



Nuclear matter, Neutron Stars & Correlations

Isaac Legred (Caltech)

INT

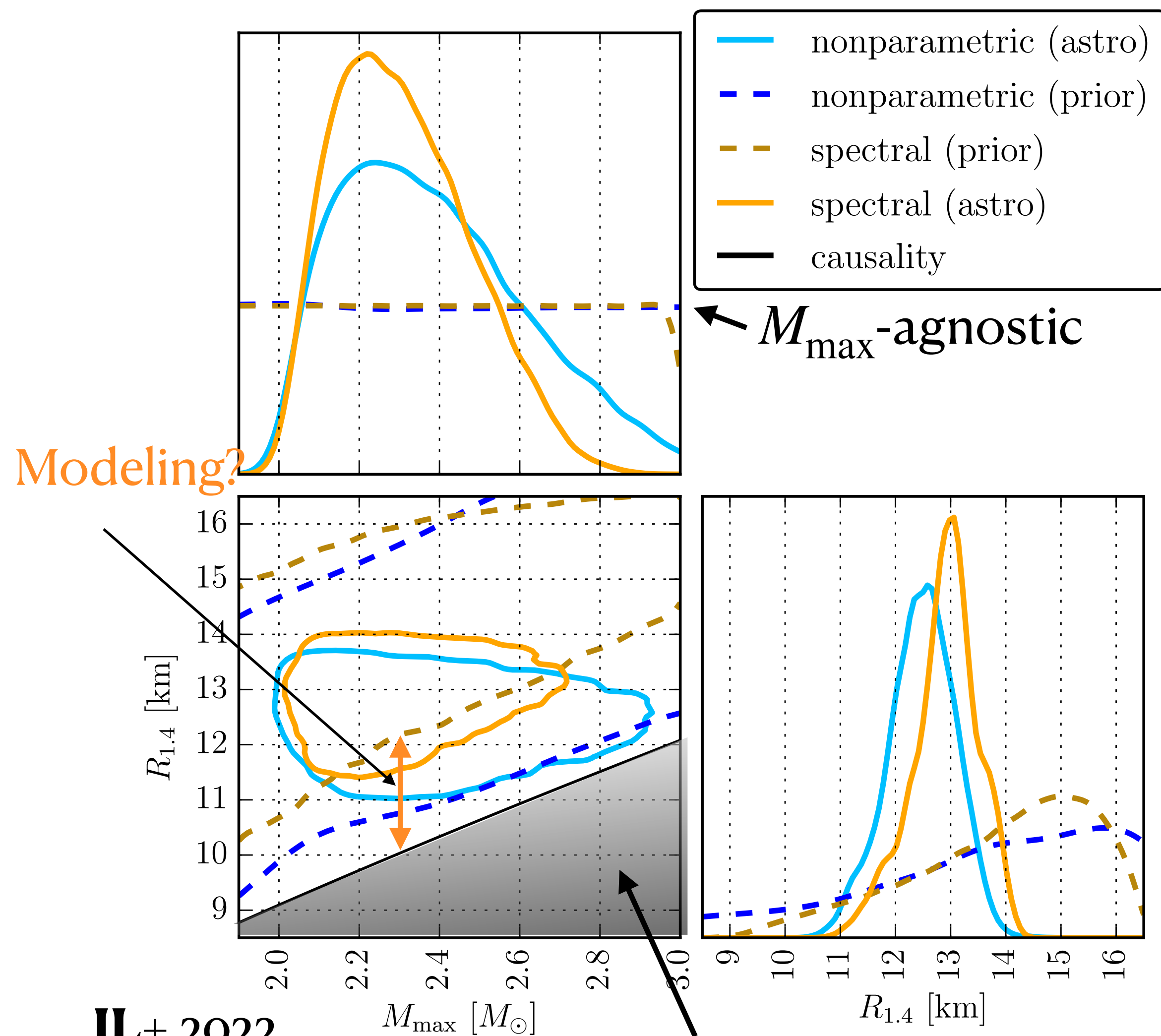
July 19, 2022

Work with: Katerina Chatziioannou,
Reed Essick, and Philippe Landry

Caltech

10.1103/PhysRevD.105.043016
<https://arxiv.org/abs/2201.06791>

Different Priors, Different Results

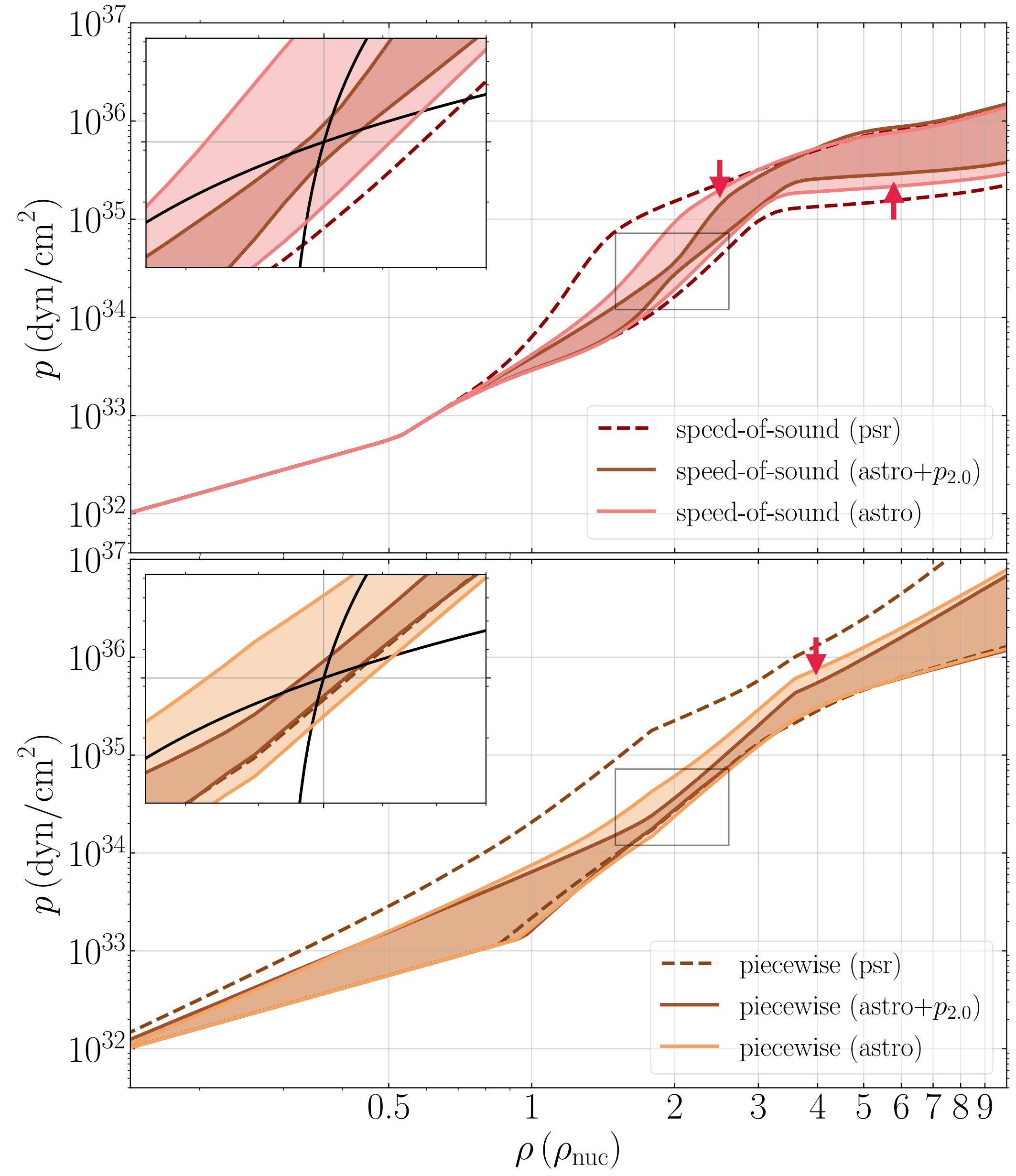
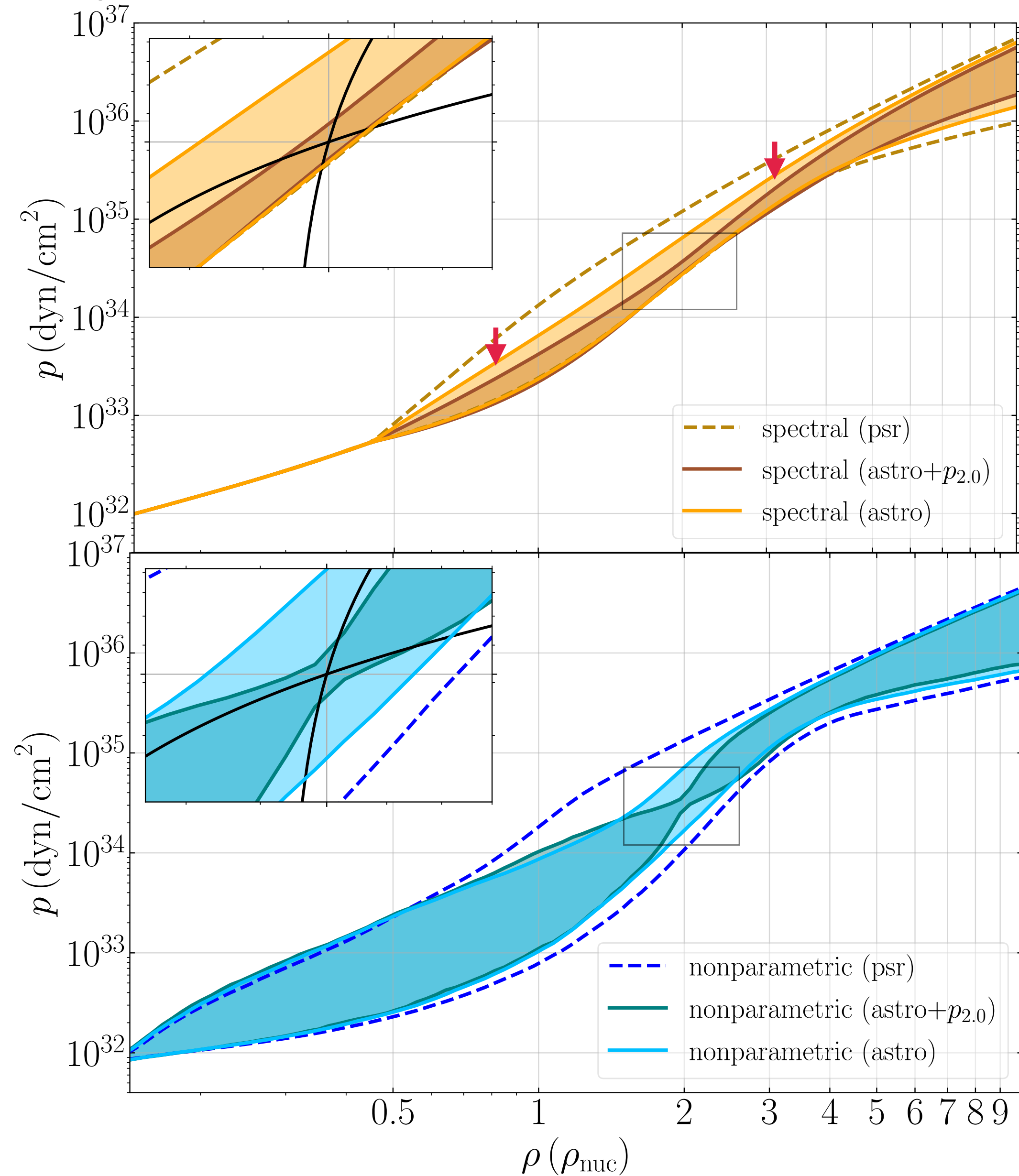


TOV maximum mass and radius of a 1.4 solar mass NS are correlated among equation of state candidates due to causality

Spectral model sees a “tighter correlation” than the Nonparametric model — not likely due to causality!

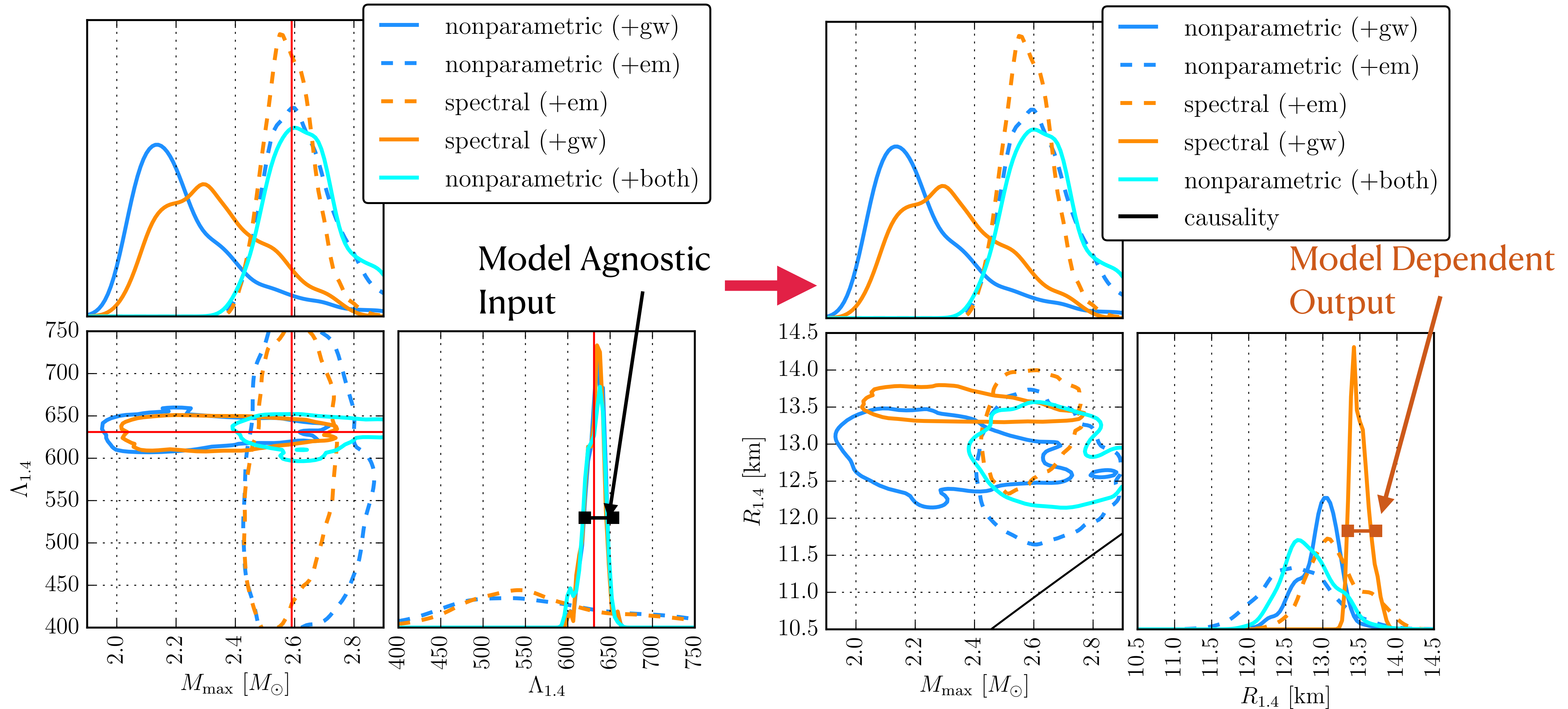
Implicit Correlations - Mock Data

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So What?

