

Moving from continuous to discrete symmetry in the 2D XY model

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Based on
[arXiv:2205.03548](https://arxiv.org/abs/2205.03548)

Outline

- Examine effects of discretization and perturbations phase structure of 2d classical XY model
- Compare discretizations of XY model using HOTRG
 - Angular discretization (Z_N clock model)
 - U(1) character expansion
- Structure of TN formulations
 - Role of core and interaction tensors
- Perturbed XY model
 - Interpolates between XY and Z_N models
 - Study effect of small perturbation on phase structure
 - How large is the effect of small perturbation?

XY model discretizations

- 2d XY model

$$H = - \sum_{\langle xy \rangle} \cos(\theta_x - \theta_y) - h \sum_x \cos(\theta_x)$$

- Angular discretization (N states) $\rightarrow \mathbb{Z}_N$ clock model

$$\theta_x \in \left\{ \frac{2\pi k}{N}; k = 0 \dots N - 1 \right\}$$

- U(1) character expansion discretization (D = 2S+1 states)

$$e^{\beta \cos(\theta_x - \theta_y)} \approx \sum_{k=-S}^S I_k(\beta) e^{ik(\theta_x - \theta_y)}$$

Tensor network simulations

- HOTRG method (Z. Y. Xie, J. Chen, M. P. Qin, J. W. Zhu, L. P. Yang, T. Xiang 2012)
 - D (N for Z_N model) = initial number of states
 - D_{cut} = maximum states (after SVD)

- Observables
 - Specific heat

$$C_V = -T \partial_T^2 F$$

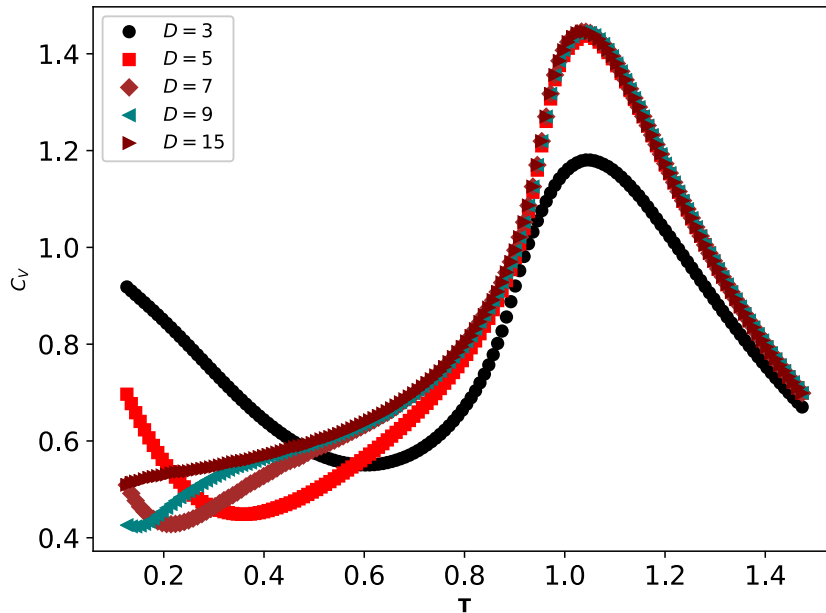
$$F = -(\ln Z) / (\beta V)$$

- Second derivative from 7-point polynomial fit
- Magnetization using impurity tensor method

$$\langle m \rangle = \frac{1}{V} \langle \sum_x \cos(\theta_x) \rangle$$

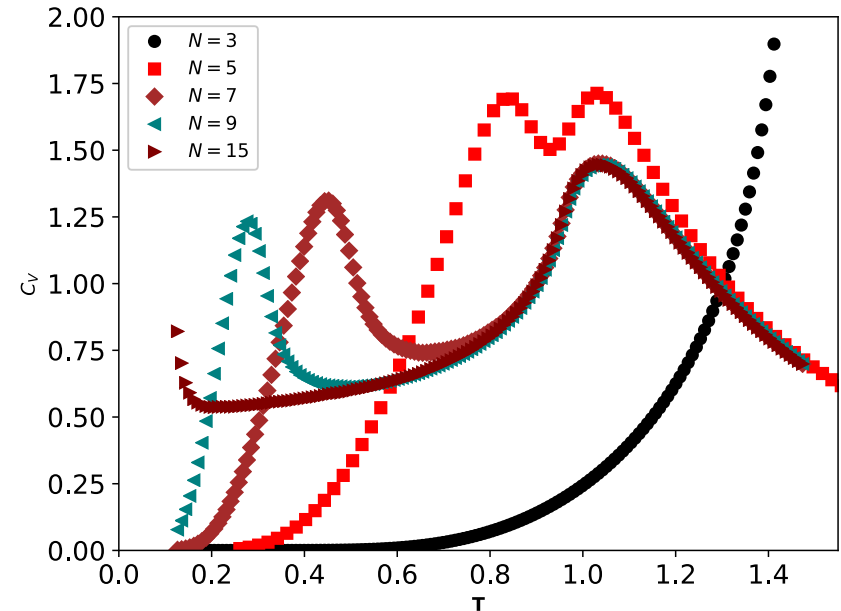
Specific heat of discretized XY model

$D_{\text{cut}} = 40$



U(1) character expansion

- Peak converges well for $D \geq 5$
- Low temperature converges well for $D \geq 15$
- Specific heat peak doesn't match phase transition (T_c), but similar

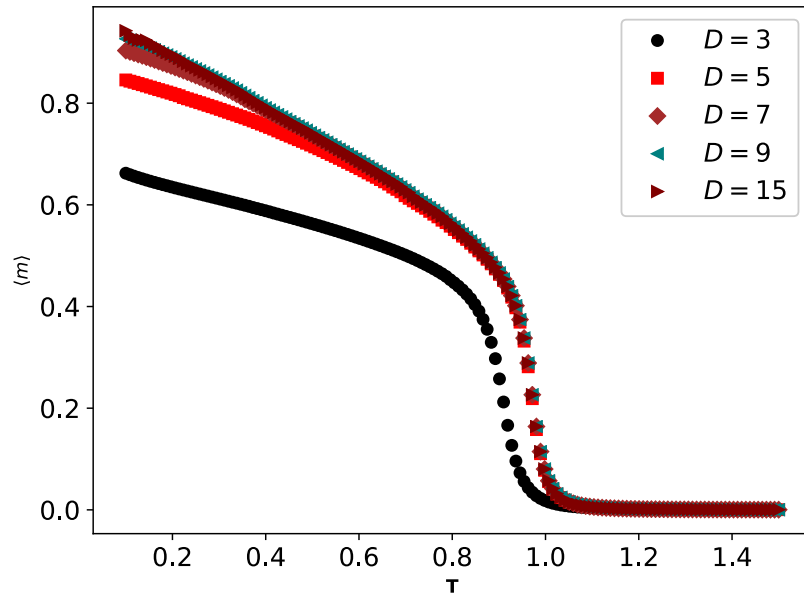


Z_N clock discretization

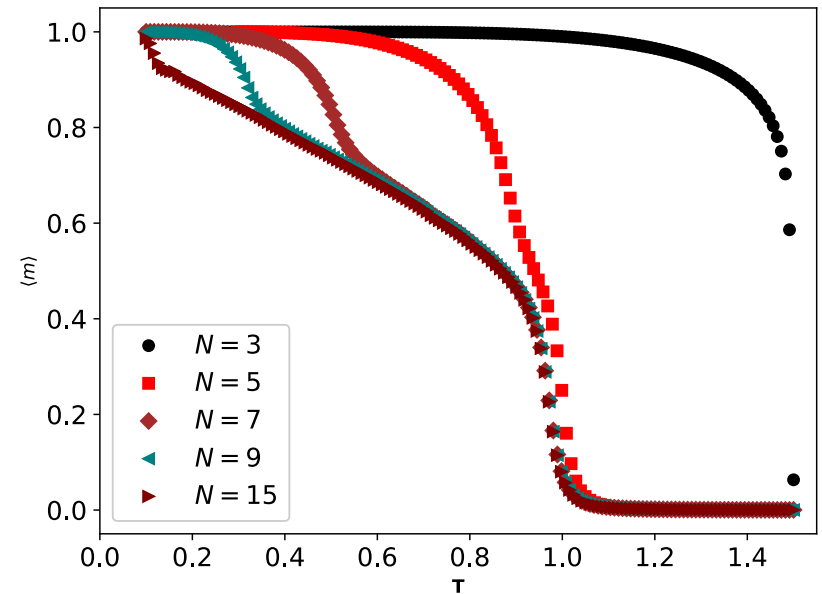
- One peak for $N \leq 4$
- Two peaks for $N \geq 5$
- Upper peak converges well for $N \geq 7$
- Larger difference from XY at low T for same number of states ($D=N$)

