

Multi-messenger nuclear & astrophysics with added crust

William G. Newton

The work presented in this talk would not be possible without an amazing team of undergraduates and Master's students:

Rebecca Preston, Amber Stinson, Lauren Balliet, Brianna Douglas, Michael Ross, Gabriel Crocombe, Blake Head, Alex Westbrooks, Sarah Cantu, Josh Sanford, Srdj Budimir, Luis Rivera, Zachary Langford

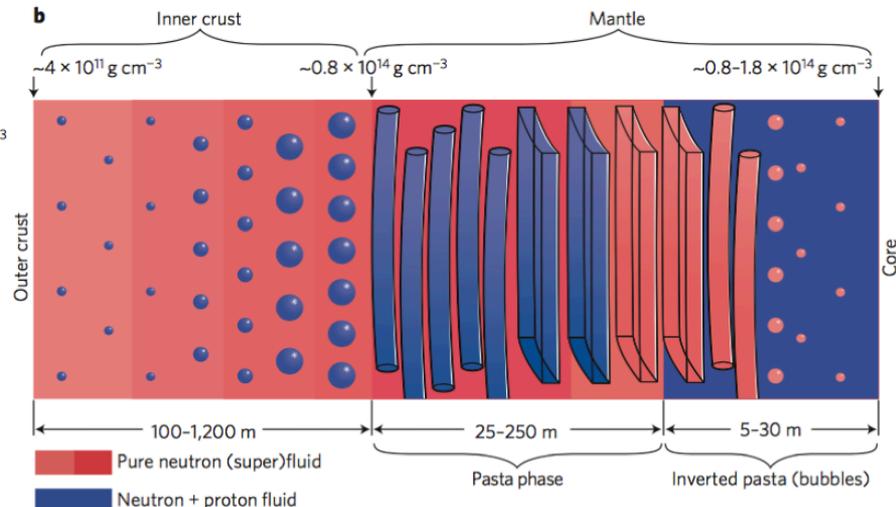
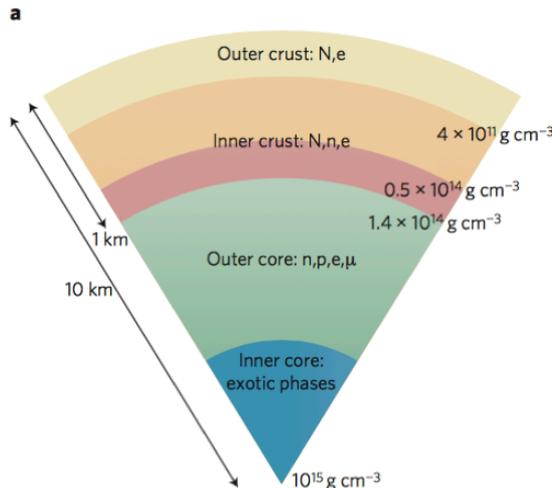
Texas A&M University-Commerce

Duncan Neill, David Tsang – University of Bath

Jirina Rikovska Stone, Alex Kaltenborn - University of Tennessee



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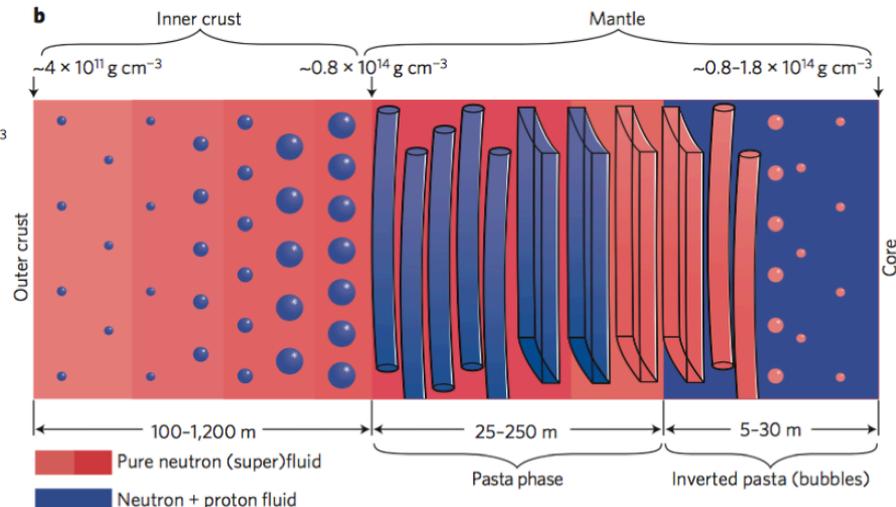
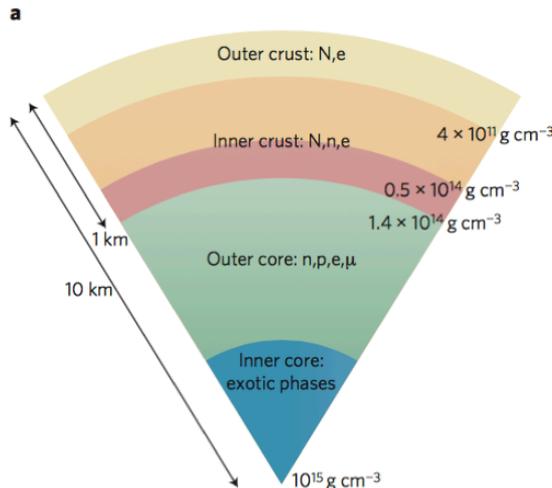
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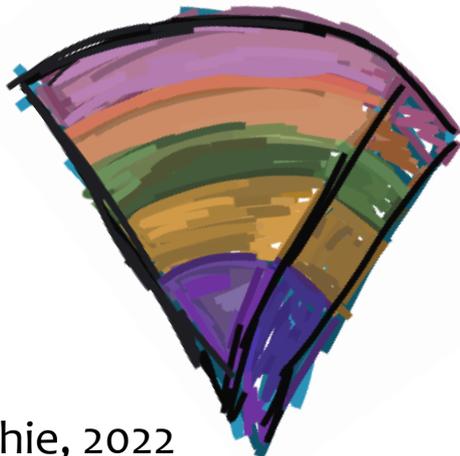
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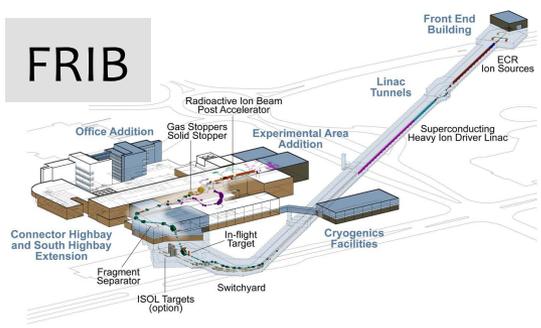


Strong, Weak, EM signals

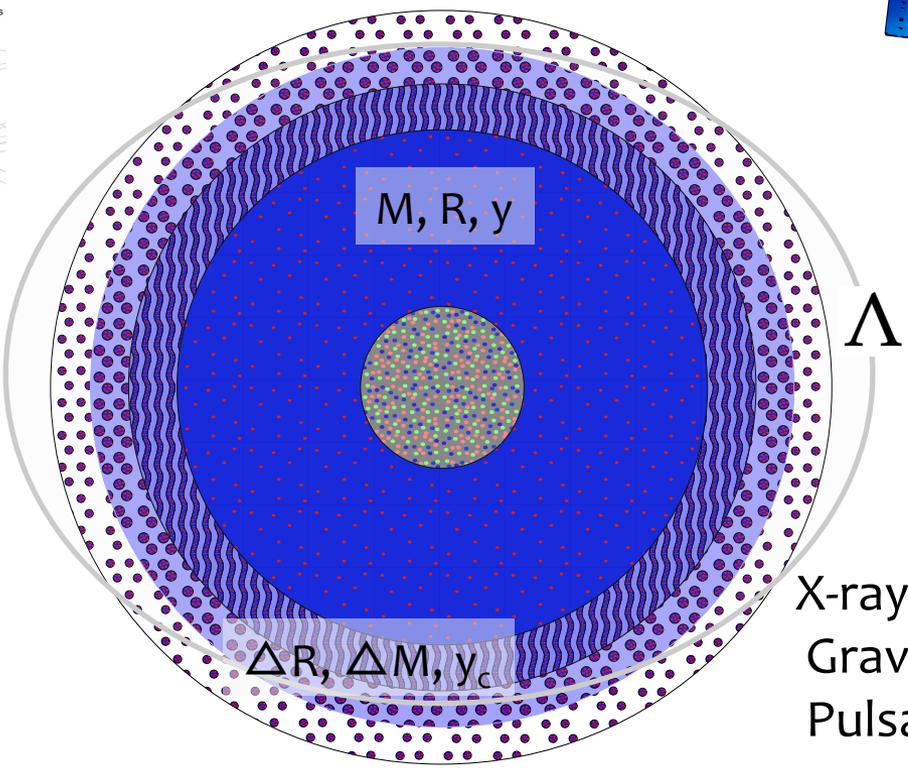
Multimessenger Nuclear & Astro Physics

Weak, EM, Grav signals

FRIB



Elliptic flow
 p/n ratios
 Pion production
 Resonance widths,
 Centroid energies
 Optical potentials
 Scattering X-sections

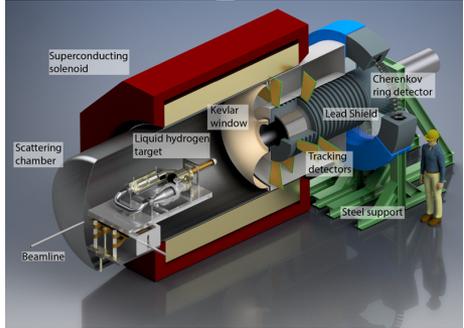


CHANDRA



NICER

X-ray flux and light curves
 Gravitational waveforms
 Pulsar timing



PREX/CREX/MREX

Computation



Randy Wong/LLNL

PARKES

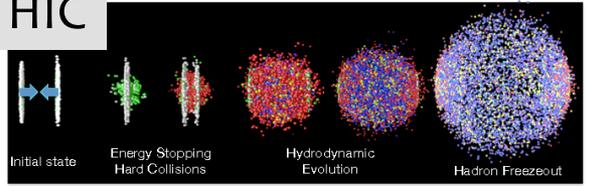


LIGO/VIRGO



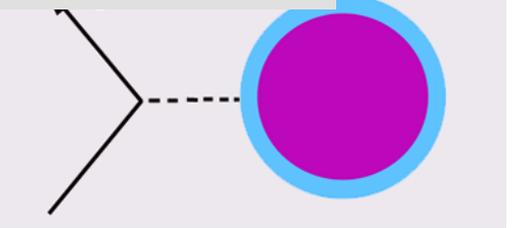
Nuclear structure/ dynamics

HIC

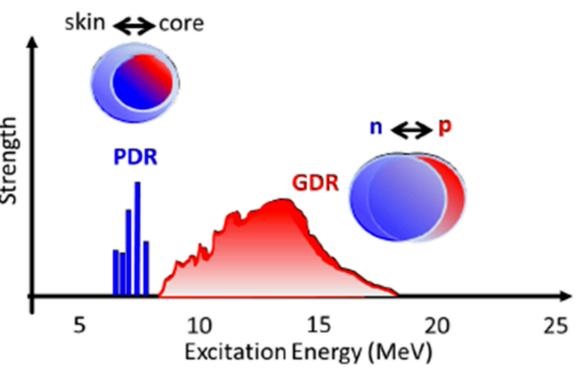


T.K.Nayak, arxiv:1201.4264

neutron skins

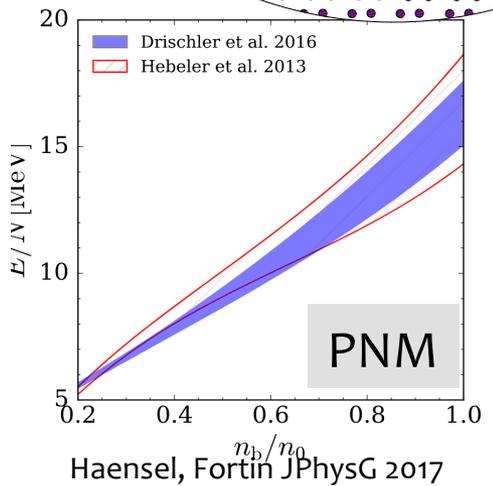
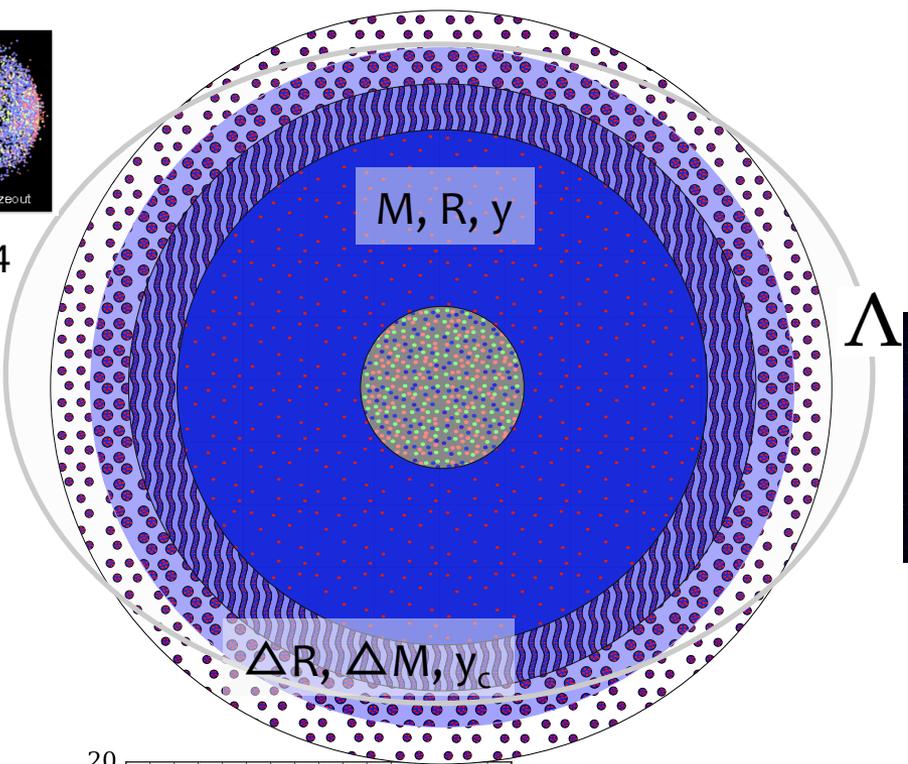


Abrahamyan+,
PRL 108, 112592 (2012)



Bracco, Lanza, Tamii,
PPNP 106, 360 (2019)

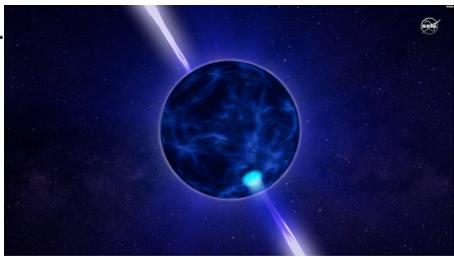
Multimessenger Nuclear & Astro Physics



Haensel, Fortin JPhysG 2017

Neutron star structure/ dynamics

Glitches, flares, cooling



Hot spots Oscillations, Crust cooling

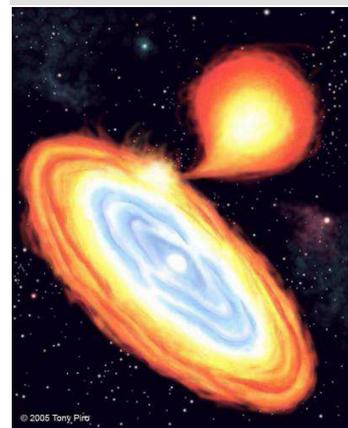
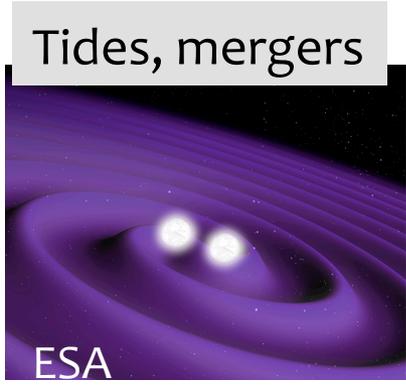
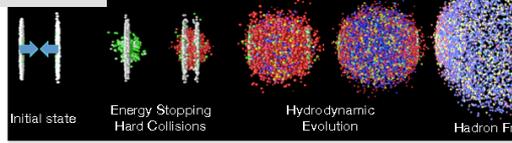


Figure: Artist's impression of a LMXB - credit Tony Piro, 2005.



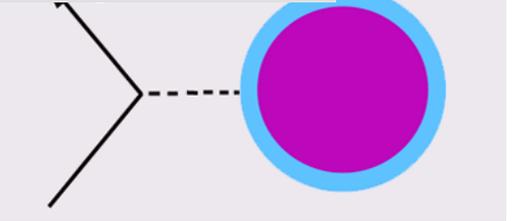
Nuclear structure/ dynamics

HIC

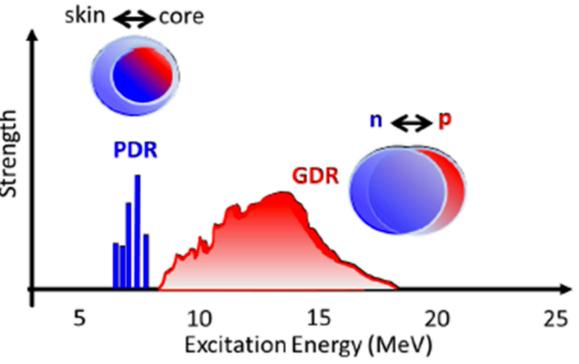


T.K.Nayak, arxiv:1201.426

neutron skins

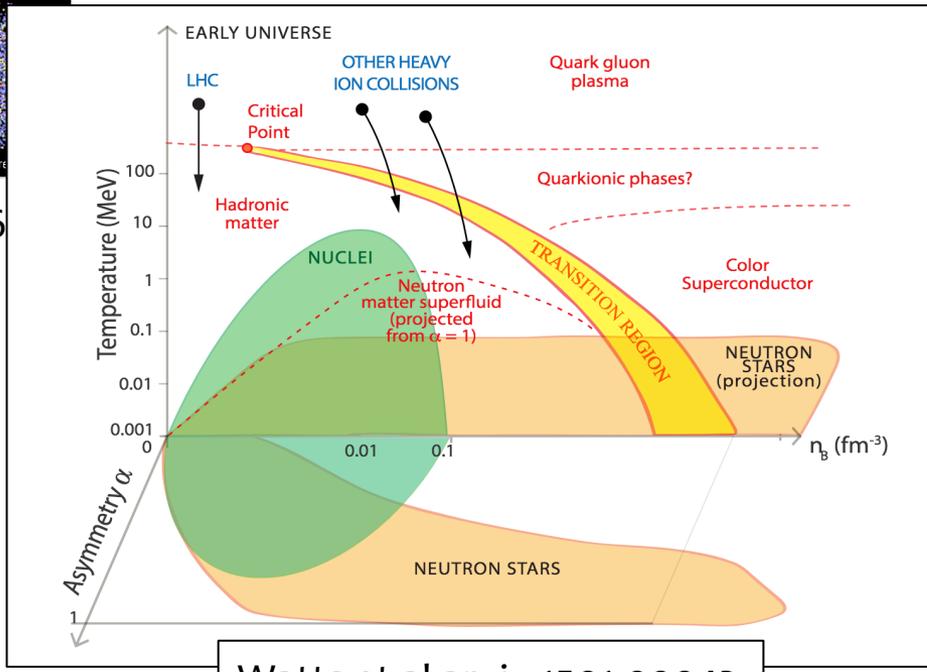


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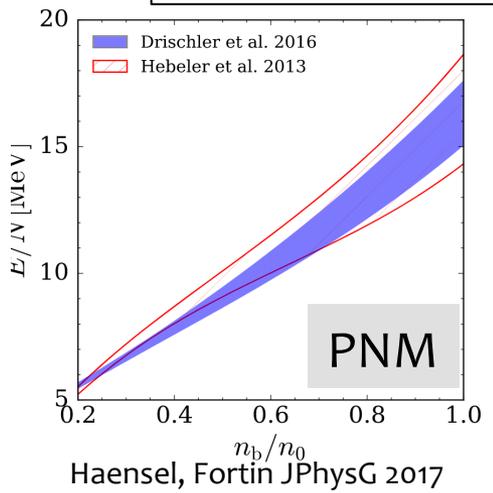


Bracco, Lanza, Tamii,
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Multimessenger Nuclear & Astro Physics



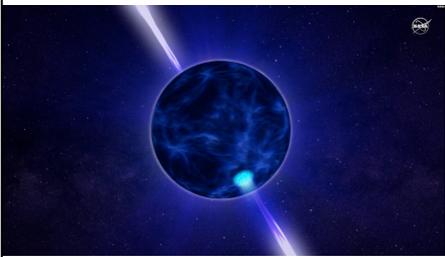
Watts et al arxiv:1501.00042



Haensel, Fortin JPhysG 2017

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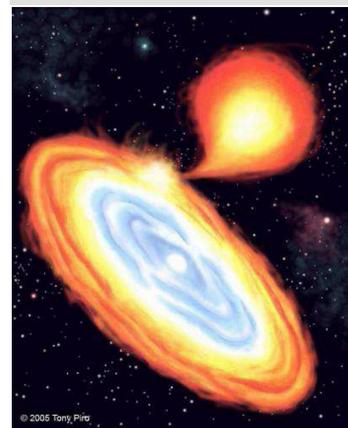
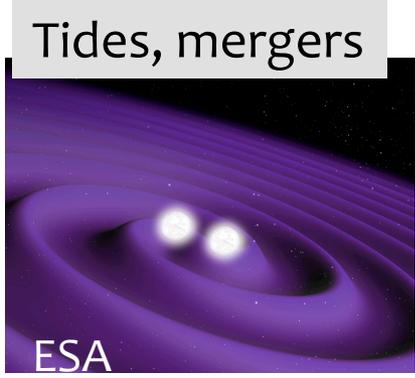


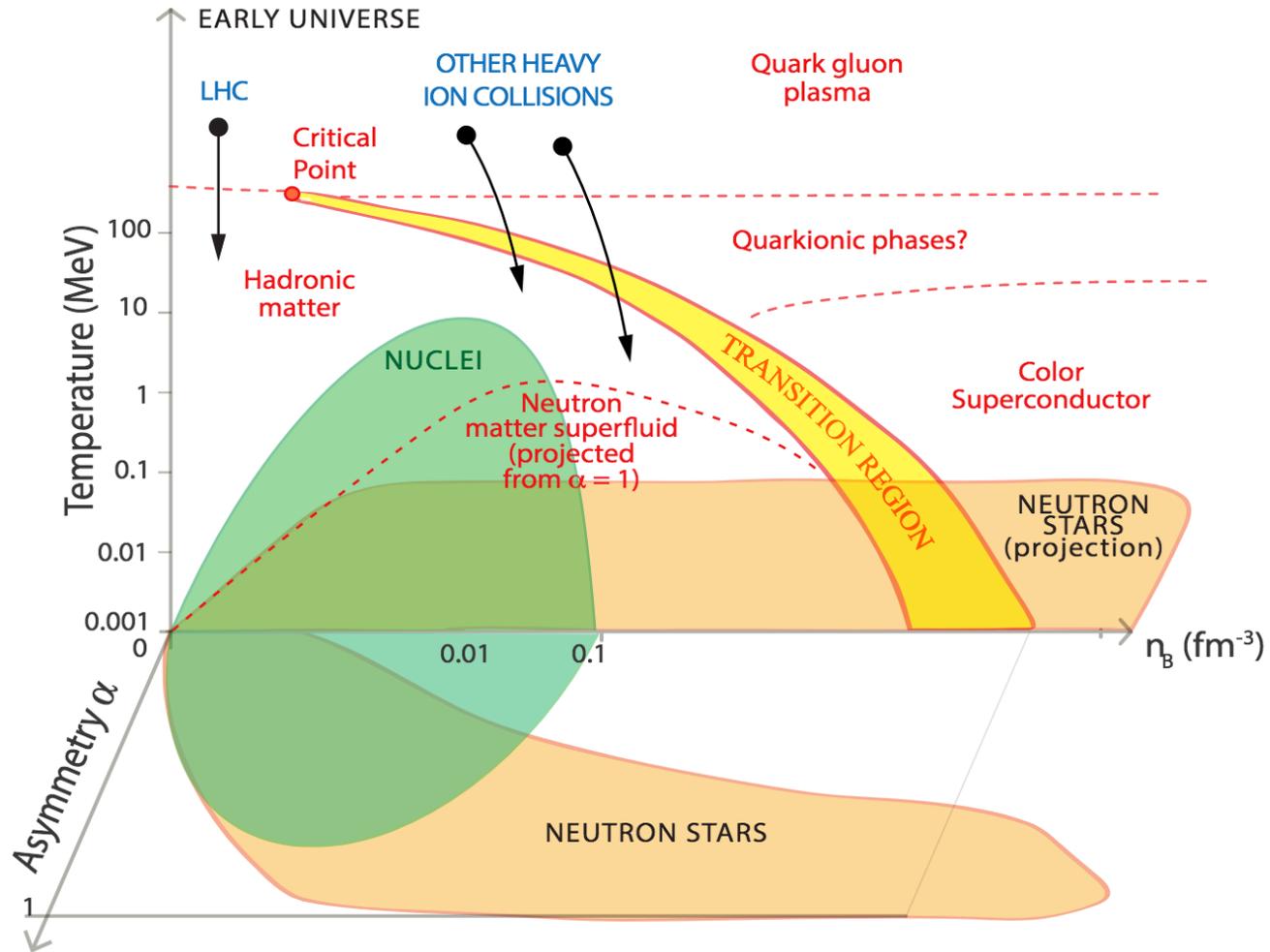
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Tides, mergers

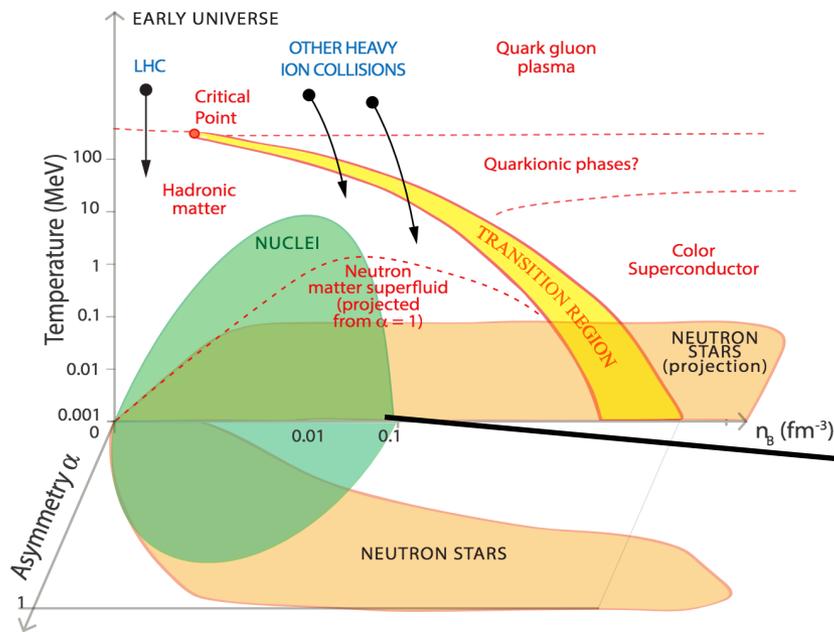
ESA

The phase diagram of nuclear matter



Watts et al arxiv:1501.00042

The phase diagram of nuclear matter to the EOS



Watts et al arxiv:1501.00042

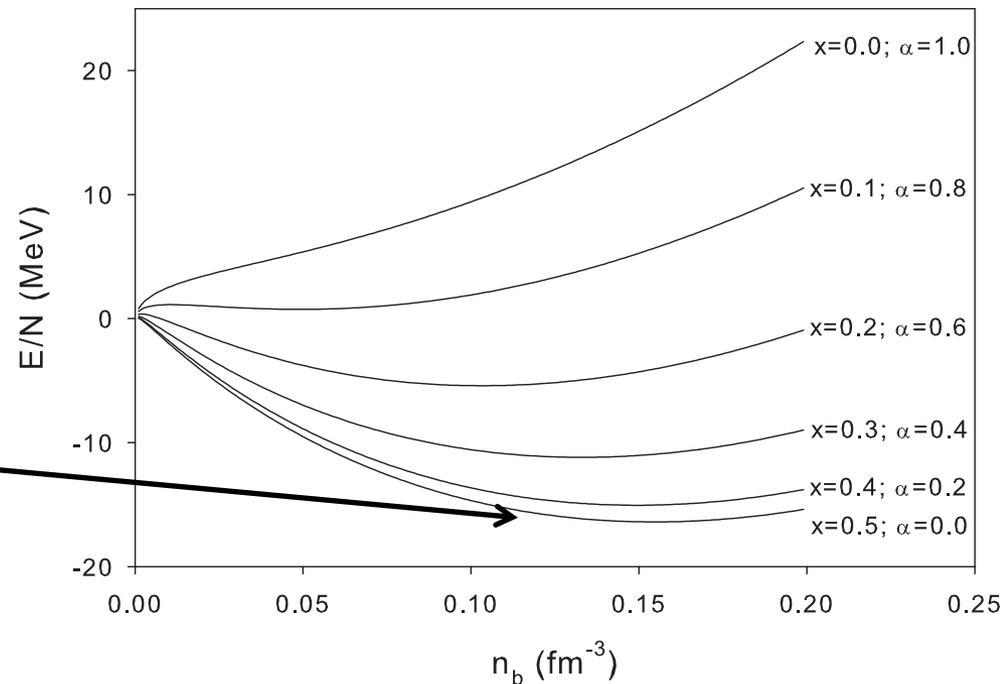
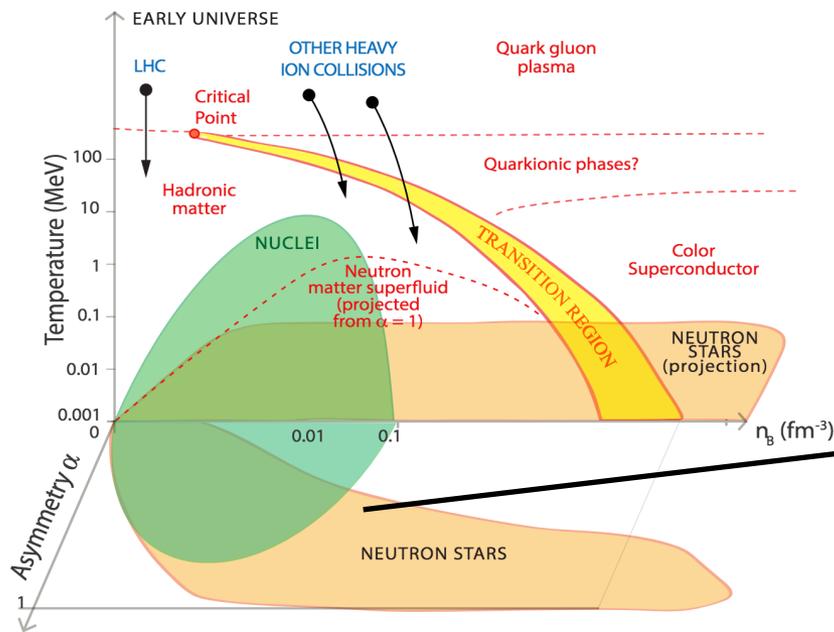


Figure: Lauren Balliet

The phase diagram of nuclear matter to the EOS



Watts et al arxiv:1501.00042

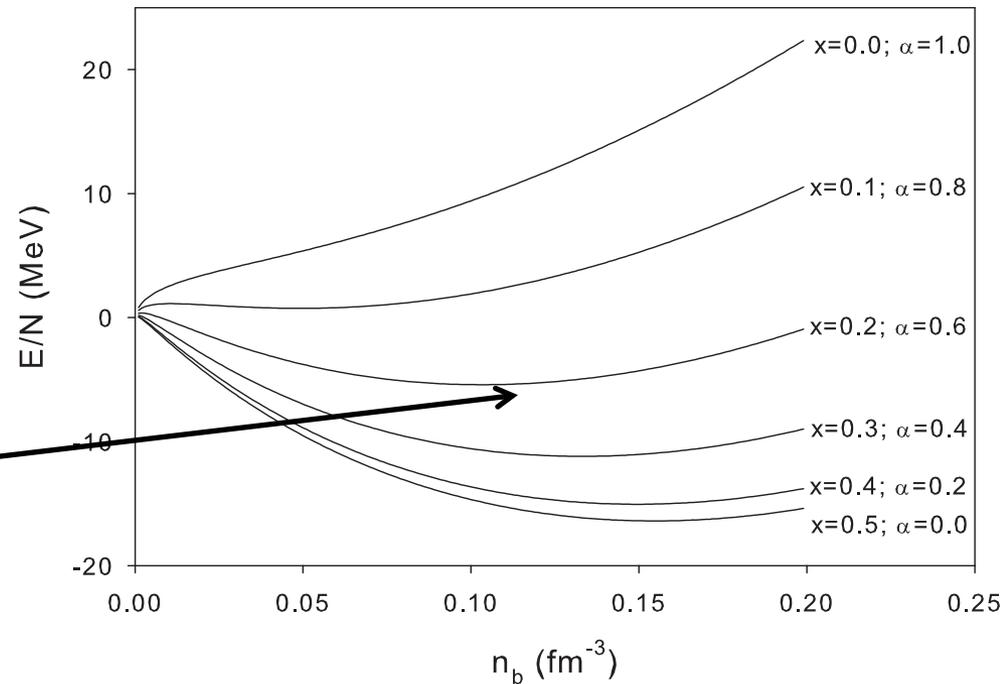
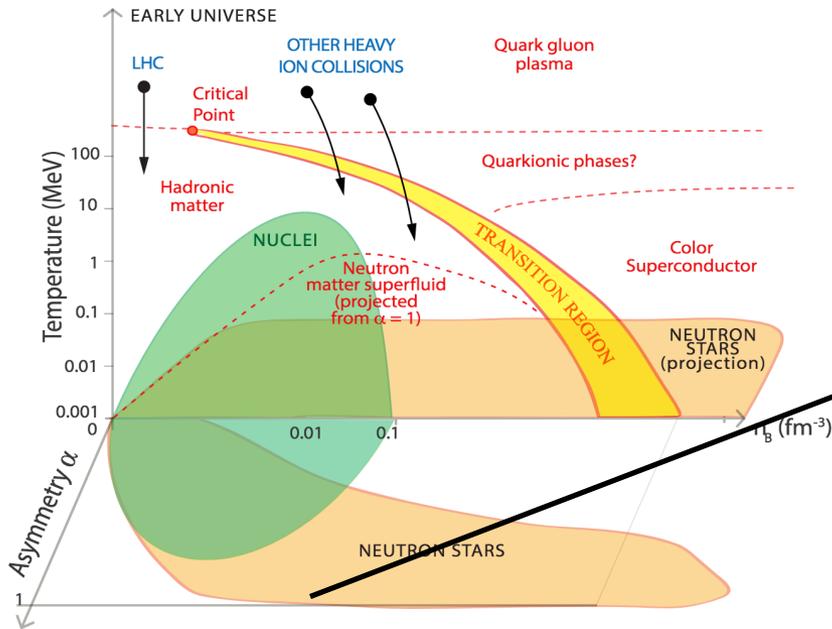


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The phase diagram of nuclear matter to the EOS



Watts et al arxiv:1501.00042

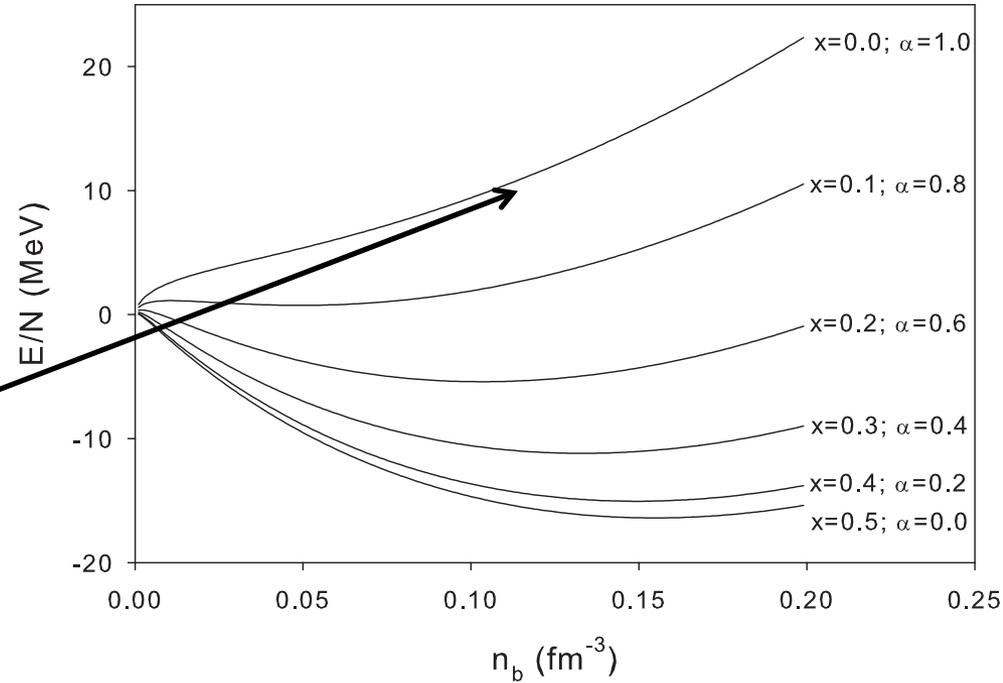
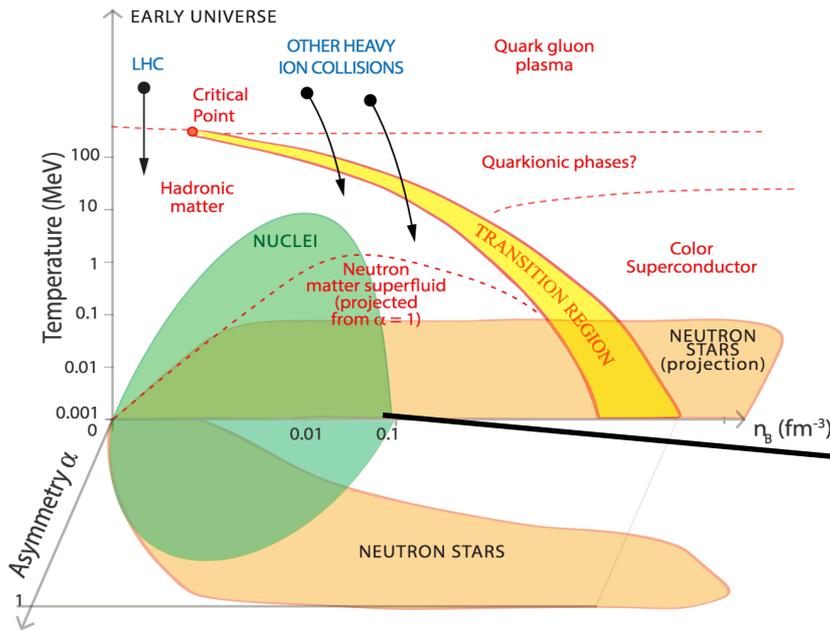


Figure: Lauren Balliet

The phase diagram of nuclear matter to the EOS



Watts et al arxiv:1501.00042

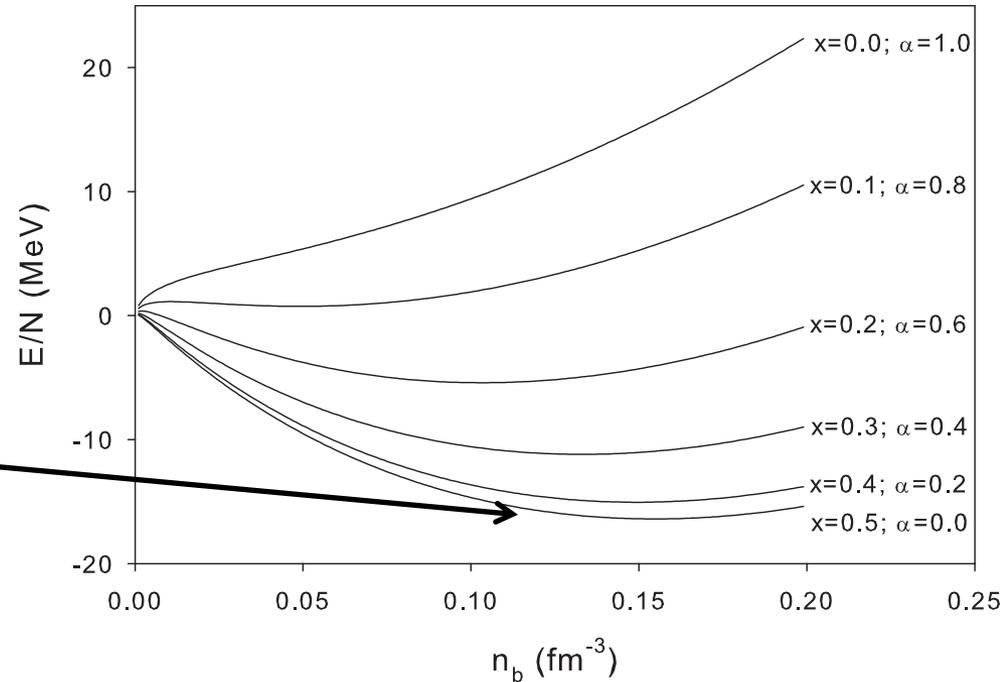
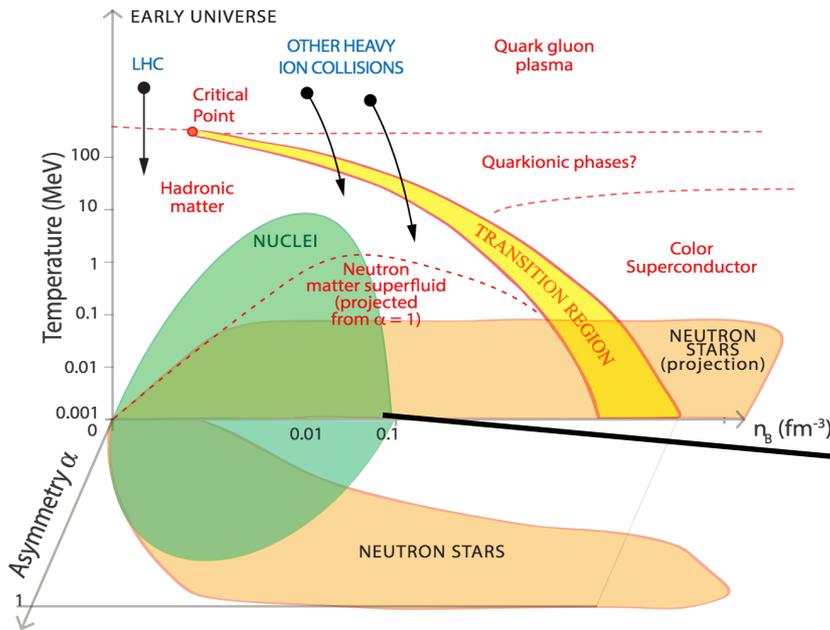


Figure: Lauren Balliet

$$E_0(\rho) = E_0(\rho_0) + \frac{K_0}{2} \left(\frac{\rho - \rho_0}{3\rho_0} \right)^2 + \frac{J_0}{6} \left(\frac{\rho - \rho_0}{3\rho_0} \right)^3,$$

The phase diagram of nuclear matter to the EOS



Watts et al arxiv:1501.00042

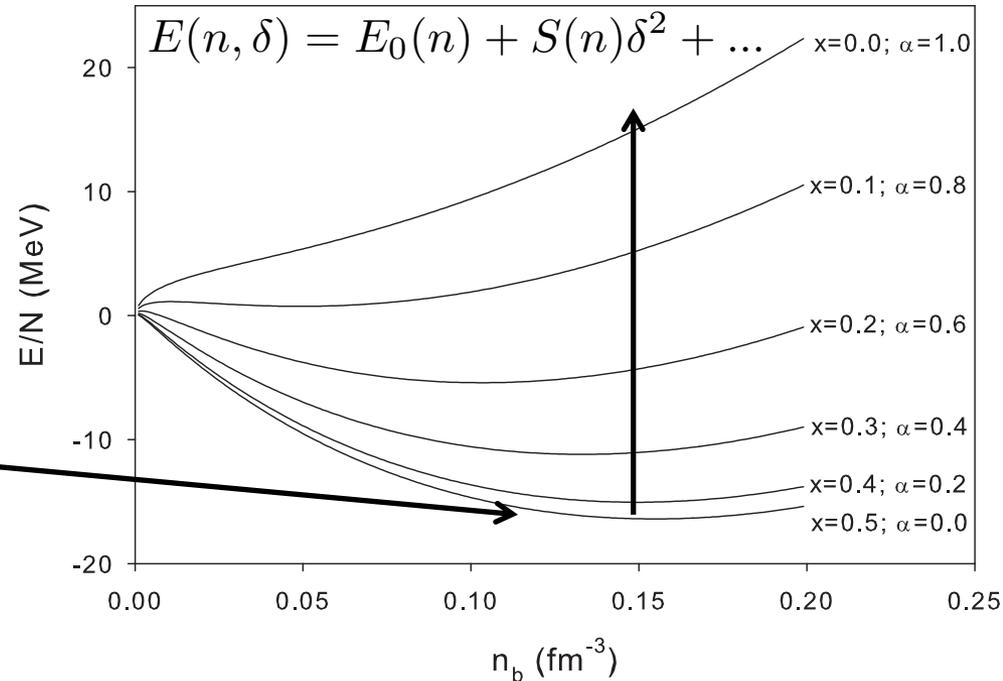
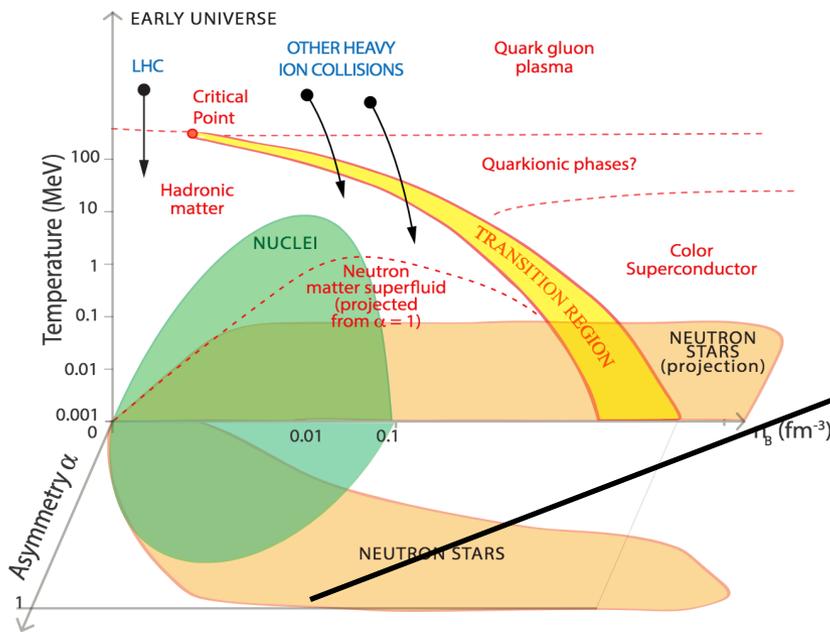


Figure: Lauren Balliet

$$E_0(\rho) = E_0(\rho_0) + \frac{K_0}{2} \left(\frac{\rho - \rho_0}{3\rho_0} \right)^2 + \frac{J_0}{6} \left(\frac{\rho - \rho_0}{3\rho_0} \right)^3,$$

The phase diagram of nuclear matter to the EOS

$$E_{\text{sym}}(\rho) = E_{\text{sym}}(\rho_0) + L\left(\frac{\rho - \rho_0}{3\rho_0}\right) + \frac{K_{\text{sym}}}{2}\left(\frac{\rho - \rho_0}{3\rho_0}\right)^2 + \frac{J_{\text{sym}}}{6}\left(\frac{\rho - \rho_0}{3\rho_0}\right)^3$$



Watts et al arxiv:1501.00042

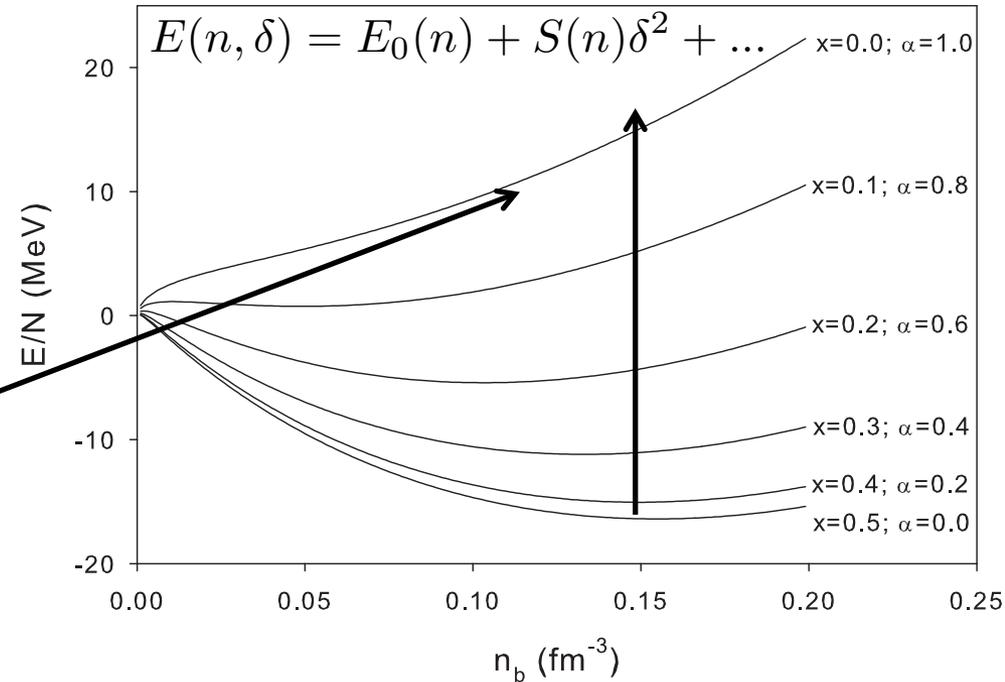


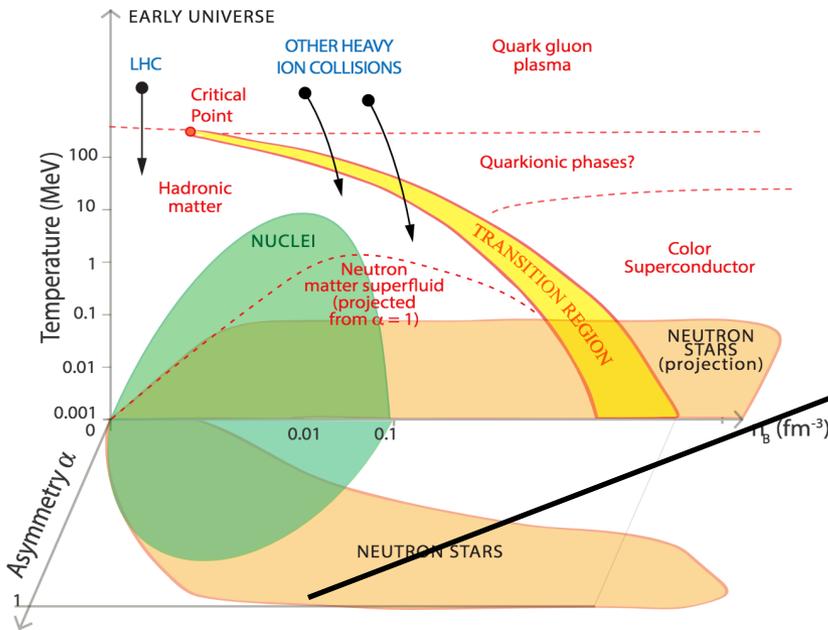
Figure: Lauren Balliet

$$E_0(\rho) = E_0(\rho_0) + \frac{K_0}{2}\left(\frac{\rho - \rho_0}{3\rho_0}\right)^2 + \frac{J_0}{6}\left(\frac{\rho - \rho_0}{3\rho_0}\right)^3,$$

Li, arxiv:2105.04629

The phase diagram of nuclear matter to the EOS

$$E_{\text{sym}}(\rho) = \cancel{E_{\text{sym}}(\rho_0)} + L\left(\frac{\rho - \rho_0}{3\rho_0}\right) + \frac{K_{\text{sym}}}{2}\left(\frac{\rho - \rho_0}{3\rho_0}\right)^2 + \cancel{\frac{J_{\text{sym}}}{6}\left(\frac{\rho - \rho_0}{3\rho_0}\right)^3} + \frac{Q_{\text{sym}}}{6}\left(\frac{\rho - \rho_0}{3\rho_0}\right)^3$$



Watts et al arxiv:1501.00042

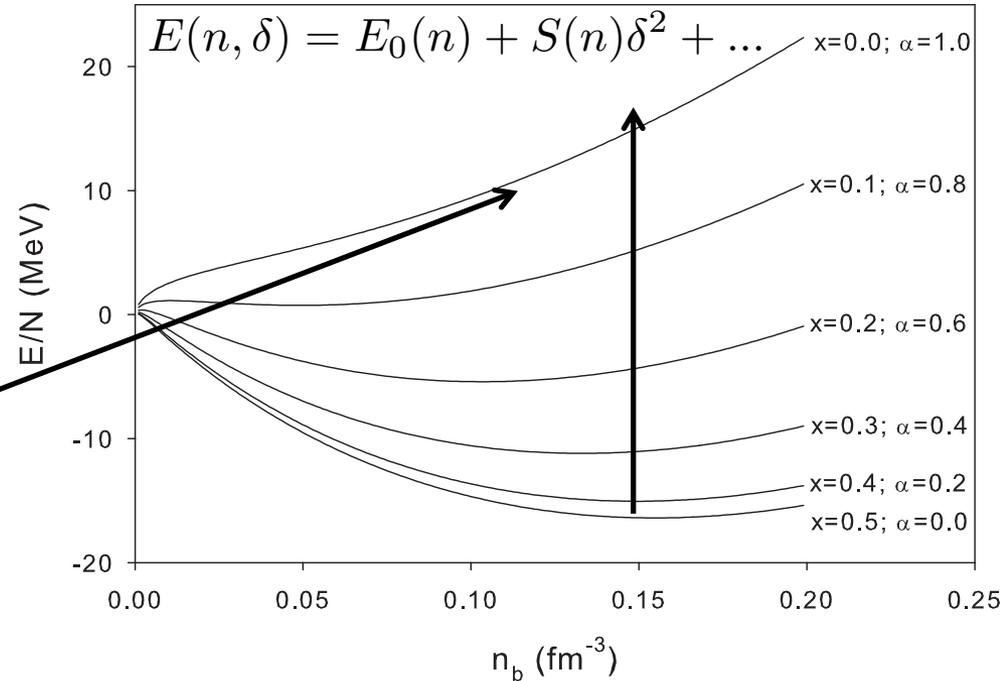
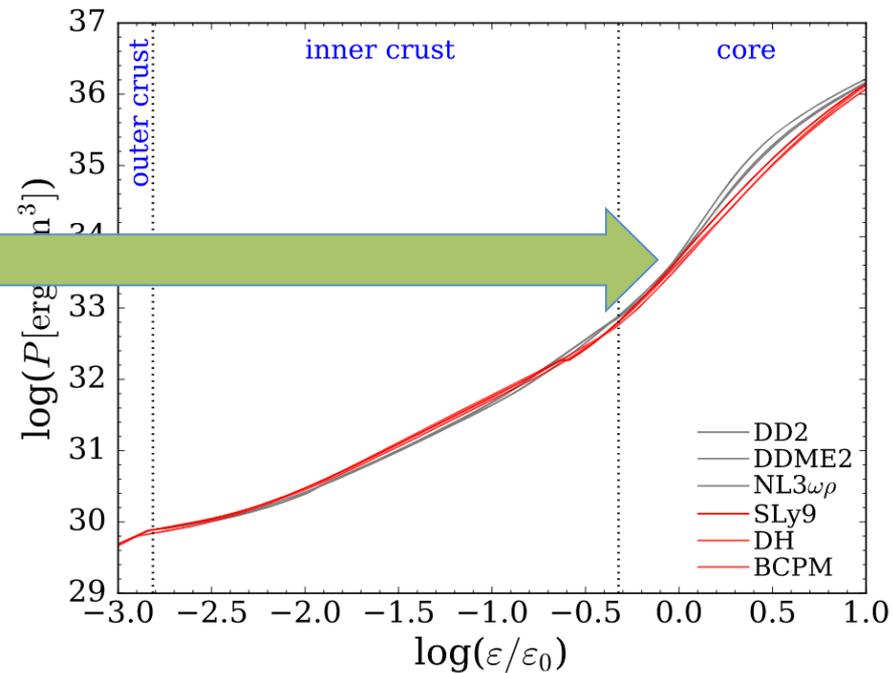
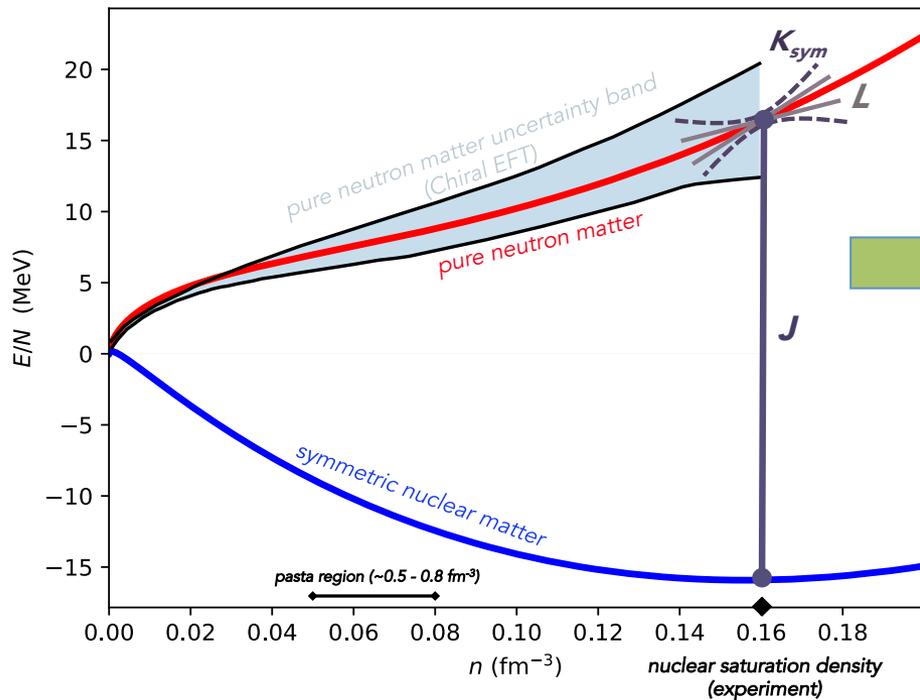


Figure: Lauren Balliet

$$E_0(\rho) = E_0(\rho_0) + \frac{K_0}{2}\left(\frac{\rho - \rho_0}{3\rho_0}\right)^2 + \frac{J_0}{6}\left(\frac{\rho - \rho_0}{3\rho_0}\right)^3,$$

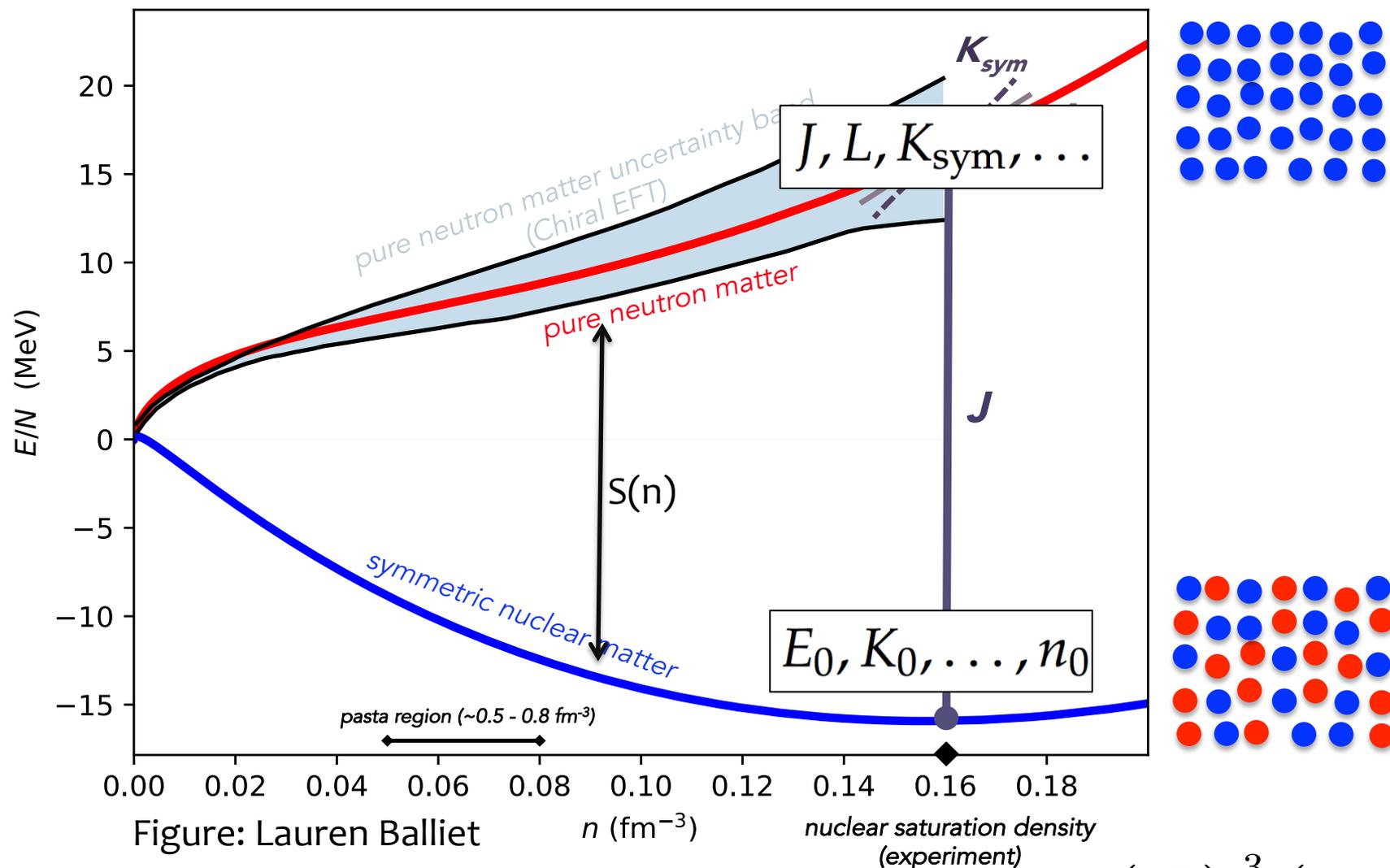
Li, arxiv:2105.04629



$$P_{\text{NS}}(n_0) \approx \frac{n_0}{3} L + 0.048 n_0 \left(\frac{J}{30} \right)^3 \left(J - \frac{4}{3} L \right)$$

