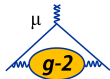


# Status of the FNAL muon $g - 2$ experiment

Alec Tewsley-Booth

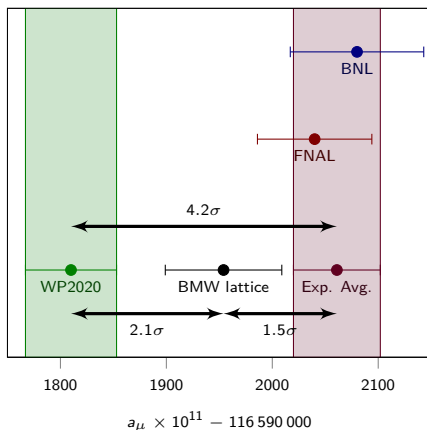
University of Kentucky, Lexington

May 11, 2023



# Spoilers (460 ppb Uncertainty)

- Fermilab measurement agrees with Brookhaven; reasonable to combine statistics-dominated measurements
- World average is  $4.2\sigma$  from the Standard Model prediction (2020 white paper)
- BMW lattice results is  $1.5\sigma$  from the world average



# The history of the muon $g$ -factor

- 1928, Dirac equation,  $g = 2$



# The history of the muon $g$ -factor

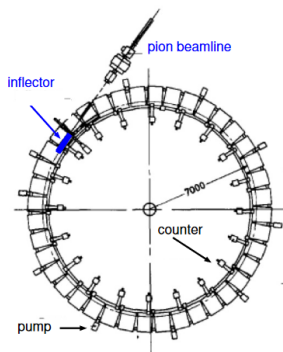
- 1928, Dirac equation,  $g = 2$
- 1947, Schwinger term,  
 $a_l = \frac{\alpha}{2\pi}$



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- 1928, Dirac equation,  $g = 2$
- 1947, Schwinger term,  

$$a_l = \frac{\alpha}{2\pi}$$
- 1962, CERN muon  $g$ -factor experiments



# The history of the muon $g$ -factor

- 1928, Dirac equation,  $g = 2$
- 1947, Schwinger term,  

$$a_l = \frac{\alpha}{2\pi}$$
- 1962, CERN muon  $g$ -factor experiments
- 2006, BNL experiment finds hints of discrepancy with Standard Model



# Spin precession equations

The “anomalous precession frequency” is the rate at which the muon’s spin and momentum accumulate relative angle.

$$\vec{\omega}_a = \vec{\omega}_s - \vec{\omega}_c = -\frac{q}{m} a_\mu \vec{B}$$

where  $a_\mu = \frac{1}{2}(g - 2)$ . With the right experiment, this rate is (*nearly*) proportional to the anomalous magnetic moment.



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$$\vec{\omega}_a = -\frac{q}{m} \left( a_\mu \vec{B} - a_\mu \frac{\gamma}{\gamma + 1} (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left[ a_\mu - \frac{1}{\gamma^2 - 1} \right] \frac{\vec{\beta} \times \vec{E}}{c} \right)$$





# Extracting $a_\mu$ from spin precession measurements

With simultaneous measurements of anomalous precession frequency and magnetic field, we determine the anomalous magnetic moment,

$a_\mu$

$$a_\mu = \frac{\omega_a}{B} \frac{m}{q}$$



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$$a_\mu = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} \frac{\mu'_p(T_r)}{\mu_e(H)} \frac{\mu_e(H)}{\mu_e} \frac{g_e}{2} \frac{m_\mu}{m_e}$$

This experiment measured the ratio of the anomalous precession frequency,  $\omega_a$ , to the precession frequency of a spherical water sample in the same volume,  $\tilde{\omega}'_p$ . Other terms from the literature are used to convert this ratio to  $a_\mu$ .  $\pm 25$  ppb



# A schematical view of the ratio $\omega_a/\tilde{\omega}'_p(r, y, \theta, t)$

$$a_\mu = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} \frac{\mu'_p(T_r)}{\mu_e(H)} \frac{\mu_e(H)}{\mu_e} \frac{g_e}{2} \frac{m_\mu}{m_e}$$

