

New Physics Searches at the

Precision Frontier

9 May, 2023

INT-23-1b, Seattle, USA

QED radiative corrections to
charged-current neutrino scattering
for accelerator neutrinos

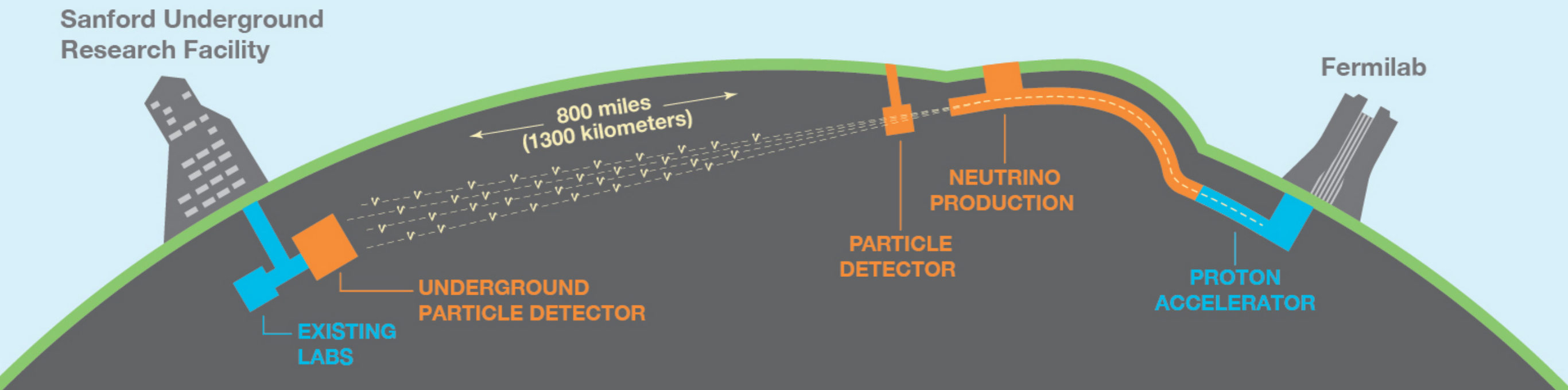


Oleksandr (Sasha) Tomalak

LA-UR-23-23974

Neutrino experiments

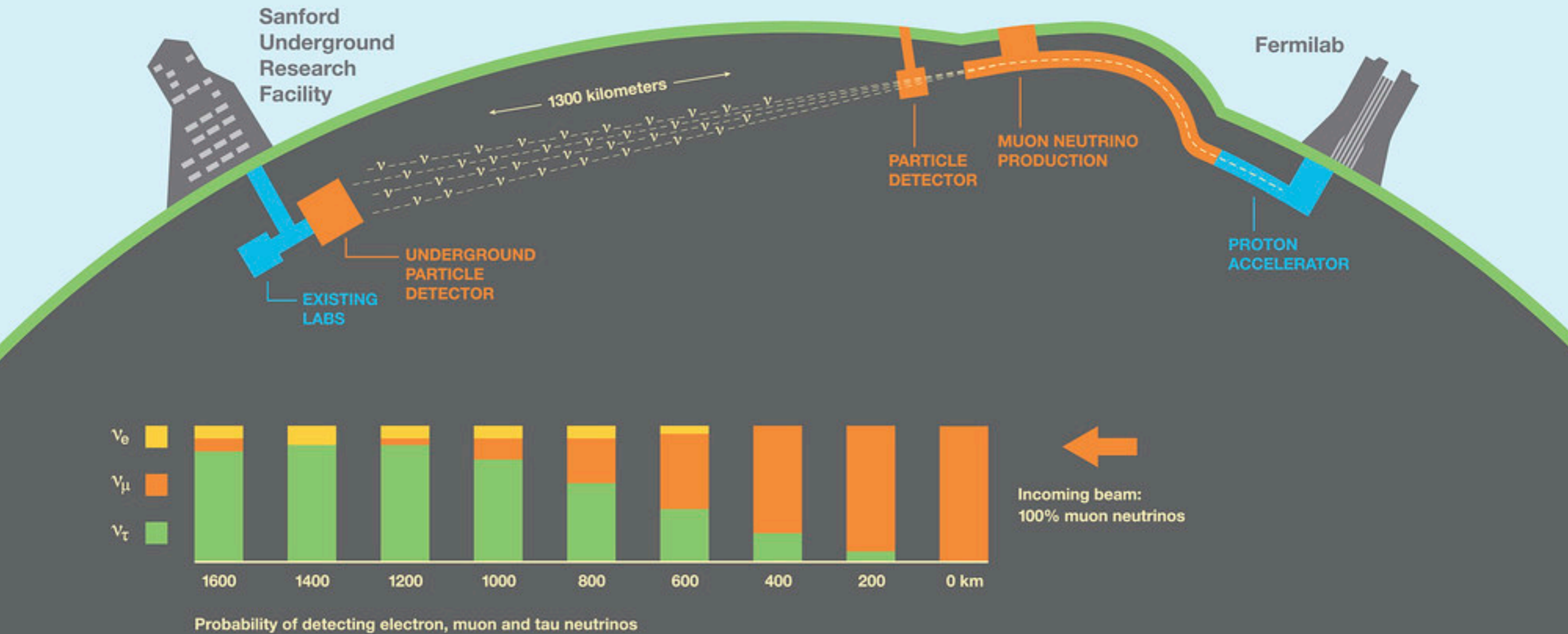
- **DUNE** and Hyper-K: leading-edge ν science experiments



- origin of matter-antimatter asymmetry δ_{CP}
- mass hierarchy and oscillation parameters PMNS matrix, Δm_{31}^2
- Grand Unified Theories proton decay
- dynamics of supernova explosion wait for one;)

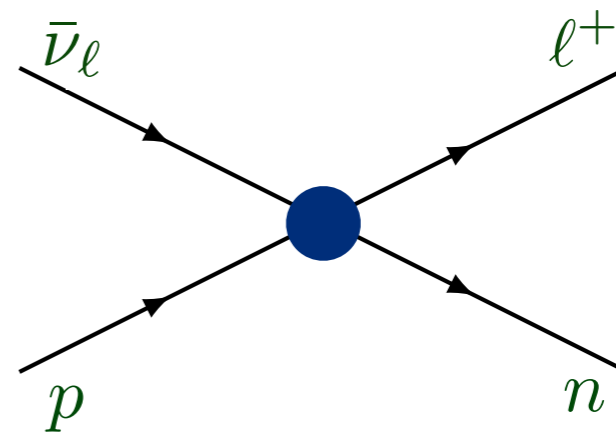
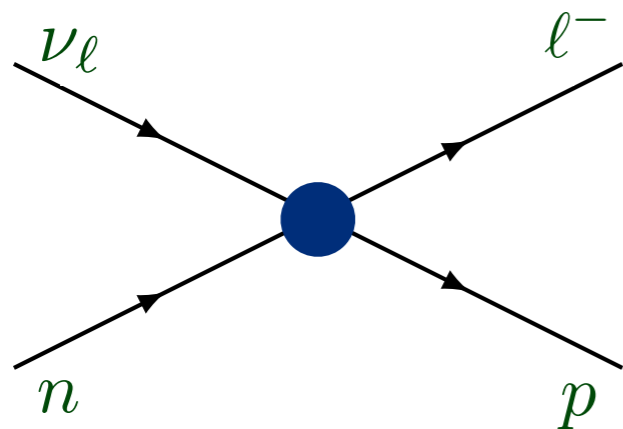
Neutrino experiments

Deep Underground Neutrino Experiment



$$N_\nu \sim \int dE_\nu \Phi_\nu (E_\nu) \times \sigma (E_\nu) \times R (E_\nu, E_\nu^{\text{rec}})$$

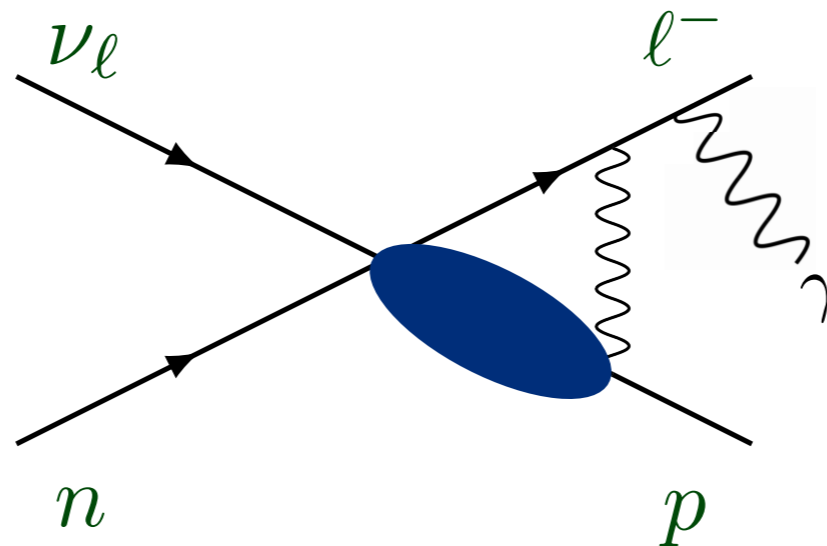
- precise neutrino physics: need in cross sections



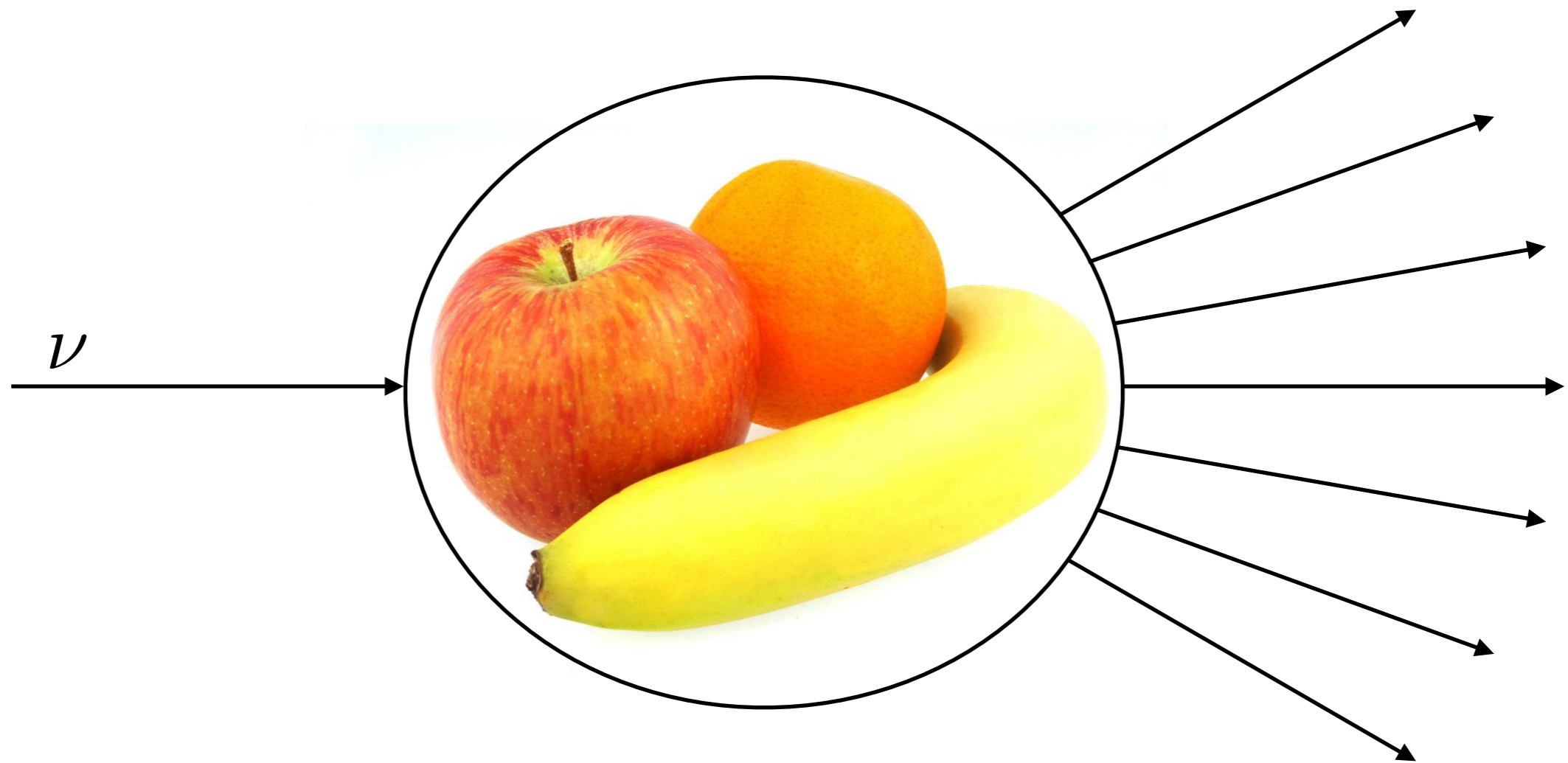
O. T., Qing Chen, Richard J. Hill and Kevin S. McFarland, Nature Commun. 13 (2022), 1, 5286

