
Progress Towards Radiative Corrections in Nuclear Electroweak Matrix Elements through the Renormalization Group

Thomas R. Richardson
N3AS
UC Berkeley & LBNL

Testing the Standard Model in Charged-
Weak Decays
INT
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In collaboration with Immo C. Reis

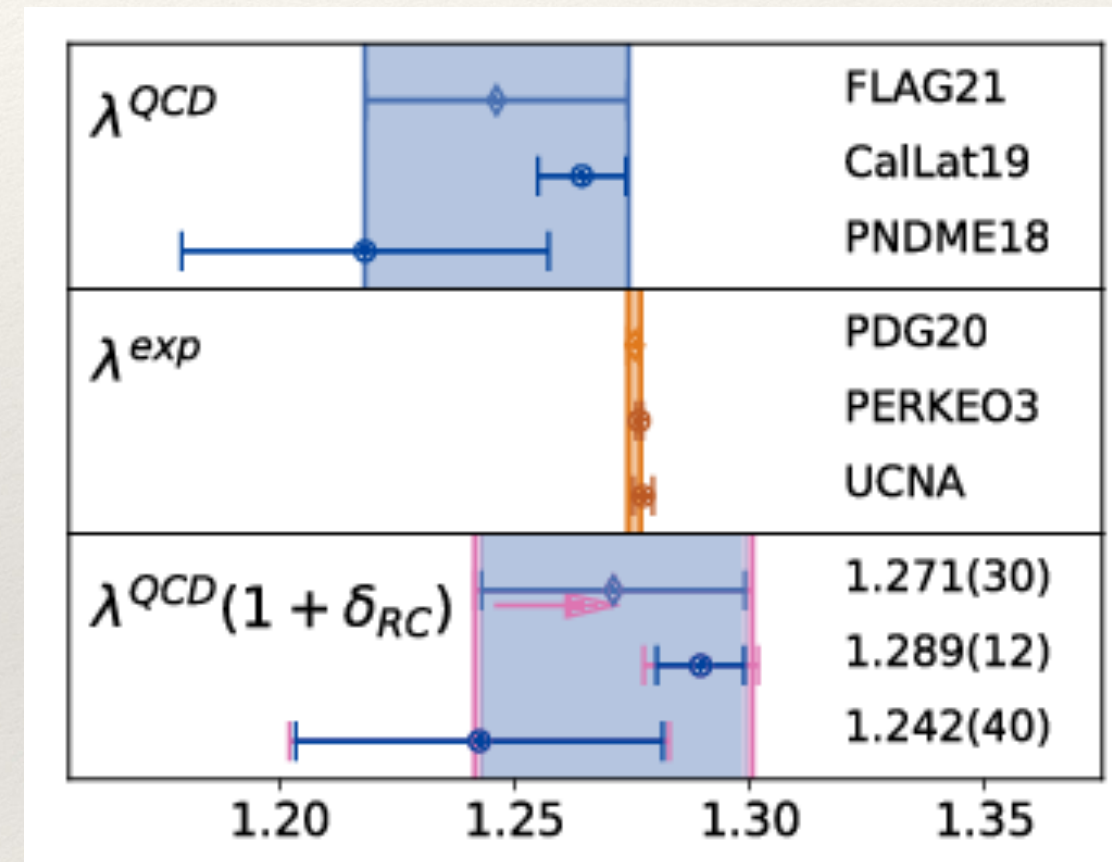
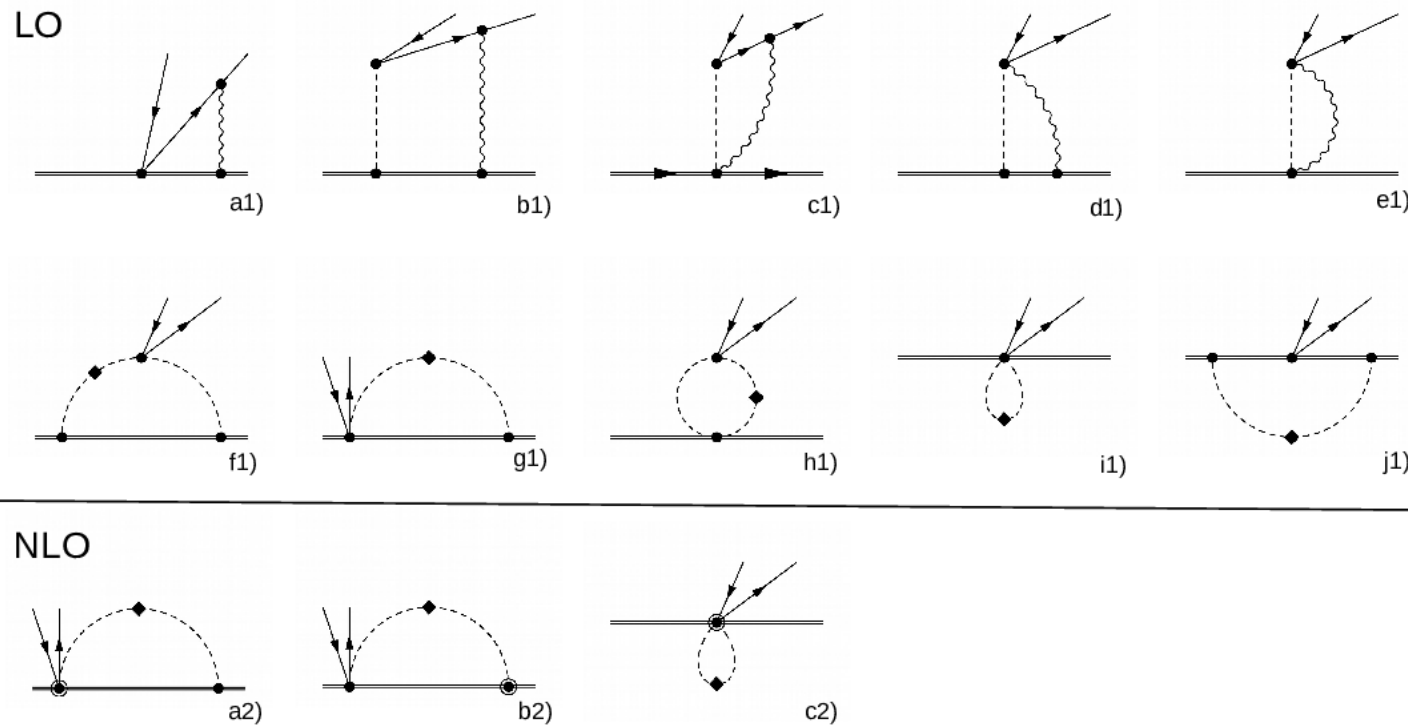
Few Body Syst. 65, 79 (2024)

PRC 111, 064011 (2025)

In preparation

INT 20R-2b: Beyond the Standard Model Physics with Nucleons and Nuclei

- ❖ Electromagnetic corrections to g_A at the % level



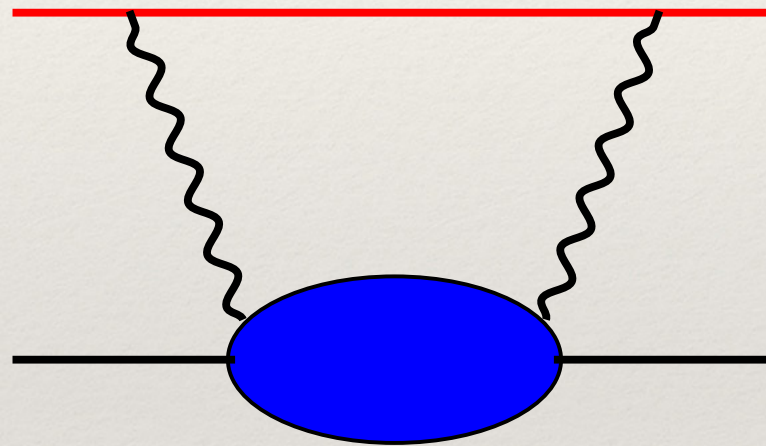
Cirigliano et al. PRL 129

- ❖ What would happen in few-body systems?

