

What can we learn about nuclear Hamiltonians from neutron star observations?

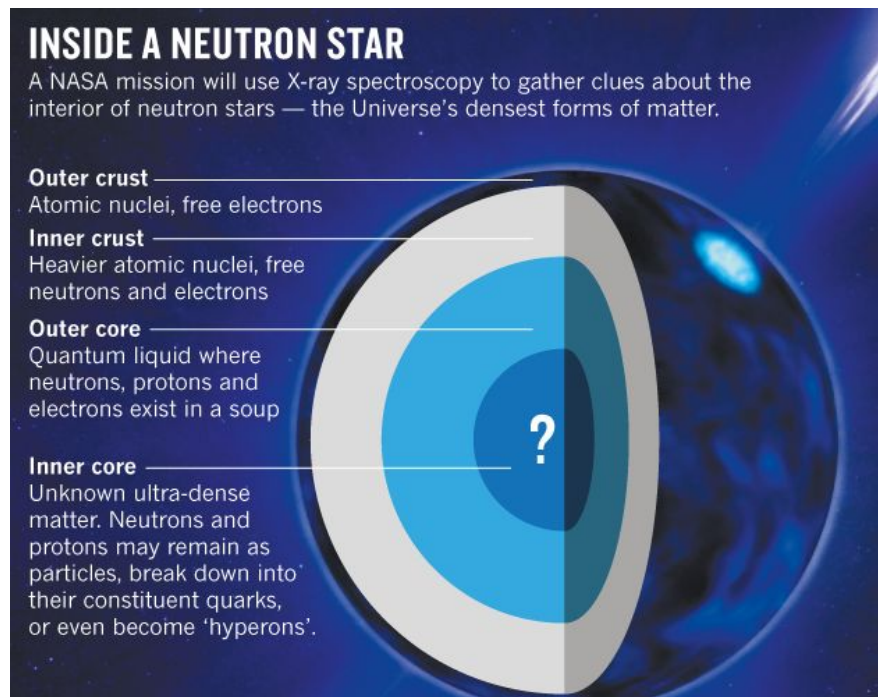
Rahul Somasundaram
Los Alamos National Laboratory

04/30/2026, Nuclear Hamiltonians for Advancing Nuclear Physics and Beyond



Neutron stars are unique nuclear systems

- Densest matter in the universe, with densities as high as $5n_{\text{sat}}$ to $8n_{\text{sat}}$
- Highly neutron rich, with the proton fraction $\sim 5\%$
- Many different observables: glitches, cooling curves etc. But Mass-Radius measurements are “clean” EOS measurements
- Exciting era of “multi-messenger” astronomy: gravitational waves, X-ray observations, electromagnetic signals, etc.



Implications for the nuclear Hamiltonian

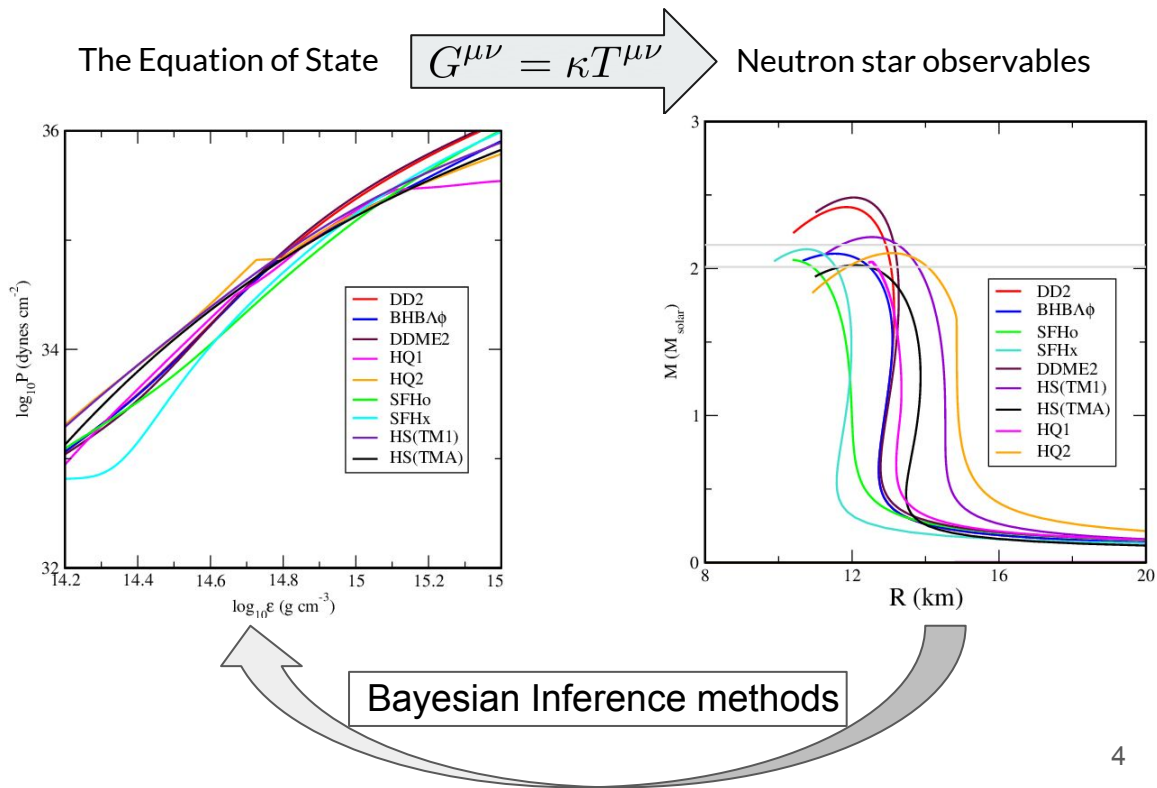
- Calibrate higher order LECs
- Understand how EFT breaks down in dense matter:
 $Q = K_F/\Lambda_B$ and K_F/Λ_C
- Probe “new” physics: Hyperon-Nucleon couplings, phase transitions, etc.

	NN	3N	4N
LO $\alpha(Q^0/\Lambda^0)$	1990 [151,152] 2 	—	—
NLO $\alpha(Q^2/\Lambda^2)$	1992 [164,165] 7 	1992,1994 [166-169] —	—
N ² LO $\alpha(Q^3/\Lambda^3)$	1992 [164,165] 0 	1994 [167,170] 2 	—
N ³ LO $\alpha(Q^4/\Lambda^4)$	2000–2002 [179-182] 12 	2008–2011 [183-185] 0 	2006 [186] 0
N ⁴ LO $\alpha(Q^5/\Lambda^5)$	2015 [188,189] 0 	2011– [190-192] ? 	?

Hebeler, Phys.Rept. 890 (2021) 1-116

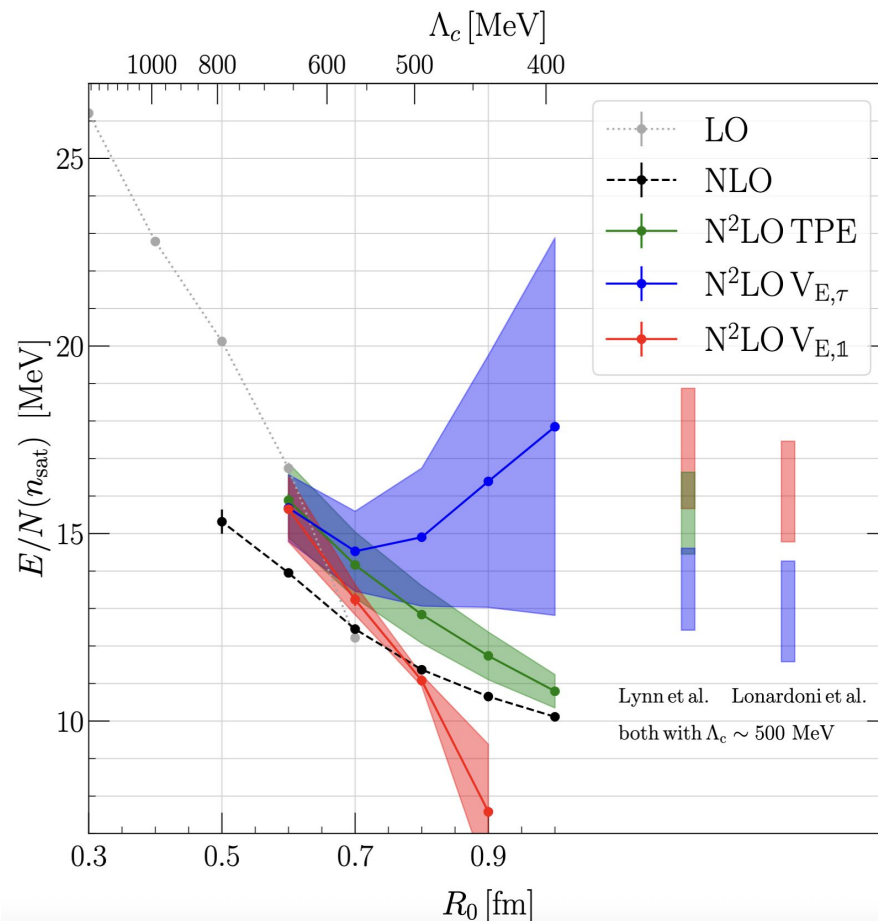
Non-trivial inverse problem

- For a given EOS, there is a single M-R curve which all neutron stars should obey
- This is an ill-defined inverse problem: crossing points, etc.
- We need an uncertainty-aware (Bayesian) framework

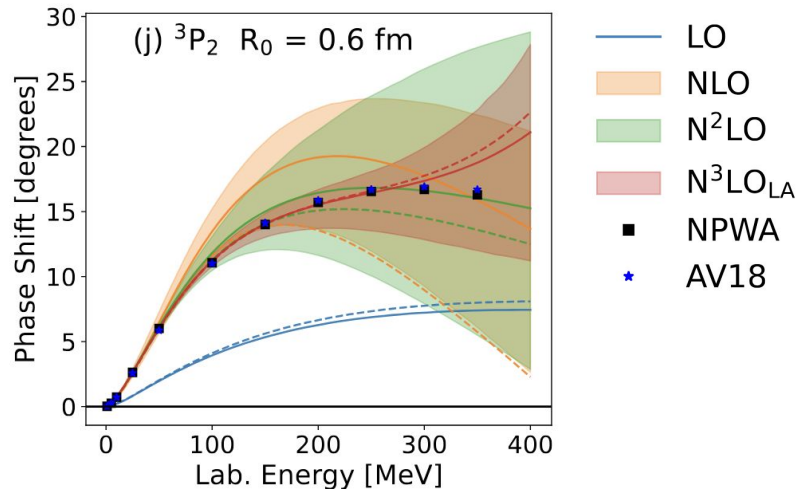
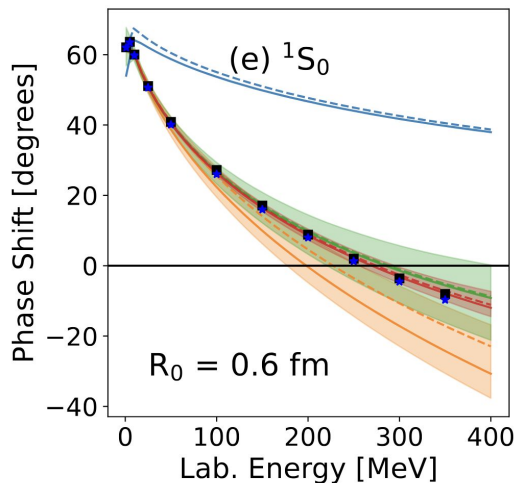


We'll focus on "High" cutoff local interactions

- Local regulators violate the Fierz rearrangement freedom. This can have a drastic effect in the 3-body sector
- Different operator structures are possible for the same 3N force, V_E . This ambiguity should vanish in the limit of infinite cutoff.



