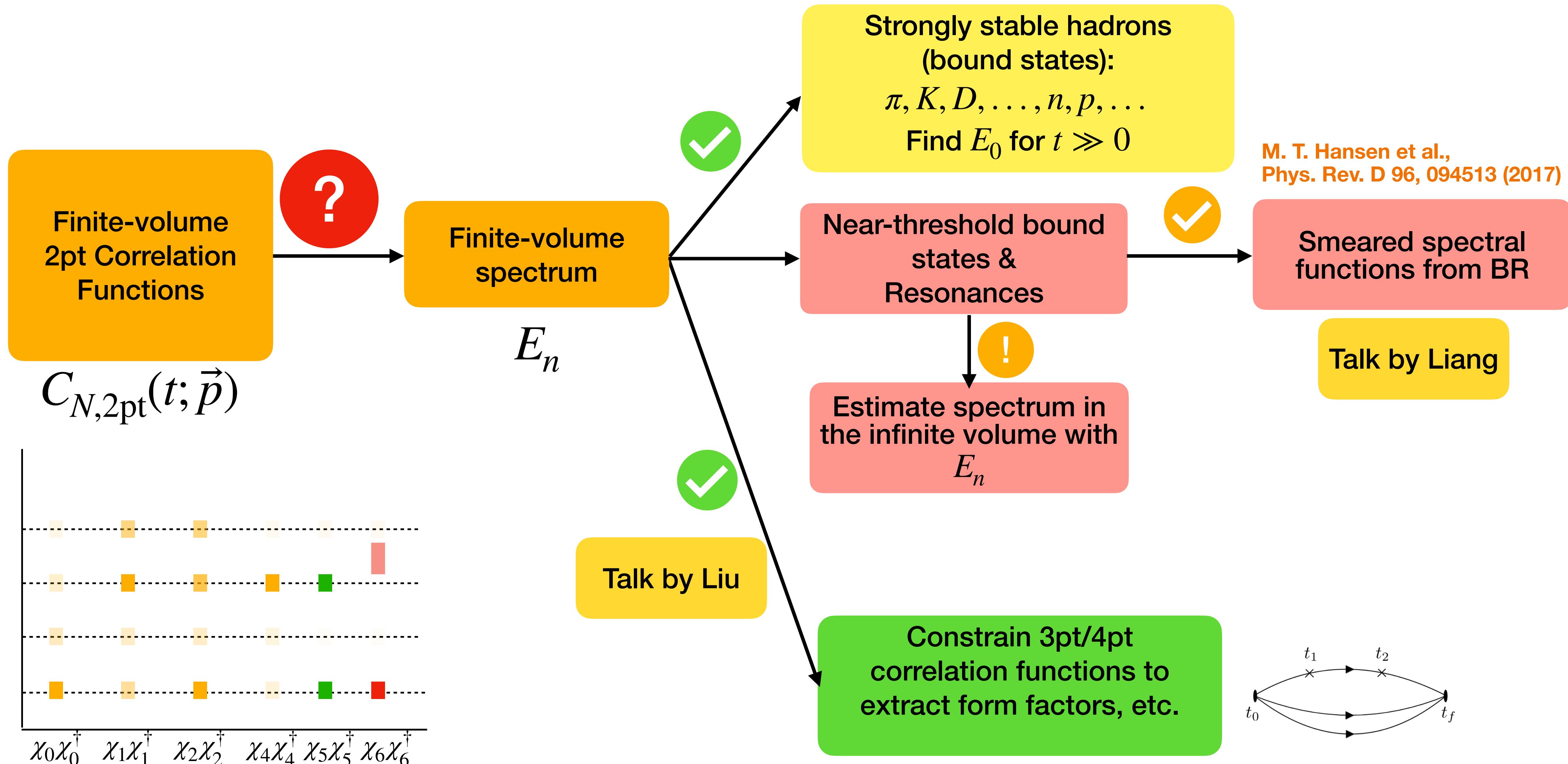
Default model
 $m(\omega) = \text{const.}$