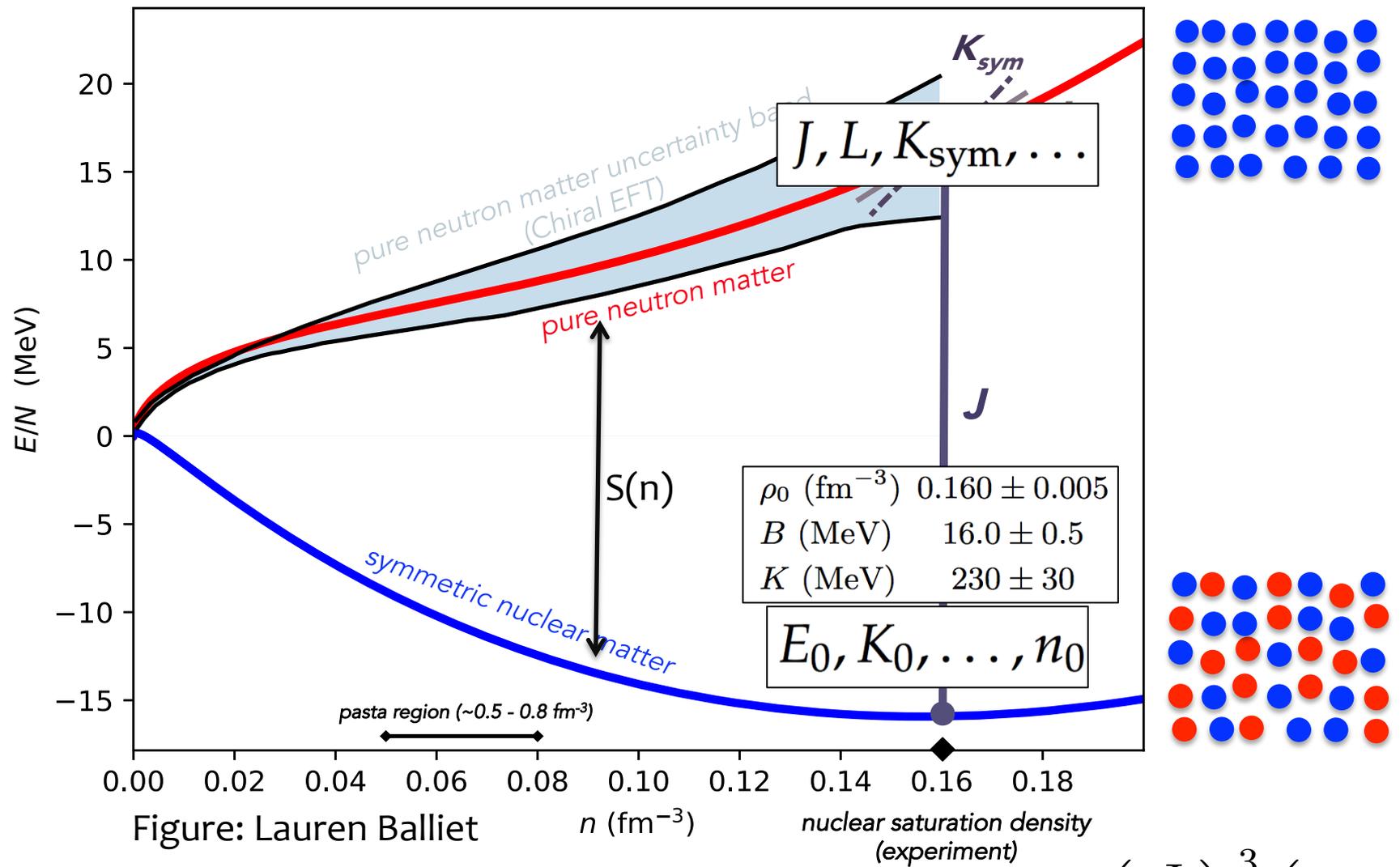
Lattimer, Prakash; astro-ph/0002232

The nuclear symmetry energy: parameterizing our ignorance in a physically meaningful way



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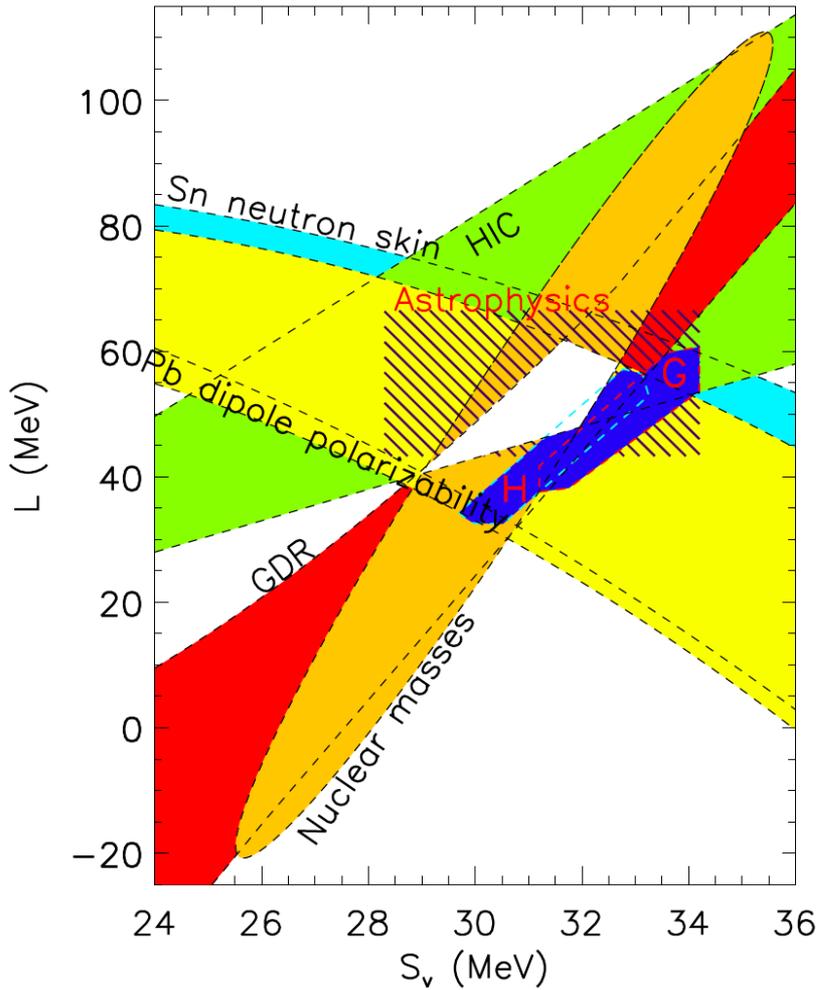
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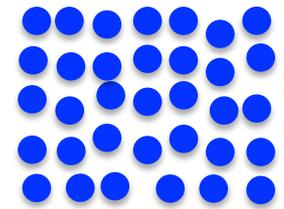
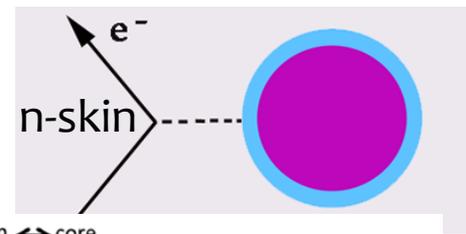
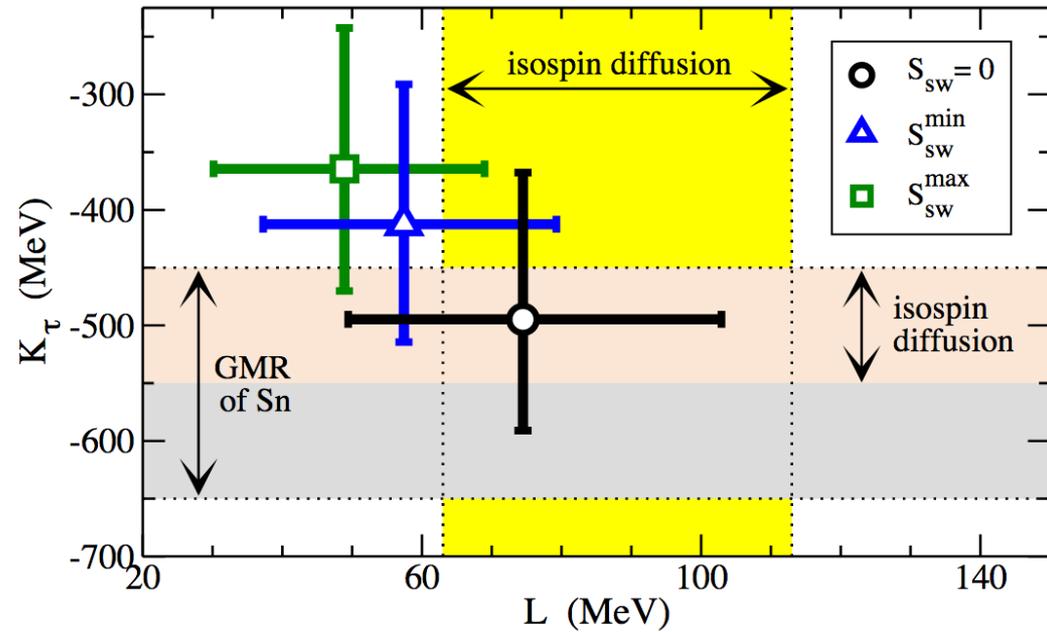
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Symmetry energy constraints

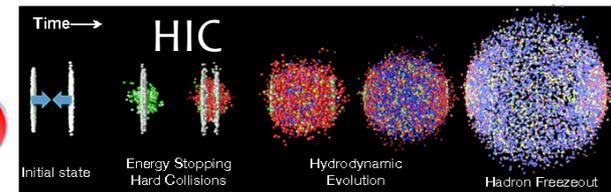
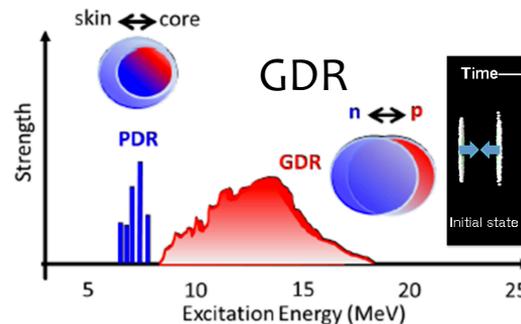
Centelles et al, arxiv:0806.2886



Lattimer, Lim ApJ771(2013)
Lattimer, Steiner EPJA50 (2013)



PNM



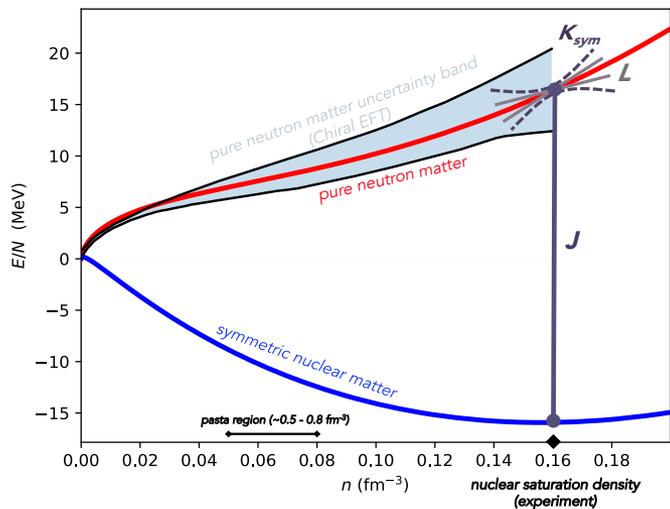
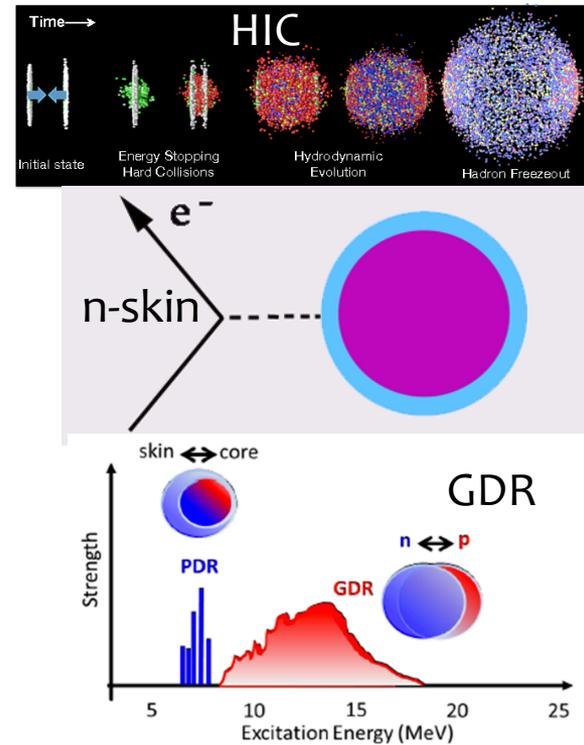
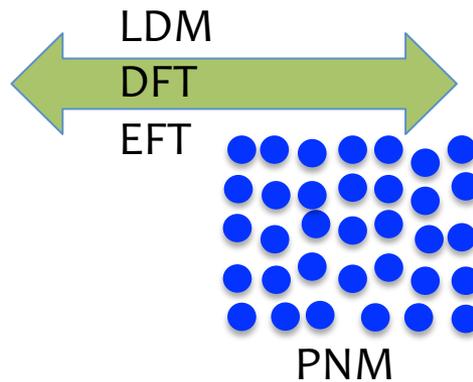
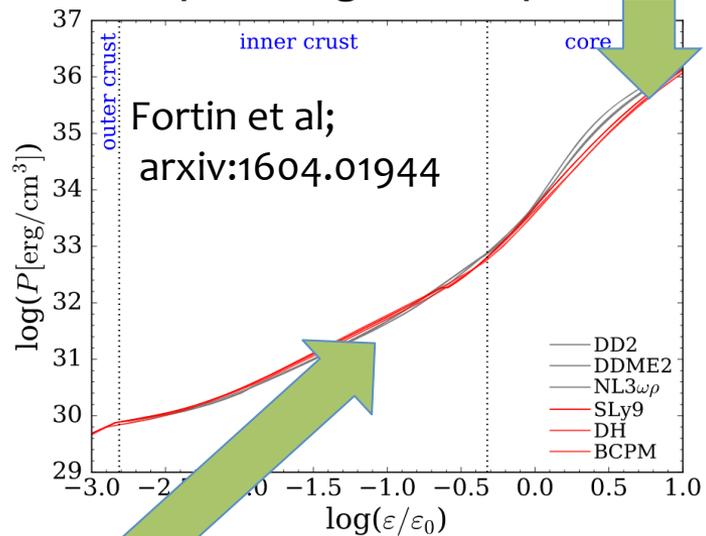


Figure: Lauren Balliet



np(Yq)
 Polytrope/speed of sound
 Spectral/gaussian processes



CLDM,TF,HF (crust), β -equilibrium matter

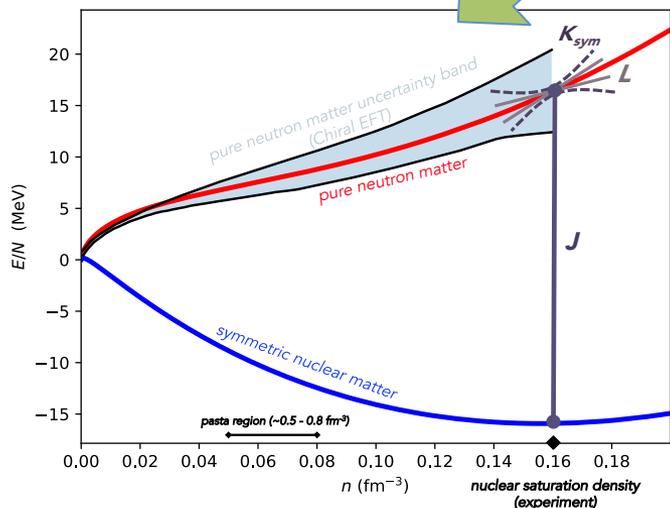
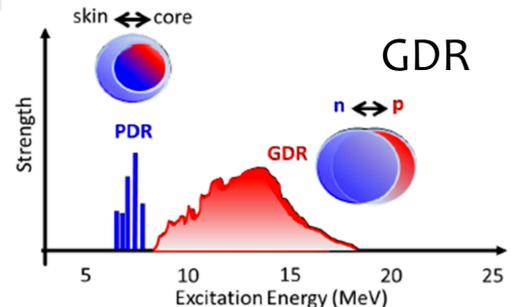
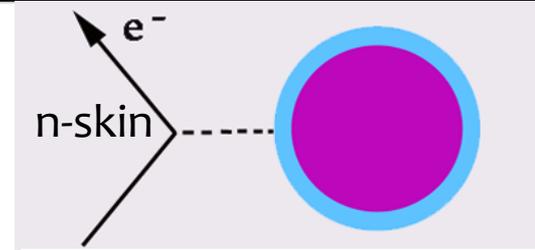
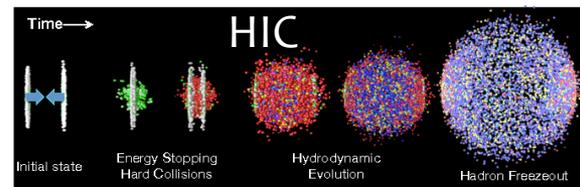
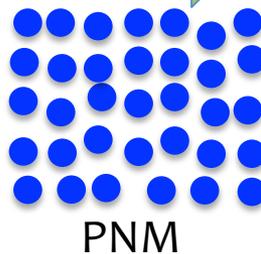
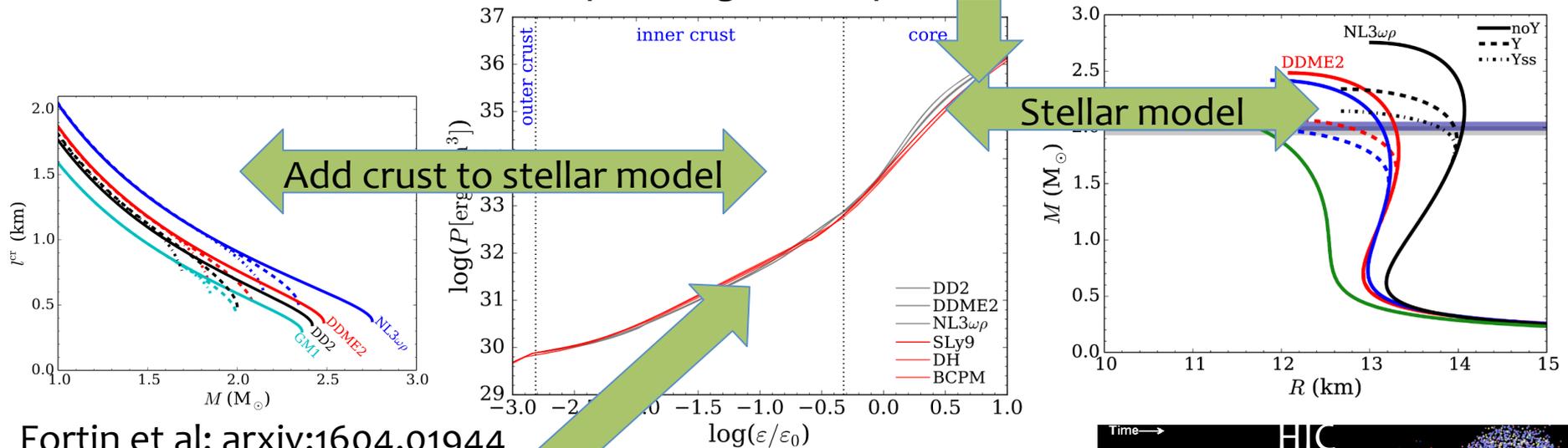


Figure: Lauren Balliet

LDM
 DFT
 EFT



np(Yq)
 Polytrope/speed of sound
 Spectral/gaussian processes

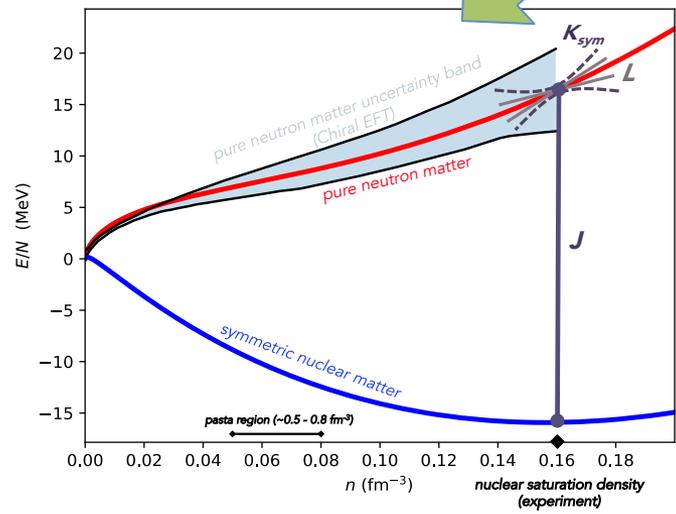


Add crust to stellar model

Stellar model

Fortin et al; arxiv:1604.01944

CLDM, TF, HF (crust), β -equilibrium matter



LDM
 DFT
 EFT

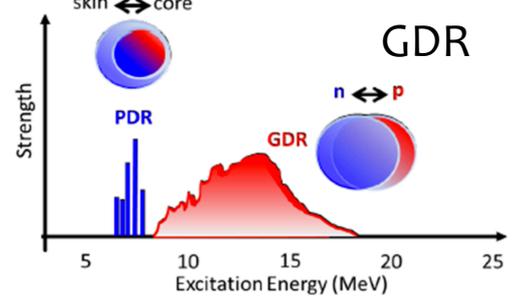
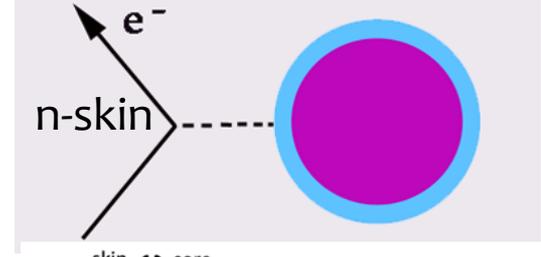
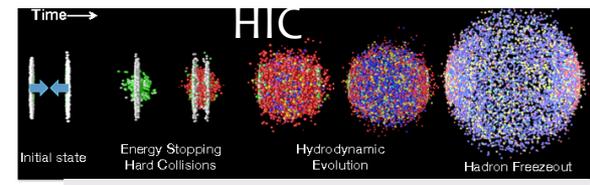
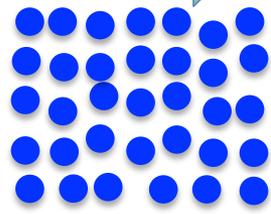
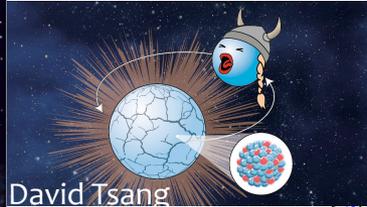
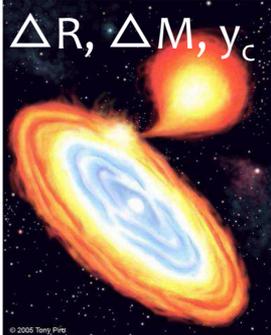


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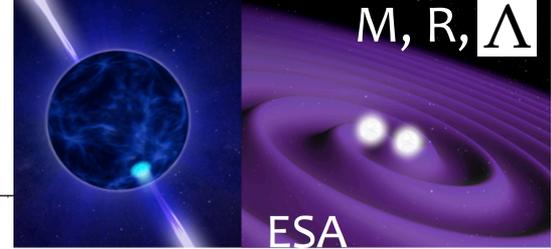
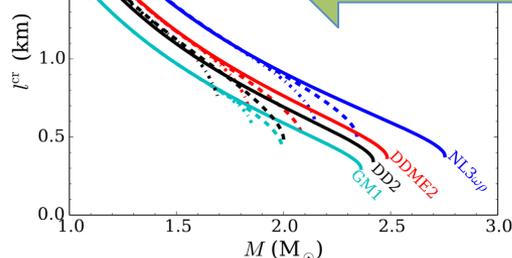
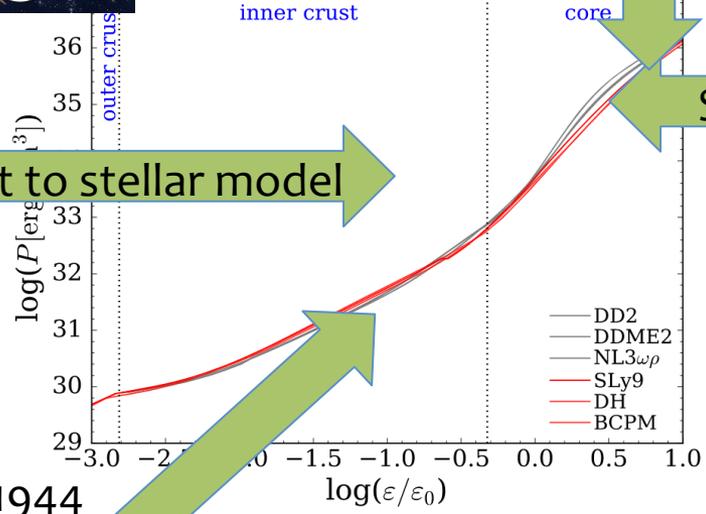


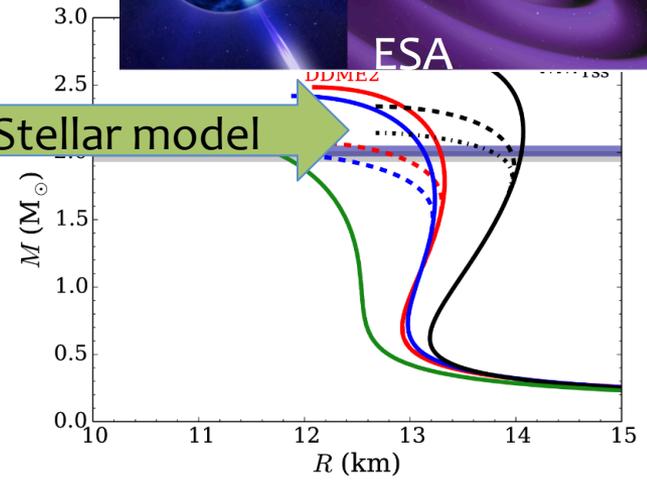
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Add crust to stellar model

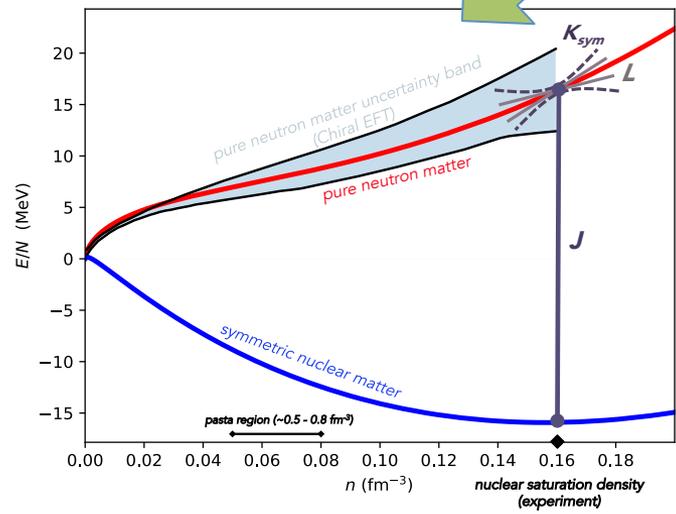


Stellar model



Fortin et al; arxiv:1604.01944

CLDM, TF, HF (crust), β -equilibrium matter



LDM
 DFT
 EFT

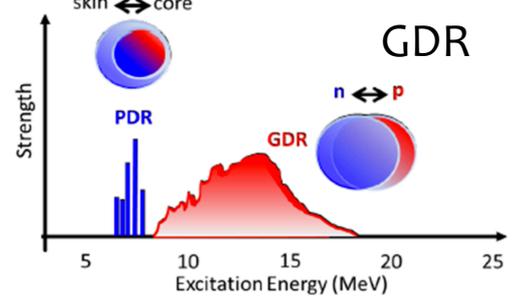
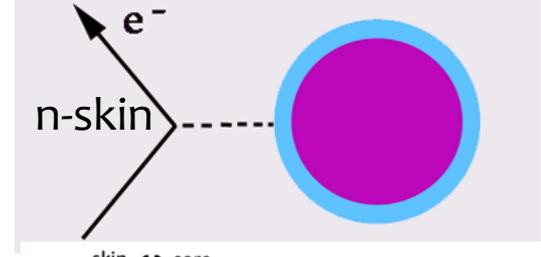
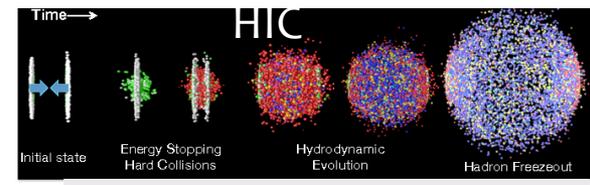
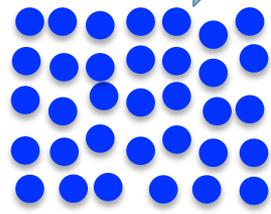
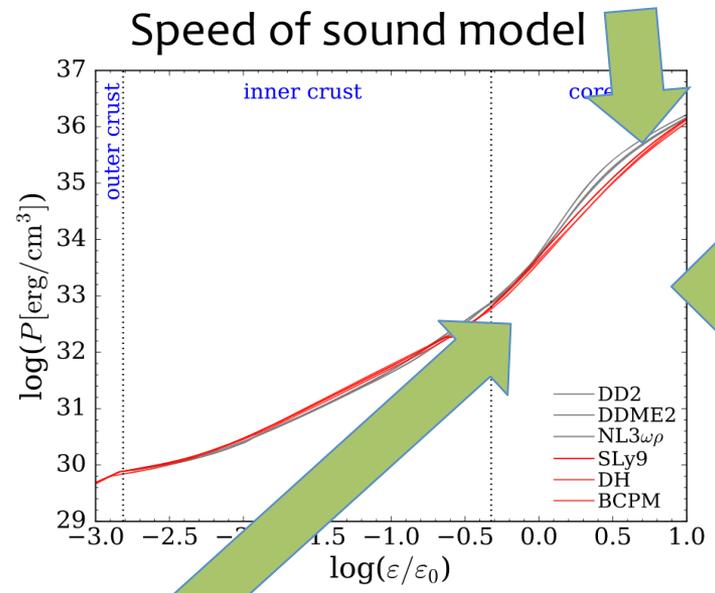
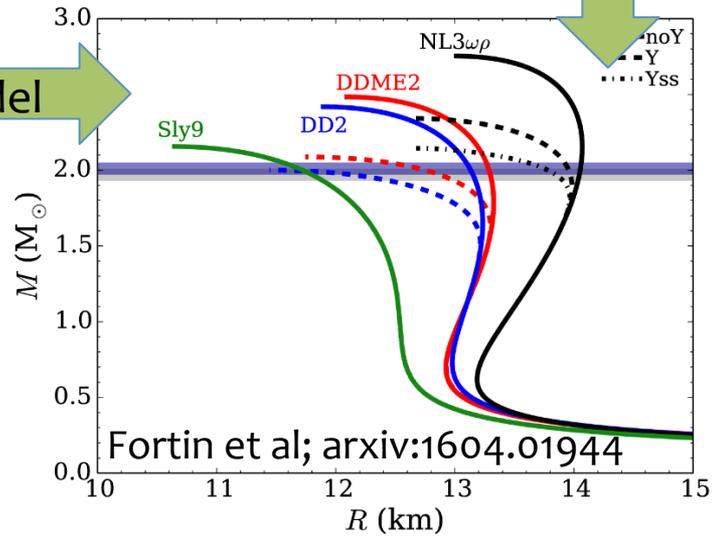


Figure: Lauren Balliet

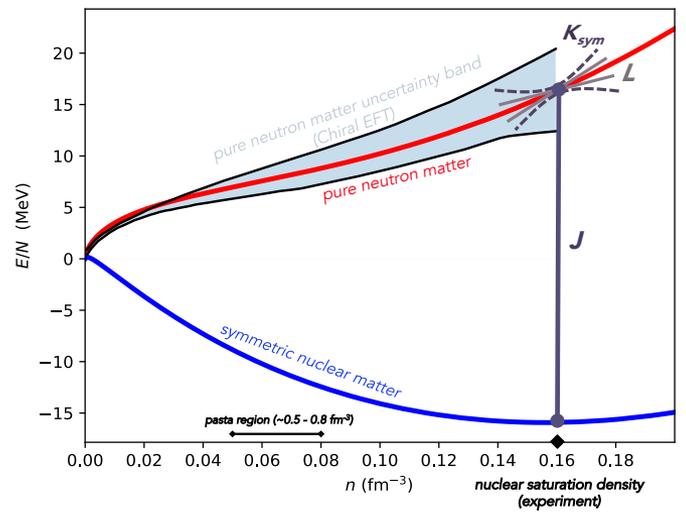
Example: Huth et al; Nature 606, 276 (2022)



Stellar model



BBP crust β -equilibrium matter



←

$$E_{\text{sym}}(\rho) = A \left(\frac{\rho}{\rho_0} \right)^{2/3} + B \left(\frac{\rho}{\rho_0} \right)^\gamma$$

χ EFT

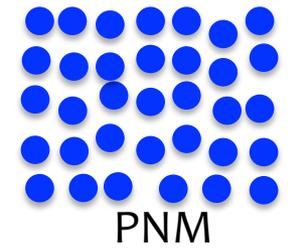
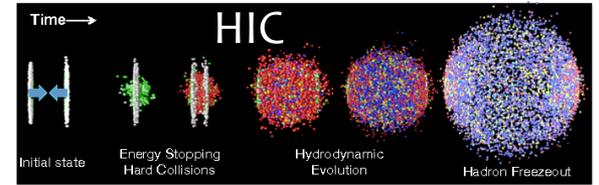
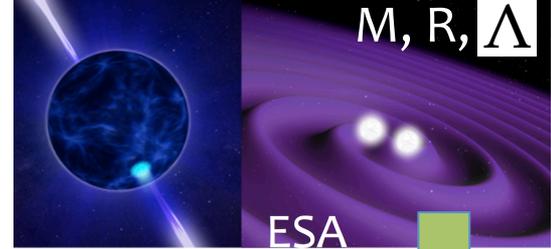
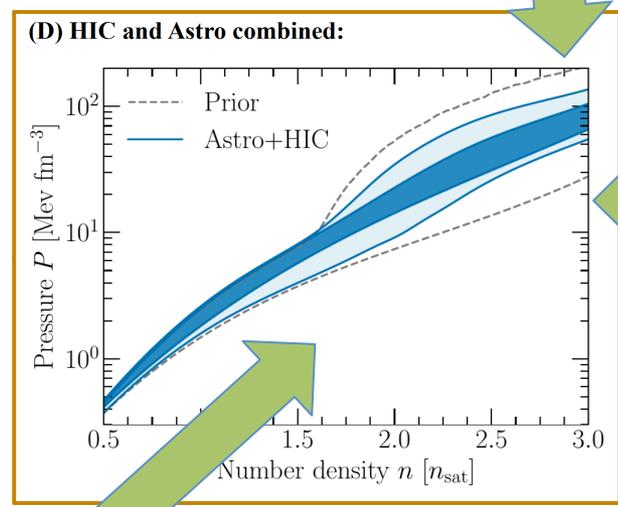


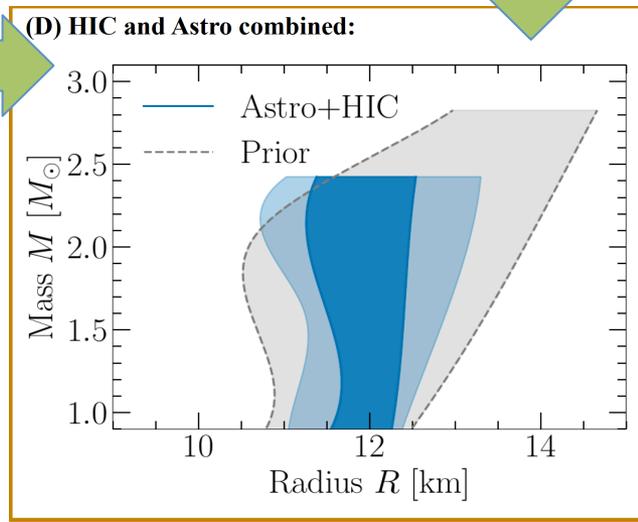
Figure: Lauren Balliet



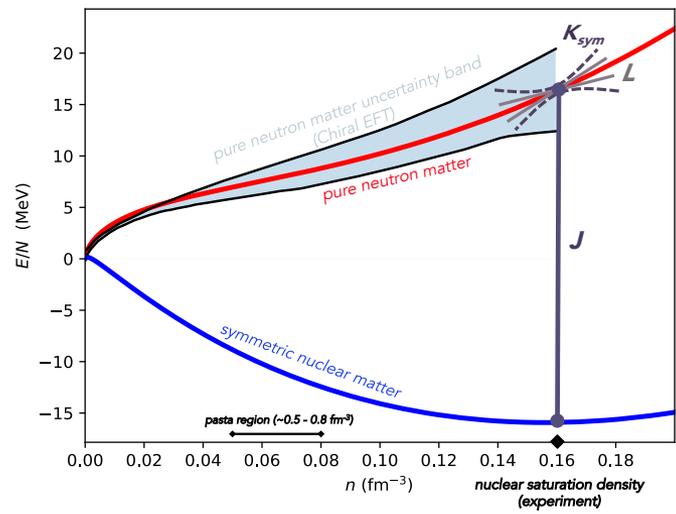
Speed of sound model



Stellar model



BBP crust β -equilibrium matter



$$E_{\text{sym}}(\rho) = A \left(\frac{\rho}{\rho_0} \right)^{2/3} + B \left(\frac{\rho}{\rho_0} \right)^{\gamma}$$

χ EFT

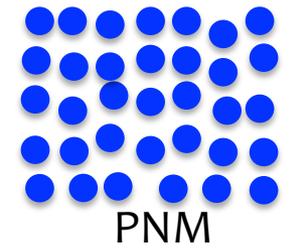
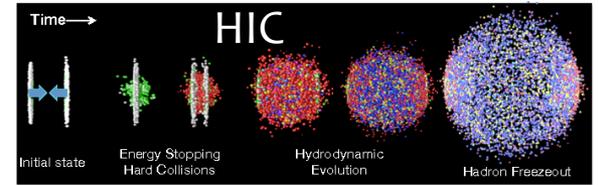


Figure: Lauren Balliet

Let's dig deeper into this strategy

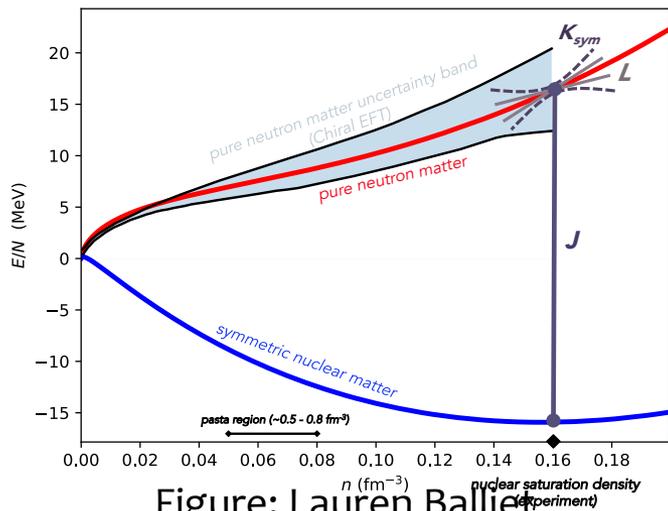
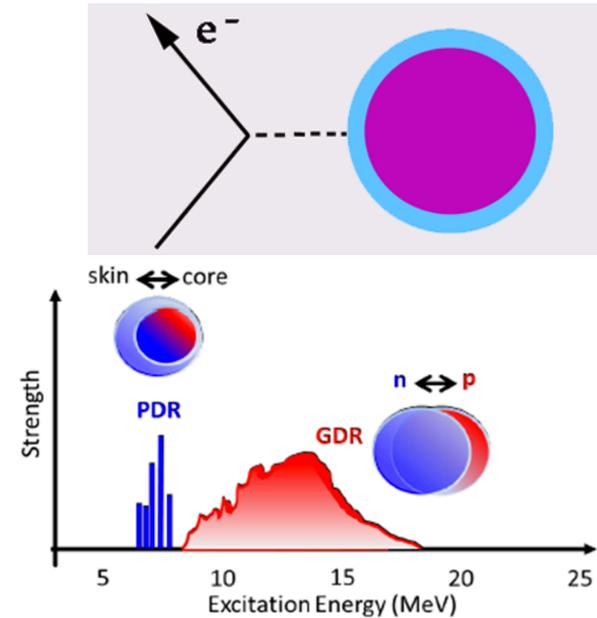
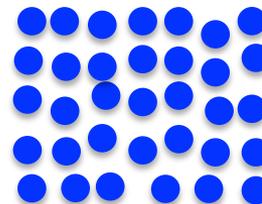
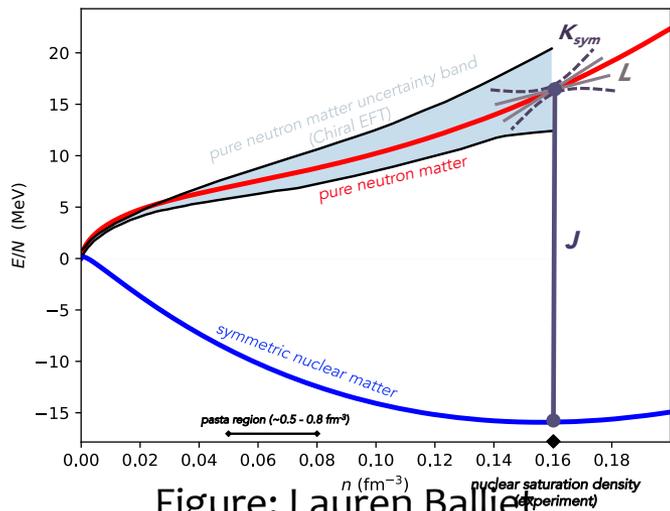


Figure: Lauren Ballet

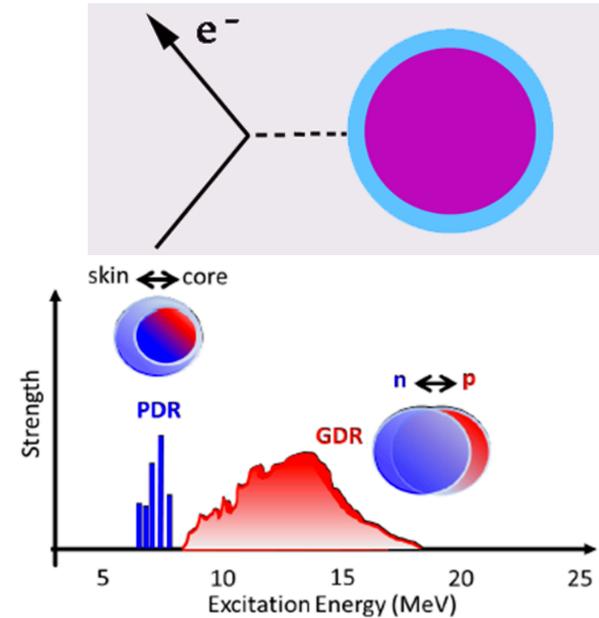
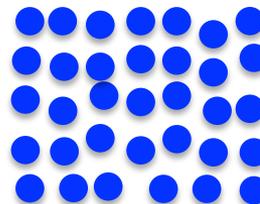
LDM
DFT
EFT

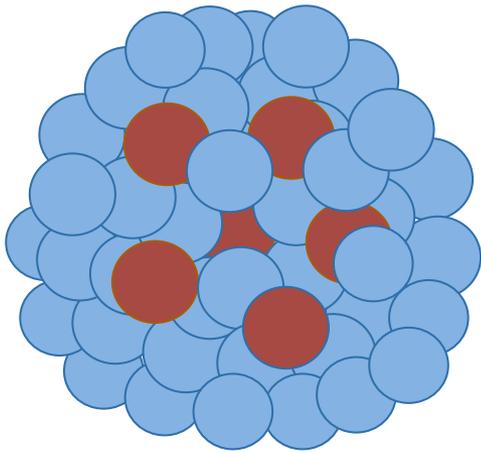


Let's dig deeper into this strategy

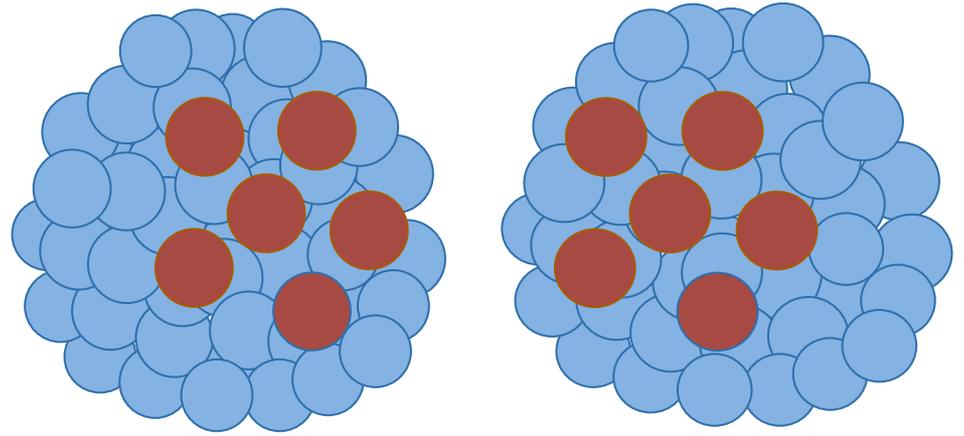


DFT

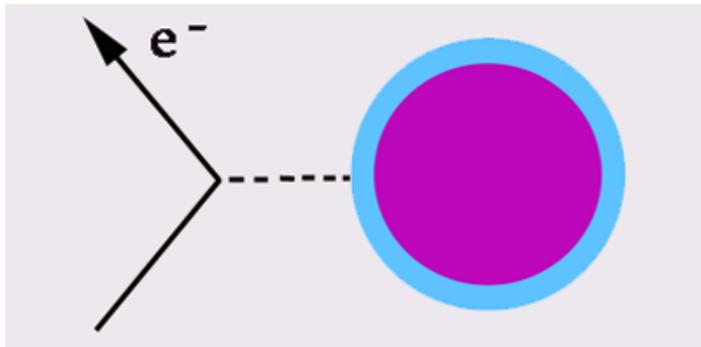




Neutron skin thickness



Collective isovector dipole excitations (PDR, GDR)

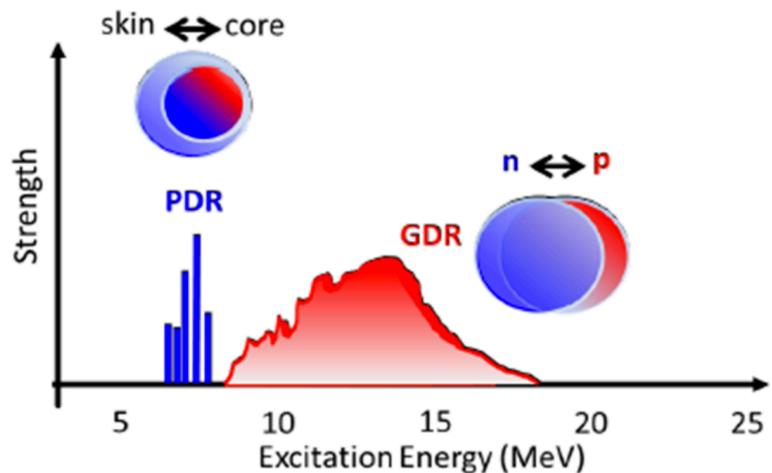


Abrahamyan+, PRL 108, 112592 (2012)

Parity-violating electron scattering

$$\text{PREX } \Delta r_{np}^{208\text{Pb}} = 0.283 \pm 0.071 \text{ fm}$$

$$\text{CREX: } \Delta r_{np}^{48\text{Ca}} = 0.121 \pm 0.035 \text{ fm}$$

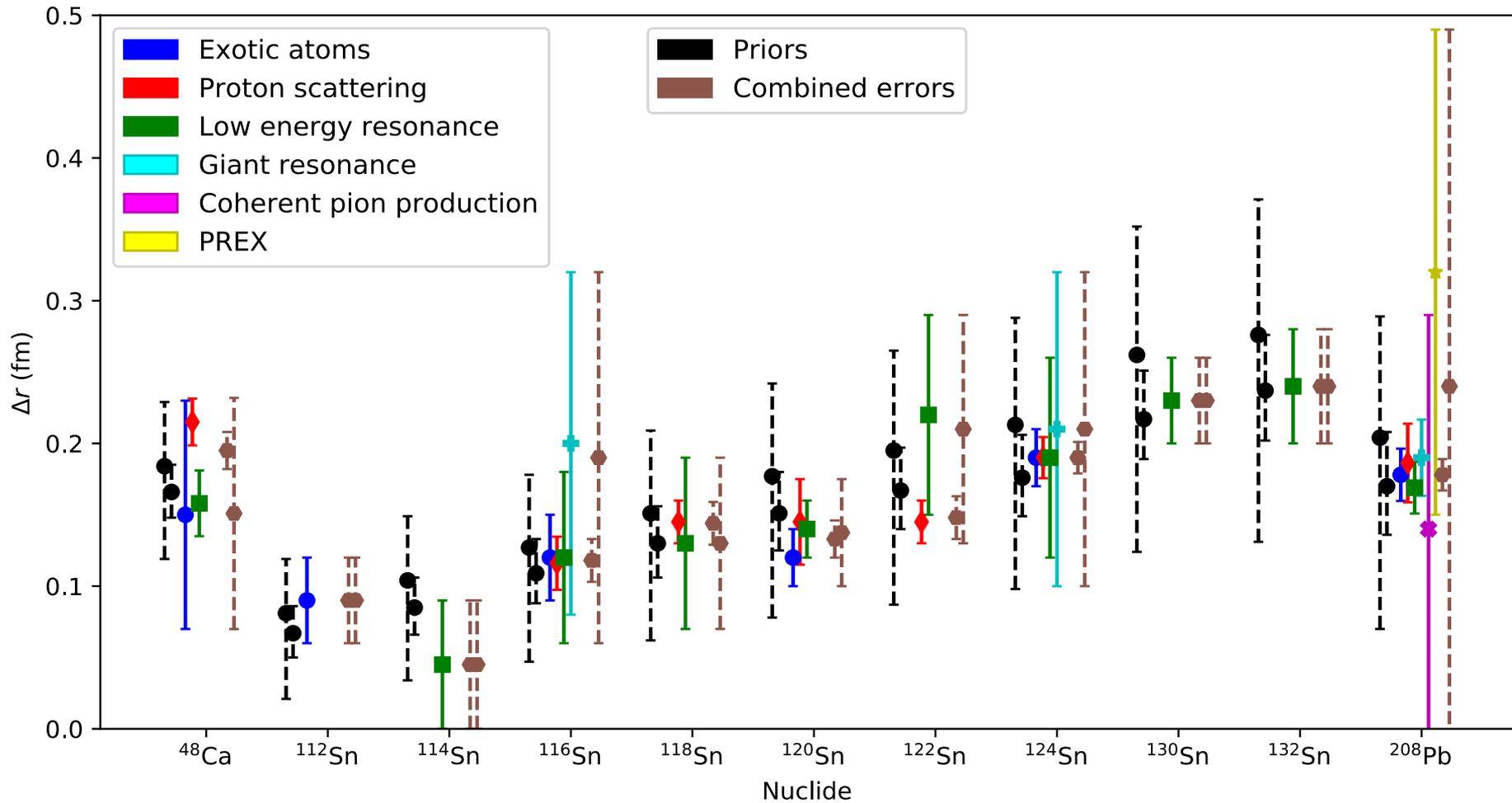


Bracco, Lanza, Tamii, PPNP 106, 360 (2019)

Proton inelastic scattering

$$\text{RCNP: } \alpha_D^{208\text{Pb}} = 20.1 \pm 0.6 \text{ fm}$$

$$\text{RCNP: } \alpha_D^{48\text{Ca}} = 2.07 \pm 0.22 \text{ fm}$$



Newton, Crocombe arxiv:2008.00042

Density Functional Theory (e.g. Skyrme)

$$\mathcal{H}_\delta = \frac{1}{4}t_0\rho^2[(2 + x_0) - (2x_0 + 1)(y_p^2 + y_n^2)]$$

Local interaction

$$\begin{aligned}\mathcal{H}_\rho &= \frac{1}{4}t_3\rho^{2+\alpha_3}[(2 + x_3) - (2x_3 + 1)(y_p^2 + y_n^2)] \\ &+ \frac{1}{4}t_4\rho^{2+\alpha_4}[(2 + x_4) - (2x_4 + 1)(y_p^2 + y_n^2)]\end{aligned}$$

Density dependent

$$\begin{aligned}\mathcal{H}_{\text{eff}} &= \frac{1}{8}\rho[t_1(2 + x_1) + t_2(2 + x_2)]\tau \\ &+ \frac{1}{8}\rho[t_1(2x_1 + 1) + t_2(2x_2 + 1)](\tau_p y_p + \tau_n y_n)\end{aligned}$$

3 body

$$\begin{aligned}\mathcal{H}_{\text{grad}} &= \frac{1}{32}(\nabla\rho)^2[3t_1(2 + x_1) - t_2(2 + x_2)] \\ &- \frac{1}{32}[3t_1(2x_1 + 1) + t_2(2x_2 + 1)][(\nabla\rho_p)^2 + (\nabla\rho_n)^2]\end{aligned}$$

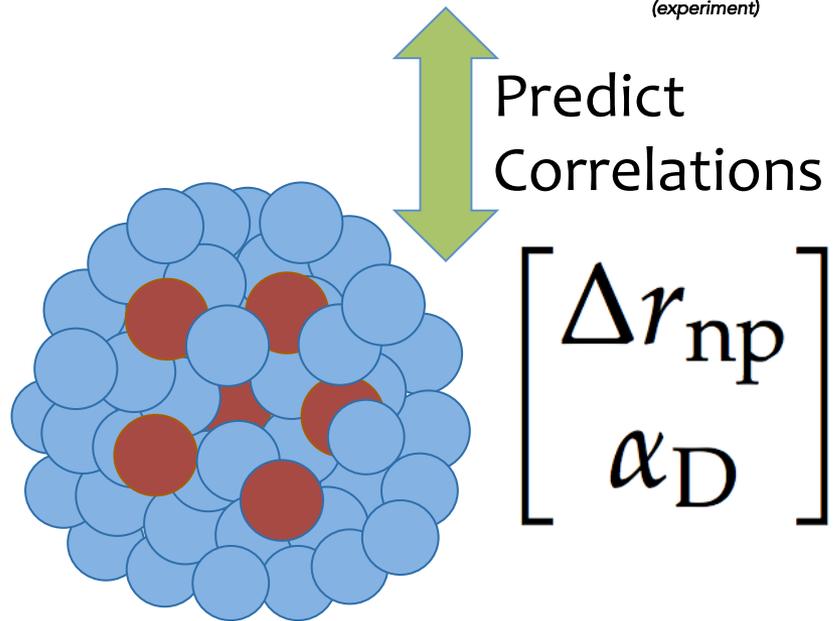
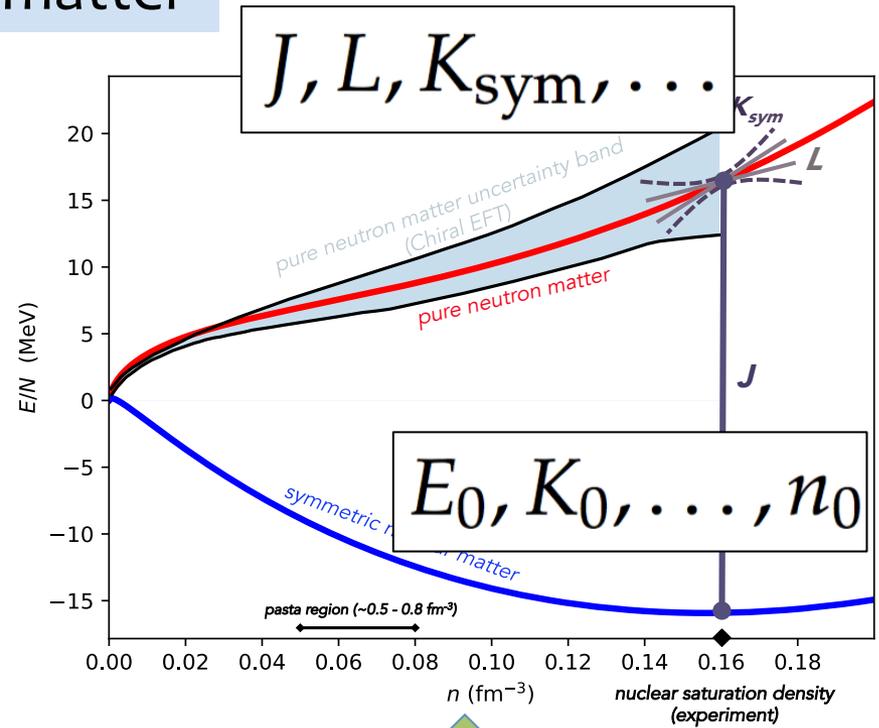
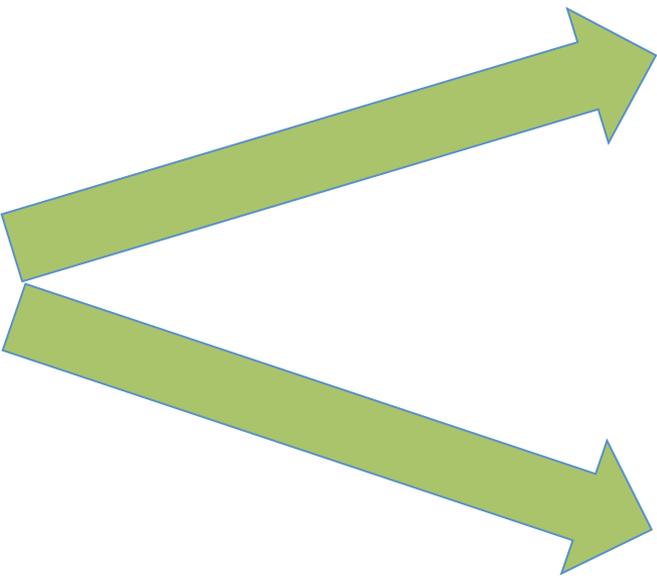
Gradient...

Used in a variational principle on total energy leads to coupled Schrödinger-like equations for the wavefunctions.
Solutions converge to ground state (Hohenberg-Kohn theorem)

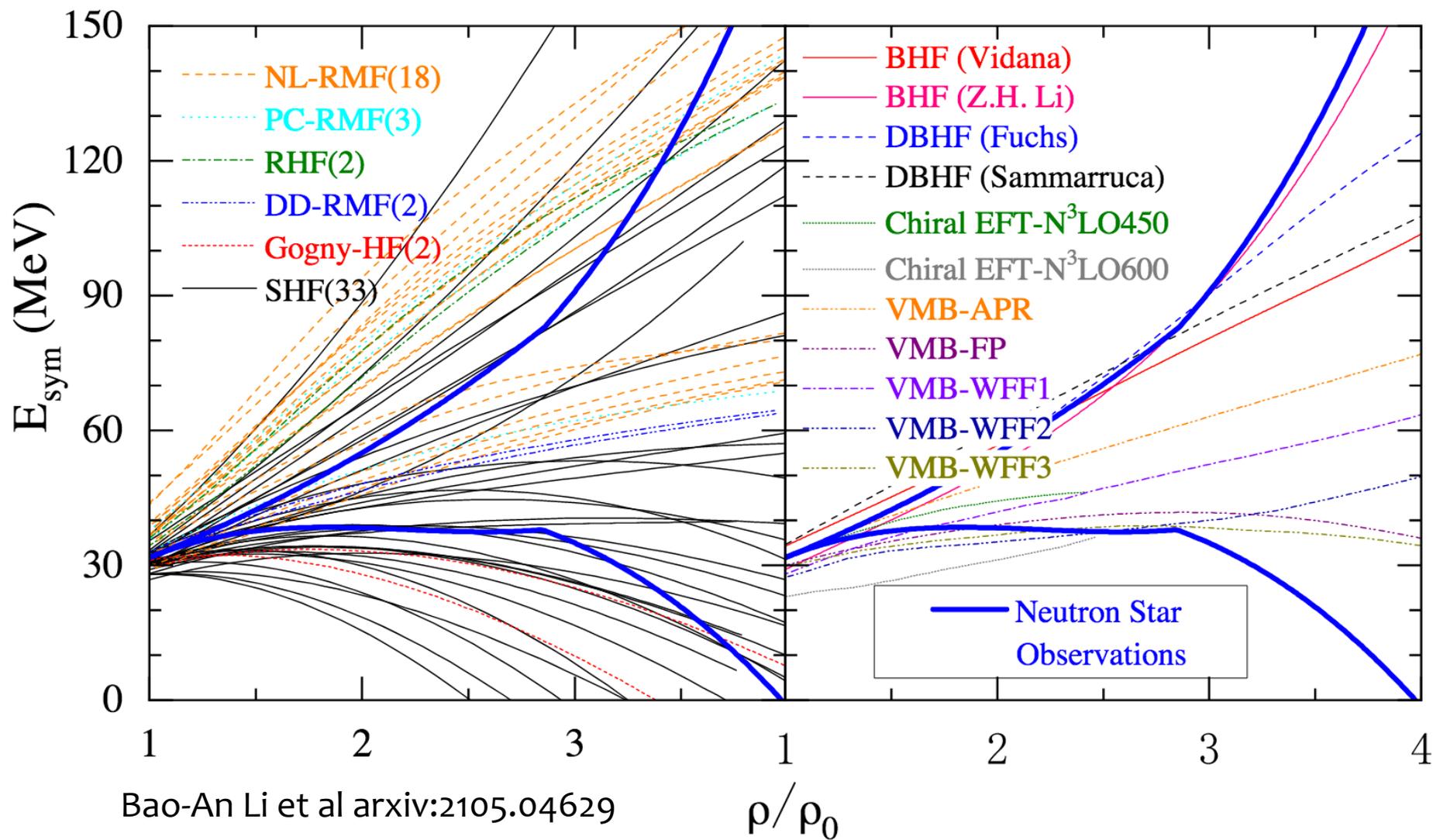
From models to nuclei and nuclear matter

Parameterization
 Based on fits to
 masses, radii,
 s.p. levels, fission barriers...