“Preliminary”



Knowing $\Lambda_{1.4} \neq$ Knowing $R_{1.4}$

Conclusions

- Phenomenological models of the nuclear equation of state can build in (often hidden) correlations due to the functional form of the EoS
- Nonparametric models (such as the Gaussian Process model), can provide more flexibility in inference of the EoS (but do not guarantee it)
- Need to be very careful when talking about translating constraints between variables, both micro and macroscopically

Inferring the EoS — In practice

- Want to establish a probability distribution on candidate equations of state given observed astrophysical data

Equation of state candidate

$$P(\varepsilon_i | d) = P(\varepsilon | d_1, d_2, \dots) \propto P(d_1, d_2, \dots | \varepsilon_i) \times \boxed{P(\varepsilon_i)}$$

Prior

Astrophysical data

Inferring the EoS — In practice

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Equation of state candidate

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Astrophysical data

Phenomenological

Parametrize a *functional form* (i.e. Spectral, Piecewise-polytrope)

Nonparametric methods, i.e. Gaussian process (GP)

Tabulated models from nuclear theory

Nonparametric: Gaussian Process

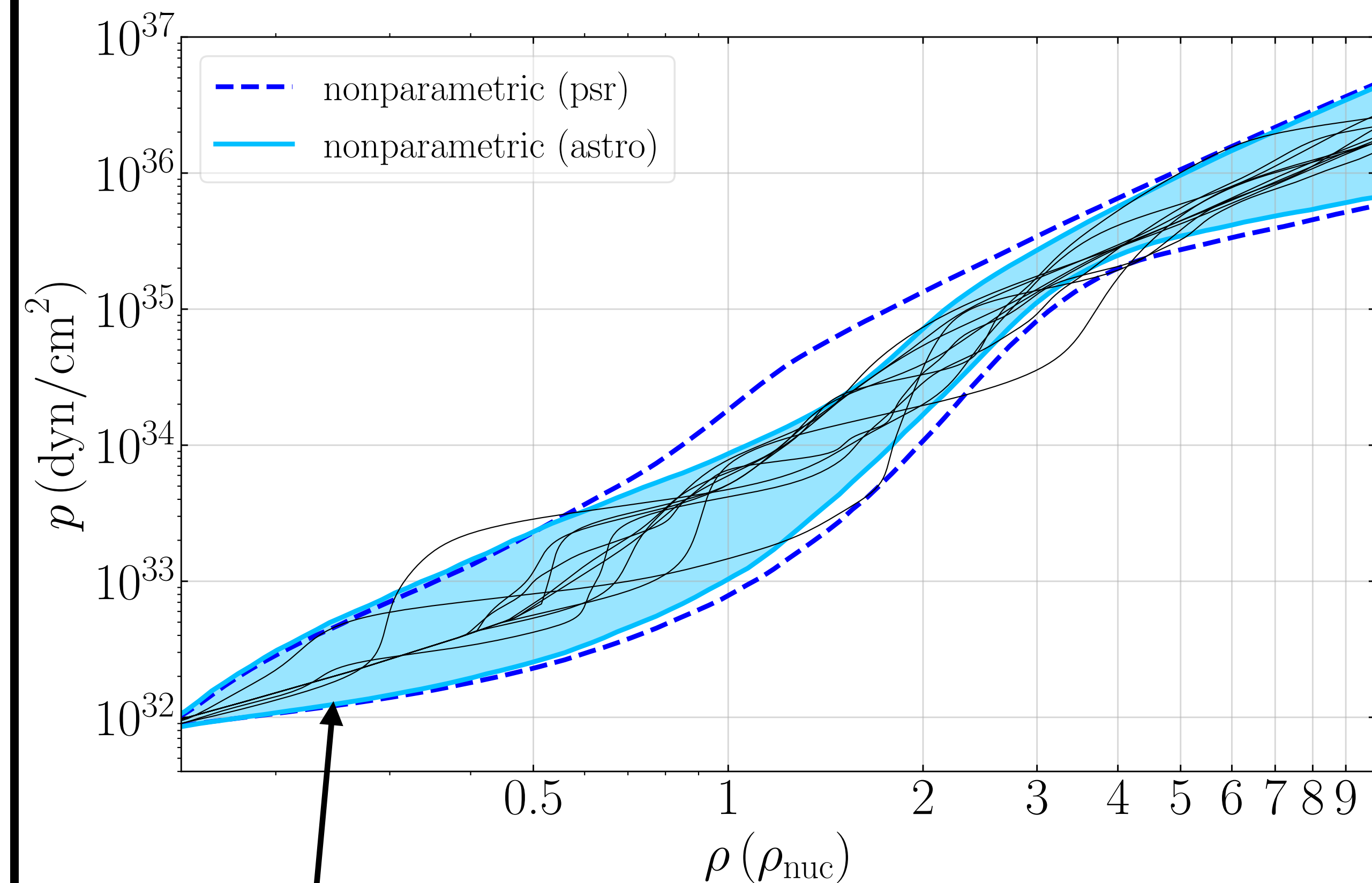
Gaussian Process Regression (Landry and Essick 2018)

Tabulate a draw $\phi(p_i) = \ln(1/c_s^2(p_i) - 1)$ @ Pressures p_i from a multivariate Gaussian distribution

Parameters for the covariance kernel are chosen to Control “shape” of EoS distribution

Model-Agnostic Prior (broadest range of models)

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90% credible interval for $p(\rho)$

psr -> just heavy pulsar mass measurements

astro -> Heaviest pulsar, 2 NICER x-ray, 2 GWs

Parametric

Spectral (Lindblom 2010)

Parametrize the adiabatic index

$$p(\rho) = \rho^{\Gamma(x)} \quad \Gamma(x) = \sum_{i=0}^n \gamma_i (\log(x))^i$$

Piecewise-polytrope (Read 2008)

A polytrope with multiple segments

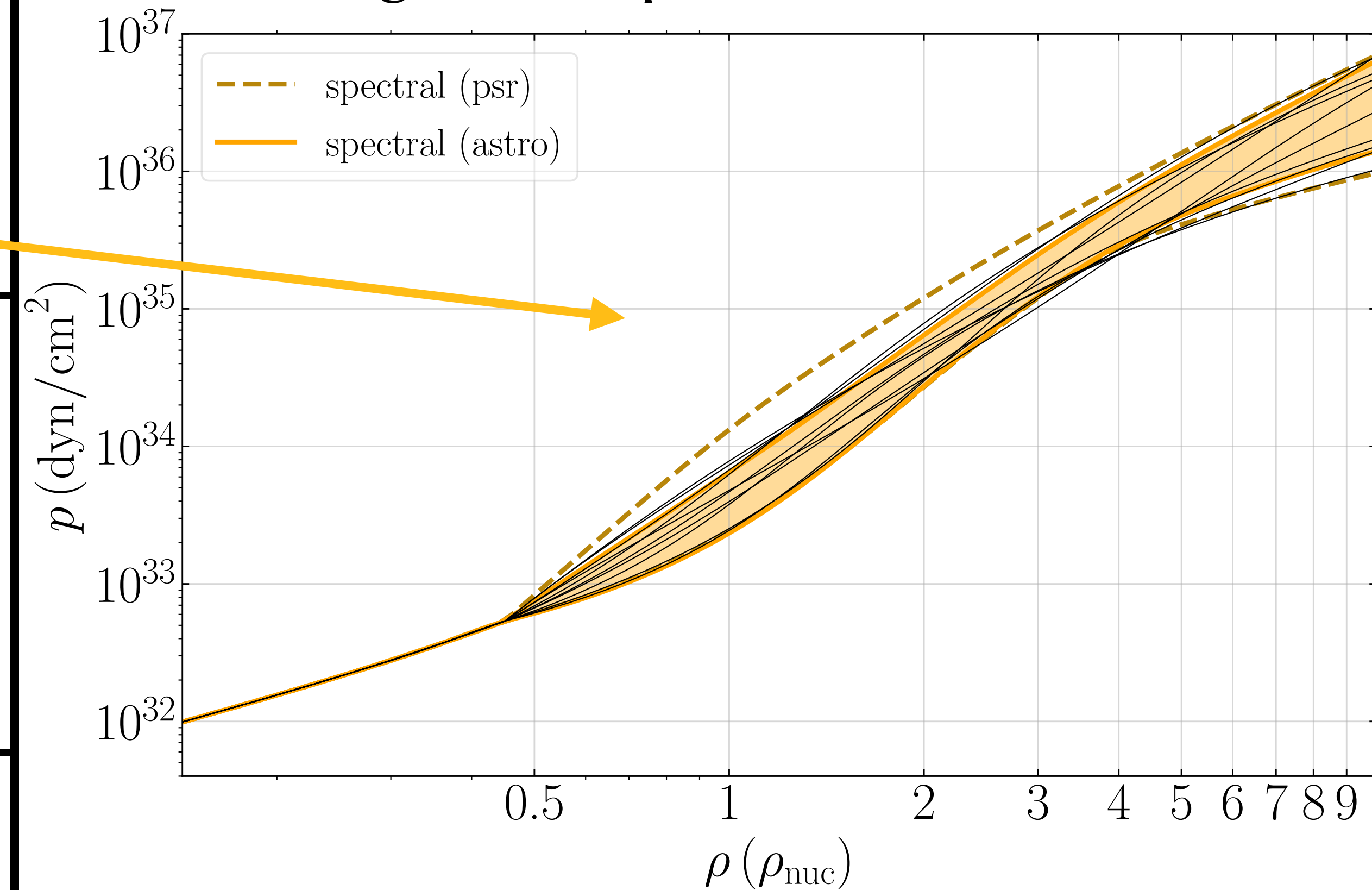
$$p(\rho) = \begin{cases} K_1 \rho^{\Gamma_1} & : \rho < \rho_1 \\ K_2 \rho^{\Gamma_2} & : \rho_1 < \rho < \rho_2 \\ K_3 \rho^{\Gamma_3} & : \rho_2 < \rho \end{cases}$$

Direct speed-of-sound (Greif 2018)

