

# Overview: Numerical Methods for Neutrino Quantum Kinetics

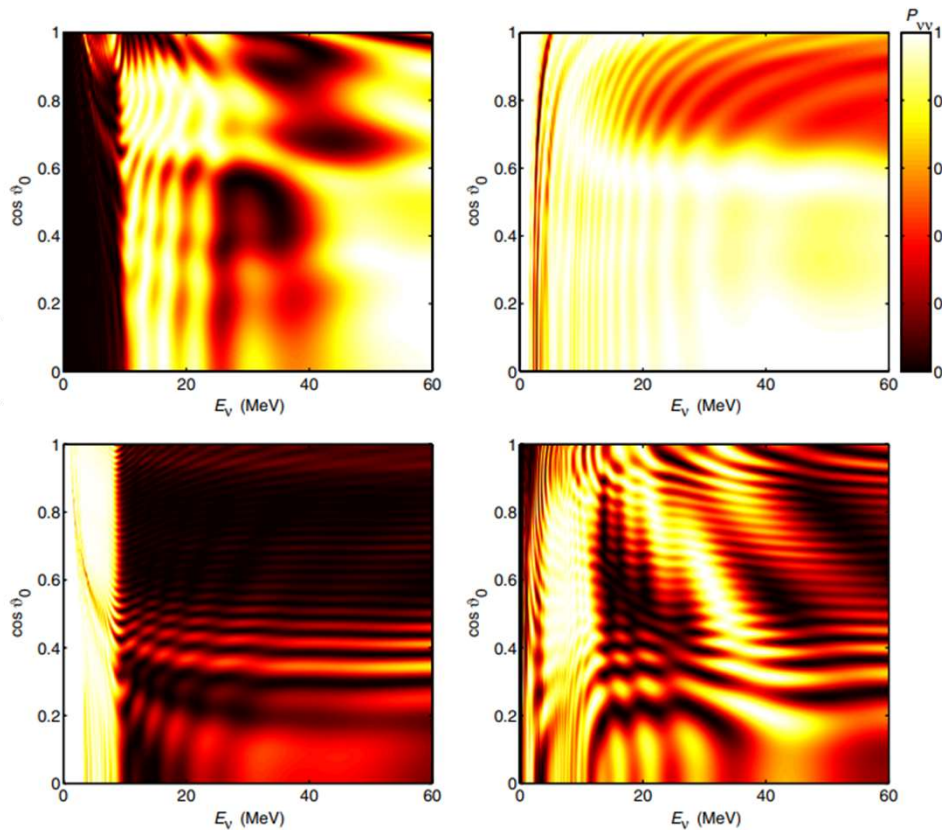
Sherwood Richers, University of Tennessee Knoxville



# The Problem

- Neutrino transport is the dominant cost of state-of-the-art simulations of core-collapse supernovae and neutron star mergers
- Neutrino flavor transformation modifies amount of heating, amount of mass ejection, and composition of ejecta
- Neutrino flavor transformation occurs on smaller length/time scales than transport

# The results are sensitive to resolution



- High-resolution 3D NSM simulations: **12.5 meters**

Kiuchi et al (2023)

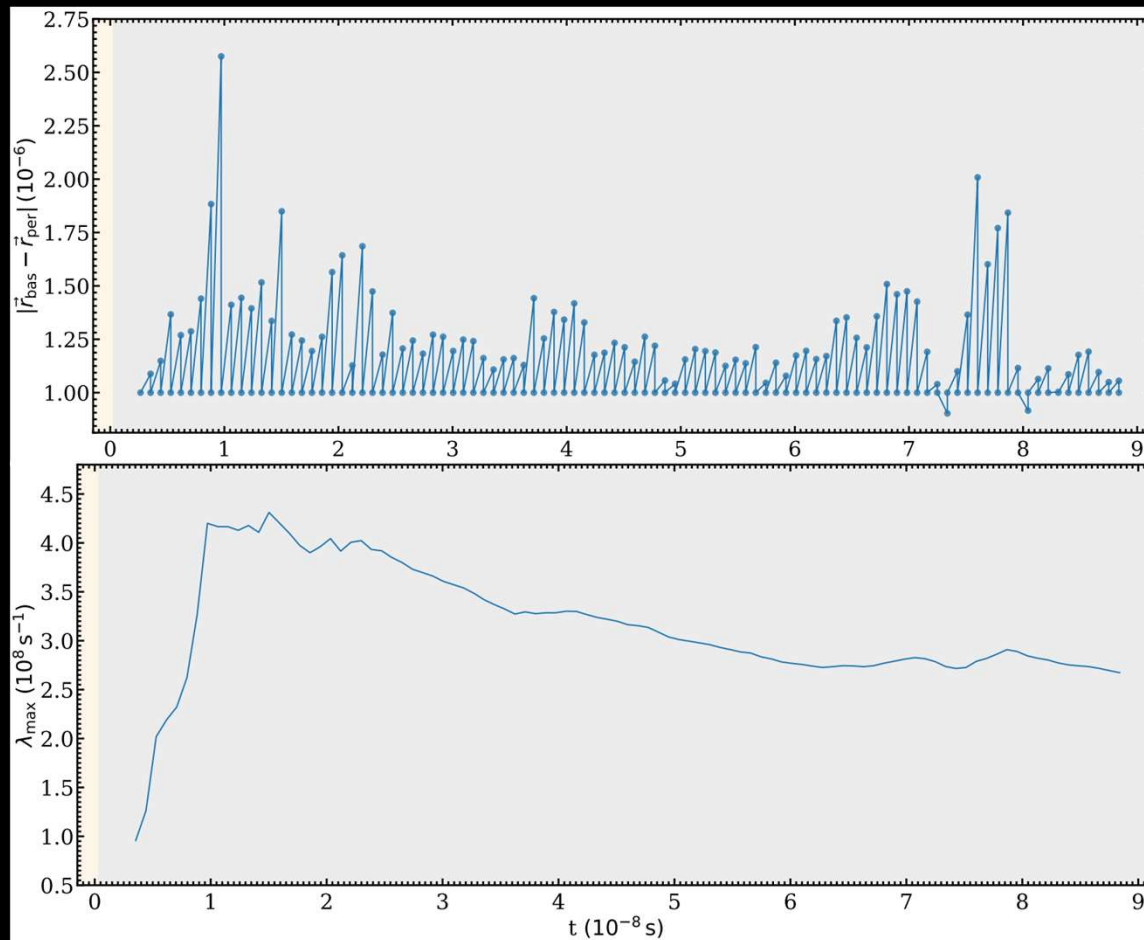
- High-resolution 2D flavor transformation: **3 m**

Nagakura (2023)

- Estimated required resolution: **0.0003 m**

Duan et al. (2006)

# One does not simply resolve the FFI.

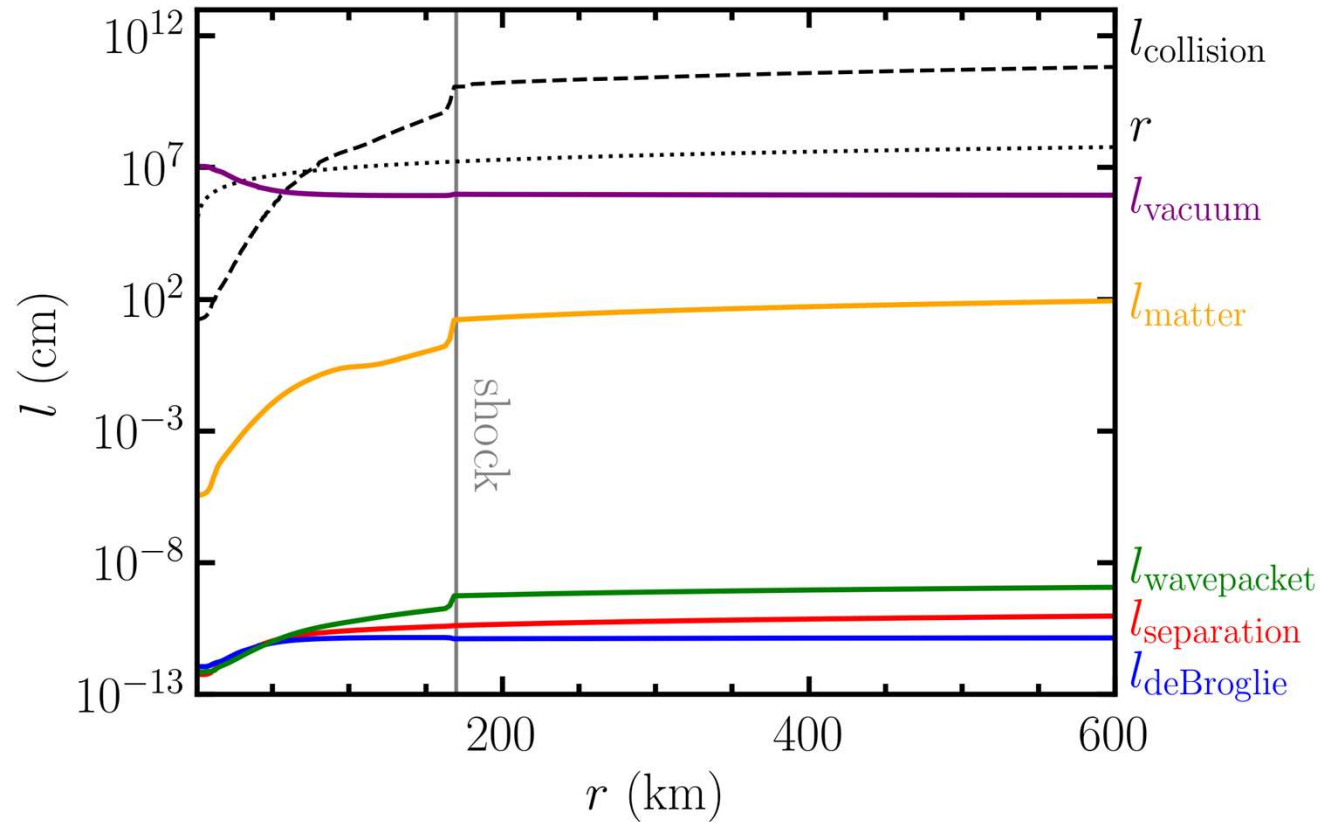


"fiducial  
trajectory"



Erick Urquilla Orellana  
(Licenciatura student @ U. El Salvador)

