# The $R_{pA}$ and $v_2$ puzzle of $D^0$ mesons in p–Pb collisions at the LHC

Zi-Wei Lin East Carolina University (ECU)

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# Outline

- The  $D^0 R_{pA}$  and  $v_2$  puzzle
- An improved multi-phase transport model
- A solution of the  $R_{pA}/v_2$  puzzle with the Cronin effect
- Effects from parton scatterings & Cronin effect
- Summary

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## The $D^0 R_{pA}$ and $v_2$ puzzle



It has been a challenge to describe both data simultaneously:

- sizable  $v_2 \rightarrow$  significant charm quark interaction with medium  $\rightarrow$  suppression of charm high  $p_T$  spectrum in pA and  $R_{pA}$  (above)
- Studies based on color glass condensate can describe D and J/ $\psi$  v<sub>2</sub>, no R<sub>pA</sub> results yet. Cheng Zhang et al. PRL (2019), PRD (2020)

## The $D^0 R_{pA}$ and $v_2$ puzzle

• Without charm quark scatterings (below),



- This was seen in an earlier study: ~ no suppression in  $R_{pA}$ , then  $v_2$  is too small. Beraudo et al. JHEP (2016)
- A simultaneous description of the R<sub>pA</sub> and v<sub>2</sub> data could disentangle different effects (*initial state correlations, cold nuclear, hot medium*) and help understand onset of collectivity & formation of parton matter or QGP

## An improved multi-phase transport model

We use a multi-phase transport (AMPT) model for this study. It was constructed as a self-contained kinetic description of heavy ion collisions:

- evolves the system from initial condition to final observables;
- particle productions of all flavors from low to high p<sub>T</sub>;
- non-equilibrium initial condition & dynamics (more important for small systems).



## Improved heavy flavor (HF) productions in AMPT

 $gg \rightarrow gg$  cross section in leading-order pQCD  $\frac{d\sigma}{dt} \sim \frac{9\pi\alpha_s^2}{2t^2}$ is divergent for massless g, for minijets (of ALL flavors).

But heavy flavor production does not need a cutoff due to heavy quark mass  $>> \Lambda_{QCD}$  (e.g. in FONLL)

$$g + g \rightarrow Q + \overline{Q}, \quad q + \overline{q} \rightarrow Q + \overline{Q}, \dots$$

- So we remove the  $p_0$  cut on HF productions Zheng et al. PRC (2020) in the two-component model HIJING for AMPT.
- Unlike HIJING, we include HF in  $\sigma_{jet}$ :  $\sigma_{jet} = \sigma_{jet}^{LF} + \sigma^{HF}$
- We also correct factor of  $\frac{1}{2}$  in certain  $\sigma_{iet}$  channels

## Improved heavy flavor (HF) productions in AMPT





- Old/public AMPT charm yield << data
- Removing p<sub>0</sub> in HF production greatly enhances charm yield
- AMPT model now well describes world data on total  $c\bar{c}$  cross section

#### Local scaling for self-consistent size dependence in AMPT

Lund symmetric string fragmentation function:  $f(z) \propto z^{-1}(1-z)^{a_L} e^{-b_L m_T^2/z}$  $b_L$  typical values (in 1/GeV<sup>2</sup>): ~ 0.58 (PYTHIA6.2), 0.9 (HIJING1.0), 0.7-0.9 (AMPT for pp)

 $b_L \sim 0.15$  is needed for string melting AMPT to describe the bulk matter at high energy AA collisions. This corresponds to a much higher string tension  $\kappa$ :  $(p_T^2) \propto \kappa \propto \frac{1}{b_L(2+a_L)}$  ZWL et al. PRC (2005)

pp and AA collisions need different values of  $\mathbf{b}_{\mathbf{L}}$ ; same for Chao Zhang et al. PRC (2019)

minijet cutoff  $\mathbf{p}_0$  (for modern PDFs, is related to  $Q_s \propto A^{1/6}$ ) Zheng et al. PRC (2020)

 $\rightarrow$  We scale them with local nuclear thickness functions:

$$b_L(s_A, s_B, s) = \frac{b_L^{p_P}}{[\sqrt{T_A(s_A)T_B(s_B)}/T_p]^{\beta(s)}}$$
  

$$p_0(s_A, s_B, s) = p_0^{p_P}(s)[\sqrt{T_A(s_A)T_B(s_B)}/T_p]^{\alpha(s)}$$
  
Chao Zhang

We fit charged hadron  $\langle p_T \rangle$  in *pp* to determine  $b_L^{pp} = 0.7$ , then used central AuAu/PbPb  $\langle p_T \rangle$  data to determine  $\alpha(s)$ ,  $\beta(s)$  versus energy  $\sqrt{s}$ 

et al. PRC (2021)

Local scaling for self-consistent size dependence in AMPT

The scaling allows AMPT to self-consistently describe the system size dependence, including centrality dependences of AuAu & PbPb and smaller systems.