We'll focus on "High" cutoff local interactions



$$\mathcal{L} \propto \prod_i \exp \left\{ -\frac{1}{2} \left(\frac{X_i^{\text{exp}} - X_i^{\text{theo}}}{\sigma_i} \right)^2 \right\}, \quad \sigma^2 = \sigma_{\text{exp}}^2 + \sigma_{\text{theo}}^2,$$

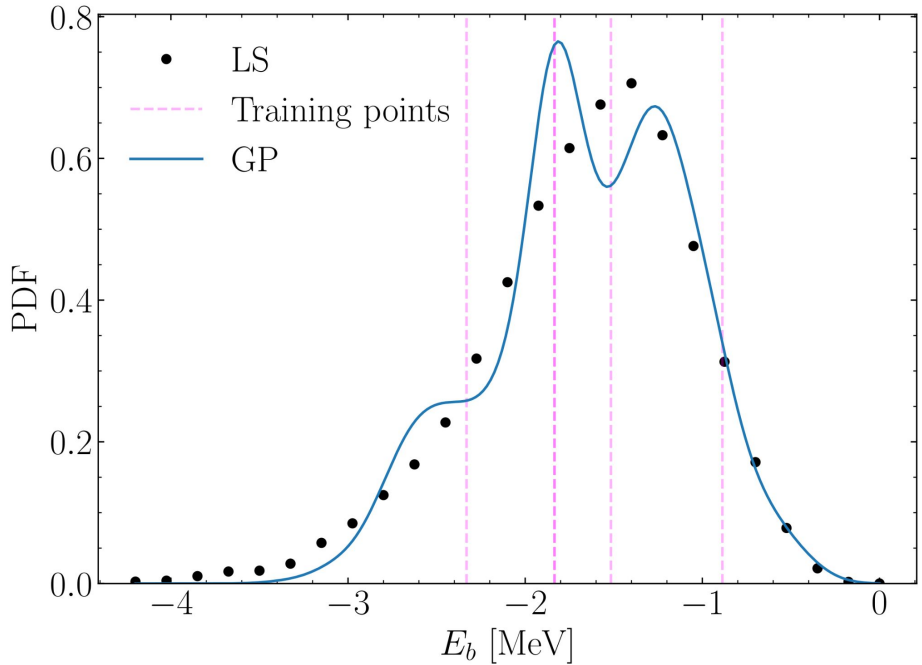
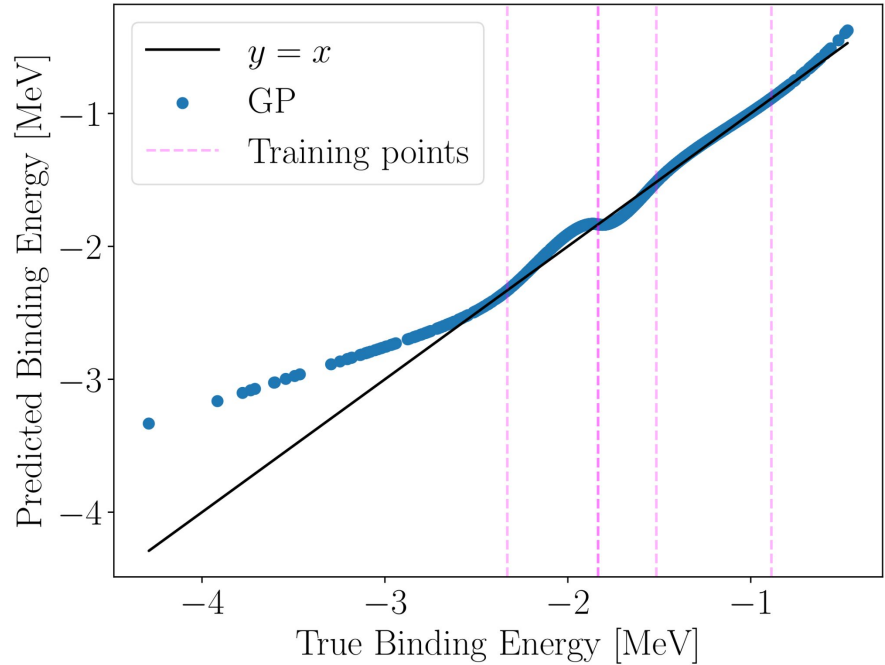
EFT Truncation error:

$$X(p) = X_{\text{ref}}(p) \sum_{n=0}^k c_n Q^n(p) \quad \Delta X_{\text{EKM}}^k = X_{\text{ref}} Q^j \max(|c_0|, |c_1|, \dots, |c_k|)$$

where $j = \max(2, k + 1)$

Emulators with scarce data: The Deuteron at leading order

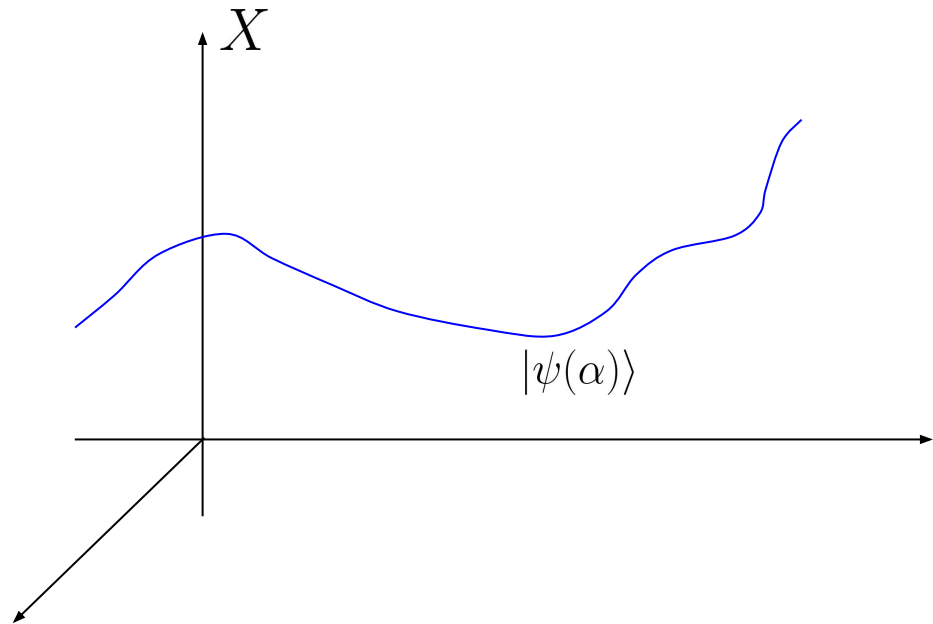
- Goal: Build emulator for QMC with ~5 - 10 training points
- A baseline GP fails to accurately interpolate and extrapolate between training points



Emulators with scarce data: Reduced basis methods

- Start with the full Schrödinger equation

$$H|\psi\rangle = E|\psi\rangle$$

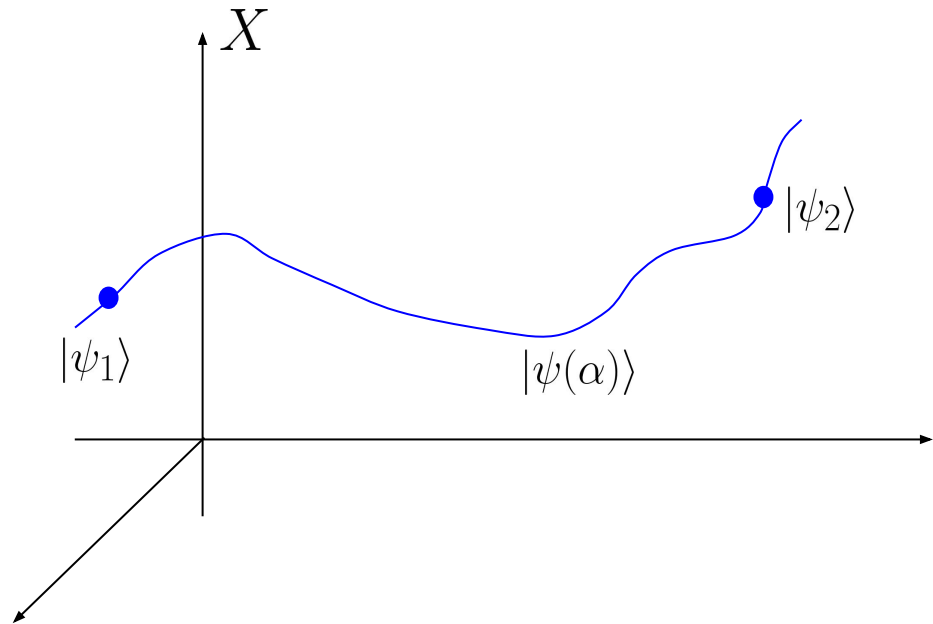


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- Compute N training functions or 'snapshots' $|\psi_j\rangle$