Magnetization of discretized XY model

$D_{\text{cut}} = 40$ $h = 10^{-4}$



U(1) character expansion



Z_N clock discretization

- Ordered phase appears in Z_N model, but not XY

Finding phase transitions

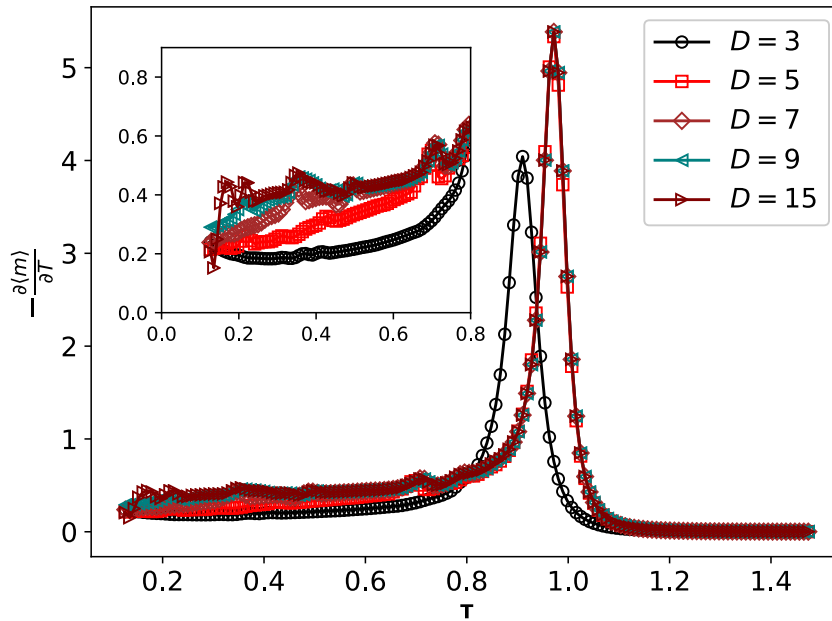
- Specific heat, magnetic susceptibility not reliable indicators of lower phase transition in Z_N ($N \geq 5$) model
- Cross derivative (Y. Chen, K. Ji, Z. Y. Xie, J. F. Yu 2020)

$$\frac{\partial^2 F}{\partial h \partial T} = - \frac{\partial \langle m \rangle}{\partial T}$$

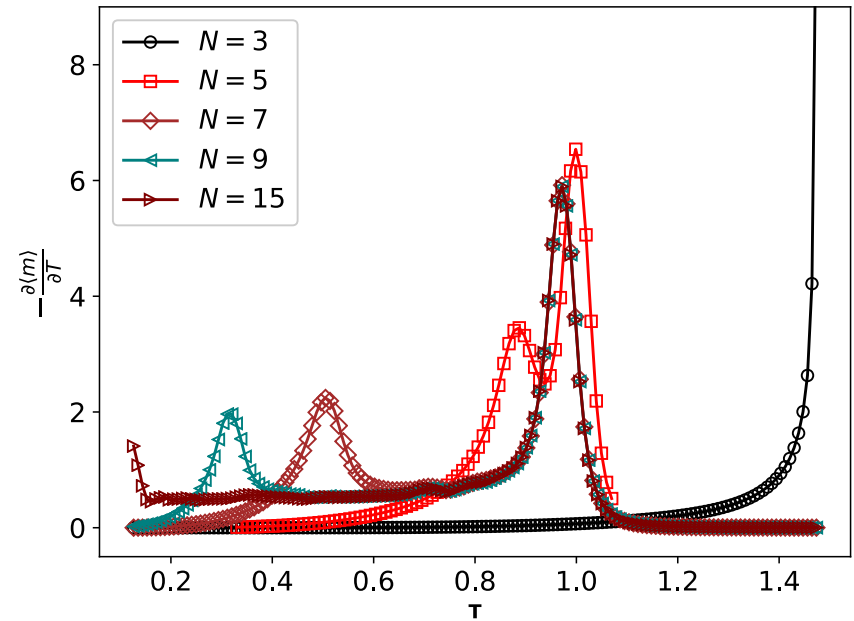
- Seems to be good indicator of phase transitions
- Obtained from fit of magnetization to polynomial in T

Cross derivative of discretized XY model

$D_{\text{cut}} = 40$ $h = 10^{-4}$



U(1) character expansion



Z_N clock discretization

- Gives clear peaks at phase transitions

Tensor network components

Spin models

- Typically start with interaction and local terms

$$H = \sum_{\langle xy \rangle} H_I(\theta_x, \theta_y) + \sum_x H_L(\theta_x)$$

- Expand interaction in some basis

$$e^{-\beta H_I(\theta_x, \theta_y)} \approx \sum_{a,b} M_{ab} f_a(\theta_x) g_b(\theta_y)$$

- Gives a core tensor of (for 2d)

$$C_{abcd} = \int d\theta e^{-\beta H_L(\theta)} f_a(\theta) g_b(\theta) f_c(\theta) g_d(\theta)$$

Tensor network components

- XY model tensor components (U(1) character basis)

$$M_{ab} = \delta_{ab} I_a(\beta)$$
$$C_{abcd} = I_{a-b+c-d}(\beta h) \xrightarrow{h \rightarrow 0} \delta_{a-b+c-d}$$

- Clock model tensor components (Z_N character basis)

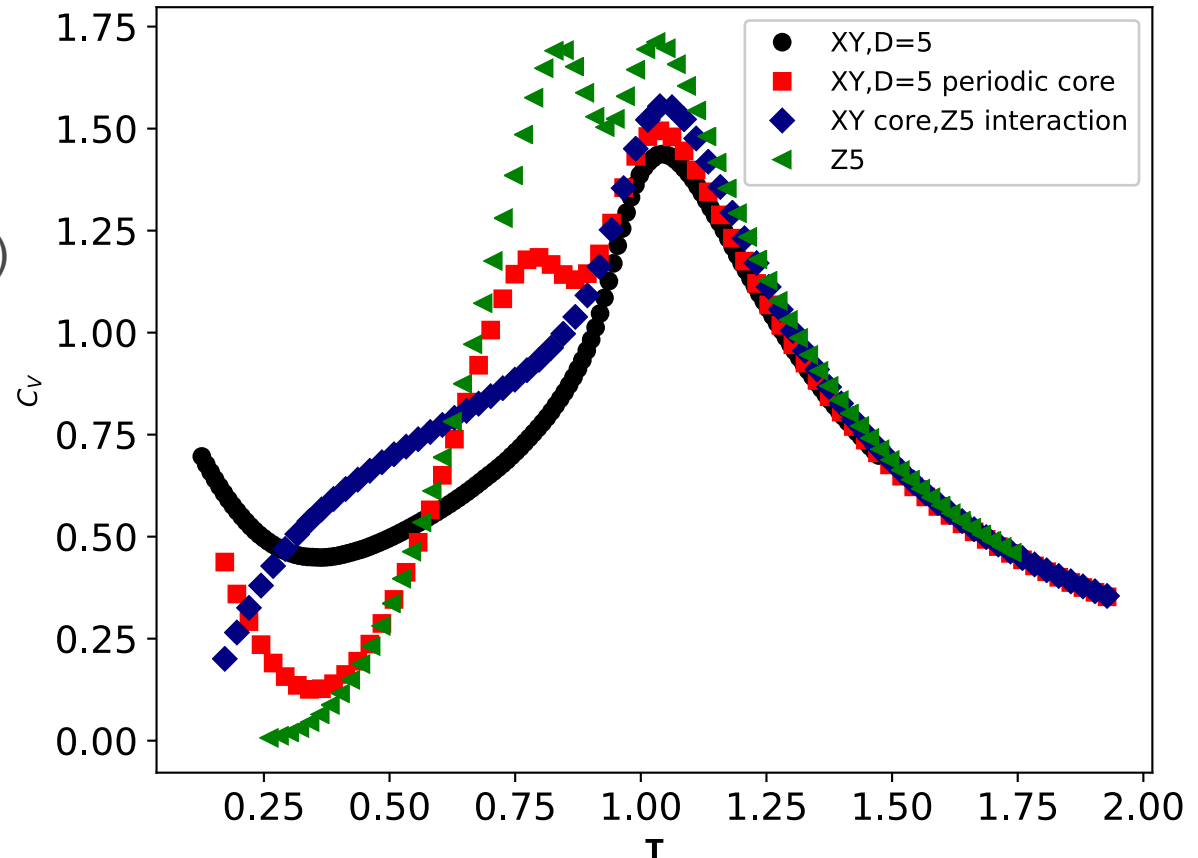
$$M_{ab} = \delta_{ab} \sum_k I_{a+kN}(\beta)$$
$$C_{abcd} = \sum_k I_{a-b+c-d+kN}(\beta h) \xrightarrow{h \rightarrow 0} \sum_k \delta_{a-b+c-d+kN}$$