Radiative corrections in charged-current elastic scattering on free nucleons

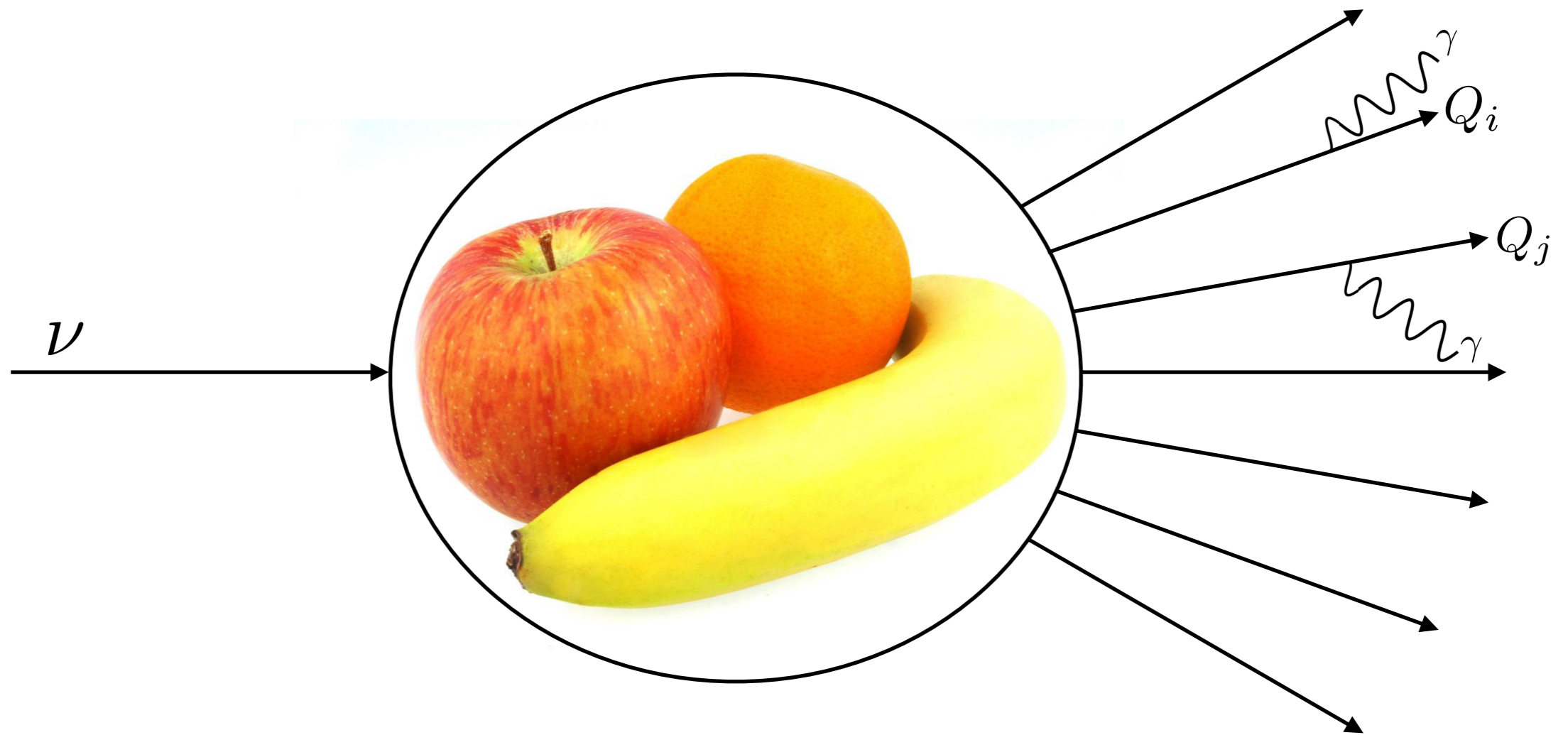
O. T., Qing Chen, Richard J. Hill, Kevin S. McFarland and Clarence Wret
 editors suggestion in Phys. Rev. D (2022)



Neutrino interactions

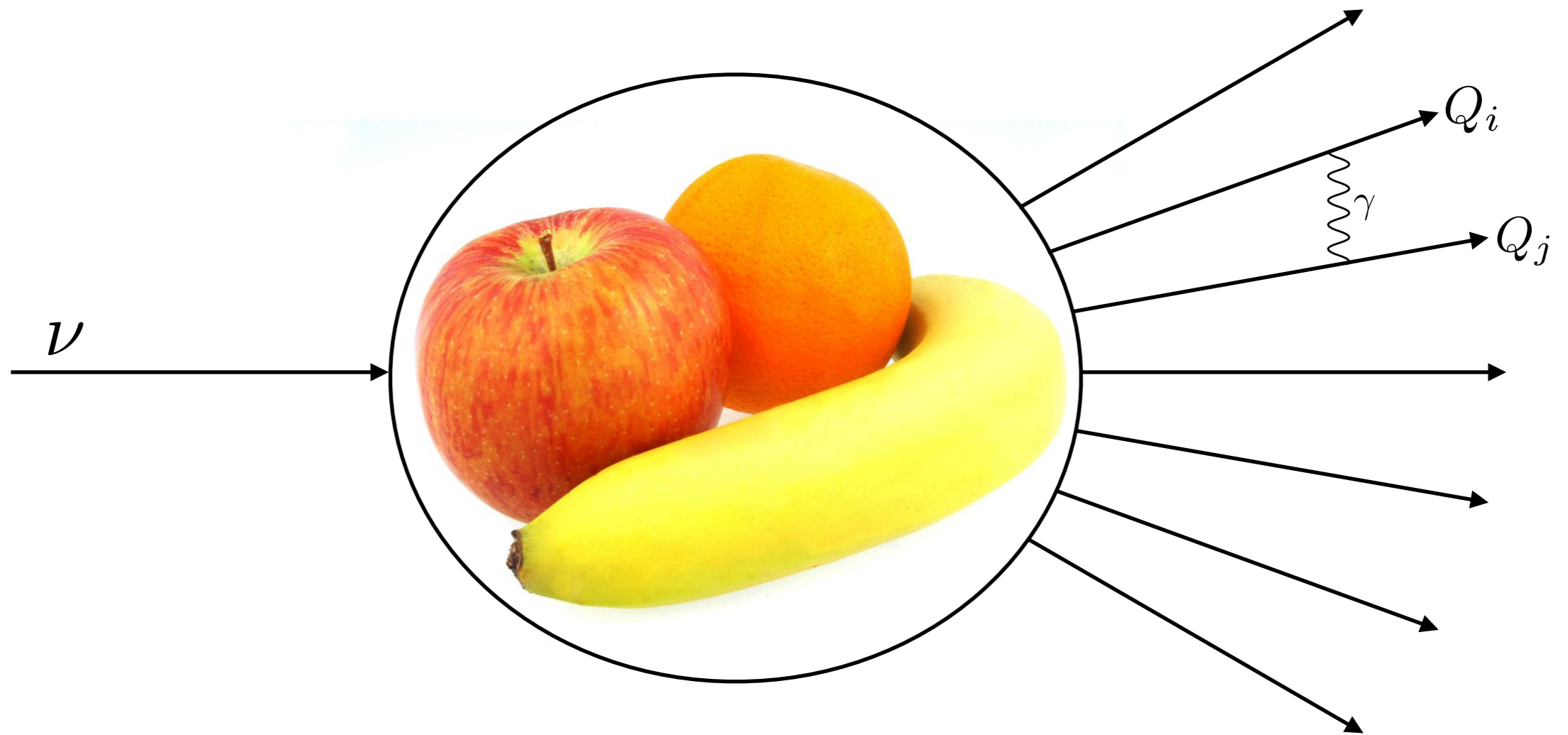


QED corrections



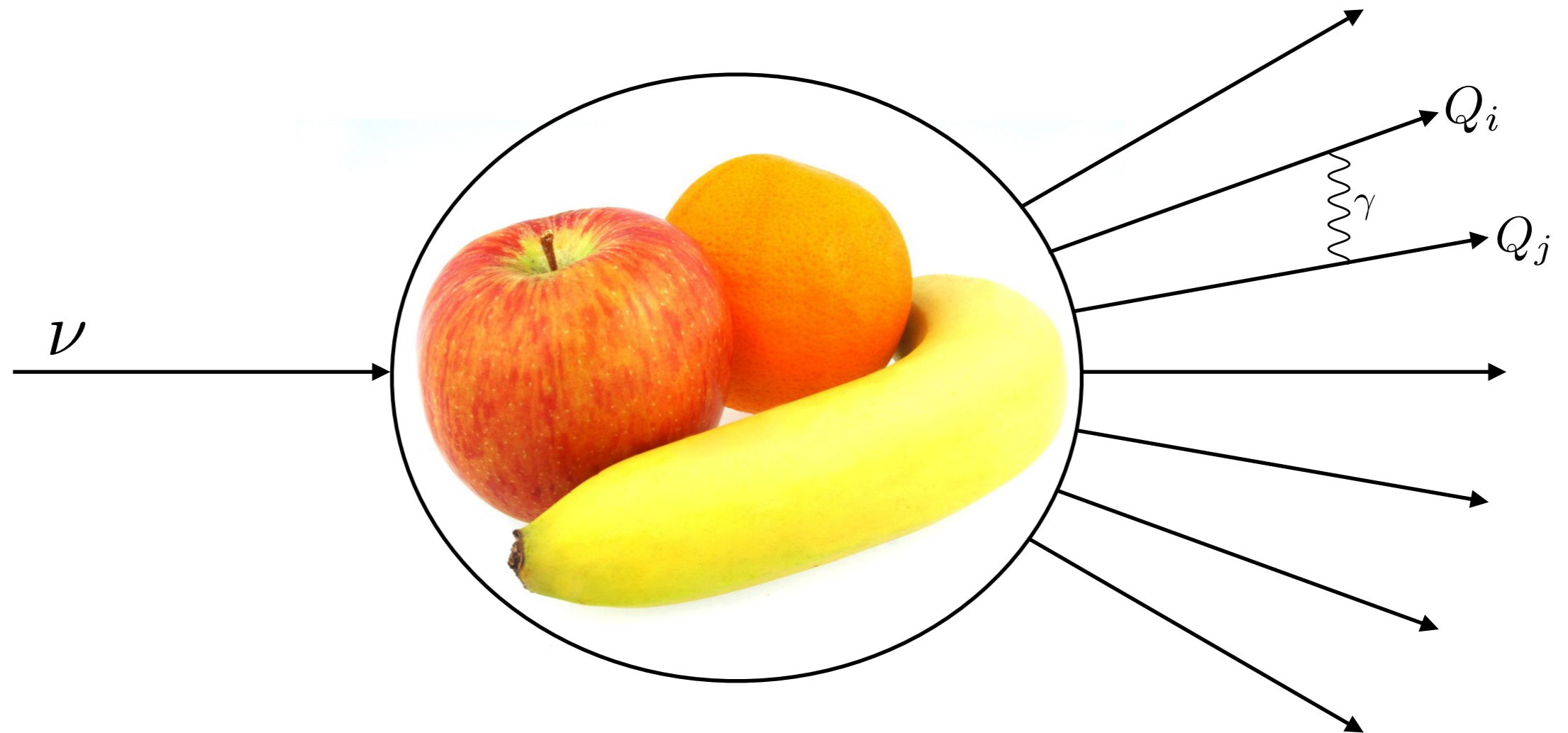
- all charged particles couple to real and virtual photons

QED corrections



- all charged particles couple to real and virtual photons

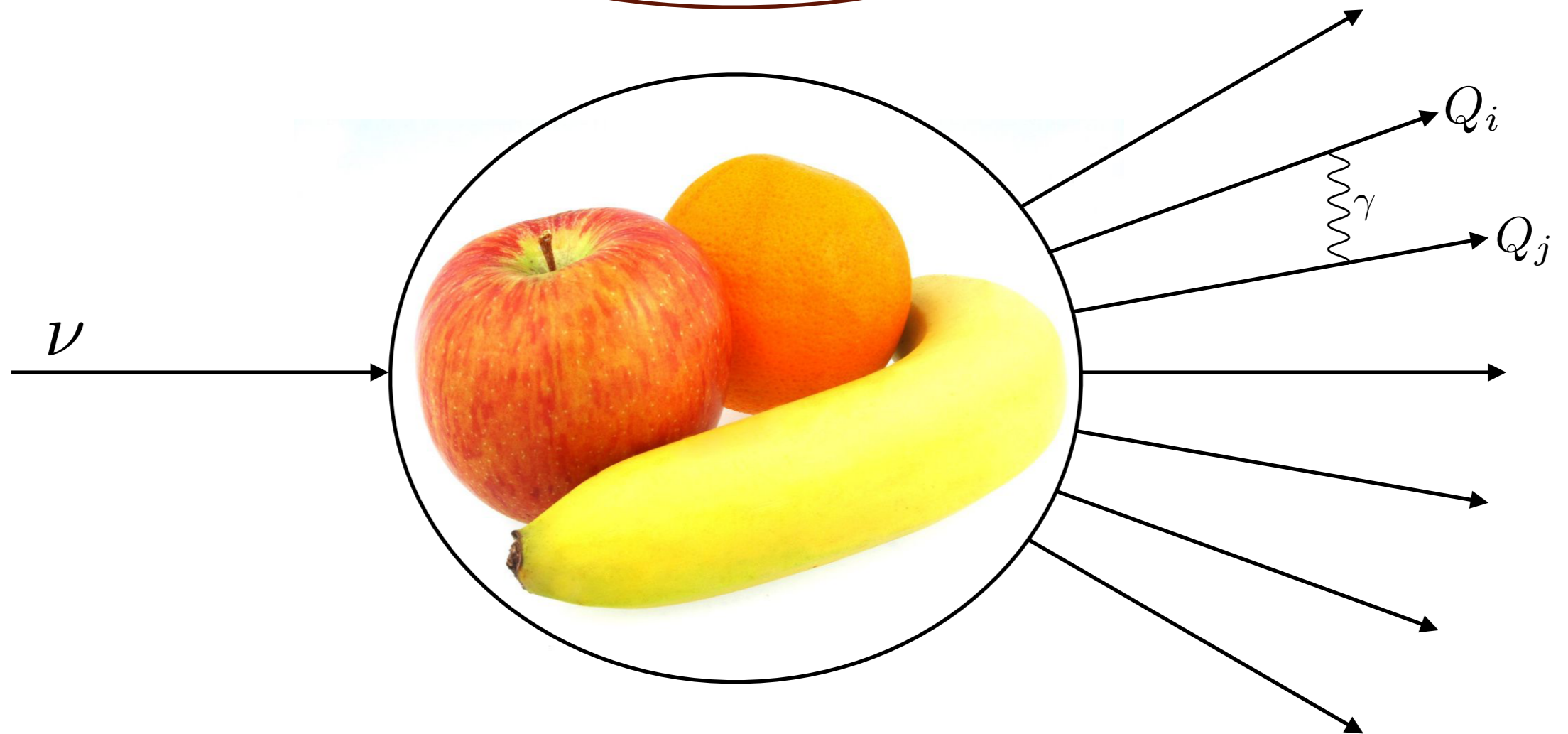
QED corrections



- $\frac{\alpha}{\pi} \sim 0.2\%$ suppression by electromagnetic coupling constant