A bump in the speed of sound before asymptotic behavior

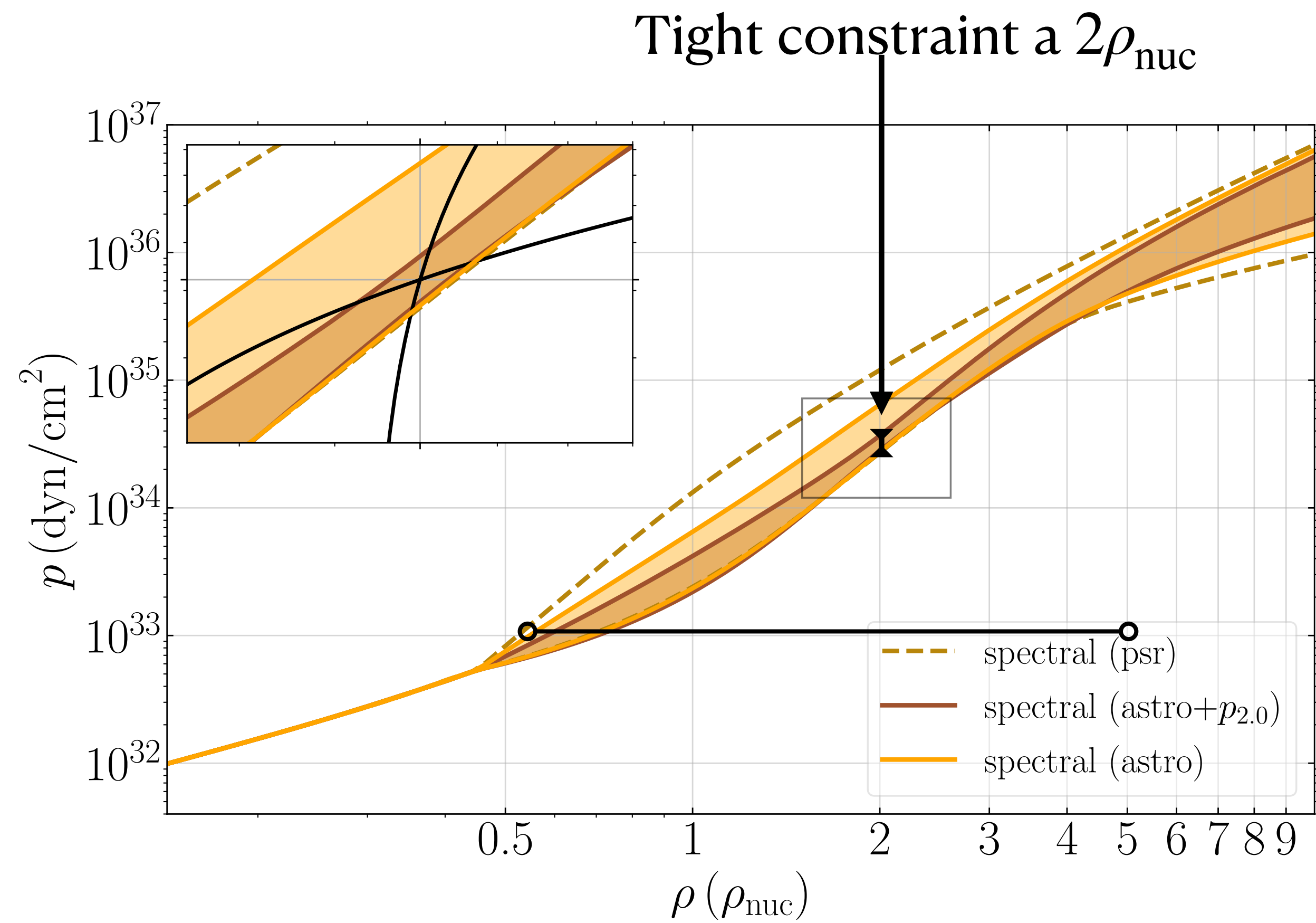
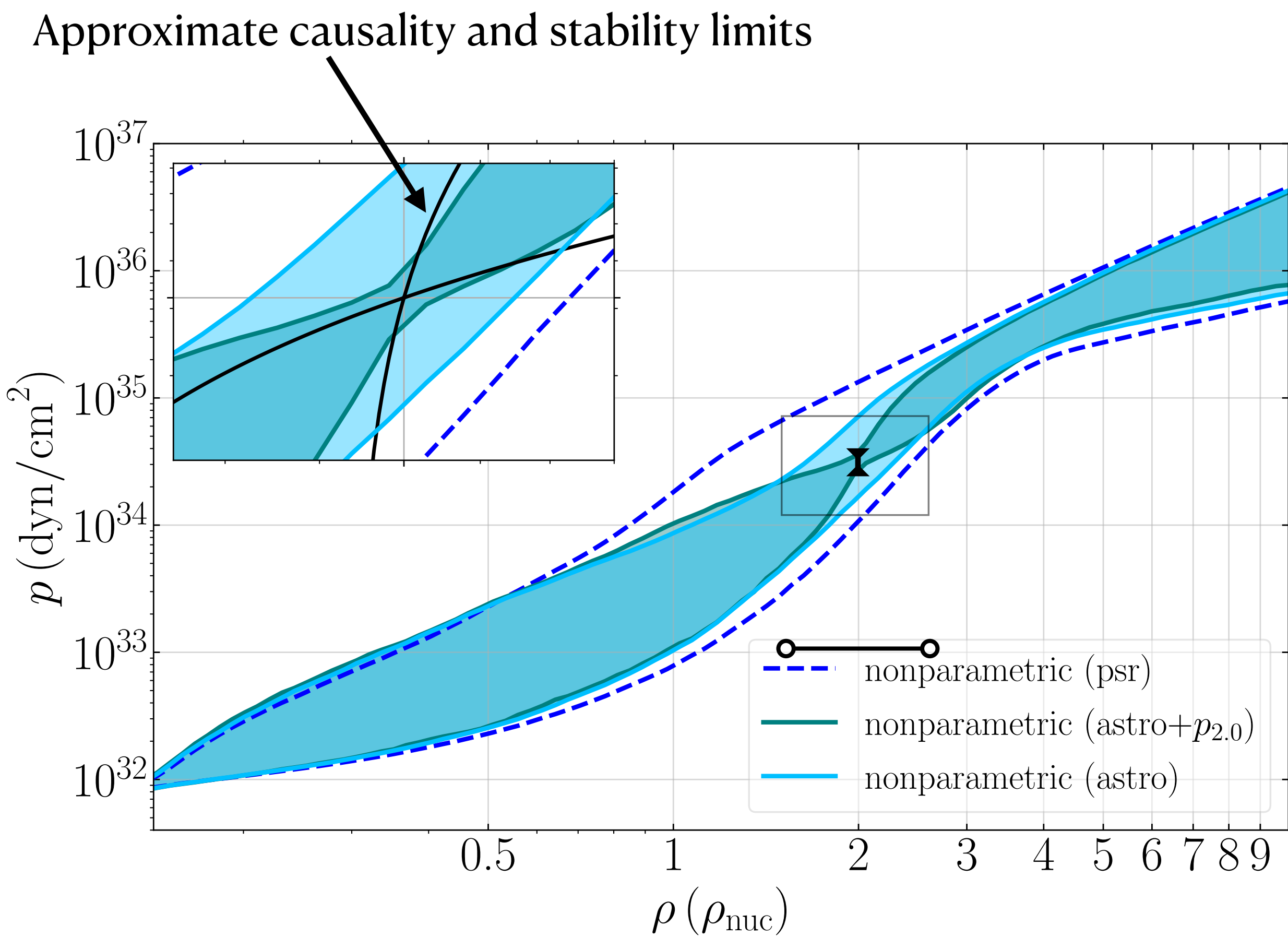
$$\frac{c_s^2(z)}{c^2} = a_1 e^{-\frac{1}{2}(z-a_2)^2/a_3^2} + a_6 + \frac{\frac{1}{3} - a_6}{1 + e^{-a_5(z-a_4)}}$$

E.g. for the Spectral Parametrization



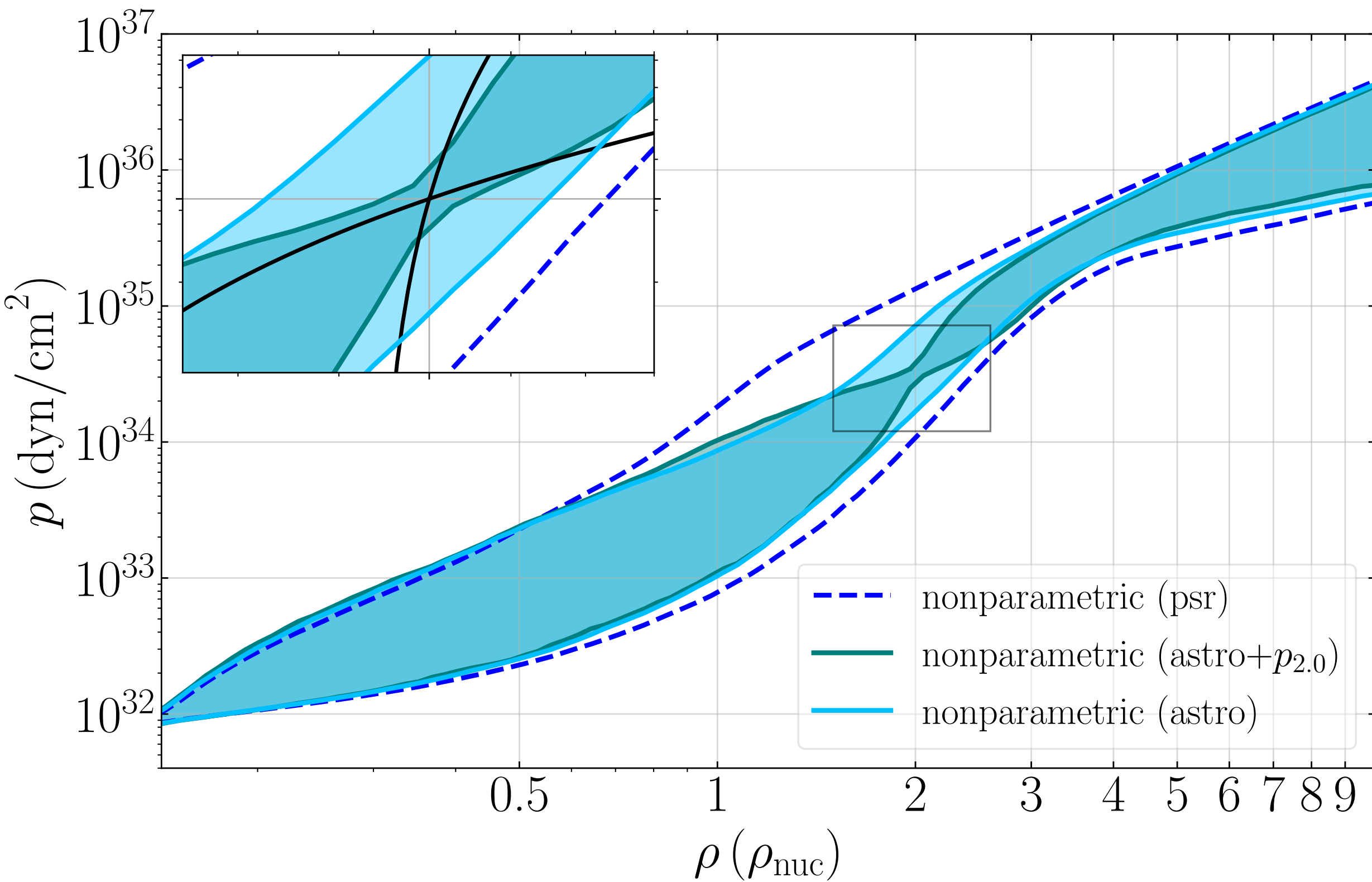
Correlations

Correlations between astro observables \Leftrightarrow Correlations between density scales



Implicit Correlations

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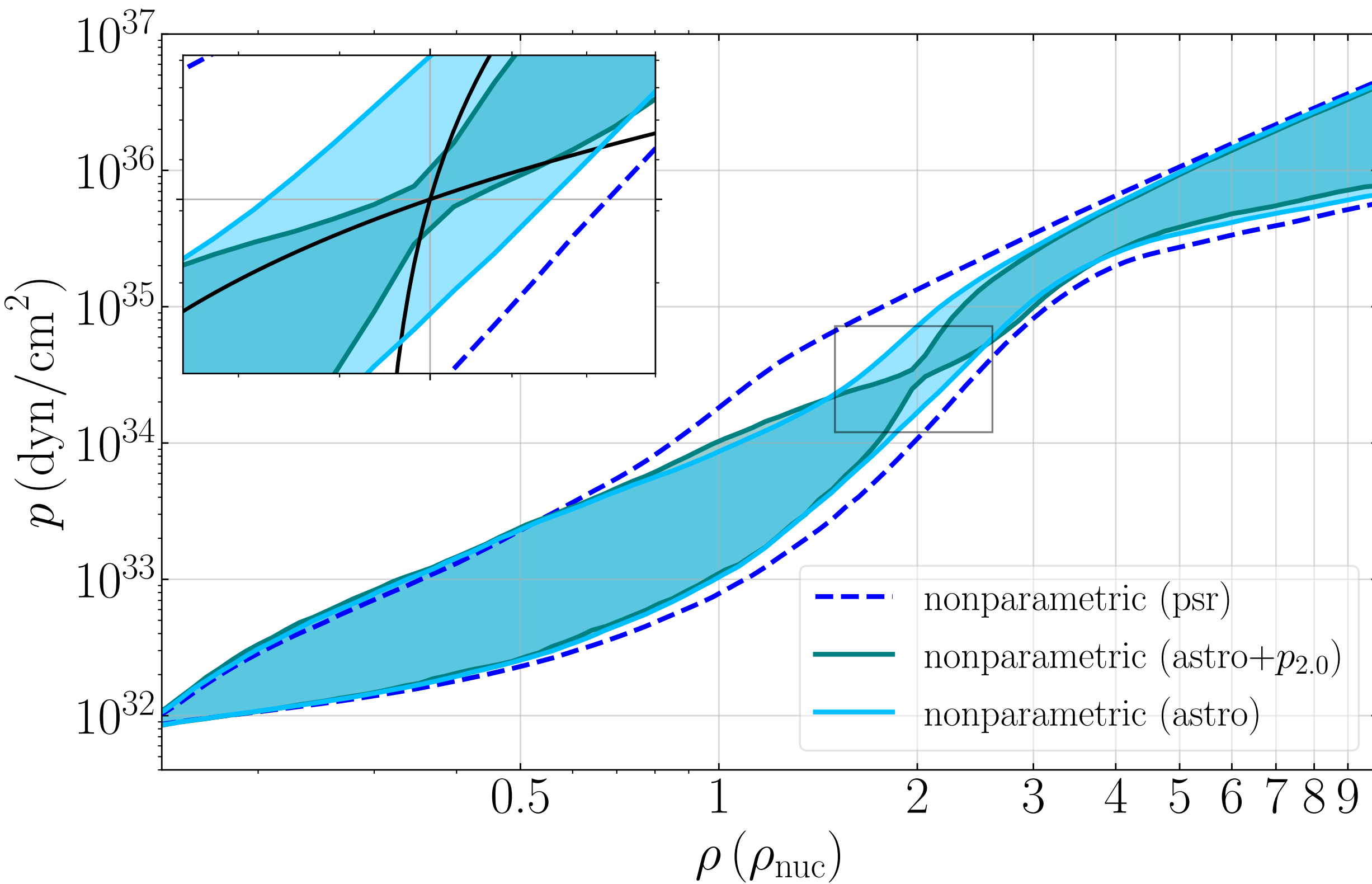
Quantifying correlations — Mutual Information

How much information is gained about other density Scales by knowing the EoS at some fixed density

$$I(p_a, p_b) \equiv \int dp_a dp_b P(p_a, p_b) \ln \left(\frac{P(p_a, p_b)}{P(p_a)P(p_b)} \right)$$

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Implicit Correlations



Quantifying correlations — Mutual Information

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Also a K-L divergence!

$$I(p_a, p_b) = \int dp_a P(p_a) \int dp_b P(p_b | p_a) \ln \left(\frac{P(p_b | p_a)}{P(p_b)} \right)$$

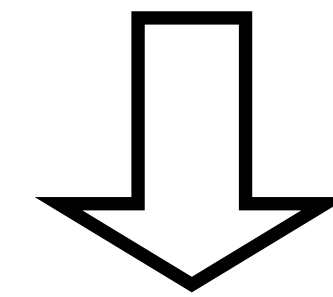
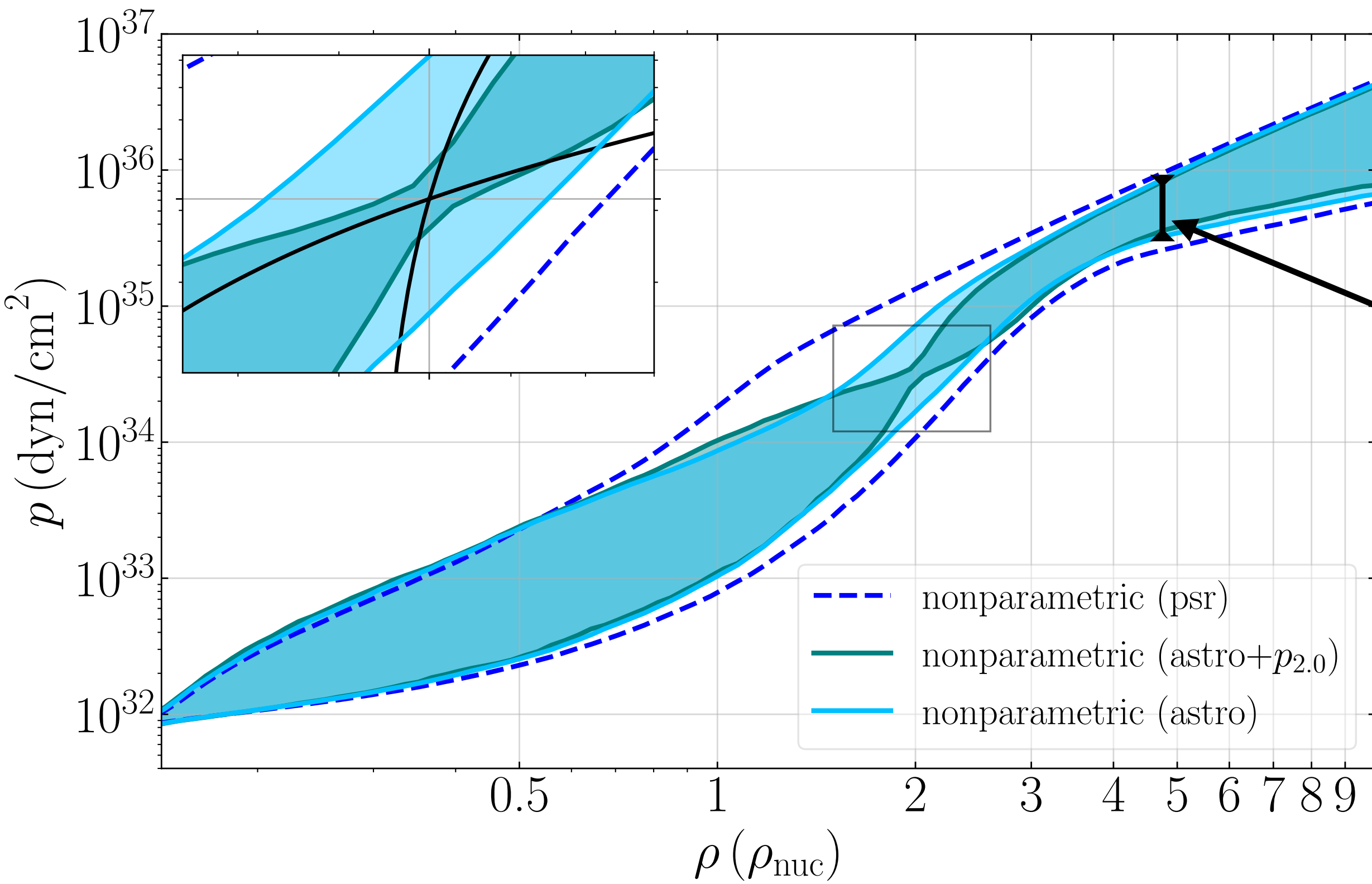
Difference in knowledge about p_b after learning p_a

