

Quantum “Dark” Matter

Surjeet Rajendran

Classical and Quantum Particle Mechanics



$$\vec{F} = m\vec{a}$$

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$$i \frac{\partial |\Psi\rangle}{\partial t} = H |\Psi\rangle$$

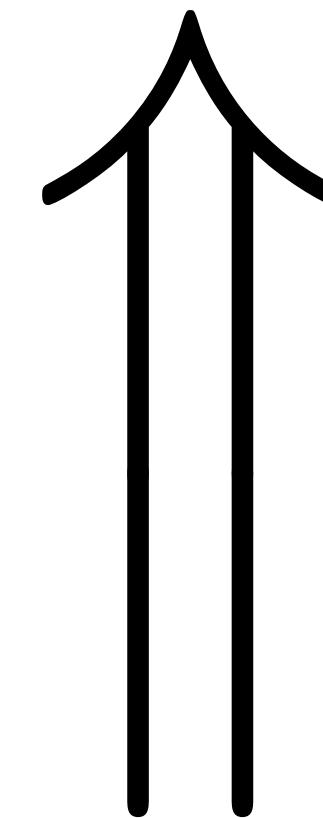


Classical and Quantum Particle Mechanics



$$\vec{F} = m\vec{a}$$

$$\frac{d\langle \Psi | \hat{p} | \Psi \rangle}{dt} = -\langle \Psi | \frac{dV}{dx} | \Psi \rangle$$



Ehrenfest's Theorem



$$i\frac{\partial|\Psi\rangle}{\partial t} = H|\Psi\rangle$$



Quantum Supremacy

Classical and Quantum Field Theory

Classical Field Theory
(e.g. Klein Gordon)

$$\frac{\partial S}{\partial \phi} = 0$$

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Schrodinger Picture of Quantum Field Theory

$|\chi(t)\rangle$

Quantum State of Fields
(e.g. in Fock states)

$\phi(x)$

**Time Independent
Operators**

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(e.g. in Fock states)

$\phi(x)$ **Time Independent Operators**

$$H = \int d^3x \mathcal{H}(\phi(x), \pi(x))$$

Time Evolution

$$i \frac{\partial |\chi(t)\rangle}{\partial t} = H |\chi(t)\rangle$$

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Time Evolution

Analog of Ehrenfest Theorem?

$$i \frac{\partial |\chi(t)\rangle}{\partial t} = H |\chi(t)\rangle$$

Schwinger Dyson Equations
Quantum Supremacy

Schwinger Dyson

$$Z = \int D\phi e^{iS[\phi]}$$

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Change Variables: $\phi \rightarrow \phi + \delta\phi \implies \delta Z = 0$

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Gauge theories?

Classical and Quantum Gauge Theories

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(e.g. General Relativity)**

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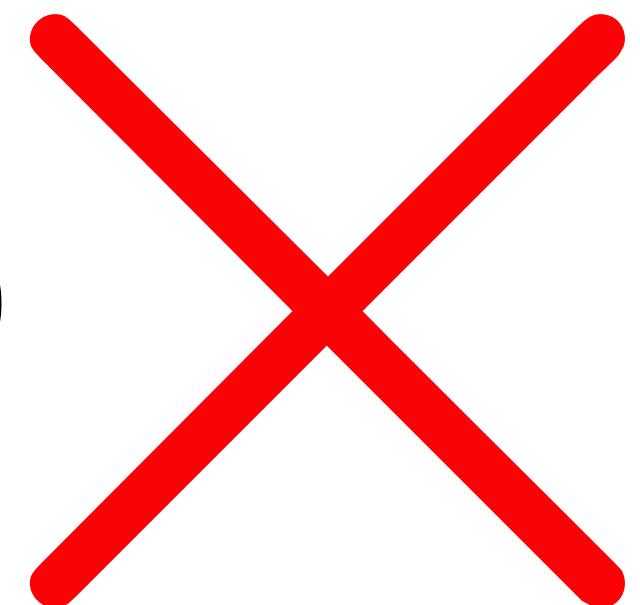
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Need to Gauge Fix to define Path Integral

