Hadrons, Type II Superconductors, and Cosmological Constant

PRD104,076010 (2021) [arXiv:2103.15768]

- Mass decomposition of hadrons from Hamiltonian and gravitational form factors
- Pressure balance of hadrons and role of trace anomaly
- Cosmological constant in Einstein's equation
- Type II superconductors
- Puzzle of pion mass and trace anomaly

June 17, 2022 INT

Motivation

Where does the proton mass come from, and how?

the proton is

masses!





The Higgs boson make the u/d quark having masses (2GeV MS-bar):

 $m_u = 2.08(9) \text{ MeV}$ $m_d = 4.73(12) \text{ MeV}$

Laiho, Lunghi, & Van de Water, Phys.Rev.D81:034503,2010

Mass (Rest Energy) from Hamiltonian

Energy momentum tensor (EMT)

$$T_{\mu\nu} = \frac{1}{4} \overline{\psi} \gamma_{(\mu} \overline{D}_{\nu)} \psi + G_{\mu\alpha} G_{\nu\alpha} - \frac{1}{4} \delta_{\mu\nu} G$$

Separate the EMT into traceless part and trace part

$$T_R^{\mu\nu} = \overline{T}_R^{\mu\nu} + \frac{1}{4}\eta^{\mu\nu}(T_\rho^\rho)_R$$

• Hamiltonian -- $H = \int d^3 \vec{x} T^{00}(x)$

X. Ji (1995)

With equation of motion (scale dependent)

$$H_m = \int d^3 \vec{x} \sum_f m_f \bar{\psi}_f \psi_f,$$

$$H_E(\mu) = \int d^3 \vec{x} \sum_f (\psi_f^{\dagger} i \vec{\alpha} \cdot \vec{D} \psi_f)_M,$$

$$pc$$

$$H_g(\mu) = \int d^3 \vec{x} \frac{1}{2} (B^2 + E^2)_M,$$

$$H_{tr} = \int d^3 \vec{x} \frac{1}{4} (T^{\mu}_{\mu})_R.$$

Quark kinetic and potential energy

Glue field energy

 $E_0 = M = \langle H_M \rangle + \langle H_E(\mu) \rangle + \langle H_g(\mu) \rangle + \langle H_{tr} \rangle,$

Proton Mass Decomposition

Lattice calculation with systematics (physical pion mass, continuum, infinite volume extrapolations, renormalization)



Y.B. Yang et al (χ QCD), PRL 121, 212001 (2018) Physic 11, 118 (2018); ScienceNews, Nov. 16 (2018)

Experimentally Measurable Decomposition

Separate the EMT into traceless part and trace part $T_R^{\mu\nu} = \overline{T}_R^{\mu\nu} + \frac{1}{4}\eta^{\mu\nu}(T_\rho^\rho)_R$ Hamiltonian -- $H = \int d^3 \vec{x} T^{00}(x)$ $H_{q}(\mu) = \int d^{3}\vec{x} \left(\frac{i}{4}\sum_{f}^{J} \bar{\psi}_{f} \gamma^{\{0} \overleftrightarrow{D}^{0\}} \psi_{f} - \frac{1}{4}T^{\mu}_{q\,\mu}\right)_{M},$ Quark momentum fraction (scale dependent) Glue momentum fraction $H_g(\mu) = \int d^3 \vec{x} \, \frac{1}{2} (B^2 + E^2)_M,$ (scale dependent) $H_{tr} = \int d^3 ec{x} \, rac{1}{4} (T^{\mu}_{\mu})_R \cdot T^{\mu}_{\mu} = \sum_f m_f ar{\psi}_f \psi_f + \left| \sum_f m_f \gamma_m(g) ar{\psi}_f \psi_f + rac{eta(g)}{2g} F^{lphaeta} F_{lphaeta}
ight|$ $E_0 = M = \langle H_{q_f}(\mu) \rangle + \langle H_q(\mu) \rangle + \langle H_{tr} \rangle,$ Rest energy -- $\langle H_{q_f}(\mu) \rangle = \frac{3}{4} \sum \langle x \rangle_f(\mu) M, \quad \langle H_g(\mu) \rangle = \frac{3}{4} \langle x \rangle_g(\mu) M,$ <x> - momentum fraction $\langle H_{tr} \rangle = \frac{1}{4}M.$ 5