Changing this analogous to adding a tight Pressure “mock-measurement”

Implicit Correlations

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$$I(p_a, p_b) = \int dp_a P(p_a) \int dp_b P(p_b | p_a) \ln \left(\frac{P(p_b | p_a)}{P(p_b)} \right)$$



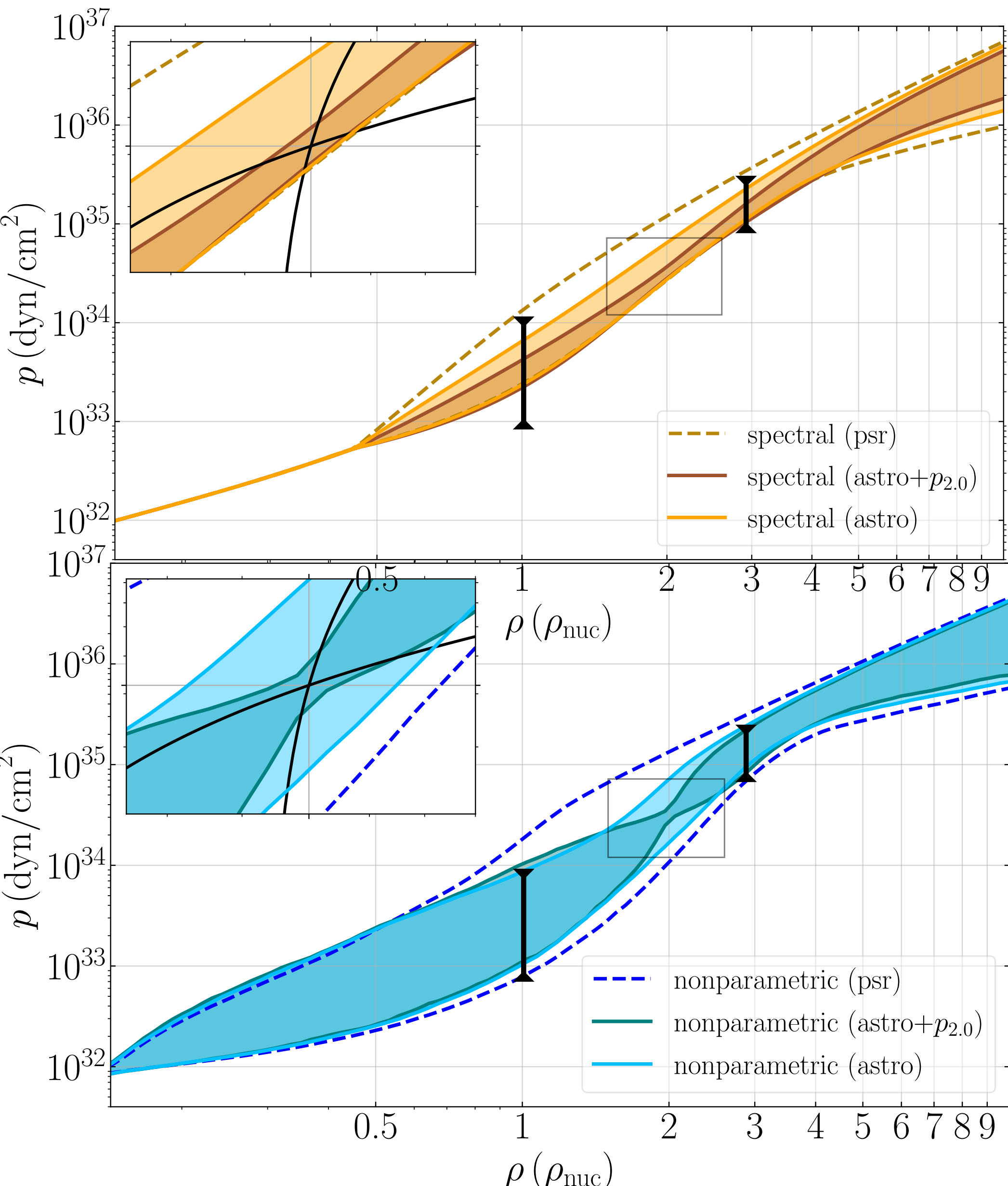
Caveats!

Scales with overall uncertainty of marginal distributions

Want to keep I small even with large entropy in
Marginal distributions $P(p_a), \dots$

Implicit Correlations

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$$I(p_a, p_b) = \int dp_a P(p_a) \int dp_b P(p_b | p_a) \ln \left(\frac{P(p_b | p_a)}{P(p_b)} \right)$$

↓ Caveats!

Scales with overall uncertainty of marginal distributions

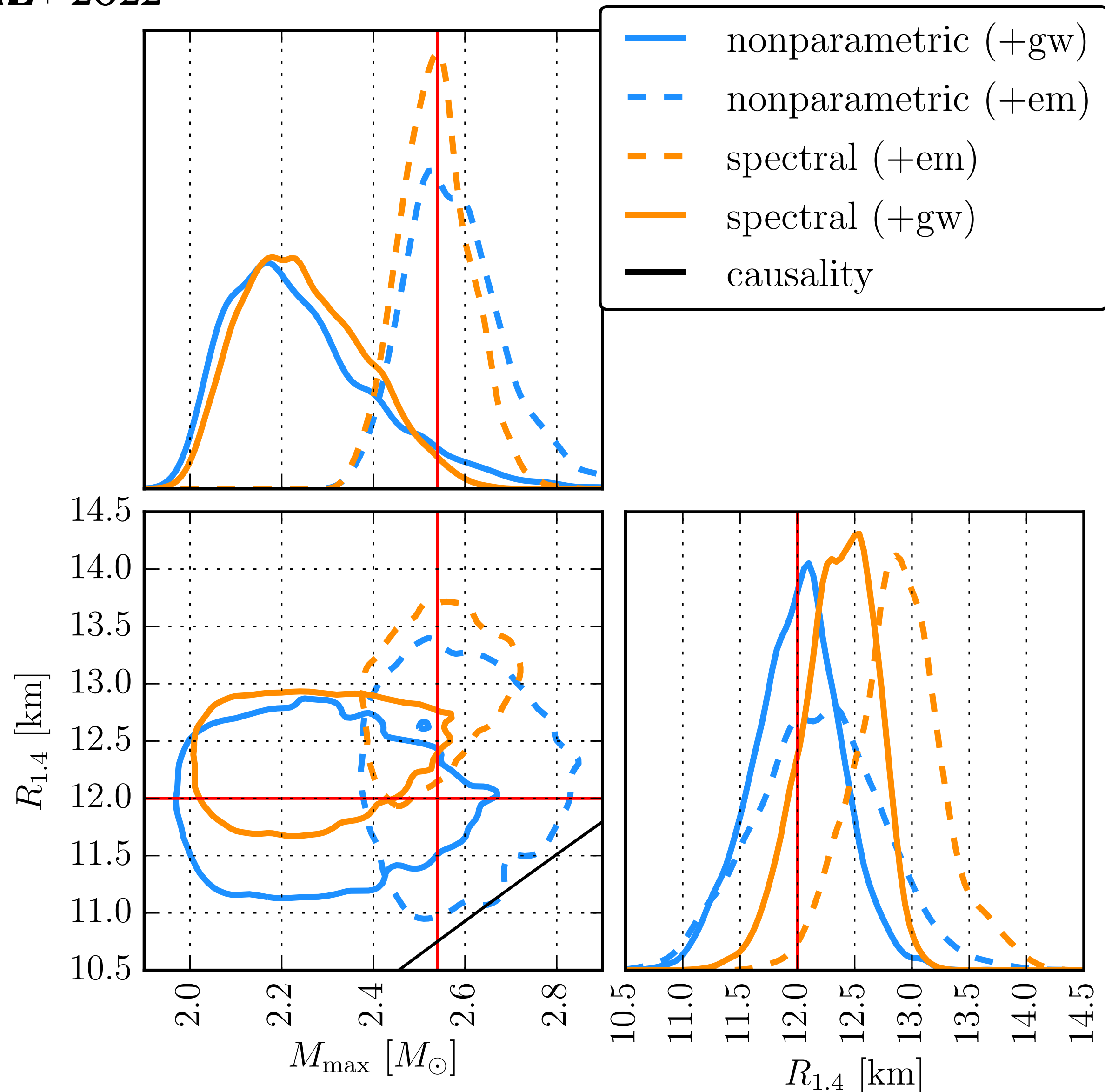
Want to keep I small even with large entropy in Marginal distributions $P(p_a), \dots$

$$I(\ln(p_{1.0}), \ln(p_{1.5}), \ln(p_{2.0}), \ln(p_{3.0}), \ln(p_{4.0}))$$

	PSR	Astro	Astro+p _{2.0}
Nonparametric	3.7	3.1	2.9
Spectral	6.6	5.5	4.7
Polytrope	5.7	4.6	3.8
Speed of sound	5.0	4.7	4.3

Simulated Astrophysical Data

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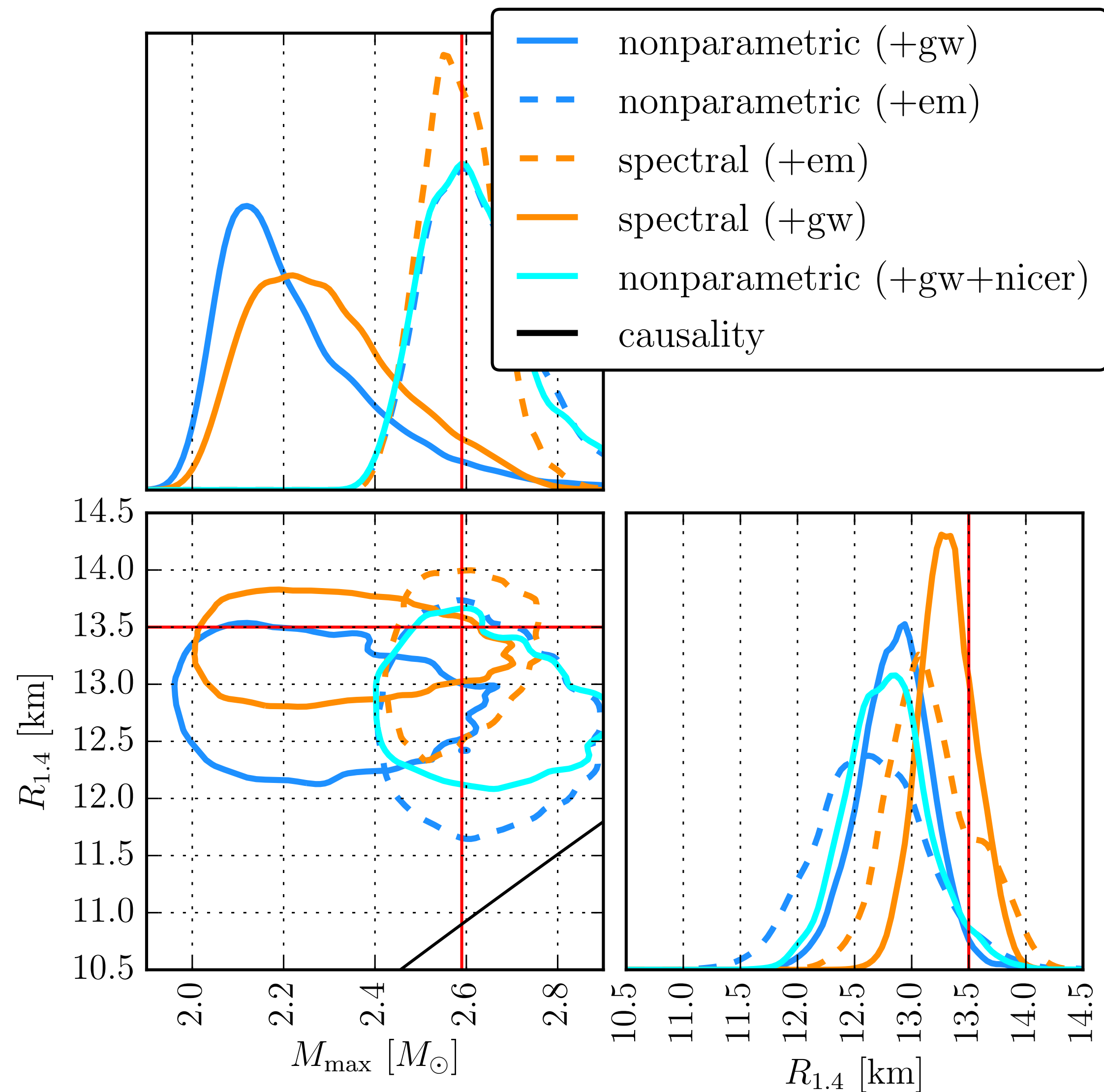


We inject gravitational-wave (gw) and x-ray-radio (em) observations on top of existing constraints

We intentionally choose an **EoS** that we expect the **Spectral** model to fail to recover

Gives a sense of tension that may arise from combining constraints using models with unphysical correlations