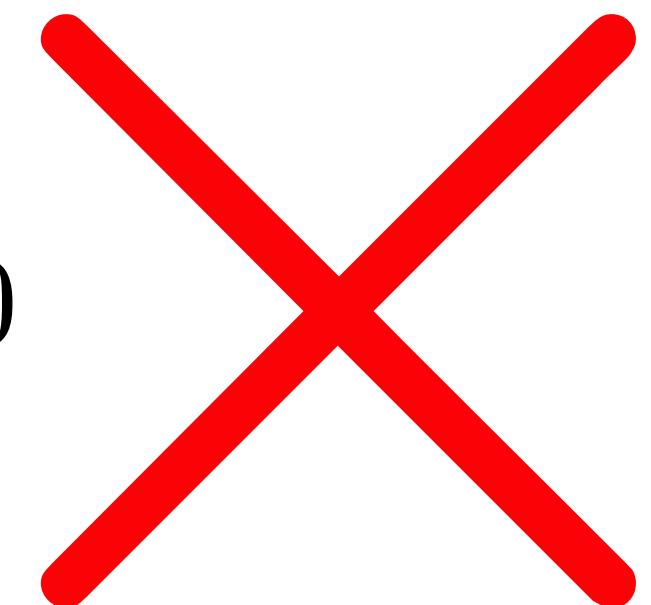
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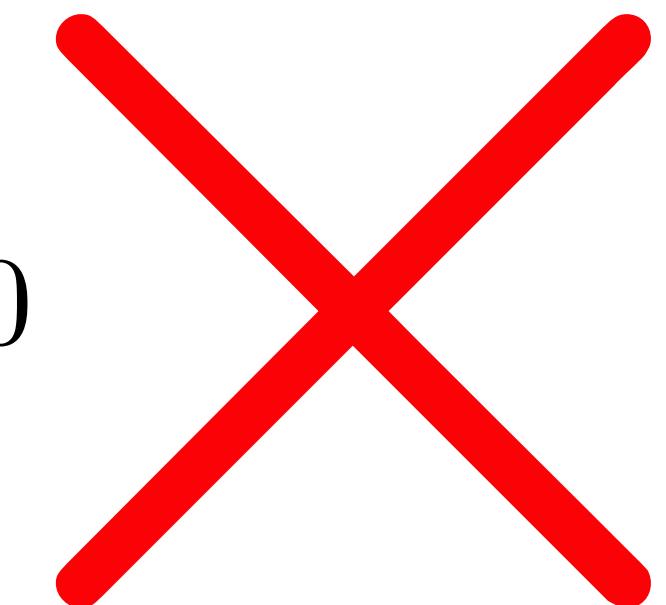
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Quantum Field Theory

Not the Same
????????

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Outline

1. Classical Mini Superspace Cosmology

2. Path Integral

3. Hamiltonian

4. Wheeler deWitt

5. Conclusions

Mini Superspace Cosmology

Homogeneous Universe

Study homogeneous space-times in General Relativity

$$g_{\mu\nu} = -N(t)^2 dt^2 + a(t)^2 (dx^2 + dy^2 + dz^2)$$

Couple this to homogenous sources of matter, for e.g. rolling scalar field ϕ

$$S = \int dt \mathcal{L} \left(a(t), \dot{a}(t), N(t), \phi(t), \dot{\phi}(t) \right)$$

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Classical Equations?

2 FRW + 1 Scalar Field

Classical Equations

$$S = \int dt \mathcal{L} \left(a(t), \dot{a}(t), N(t), \phi(t), \dot{\phi}(t) \right)$$

$$\frac{\partial S}{\partial a} = 0 \implies \partial_t \left(\frac{\partial \mathcal{L}}{\partial \dot{a}} \right) - \frac{\partial \mathcal{L}}{\partial a} = 0 \text{ i.e. } \left(\frac{\ddot{a}}{a} + \dots = 0 \right)$$

2nd FRW Equation

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Klein Gordon

$$\frac{\partial S}{\partial N} = 0 \implies \frac{\partial \mathcal{L}}{\partial N} = 0 \text{ i.e. } \left(\left(\frac{\dot{a}}{a} \right)^2 + \dots = 0 \right)$$

