

SUBA-Jet

A New Coherent Jet Energy Loss Model For Heavy Ion Collisions

Alexander Lind

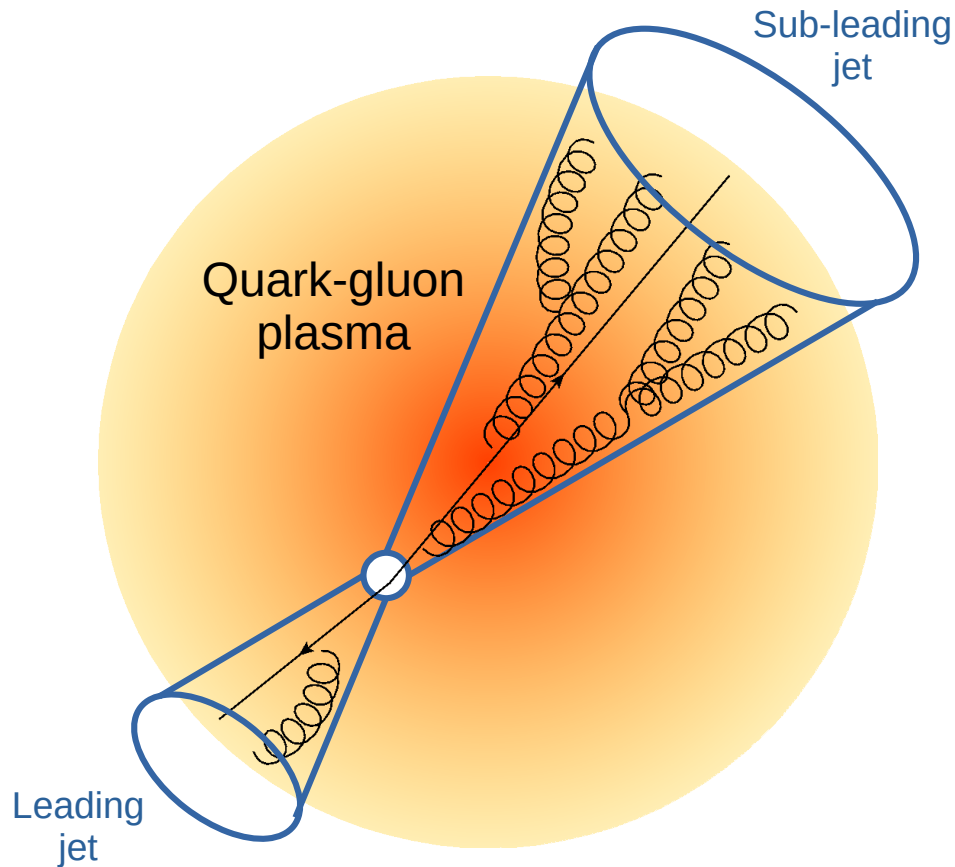
with Iurii Karpenko, Joerg Aichelin, Pol-Bernard Gossiaux,
Martin Rohrmoser, and Klaus Werner



INT Workshop 21r-2b
Probing QCD at High Energy and Density with Jets
Institute for Nuclear Theory, University of Washington
Seattle, Washington, USA
20 October 2023



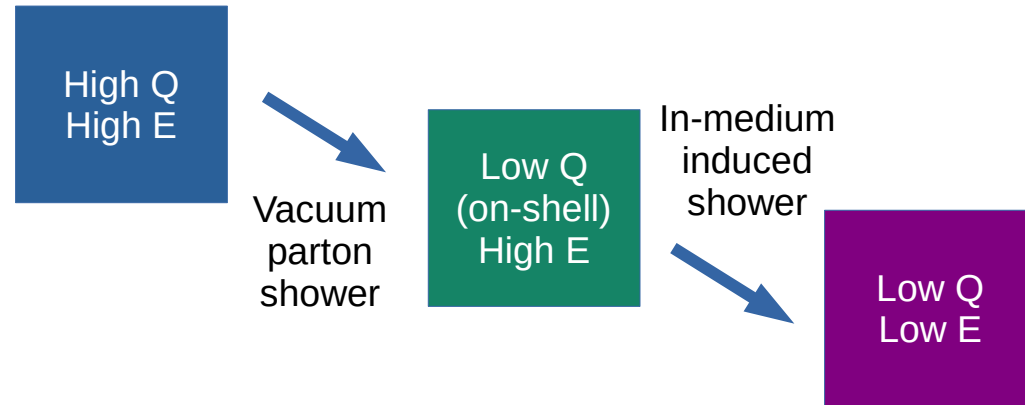
Jets in Heavy Ion Collisions



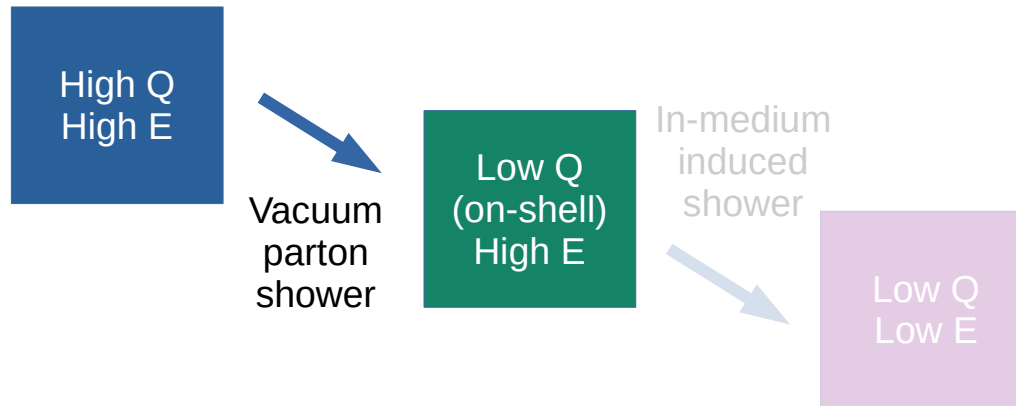
- Interactions between jet partons and the QGP medium leads to modifications of jet properties

→ **Jet Energy Loss / Quenching**

- **SUBA-Jet:**
Monte Carlo for jet energy loss in heavy ion collisions



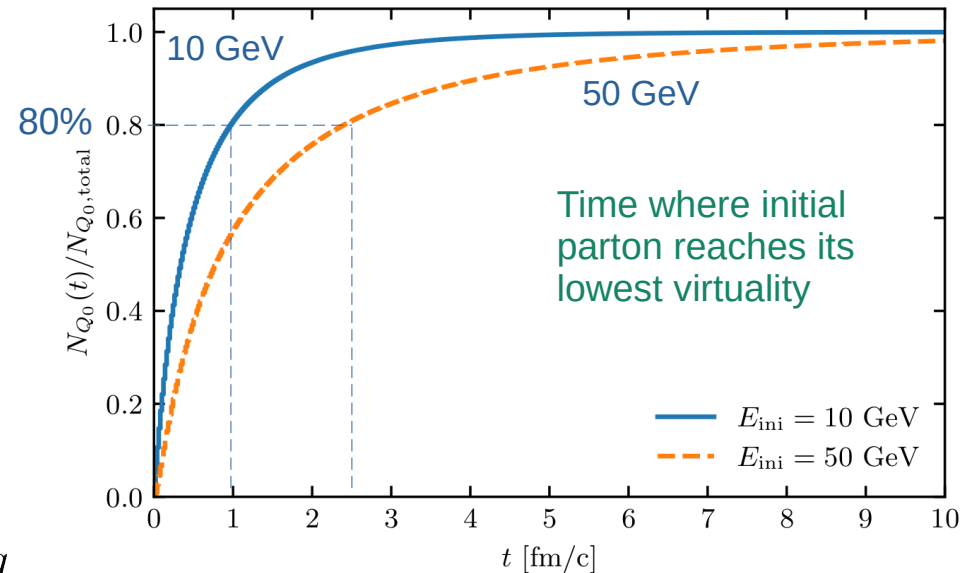
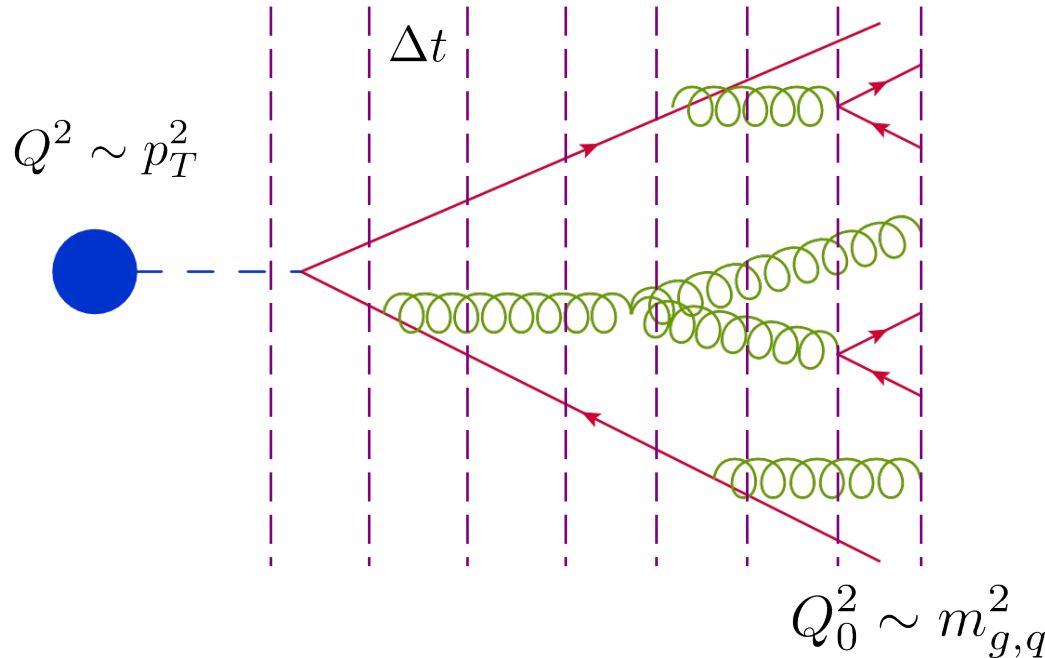
High Virtuality Regime



Vacuum Parton Shower

- Monte Carlo of a vacuum parton shower originally developed by Martin Rohrmoser
- Evolution according to the DGLAP equations from high virtuality $Q_{\max} \sim p_T$ to low virtuality Q_0
- Time evolution split into time steps, mean life time

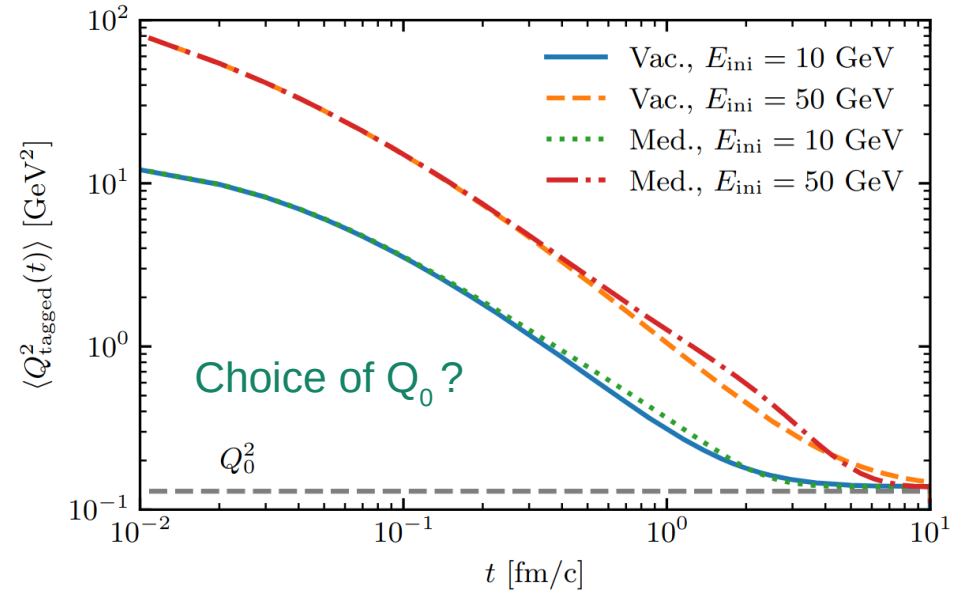
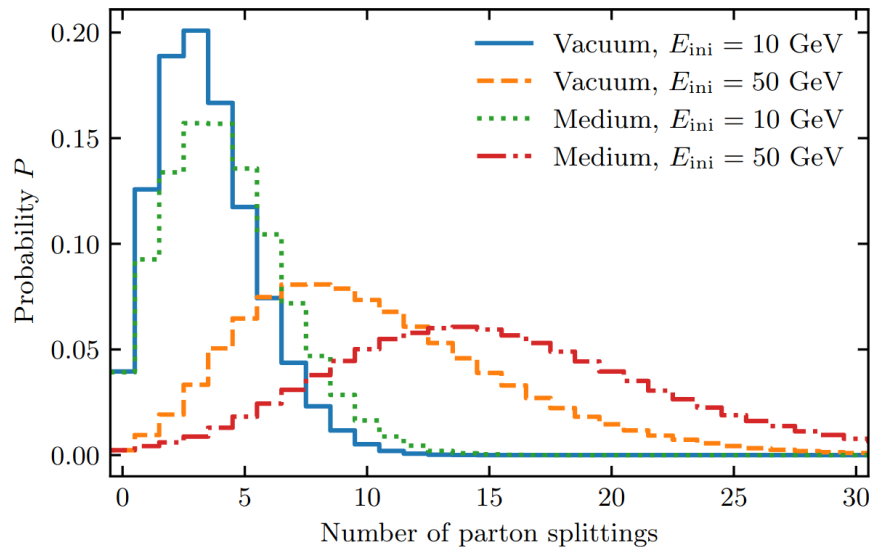
$$\Delta t = \tau = \frac{E}{Q^2}$$



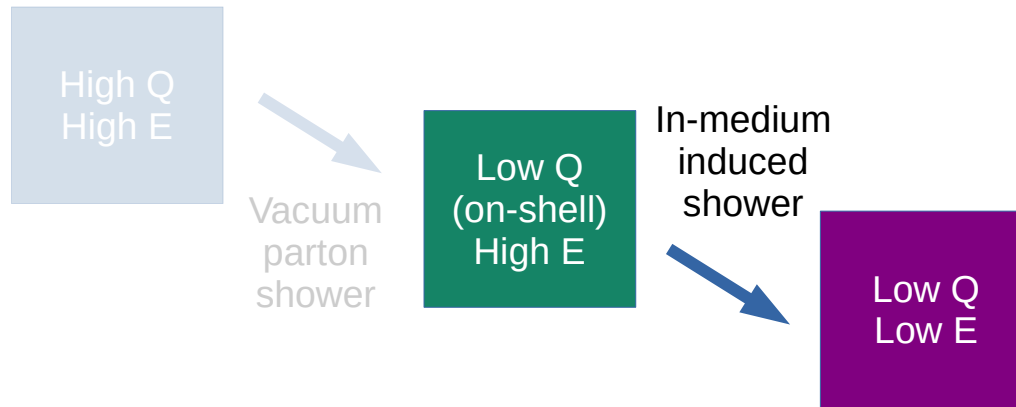
“Vacuum” Parton Shower in Medium

- Medium interactions for high Q regime resulting in virtuality increase, similar to YaJEM (T. Renk, 2008)

$$\frac{dQ^2}{dt} = \hat{q}(T)$$

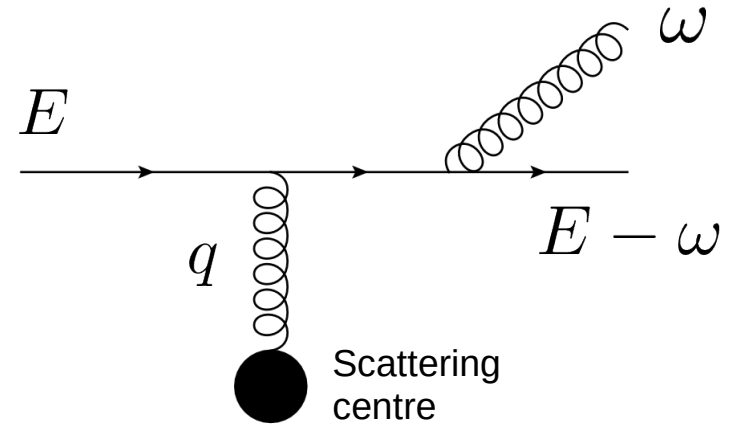


Low Virtuality Regime



Medium-Induced Single Radiation

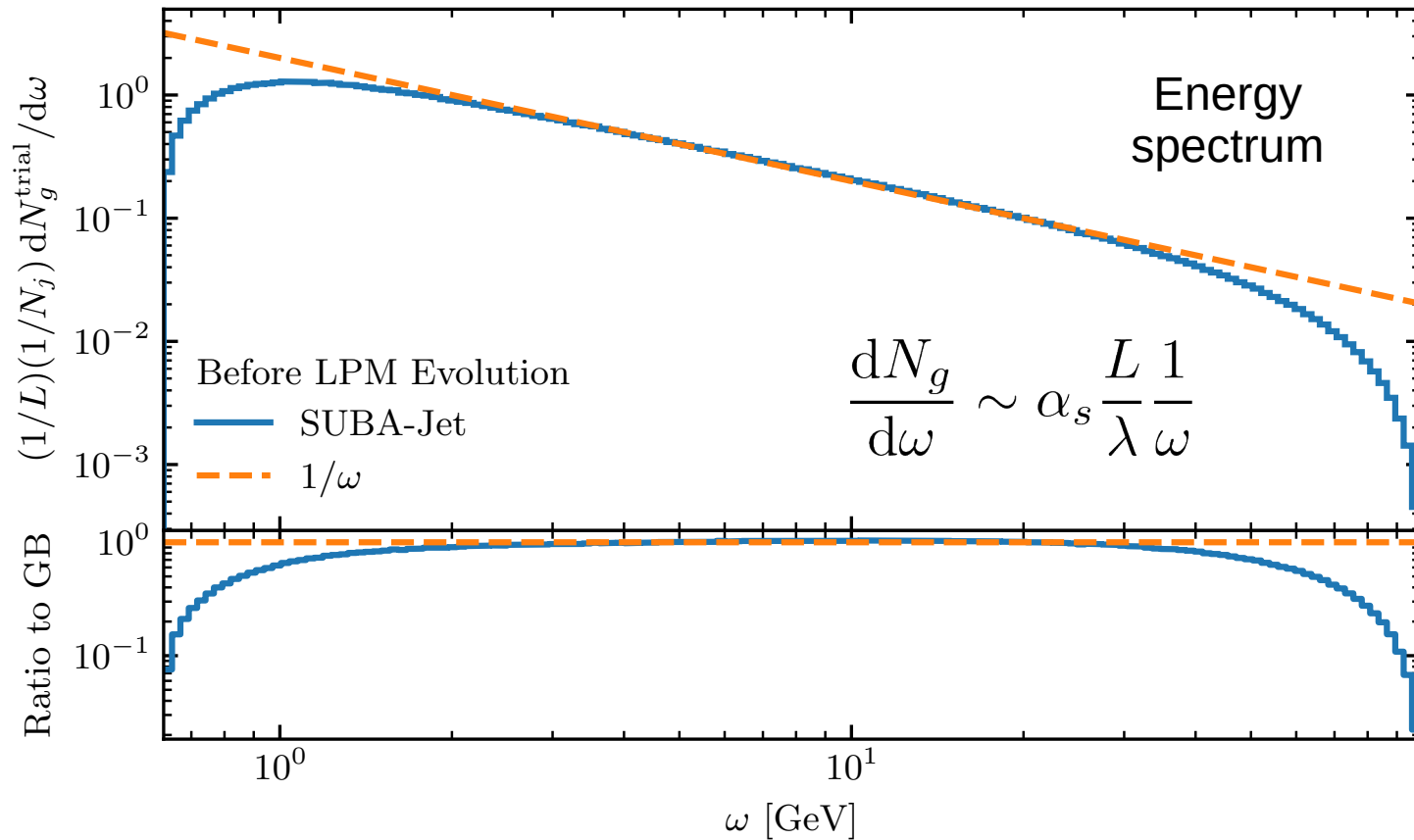
- **Inelastic collision:**
Single gluon emission from single medium scattering
- **Original result from Gunion-Bertsch (1982)**
Generalised to massive case by Aichelin, Gossiaux, Gousset (2014)
- **Initial Gunion-Bertsch seed:** i.e. radiation of a **preformed gluon** from a single scattering (Each parton can generate a number of preformed gluons)
- Gunion-Bertsch cross-section from scalar QCD