# How hard could it be?

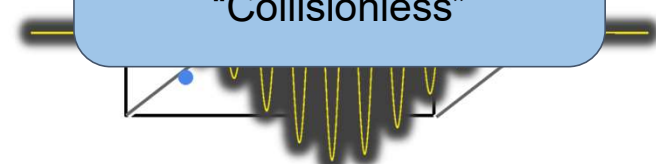


$(1e7)^4 = 1e28$

**Neutrino Decoupling**  
size > mean free path

**Classical Scattering**  
"Instantaneous" collisions

**Quantum Kinetics**  
Flavor changing is fast!  
"Collisionless"



# Theory of Neutrino Quantum Kinetics

$$\frac{\partial f_{ab}}{\partial t} + c\Omega \cdot \nabla f_{ab} = \mathcal{C}_{ab} - \frac{i}{\hbar} [\mathcal{H}, f]_{ab}$$

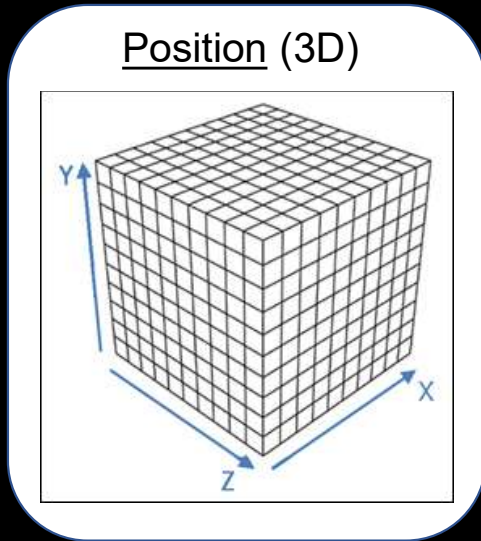
Transport
Collision
Flavor

Vlasenko+ (2014)

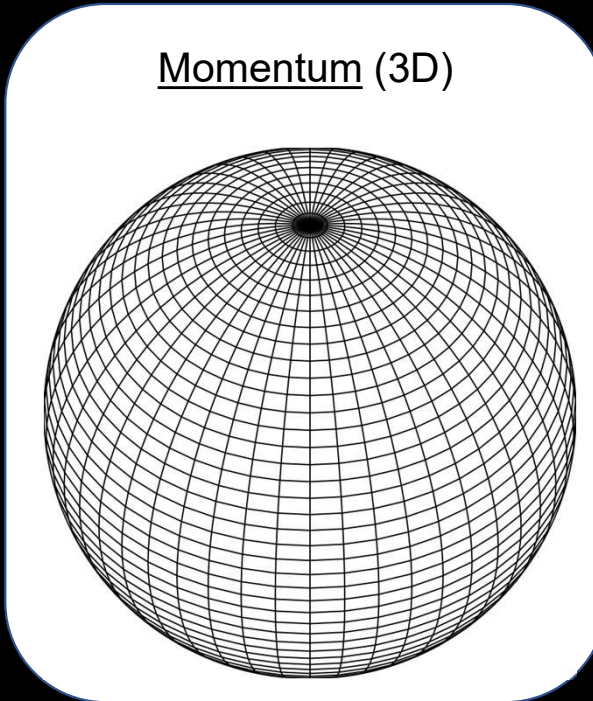
Volpe (2015)

Blaschke &  
Cirigliano (2016)

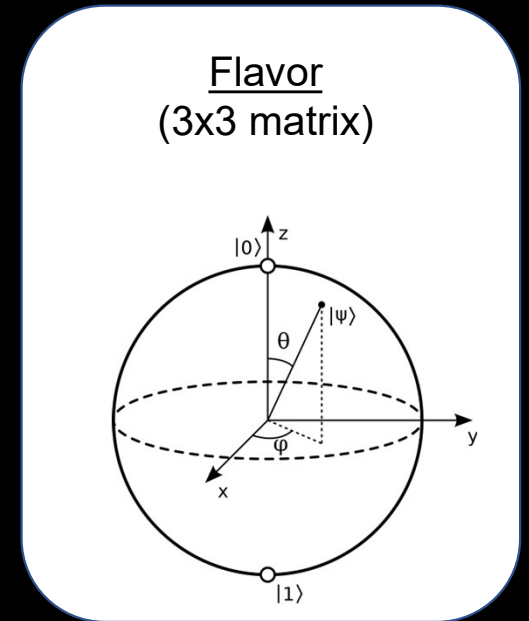
$f_{ab} =$



x



x



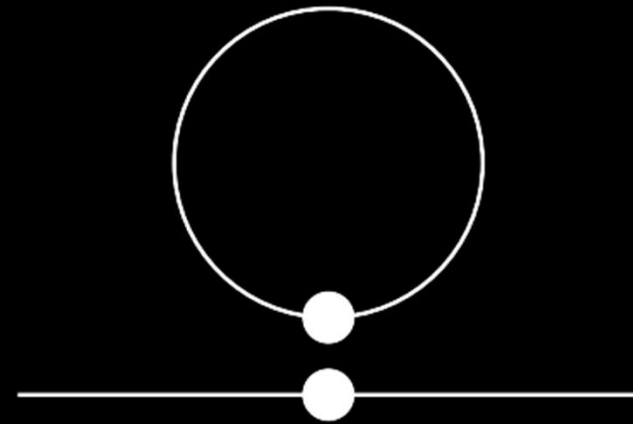
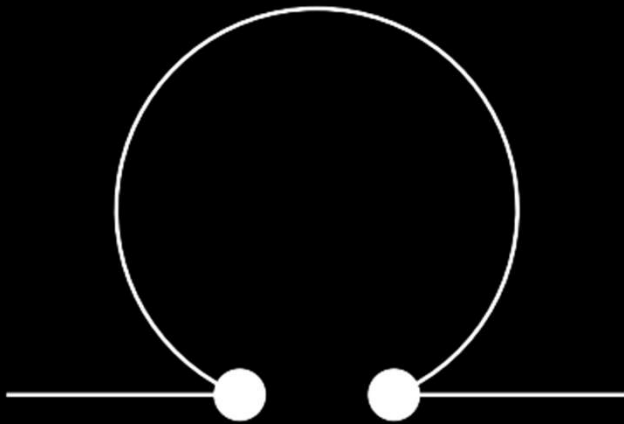
$$\frac{\partial f_{ab}}{\partial t} + c\boldsymbol{\Omega} \cdot \nabla f_{ab} = \mathcal{C}_{ab} - \frac{i}{\hbar} [\mathcal{H}, f]_{ab}$$

Flavor

Vlasenko+ (2014)

Volpe (2015)

Blaschke &  
Cirigliano (2016)





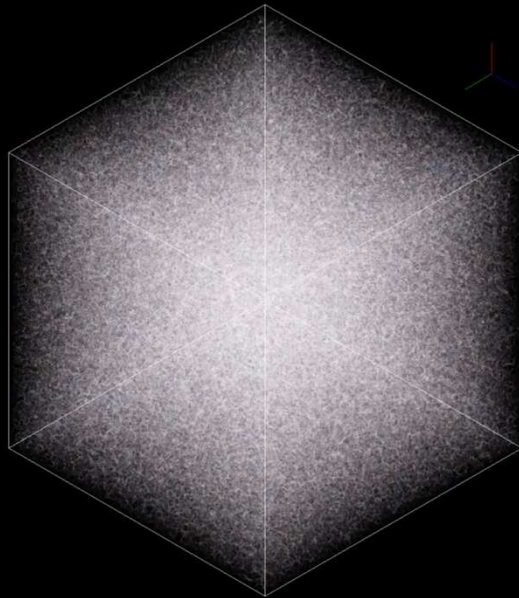
# What does the FFI look like?



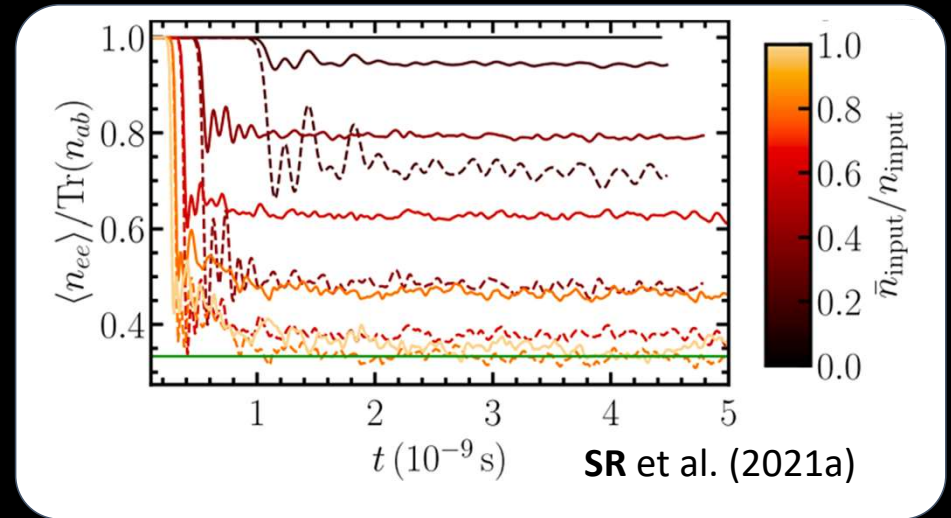
Nicole Ford  
(Berkeley post-bac  
→ McGill grad student)

$t = 0.0000$  ns

8 cm



SR et al. (2021b)



Amount of flavor transformation  
depends on the angular distribution.