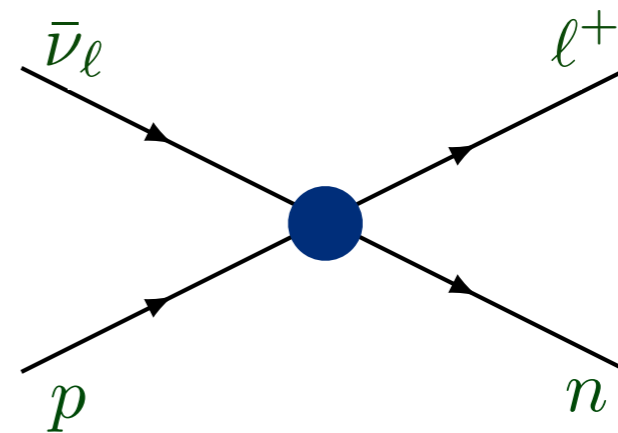
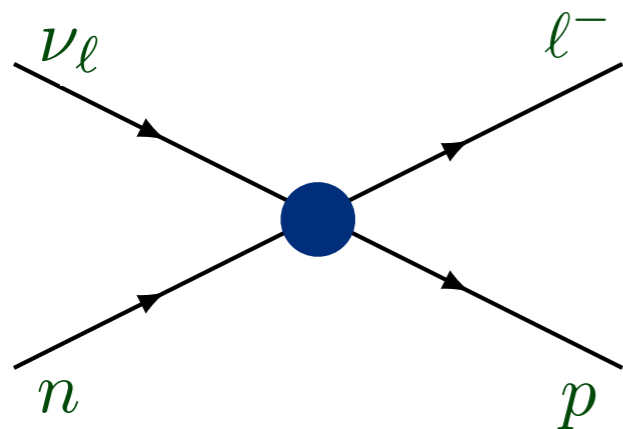
QED corrections

$$m_e \ll m_\mu \ll E_\nu$$



$$\frac{\alpha}{\pi} \sim 0.2 \% \text{ multiplied by } \ln \frac{E_\nu}{m_e} \sim 6 - 10 \text{ or } \ln^2 \frac{E_\nu}{m_e} \sim 36 - 100$$

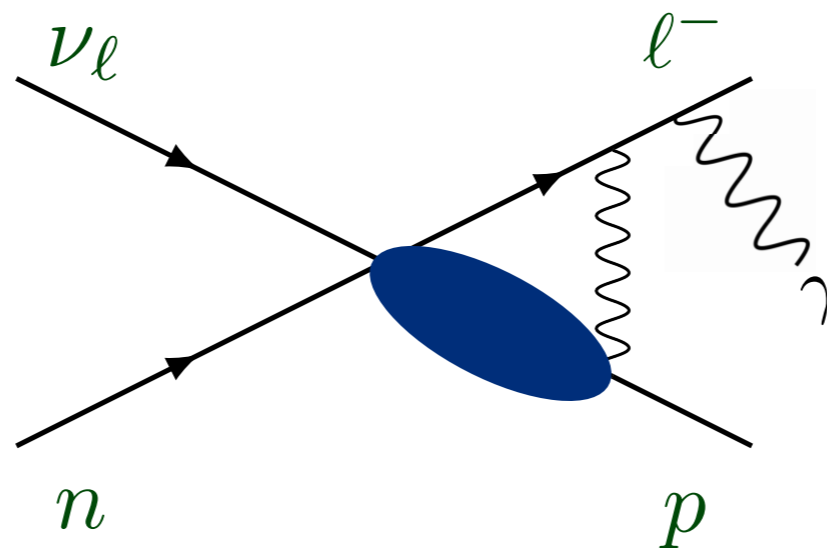
- scale separation introduces large flavor-dependent QED logarithms



O. T., Qing Chen, Richard J. Hill and Kevin S. McFarland, Nature Commun. 13 (2022), 1, 5286

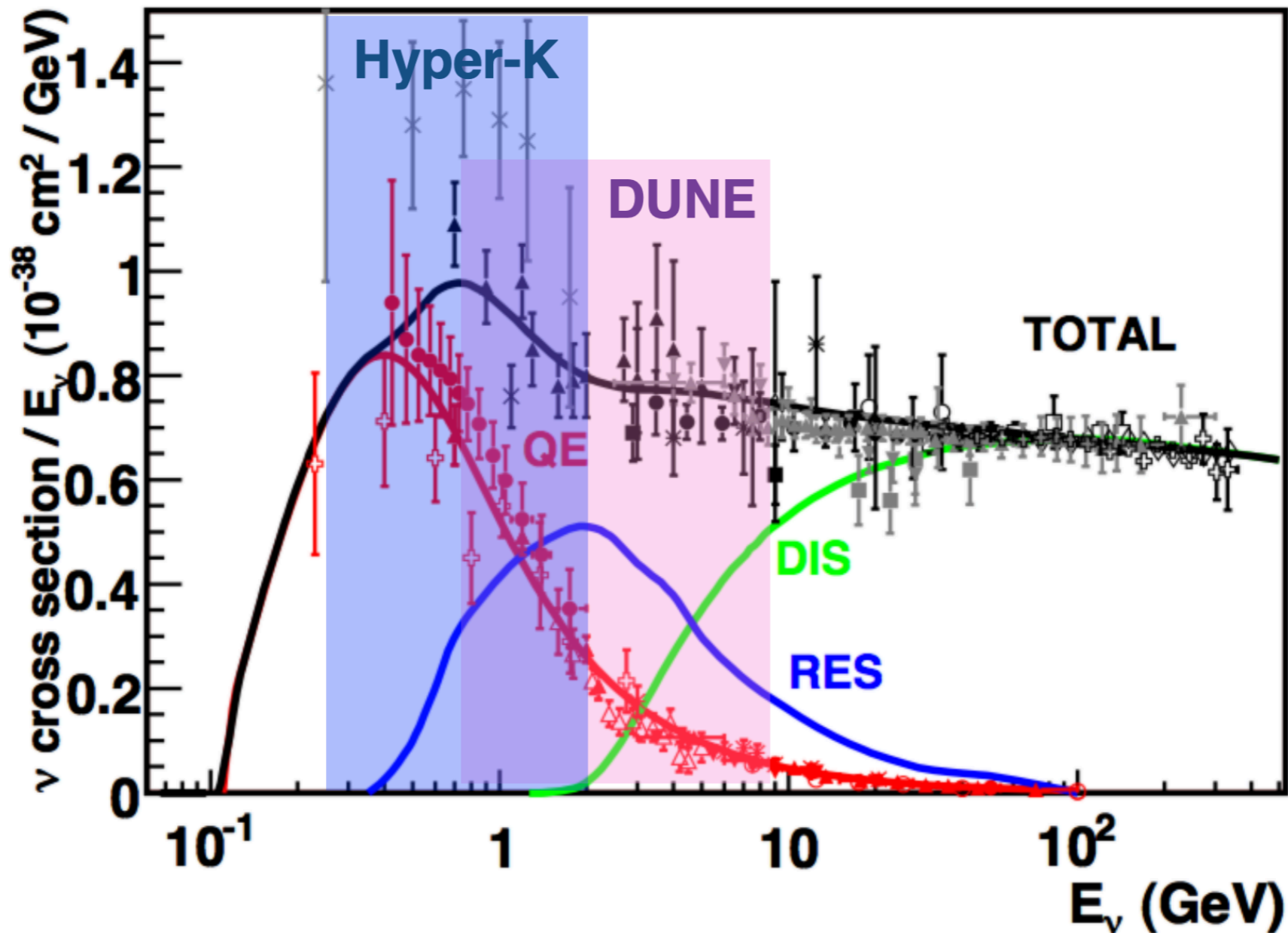
Radiative corrections in charged-current elastic scattering on free nucleons

O. T., Qing Chen, Richard J. Hill, Kevin S. McFarland and Clarence Wret
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CCQE. Why should we care?

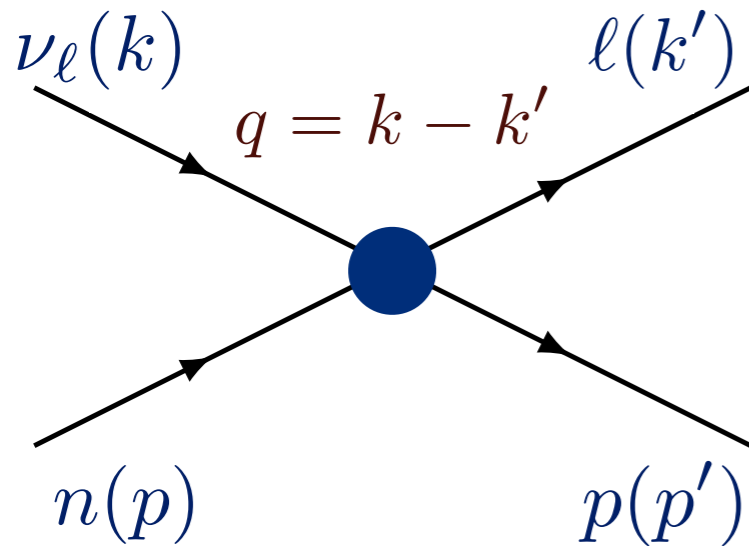
- neutrino-nucleus cross sections and future accelerator-based fluxes



Formaggio
and Zeller
(2013)

- basic process: bulk of events at Hyper-K and DUNE
- channel for reconstruction of neutrino energy

CCQE scattering on free nucleon



neutrino energy

$$E_\nu$$

momentum transfer

$$Q^2 = -q^2$$

contact interaction at GeV energies

- assuming isospin symmetry, nucleon current:

$$\Gamma^\mu(Q^2) = \langle p | \bar{u} (\gamma^\mu - \gamma^\mu \gamma_5) d | n \rangle$$

$$\Gamma^\mu(Q^2) = \gamma^\mu F_D^V(Q^2) + \frac{i\sigma^{\mu\nu} q_\nu}{2M} F_P^V(Q^2) + \gamma^\mu \gamma_5 F_A(Q^2) + \frac{q^\mu}{M} \gamma_5 F_P(Q^2)$$

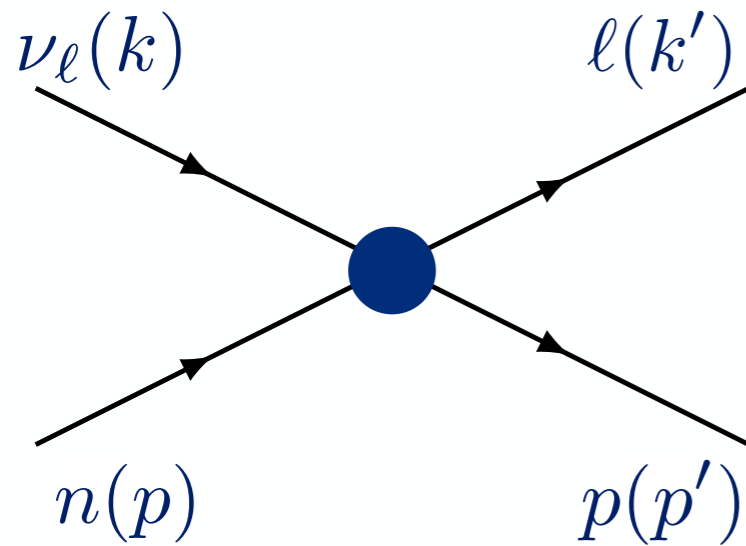
form factors: isovector Dirac and Pauli axial and pseudoscalar

$$F_{D,P}^V = F_{D,P}^p - F_{D,P}^n$$

tree-level amplitude

$$T = \frac{G_F V_{ud}}{\sqrt{2}} (\bar{\ell}(k') \gamma_\mu (1 - \gamma_5) \nu_\ell(k)) (\bar{p}(p') \Gamma^\mu(Q^2) n(p))$$

CCQE scattering on free nucleon



$$\nu = E_\nu/M - \tau - r^2$$

$$r = \frac{m_\ell}{2M} \quad \tau = \frac{Q^2}{4M^2}$$

unpolarized cross section

$$\frac{d\sigma}{dQ^2} \sim \frac{M^2}{E_\nu^2} \left((\tau + r^2) A(Q^2) - \nu B(Q^2) + \frac{\nu^2}{1 + \tau} C(Q^2) \right)$$

Llewellyn Smith (1972)

- structure-dependent functions

$$A = \tau (G_M^V)^2 - (G_E^V)^2 + (1 + \tau) F_A^2 - r^2 \left((G_M^V)^2 + F_A^2 - \underline{4\tau F_P^2 + 4F_A F_P} \right)$$

$$B = \pm 4\tau F_A G_M^V$$

$$C = \tau (G_M^V)^2 + (G_E^V)^2 + (1 + \tau) F_A^2$$

- **pseudoscalar** form factor contribution is suppressed by lepton mass
- cross section is sensitive to both **vector** and **axial** contributions

Elastic scattering on free nucleon

- only 3 experiments performed with deuterium bubble chamber
- direct access to form-factor shape

ANL 1982: 1737 events

BNL 1981: 1138 events

FNAL 1983: 362 events

world data: ~3200 events



Fermilab bubble chamber, Richard Drew

- axial form factor extracted based on electromagnetic structure

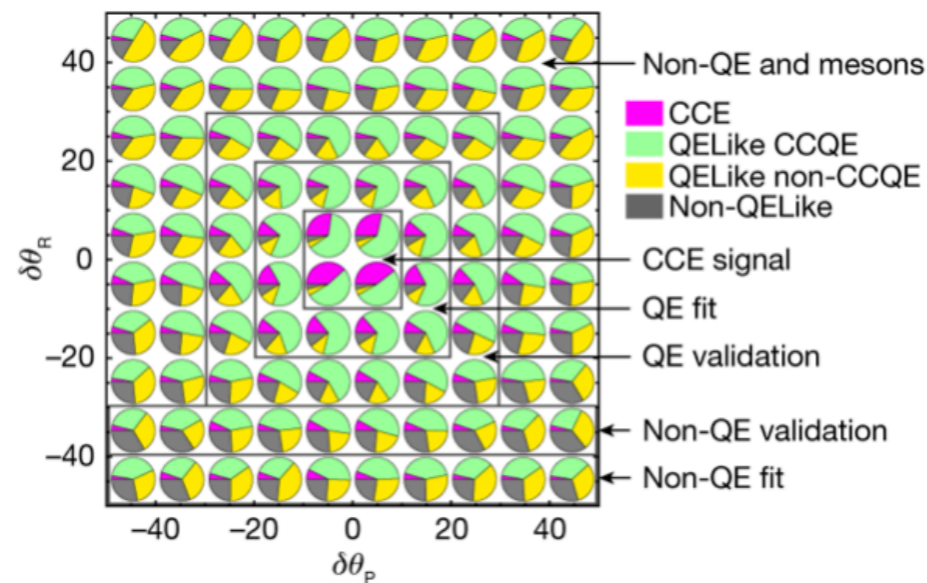
A.S. Meyer, M. Betancourt, R. Gran and R.J. Hill (2016)

MINERvA result with free protons

- idea of scattering on molecular hydrogen realized !!!