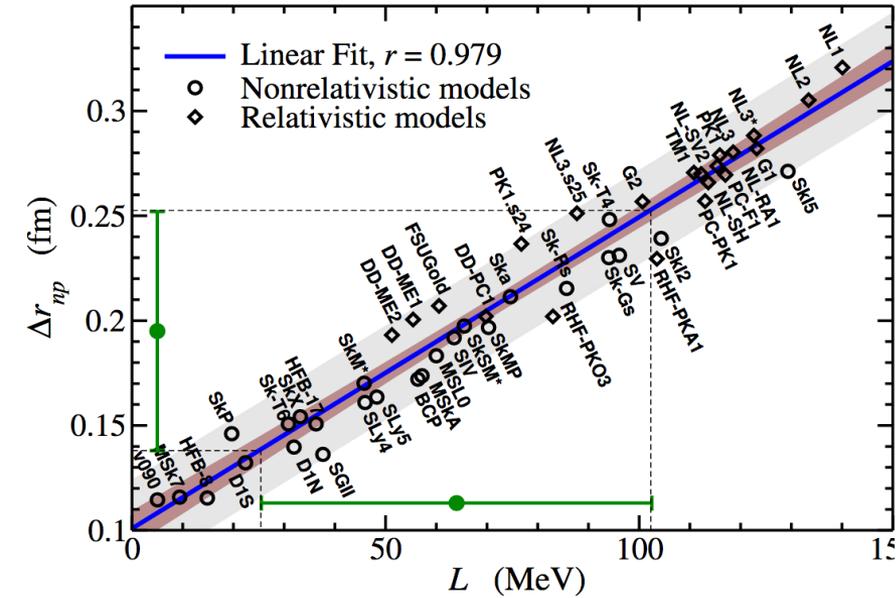
- t_0
- \vdots
- t_4
- x_0
- \vdots
- x_4
- α_3
- α_4



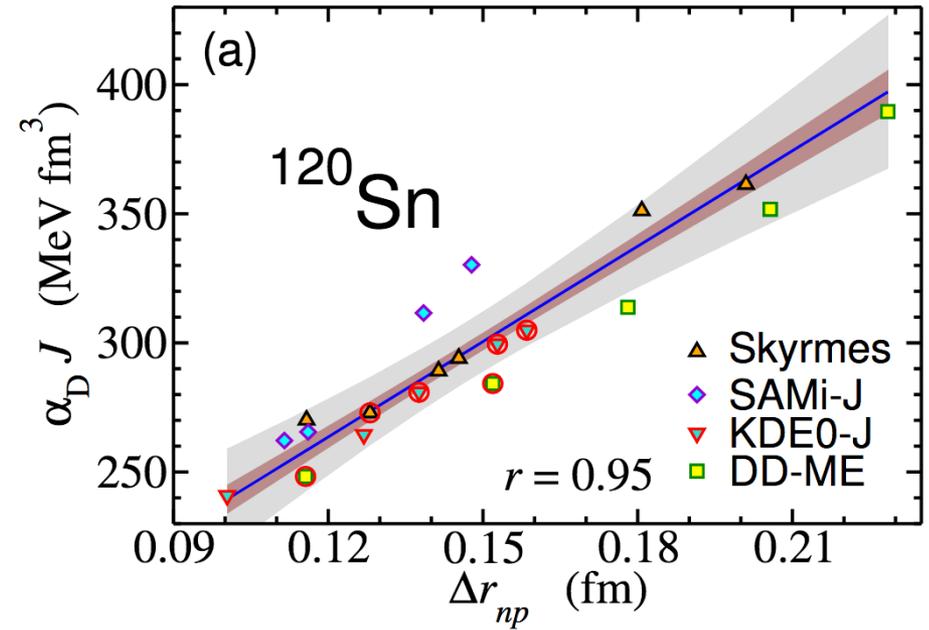
Wide range of symmetry energy dependence



Correlations are revealed between nuclear matter parameters and nuclear properties



Roca-Maza et al, arxiv:1103.1762



Roca-Maza et al, arxiv: 1510.01874

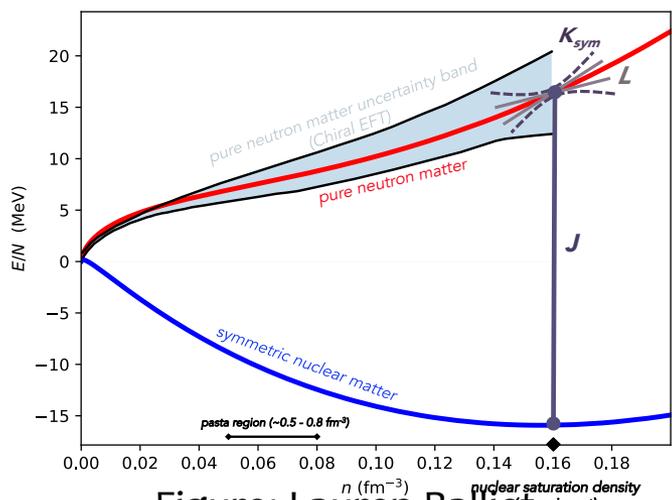
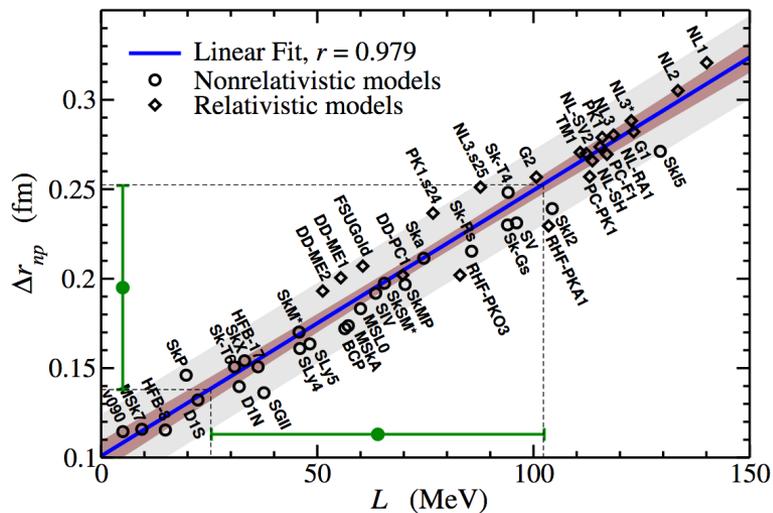
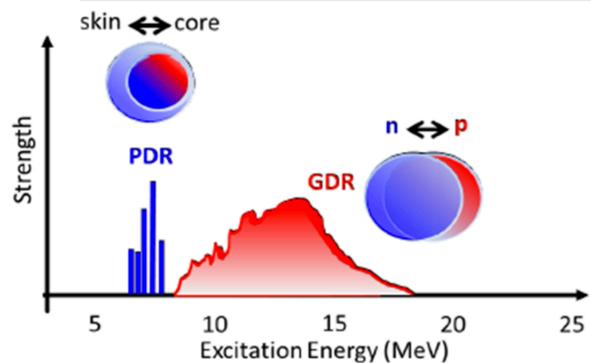
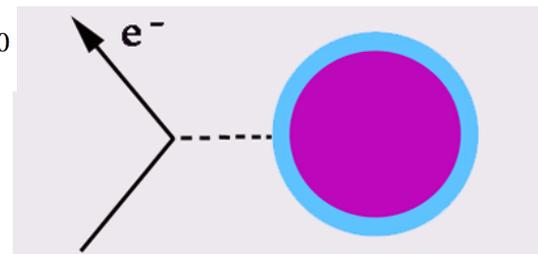
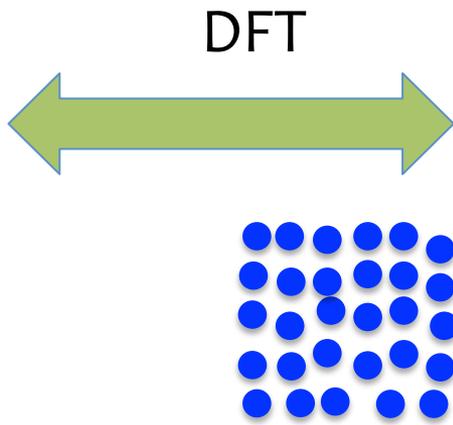
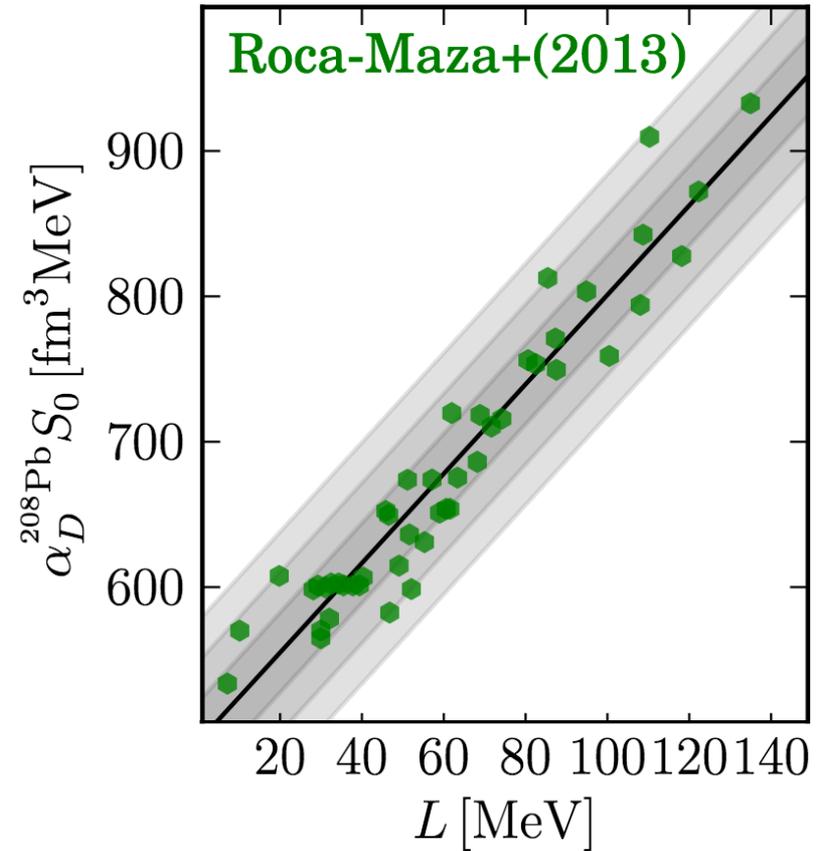
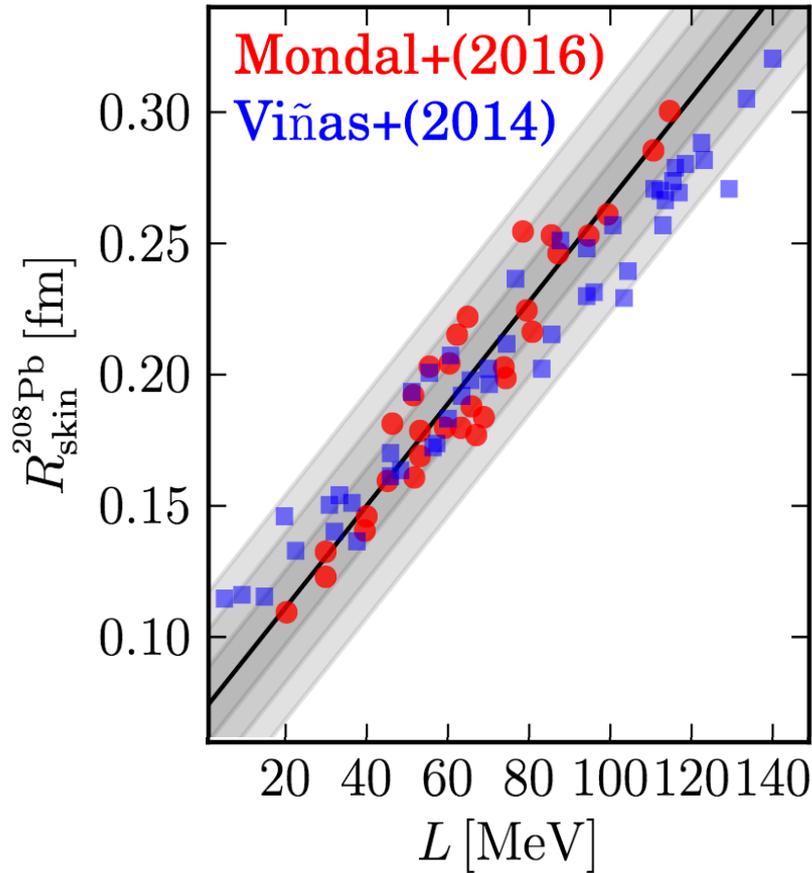


Figure: Lauren Balliet

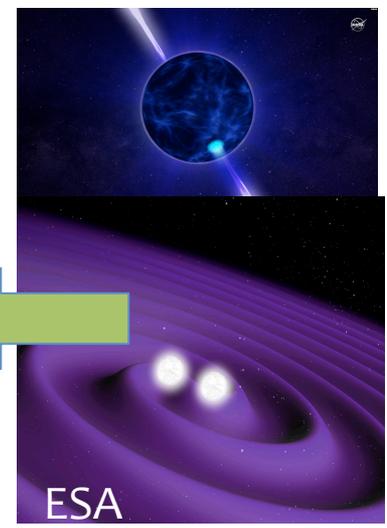
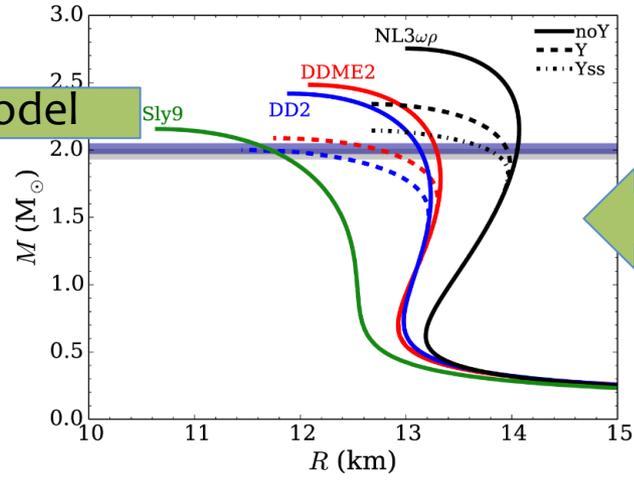
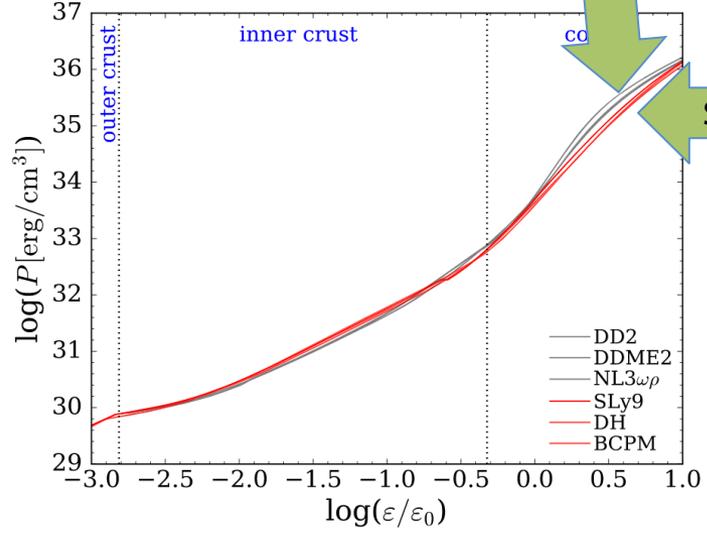


Correlations are revealed between nuclear matter parameters and nuclear properties

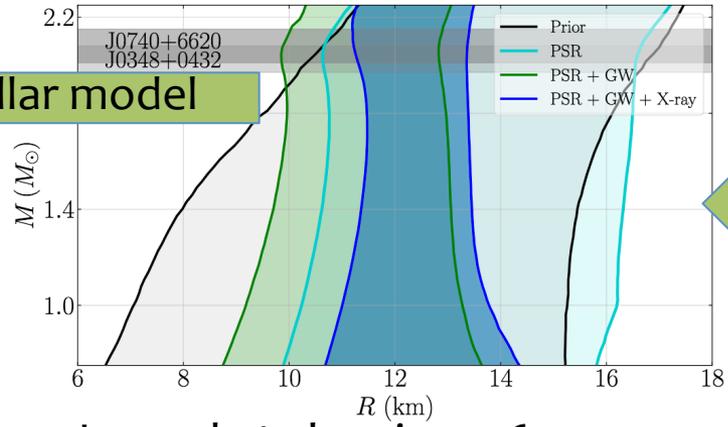
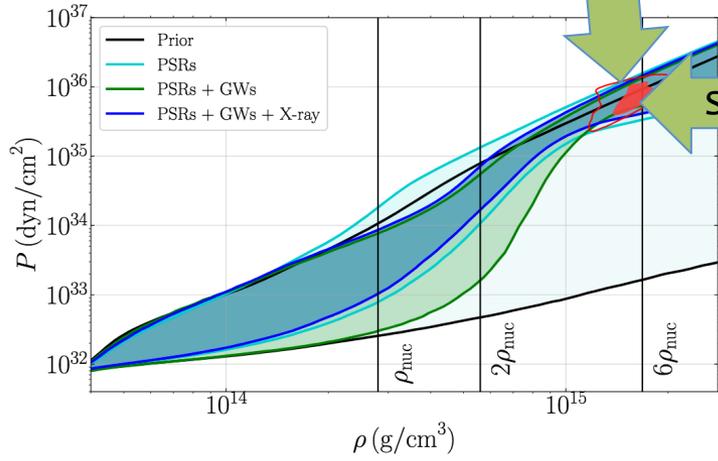


Essick et al: 2107.05528

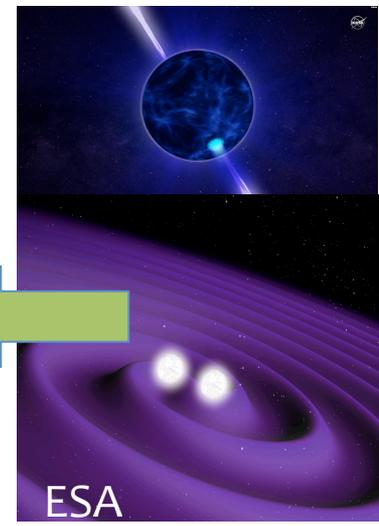
Gaussian processes



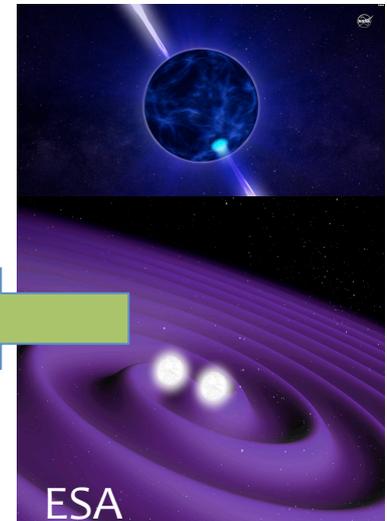
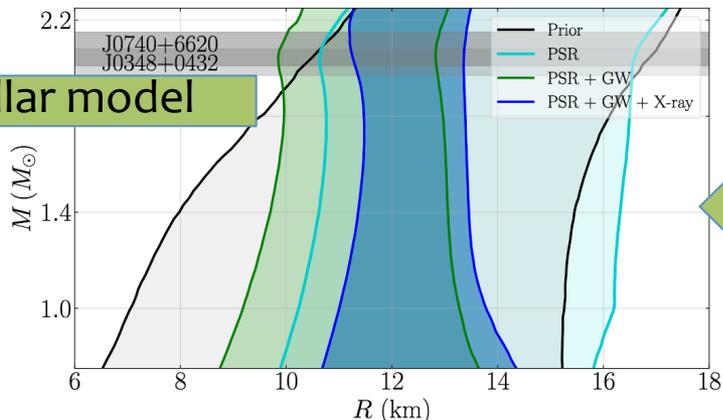
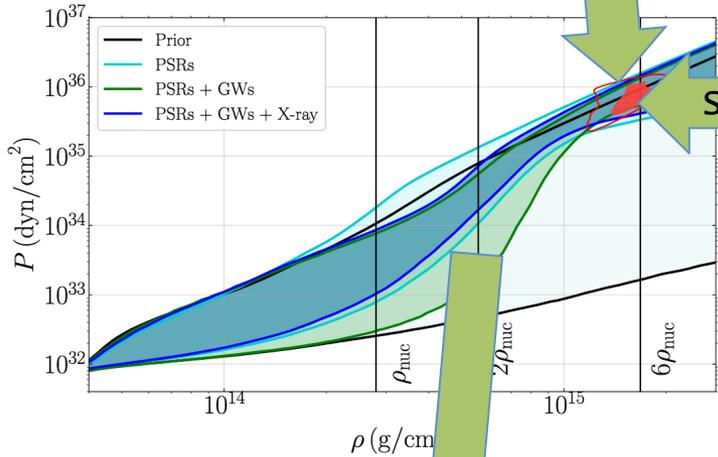
Gaussian processes



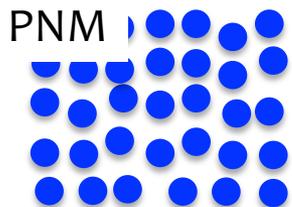
Legred et al arxiv:2106.05313



Gaussian processes



Legred et al arxiv:2106.05313
Essick et al: 2107.05528



BBP; Reconstruct Nuclear matter

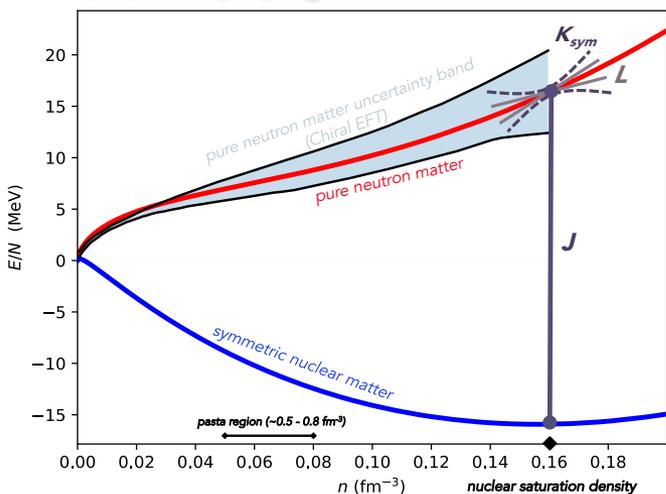
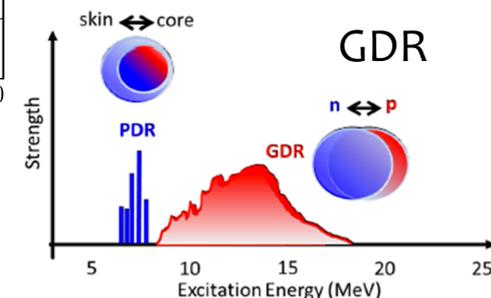
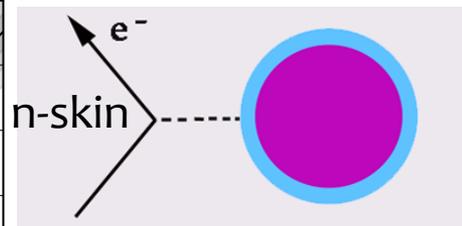
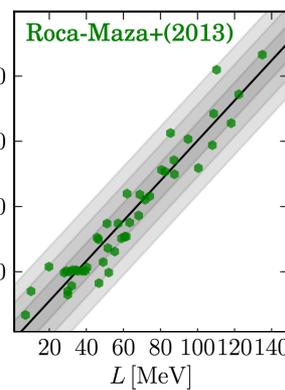
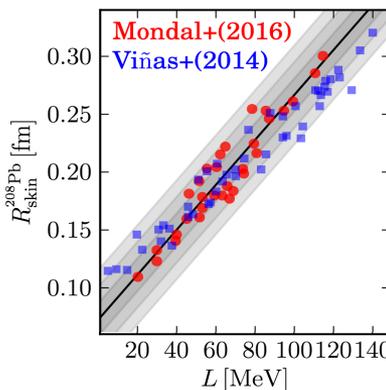
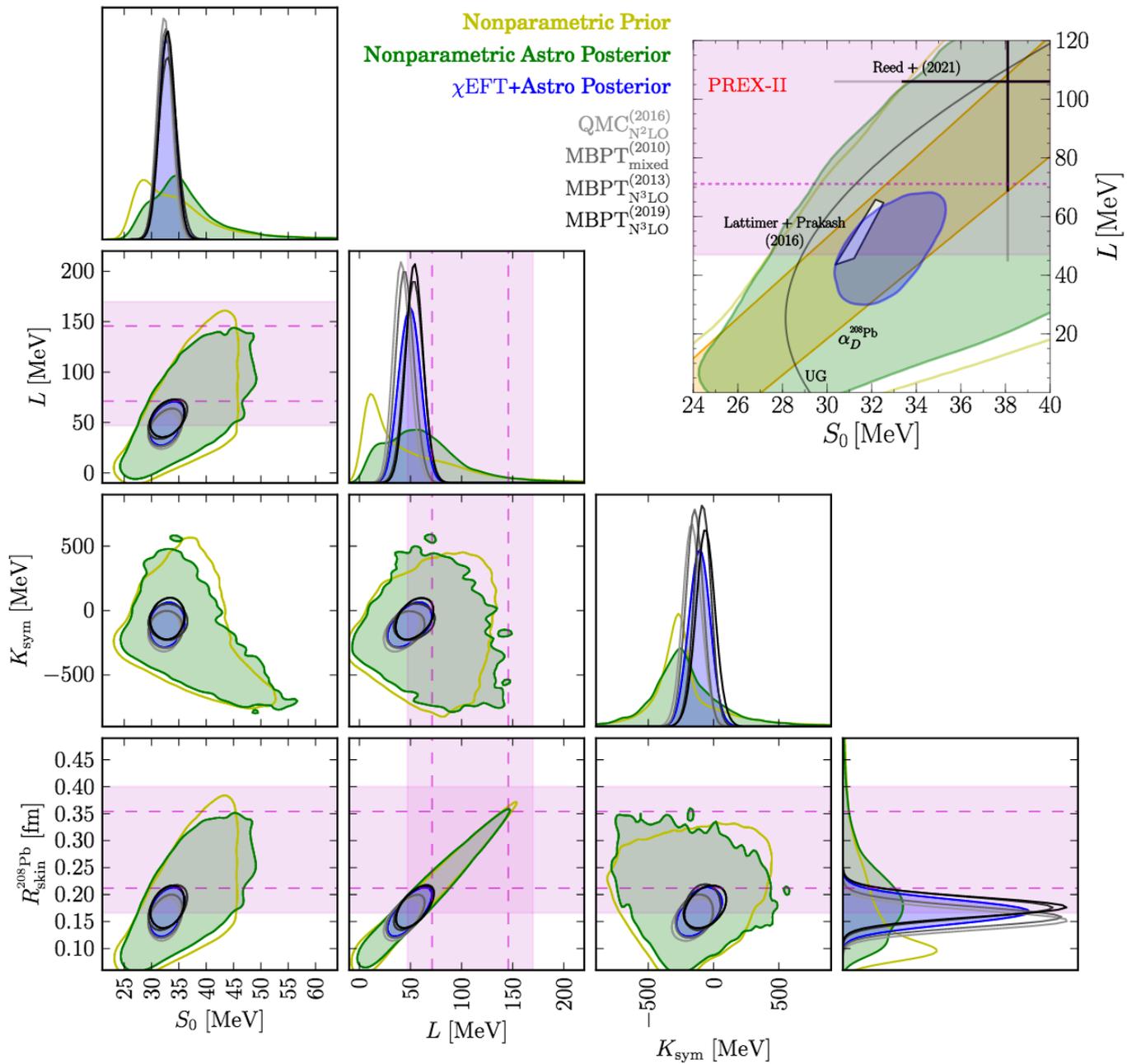


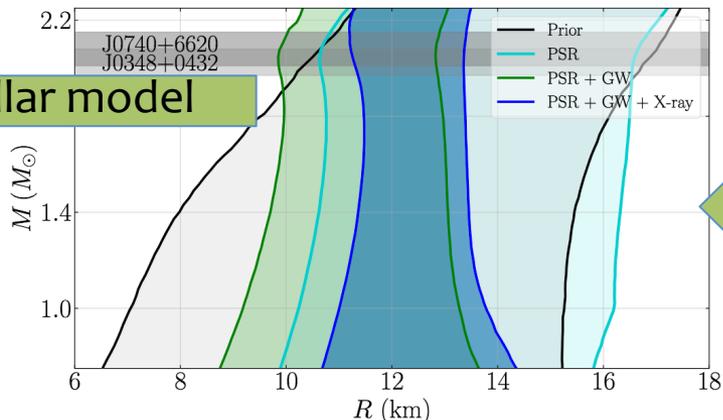
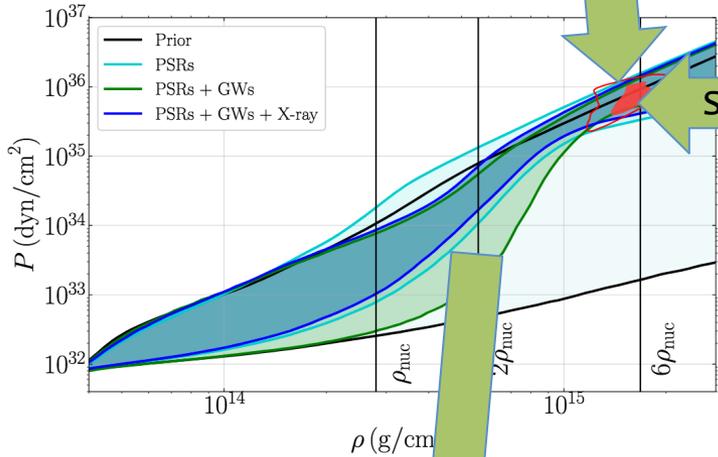
Figure: Lauren Balliet



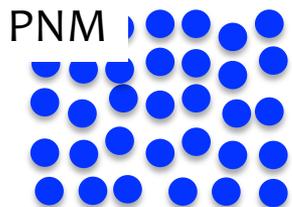
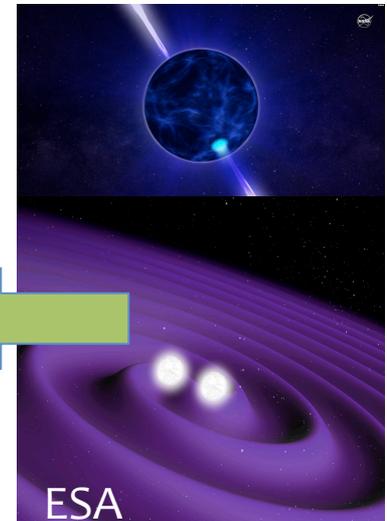


Essick et al: 2107.05528

Gaussian processes



Legred et al arxiv:2106.05313
Essick et al: 2107.05528



BBP; Reconstruct Nuclear matter

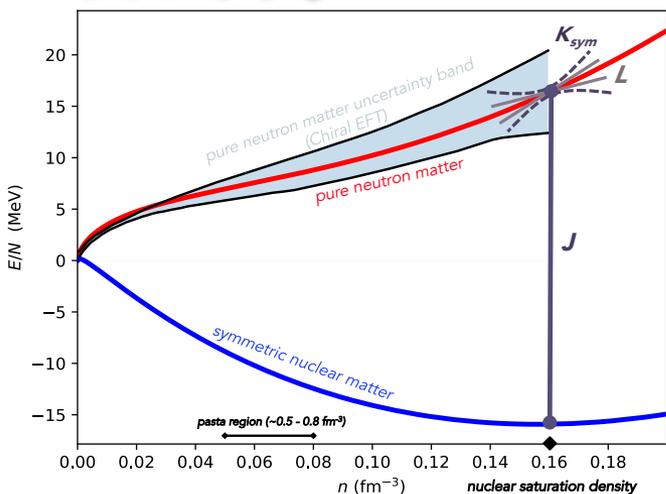
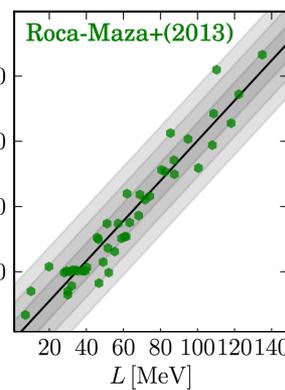
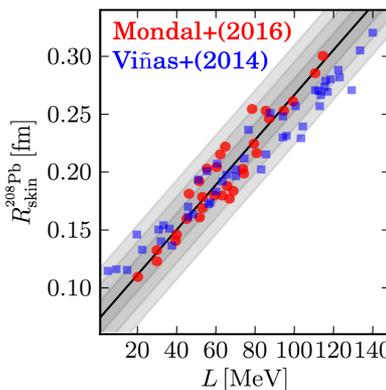
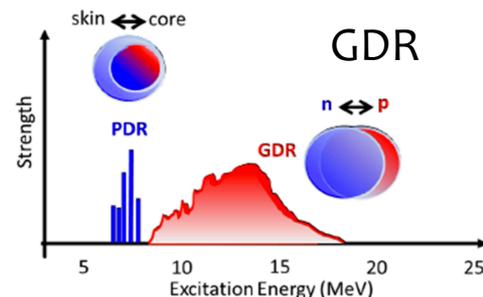
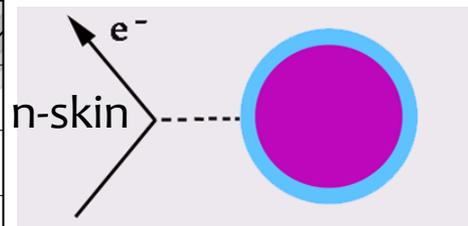


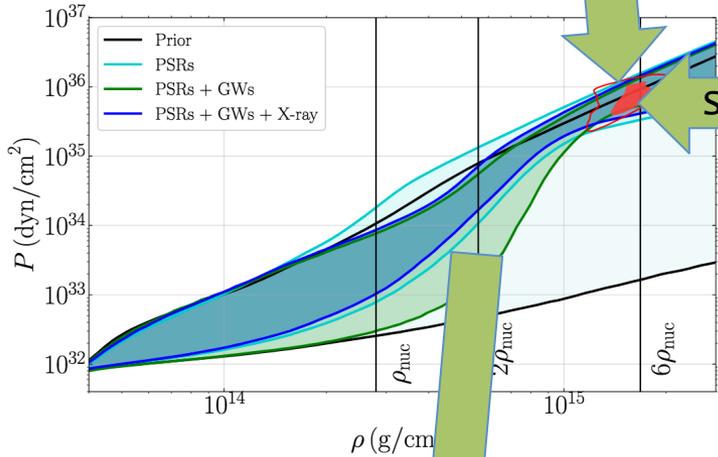
Figure: Lauren Balliet



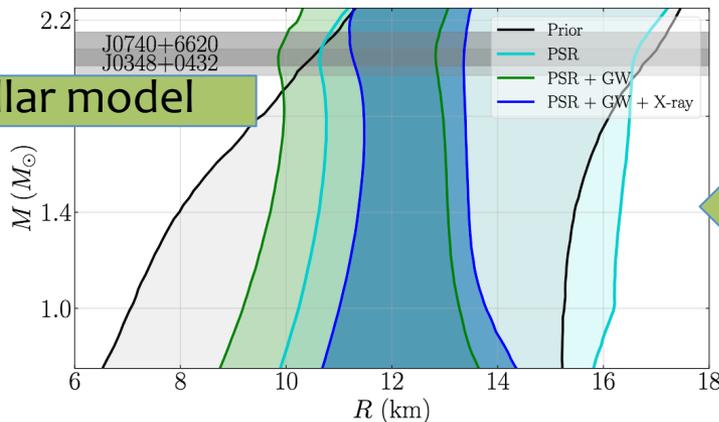
Use "empirical" correlations



Gaussian processes

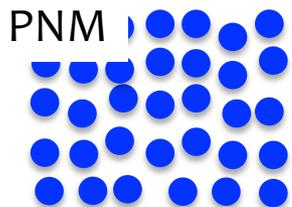
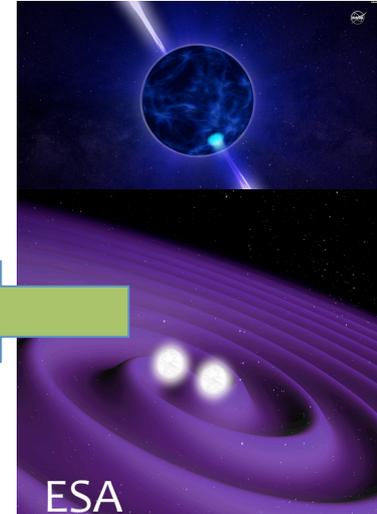


stellar model



Legred et al arxiv:2106.05313

Essick et al: 2107.05528



BBP; Reconstruct Nuclear matter

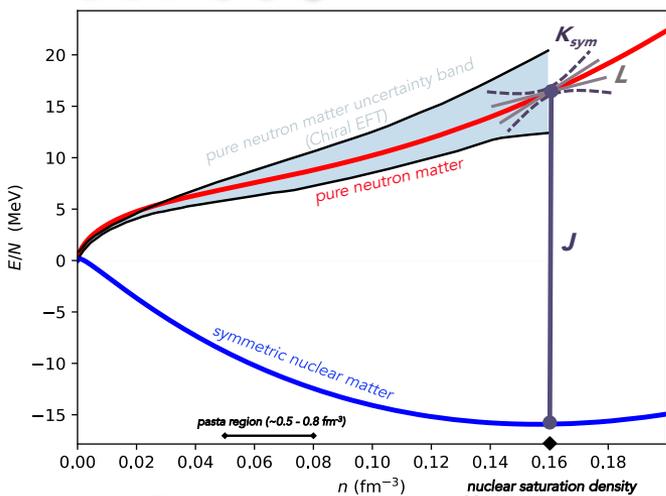
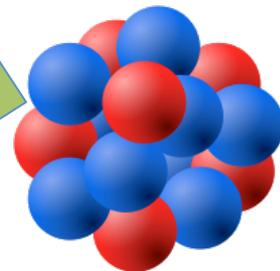
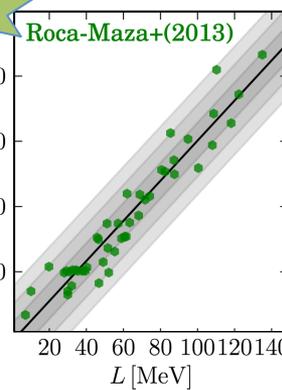
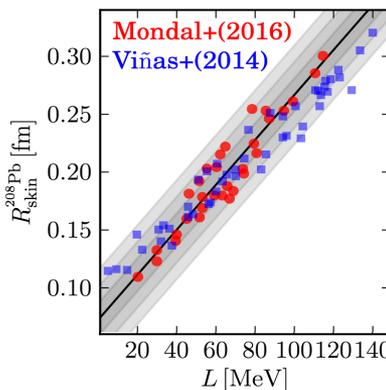
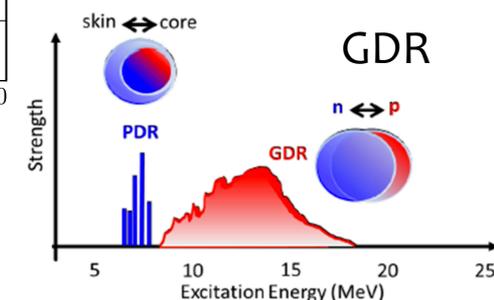
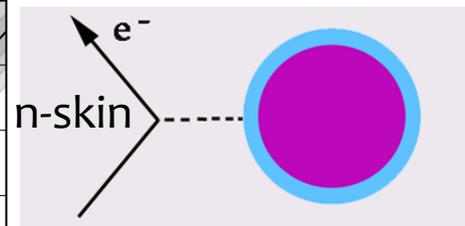


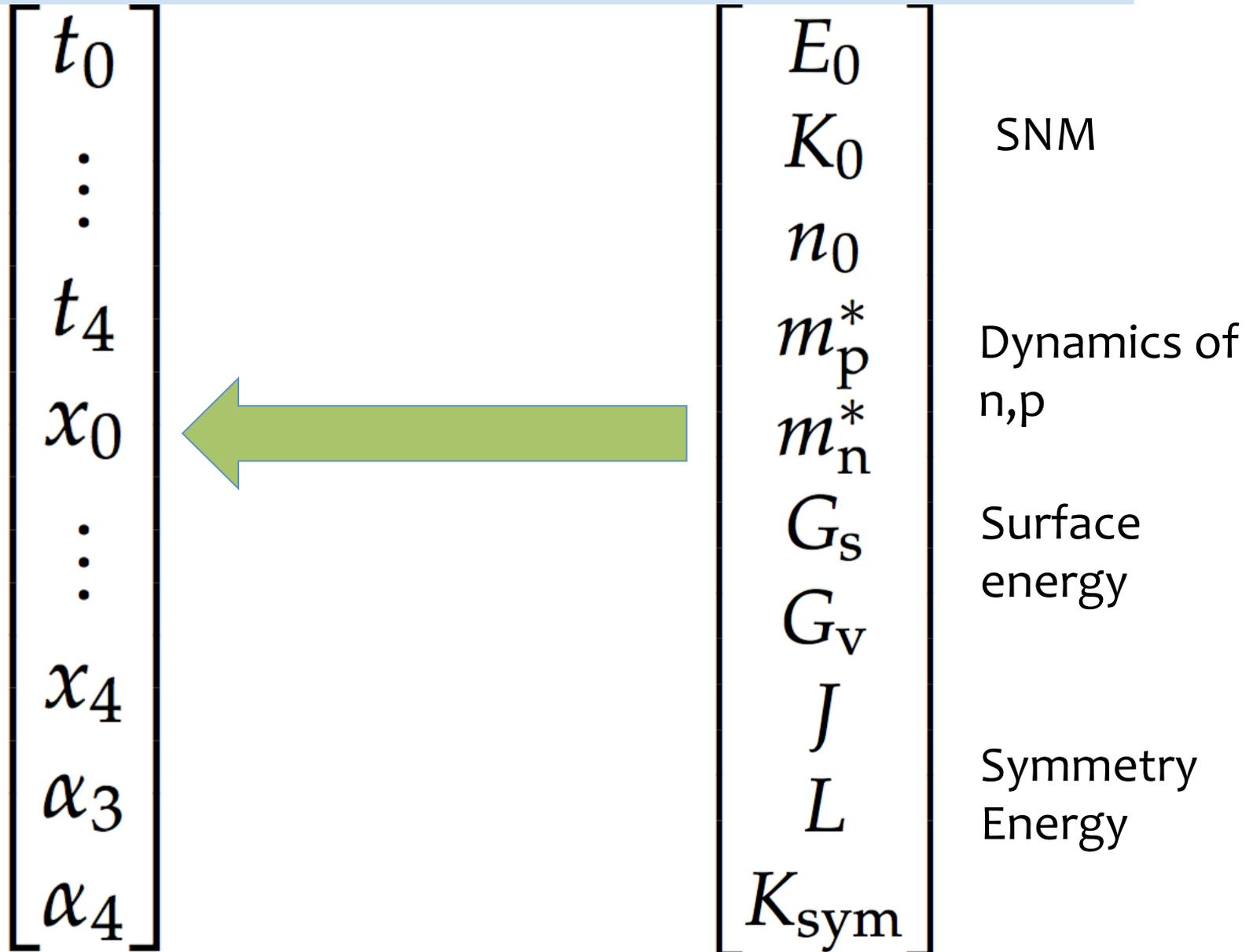
Figure: Lauren Balliet



Use "empirical" correlations



More systematic: map nuclear matter parameters to model parameters and systematically generate models



The overarching strategy

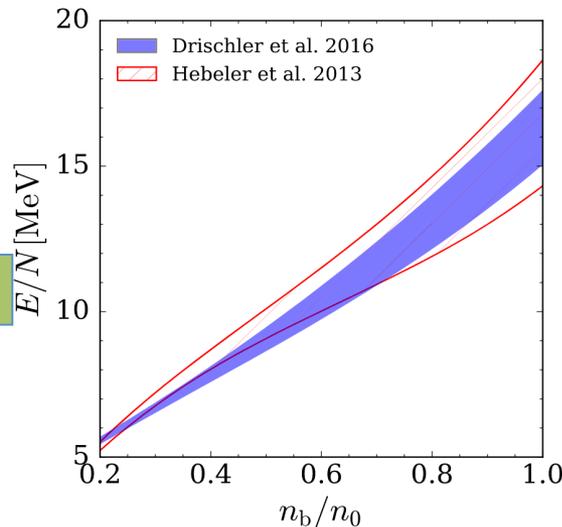
- Choose an EDF with enough degrees of freedom to mitigate the influence of choosing that EDF rather than any other.
- Prepare ensembles of parameterizations of the EDF that distributed over a wide range of the space of nuclear matter parameters (Priors)
- Choose methods of modeling nuclei and neutron star crust which account for as much physics as possible in as reasonable way as possible while being computationally expeditious (10,000s-100,000s models will need to be sampled)
- Use ensemble to calculate nuclear observables, unified crust-core EOS and astro observables
- Add data, construct Likelihoods -> MCMC sampling of posterior probability distribution of the EOSs

More systematic: map nuclear matter parameters to model parameters and systematically generate models

$$\begin{bmatrix} x_0 \\ x_3 \\ x_4 \end{bmatrix} \longleftarrow \begin{bmatrix} J \\ L \\ K_{\text{sym}} \end{bmatrix}$$

$$\begin{bmatrix} t_0 \\ \vdots \\ t_4 \\ x_1 \\ x_2 \end{bmatrix} \longleftarrow \begin{bmatrix} E_0 \\ K_0 \\ n_0 \\ m_p^* \\ m_n^* \\ G_s \\ G_v \end{bmatrix} \begin{bmatrix} x_0 \\ x_3 \\ x_4 \end{bmatrix}$$

$$\begin{bmatrix} \alpha_3 \\ \alpha_4 \end{bmatrix} \longleftarrow$$



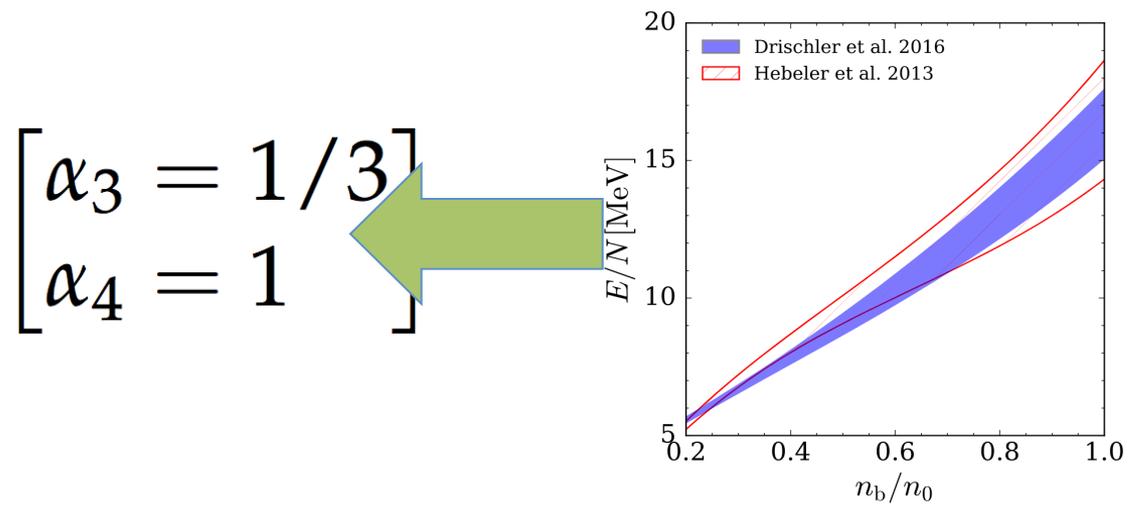
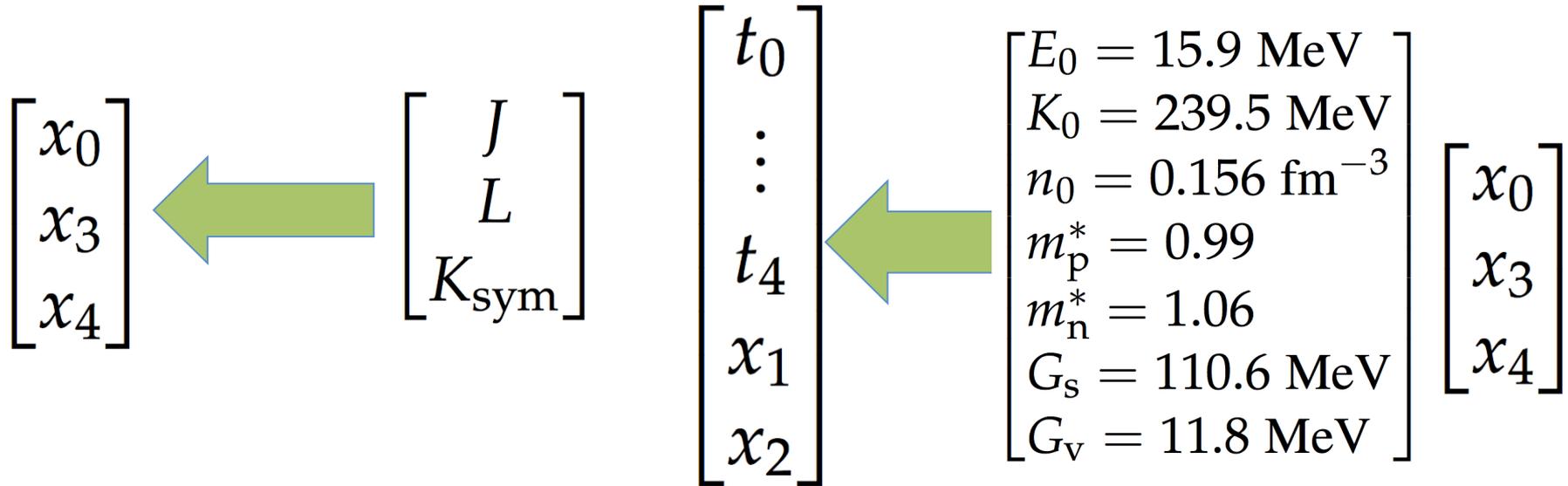
Haensel, Fortin JPhysG 2017
Lim, Holt arXiv:1702.02898

Nuclear masses,
giant resonances

ρ_0 (fm^{-3})	0.160 ± 0.005
B (MeV)	16.0 ± 0.5
K (MeV)	230 ± 30

Lim, Holt arXiv:1702.02898

More systematic: map nuclear matter parameters to model parameters and systematically generate models



Haensel, Fortin JPhysG 2017
Lim, Holt arXiv:1702.02898

Nuclear masses, giant resonances

ρ_0 (fm ⁻³)	0.160 ± 0.005
B (MeV)	16.0 ± 0.5
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Potential sources of systematic error

Had to choose an EDF (Skyrme). Enough degrees of freedom?
Can add more (Q_{sym})

Symmetric nuclear matter and gradient parameters held fixed; extending inference to those parameters may change posteriors

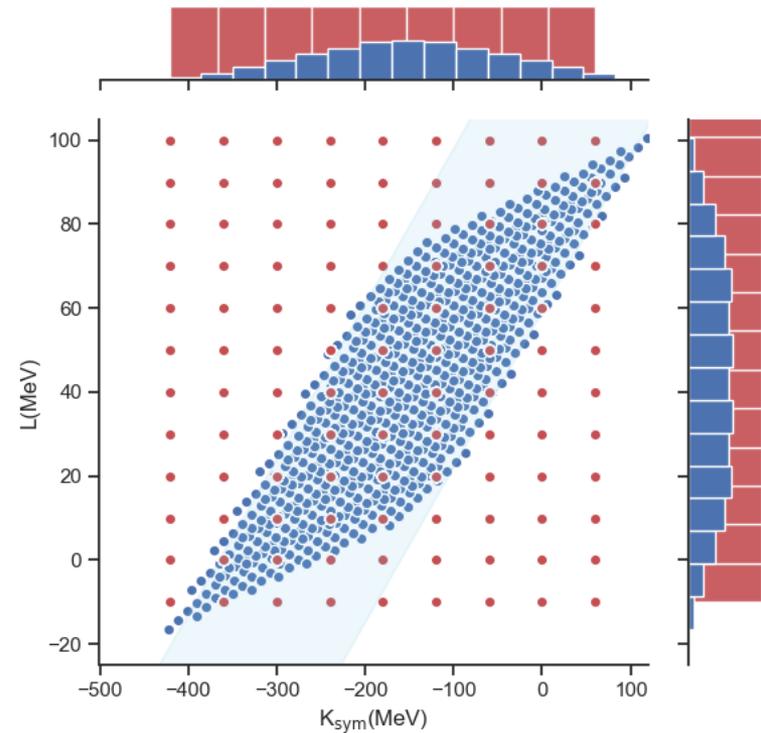
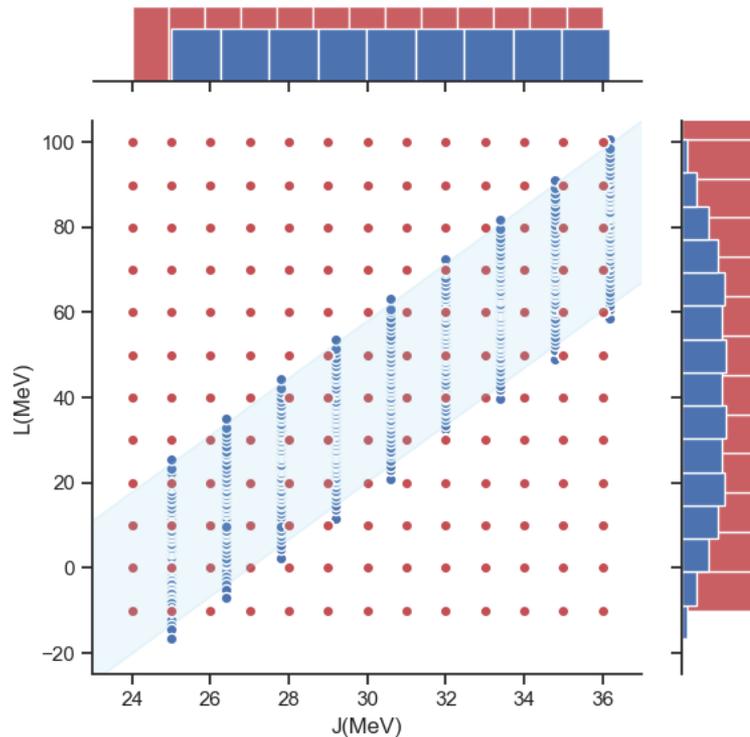
Priors

$$P(J, L, K_{sym})$$

A uniform grid of Skyrme models

Red – Uninformative priors

Blue – Pure neutron matter priors (Fermi liquid theory) [Holt&Lim PLB 2018](#)



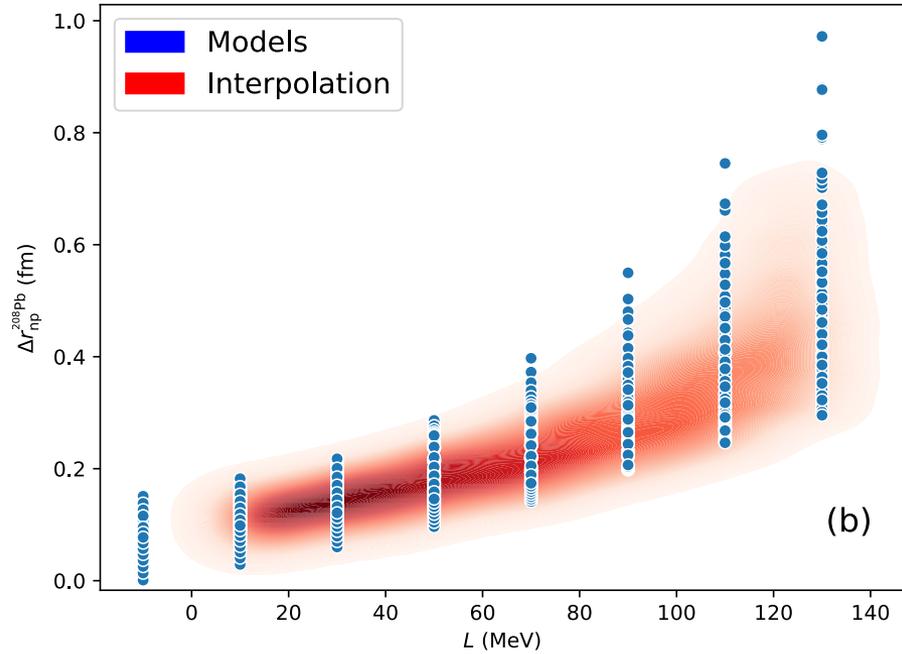
J : 25-43 MeV

L : 5-145 MeV

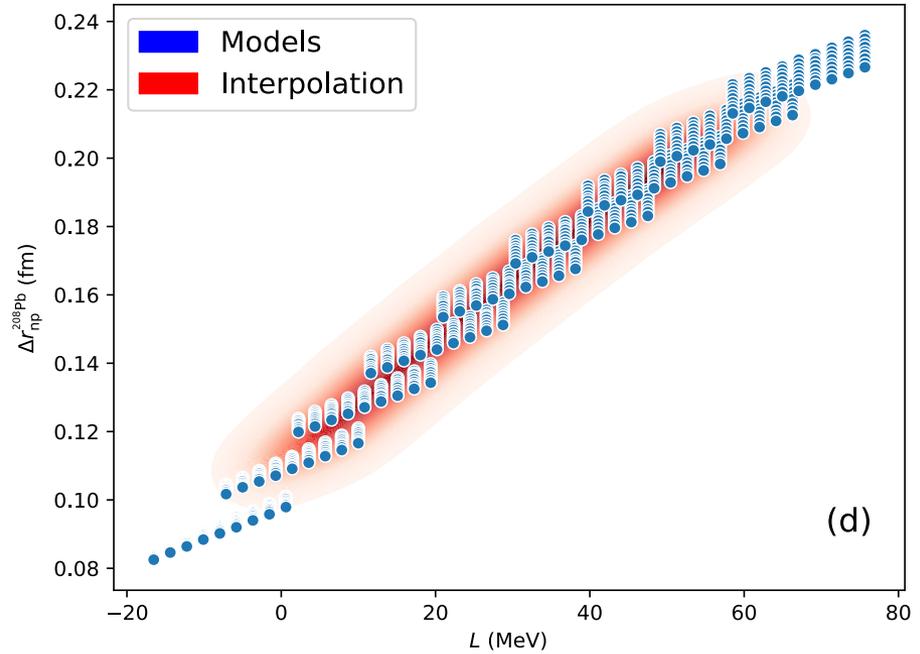
K_{sym} : -500- +200 MeV

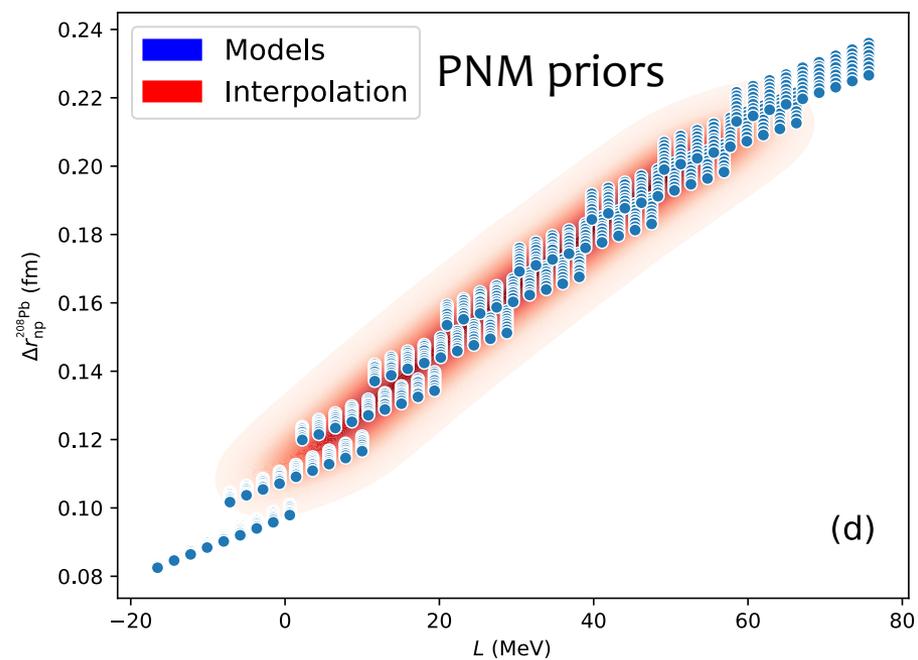
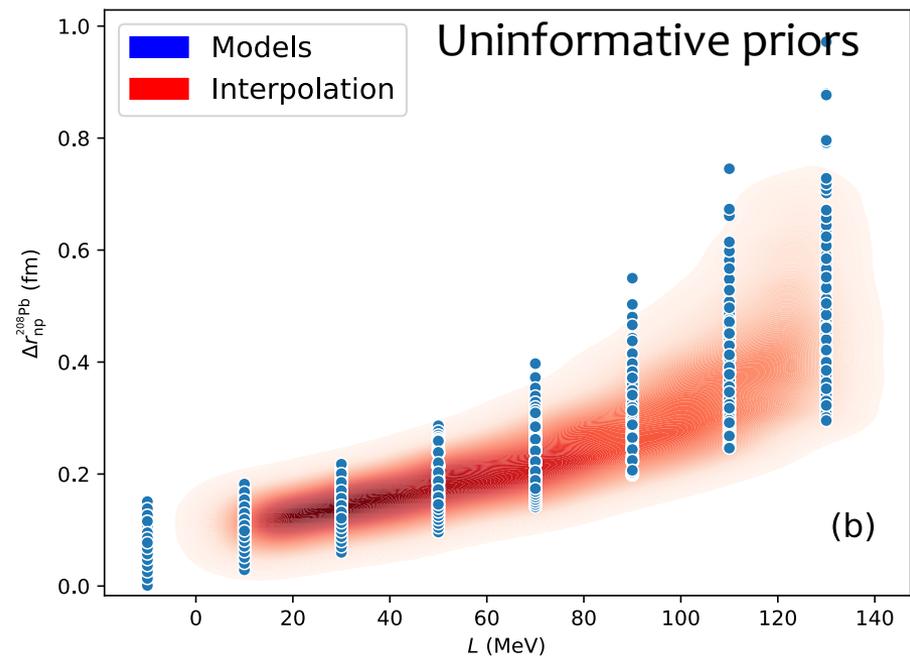
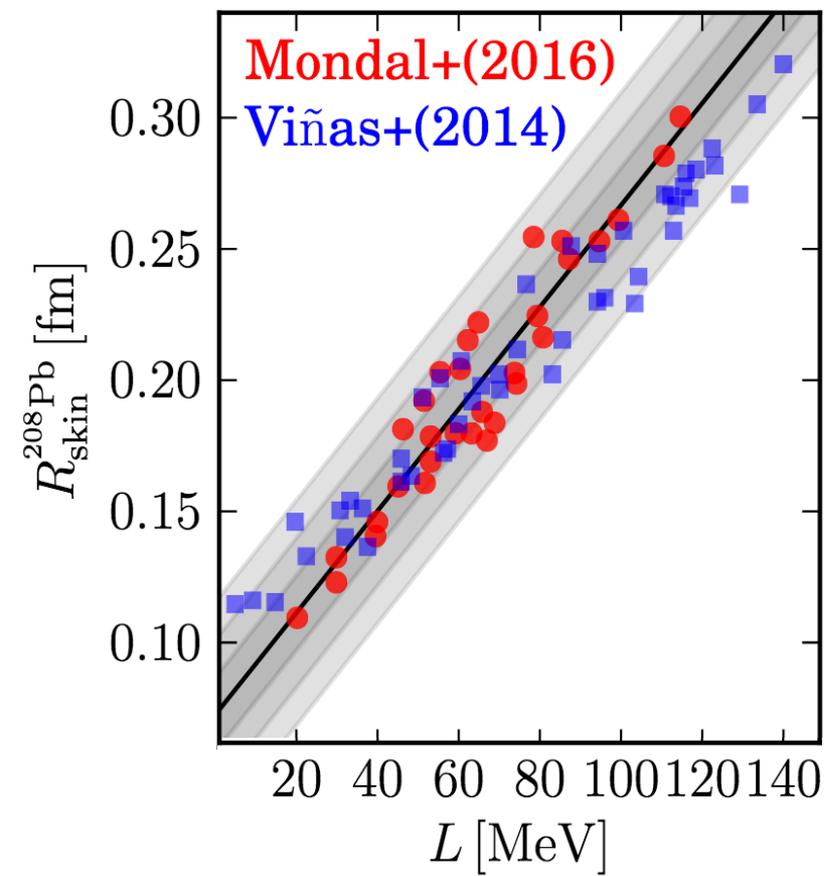
Starting from a set of systematically generated EDFs with minimal symmetry energy assumptions

