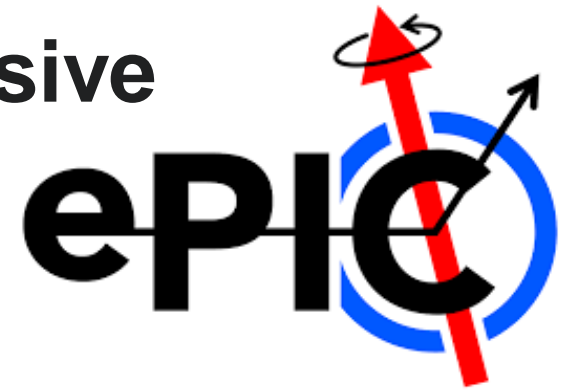
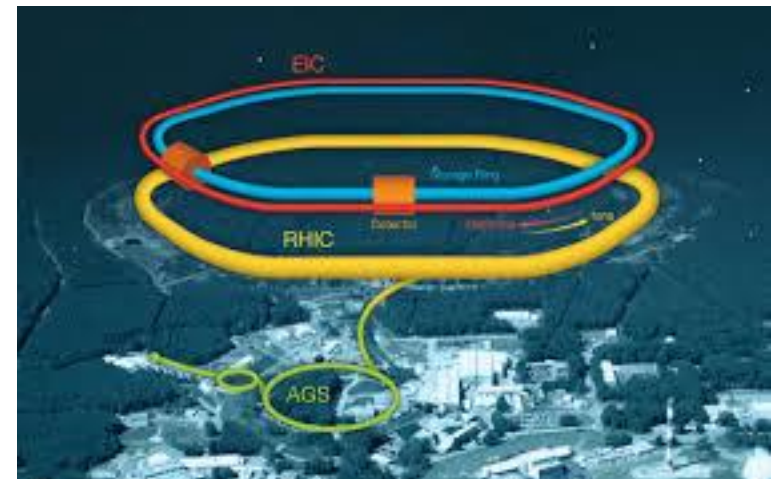
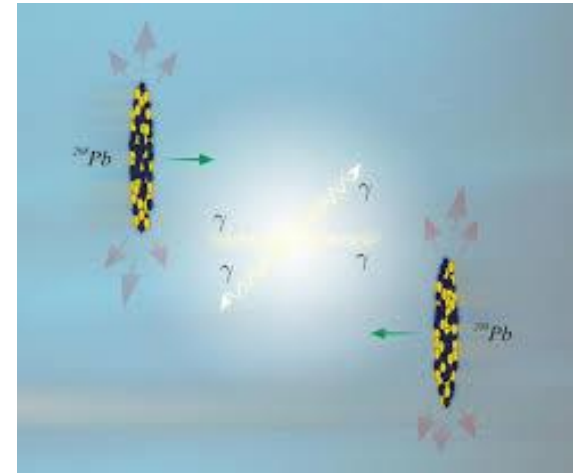


# Connections between an exclusive program at EIC and UPCs



Spencer Klein, LBNL for  
the ePIC Collaboration

- Ultra-peripheral collisions
- Exclusive interactions at the EIC
  - The ePIC detector (forward)
  - Exclusive interactions in ePIC
- Cross-sections
- $d\sigma/dt$  and gluon spatial distributions
- Reggeons and exotica
- Backward production
- UPCs in the EIC era
- Conclusions



# Ultra-peripheral collisions (UPCs)

- Heavy nuclei carry strong electric and magnetic fields
  - Fields are perpendicular -> nearly-real virtual photon field
    - $E_{\max} = \gamma hc/b$
  - Photonuclear interactions
    - Two-photon interactions also occur, but less relevant here
- Most visible when  $b > \sim 2R_A$ , so there are no hadronic interactions;
  - We also see coherent  $J/\psi$  photoproduction in peripheral nuclear collisions

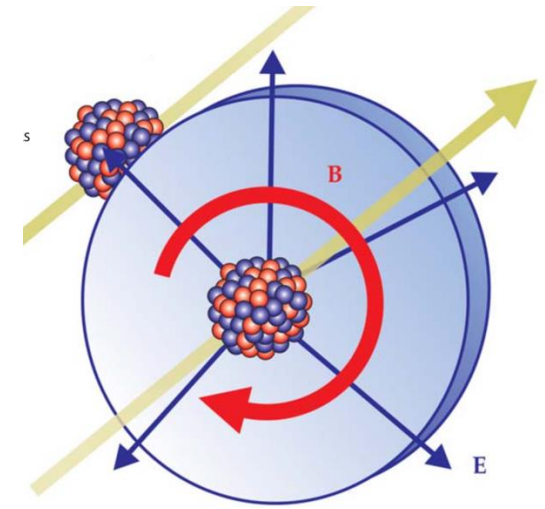
Energy	AuAu RHIC	pp RHIC	PbPb LHC	pp LHC
Photon energy (target frame)	0.6 TeV	~12 TeV	500 TeV	~5,000 TeV
CM Energy $W_{\gamma p}$	24 GeV	~80 GeV	700 GeV	~3000 GeV
Max $\gamma\gamma$ Energy	6 GeV	~100 GeV	200 GeV	~1400 GeV

\*LHC at full energy  $\sqrt{s}=14$  TeV/5.6 TeV

The energy frontier for photon physics!