Radiative Corrections in Nuclear Physics

Spectroscopy

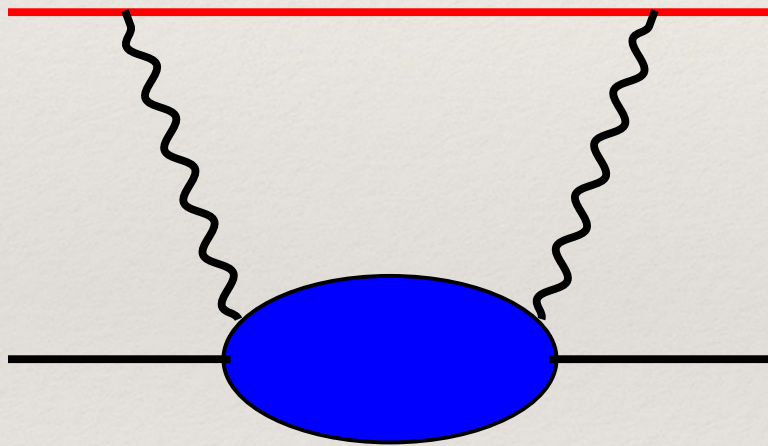
Lepton Scattering



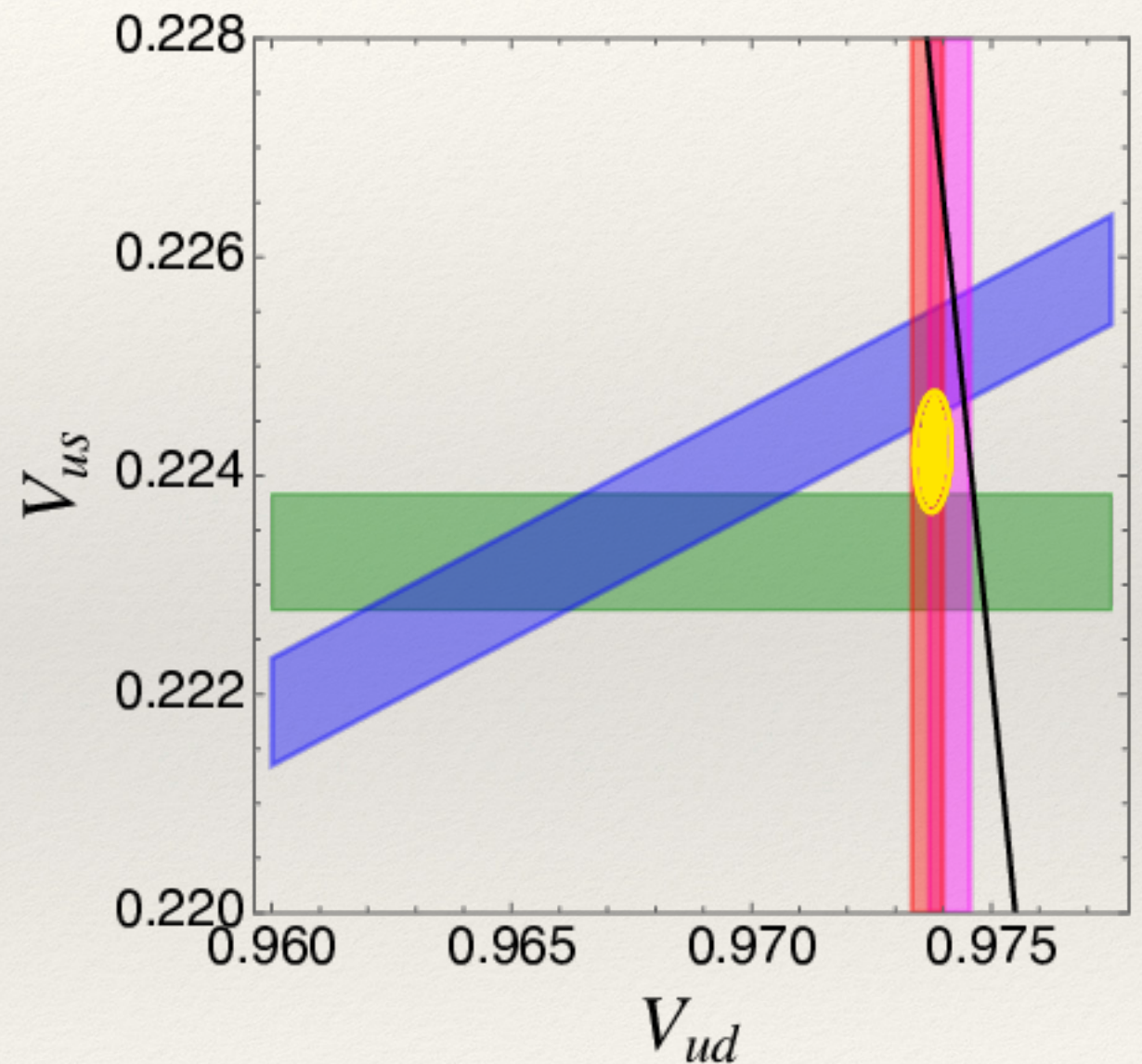
Beta Decay

Radiative Corrections in Nuclear Physics

Beta Decay



Cirigliano et al. PLB 838



$K \rightarrow \pi l \nu$

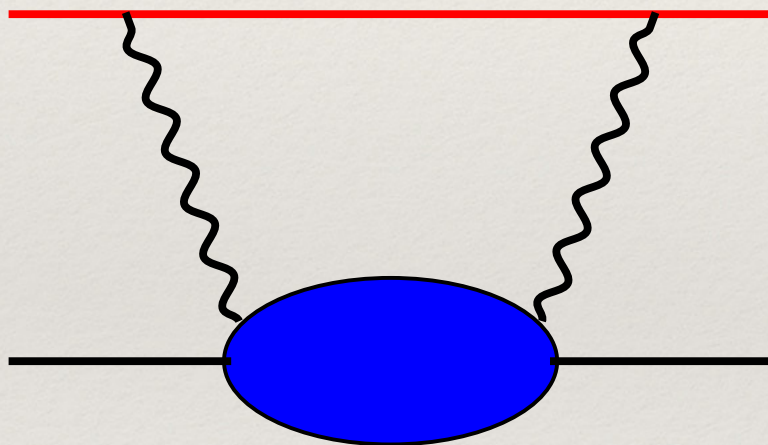
neutron

$K \rightarrow \mu \nu / \pi \rightarrow \mu \nu$

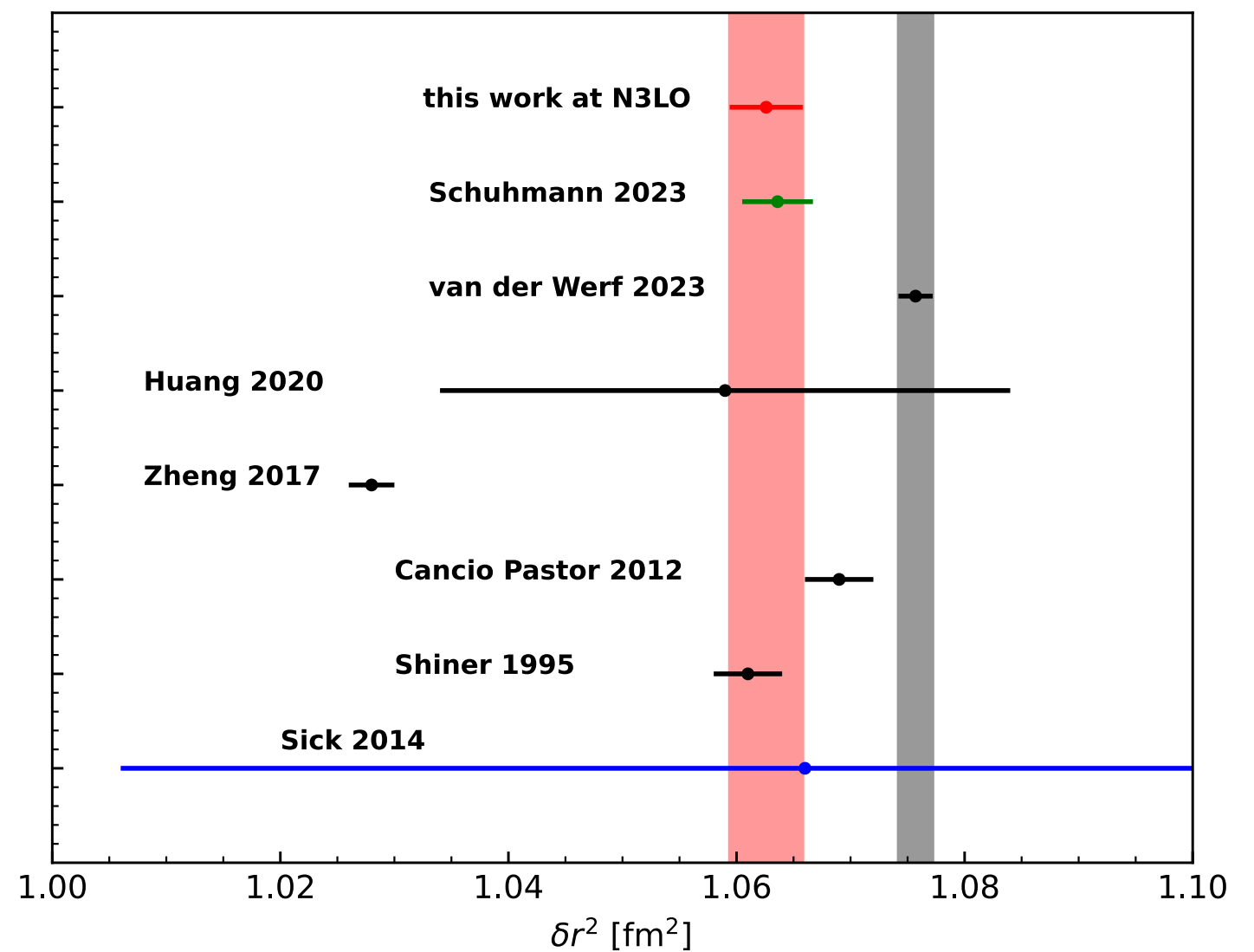
$0^+ \rightarrow 0^+$

Radiative Corrections in Nuclear Physics

Spectroscopy



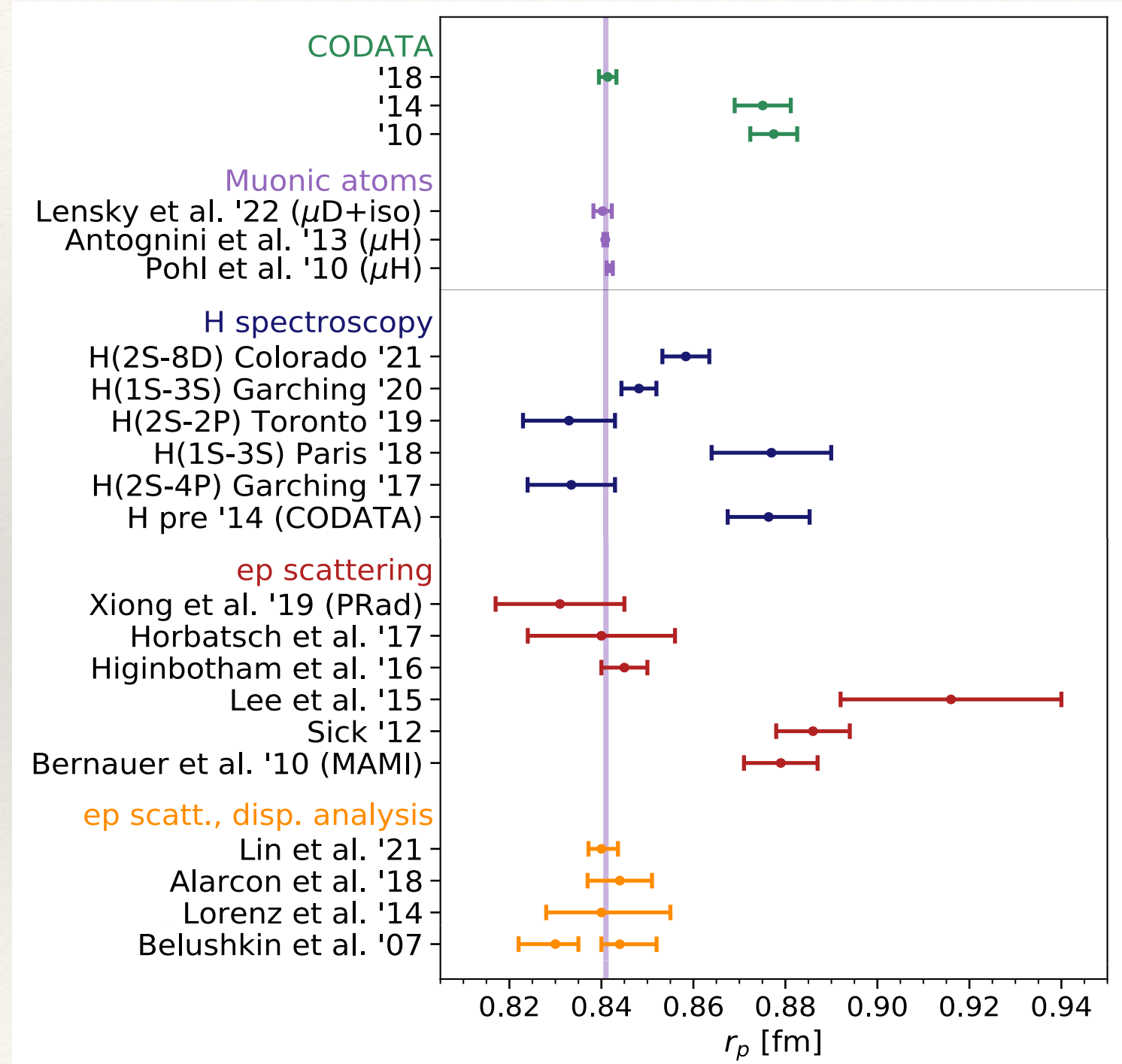
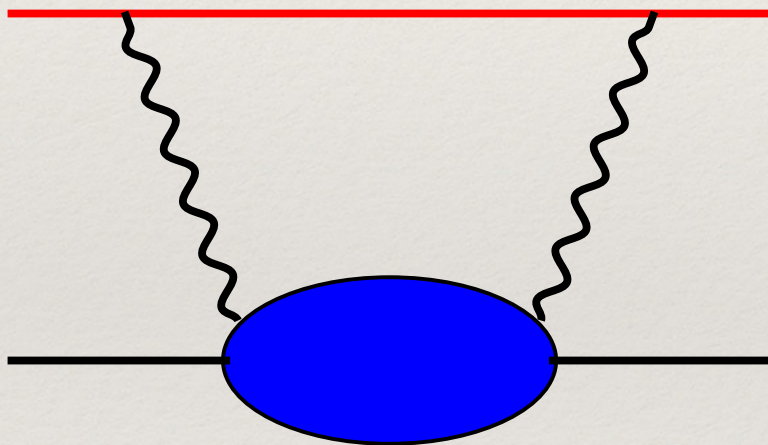
$$r^2(^3\text{He}) - r^2(^4\text{He})$$



Li Muli, TRR, Bacca PRL 134

Radiative Corrections in Nuclear Physics

Lepton Scattering