- Can mix core and interaction tensors from different models

Specific heat for mixed XY/ Z_N models

$D_{\text{cut}} = 40$

- 2 peaks appear with periodic core (Z_5 and XY+periodic core)
- Only one peak apparent with XY core
- Phase structure determined by symmetry of core tensor (Y. Meurice 2019)



Discretization as a choice of basis

- Core, interaction decomposition not unique
- For example, can move local term into interaction

$$H'_I(\theta_x, \theta_y) = H_I(\theta_x, \theta_y) + \frac{1}{2d} [H_L(\theta_x) + H_L(\theta_y)] \quad e^{-\beta H'_I(\theta_x, \theta_y)} \approx \sum_{a,b} M'_{ab} f_a(\theta_x) g_b(\theta_y)$$

- Core tensor then determined only by basis $C'_{abcd} = \int d\theta f_a(\theta) g_b(\theta) f_c(\theta) g_d(\theta)$
- Z_N clock model as approximation to interaction term using a rectangular step function basis, each over distinct angular range of $2\pi/N$, zero elsewhere
 - Evaluating interaction at center of each interval reproduces clock model
- Could choose other sets of basis functions, expansion approximations
- Evaluating error in discrete approximation to interaction term may help determine which basis is best (does this consider symmetries?)
 - RMS error for Z_N basis is larger than $U(1)$ character basis
- Could use this to optimize basis for given model

Moving between XY and Z_N models

- Previously looked at Z_N for increasing N
- Instead consider fixed N
 - interpolate between periodic, non-periodic core tensor

$$C_{abcd} = \sum_k \alpha_k(h_N) I_{a-b+c-d+kN}(\beta h)$$

- Perturbed XY model (Jose, Kadanoff, Kirkpatrick, Nelson 1977)

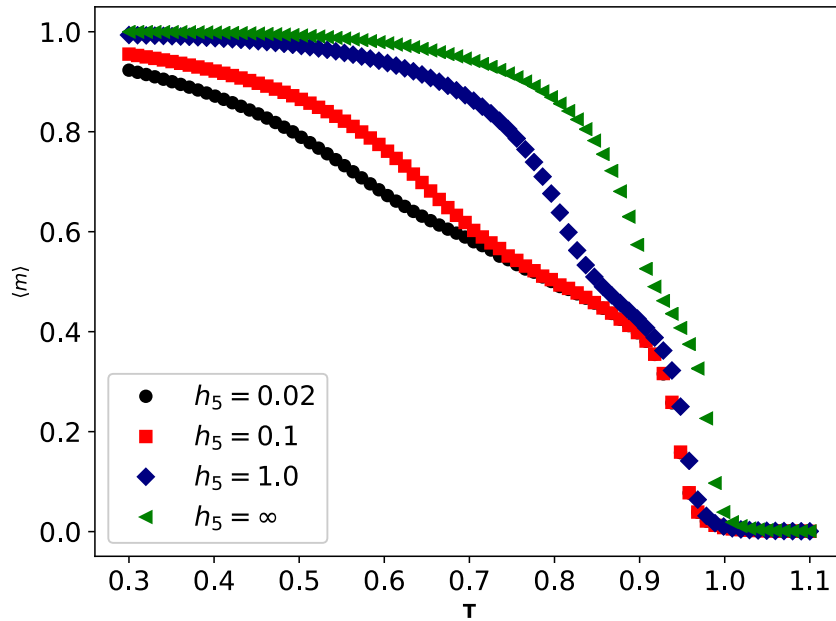
$$H = - \sum_{\langle xy \rangle} \cos(\theta_x - \theta_y) - h \sum_x \cos(\theta_x) - h_N \sum_x \cos(N\theta_x)$$

$$\alpha_k(h_N) \propto \frac{I_k(\beta h_N)}{I_0(\beta h_N)} = \begin{cases} \delta_k : h_N = 0 \text{ (XY)} \\ 1 : h_N \rightarrow \infty \text{ (ZN)} \end{cases}$$

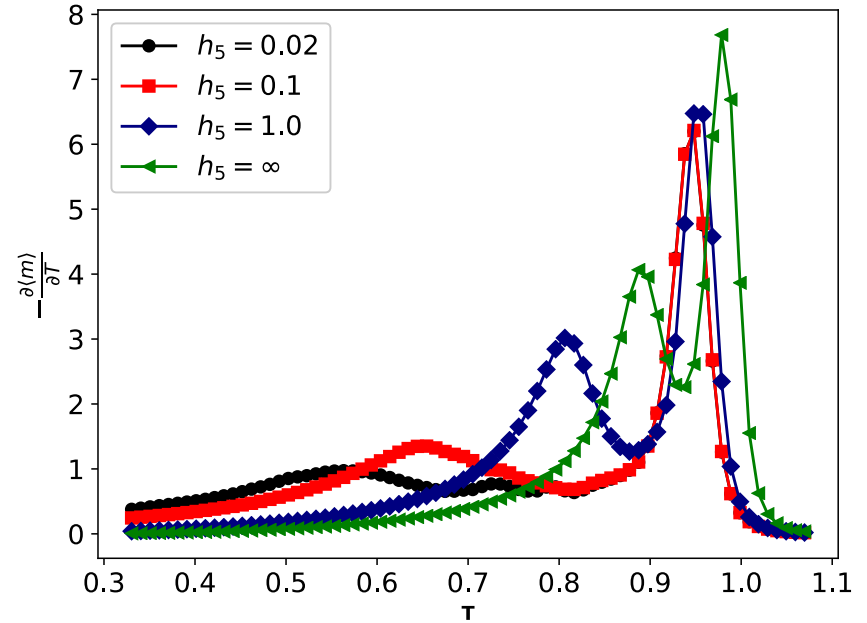
- Interested in behavior of lower phase transition for small h_N