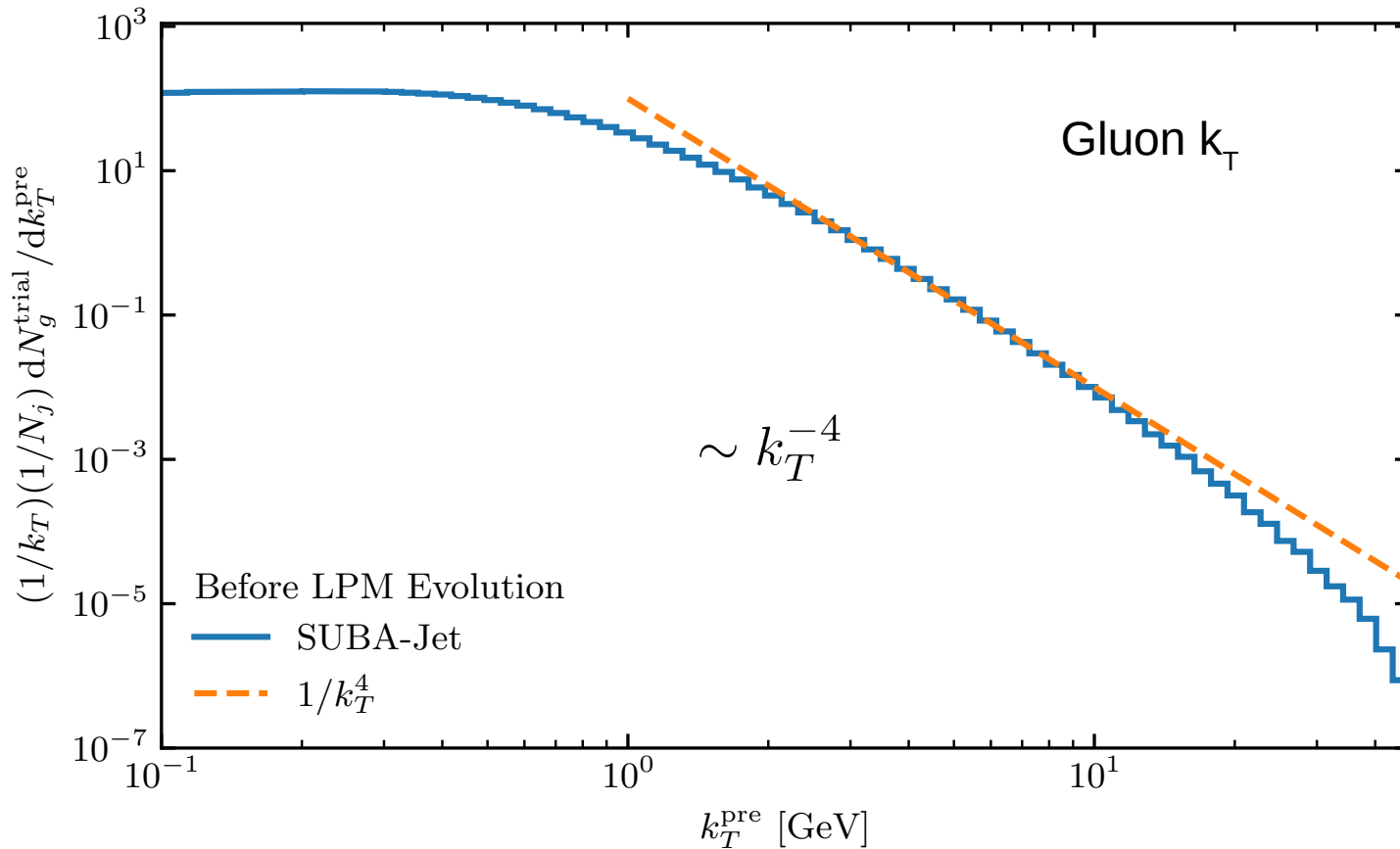
$$\frac{d\sigma^{Qq \rightarrow Qqg}}{dx d^2k_T d^2l_t} = \frac{d\sigma_{\text{el}}}{d^2l_t} P_g(x, k_T, l_T) \theta(\Delta)$$

$$\frac{d\sigma_{\text{el}}}{d^2l_t} \sim \frac{8\alpha_s^2}{9(l_T^2 + \mu^2)^2}$$

Medium-Induced Single Radiation



Medium-Induced Single Radiation



Coherency and the LPM Effect

- The formation of the radiated gluon is a quantum mechanical process

Formation time: $t_f \sim \sqrt{\frac{\omega}{\hat{q}}}$

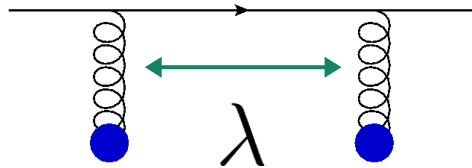
- Coherence effects:
Landau-Pomeranchuk-Migdal (LPM) effect

- Have to take into account multiple scatterings with the medium during the formation time

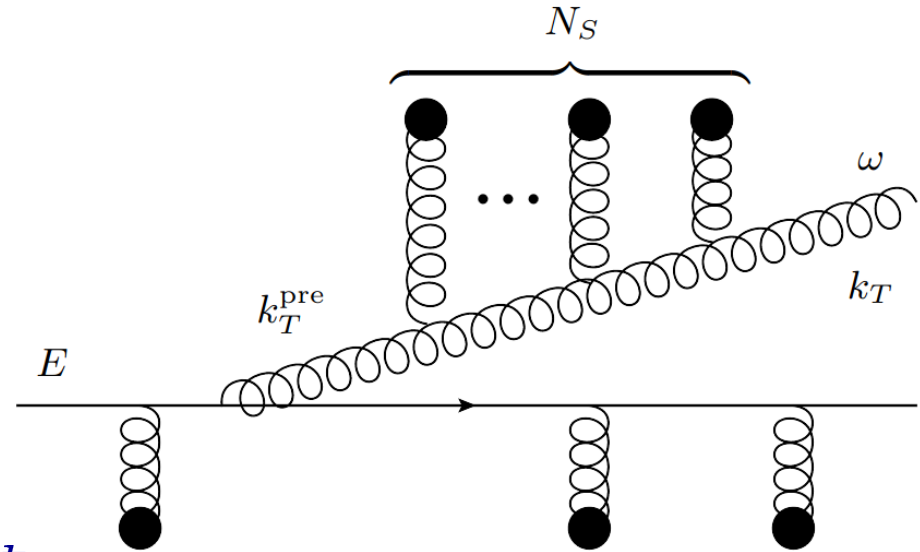
ω = gluon energy

\hat{q} = medium modifications

$$N_s = \frac{t_f}{\lambda}$$



$$\lambda \simeq \frac{\hbar c}{\alpha_s T}$$



L = path length of medium

Implementation of the LPM Effect

- At each timestep:

- Elastic scattering with prob. $\Gamma_{\text{el}}\Delta t$

$$\Gamma_{\text{el}}^q = \left(1 + \frac{N_f}{N}\right) \frac{(N^2 - 1)T^3}{\pi\hbar c} \frac{4\alpha_s^2}{\mu^2}$$

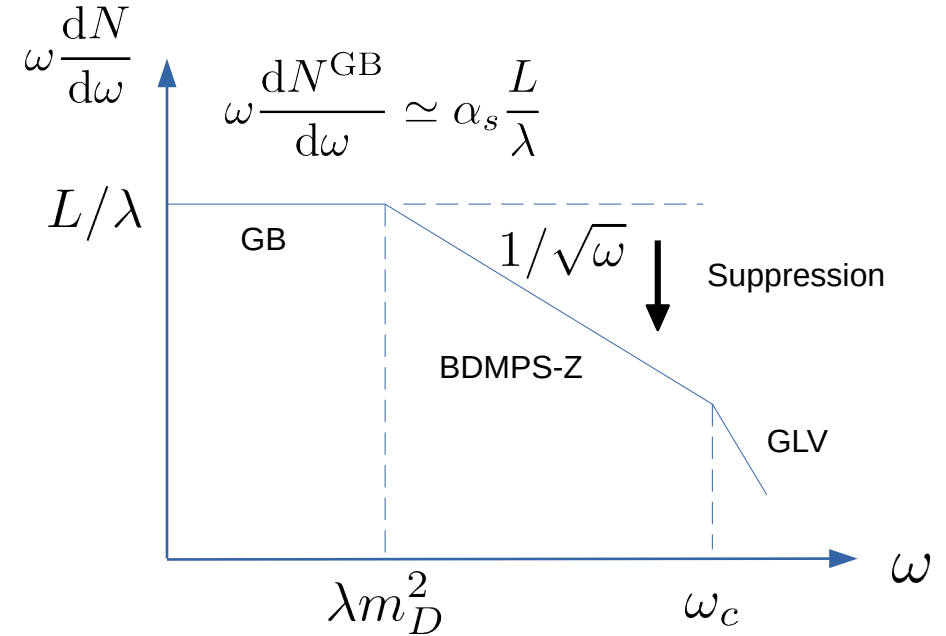
- Radiation of preformed gluon with prob. $\Gamma_{\text{inel}}\Delta t$

- BDMPS-Z spectrum at intermediate energies achieved by suppressing GB seed by

$$1/N_s$$

Like in Zapp, Stachel, Wiedemann, JHEP 07 (2011), 118

Radiation energy spectrum:



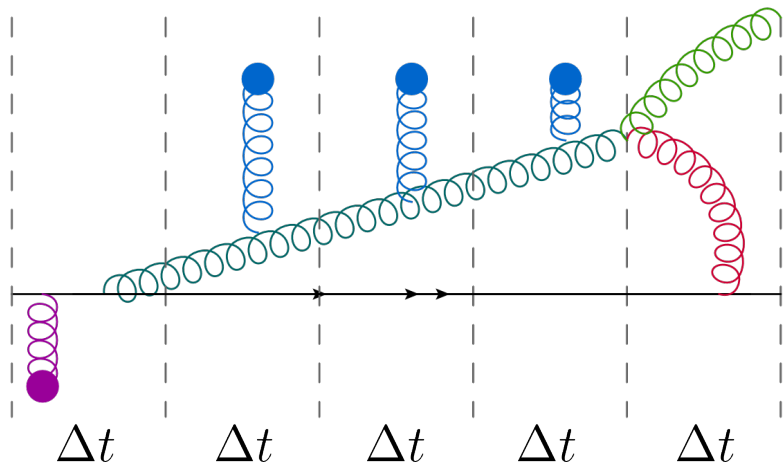
$$\omega \frac{dN^{\text{BDMPS-Z}}}{d\omega} \simeq \alpha_s \sqrt{\frac{\hat{q} L^2}{\omega}}$$

The Algorithm

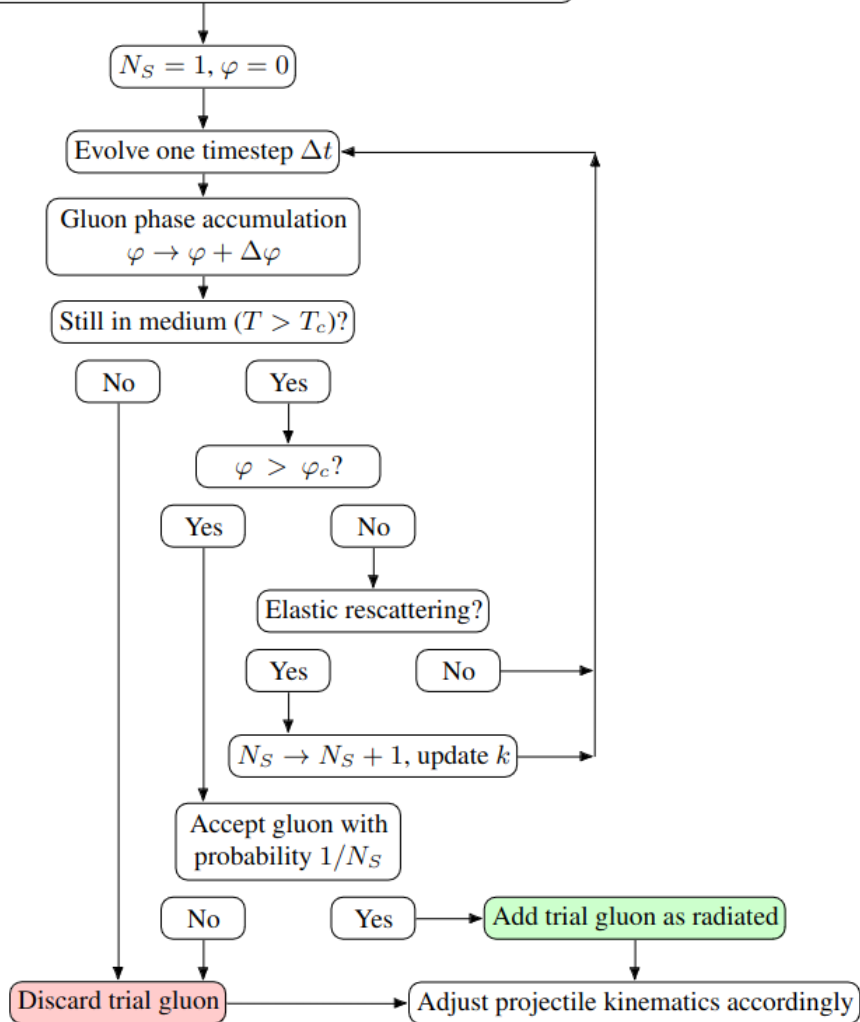
Flow diagram:

Algorithm for the coherent medium-induced gluon radiation in our model

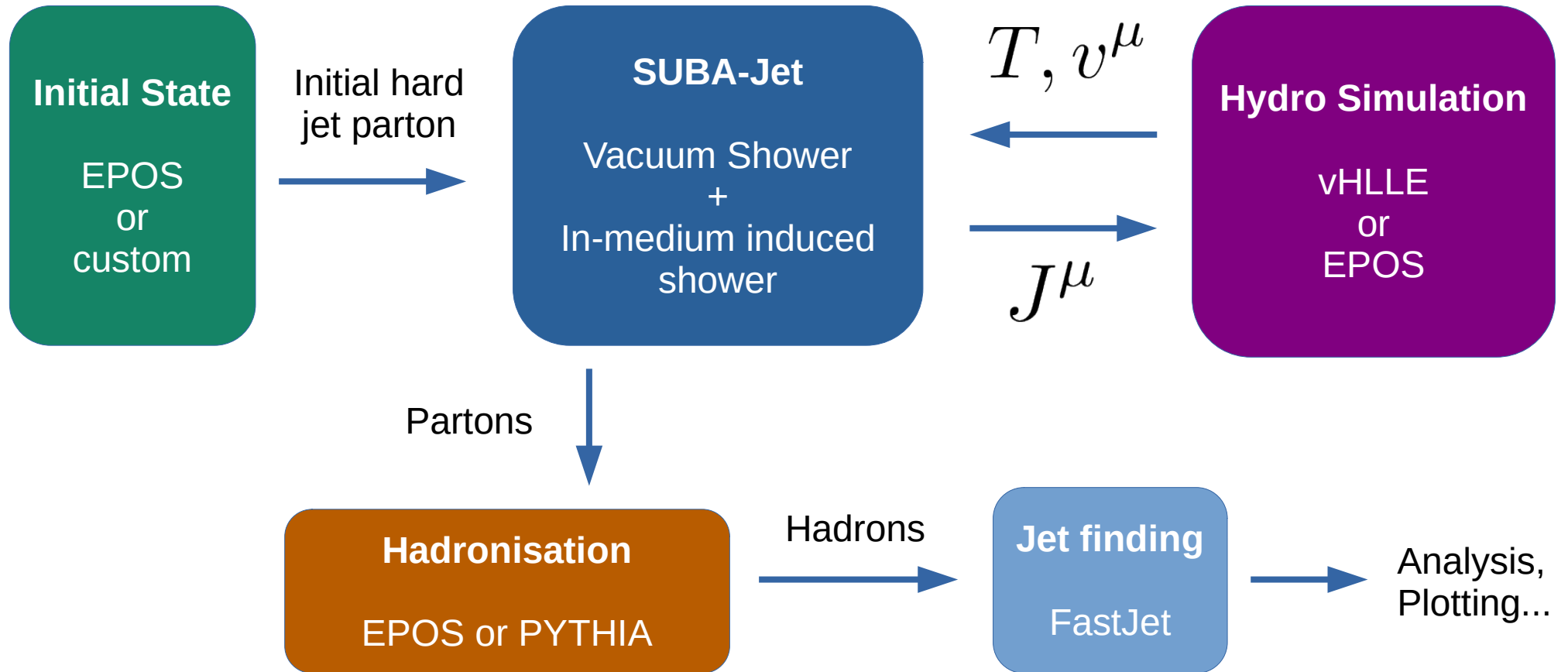
Various parameters and settings can be changed and tuned to compare distributions