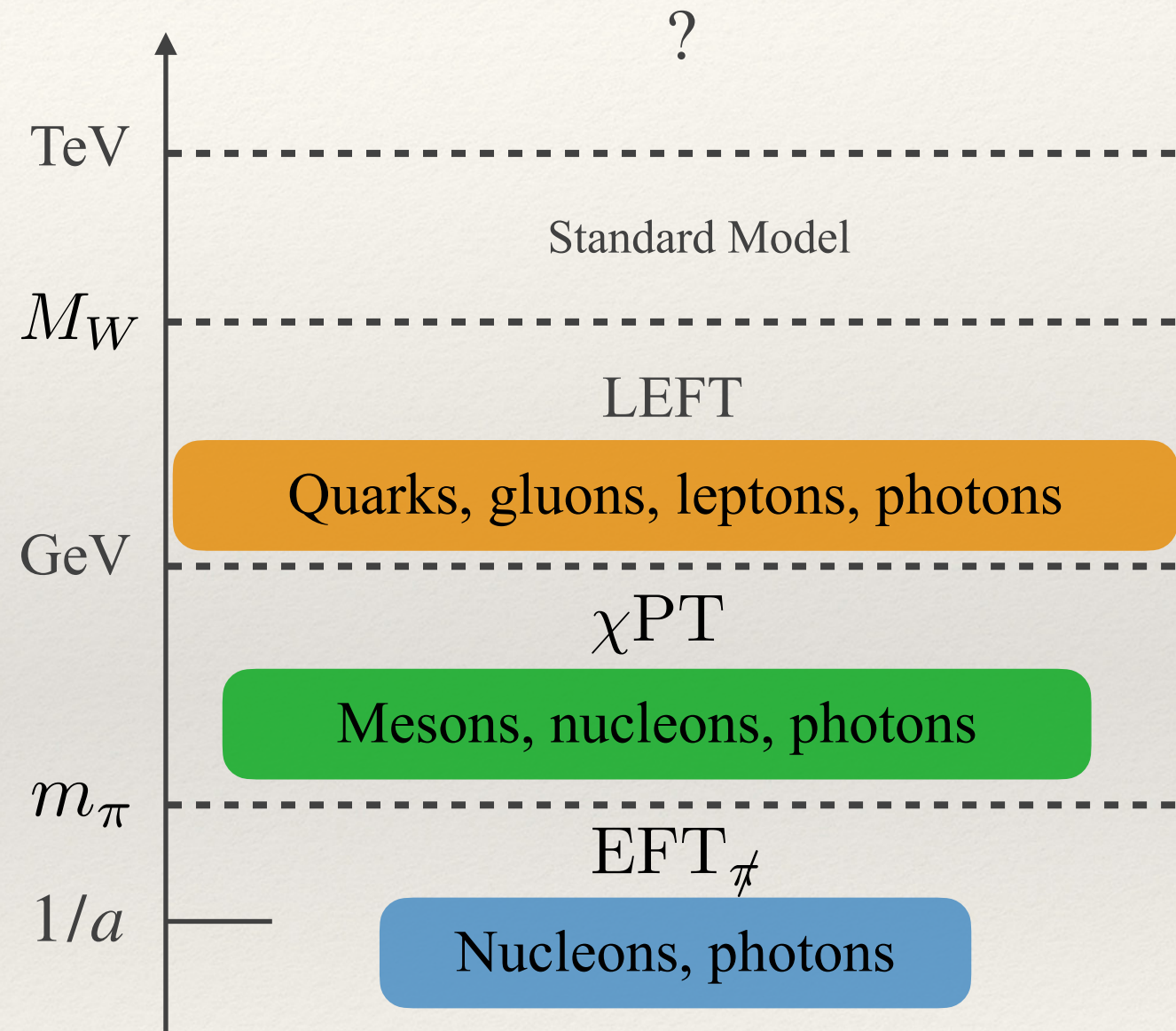


Effective Field Theory Tower

Goal: Low energy, high precision, universal

Effective field theories should...

1. Have all particles that are near the mass-shell
2. Have a homogeneous power counting governed by a single ratio of scales
3. Be renormalization group invariance up to the order we are working
4. Preserve relevant symmetries (gauge, discrete, internal)
5. Make life easier



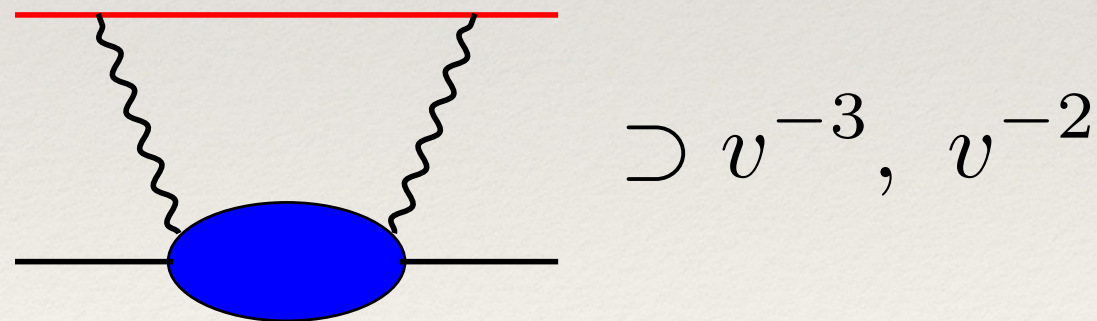
Nonrelativistic EFT with Virtual Photons

- ❖ Energy and momenta are different but correlated scales

$$v \sim p/M = \sqrt{E/M}$$

potential region : $(E, \vec{p}) \sim (Mv^2, Mv)$

- ❖ But our prototype diagram doesn't scale with a single power of v ?