Mass Decomposition from Hamiltonian



$$f_f^H = \langle H_q \rangle / M = \frac{3}{4} \langle x \rangle_f(\mu), \quad f_g^H = \langle H_g \rangle / M = \frac{3}{4} \langle x \rangle_g(\mu),$$
$$f_{\pi N}^N = \frac{1}{4} \frac{\sigma_{\pi N}}{M}, \quad f_s^N = \frac{1}{4} \frac{\sigma_s}{M}, \quad f_{\text{trace anomaly}} = \frac{1}{4} \frac{\langle H_{\text{ta}} \rangle}{M}$$

Y.B. Yang et al (χ QCD) [arXiv: 1511.15089]

Momentum fractions from CT18 (T.J. Hou et al, PRD, arXiv:1912.10053) at $\mu = 2$ GeV and 250 GeV.

Mass from Gravitational FF

- Gravitational Form factors from the EMT matrix elements $\langle P'|(T_{q,g}^{\mu\nu})_R(\mu)|P\rangle/2M_N = \bar{u}(P')[T_{1_{q,g}}(q^2,\mu)\gamma^{(\mu}\bar{P}^{\nu)} + T_{2_{q,g}}(q^2,\mu)\frac{\bar{P}^{(\mu}i\sigma^{\nu)\alpha}q_{\alpha}}{2M_N}$ $+ D_{q,g}(q^2,\mu)\frac{q^{\mu}q^{\nu} - g^{\mu\nu}q^2}{M_N} + \bar{C}_{q,g}(q^2,\mu)M_N\eta^{\mu\nu}]u(P)$
 - $-T_1$ and T_2 : momentum and angular momentum [Ji]
 - D term: deformation of space = elastic property [Polyakov]
 - C-bar term: pressure-volume work [Lorce, Liu]

$$\begin{split} M_{N}(q,g) &= \langle P|(T_{q,g}^{00})_{\mathrm{M}}(\mu)|P\rangle|_{\vec{P}=0}/2M_{N} = \langle x \rangle_{q,g}(\mu)M_{N} + \bar{C}_{q,g}(0,\mu)M_{N} \\ \langle P|(T_{q,g}^{ii})_{\mathrm{M}}(\mu)|P\rangle|_{\vec{P}=0}/2M_{N} = -3\bar{C}_{q,g}(0,\mu)M_{N} \\ \bar{C}_{q} + \bar{C}_{g} &= \frac{1}{4} \left(\sum_{f} f_{f}^{N} + f_{a}^{N} \right) - \frac{1}{4} \left(\langle x \rangle_{q} + \langle x \rangle_{q} \right) = 0 \quad \partial_{\nu}T^{\mu\nu} = 0 \\ M_{N} &= \frac{3}{4} (\langle x \rangle_{q}(\mu) + \langle x \rangle_{g}(\mu))M_{N} + \frac{1}{4} (\sum_{f} f_{f}^{N} + f_{a}^{N})M_{N} \quad \begin{array}{c} \text{Same as from} \\ \text{Hamiltonian} \end{array} \end{split}$$

Mass and Pressure from Gravitational FF

• Mass $M_N = M_N(q) + M_N(g)$

 $M_N(q,g) = \langle P|(T_{q,g}^{00})_{\mathcal{M}}(\mu)|P\rangle|_{\vec{P}=0}/2M_N = \langle x\rangle_{q,g}(\mu)M_N + \bar{C}_{q,g}(0,\mu)M_N$

Note: Being scale dependent, separate quark and glue T⁰⁰ are renormalized and mixed.