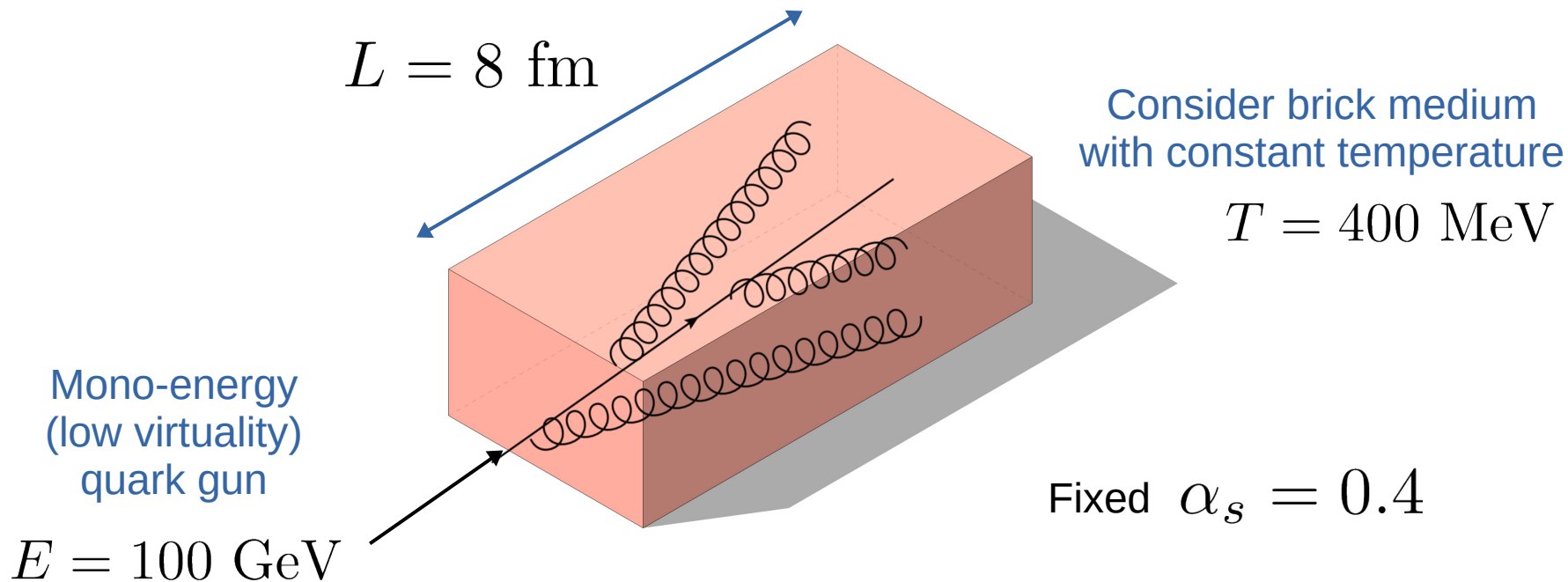
Trial (preformed) incoherent gluon formation according to GB seed



The Monte Carlo

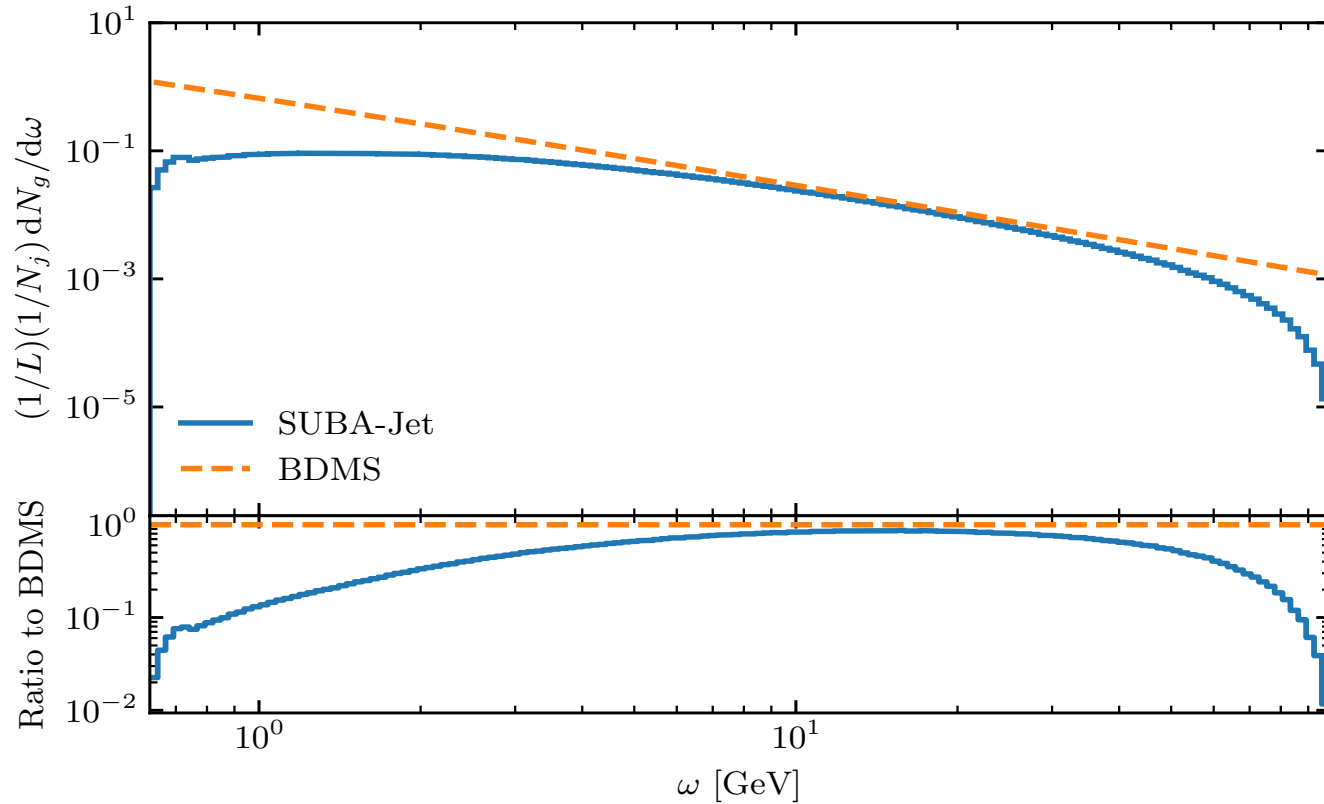


First Results



**We want to reproduce theoretical expectation
and check effect of model parameters**

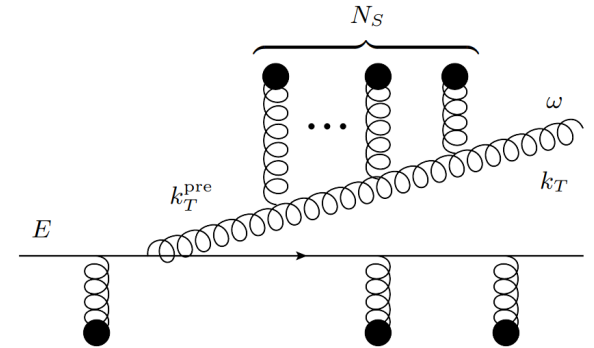
Reproduction of the BDMPS-Z Limit



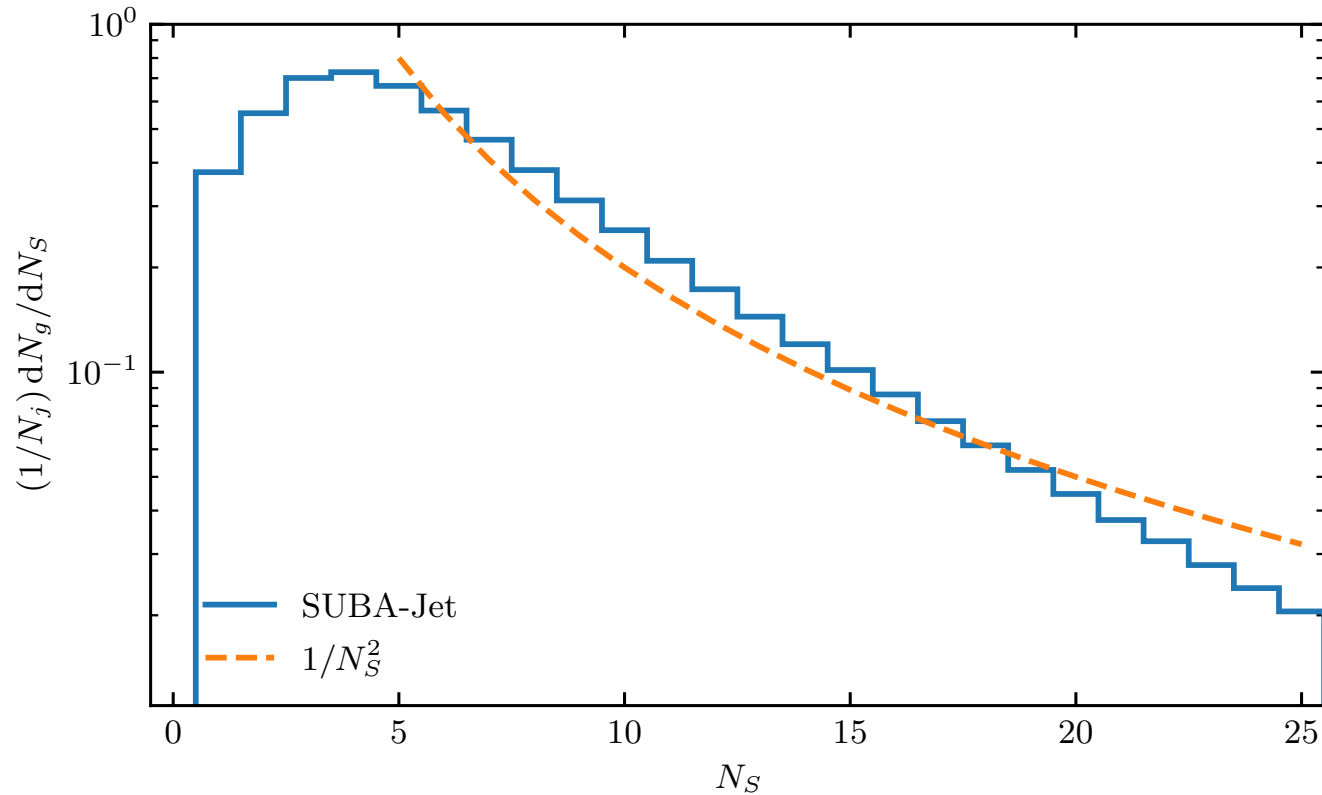
Glucn energy ω spectrum

Reproduces BDMPS-Z for intermediate energies

$$\frac{dN}{d\omega} \sim \frac{1}{\omega^{3/2}}$$



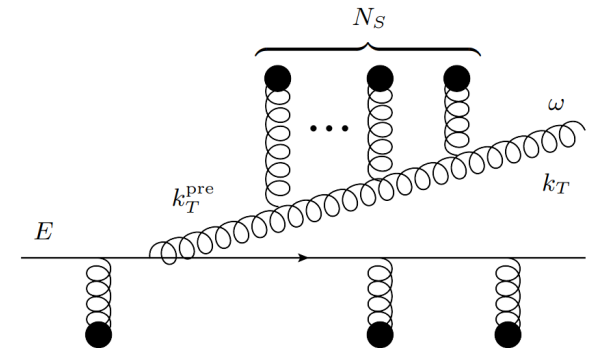
Reproduction of the BDMPS-Z Limit



Number of elastic scatterings N_S

Reproduces BDMPS-Z expectation

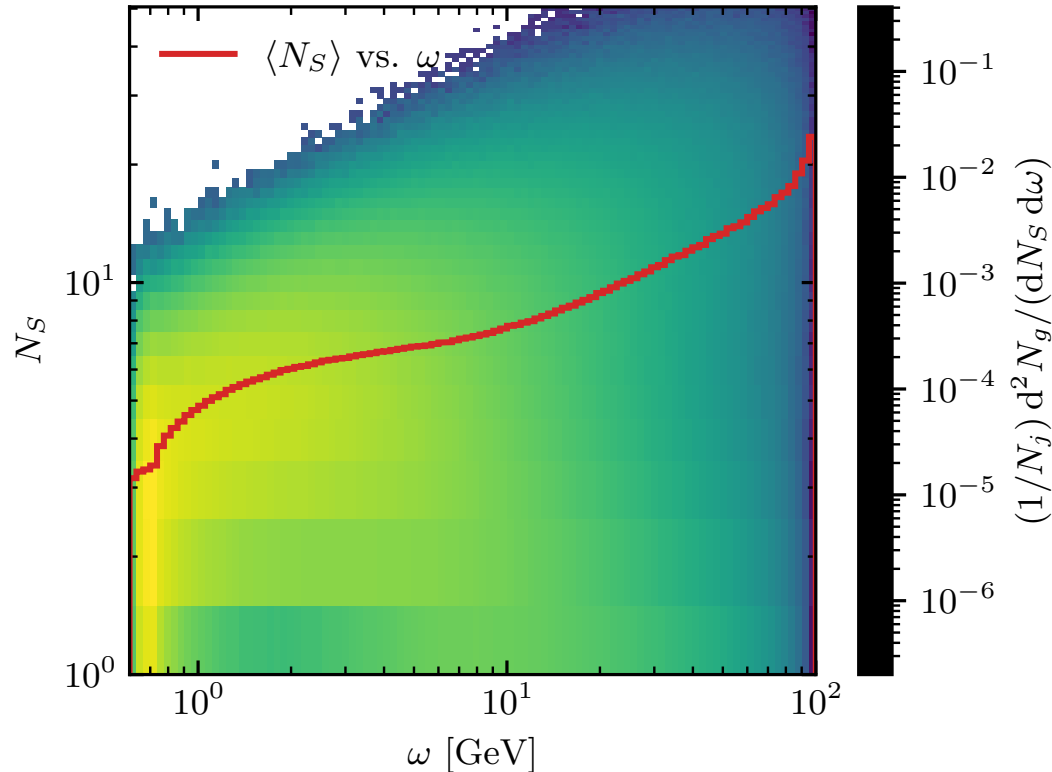
$$\frac{dN}{dN_S} \sim \frac{1}{N_S^2}$$



Reproduction of the 3 Regimes

$E = 100 \text{ GeV}$, $L = 8 \text{ fm}$, $m_q \rightarrow \infty$

$k_T^{\text{pre}} < c_{\text{colin}} \omega$, $\Delta\phi = (2P_Q \cdot k/E_Q)\Delta t$, Energy conservation