Nucleons vs. $t\bar{t}$

NRQCD

NN EFT

Heavy potential
quarks

Heavy potential
nucleons

Massless gluons

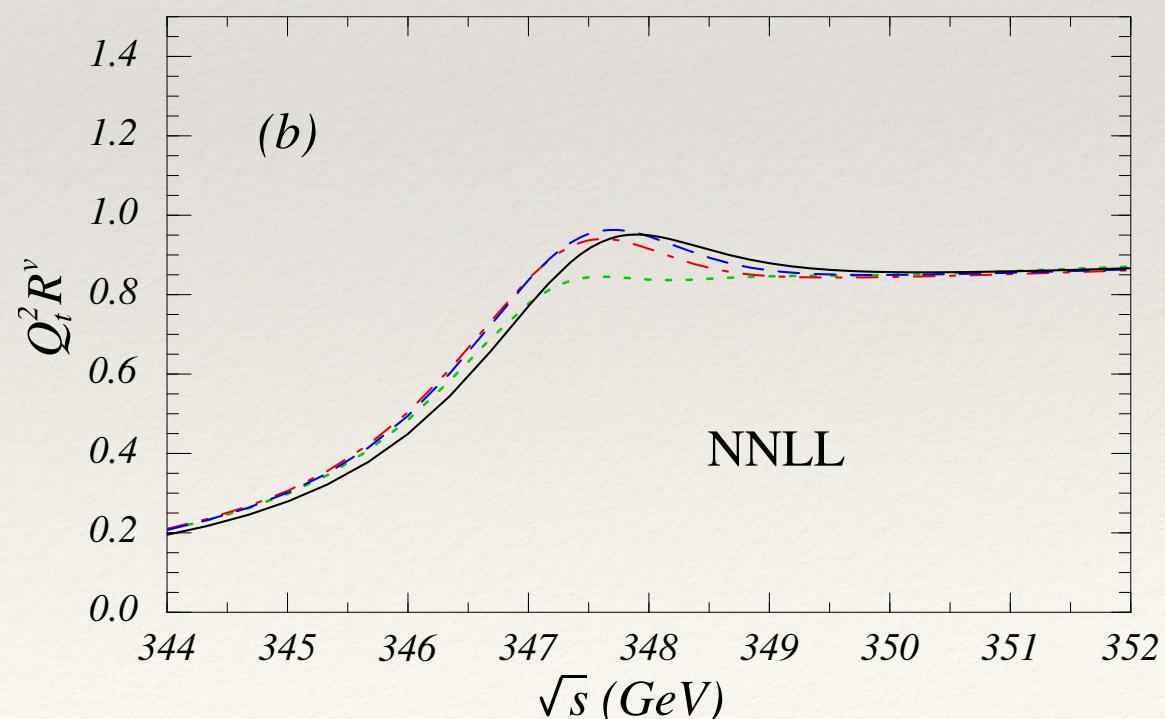
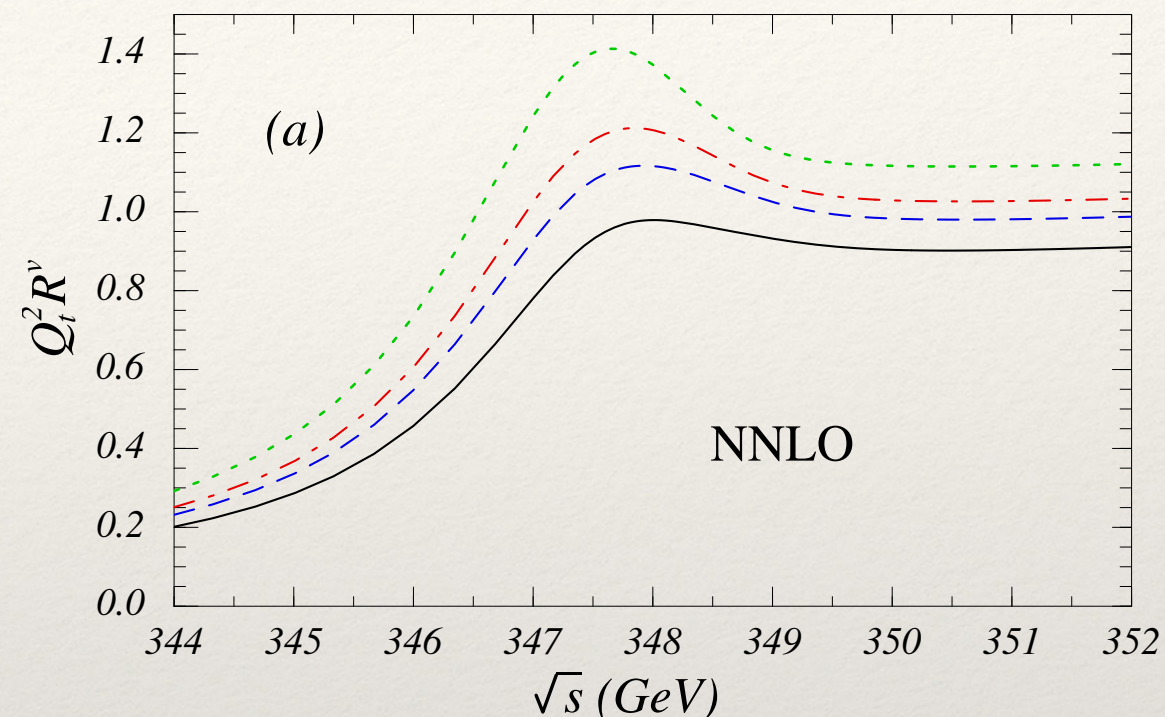
Massless photons
(light pions)

$$p \sim m\alpha$$

$$p \sim 1/a$$

$$E \sim m\alpha^2$$

$$E \sim 1/M_N a^2$$



Nonrelativistic EFT with Virtual Photons

- ❖ Energy and momenta are different but correlated scales

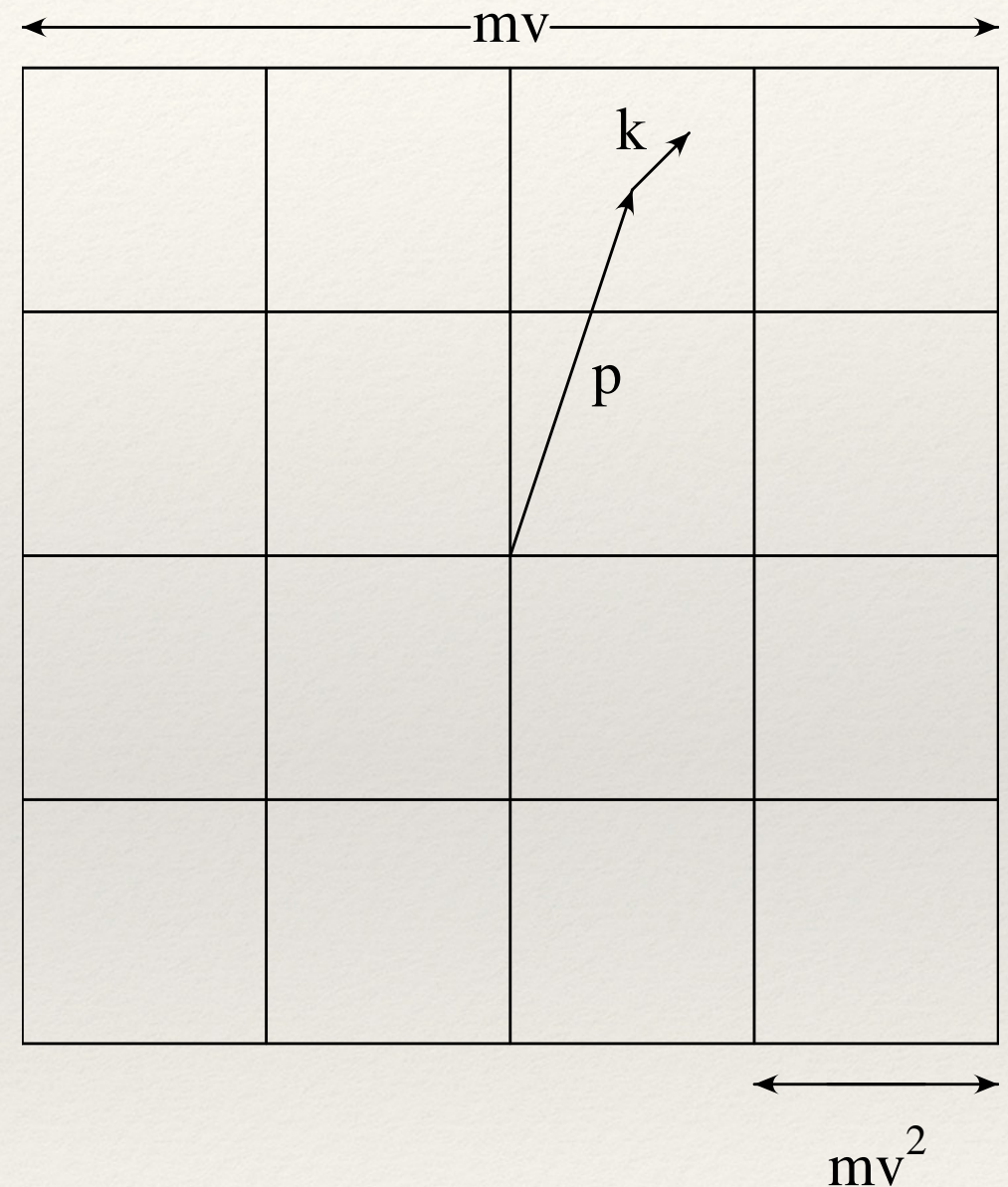
$$v \sim p/M = \sqrt{E/M}$$

- ❖ Homogeneous power counting requires mode separation

$$\text{potential} \sim (Mv^2, Mv) : N_{\mathbf{p}}$$

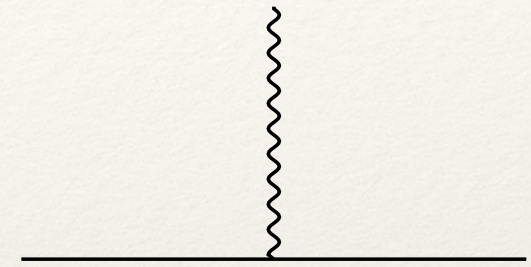
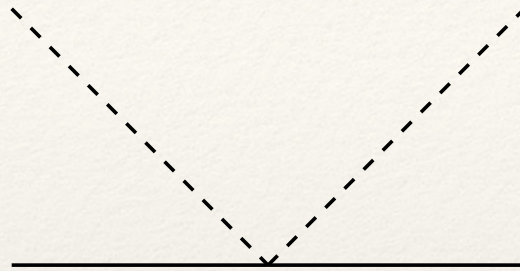
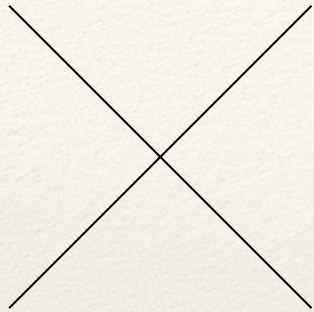
$$\text{soft} \sim (Mv, Mv) : A_p$$

$$\text{ultrasoft} \sim (Mv^2, Mv^2) : A$$



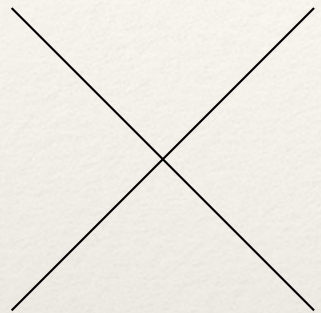
Caswell, Lepage, Labelle, Luke, Manohar,
Rothstein, Savage, Grißhammer, Grinstein,
Stewart, Hoang, Pineda, Soto, Brambilla, Vairo

