

Finite temperature effects in binary neutron star simulations

[Rivieccio, Nadal-Matosas, Rios & Ruiz, *ApJ* **987**, 67 \(2025\)](#)

[D Guerra, Rios et al arxiv:2512.05118](#)

[Kochankovski, Rozalén, Rios & Ramos, *in prep*](#)



Arnau Rios Huguet

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Institute of Cosmos Sciences
Universitat de Barcelona
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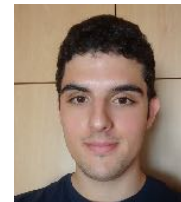
Guerra



Rivieccio



Rozalén



Kochankovski



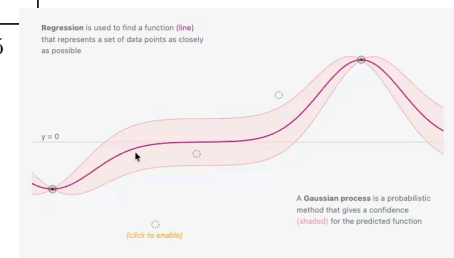
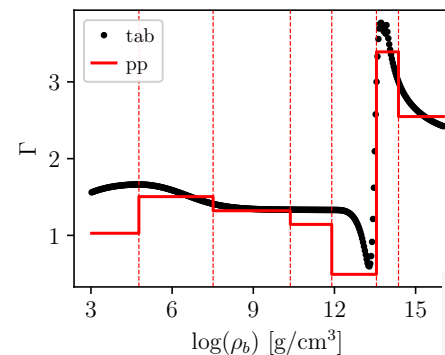
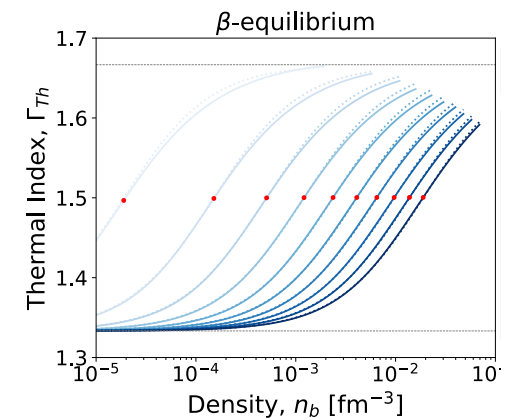
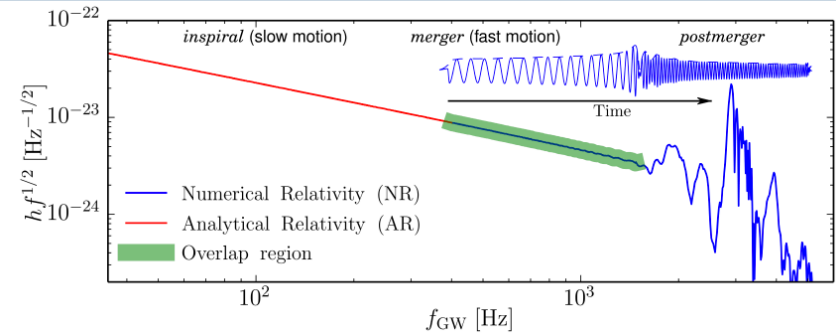
INT Finite T
15 June 2026

- **Motivation**

- **Virial approximation**

- **Thermal effects in BNS**

- **Gaussian processes for finite T**

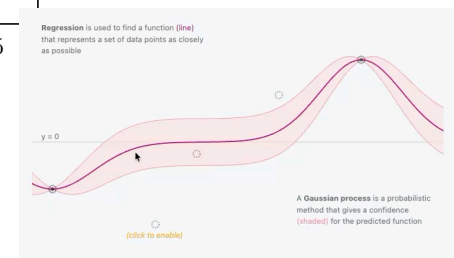
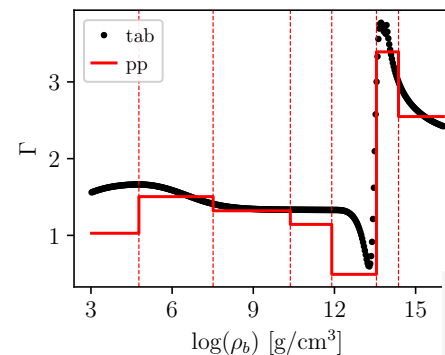
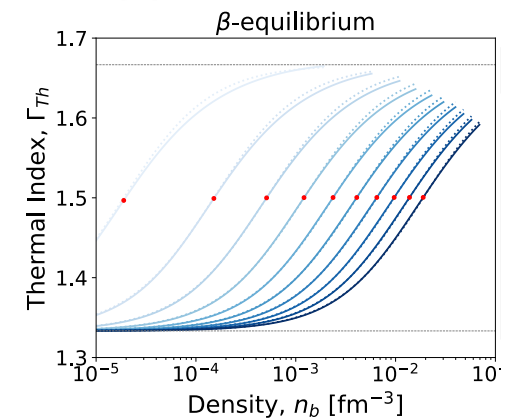
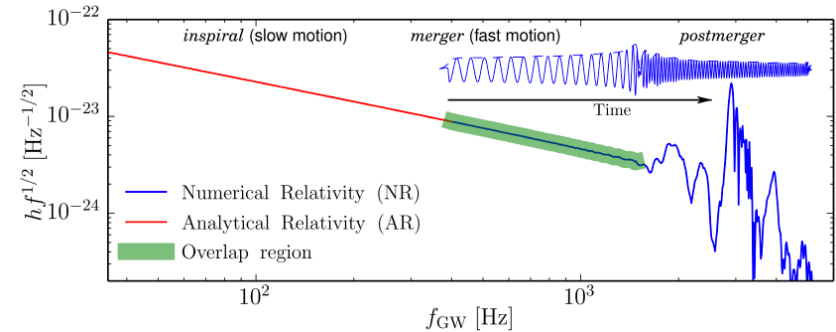


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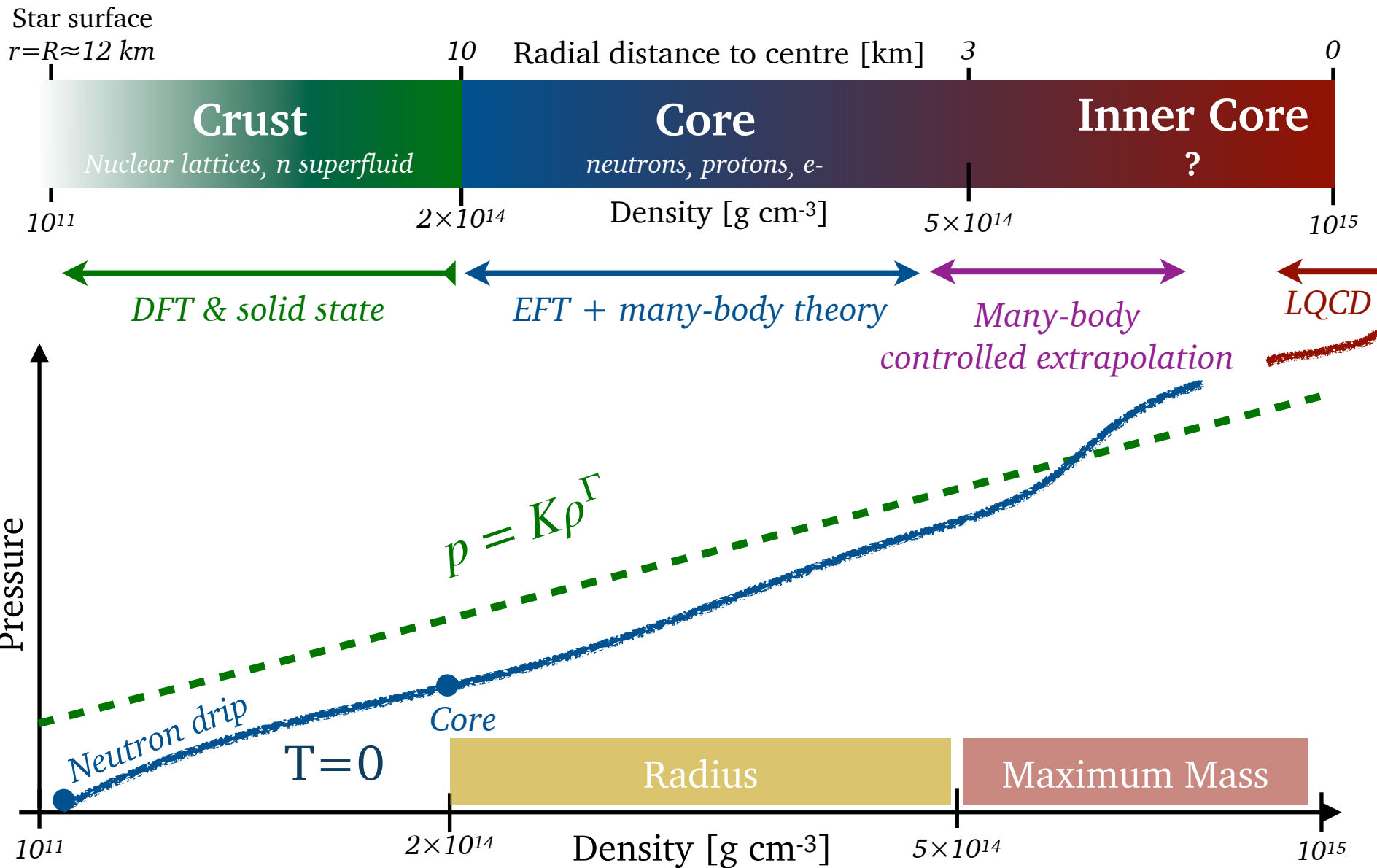


*Thermodynamical Properties of Nuclear Matter
from a Self-Consistent Green's Function Approach*

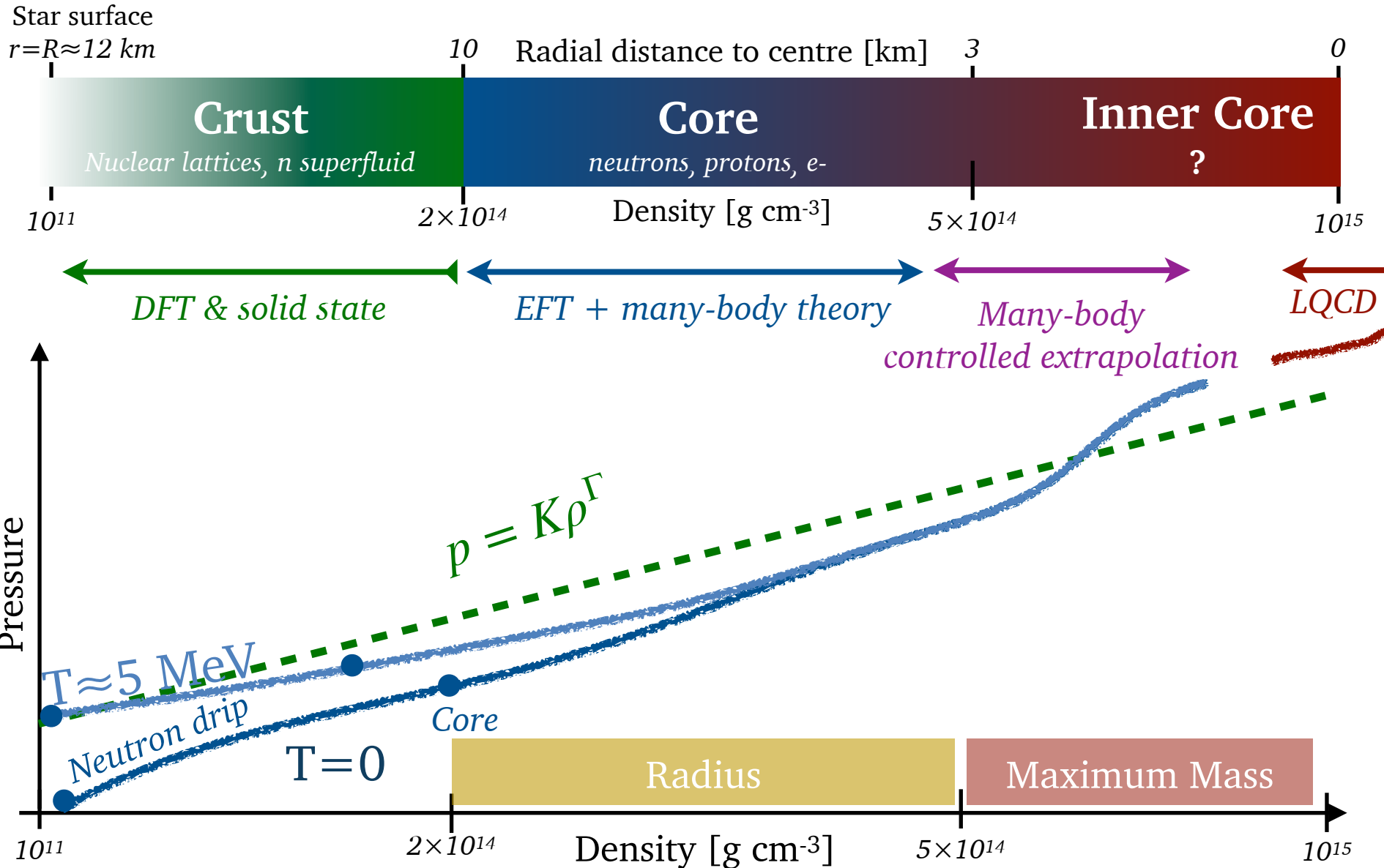
Memòria presentada per
Arnau Rios Huguet
per optar al títol de
Doctor en Ciències Físiques.
Barcelona, 23 febrer de 2007.

Programa de doctorat
Física Avançada
Bienni 2002-2004,
del Departament d'Estructura
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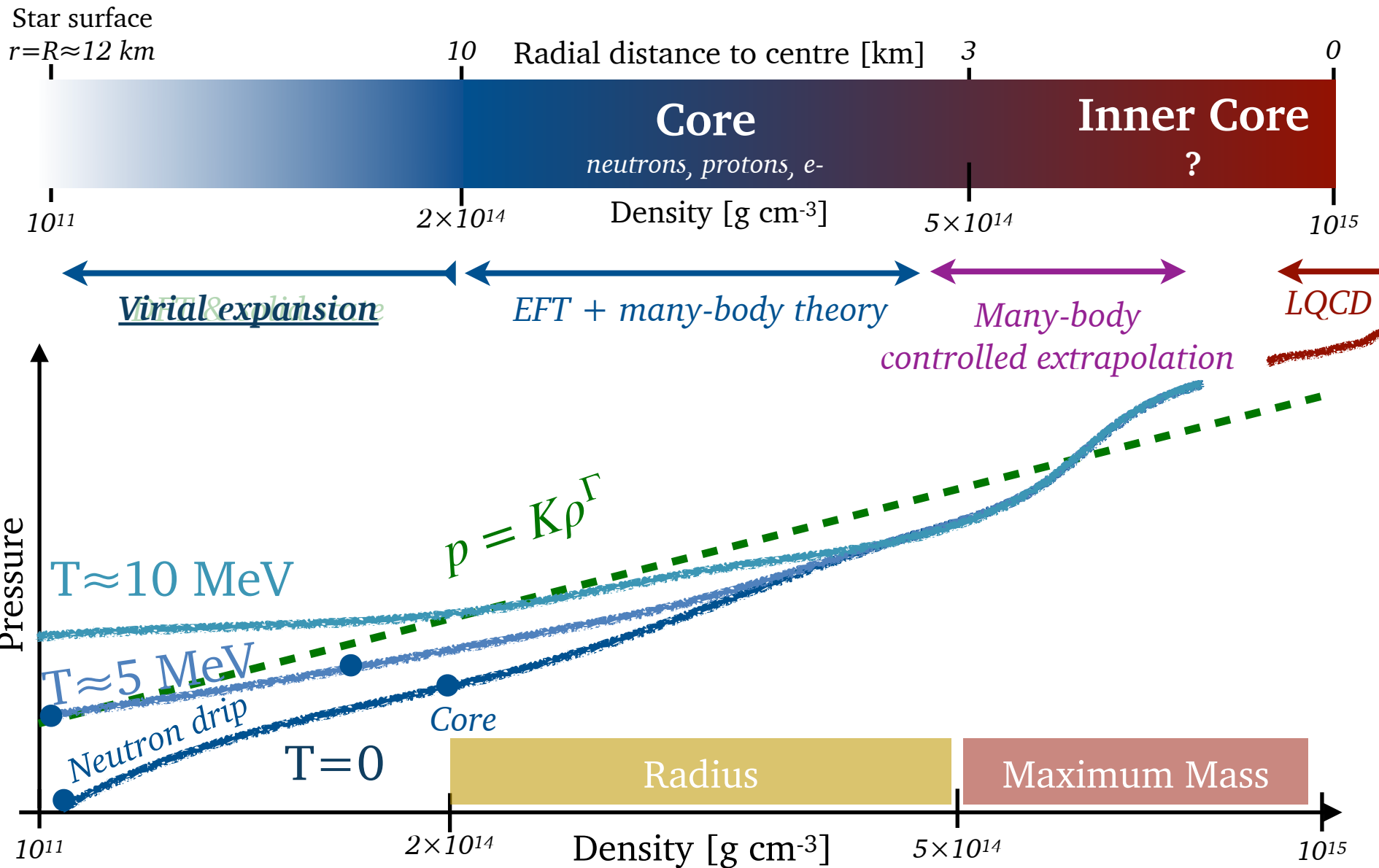
Neutron-star modelling



Neutron-star modelling



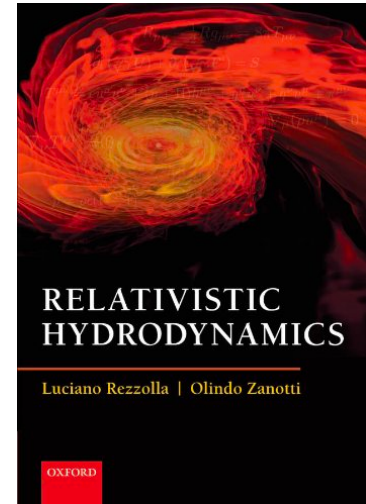
Neutron-star modelling



Simulations require **coupled dynamics** of:

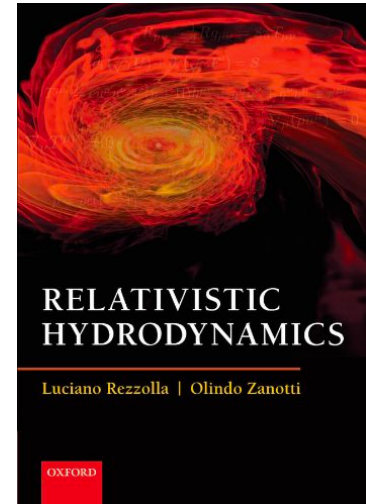
- A fluid which is essentially **cold** (eg NSs before merger)
- A fluid **heated by shocks**, with kinetic energy dissipated into internal energy

$$P(\epsilon, T) = P(\epsilon, T = 0) + P_{\text{th}}(\epsilon_{\text{th}})$$



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Cold part

(Piecewise polytrope)

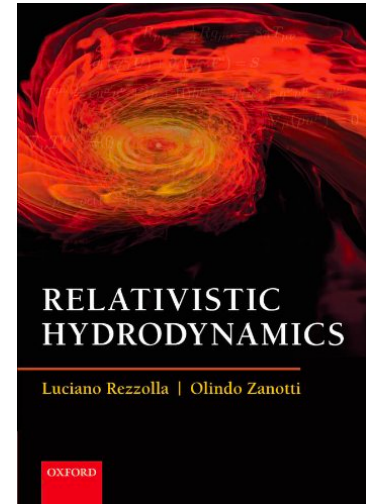
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No thermal energy

Read et al, *Phys Rev D* **79** 124032 (2009)

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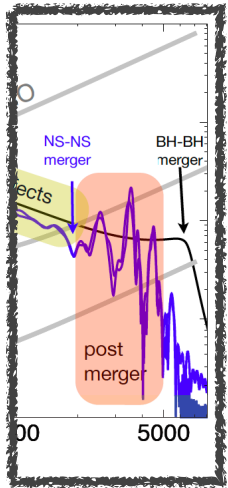
Thermal index

$$P_{\text{th}}(\epsilon_{\text{th}}) \equiv [\Gamma_{\text{th}} - 1] \epsilon_{\text{th}}$$

$$\Gamma_{\text{th}}(\epsilon_{\text{th}}) = 1 + \frac{P_{\text{th}}(\epsilon_{\text{th}})}{\epsilon_{\text{th}}}$$

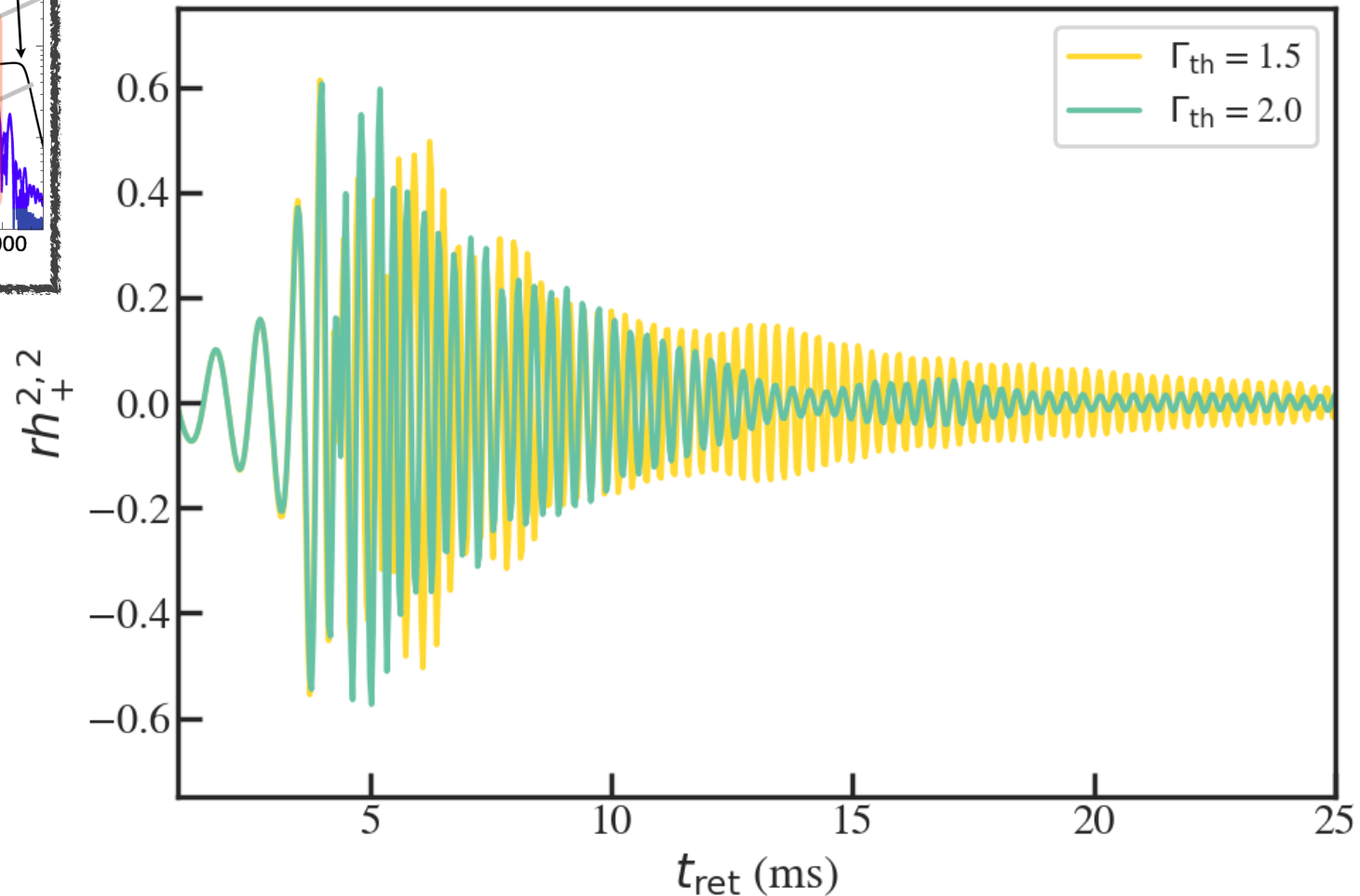
$$\epsilon_{\text{th}} = \epsilon(T) - \epsilon_0$$

Post-merger: from Γ_{th} to GWs



$1.4M_{\odot}+1.4M_{\odot}$ GW spectrum

Same cold component, different constant Γ_{th}



Cold part

Non-relativistic

$$\Gamma_c = \frac{5}{3}$$



Thermal part

$$\Gamma_{\text{th}} = \frac{5}{3}$$

Relativistic

$$\Gamma_c = \frac{4}{3}$$



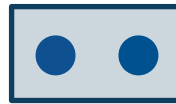
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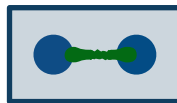


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Unitary

$$k_{FaS} \gg 1$$

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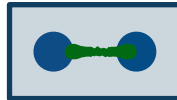