Simulated Astrophysical Data



We inject gravitational-wave (gw) and x-ray-radio (em) observations on top of existing constraints

Inverse Problem : Spectral Eos \rightarrow NP analysis

Slow convergence, but no bias

Modified parametric priors

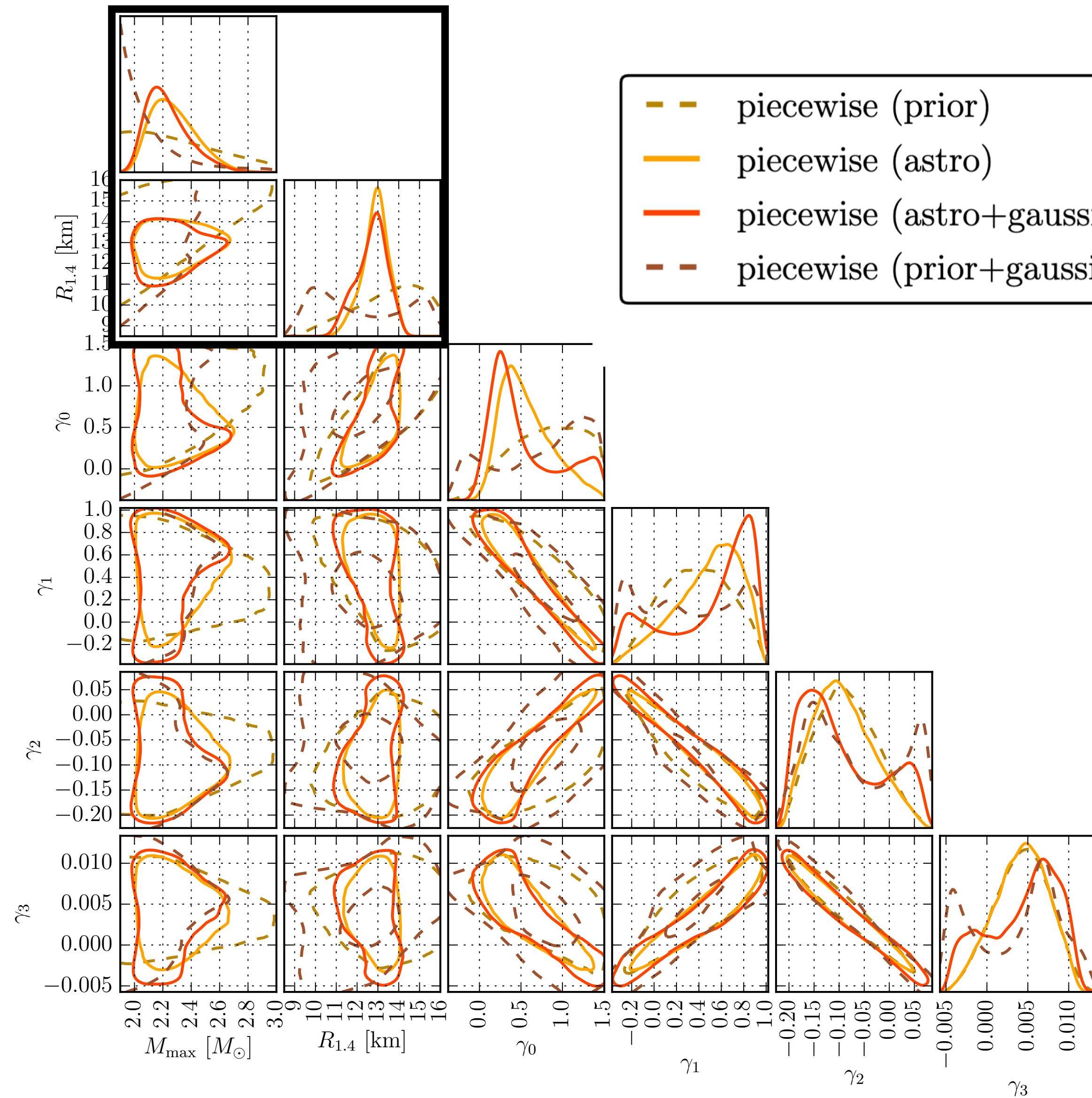
Why not just modify the parametric models to get more flexibility?

E.g.

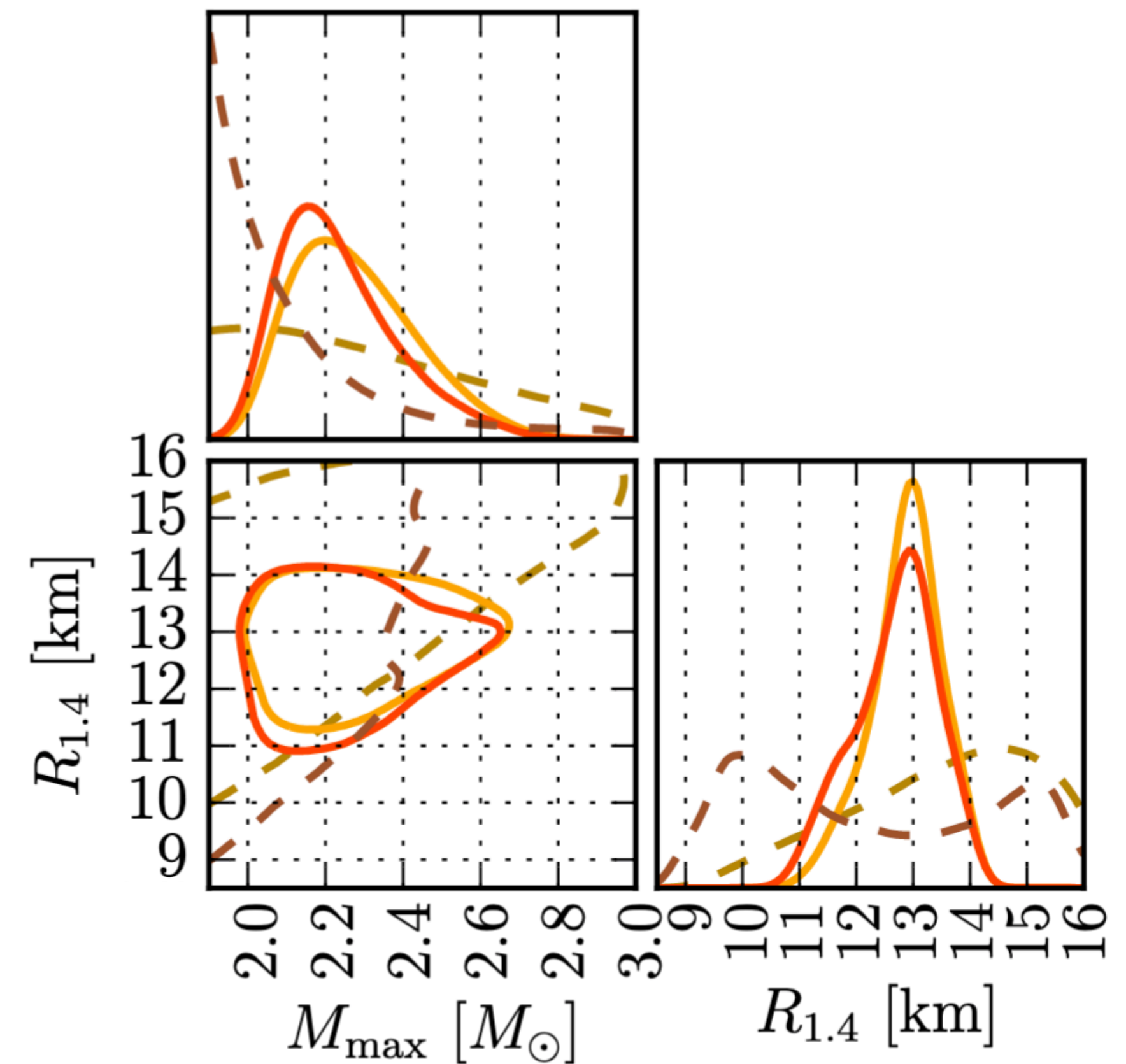
$$p(\rho) = \rho^\Gamma; \quad \Gamma(p) = \sum_{i=0}^3 \gamma_i \log(p/p_0)^i + \text{more terms}$$

Modified parametric priors

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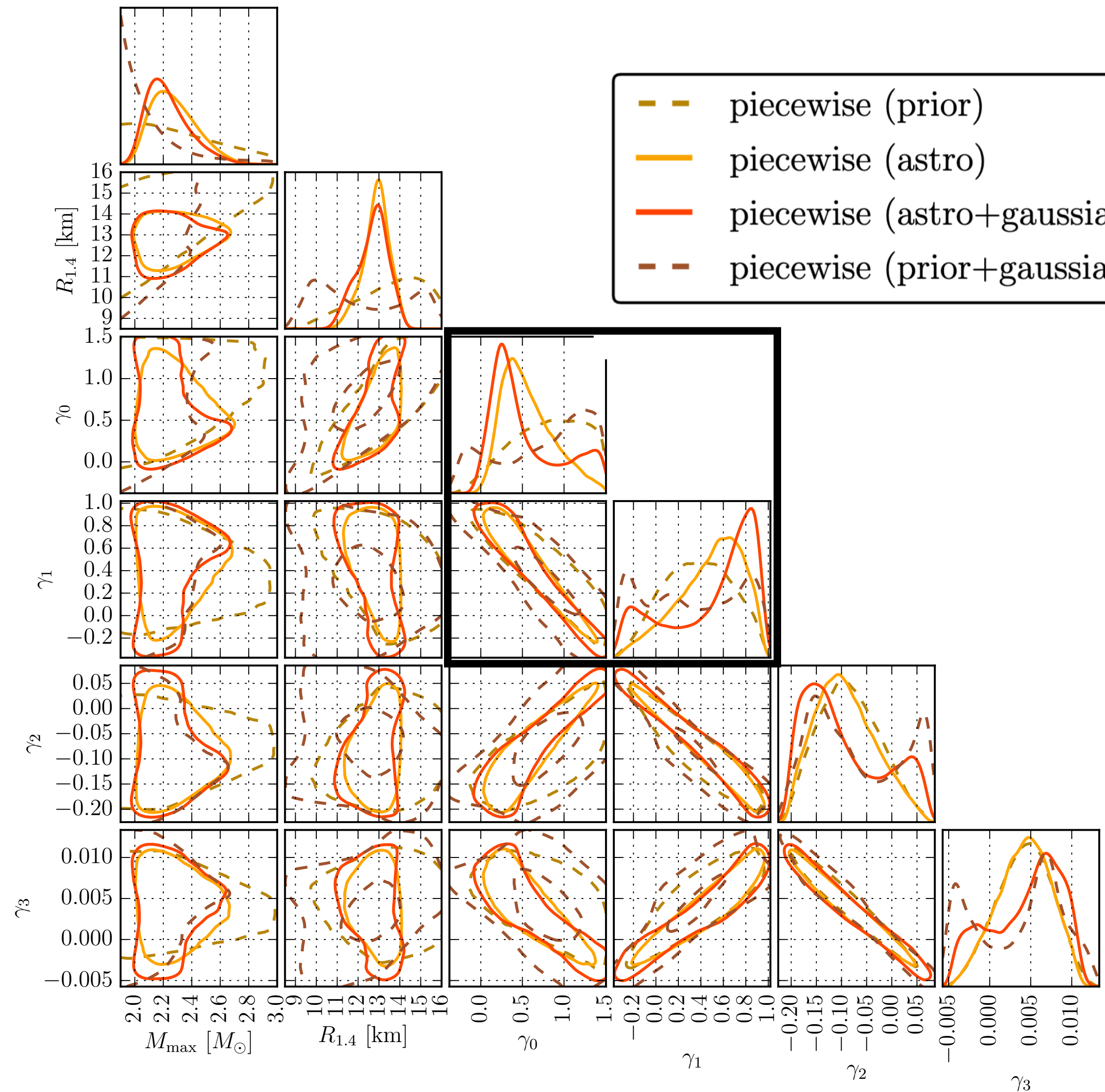


We find changing the
Prior on parameters doesn't
Remove the correlations

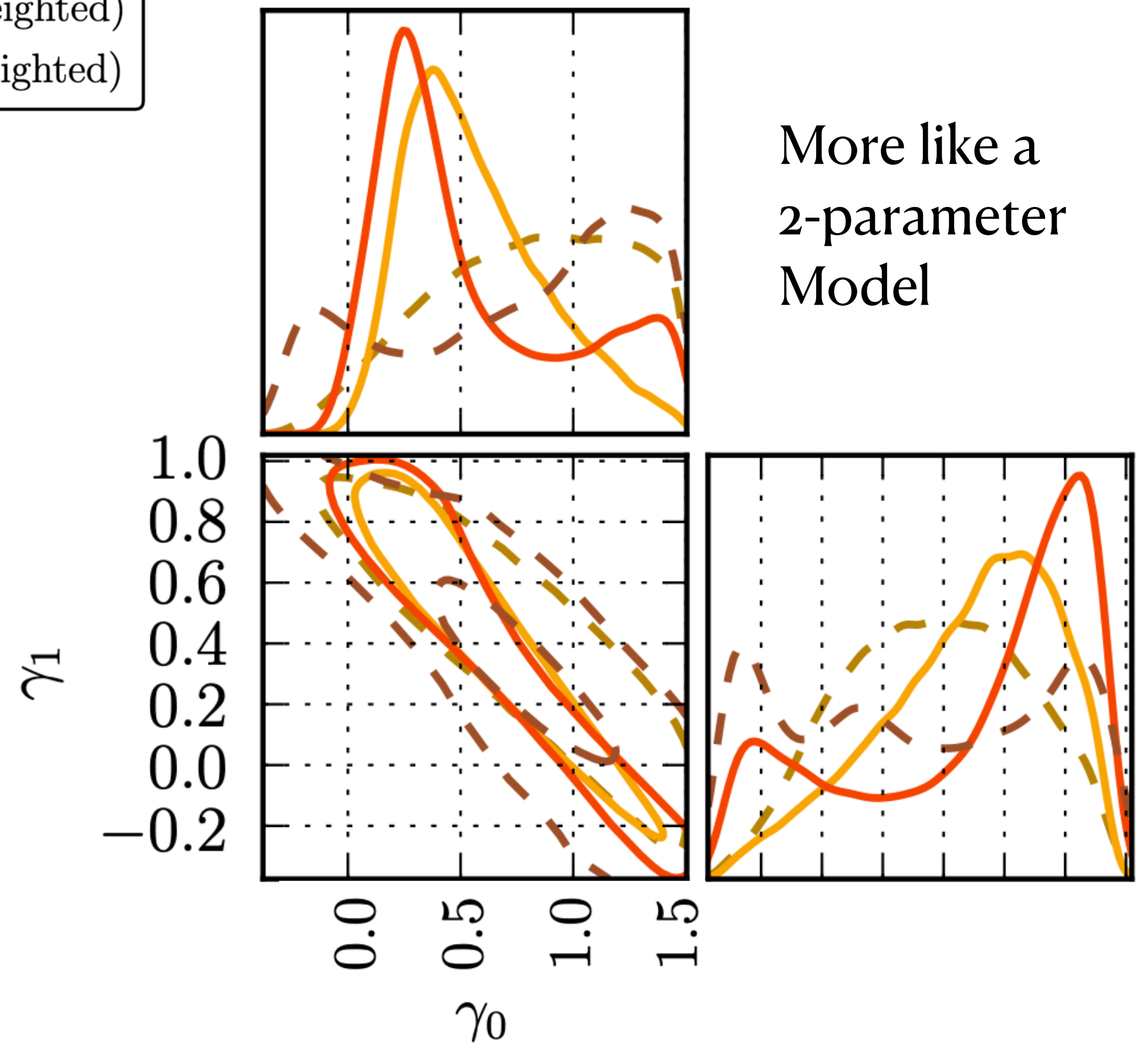


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Modified parametric priors



With the current data we only
Really only infer γ_0 & γ_1



Modified parametric priors

Why not just modify the parametric models to get more flexibility?