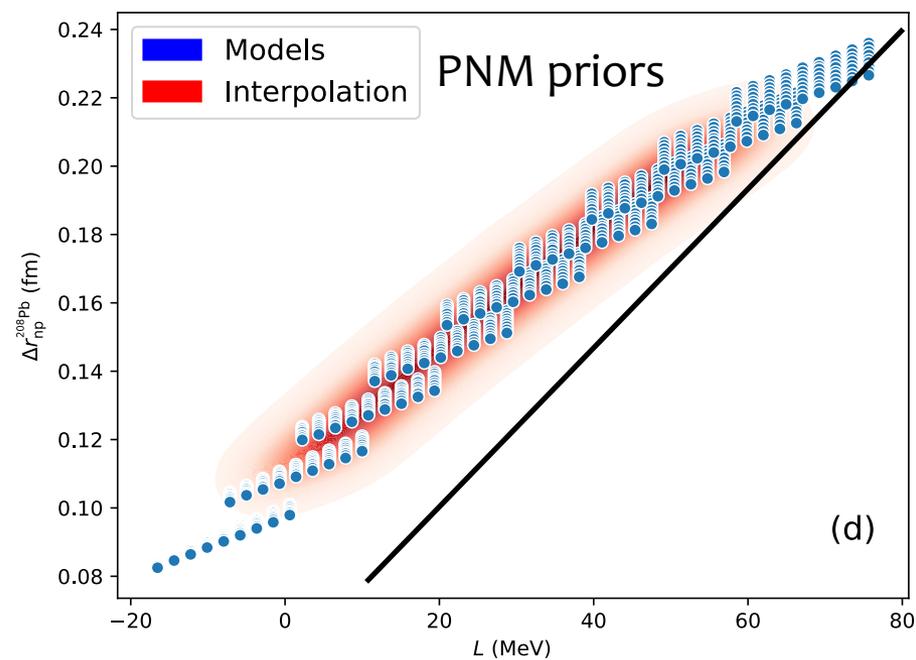
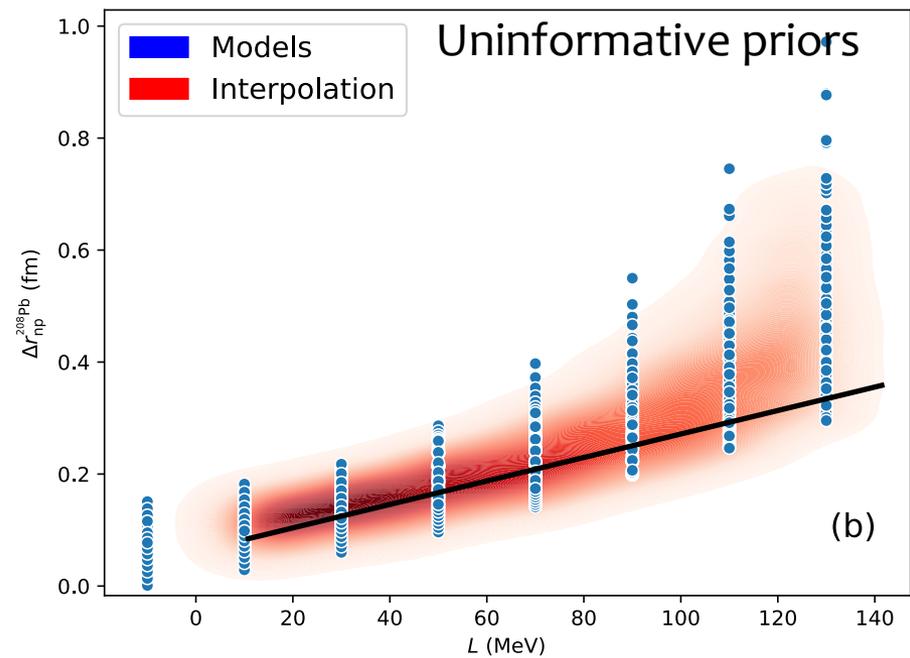
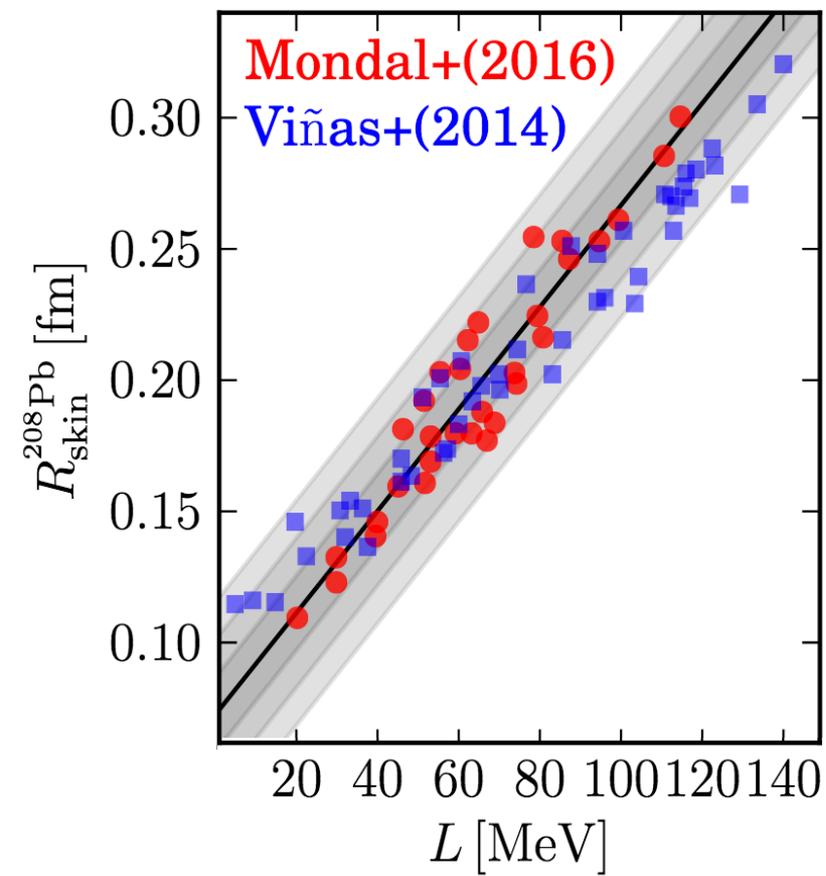
Uninformative priors

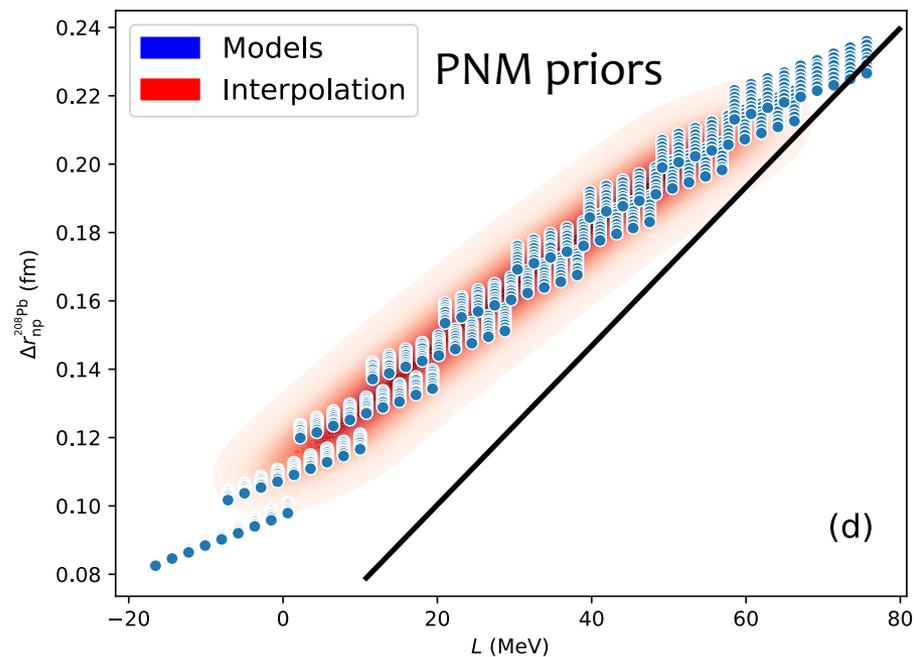
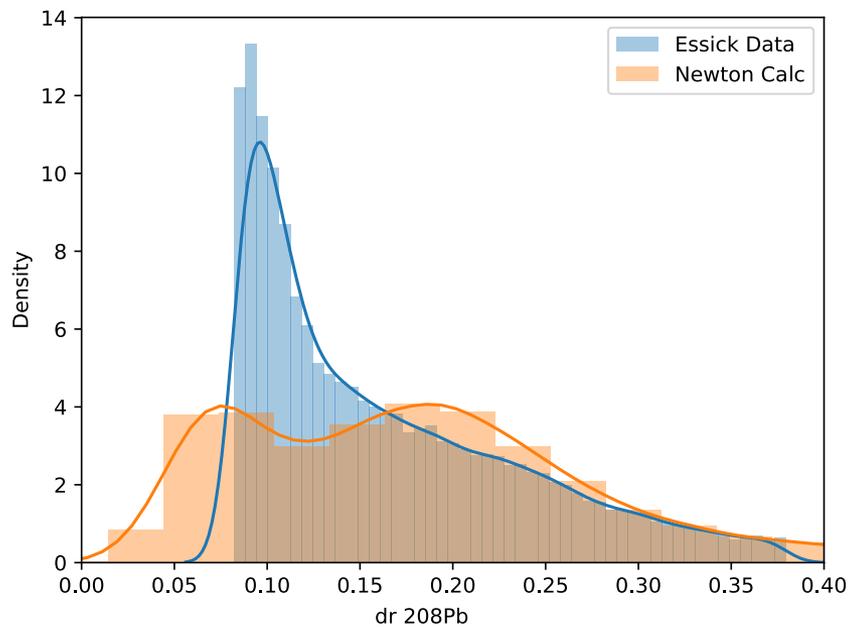
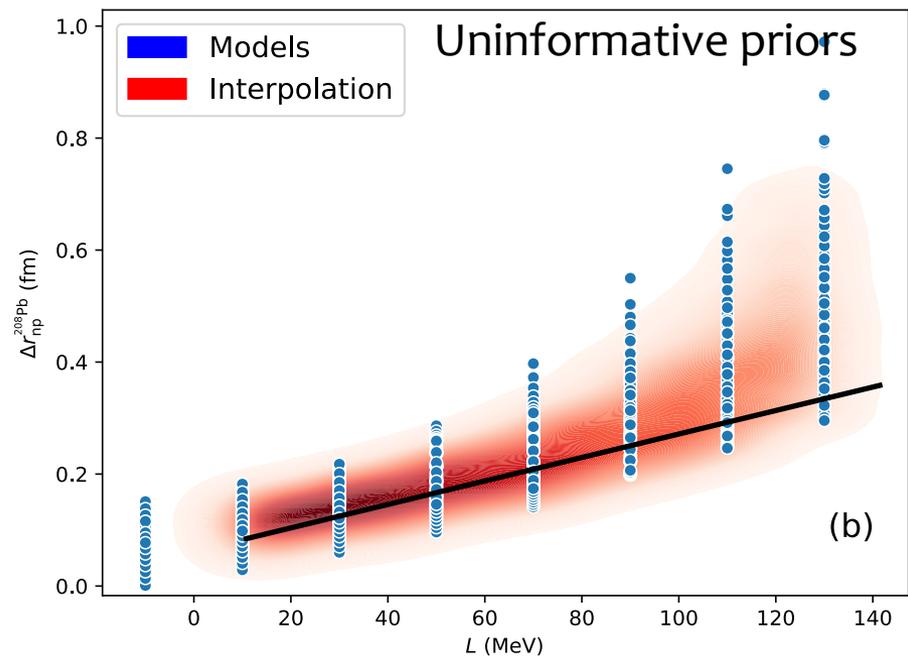
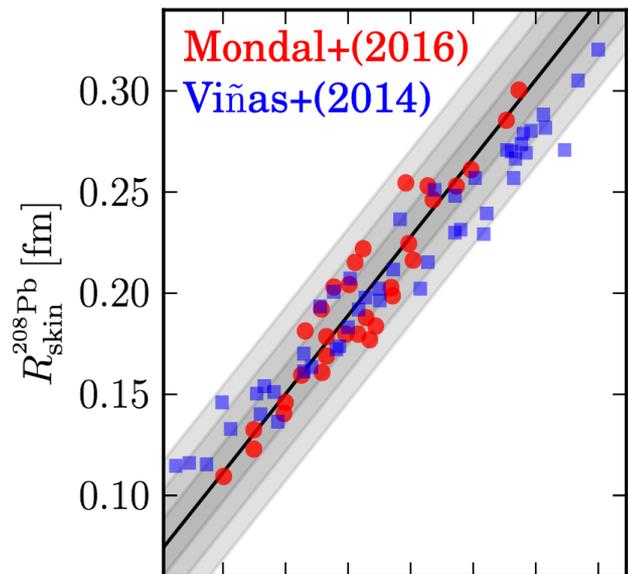


PNM priors









Rebecca Preston, yesterday

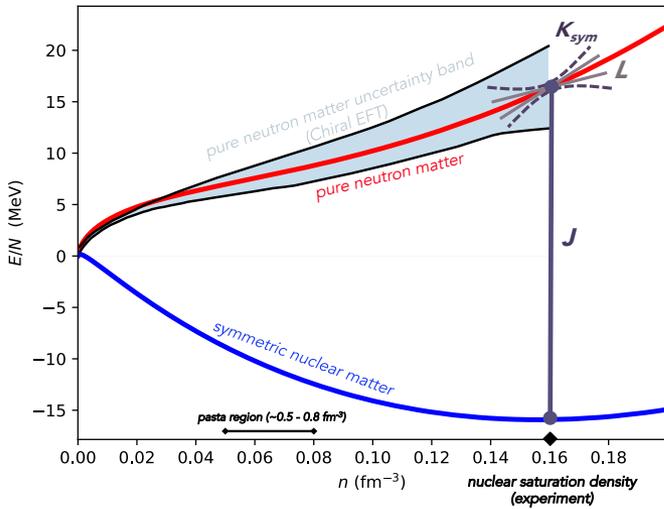
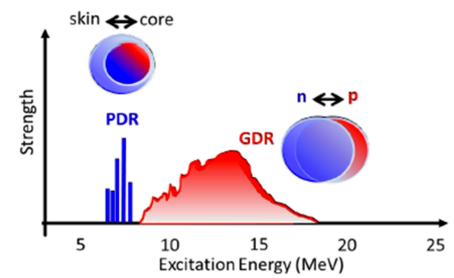
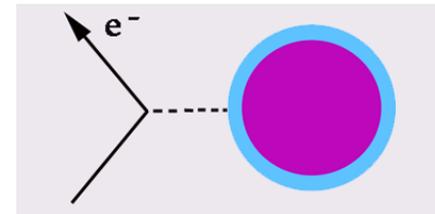
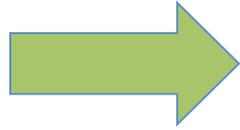


Figure: Lauren Balliet

$$P(J, L, K_{\text{sym}} | \mathcal{D})$$



$$\begin{bmatrix} J \\ L \\ K_{\text{sym}} \end{bmatrix}$$

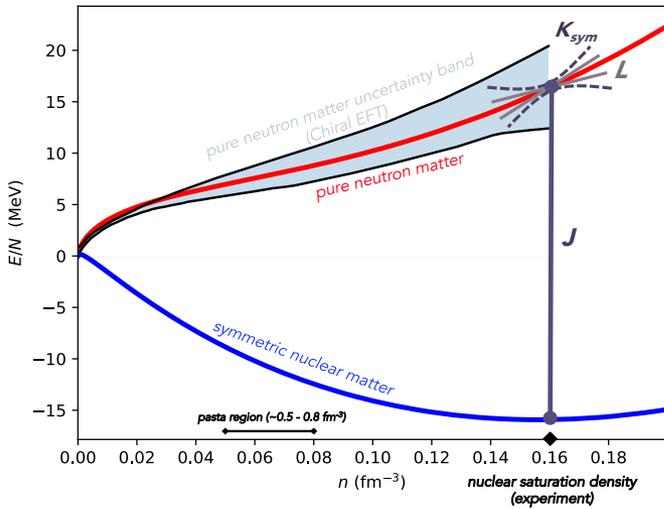


$$\begin{bmatrix} x_0 \\ x_3 \\ x_4 \end{bmatrix}$$



$$\Delta r$$

Skyrme Hartree-Fock
SkyrmeRPA Comp Phys Comms, 184, (2013)



$$P(J, L, K_{\text{sym}} | \mathcal{D})$$

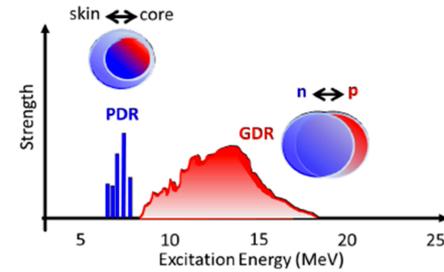
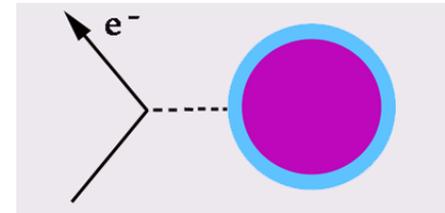
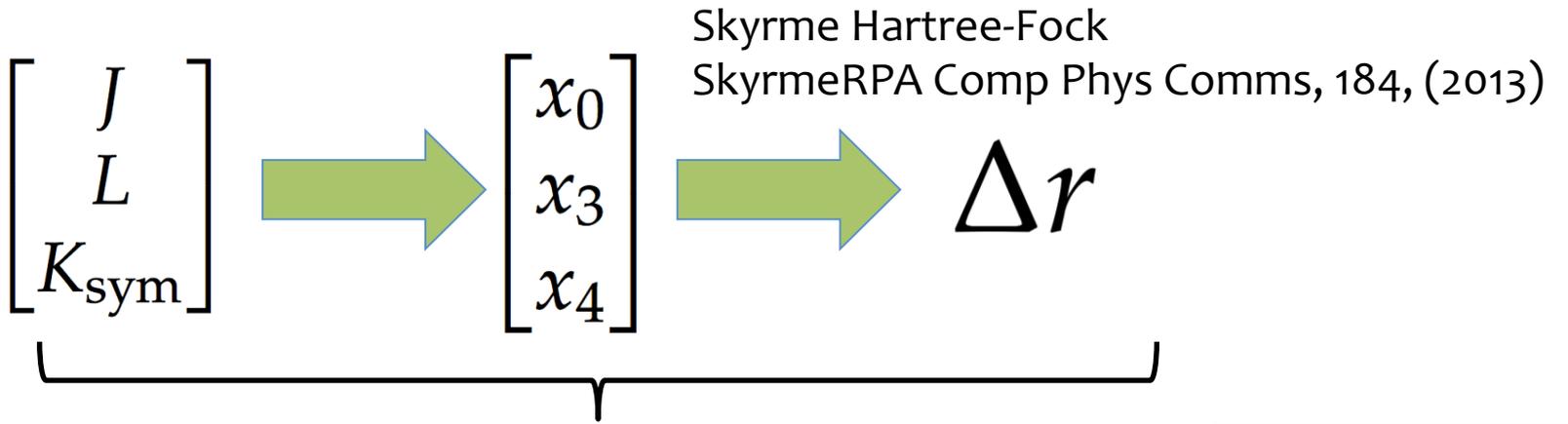


Figure: Lauren Balliet



$$P(\Delta r | J, L, K_{\text{sym}})$$

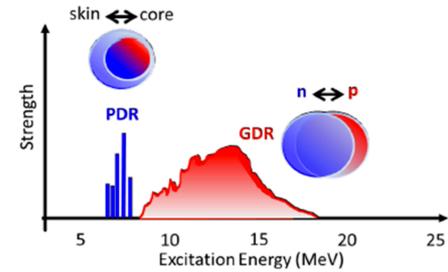
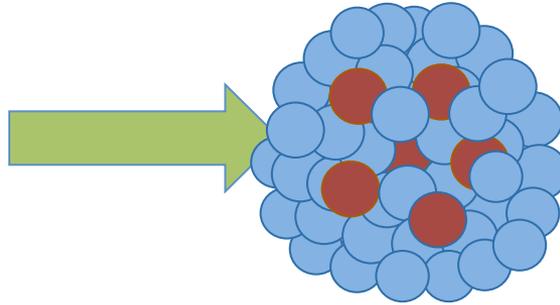
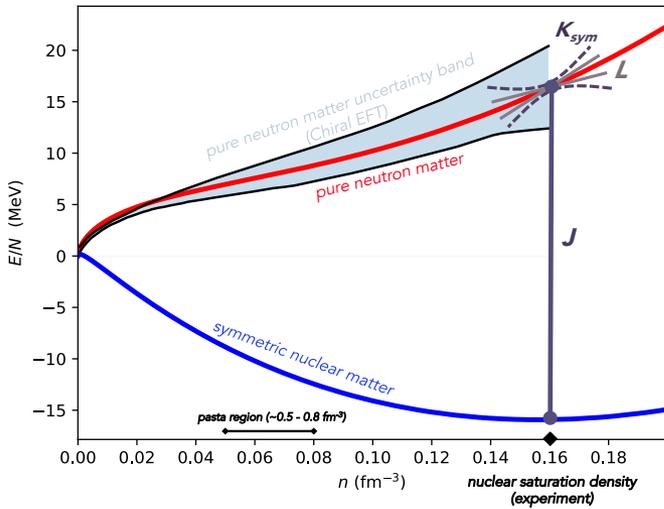
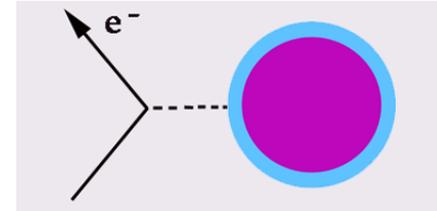
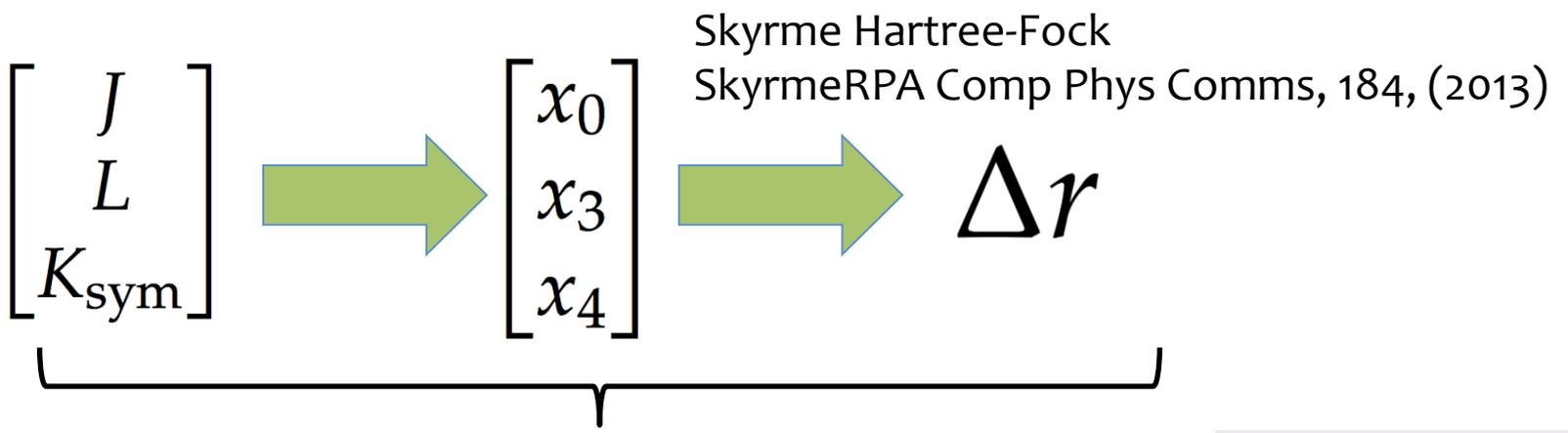
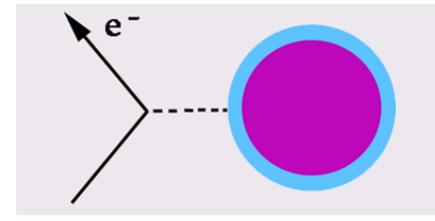
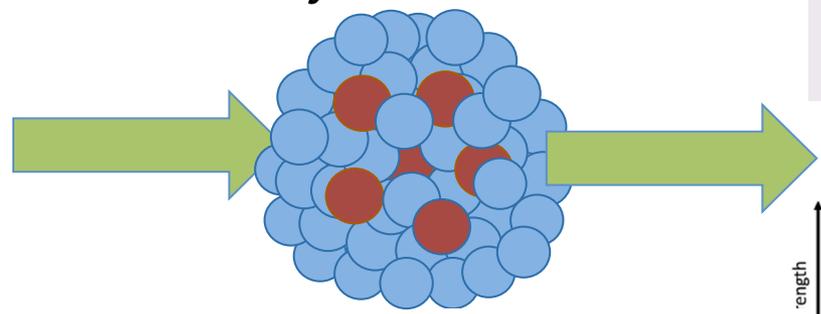
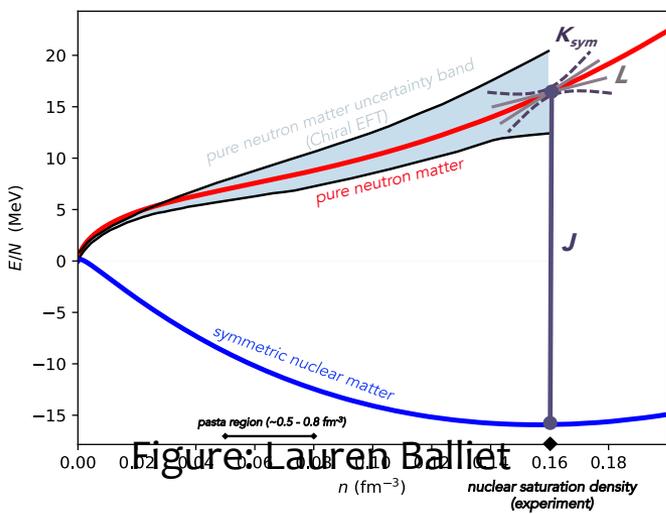


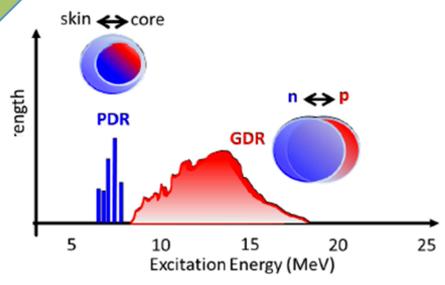
Figure: Lauren Balliet



$$P(\Delta r | J, L, K_{\text{sym}})$$



$$P(\mathcal{D} | \Delta r)$$



$$= \exp \left[- \frac{(\Delta r_{\text{np}}^{\text{data}} - \Delta r_{\text{np}}(J, L, K_{\text{sym}}))^2}{2\sigma_{\text{data}}^2} \right]$$

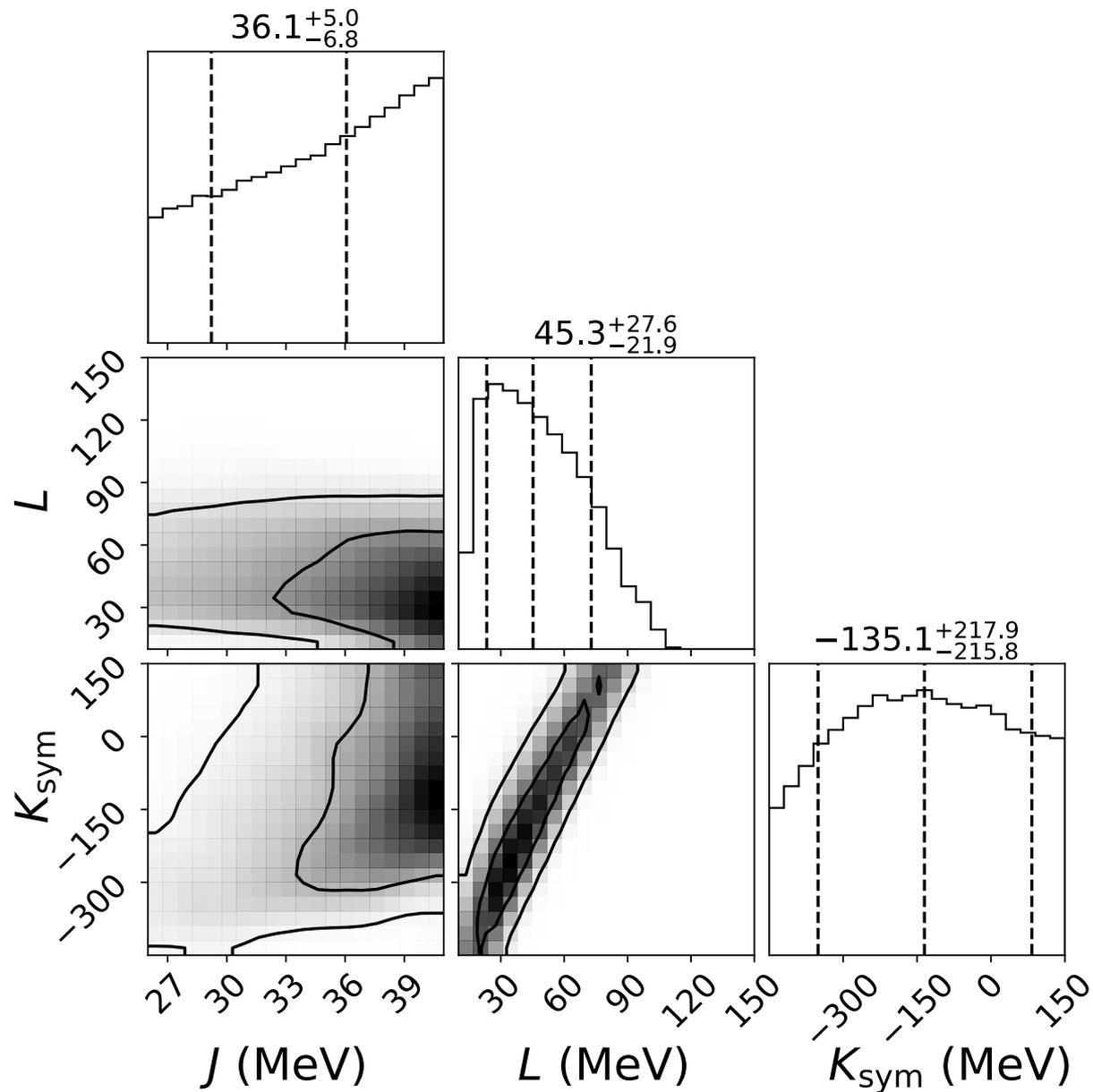
Potential sources of systematic error

Had to choose an EDF (Skyrme). Enough degrees of freedom?
Can add more (Q_{sym})

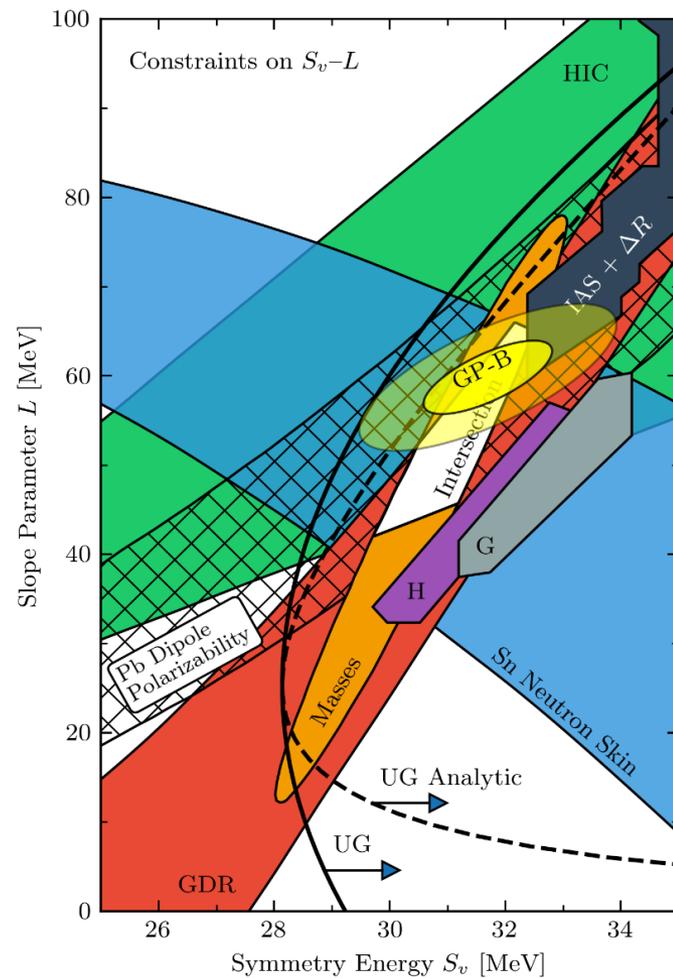
Symmetric nuclear matter and gradient parameters held fixed; extending inference to those parameters may change posteriors

We're usually not directly modeling nuclear observables - but in some cases we could (e.g. weak form factor) and thus improve consistency

Neutron skins: Sn

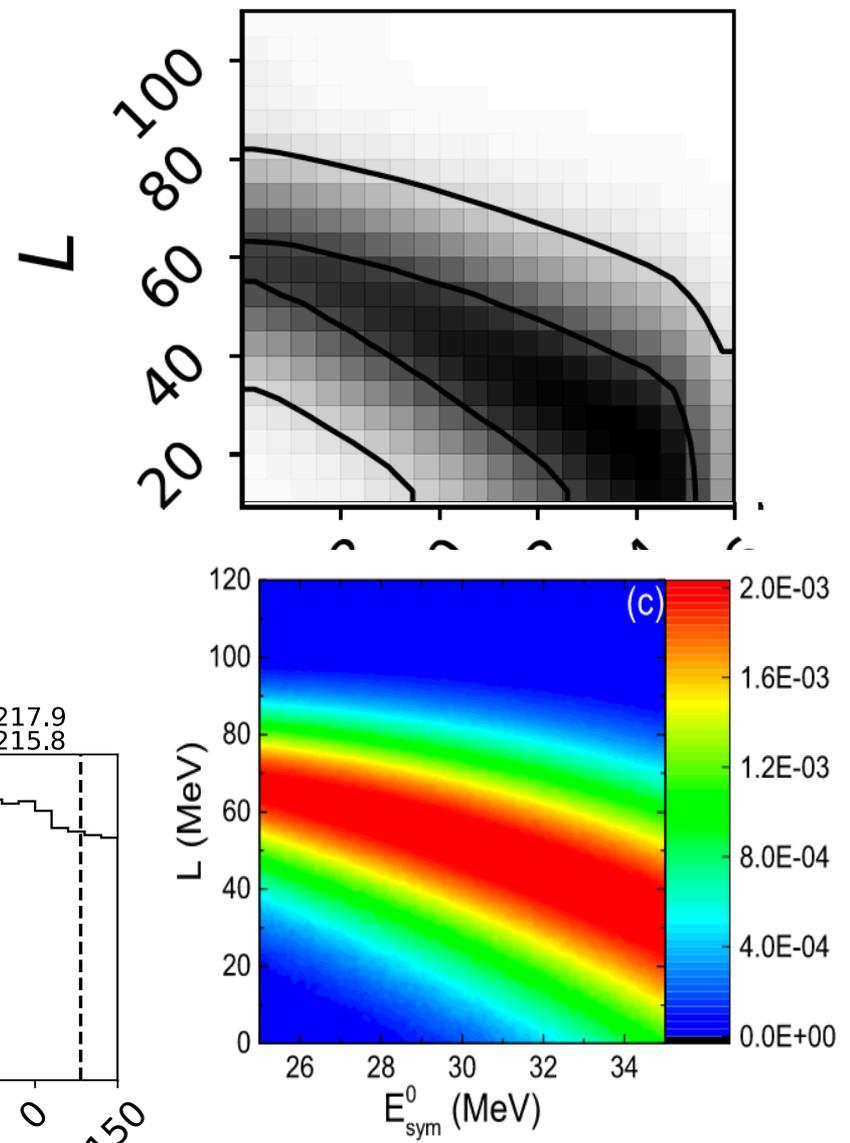
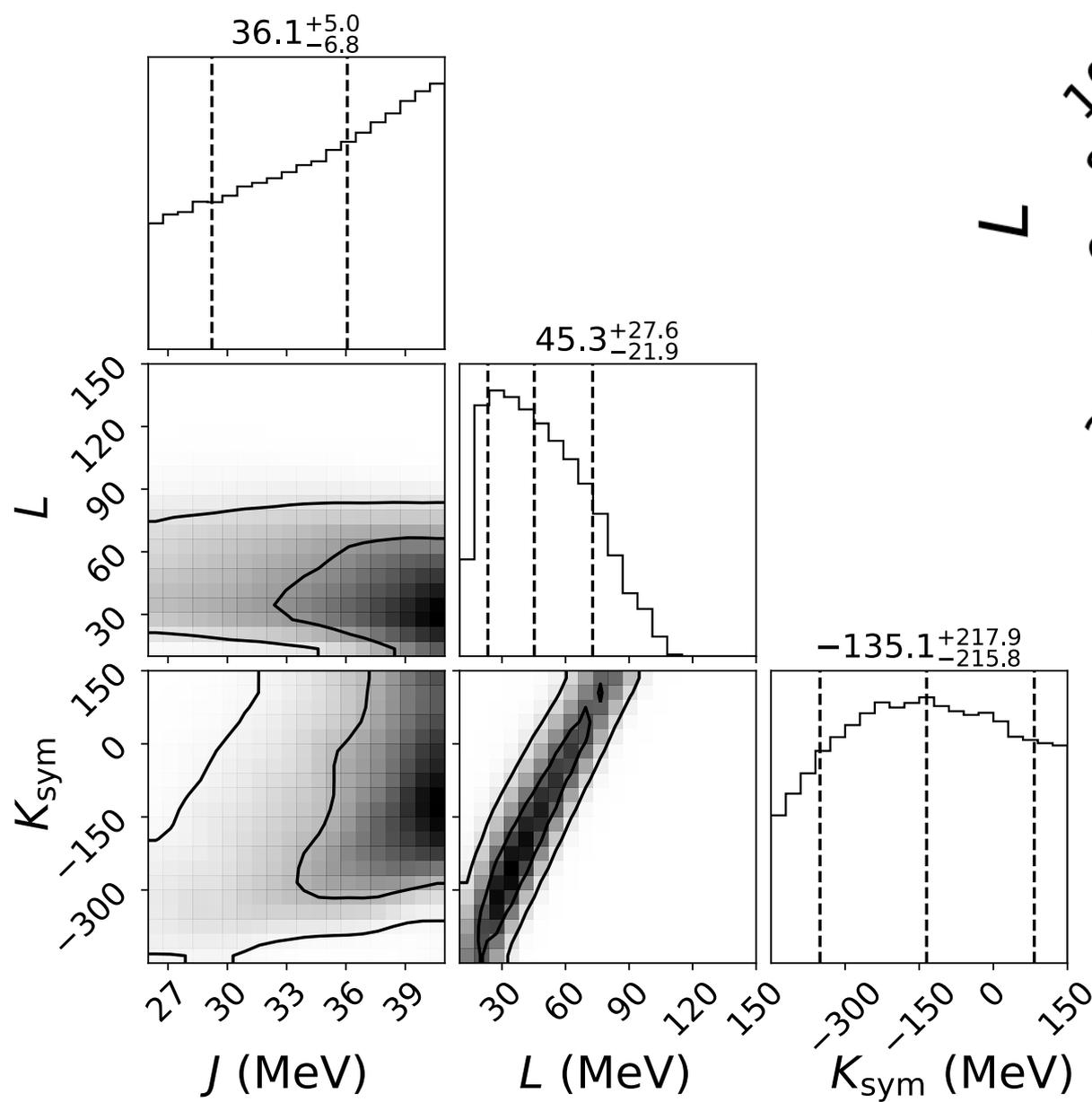


Drischler, arxiv:2004.07232

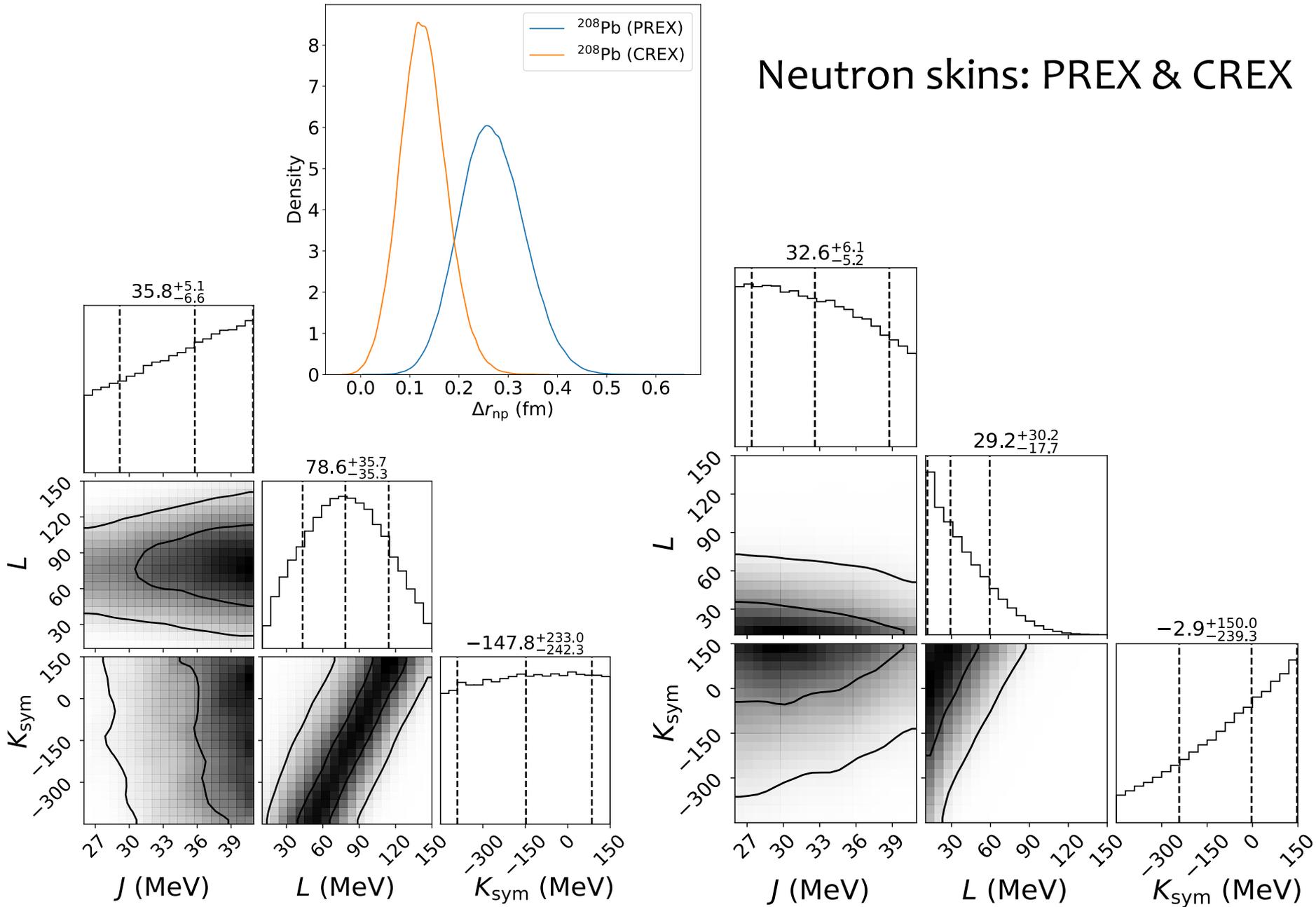


Li, Xie, Xu arxiv:2007.07669

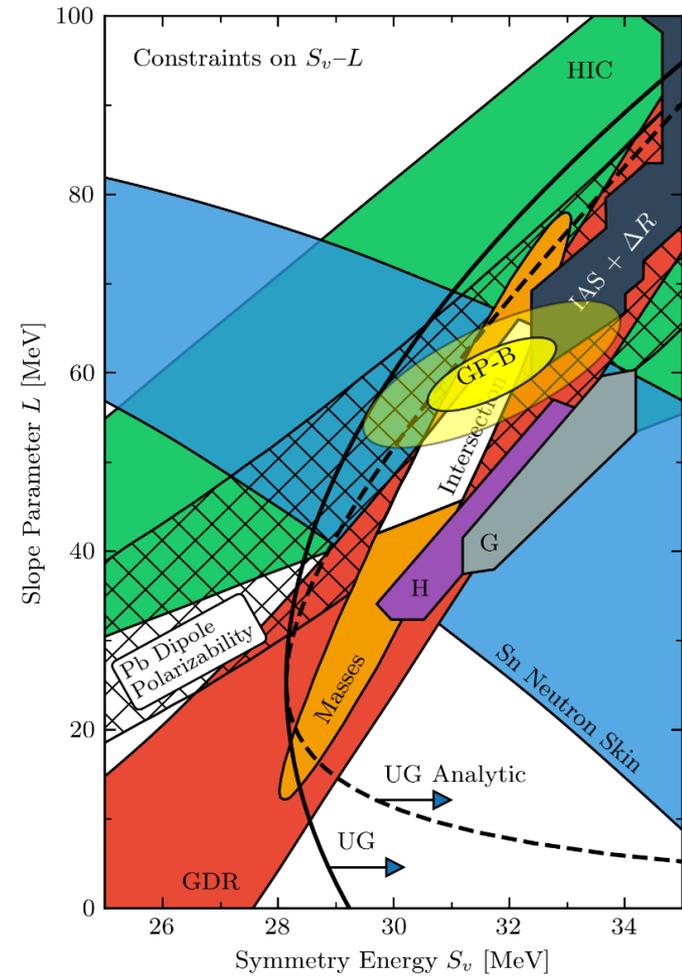
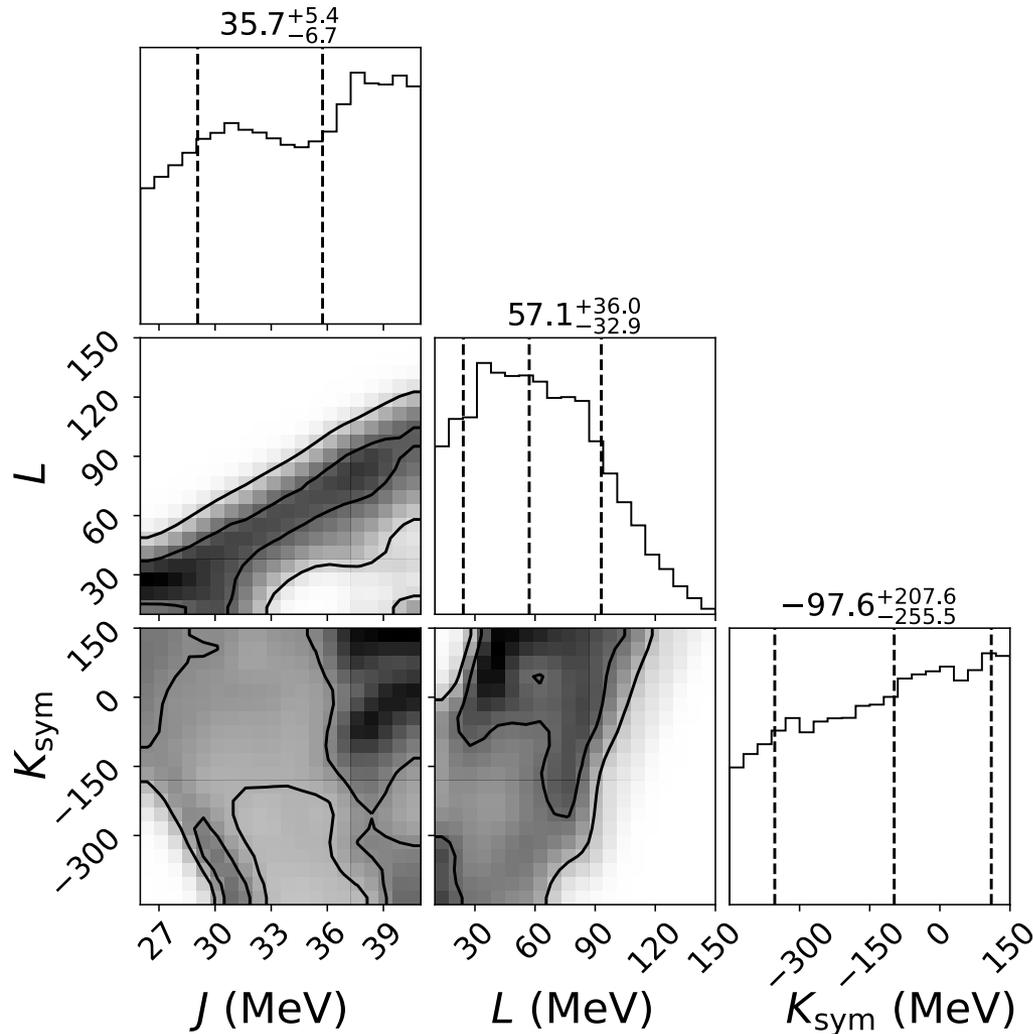
Neutron skins: Sn



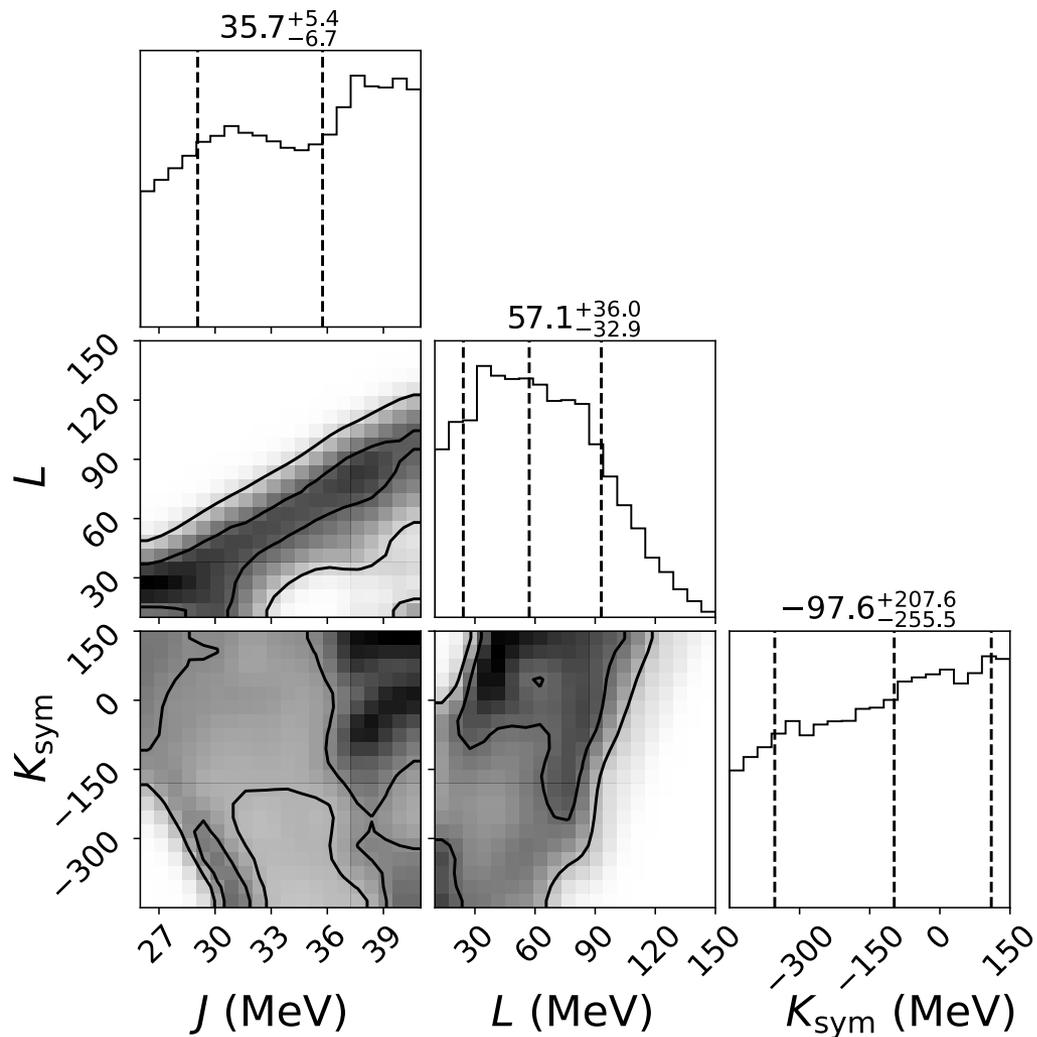
Neutron skins: PREX & CREX



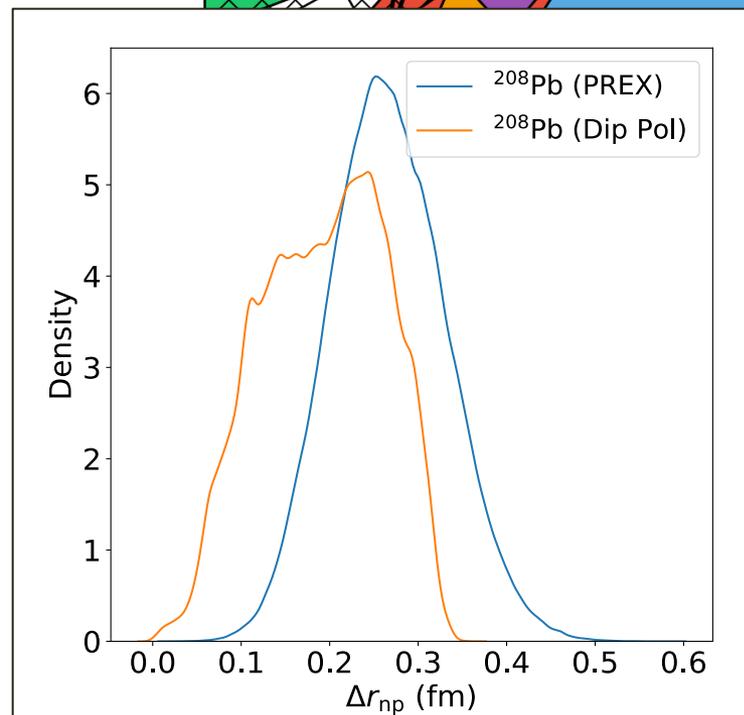
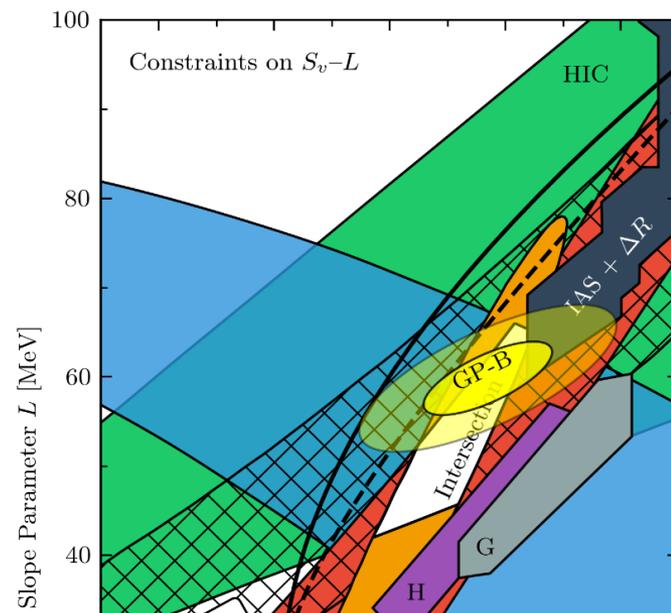
Pb208, Ca48 Dipole Polarizability



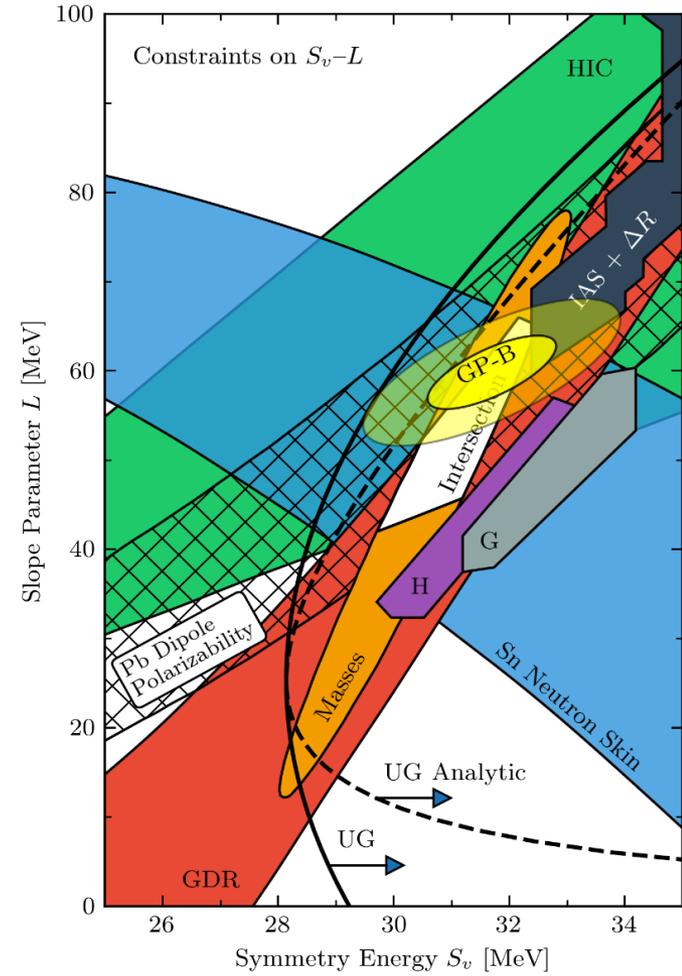
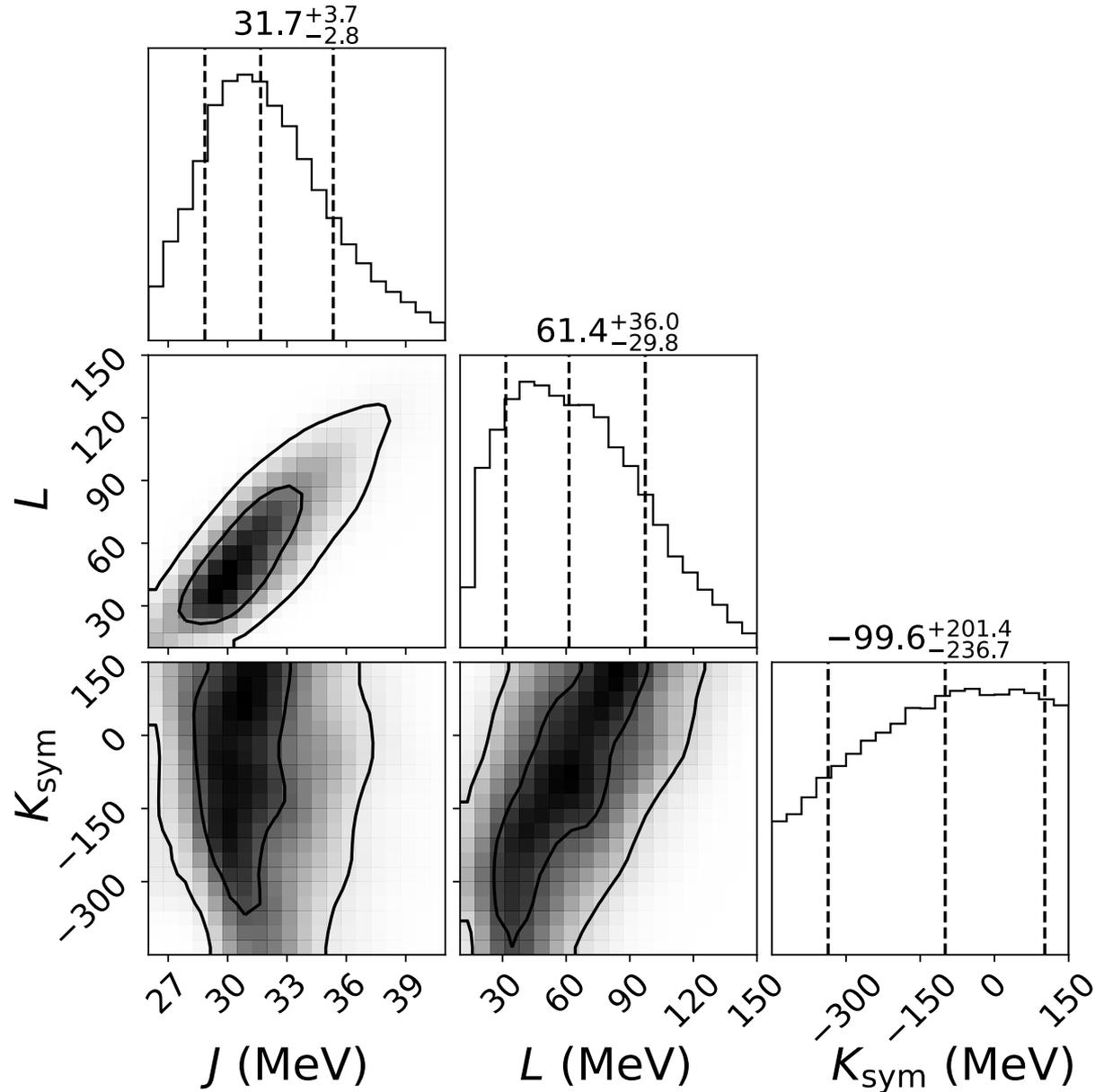
Pb208, Ca48 Dipole Polarizability