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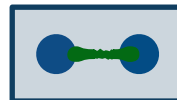
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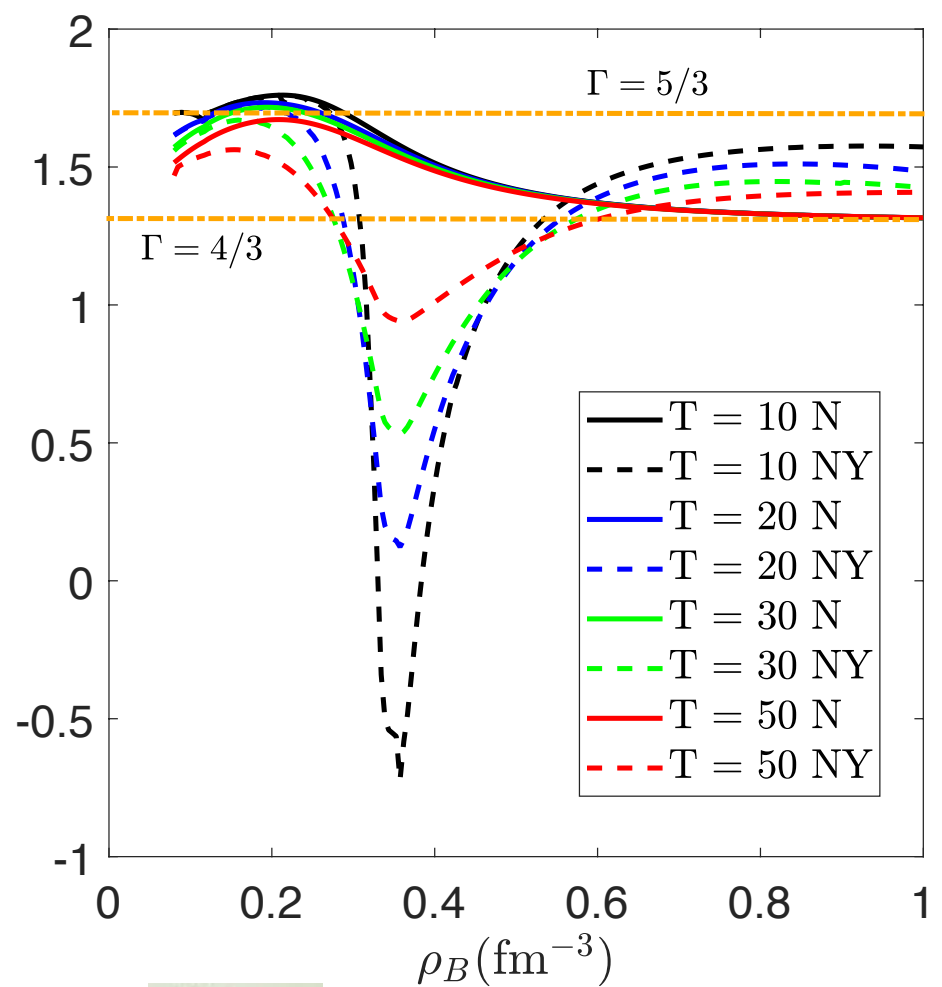
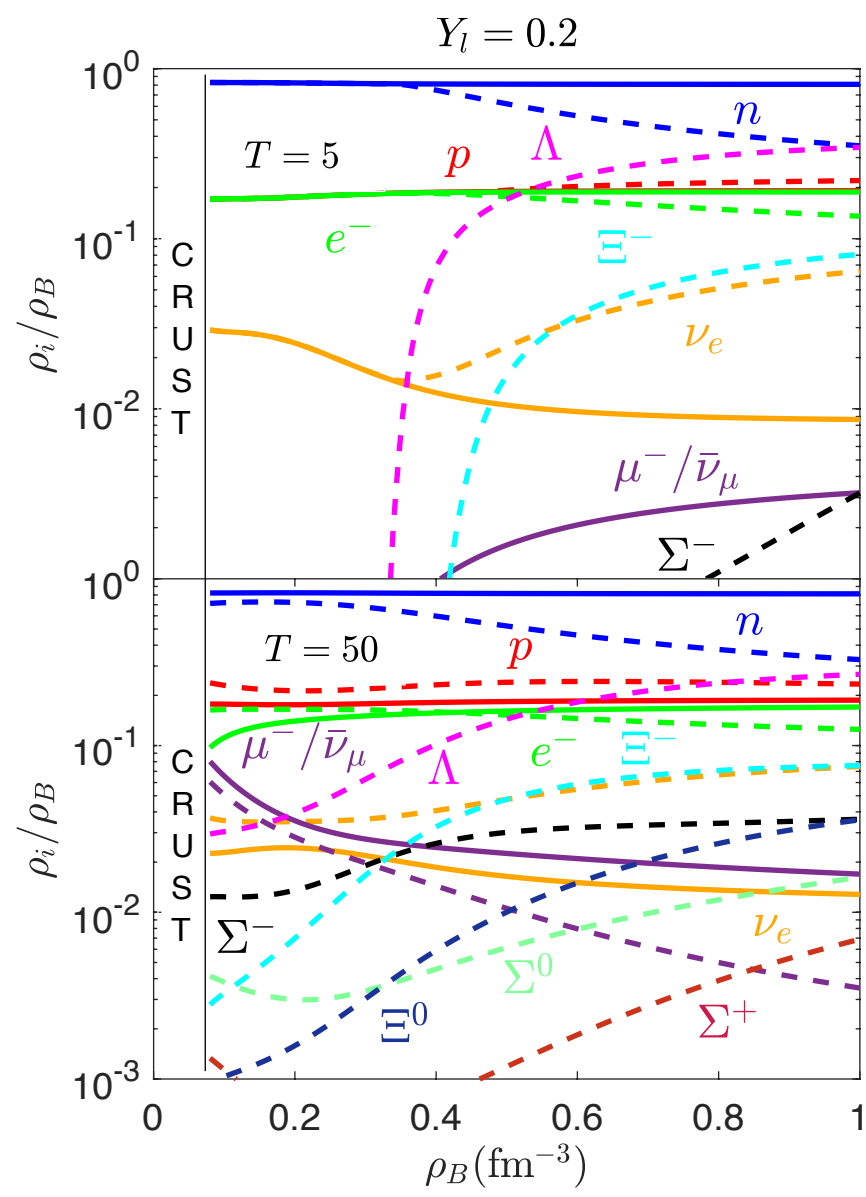
Mean-field/Fermi liquid

$$\Gamma_c = \Gamma_c(\epsilon)$$

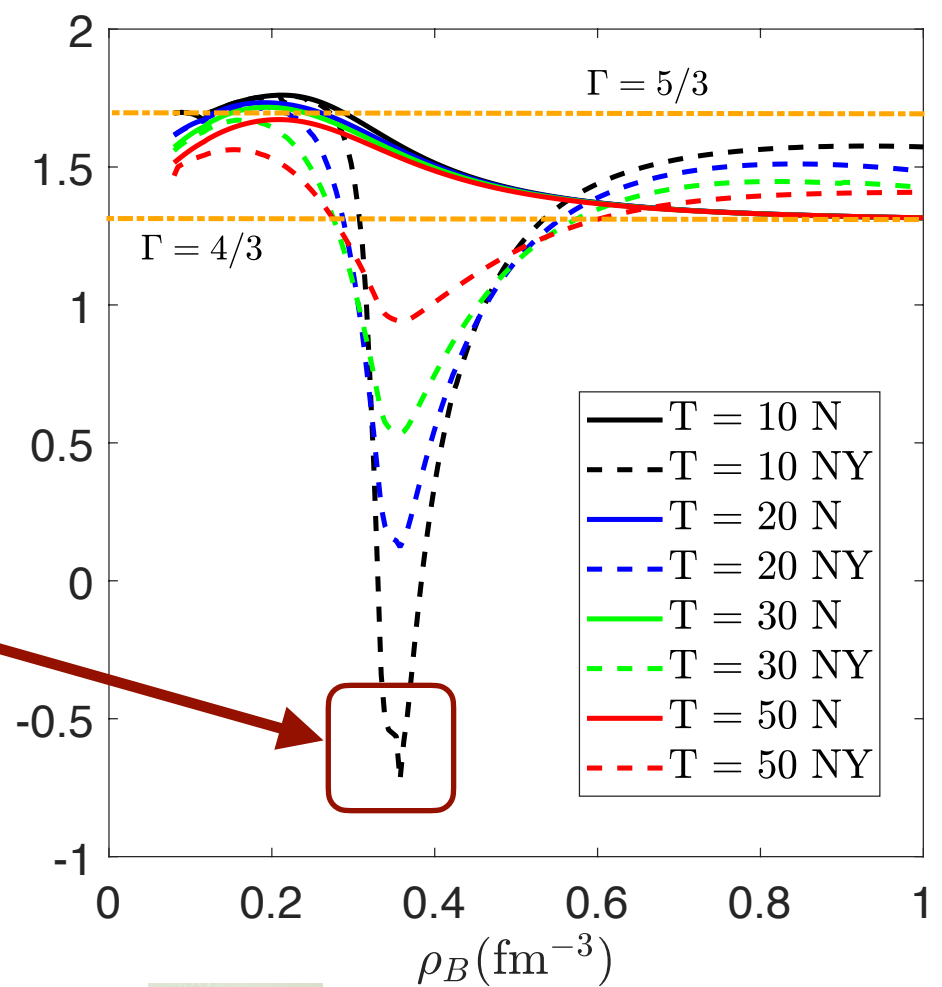
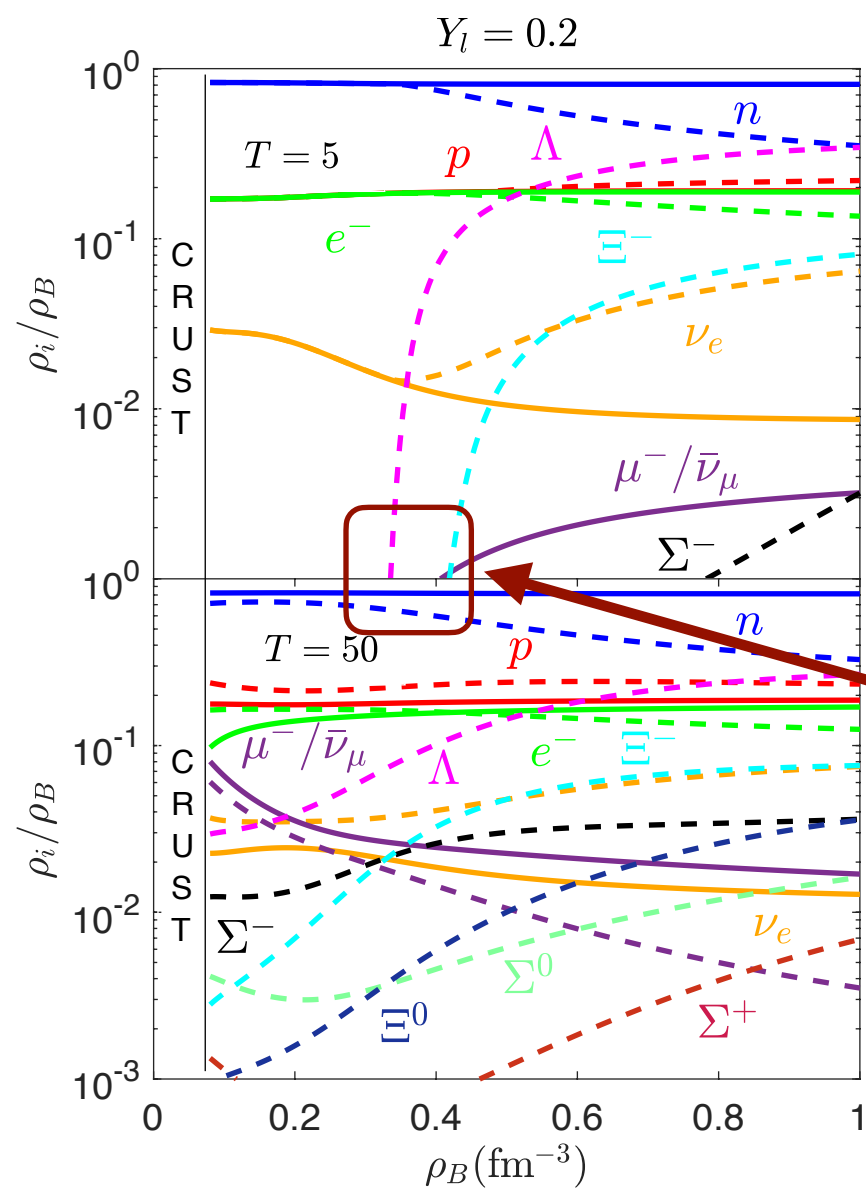


$$\Gamma_{\text{th}}^{m^*} = \frac{5}{3} - \frac{n}{m^*} \frac{\partial m^*}{\partial \rho}$$

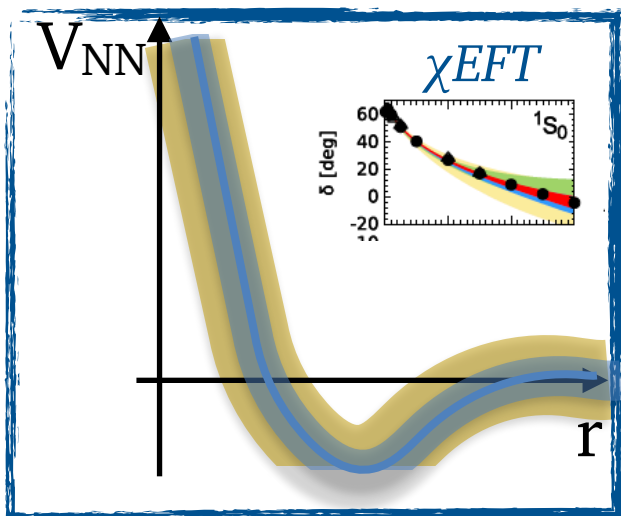
Multicomponent systems: hyperons



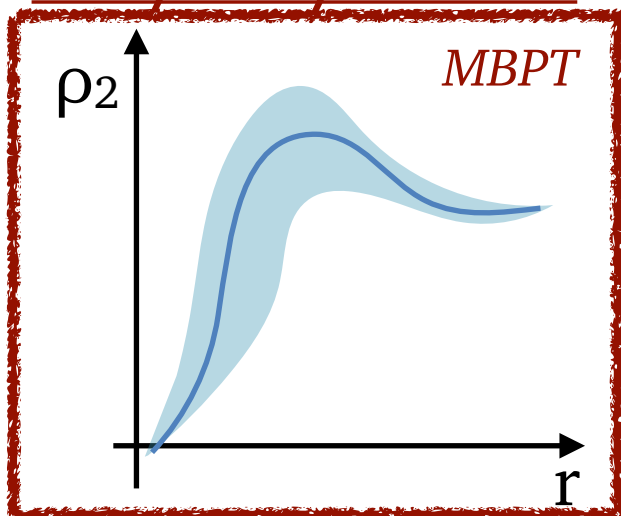
Multicomponent systems: hyperons



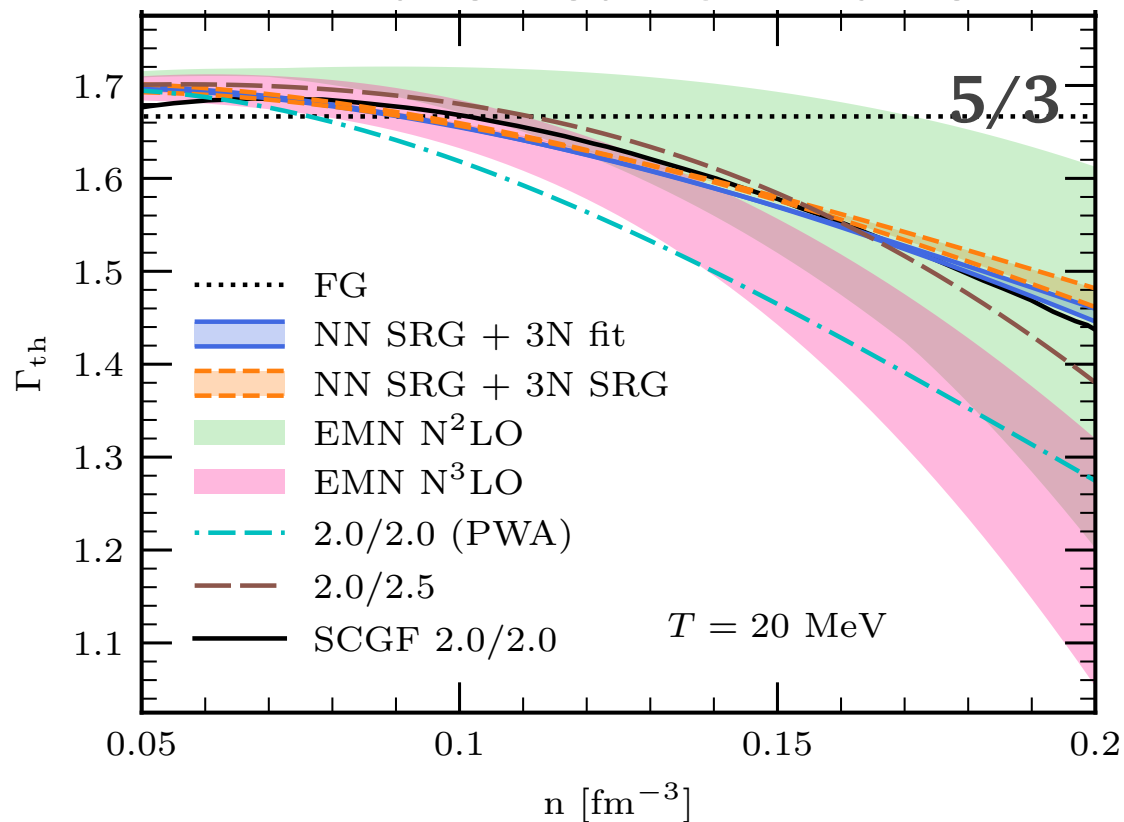
Hamiltonian



Many-body method

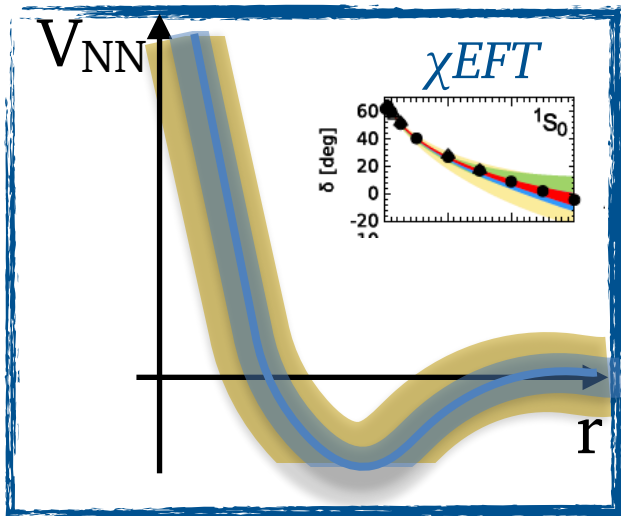


Pure neutron matter

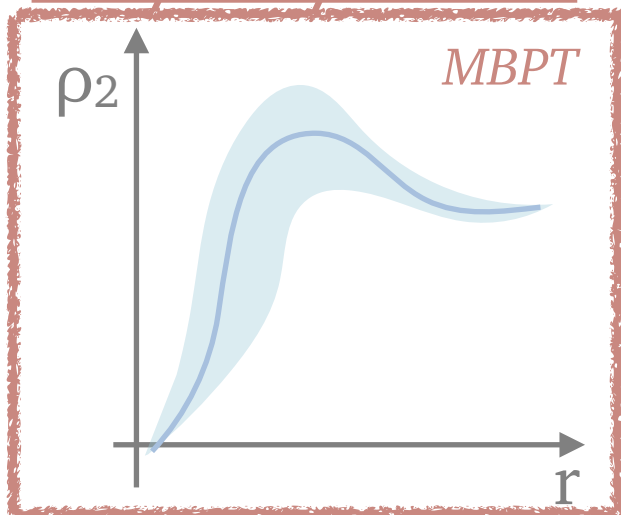


Keller, Wellenhofer, Hebeler & Schwenk *Phys Rev C* **103**, 055806 (2021)

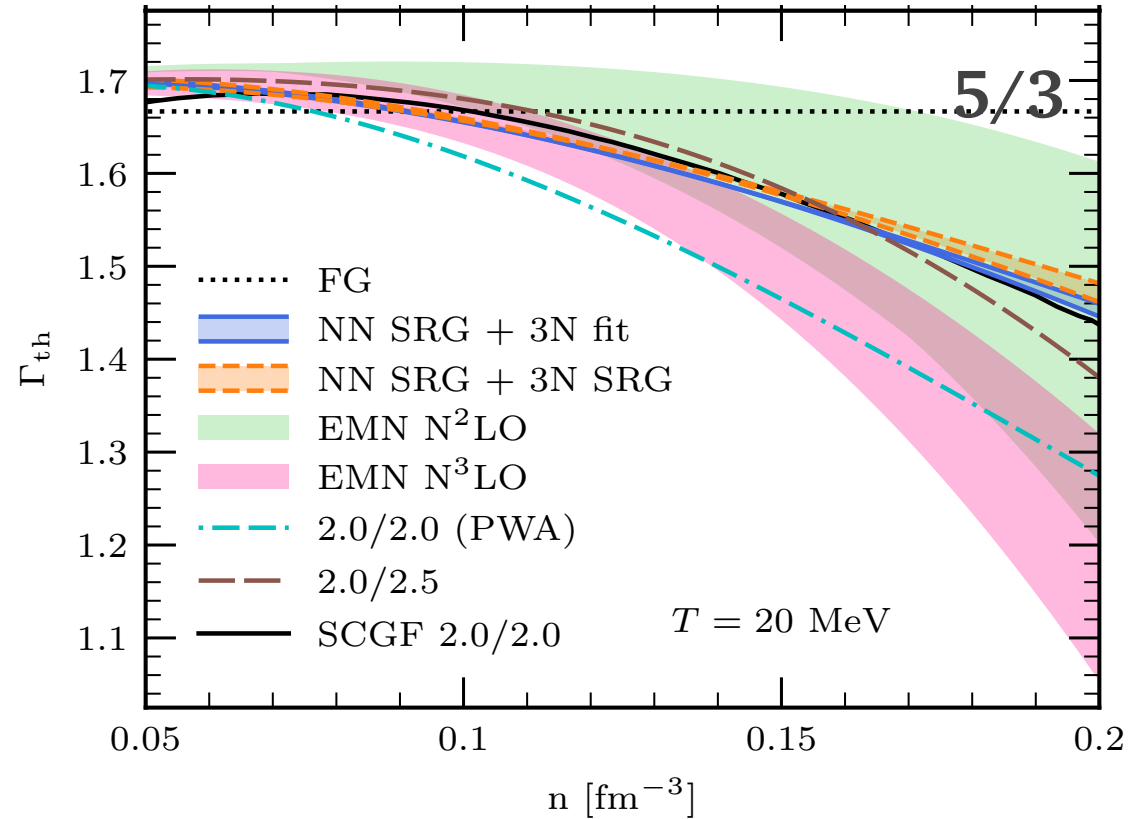
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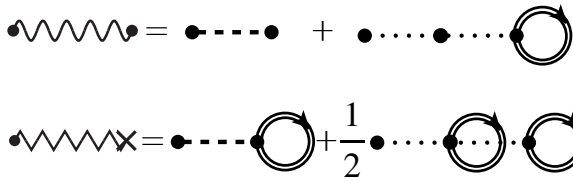


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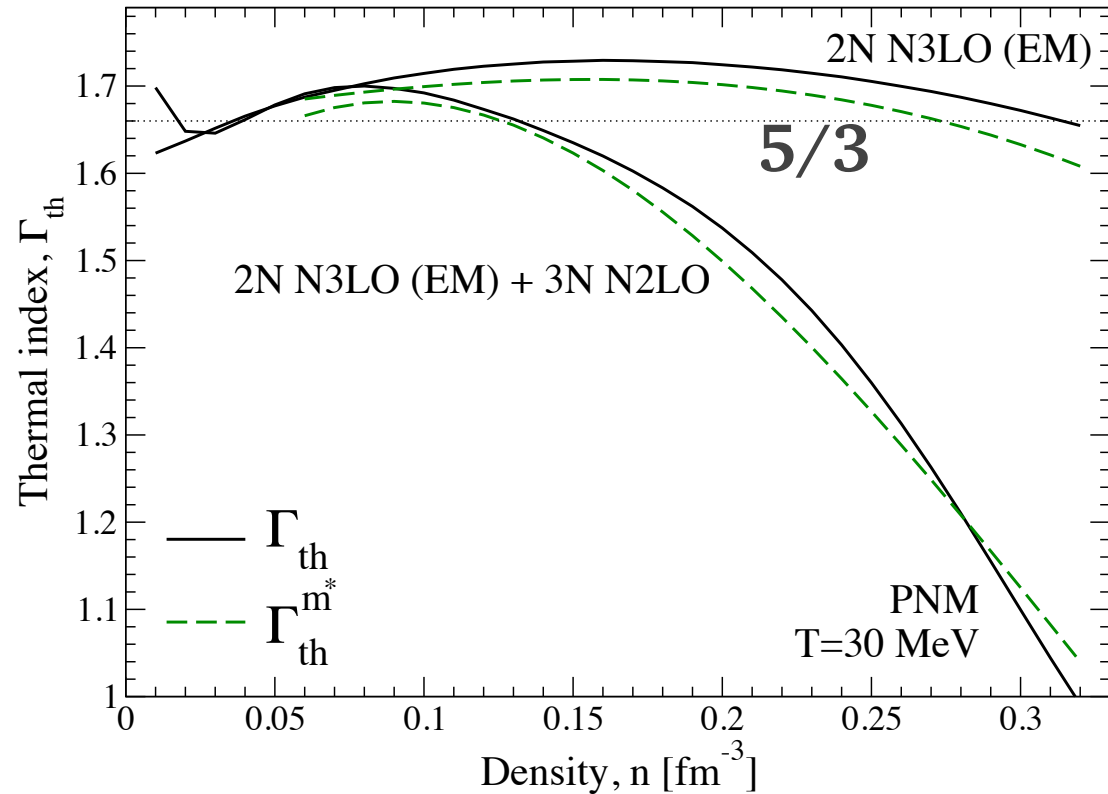


Keller, Wellenhofer, Hebeler & Schwenk *Phys Rev C* **103**, 055806 (2021)