Models are either

- (1) fine-tuned => extending them without breaking is difficult (spectral + speed of sound)
- (2) Need overhaul-type improvements (piecewise-polytrope + speed of sound)

This is already being done!

i.e. Steiner+ 2016 -> better piecewise-polytrope models

But... Extensions are nontrivial.

Best to understand limitations of each model while using it

Not all Correlations are Bad!

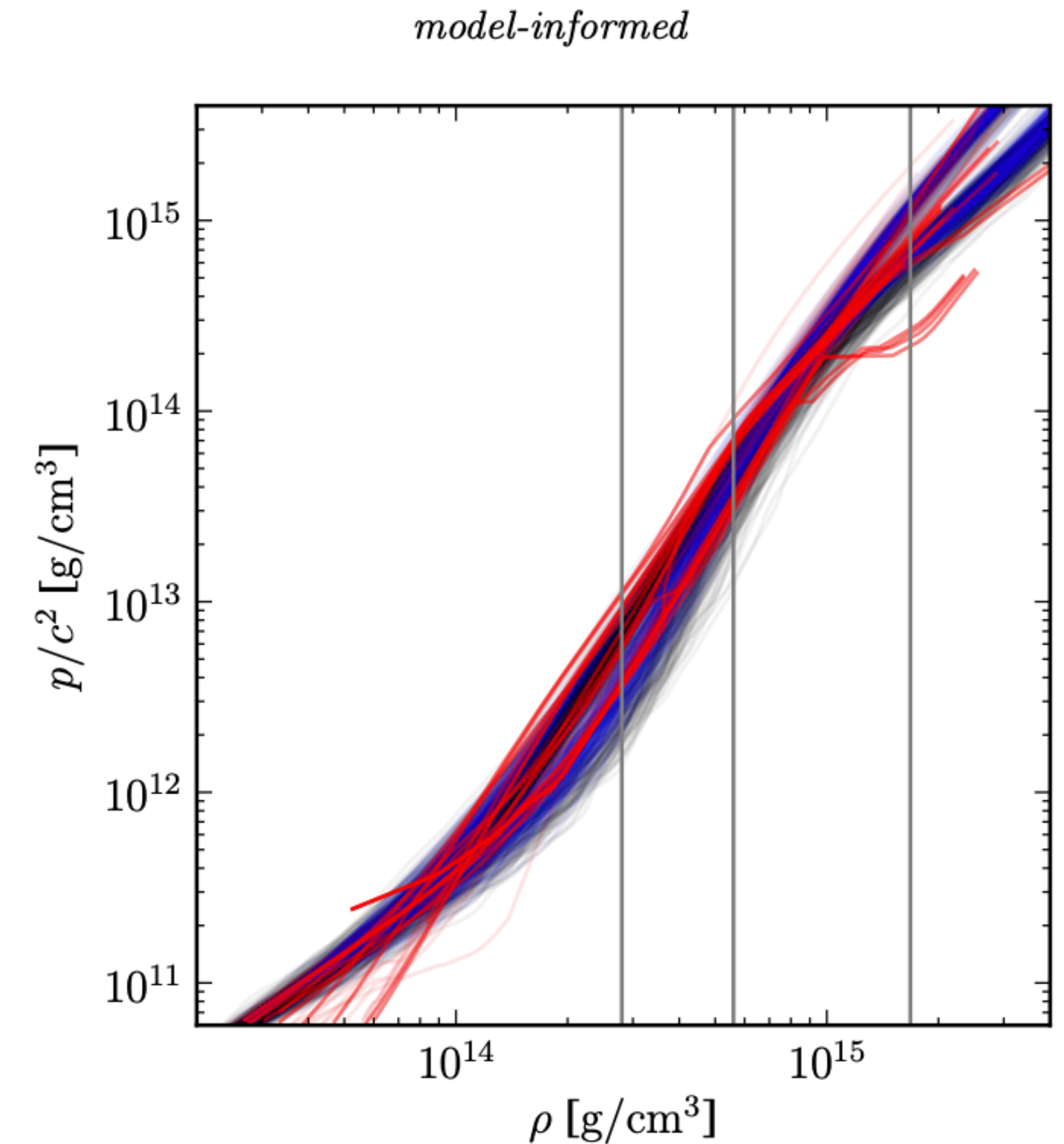
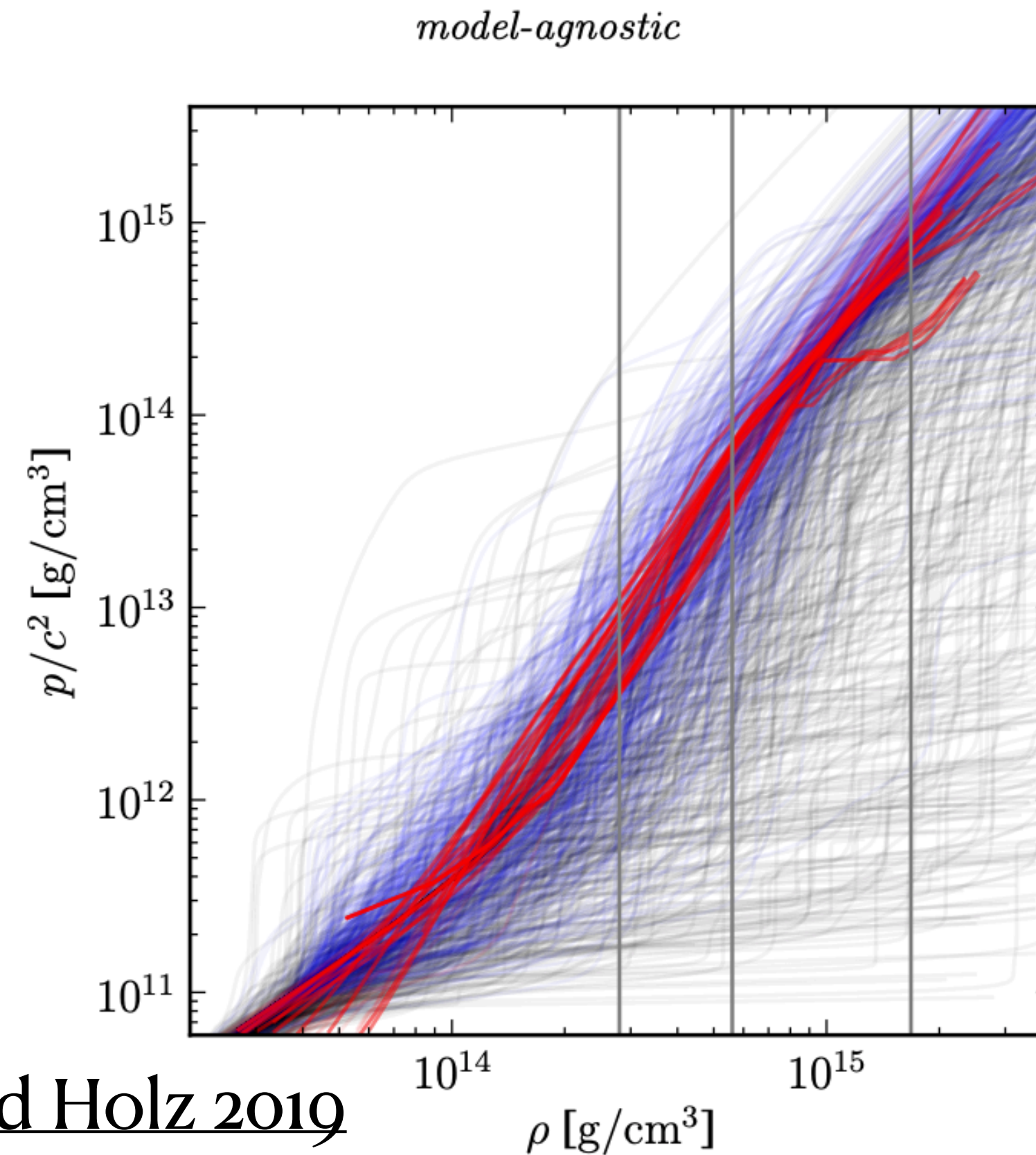
Physical theories have correlations between quantities “ $F=ma$ ”

Correlated

Goal is to give flexibility in the choice of correlations

See e.g. Miller+2021 : GP with
“tight” correlations

Eventually one should *infer* the
correlations



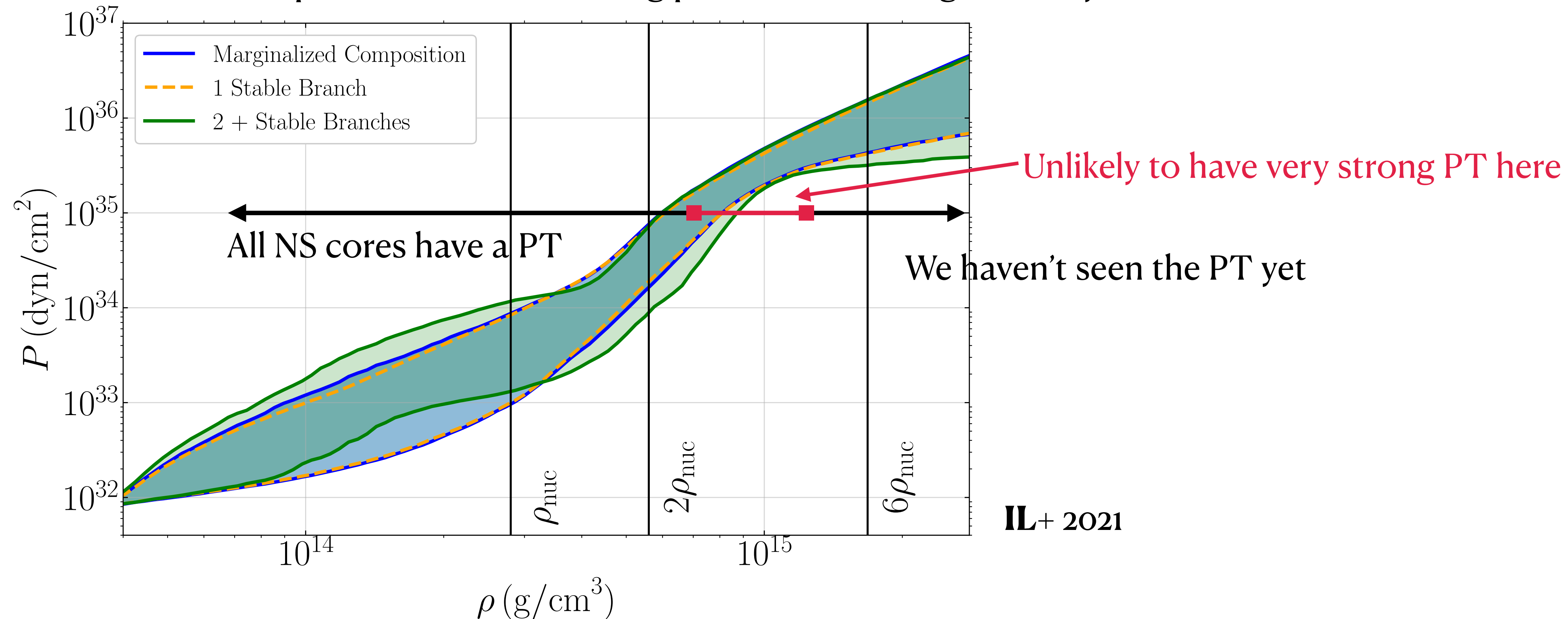
Essick, Landry, and Holz 2019

(Backup) : Strong Phase Transitions

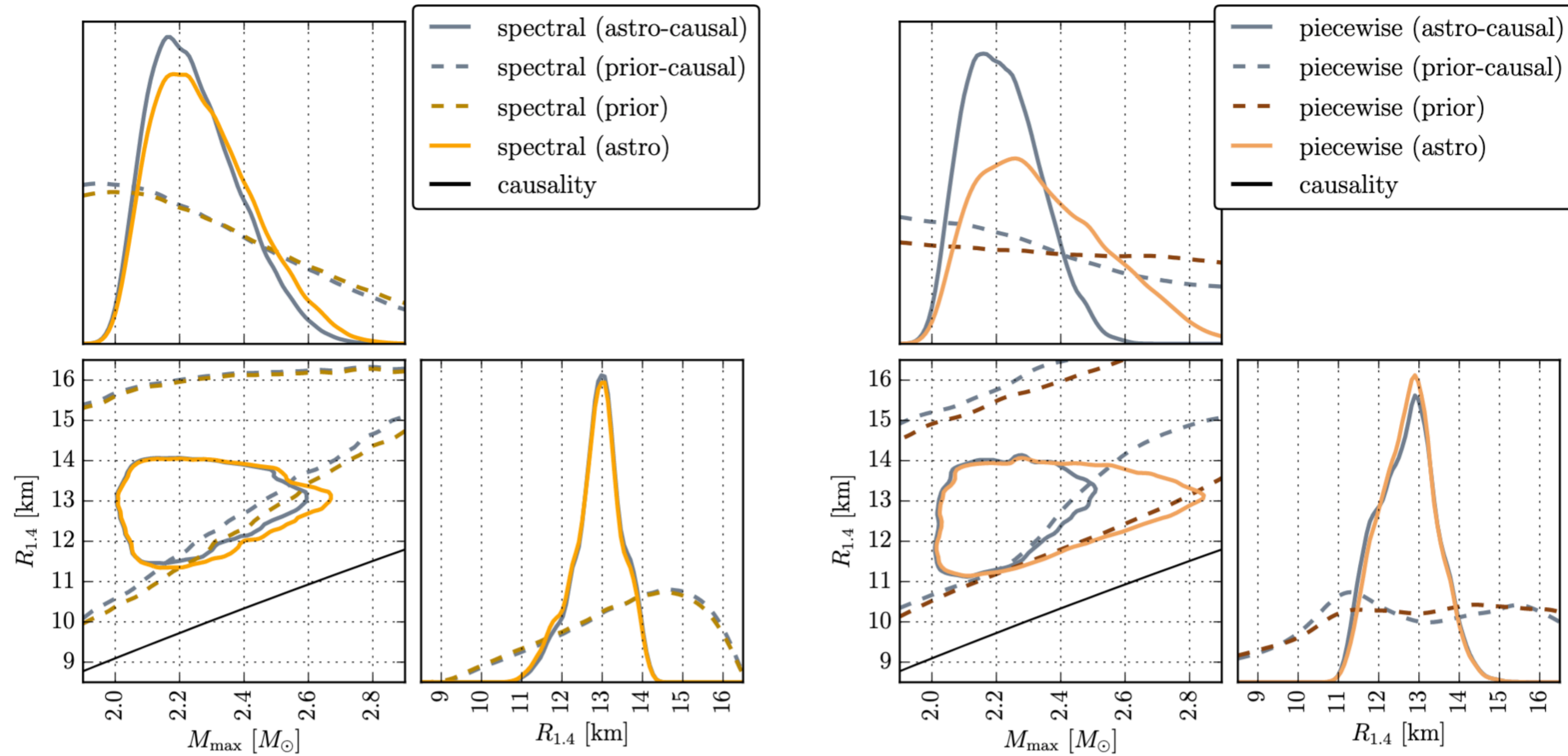
Parametric models struggle to model phase transitions