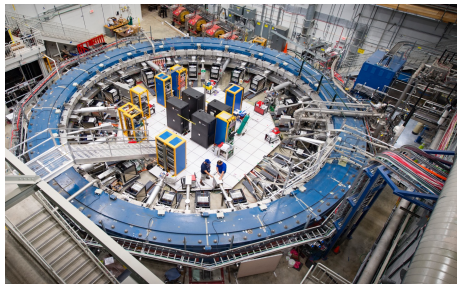
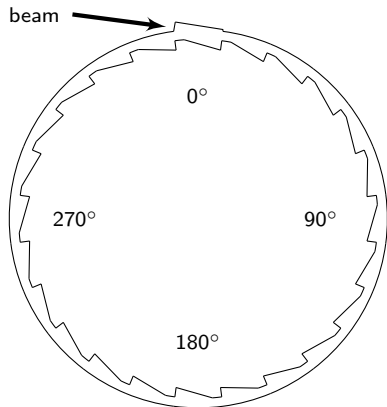
$$\vec{\omega}_a = -\frac{q}{m} \left( a_\mu \vec{B} - a_\mu \frac{\gamma}{\gamma + 1} (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left[ a_\mu - \frac{1}{\gamma^2 - 1} \right] \frac{\vec{\beta} \times \vec{E}}{c} \right)$$

Insert equation for  $\omega_a$  into the measured ratio in equation for  $a_\mu$ , collect correction terms to write a schematical equation:

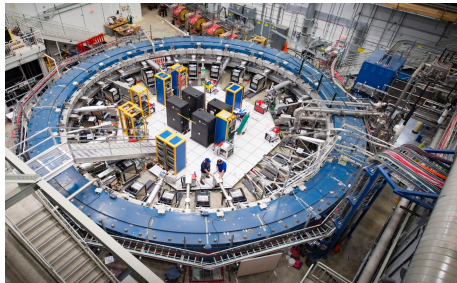
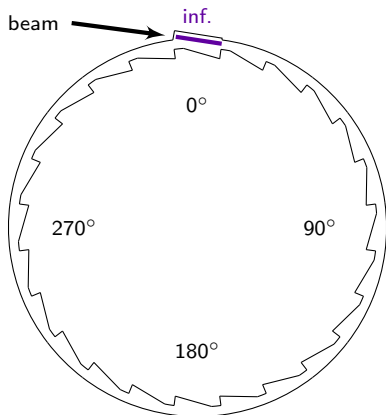
$$\frac{\omega_a}{\tilde{\omega}'_p} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$



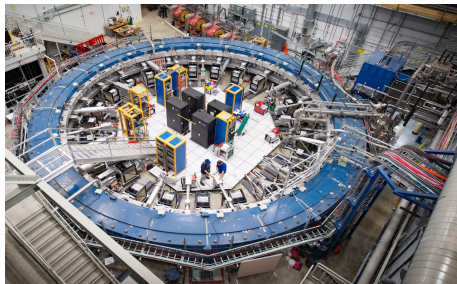
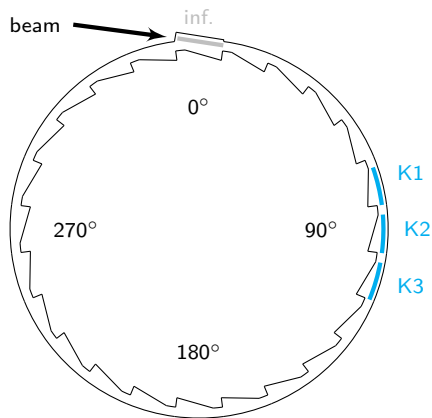
# The Muon $g - 2$ storage ring