Velocity EFT Lagrangian



$$\begin{aligned}
 \mathcal{L} = & \sum_p N_p^\dagger \left(iD_0 - \frac{(p - iD)^2}{2M_N} \right) N_p - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\
 & + \sum_p |p^\mu A_p^\nu - p^\nu A_p^\mu|^2 - \sum_{p', p} V(p', p) \\
 & - \frac{4\pi\alpha}{2M_N} \sum_{q, q', p, p'} A_{q'} \cdot A_q N_{p'}^\dagger Q N_p \\
 & + \frac{e}{2M_N} \epsilon^{ijk} (\nabla^j A^k) \sum_p N_p^\dagger \sigma^i [\kappa_0 + \kappa_1 \tau^3] N_p ,
 \end{aligned}$$

Neutron-Proton Potential



$$V_{pn} = \sum_{v=-1} \sum_{p',p} V_{abcd}^{(v)}(p',p) p_{p',a}^\dagger p_{p,b} n_{-p',c}^\dagger n_{-p,d}$$

$$V_{abcd}^{(-1)} = C_{0,pn}^{(S=1)} P_{ab,cd}^{(1)} + C_{0,pn}^{(S=0)} P_{ab,cd}^{(0)}$$

$$V_{abcd}^{(0)} = \frac{1}{2} (p'^2 + p^2) \left[C_{2,pn}^{(S=1)} P_{ab,cd}^{(1)} + C_{2,pn}^{(S=0)} P_{ab,cd}^{(0)} \right]$$

$$V_{abcd}^{(1)} = \frac{1}{4} (p'^2 + p^2)^2 \left[C_{4,pn}^{(S=1)} P_{ab,cd}^{(1)} + C_{4,pn}^{(S=0)} P_{ab,cd}^{(0)} \right]$$

Proton-proton potential

$$V_{pp} = \sum_{v=-2} \sum_{\mathbf{p}', \mathbf{p}} V_{abcd}^{(v)}(\mathbf{p}', \mathbf{p}) p_{\mathbf{p}', a}^\dagger p_{-\mathbf{p}', c}^\dagger p_{-\mathbf{p}d} p_{\mathbf{p}, b}$$

$$V_{abcd}^{(-2)}(\mathbf{p}', \mathbf{p}) = \frac{4\pi\alpha}{(\mathbf{p}' - \mathbf{p})^2} \delta_{ab} \delta_{cd}$$

$$V_{abcd}^{(0)}(\mathbf{p}', \mathbf{p}) = \frac{1}{4} C_0^{(pp)} \sigma_{ac}^2 \sigma_{db}^2$$

“Matching” to AV18

- ❖ The LECs need to be matched at $\mu = m_\pi$ at $O(\alpha^0)$

$$C_0^{(s)}(\mu_{\text{match}} = m_\pi) = \frac{4\pi a_s}{M_N}$$

$$C_2^{(s)}(\mu_{\text{match}} = m_\pi) = \frac{2\pi a_s^2 r_s}{M_N} \qquad C_4^{(s)}(\mu_{\text{match}} = m_\pi) = \frac{\pi a_s^3 r_s^2}{M_N}$$

- ❖ Use Argonne v18 with $\alpha = 0$ as proxy for lattice QCD—“integrate out the pion”

$$a_0 = -23.084 \text{ fm}$$

$$r_0 = 2.703 \text{ fm}$$

$$a_{pp} = -17.164 \text{ fm}$$

$$a_1 = 5.402 \text{ fm}$$

$$r_1 = 1.752 \text{ fm}$$

$$r_{pp} = 2.865 \text{ fm}$$

- ❖ Next: Run μ from $m_\pi \rightarrow p$ at $O(\alpha)$ and calculate matrix elements

Summing Logs in Bound States

1. Obtain anomalous dimensions for the couplings in perturbation theory
2. Integrate the RG equations from matching scale to typical bound state velocity
3. Calculate bound state matrix elements

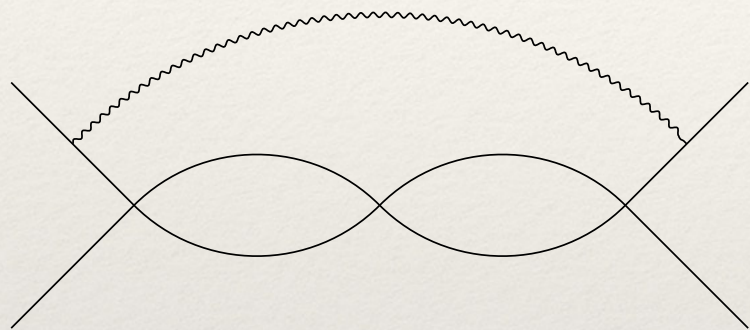
$$L \sim \log(m_L/m_H)$$

LO	1		
NLO	αL	α	
NNLO	$\alpha^2 L^2$	$\alpha^2 L$	α^2
...	...	...	...
$N^k\text{LO}$	$\alpha^k L^k$	$\alpha^k L^{k-1}$	$\alpha^k L^{k-2}$

Cohen arXiv:1903.03622

Loop Corrections and the Velocity Renormalization Group

- ❖ Perturbation theory generates two (possibly) large logarithms



$$p^2 C_0^3 \left[\lambda_S \log \frac{\mu}{p} + \lambda_{US} \log \frac{\mu}{E} \right]$$

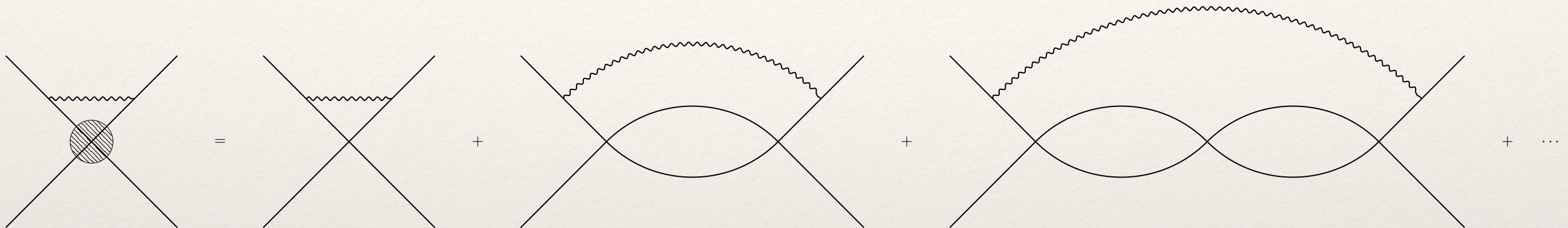
- ❖ Introduce two correlated scales in dimensional regularization

$$\mu_S = M_N \nu \qquad \mu_U = M_N \nu^2$$

- ❖ Run in ν —sum soft and ultrasoft logarithms simultaneously

$$\nu \frac{dV}{d\nu} = \gamma_S + 2\gamma_U \qquad \mu_U \frac{dV}{d\mu_U} = \gamma_U \qquad \mu_S \frac{dV}{d\mu_S} = \gamma_S$$

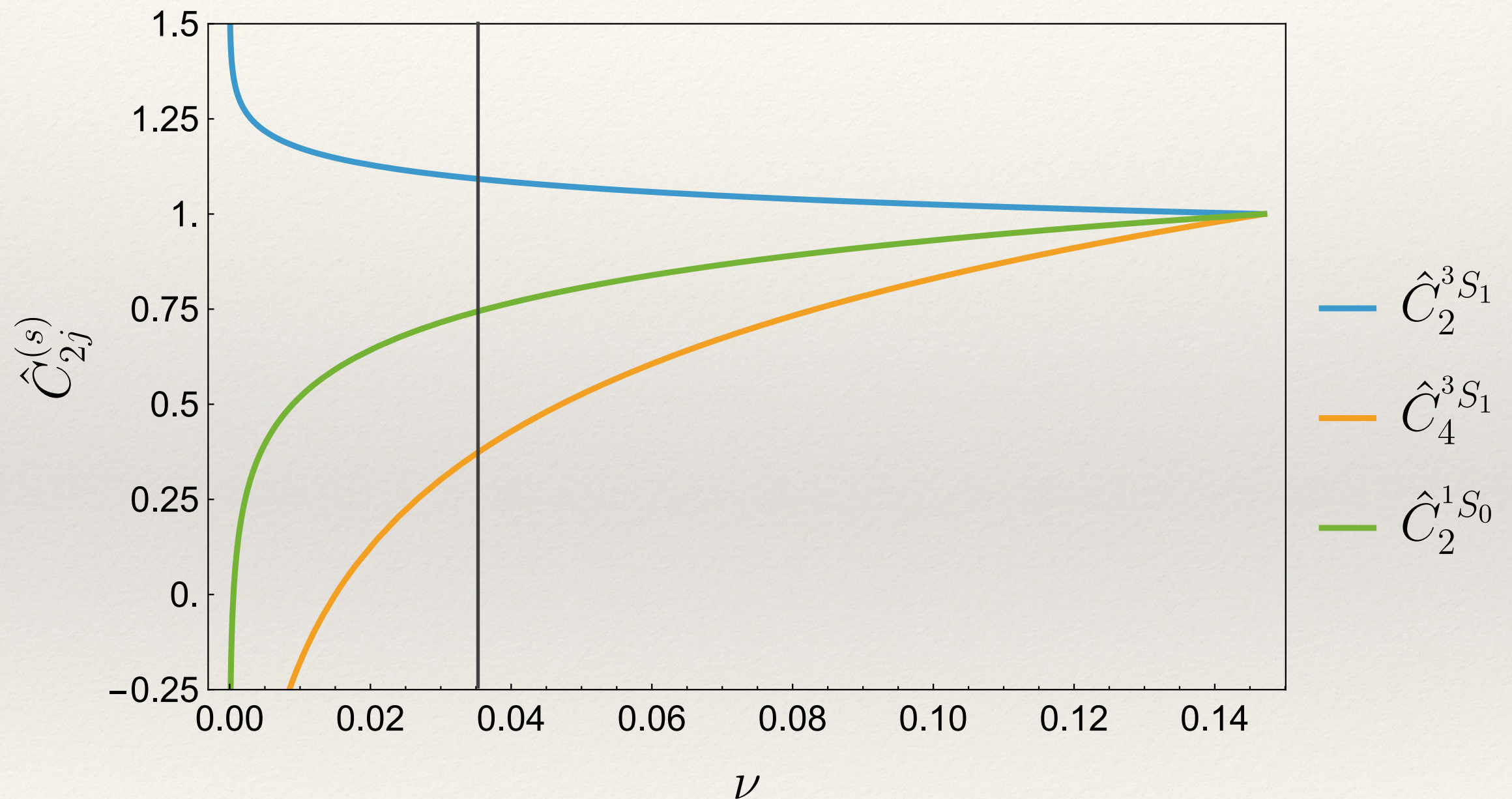
Velocity RG Solution



$$C_2(\nu) = C_2 \left(\frac{m_\pi}{M_N} \right) - \frac{15}{4} \left(\frac{M_N}{4\pi} \right)^2 C_0^3 \log \left(\frac{\alpha(M_N \nu^2)}{\alpha(m_\pi^2/M_N)} \right) ,$$