Perturbed model N=5

$D_{\text{cut}} = 40$ $h = 1e-5$



Magnetization



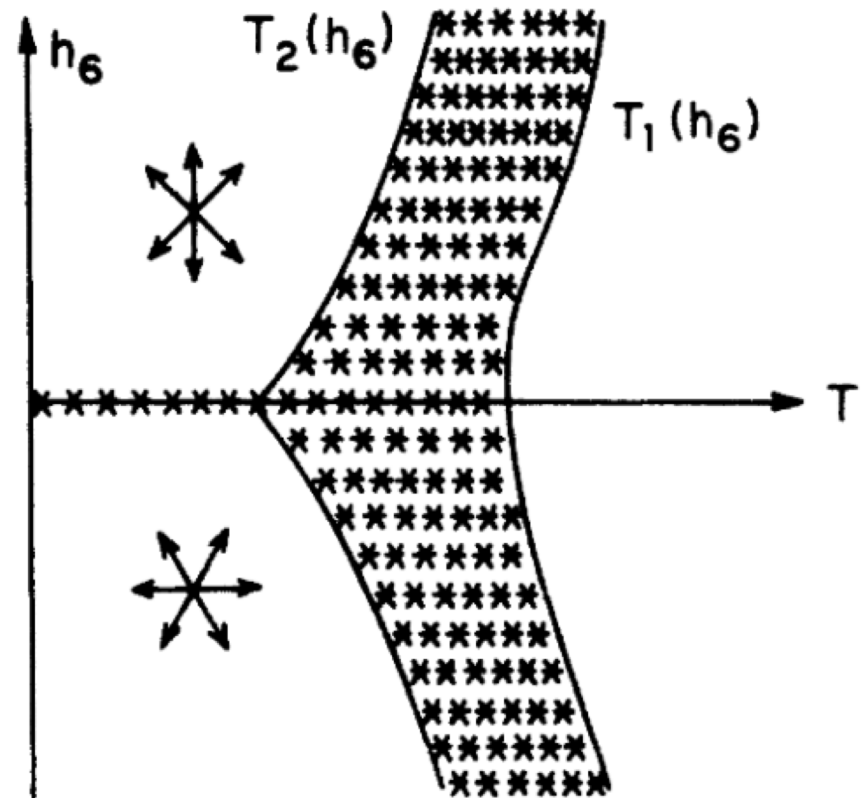
Cross derivative

Perturbed XY model RG analysis

- RG analysis using Villain approximation in JKKN paper

$$1 - \cos(\theta_x - \theta_y) \approx (\theta_x - \theta_y)^2 / 2$$

- For fixed $h_N > 0$, found
 - 2 critical temperatures for $N > 4$
 - At $N=4$ collapse to single T_c
- Value of $T_2(h_N \rightarrow 0)$ not clear from RG analysis
- Interested if non-zero, and how large
- If large, infinitesimal perturbation has large effect on phase diagram



(Jose, Kadanoff, Kirkpatrick, Nelson 1977)

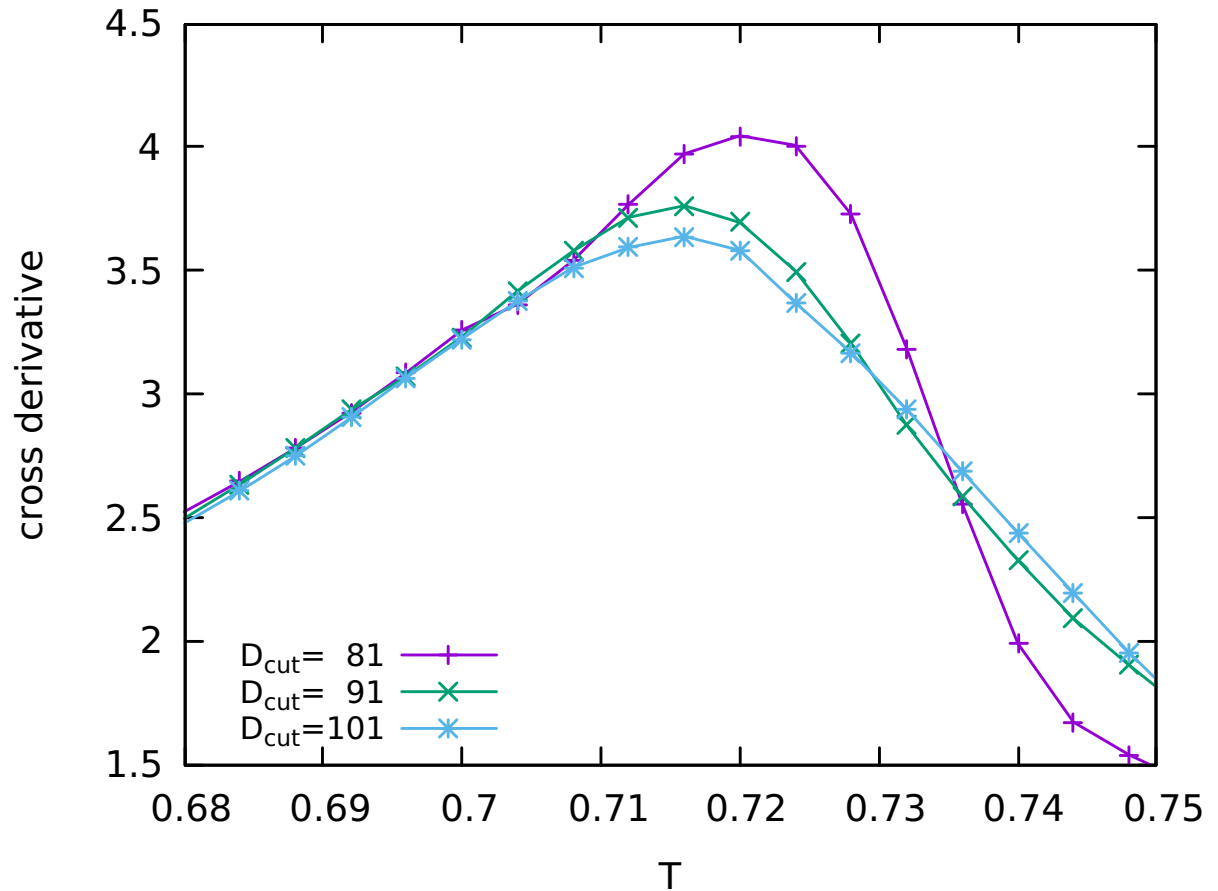
Determining peaks of cross derivative

- Calculated $\langle m \rangle$ with spacing $\Delta T = 0.004$
- Fit 11 consecutive points to 4th order polynomial
 - Plots of cross derivative from derivative of polynomial at center of interval
 - Peak location from zero of 2nd derivative of polynomial
- Consider only peak fits where peak is well within interval
- Peak temperature and height averaged over fit intervals