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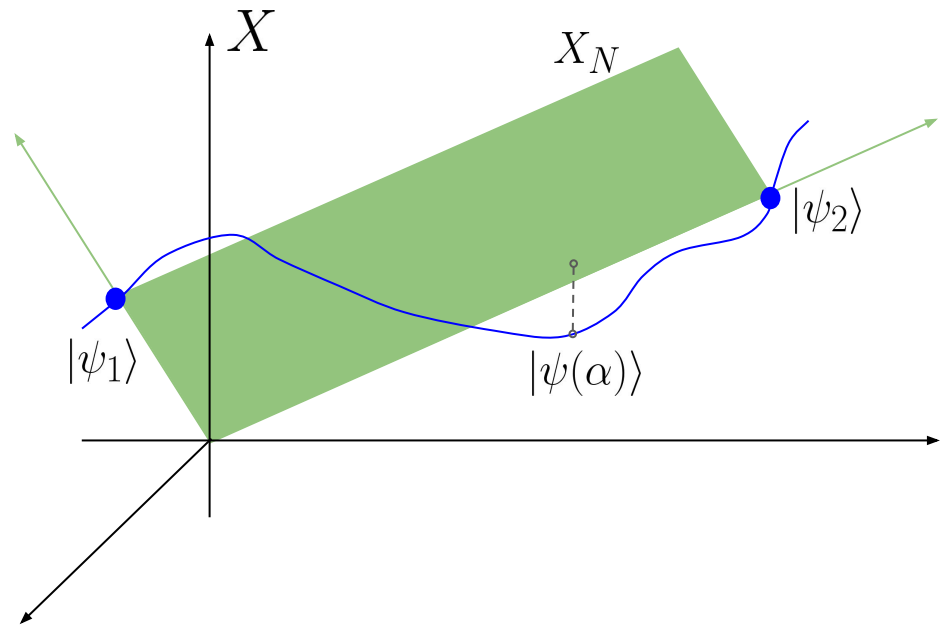
$$H|\psi\rangle = E|\psi\rangle$$

- Compute N training functions or 'snapshots' $|\psi_j\rangle$
- Project the Hamiltonian into the reduced space spanned by $|\psi_j\rangle$

Mathematically, this corresponds to computing the matrix

$$M_{ij} \equiv \langle \psi_i | H | \psi_j \rangle$$

Petrov-Galerkin projection method



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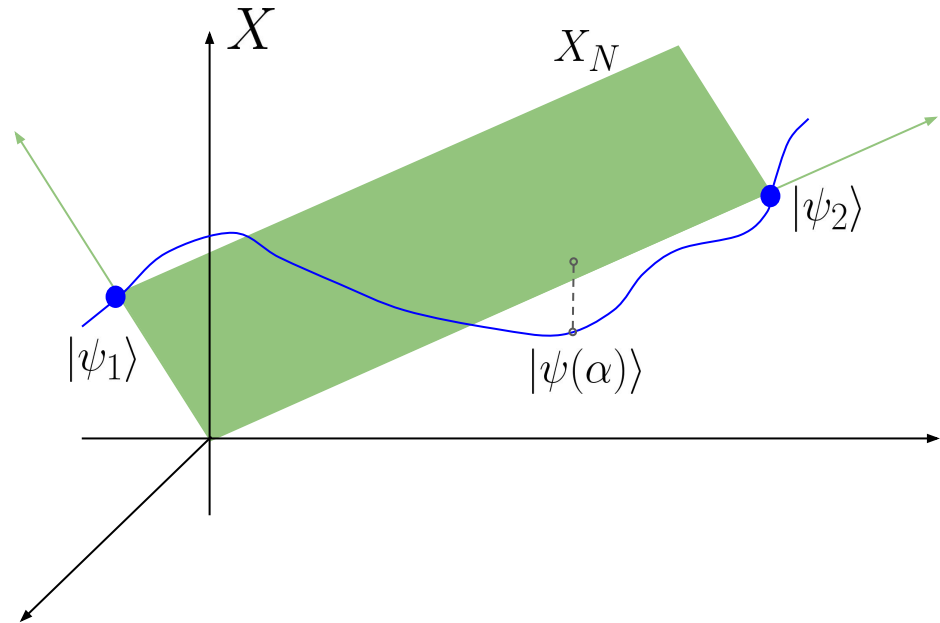
Mathematically, this corresponds to computing the matrix

$$M_{ij} \equiv \langle \psi_i | H | \psi_j \rangle$$

- For QMC, M_{ij} is dominated by stochastic noise and cannot be calculated. We therefore implemented a Petrov-Galerkin projection method for this problem:

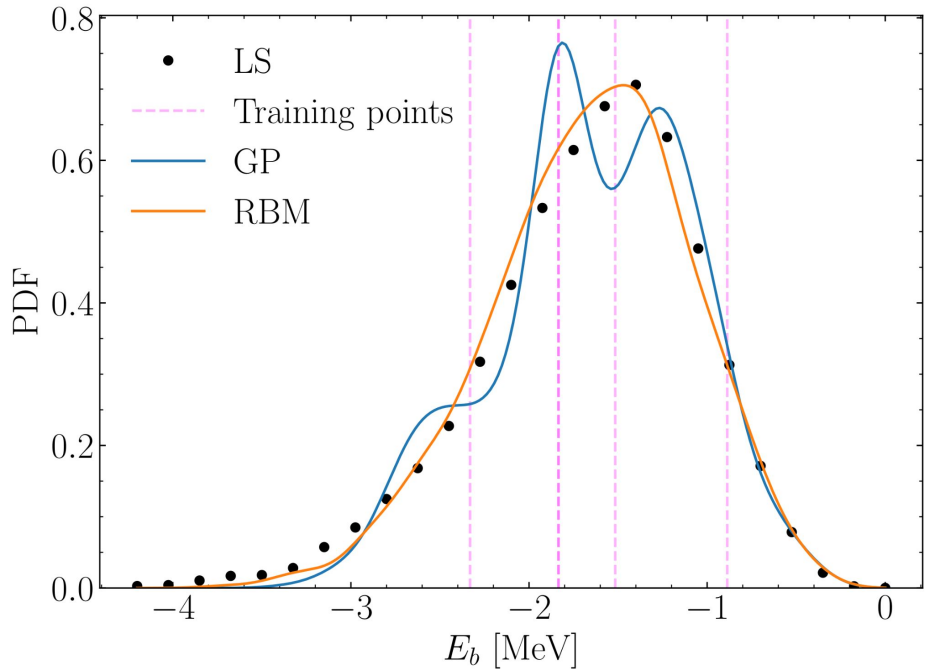
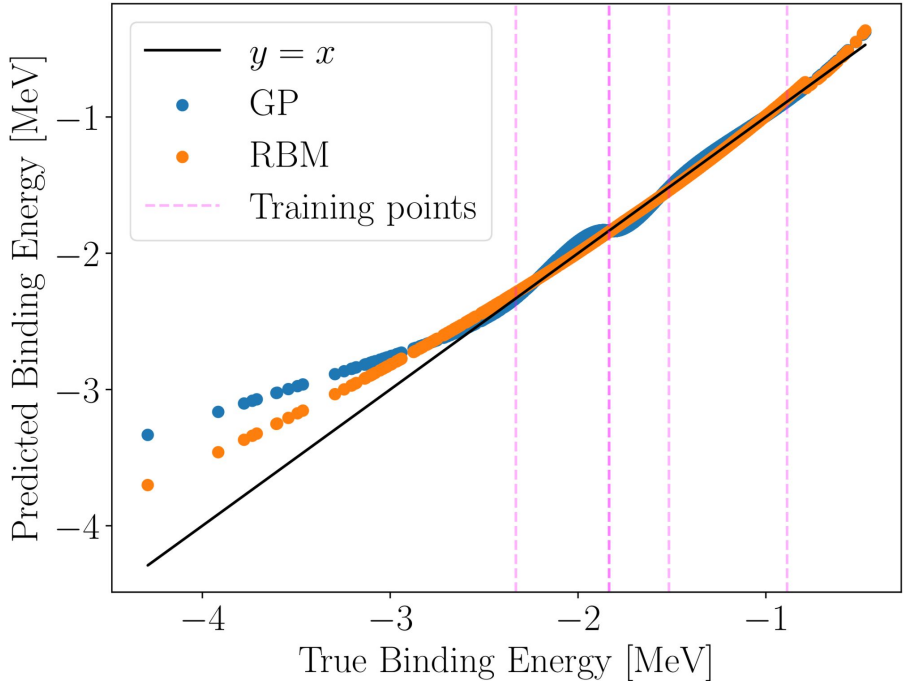
$$\tilde{M}_{ij} = \langle \psi_i^T | H | \psi_j \rangle$$

Petrov-Galerkin projection method



Emulators with scarce data: Reduced basis methods

- This RBM outperforms the GP
- This RBM is capable of interpolating but fails to extrapolate away from training points. Note that this is NOT Eigenvector Continuation

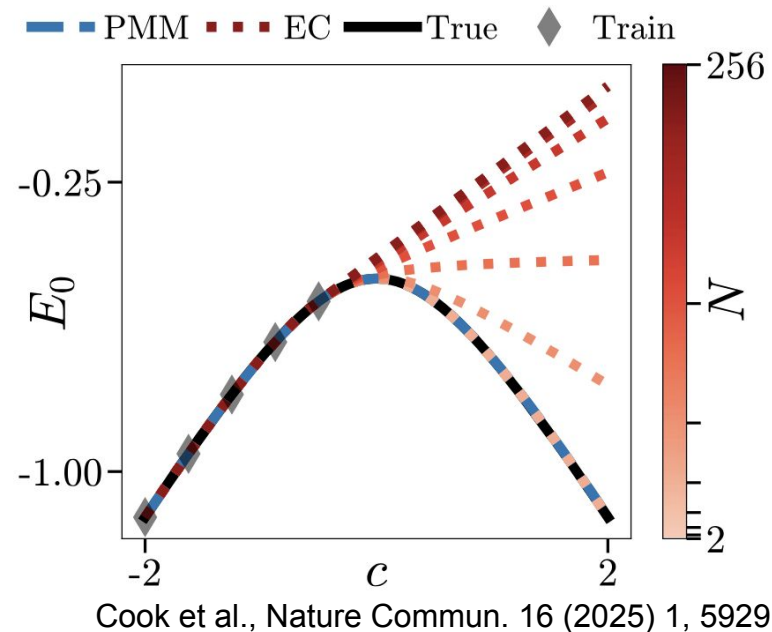


Emulators with scarce data: Hybrid models

- Combine elements of RBMs with data-driven emulators
- We employ the recently proposed parametric matrix models