**1st FRW Equation
(Hubble)**

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Classical Solutions are over-constrained (3 equations for 2 variables)

Solve: Klein Gordon + 2nd FRW with boundary condition from 1st FRW

Path Integral

Path Integral

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To get finite path integral, need to gauge fix $N(t)$

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$\dot{N} = 0$, enforce via Lagrange multiplier

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$$\dot{N} = 0 \implies N(t_2) = N_0, N(t_1) = N_0$$

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Can show that path integral over $a(t)$, $\phi(t)$ are finite

$$\langle \phi_f, a_f | T(t_2; t_1) | \phi_i, a_i \rangle = e^{i P(\phi_f, a_f, t_2; \phi_i, a_i, t_1; N_0)}$$

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Path Integral explicitly depends upon random gauge choice No

Path integral obviously not gauge invariant

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Is this a problem?

Gauge Invariant Physics

$$\langle \phi_f, a_f | T(t_2; t_1) | \phi_i, a_i \rangle$$

Path integral defines the time evolution operator

$$g_{\mu\nu} = -N(t)^2 dt^2 + a(t)^2 (dx^2 + dy^2 + dz^2)$$

Time coordinate is a gauge choice in General Relativity

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Perfectly Reasonable for the time evolution operator to be gauge dependent

We need gauge invariant physics. What does this mean?

Gauge Invariant Physics

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What happens if $N_0 \rightarrow \tilde{N}_0$?

Gauge Invariant Physics

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Can Explicitly Show: $\langle \phi_f, a_f | T(t_2; t_1) | \phi_i, a_i \rangle = \langle \phi_f, a_f | T(\tilde{t}_2; \tilde{t}_1) | \phi_i, a_i \rangle$

$$\text{where } \tilde{t} = \frac{N_0}{\tilde{N}_0} t$$

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$$|\phi_i, a_i\rangle \rightarrow \sum_{\phi_f, a_f} c_{\phi_f, a_f}(t_2, t_1) |\phi_f, a_f\rangle = \sum_{\phi_f, a_f} c_{\phi_f, a_f}(\tilde{t}_2, \tilde{t}_1) |\phi_f, a_f\rangle$$

$$|\Psi(t_1)\rangle = |\tilde{\Psi}(\tilde{t}_1)\rangle = |\Sigma\rangle \rightarrow |\Psi(t_2)\rangle = |\tilde{\Psi}(\tilde{t}_2)\rangle = |\Omega\rangle$$

Gauge Invariant Physics

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Fix manifold (\mathbf{R}^1).

Pick any choice of time co-ordinate on the manifold

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General Gauge: $N(t) = f(t)$

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Where $N_0 dt = f(\tilde{t}) d\tilde{t}$

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Fix manifold (\mathbf{R}^4).

Pick any choice of time co-ordinate on the manifold

General Gauge: $N(t) = f(t)$

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Moral: Gauge invariant physics from gauge dependent path integral
Reasonable since time is gauge choice!

Schwinger Dyson Equations

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$$\langle \Psi | \partial_t \left(\frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) - \frac{\partial \mathcal{L}}{\partial \phi} | \Psi \rangle = 0 \quad \text{Klein Gordon}$$

Heisenberg Picture

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$$\langle \Psi | \frac{\partial \lambda}{\partial t} - \frac{\partial \mathcal{L}}{\partial N} | \Psi \rangle = 0 \quad \text{Tells you how } \lambda \text{ evolves in the path integral - not 1st FRW Equation????}$$

Heisenberg Picture

Hamiltonian

Hamiltonian Construction

$$g_{\mu\nu} = -N(t)^2 dt^2 + a(t)^2 (dx^2 + dy^2 + dz^2)$$

Physical Degrees of freedom: $a(t), \phi(t)$
Gauge Freedom: $N(t)$

$$S = \int dt \mathcal{L} \left(a(t), \dot{a}(t), N(t), \phi(t), \dot{\phi}(t) \right)$$

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Quantize These

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Quantize These

$$\Pi_N = \frac{\partial \mathcal{L}}{\partial \dot{N}} = 0$$

What to do with N ?