Cross derivative peak example

$h_5 = 0.2$ $h = 2e-7$

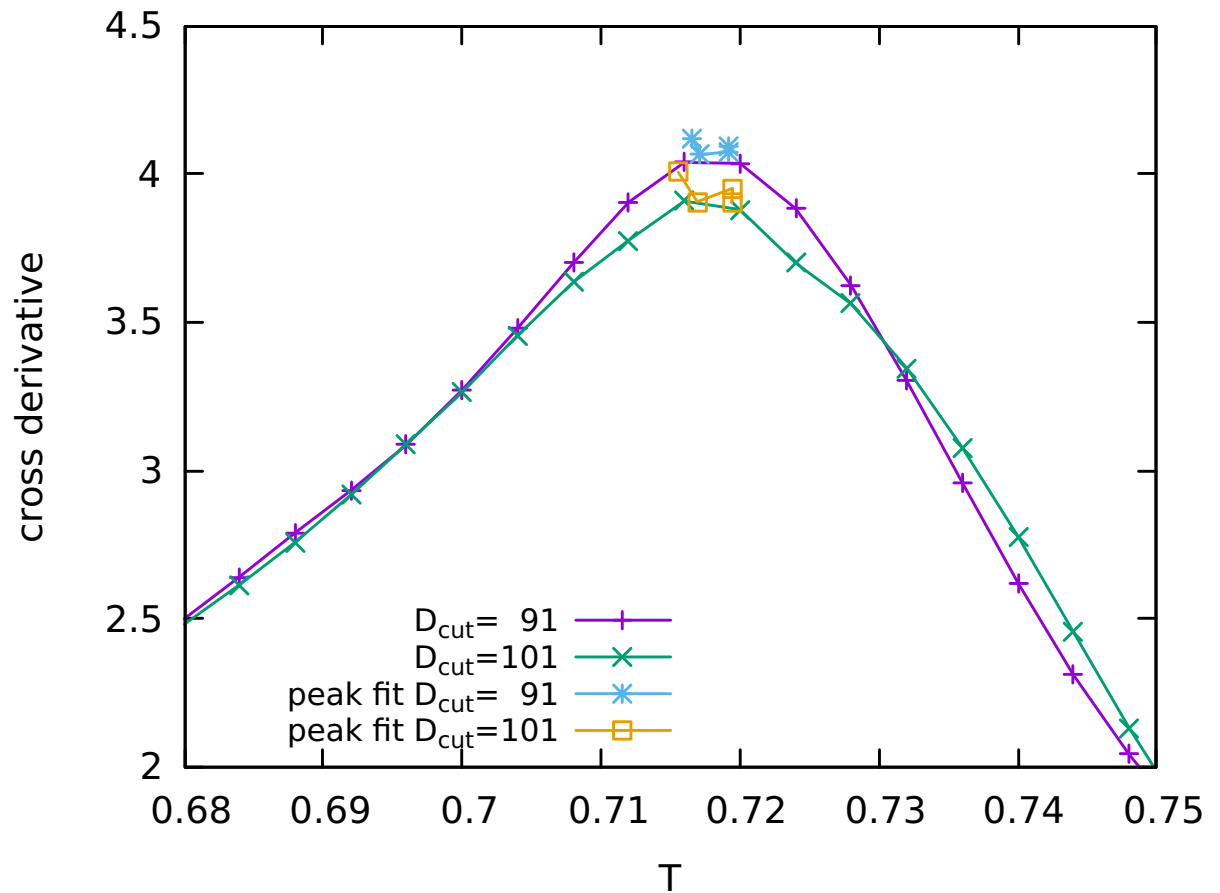
- $D_{\text{cut}}=81$ not enough
- Peak temperature consistent for $D_{\text{cut}}=91, 101$
- $D_{\text{cut}}=91$ seems sufficient for T_c
- Peak height still moving lower a bit
- Lowest $h_5 = 0.2$, simulations became unstable around $h=1e-7$ for $h_5 < 0.2$



Cross derivative peak fits

$h_5 = 0.2$ $h = 1e-7$

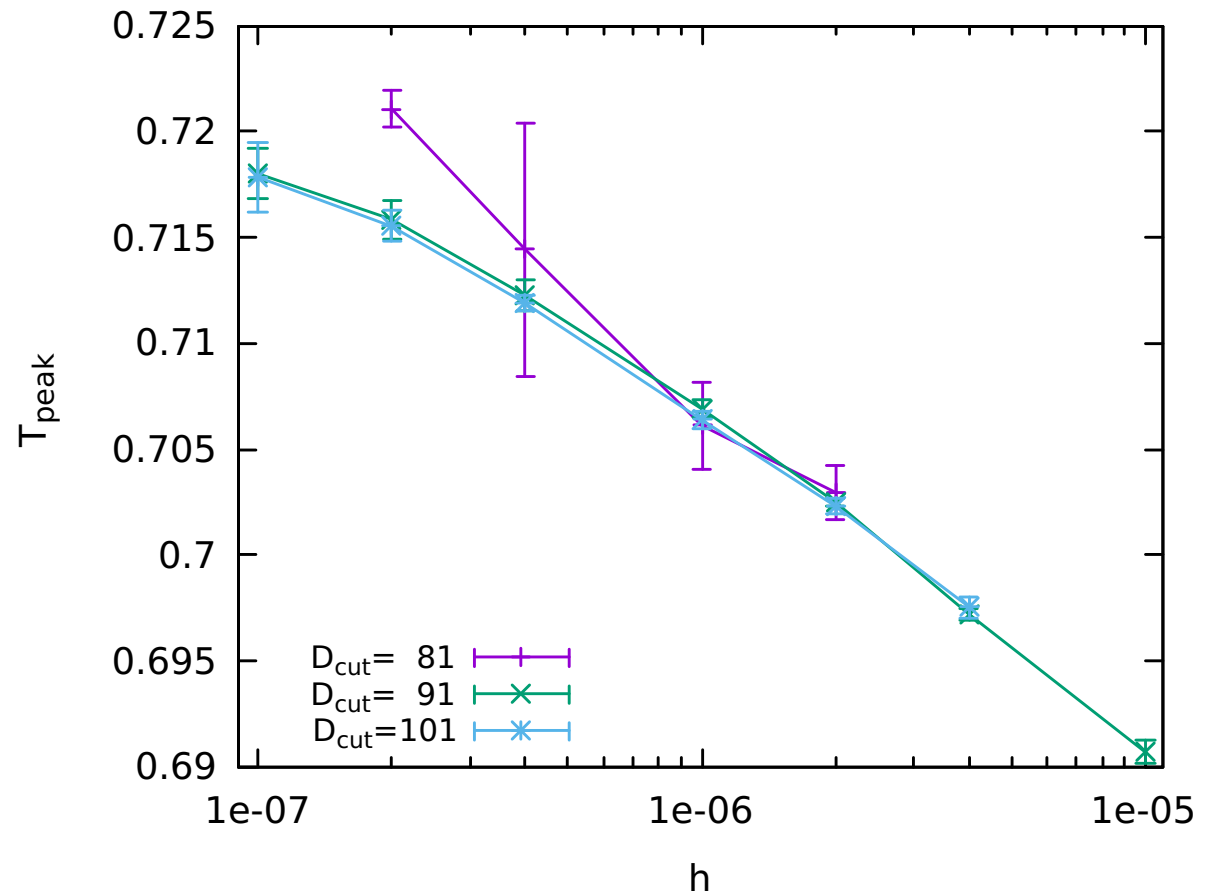
- ~4 points fell well within fit interval
- Peak determined from average, error from variance
- Error is systematic error of fit
- For peak height additional error from D_{cut}