What are \bar{C}_q and \bar{C}_g ? $\langle P|(T_{q,g}^{ii})_{\mathrm{M}}(\mu)|P\rangle|_{\vec{P}=0}/2M_N = -3\bar{C}_{q,g}(0,\mu)M_N$ $3\bar{C}_{q,g}(0,\mu)M_N = [\langle P|\eta_{\mu\nu}(T_{q,g}^{\mu\nu})_{RM}|P\rangle - \langle P|(T_{q,g}^{00})_{RM}(\mu)|P\rangle]/2M_N$ $\bar{C}_q(0,\mu) = \frac{1}{4}\sum_f (f_f^N - \langle x \rangle_f(\mu)), \ \bar{C}_g(0,\mu) = \frac{1}{4}(f_a^N - \langle x \rangle_g(\mu))$

Therefore,

 $M_N = \frac{3}{4} \left(\langle x \rangle_q(\mu) + \langle x \rangle_g(\mu) \right) M_N + \frac{1}{4} \left(\sum_f f_f^N + f_a^N \right) M_N$

the same as from the Hamiltonian.

Role of Trace Anomaly --String tension in charmonium

- Heavy quarkonium is confined by a linear potential.
- Constant vacuum energy density and flux tube

 $V(r) = |\epsilon_{vac}| A r = \sigma r$

Infinitely heavy quark with Wilson loop

 $V(r) + r \frac{\partial V(r)}{\partial r} = \frac{\langle \frac{\beta}{2g} (\int d^3 \vec{x} F^2) W_L(r, T) \rangle}{\langle W_L(r, T) \rangle}.$ For charmonium $2\sigma \langle r \rangle = \langle H_\beta \rangle_{\bar{c}c} = \frac{\langle \bar{c}c | \frac{\beta}{2g} \int d^3 \vec{x} F^2 | \bar{c}c \rangle}{\langle \bar{c}c | \bar{c}c \rangle}.$

Dosch, Nachtmann, Rueter - 9503386; Rothe - 9504102

$$\langle H_{\beta} \rangle_{\bar{c}c} = M_{\bar{c}c} - (1 + \gamma_m) \langle H_m \rangle_{\bar{c}c}.$$

- Lattice calculation of charmonium (W. Sun et al.,2012.06228) $\langle H_{\beta} \rangle_{\bar{c}c} = 199 \,\mathrm{MeV} \rightarrow \sigma = 0.153 \,\mathrm{GeV^2}$
- Cornell potential fit of charmonium $\rightarrow \sigma = 0.164(11) \text{ GeV}^2$

Mateu:2018zym

Trace Anomaly and Gluon Condensate

What is trace anomaly? What dynamical role does it play, if any?

- Note $\langle P|(T_{q,g}^{ii})_{\mathrm{M}}(\mu)|P\rangle|_{\vec{P}=0}/2M_N = -3\bar{C}_{q,g}(0,\mu)M_N$ • $\frac{1}{3}\langle P|(T_{q,g}^{ii})|P\rangle|_{\vec{P}=0}$ is pressure-volume work
- The pressure balance equation



 $PV = -\frac{dE}{dV}V = -(\bar{C}_q + \bar{C}_g) = \frac{1}{4}(\langle x \rangle_q + \langle x \rangle_q) - \frac{1}{4}(\sum_f f_f^N + f_a^N) = 0$

Nucleon is a bubble in the sea of gluon condensate, where

$$\langle H_a \rangle = -\epsilon_{vac} V, \\ \epsilon_{vac} = \frac{\beta(g)}{2g} \langle 0 | F^{\alpha\beta} F_{\alpha\beta} | 0 \rangle < 0$$

$$\begin{split} E_0 &= E_T + E_S, \\ E_S &= \frac{1}{4} [\langle H_m \rangle + \langle H_a \rangle] \propto V, \\ E_T &= \langle H_{q_f}(\mu) \rangle + \langle H_g(\mu) \rangle \propto V^{-1/3} \\ \end{split}$$
 (MIT bag model, E₀ = BV+ $\Sigma_{a,a}$ /R)