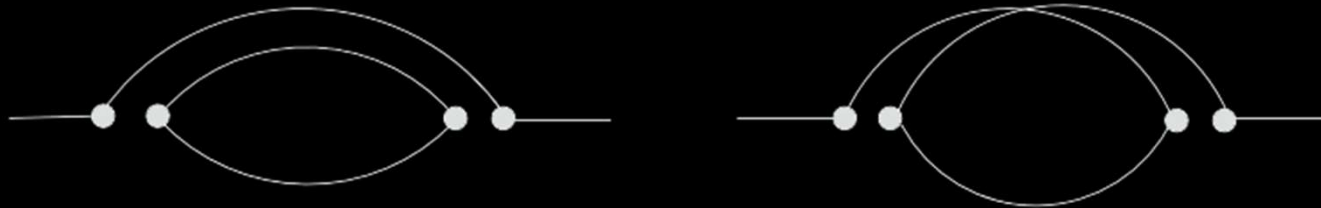
$$\frac{\partial f_{ab}}{\partial t} + c\Omega \cdot \nabla f_{ab} = \boxed{\mathcal{C}_{ab}} - \frac{i}{\hbar} [\mathcal{H}, f]_{ab}$$

Vlasenko+ (2014)

Volpe (2015)

Blaschke &  
Cirigliano (2016)

“The Supernova Problem”



### Neutrino Transport Reviews

Bruenn (1985)

Burrows, Reddy, Thompson (2007)

Mezzacappa (2022)

### Combining with one-loop effects

Cherry (2012)

Vlasenko (2017)

Vlasenko & McLaughlin (2018)

SR et al. (2019)

Shalgar & Tamborra (2020, 2022)

Johns (2021)

Martin et al. (2021)

Sasaki et al. (2021)

Nagakura (2022)

Hansen et al. (2022)

Johns & Xiong (2022)

Kato & Nagakura (2022)

Padilla-Gay et al. (2022)

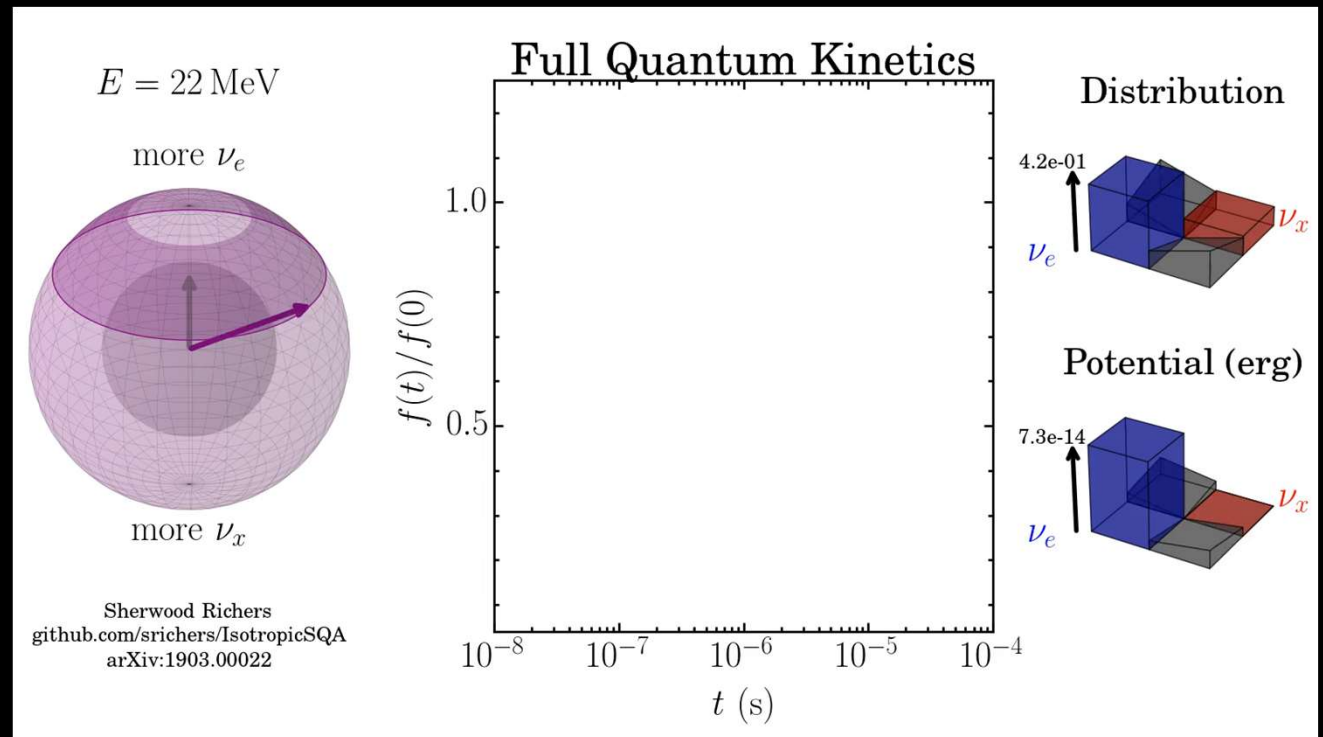
Kato, Nagakura, & Zaizen (2023)

Lin & Duan (2023)

Xiong et al. (2023)

$$\frac{\partial f_{ab}}{\partial t} + c\Omega \cdot \nabla f_{ab} = C_{ab} - \frac{i}{\hbar} [\mathcal{H}, f]_{ab}$$

Oscillations and collisions  
are not generally separable



Richers+ (2019)

# Collisional Processes

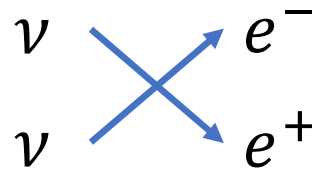
Abs. & Emis.



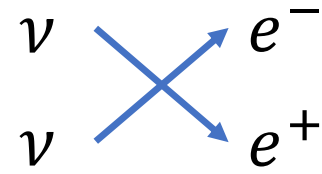
Scattering



Pair Annihilation  
Bremsstrahlung



4-neutrino  
Processes



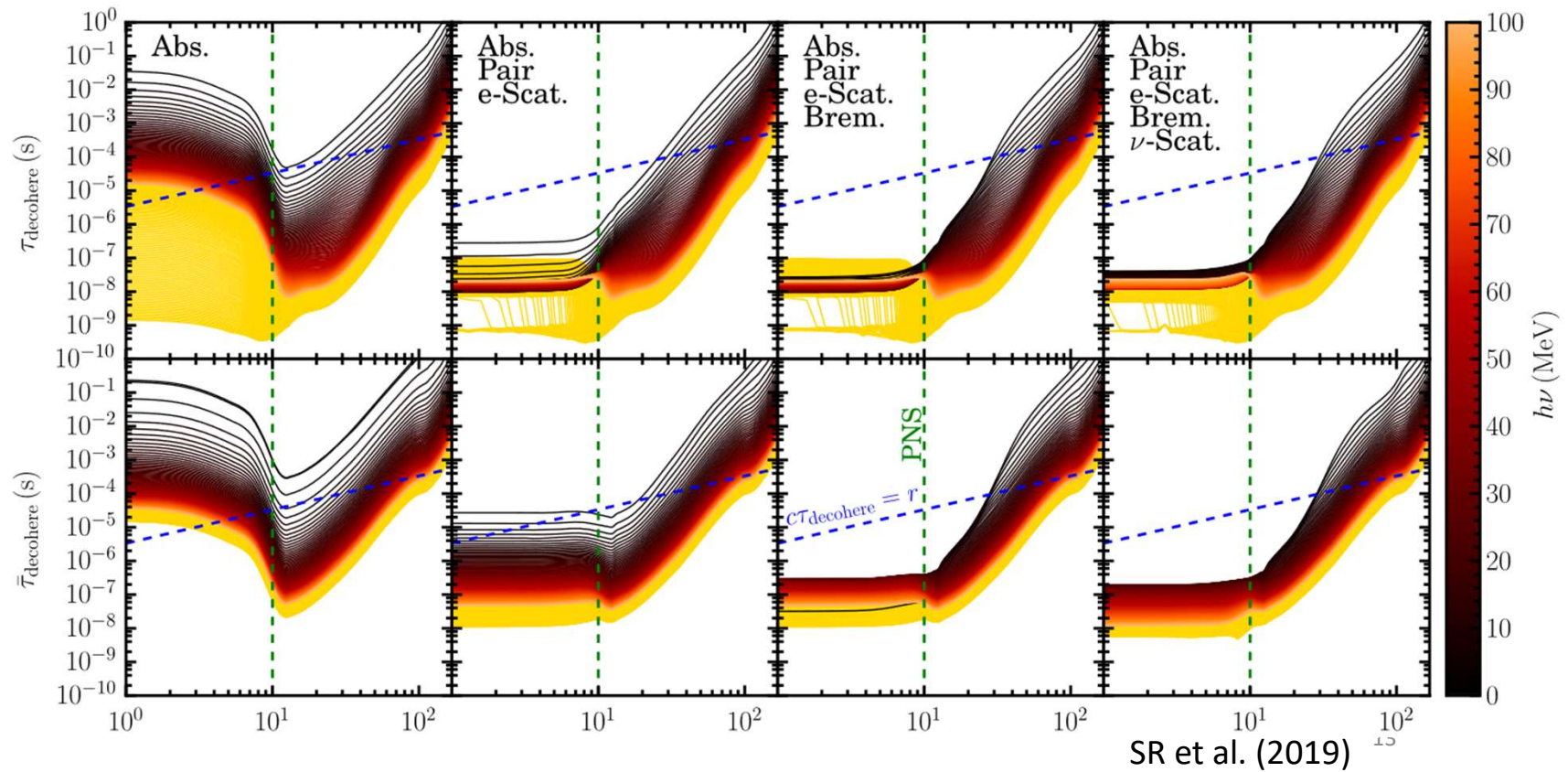
$$\Pi_{ab}^+ = \int \frac{d^3\nu_1'}{c^4} \langle R \rangle_{ab}^+ f'_{1ab},$$

$$\Pi_{ab}^- = \int \frac{d^3\nu_1'}{c^4} \langle R \rangle_{ab}^- (\delta_{ab} - f'_{1ab}),$$

$$R_{(\nu_a)}^+ = \sum_c (1 + \delta_{ac}) \int \frac{d^3\nu_2'}{c^3} \frac{d^3\nu_3'}{c^3} \\ \times r_{(p_1+p_3 \rightarrow p+p_2)} (1 - f'_{2cc}) f'_{3cc},$$