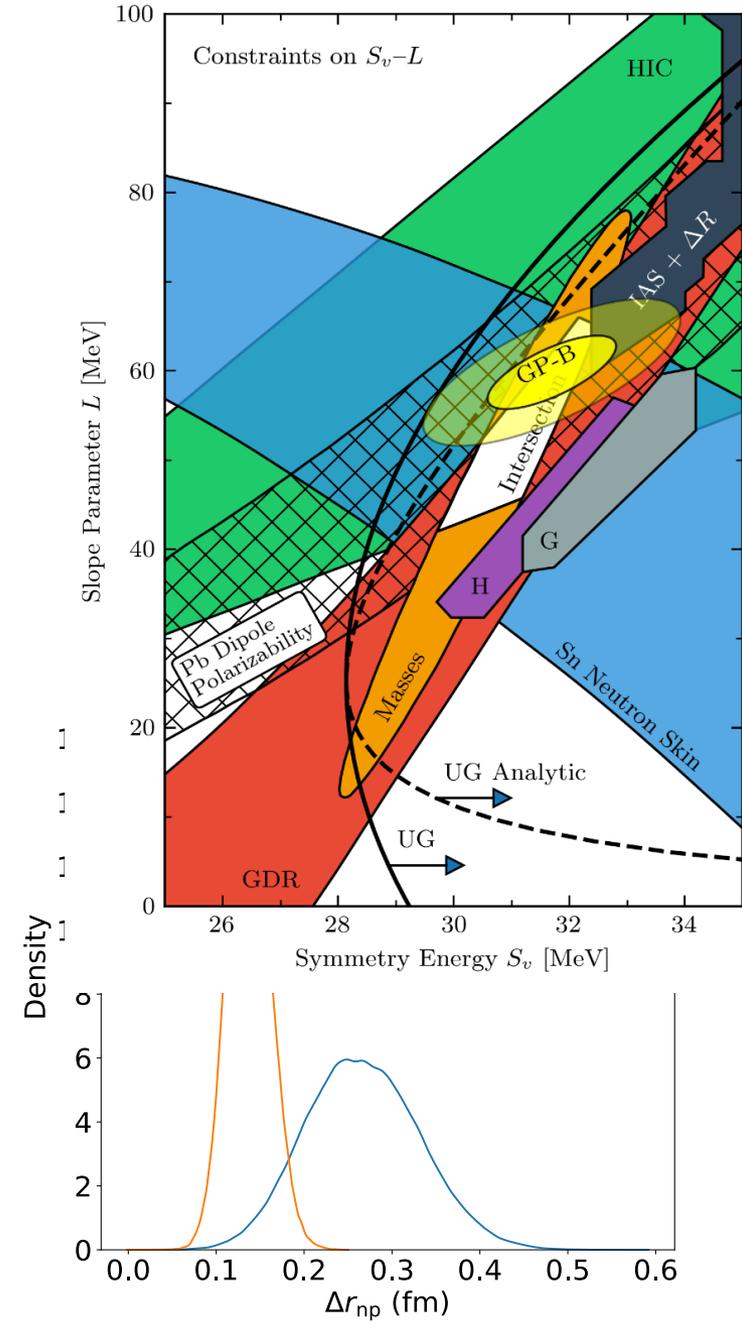
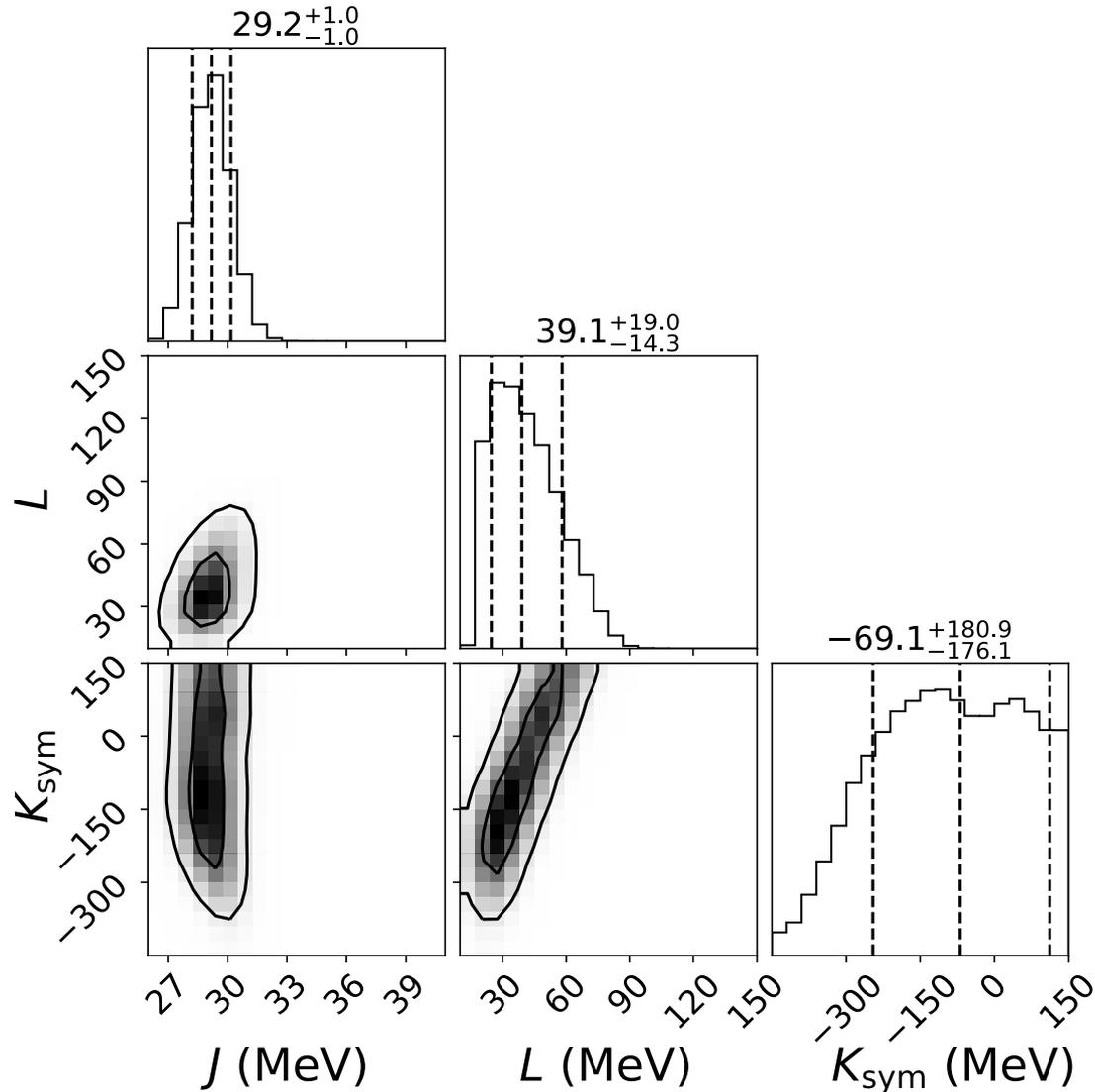
Drischler, arxiv:2004.07232



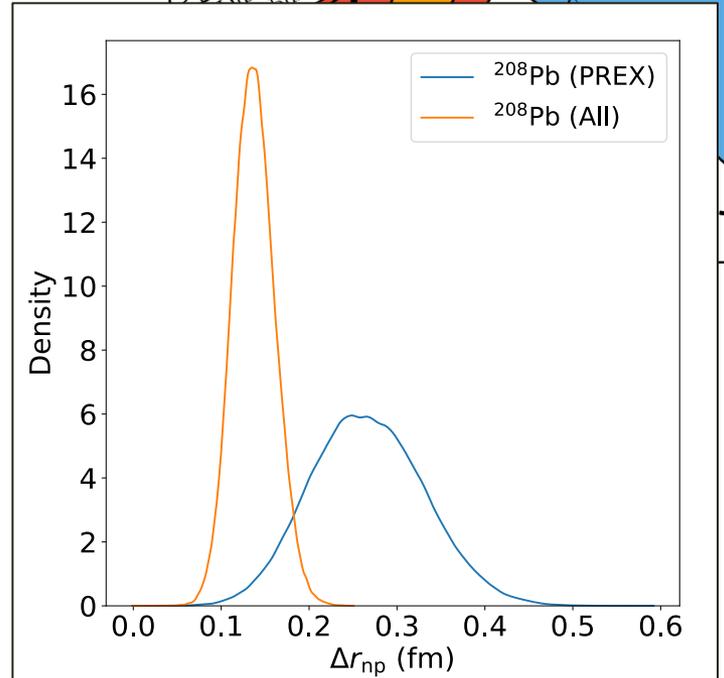
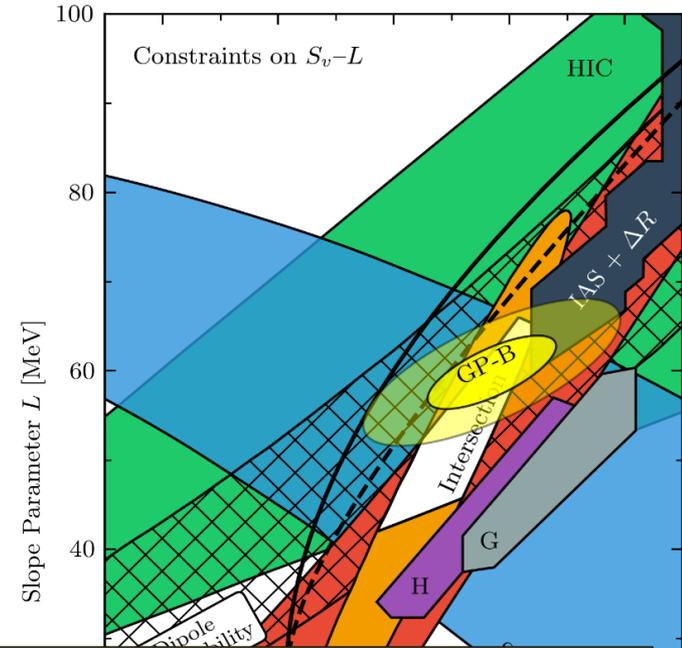
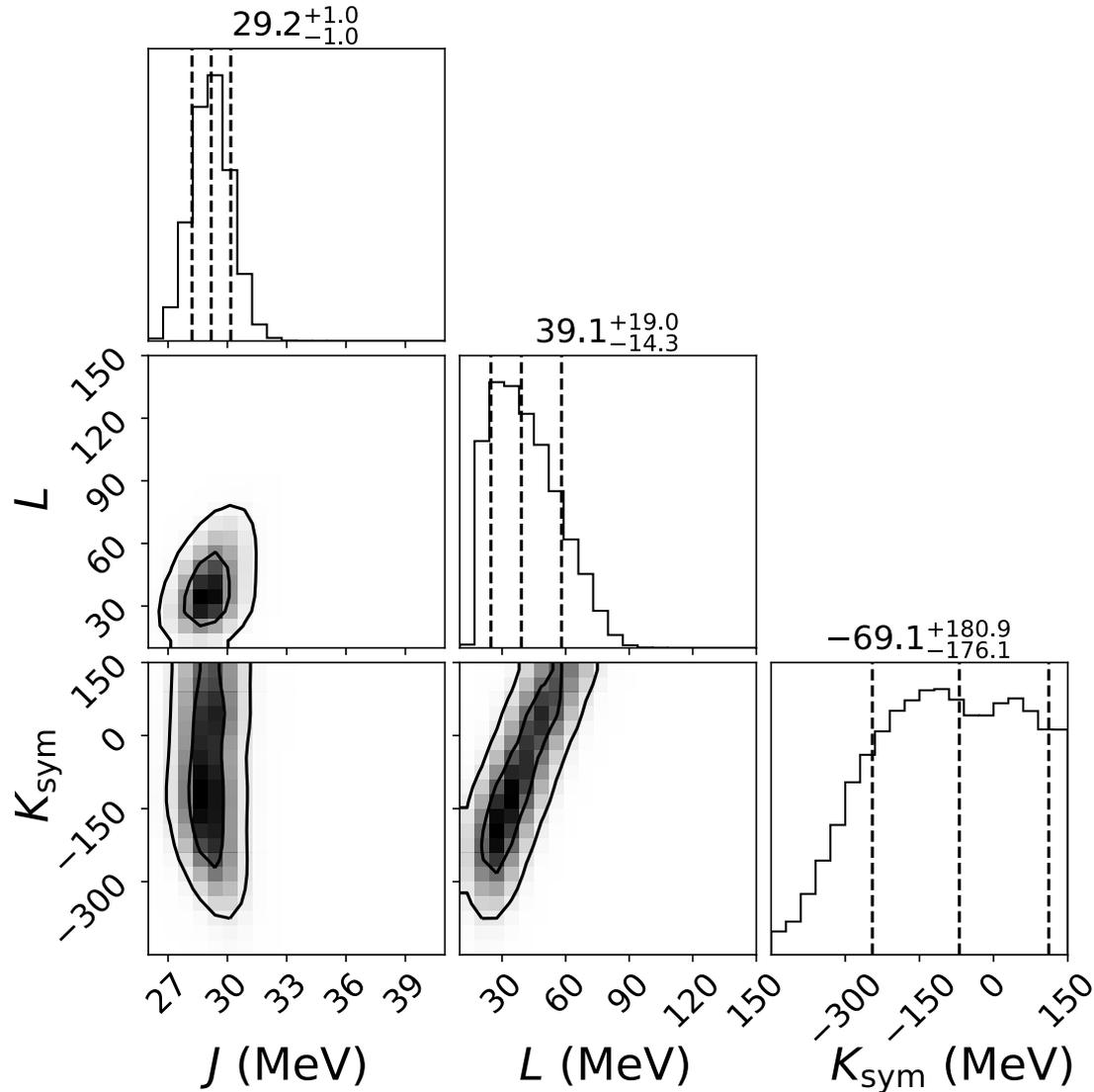
Binding energies/charge radii

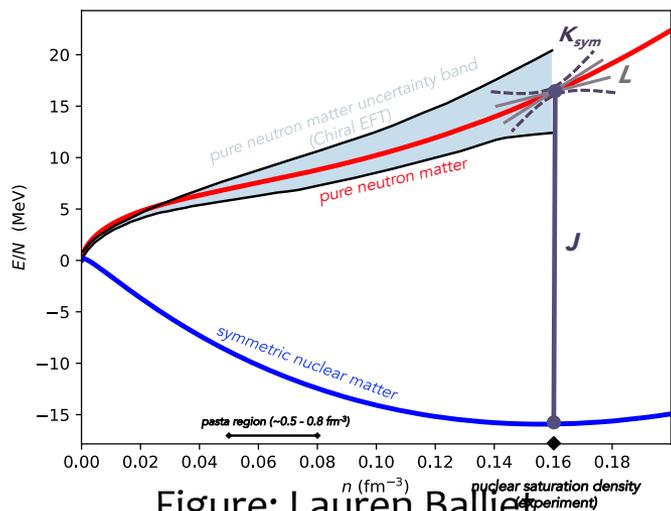


Skins (no PREX)+dip pol+BE+rp

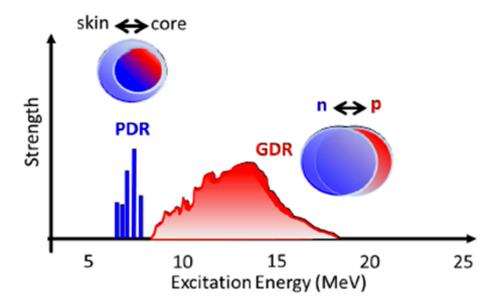
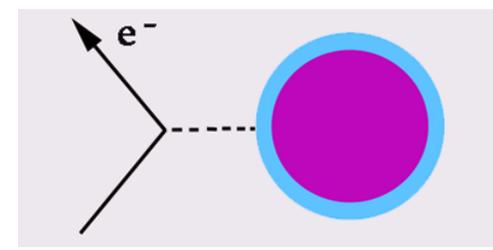
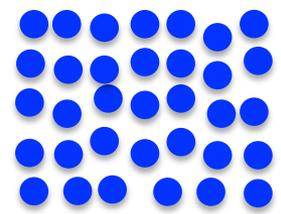


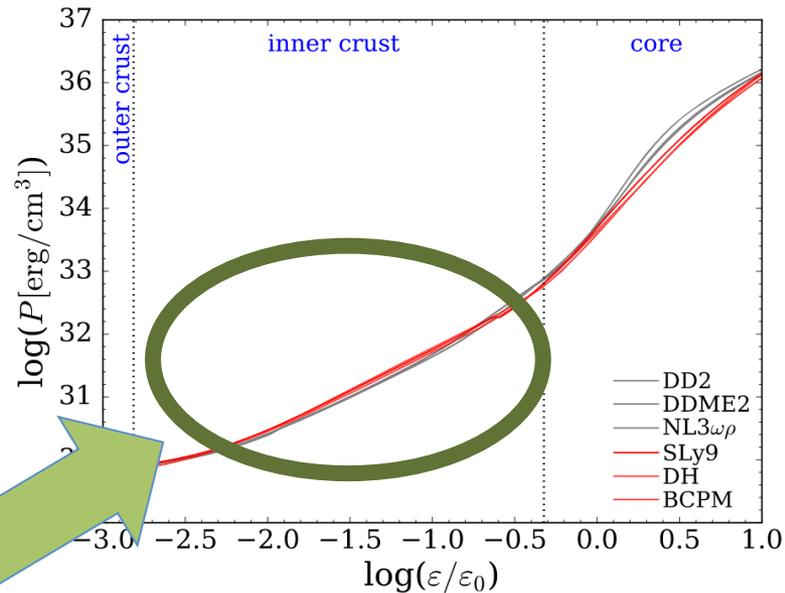
Skins (no PREX)+dip pol+BE+rp





$$P(J, L, K_{\text{sym}} | \mathcal{D})$$





DFT+CLDM

$$P(J, L, K_{\text{sym}} | \mathcal{D})$$

DFT

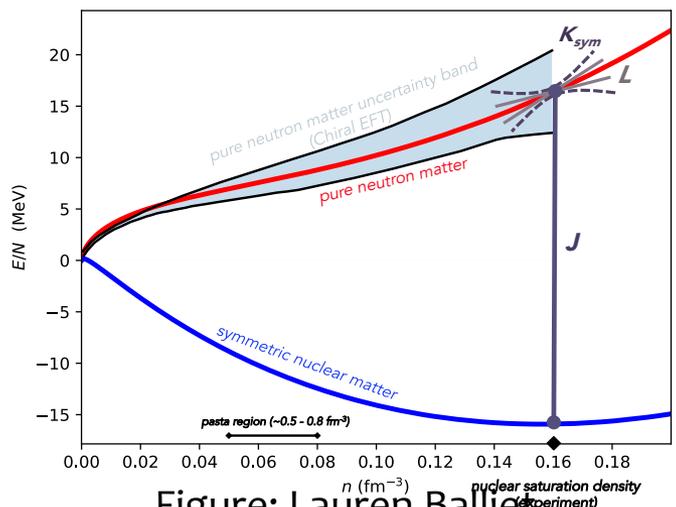
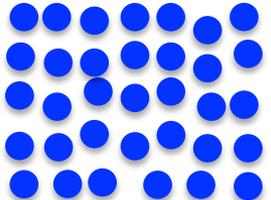
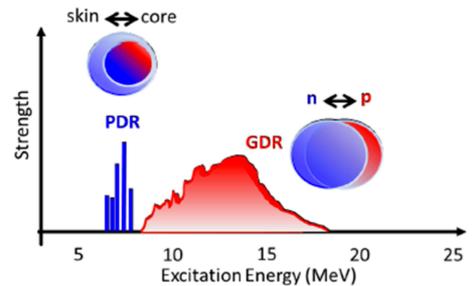
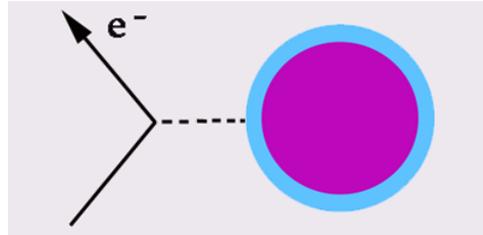


Figure: Lauren Balliet



$P(\text{Fraction of pasta} | \mathcal{D})$

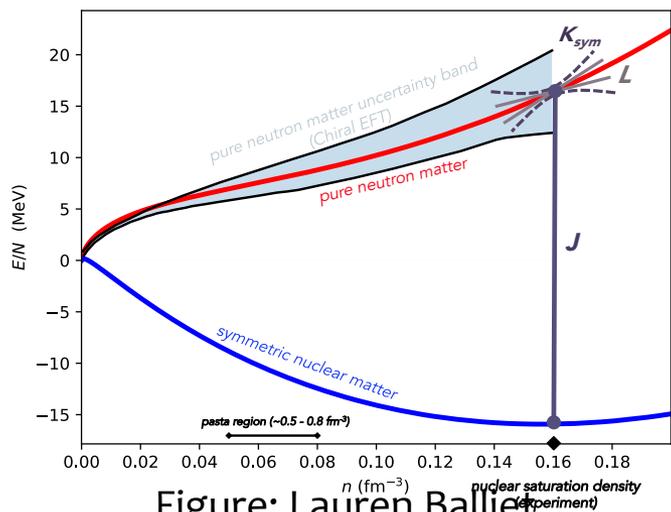
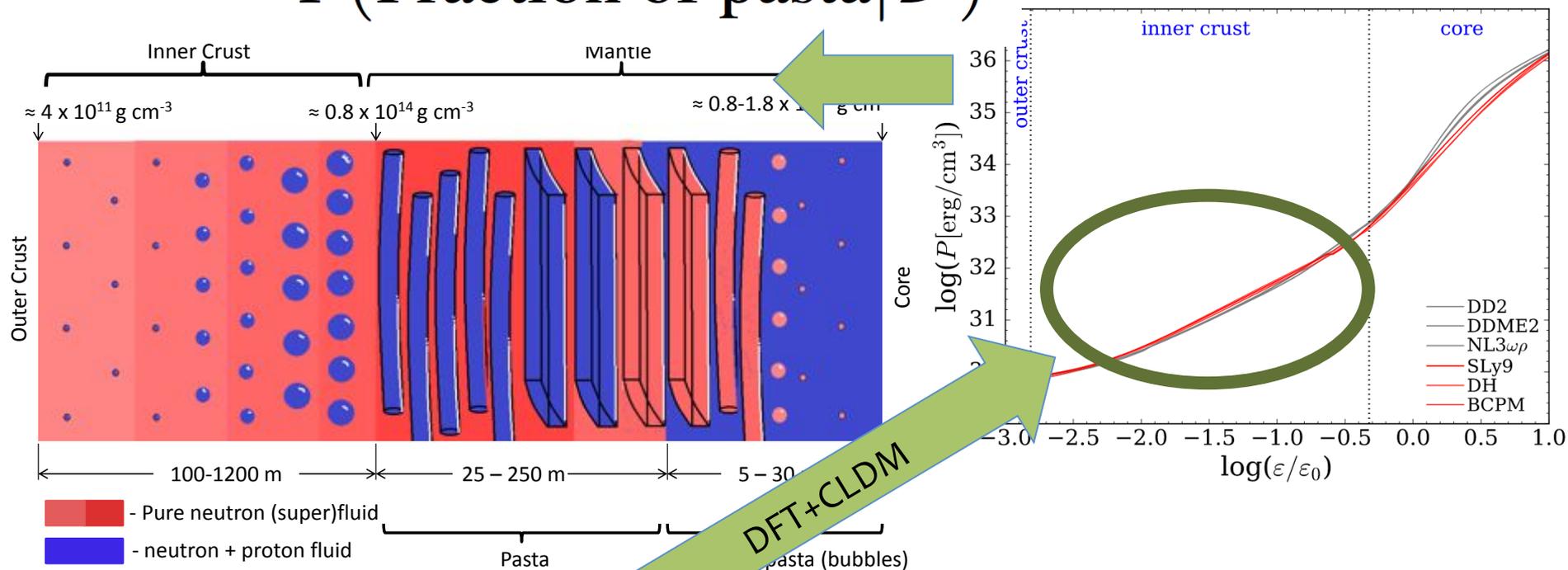
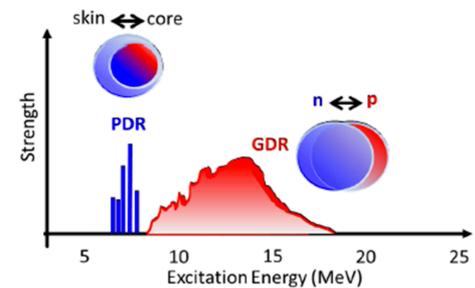
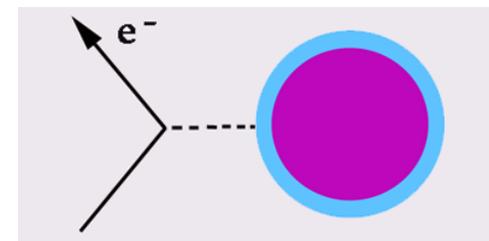
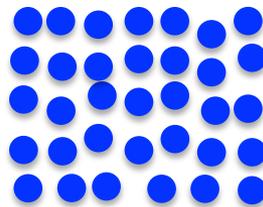


Figure: Lauren Ballot

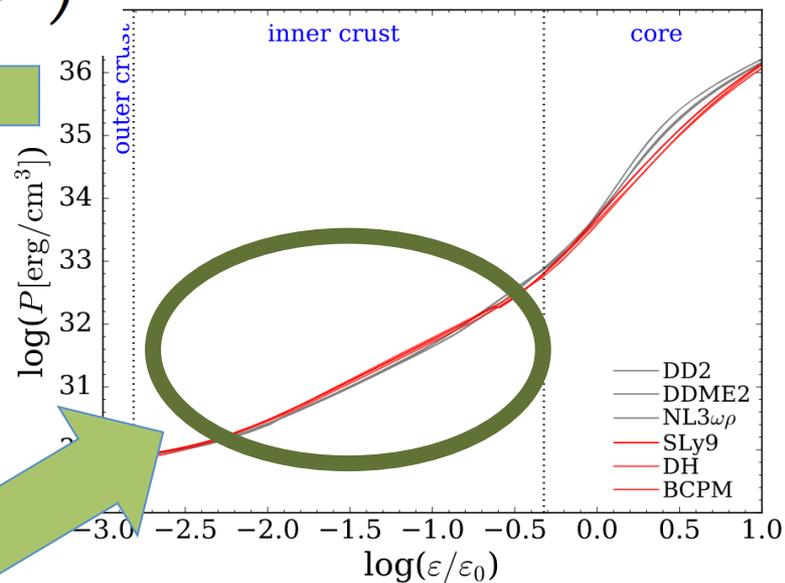
$P(J, L, K_{\text{sym}} | \mathcal{D})$



$P(\text{Fraction of pasta}|\mathcal{D})$



Noa Fritschie, 2022



DFT+CLDM

$P(J, L, K_{\text{sym}}|\mathcal{D})$

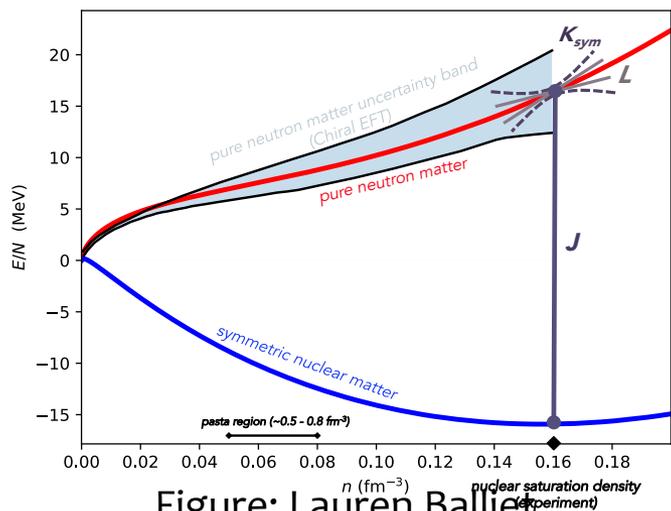
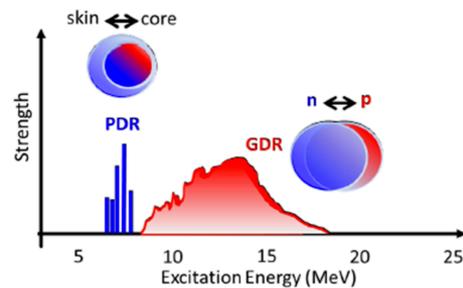
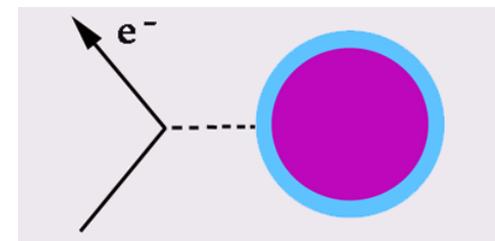
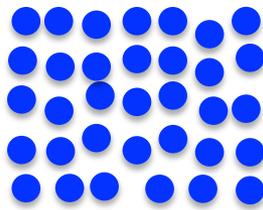


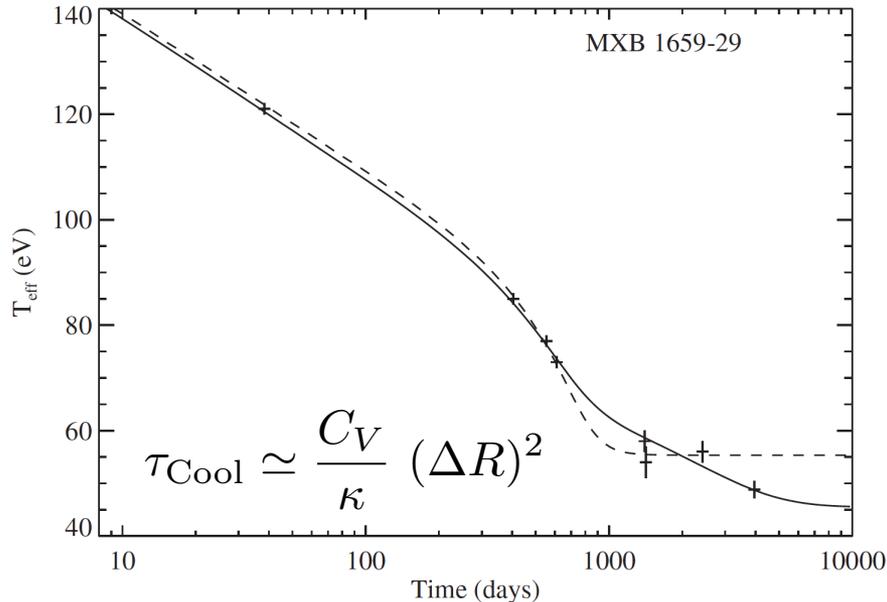
Figure: Lauren Ballett

DFT

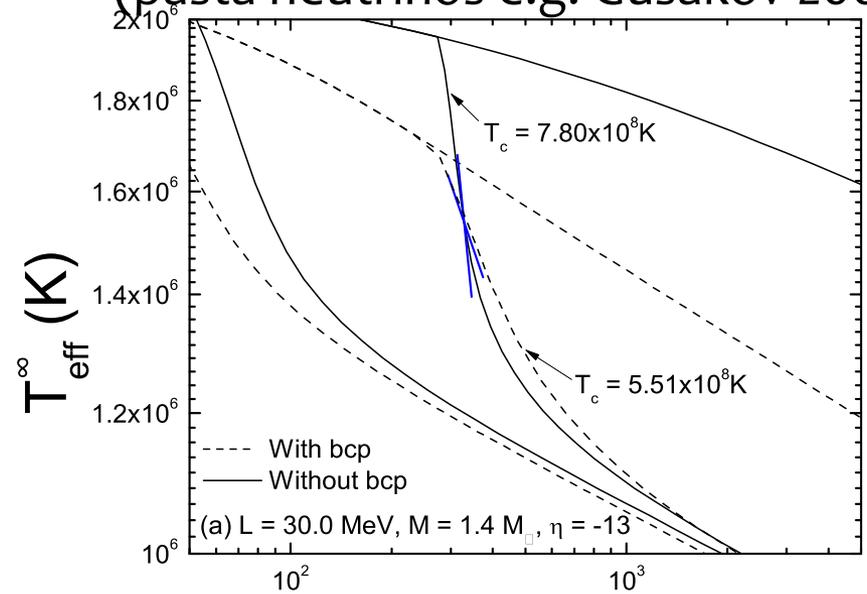


A number of ways pasta and the physics of the crust-core boundary leave signatures on observables have tentatively been put forward

Horowitz+, PRL114, 031102 (2015)
Cooling of accreted crusts



Newton+, ApJL 779, 4 (2013)
Cooling in isolated neutron stars
(pasta neutrinos e.g. Gusakov 2004)



Magnetic field evolution - Pons Nature Physics, 9, 7, 431-434 (2013)

Mountains on neutron stars - Gearheart, Newton, Li, MNRAS 418 (2011)

Crust oscillations... - Gearheart+, MNRAS 418 (2011)

... leading to resonant shattering – Neill+, MNRAS 504, 2021

Pulsar glitches – Graber+, Apj 865, 23 (2018)

Evolution of r-modes – Wen+Phys Rev C, 85, 025801 (2012)

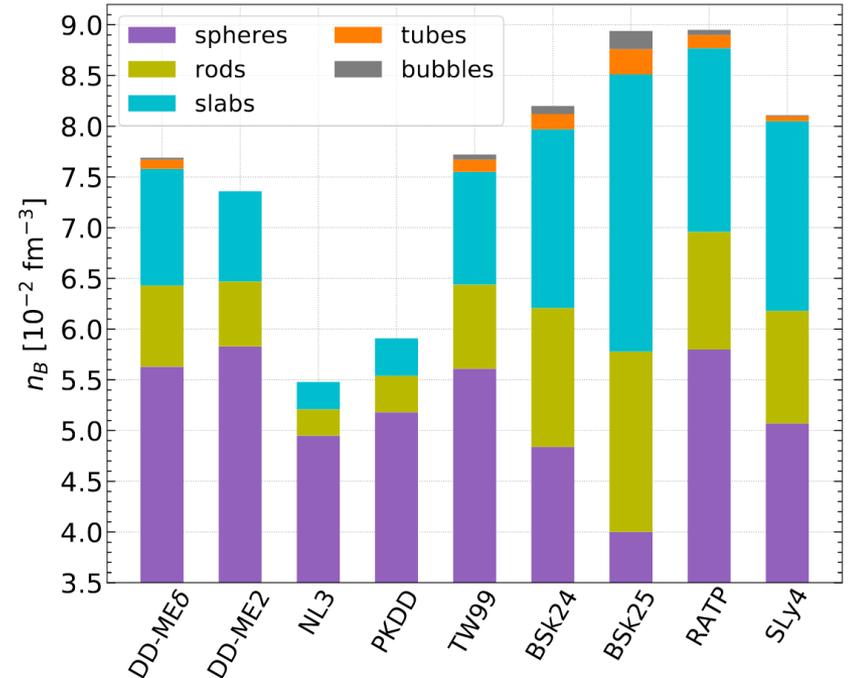
Vidana Phys Rev C 85, 045808 (2012)

The amount of crust and pasta is highly nuclear-EOS dependent

Model	ρ_{tt} [fm ⁻³]	$Y_{p,tt}$	P_{tt} [MeVfm ⁻³]	ρ_{td} [fm ⁻³]	$Y_{p,td}$	P_{td} [MeVfm ⁻³]
Microscopic						
BHF-1	0.061	0.023	0.193			
BHF-1 _{para}	0.083	0.026	0.400			
BHF-2	0.078	0.027	0.370			
BHF-2 _{para}	0.094	0.028	0.571			
Skyrme						
Bsk14	0.090	0.033	0.483	0.081	0.030	0.381
Bsk16	0.096	0.037	0.502	0.087	0.035	0.402
Bsk17	0.095	0.036	0.499	0.086	0.034	0.397
G _σ	0.063	0.013	0.278	0.054	0.010	0.172
R _σ	0.067	0.014	0.312	0.058	0.012	0.202
LNS	0.088	0.031	0.614	0.077	0.028	0.469
NRAPR	0.083	0.034	0.545	0.073	0.030	0.413
RATP	0.097	0.037	0.500	0.086	0.034	0.390
SV	0.071	0.021	0.372	0.061	0.016	0.235
SGII	0.086	0.026	0.401	0.077	0.024	0.311
SkI2	0.064	0.014	0.291	0.054	0.011	0.170
SkI3	0.071	0.022	0.363	0.062	0.018	0.244
SkI4	0.081	0.024	0.332	0.072	0.021	0.234
SkI5	0.061	0.014	0.271	0.051	0.010	0.149
SkI6	0.082	0.026	0.352	0.073	0.024	0.257
SkMP	0.072	0.020	0.357	0.062	0.017	0.241
SkO	0.073	0.020	0.413	0.062	0.017	0.270
Sly230a	0.090	0.039	0.404	0.081	0.037	0.319
Sly230b	0.089	0.038	0.462	0.080	0.036	0.362
SLy4	0.089	0.038	0.461	0.080	0.036	0.361
SLy10	0.091	0.042	0.447	0.083	0.041	0.369

Force	BBP	SKM	FPS
R	10.49	10.78	10.79
ΔM_c	0.0299 (2.07%)	0.0122 (0.84%)	0.0125 (0.86%)
ΔM_d	0.0242 (1.67%)	0.0103 (0.71%)	0.0084 (0.58%)
ΔM_n	0.0062 (0.43%)
ΔM_{dn}	0.0051 (0.35%)
I	61.56	60.89	62.57
ΔI_c	2.74 (4.45%)	1.21 (1.99%)	1.22 (1.94%)
ΔI_d	2.22 (3.60%)	1.02 (1.68%)	0.82 (1.32%)
ΔI_n	0.59 (0.94%)
ΔI_{dn}	0.48 (0.77%)

Lorenz et al PRL70 (1993)

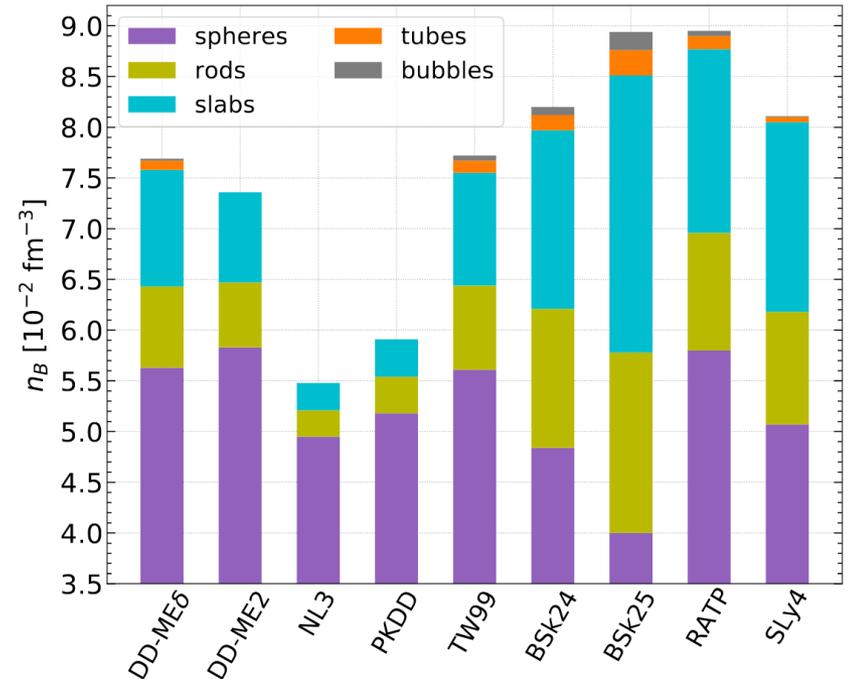


The amount of crust and pasta is highly nuclear-EOS dependent

Model	ρ_{tt} [fm ⁻³]	$Y_{p,tt}$	P_{tt} [MeVfm ⁻³]	ρ_{td} [fm ⁻³]	$Y_{p,td}$	P_{td} [MeVfm ⁻³]
Microscopic						
BHF-1	0.061	0.023	0.193			
BHF-1 _{para}	0.083	0.026	0.400			
BHF-2	0.078	0.027	0.370			
BHF-2 _{para}	0.094	0.028	0.571			
Skyrme						
Bsk14	0.090	0.033	0.483	0.081	0.030	0.381
Bsk16	0.096	0.037	0.502	0.087	0.035	0.402
Bsk17	0.095	0.036	0.499	0.086	0.034	0.397
G _σ	0.063	0.013	0.278	0.054	0.010	0.172
R _σ	0.067	0.014	0.312	0.058	0.012	0.202
LNS	0.088	0.031	0.614	0.077	0.028	0.469
NRAPR	0.083	0.034	0.545	0.073	0.030	0.413
RATP	0.097	0.037	0.500	0.086	0.034	0.390
SV	0.071	0.021	0.372	0.061	0.016	0.235
SGII	0.086	0.026	0.401	0.077	0.024	0.311
SkI2	0.064	0.014	0.291	0.054	0.011	0.170
SkI3	0.071	0.022	0.363	0.062	0.018	0.244
SkI4	0.081	0.024	0.332	0.072	0.021	0.234
SkI5	0.061	0.014	0.271	0.051	0.010	0.149
SkI6	0.082	0.026	0.352	0.073	0.024	0.257
SkMP	0.072	0.020	0.357	0.062	0.017	0.241
SkO	0.073	0.020	0.413	0.062	0.017	0.270
Sly230a	0.090	0.039	0.404	0.081	0.037	0.319
Sly230b	0.089	0.038	0.462	0.080	0.036	0.362
SLy4	0.089	0.038	0.461	0.080	0.036	0.361
SLy10	0.091	0.042	0.447	0.083	0.041	0.369

Force	BBP	SKM	FPS
R	10.49	10.78	10.70
ΔM_c	0.0299 (2.07%)	0.0122 (0.84%)	0.0125 (0.86%)
ΔM_d	0.0242 (1.67%)	0.0103 (0.71%)	0.0084 (0.58%)
ΔM_n	0.0062 (0.43%)
ΔM_{dn}	0.0051 (0.35%)
I	61.56	60.89	62.57
ΔI_c	2.74 (4.45%)	1.21 (1.99%)	1.22 (1.94%)
ΔI_d	2.22 (3.60%)	1.02 (1.68%)	0.82 (1.32%)
ΔI_n	0.59 (0.94%)
ΔI_{dn}	0.48 (0.77%)

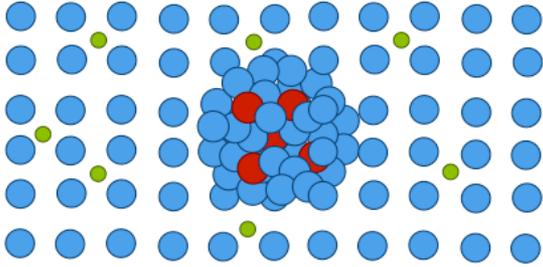
Lorenz et al PRL70 (1993)



Modeling the crust

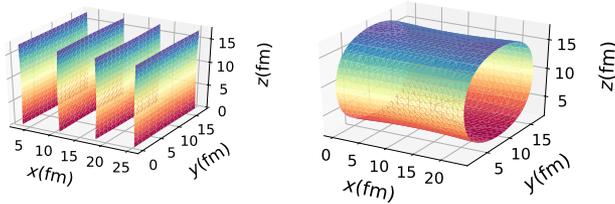
3D Skyrme HF:

n,p degrees of freedom



Newton+ arxiv:2104.11835

Pictures: Lauren Balliet



$$\mathcal{H}_\delta + \mathcal{H}_\rho + \mathcal{H}_{\text{eff}} + \mathcal{H}_{\text{grad}} + \mathcal{H}_{\text{Coul}}$$

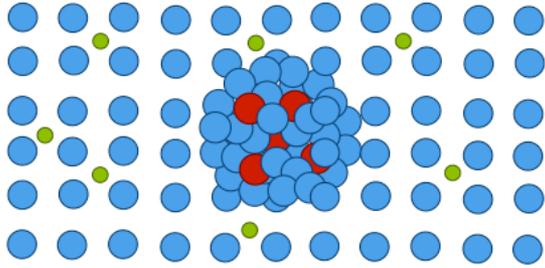
Nuclear EDF: Bulk+Gradient

Specific model: Skyrme

$$\begin{aligned} \mathcal{H}_\rho = & \frac{1}{4} t_3 \rho^{2+\alpha_3} [(2 + x_3) - (2x_3 + 1)(y_p^2 + y_n^2)] \\ & + \frac{1}{4} t_4 \rho^{2+\alpha_4} [(2 + x_4) - (2x_4 + 1)(y_p^2 + y_n^2)] \end{aligned}$$

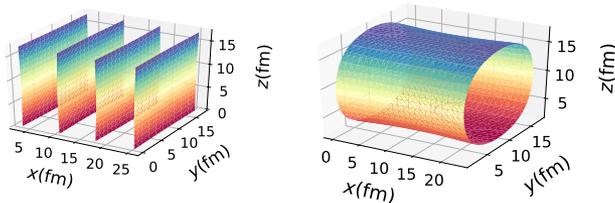
Modeling the crust

3D Skyrme HF:
n,p degrees of freedom



Newton+ arxiv:2104.11835

Pictures: Lauren Balliet



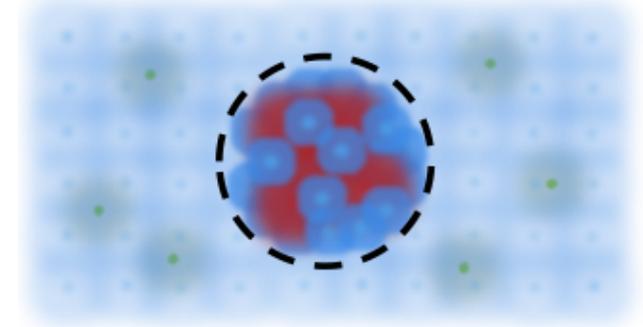
$$\mathcal{H}_\delta + \mathcal{H}_\rho + \mathcal{H}_{\text{eff}} + \mathcal{H}_{\text{grad}} + \mathcal{H}_{\text{Coul}}$$

Nuclear EDF: Bulk+Gradient
Specific model: Skyrme

$$\mathcal{H}_\rho = \frac{1}{4} t_3 \rho^{2+\alpha_3} [(2 + x_3) - (2x_3 + 1)(y_p^2 + y_n^2)]$$

$$+ \frac{1}{4} t_4 \rho^{2+\alpha_4} [(2 + x_4) - (2x_4 + 1)(y_p^2 + y_n^2)]'$$

CLDM: Bulk fluid and surface
degrees of freedom



Newton et al arxiv: 1110.4043
Balliet+; arxiv:2009.07696

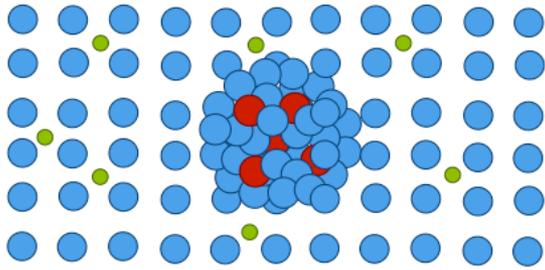
$$\mathcal{H}_\delta + \mathcal{H}_\rho + \mathcal{H}_{\text{eff}} \quad \sigma(y_p)$$

Nuclear EDF: Bulk +
separate surface energy function
specific model: LLPR 1985

$$\sigma_s(y_p) = \sigma_0 \frac{2^{p+1} + b}{\frac{1}{y_p^p} + b + \frac{1}{(1-y_p)^p}}$$

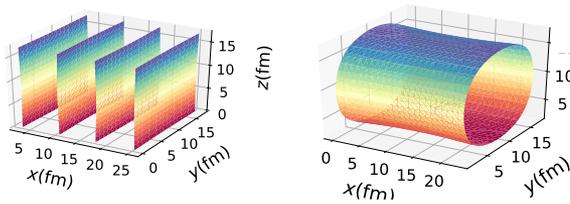
Modeling the crust

3D Skyrme HF:
n,p degrees of freedom

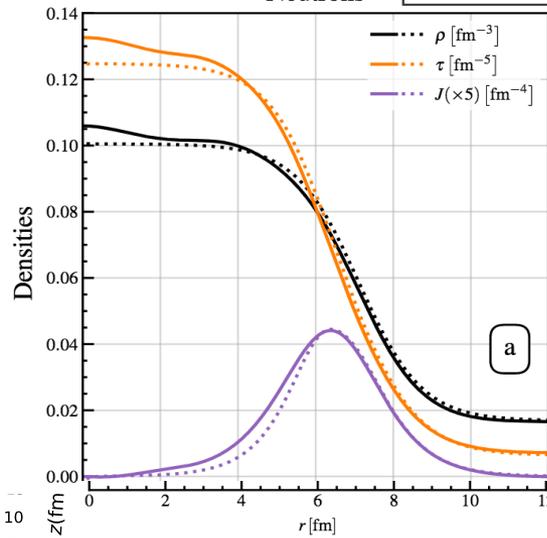


Newton+ arxiv:2104.11835

Pictures: Lauren Balliet

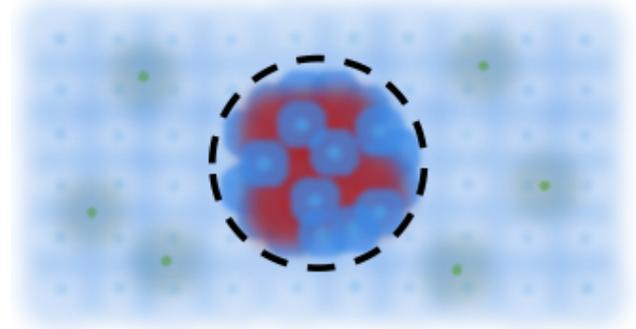


Thomas-Fermi: density profile
degree of freedom



Shelley, arxiv:2011.10511

CLDM: Bulk fluid and surface
degrees of freedom



Newton et al arxiv: 1110.4043
Balliet+; arxiv:2009.07696

$$\mathcal{H}_\delta + \mathcal{H}_\rho + \mathcal{H}_{\text{eff}} + \mathcal{H}_{\text{grad}} + \mathcal{H}_{\text{Coul}}$$

Nuclear EDF: Bulk+Gradient
Specific model: Skyrme

$$\mathcal{H}_\rho = \frac{1}{4} t_3 \rho^{2+\alpha_3} [(2 + x_3) - (2x_3 + 1)(y_p^2 + y_n^2)] \\ + \frac{1}{4} t_4 \rho^{2+\alpha_4} [(2 + x_4) - (2x_4 + 1)(y_p^2 + y_n^2)]'$$

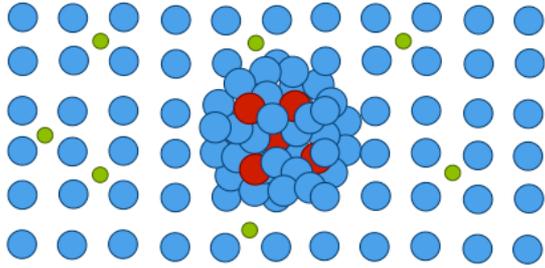
$$\mathcal{H}_\delta + \mathcal{H}_\rho + \mathcal{H}_{\text{eff}} \quad \sigma(y_p)$$

Nuclear EDF: Bulk +
separate surface energy function
specific model: LLPR 1985

$$\sigma_s(y_p) = \sigma_0 \frac{2^{p+1} + b}{\frac{1}{y_p^p} + b + \frac{1}{(1-y_p)^p}}$$

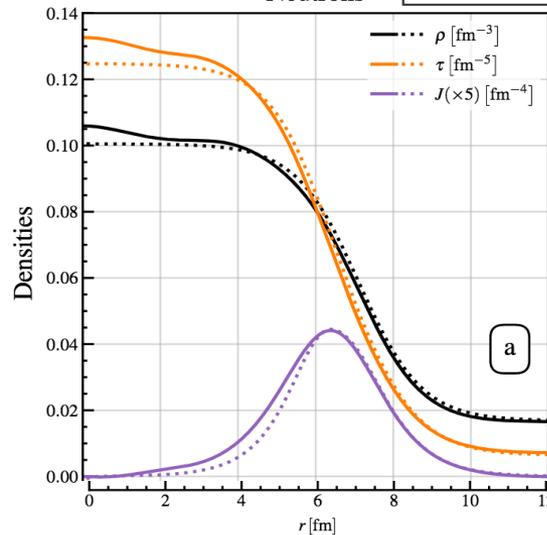
Modeling the crust

3D Skyrme HF:
n,p degrees of freedom



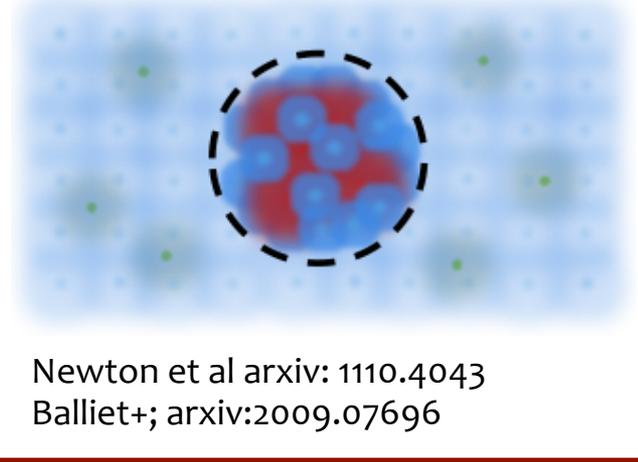
Newton+ arxiv:2104.11835
Pictures: Lauren Balliet

Thomas-Fermi: density profile
degree of freedom



Shelley, arxiv:2011.10511

CLDM: Bulk fluid and surface
degrees of freedom



Newton et al arxiv: 1110.4043
Balliet+; arxiv:2009.07696

$$\mathcal{H}_\delta + \mathcal{H}_\rho + \mathcal{H}_{\text{eff}} + \mathcal{H}_{\text{grad}} + \mathcal{H}_{\text{Coul}}$$

Nuclear EDF: Bulk+Gradient
Specific model: Skyrme

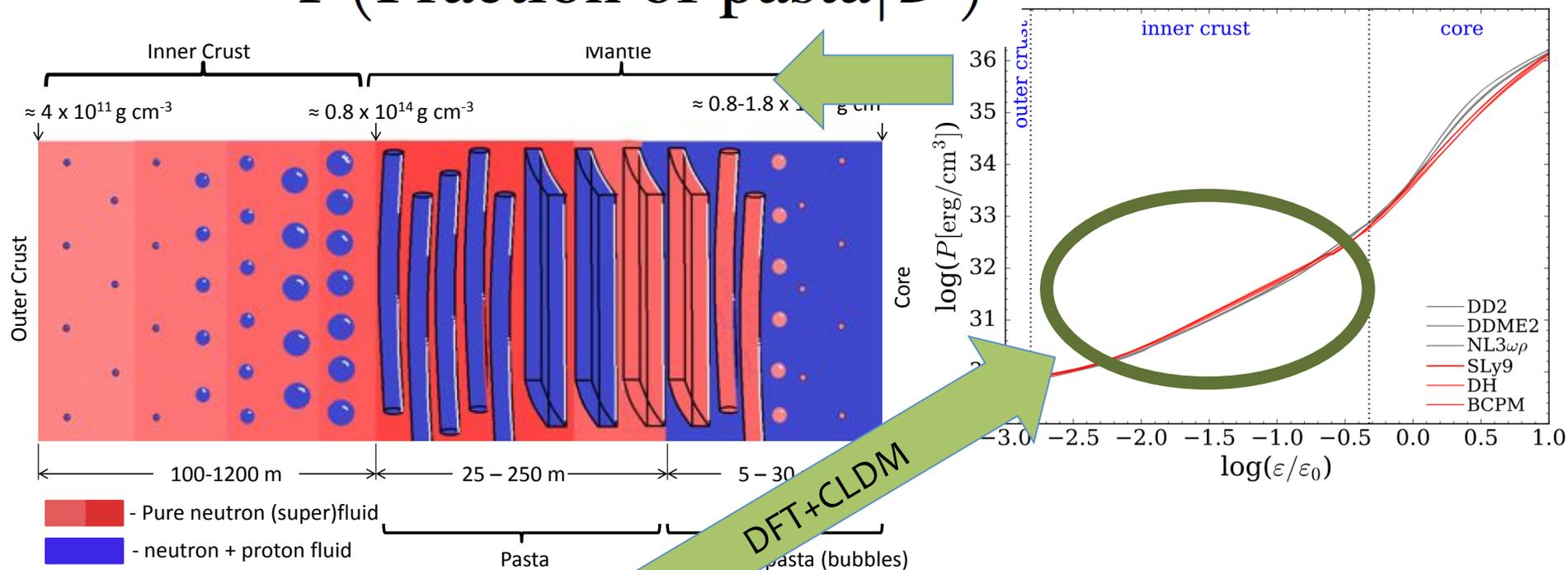
$$\mathcal{H}_\rho = \frac{1}{4} t_3 \rho^{2+\alpha_3} [(2 + x_3) - (2x_3 + 1)(y_p^2 + y_n^2)] \\ + \frac{1}{4} t_4 \rho^{2+\alpha_4} [(2 + x_4) - (2x_4 + 1)(y_p^2 + y_n^2)]'$$

$$\mathcal{H}_\delta + \mathcal{H}_\rho + \mathcal{H}_{\text{eff}} \quad \sigma(y_p)$$

Nuclear EDF: Bulk +
separate surface energy function
specific model: LLPR 1985

$$\sigma_s(y_p) = \sigma_0 \frac{2^{p+1} + b}{\frac{1}{y_p^p} + b + \frac{1}{(1-y_p)^p}}$$

$P(\text{Fraction of pasta} | \mathcal{D})$



DFT+CLDM

$P(J, L, K_{\text{sym}} | \mathcal{D})$

DFT

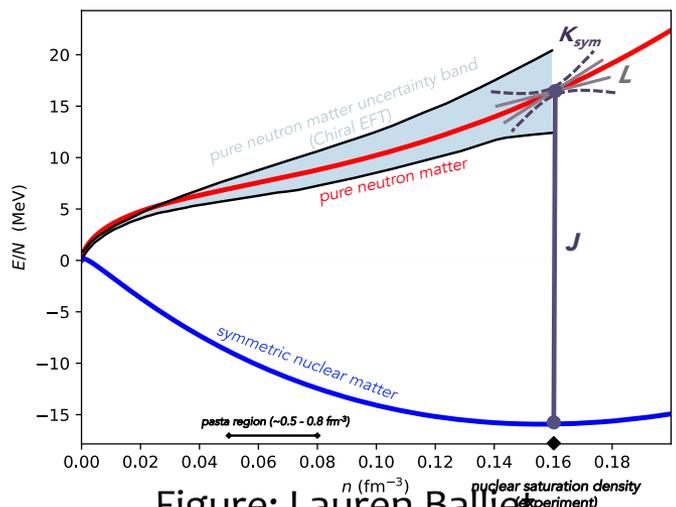
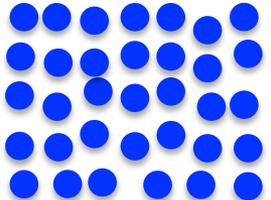
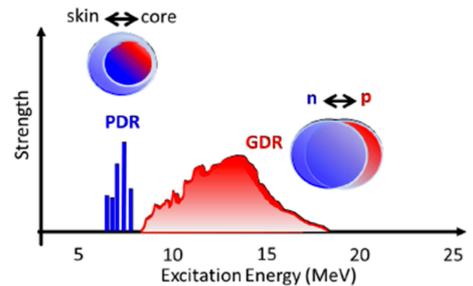
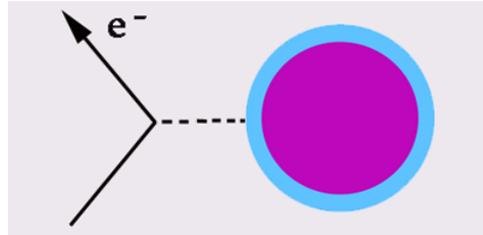
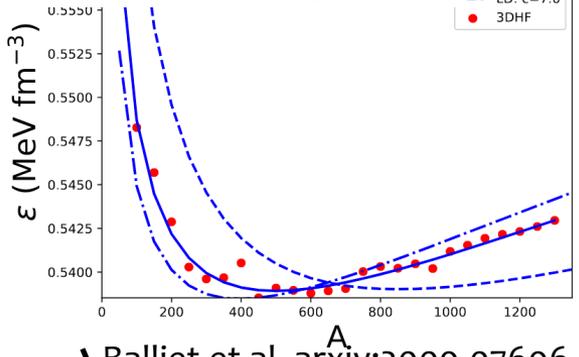
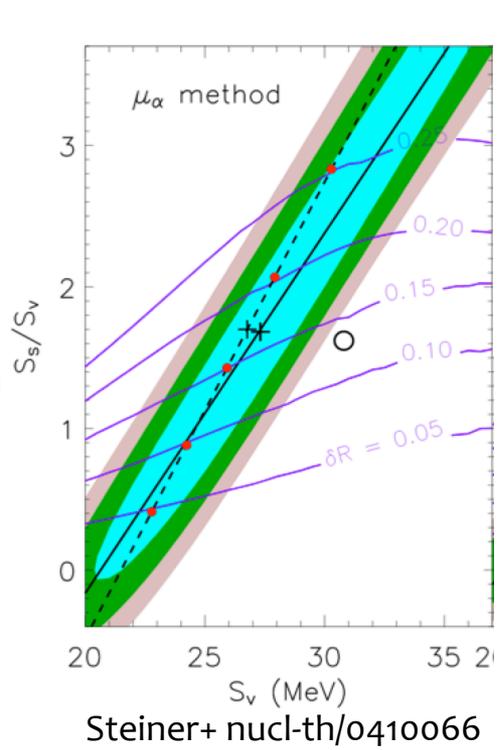


Figure: Lauren Balliet



Correlations from mass fits, fits to 3DHF



Balliet et al, arxiv:2009.07696

$$P(\sigma_0, \sigma_c, c, p)$$

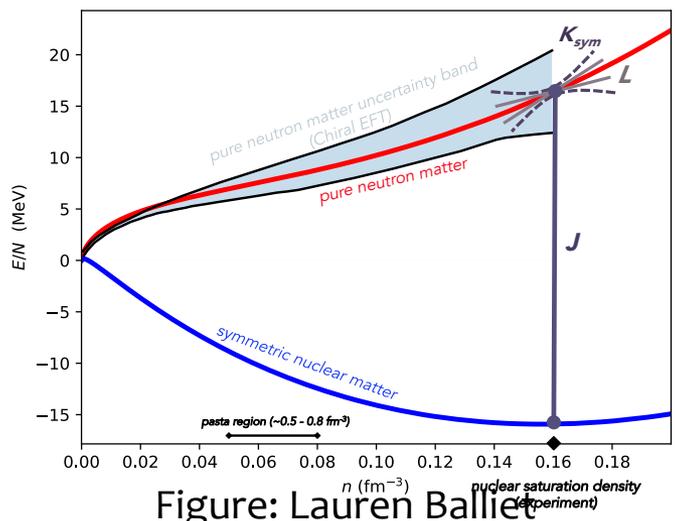
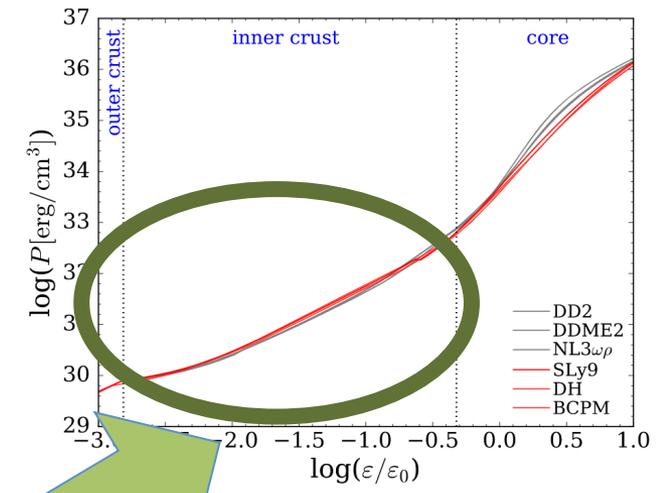
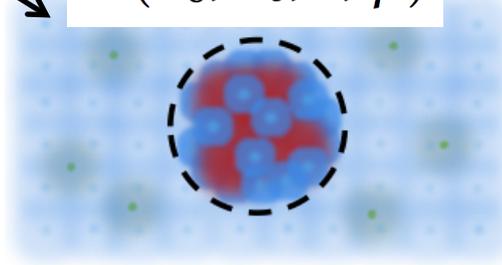
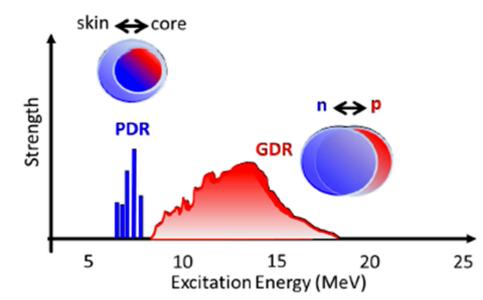
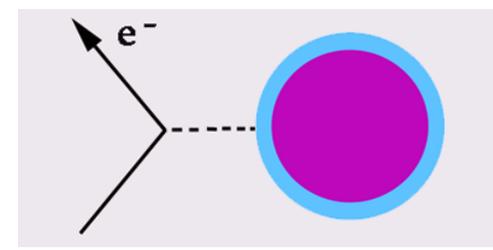
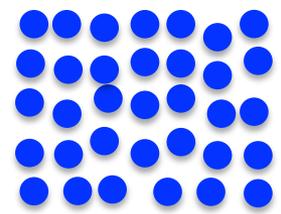