Lower peak temperature

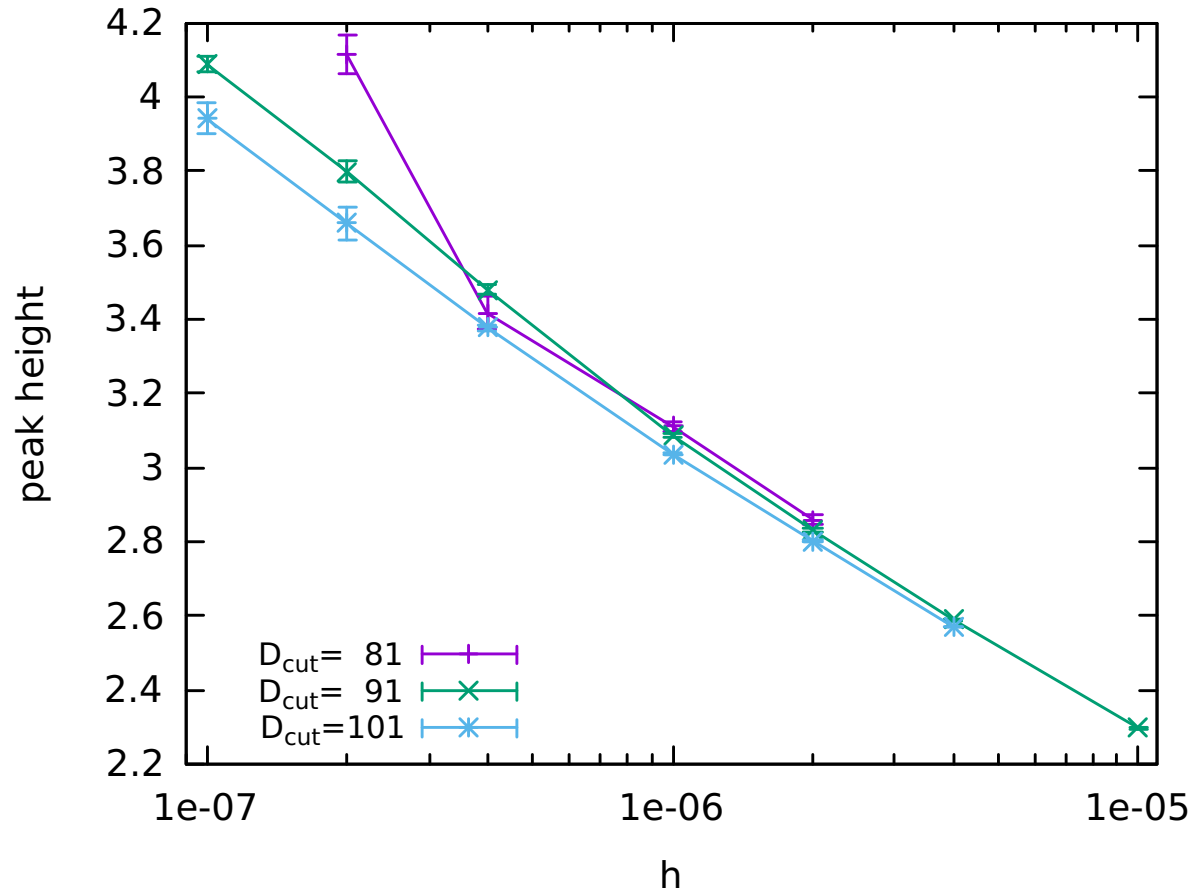
$h_5 = 0.2$

- $D_{\text{cut}} = 81$ not accurate enough at small h
- $D_{\text{cut}} = 91$ and 101 agree within errors
- Used $D_{\text{cut}} = 91$ for most of remaining results



Lower peak height

$h_5 = 0.2$

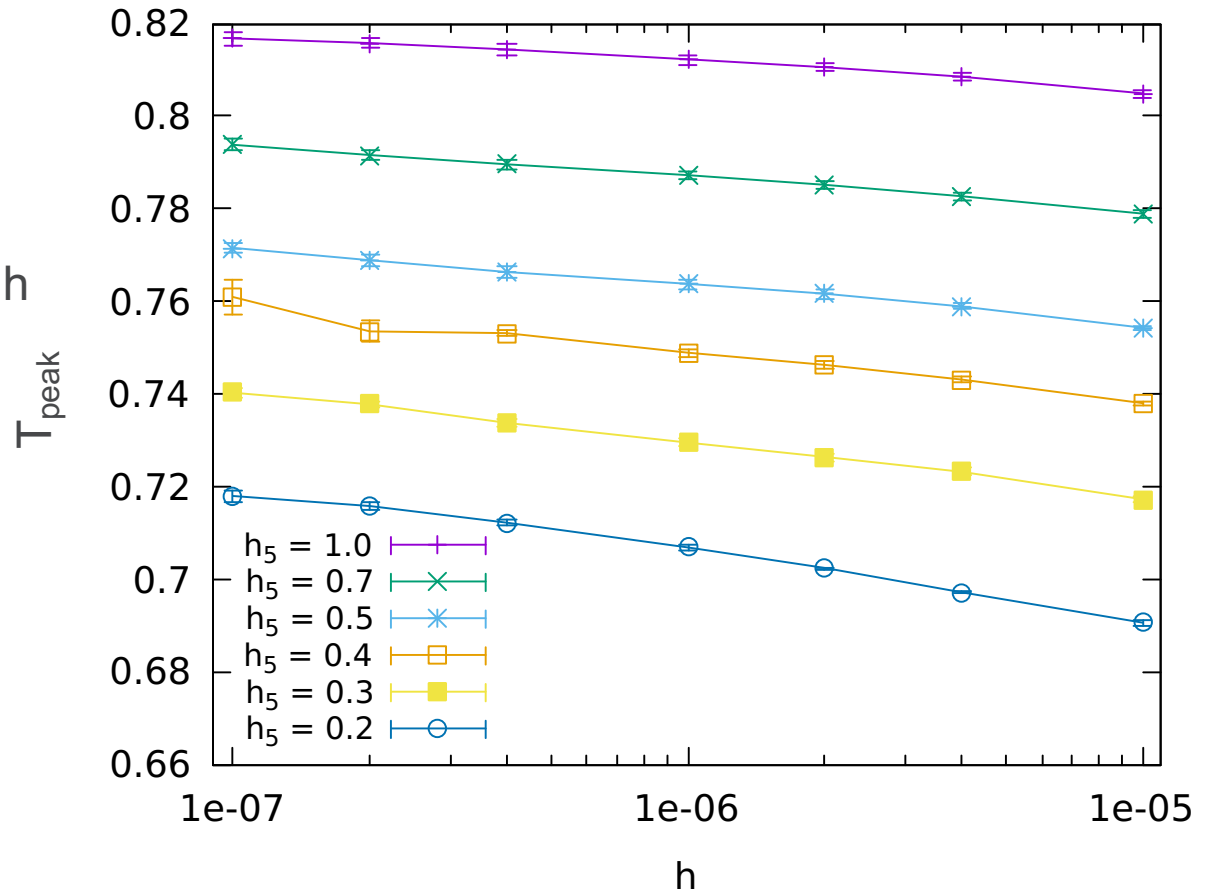


Lower peak temperature versus h

- T_{peak} rises at lower h,
falls at lower h_5
- Fit T_{peak} to power-law in h

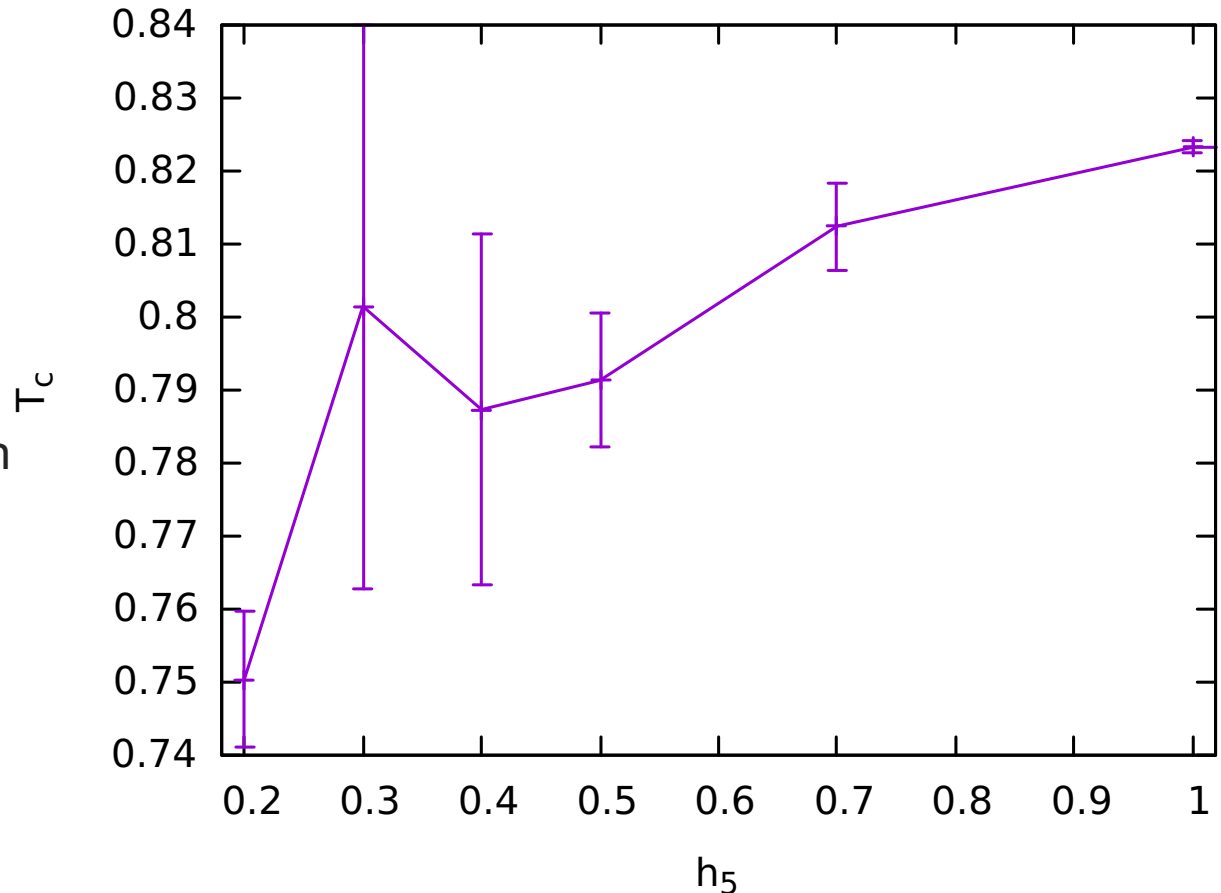
$$T_c(h_5) - T_{peak} \propto h^\alpha$$

to extract $T_c(h_5)$



Lower critical temperature versus h_5

- Error bars too large for power-law fit
- Trending down at lower h_5
- Still far from 0, but unable to do reliable extrapolation with this approach