Hamiltonian Construction

$$\Pi_a = \frac{\partial \mathcal{L}}{\partial \dot{a}}, \Pi_\phi = \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \quad \text{Quantize These}$$

$$[a, \Pi_a] = [\phi, \Pi_\phi] = i$$

$$a = A + A^\dagger, \Pi_a = i(A - A^\dagger) \quad \phi = B + B^\dagger, \Pi_\phi = i(B - B^\dagger)$$

Hamiltonian Construction

$$\Pi_a = \frac{\partial \mathcal{L}}{\partial \dot{a}}, \Pi_\phi = \frac{\partial \mathcal{L}}{\partial \dot{\phi}}$$

Quantize These

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Fully non-linear General Relativity - no “free” theory with “free” kinetic term

But, can still construct Hilbert space with Fock states of A, B - these are operator level statements independent of kinetic terms of the theory

$$A|0\rangle = 0, A^\dagger|0\rangle = |1\rangle \text{ etc.}$$

Hamiltonian Construction

$$S = \int dt \mathcal{L} \left(a(t), \dot{a}(t), N(t), \phi(t), \dot{\phi}(t) \right)$$

Construct Canonical Hamiltonian from this Lagrangian

$$H_N = N H_0 (a, \Pi_a, \phi, \Pi_\phi)$$

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**What is N ? Different values of N yield different Hamiltonians
Different physics? Gauge Invariance?**

**If N is a non-trivial operator on the Fock space, no way to make physics gauge invariant
Possible choice: N is a c-number
But still, different choices of N yield different Hamiltonians!**

Schrodinger Equation

$$i \frac{\partial |\chi(t)\rangle}{\partial t} = NH_0|\chi(t)\rangle$$

N is a c-number, different choices of N yield different Hamiltonians

Schrodinger Equation

$$i \frac{\partial |\chi(t)\rangle}{\partial t} = NH_0 |\chi(t)\rangle$$

N is a c-number, different choices of N yield different Hamiltonians

$$i \frac{1}{N} \frac{\partial |\chi(t)\rangle}{\partial t} = H_0 |\chi(t)\rangle \quad d\tilde{t} = N dt$$

Schrodinger Equation

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N is a c-number, different choices of N yield different Hamiltonians

$$i \frac{1}{N} \frac{\partial |\chi(t)\rangle}{\partial t} = H_0 |\chi(t)\rangle \quad d\tilde{t} = N dt$$

**Different choices of N correspond to different choices of time co-ordinate.
Gauge invariant physics!**

$$\langle \phi_f, a_f | T(t_2; t_1) | \phi_i, a_i \rangle = \int_{\phi(t_1) = \phi_i, a(t_1) = a_i}^{\phi(t_2) = \phi_f, a(t_2) = a_f} D\phi Da e^{i \int_{t=t_1}^{t=t_2} (\tilde{\mathcal{L}}(N_0))}$$

Schrodinger Equation

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Consistent with Path Integral - you just pick N

Consequence of Schrodinger Equation

$$i \frac{\partial |\chi(t)\rangle}{\partial t} = N H_0 |\chi(t)\rangle \implies \frac{d \langle \chi(t) | H_0 | \chi(t) \rangle}{dt} = 0$$

Identity, similar to Ehrenfest and Schwinger-Dyson

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Thus: $\frac{d \langle \chi | a^3 \frac{\partial \mathcal{L}}{\partial N} | \chi \rangle}{dt} = 0$

**This is almost the 1st FRW equation - but not quite.
1st FRW equation only satisfied up to overall constant**

Initial State

$$a = A + A^\dagger, \Pi_a = i(A - A^\dagger) \quad \phi = B + B^\dagger, \Pi_\phi = i(B - B^\dagger)$$

Create quantum states of a, ϕ

$$|\chi\rangle = f(A, A^\dagger, B, B^\dagger) |0\rangle$$

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Choose $\langle\chi|a^3 \frac{\partial\mathcal{L}}{\partial N}|\chi\rangle = 0 \implies$ 1st FRW holds $\left(\frac{d\langle\chi|a^3 \frac{\partial\mathcal{L}}{\partial N}|\chi\rangle}{dt} = 0\right)$