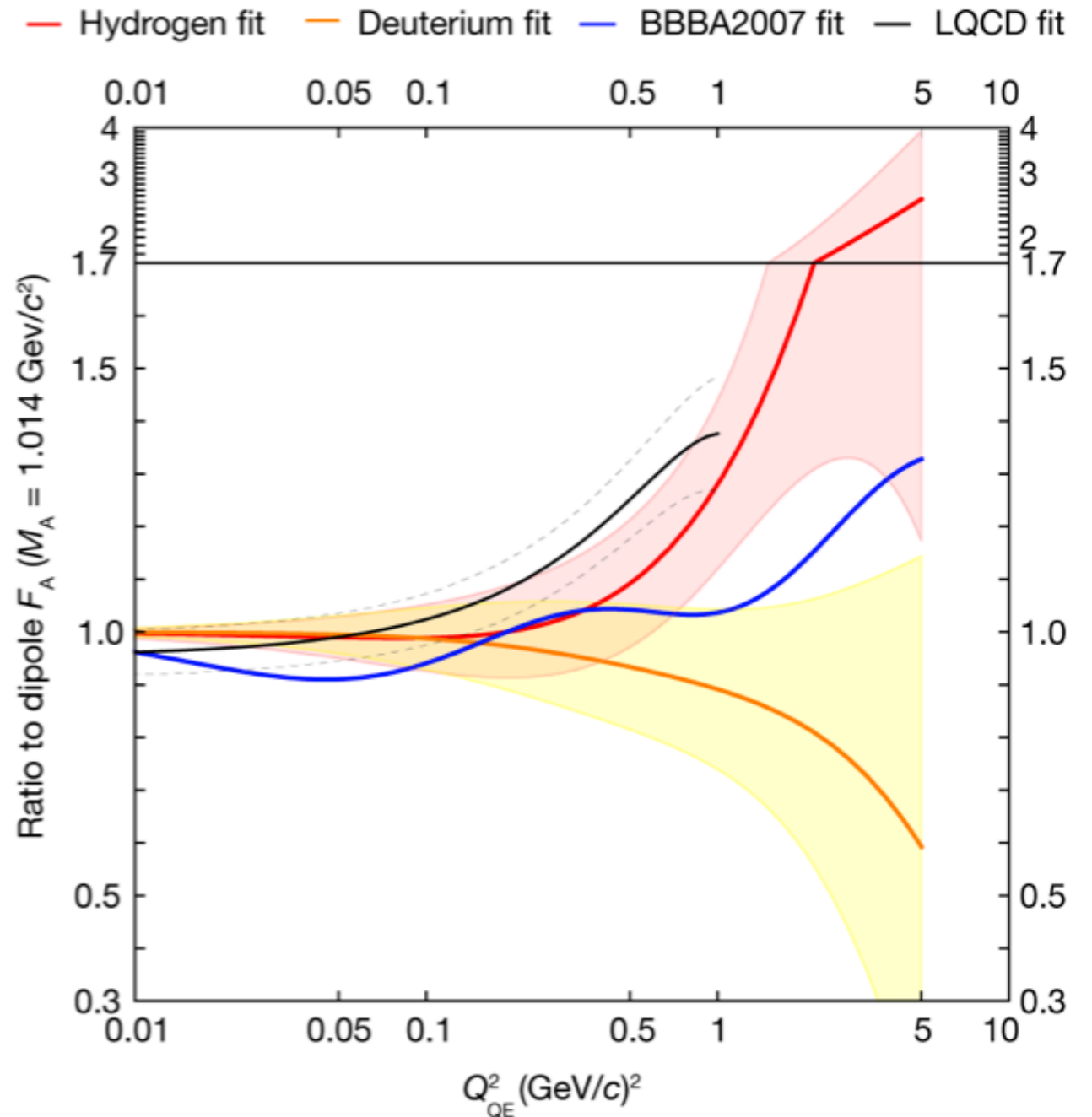


muon kinematic selection



5580 events over
12500 background

background nuclear events
constrained by scattering of ν

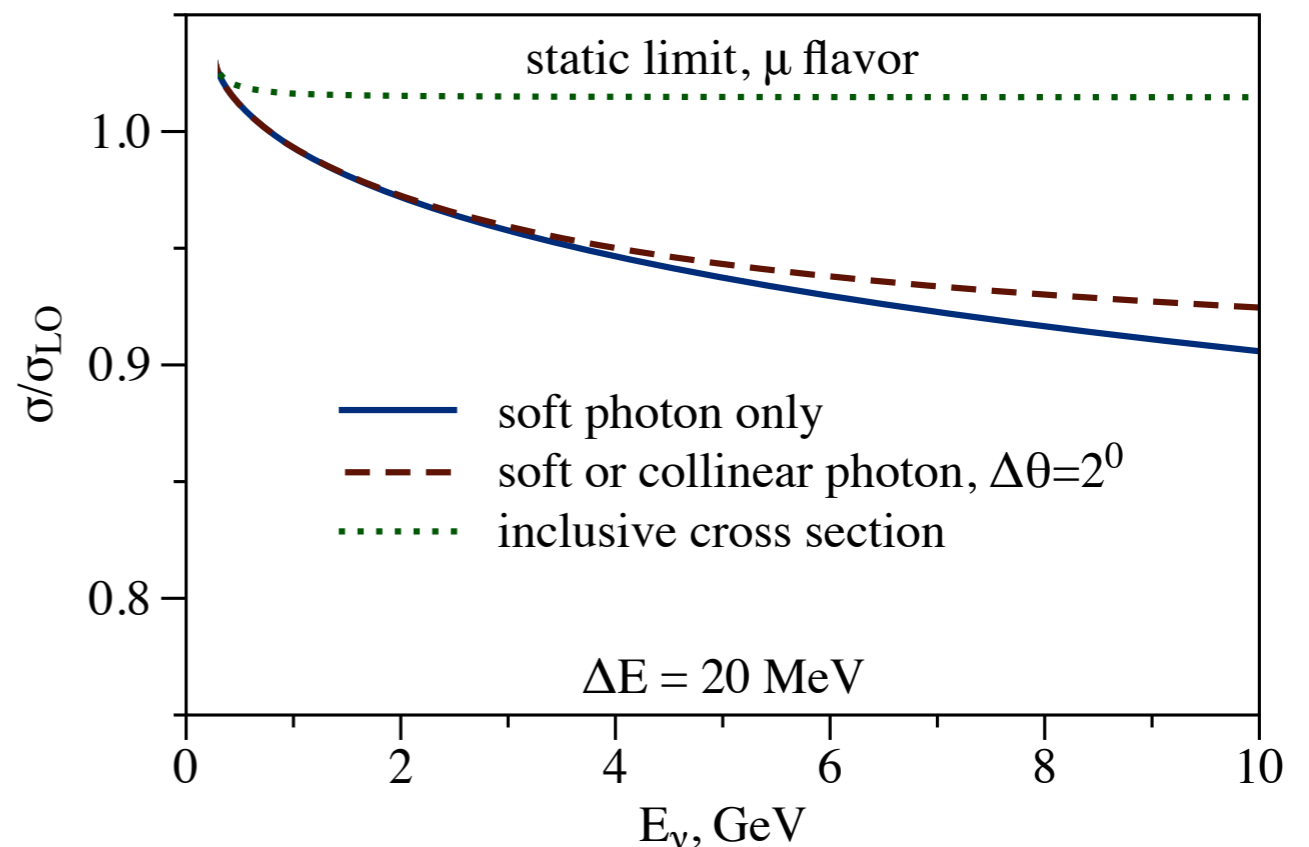
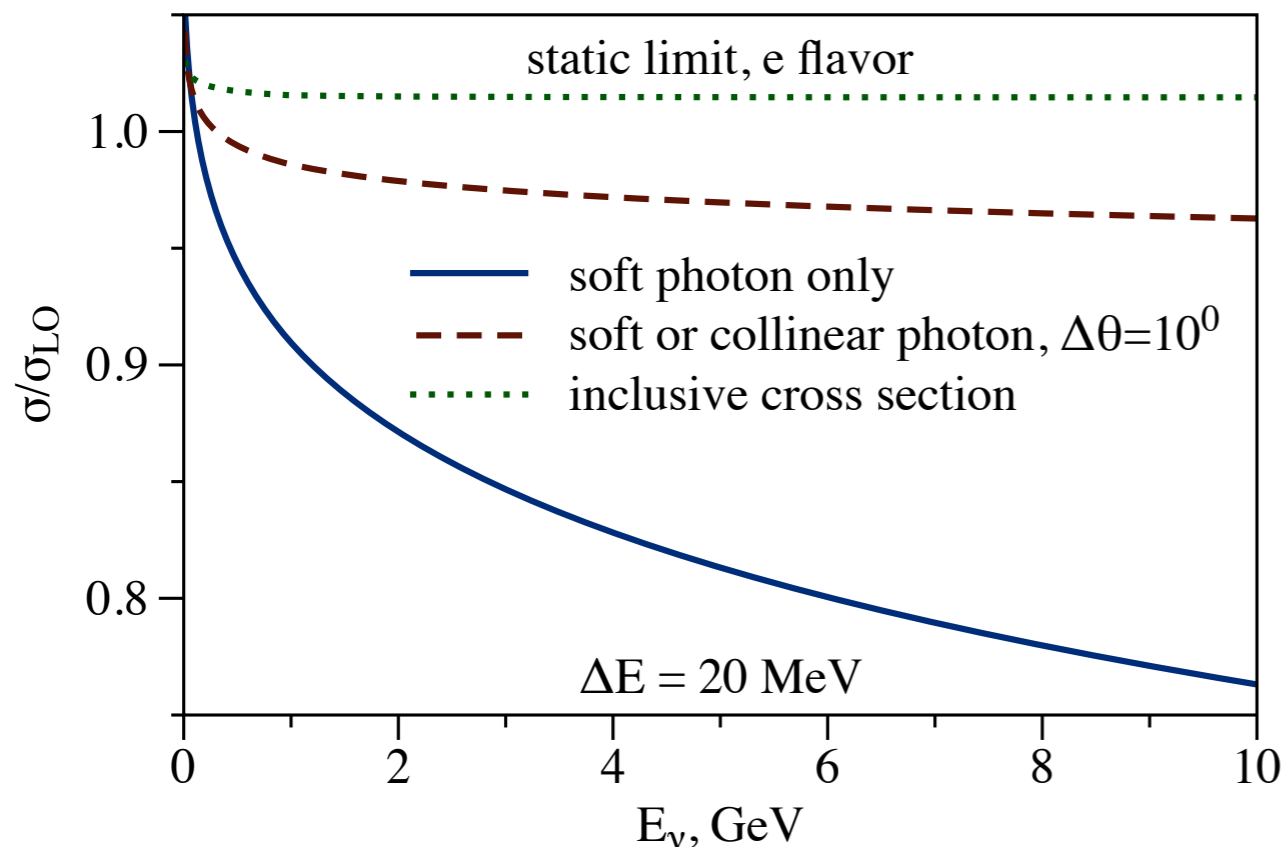


- 1st measurement of axial form factor on “free” protons $\bar{\nu}_\mu p \rightarrow \mu^+ n$

T. Cai et al., MINERvA Collaboration, Nature (2023), 614, 48-53

Static nucleon limit

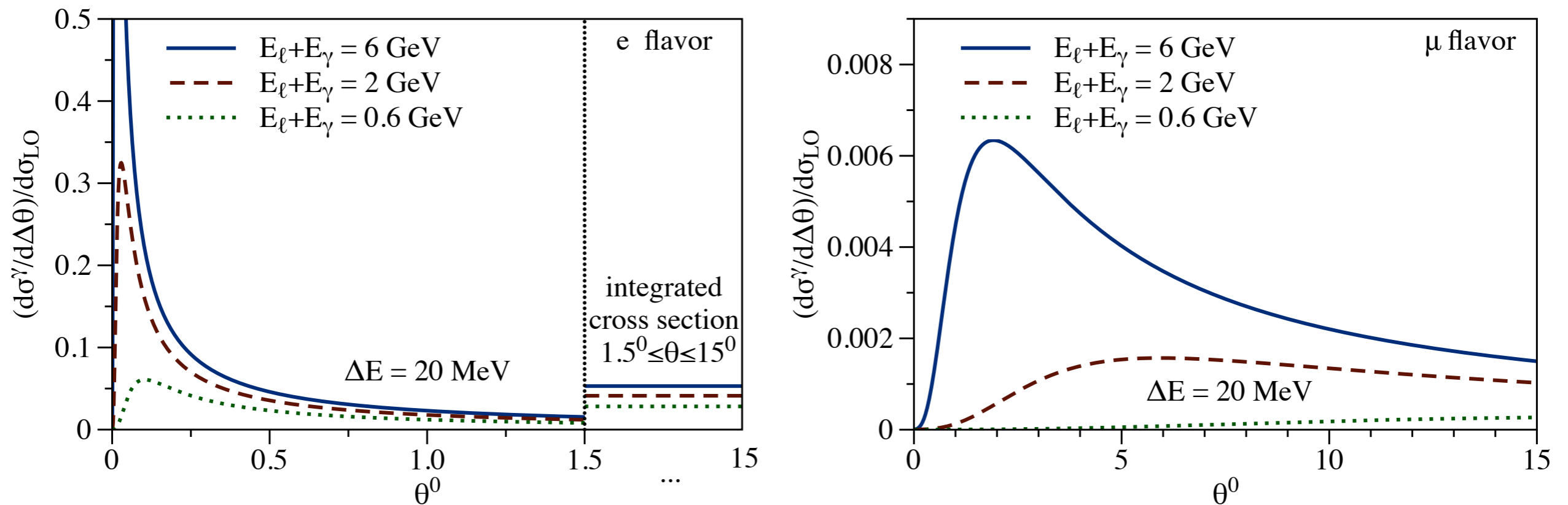
- formal limit of infinitely heavy nucleus $m_\ell \ll E_\ell \ll M$
- provides correct soft and collinear logarithms
- soft-photon energy < 20 MeV, jet size: 10° for electron and 2° for muon