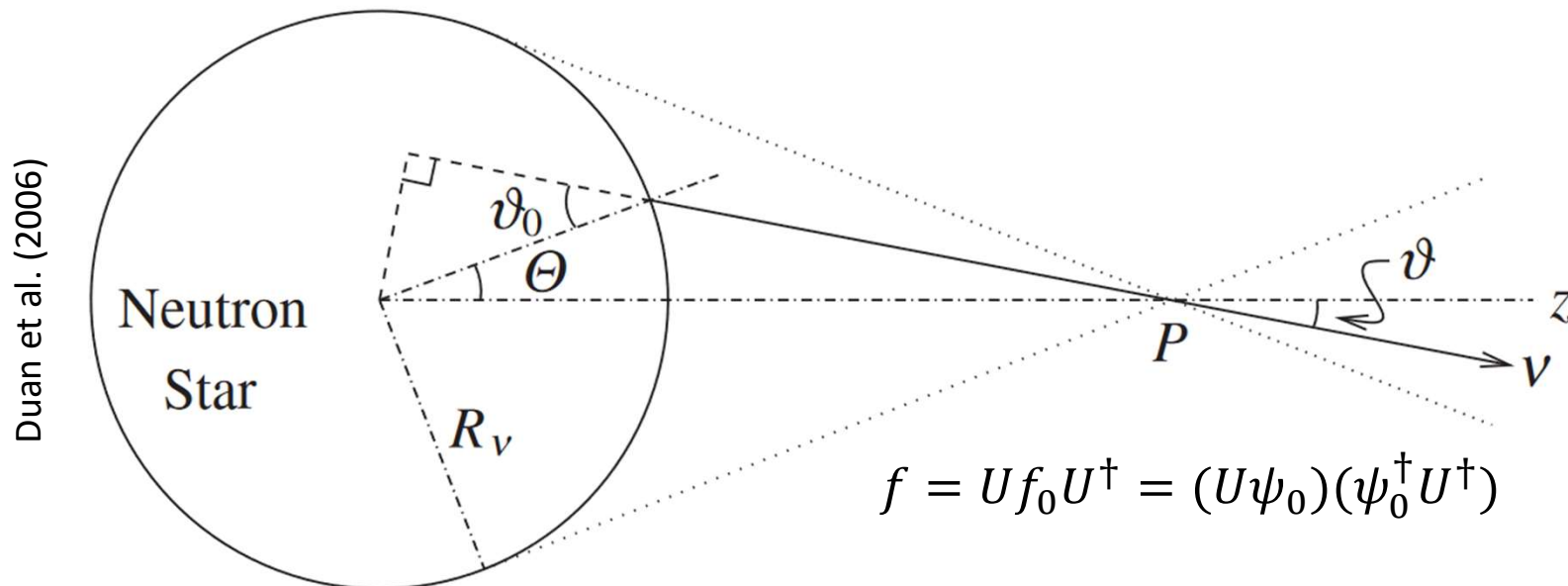
$$R_{(\nu_a)}^- = \sum_c (1 + \delta_{ac}) \int \frac{d^3\nu_2'}{c^3} \frac{d^3\nu_3'}{c^3} \\ \times r_{(p+p_2 \rightarrow p_1+p_3)} f'_{2cc} (1 - f'_{3cc})$$

# Multiple collision processes matter



# Simulation of Neutrino Quantum Kinetics

# Bulb Model (Dirichlet boundary conditions)



→ Numerical demonstrations of collective oscillations, MNR, Halo effect

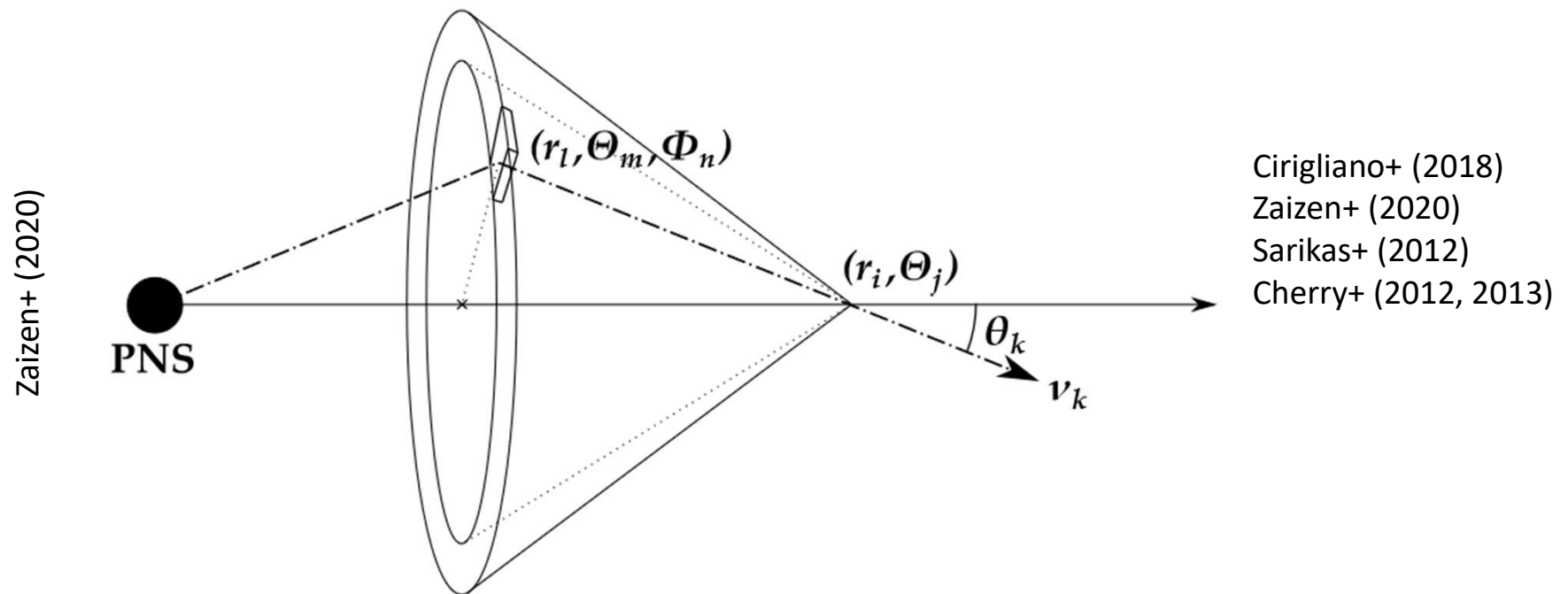
(see also Galais+2012, Malkus+2012, Tian+2017, many more)

Evolve  $U$  instead of  $f$  (unitary operator)

Evolve OUTWARD → 1+0 dimensional

(Single-angle approximation)

# Bulb Model (Halo Effect)





$$\frac{\partial f_{ab}}{\partial t} + c\Omega \cdot \nabla f_{ab} = \mathcal{C}_{ab} - \frac{i}{\hbar} [\mathcal{H}, f]_{ab}$$

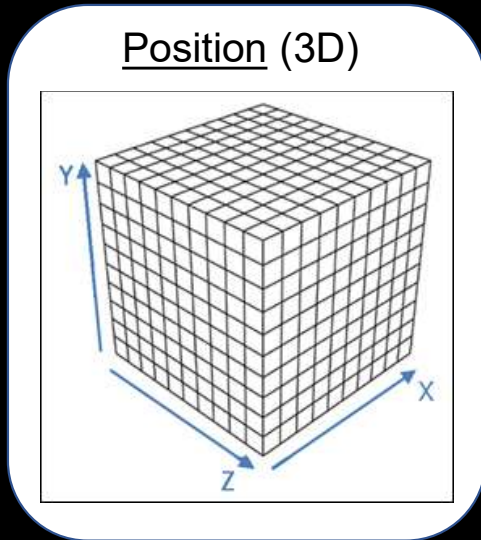
Transport
Collision
Flavor

Vlasenko+ (2014)

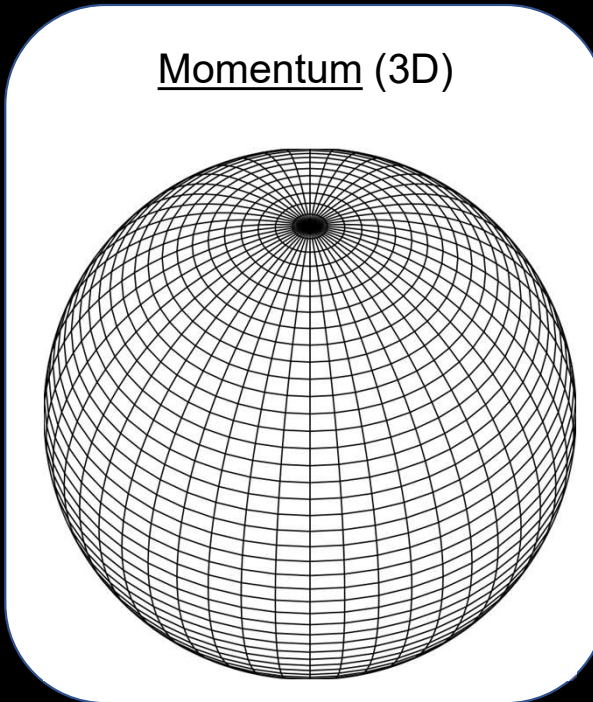
Volpe (2015)

Blaschke &  
Cirigliano (2016)