Double differential plot
in N_S and ω

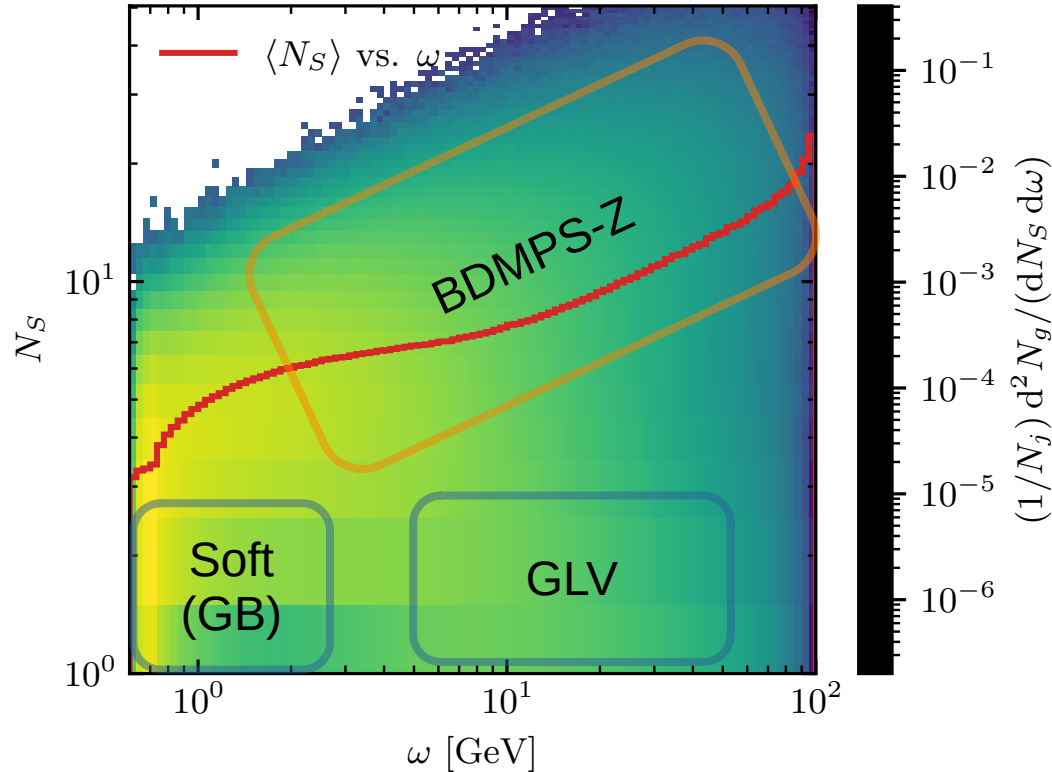
Red line: $\langle N_S \rangle$ vs. ω

$$N_S \sim t_f \sim \sqrt{\omega}$$

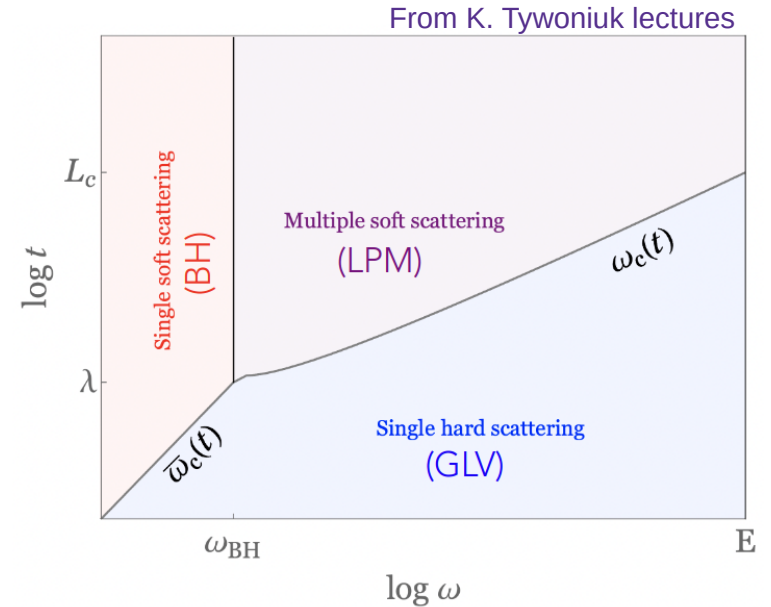
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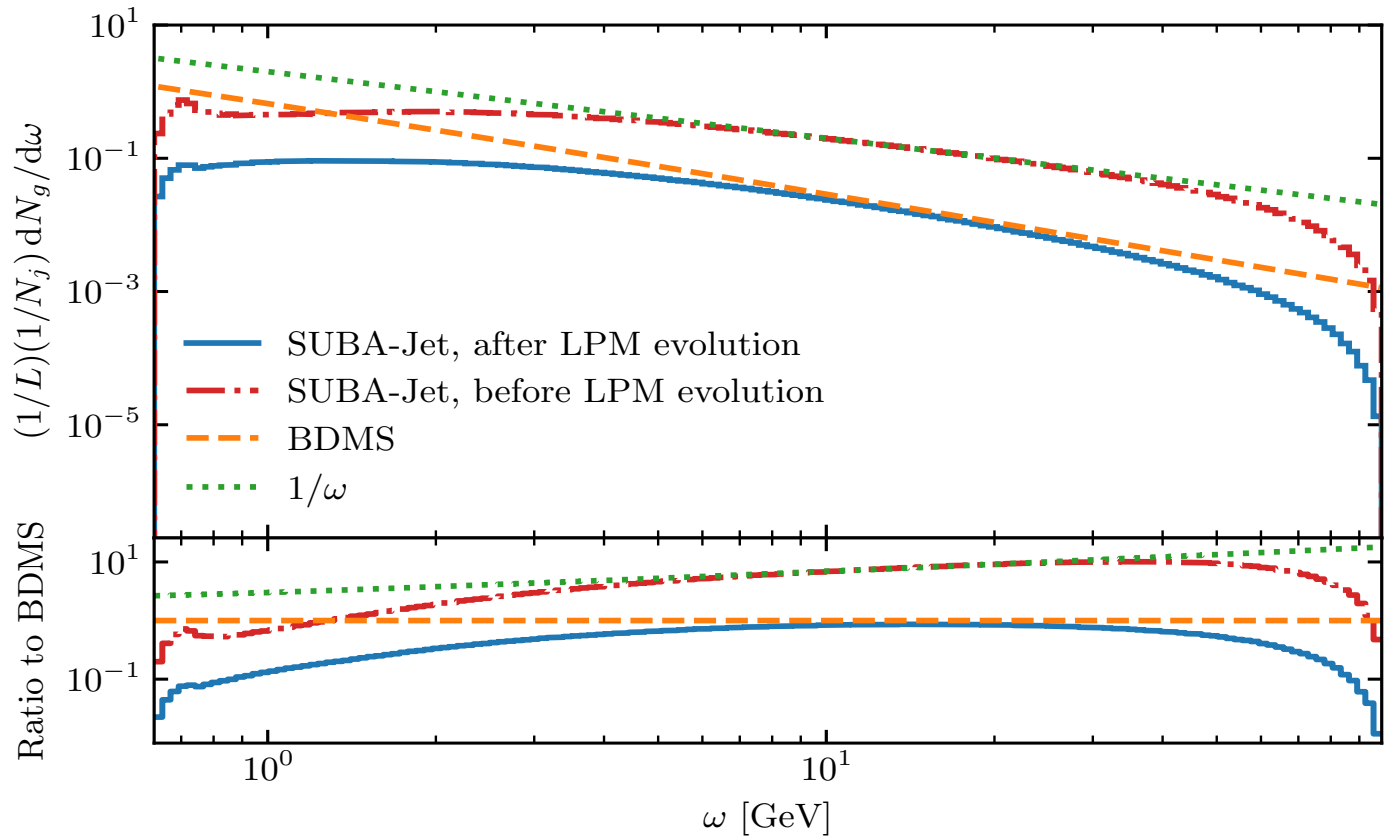
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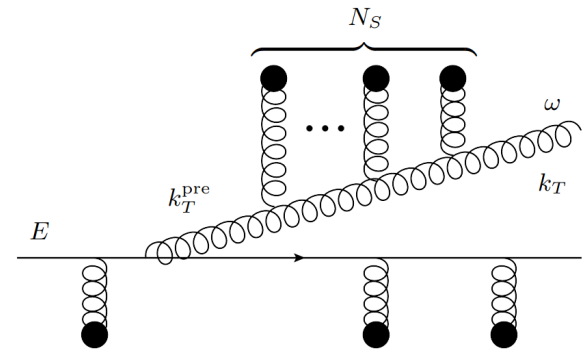
Convolution of different distributions



The Role of the Phase Accumulation

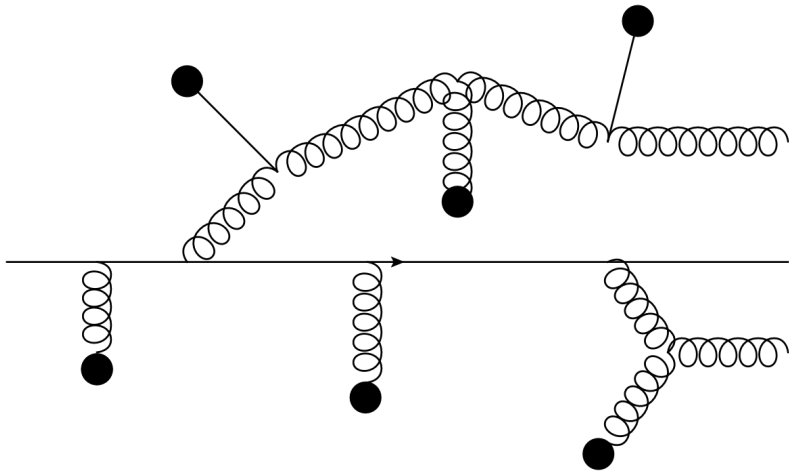


Effect of the phase accumulation on the energy spectrum



The Role of the Phase Accumulation

Choice of phase accumulation
of the preformed (trial) gluons:



- **More general formula:**

$$\Delta\varphi = \frac{2P_Q \cdot k}{E_Q} \Delta t$$

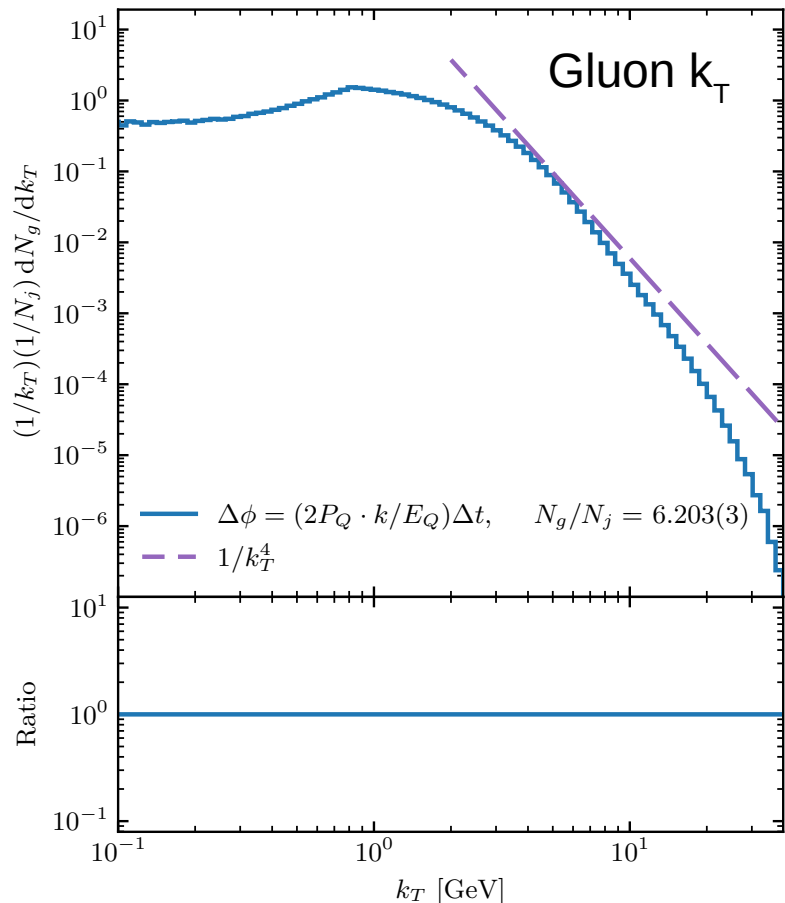
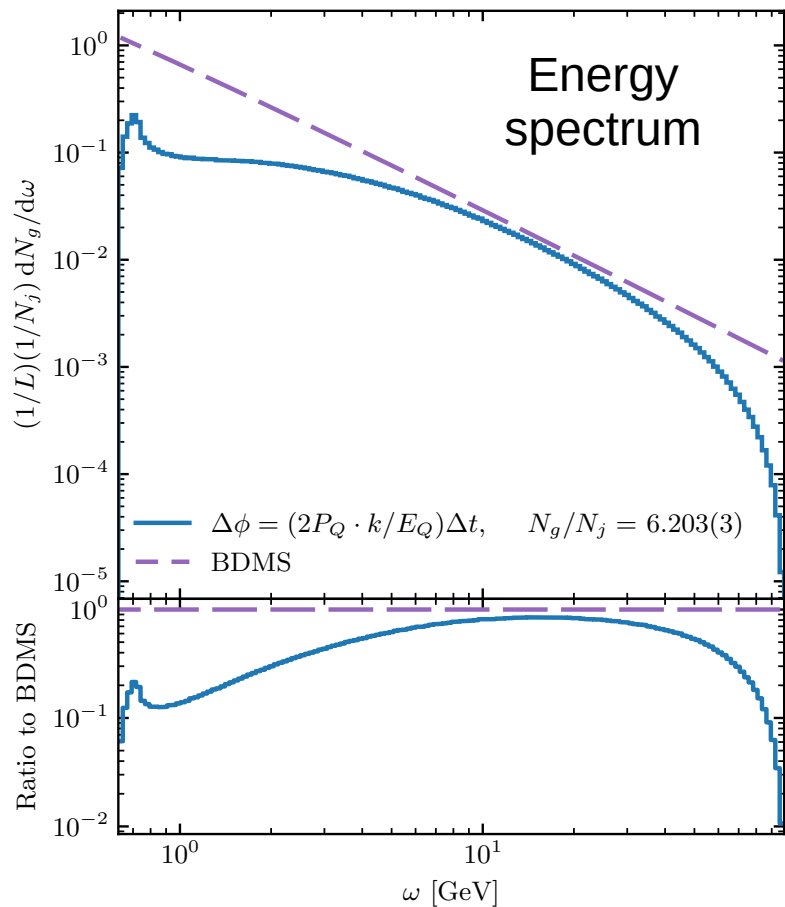
- **What is used in JEWEL:**

$$\Delta\varphi = \frac{k_T^2}{\omega} \Delta t$$

- **Including thermal gluon mass:**

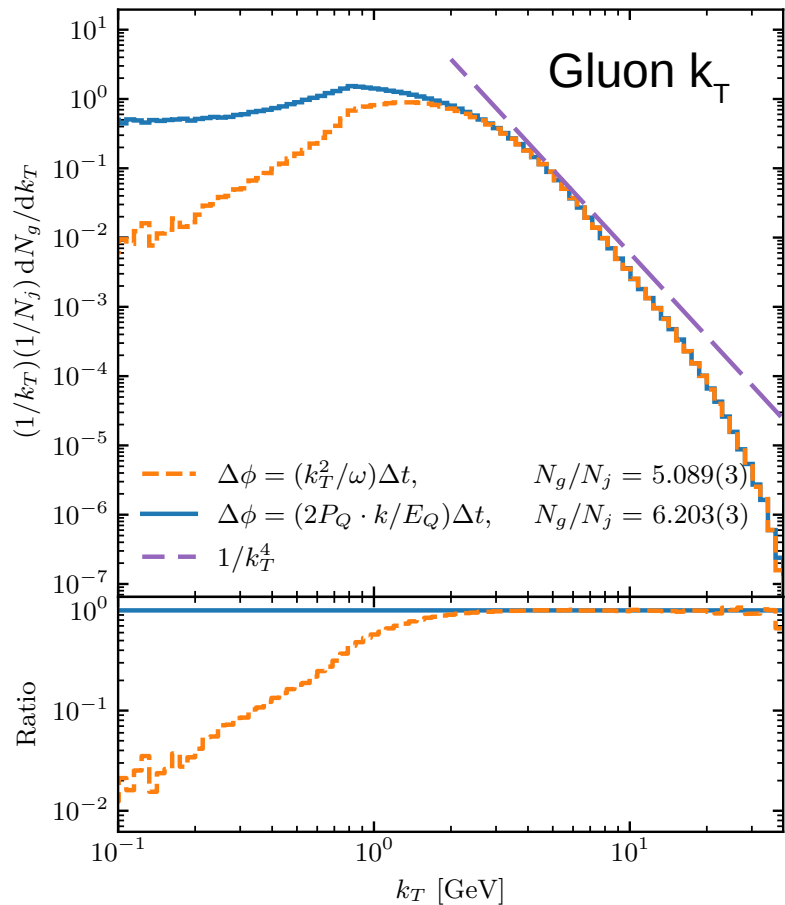
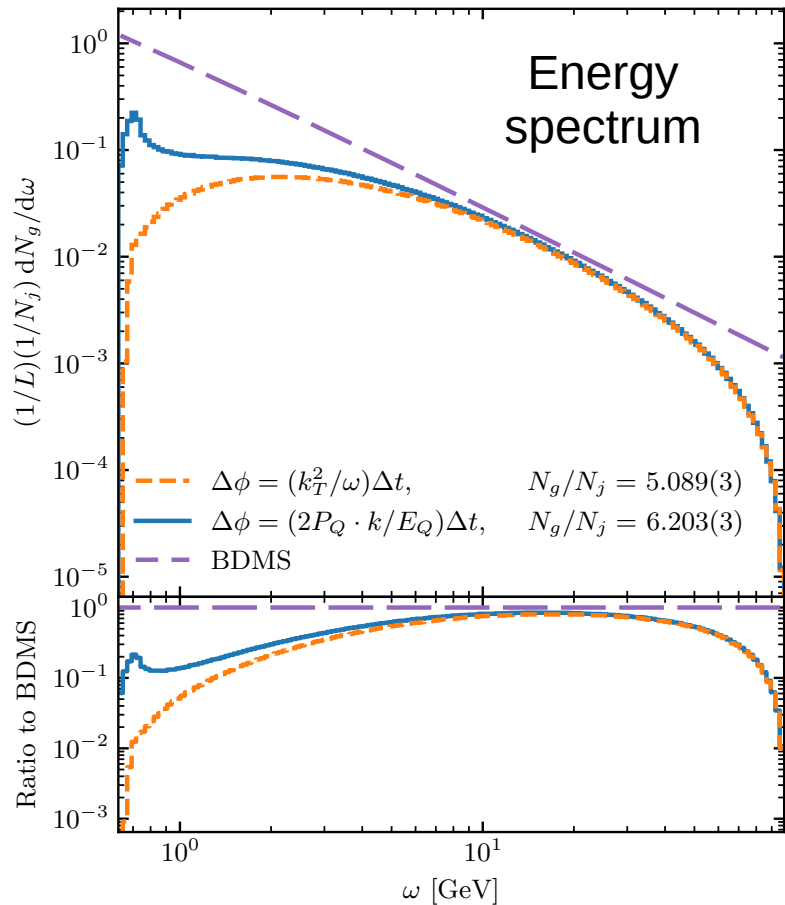
$$\Delta\varphi = \frac{m_g^2 + k_T^2}{\omega} \Delta t$$

The Role of the Phase Accumulation



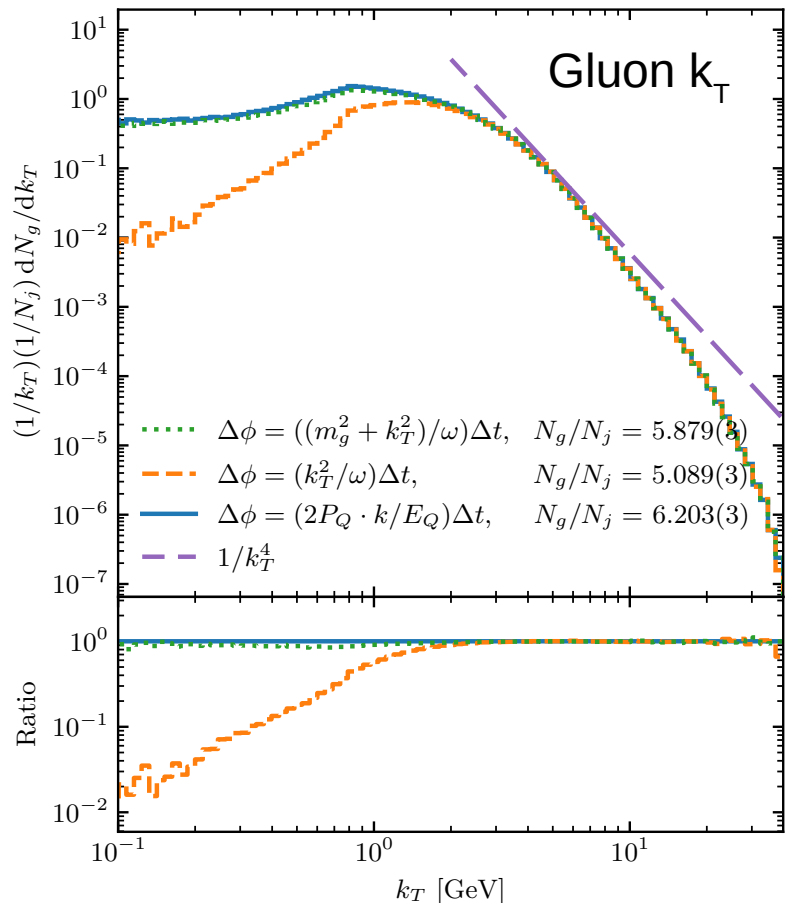
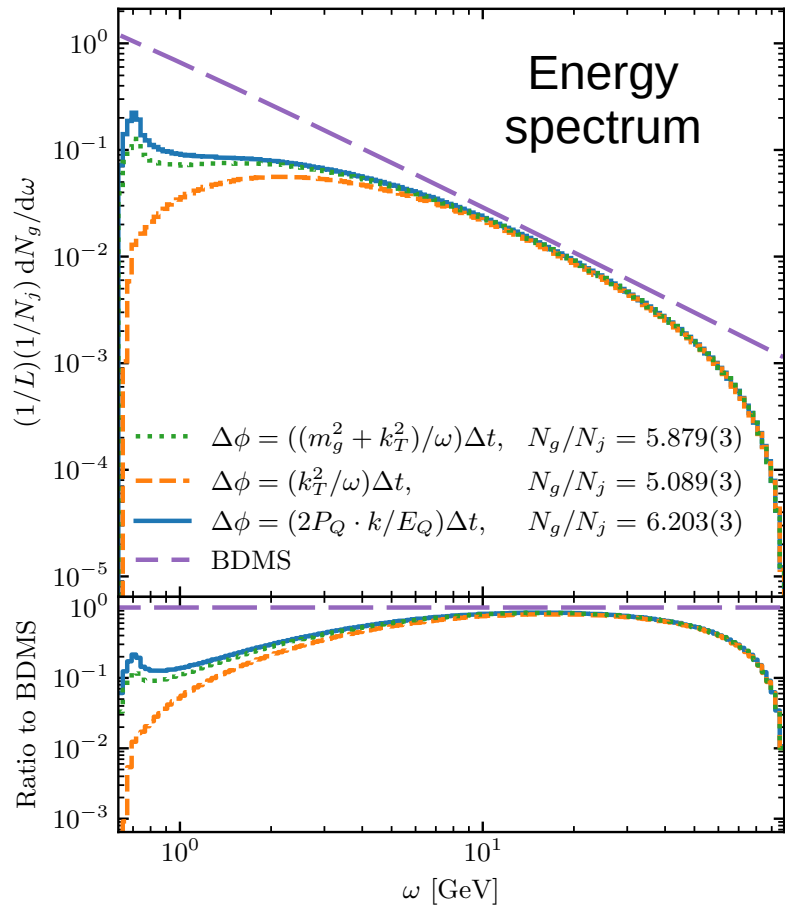
The Role of the Phase Accumulation