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In classical physics, we demand 2 FRW + 1 KG equation to hold - so we restrict initial conditions to obey 1st FRW

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Quantum Dynamics

$$i \frac{\partial |\chi(t)\rangle}{\partial t} = N H_0 |\chi(t)\rangle$$

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Quantum Dynamics

$$i \frac{\partial |\chi(t)\rangle}{\partial t} = NH_0|\chi(t)\rangle$$

First order ODE - no issue with time evolving

$$\langle \chi | a^3 \frac{\partial \mathcal{L}}{\partial N} | \chi \rangle \neq 0$$

Violating 1st FRW

$$\langle \chi | a^3 \frac{\partial \mathcal{L}}{\partial N} | \chi \rangle = c$$

Quantum Dynamics: $i \frac{\partial |\chi(t)\rangle}{\partial t} = NH_0 |\chi(t)\rangle$

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Quantum Dynamics: $i \frac{\partial |\chi(t)\rangle}{\partial t} = NH_0 |\chi(t)\rangle$

Implied Classical Dynamics

$$\langle \chi | \partial_t \left(\frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) - \frac{\partial \mathcal{L}}{\partial \phi} | \chi \rangle = 0$$

2nd FRW Equation

$$\langle \chi | \partial_t \left(\frac{\partial \mathcal{L}}{\partial \dot{a}} \right) - \frac{\partial \mathcal{L}}{\partial a} | \chi \rangle = 0$$

Klein Gordon

$$\langle \chi | \frac{\partial \mathcal{L}}{\partial N} | \chi \rangle = \langle \chi | \frac{c}{a^3} | \chi \rangle$$

1st FRW but with “Dark” Matter

Quantum “Dark” Matter

$$\langle \chi | a^3 \frac{\partial \mathcal{L}}{\partial N} | \chi \rangle = c$$

$$a = A + A^\dagger, \Pi_a = i(A - A^\dagger) \quad \phi = B + B^\dagger, \Pi_\phi = i(B - B^\dagger)$$

Create quantum states of a, ϕ

Quantum Dynamics: Just 1 first order ODE (Schrodinger)

No reason to constrain initial state!

Failure manifests classically as “dark” matter - though no real particle excitation there. Conservation implies super-selection sector.

Can be positive or negative!

Wheeler deWitt

Can we get all of Einstein's Equations?

$$i\frac{\partial|\chi(t)\rangle}{\partial t} = NH_0|\chi(t)\rangle$$

Want $\langle\chi|a^3\frac{\partial\mathcal{L}}{\partial N}|\chi\rangle = 0$

Hamiltonian changes with N

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Universe is in an energy eigenstate

But... Time??

Path Integral Version

$$\langle \phi_f, a_f | T(t_2; t_1) | \phi_i, a_i \rangle = e^{iP(\phi_f, a_f, t_2; \phi_i, a_i, t_1; N_0)}$$

Gauge fixed path integral implies we lose 1st FRW equation

Path Integral explicitly depends upon random gauge choice No

Path integral obviously not gauge invariant

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Integrate over all N_0

Unsurprisingly, this yields infinity

Demanding gauge invariant path integral implies only possible states are static (Wheeler deWitt)

Conclusions

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Quantum Supremacy

Wheeler deWitt

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Quantum Supremacy

Gauge Dependent Time Evolution
With Gauge Invariant Physics

1st FRW Equation only true up
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Quantum “Dark” Matter

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Wheeler deWitt

Gauge Independent Time
Evolution, Infinite path integral

All Equations Hold...but...

Only static states...no time!

Conclusions

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Quantum Supremacy

Gauge Dependent Time Evolution
With Gauge Invariant Physics

1st FRW Equation only true up
to constant

Quantum “Dark” Matter

Full General Relativity , similar issues in Electromagnetism

Wheeler deWitt

Gauge Independent Time
Evolution, Infinite path integral

All Equations Hold...but...

Only static states...no time!