2N & 3N forces



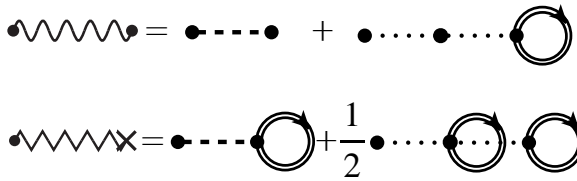
Pure neutron matter



Carbone & Schwenk *Phys Rev C* **100** 025805 (2019)

Carbone, Polls & Rios, *Phys Rev C* **88** 044302 (2014)
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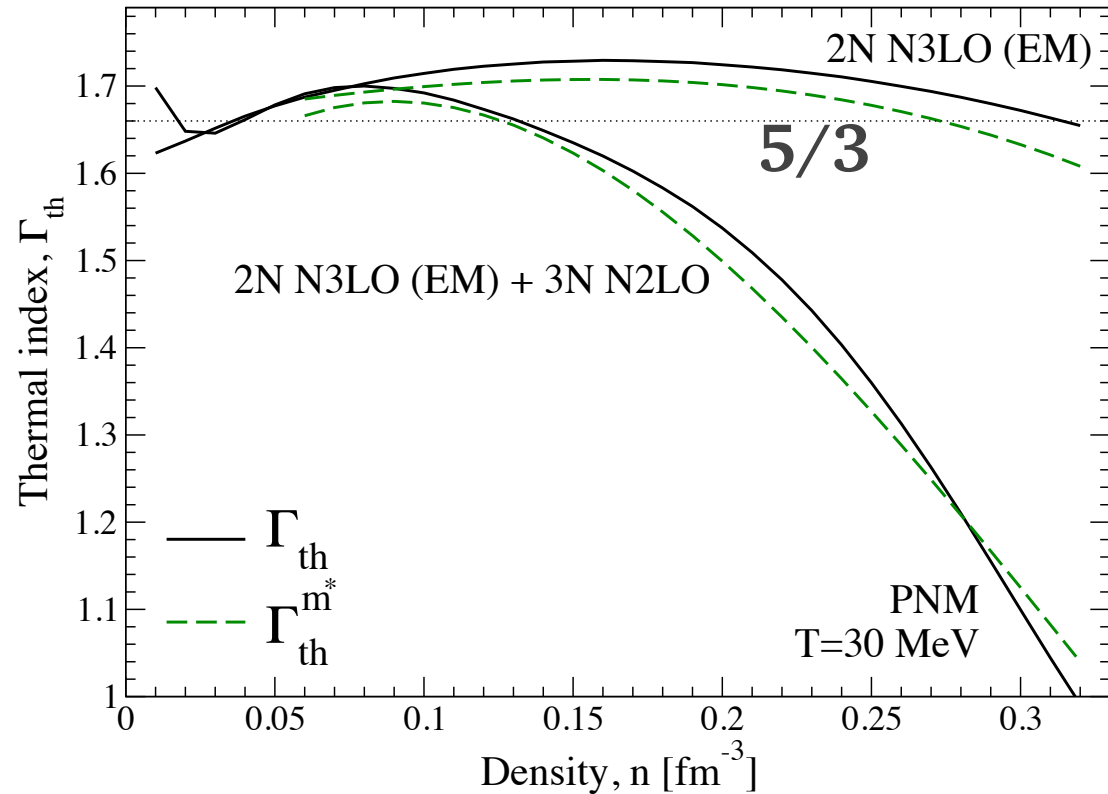
2N & 3N forces



In-medium interaction



Pure neutron matter

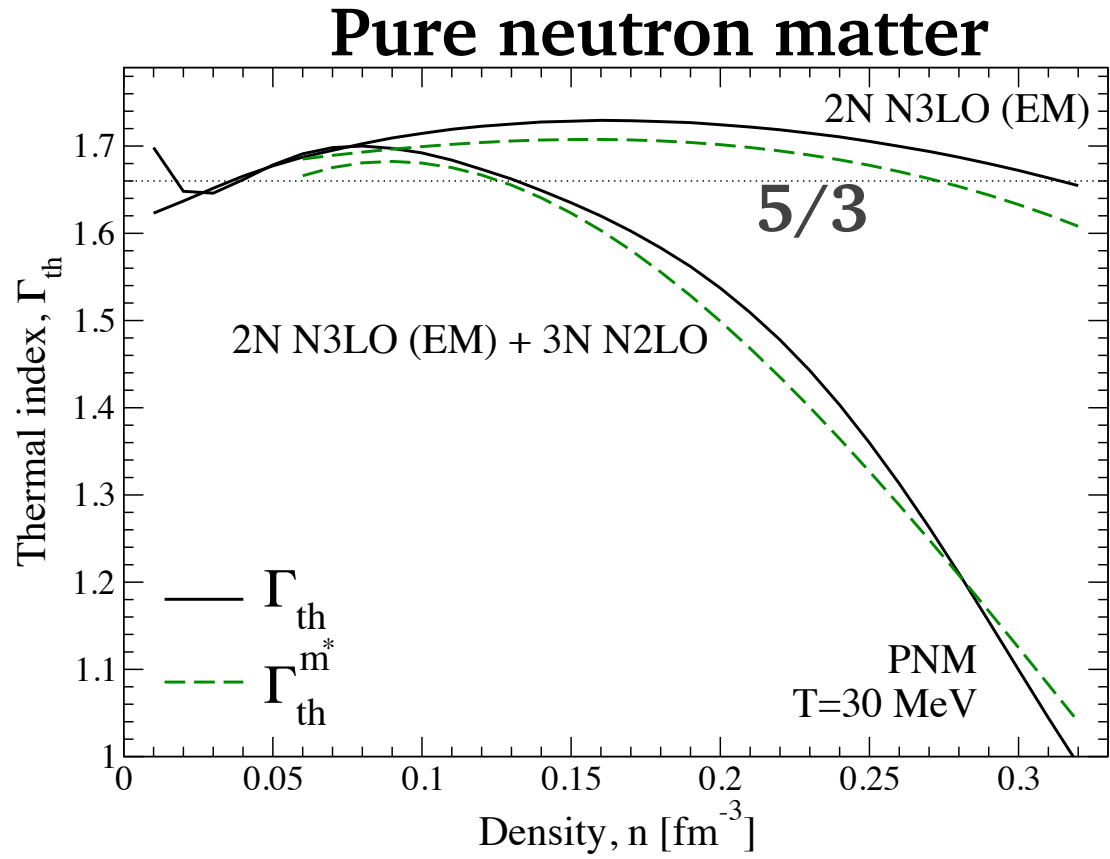
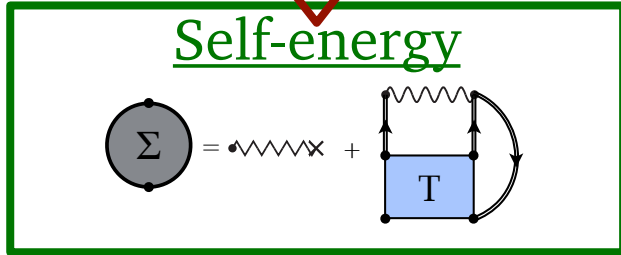
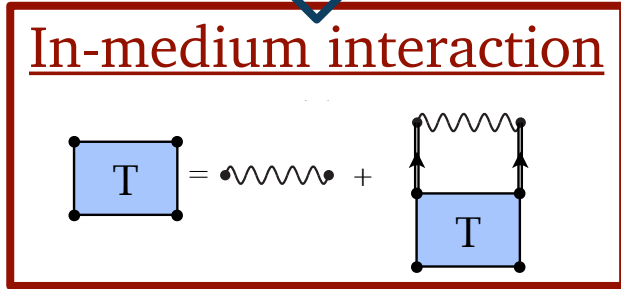
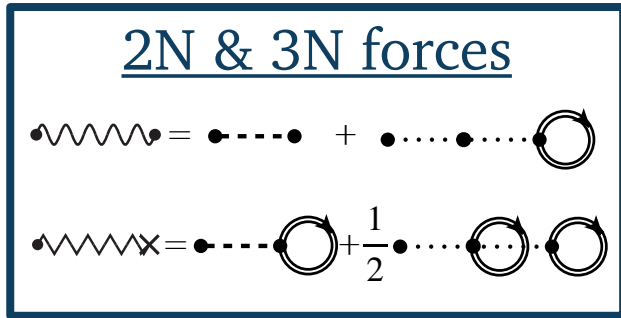


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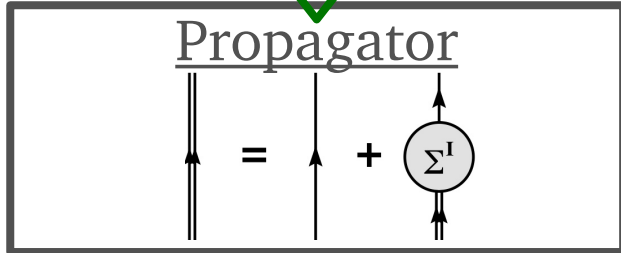
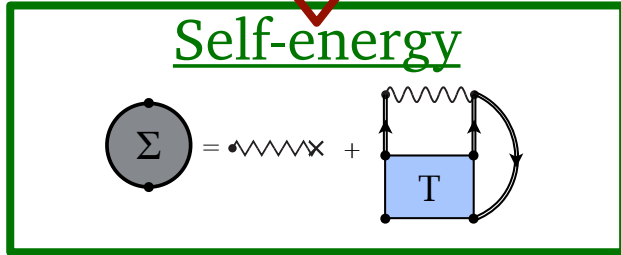
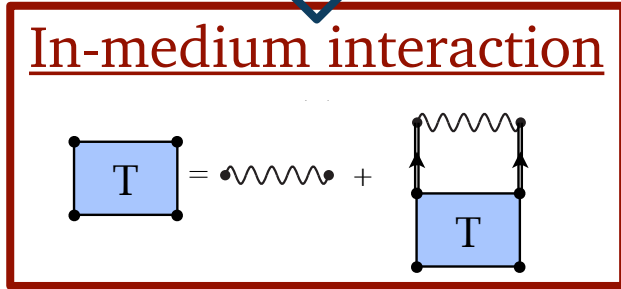
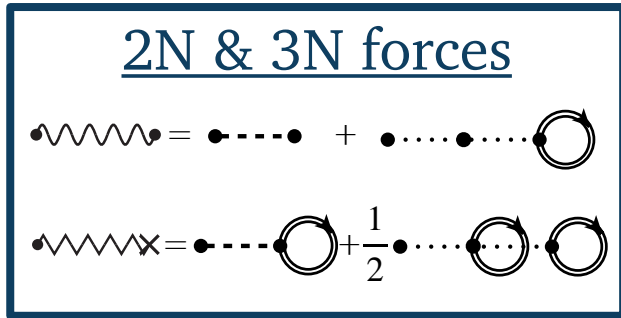
Ab initio thermal index



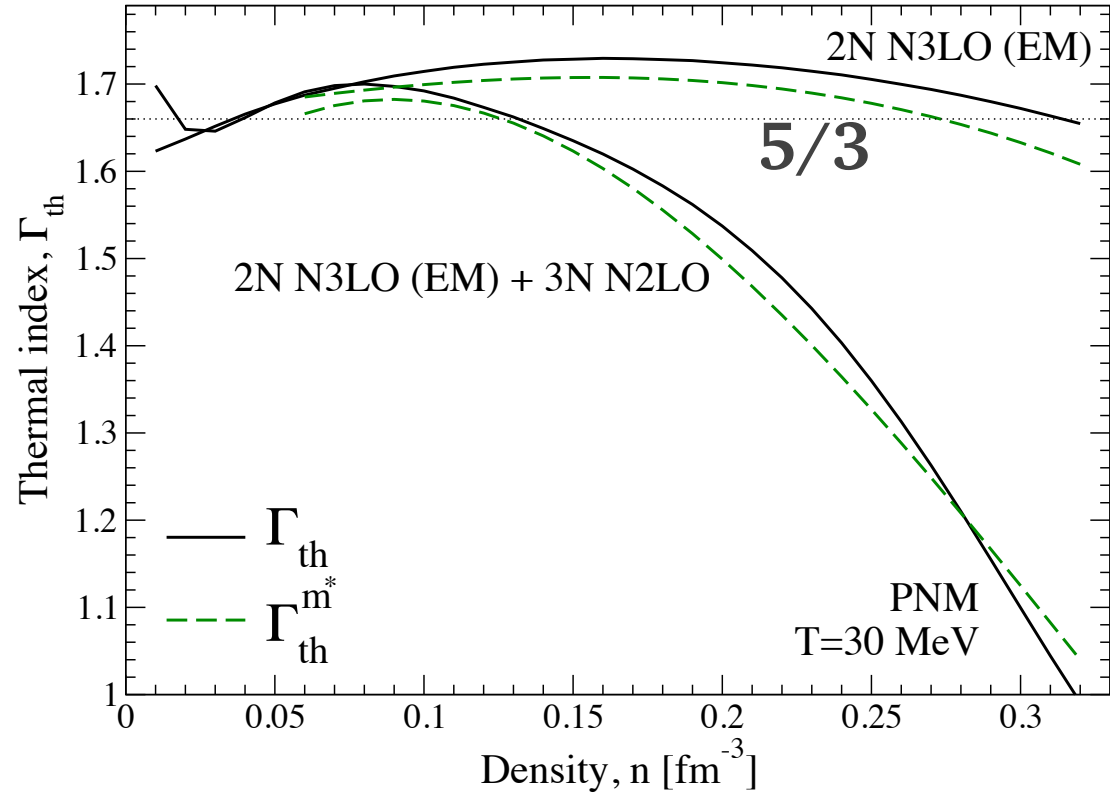
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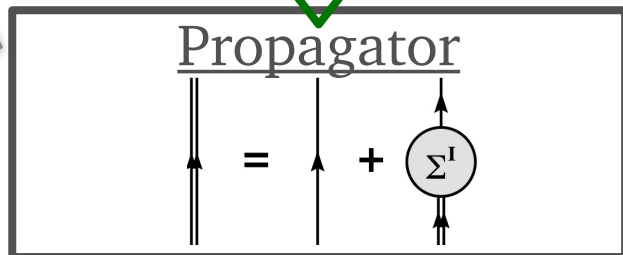
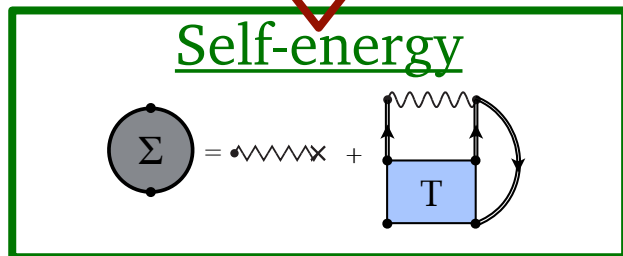
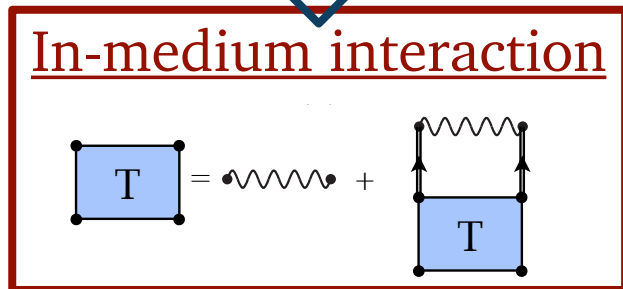
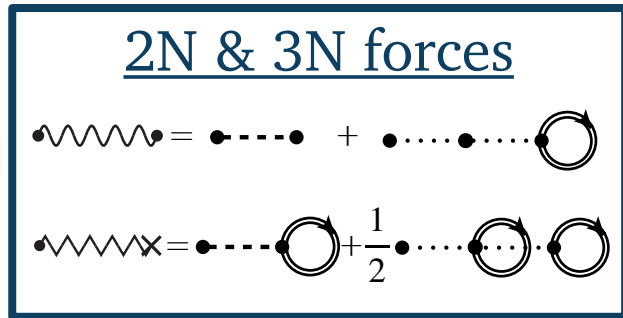
Pure neutron matter



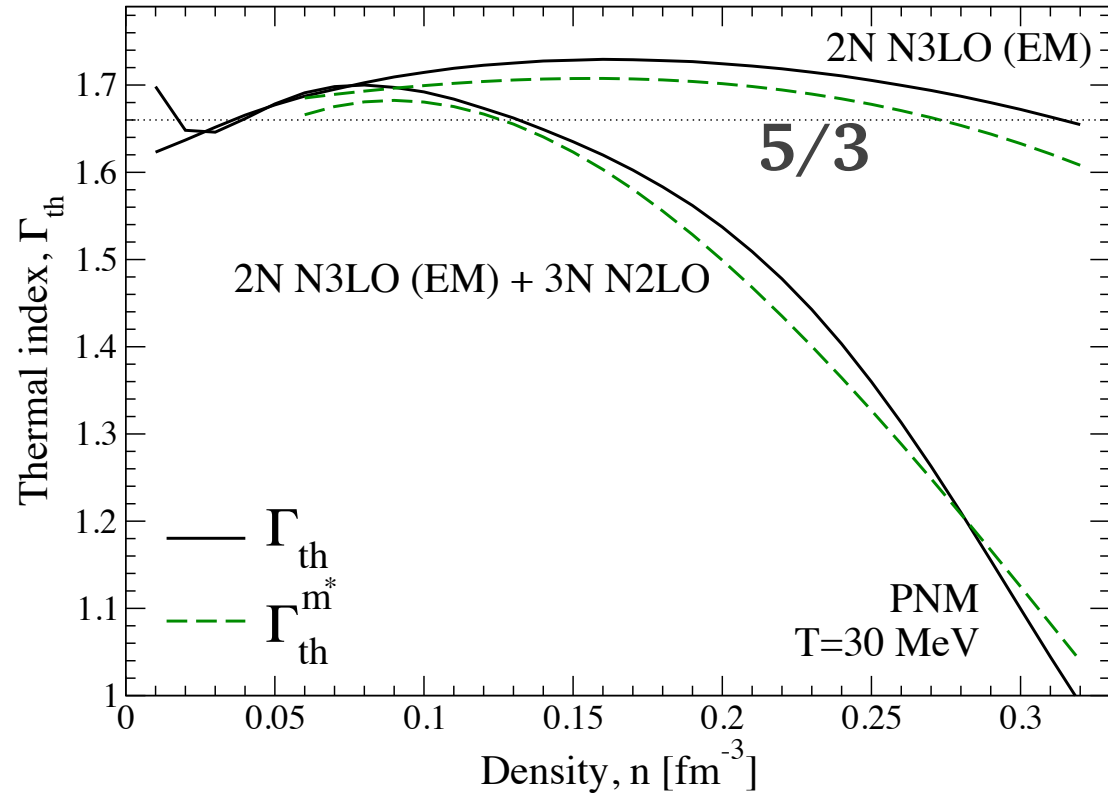
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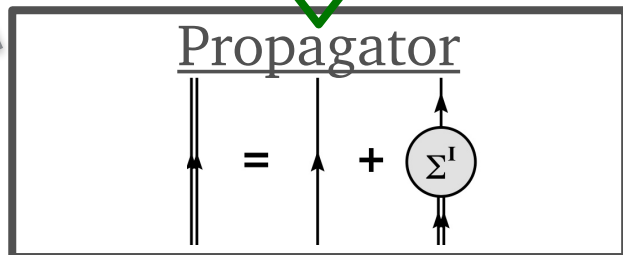
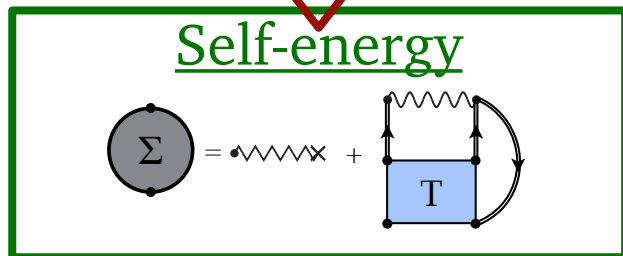
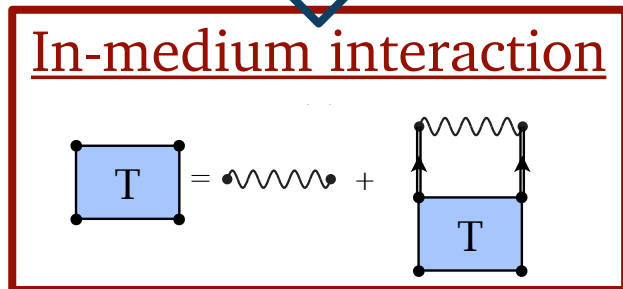
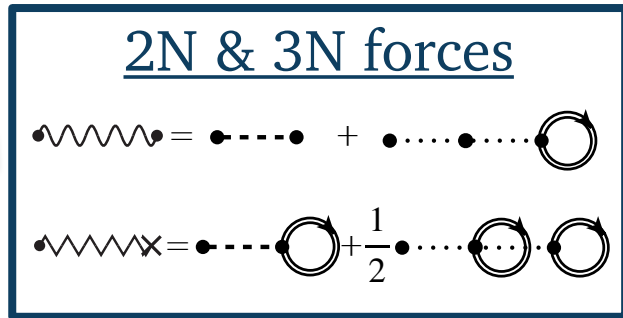
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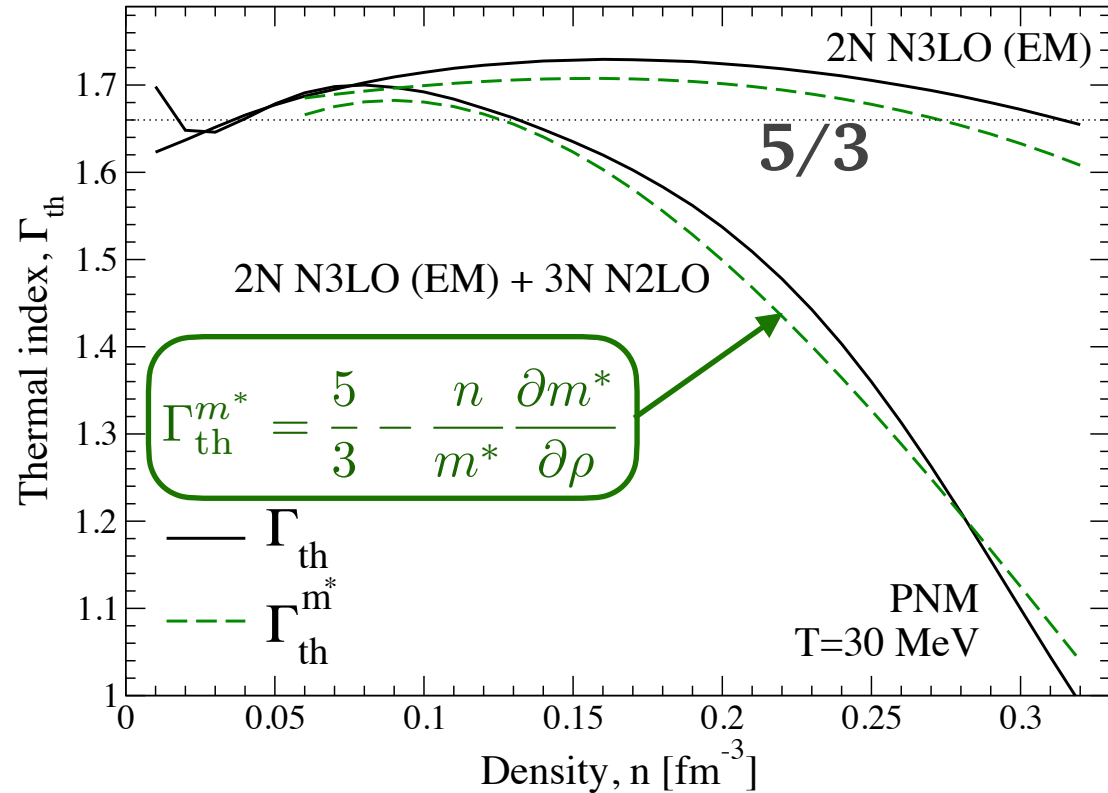
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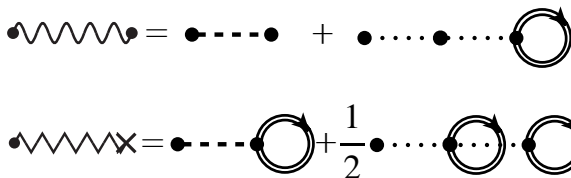
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Self-Consistent Green's Functions

 (ρ, T)

2N & 3N forces

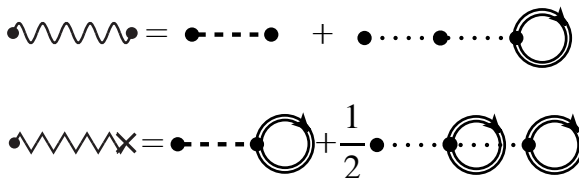


Carbone, Rios & Polls PRC **88** 044302 (2013);
 PRC **90**, 054322 (2014);
 Carbone PhD Thesis

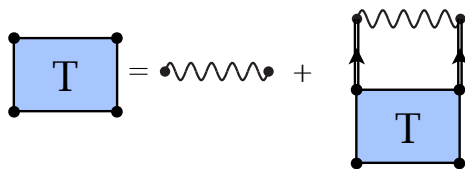
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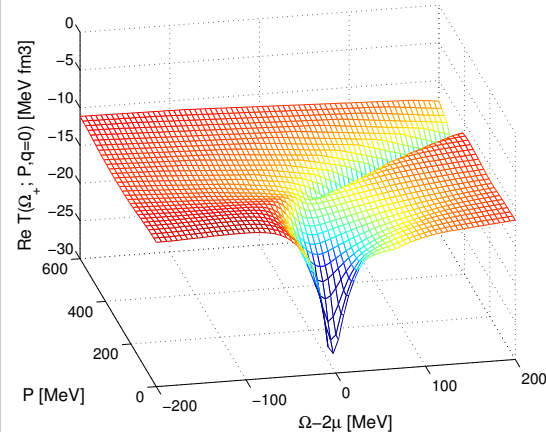