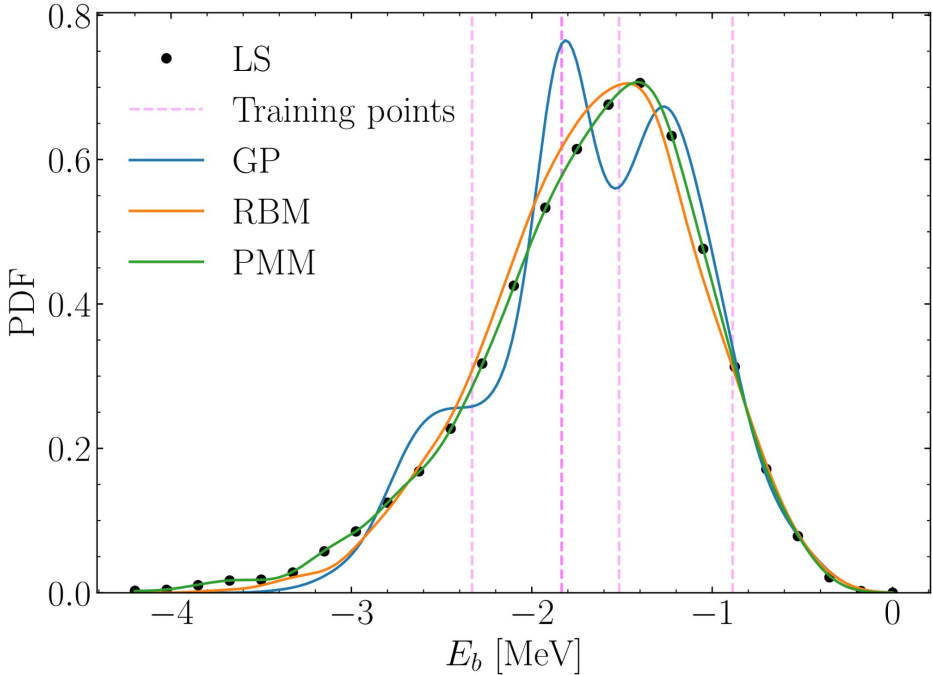
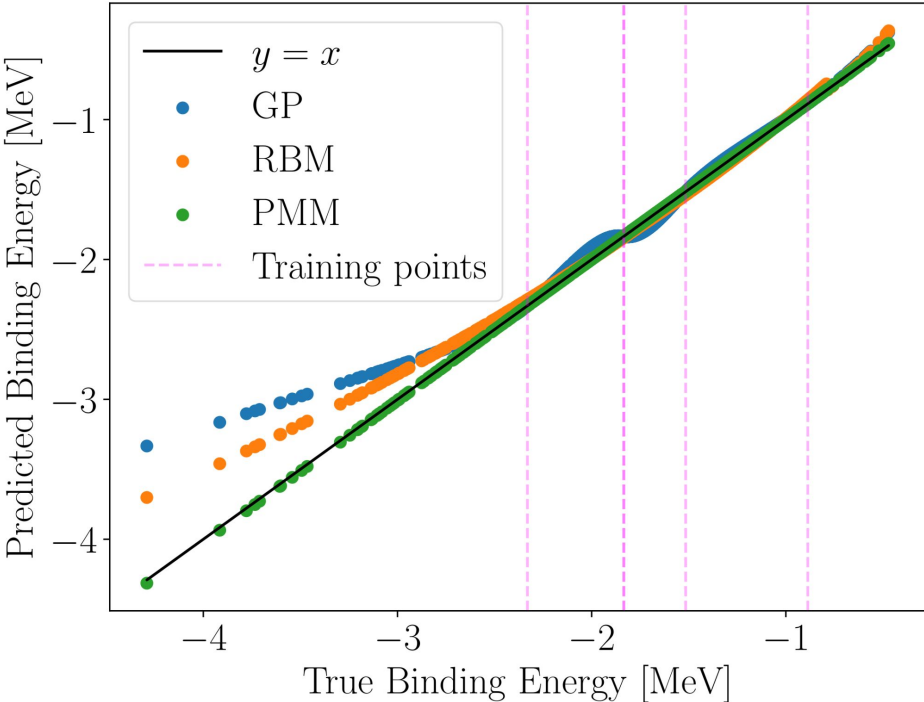
$$M(\vec{\alpha}) = M_0 + \sum_i \alpha_i M_i$$

- The form of the reduced subspace matrix is inspired by RBMs. However, we do not directly compute the projections, i.e. we do not compute the subspace matrix elements
- Instead they are learned in some manner from the data



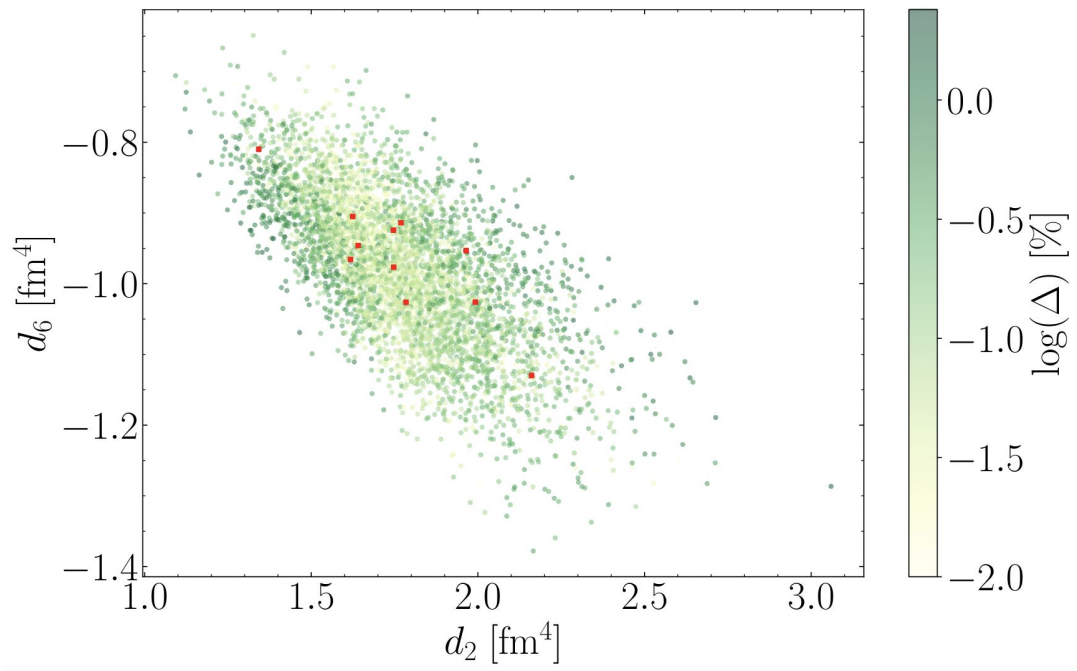
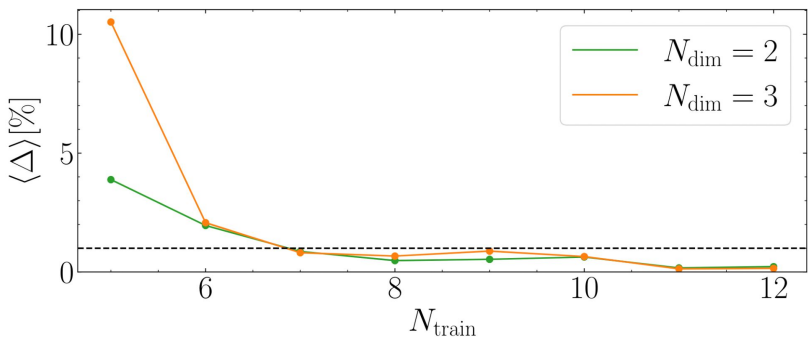
Emulators with scarce data: Hybrid models

- The PMM outperforms both the GP and our RBM
- It interpolates well but also gives excellent results for extrapolation!



Emulators with scarce data: Hybrid models. The Deuteron at N²LO

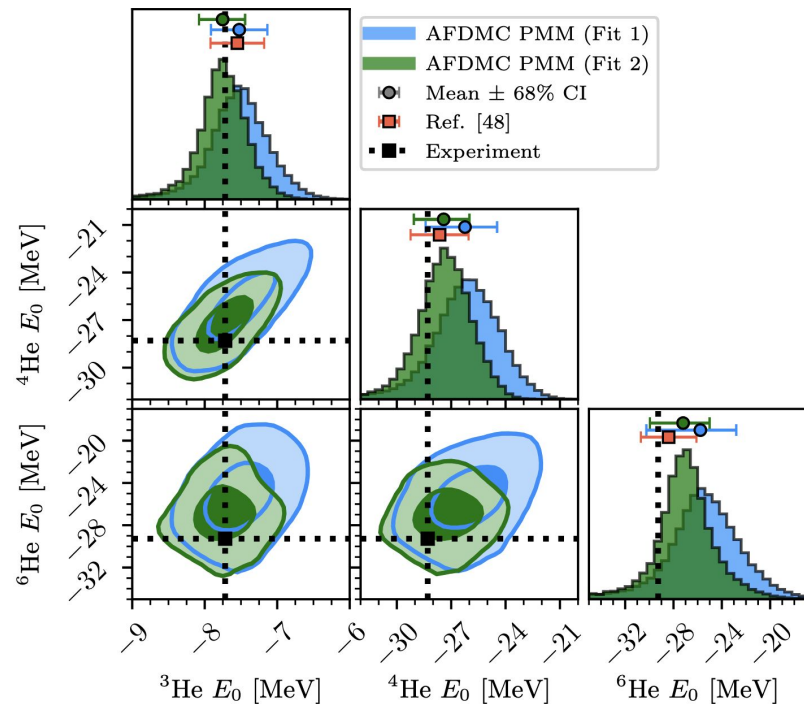
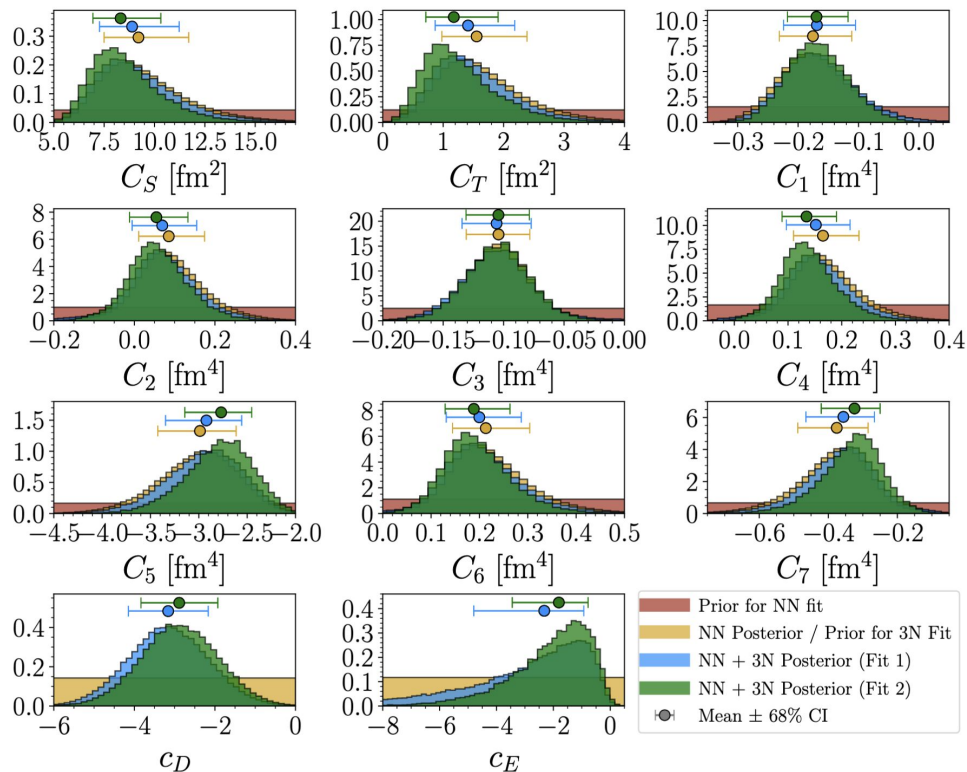
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- We found that our PMM generalises very well to at least 4 - 5 dimensional parameter spaces.