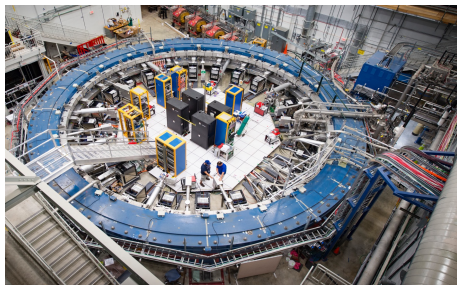
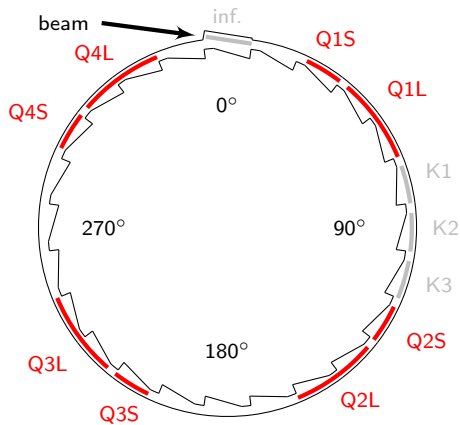
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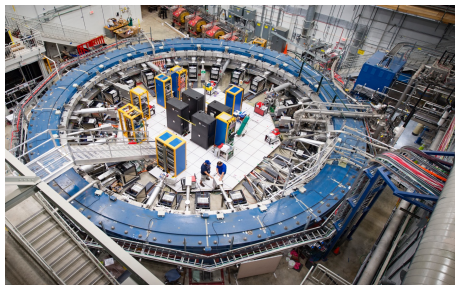
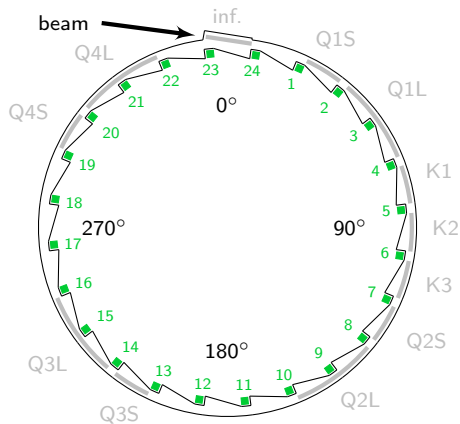
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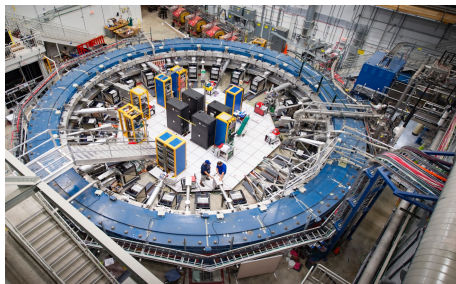
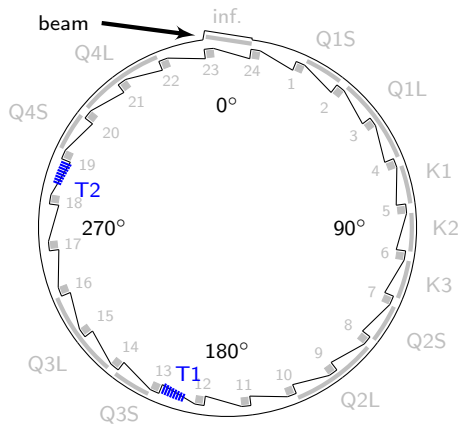


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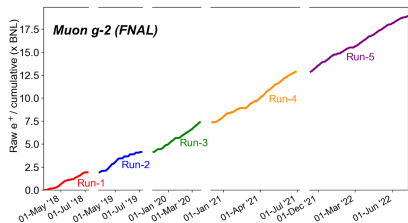




# The Muon $g - 2$ storage ring



# What's new in Run-2/3?



- Broken resistors in Run-1 replaced, fixed quad charging time
- Kicker system upgraded to kick at full spec
- Magnet insulation to stabilize temperature-dependent effects
- Heroic transient measurement efforts



# Collecting Uncertainties

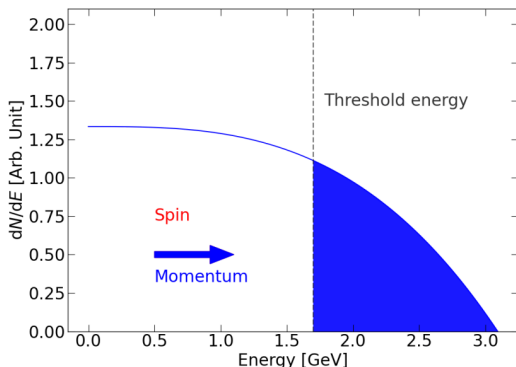
Quantity	Correction (ppb)	Uncertainty(ppb)
$\omega_a$ (stats)	—	434
$\omega_a$ (syst)	—	56
$C_e$	489	53
$C_p$	180	13
$C_{ml}$	-11	5
$C_{pa}$	-158	75
$f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle$	—	56
$B_k$	-27	37
$B_q$	-17	92
$\mu_p^i(34.7^\circ)/\mu_e$	—	10
$m_\mu/m_e$	—	22
$g_e/2$	—	0
Systematic total	—	157
Fundamental factors	—	25
Totals	544	462