In-medium interaction



Carbone, Rios & Polls PRC **88** 044302 (2013);
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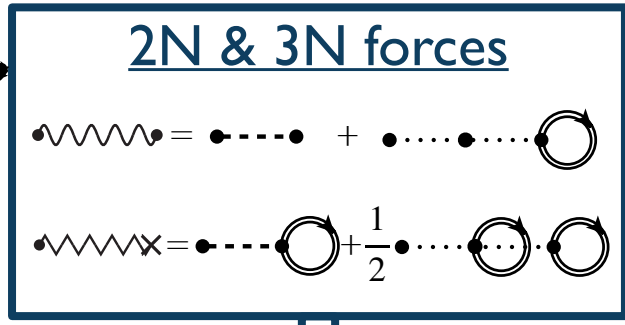
T-matrix at $T=5$ MeV



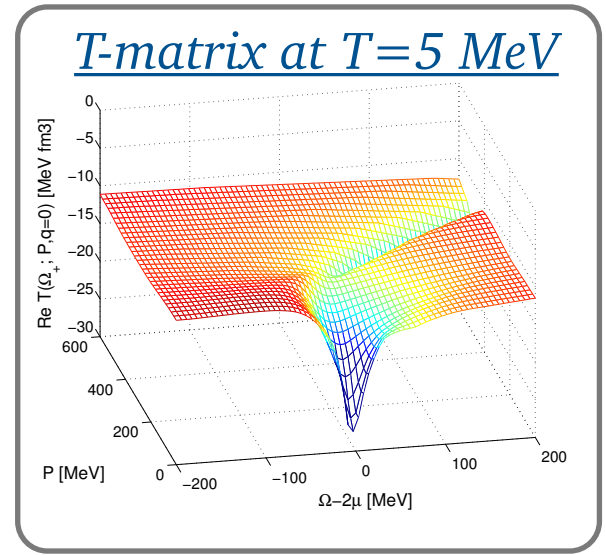
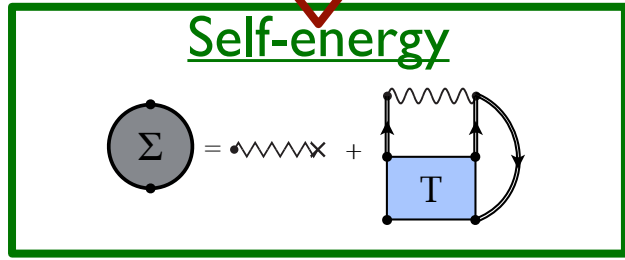
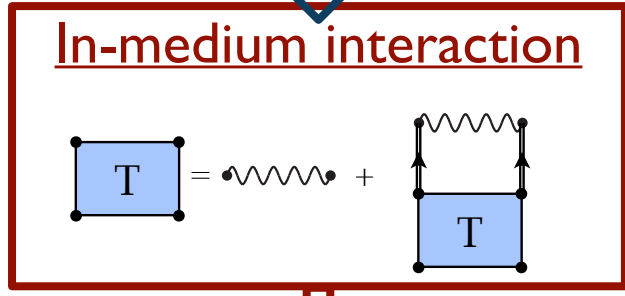
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 Dewulf *et al.*, PRL **90** 152501 (2003)
 Frick & Muther, PRC **68** 034310 (2003)
 Rios, PhD Thesis, U. Barcelona (2007)
 Soma & Bozek, PRC **78** 054003 (2008)
 Rios & Soma PRL **108** 012501 (2012)

Self-Consistent Green's Functions

(ρ, T)



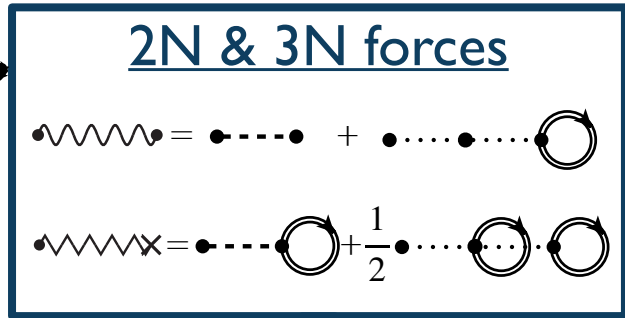
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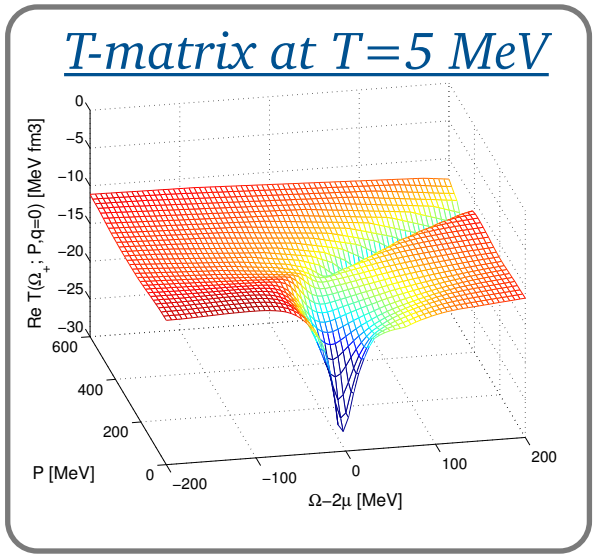
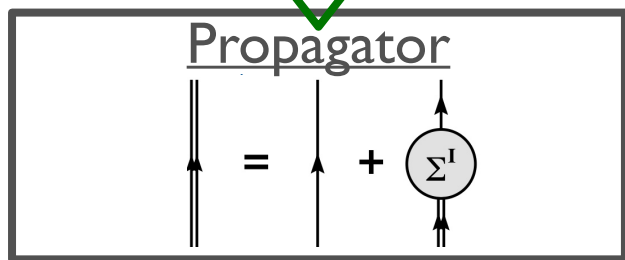
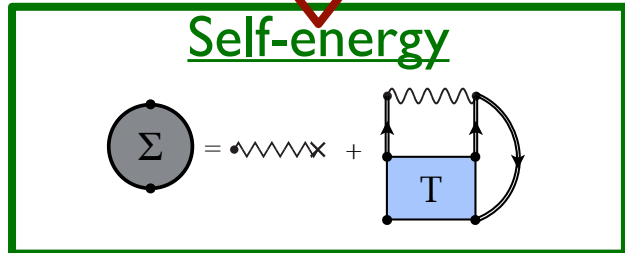
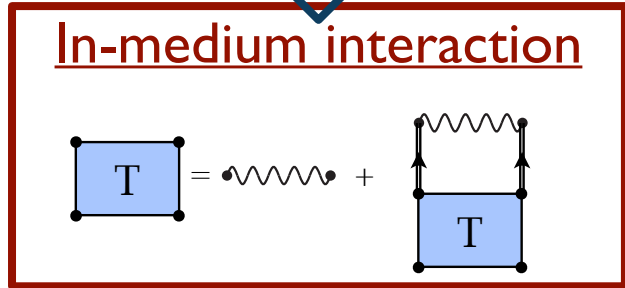
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Self-Consistent Green's Functions

(ρ, T)



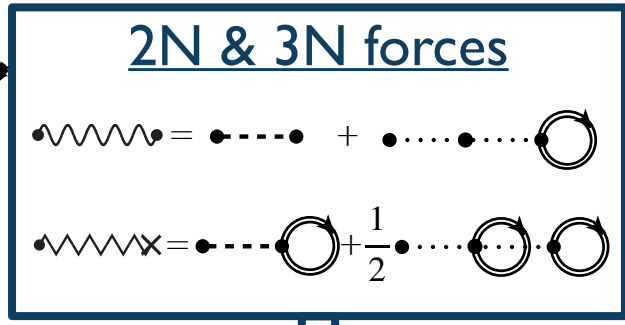
Carbone, Rios & Polls PRC **88** 044302 (2013);
 PRC **90**, 054322 (2014);
 Carbone PhD Thesis



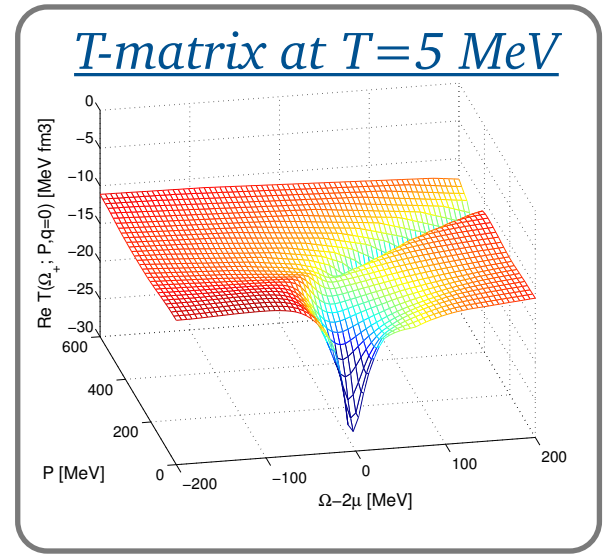
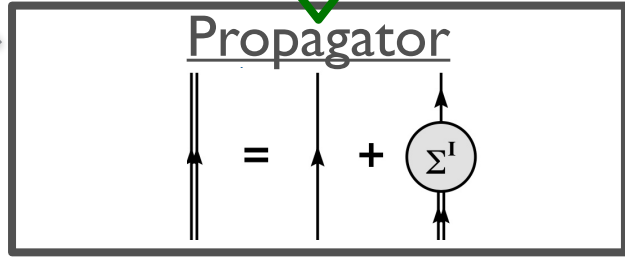
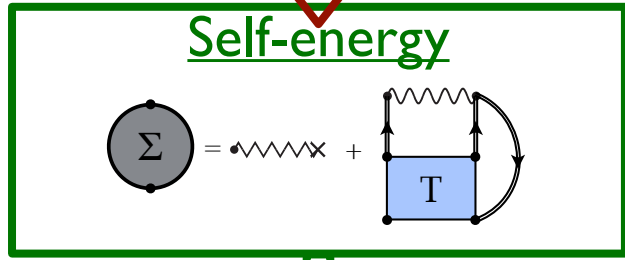
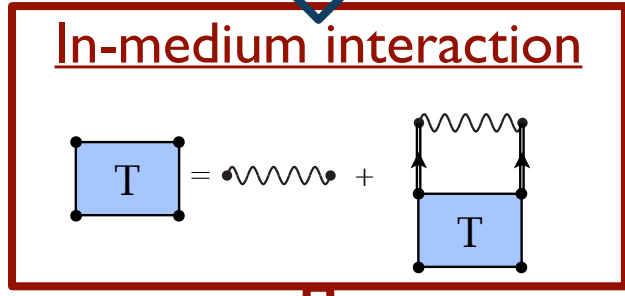
Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
 Alm et al., PRC **53** 2181 (1996)
 Dewulf et al., PRL **90** 152501 (2003)
 Frick & Muther, PRC **68** 034310 (2003)
 Rios, PhD Thesis, U. Barcelona (2007)
 Soma & Bozek, PRC **78** 054003 (2008)
 Rios & Soma PRL **108** 012501 (2012)

Self-Consistent Green's Functions

(ρ, T)



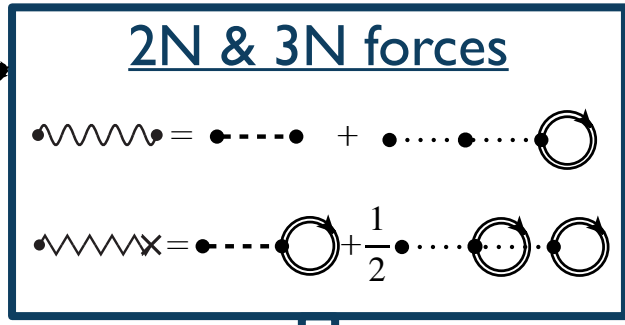
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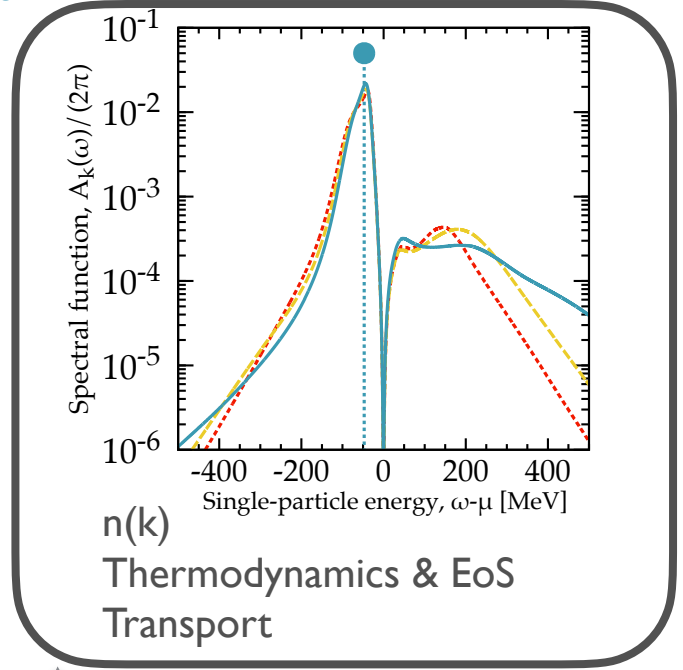
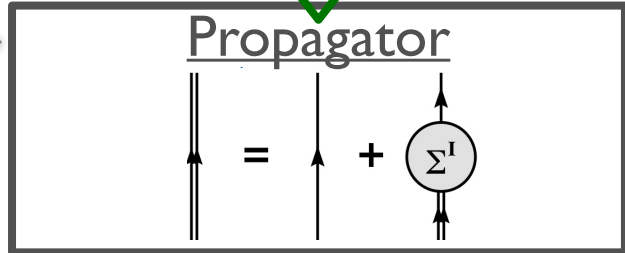
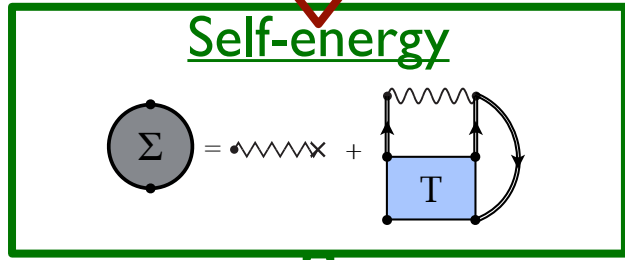
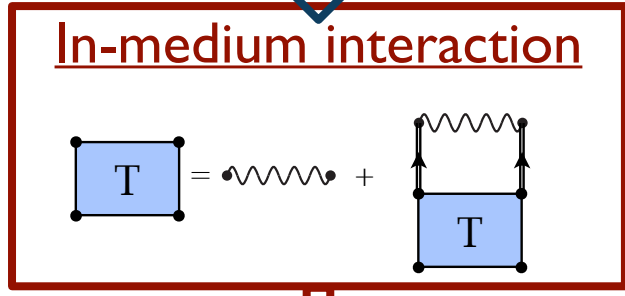
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Self-Consistent Green's Functions

(ρ, T)



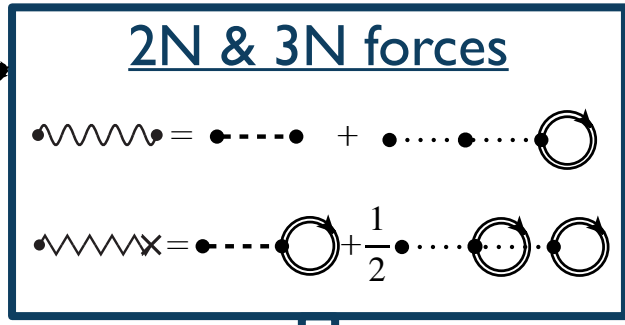
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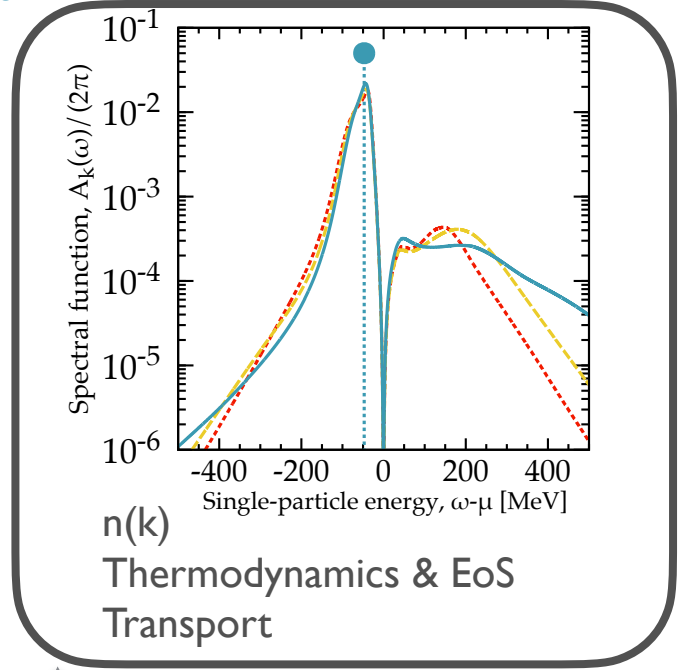
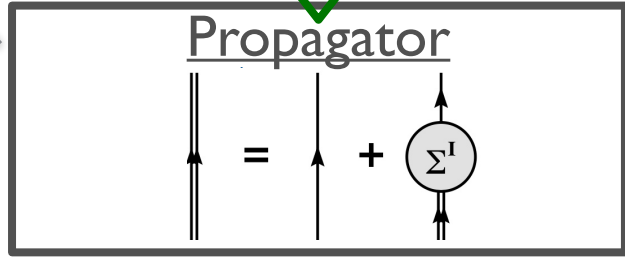
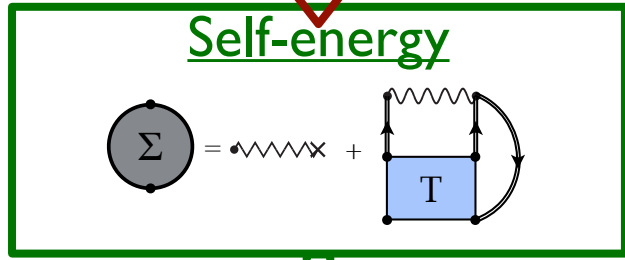
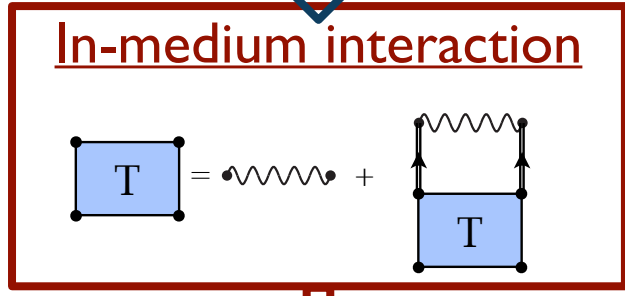
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Self-Consistent Green's Functions

(ρ, T)



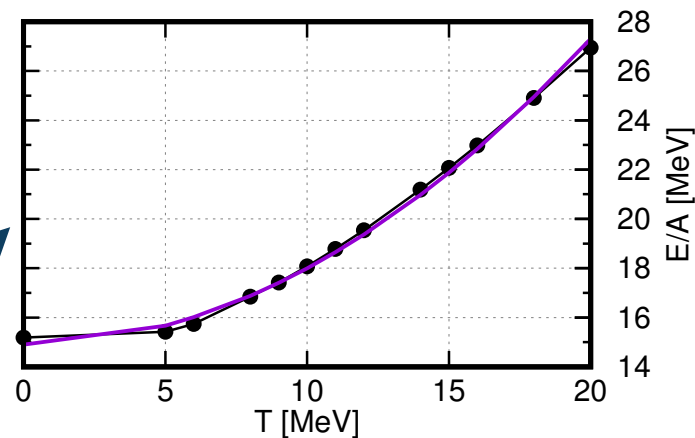
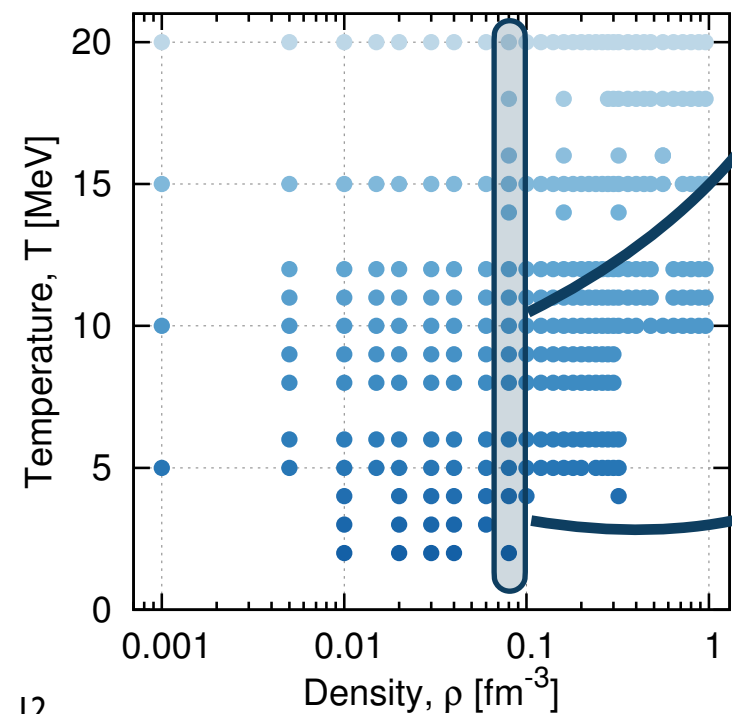
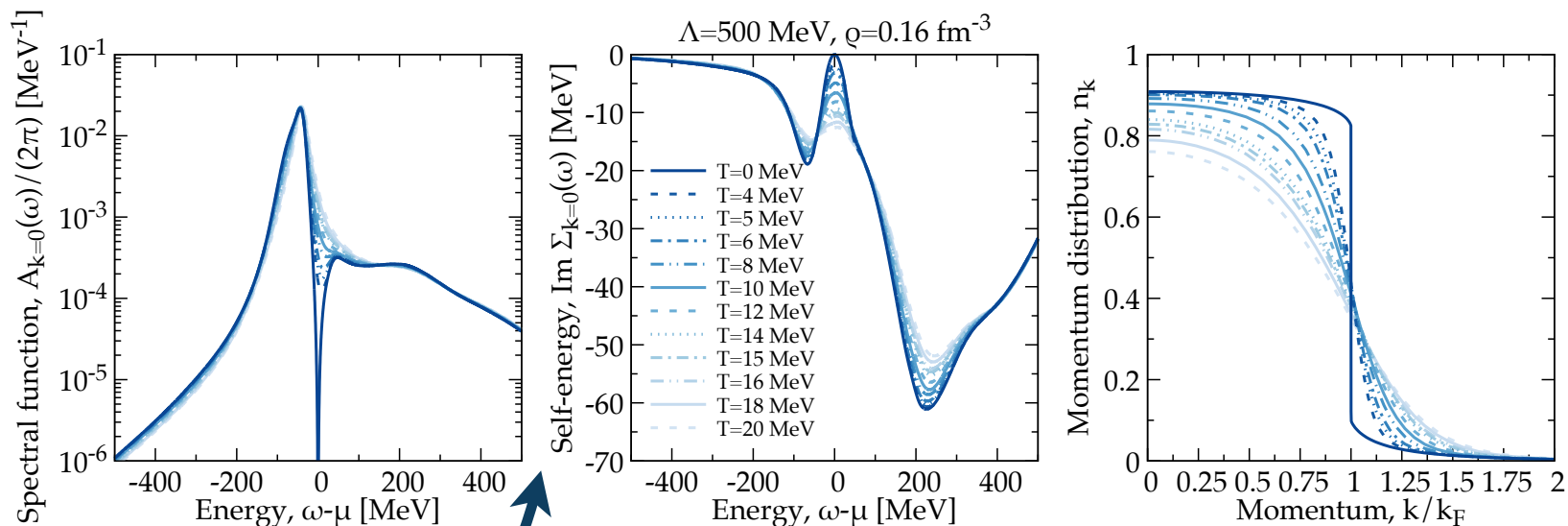
Carbone, Rios & Polls PRC **88** 044302 (2013);
 PRC **90**, 054322 (2014);
 Carbone PhD Thesis