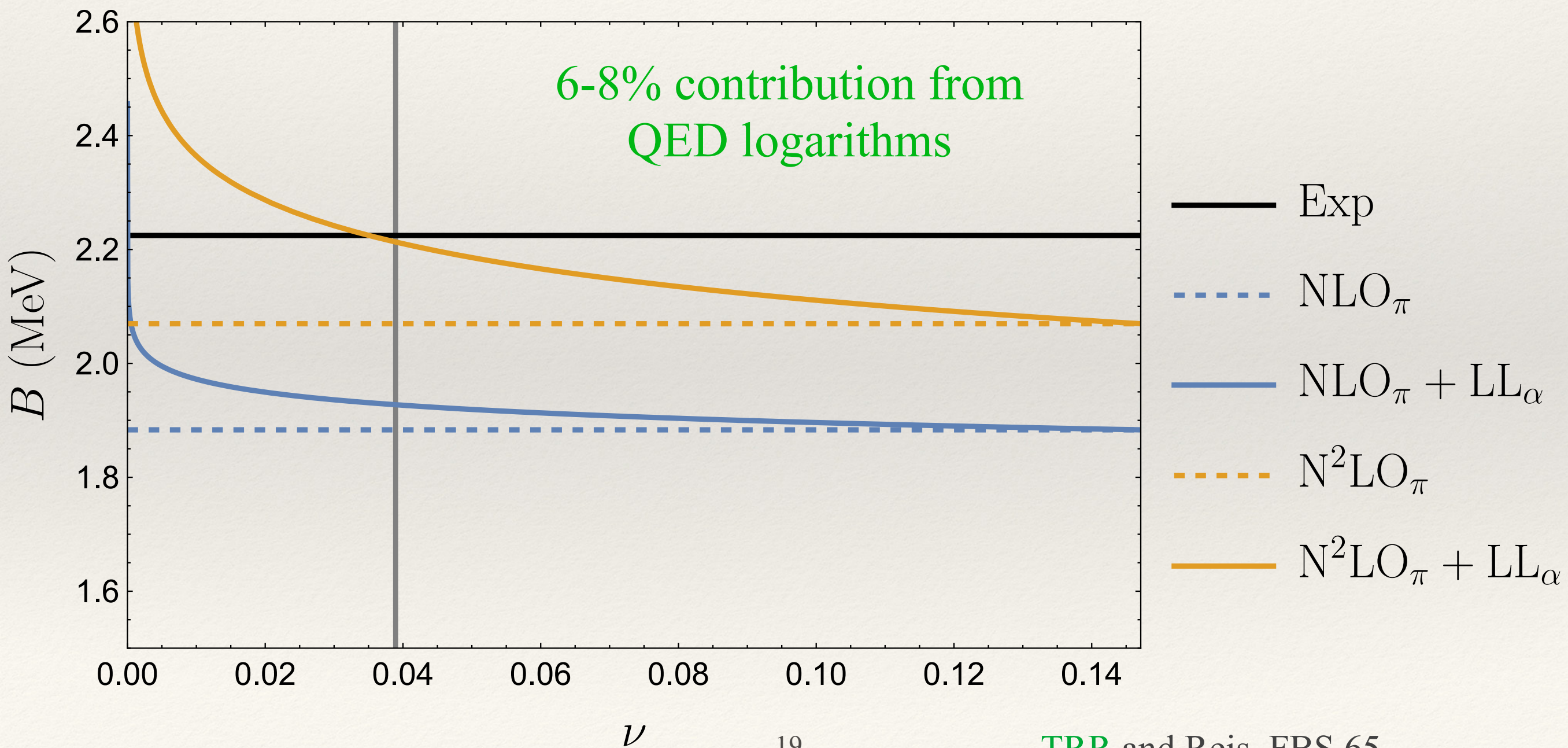
$$C_4(\nu) = C_4 \left(\frac{m_\pi}{M_N} \right) + 4 \left(\frac{M_N}{4\pi} \right)^4 C_0^5 \log \left(\frac{\alpha(M_N \nu^2)}{\alpha(m_\pi^2/M_N)} \right)$$

Neutron-proton potential



Deuteron Binding Energy

$$B = \frac{1}{M_N} \left(\frac{4\pi}{M_N C_0} \right)^2 + \frac{1}{2\pi} C_2 \left(\frac{4\pi}{M_N C_0} \right)^5 + \frac{7}{16\pi^2} M_N C_2^2 \left(\frac{4\pi}{M_N C_0} \right)^8 - \frac{1}{2\pi} C_4 \left(\frac{4\pi}{M_N C_0} \right)^7$$



Electroweak Matrix Elements: A Case Study in Proton-Proton Fusion

- ❖ Typical proton momentum in the sun $p \sim 2.5$ MeV, $E_e \sim 1$ MeV

$$p < a_{pp}^{-1} \ll \sqrt{2M_r B_d} \ll m_\pi$$

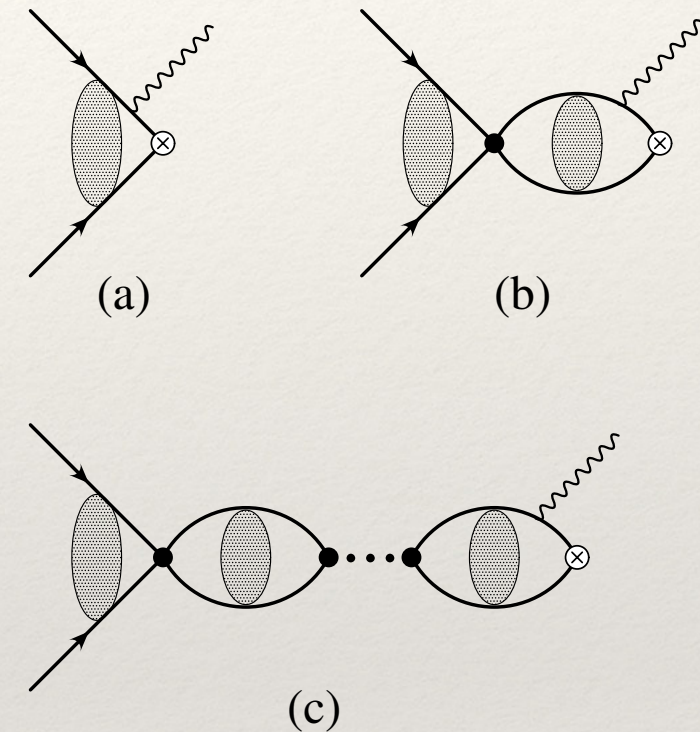
- ❖ Coulomb interaction is nonperturbative for $v \sim \alpha$

$$A(p) \sim \int_{\vec{k}} \frac{1}{\mathbf{k}^2 + 2M_r B_d} \psi_{\mathbf{p}}(\mathbf{k}) \quad B(p) \sim \int_{\mathbf{k}_1, \mathbf{k}_2} \frac{\langle \mathbf{k}_1 | G_C | \mathbf{k}_2 \rangle}{\mathbf{k}_1^2 + 2M_r B_d}$$

$$T \sim A + \frac{C_0^{(pp)}}{1 - C_0^{(pp)} J_0(p)} \lambda(p) B(p)$$

$$J_0(p) = i \frac{M_p p}{4\pi} + \frac{M_p^2 \alpha}{4\pi} \log \left(\frac{\mu^2}{-4p^2} \right) - \frac{M_p^2 \alpha}{4\pi} \left[\psi \left(i \frac{M_p \alpha}{2p} \right) + \gamma_E \right]$$

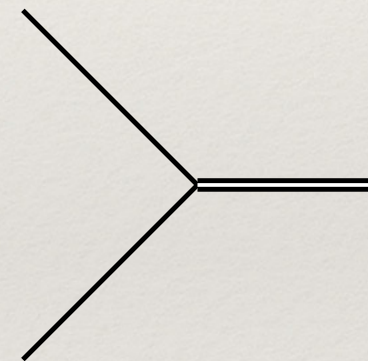
- ❖ Want to recover from an EFT perspective—expand in $p/\sqrt{2M_r B_d}$ and resum what has to be resummed



Matching

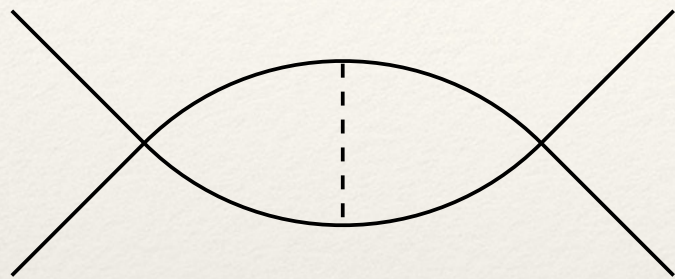
- ❖ Treat deuteron like a fundamental particle
- ❖ Wilson coefficient given by hard part ($k_0 \sim \gamma_d, \mathbf{k} \sim \gamma_d$) of pionless result—perturbative in α and to all orders in $C_0^{(pp)}$

$$\mathcal{L} \supset \frac{1}{2} g_{\text{eff}} L_i d^{i\dagger} (p \sigma^2 p)$$