Mass-Pressure Correspondence

Mass
$$M_N = \frac{3}{4} (\langle x \rangle_q(\mu) + \langle x \rangle_g(\mu)) M_N + \frac{1}{4} (\sum_f f_f^N + f_a^N) M_N$$

Pressure-volume work

 $PV = \frac{1}{4} (\langle x \rangle_q(\mu) + \langle x \rangle_g(\mu)) M_N - \frac{1}{4} (\sum_f f_f^N + f_a^N) M_N$

Virial theorem

D dimension

- Mass
$$M = E_T + E_S, \quad E_T = \langle \overline{T}^{00} \rangle = \frac{D-1}{D}M, \quad E_S = \frac{1}{D} \langle T^{\mu}_{\mu} \rangle = \frac{1}{D}M$$

- Pressure $PV = \frac{1}{D-1} \langle T^{ii} \rangle = -\frac{1}{D-1} [\langle T^{\mu}_{\mu} \rangle - \langle T^{00} \rangle]$ = $-\frac{1}{D-1} [DE_s - E_T - E_s] = -E_s + \frac{1}{D-1} E_T = 0$

Mass-Pressure Correspondence

Other mass decomposition formulae

 $- \operatorname{Trace}_{\frac{\langle P| \int d^3 \vec{x} \, T^{\mu}_{\mu}(x) |P \rangle}{\langle P|P \rangle}|_{\vec{P}=0}} = M_N$

$$T^{\mu}_{\mu} = \sum_{f} m_{f} \bar{\psi}_{f} \psi_{f} + \left[\sum_{f} m_{f} \gamma_{m}(g) \bar{\psi}_{f} \psi_{f} + \frac{\beta(g)}{2g} F^{\alpha\beta} F_{\alpha\beta} \right]$$

No kinetic energy, no connection to pressure

Gravitational FF without C

 $M_N = (\langle x \rangle_q(\mu) + \langle x \rangle_g(\mu))M_N$

No potential energy, no relation to pressure

Virial Theorem

- Coulomb potential: H = T + V, 2 < T > = < V >
- $E = -\frac{1}{2}$ <T>, $E = \frac{1}{2}$ <V> are not good physics explanation of the decomposition of the binding energy.

Generic bound state energy and pressure correspondence Hadronic model mass – pressure relation

- MIT bag model ($E_0 = BV + \Sigma_{q,g}/R$), pressure: $\partial E_0/\partial R = 0$

- Skyrmion: Derrick's theorem $(r \rightarrow \lambda r)$ $\mathcal{L}_2 = f_{\pi}^2 Tr \left[\partial_{\mu} U^{\dagger} \partial^{\mu} U\right] \longrightarrow \lambda$ $\mathcal{L}_4 = \frac{1}{e^2} Tr \left[U^{\dagger} \partial_{\mu} U, U^{\dagger} \partial_{\nu} U\right]^2 \longrightarrow 1/\lambda$
- Potential models: kinetic energy ↘ r, potential energy ↗ r.

Feynman's quote

From the very beginning of his first-ever lecture comes this timeless gem (mentioned in Daniel Bor's excellent <u>The</u> Ravenous Brain) that set the tone for both Feynman's academic contribution and his broader cultural legacy: If, in some cataclysm, all of scientific knowledge were to be destroyed, and only one sentence passed on to the next generation of creatures, what statement would contain the most information in the fewest words?



I believe it is the atomic hypothesis that all things are made of atoms — little particles that move around in perpetual motion, attracting each other when they are a little distance apart, but repelling upon being squeezed into one another. In that one sentence, you will see, there is an enormous amount of information about the world, if just a little imagination and thinking are applied.