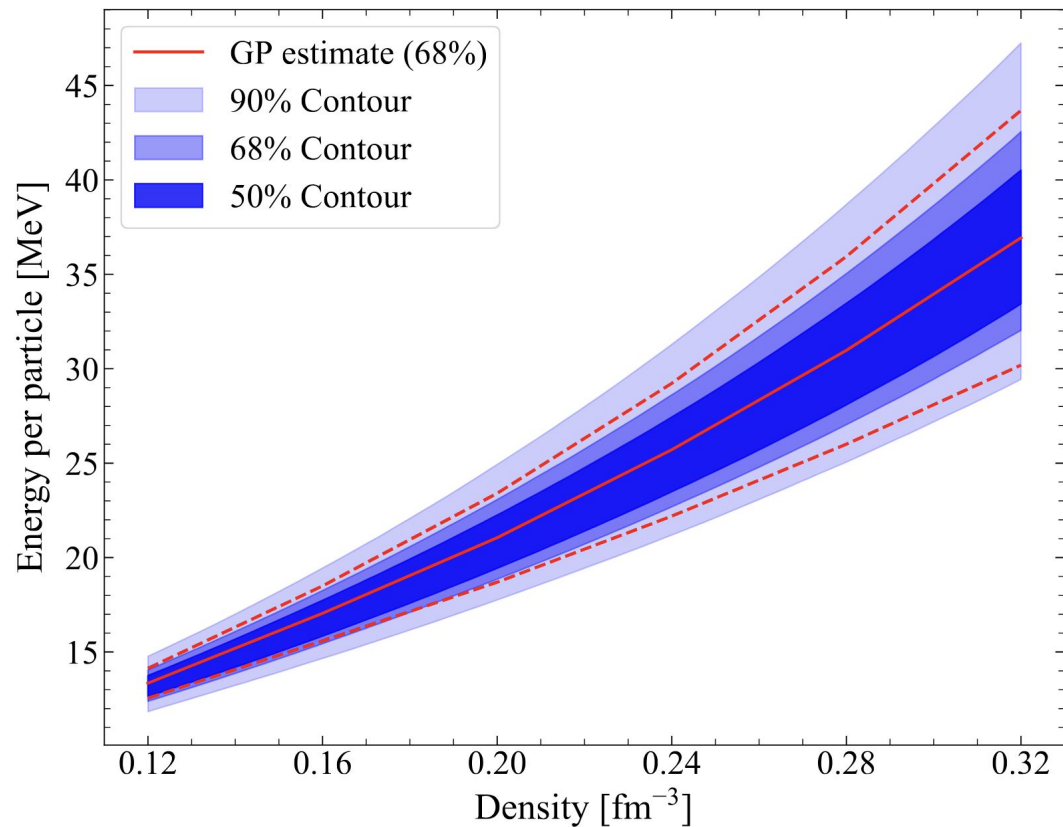
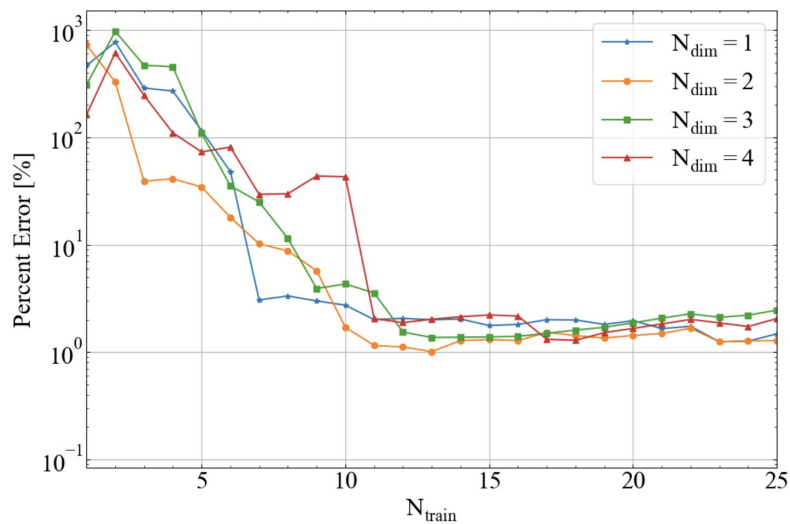
Using PMMs for light nuclei

$$P(\alpha_{3N}, \alpha_{2N} | D) \propto \mathcal{L}(\alpha_{3N}, \alpha_{2N} | D_L) \Pi(\alpha_{3N}) P(\alpha_{2N} | D_S)$$



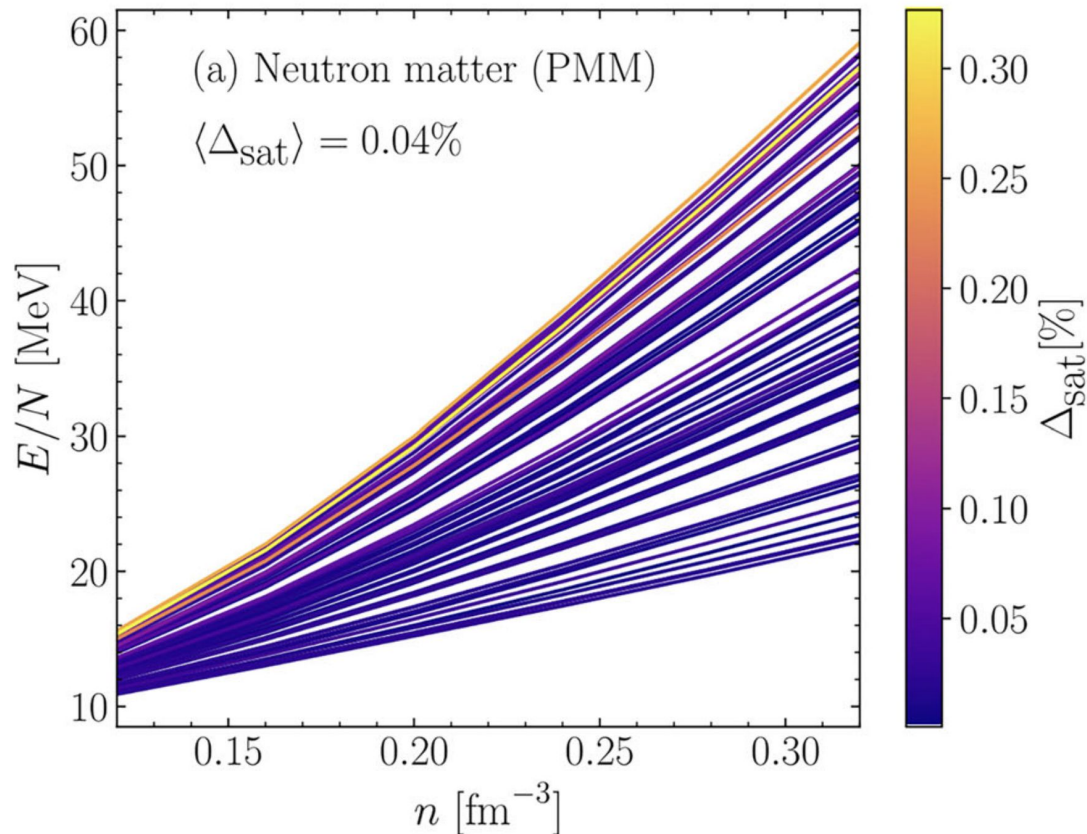
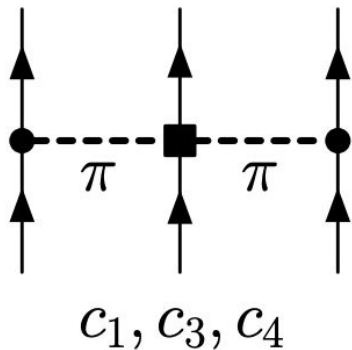
Using PMMs for neutron matter

Joint posterior distribution on 6 two-body spectral LECs propagated to pure neutron matter.



Using PMMs for neutron matter (MBPT calculations)

- N²LO interaction of Entem, Machleidt, and Nosyk. $\Lambda_C = 450$ MeV
- We kept NN part fixed. There are then two parameters that determine the 3N sector in PNM



Impact of neutron star data

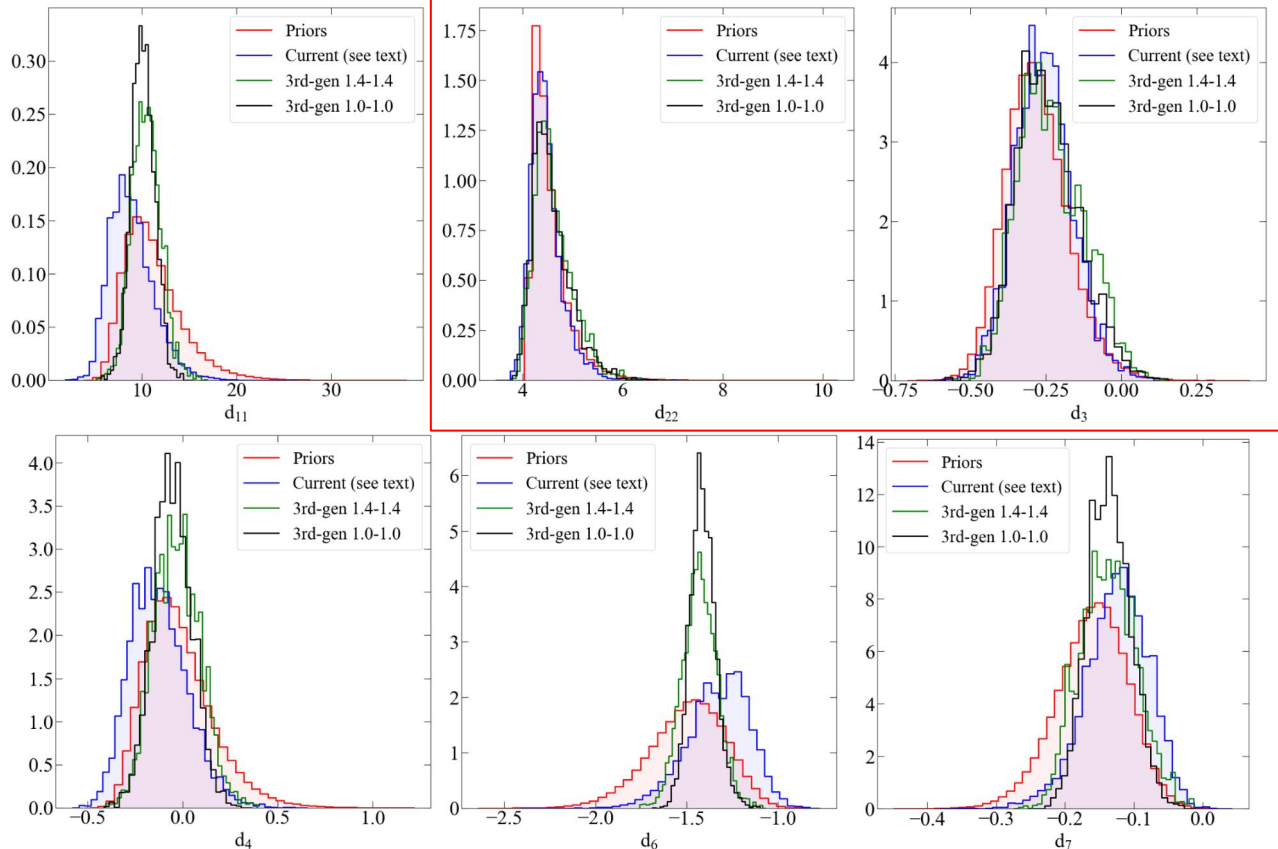
Astrophysical data includes:

- GW170817
- NICER
 - PSR J0030
 - PSR J0740
 - PSR J0437
 - PSR J0614

- Radio observations ($M_{\text{TOV}} > 2 M_{\text{sun}}$)
- Simulated next-generation events from Cosmic Explorer and Einstein Telescope

Largely, we see consistency between astrophysical data and np scattering

S-wave

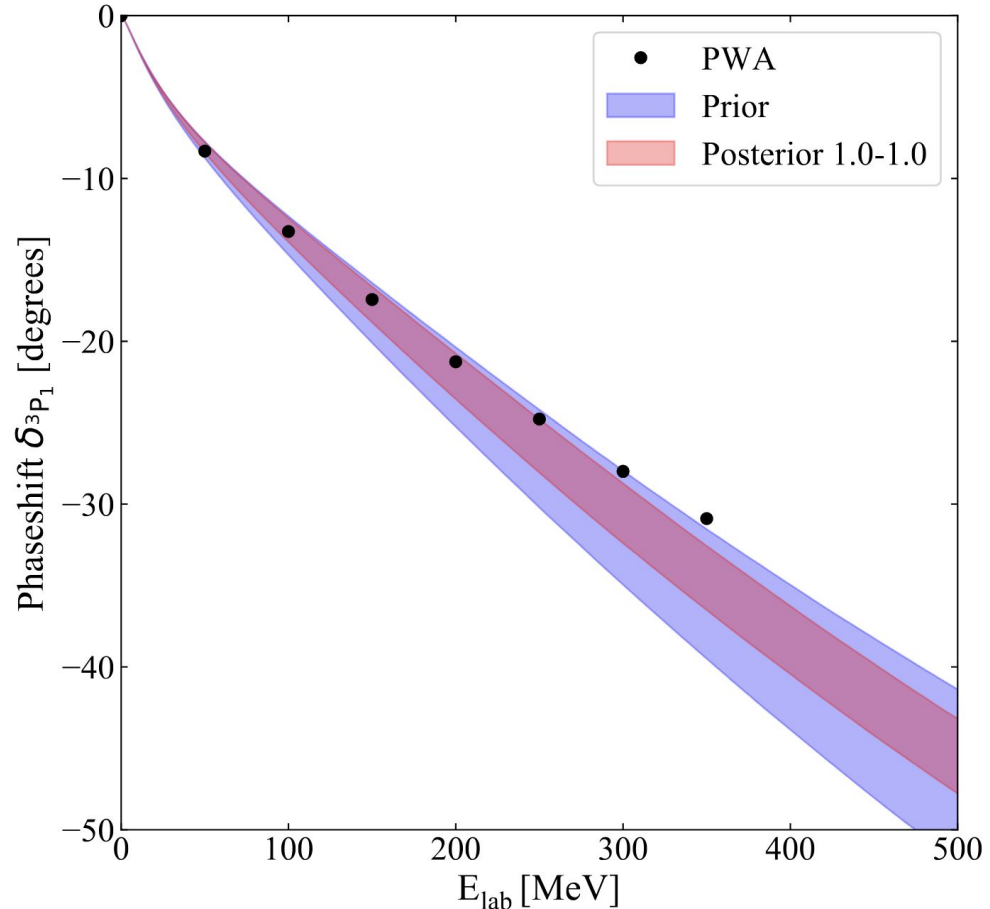


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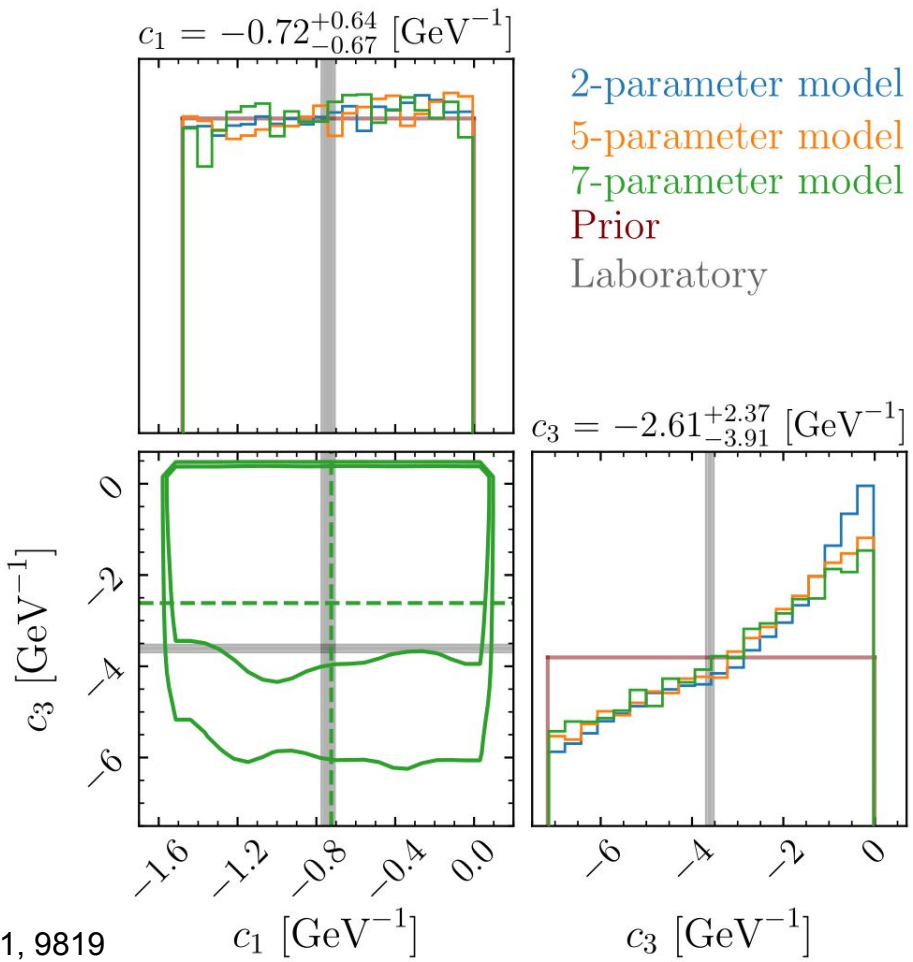
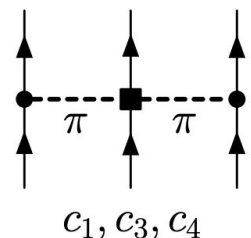
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Impact of neutron star data (EMN, 450 MeV)

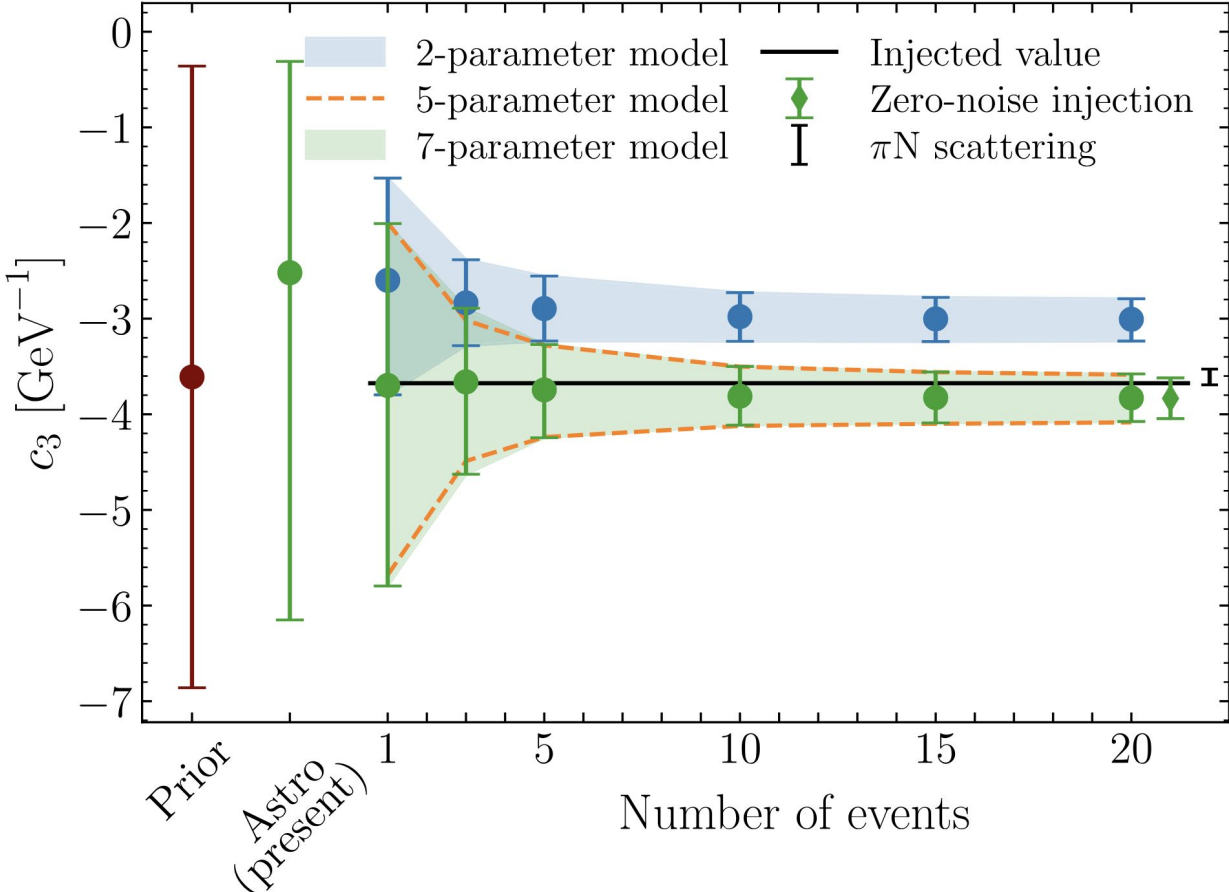
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2-parameter model
 5-parameter model
 7-parameter model
 Prior
 Laboratory