- flavor-dependent effect, same for $\nu_\ell n \rightarrow \ell^- p$ vs $\bar{\nu}_\ell p \rightarrow \ell^+ n$
- collinear observable: cancellation of virtual vs real logs
- inclusive observables (+ γ): few % level, flavor independent

Electron vs muon jets

- factorization for radiation of collinear photons
- cone angle is defined to lepton direction
- photons of energy > 20 MeV, fixed energy in the cone



- flavor-dependent effect, same for $\nu_\ell n \rightarrow \ell^- p$ vs $\bar{\nu}_\ell p \rightarrow \ell^+ n$
- forward-peaked radiation for electron flavor
- negligible radiation for muons with shifted peak position

Factorization approach

- cross section is given by **factorization formula**

$$d\sigma \sim S \left(\frac{\Delta E}{\mu} \right) J \left(\frac{m_\ell}{\mu} \right) H \left(\frac{M}{\mu} \right)$$

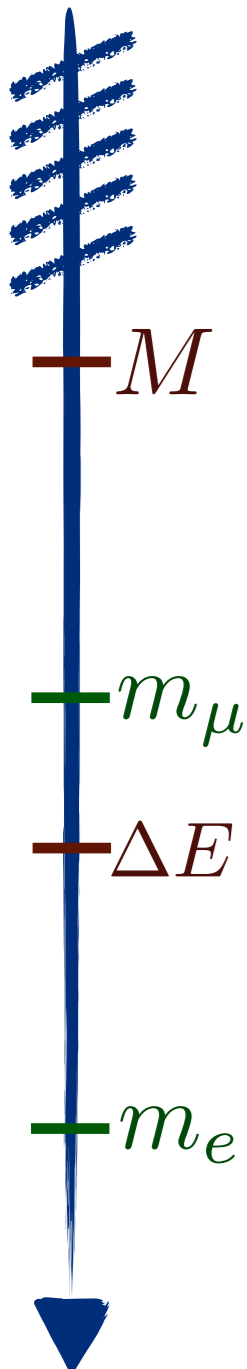
- determine **hard function** at hard scale by matching experiment or **hadronic model** to the theory with heavy nucleon

- **soft and collinear functions** are evaluated **perturbatively**

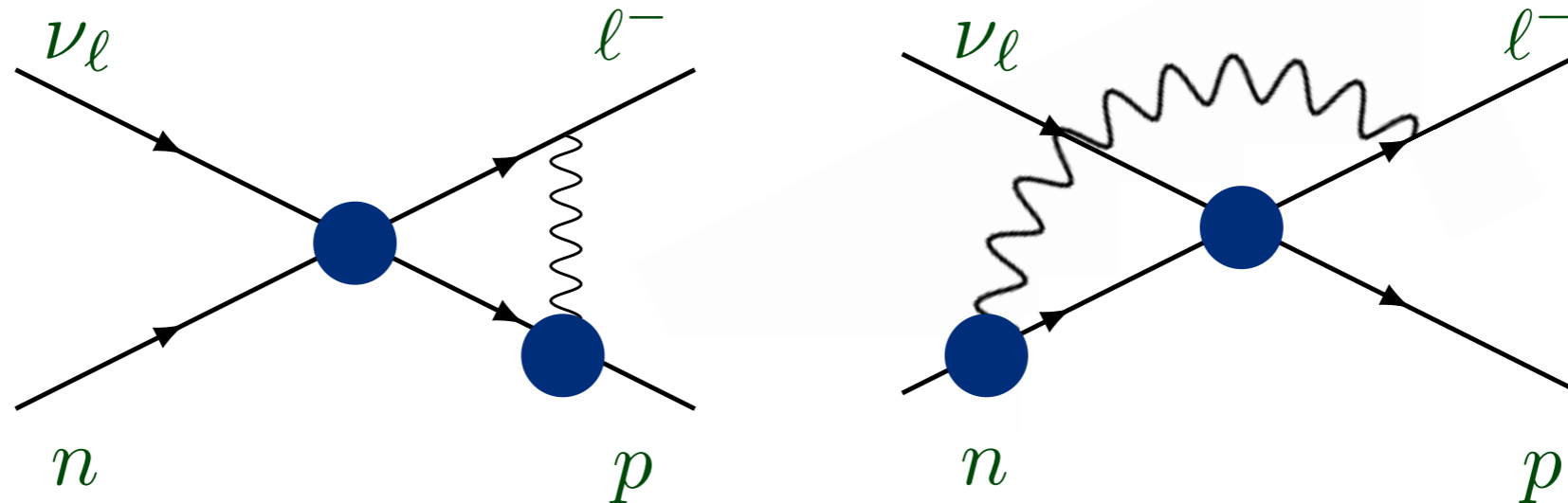
SCET power expansion parameter

$$\lambda \sim \frac{m_\mu^2}{E_\nu^2} \sim (\Delta\theta)^2 \sim \frac{\Delta E}{E_\nu}$$

$\ln \lambda$ enhancements



Hadronic model at GeV scale



- exchange of photon between the charged lepton and nucleons
- assume **onshell form** for each interaction with dipole form factors
discussed for neutrino-nucleon scattering: Graczyk (2013)
- add **self energy** for charged particles
- reproduce soft and collinear regions of SCET

- best determination of hard function

Factorization approach

- cross section is given by **factorization formula**

$$d\sigma \sim S \left(\frac{\Delta E}{\mu} \right) J \left(\frac{m_\ell}{\mu} \right) H \left(\frac{M}{\mu} \right)$$

- determine **hard function** at hard scale by matching experiment or **hadronic model** to the theory with heavy nucleon

- **RGE evolution** of the hard function to scales $\Delta E, m_\ell$

- **soft and collinear functions** are evaluated **perturbatively**

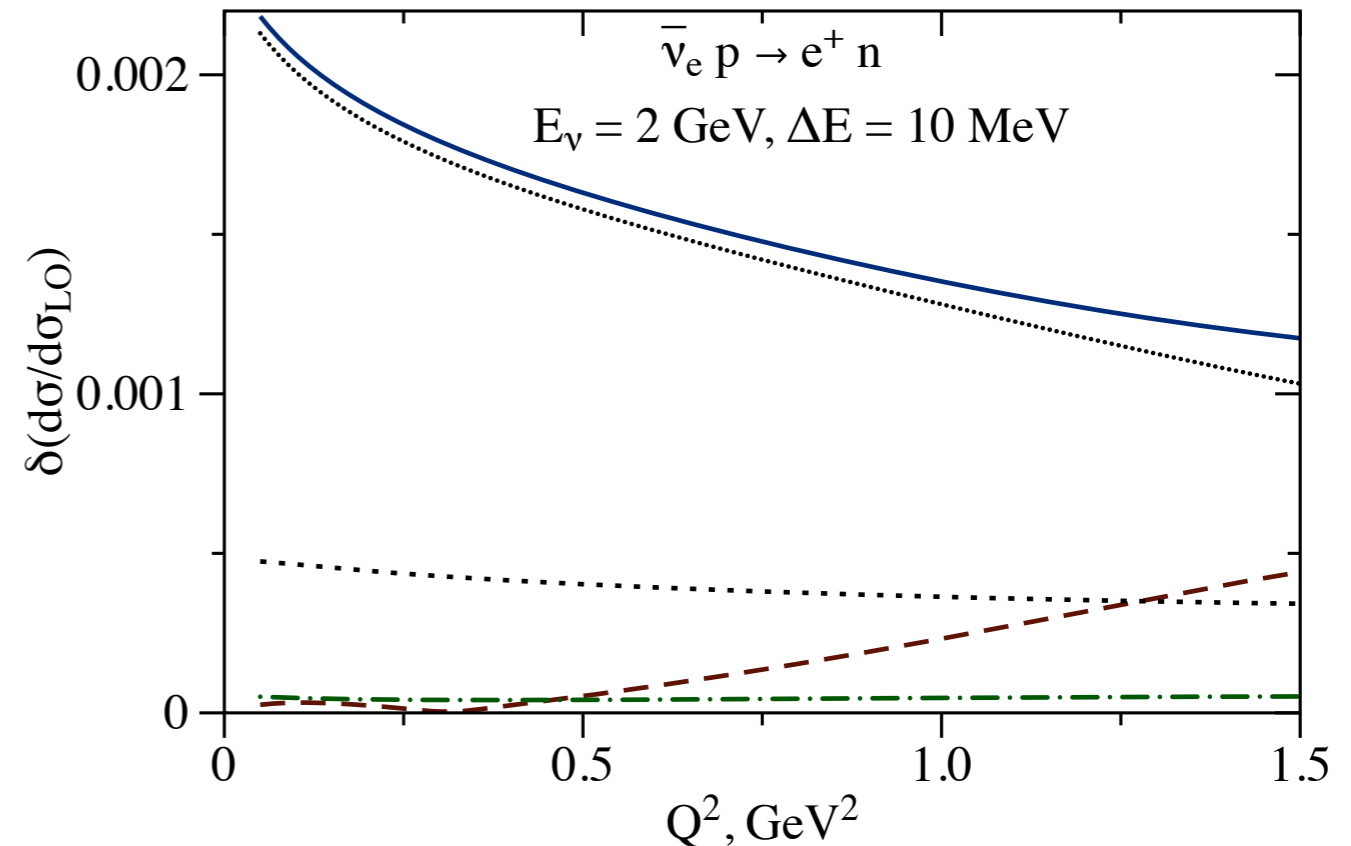
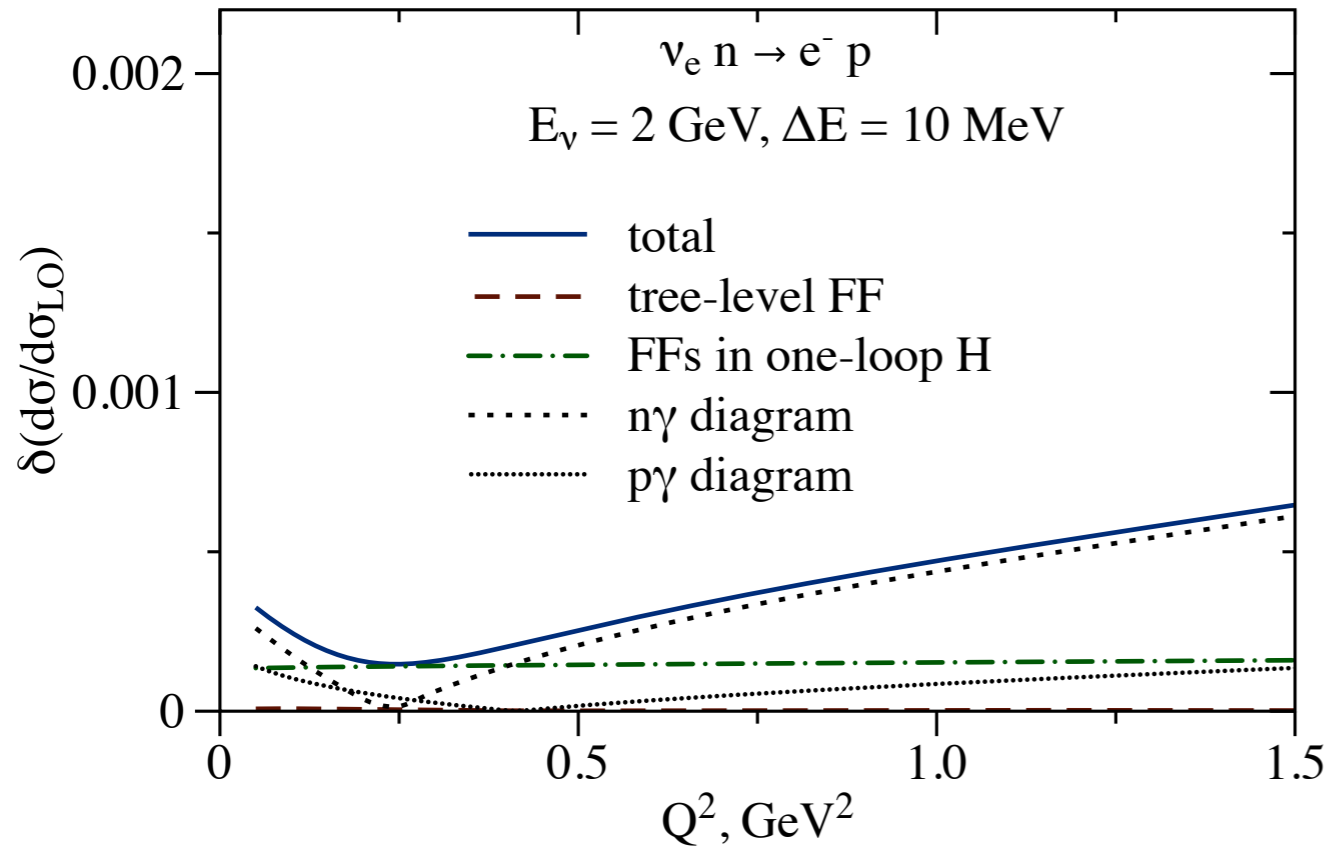
- calculate cross section at low energies accounting for **all large logs**

ep scattering with soft radiation only: Richard J. Hill (2016)

- **soft and collinear functions** determined **analytically**
- **hard function** describes physics at GeV energies

Error budget

- uncertainties from hard function



Meyer, Betancourt, Gran and Hill (2016)