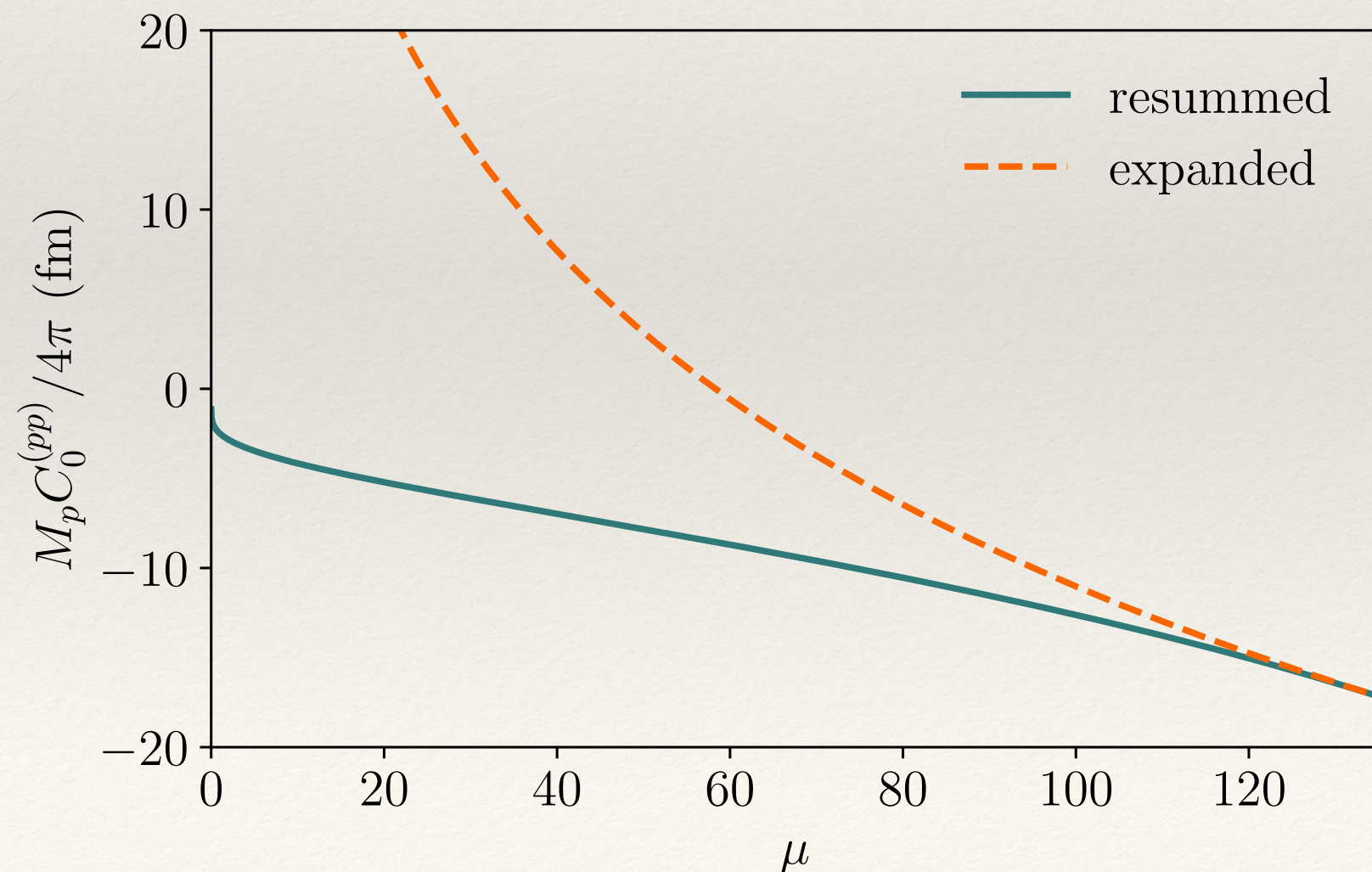
$$g_{\text{eff}}(\mu = \gamma_d) = \sqrt{\frac{8\pi}{\gamma_d^3}} \left(1 - \frac{M_p C_0^{(pp)}(\gamma_d)}{4\pi} \gamma_d + \frac{M_p \alpha(\gamma_d)}{\gamma_d} \right)$$

$C_0^{(pp)}$ Anomalous Dimension

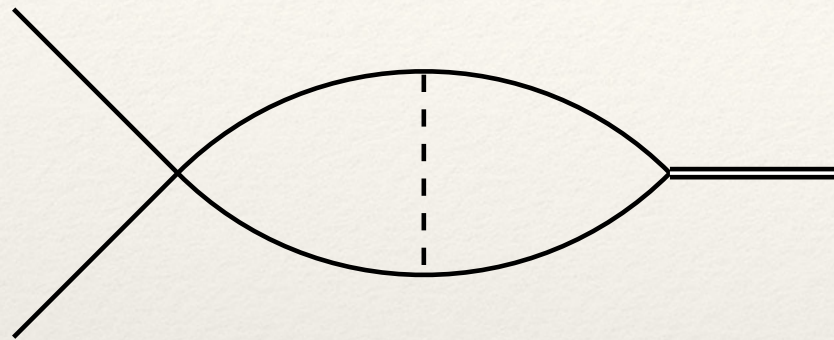


$$\beta(C_0^{(pp)}) = -\frac{M_p^2 \alpha}{2\pi} C_0^{(pp)2}$$

$$C_0^{(pp)} = \frac{C_0^{(pp)}(m_\pi)}{1 + \frac{3M_p^2}{4} C_0^{(pp)}(m_\pi) \log\left(\frac{\alpha(\mu)}{\alpha(m_\pi)}\right)}$$

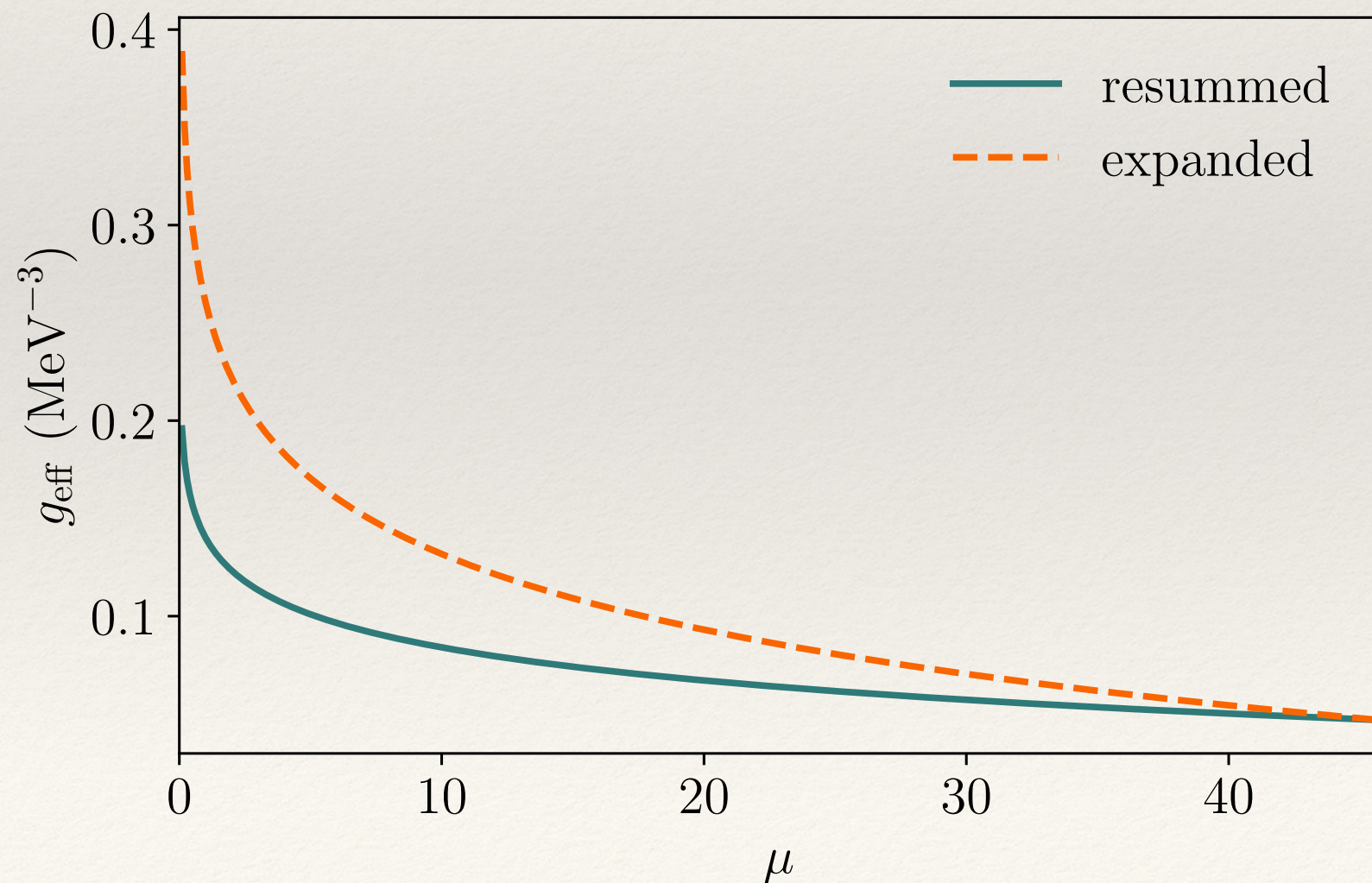


g_{eff} Anomalous Dimension

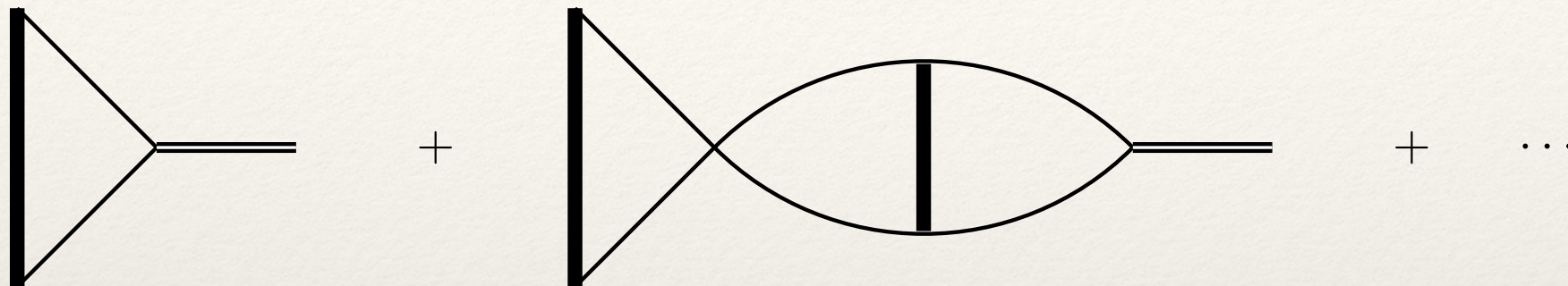


$$\beta(g_{\text{eff}}) = \frac{M_p^2 \alpha}{2\pi} C_0^{(pp)} g_{\text{eff}}$$

$$g_{\text{eff}}(\mu) = g(\gamma_d) \left[1 + \frac{3M_p^2}{4} C_0^{(pp)} \log \frac{\alpha(\mu)}{\alpha(\gamma_d)} \right]$$