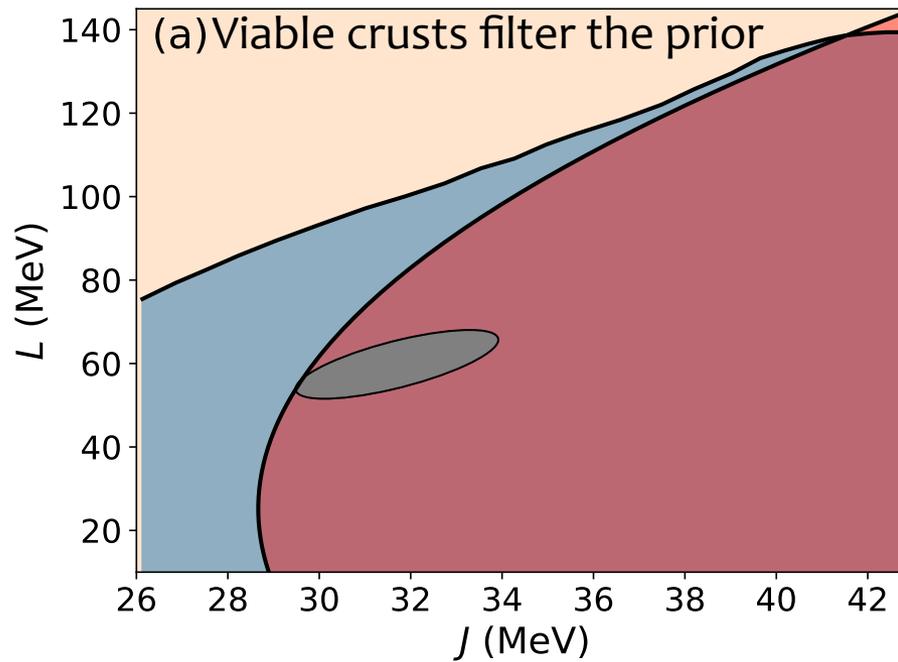
Figure: Lauren Balliet

$$P(J, L, K_{sym} | \mathcal{D})$$

DFT



Priors: crust models



Prior	σ_0	p	c	$\sigma_{0,c}$
Surf:Unif	0.8-1.3	2.0-4.0	2.0-7.0	0.0-1.0
Surf:Fit	1.1	3.8 ± 0.02	4.5	0.6

Newton et al arxiv: 1110.4043

Balliet+; arxiv:2009.07696

(Unitary gas bounds: Tews et al, arxiv:1611.07133 PNM: Drischler, arxiv:2004.07232)

Potential sources of systematic error

Had to choose an EDF (Skyrme). Enough degrees of freedom?
Can add more (Q_{sym})

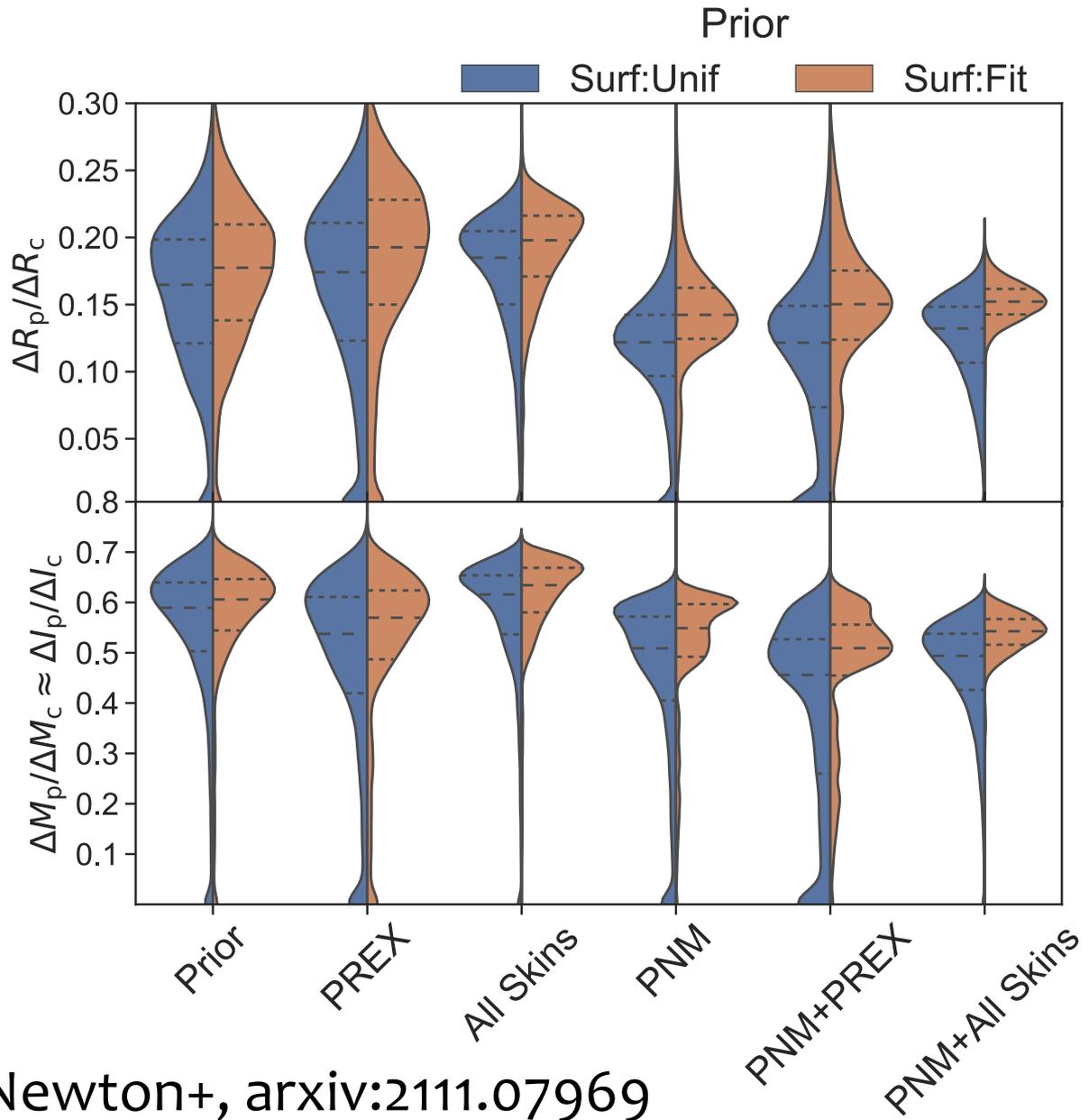
Symmetric nuclear matter and gradient parameters held fixed; extending inference to those parameters may change posteriors

We're usually not directly modeling nuclear observables - but in some cases we can (e.g. weak form factor) and thus improve consistency

We've chosen a model of crust (CLDM) different to modeling of the nuclei (1D SHF+RPA)

Need more information to constrain surface parameters of crust model (mass fits, semi-infinite nuclear matter, Thomas-Fermi calculations)

Results: relative thickness and mass of pasta



$$\frac{\Delta R_p}{\Delta R_c} \approx \frac{\mu_{cc} - \mu_p}{\mu_{cc} - \mu_0}$$

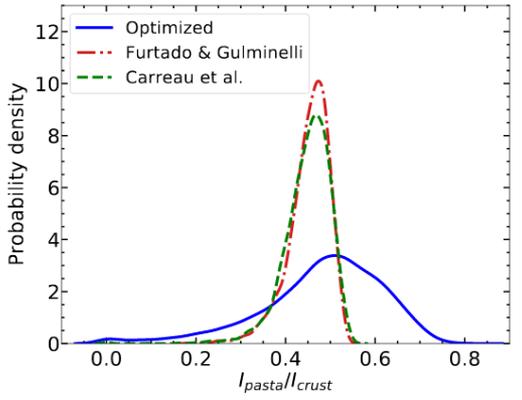
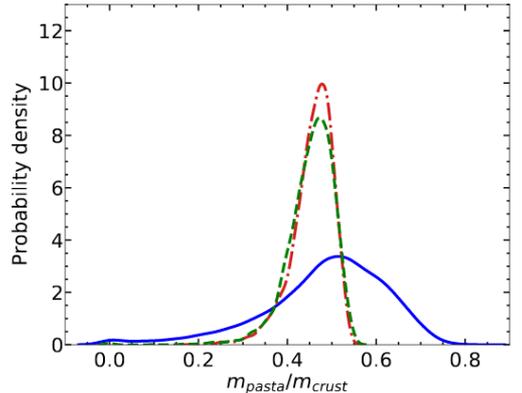
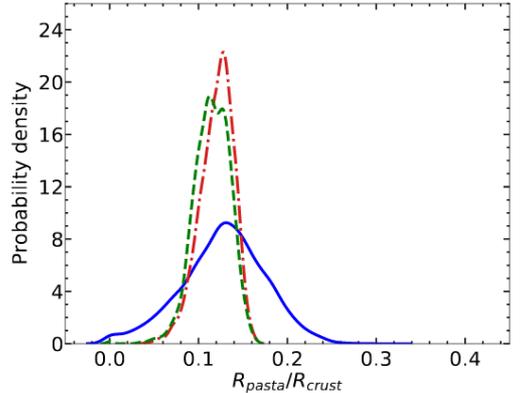
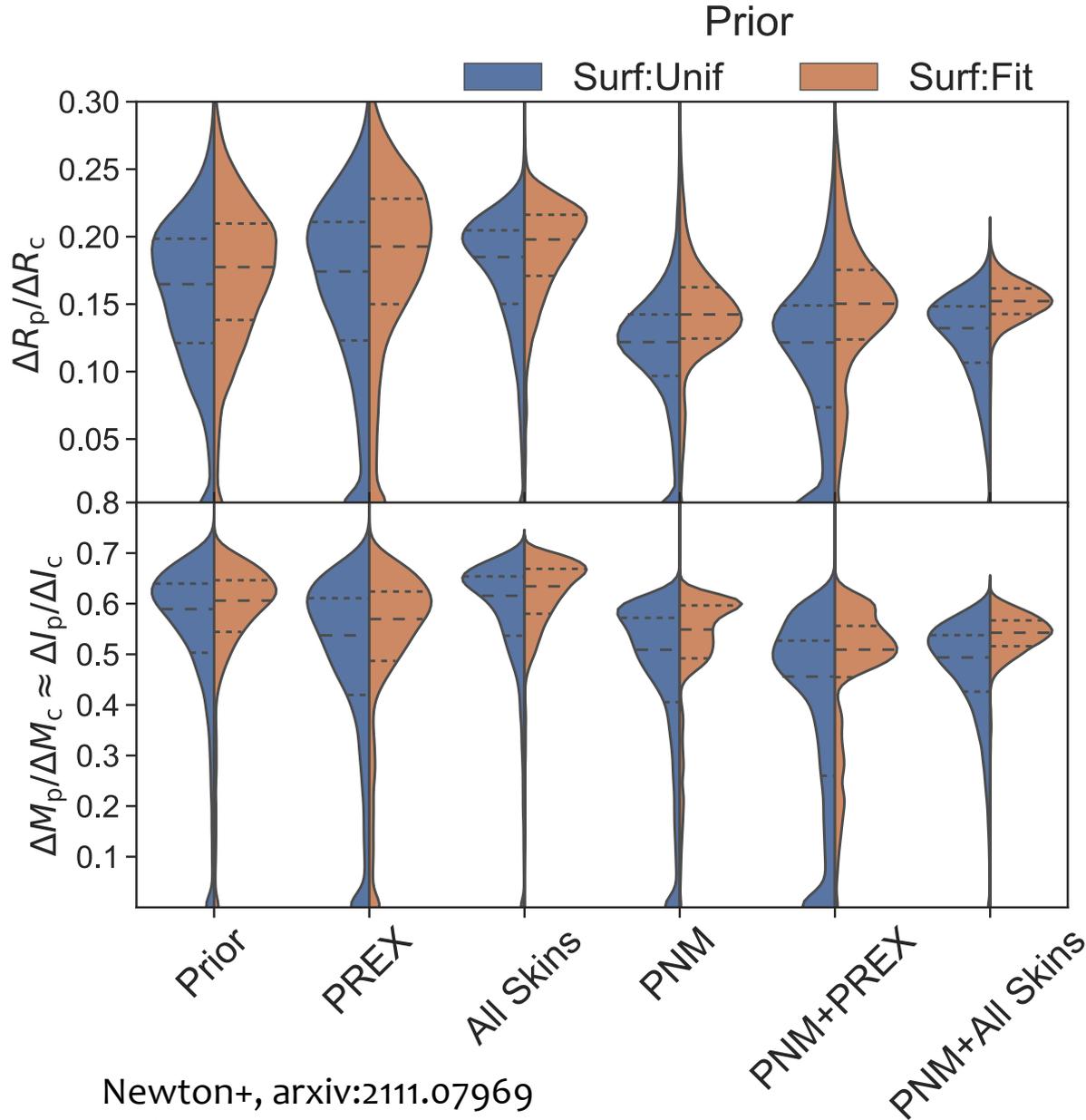
$$\frac{\Delta M_p}{\Delta M_c} \approx \frac{P_p}{P_{cc}}$$

Combining our best experimental and computational data:

$$\Delta R_p / \Delta R_c = 0.132^{+0.023}_{-0.041}$$

$$\Delta M_p / \Delta M_c \approx \Delta I_p / \Delta I_c = 0.49^{+0.06}_{-0.11}$$

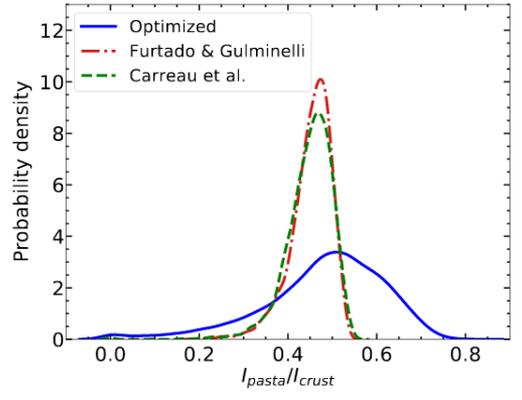
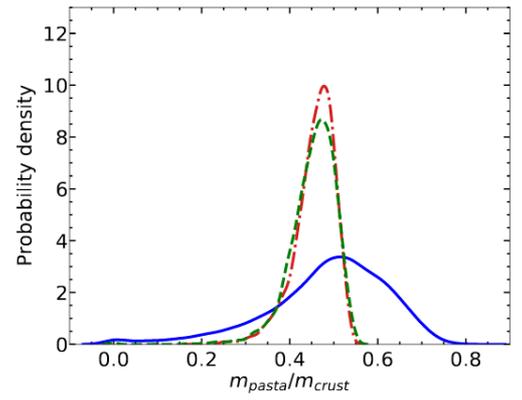
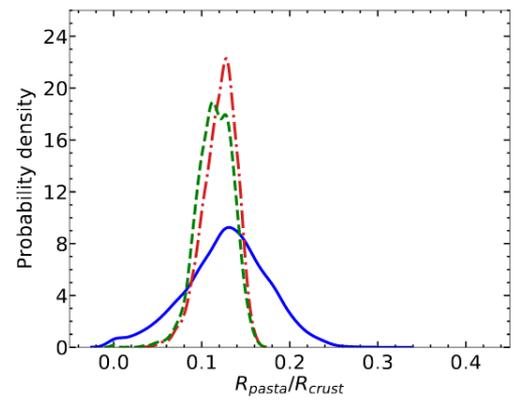
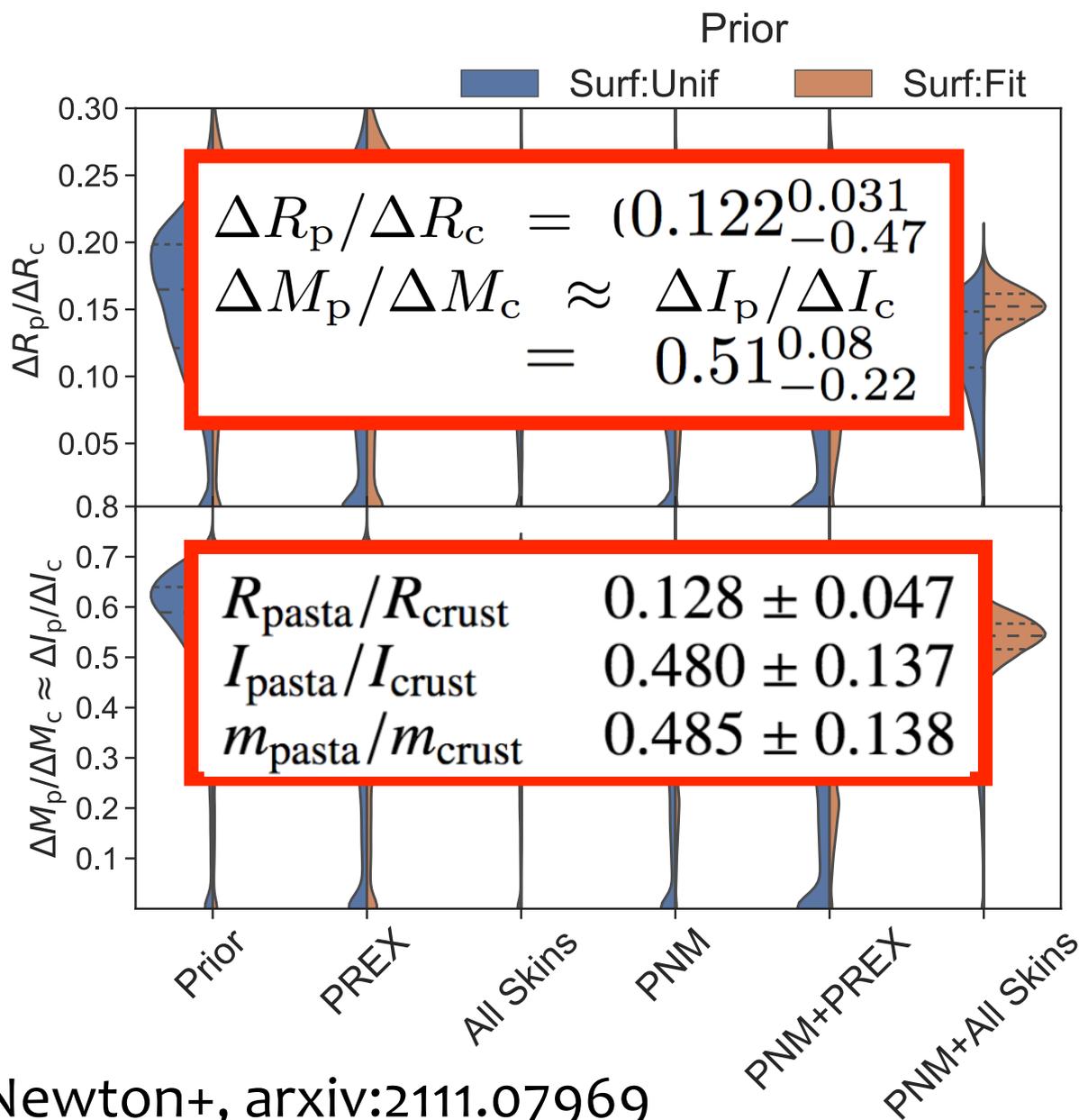
Relative thickness and mass of pasta: agreement with other studies



Newton+, arxiv:2111.07969
 Balliet+, arxiv:2009.07696

Dinh Thi+ arxiv: 2109.13638

Relative thickness and mass of pasta: agreement with other studies

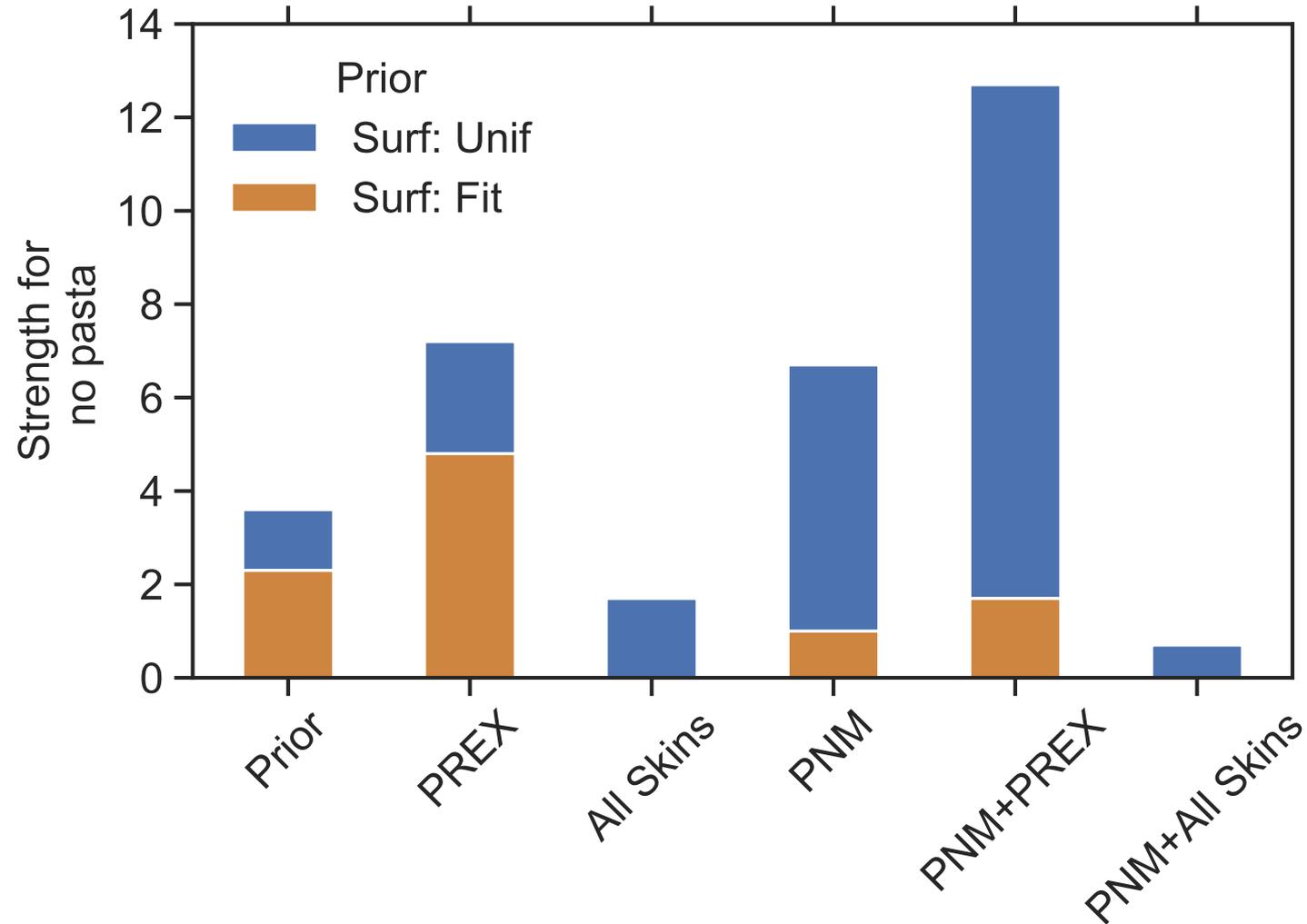


Newton+, arxiv:2111.07969

Balliet+, arxiv:2009.07696

Dinh Thi+ arxiv: 2109.13638

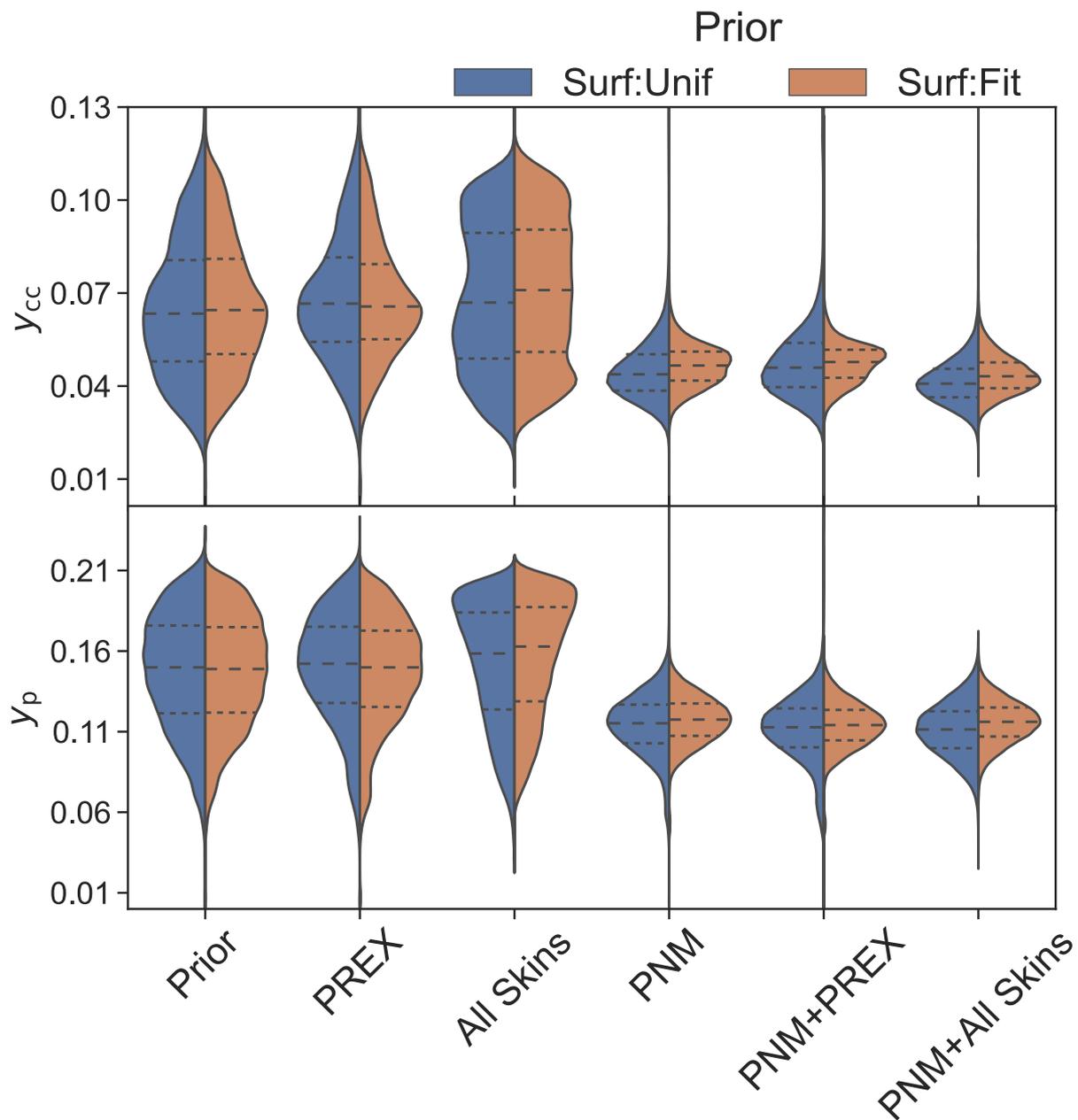
There's a non-negligible range of models that predicts no pasta



Newton+, arxiv:2111.07969

Balliet+, arxiv:2009.07696

Proton fractions



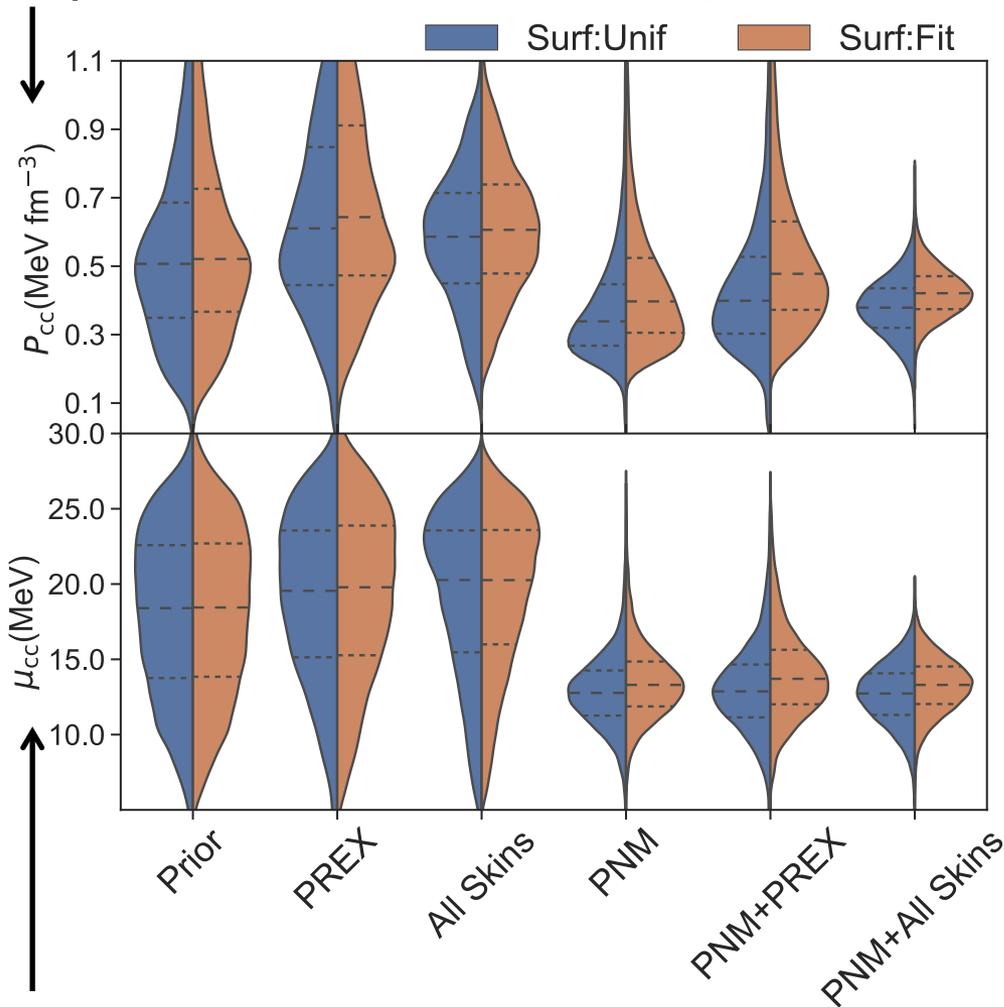
$$y_p = 0.111^{+0.017}_{-0.17}$$

$$y_{cc} = 0.041^{+0.007}_{-0.006}$$

Crust-core transition pressure and chemical potential

Responsible for crust *mass*

Prior



$$P_{cc} = 0.38(0.42)^{+0.08(0.07)}_{-0.09(0.07)}$$

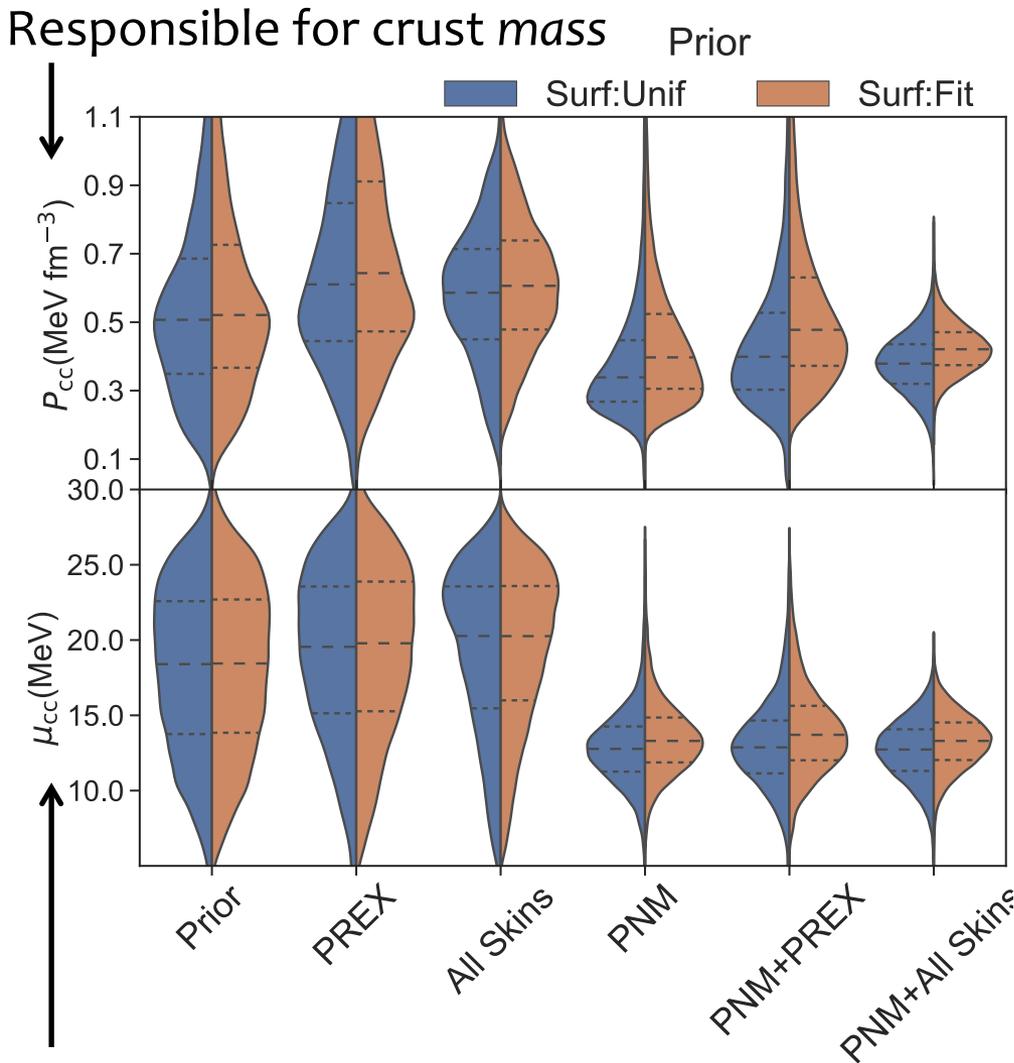
$$\mu_{cc} = 12.7(13.3)^{+2.0(1.8)}_{-2.1(1.9)}$$

Responsible for crust *thickness*

Newton+, arxiv:2111.07969

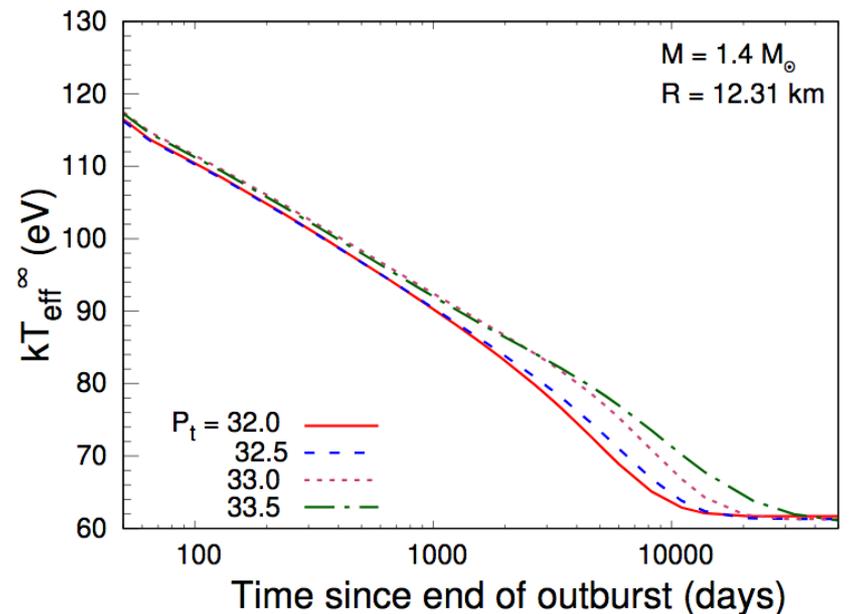
Balliet+, arxiv:2009.07696

Crust-core transition pressure and chemical potential



Lalit+ arxiv:1906.01535

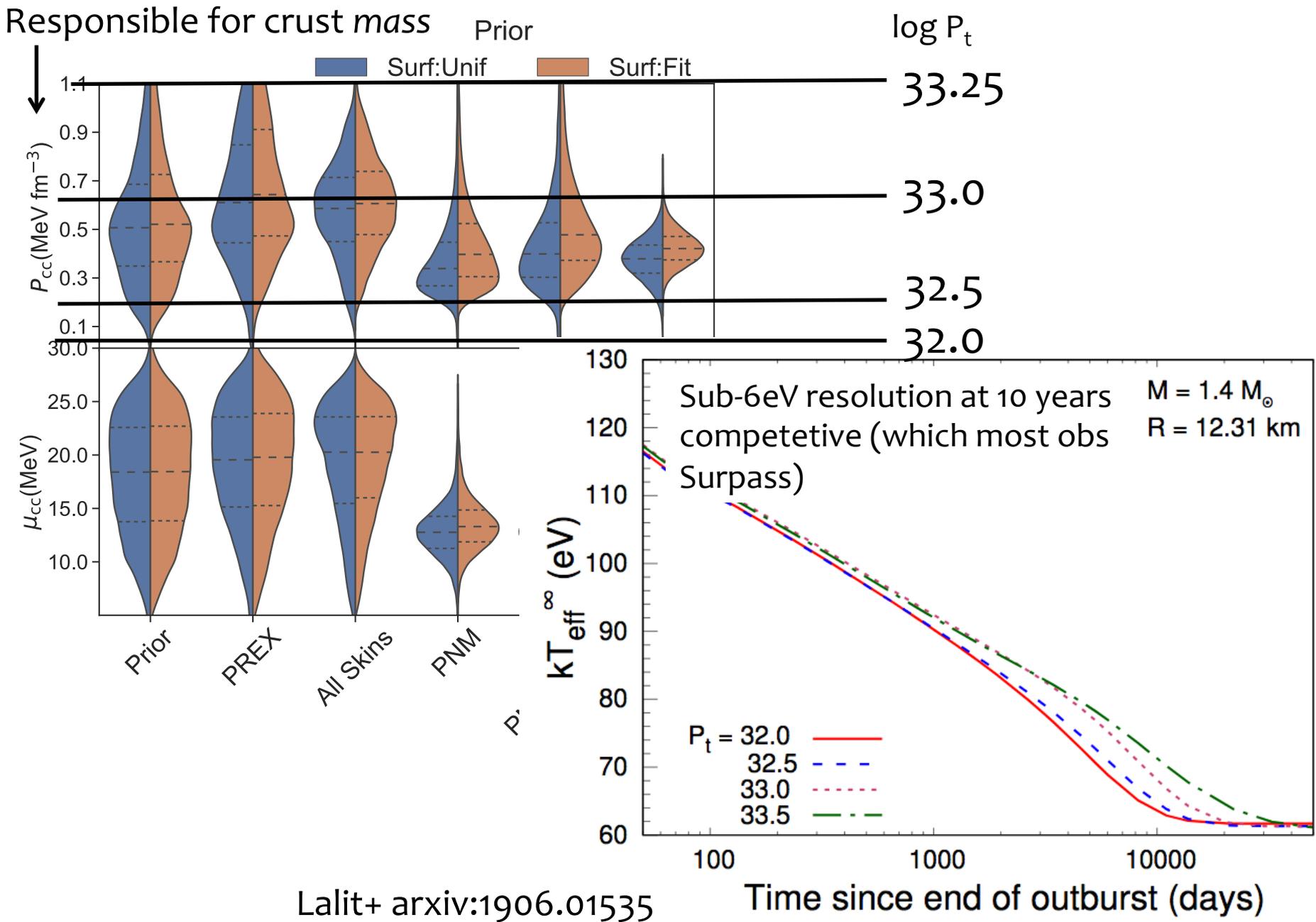
How much does the cooling light curve care about the base of the crust?



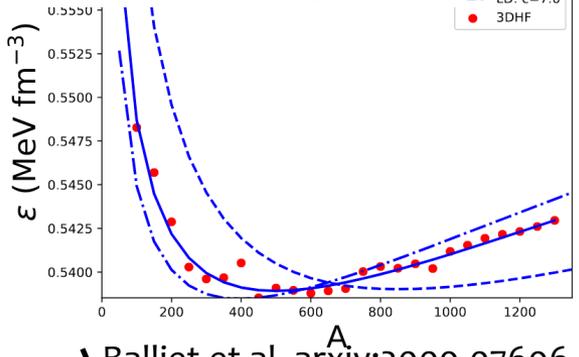
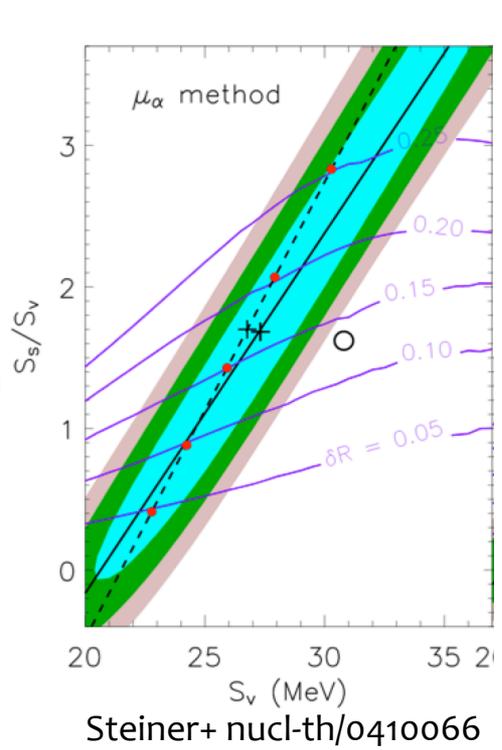
Newton+, arxiv:2111.07969

Balliet+, arxiv:2009.07696

Crust-core transition pressure and chemical potential



Correlations from mass fits, fits to 3DHF



Balliet et al, arxiv:2009.07696

$$P(\sigma_0, \sigma_c, c, p)$$

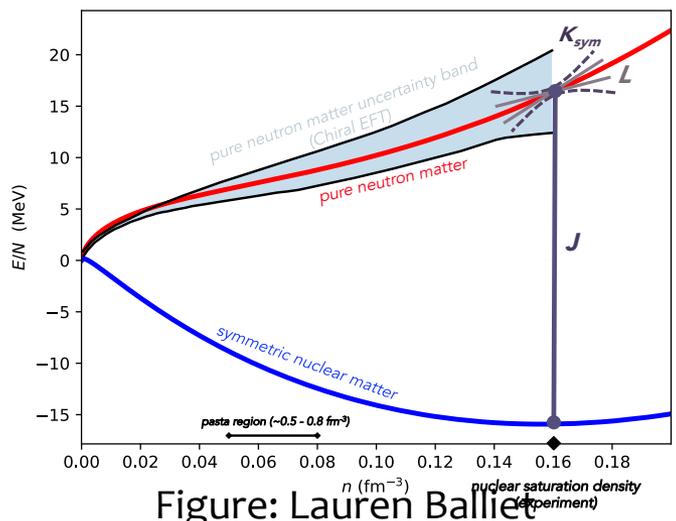
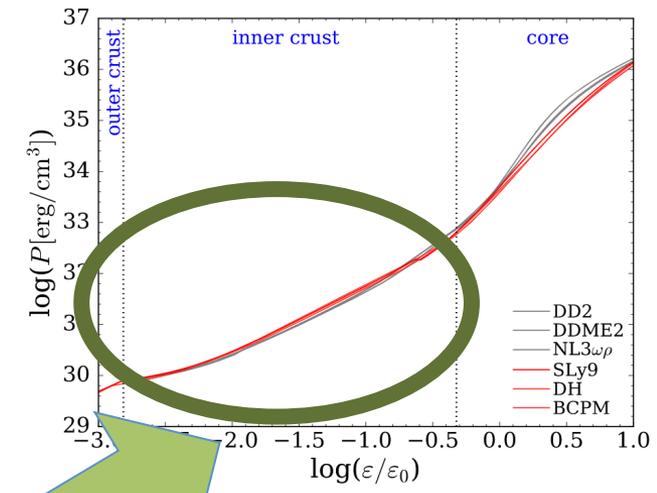
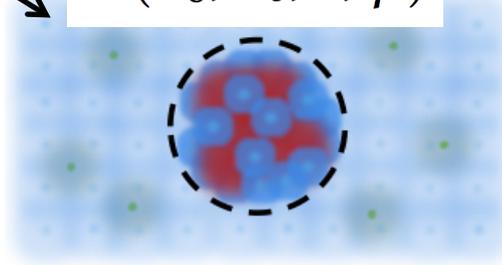
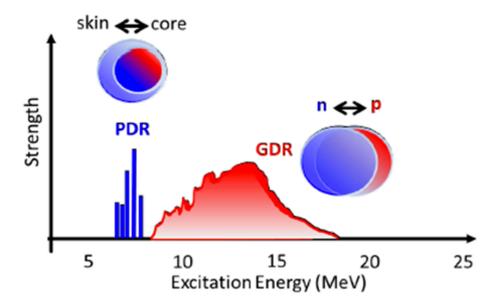
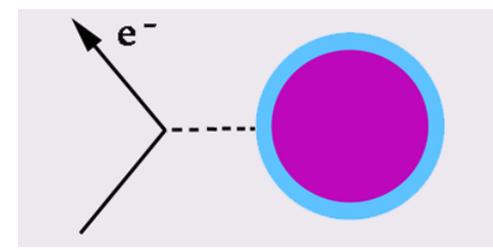
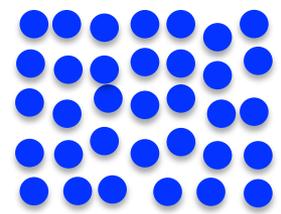
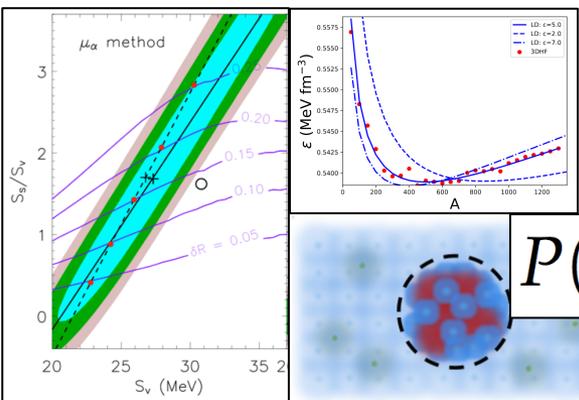
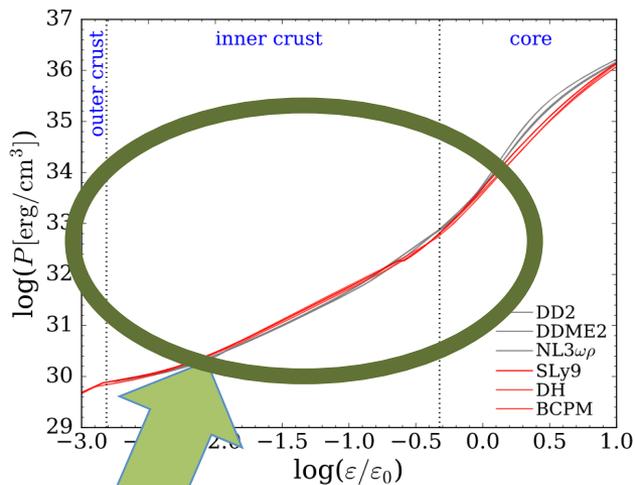


Figure: Lauren Balliet

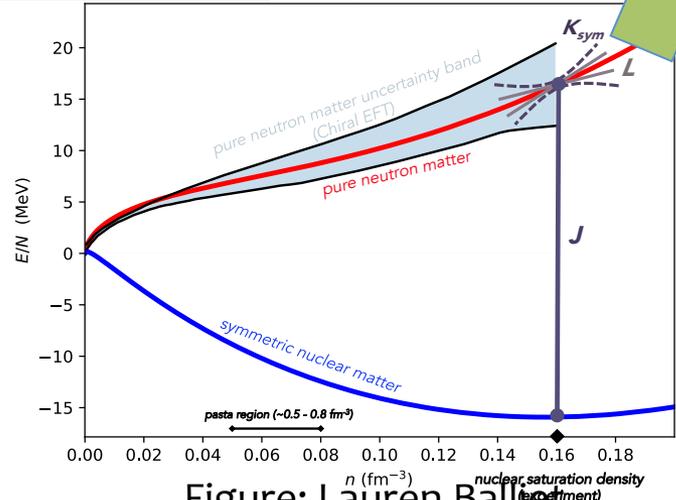
$$P(J, L, K_{\text{sym}} | \mathcal{D})$$

DFT





$$P(\text{crust+outer core}(J, L, K_{\text{sym}}) | \mathcal{D}_{\text{nuc}}; \sigma_0, \sigma_c, c, \rho)$$



$$P(J, L, K_{\text{sym}} | \mathcal{D})$$

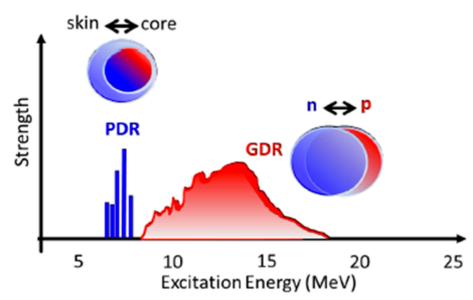
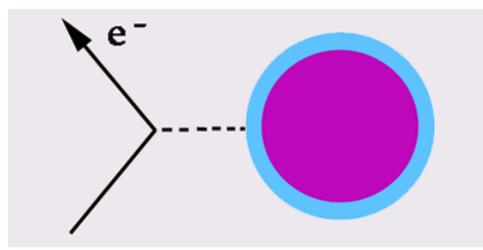
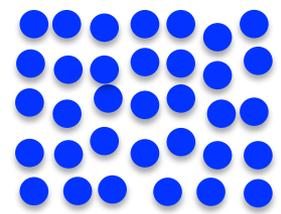
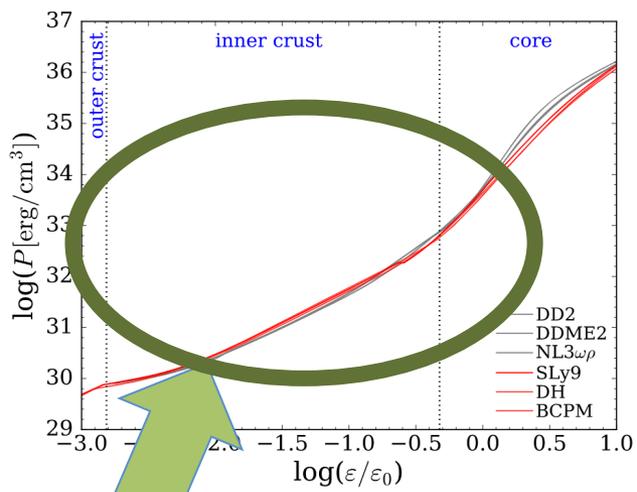
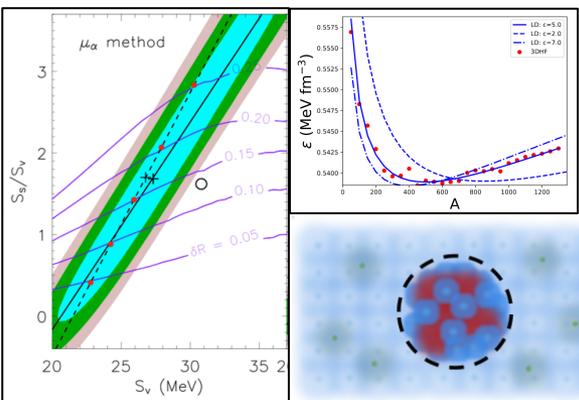
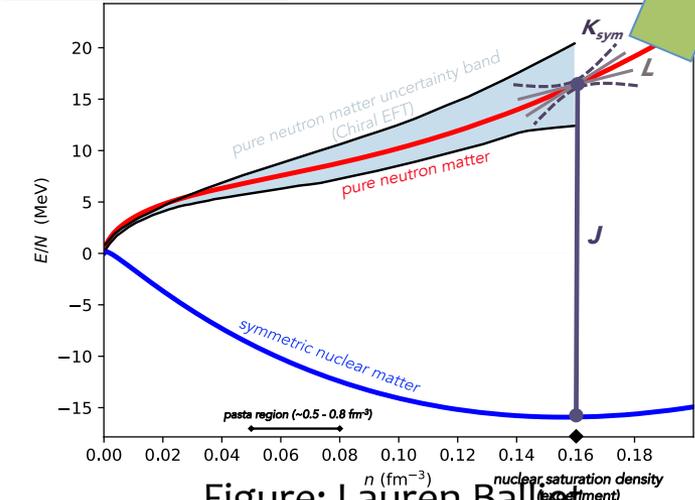


Figure: Lauren Ballet



$$P(\text{crust+outer core}(J, L, K_{\text{sym}}) | \mathcal{D}_{\text{nuc}})$$

CLDM



$$P(J, L, K_{\text{sym}} | \mathcal{D})$$

DFT

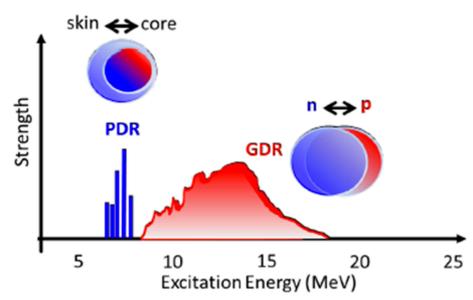
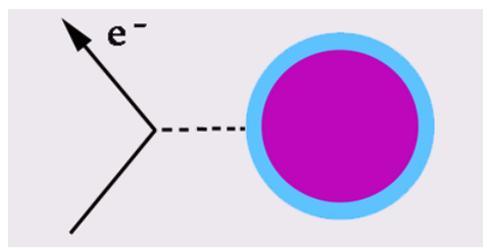
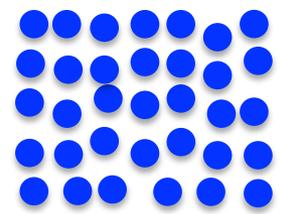
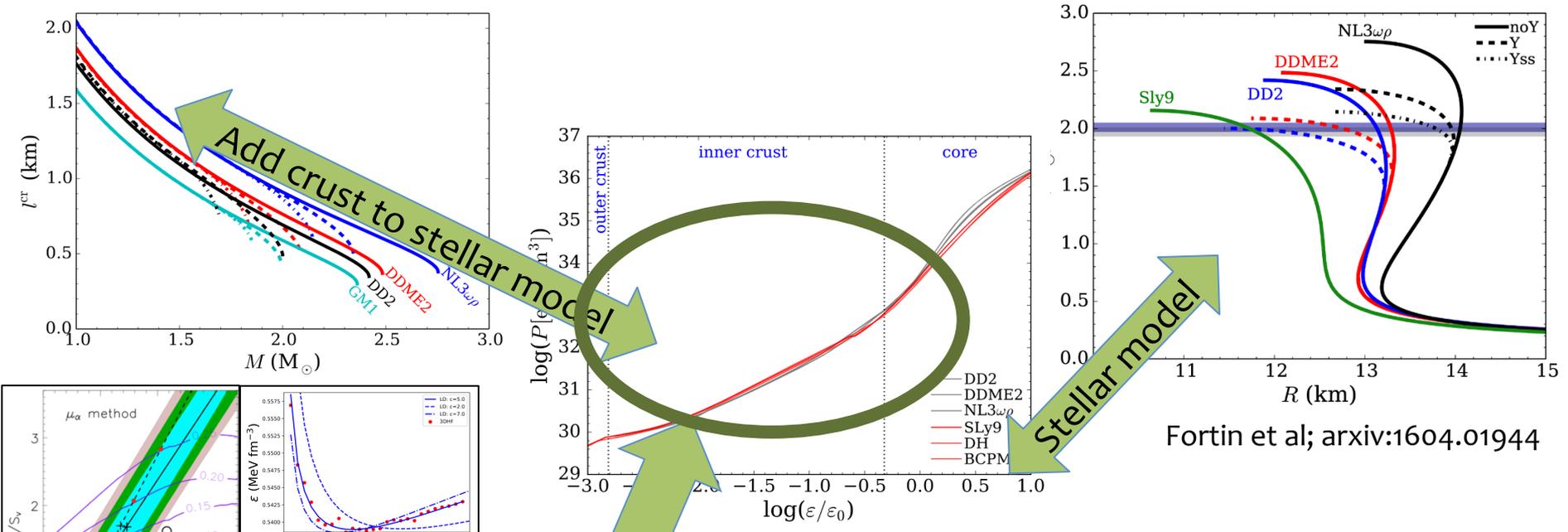
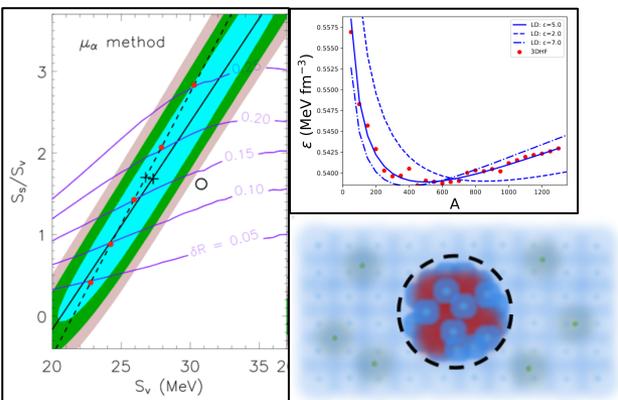


Figure: Lauren Ballet



Fortin et al; arxiv:1604.01944



CLDM

$$P(\text{crust+outer core}(J, L, K_{\text{sym}}) | \mathcal{D}_{\text{nuc}})$$

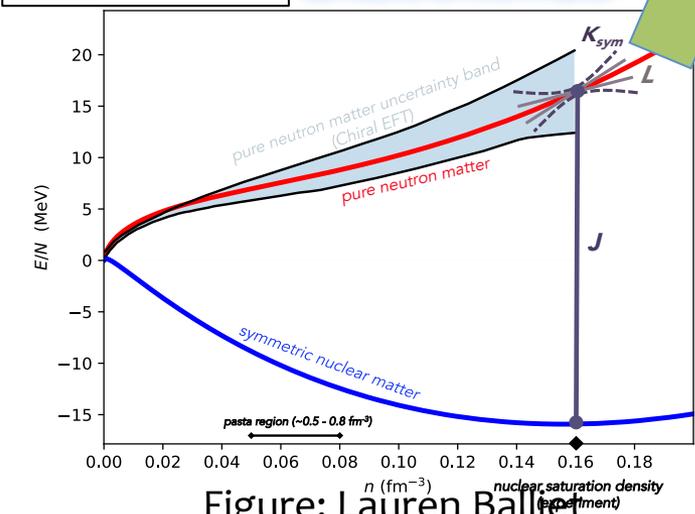
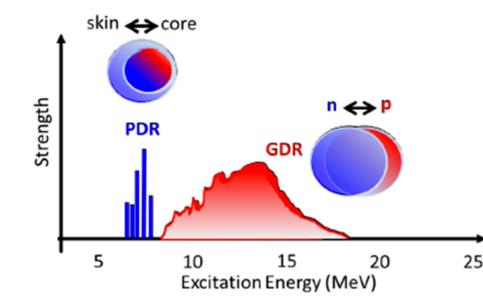
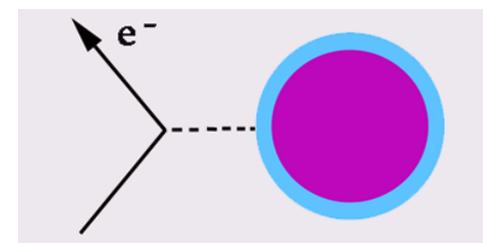
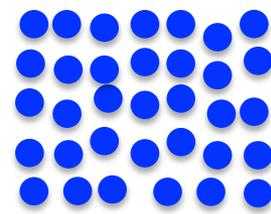
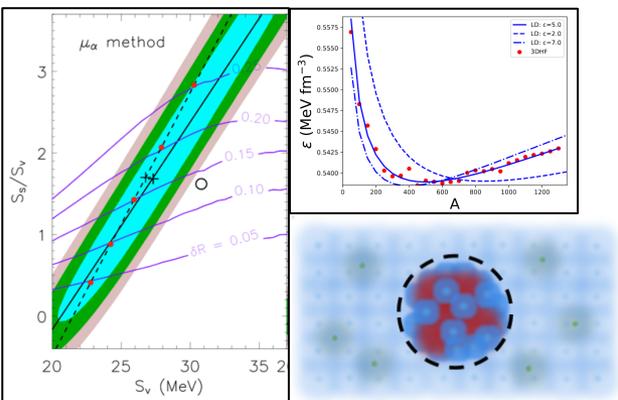
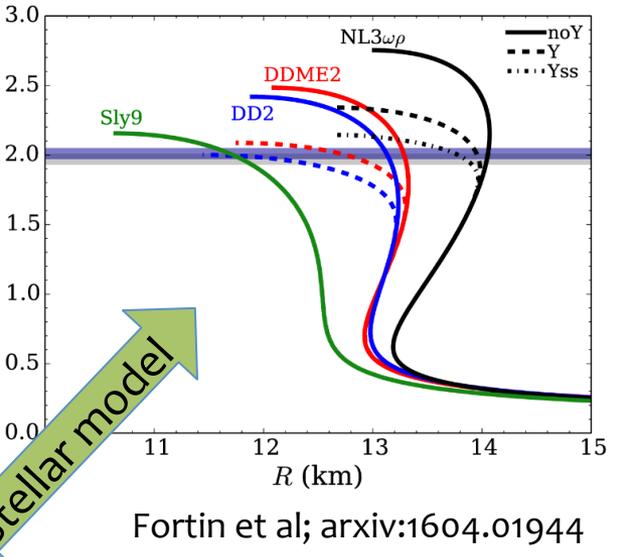
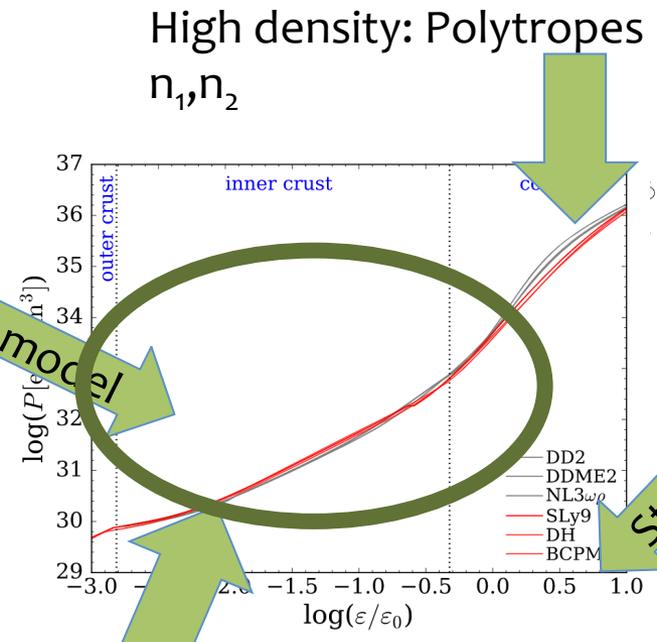
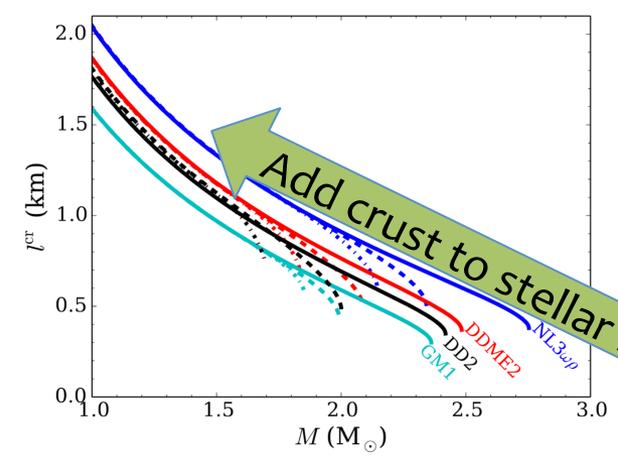