- nucleon form factors

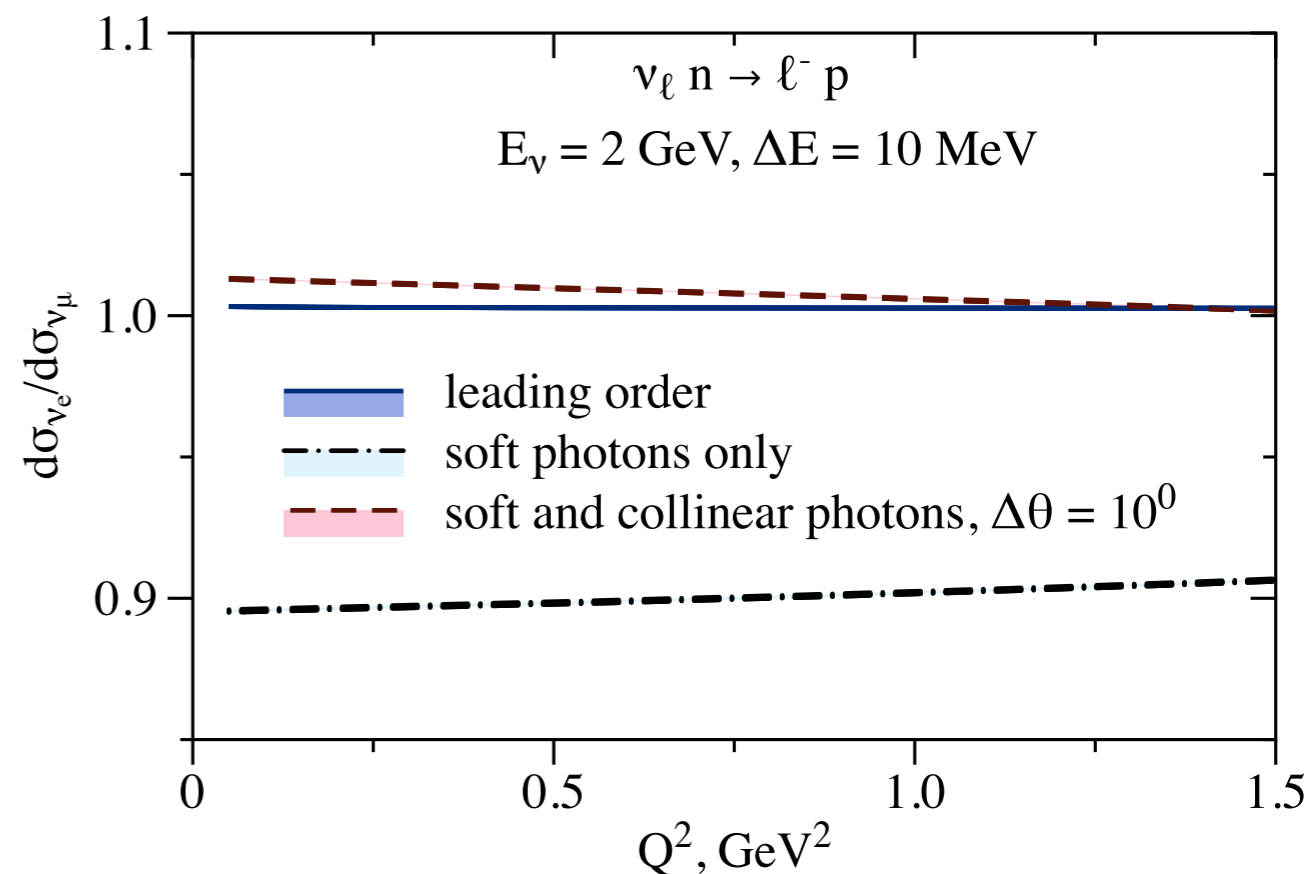
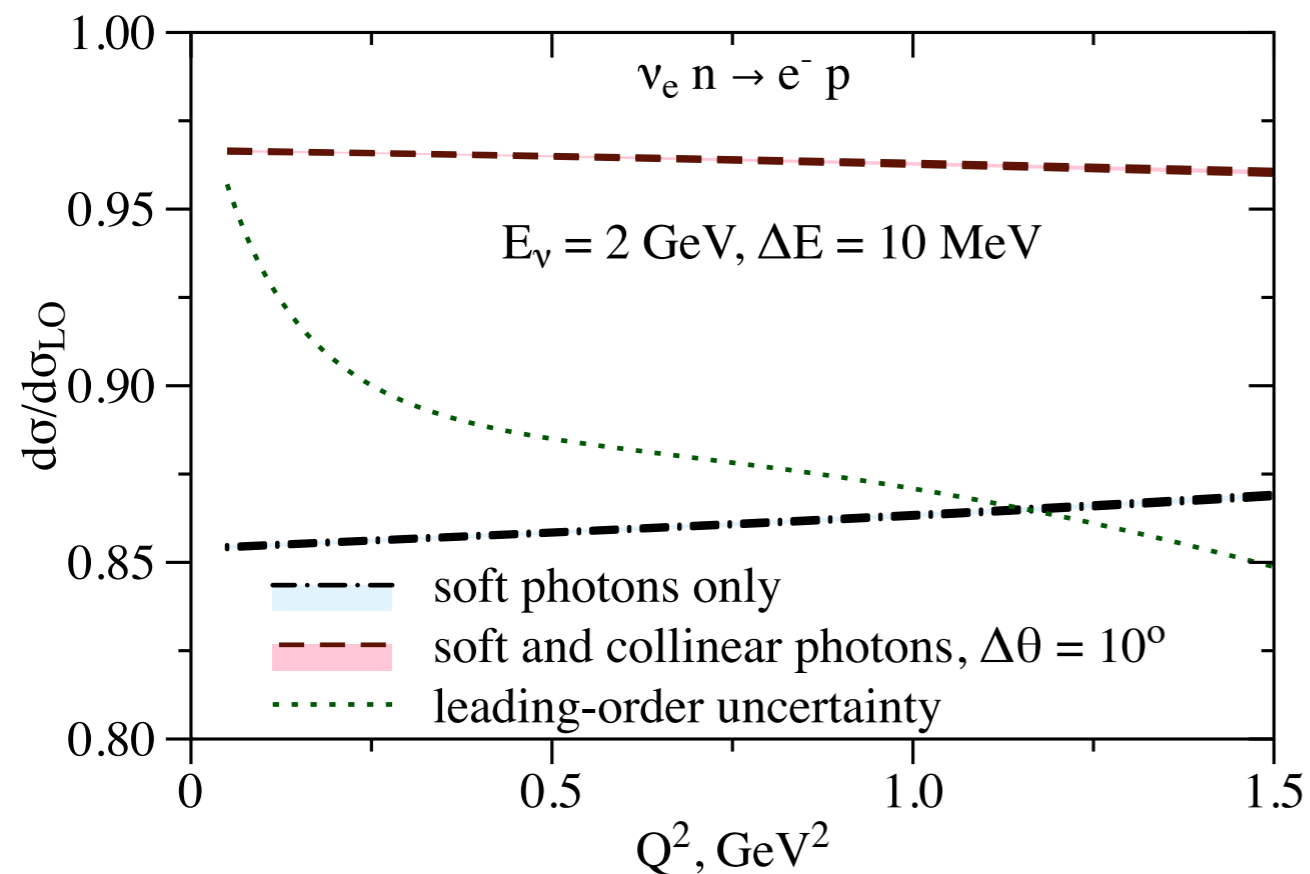
Kaushik Borah, Gabriel Lee, Richard J. Hill and O.T. (2020)

- add perturbative uncertainty by variation of scale

- uncertainty of permille level for the ratio to LO result

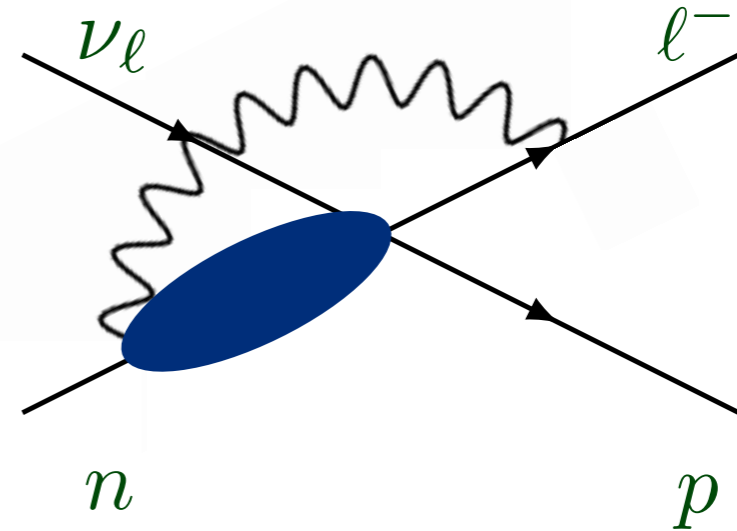
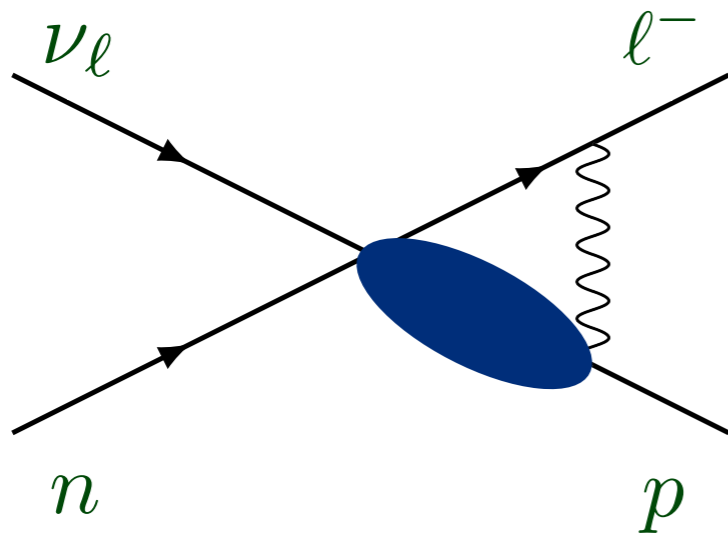
Exclusive observables