- Dominated by statistics
- Largest systematics uncertainties are phase acceptance and quadrupole transient



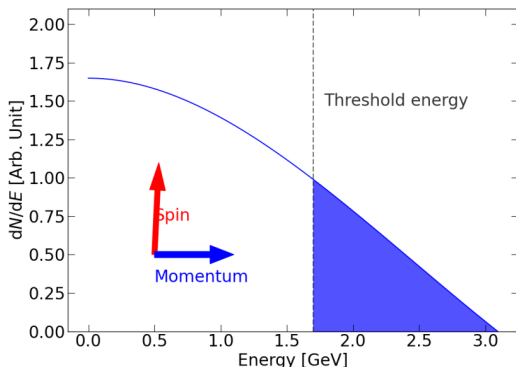
# Measuring the precession frequency with positrons

- High-momentum positrons preferentially emitted along the muon's spin
- In lab frame, more high-energy positrons emitted when spin parallel to momentum



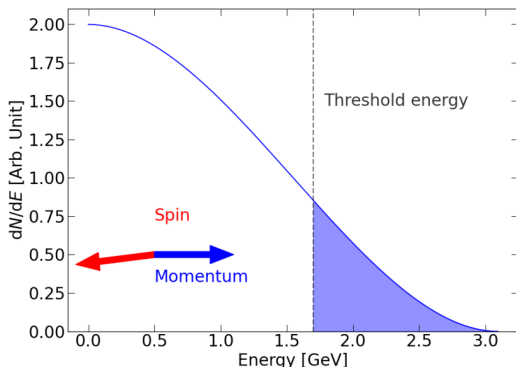
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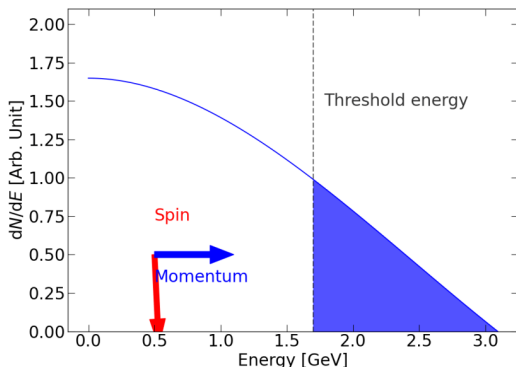
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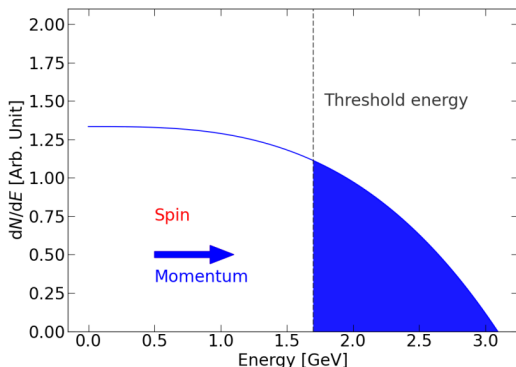
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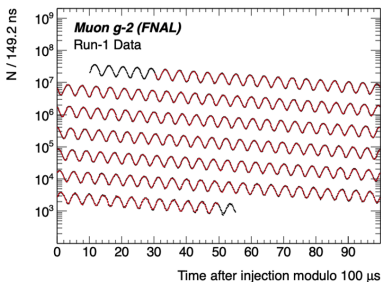




# Measuring the precession frequency with positrons

- High-momentum positrons preferentially emitted along the muon's spin
- In lab frame, more high-energy positrons emitted when spin parallel to momentum
- Modulation of positron energy spectrum encodes anomalous precession frequency

$$\frac{\omega_a}{\omega'_p} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$