Chao Zhang et al. PRC (2021)

Centrality dependence of  $< p_T >$  is now reasonable, much better than public AMPT (v2.26t9)

## Structure of improved AMPT (String Melting version)



The AMPT model used in this  $R_{pA}/v_2$  study contains all these improvements

## A solution of the $R_{pA}/v_2$ puzzle with the Cronin effect

We implement the Cronin effect on initial charm quarks by broadening  $c\bar{c} p_T$  with a random  $k_T$  sampled from

Mangano et al. NPB (1993) Vogt, PRC (2018, 2021)

 $w = w_0 \sqrt{1 + (n_{\text{coll}} - i)\delta}$  grows with # of NN collisions of the wounded nucleon(s).

 $f(\vec{k_{\rm T}}) = \frac{1}{\pi w^2} e^{-k_{\rm T}^2/w^2}$ 



Full model, with Cronin effect at  $\delta=7$ ,  $\sigma_{LQ}=0.5$ mb (for scatterings among u/d/s quarks),  $\sigma_{HQ}=1.5$ mb (for scatterings of charm quarks with other partons), can describe both  $R_{pA}$  and  $v_2$  data of  $D^0$  mesons

### A solution of the $R_{pA}/v_2$ puzzle with the Cronin effect

Without the Cronin effect ( $\delta$ =0): if we get sizable v<sub>2</sub>, then D<sup>0</sup> R<sub>pA</sub> is underestimated due to charm scatterings with the medium (via  $\sigma_{HQ}$ ).



Black curve vs blue curve (*both at*  $\sigma_{HQ}=1.5mb$ ): the Cronin effect significantly increases charm  $R_{pA}$  at moderate/high  $p_T$  but modestly decreases charm  $v_2$ 

#### More on the Cronin effect

Often considered as transverse momentum broadening of a produced parton from a hard process due to multiple scatterings of initial parton(s) in the nucleus

• We take the k<sub>T</sub> width as 
$$w = w_0 \sqrt{1 + (n_{\rm coll} - i)\delta}$$

grows with  $n_{coll}$ : # of NN collisions of the wounded nucleon(s), *i*=1 for  $c\bar{c}$  produced from the radiation of 1 wounded nucleon, =2 for  $c\bar{c}$  produced from the collision of 2 wounded nucleons, This way,  $w=w_0$  for pp collisions.

$$w_0 = (0.35 \text{ GeV}/c) \sqrt{b_{\rm L}^0 (2 + a_{\rm L}^0)/b_{\rm L}/(2 + a_{\rm L})} \propto \sqrt{\kappa}$$
  
motivated by  $\kappa \propto \frac{1}{b_{\rm L}(2 + a_{\rm L})}$  for Lund string fragmentation.

Kopeliovich et al. PRL (2002) Kharzeev et al. PRD (2003) Vitev et al. PRD (2006) Accardi, hep-ph/0212148

## Effects from parton scatterings & Cronin effect

Test results for charm quarks:

- parton scatterings are mostly responsible for generating charm v<sub>2</sub>
- the Cronin effect modestly decreases charm v<sub>2</sub>

- parton scatterings significantly suppress charm spectra at moderate/high p<sub>T</sub>
- the Cronin effect significantly increases charm spectra at moderate/high p<sub>T</sub>



#### At 5.02 TeV, the full model also reasonably describes



15

The full model at 8.16 TeV

at the same  $\sigma_{LQ}=0.5mb$ or a smaller  $\sigma_{LO}=0.3mb$  :



This change of  $\sigma_{LQ}$  has ~no effect on D<sup>0</sup> R<sub>pA</sub> or v<sub>2</sub>:



## Summary

We have studied p-Pb collisions at LHC energies with an improved multi-phase transport model.

Including a strong Cronin effect allows a simultaneous description of the D<sup>0</sup> meson  $R_{pA}$  and  $v_2$  data (at  $p_T \le 8 \text{ GeV/c}$ )

Parton scatterings significantly suppress charm spectra at moderate/ high  $p_T$ , Cronin effect significantly increases charm spectra at moderate/high  $p_T$ and thus compensates for the effect from parton scatterings

Charm  $v_2$  is found to be mostly generated by charm quark scatterings, Cronin effect modestly decreases the charm quark or meson  $v_2$ 

The Cronin effect is expected to grow with the system size, so this implies the importance of including the Cronin effect in heavy flavor studies (especially  $R_{AA}$ ) in large systems