Effects at low energy & low k_T



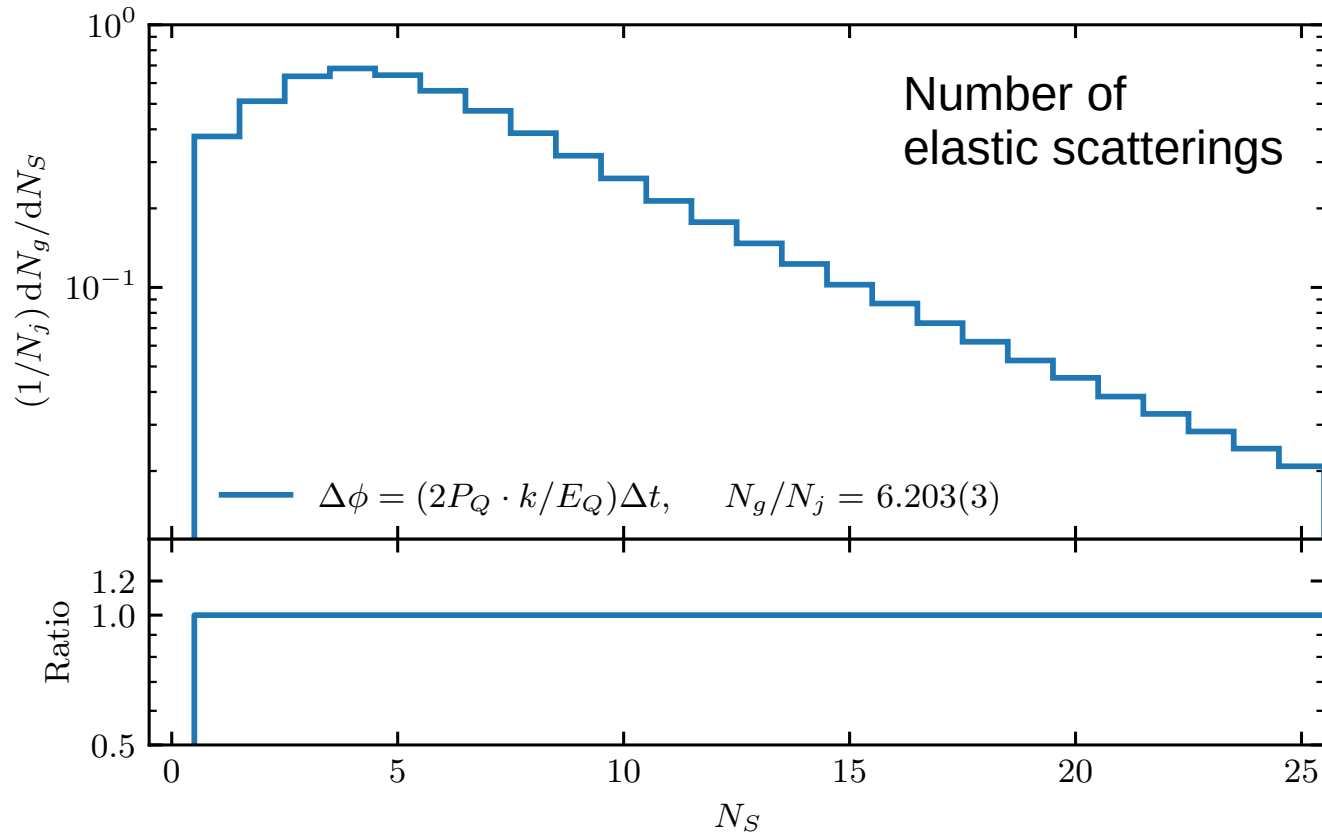
The Role of the Phase Accumulation

Effects at low energy & low k_T

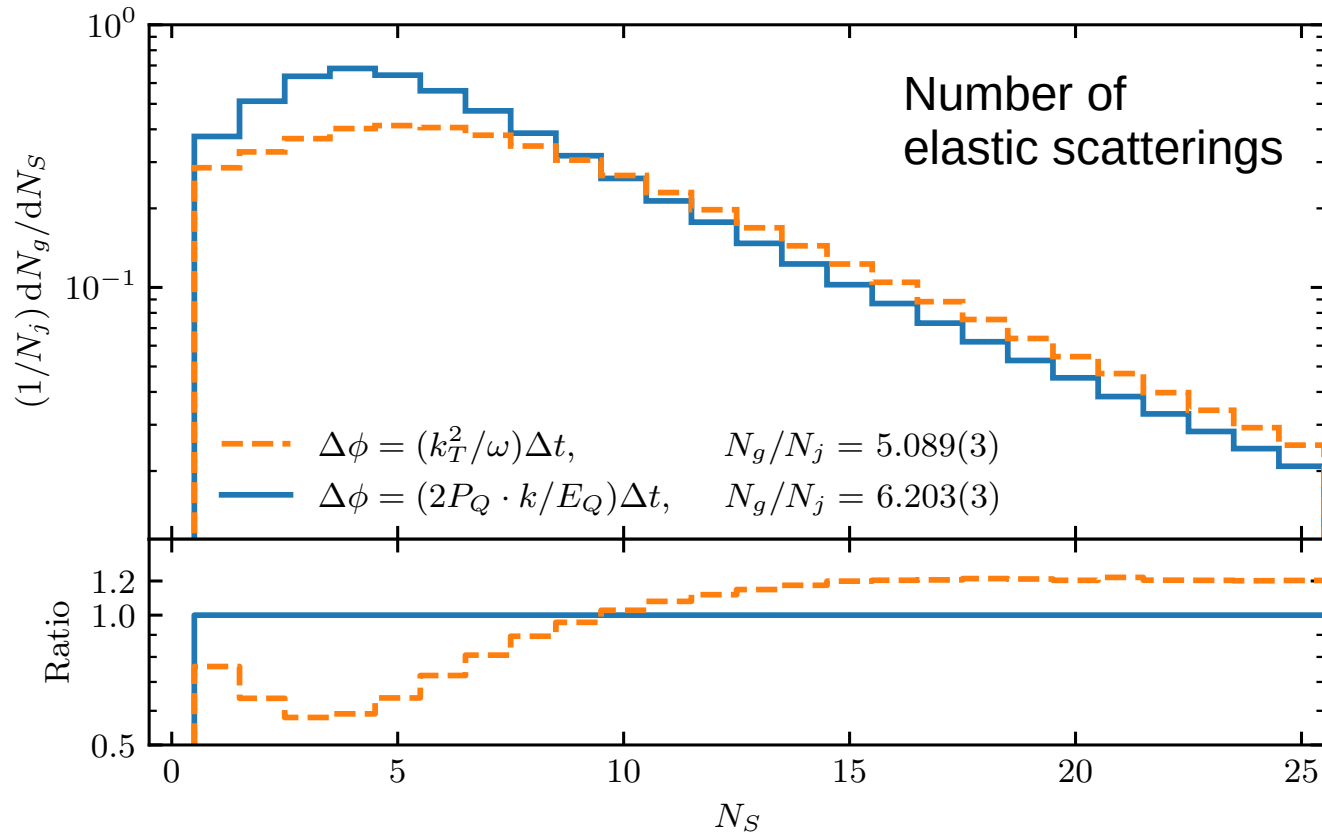


Number of radiated gluons per jet $\sim 5 - 6$

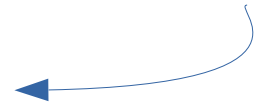
The Role of the Phase Accumulation



The Role of the Phase Accumulation

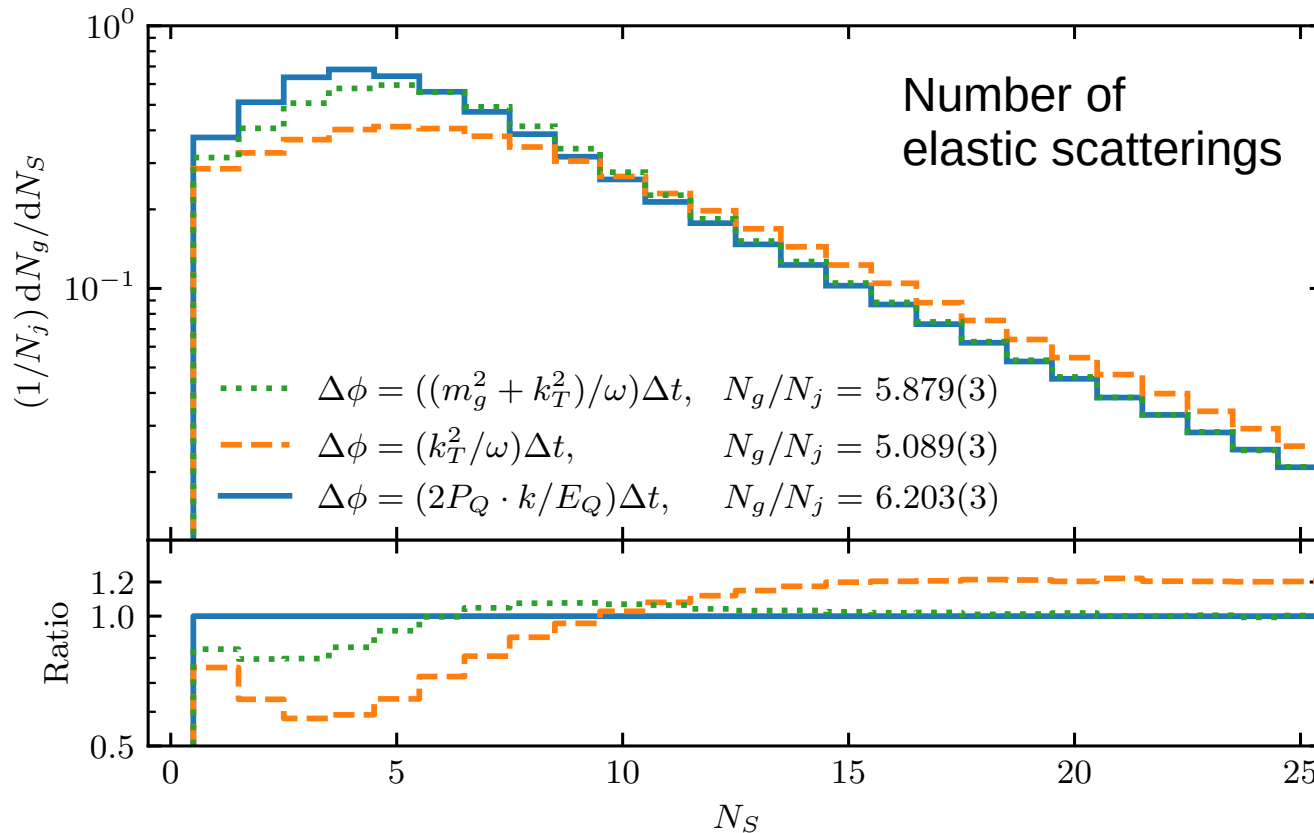
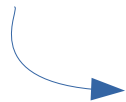


When neglecting the gluon mass in the phase accumulation, a larger path length is required to have a comparable overall number of radiations

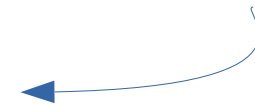


The Role of the Phase Accumulation

Effects at low N_S

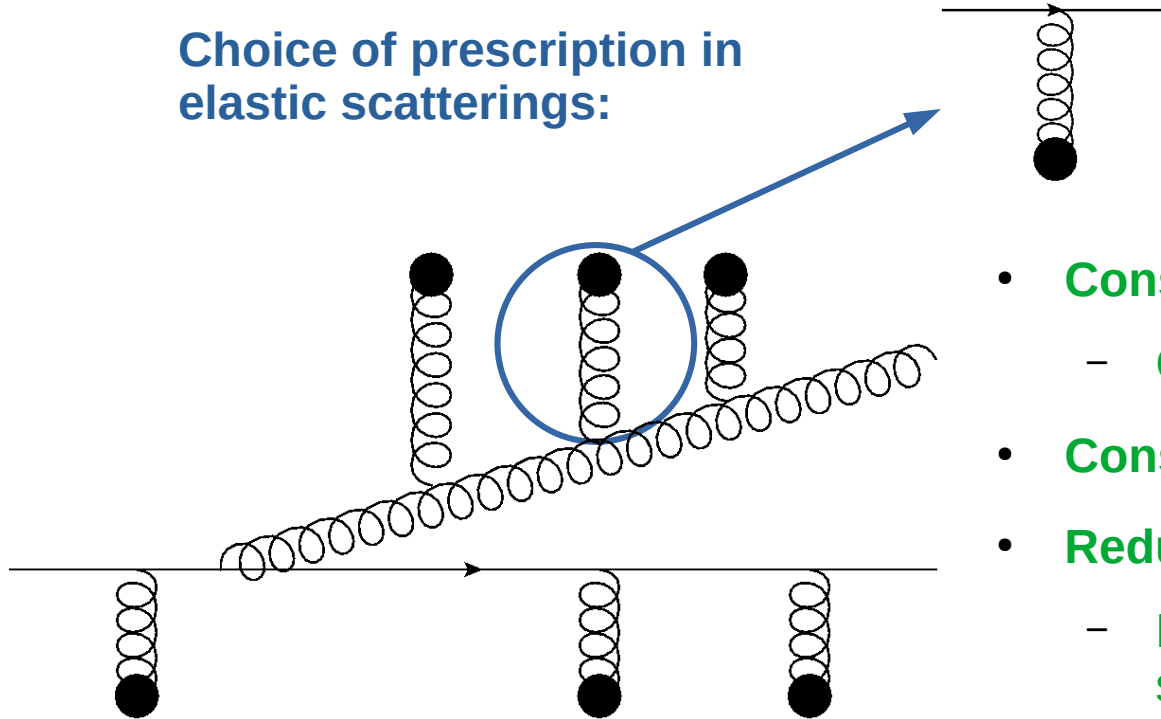


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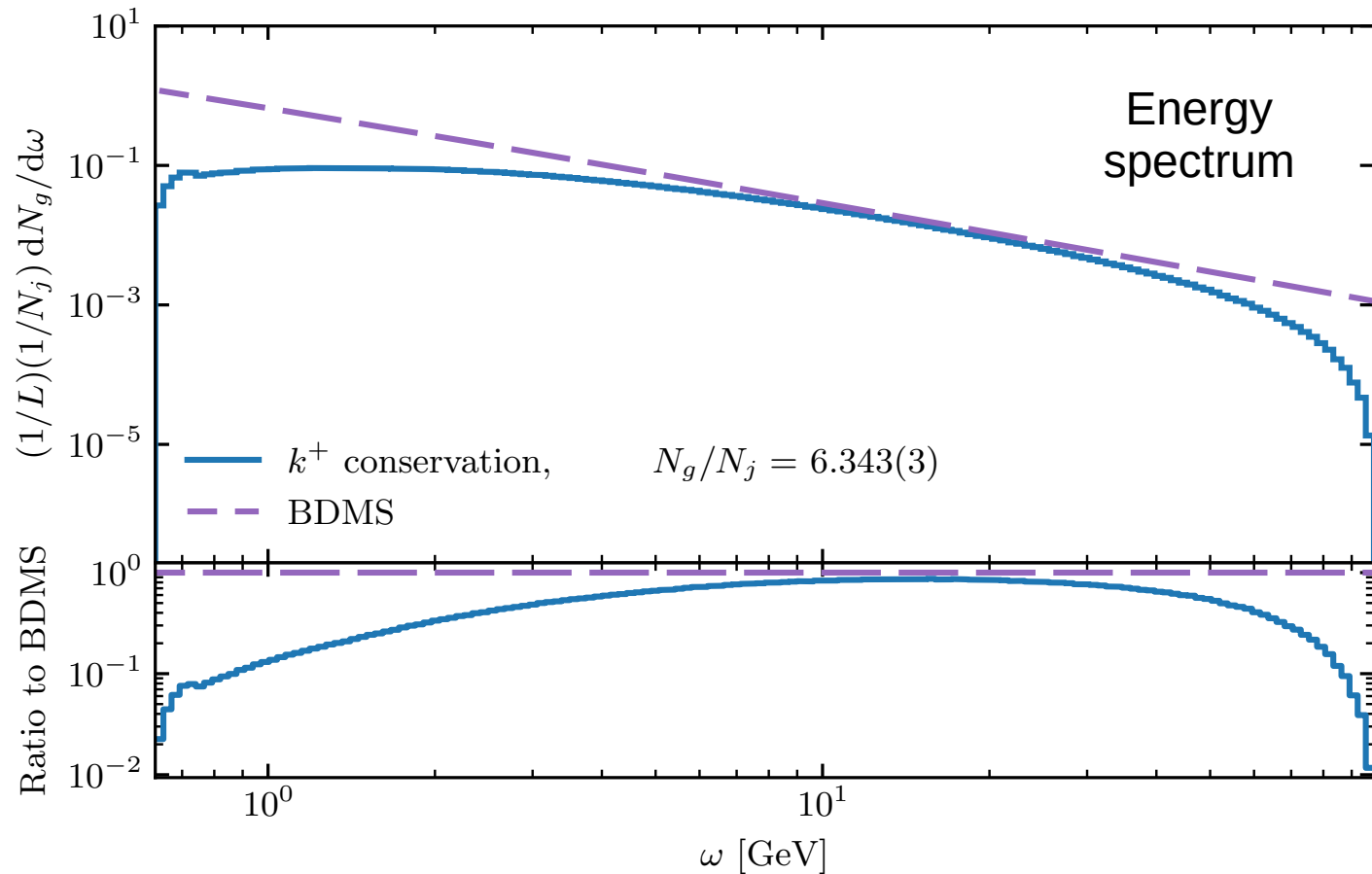
The Role of the Elastic Scatterings