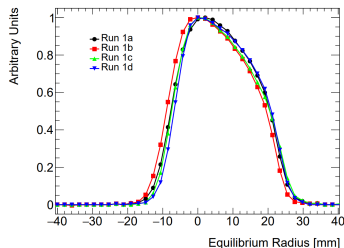


Stat.:  $\pm 434 \text{ ppb}$ , Syst.:  $\pm 56 \text{ ppb}$



The corrections to  $\omega_a$ 

$$\frac{\omega_a}{\omega_p'} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$



$$c_e = 2n(1 - n)\beta_0^2 \frac{\langle x_e^2 \rangle}{R_0^2}$$

$489 \pm 53$  ppb

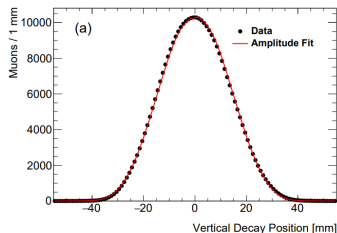
$c_e$  E-field correction minimized using magic momentum, remainder corrected by studying mean radius which is correlated to momentum dispersion

$$\frac{q}{m} \left[ a_\mu - \frac{1}{\gamma^2 - 1} \right] \frac{\vec{\beta} \times \vec{E}}{c}$$



The corrections to  $\omega_a$ 

$$\frac{\omega_a}{\omega'_p} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$



$$c_p = \frac{n}{2} \frac{\langle y^2 \rangle}{R_0^2}$$

$180 \pm 13$  ppb

$c_e$  The E-field

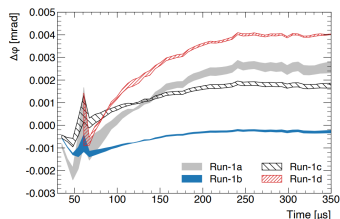
$c_p$  Pitch correction from the small vertical component of muons' momentum, calculated by studying vertical position distribution from trackers

$$\frac{q}{m} a_\mu \frac{\gamma}{\gamma + 1} (\vec{\beta} \cdot \vec{B}) \vec{\beta}$$



# The corrections to $\omega_a$

$$\frac{\omega_a}{\omega_p} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$



$$\frac{d\phi_0}{dt} = \frac{d\phi_0}{d\langle p \rangle} \frac{d\langle p \rangle}{dt}$$

$$-11 \pm 5 \text{ ppb}$$

$c_e$  The E-field

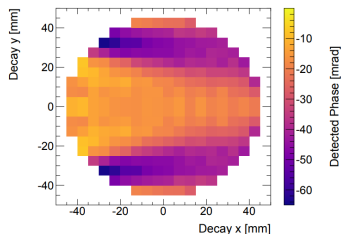
$c_p$  Pitch correction

$c_{ml}$  Muon loss correction from initial phase-momentum correlation in muons, as muons are lost in time, there is time dependent change in phase



The corrections to  $\omega_a$ 

$$\frac{\omega_a}{\omega_p} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_{\mu} \rangle (1 + B_{qt} + B_{kick})}$$



$-158 \pm 75$  ppb

$c_e$  The E-field

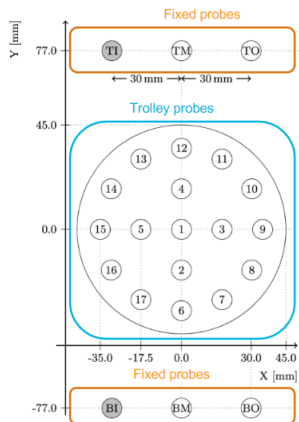
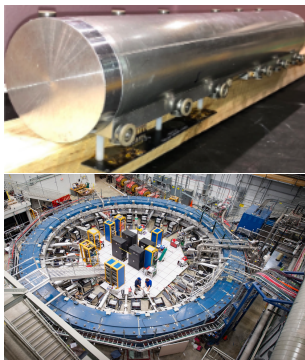
$c_p$  Pitch correction

$c_{ml}$  Muon loss correction

$c_{pa}$  Phase acceptance correction caused by decay-position dependence of positron phase, early-to-late beam motion modulation leads to time-dependent phase