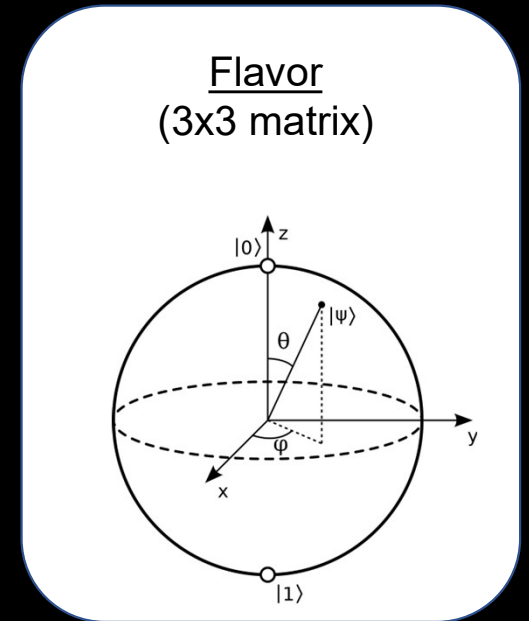
$f_{ab} =$



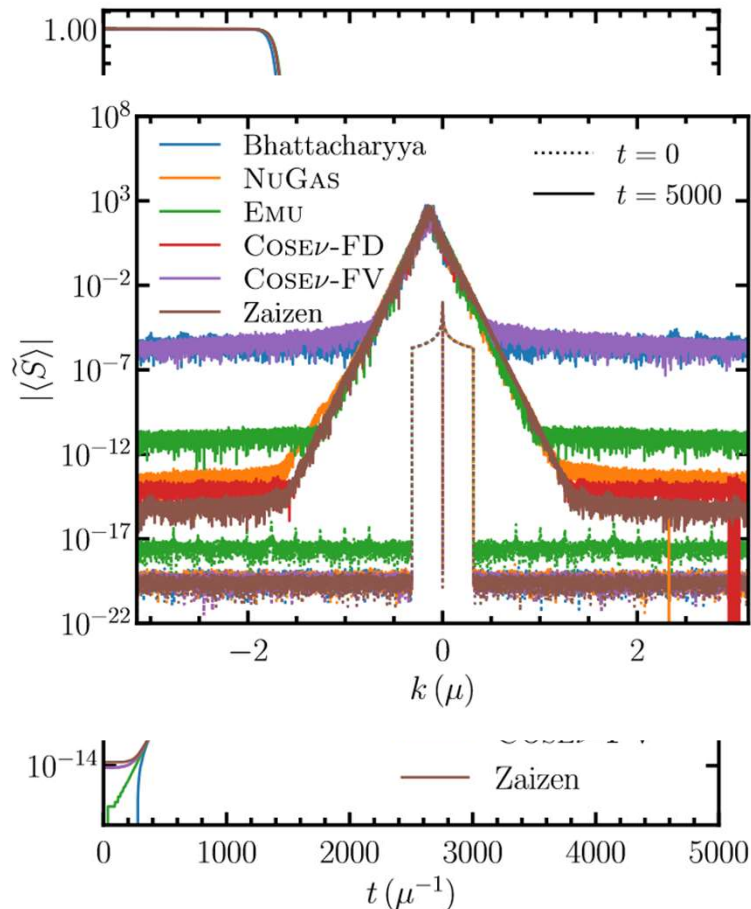
x



x



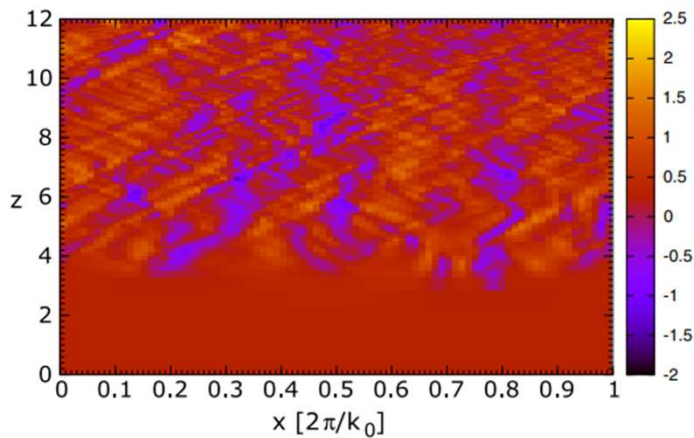
# General Features of the FFI



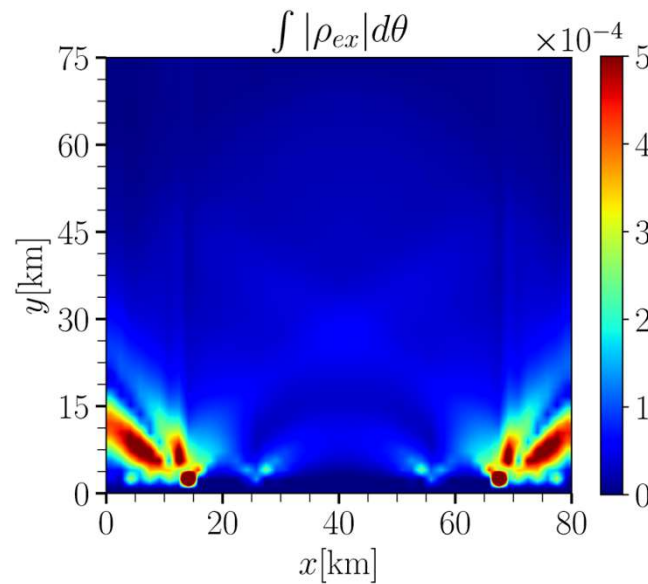
SR+ (2022), following many other works

1. Exponential growth of perturbations  
Sawyer (2005), Dasgupta, Sen, Mirizzi, Morinaga, Padilla-Gay, Abbar, Xiong, Wu, Bhattacharyya, Zaizen, George, Duan, Sigl, Capozzi, Shalgar, Raffelt, Chakraborty, Kato ... [many contributions]
2. Complete mixing within “ELN Crossing”, incomplete elsewhere to preserve lepton #  
Bhattacharyya & Dasgupta (2021)
3. Modes spreading to exponential distribution.  
SR et al. (2021)
4. Coherent post-saturation flavor wave  
Duan et al. (2021)
5. Non-trivial interplay with collisions  
Padilla-Gay, Shalgar, Johns, Xiong, Sasaki, Sigl, Tamborra, Hansen, Martin, SR

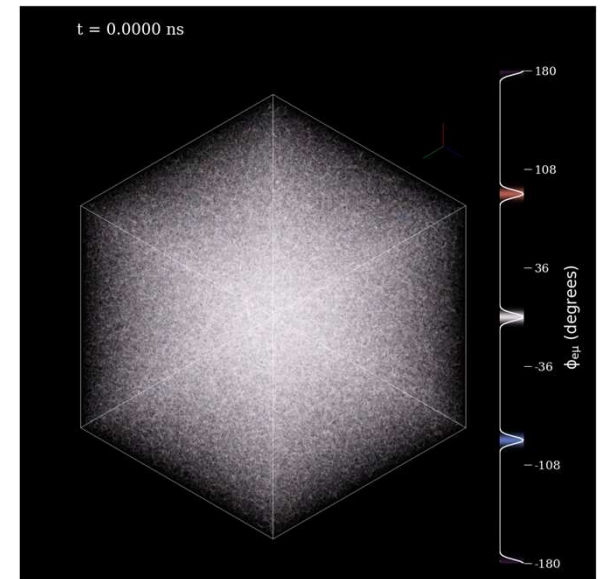
# Multiple dimensions allow broken symmetries



Mirizzi+ (2015)



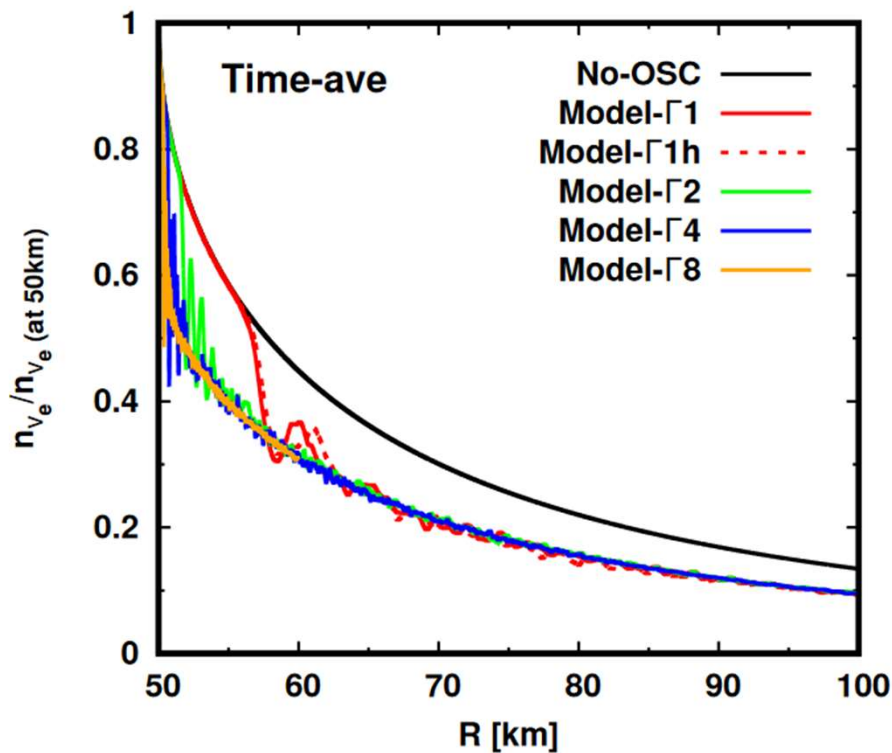
Padilla-Gay+ (2020)



SR et al. (2021b)

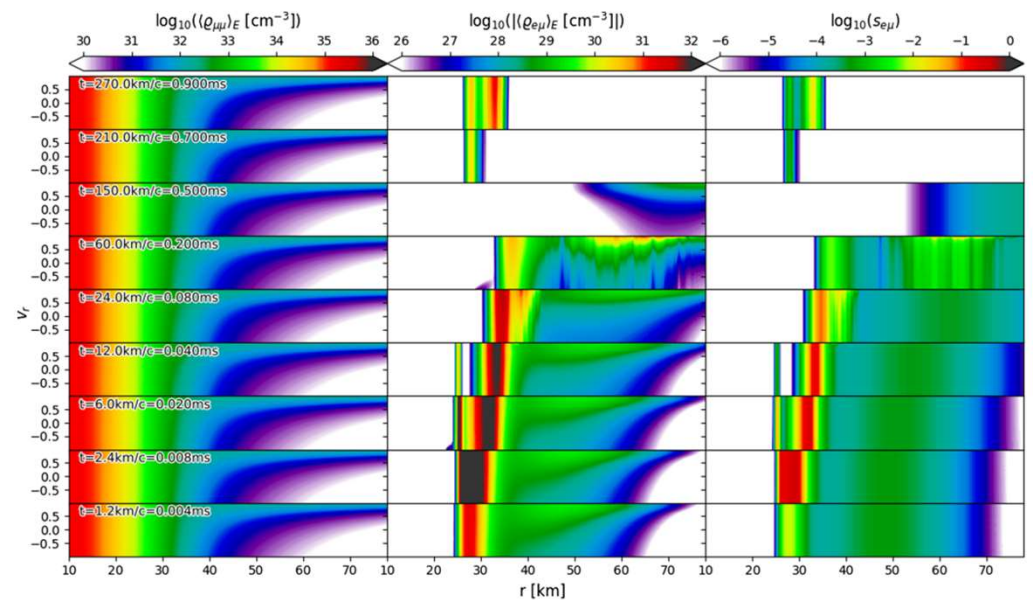
Local FFI in 3D is similar to well-constructed 1D model

# Reduced coupling enables global analysis



Nagakura & Zaizen (2022)