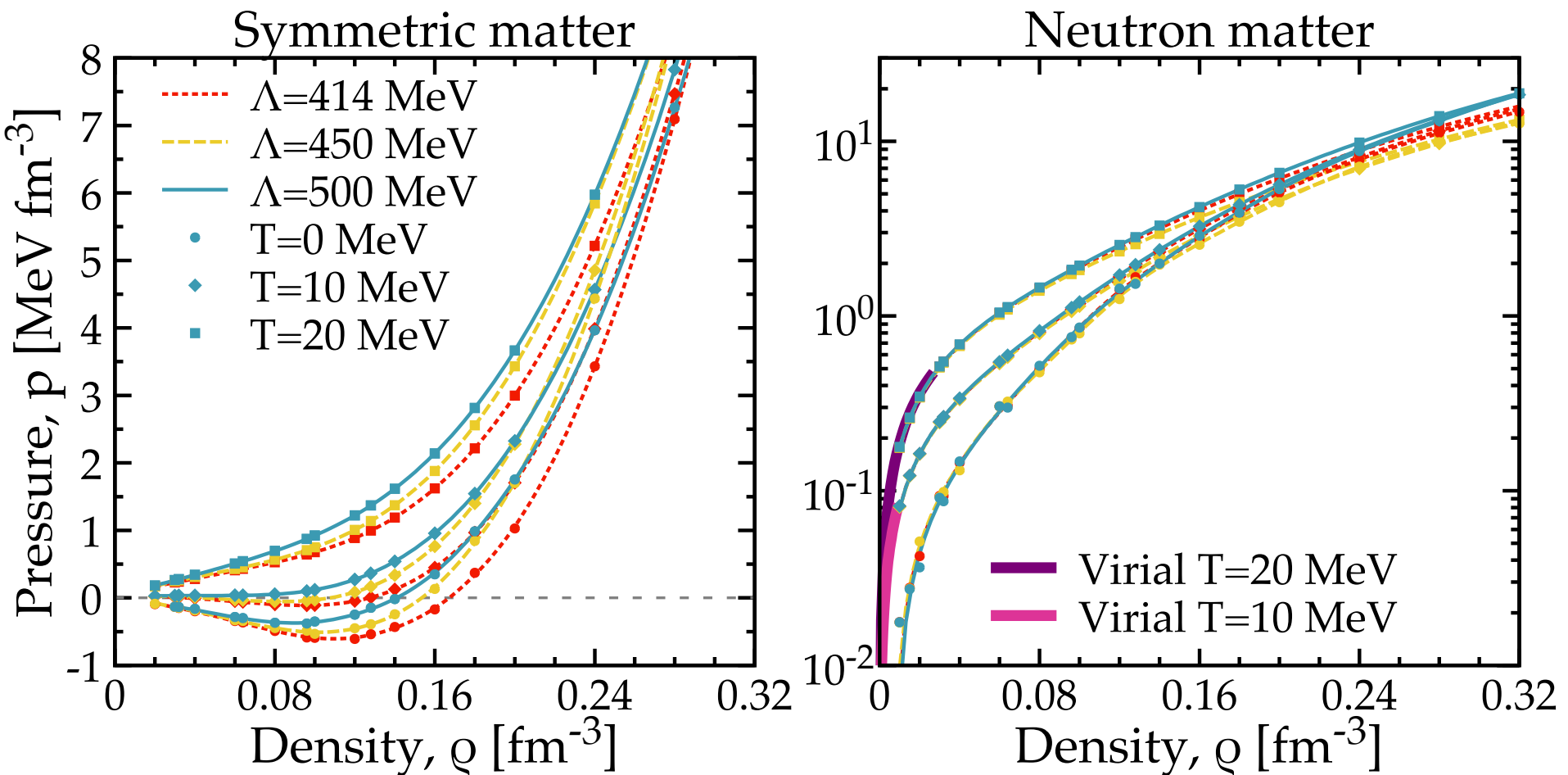


- Off-shell ✓
- Matsubara formalism ✓
- Φ -derivable ✓

Ramos, Polls & Dickhoff, NPA **503** 1 (1989)
 Alm et al., PRC **53** 2181 (1996)
 Dewulf et al., PRL **90** 152501 (2003)
 Frick & Muther, PRC **68** 034310 (2003)
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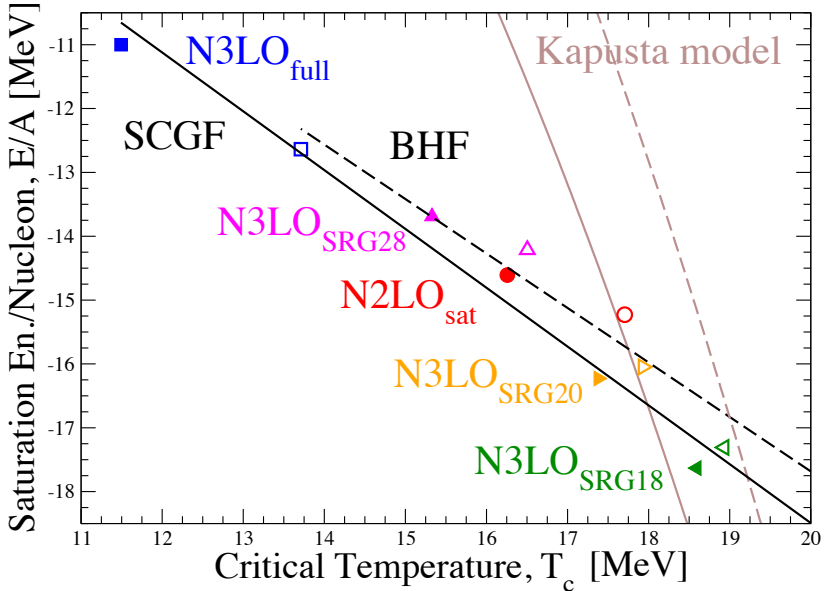
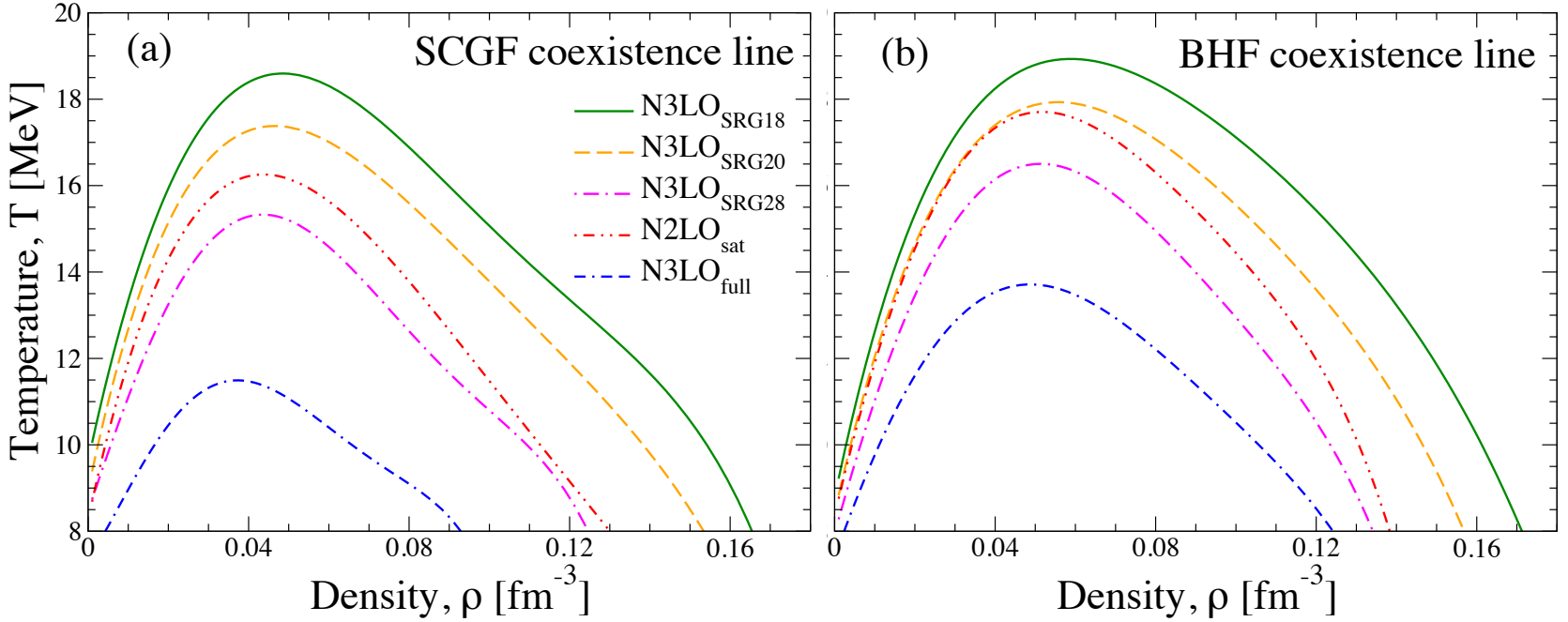
Zero temperature extrapolation





- From a finite set of points to functions?

Liquid-gas phase transition



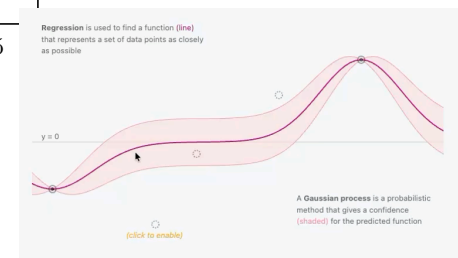
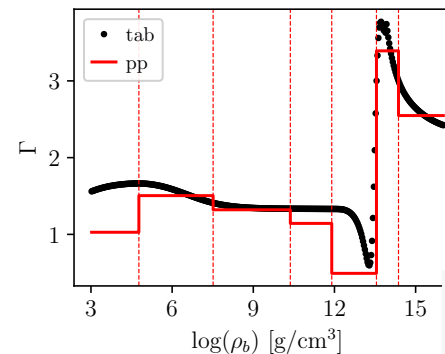
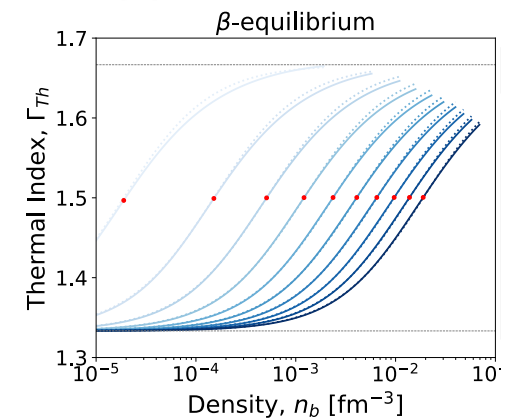
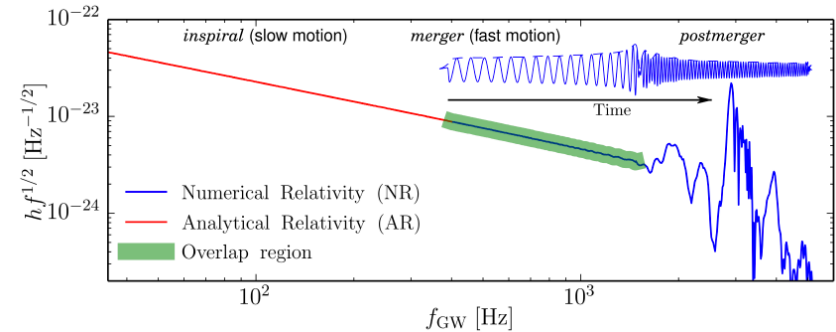
SCGF	ρ_c (fm ⁻³)	T_c (MeV)	ρ_0 (fm ⁻³)	$\frac{E_0}{A}$ (MeV)	$\frac{m_0^*}{m}$
N3LO _{SRG18}	0.048	18.6	0.19	-17.6	0.83
N3LO _{SRG20}	0.047	17.4	0.19	-16.2	0.84
N3LO _{SRG28}	0.043	15.3	0.16	-13.7	0.88
N2LO _{sat}	0.043	16.3	0.15	-14.6	0.90
N3LO _{full}	0.038	11.5	0.14	-11.0	0.84
BHF	ρ_c (fm ⁻³)	T_c (MeV)	ρ_0 (fm ⁻³)	$\frac{E_0}{A}$ (MeV)	$\frac{m_0^*}{m}$
N3LO _{SRG18}	0.058	18.9	0.19	-17.3	0.70
N3LO _{SRG20}	0.056	17.9	0.18	-16.0	0.73
N3LO _{SRG28}	0.051	16.5	0.16	-14.2	0.79
N2LO _{sat}	0.051	17.7	0.16	-15.2	0.82
N3LO _{full}	0.048	13.7	0.16	-12.6	0.76

- Motivation

- **Virial approximation**

- Thermal effects in BNS

- Gaussian processes for finite T

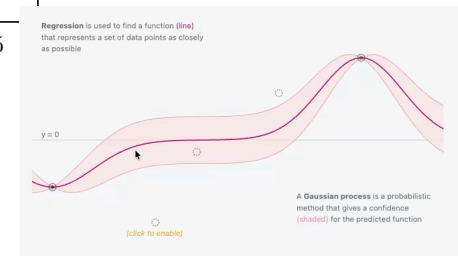
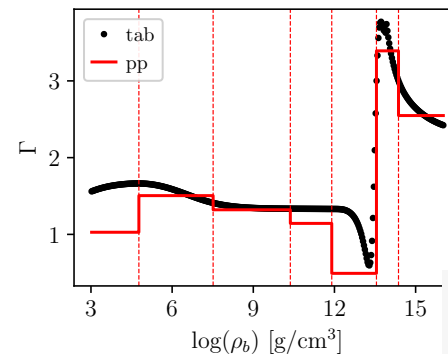
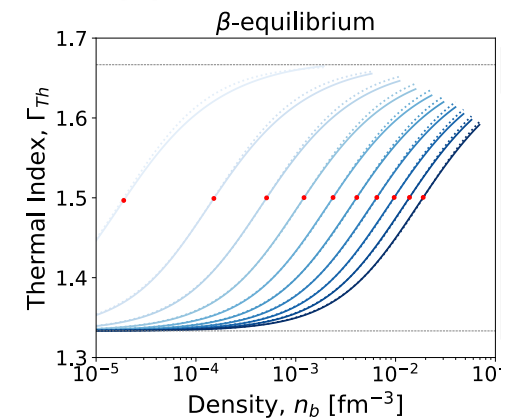
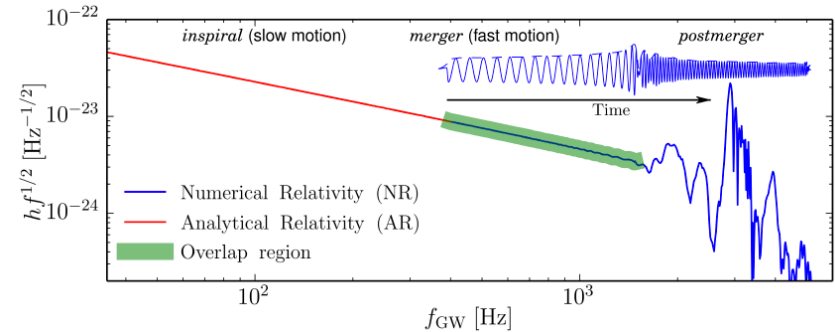


- Motivation

- **Virial approximation**

- Thermal effects in BNS

- Gaussian processes for finite T



Expansion

$$P = \frac{2T}{\lambda^3} [z + z^2 b_2 + z^3 b_3 + O(z^4)]$$

$$n = \frac{2}{\lambda^3} [z + 2z^2 b_2 + 3z^3 b_3]$$

$$\epsilon = \frac{3}{2}P + \frac{2T^2}{\lambda^3} [z^2 b'_2 + z^3 b'_3]$$

$$z = e^{\mu/T} \quad b'_n = \frac{\partial b_n}{\partial T}$$

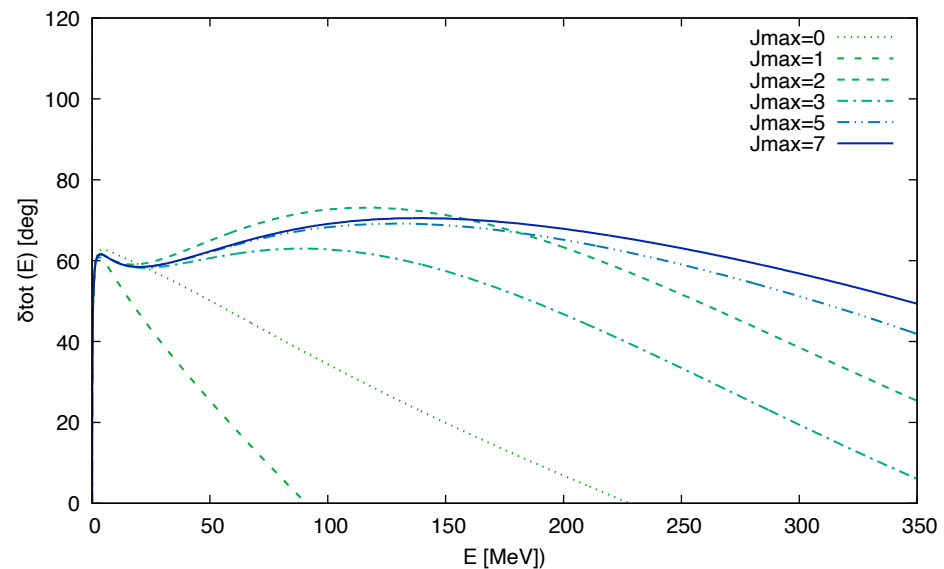
- Systematic expansion
- Improvable order-by-order

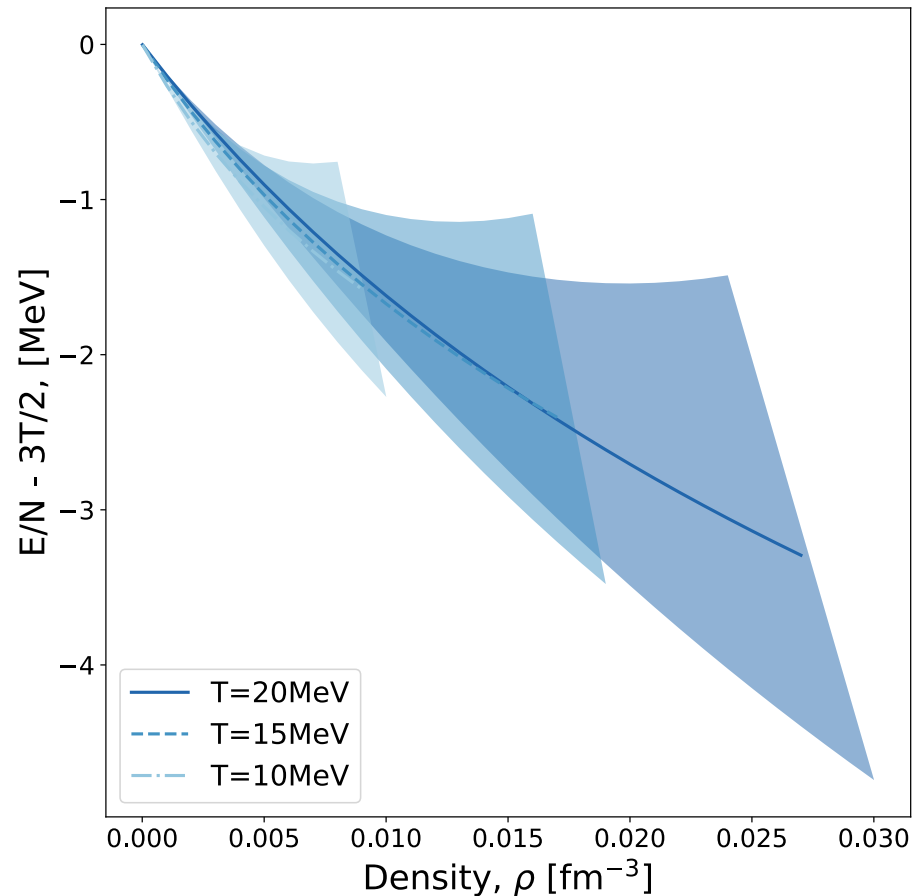
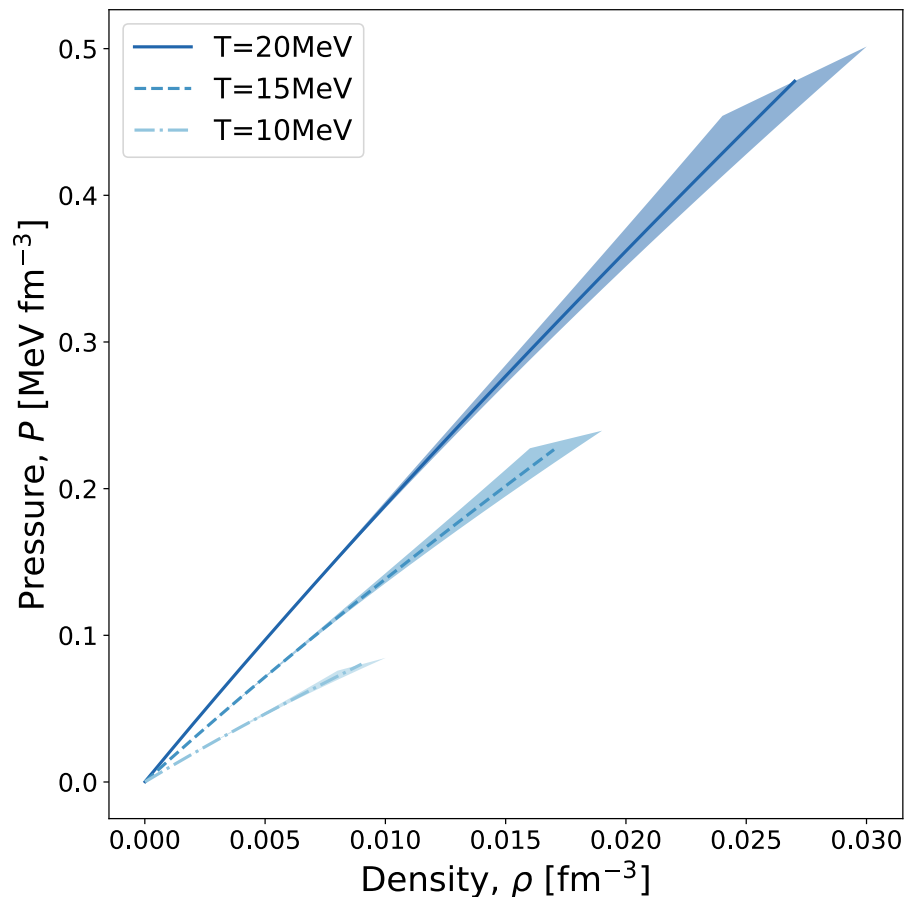
Virial coefficients

$$b_2(T) = \frac{1}{\sqrt{2\pi T}} \int_0^\infty dE e^{-E/T} \delta(E) - 2^{-5/2}$$

$$b_3(T) = ?$$

Drut et al. Unitary gas





Horowitz & Schwenk, *Phys. Lett. B* **638** 153159 (2006); *Nuc. Phys. A* **776** 055079 (2006)
 Riviuccio, Nadal-Matosas, Rios & Ruiz, *ApJ* **987**, 67 (2025)

- Use $\pm b_3 = 3^{-5/2} - 0.5b_2$ for uncertainty
- Truncate at $z=0.3$
- Granada phase-shifts for calculations

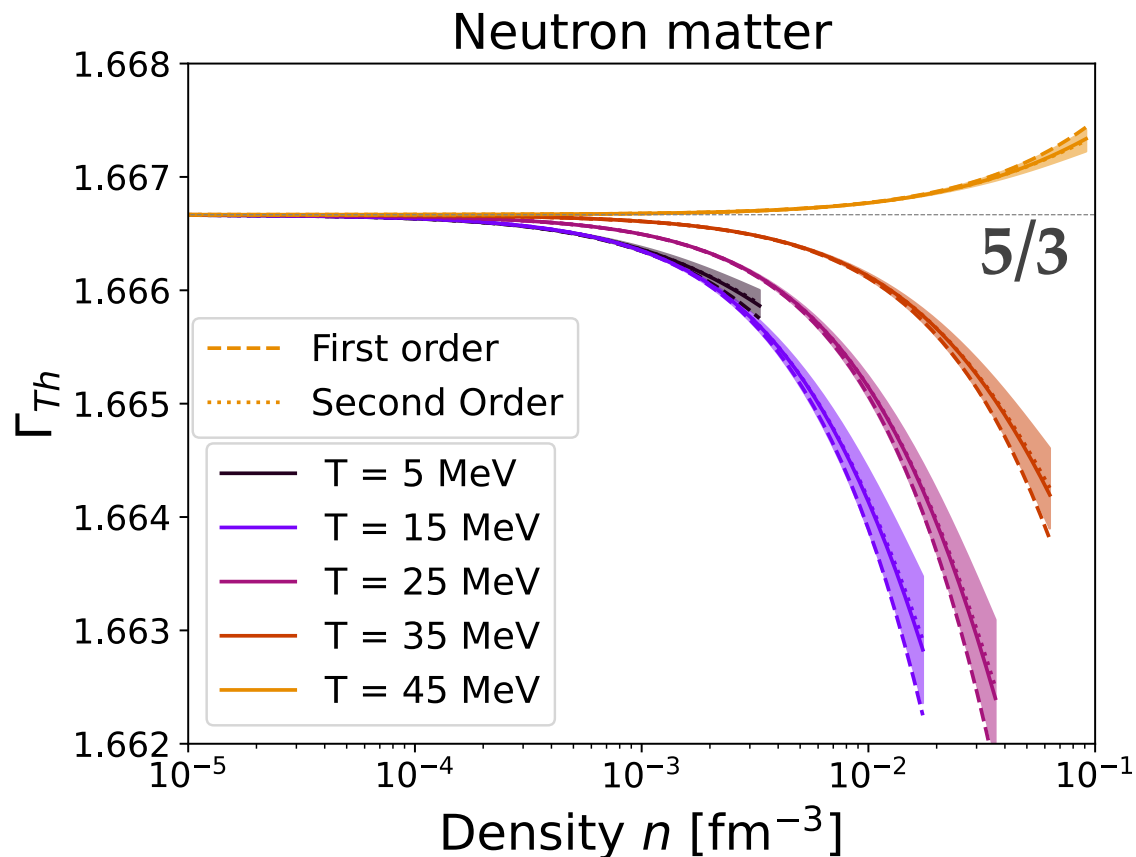
Thermal index: pure neutron matter

$$\Gamma_{th} = 1 + \frac{P_{th}}{\epsilon_{th}}$$

$$\Gamma_{th} = \frac{5}{3} + \Gamma_1 z + \Gamma_2 z^2$$

$$\Gamma_1 = -\frac{4}{9} T b'_2 \leq 0$$

$$\Gamma_2 = \frac{2}{3} T b'_2 + b_2 - \frac{b'_3}{b'_2}$$

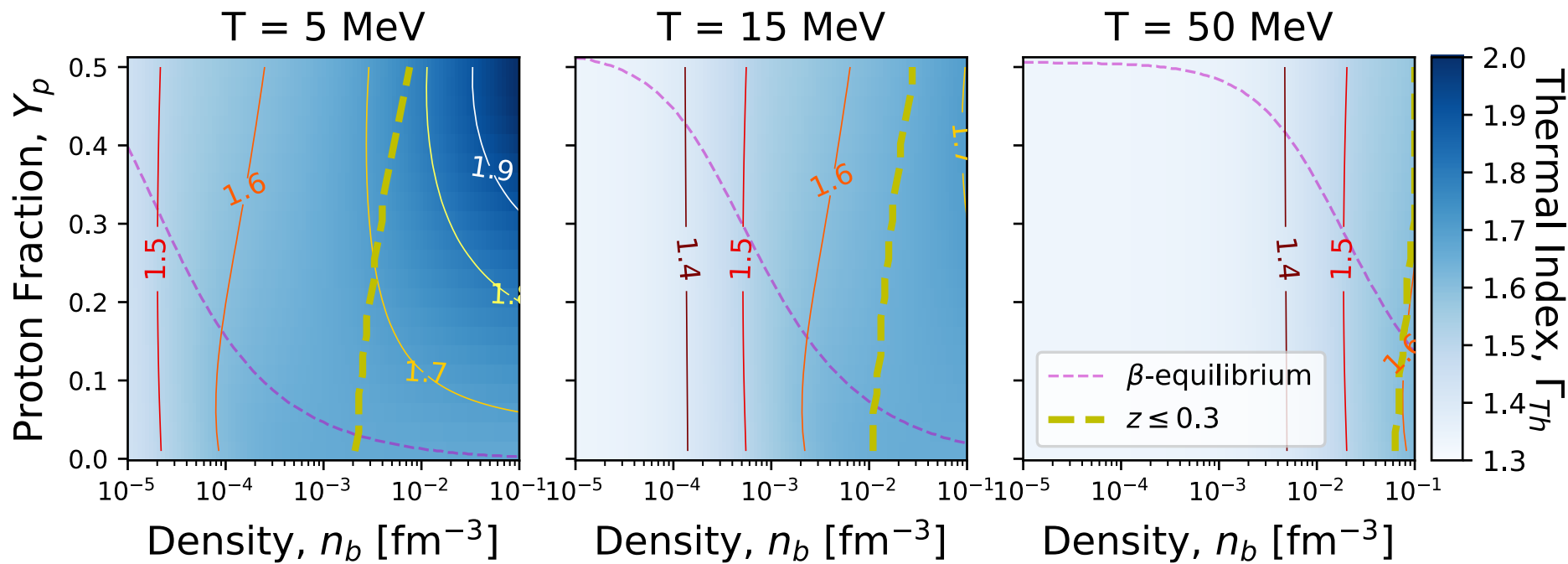


G Riviuccio @ UV



- **Small variation: 0.3%**
- **Small uncertainty: 0.1%**
- **T** dependence mild
- Closeness to **unitary gas**