- cancellation of uncertainties from hard function for e/μ and ratio to LO

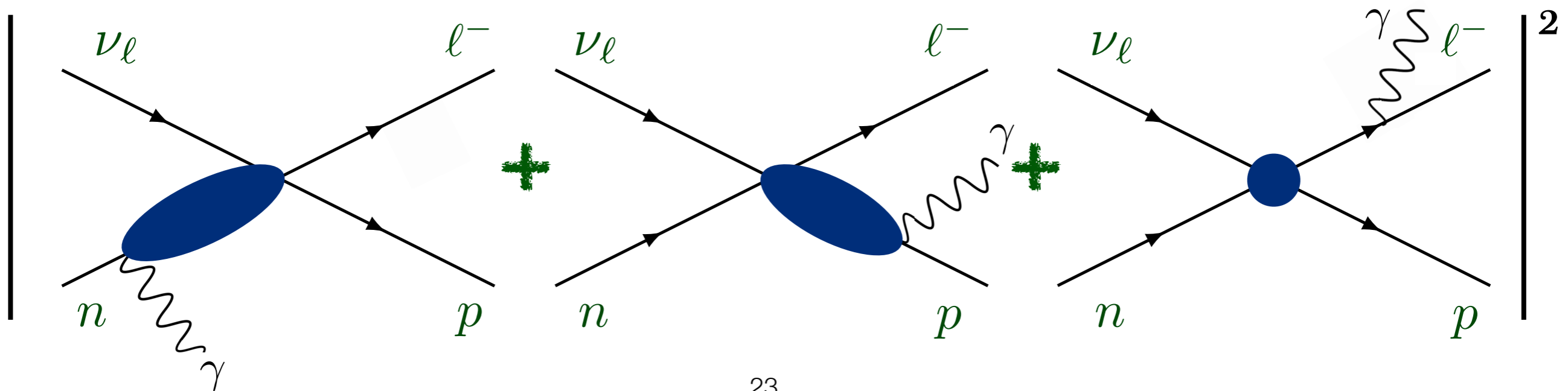


- ratios: cancellation of uncertainty from hard function

Inclusive observables

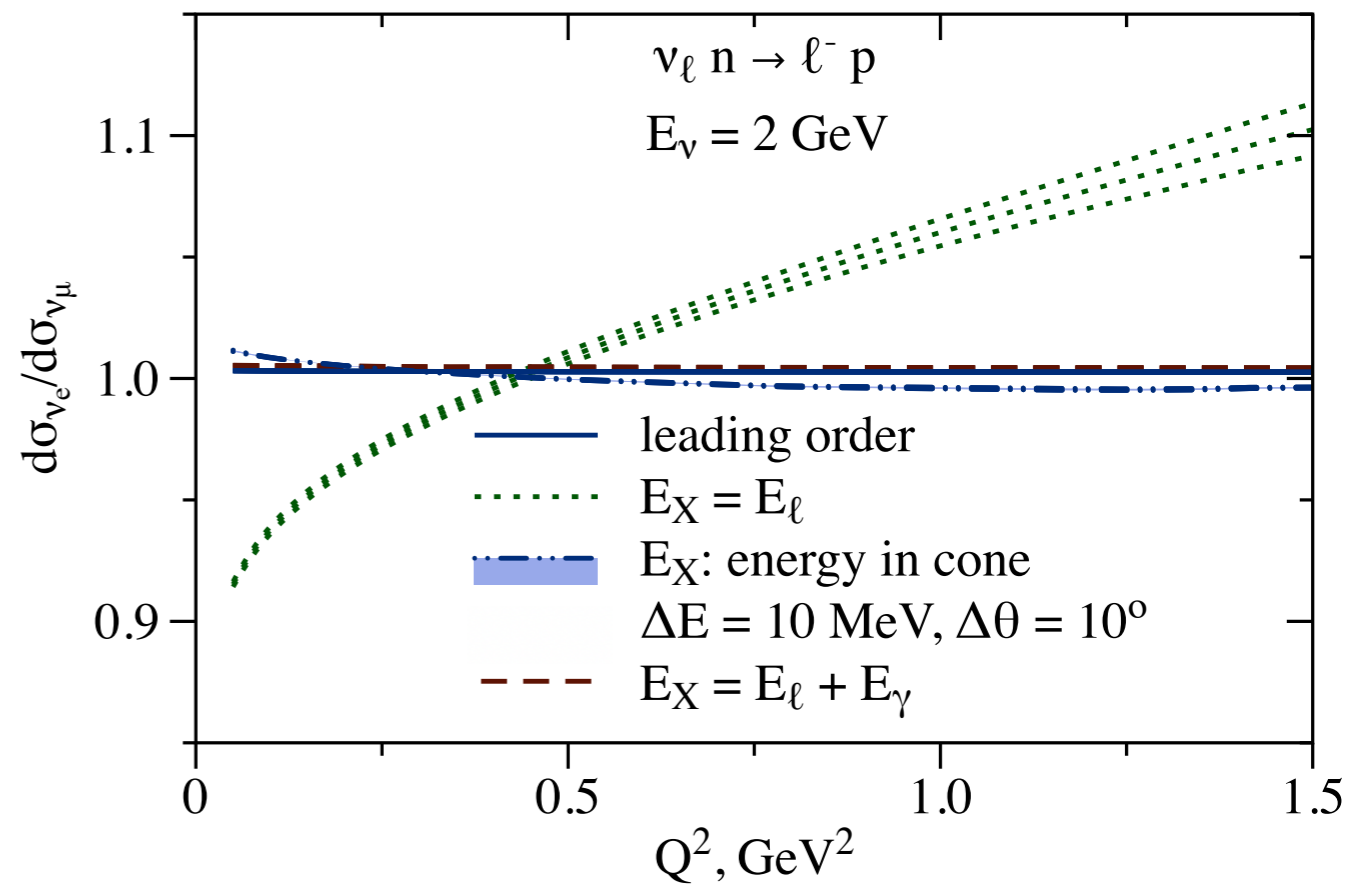
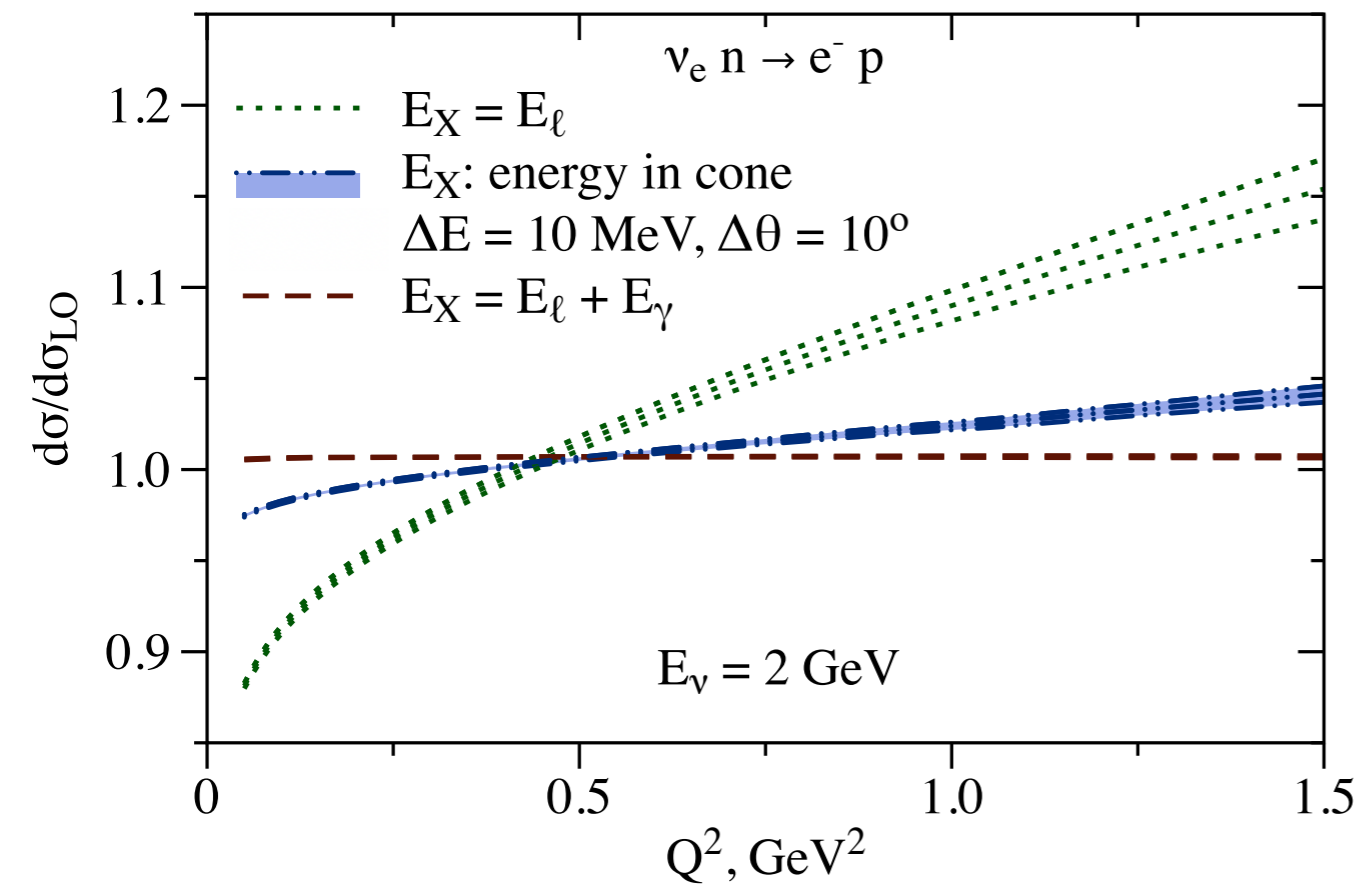


- the same gauge-invariant model for the real radiation
- arbitrary hard photons are part of the observable



Inclusive observables

- kinematics $Q^2 = 2M(E_\nu - E_X)$ is reconstructed with 3 different E_X



- dependence on reconstruction of kinematics and cuts
- predict σ_{ν_e} from σ_{ν_μ} measurements with neutrino beam



Electron/muon ratio

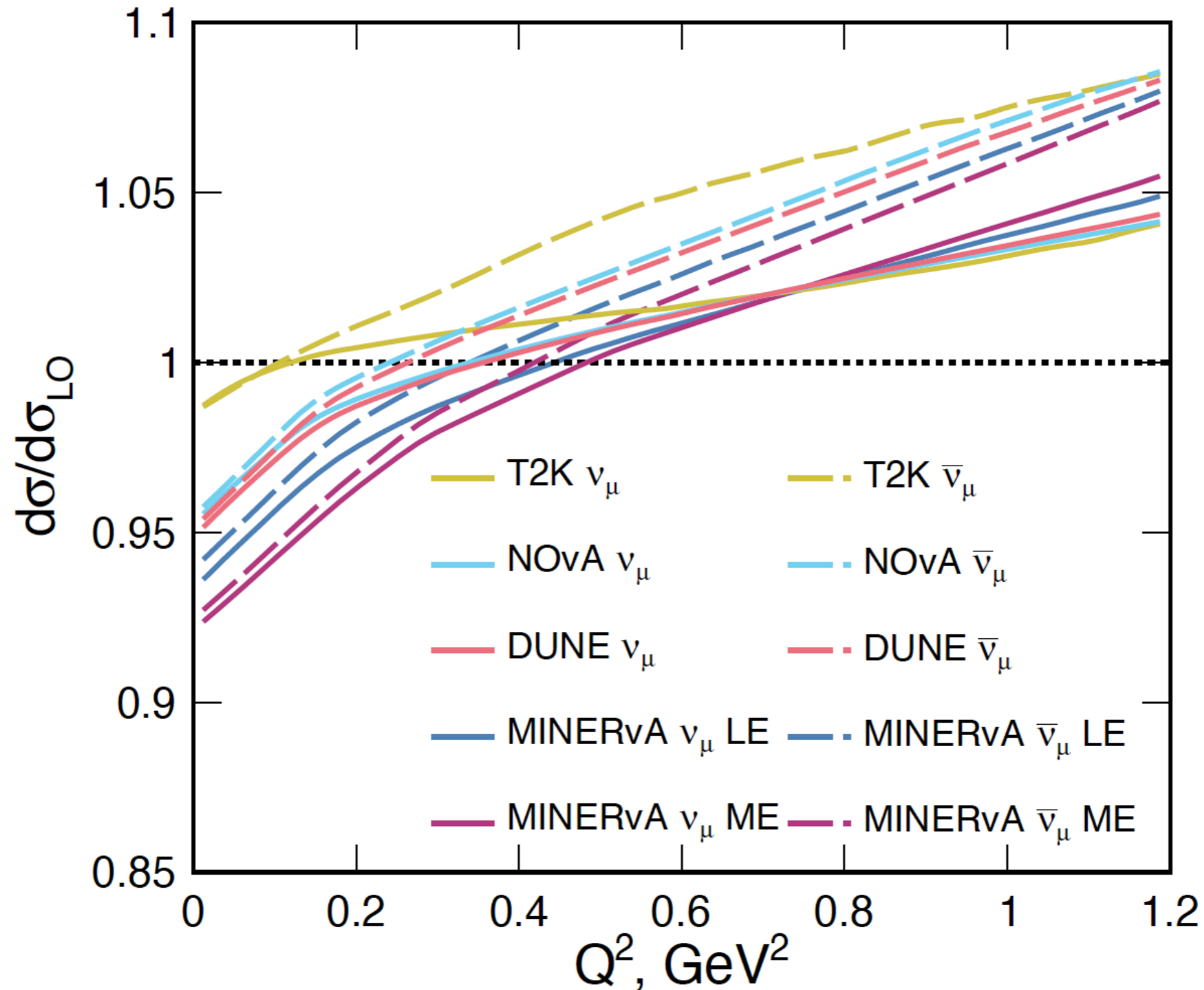
	E_ν , GeV		$\left(\frac{\sigma_e}{\sigma_\mu} - 1\right)_{\text{LO}}$, %	$\frac{\sigma_e}{\sigma_\mu} - 1$, %
T2K/HyperK	0.6	ν	2.47 ± 0.06	$2.84 \pm 0.06 \pm 0.37$
		$\bar{\nu}$	2.04 ± 0.08	$1.84 \pm 0.08 \pm 0.20$
NOvA/DUNE	2.0	ν	0.322 ± 0.006	$0.54 \pm 0.01 \pm 0.22$
		$\bar{\nu}$	0.394 ± 0.003	$0.20 \pm 0.01 \pm 0.19$

TABLE II: Inclusive electron-to-muon cross-section ratios for neutrinos and antineutrinos without kinematic cuts. Uncertainties at leading order are from vector and axial nucleon form factors. For the final result, we include an additional hadronic uncertainty from the one-loop correction to the first uncertainty, and provide a second uncertainty as the magnitude of the radiative correction.

$$\frac{\sigma(m_\ell \rightarrow 0)}{\sigma(m_\ell = 0)} \approx 1 + Am_\ell^2 + \alpha Bm_\ell^2 \ln m_\ell$$

- inclusive cross sections and flavor ratios determined by KLN
- nuclear effects: suppressed by expansion parameters squared