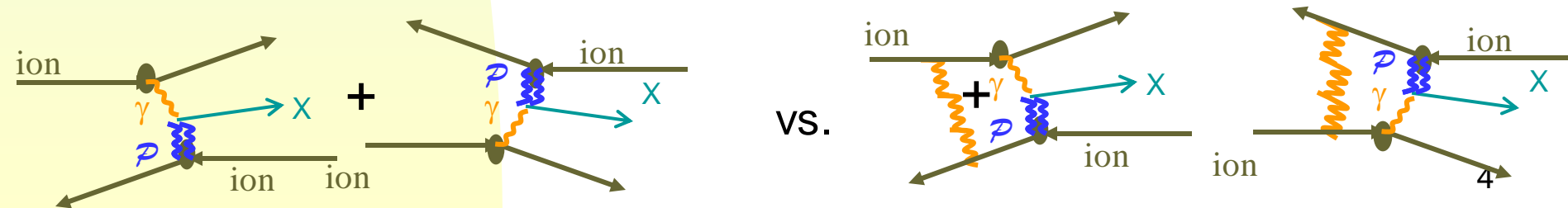
# UPCs – good and bad

- The energy frontier for electromagnetic probes
  - Maximum CM energy  $W_{\gamma p} \sim 3 \text{ TeV}$  for pp at the LHC
    - $\sim 10$  times higher than HERA
  - Probe parton distributions in proton and heavy-ions down to
    - Bjorken-x down to a few  $10^{-6}$  at moderate  $Q^2$
- Electromagnetic probes have  $\alpha_{EM} \sim 1/137$ , so are less affected by multiple interactions than hadronic interactions
  - Exclusive interactions
- Bidirectional photon beams
- $Z\alpha \sim 0.6$  for lead  $\rightarrow$  multiple interactions with a single ion pair.
  - E. g. vector meson production + nuclear excitation or 2 vector mesons
  - Useful for tagging the impact parameter vector, but we cannot select pure single-photon exchange events



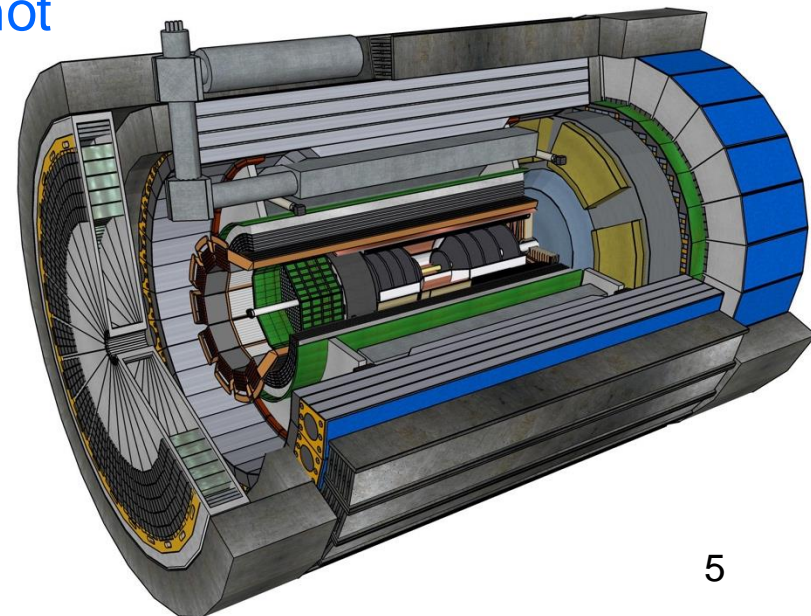
# Bidirectional photon beams

- In pp/AA collisions, either nucleus can emit the photon
  - In pA, photon usually comes from the heavy nucleus
- In coherent reactions, the 2 possibilities are indistinguishable, so amplitudes add, and interfere destructively
  - $\sigma \rightarrow 0$  as  $p_T \rightarrow 0$  at  $y=0$
- 2 directions have different photon energies and Bjorken-x:
  - $k = M_V/2 \exp(\pm y)$  and  $x m_p = M_V/2 \gamma_{\text{beam}} m_p \exp(\mp y)$
- To find  $\sigma(k)$  requires selecting events with different photon spectra
  - Additional photons  $\rightarrow$  Different impact-parameter distributions
  - Events with and w/o nuclear excitation
  - Systems of linear equations  $\rightarrow$  solvable, at a cost in uncertainty