Finite-volume spectrum from lattice QCD

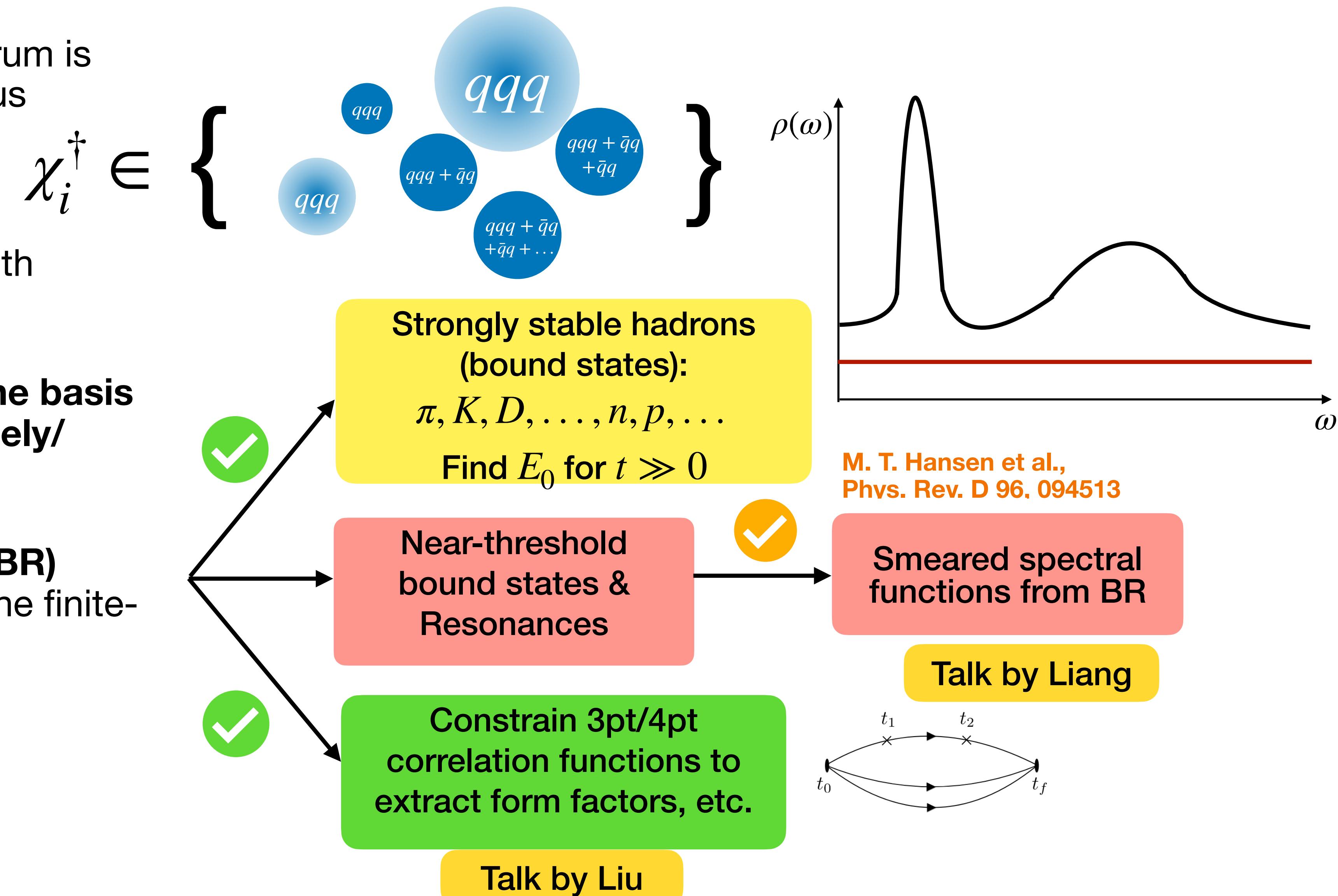


Finite-volume spectrum from lattice QCD with BR



Conclusion and outlook

- Extraction of finite-volume spectrum is important in lattice QCD for various physics topics.
- We don't **always** have to use **many interpolating operators** with **high computational cost**.
 - Theories can help us **choose the basis of interpolating operators wisely/efficiently**.
 - **The Bayesian Reconstruction (BR) method can be used** to extract the finite-volume spectrum $\rho(\omega)$.
 - More on-going tests of BR
 - To be continued ...



Thanks for your attention!