→ FFI can modify CCSN and NSM outcomes



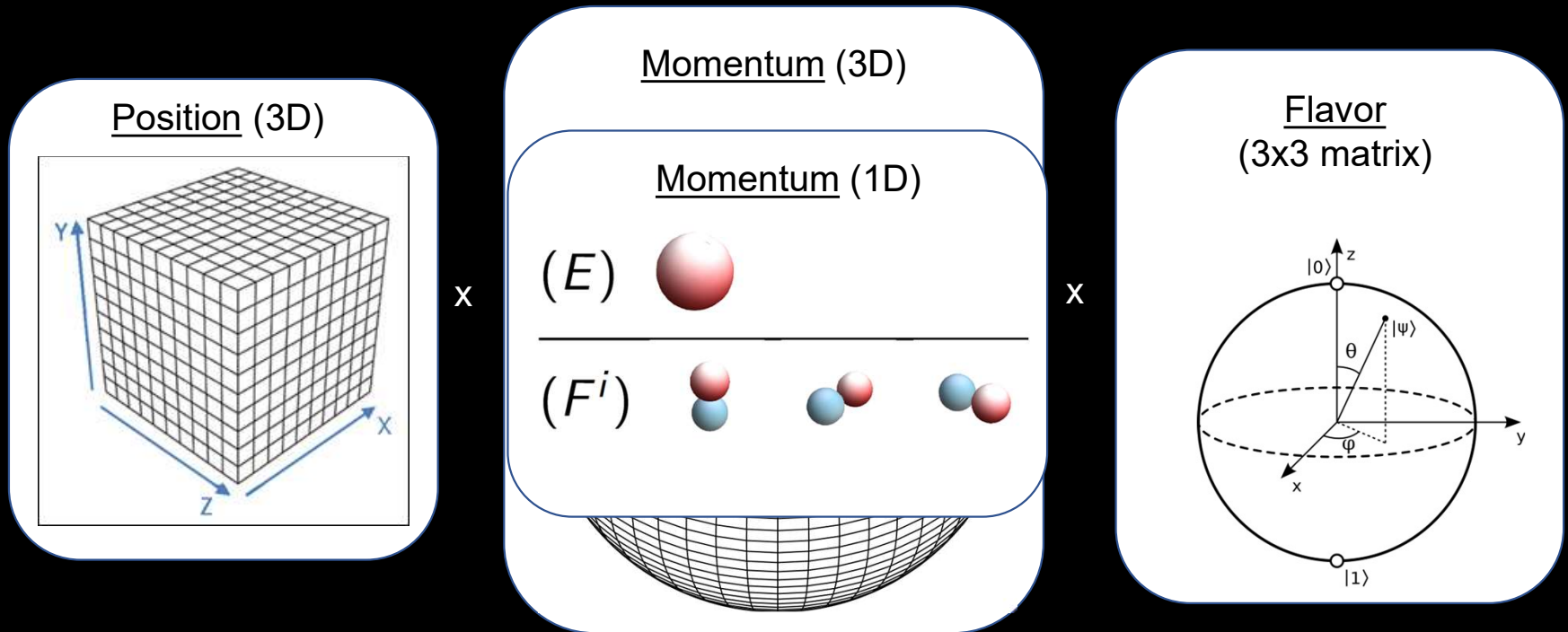
Xiong+ (2023)

Collisional instability significantly modifies heavy lepton neutrinos in CCSNe

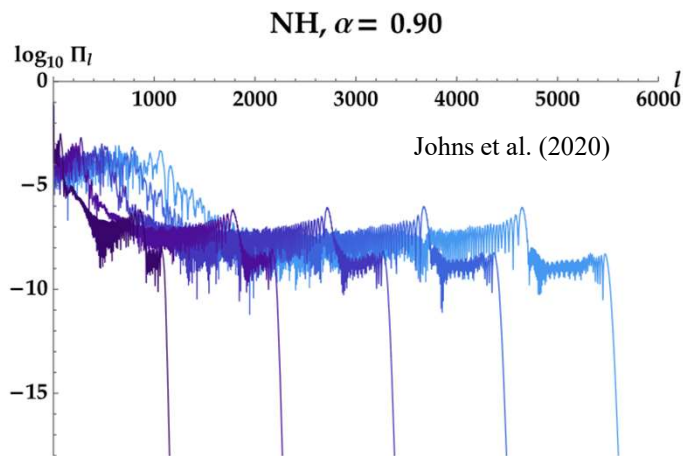
$$\frac{\partial f_{ab}}{\partial t} + c\Omega \partial_\alpha T_{ab}^{\alpha\beta} = \mathcal{C}_{ab} - \frac{i}{\hbar} \left[ \mathcal{H}_{\alpha,ab}, T_{ab}^{\alpha\beta} \right]$$

Transport
Collision
Flavor

$f_{ab} =$

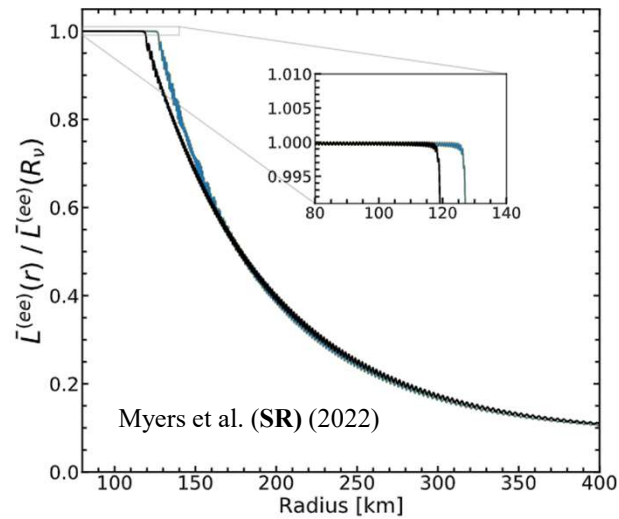


# Moments are fast, but face difficulties



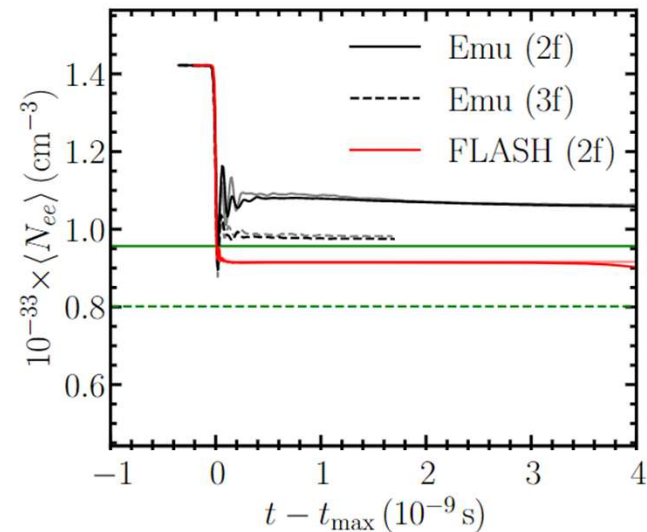
Moments with  $k=0$  FFI

Power goes to high- $l$  moments



Moments with bulb model

Reproduction of collective oscillations



Grohs, SR et al. (2022)

30x cheaper than Emu



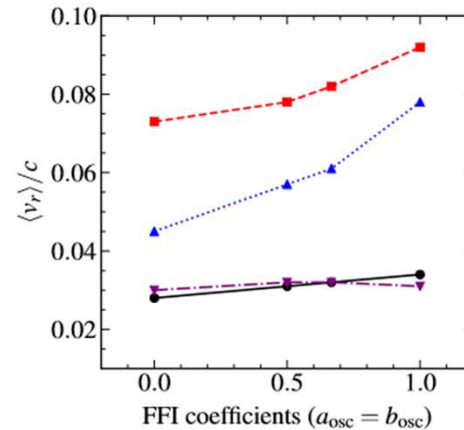
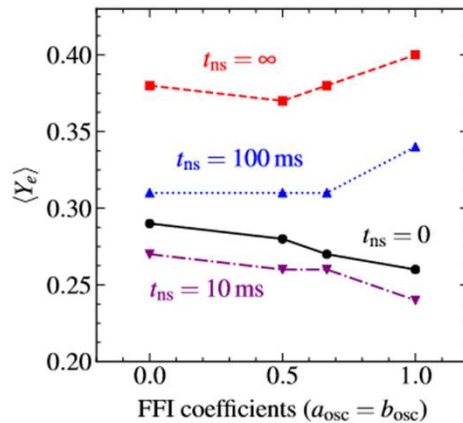
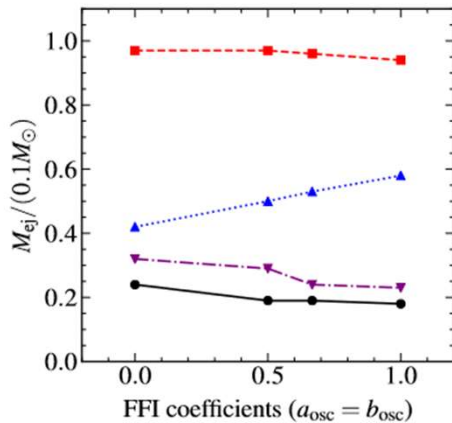
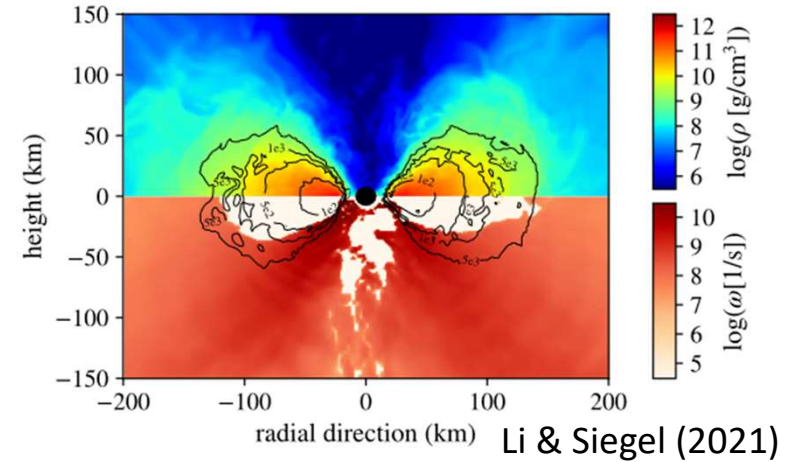
# Effective models probe sensitivity