Thermal index: npe matter



Baryon matter

$$P_{th}^{nuc} = \frac{2T}{\lambda^3} [z_n + z_p + (z_n^2 + z_p^2)b_n + 2z_n z_p b_{np}]$$

Leptons & photons ultrarelativistic

$$\Gamma_{th} = \frac{4}{3}$$

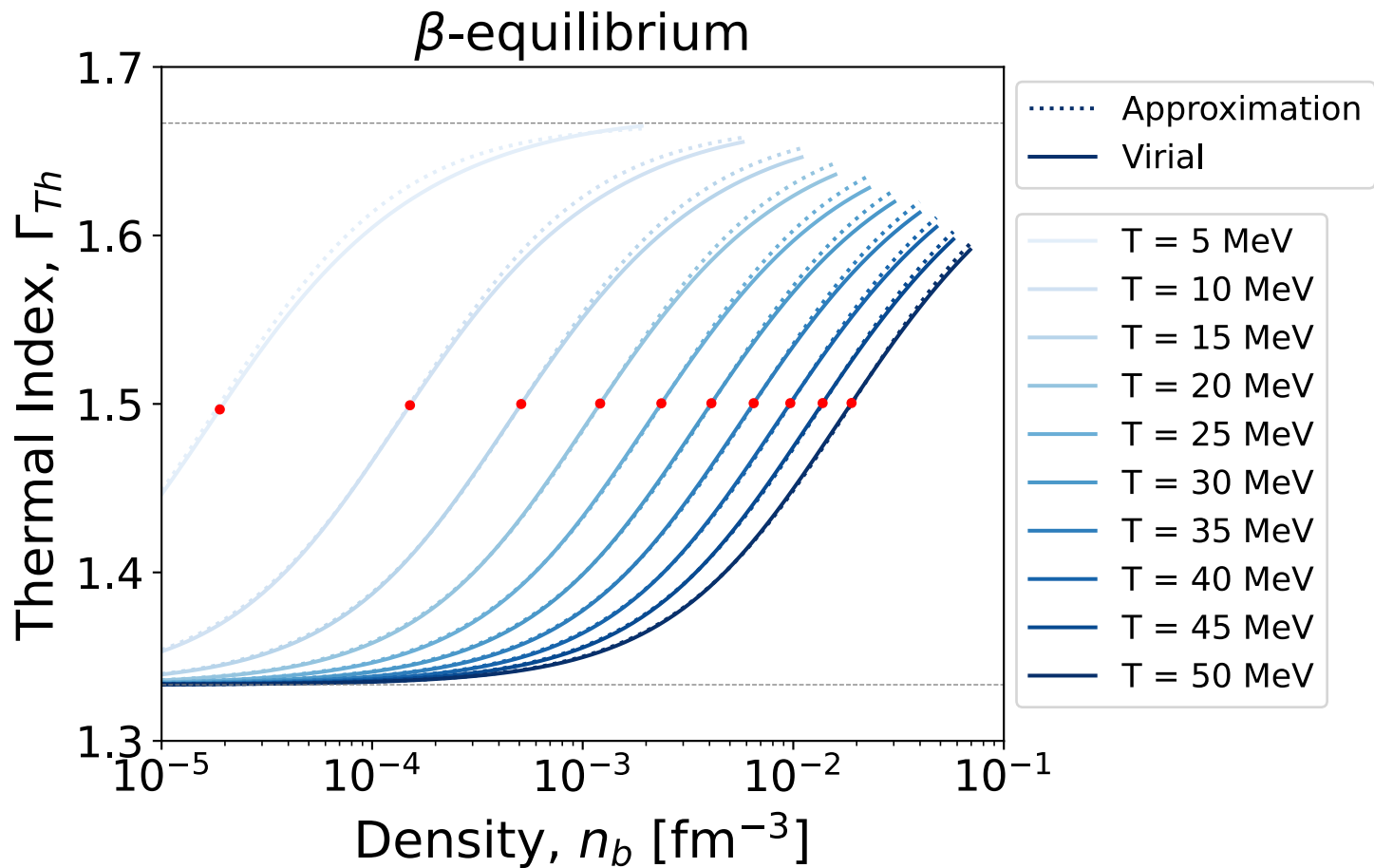
β equilibrium

$$\mu_n = \mu_p + \mu_e$$

Thermal index

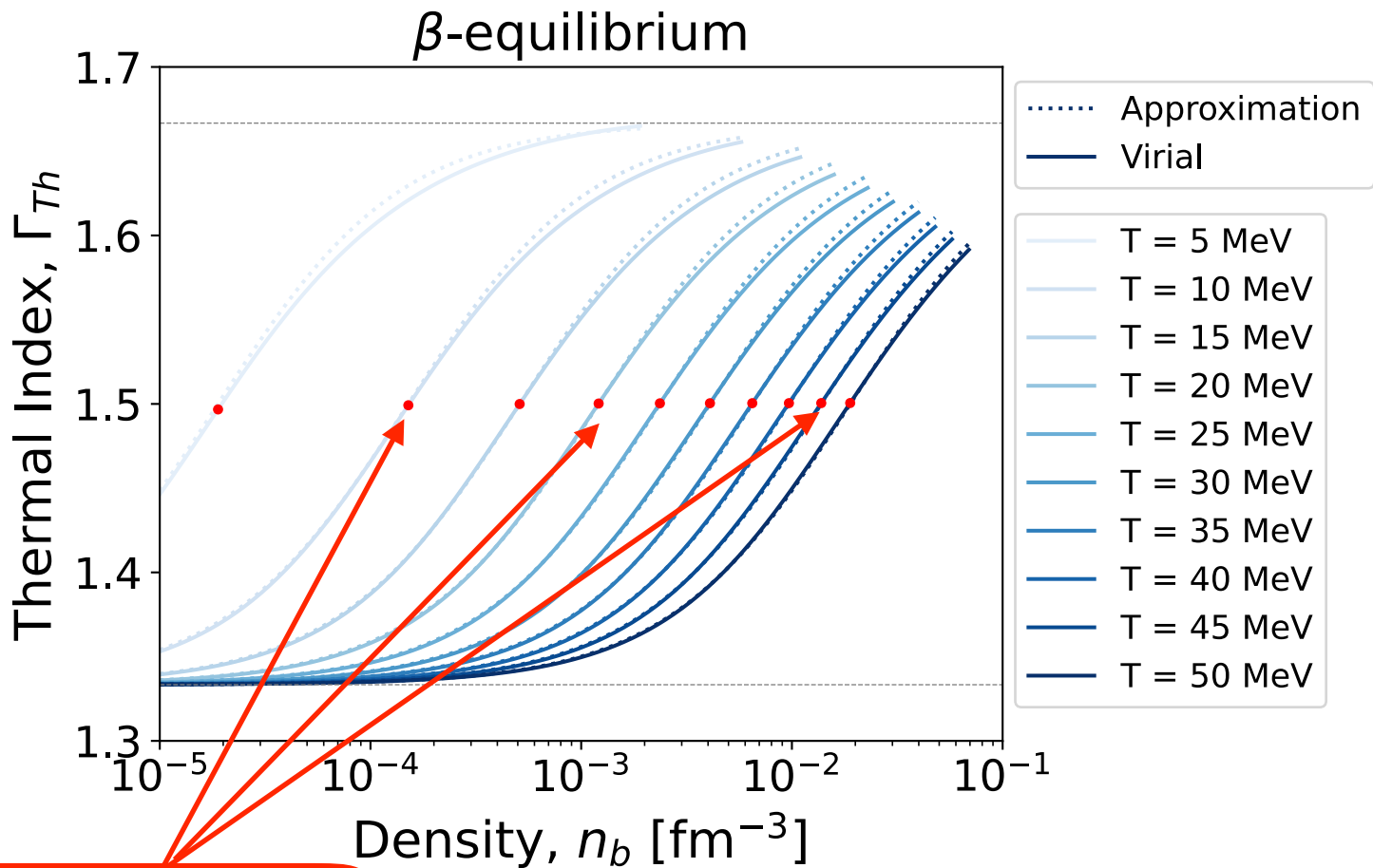
$$\Gamma_{th} = 1 + \frac{P_{th}^{nuc} + P_{th}^{lep} + P_{th}^{\gamma}}{\epsilon_{th}^{nuc} + \epsilon_{th}^{lep} + \epsilon_{th}^{\gamma}}$$

Thermal index: parametrization



$$\Gamma_{th} = \frac{4}{3} + \frac{1}{3} \frac{n_b}{n_b + n_{inf}}$$

Thermal index: parametrization



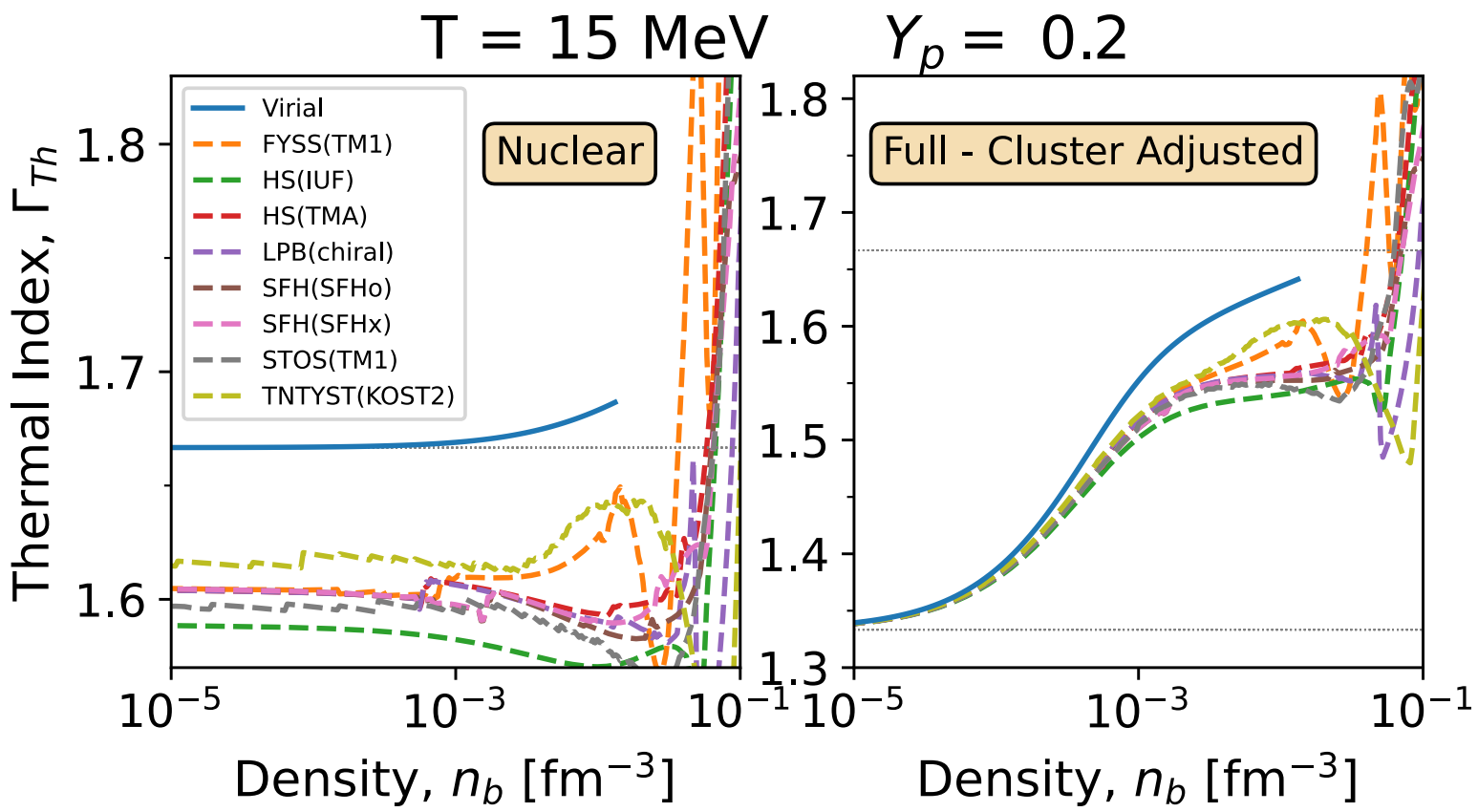
Regime change

$$P_{th}^{nuc} \approx P_{th}^{lep}$$

$$n_{inf} = 1.5 \times 10^{-4} \left(\frac{T}{10 \text{ MeV}} \right)^3 \text{ fm}^{-3}$$

$$\Gamma_{th} = \frac{4}{3} + \frac{1}{3} \frac{n_b}{n_b + n_{inf}}$$

Full tabulated comparison



Tabulated

ρ	T	ϵ	P	Γ_{th}

- **Clusters** are nuclear-only contributions
- Irrelevant at $T > 15 \text{ MeV}$

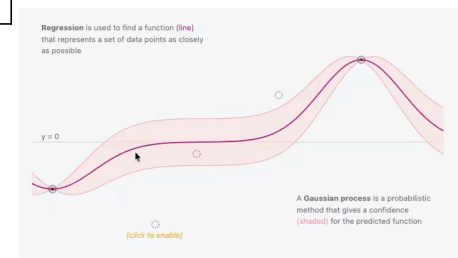
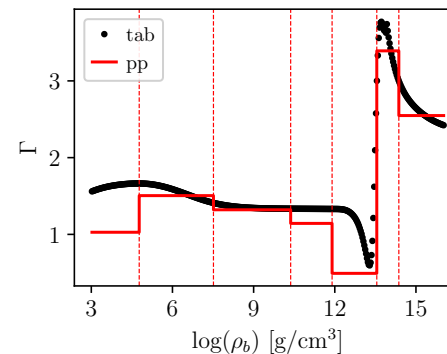
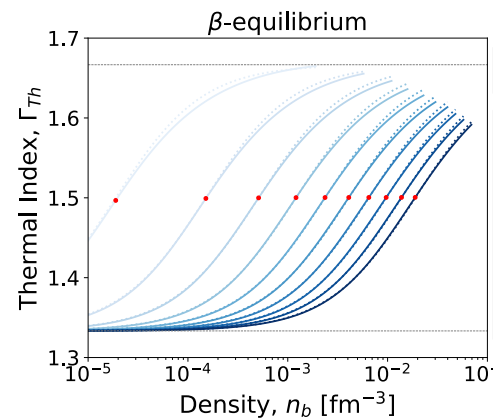
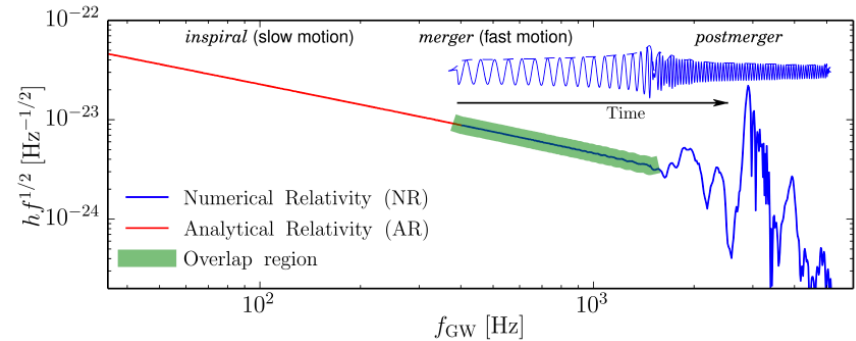
<https://github.com/Gpeppee/Virial-EOS.git>

- Motivation

- Virial approximation

- Thermal effects in BNS

- Gaussian processes for finite T

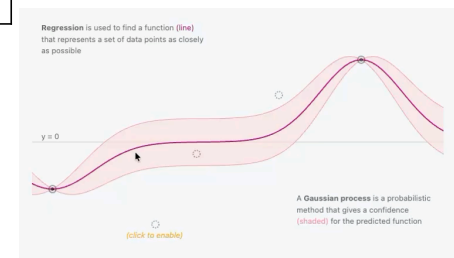
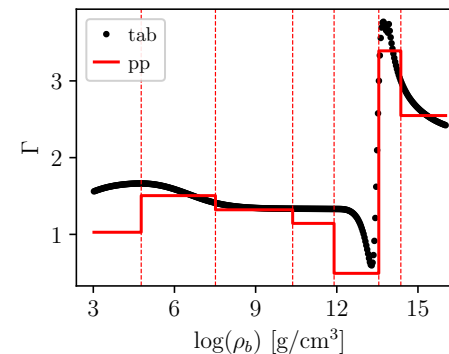
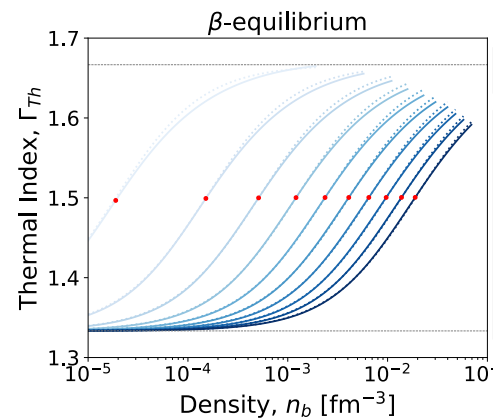
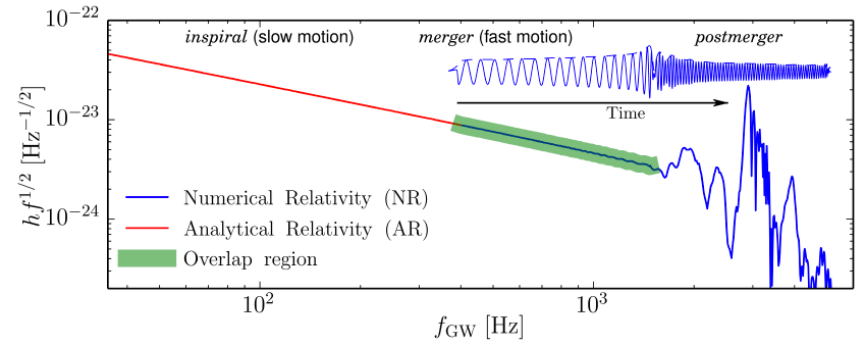


- Motivation

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On the treatment of thermal effects in the equation of state on neutron star merger remnants

Davide Guerra^{1,*}, Milton Ruiz¹, Michele Pasquali², Pablo Cerdá-Durán^{1,3}, Arnau Rios^{4,5} and José A. Font^{1,3}

¹Departamento de Astronomía y Astrofísica, Universitat de València, Dr. Moliner 50, 46100, Burjassot (València), Spain

²Department of Chemistry, Life Sciences and Environmental Sustainability, University of Parma, 43124 Parma, Italy

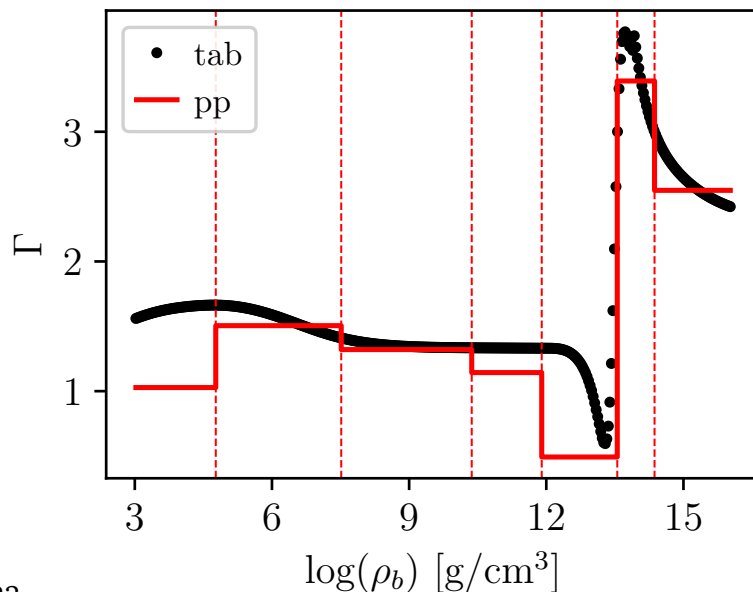
³Observatori Astronòmic, Universitat de València, C/ Catedrático José Beltrán 2, 46980, Paterna (València), Spain