# The electron-ion collider & ePIC

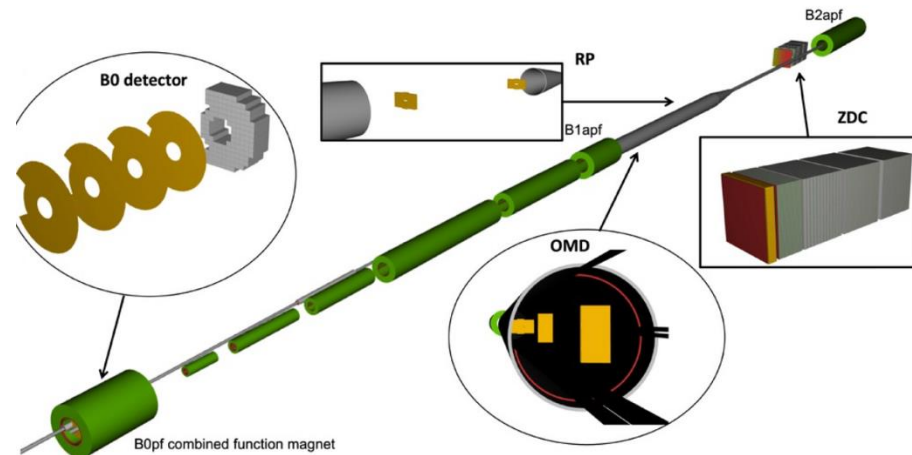
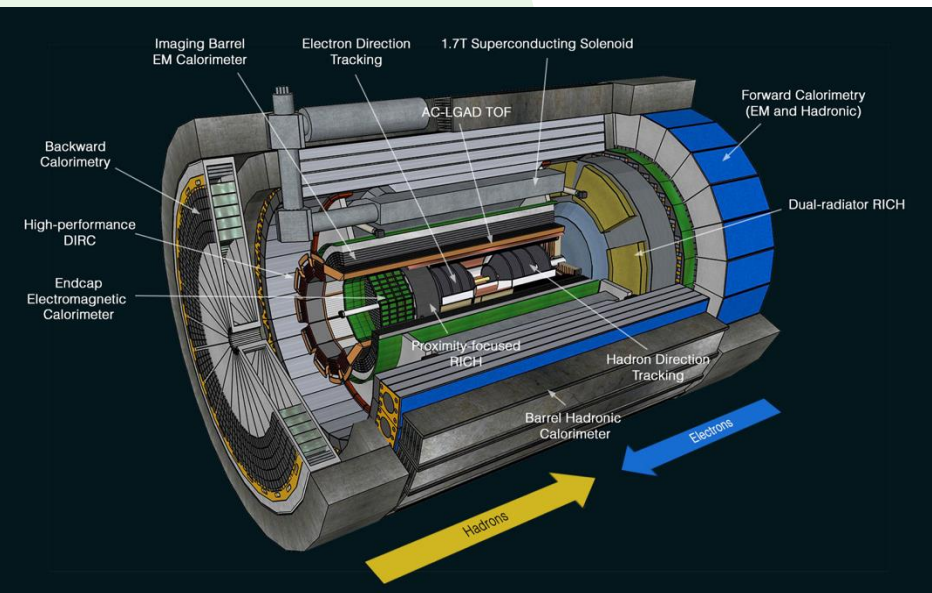
- High luminosity  $ep/eA$  collisions
- Photons with a wide range of virtuality
  - Observe scattered electron to determine photon energy and  $Q^2$
- Detector optimized for  $\gamma^*p/\gamma^*A$  collisions
  - Near  $4\pi$  acceptance
  - Good forward instrumentation to determine if nucleus dissociated or not
- Precision measurements down to Bjorken- $x \sim 10^{-4}$ 
  - Less energy reach than UPCs at the EIC, but more precision