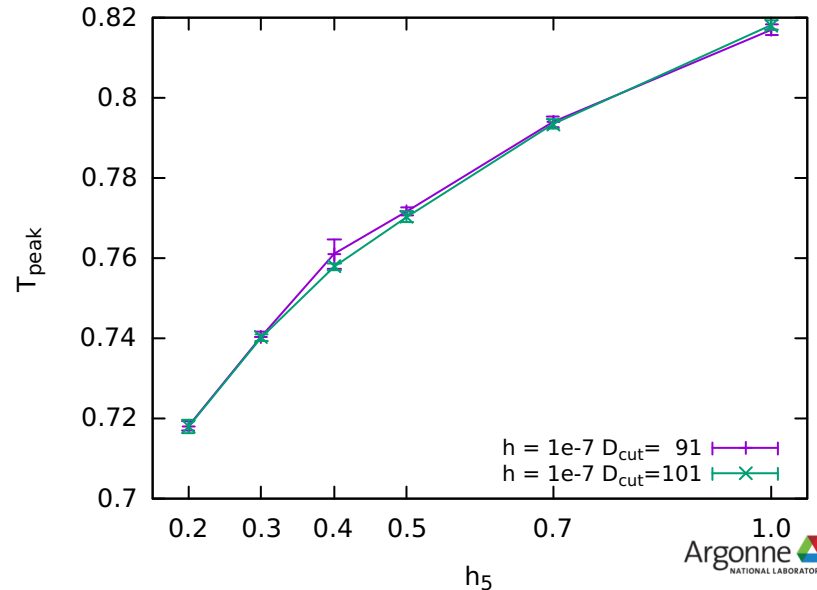
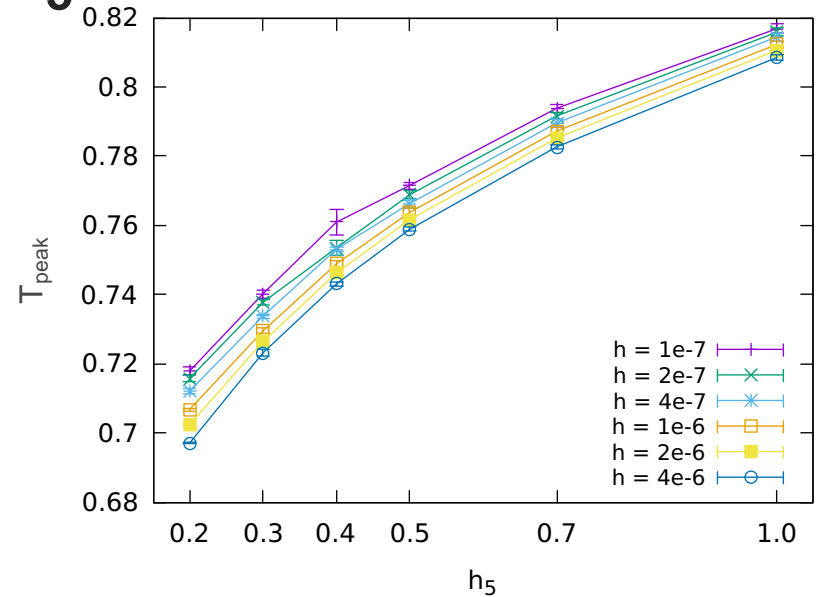


Peak temperature versus h_5

- First fit to h_5 at fixed h
- power-law in h_5

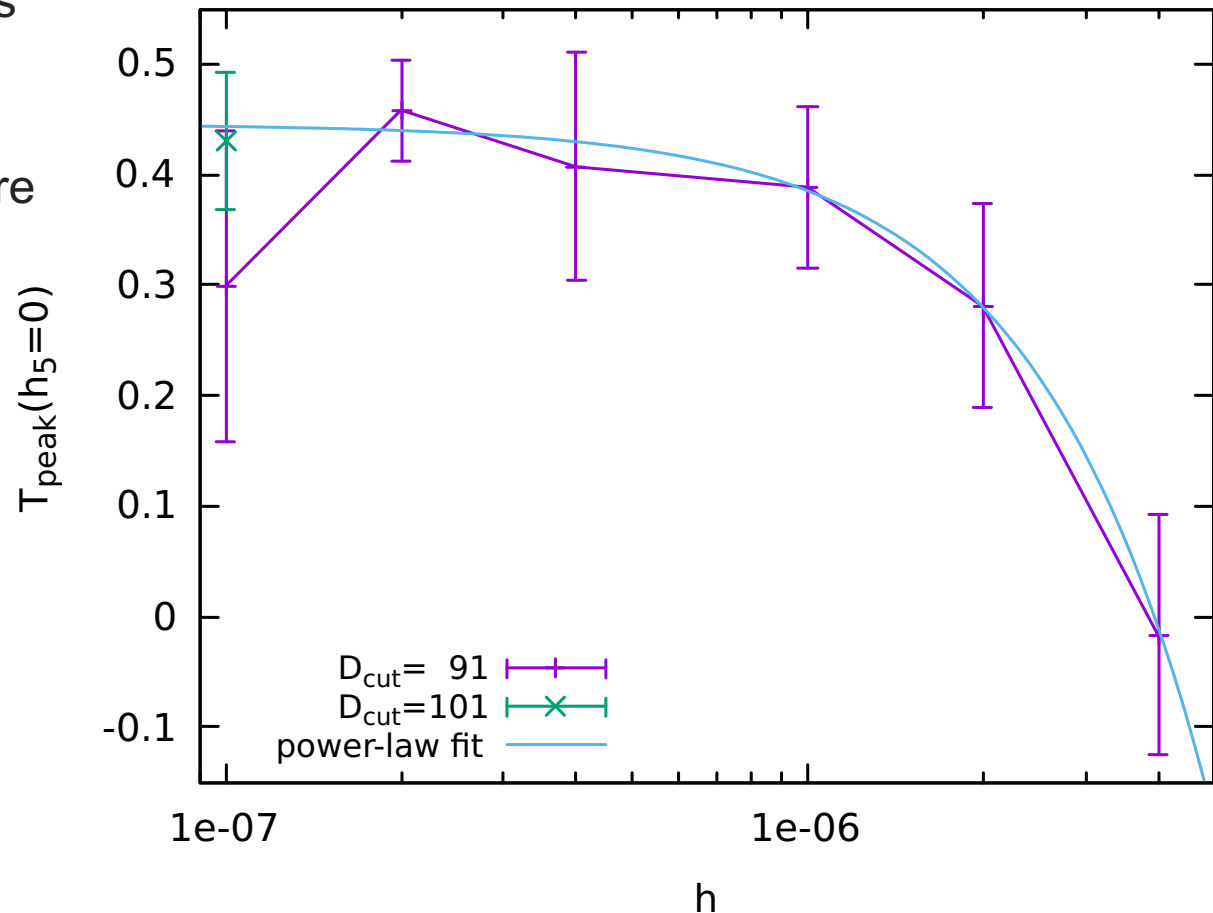
$$T_{peak}(h_5) - T_{peak}(0) \propto h_5^\alpha$$