Trace Anomaly and Cosmological Constant

Vacuum energy density is indeed a constant which is analogous to the cosmological constant in the g^{µν} term as Einstein introduced for a static universe.

$$R_{\mu\nu} + \frac{1}{2}Rg_{\mu\nu} - \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu} \qquad \Lambda = 4\pi G\rho$$

Friedman equation for the accelerating expansion of the universe

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P) + \frac{\Lambda}{3} \qquad \qquad \rho_{\text{vac}} = \frac{\Lambda}{8\pi G}$$
$$P_{\text{vac}} = -\frac{\Lambda}{8\pi G}$$

Type II Superconductor



Ginzburg–Landau equations

$$lpha\psi+eta|\psi|^2\psi+rac{1}{2m}(-i\hbar
abla-2e{f A})^2\psi=0$$

$$abla imes {f B} = \mu_0 {f j} \hspace{0.2cm} ; \hspace{0.2cm} {f j} = rac{2e}{m} \operatorname{Re} \{ \psi^* \left(-i \hbar
abla - 2e {f A}
ight) \psi \}$$

$$|\psi|^2 = n_s$$

Physics of type I and II superconductors

• "London Penetration Depth" λ_{L}

is the e-fold decay length of the magnetic field from the superconductor skin due to the Meissner effect (in the range of 10 to $10^3~\rm nm)$

"Coherence Length" ξ

the average size of Cooper in the superconductor (in the range of 10 to 100 nm, I.e. much larger than the interatomic distance typically of 0.1 to 0.3 nm.



Nb 32 39

Ginzburg-Landau Parameter κ

$$\kappa = \frac{\lambda_L}{\xi} \Rightarrow \begin{cases} k < \frac{1}{\sqrt{2}} \Leftrightarrow \text{ type I} \\ k > \frac{1}{\sqrt{2}} \Leftrightarrow \text{ type II} \end{cases} \begin{bmatrix} \text{material In Pb Sn} \\ \lambda_L \text{ [nm]} & 24 & 32 & \approx 30 \\ \hline \xi \text{ [nm]} & 360 & 510 & \approx 170 \end{cases}$$

London penetration depth

$$\Lambda_L = \sqrt{rac{m}{4\mu_0 e^2 \psi_0^2}}$$

Coherent length ξ

$$\kappa = \lambda_L / \xi$$

Energetics and Pressure

Normal Phase

Type II superconductor

 $F = F_s + F_B + F_{sc}$ $F_s = \text{ cost of condensation energy}$ $F_B = \int dv B^2 / 2\mu_0 \text{ (magnetic energy)}$ $F_{sc} = 1/2 \int dv \lambda_L^2 J_s \cdot J_s \text{ (supercurrent kinetic energy)}$ Superconducting Phase

• Variational model (J.R. Clem, Jour. Low Temp. Phys. 18, 5/6 (1975)) $\frac{|\psi|^2}{n_0} = \frac{n_s}{n_0} = \frac{\rho^2}{\rho^2 + R^2} \xrightarrow{\rho \to \infty} 1$ $\frac{1}{\sqrt{2}H_c} \frac{E}{l} = \phi_0 H'_c / 4\pi \text{ where } \phi_0 = hc/2e, \ \sqrt{2}H_c = \kappa \phi_0 / 2\pi \lambda_L^2$ $H'_c = \kappa R' / 8 + 1 / 8\kappa + K_0(R') / 2\kappa R' K_1(R'), \text{ where } R' = R / \lambda_L$ $F_s \qquad F_B + F_{sc}$

SC, Hadrons, Cosmos

Type II Superconductor

 $P_s = -\frac{\partial F_s}{\partial V} < 0, \quad P_{B+sc} = -\frac{\partial F_B + F_{sc}}{\partial V} > 0$ Hadrons

$$P_{tr} = -\frac{\partial E_S}{\partial V} < 0, \quad P_{q+g} = -\frac{\partial E_T}{\partial V} > 0$$

- Cosmos $\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3P) + \frac{\Lambda}{3}, \quad \Lambda > 0 \longrightarrow \frac{\ddot{a}}{a} > 0$
- The common theme is the existence of a condensate.
- Hadrons: condensates from breaking of conformal and chiral symmetries. SC: Cooper pair condensate from gauge symmetry breaking. Cosmos: ?