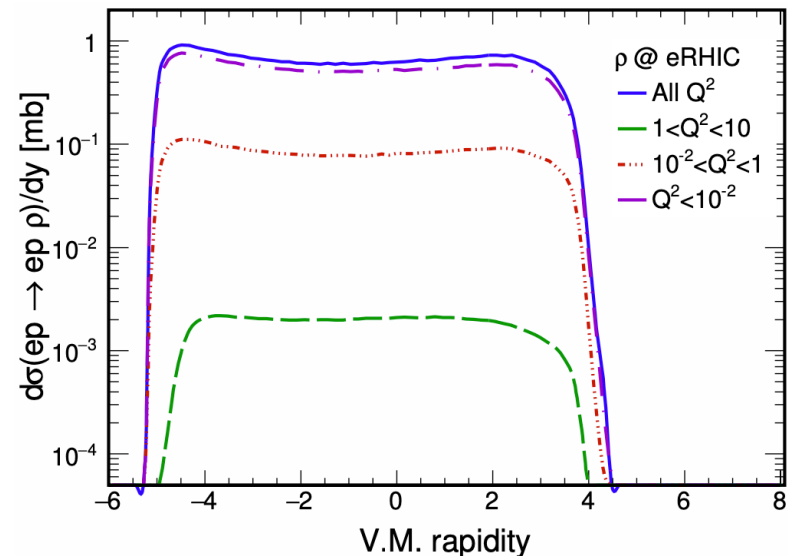
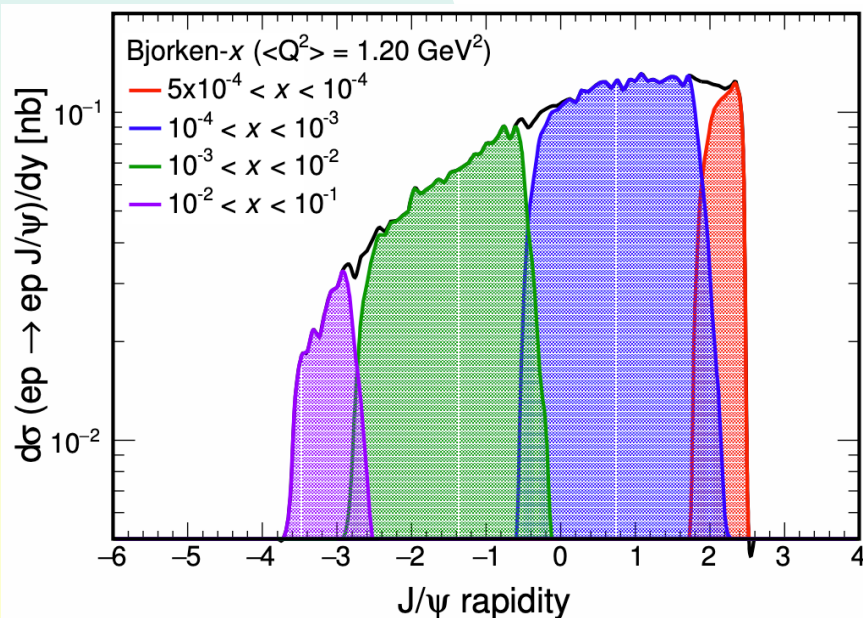
# The ePIC detector

- The central region ( $|y| < 4$ ) See Olga Evdokimov's talk
- Low  $Q^2$  electron tagger determine photon  $E$ ,  $Q^2$
- Forward detectors
  - B0 tracker & calorimeter ( $4.6 < \eta < 5.9$ )
  - Roman pots and Off-Momentum Detector detect scattered protons
  - Zero Degree Calorimeter for photons and neutrons
- Big forward question: did the nucleus break up, or not?



# Energy and rapidity

- For exclusive interactions, energy and rapidity are related
  - Photon energy  $K=M_X/2 \exp(y)$
  - Bjorken- $x$ :  $x=M_X/\gamma M_p \exp(-y)$
- Wide energy coverage requires a wide rapidity range
  - For vector mesons, need  $\sim +1$  unit of pseudorapidity coverage to cover a given rapidity range.