Piecewise-polytope models with variable stitching densities may be able to -> need fine tuning

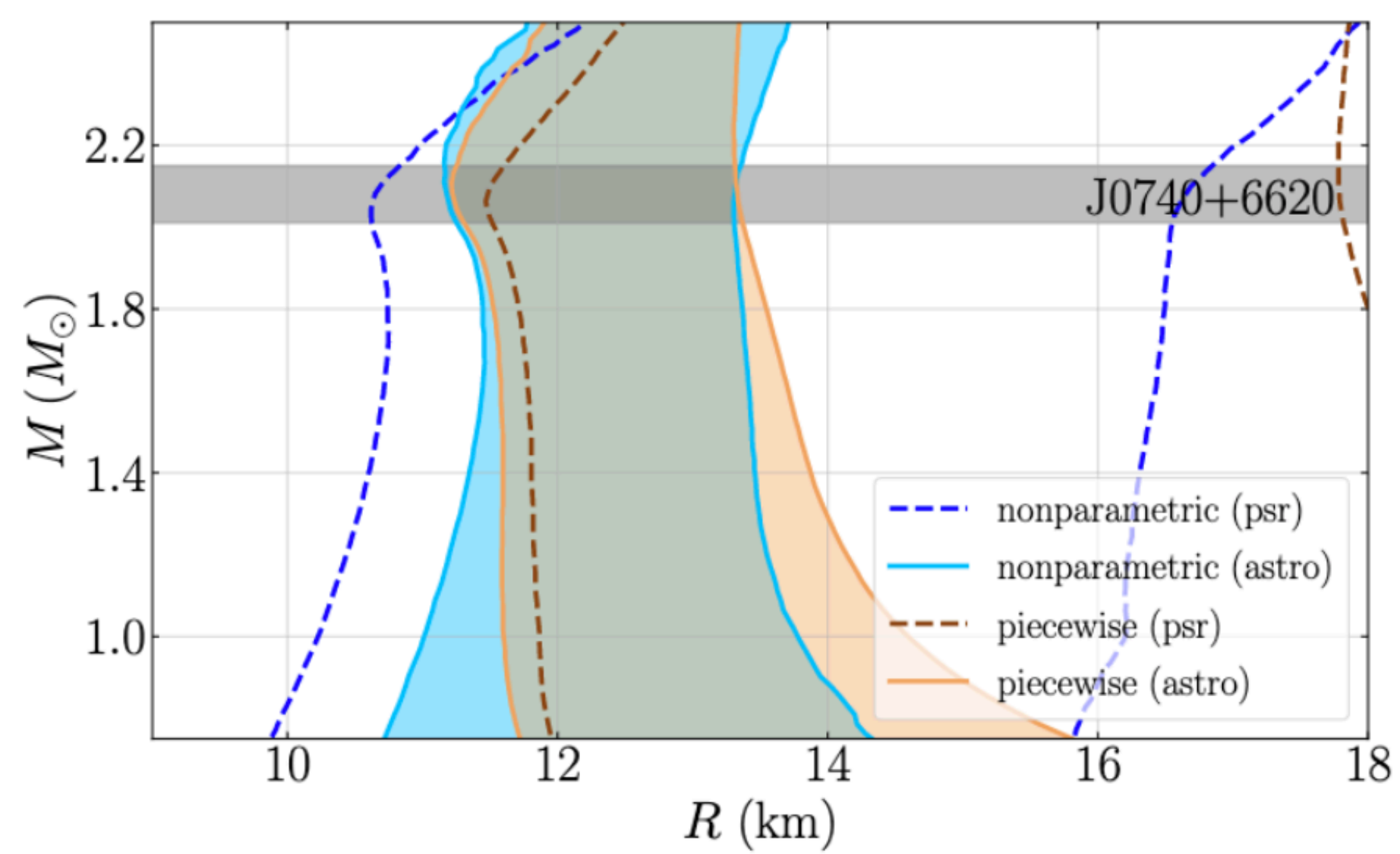
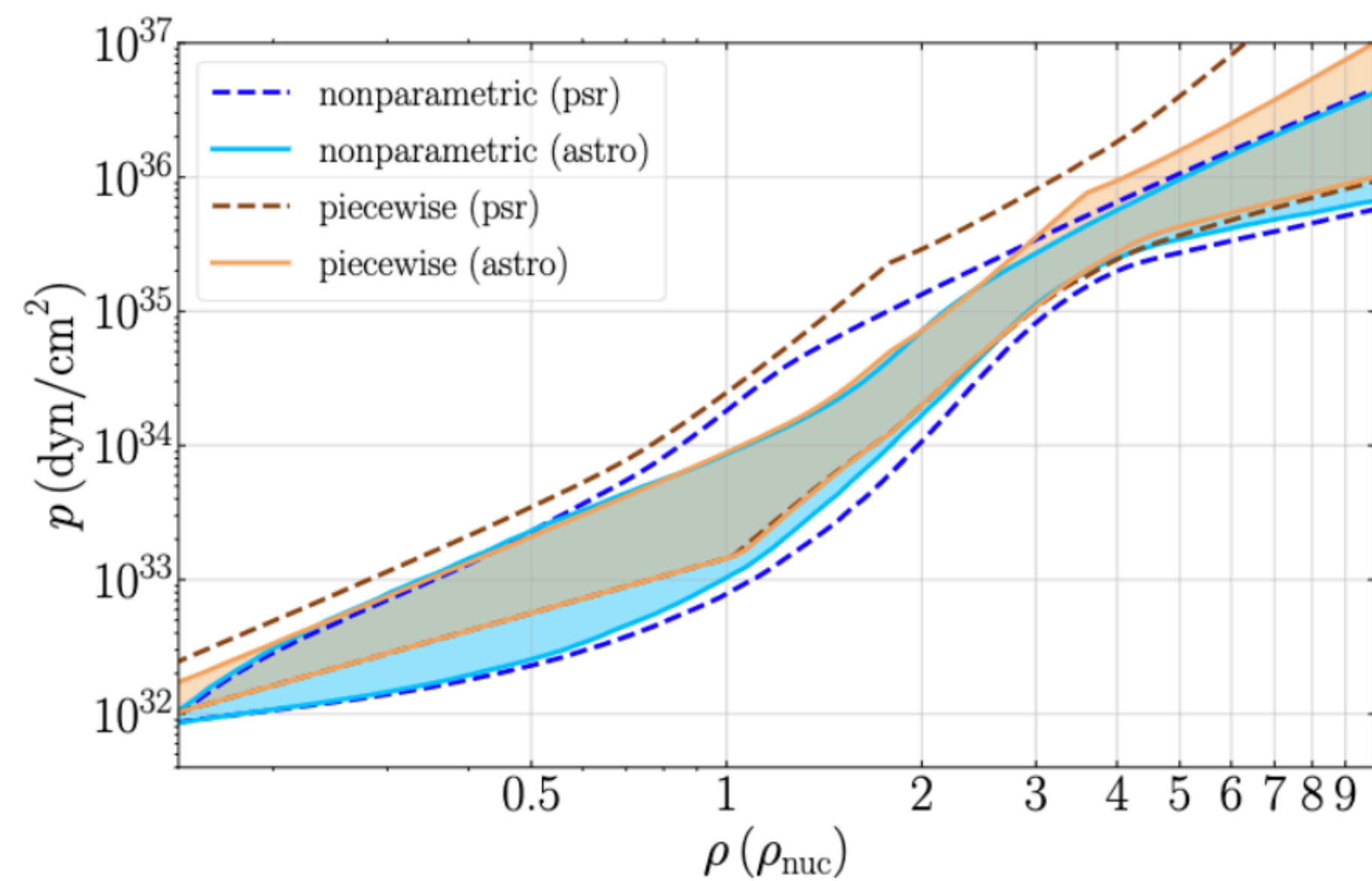
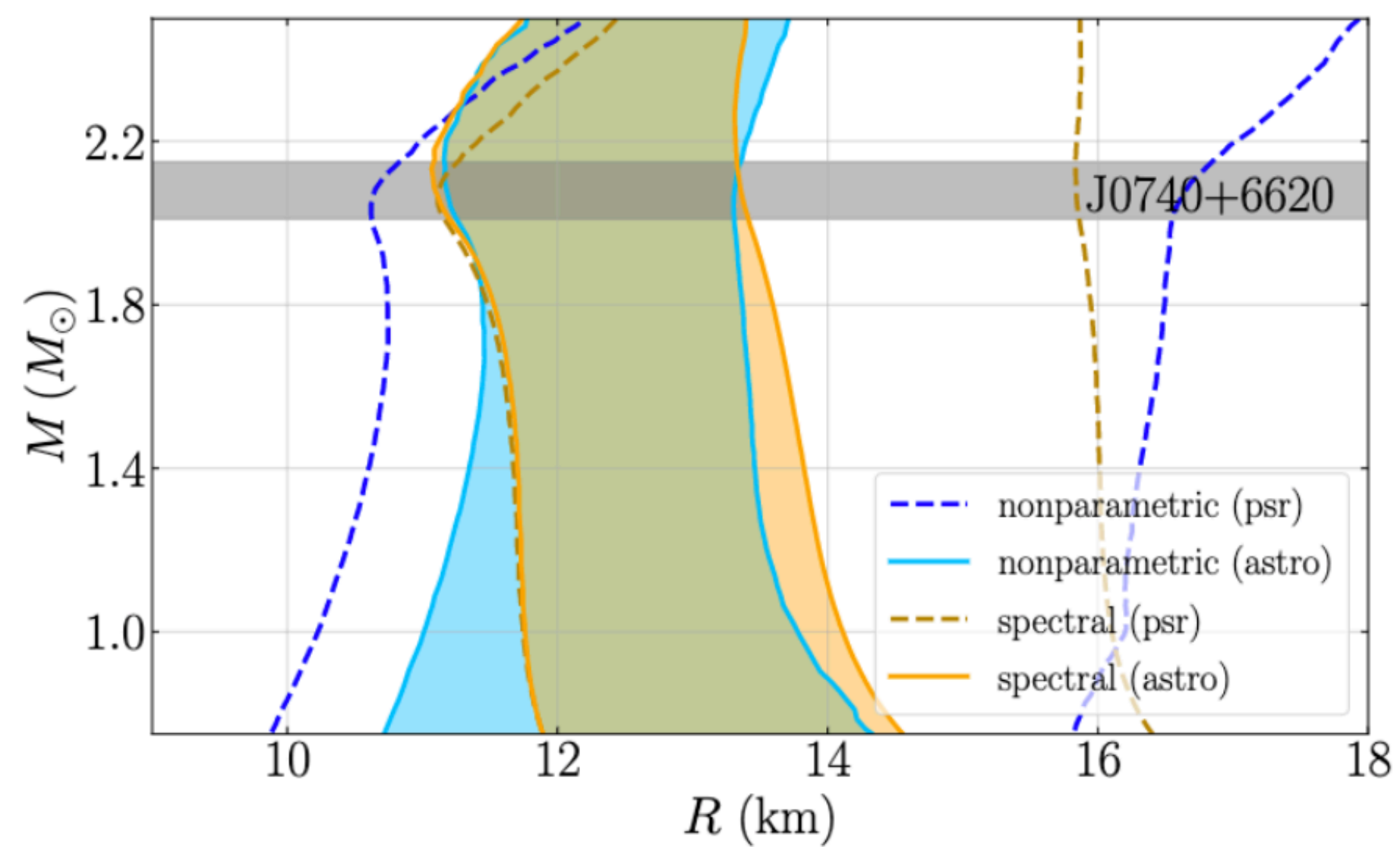
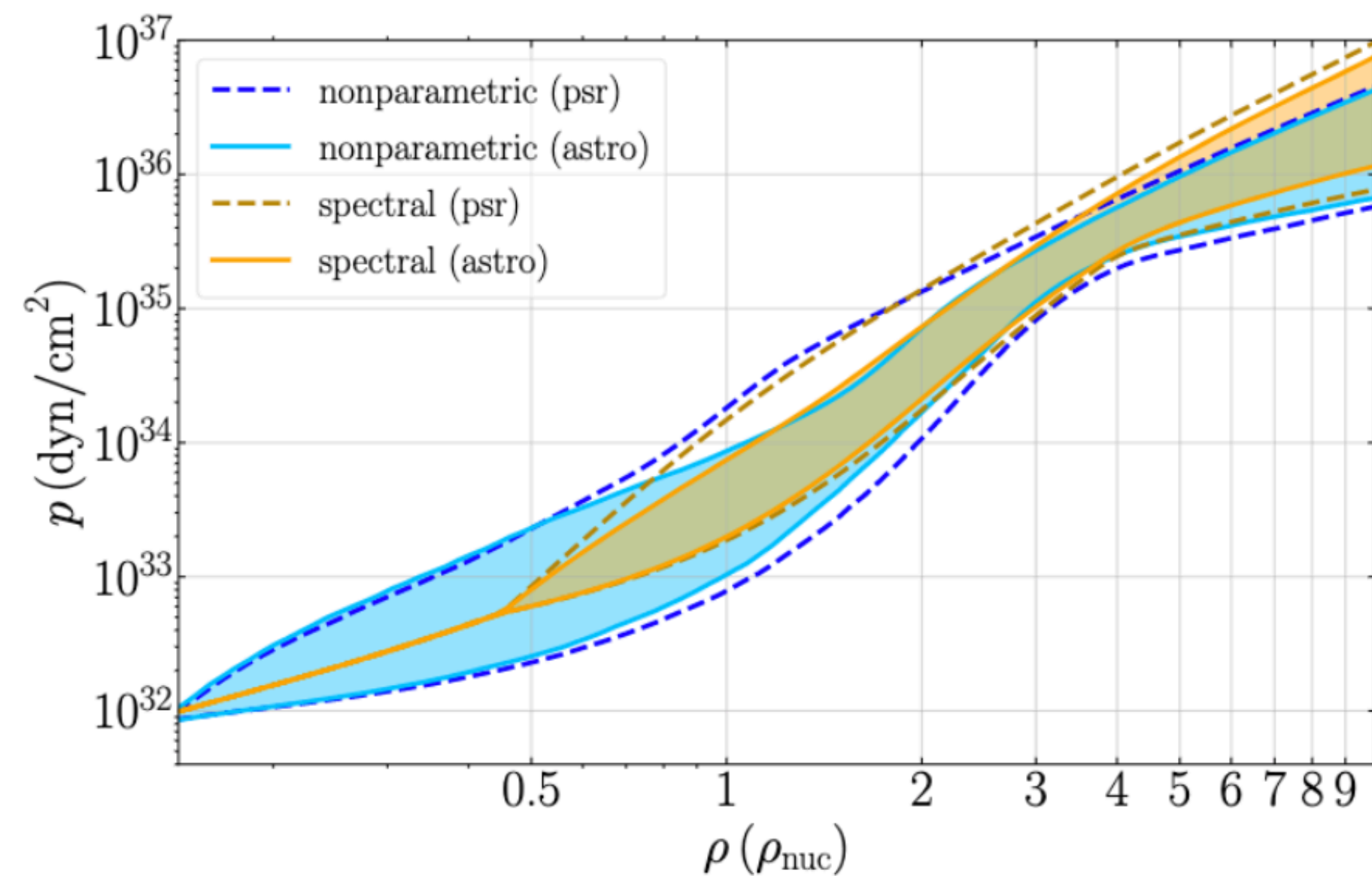
The GP model can produce EoSs mimicking phase transitions generically