- Then extrapolate $T_{peak}(h_5 \rightarrow 0)$ to $h = 0$



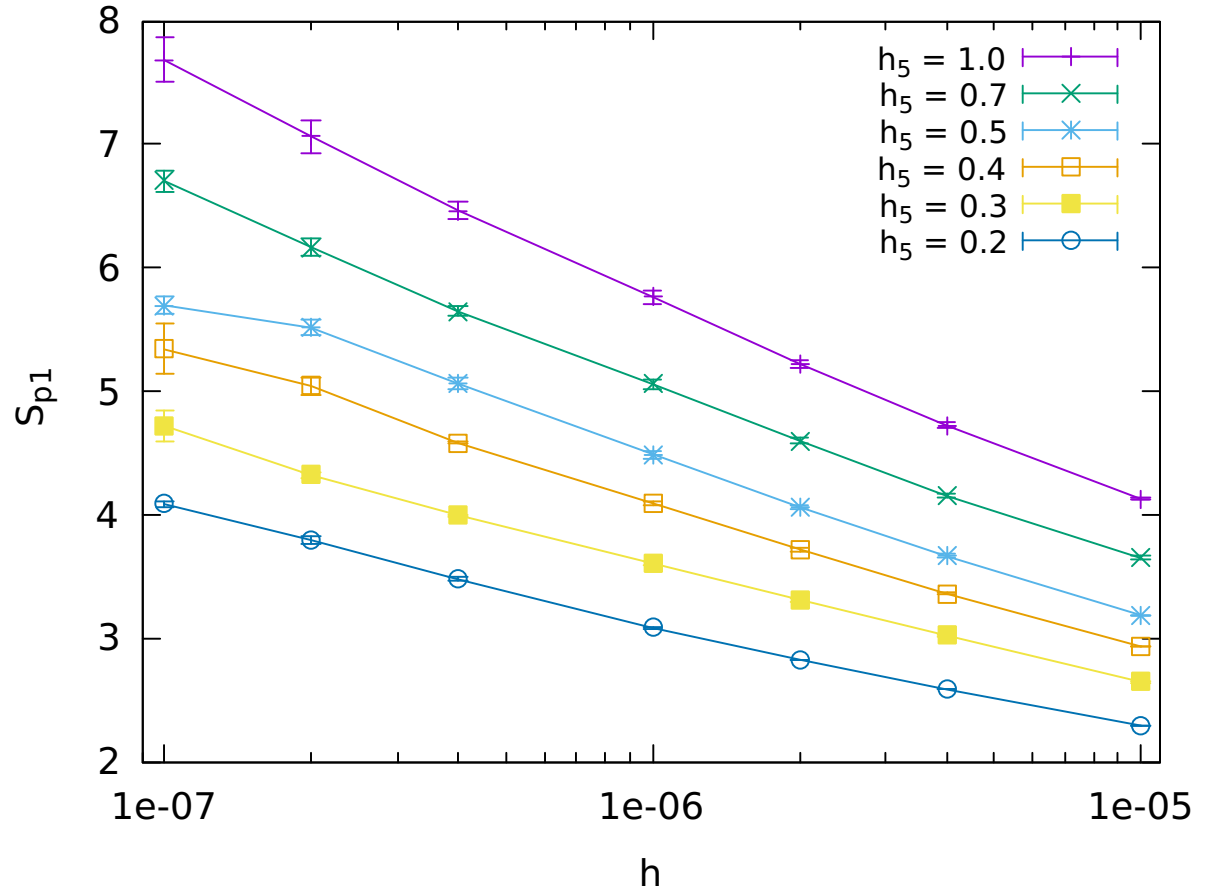
Extrapolated peak temperature versus h

- Power-law fit extrapolates to $T_{\text{peak}}(h_5 \rightarrow 0) = 0.44(3)$
- About half the temperature of XY $T_c \sim 0.89$
- Small perturbation seems to have large effect on phase diagram



Peak height

- Need to check that system is critical (peak height diverges) as $h \rightarrow 0$
- Fit peak heights to power-law
 $height \propto h^{-\gamma}$
- Check if gamma is nonzero as $h_5 \rightarrow 0$

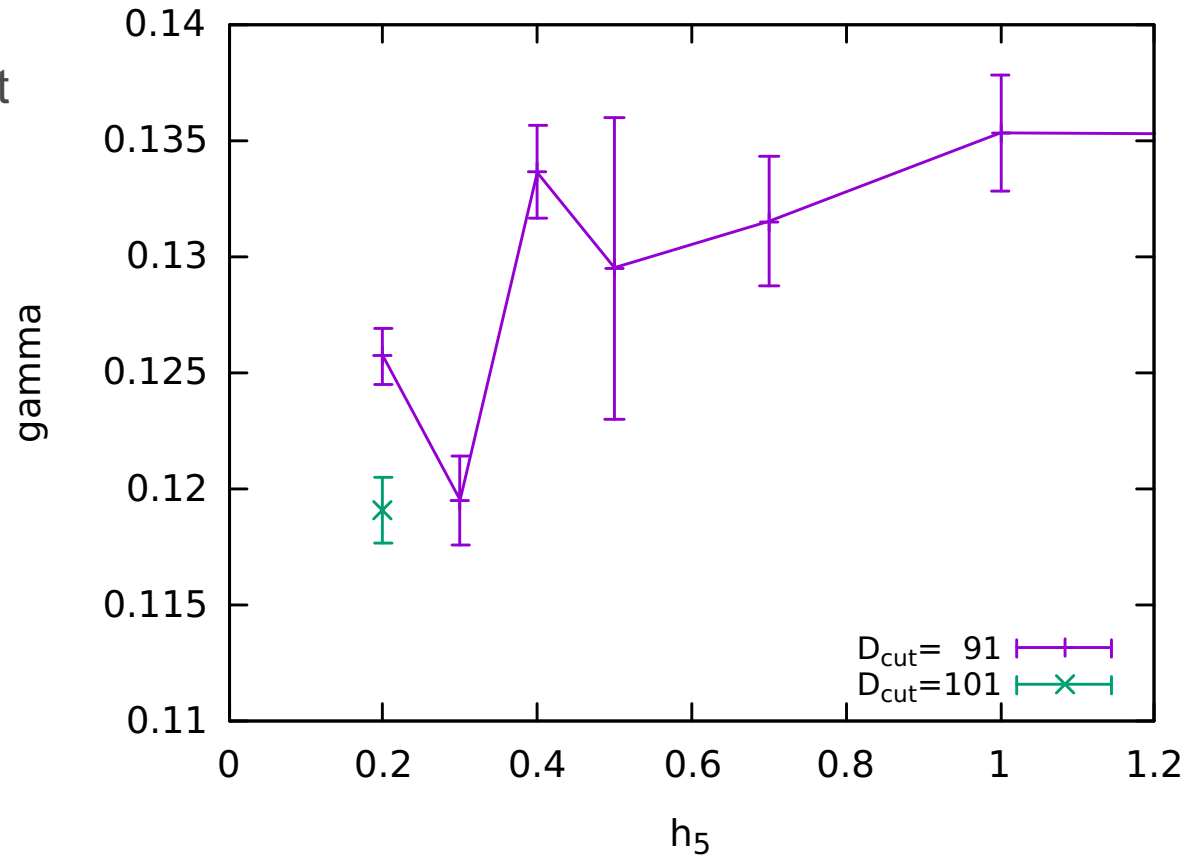


Peak height scaling

- Power-law exponent

$$\text{height} \propto h^{-\gamma}$$

- Trending down as h_5 decreases, but still far from 0
- Consistent with critical transition as $h_5 \rightarrow 0$
- Need smaller h_5 to verify



Summary

- XY truncation using U(1) character expansion has less effect on low energy behavior than angular discretization
- Expansion basis for TN construction determines core tensor symmetries, which in turn has large effect on phase structure
- Choosing basis that preserves symmetries
 - Does that also minimize error in truncated interaction expansion?
- Studied location of lower-T phase transition in perturbed XY model as function of h_5
 - consistent with non-zero T_c ($\sim 0.5 T_{XY}$) as $h_5 \rightarrow 0$
- Small symmetry-breaking perturbations can have large effect on phase structure