Pion Mass Puzzle

Pion mass in terms of trace of EMT

 $m_{\pi} = m \langle \pi | \bar{\psi} \psi | \pi \rangle + \langle \pi | \frac{\beta}{2g} G_{\mu\nu}^2 + m \gamma_m \bar{\psi} \psi | \pi \rangle$

Gellmann-Oakes-Renner relation

 $m_{\pi}^2 = -2m\langle \bar{\psi}\psi \rangle / f_{\pi}^2, \ \ m_{\pi}^2 \propto m$



■ $\langle \pi | \bar{\psi} \psi | \pi \rangle \propto 1/\sqrt{m}$ But, why should the trace anomaly be proportional to \sqrt{m} ? V → 0 ?



Y.B. Yang et al. (χQCD), PRD (2015); 1405.4440

Trace anomaly Distribution



It changes sign in pion so that the integral approaches zero at the chiral limit.

Pion as a Ring-shaped Type II Superconductor



A. Groeger et al., PRL 90, 237004 (2003)



Pion with shell of positive trace anomaly and an inner core with negative trace anomaly. Niobium, normal conducting vortex ring around a superconducting region,

Pion trace anomaly FF



$$m_{\pi} = 340 \text{ MeV}$$

X.B. Tong, J.P. Ma and F. Yuan arXiv:2203.13493

Summary and Challenges

- From femto-scale to micro-scale to that of the cosmos, Nature seems to choose the same mechanism for confinement or acceleration.
- $m_q \leftarrow Higgs mechanism$
- Quark condensate \leftarrow chiral symmetry breaking (restoration at T and μ)
- Trace anomaly (confinement) ← conformal symmetry breaking (conformal phases with multi-flavors and SU(N); finite T > T_c)
- Chiral symmetry breaking and conformal symmetry breaking are linked in the case of the pion trace anomaly distribution.
- String theory invented in hadron physic finds its home in quantum gravity.
- Cosmological constant introduced in general relativity is relevant to hadron physics.
- Challenges for EIC and COMPASS is to measure the trace anomaly form factors for the proton and, particularly, the pion.
- Glue condensate is an order parameter for confinement deconfinement transition.

Virial Theorem

- D dimension
- Mass

$$M = E_T + E_S, \quad E_T = \langle \overline{T}^{00} \rangle = \frac{D-1}{D}M, \quad E_S = \frac{1}{D} \langle T^{\mu}_{\mu} \rangle = \frac{1}{D}M$$

Pressure

$$PV = \frac{1}{D-1} \langle T^{ii} \rangle = -\frac{1}{D-1} [\langle T^{\mu}_{\mu} \rangle - \langle T^{00} \rangle]$$
$$= -\frac{1}{D-1} [DE_s - E_T - E_s] = -E_s + \frac{1}{D-1} E_T = 0$$

Trace of EMT (1/4 of Hadron Mass)

- Trace of EMT scalar, frame independent, RG invariant $T^{\mu}_{\mu} = \sum_{f} m_{f} \bar{\psi}_{f} \psi_{f} + \left[\sum_{f} m_{f} \gamma_{m}(g) \bar{\psi}_{f} \psi_{f} + \frac{\beta(g)}{2g} F^{\alpha\beta} F_{\alpha\beta} \right]$
- Lattice calculation of of quark condensate
 - Y.B. Yang et al (χ QCD) [arXiv: 1511.15089]
 - Overlap fermion ($Z_m Z_s = 1$)
 - 3 lattices (one at physical m_{π}), systematics (volume, continuum)



 πN sigma term $\sigma_{\pi N} = \frac{m_u + m_d}{2} \langle P | \bar{u}u + \bar{d}d | P \rangle$