Choice of prescription in elastic scatterings:



- **Conserve k^+ ?**
 - Considered by BDMPS-Z
- **Conserve energy?**
- **Reduce energy?**
 - Energy gain by the medium parton is subtracted from the projectile parton

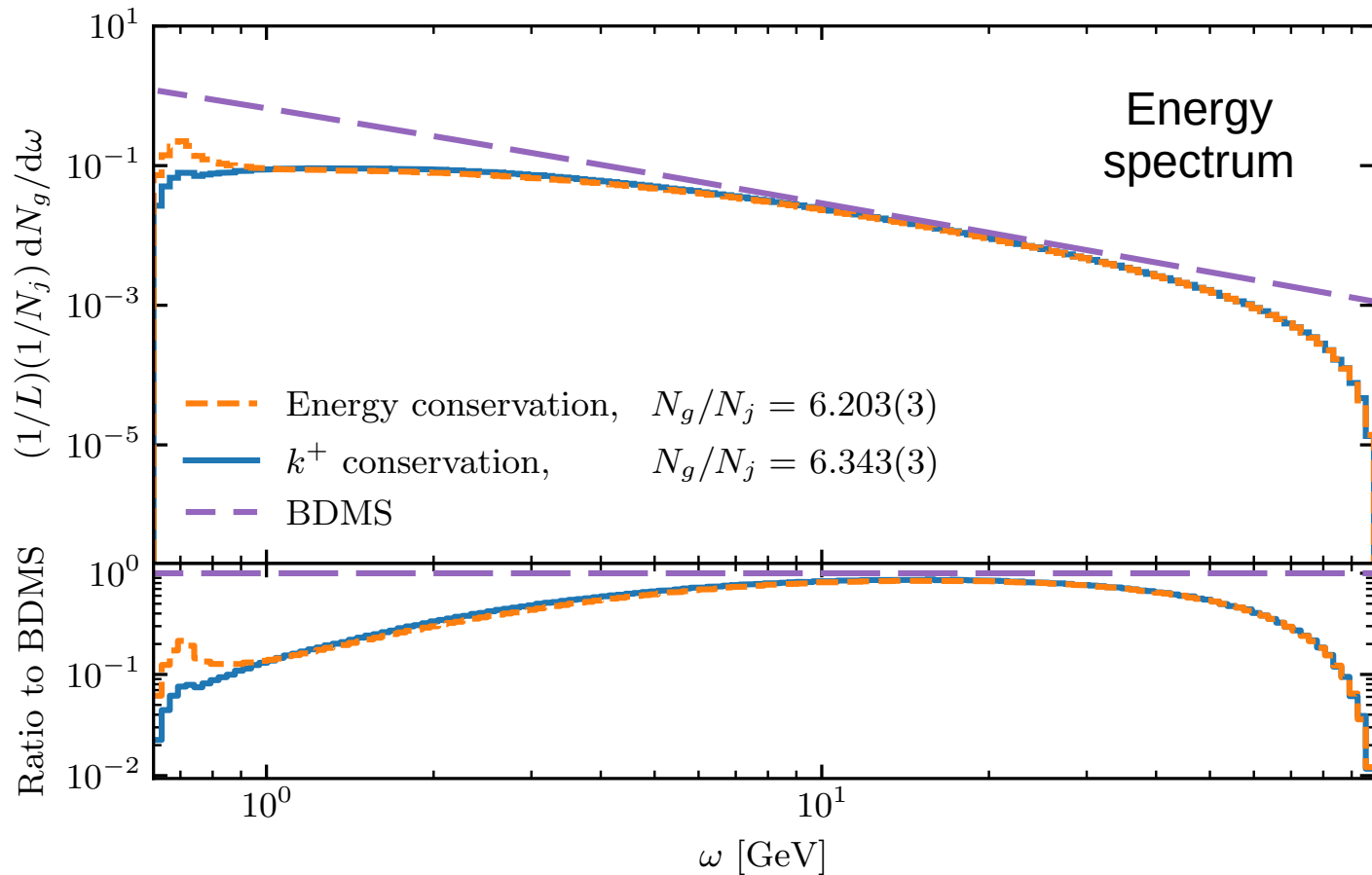
The Role of the Elastic Scatterings



Same BDMS behaviour
at intermediate energies

Difference at small
energies

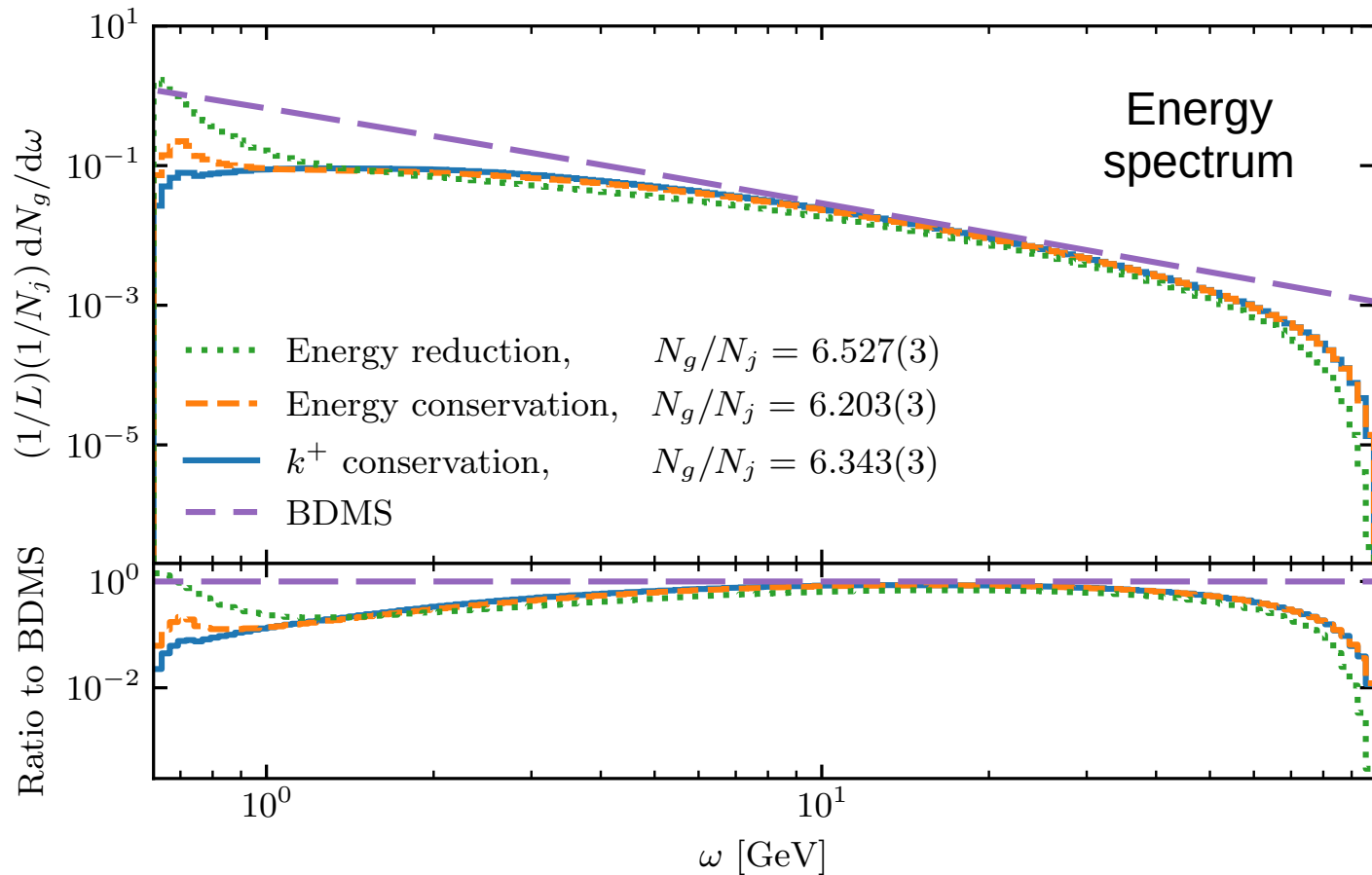
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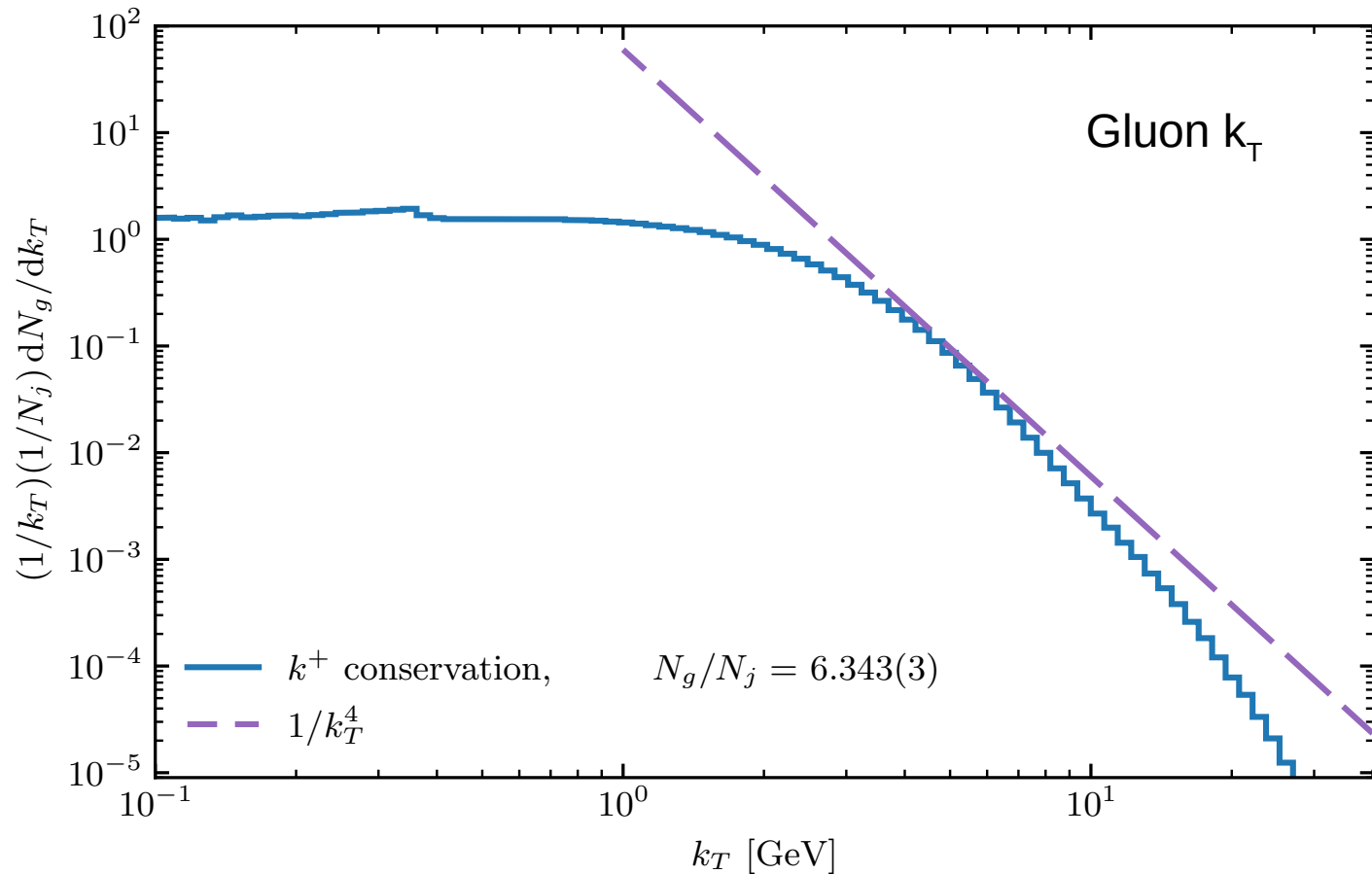
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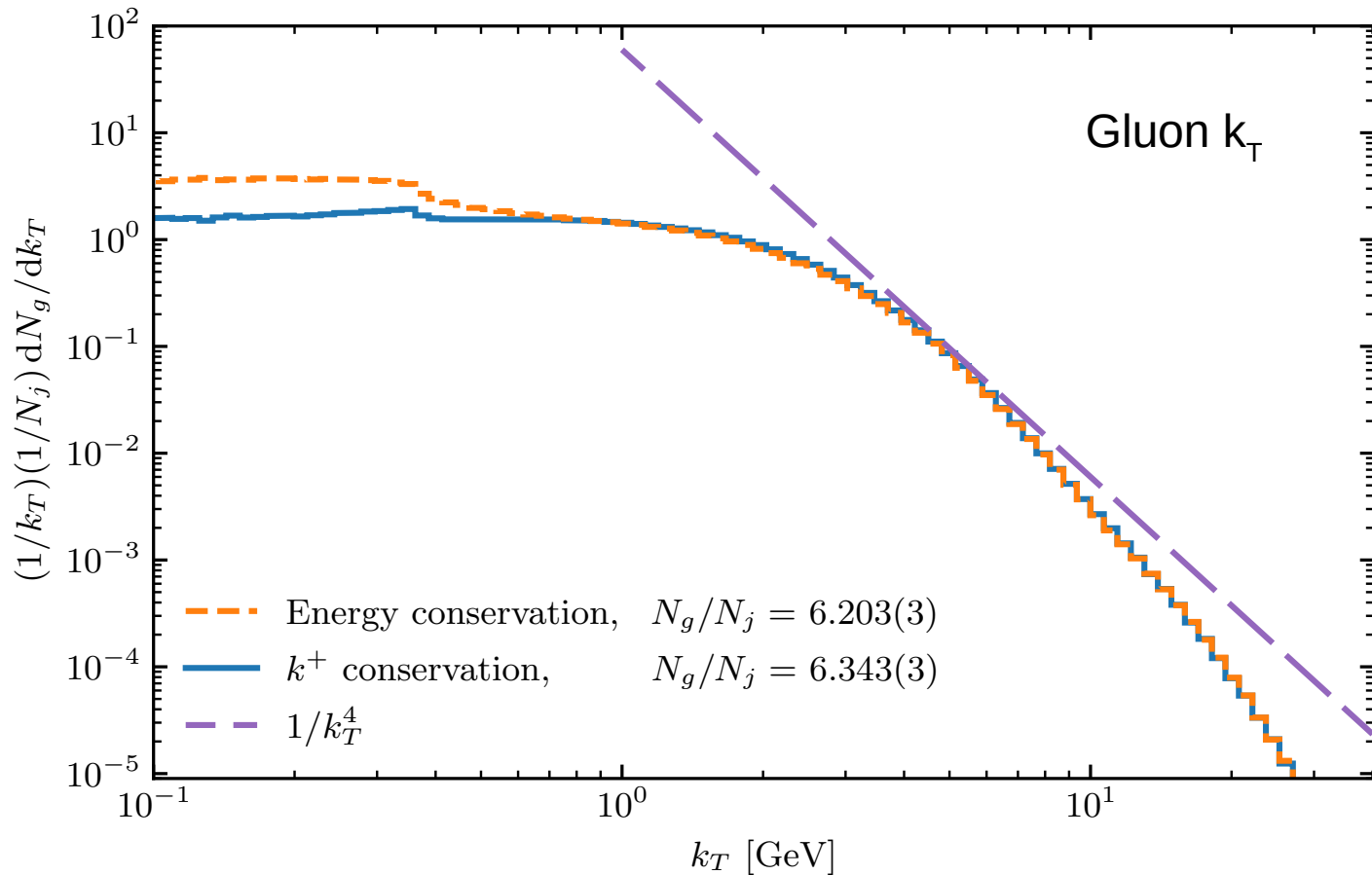
Difference at small
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The Role of the Elastic Scatterings