Figure: Lauren Ballet

$$P(J, L, K_{\text{sym}} | \mathcal{D})$$

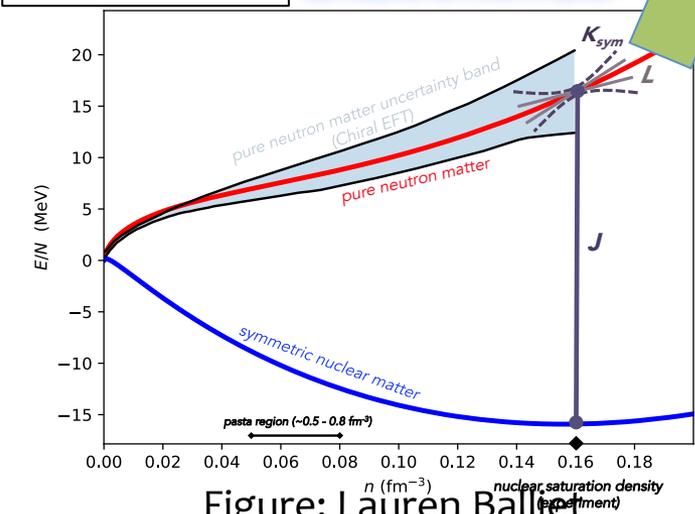
DFT





$$P(\text{crust+outer core}(J, L, K_{\text{sym}}) | \mathcal{D}_{\text{nuc}})$$

CLDM



$$P(J, L, K_{\text{sym}} | \mathcal{D})$$

DFT

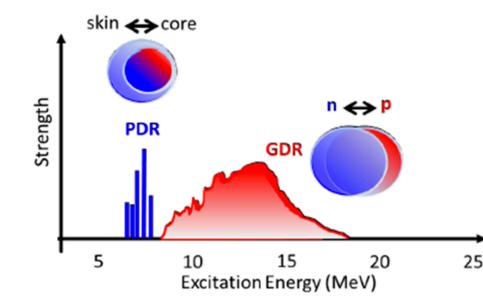
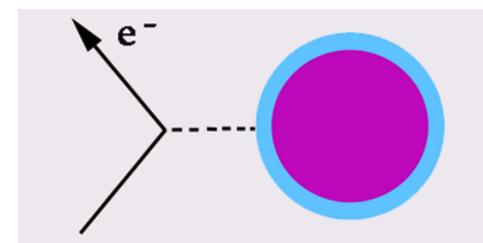
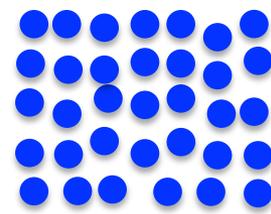
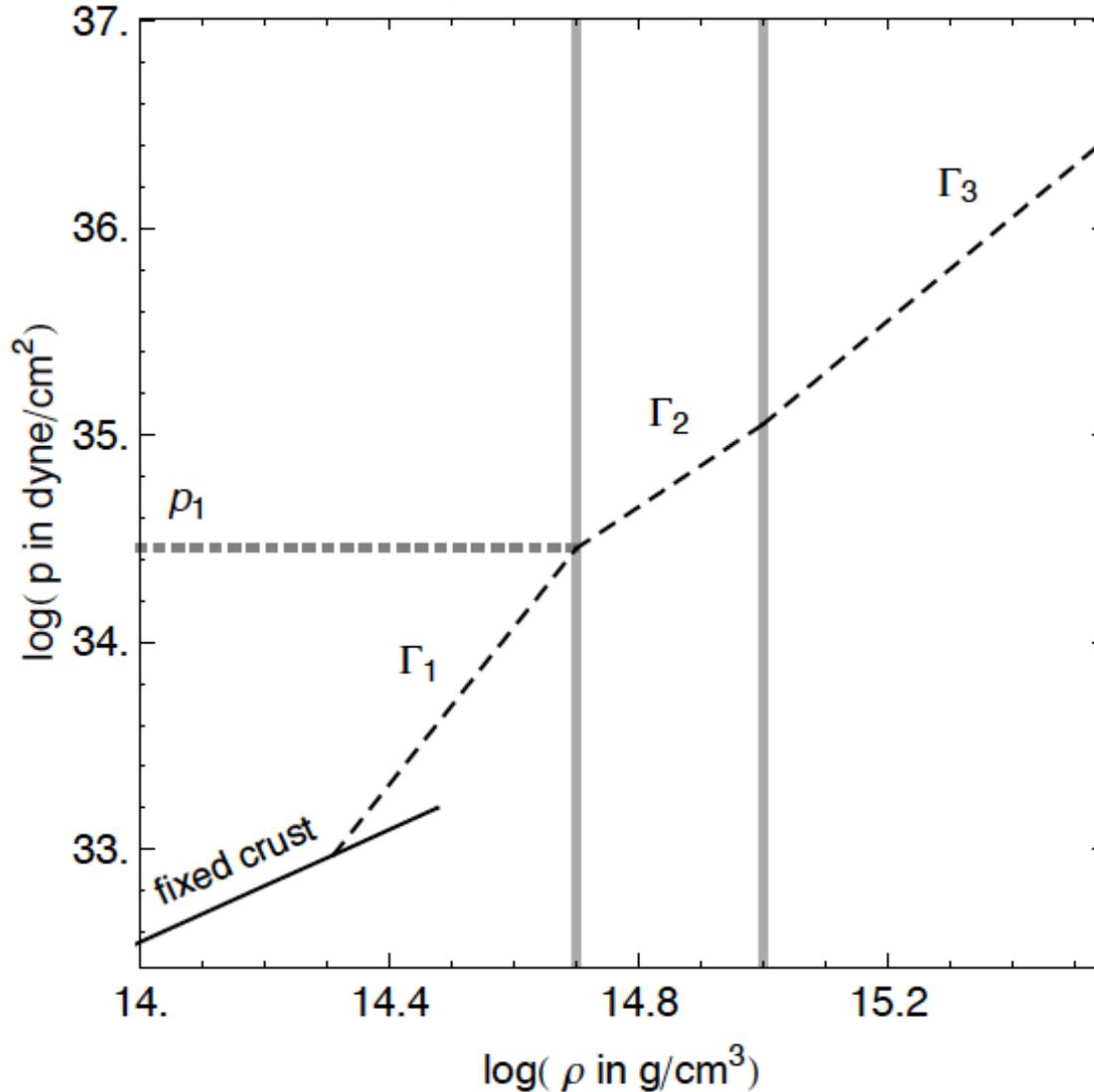


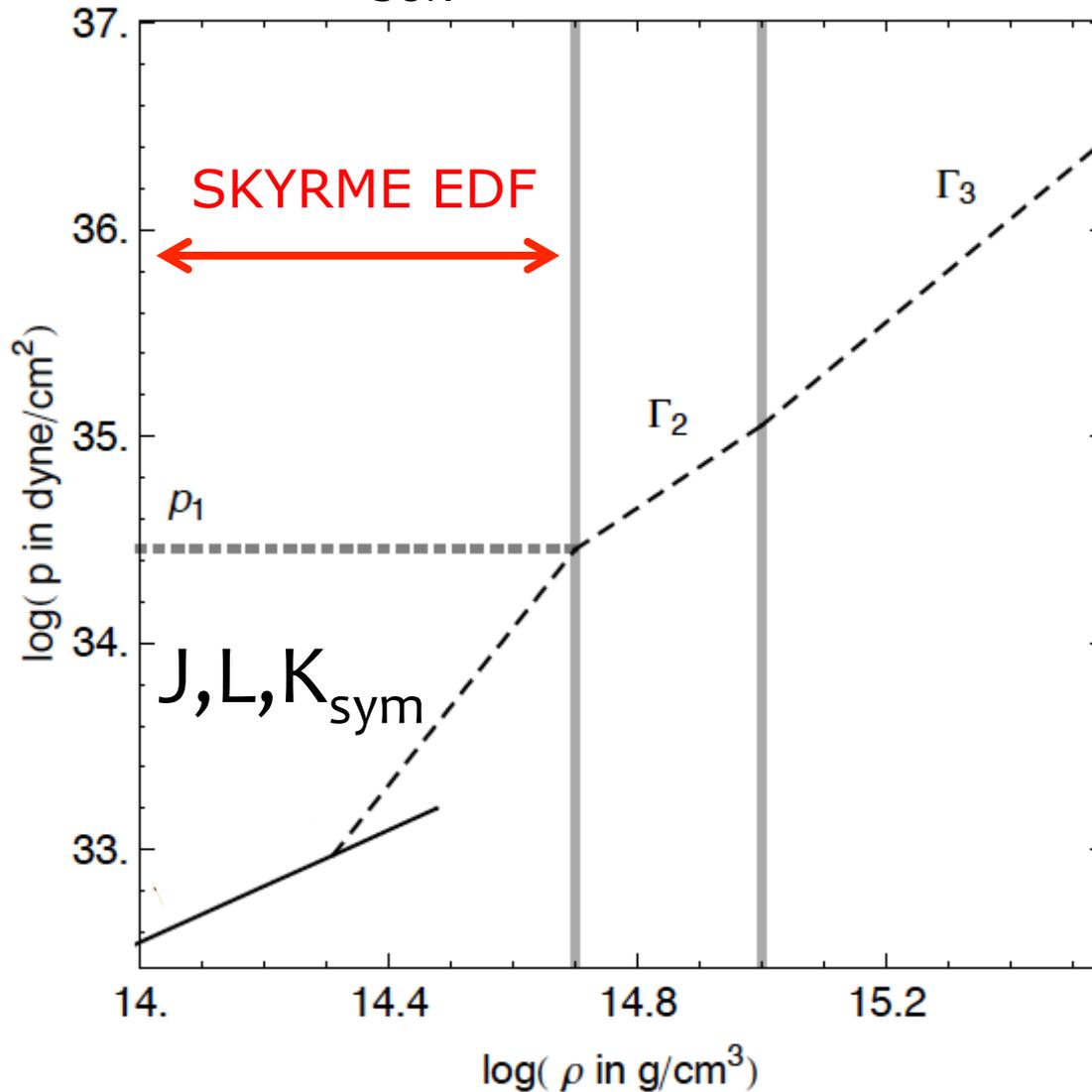
Figure: Lauren Ballet

High density EOS: piecewise polytrope tuned to give max masses $> 2.0 M_{\text{SUN}}$ up until causality is violated



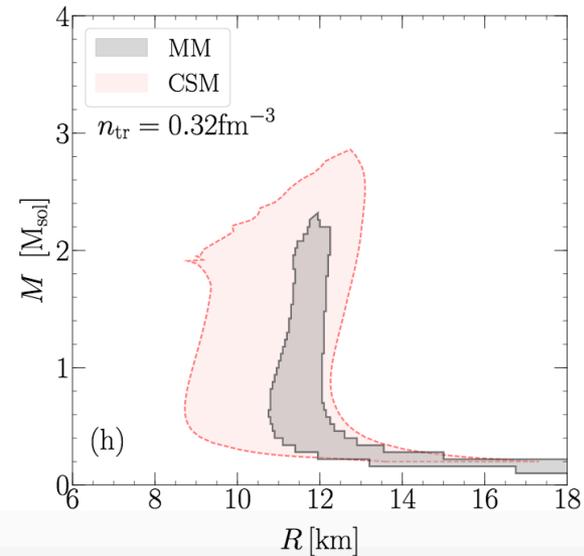
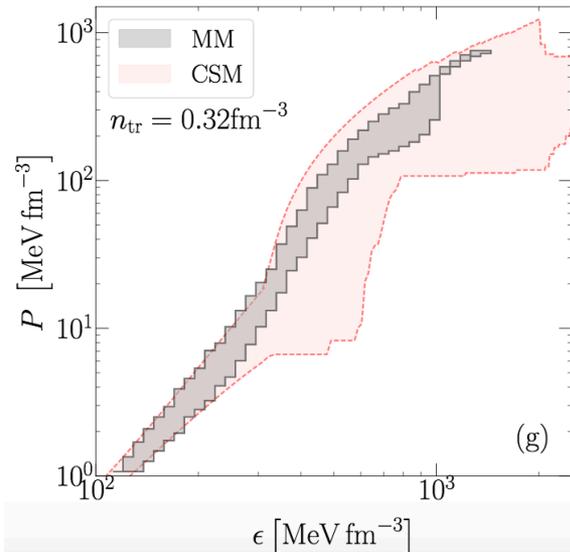
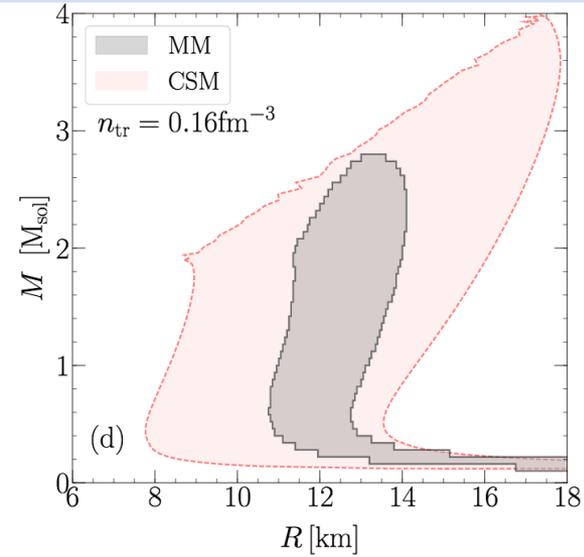
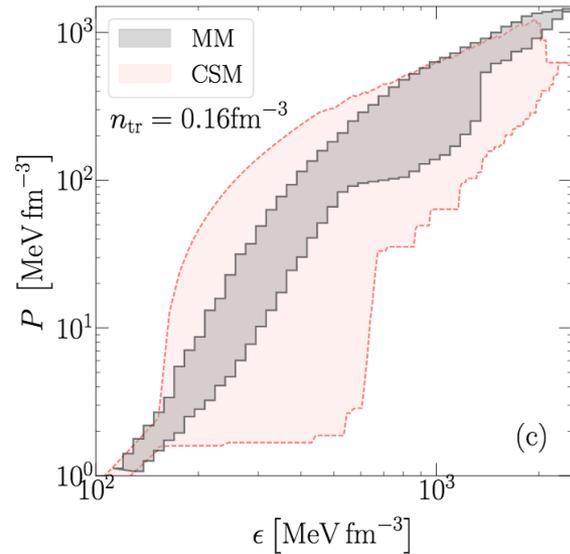
Read+, arxiv:0812.2163; see also works by Steiner, Lattimer, Özel

High density EOS: piecewise polytrope tuned to give max masses $> 2.0 M_{\text{SUN}}$ up until causality is violated



Read+, arxiv:0812.2163; see also works by Steiner, Lattimer, Özel...

Polytropes versus continuing the nuclear matter comparison/ extrapolating EDF to arbitrarily high



Potential sources of systematic error

Had to choose an EDF (Skyrme). Enough degrees of freedom?
Can add more (Q_{sym})

Symmetric nuclear matter and gradient parameters held fixed; extending inference to those parameters may change posteriors

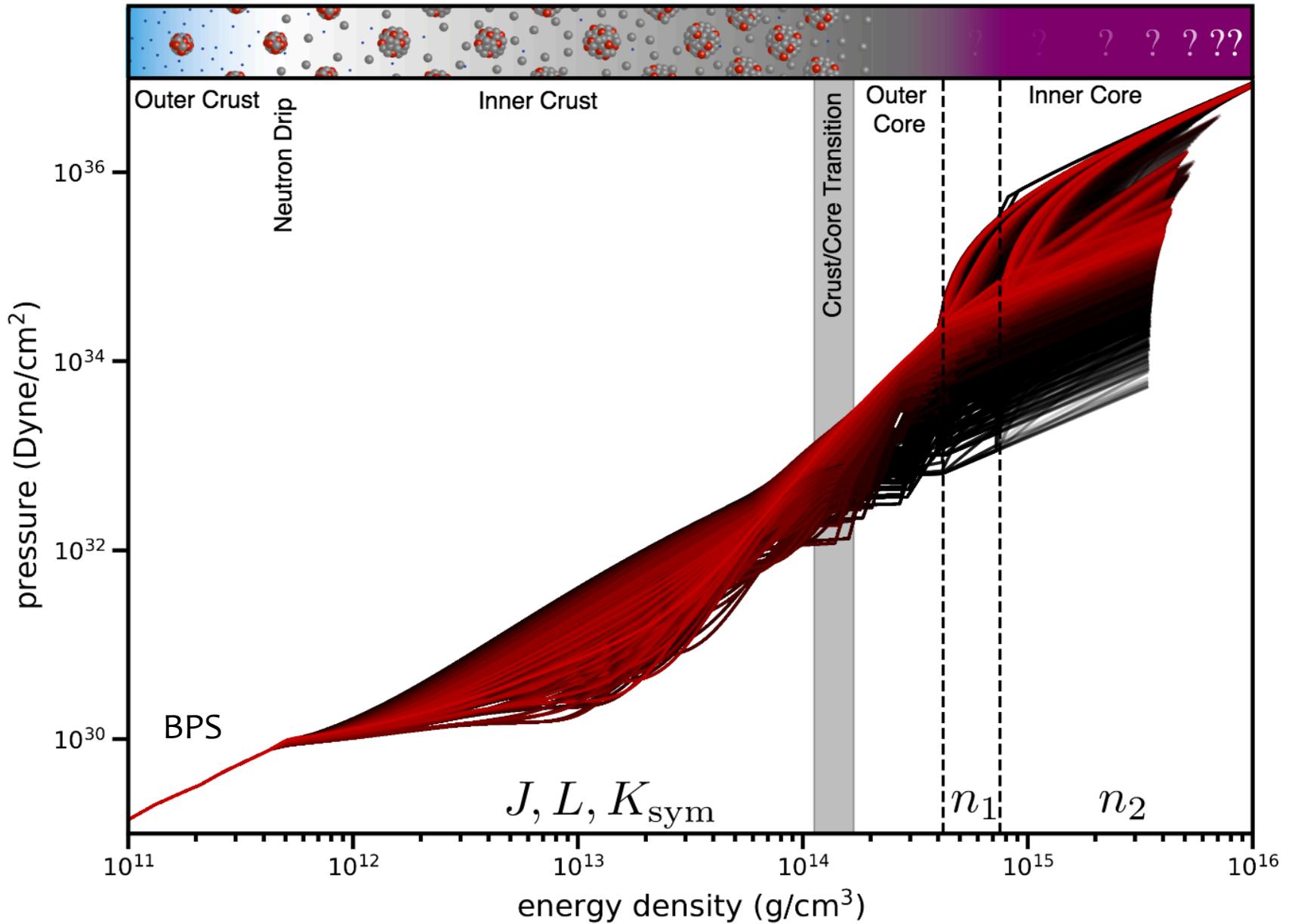
We're usually not directly modeling nuclear observables - but in some cases we can (e.g. weak form factor) and thus improve consistency

We've chosen a model of crust (CLDM) different to modeling of the nuclei (1D SHF+RPA)

Need more information to constrain surface parameters of crust model (mass fits, semi-infinite nuclear matter, Thomas-Fermi calculations)

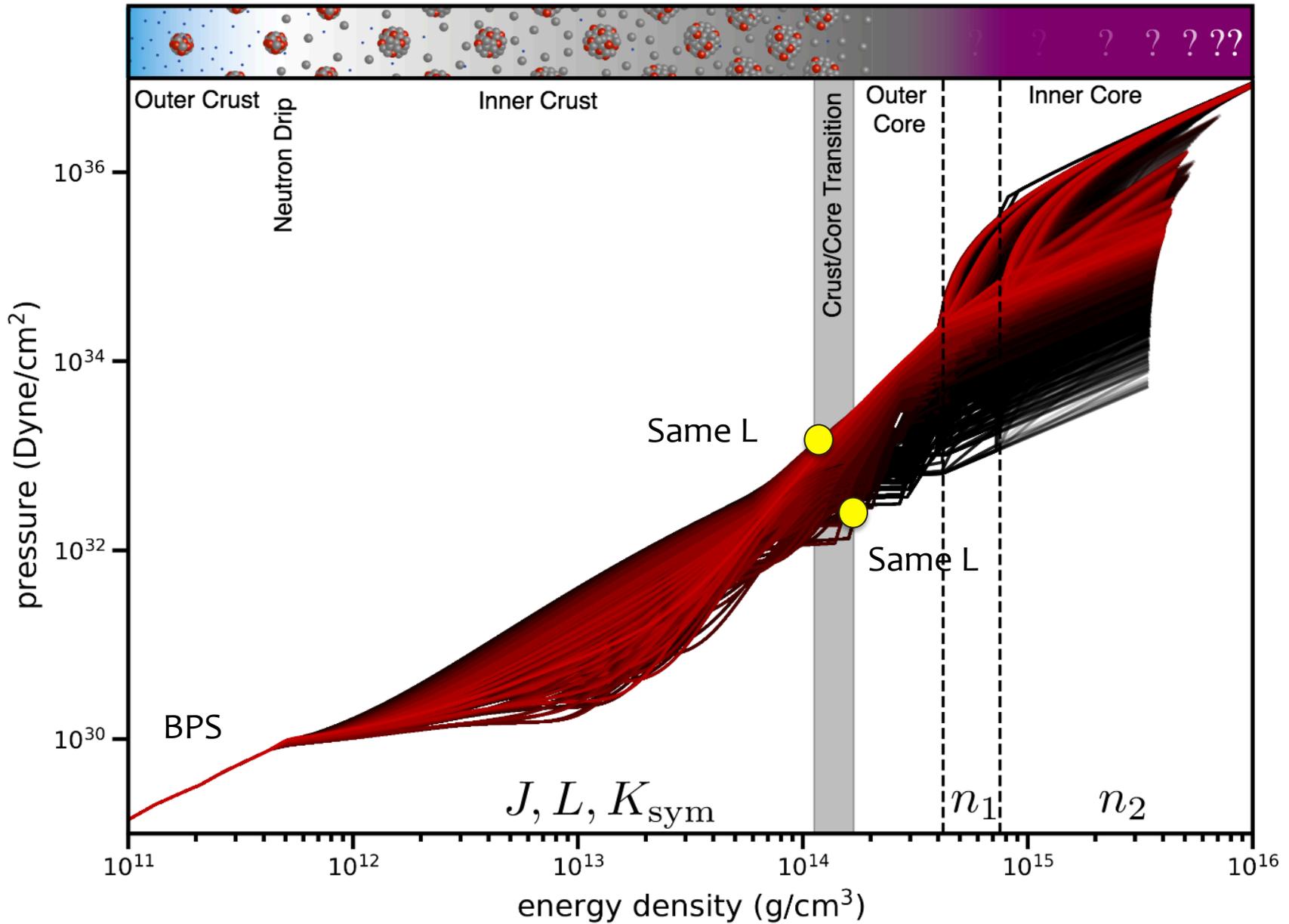
2 polytropes in the core is bare minimum: can be more sophisticated

Sample of our resulting equations of state



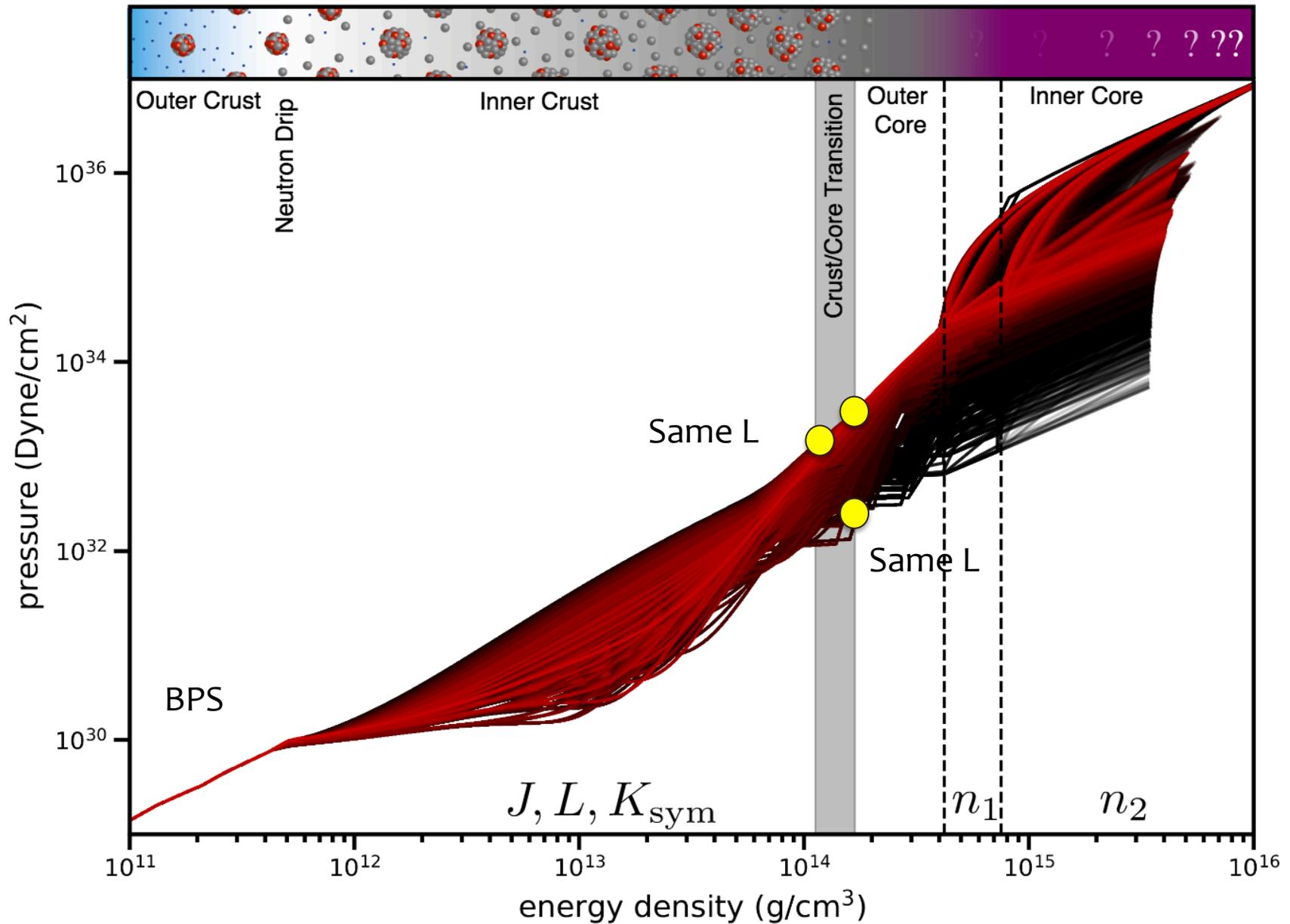
Neill, Preston, Tsang, Newton in prep

By the way, about the crust core transition...



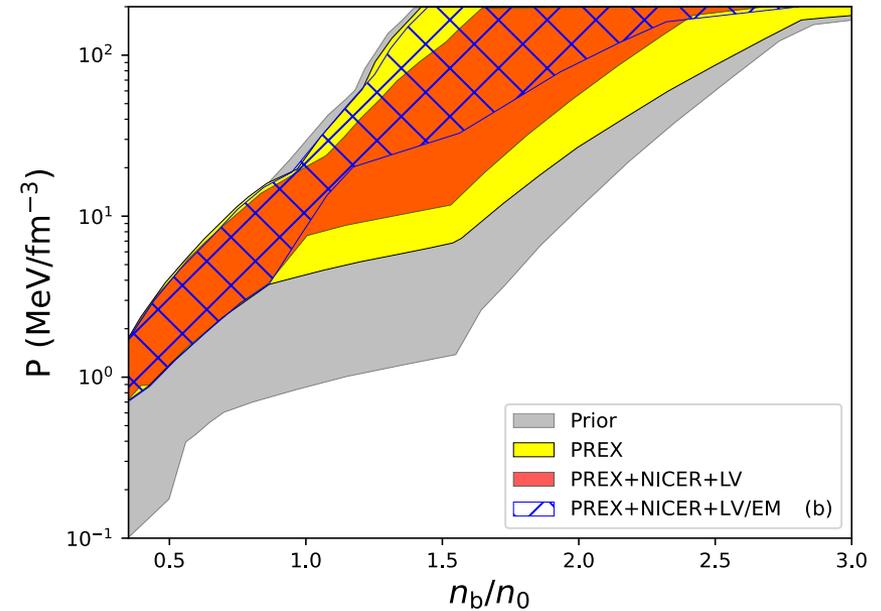
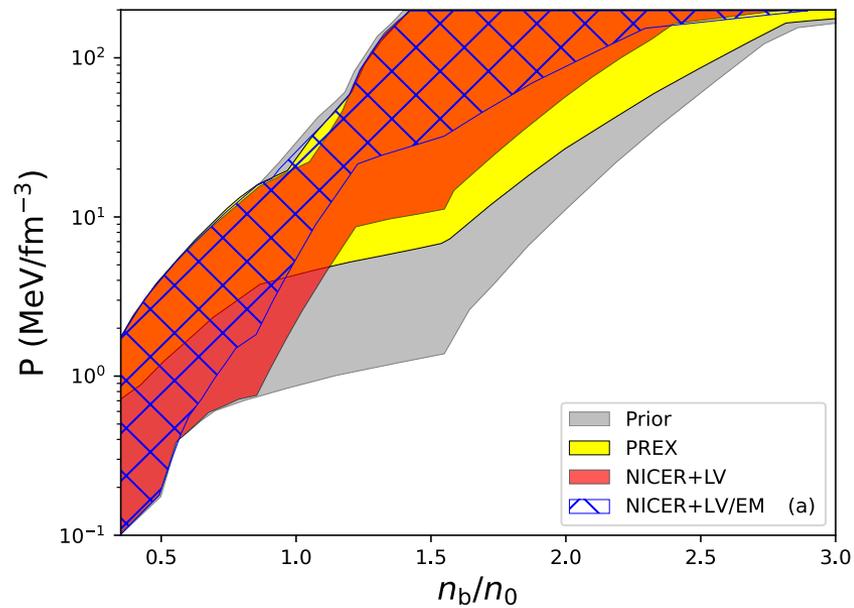
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By the way, about the crust core transition...



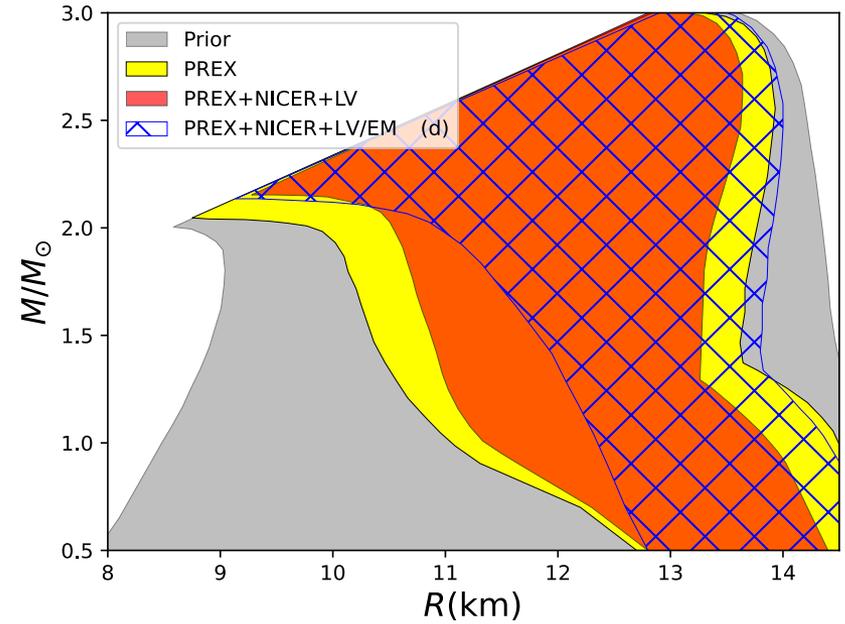
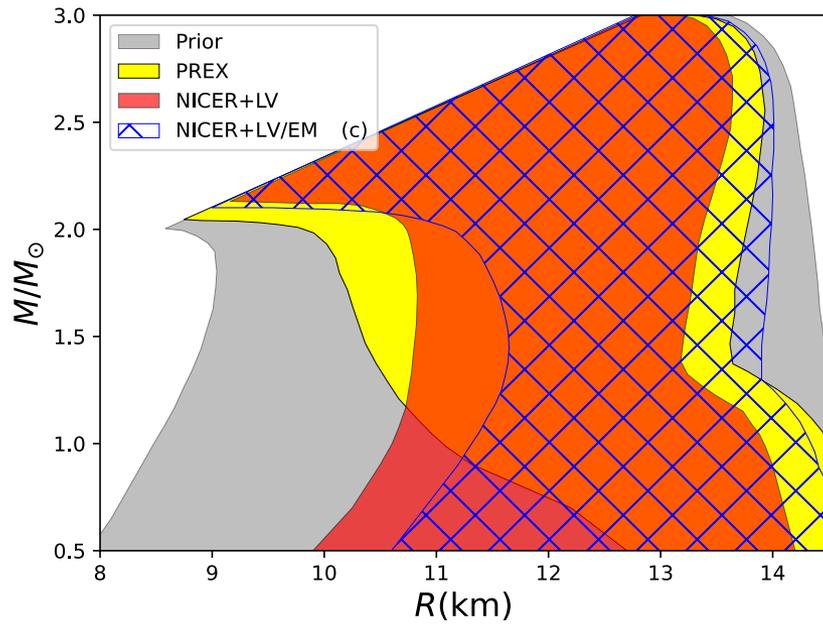
Neill, Preston, Tsang, Newton in prep

Nuclear *and* astro data can be brought to bear with greater consistency



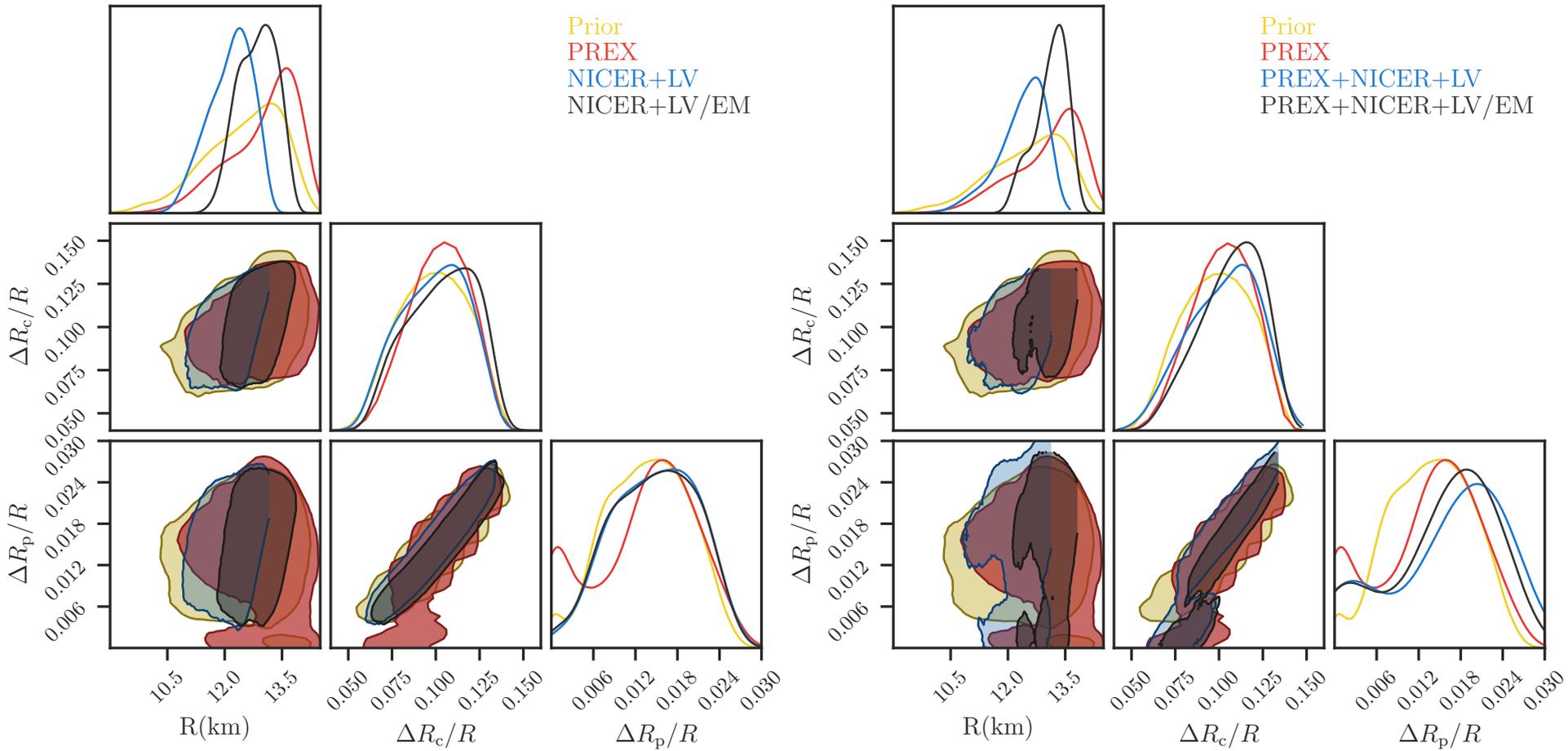
Newton et al; arxiv:2112.12108

Nuclear *and* astro data can be brought to bear with greater consistency



Newton et al; arxiv:2112.12108

Radial extents

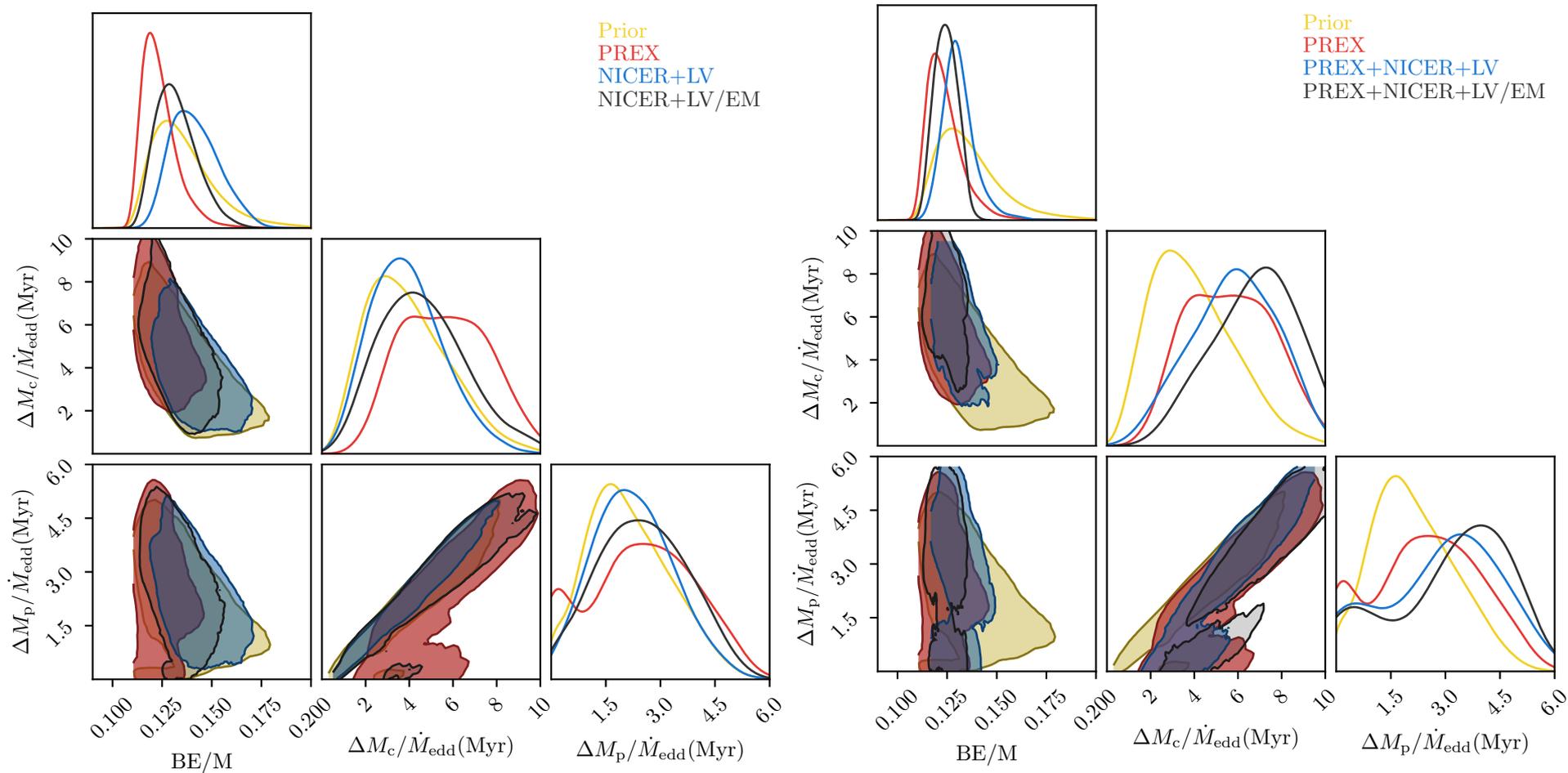


Newton et al; arxiv:2112.12108

12km star: With just NICER/LIGO data, crust can contribute 0.96-1.8 km
c.f. uncertainty from different ways of matching EoS $\approx 0.7\text{km}$; Fortin et al arxiv: 1604.01944

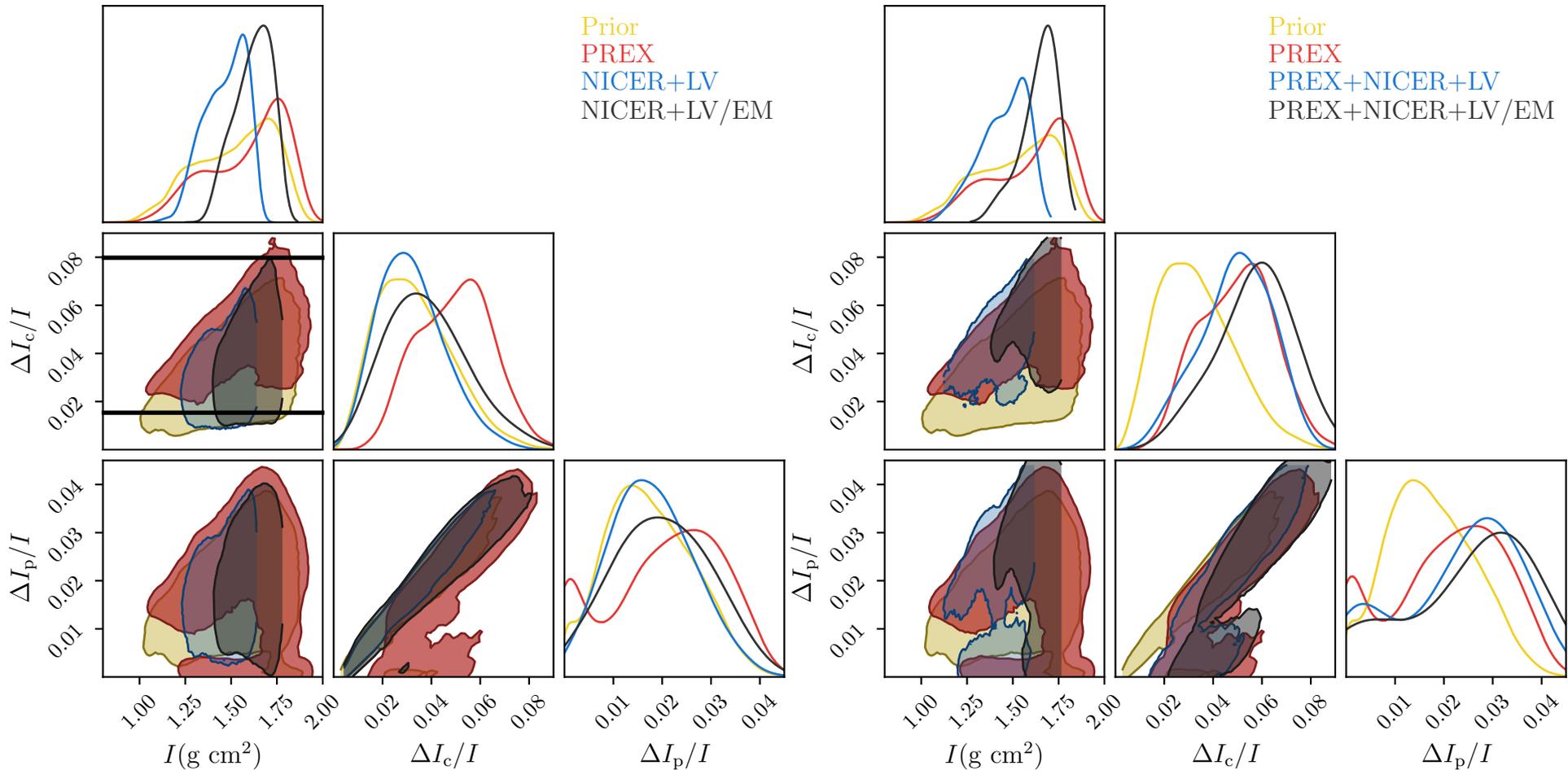
Linking crust and core models, nuclear and astro data: $M=1.4M_{\text{SUN}}$

Masses



Newton et al; arxiv:2112.12108

Moments of inertias

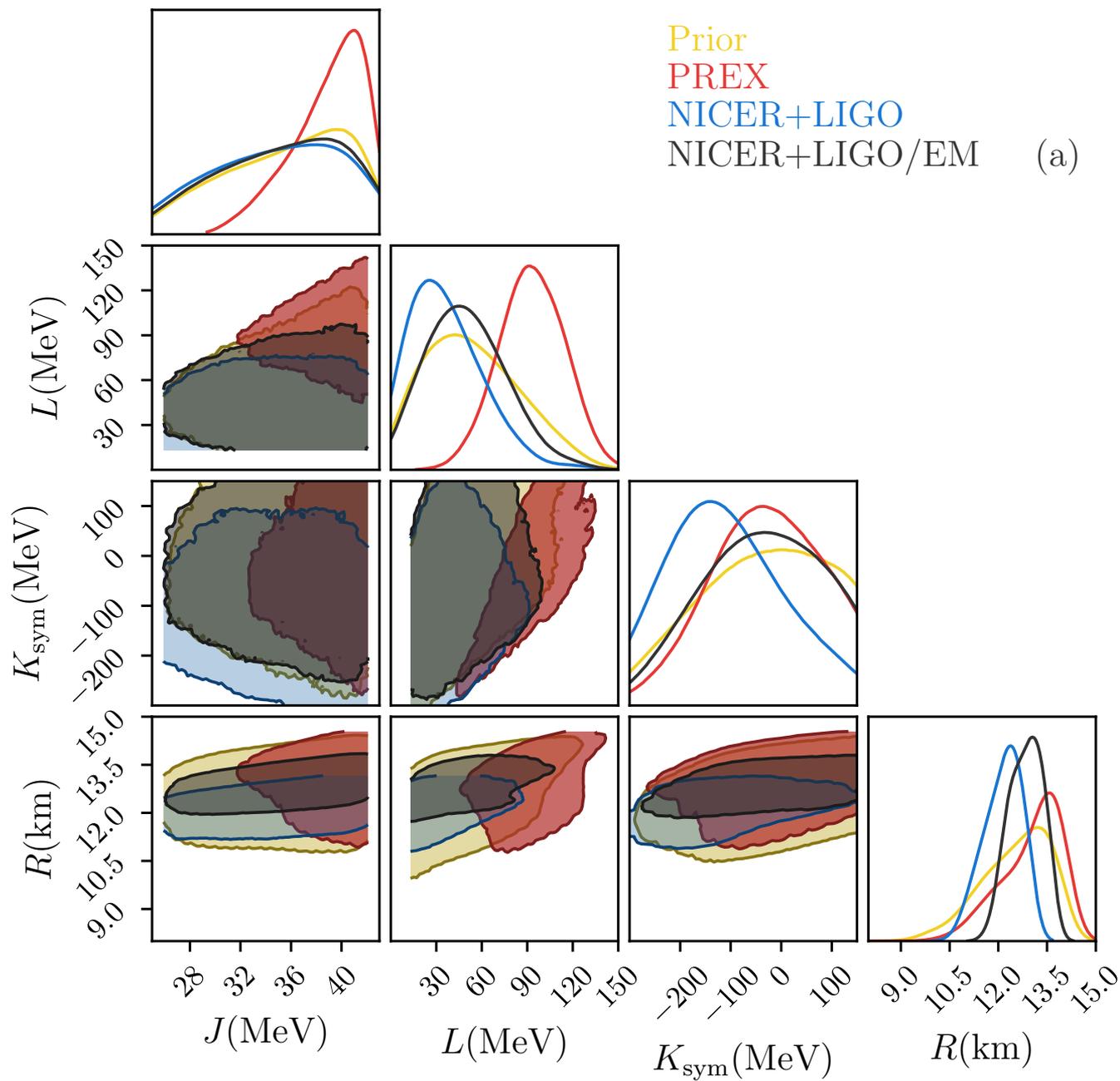


Newton et al; arxiv:2112.12108

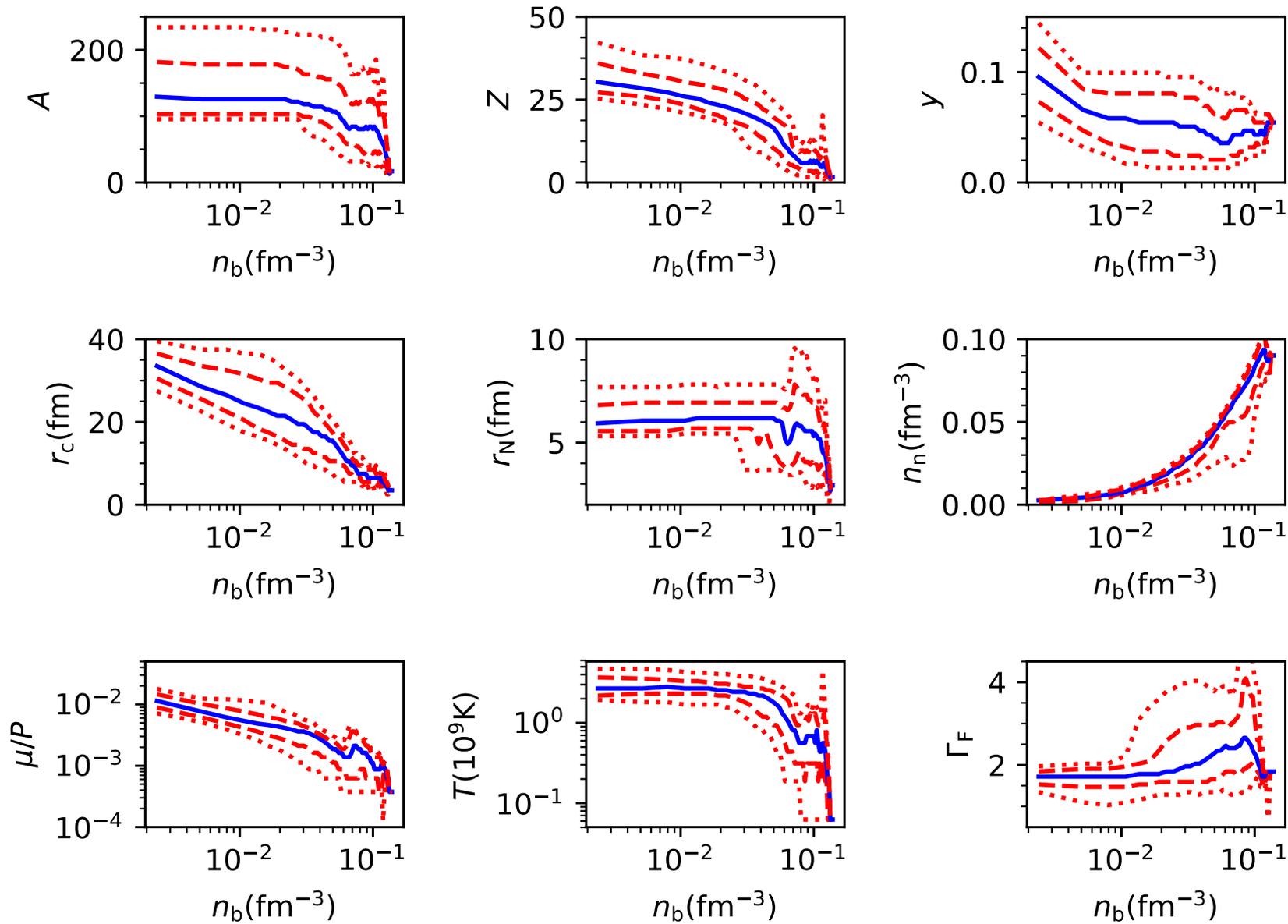
Glitches: Mol of fraction of crust must exceed 0.016 (0.08 when entrainment is high)

(crust is not/maybe enough Andersson arxiv:1207.0633/Piekarewicz arxiv:1404.2660)

Symmetry energy constraints from nuclear and astro

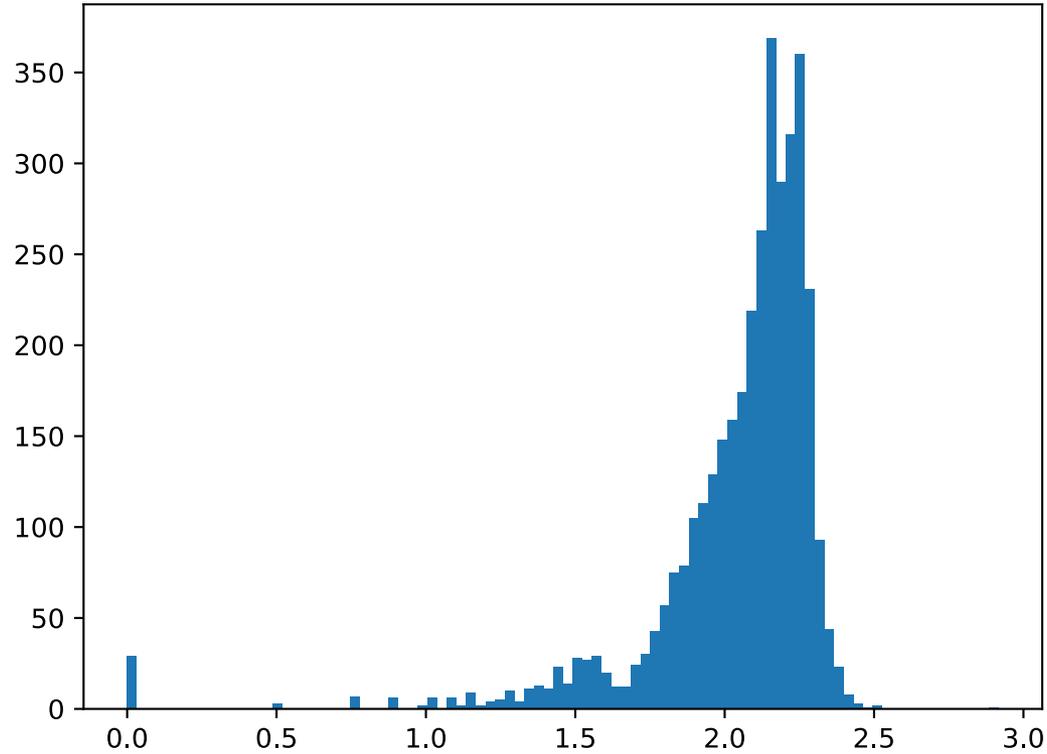


Crust Composition: Uninformative Priors



Shear modulus is likely overestimated by *at least* a factor of two in the pasta layers

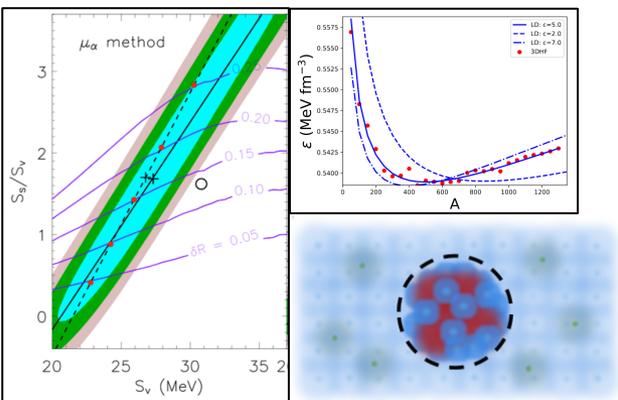
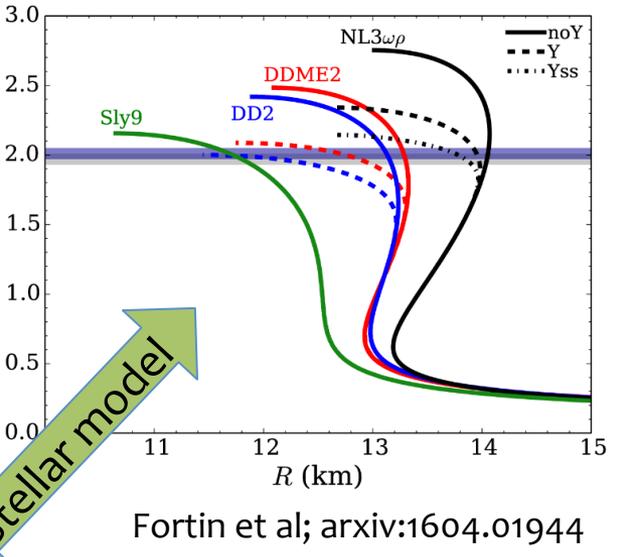
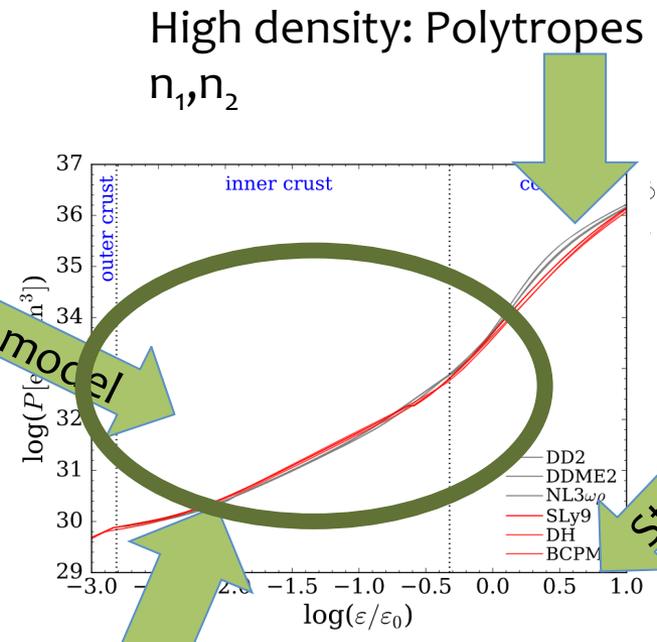
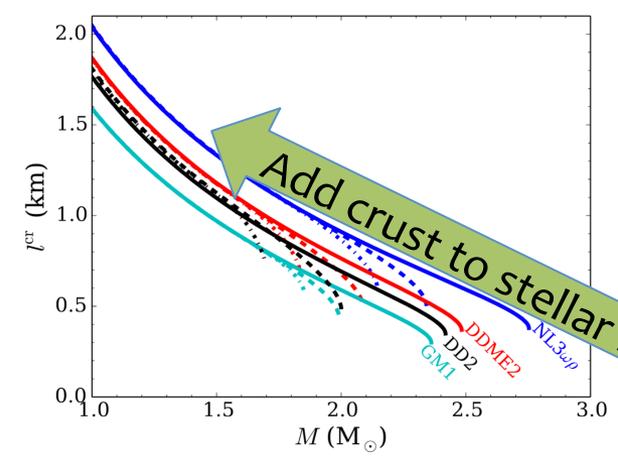
Number of EOSs



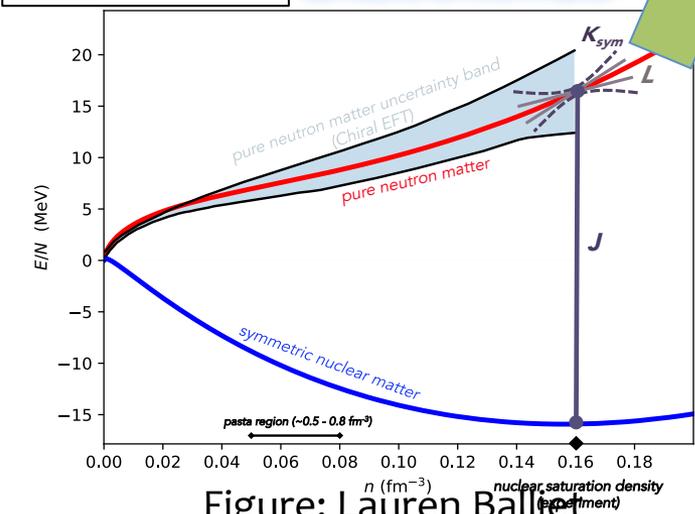
$$c_{33} = 3E_{\text{surf}}$$

$$E_{\text{surf}} = 2E_{\text{Coul}}$$

$$\mu = \frac{0.1194}{1 + 0.595(\Gamma_0/\Gamma)^2} \frac{n_i(Ze)^2}{a} \quad / \quad \mu_V \simeq \frac{c_{33}}{15}$$



$$P(\text{crust+outer core}(J, L, K_{\text{sym}}) | \mathcal{D}_{\text{nuc}})$$



$$P(J, L, K_{\text{sym}} | \mathcal{D})$$

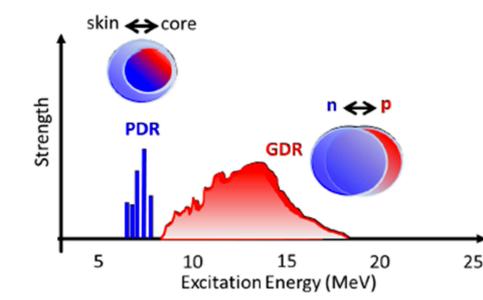
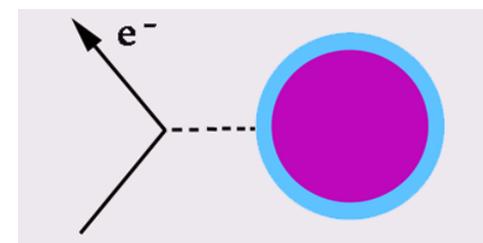
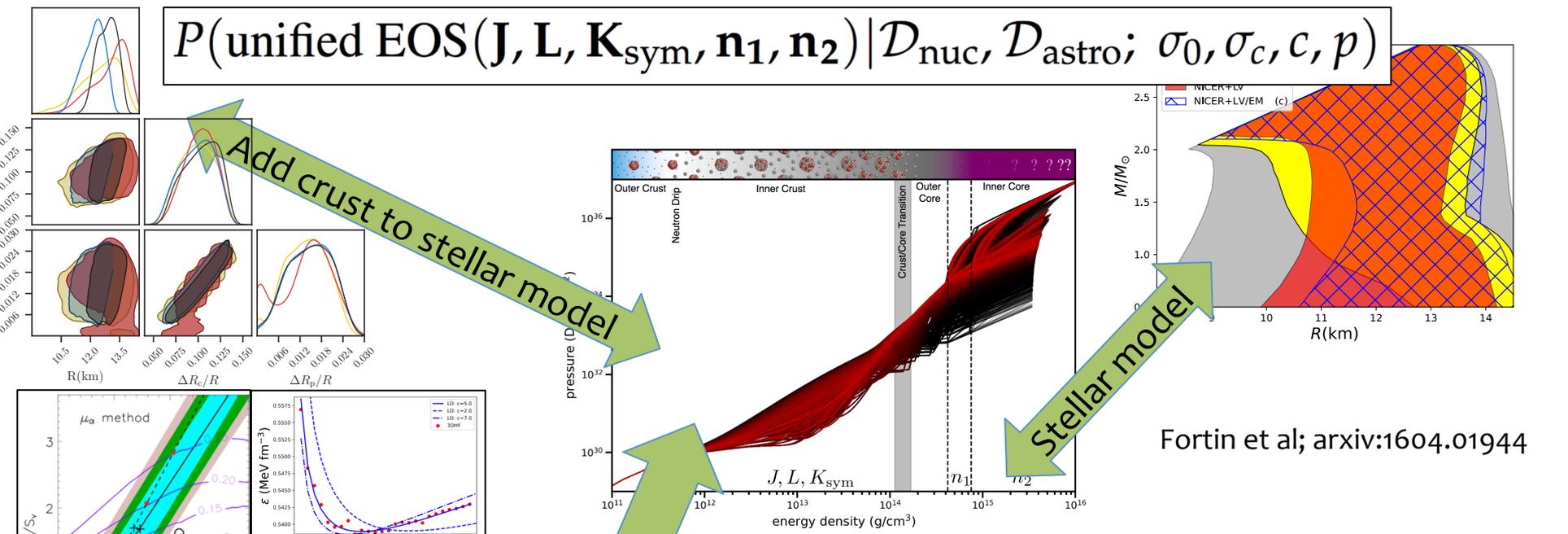