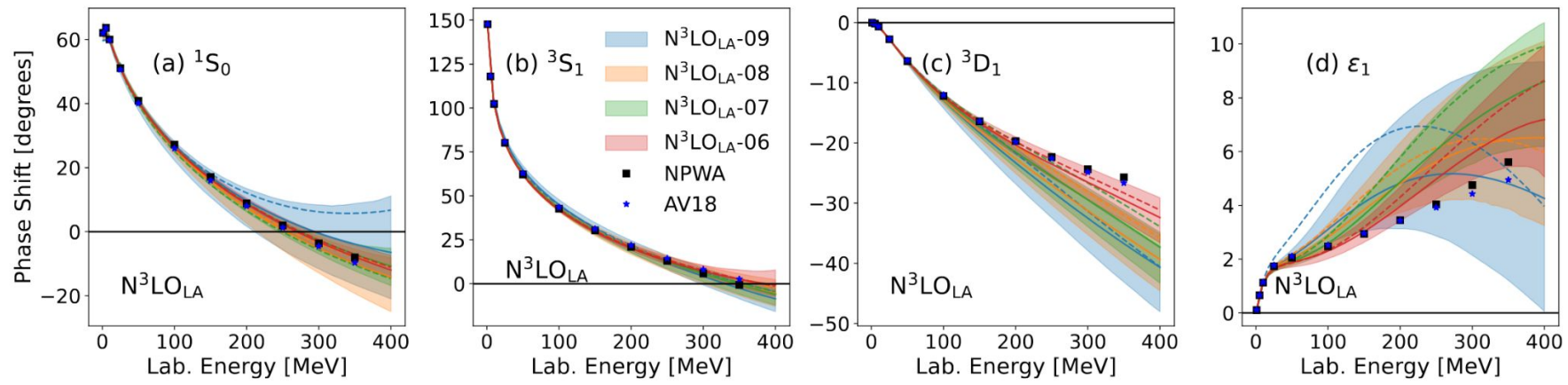
Impact of neutron star data (EMN, 450 MeV)



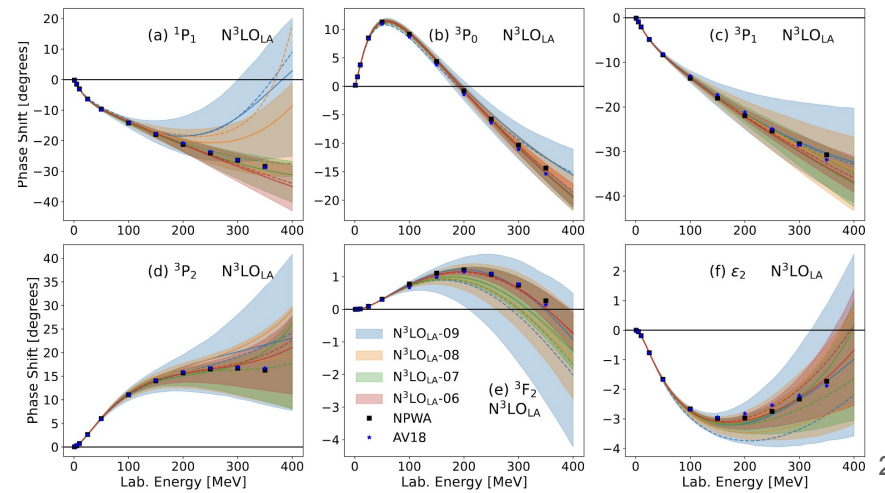
Impact of neutron star data (EMN, 450 MeV)



Local interactions at N³LO

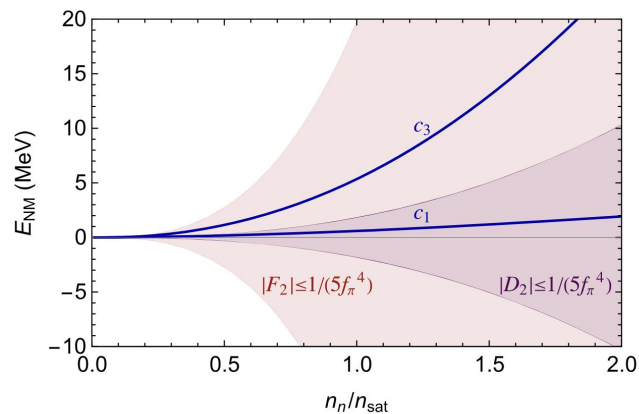


Work in progress: Extend calculations to maximally local N³LO interactions. This includes 4 non-local operators that need to be treated perturbatively.



Outlook

- Infer the breakdown scale directly in dense nuclear matter. Study cutoff dependency and compare local vs non-local interactions
- Study RG invariance of 3N forces in dense matter



Cirigliano et al., Phys.Rev.Lett. 135 (2025) 2, 022501

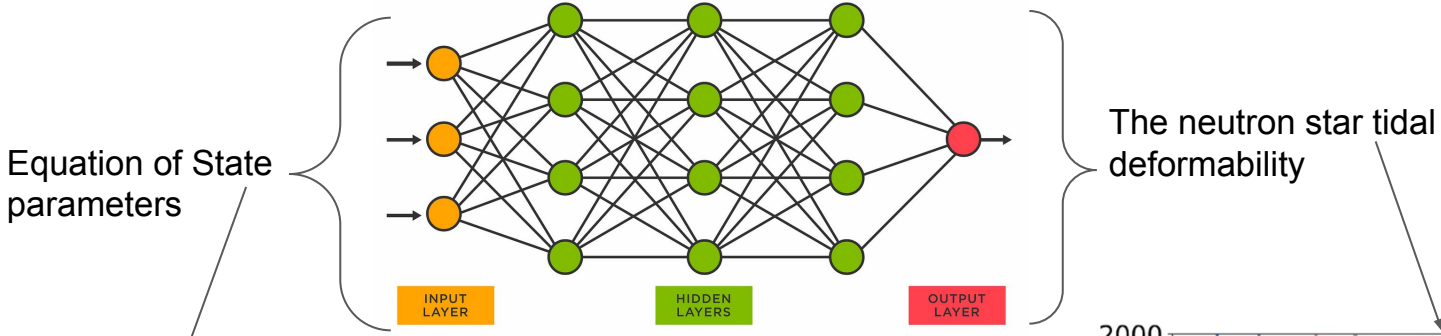
Challenges

- What other observables should we be focusing on? Can we study multi-nucleon scattering?
- How do we sample large parameter spaces (at high orders) and how much data should we include?
- Can we “sample” over the space of operators to study different power counting schemes?

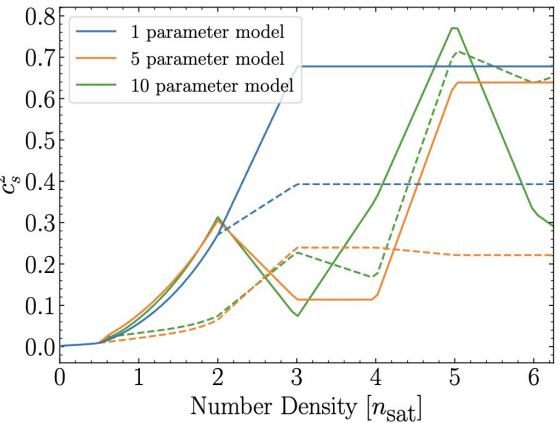
Backup Slides

Emulating the TOV equations: Deep ensembles of Multilayer Perceptrons

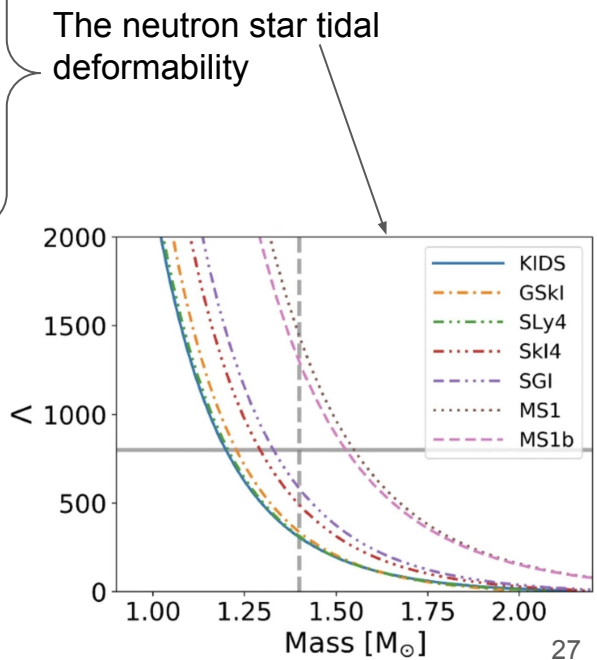
Multilayer Perceptrons (MLP) are the simplest, dense, feedforward neural networks. We use the method of deep ensembles where we use a set of 100 MLPs



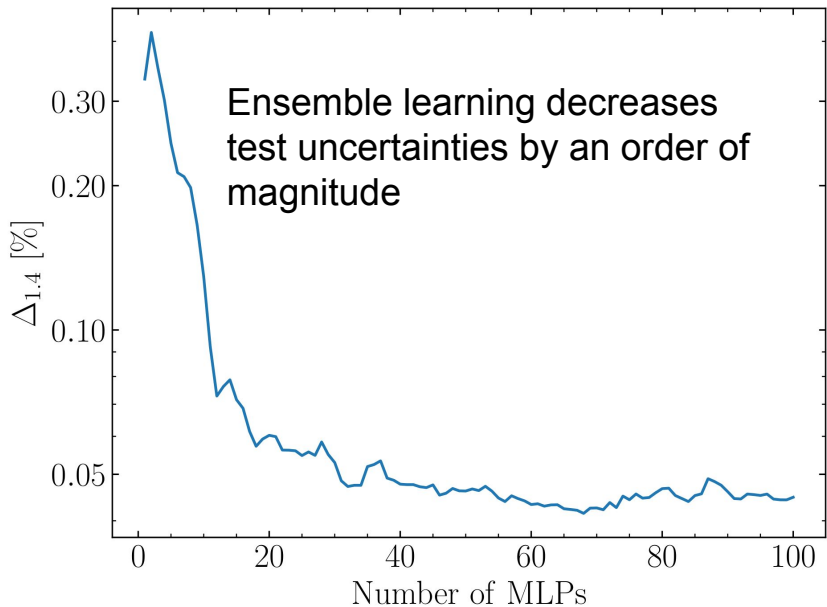
(x100)