Strangeness sigma term

$$\sigma_s = m_s \langle P | \bar{s}s | P \rangle$$

$$f_{\pi N}^{N} = \frac{\sigma_{\pi N}}{M_{N}}, \qquad f_{s}^{N} = \frac{\sigma_{s}}{M_{N}}$$
 6

Mass and Pressure from Gravitational FF

• Mass
$$M_N = M_N(q) + M_N(g)$$

 $M_N(q,g) = \langle P | (T_{q,g}^{00})_{\mathcal{M}}(\mu) | P \rangle |_{\vec{P}=0} / 2M_N = \langle x \rangle_{q,g}(\mu) M_N + \bar{C}_{q,g}(0,\mu) M_N$

Note: Being scale dependent, separate quark and glue T⁰⁰ are renormalized and mixed.

• What are \bar{C}_q and \bar{C}_g ? $\langle P|(T_{q,g}^{ii})_{\mathrm{M}}(\mu)|P\rangle|_{\vec{P}=0}/2M_N = -3\bar{C}_{q,g}(0,\mu)M_N$ $3\bar{C}_{q,g}(0,\mu)M_N = [\langle P|\eta_{\mu\nu}(T_{q,g}^{\mu\nu})_{RM}|P\rangle - \langle P|(T_{q,g}^{00})_{RM}(\mu)|P\rangle]/2M_N$ $\bar{C}_q(0,\mu) = \frac{1}{4}\sum_f (f_f^N - \langle x \rangle_f(\mu)), \ \bar{C}_g(0,\mu) = \frac{1}{4}(f_a^N - \langle x \rangle_g(\mu))$

• Therefore,

 $M_N = \frac{3}{4} \left(\langle x \rangle_q(\mu) + \langle x \rangle_g(\mu) \right) M_N + \frac{1}{4} \left(\sum_f f_f^N + f_a^N \right) M_N$

the same as from the Hamiltonian.

Trace Anomaly and Cosmological Constant

Pressure of anomaly: $d\langle H_a \rangle = -P_{vac} \, dV (dQ = T \, dS = 0), \qquad P_{vac} = -|\epsilon_{vac}| < 0$ Quark and glue energy $\langle H_E(\mu) \rangle + \langle H_a(\mu) \rangle \propto V^P$ Volume dependence of total rest energy $E_0 = |\epsilon_{vac}|V + \epsilon_{mat}V^p$ $\frac{dE_0}{dV} = -P_{vac} - P_k = |\epsilon_{vac}| + p \epsilon_{mat} V^{p-1} = 0$ • $E_0 = d E_S (d=4) \implies p = -1/3$ (MIT bag model, $E_0 = BV + \Sigma_{q,q}/R$) Rest energy as the sum of scalar trace and tensor traceless parts $E_0 = E_T + E_S$

$$E_T = \langle H_{q_f}(\mu) \rangle + \langle H_g(\mu) \rangle = \frac{3}{4} \left[\sum_f \langle x \rangle_f(\mu) + \langle x \rangle_g(\mu) \right] M,$$

$$E_S = \frac{1}{4} [\langle H_m \rangle + \langle H_a \rangle]$$

Pion Mass Puzzle

• Pion mass in terms of trace of EMT

 $m_{\pi} = m \langle \pi | \bar{\psi} \psi | \pi \rangle + \langle \pi | \frac{\beta}{2g} G_{\mu\nu}^2 + m \gamma_m \bar{\psi} \psi | \pi \rangle$

Gellmann-Oakes-Renner relation

 $m_{\pi}^2 = -2m\langle \bar{\psi}\psi \rangle / f_{\pi}^2, \ \ m_{\pi}^2 \propto m$



• $\langle \pi | \bar{\psi} \psi | \pi \rangle \propto 1/\sqrt{m}$ But, why should the trace anomaly be proportional to ? V \rightarrow 0 ?



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