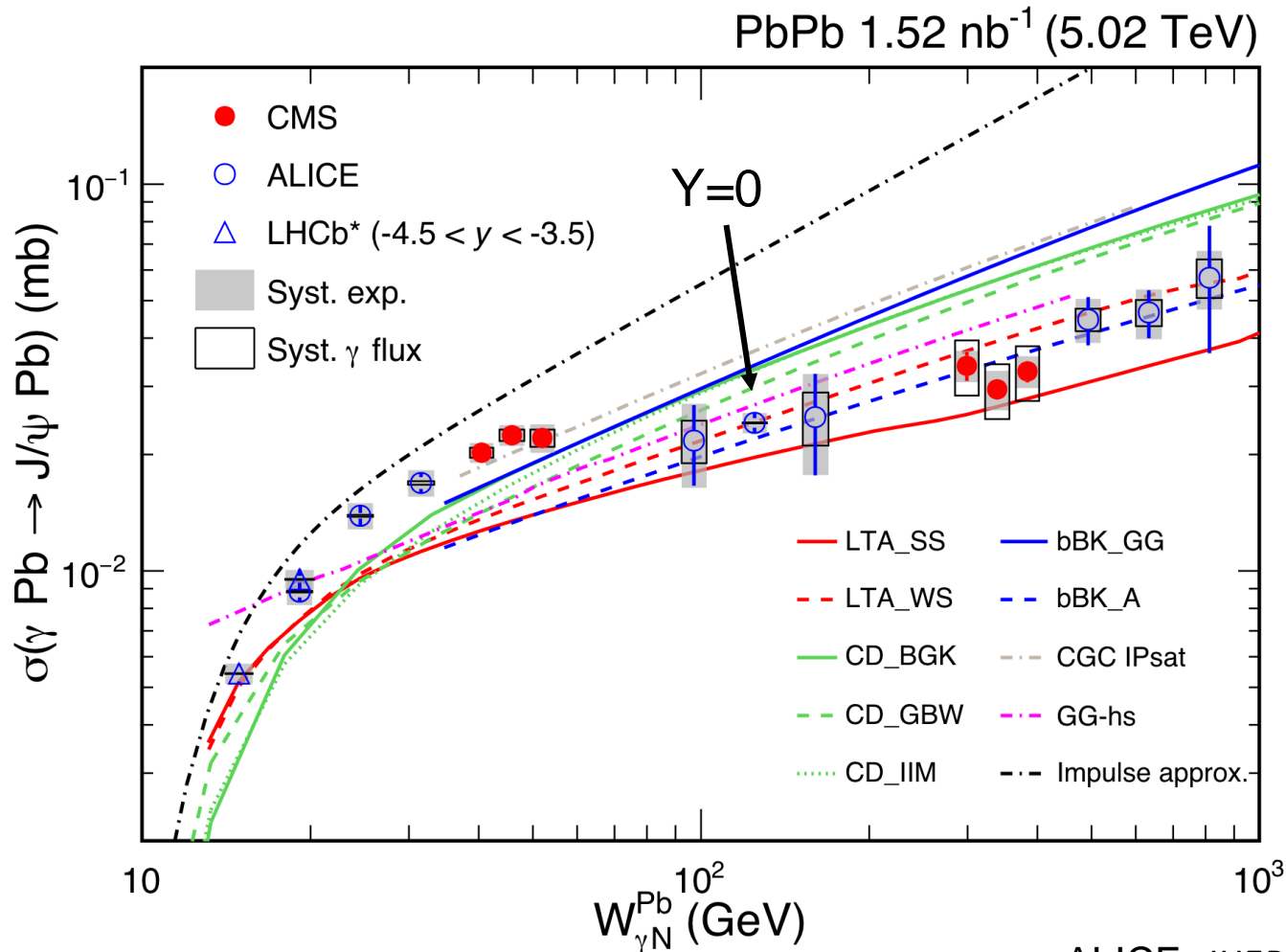
# Energy and Rapidity in UPCs

- AuAu/PbPb collisions are symmetric  $\rightarrow$  either nucleus can emit the photon  $\rightarrow$  bidirectional ambiguity
  - Photon energy  $K = M_x/2 \exp(\pm y)$
  - Bjorken- $x$ :  $x = M_x/\gamma M_p \exp(\mp y)$
- Total amplitude is sum of both directions. Away from  $y=0$ ,  $p_T=0$ , interference is small  $\rightarrow$  can directly use cross-sections.
- The cross-section at a given  $y \neq 0$  is the sum of two directional cross-sections, with different energies.
- The solution is to use measurements with two different photon spectra, so different energies, i. e. with two different cross-section ratios
  - Two different impact parameter distributions



# J/ψ cross-sections vs. energy

- $\sigma \sim W^\delta$  continues up to  $W_{\gamma p} \sim 1$  TeV
- Some wiggles -> tension between analyses?



# Coherent and Incoherent Photoproduction: a quantum view

- The Good-Walker formalism links coherent and incoherent production to the average nuclear configuration and event-by-event fluctuations respectively
  - Configuration = position of nucleons, gluonic hot spots etc.
- Coherent: Nucleus remains in ground state, so sum the amplitudes, then square -> average over different configurations
- Incoherent = Total – coherent; total: square, then sum cross-sections for different configurations

$$\frac{d\sigma_{\text{tot}}}{dt} = \frac{1}{16\pi} \left\langle |A(K, \Omega)|^2 \right\rangle \quad \text{Average cross-sections } (\Omega)$$

$$\frac{d\sigma_{\text{coh}}}{dt} = \frac{1}{16\pi} |\langle A(K, \Omega) \rangle|^2 \quad \text{Average amplitudes } (\Omega)$$

$$\frac{d\sigma_{\text{inc}}}{dt} = \frac{1}{16\pi} \left( \left\langle |A(K, \Omega)|^2 \right\rangle - |\langle A(K, \Omega) \rangle|^2 \right) \quad \text{Incoherent is difference}$$