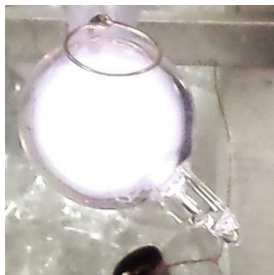
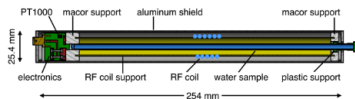
# The magnetic field measurement systems



# The field calibration chain

$$\frac{\omega_a}{\omega_p'} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_{\mu} \rangle (1 + B_{qt} + B_{kick})}$$

- Absolute calibration done at ANL, cross-checked with  $^3\text{He}$



$\pm 32$  ppb

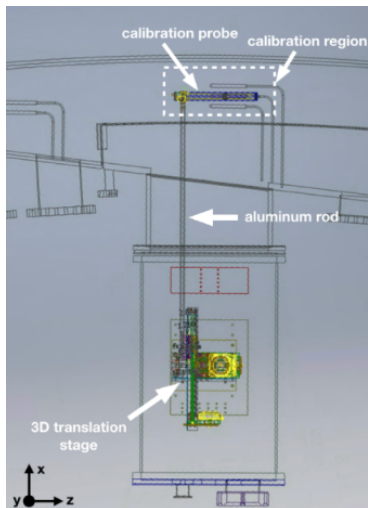


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- Absolute calibration done at ANL, cross-checked with  $^3\text{He}$
- Plunging probe calibrates trolley

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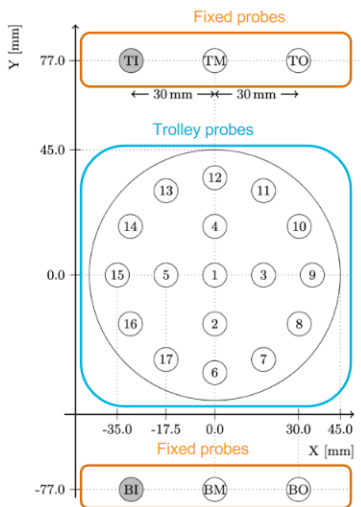


# The field calibration chain

$$\frac{\omega_a}{\omega'_p} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$

- Absolute calibration done at ANL, cross-checked with  $^3\text{He}$
- Plunging probe calibrates trolley
- Trolley synchronizes with fixed probes

$\pm 32$  ppb

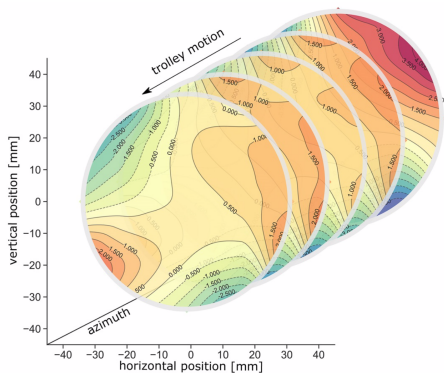


## Getting field maps from the trolley

$$\frac{\omega_a}{\omega_p'} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$

- Trolley has an array of 17 NMR probes, travels around ring, takes measurements at  $\sim 8000$  azimuthal locations
- Measurements in same volume muons are stored in

$-13 \pm 25$  ppb



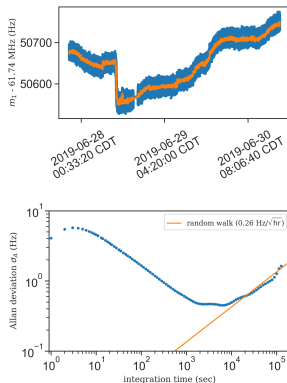
# Interpolating between trolley runs with the fixed probes

$$\frac{\omega_a}{\omega'_p} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$

- Fixed probes synchronized to the trolley during the trolley runs
- Fixed probes track field between trolley runs, additional data points to interpolate field