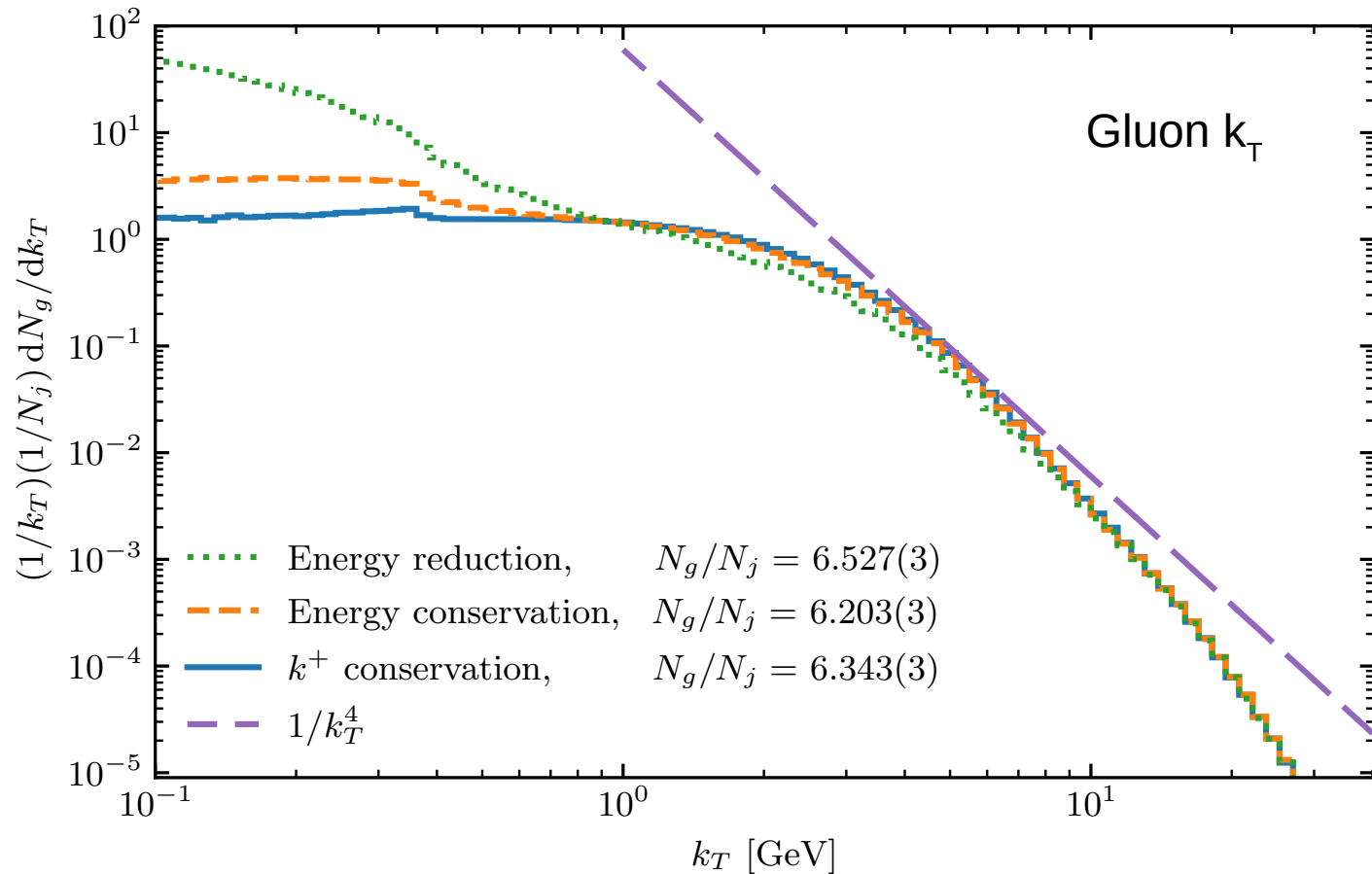
Large difference at small k_T

The Role of the Elastic Scatterings



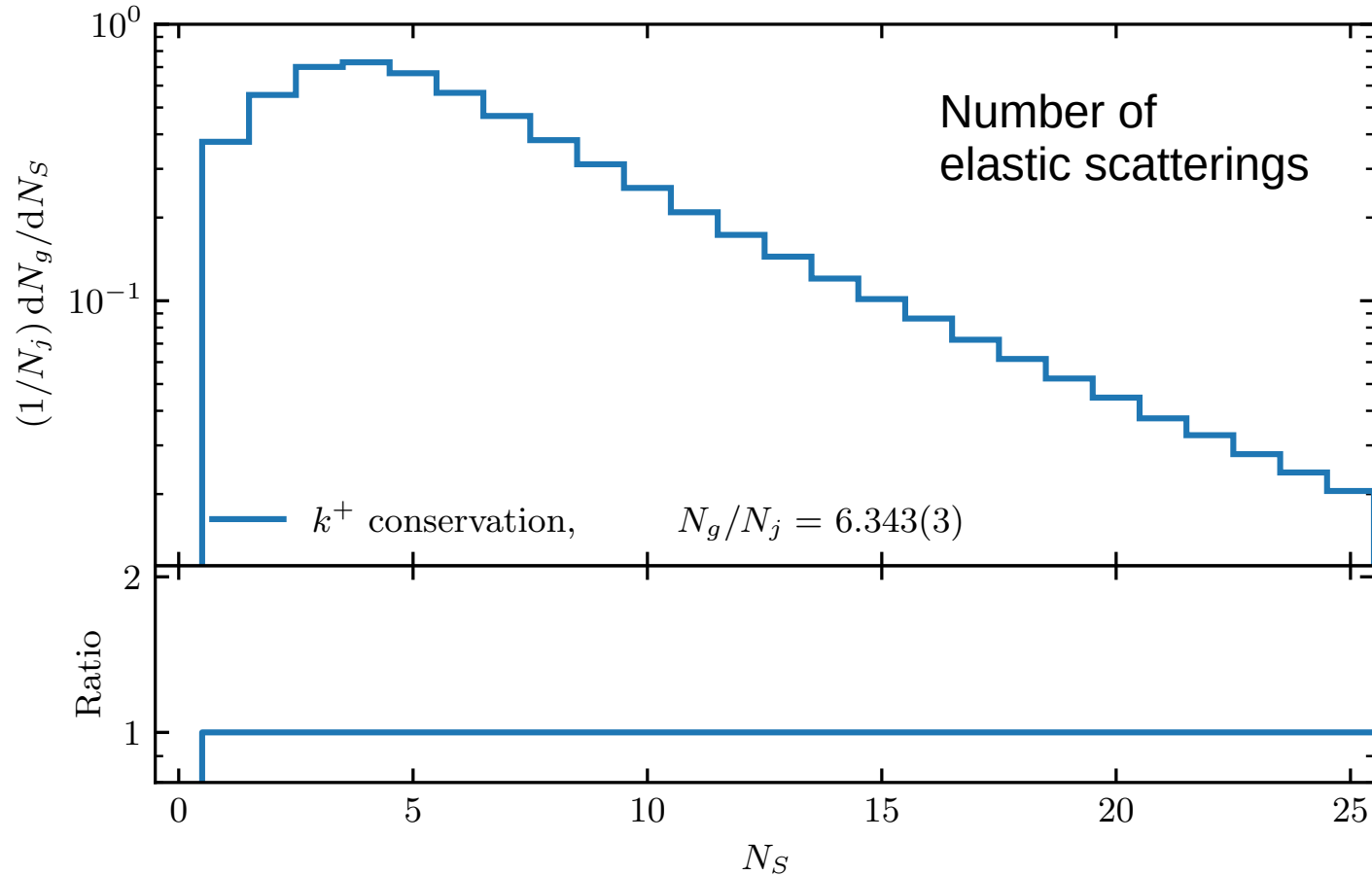
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The Role of the Elastic Scatterings

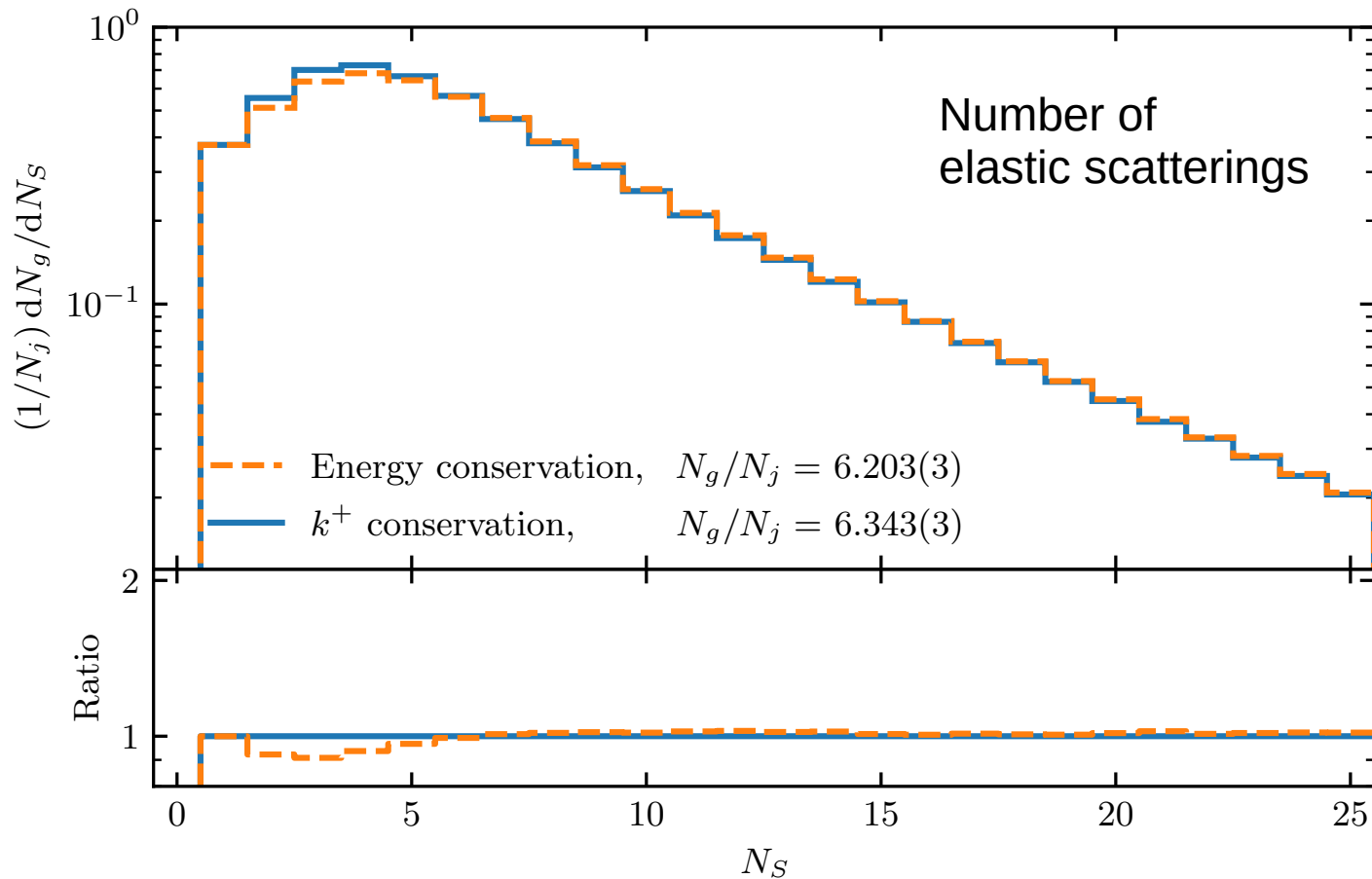


Large difference at small k_T

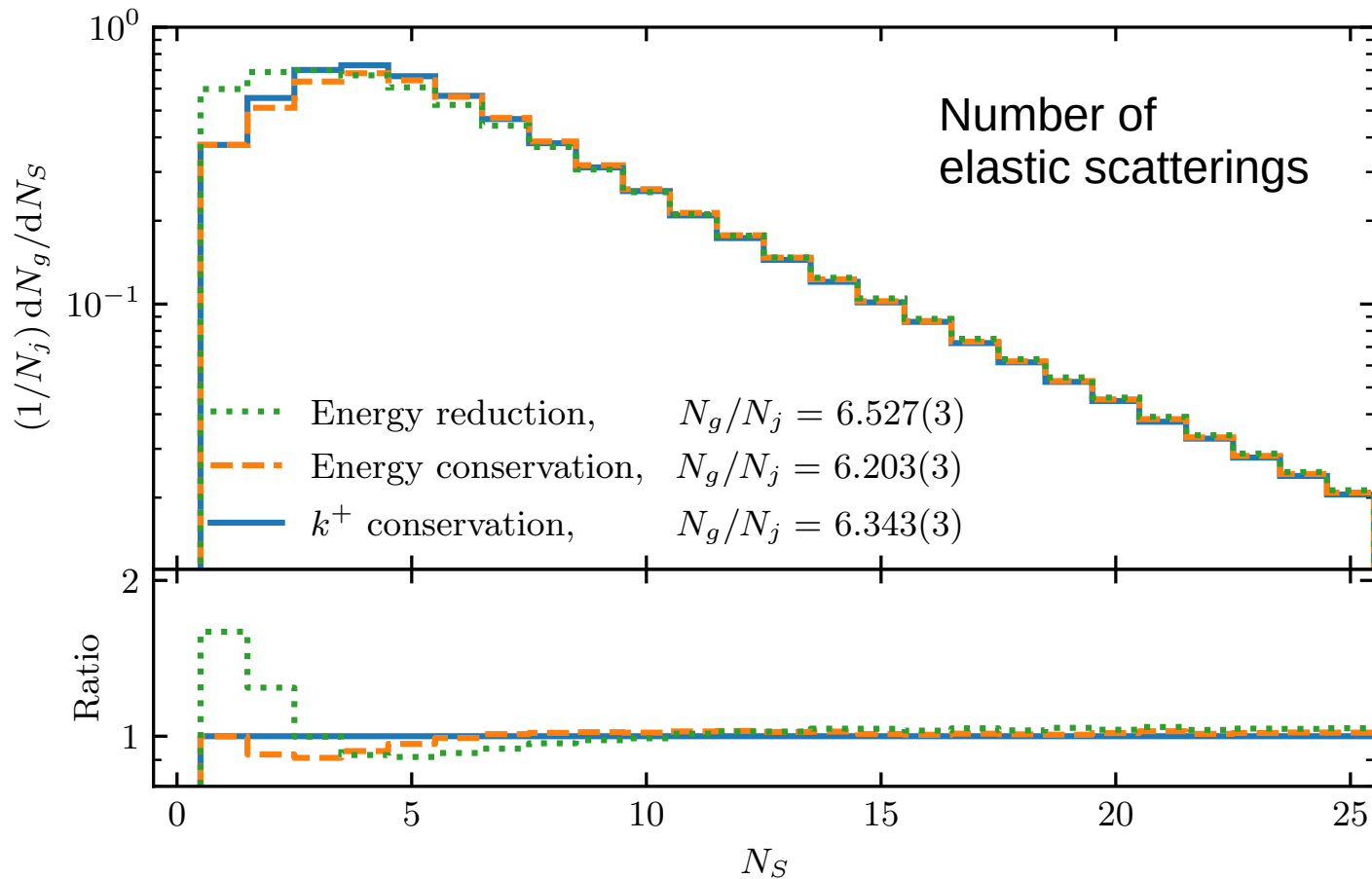
The Role of the Elastic Scatterings



The Role of the Elastic Scatterings

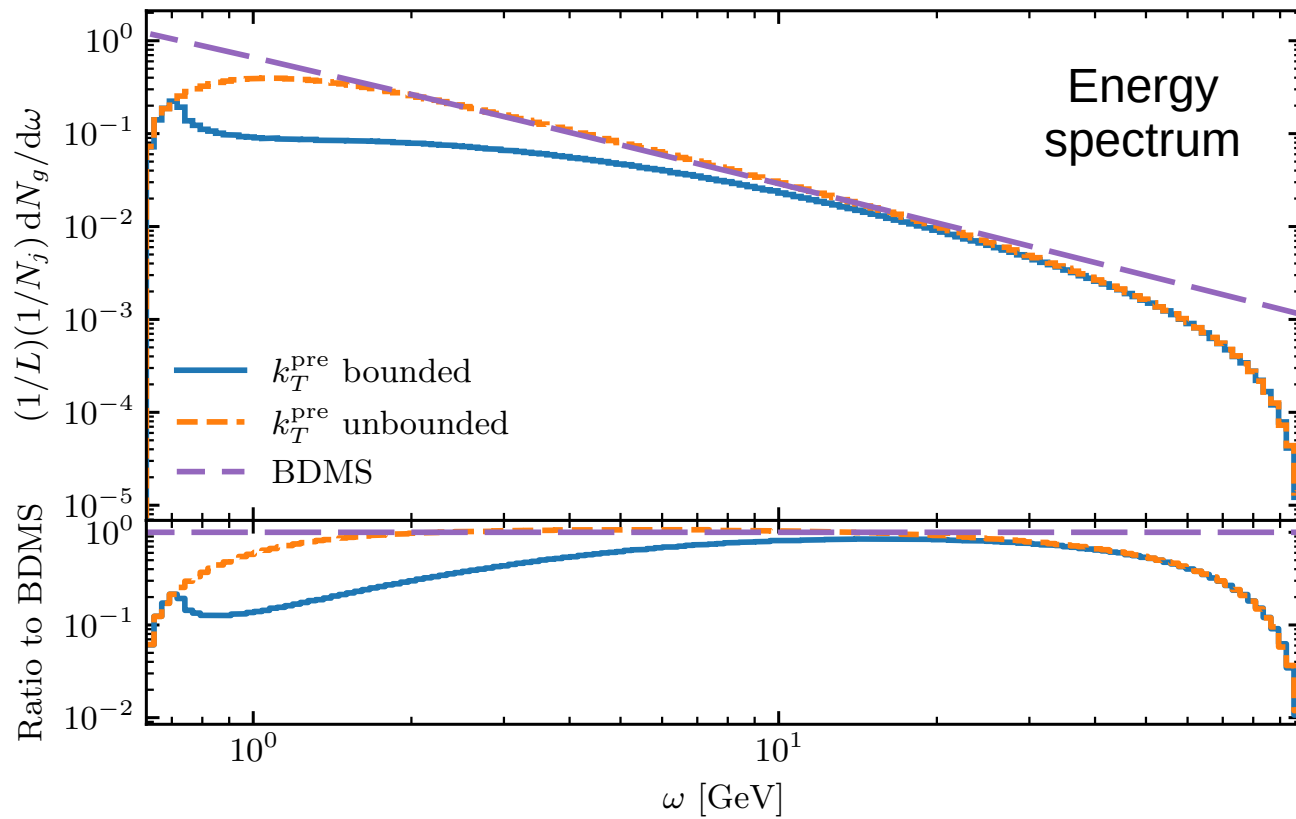


The Role of the Elastic Scatterings



The energy reduction case is larger at $N_s = 1$
→ Larger probability of emission

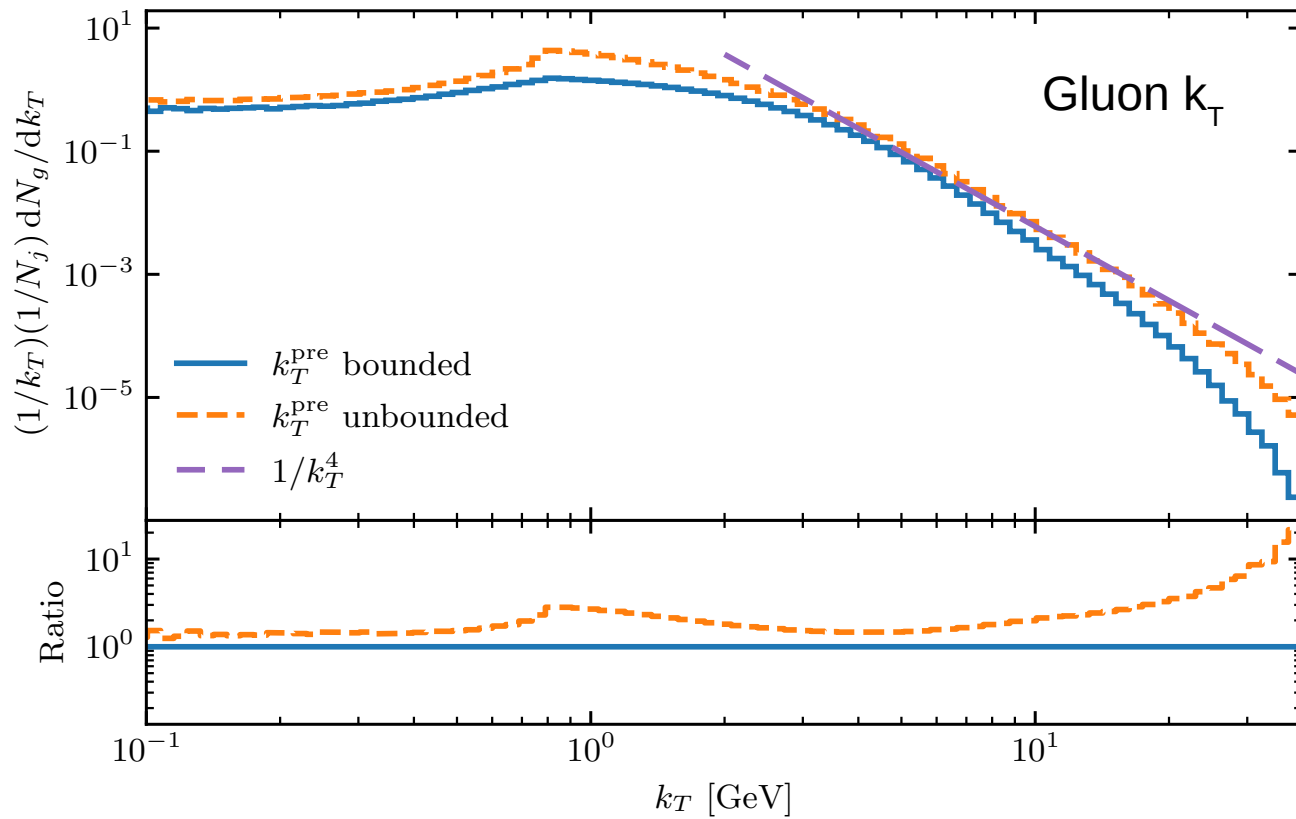
The Role of the Colinearity Hypothesis



Colinearity hypothesis

$$k_T \ll \omega$$

The Role of the Colinearity Hypothesis



Colinearity hypothesis

$$k_T \ll \omega$$

Looking Forward: Towards More Realism

Next step:

- Interface with vHLLE to get hydro evolution of the medium
- Running strong coupling in elastic scatterings
- Start with high virtuality partons
- Sampling of initial parton p_T

$$\frac{d\sigma}{dp_T} \sim p_T^{-6.5}$$

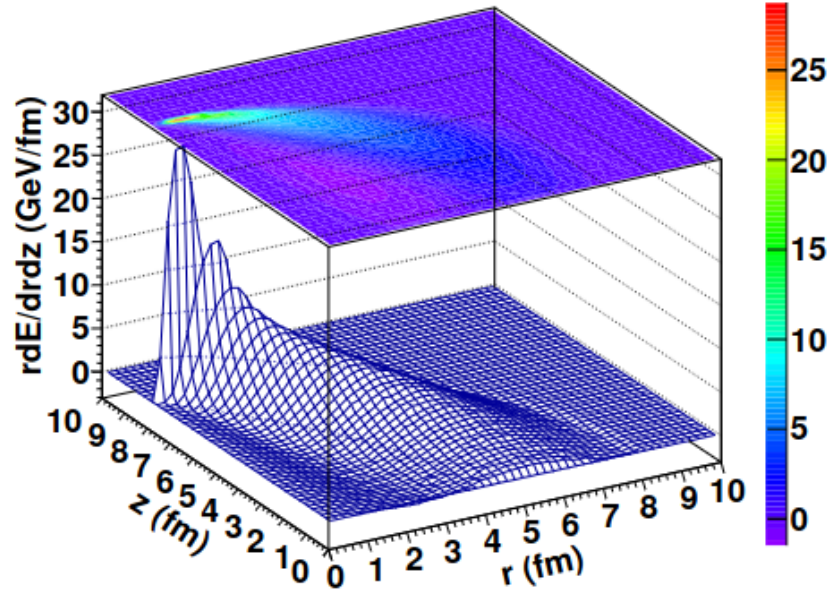
- Run with hadronisation and jet finding



Looking Forward: Effect on the Medium

The jet also affects the medium

(b) $t=8 \text{ fm}/c$



'Wake wave'
in the medium
due to the jet

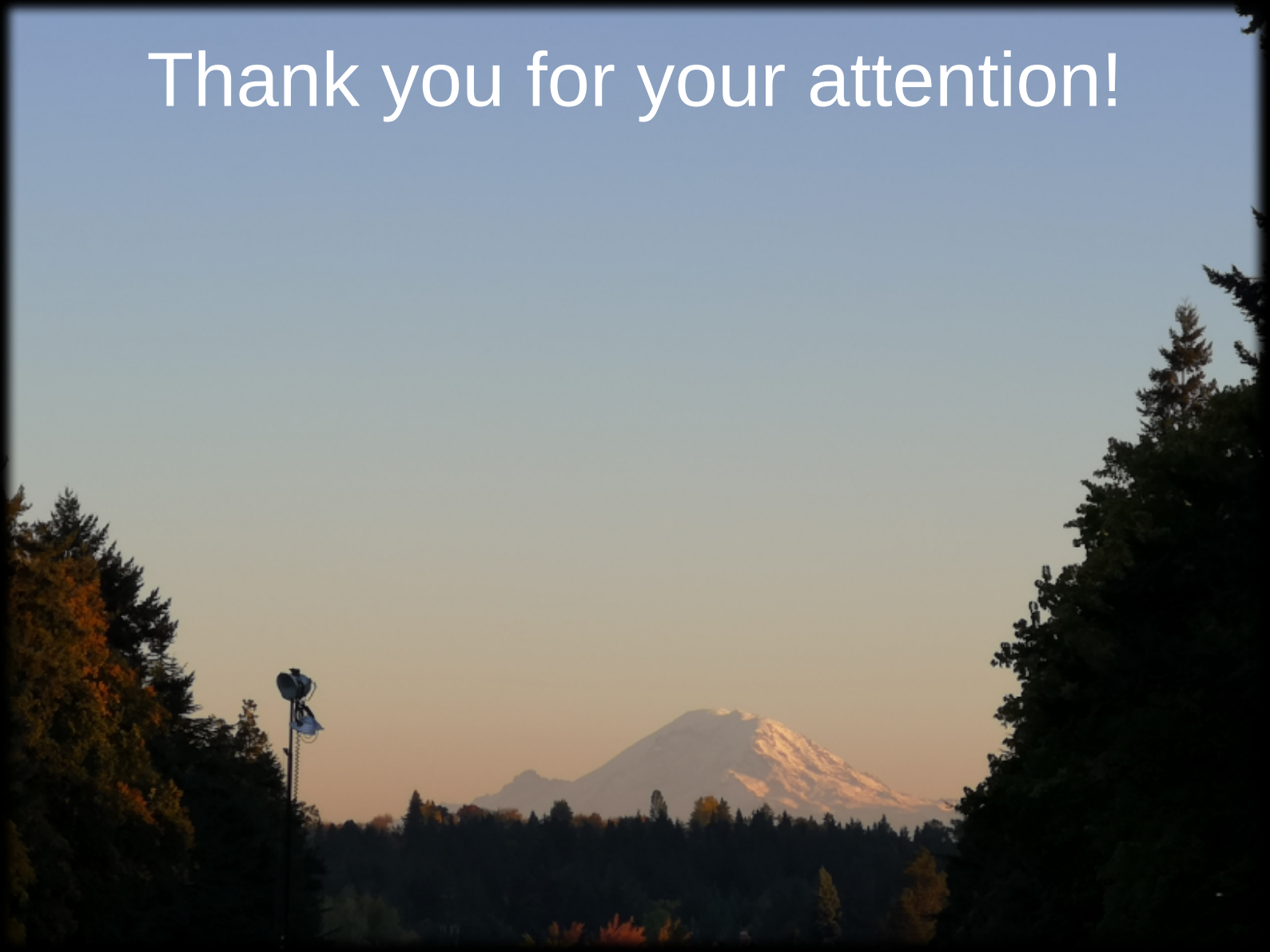


Xin-Nian Wang's talk from Monday

Summary

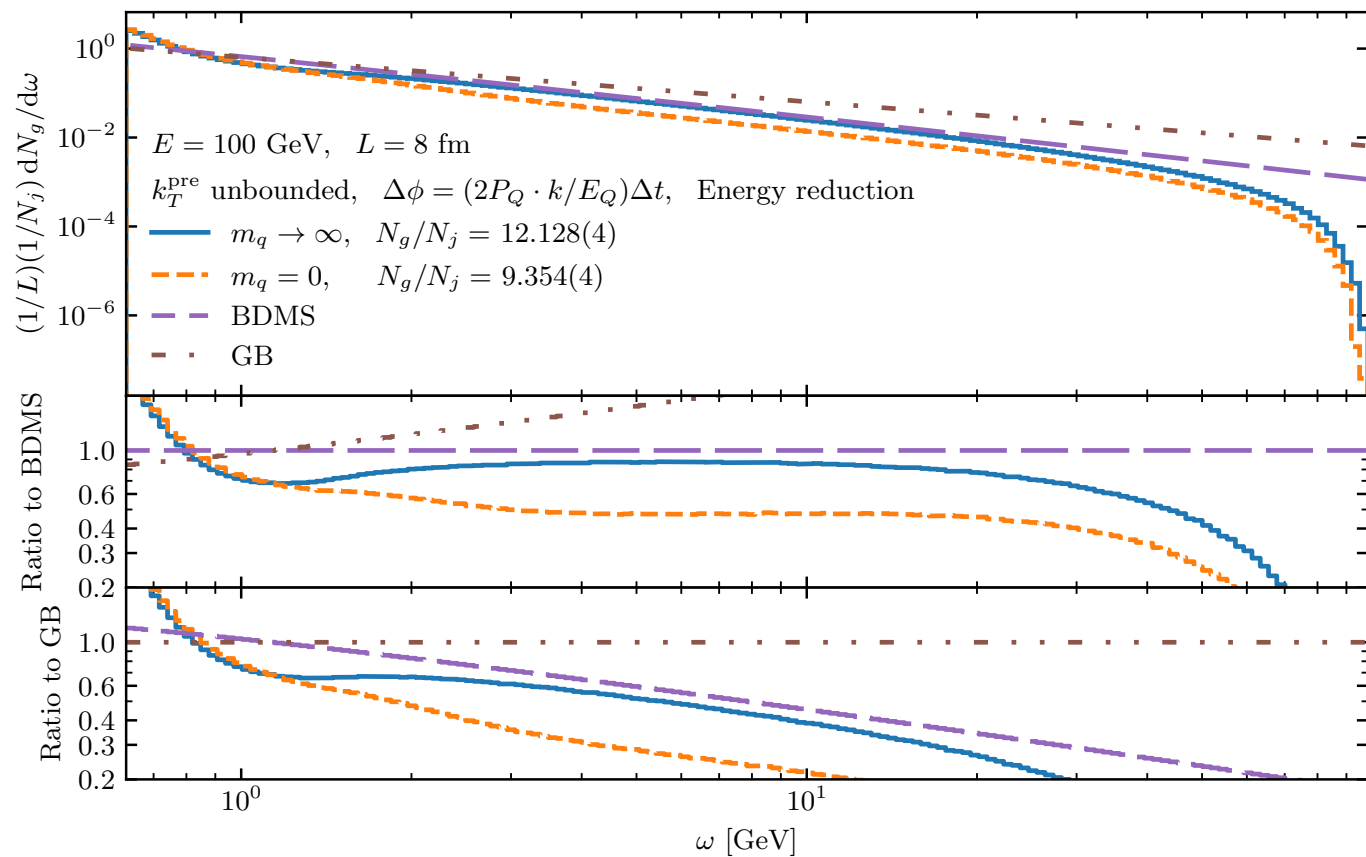
- We have presented a new model for jet energy loss in heavy ion collisions
- Implementation in a Monte Carlo framework
- **1st step done:**
 - Reproduction of the BDMS radiation energy spectrum
 - Shown effects of different model assumptions
- **2nd step:** First results with hydro evolution interface to vHLLE
- **3rd step:** Implementation within the new EPOS4
 - **EPOS4+JETS** – Initial state, hydro, and hadronisation from EPOS4

Thank you for your attention!



Backup Slides

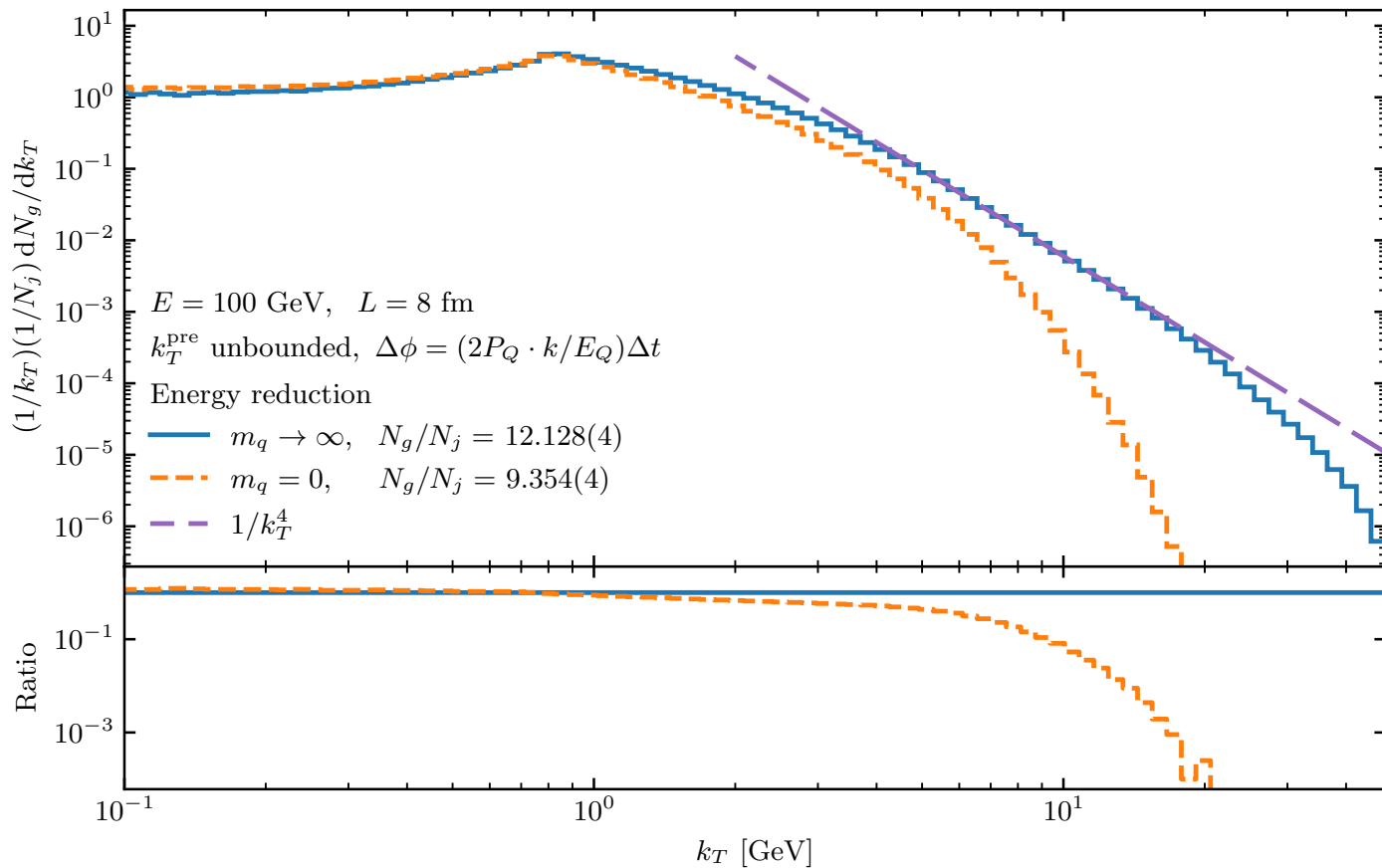
The Role of Scattering Centre mass m_q



Energy spectrum

Effect of mass of scattering centre in the initial GB seed

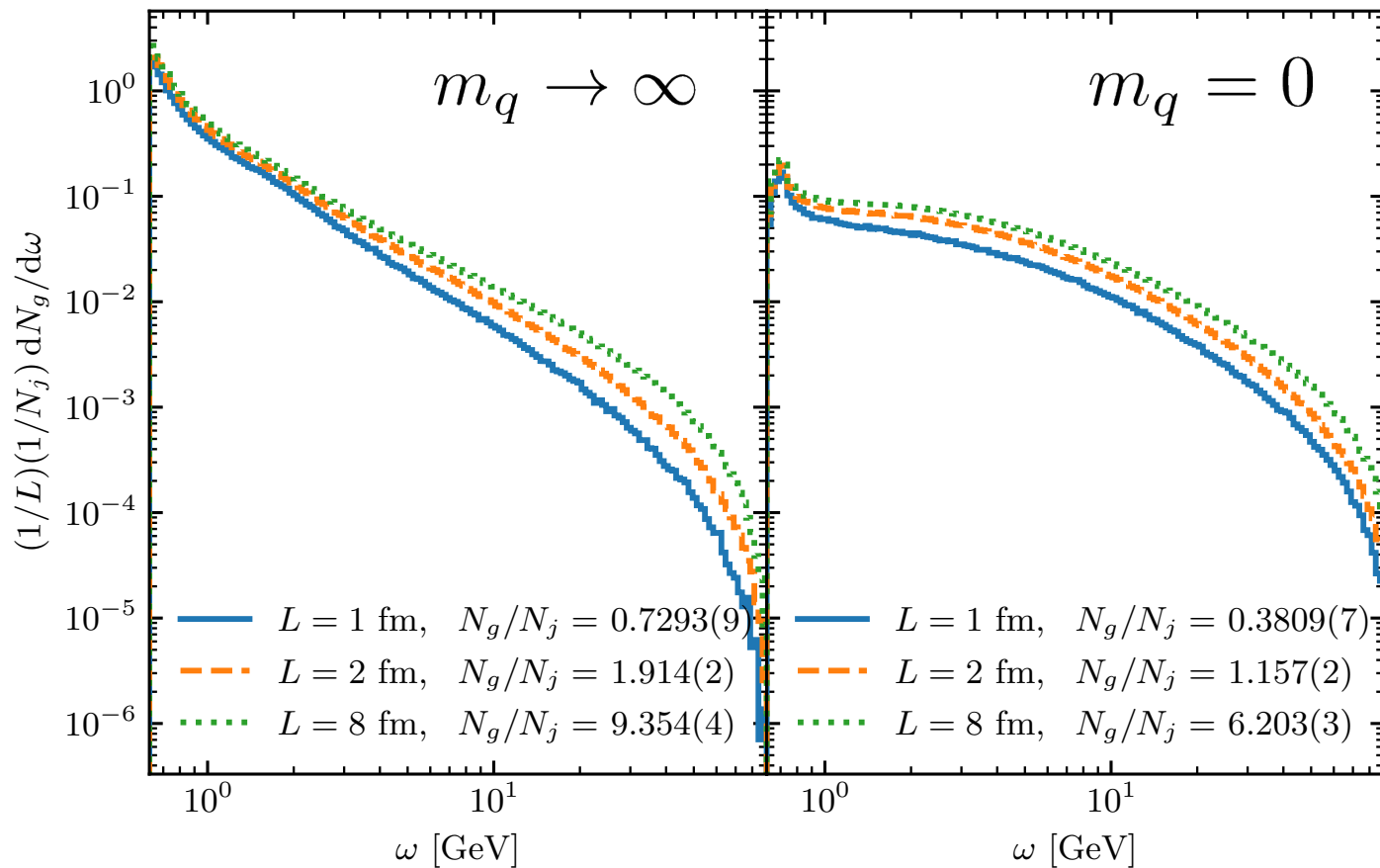
The Role of Scattering Centre mass m_q



Gluon k_T

Effect of mass of scattering centre in the initial GB seed

Effect of Path Length



Energy spectrum
for different path lengths
(medium sizes)