# Coherence in Good-Walker

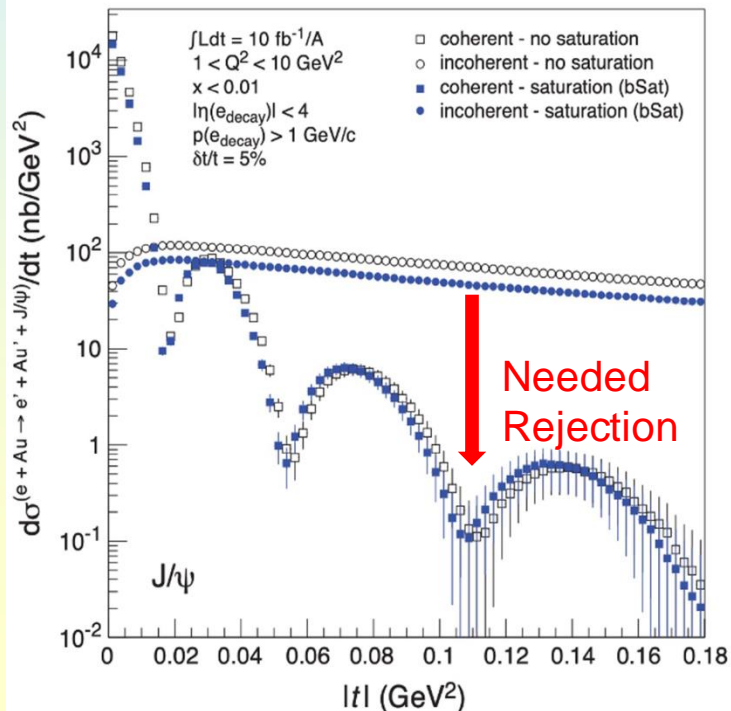
- Coherent production  $\Leftrightarrow$  Target remains in the ground state
  - $\rightarrow d\sigma/dt$  probes transverse distribution of scatterers
- Incoherent production  $\Leftrightarrow$  Target is excited/dissociated
  - Cross-section probes event-by-event target fluctuations
- But... we observe coherent production accompanied by mutual Coulomb excitation, and in peripheral heavy ion collisions
  - Here, coherent  $\Leftrightarrow$  the amplitudes from the nuclei add in-phase
    - $\sigma_{\text{coherent}} = |\sum_i A_i k \exp(ikb)|^2$
  - Something is missing/problematic from Good-Walker
    - How coherent is coherent enough?
      - A soft bremsstrahlung photon can be added to any reaction
- Use caution in interpretation, especially in relating incoherent production to target fluctuations

# $d\sigma/dt$ and the transverse distribution of gluons in protons/nuclei from coherent production

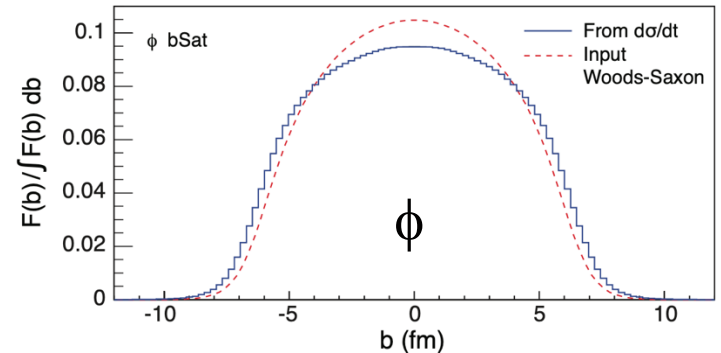
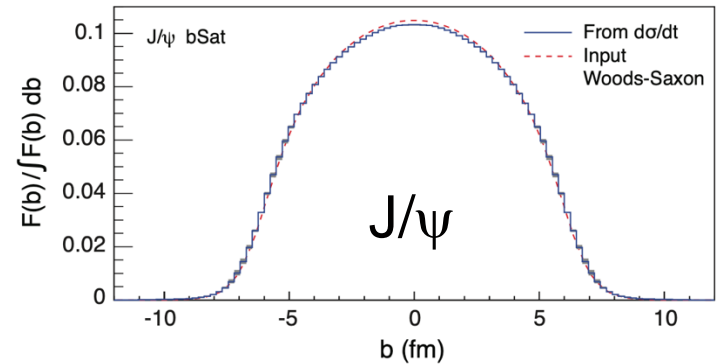
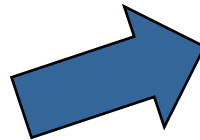
- Position (within nucleus) and  $p_T$  are conjugate variables
  - $F(b)$ , the transverse distribution of scatterers in a target, is the 2-d Fourier transform of  $ds/dp_T$

$$F(b) \propto \frac{1}{2\pi} \int_0^\infty dp_T p_T J_0(bp_T) \sqrt{\frac{d\sigma}{dt}} \quad \text{*must flip sign at each diffractive minimum}$$