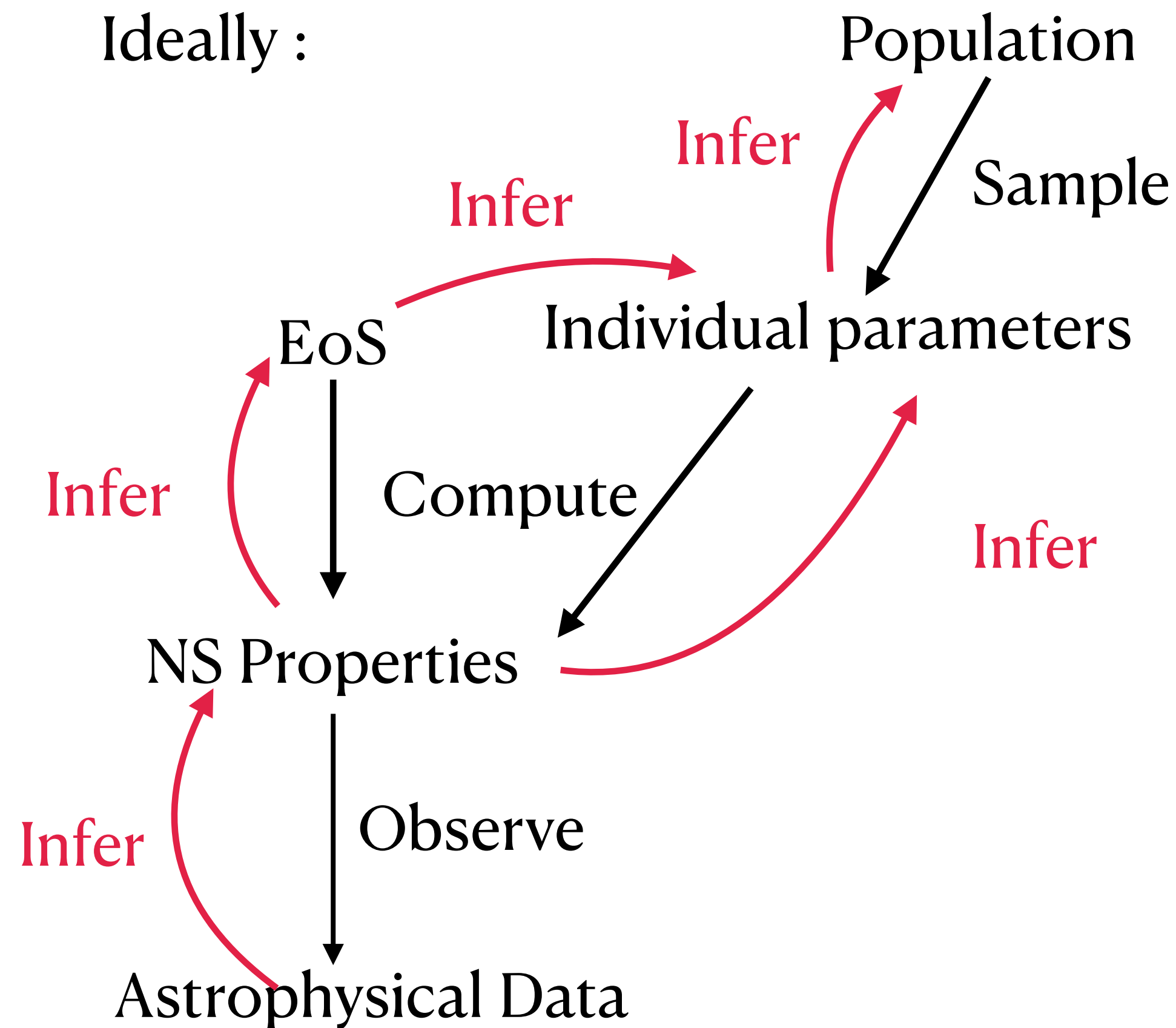
(Backup): Causality in Parametric Models



(Backup): More Parametric Results



Astrophysical Data (Brief Aside)



Lots of details

- (1) Selection Effects
- (2) Interpreting data (GW waveforms, x-ray pulse profiles)
- (3) Poorly characterized population of NSs

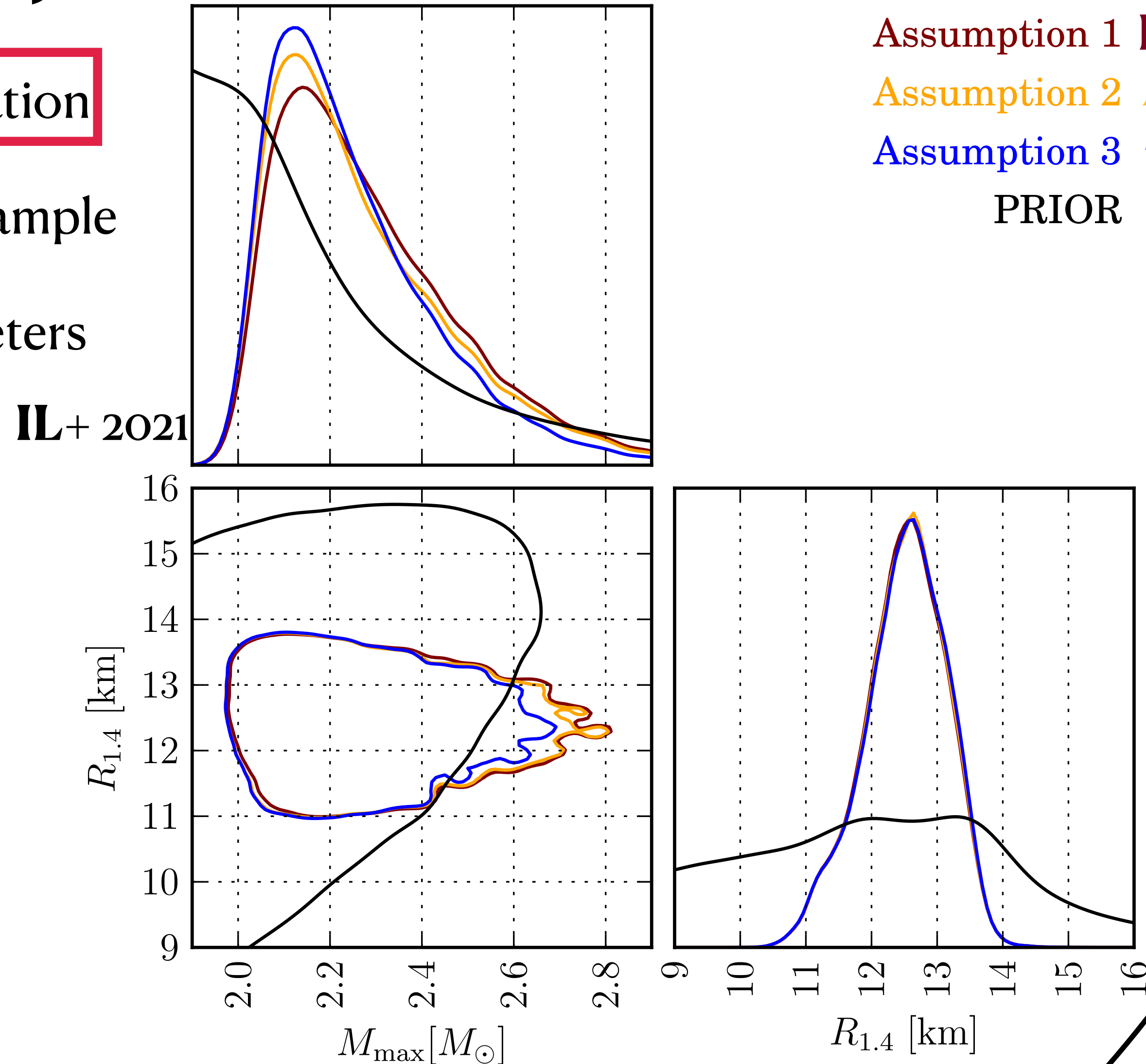
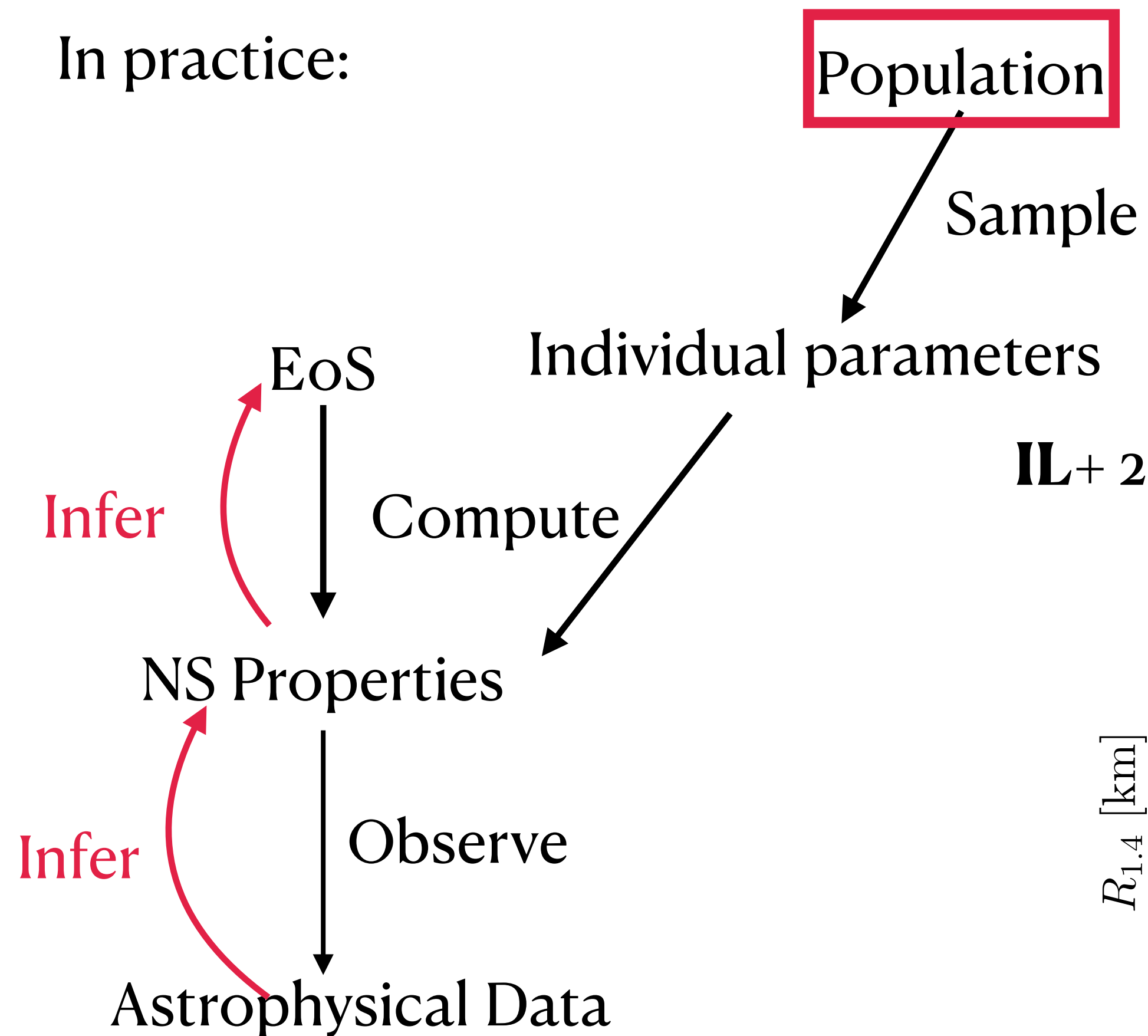
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$$P(d_1, d_2, \dots | \varepsilon_i) = P(d_1 | \varepsilon_i) \times P(d_2 | \varepsilon_i)) \dots$$

$$P(d_1 | \varepsilon_i) = \sum_{\text{NS sources}} P(d_1 | \text{NS source}) \times P(\text{NS Source} | \varepsilon_i, (\text{Population Model}))$$

Astrophysical Data (Brief Aside)

In practice:



$$P(d_1 | \varepsilon_i) = \sum_{\text{NS sources}} P(d_1 | \text{NS source}) \times P(\text{NS Source} | \varepsilon_i, (\text{Population Model}))$$

PRIOR

PSR+GW+J0030

PSR+GW+J0030+J0740

(Backup): Corner Plot