$$P(\epsilon, T) = P(\epsilon, T = 0) + P_{th}(\epsilon_{th})$$



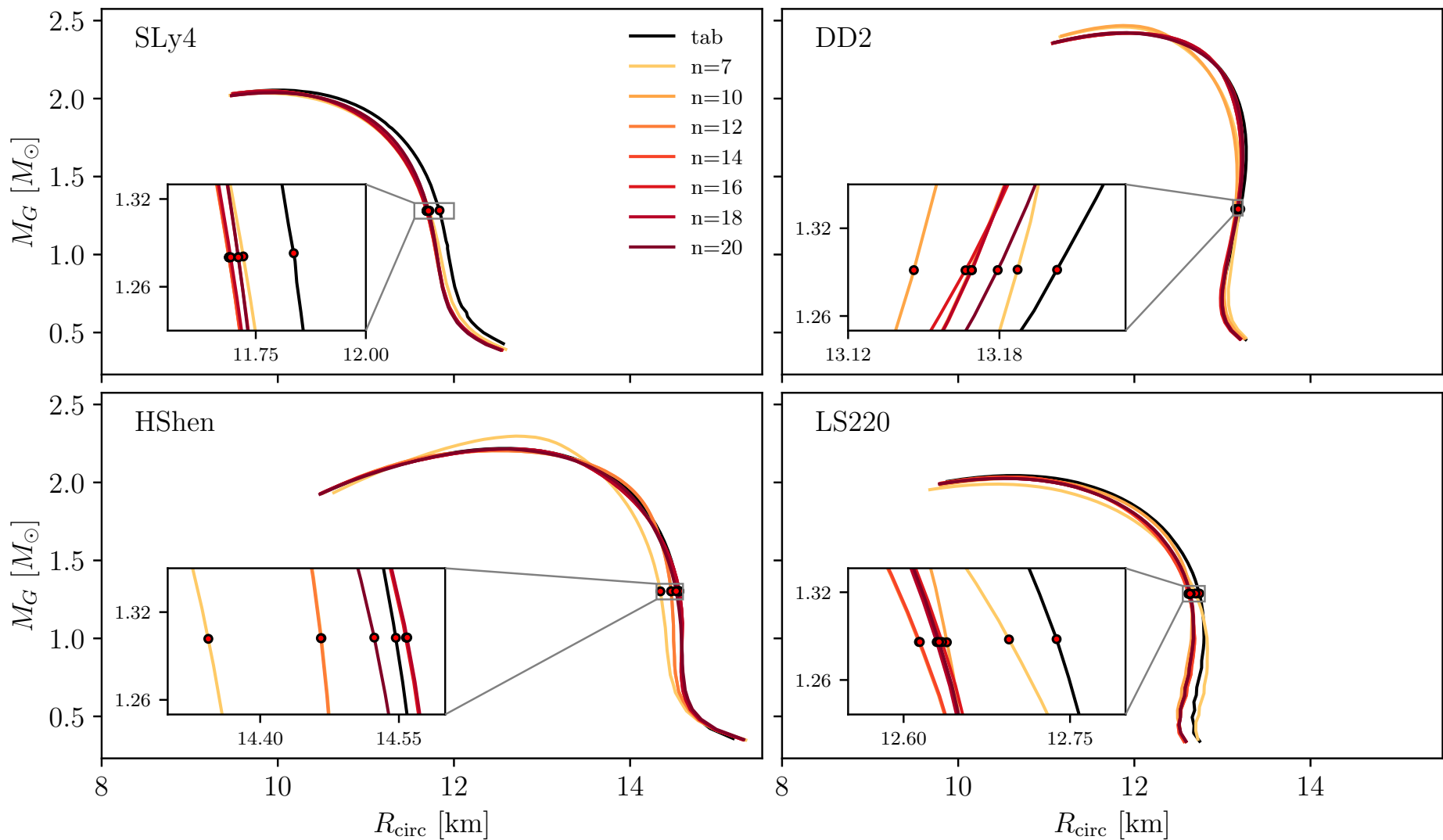
Polytropic / Piecewise

$$Q_{th}(\rho, T) = K_{th} \rho^{\Gamma_{th}}$$

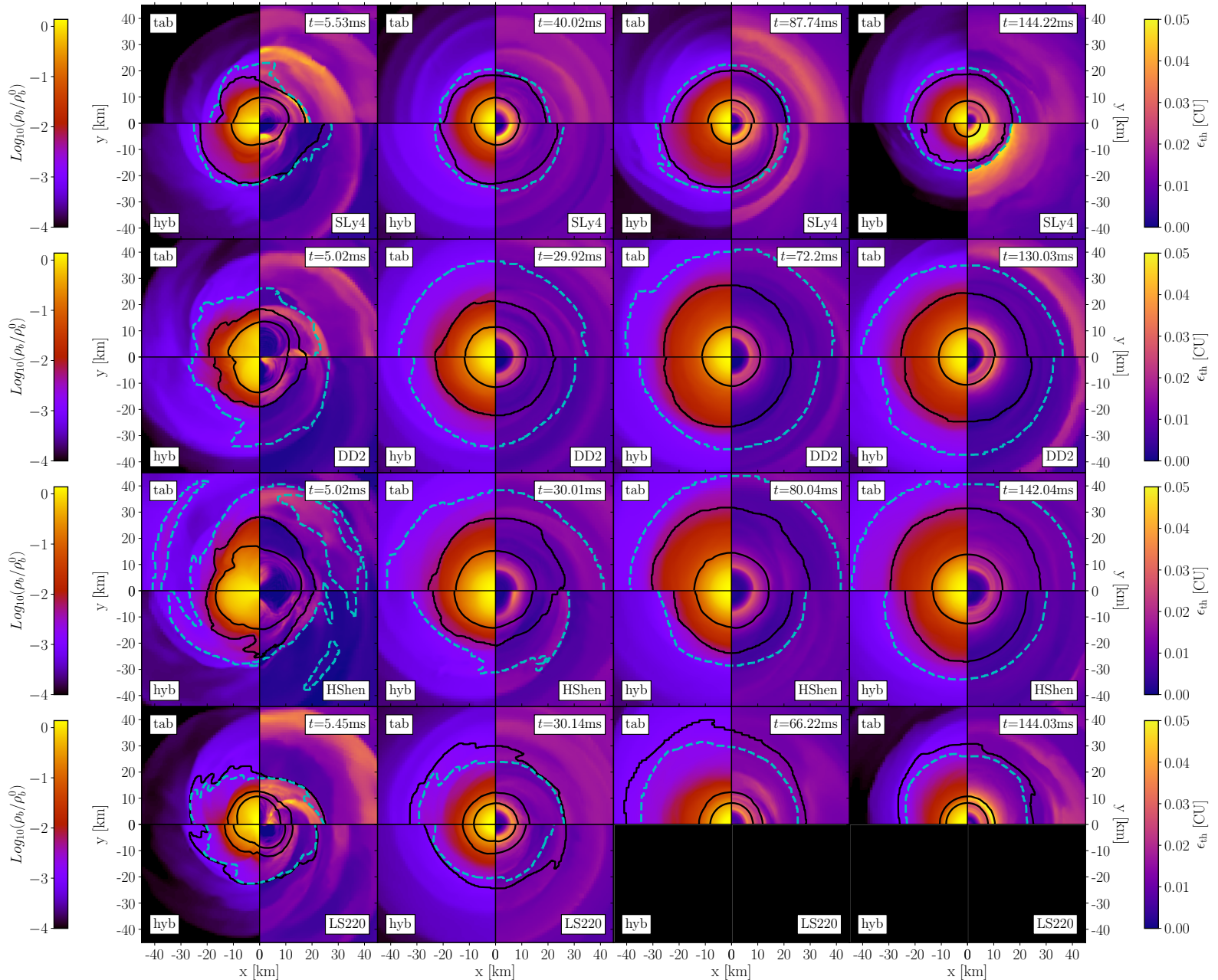


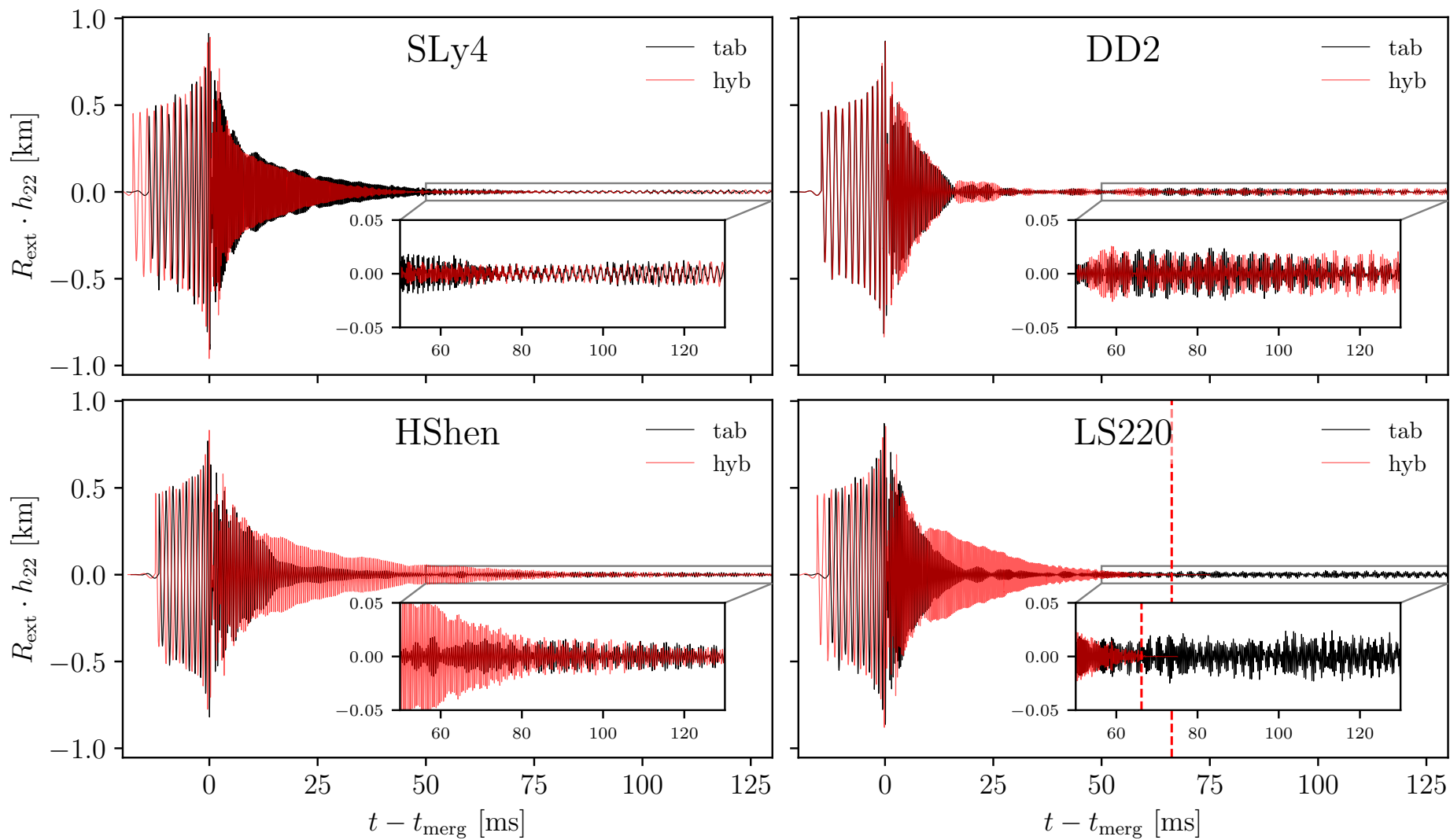
Tabulated

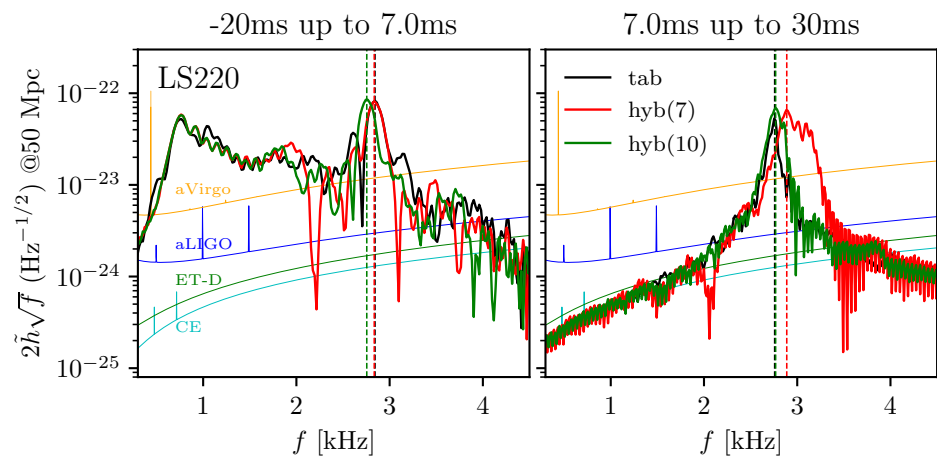
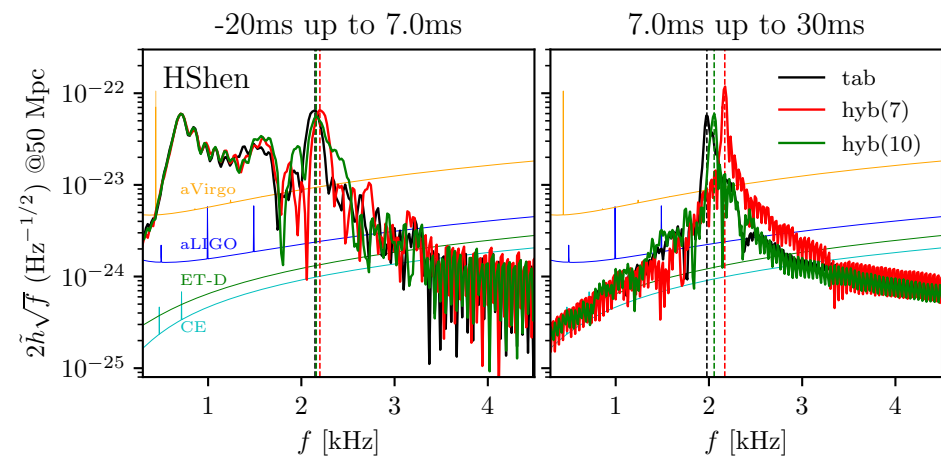
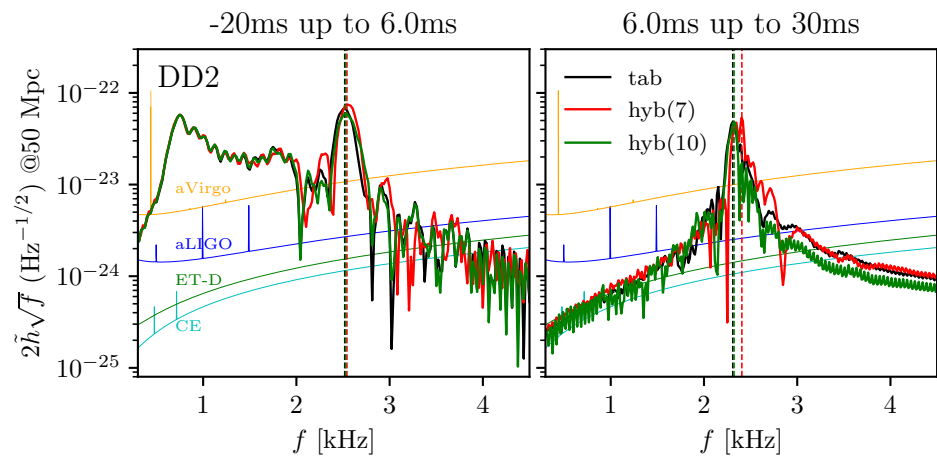
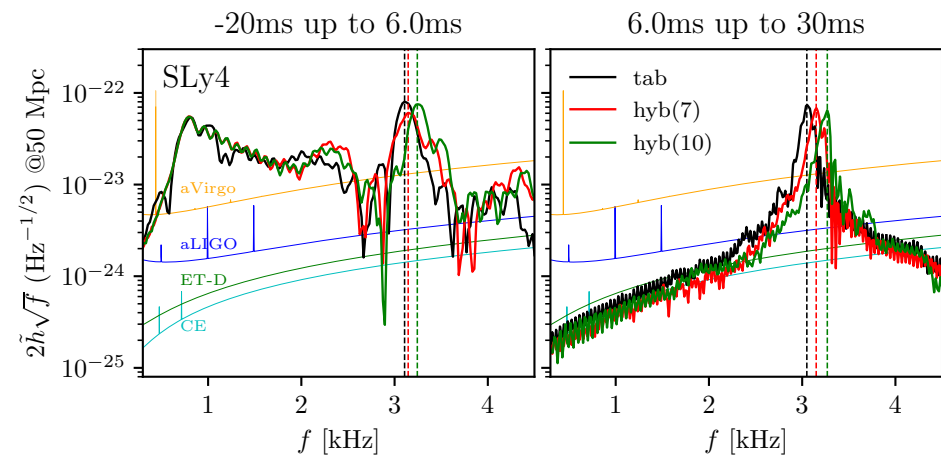
ρ	T	ϵ	P	Γ_{th}



Density and temperature

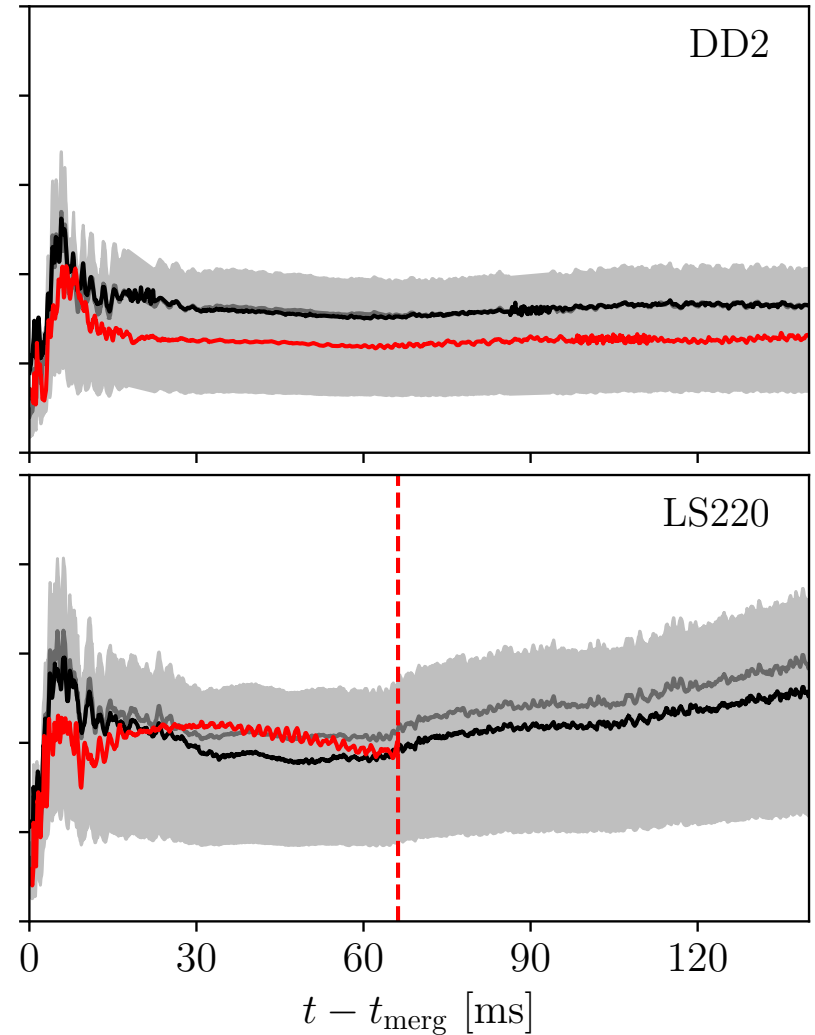
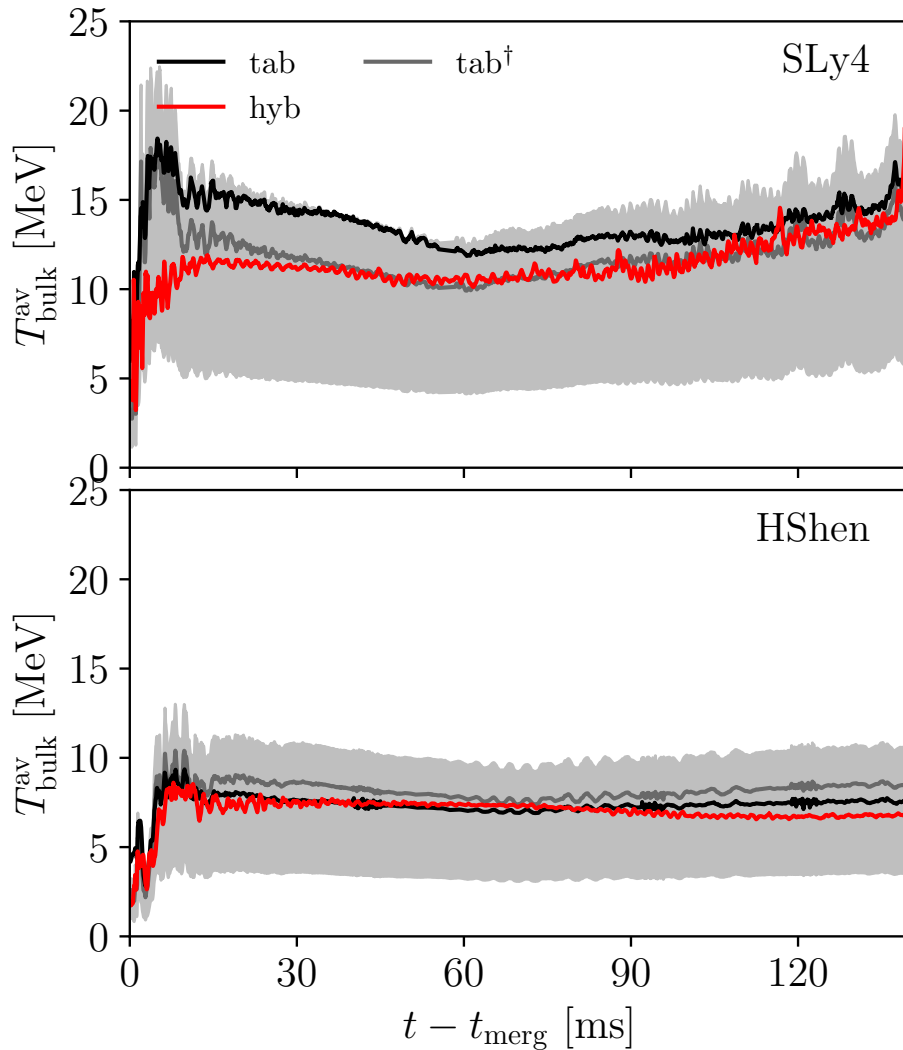




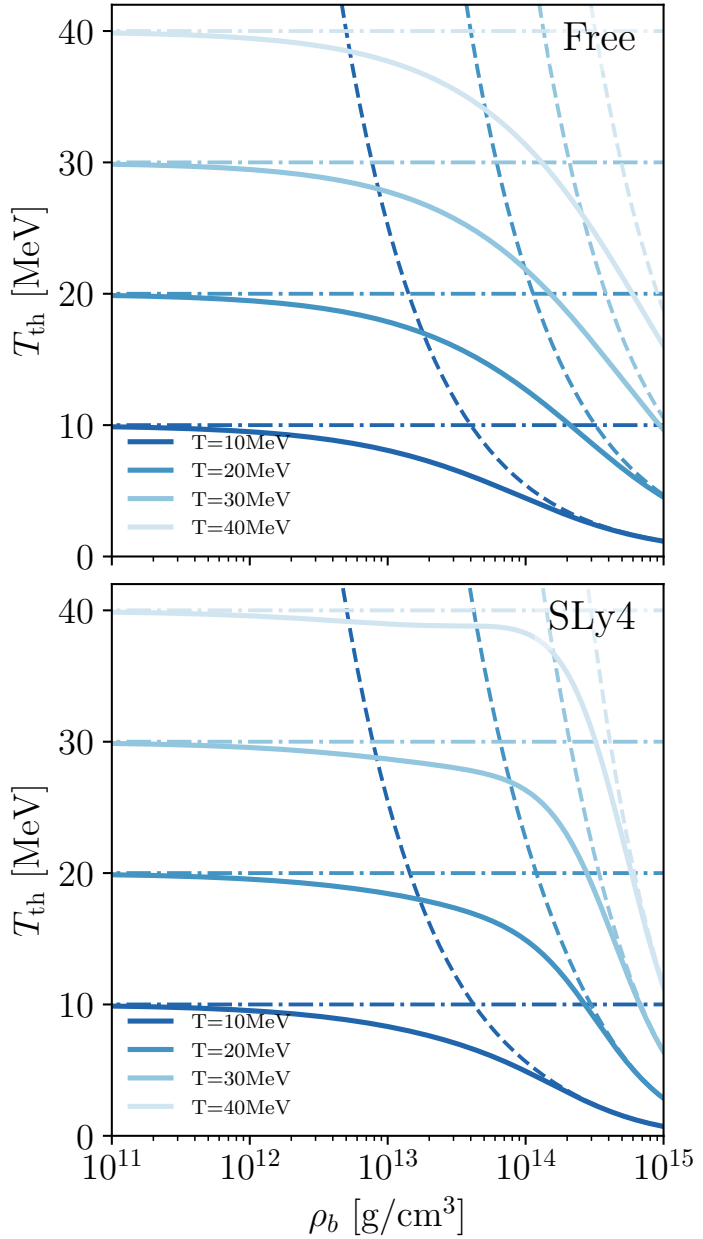


Temperature estimate

$$T_{hyb} = (\Gamma_{th} - 1)\epsilon_{hyb}$$



Estimate error in T



$$e(\rho, T) = e(\rho, T = 0) + e_{th}(\rho, T)$$

$$p(\rho, T) = p(\rho, T = 0) + (\Gamma_{th} - 1) \rho e_{th}(\rho, T)$$

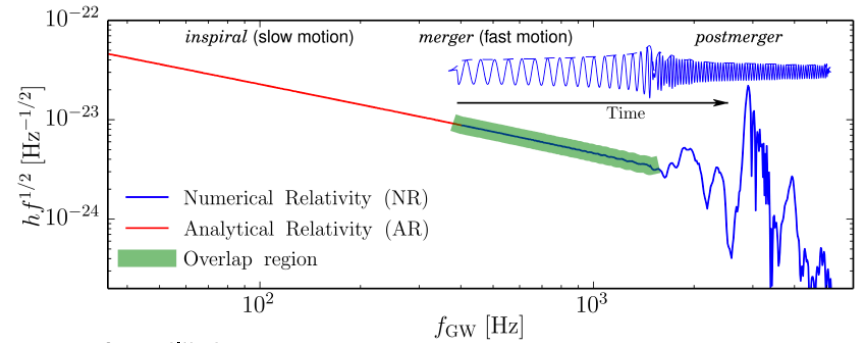
$$T_{th} = (\Gamma_{th} - 1) e_{th}(n, T)$$

Ideal Gas $T_{th} = \left(\frac{5}{3} - 1\right) \frac{3}{2} T = T$

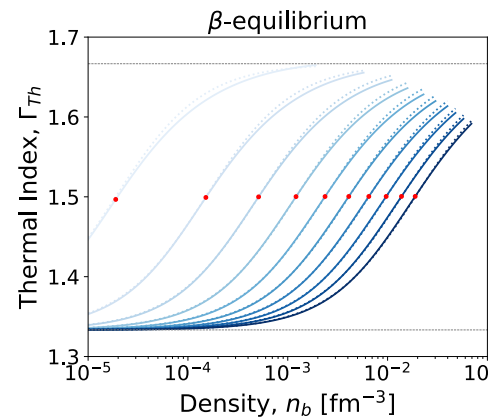
Degenerate Gas

$$T_{th} = \frac{\pi^2 T^2}{6 \varepsilon_F^*} \left[1 - \frac{3}{2} \frac{n}{m_n^*} \frac{\partial m_n^*}{\partial n} \right]$$