- Sensitive to shadowing; major focus of EIC White Paper



Fourier Transform

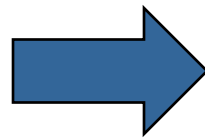
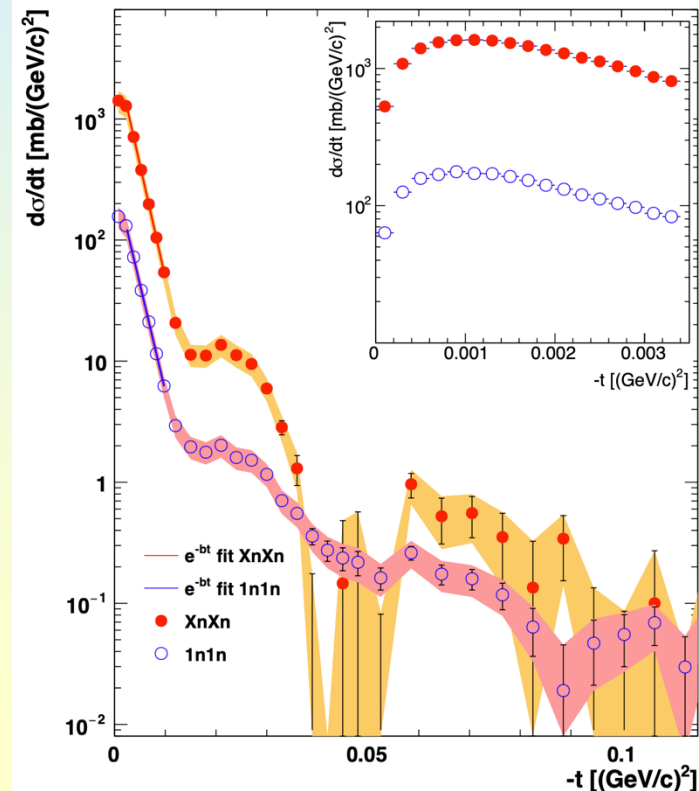


# Difficulties in measuring $d\sigma/dt$

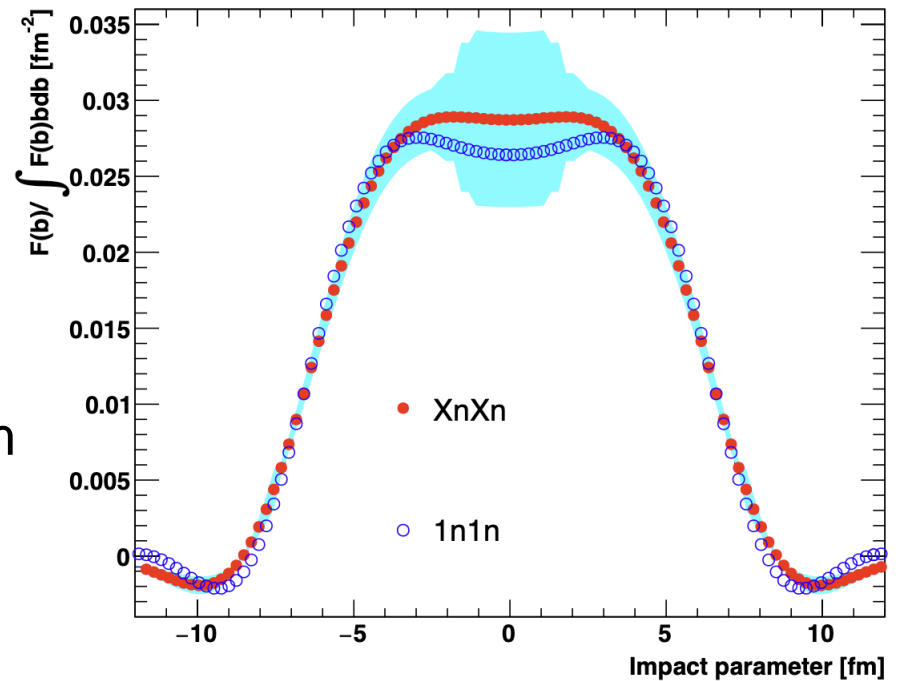
- Resolution fills in the diffractive dips
- In UPCs
  - The photon flux must be removed by deconvolution
  - Limited  $p_T$  reach creates windowing artifacts
    - May be alleviated with ALICE Run 3 data
- At the EIC
  - Resolution is an issue, especially for the electron.
  - Momentum transfer is  $\ll$  electron energy
  - Beam energy spread must be considered
- If the diffractive dips are filled in, they cannot be so well localized, and  $F(b)$  becomes less precise

# STAR transverse distribution measurements

- Fit incoherent contribution at large  $|t|$  and subtract
  - Use a dipole form factor for scattering off a single nucleon
    - Not related to event-by-event fluctuations
- Vector sum of 'Pomeron'  $p_T$ , photon  $p_T$  and resolution
- L