Figure: Lauren Ballet

$$P(\text{unified EOS}(J, L, K_{\text{sym}}, n_1, n_2) | \mathcal{D}_{\text{nuc}}, \mathcal{D}_{\text{astro}}; \sigma_0, \sigma_c, c, \rho)$$



$$P(\text{crust+outer core}(J, L, K_{\text{sym}}) | \mathcal{D}_{\text{nuc}})$$

$$P(J, L, K_{\text{sym}} | \mathcal{D})$$

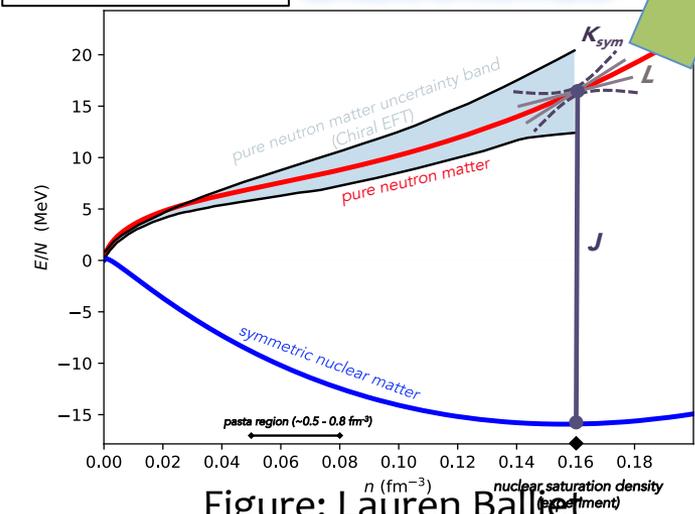
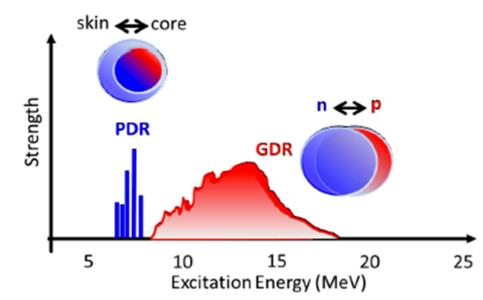
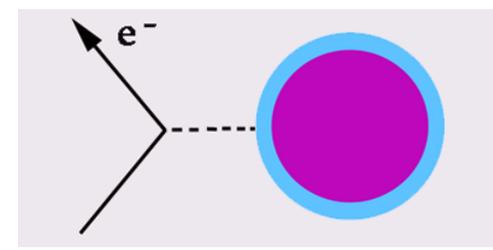
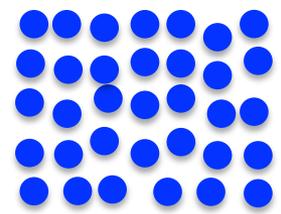
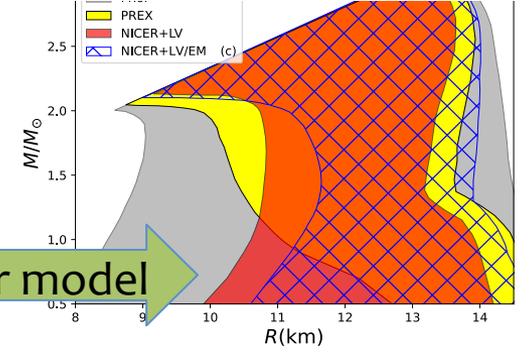
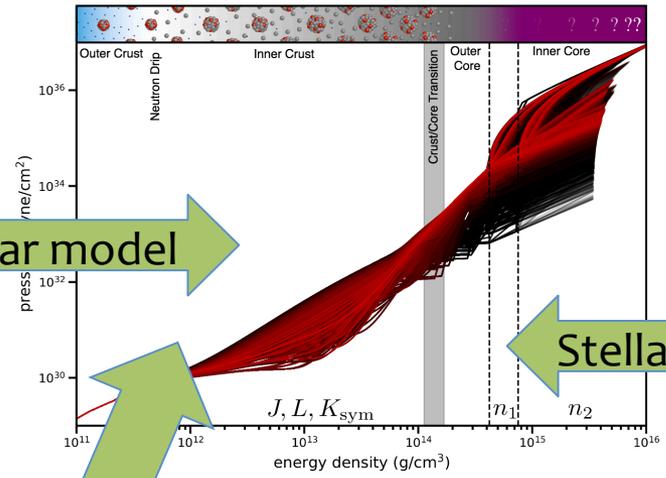
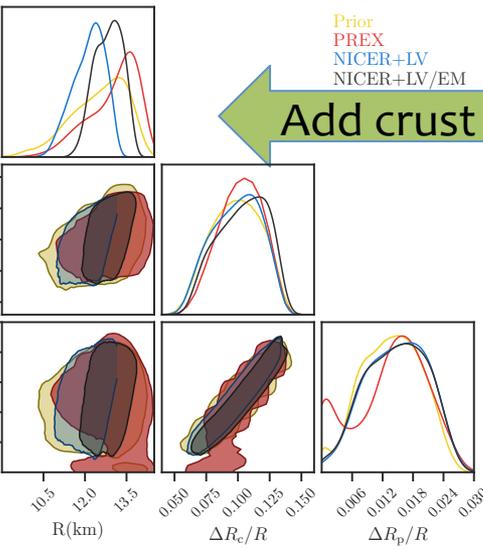
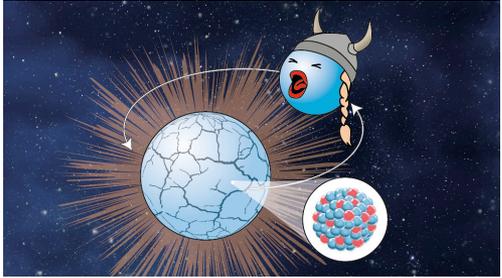
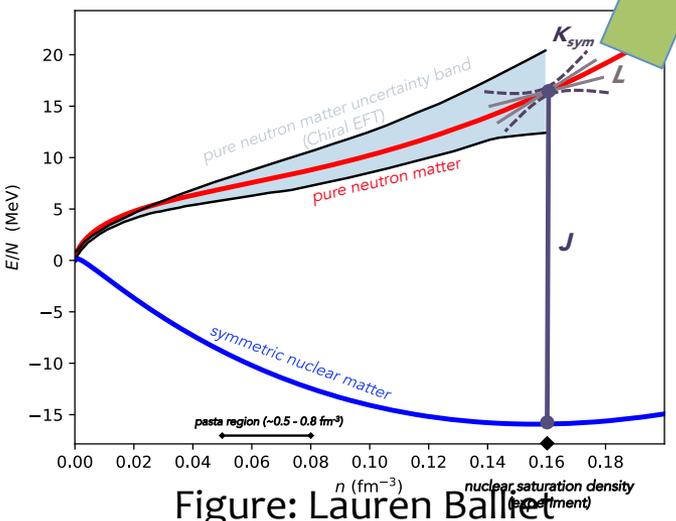


Figure: Lauren Ballet





$$P(\text{crust+outer core}(J, L, K_{\text{sym}}) | \mathcal{D}_{\text{nuc}})$$



$$P(J, L, K_{\text{sym}} | \mathcal{D})$$

DFT

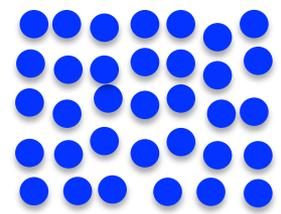
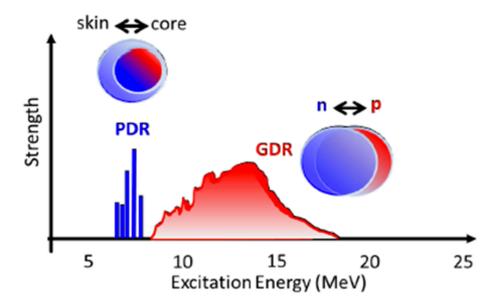
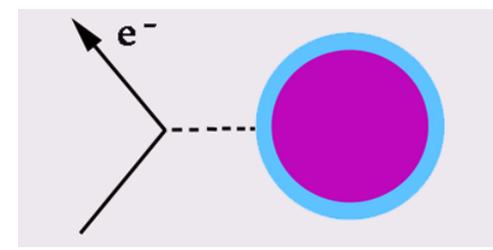
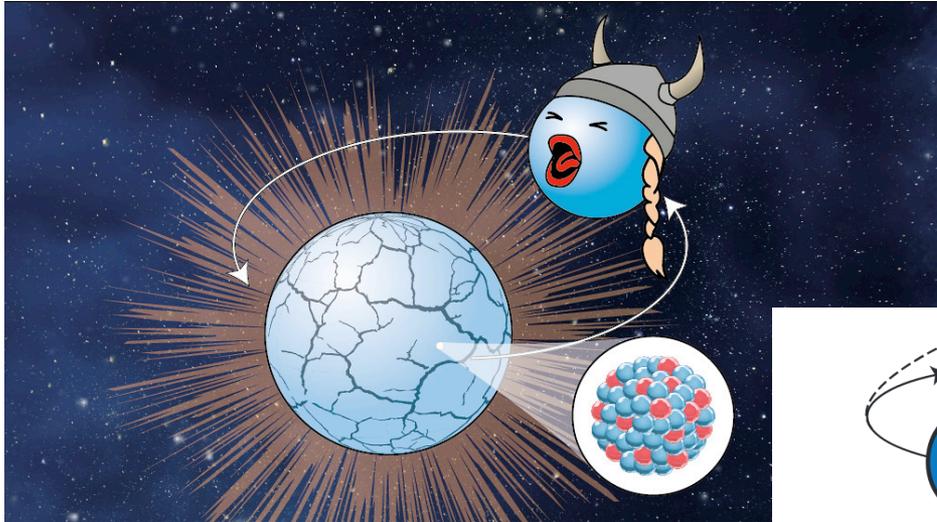


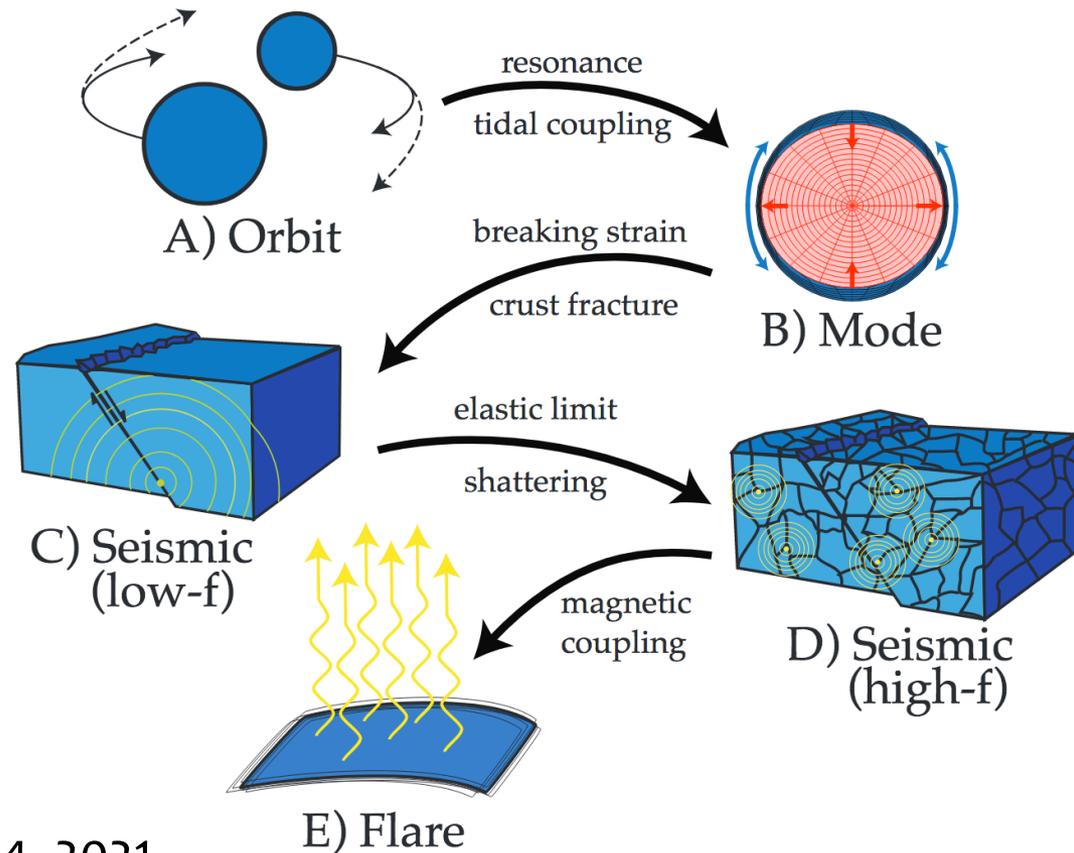
Figure: Lauren Ballet

Resonant shattering flares: combining crust and core observables



Picture: David Tsang

(See David's talk next Wednesday!)



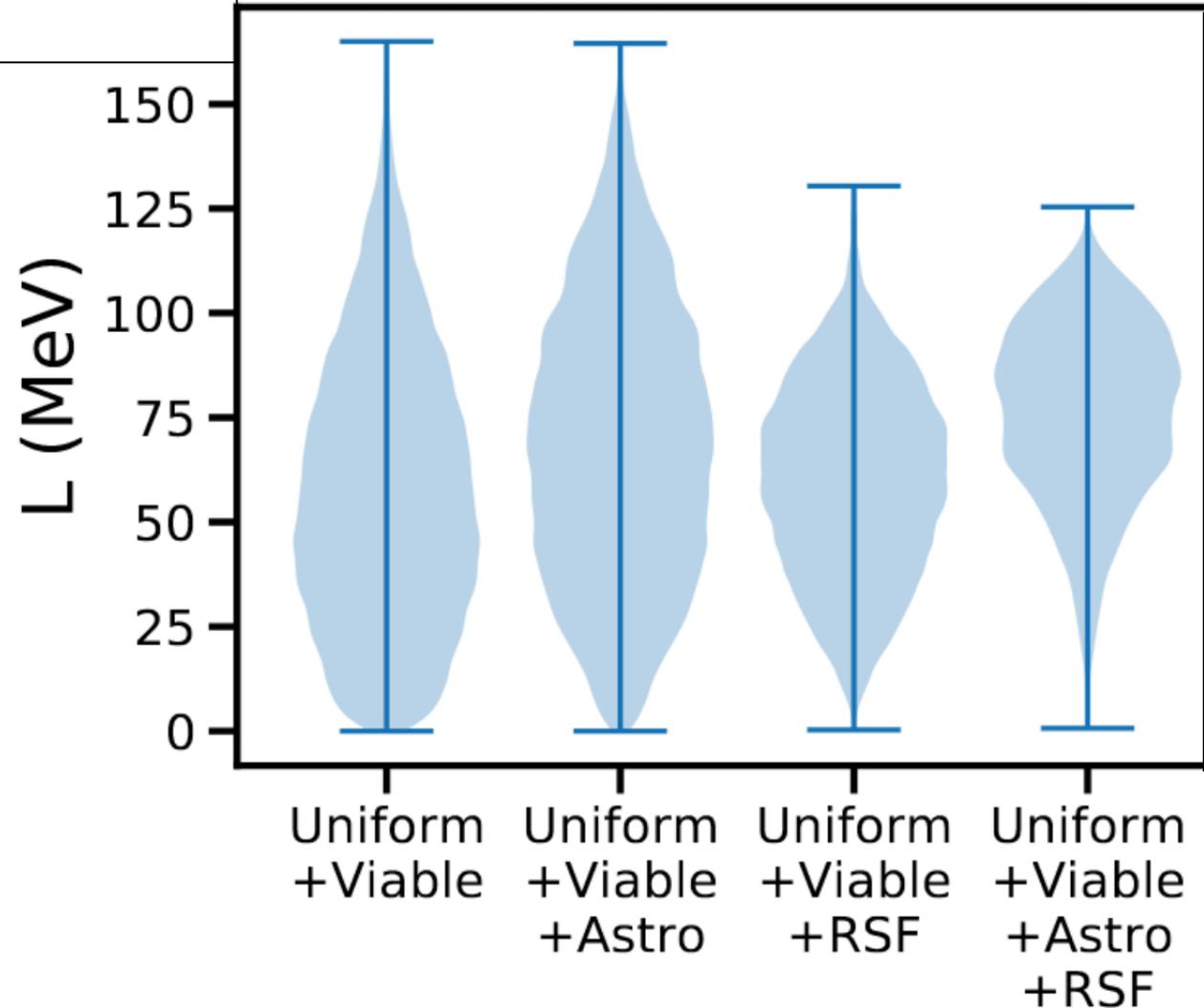
D.Tsang, *Apj* 777, 2013

Neill, Newton & Tsang, *MNRAS* 504, 2021

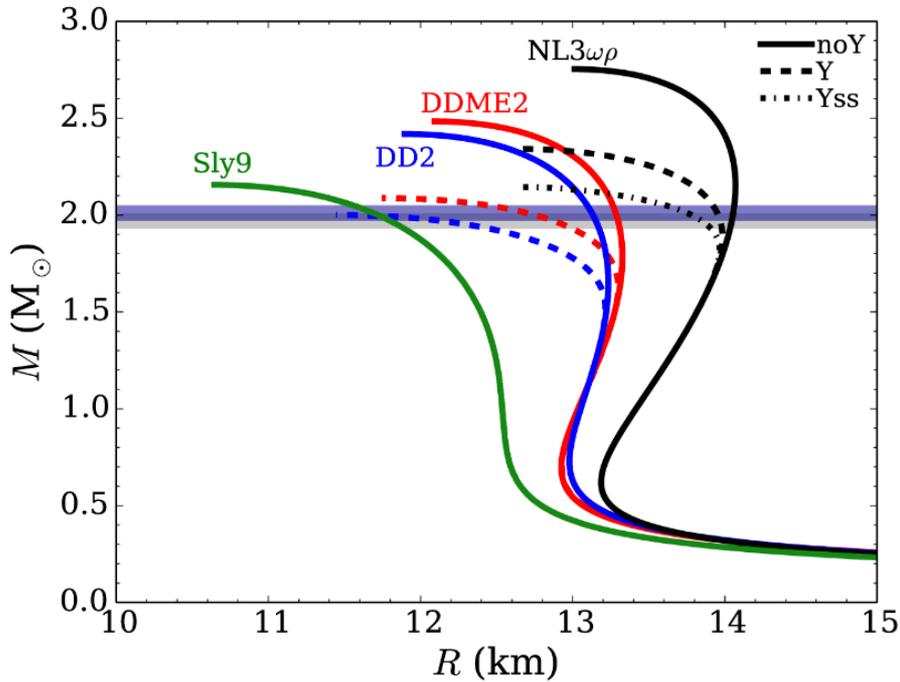
Neill, Preston, Tsang, Newton in prep

Resonant shattering flares: combining crust and core observables

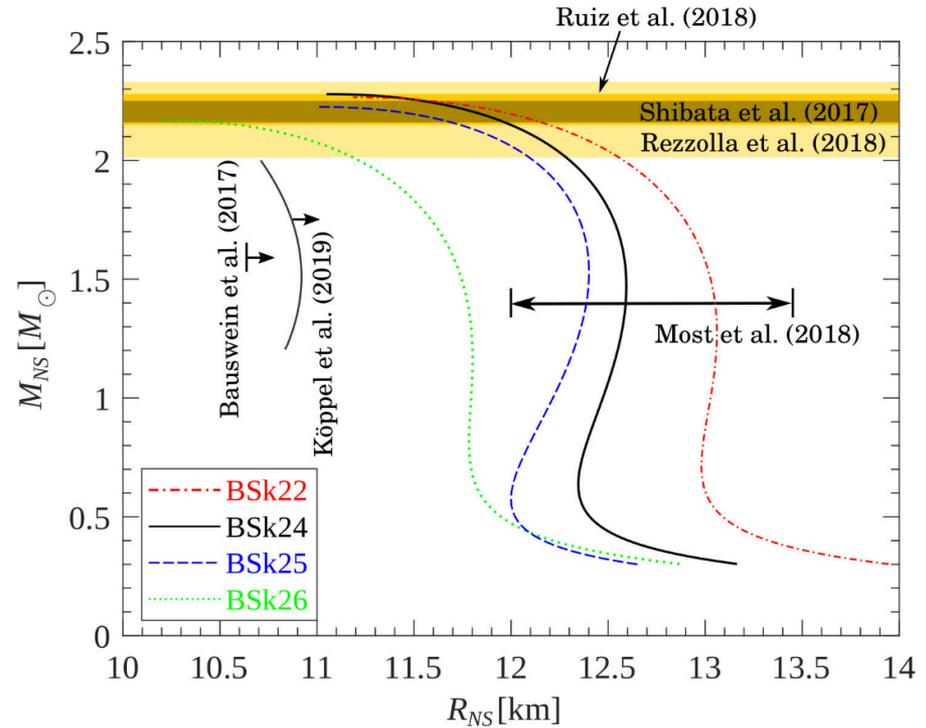
(See David's talk next Wednesday!)



Complementary approach: Precision model fits

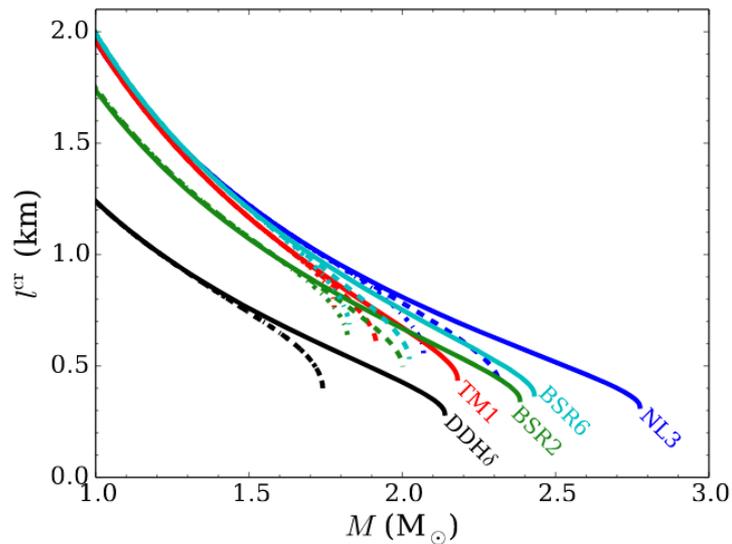
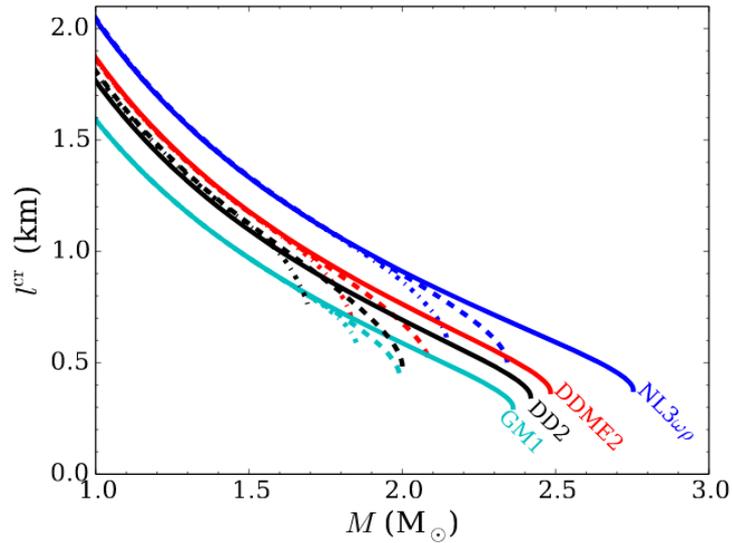


Fortin et al; arxiv:1604.01944

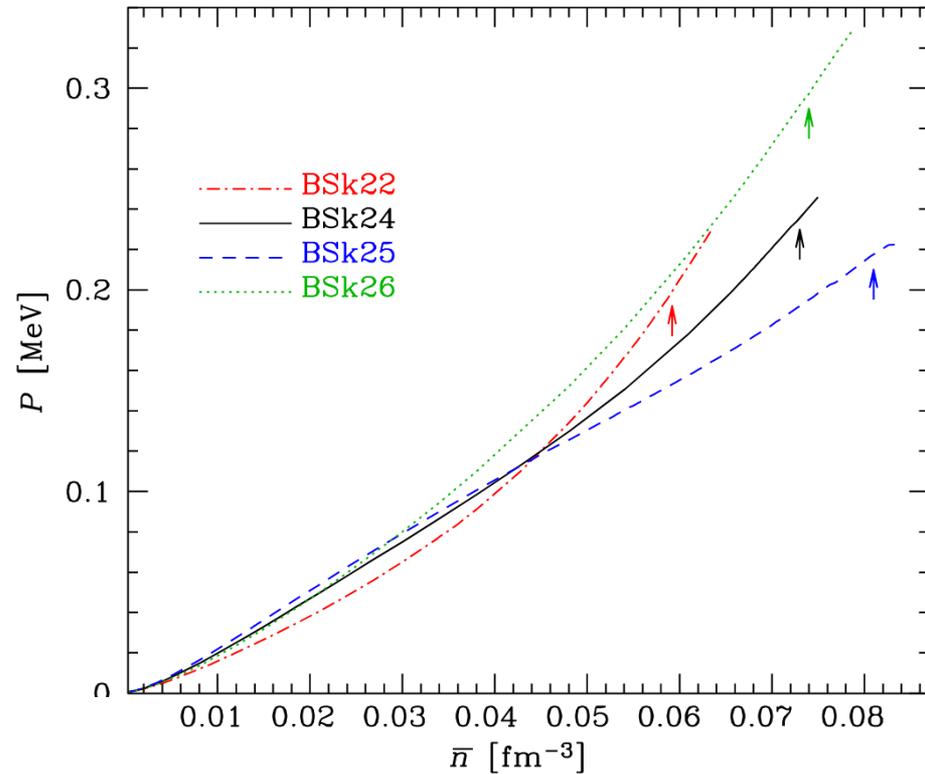


Pearson et al MNRAS481,2994–3026 (2018)
 Chamel et al arxiv:1904.12477

Complementary approach: Precision model fits

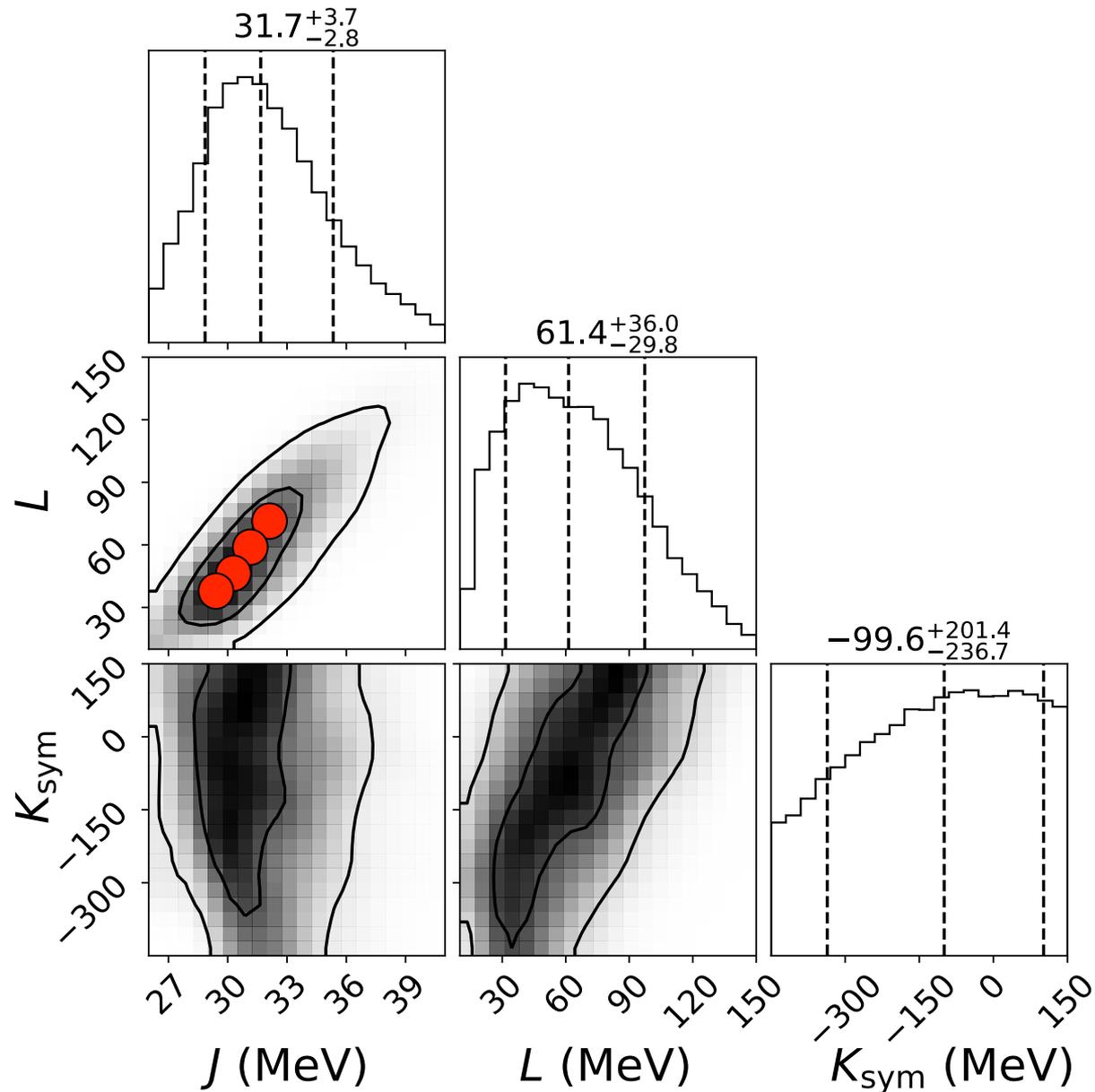


Fortin et al; arxiv:1604.01944

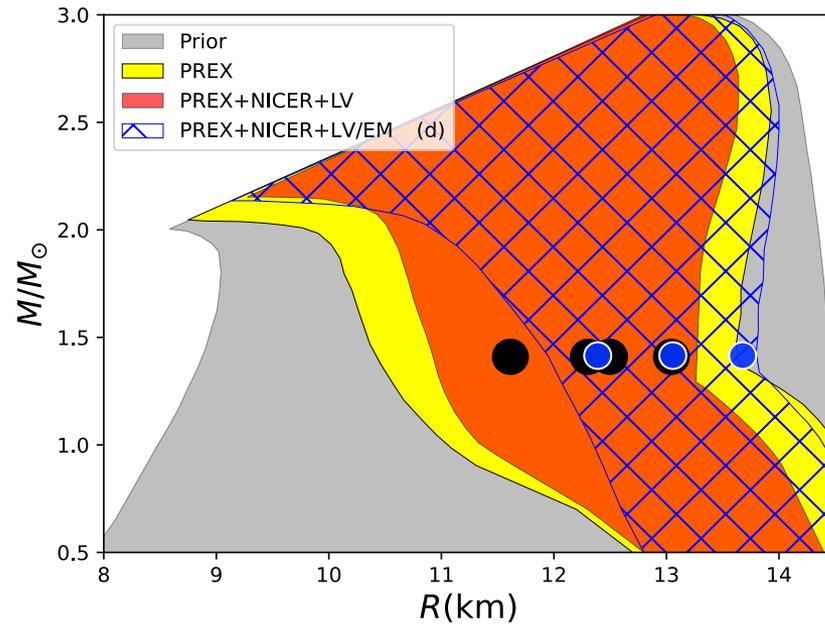
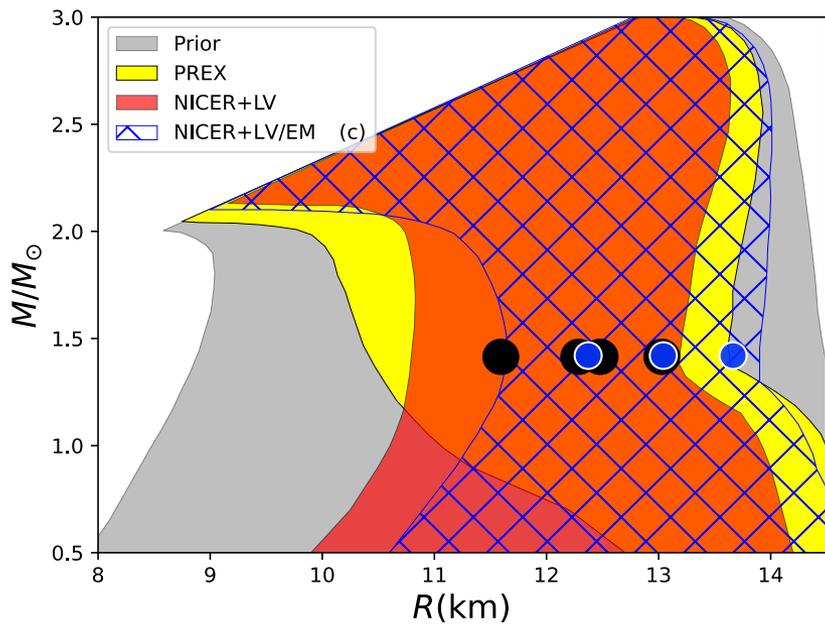


Pearson et al MNRAS481,2994–3026 (2018)
Chamel et al arxiv:1904.12477

Complementary approach: Precision model fits



Complementary approach: Precision model fits



Newton et al; arxiv:2112.12108

Some thoughts

Attempt at a framework to connect neutron star bulk and crust observables with nuclear data as consistently as possible

Consistency: eliminate systematic errors that may arise when different models are used to propagate information across multiple physical characteristics

Many rich astro datasets require crust modeling to interpret. We want to bring these into our multimessenger club

Not a replacement for precision modeling

Radius measurements are not going to be able to ignore the crust too much longer

Assumptions/limitations

Had to choose an EDF (Skyrme). Enough degrees of freedom?

Can add more (Q_{sym})

Symmetric nuclear matter and gradient parameters held fixed, underestimate model

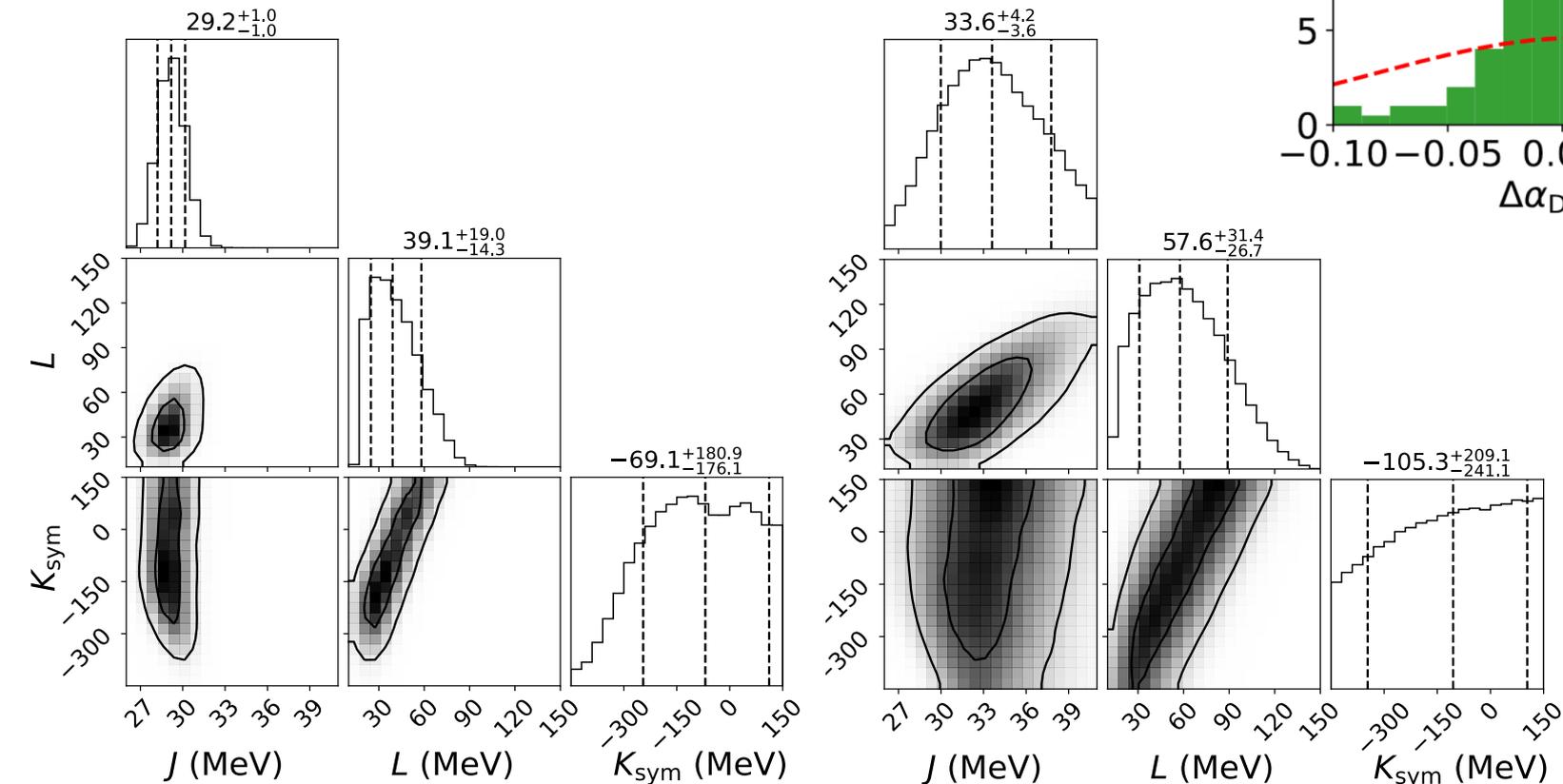
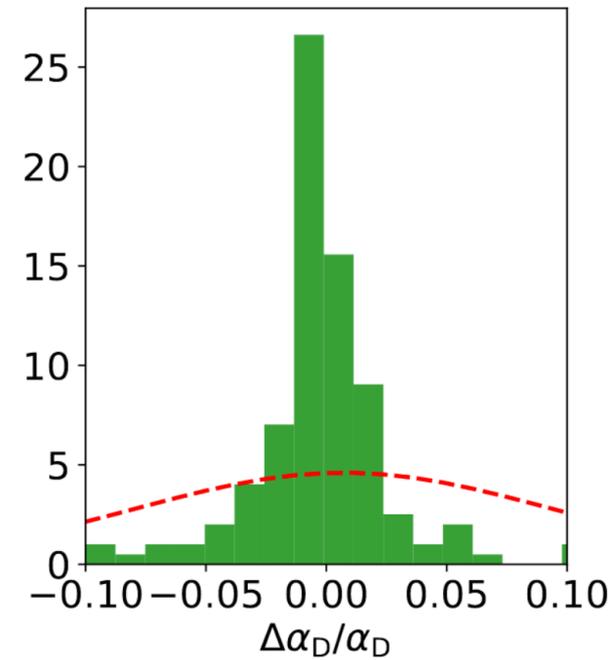
Need more information to constrain surface parameters of crust model (mass fits, semi-infinite nuclear matter, Thomas-Fermi calculations)

2 polytropes in the core is bare minimum: can be more sophisticated

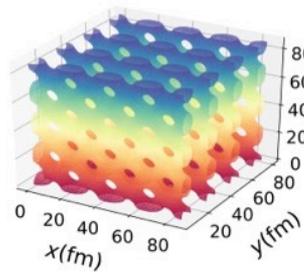
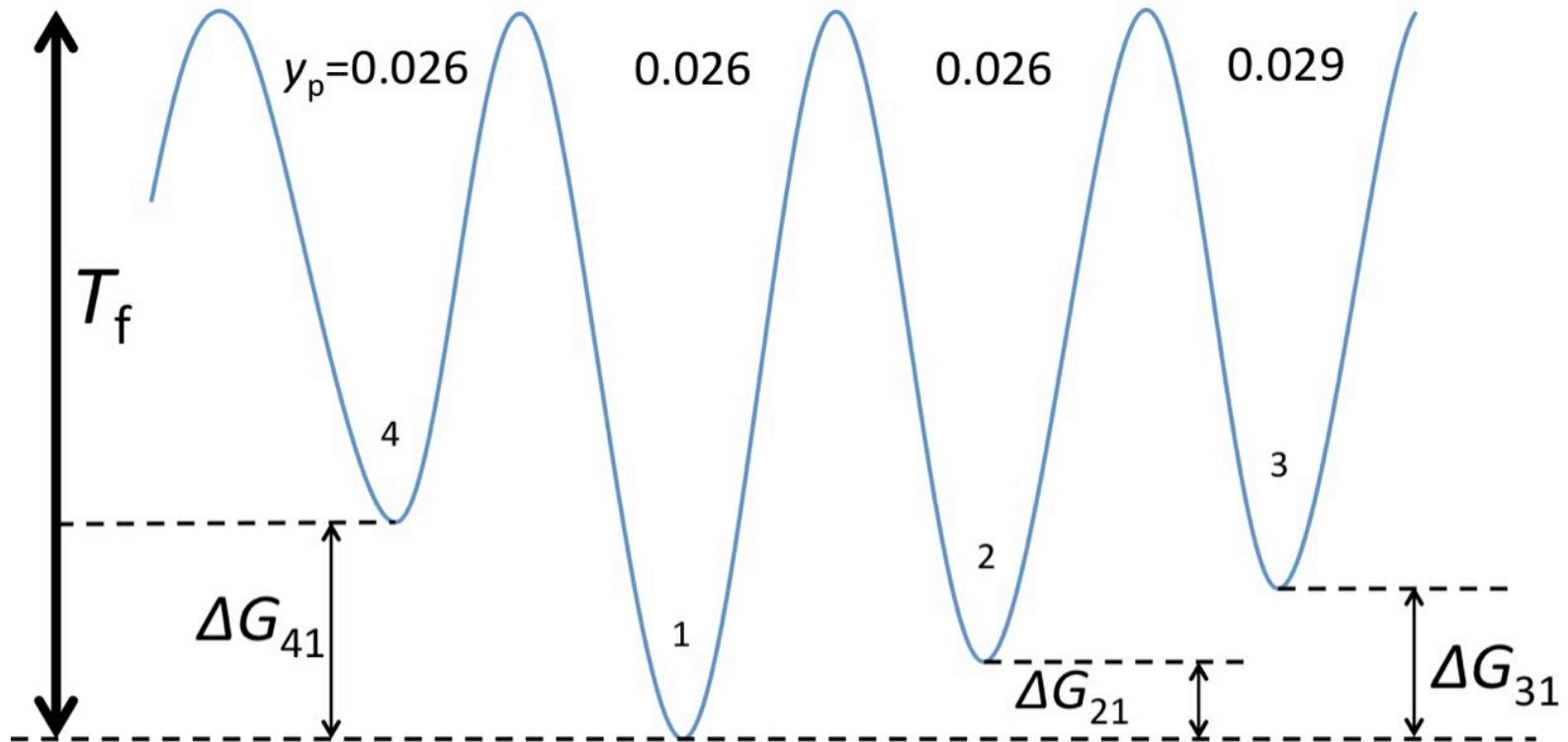
We're usually not directly modeling nuclear observables - but in some cases we can (e.g. weak form factor) and thus improve consistency

What about the other parameters?

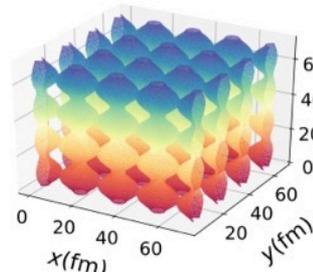
- Refit a subset of our 1000 Skyrmes using simulated annealing method
- Fit the resulting differences in observables with Gaussian – a conservative estimate



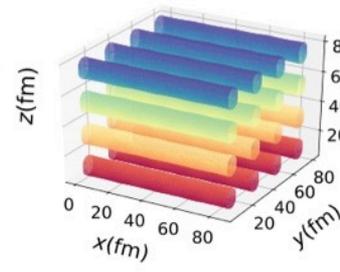
Pasta: a complex, glassy system. Multiple shapes coexist in microscopic domains. *Should* affect transport properties. How much is there?



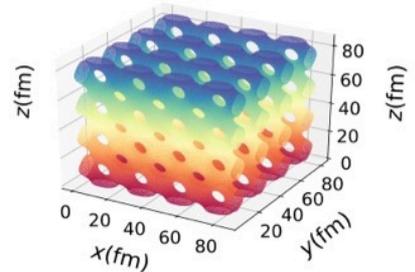
22%



30%



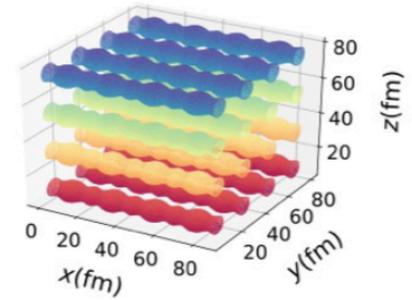
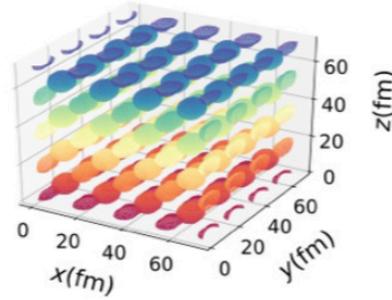
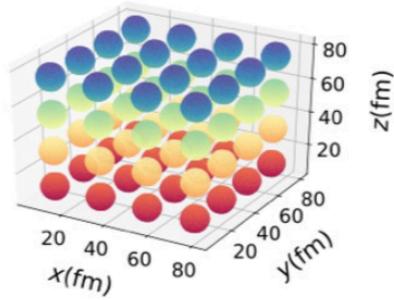
25%



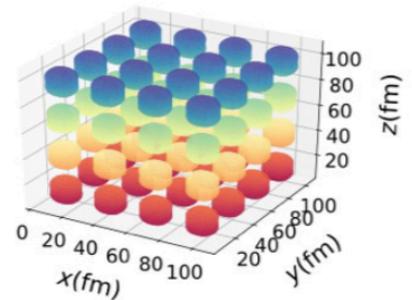
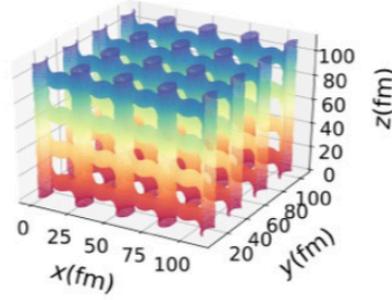
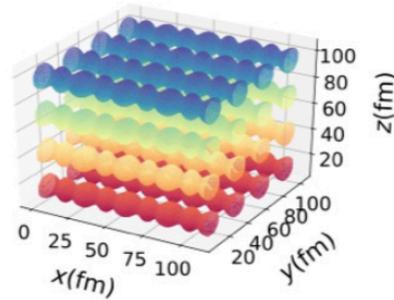
23%

Increasing energy of local minimum \rightarrow

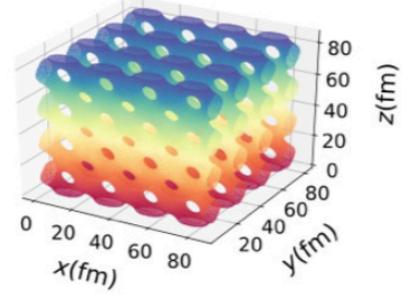
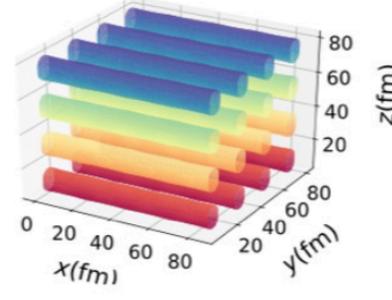
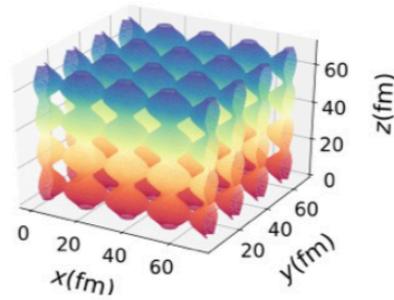
0.04-
0.05



0.05-
0.06

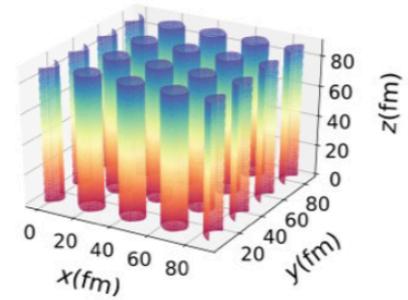
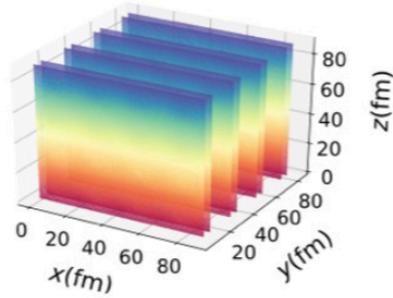
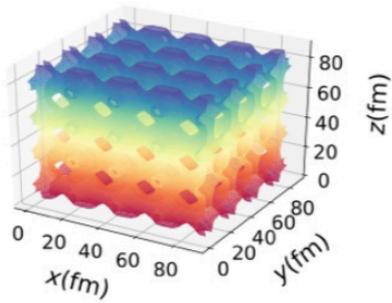


0.06-
0.07

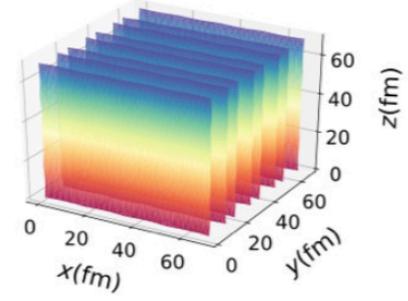
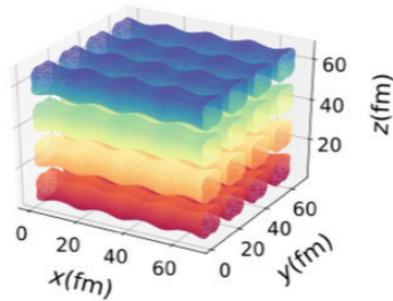
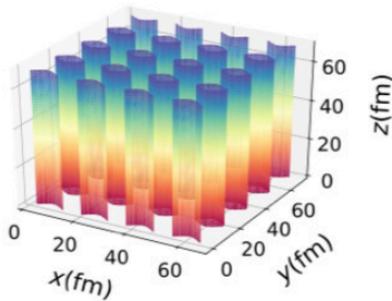
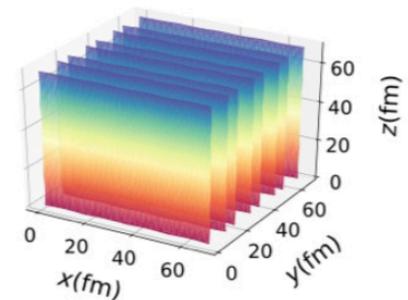
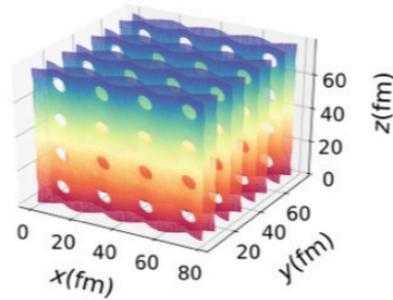
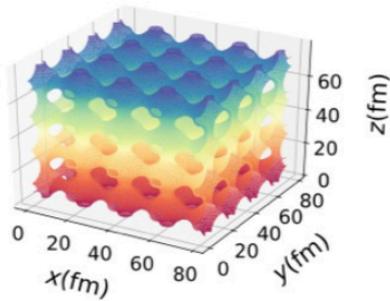


fm^{-3}

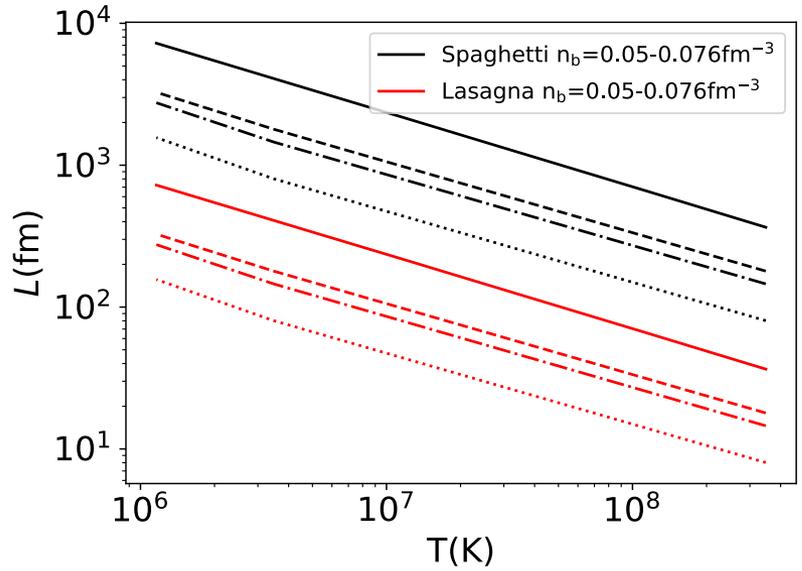
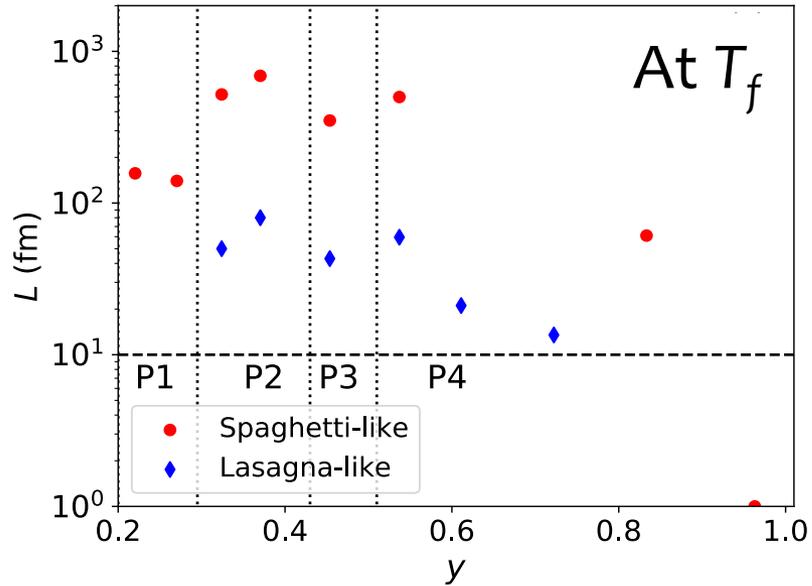
Increasing energy of local minimum \rightarrow



0.07-
0.09
 fm^{-3}



Possible sources of resistivity



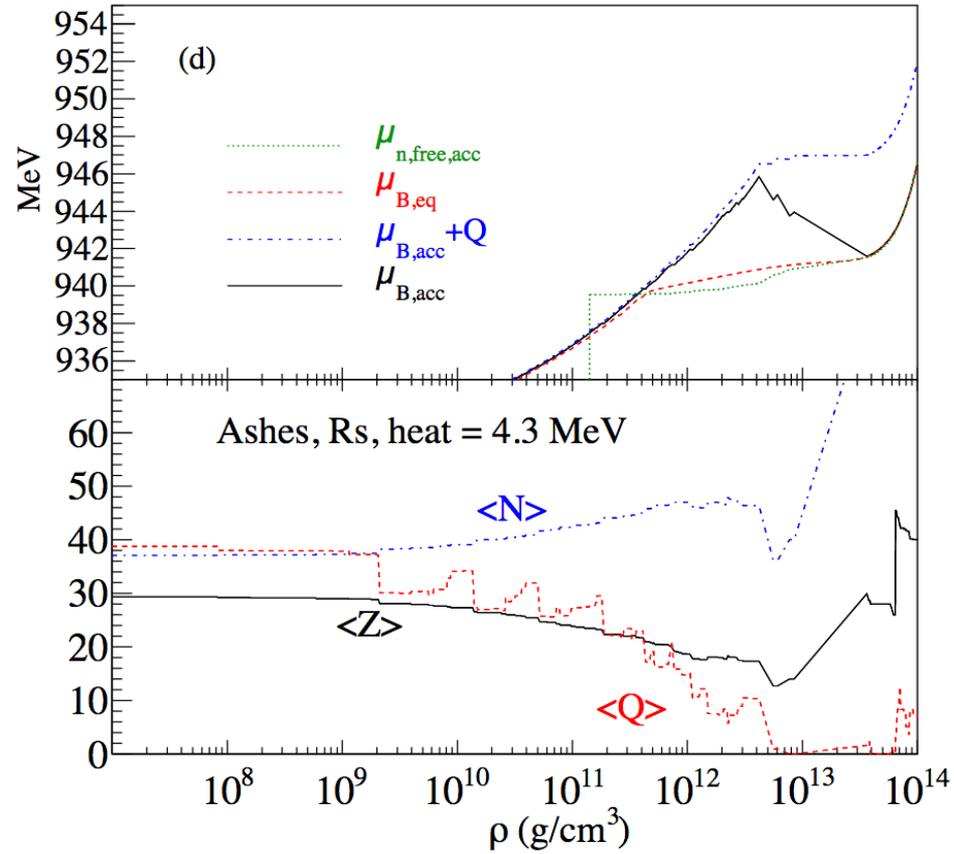
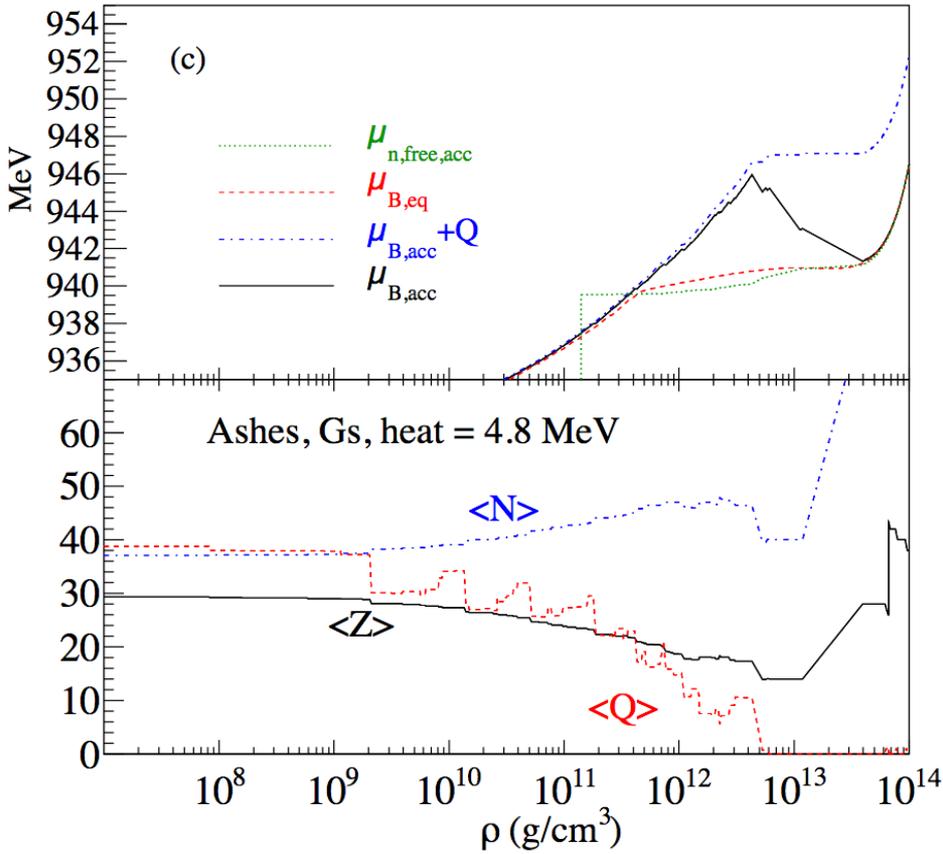
Electron scattering off domain boundaries
(annealing may lead to most energetically favorable domain growing)



Electron scattering off disordered pasta
(temperature dependent)

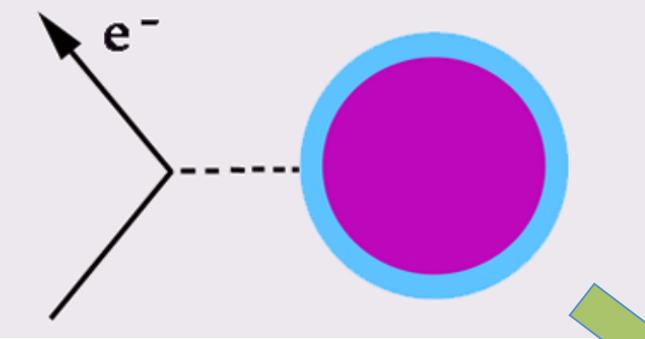


Accreted Crusts: Deep crustal heating, impurity (J,L,K_{sym})



Steiner, arxiv:1202.3378

Partially accreted: Sulieman et al, arxiv:2203.14735



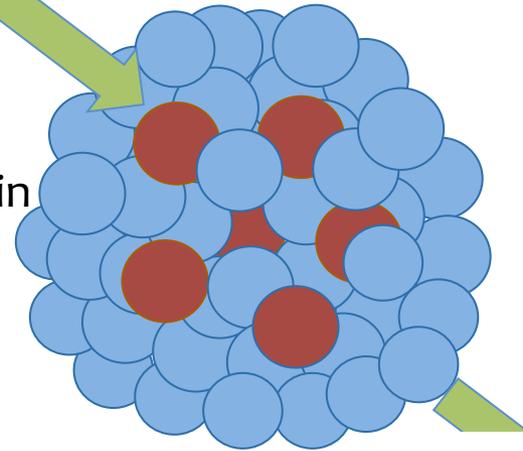
Scattering cross sections

Weak charge density

Set of EDFs

n-p densities

neutron skin



Us: completely different set of EDFs