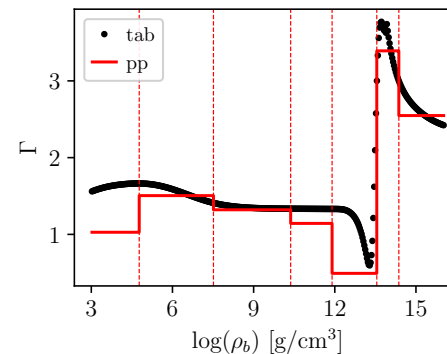
- Motivation



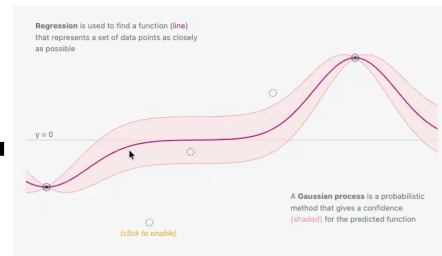
- Virial approximation



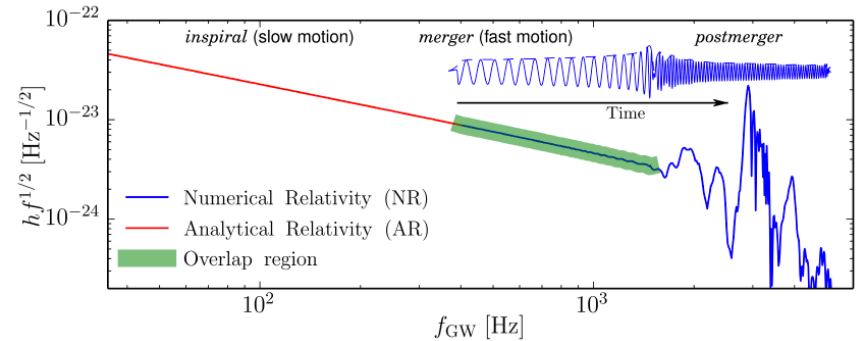
- Thermal effects in BNS



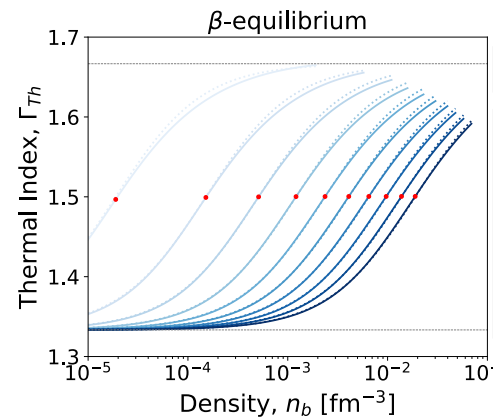
- Gaussian processes for finite T



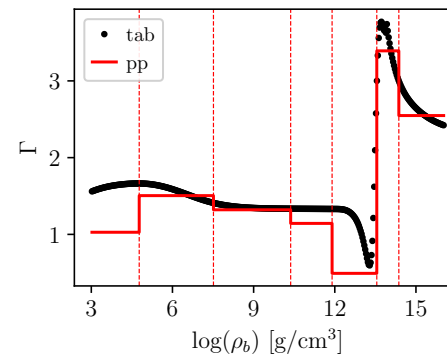
- Motivation



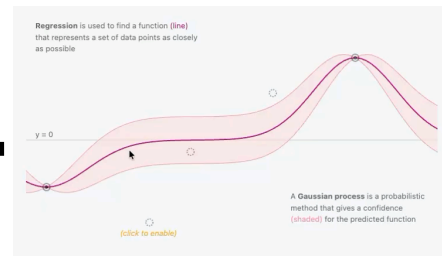
- Virial approximation



- Thermal effects in BNS



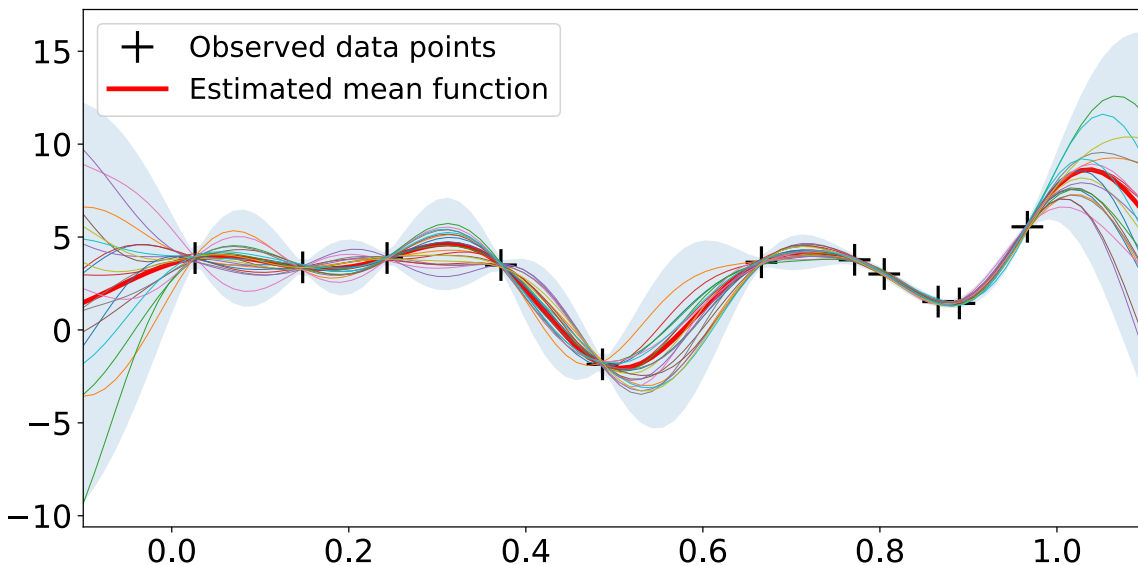
- Gaussian processes for finite T



If we denote a real-valued function $f(x)$ over an input space X , then saying that f is distributed as a *Gaussian Process* means that for any finite set of inputs $\{x_1, x_2, \dots, x_n\} \subset X$, the vector of function values

$$\mathbf{f} = (f(x_1), f(x_2), \dots, f(x_n))^{\top}$$

follows a multivariate Gaussian distribution.



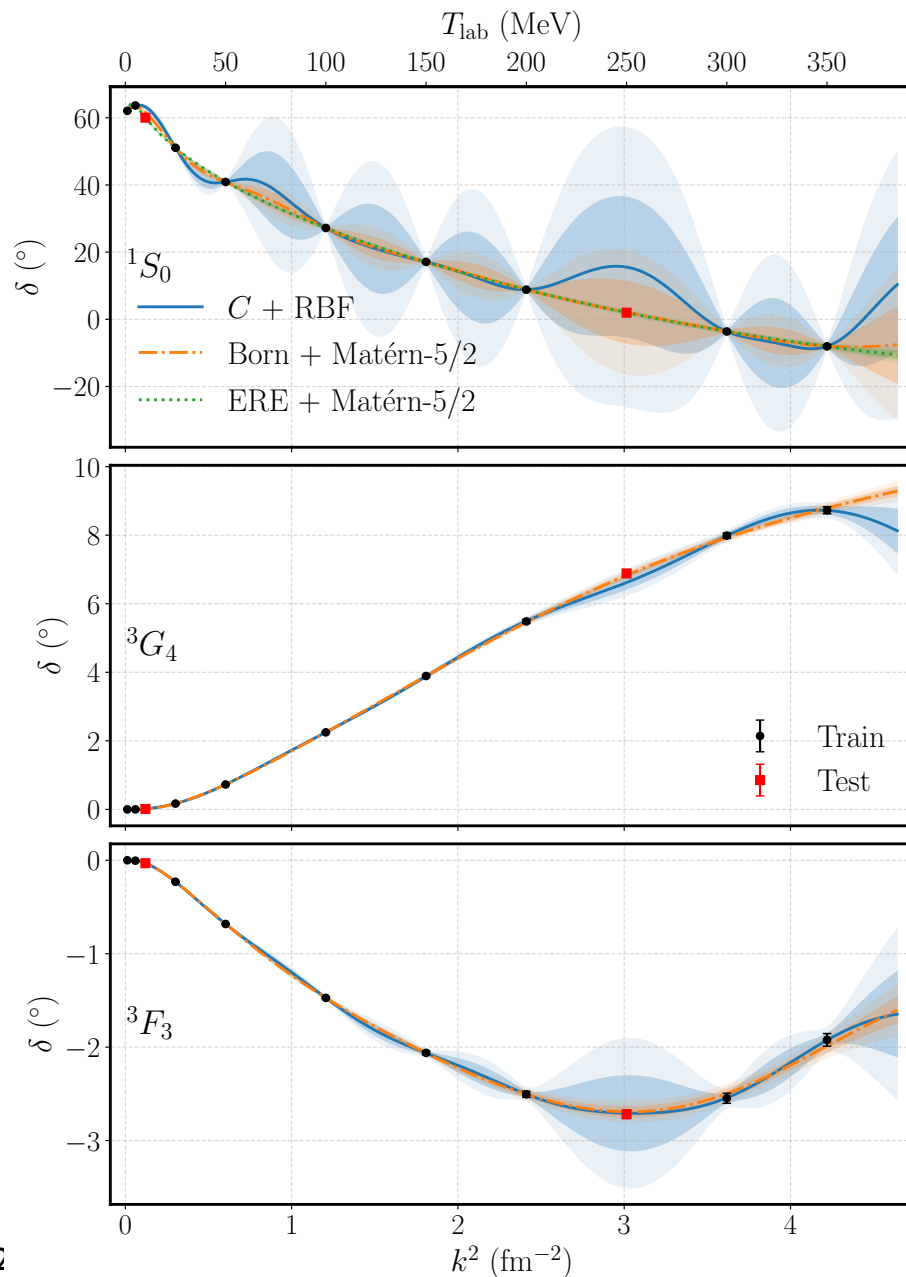
$$f(\mathbf{x}) \sim \mathcal{GP}(m(\mathbf{x}), k(\mathbf{x}, \mathbf{x}'))$$

mean

$$m(\mathbf{x}) = \mathbb{E}[f(\mathbf{x})]$$

kernel

$$k(\mathbf{x}, \mathbf{x}') = \mathbb{E}[(f(\mathbf{x}) - m(\mathbf{x})) (f(\mathbf{x}') - m(\mathbf{x}'))]$$



GP pros

- Guaranteed to reproduce data
- Bayesian interpolator
- Errors accounted for
- Differentiable model

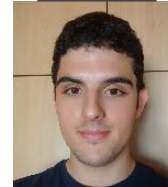
GP cons

- Extrapolation “misguided”

$$f(\mathbf{x}) \sim \mathcal{GP}(m(\mathbf{x}), k(\mathbf{x}, \mathbf{x}'))$$

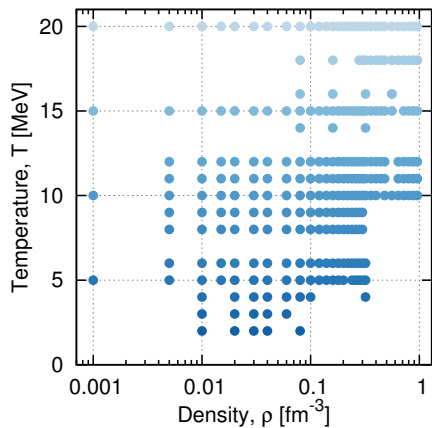
$$m(\mathbf{x}) = \mathbb{E}[f(\mathbf{x})] \approx C$$

Rozalén



Kochankovski





Mean model for EoS

$$\begin{aligned}
 m_{\text{par}}(\rho_B, T, Y_q) = & (e_0(\rho_B) + f_{0,\text{th}}(\rho_B, T)) \\
 & + (e_{\text{sym}}(\rho_B) + f_{\text{sym,th}}(\rho_B, T)) \delta^2 \\
 & + m_{\text{el}}(\rho_B, T, Y_q)
 \end{aligned}$$

Cold part

$$\chi = \frac{\rho_B - n_0}{3n_0}$$

$$e_0(\rho_B) = E_{\text{sat}} + \frac{1}{2}K_0 \chi^2 + \frac{1}{6}Q_0 \chi^3$$

$$e_{\text{sym}}(\rho_B) = J + L \chi + \frac{1}{2}K_{\text{sym}} \chi^2$$

Trainable parameters

$$E_{\text{sat}}, K_0, Q_0, J, L, K_{\text{sym}}, \alpha_0, \alpha_n, \beta_0, \beta_n$$

Thermal part

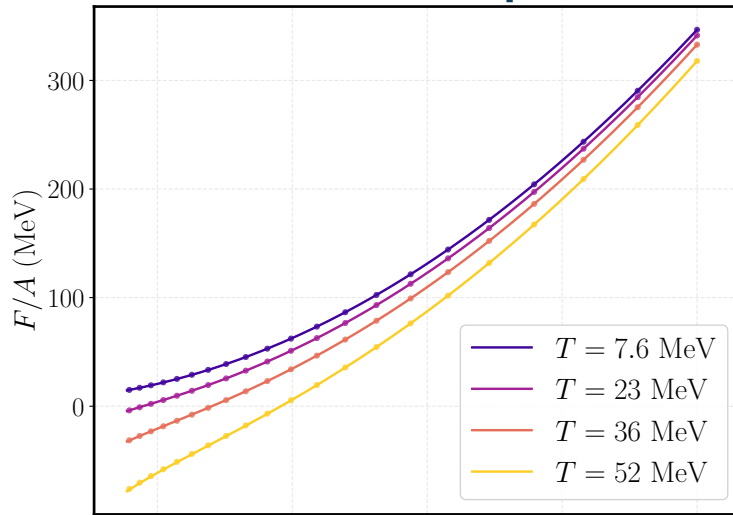
$$f_{0,\text{th}}(\rho_B, T) = -\frac{a_0(\rho_B) T^2}{1 + b_0(\rho_B) T}$$

$$f_{\text{sym,th}}(\rho_B, T) = -\frac{a_{\text{sym}}(\rho_B) T^2}{1 + b_S(\rho_B) T}$$

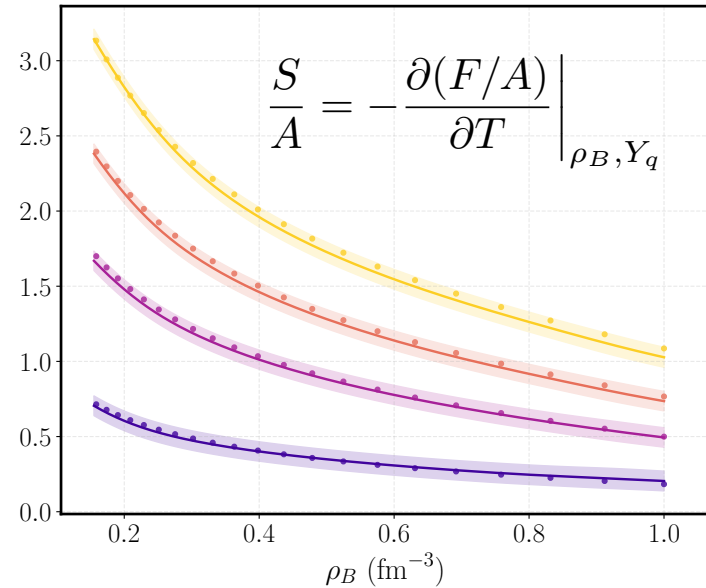
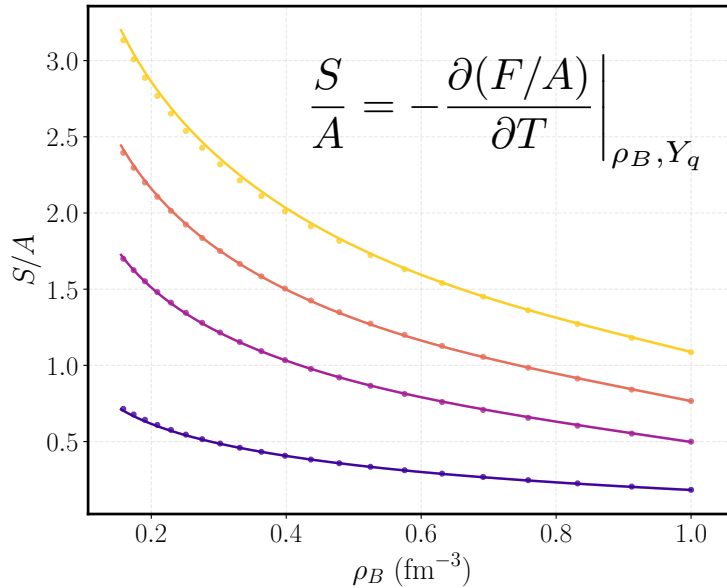
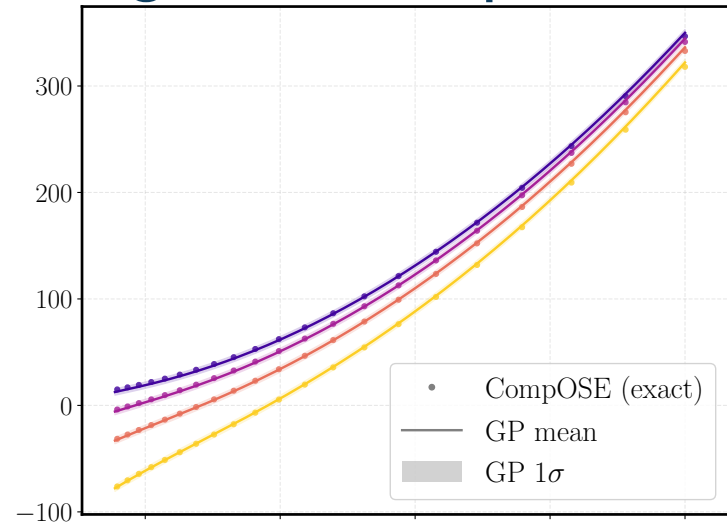
$$a_\tau(\rho) = \frac{\pi^2}{2} \frac{m_\tau^*(\rho)/m}{\varepsilon_{F,\tau}(\rho)}$$

$$\frac{m_\tau^*(\rho)}{m} = \frac{1}{1 + \alpha_\tau \rho / \rho_0}$$

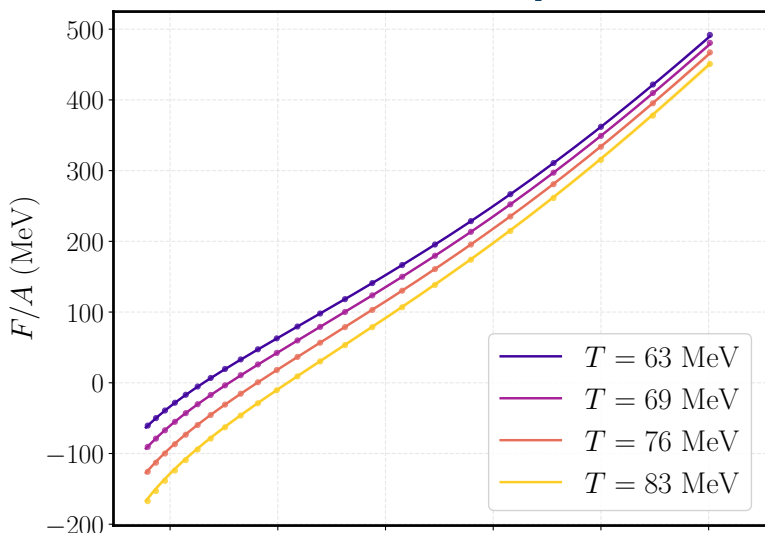
Informed interpolation



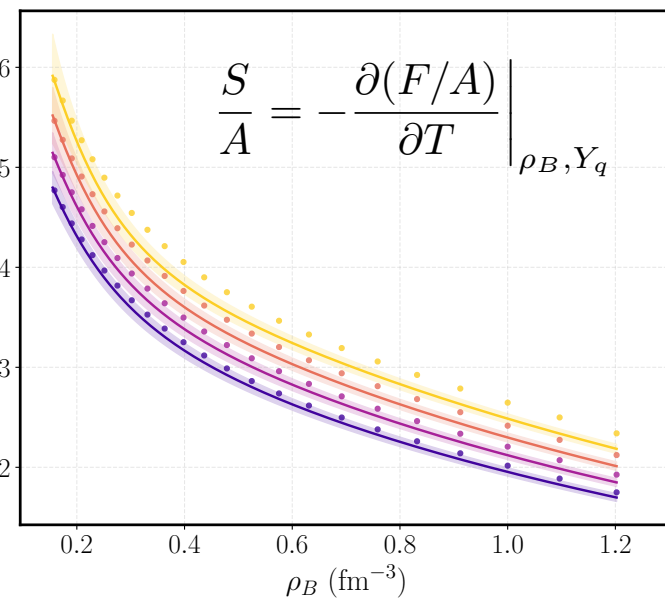
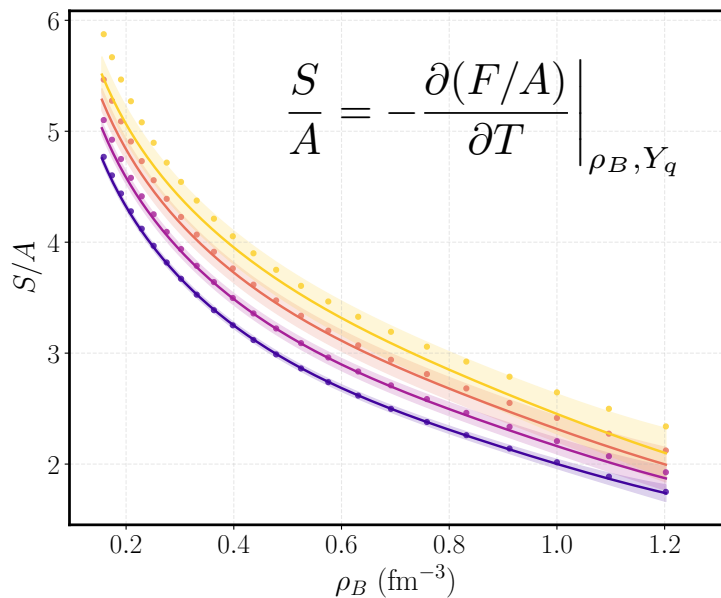
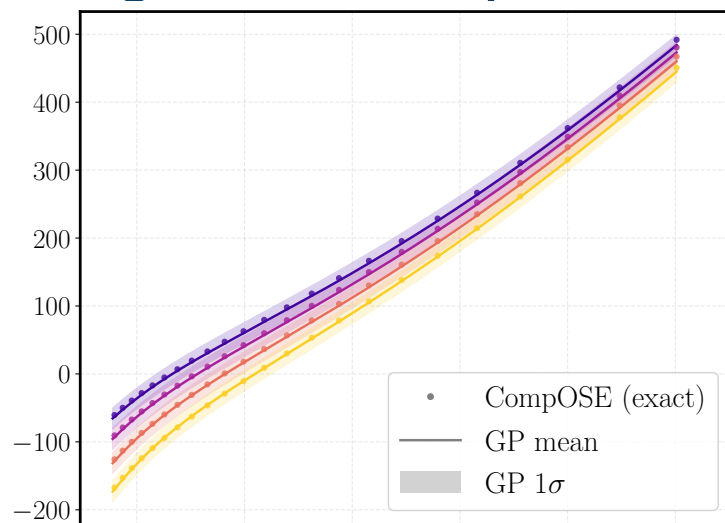
Agnostic interpolation



Informed extrapolation



Agnostic extrapolation



- Ab initio **finite temperature** effects can be simulated & quantified
- **ML** GP techniques can help **interpolate** & **extrapolate** (if physics guided)
- **Challenges**

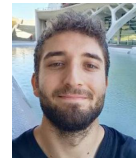
- **Uncertainty quantification**

- **Practical simulability**

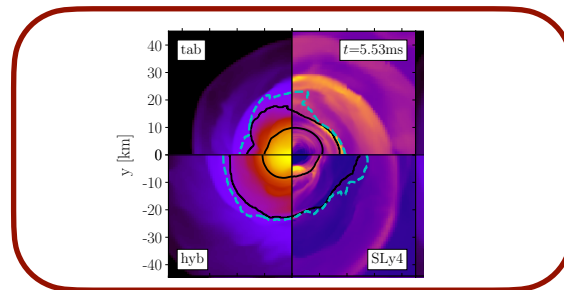
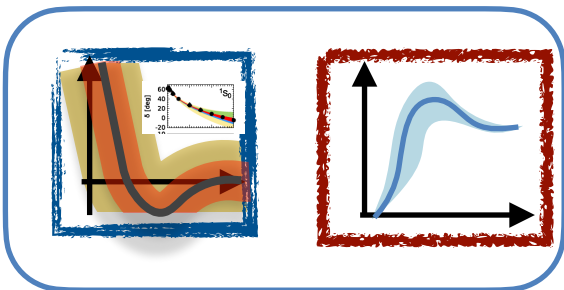
- **Internal NS structure: modes? Phases?**

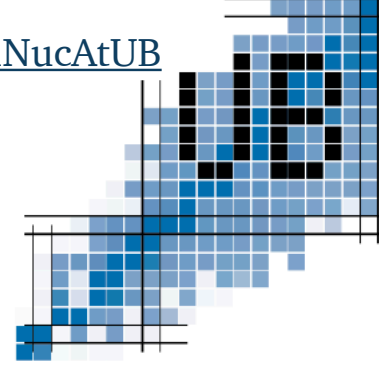


Márquez



G Riveccio





Thank you!

arnau.rios@icc.ub.edu

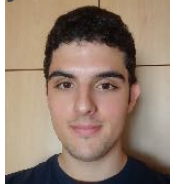
<https://sites.google.com/view/arnaorios/>



À Ramos



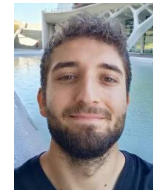
J Rozalén H Kochankovski



D Guerra



G Riviuccio



T Font



P Cerdá-Duran



M Ruiz



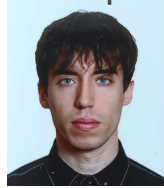
A Nadal



R Bondarescu



R Márquez



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