We generated 200K samples, randomly split into training and test data

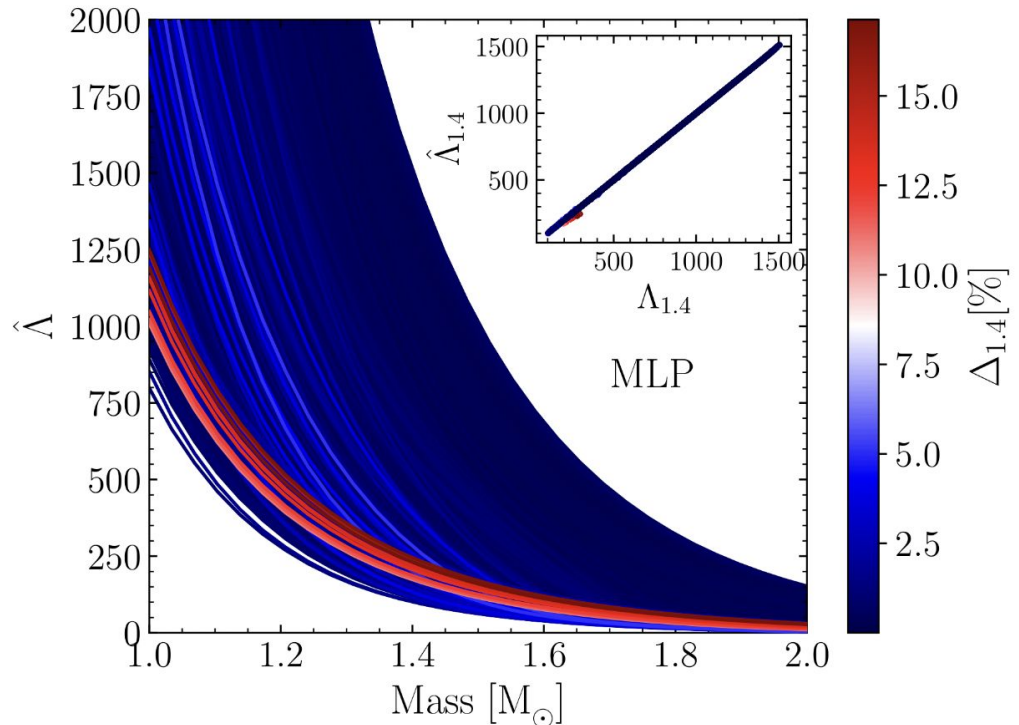


Deep ensembles of Multilayer Perceptrons



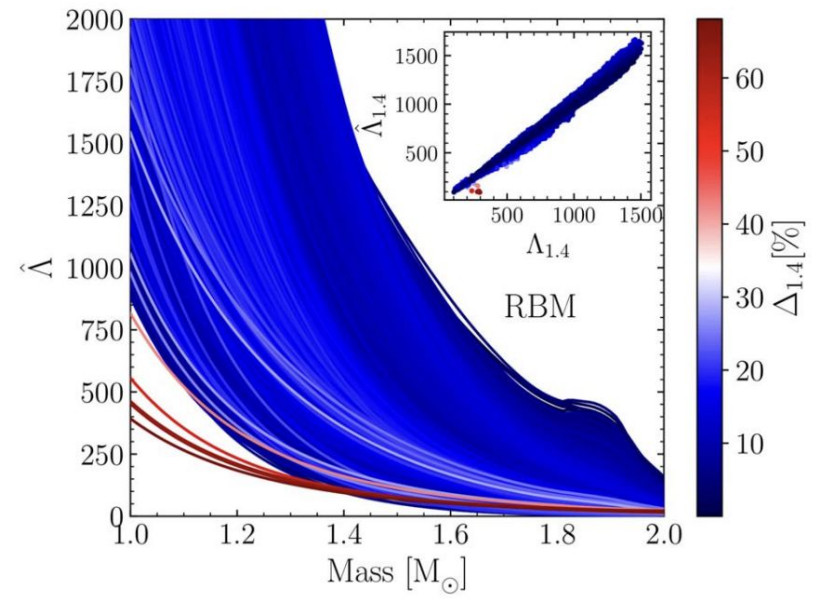
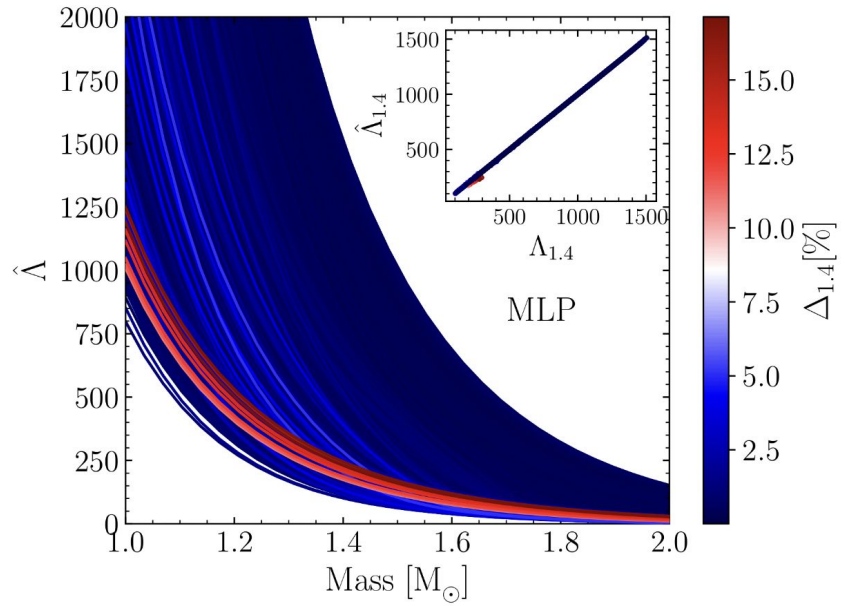
- Average uncertainty on test samples is 0.04 %
- 0.01% of test samples are outliers (uncertainty > 10%)

$$\Delta[\%] \equiv \left| \frac{\Lambda - \hat{\Lambda}}{\Lambda} \right| \times 100$$



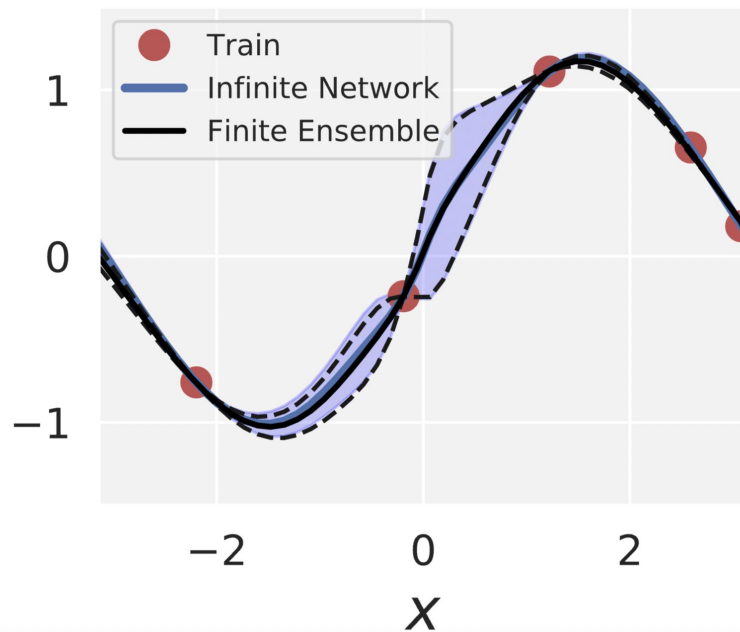
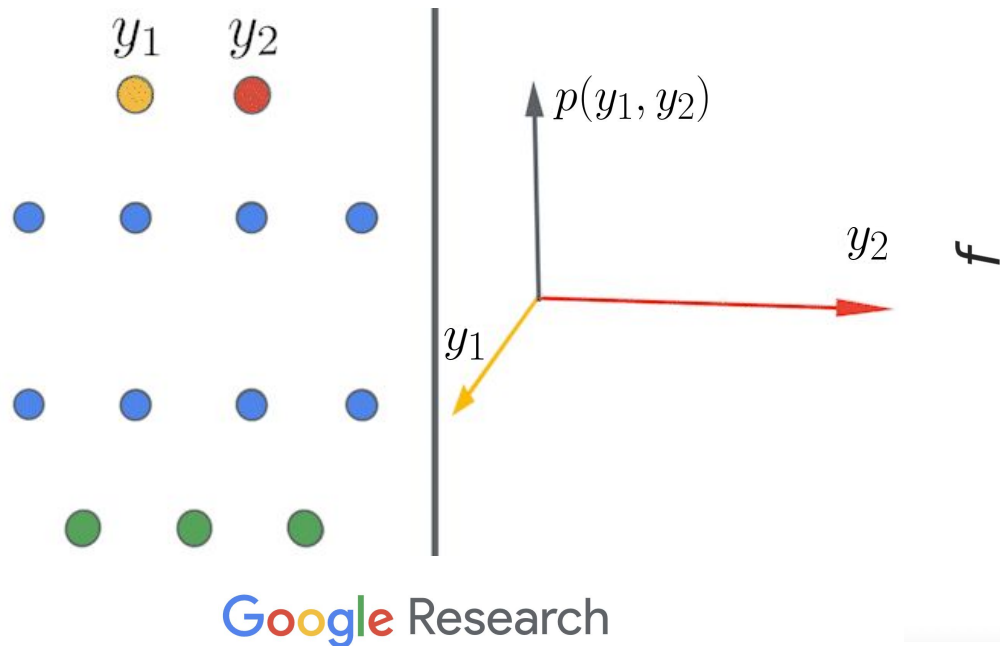
Deep ensemble versus RBMs

Overall, our deep ensemble also performs better on test samples, with better validation scores.



Uncertainty quantification using deep ensembles

- In the limit of infinite neural network width, deep ensembles are mathematically equivalent to Gaussian Processes



Uncertainty quantification using deep ensembles

We use MLPs with a width of 64 neurons.
Therefore, our predictions are not exactly Gaussian distributed