$\pm(24 \text{ to } 44) \text{ ppb}$

Super-long Trolley Runs

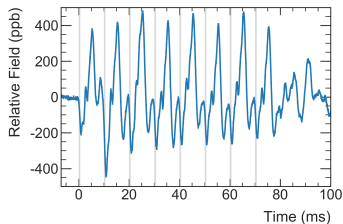


# The fast transient systematics

- Electric quadrupoles vibrate generating eddy currents that perturb the field
- Perturbation too fast for fixed probes to pick up normally, shielded by the vacuum chambers
- Measurements made with special probe to determine effect of perturbation

$-17 \pm 92$  ppb

$$\frac{\omega_a}{\omega_p'} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho \mu \rangle (1 + B_{qt} + B_{kick})}$$

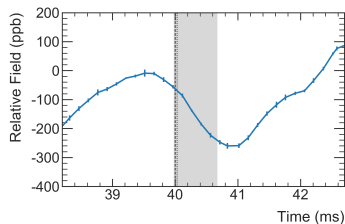


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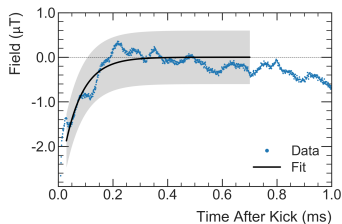


# The fast transient systematics

- Pulsed magnetic kick generates large eddy currents, perturb the magnetic field
- Like quad transient, too fast to be picked up by regular analysis
- Used optical Faraday magnetometer to study the transient

$-27 \pm 37$  ppb

$$\frac{\omega_a}{\omega_p'} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho \mu \rangle (1 + B_{qt} + B_{kick})}$$

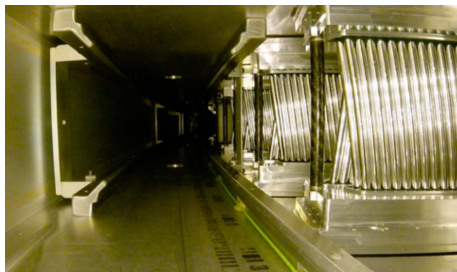


# Straw trackers

$$\frac{\omega_a}{\omega'_p} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$

- Straw trackers track decay positron path
- Tracks can be extrapolated back to muon decay position
- Decay vertices used to determine muon distribution in storage region.

$(-3 \text{ to } 1) \pm (11 \text{ to } 20) \text{ ppb}$

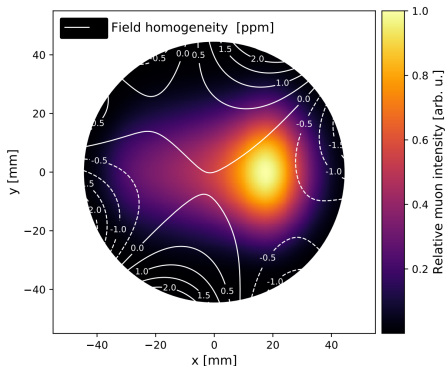


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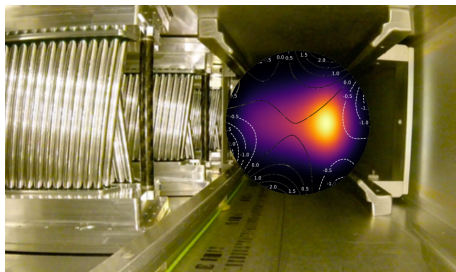


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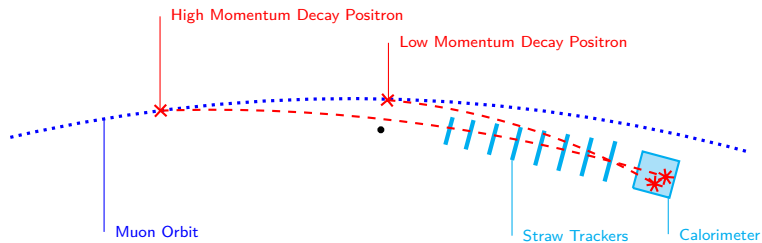
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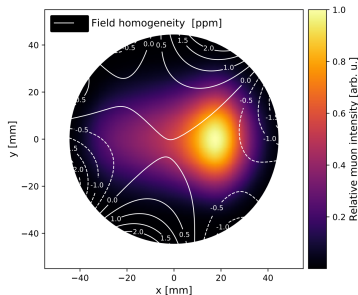
# Straw trackers