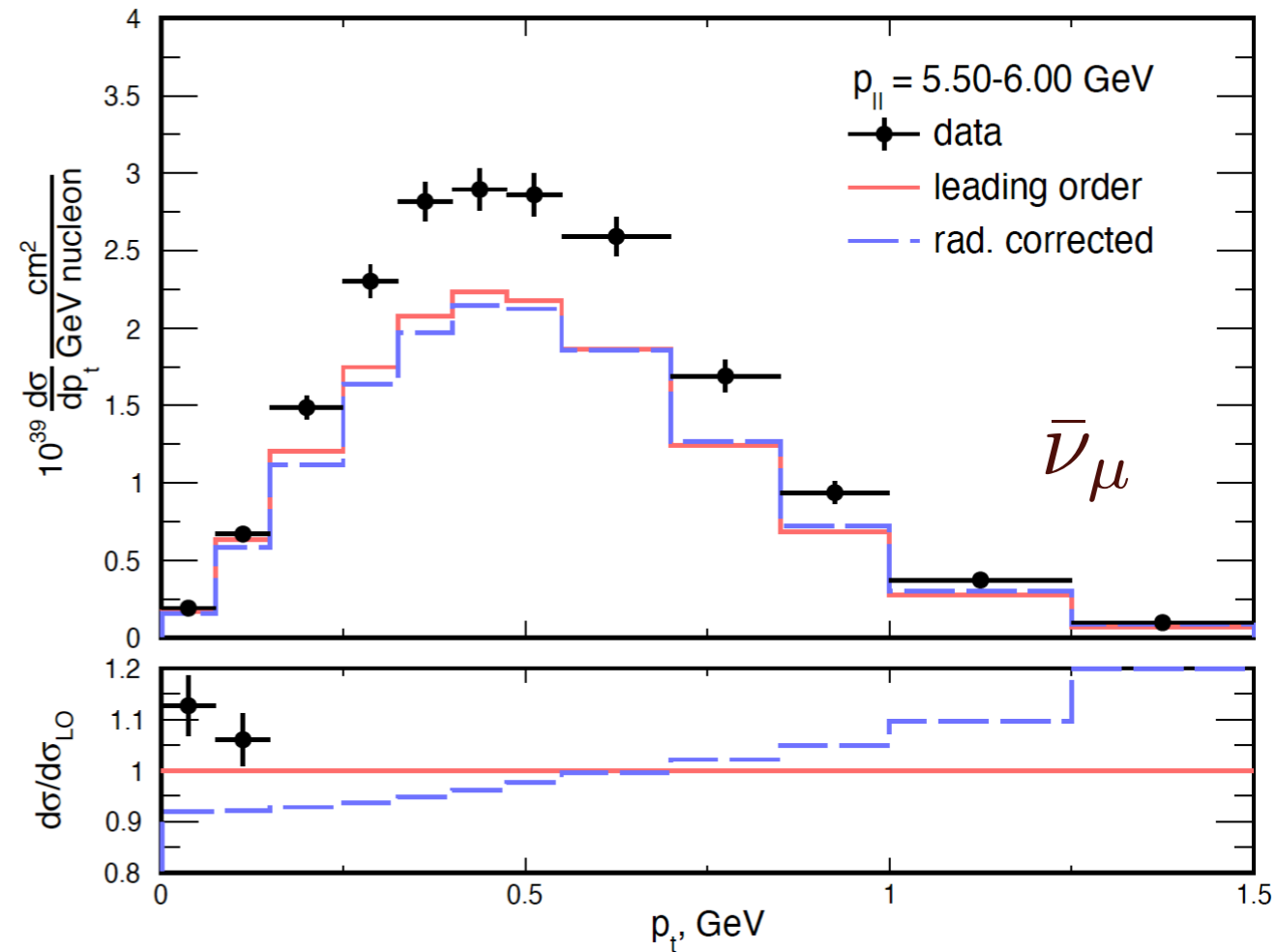
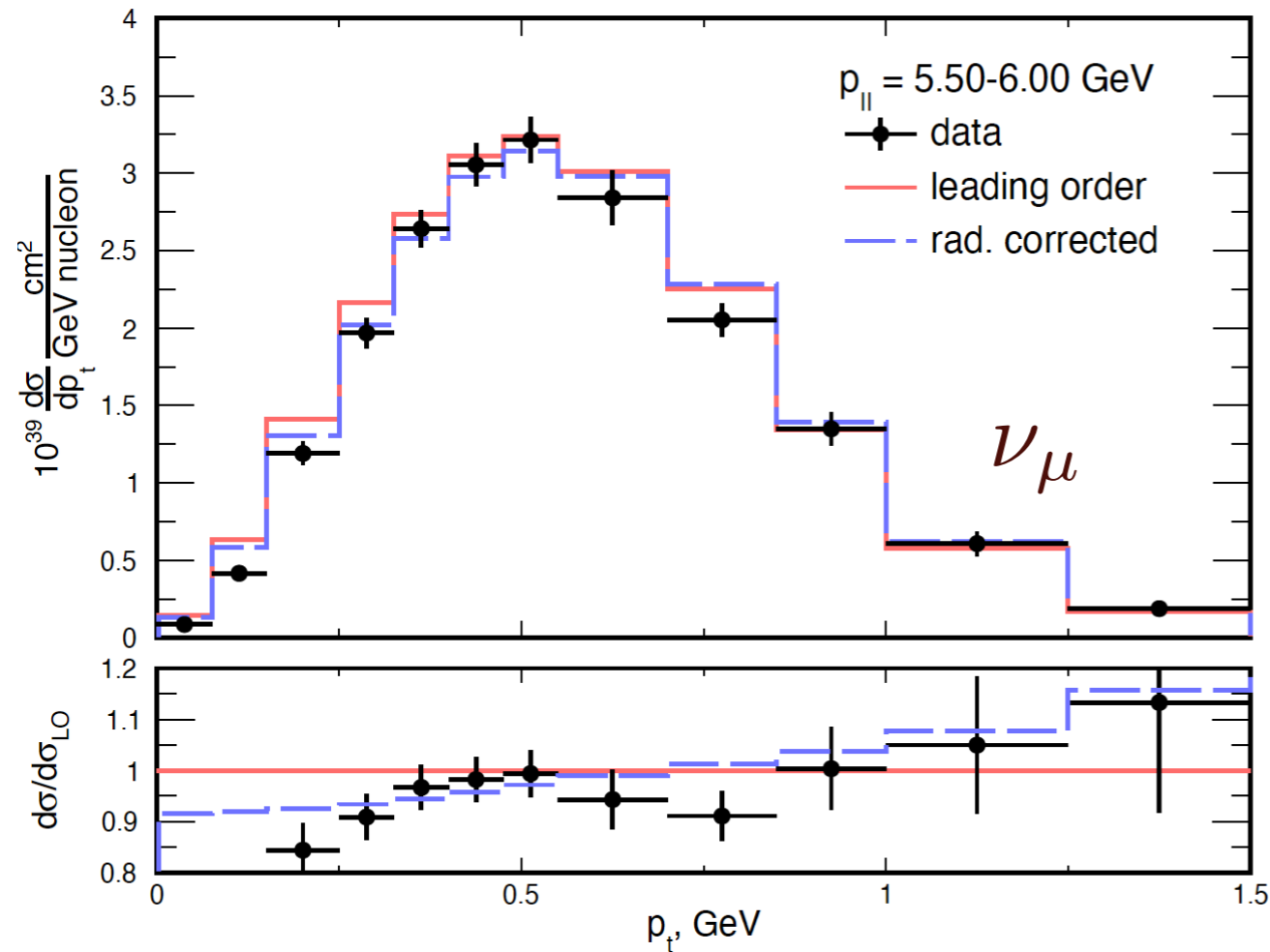
Comparison to data



- lepton energy spectra with lepton kinematics
- NEUT generator + flux averaging

Comparison to data

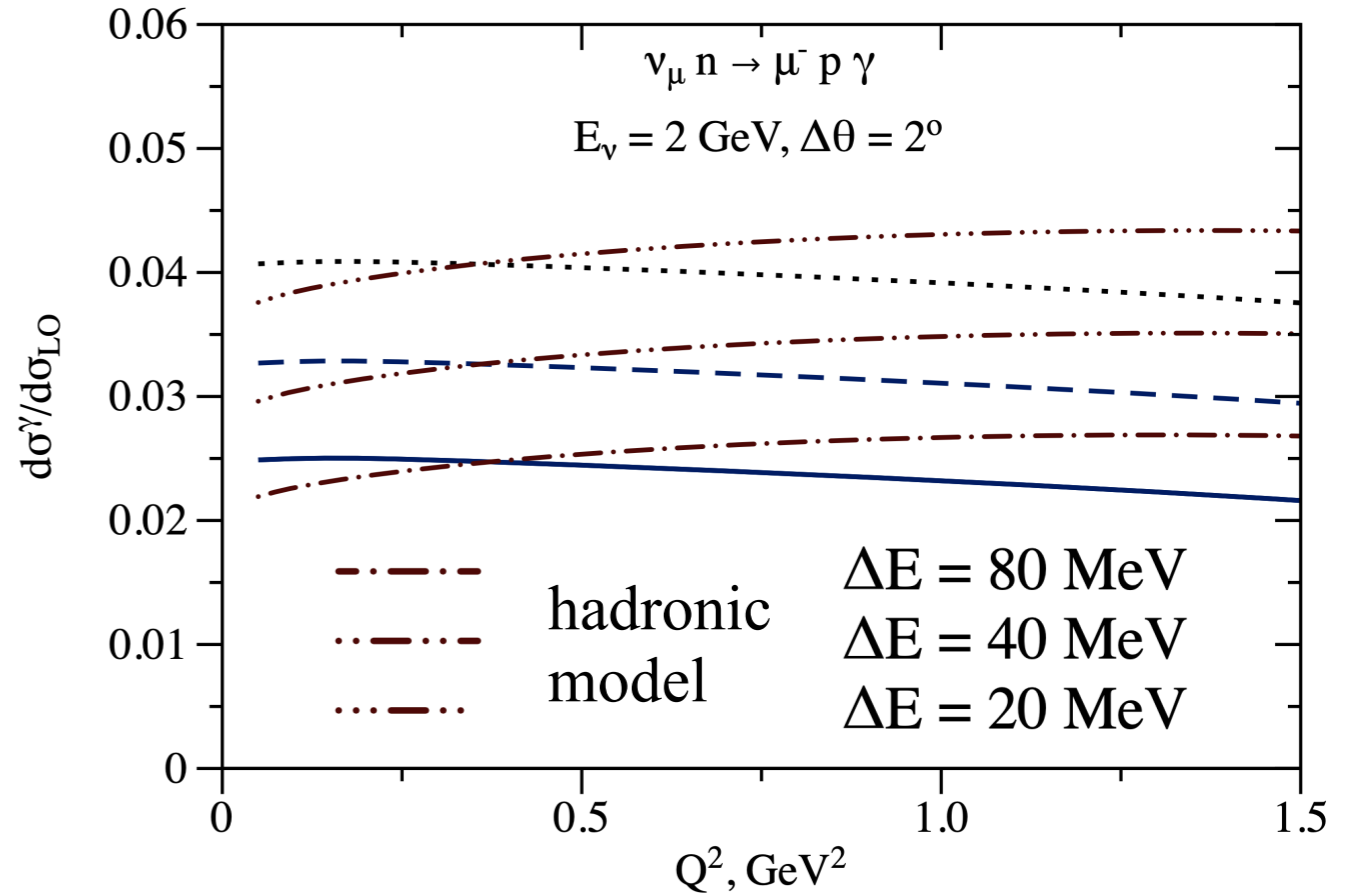
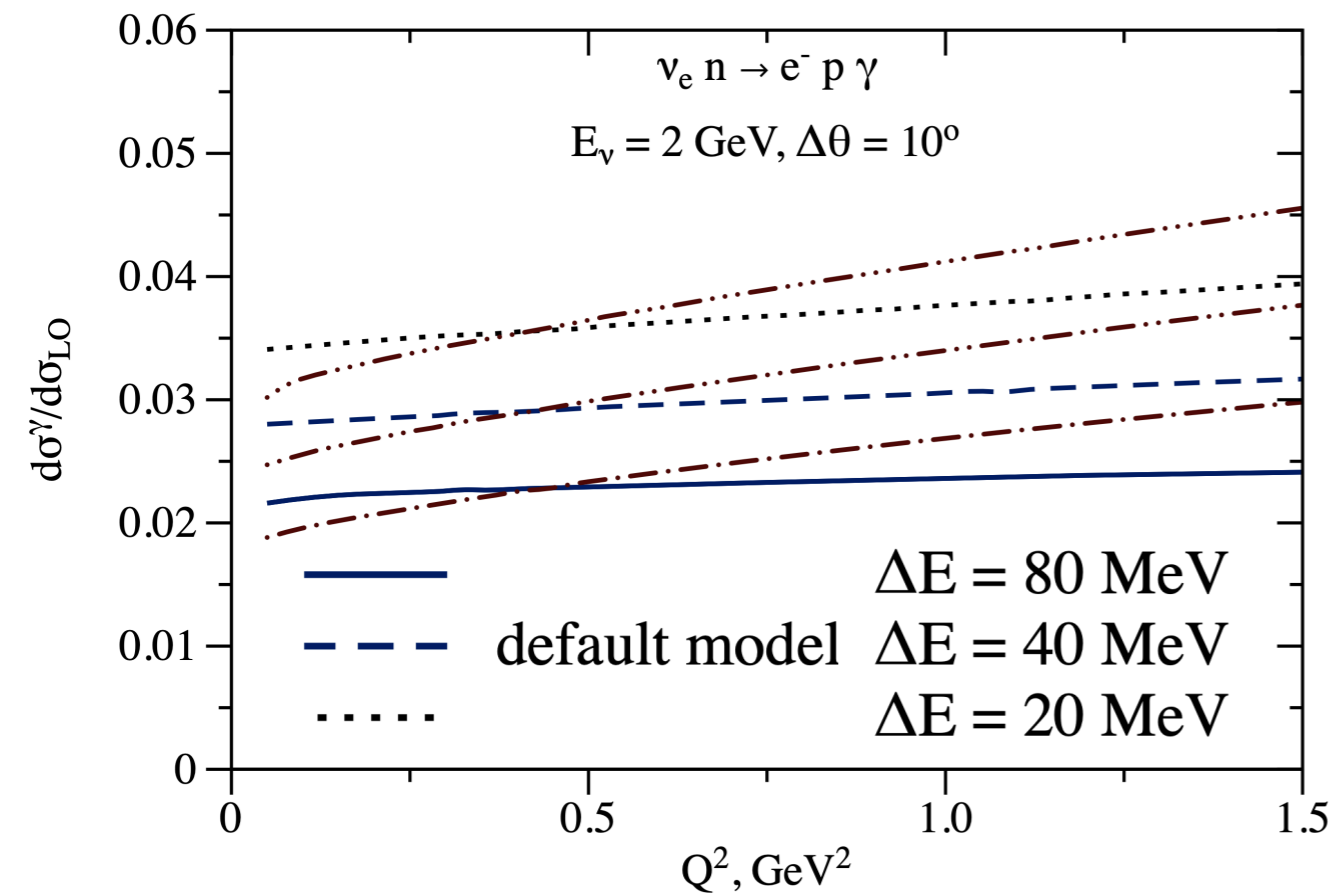
- medium-energy flux data from MINERvA@FERMILAB



- electron flavor: measurements are uncertain
- muon flavor: comparable to experimental precision

Radiation of hard photons

- model-dependent description for radiation of hard photons



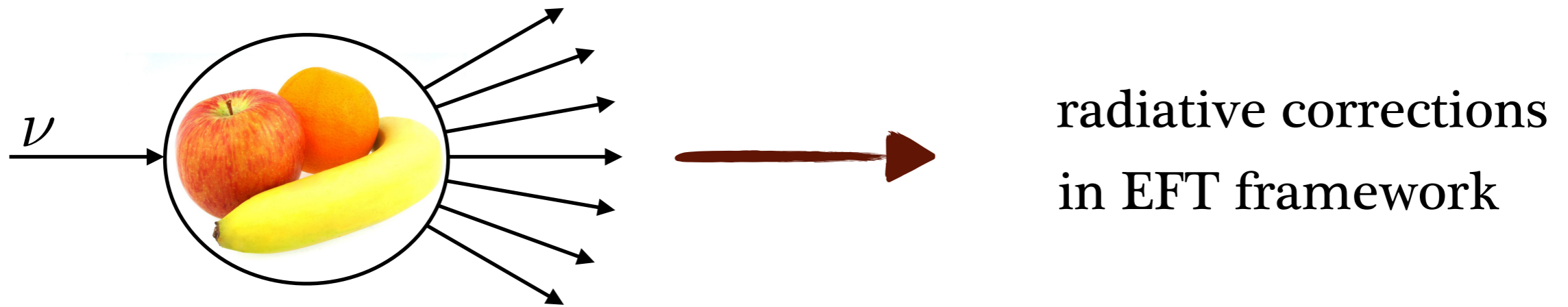
“Blunden calculation”

- photon energies are above 20, 40, and 80 MeV: default vs “SIFF”

“hadronic model”

- % -level radiation of non-collinear hard photons
- 10^{-4} flavor misidentification rate for NO ν A&T2K kinematics

Conclusions



- radiative corrections to neutrino-nucleon cross sections formulated in factorization framework
- charged-current elastic electron vs muon cross-section ratios evaluated from theory with sub-percent uncertainty

Thanks for your attention !!!