EFT Amplitude

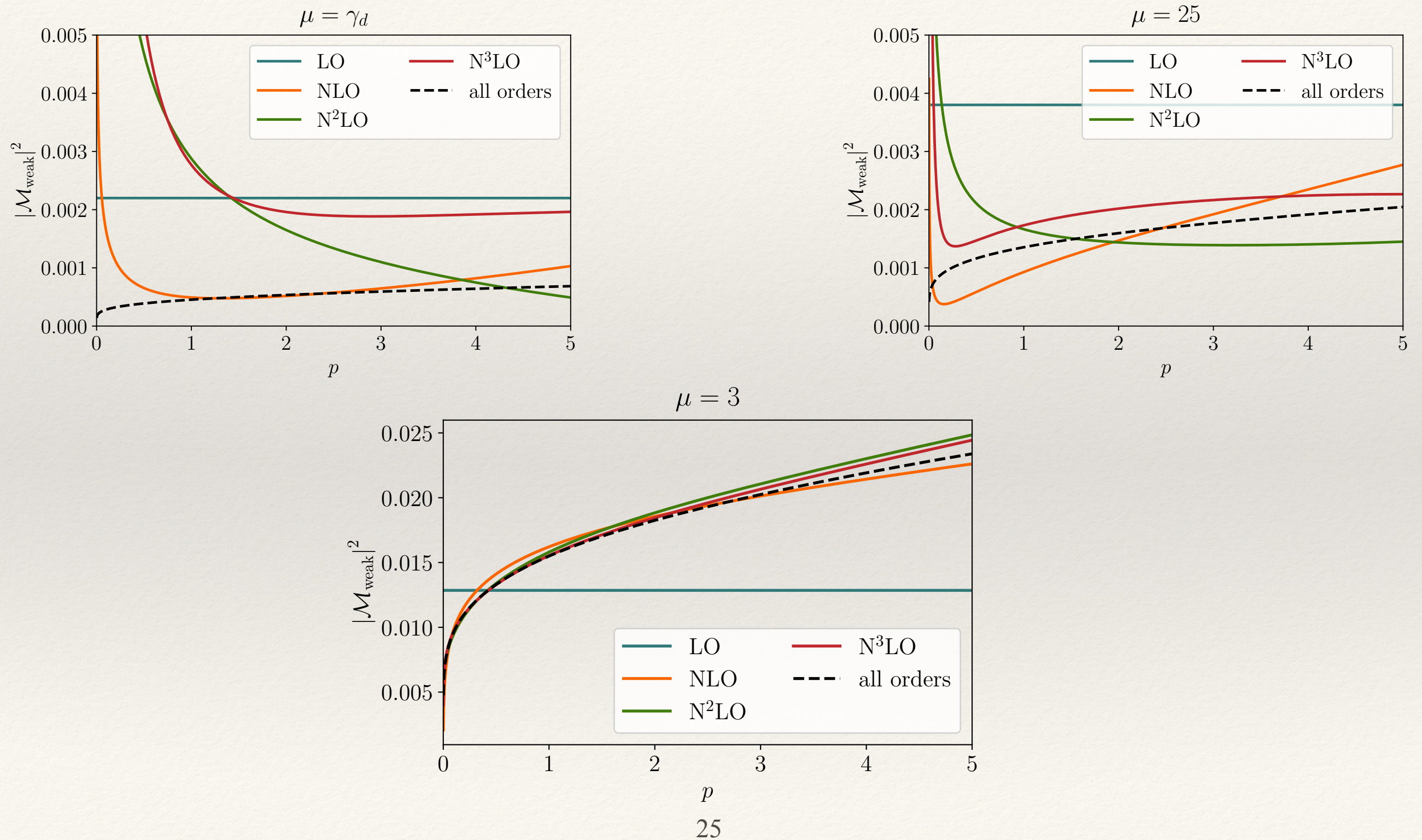


$$\begin{aligned} T &= g_{\text{eff}} \lambda(p) \left[1 + C_0 J_0 + (C_0 J_0)^2 + \dots \right] \\ &= g_{\text{eff}} \lambda(p) \frac{1}{1 - C_0 J_0} \end{aligned}$$

❖ Expected from EFT principles

$$\langle \mathcal{O}_{\text{full}} \rangle = C \langle \mathcal{O}_{\text{EFT}} \rangle$$

RG Improving the Amplitude



Uncertainty Estimate

- ❖ Matching $< 0.5 \%$

δ_{NS} goes in the matching

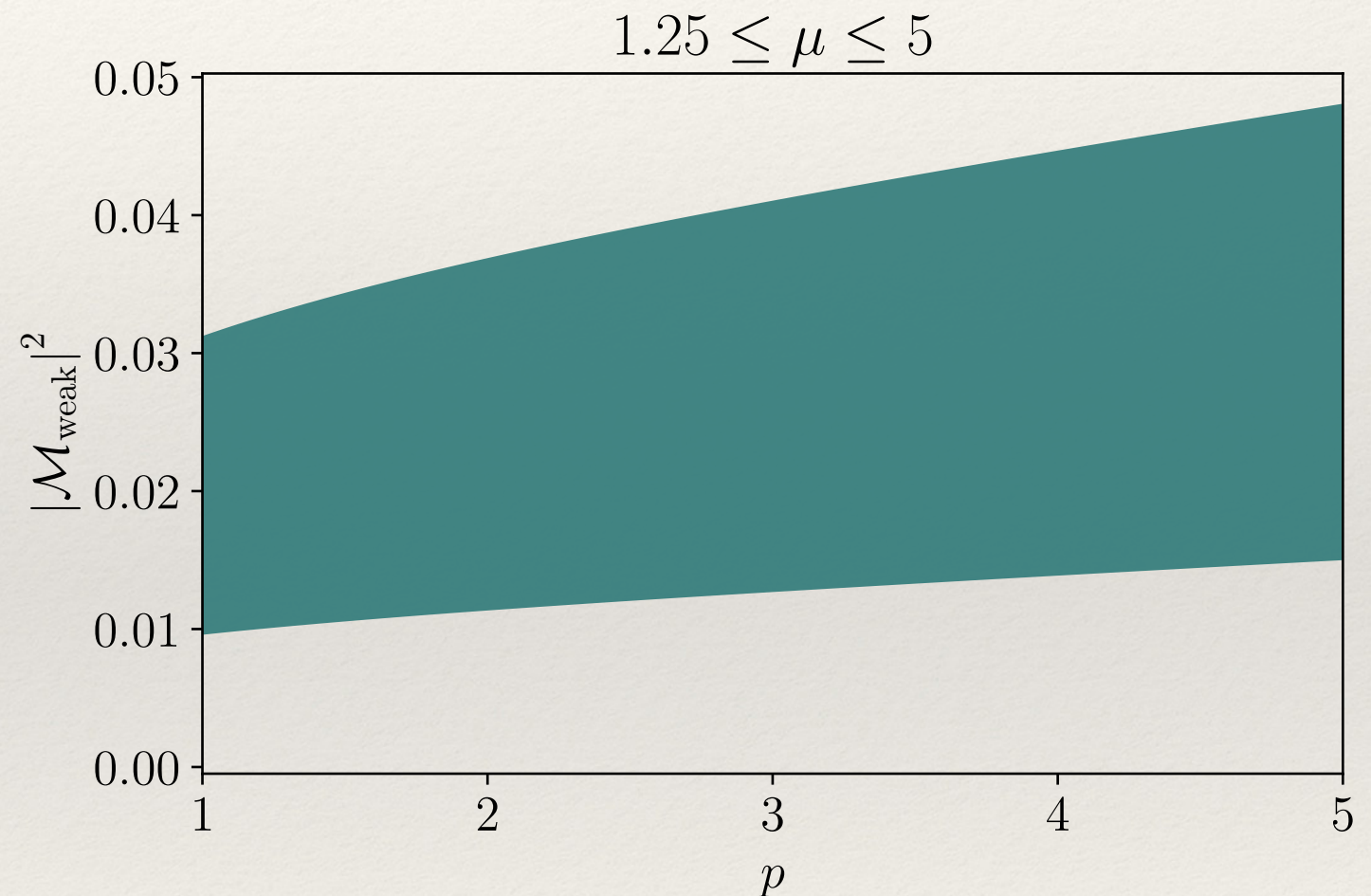
Combes et al. PRC 110

- ❖ Ultrasoft $< 0.1 \%$

- ❖ Soft $\implies$ Fermi function and other 0.5%

Hill and Plestid PRL 133, PRD 109

- ❖ Scale variation $\sim 30 \%$



Summary

- ❖ Nucleons + photons $\approx$ heavy quarkonium + gluons
- ❖ LL_α improved two nucleon interaction—few percent effect from QED logs in deuteron binding
- ❖ 30% uncertainty in pp fusion—improve with NLL_α ? Power suppressed potentials?
- ❖ Embed correlated running in the superallowed framework—are intranuclear corrections understood?
Cirigliano et al. PRL 133
- ❖ Revisit corrections in tritium—KATRIN, Project 8