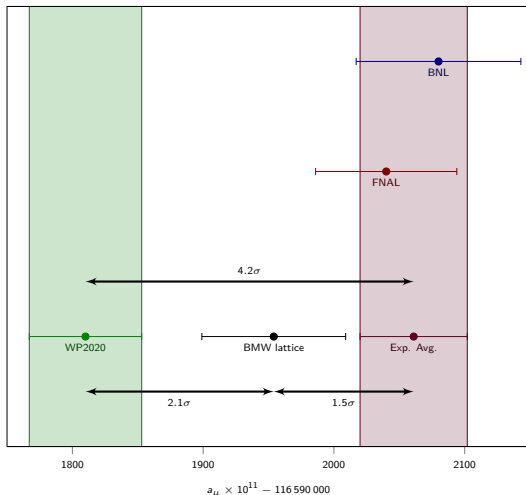
# Combining the field maps and muon distribution

$$\frac{\omega_a}{\omega'_p} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$

- Magnetic field maps weighted by muon distribution determined by trackers
- Trackers measure at two locations storage ring, use beam dynamics simulations to extrapolate distribution around ring



## Recap (460 ppb Uncertainty)



# Acknowledgments

Thanks to:

- University of Kentucky
- Muon  $g - 2$  field team
- Muon  $g - 2$  team in general
- Fermilab and Argonne

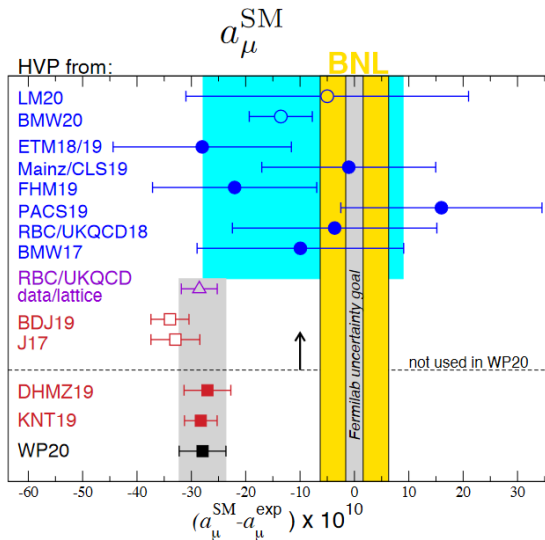


# Backup Slides



# Lattice QCD Results

- Large range of lattice values that agree with both 2020 white paper and experiment
- The BMW20 calculation still needs to be vetted by independent lattice groups
- More lattice results are upcoming, worth keeping an eye on



# Blinding

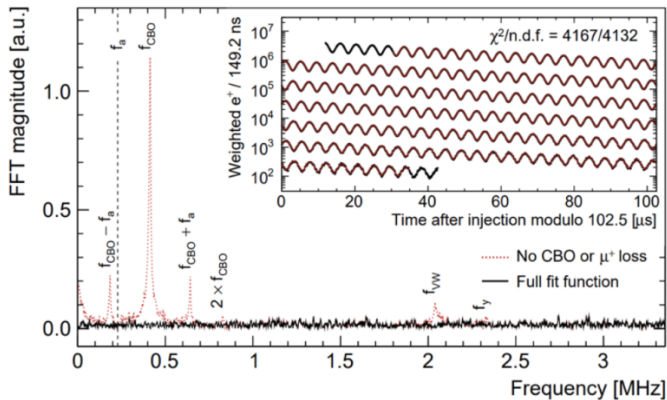
$$\frac{\omega_a}{\omega'_p} = \frac{f_{\text{clock}} \omega_{a,\text{meas}} (1 + c_e + c_p + c_{ml} + c_{pa})}{f_{\text{field}} \langle \omega_p \otimes \rho_\mu \rangle (1 + B_{qt} + B_{kick})}$$

- Experiment blinded at  $\omega_a$  clock
- Blinders not part of Muon  $g - 2$ , Fermilab personnel
- Clock frequencies kept in sealed envelopes until unblinding





# The frequencies in the measurement



# The Muon $g - 2$ storage ring