Li & Siegel 2021: [k=0 dispersion] FFI → more neutron-rich outflow

Just+2022: [Polynomial] FFI → modest neutron-rich enhancement

Fernández+2022: [Opt. depth] FFI → Long-lived HMNS can reverse neutron enhancement

Ehring+2023: [Density cutoff] FFI can help (low-mass) or hinder (high-mass) CCSN explosion



Fernández, SR et al (2022)

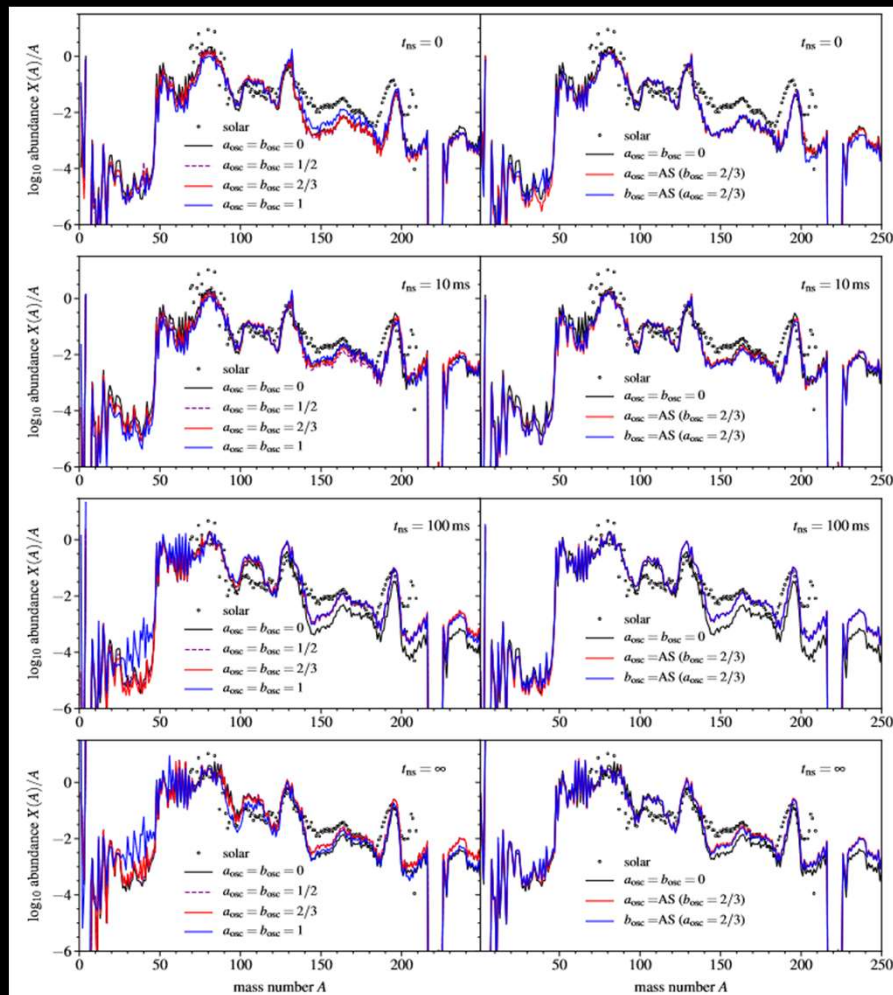
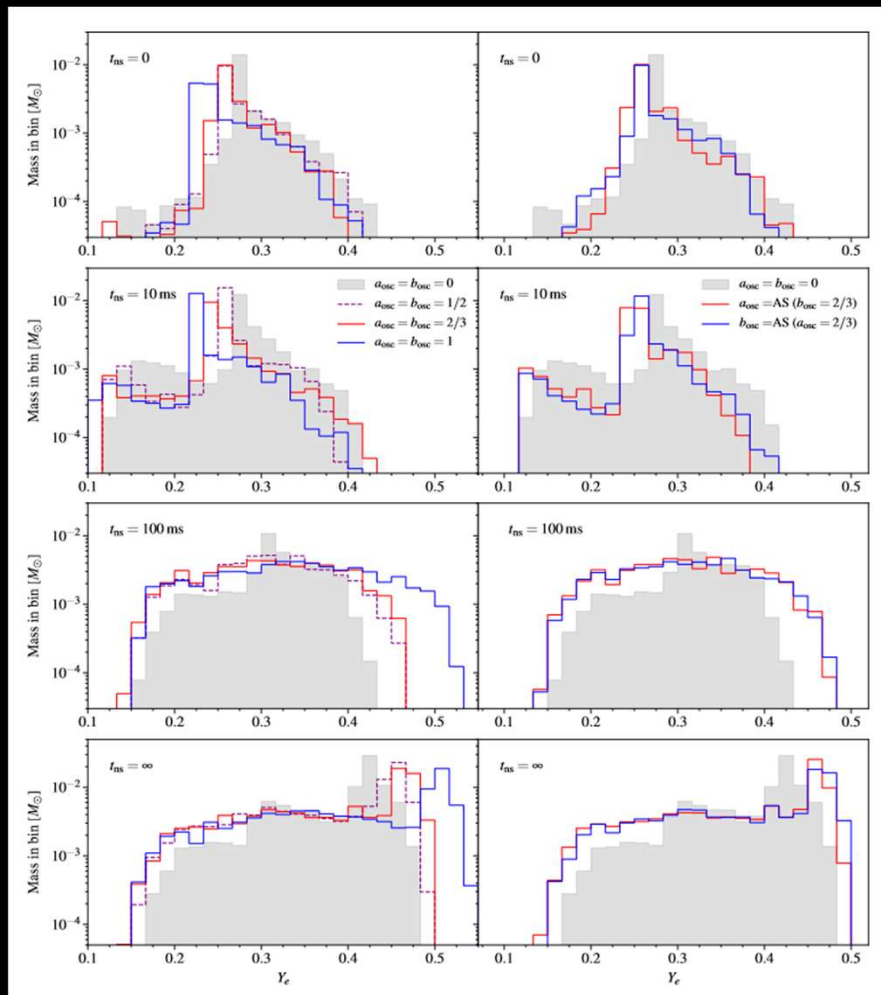
Multi-angle effective models

Bhattacharyya & Dasgupta (2022)

Zaizen & Nagakura (2023)



# Expect FFI to have a moderate impact on outflows



# Replacing Simulation with Machine Learning

## Invariance

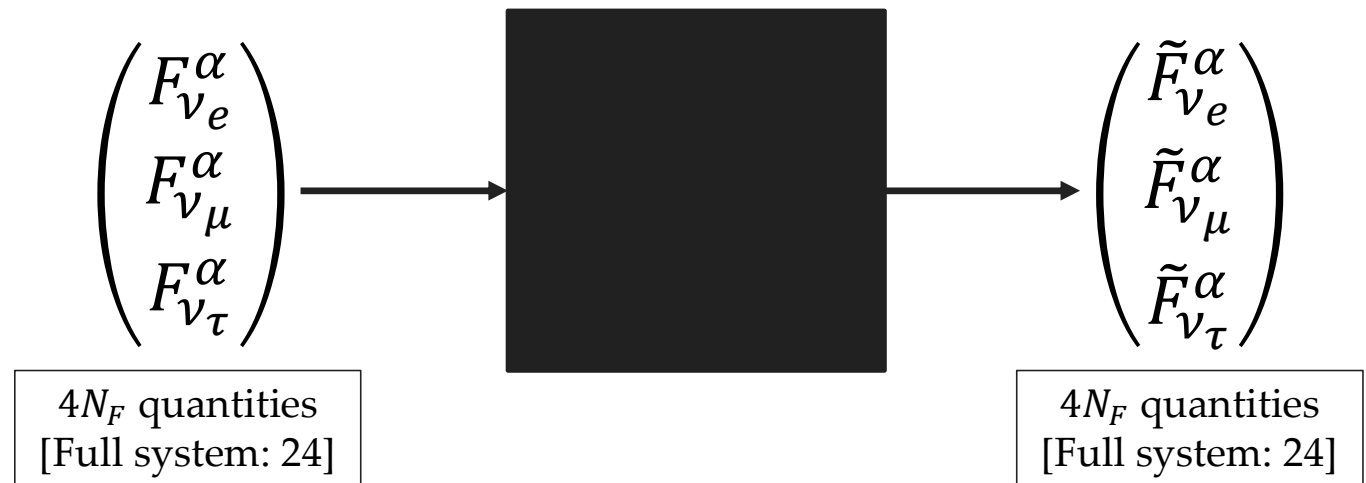
- Rotation (Lorentz)
- $v_i \leftrightarrow v_j$
- $v_i \leftrightarrow \bar{v}_i$

## Conservation

- $\sum_i F_{v_i}^\alpha$  and  $\sum_i F_{\bar{v}_i}^\alpha$
- $\sum_i (F_{v_i}^t - F_{\bar{v}_i}^t)$

## Other

- Do exactly nothing when stable
- $\tilde{F}$  must be stable
- Flux factor  $< 1$
- Positive density



# Replacing Simulation with Machine Learning

## Invariance

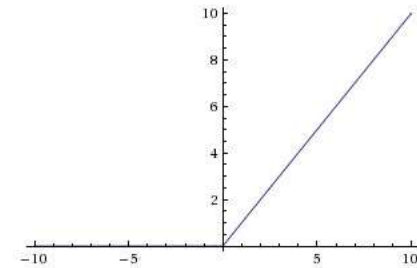
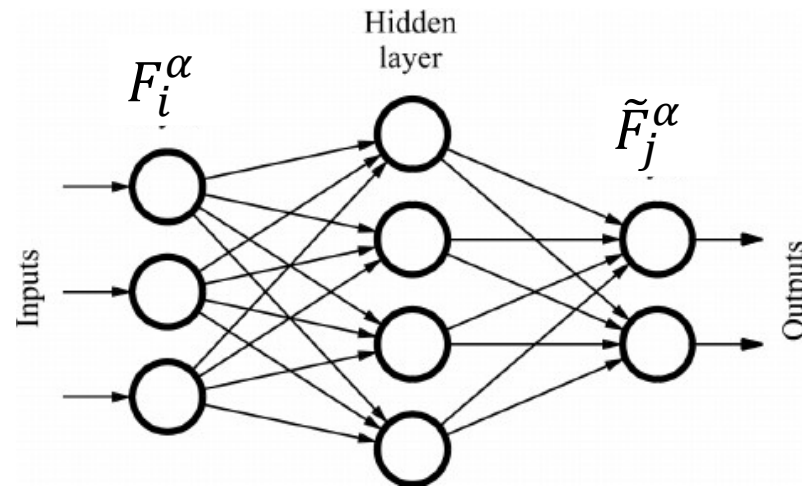
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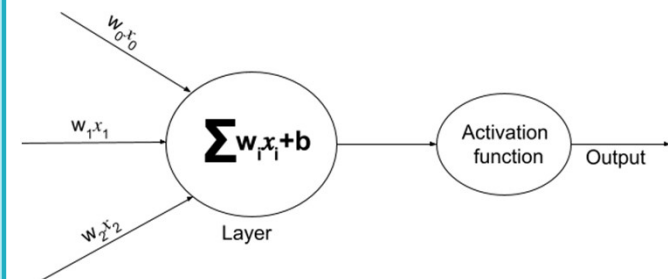
## Other

- Flux factor  $< 1$
- Positive density



## Feed-forward NN

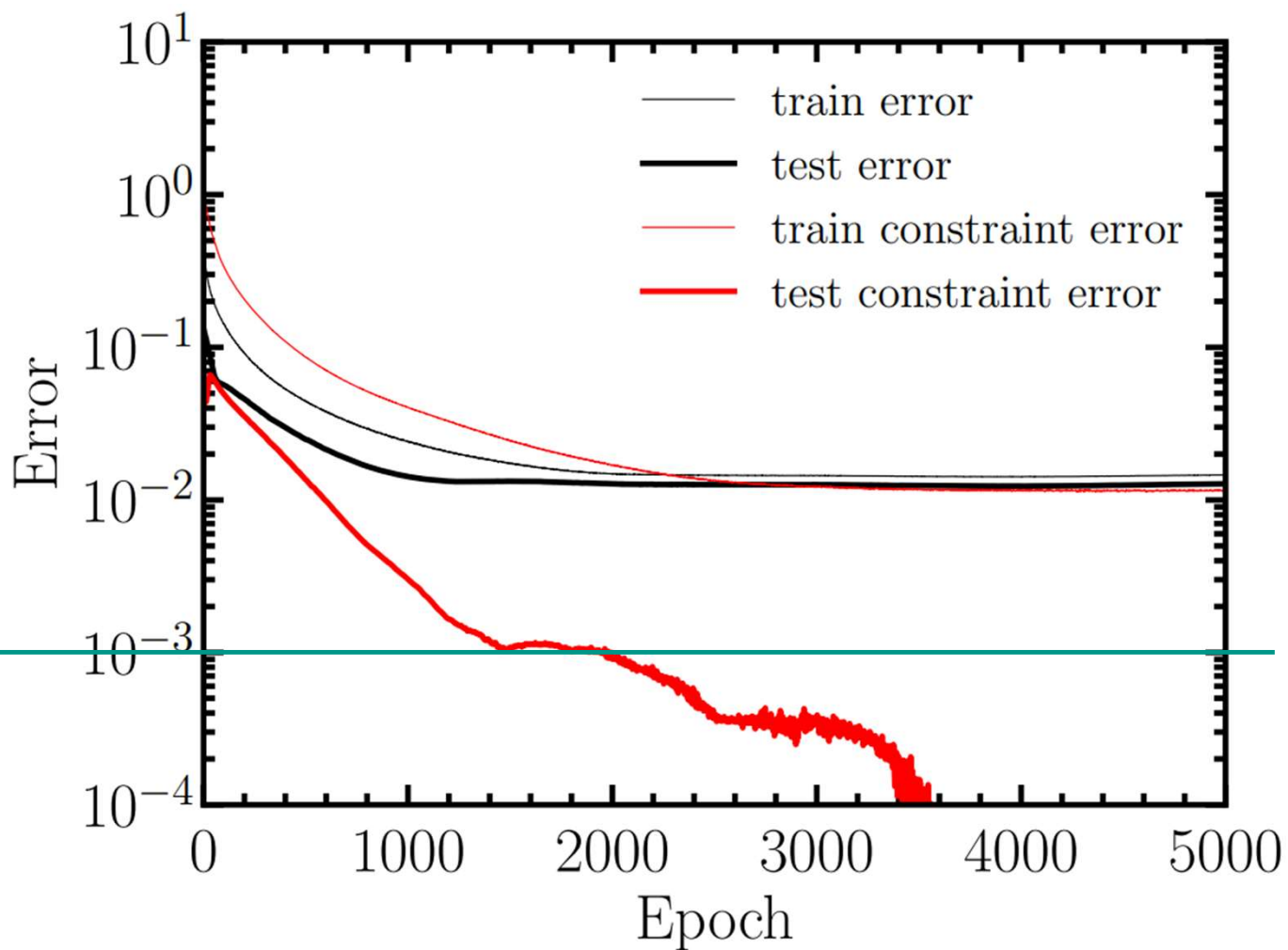
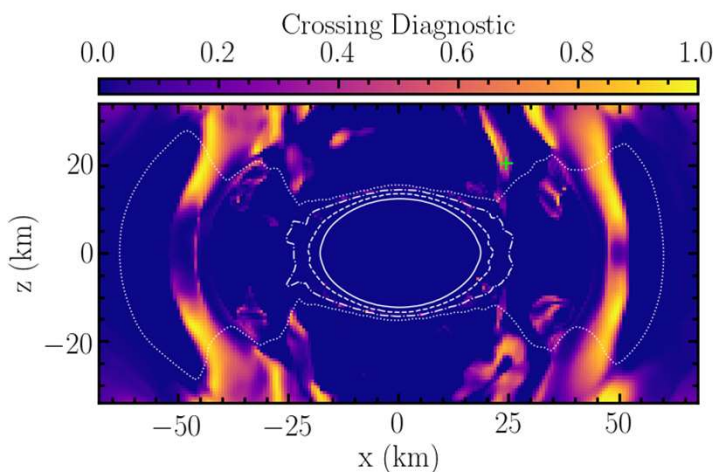
- 300 samples (plus augmentation)
- 1 hidden layer, 32 wide
- ADAM optimizer
- 50% dropout probability
- Batch Normalization
- ReLU activation
- Training time: 1 minute



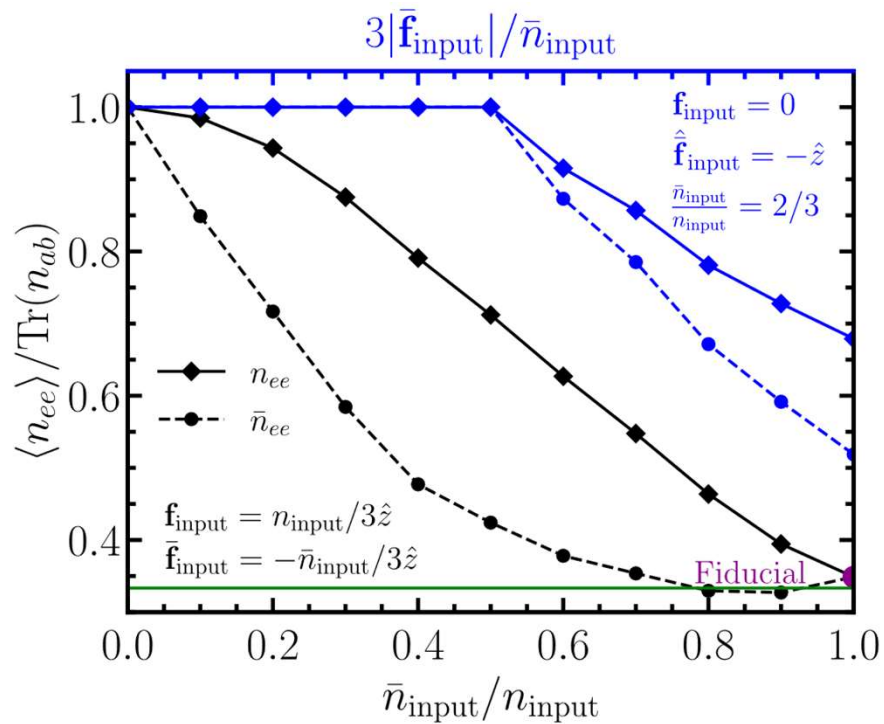
# Replacing Simulation with Machine Learning

On average:

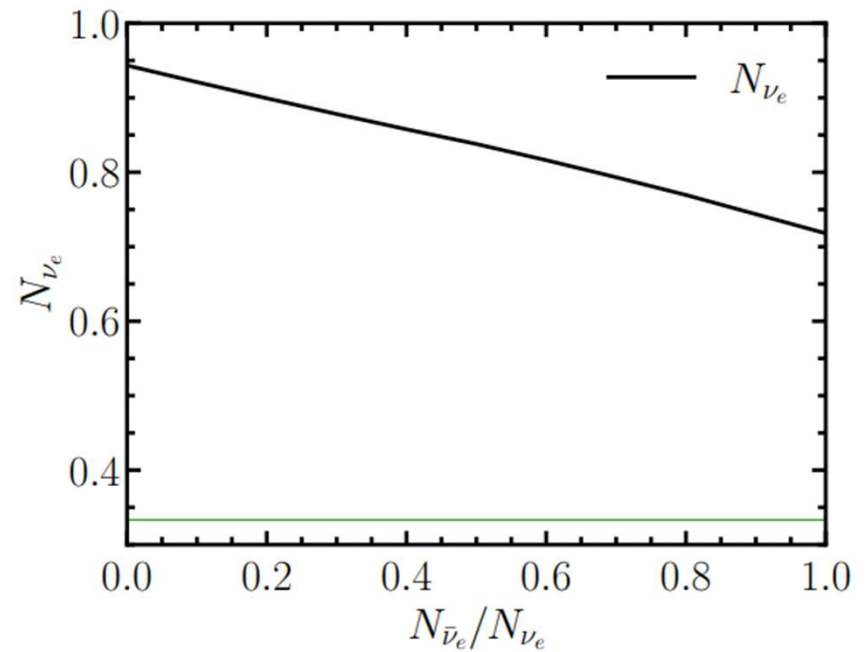
- Percent-level accuracy
- Learned constraints



# A Warning about Parameter Space



SR et al. (2021) using Emu



Neural Network

Results generalize poorly outside the training data

# The Future

Better Effective Models

Larger exact calculations

More complete microphysics

Better approximate methods

Expect rapid development in coming years!

General adoption into dynamical models  
requires reliable predictions