Fourier Transform





# Low-x VM production in eA in ePIC

- More saturation expected for light mesons
  - $\phi$  (light) and  $J/\psi$  (heavy) are featured in EIC studies
    - $\phi$  is particular problem because the  $K^\pm$  daughters are so soft
      - $p=135$  MeV/c in  $\phi$  rest frame;  $\beta \sim 0.2$  so  $dE/dx$  is large
  - Consider  $\rho$  as a replacement
- Usually cannot see outgoing ion
  - Some protons observable in Roman pot detectors, etc.
- Even if ion is observed,  $t$  is difference of large numbers
  - Beam spread, measurement errors
- Multiple  $t$ -measurement methods considered

- Method Exact (E):
- Method Approximate (A) (UPCs)
- Method with **exclusivity corrected** (L):

$$-t = -(\mathbf{p}_e - \mathbf{p}_{e'} - \mathbf{p}_{VM})^2 = -(\mathbf{p}_{A'} - \mathbf{p}_A)^2$$

$$-t = (\mathbf{p}_{T,e'} + \mathbf{p}_{T,VM})^2$$

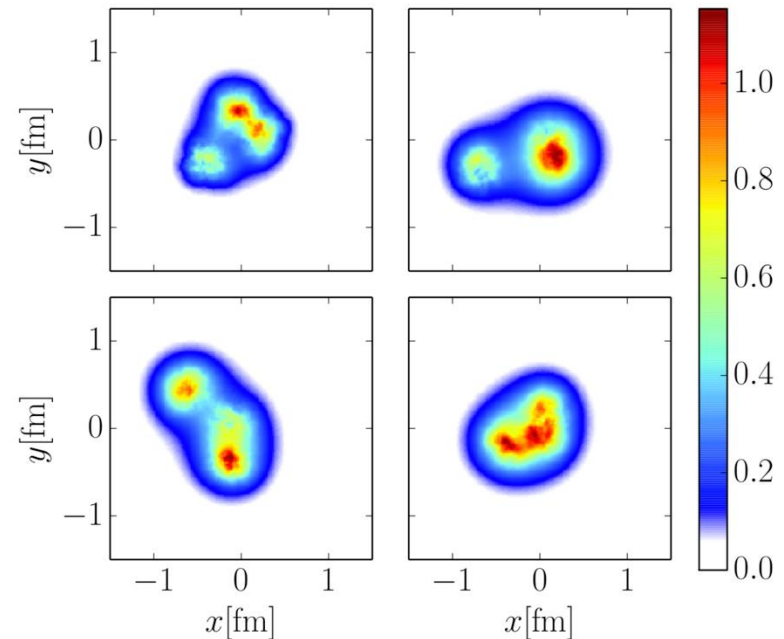
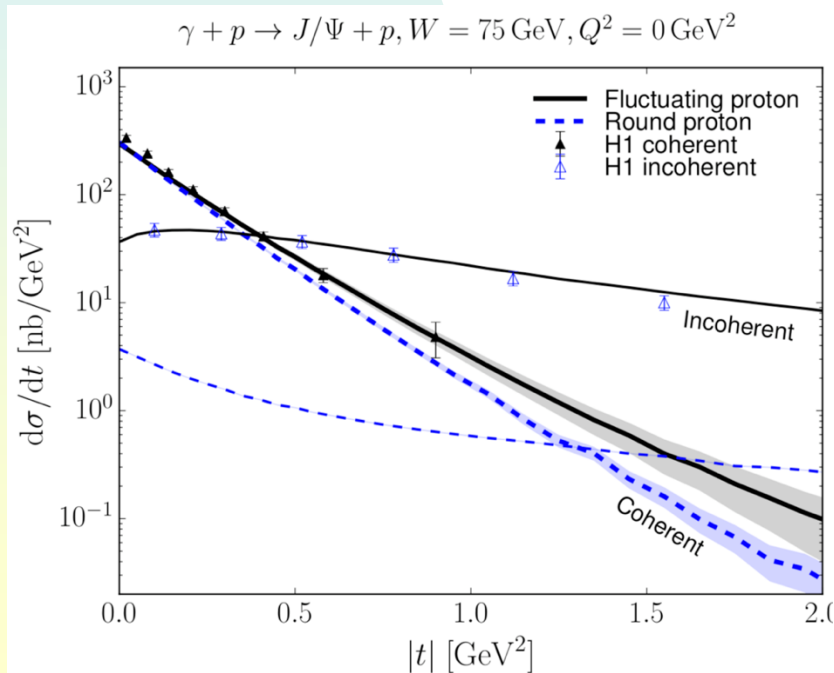
$$-t = -(\mathbf{p}_{A',\text{corr}} - \mathbf{p}_A)^2,$$

where  $\mathbf{p}_{A',\text{corr}}$  is constrained by exclusive reaction.



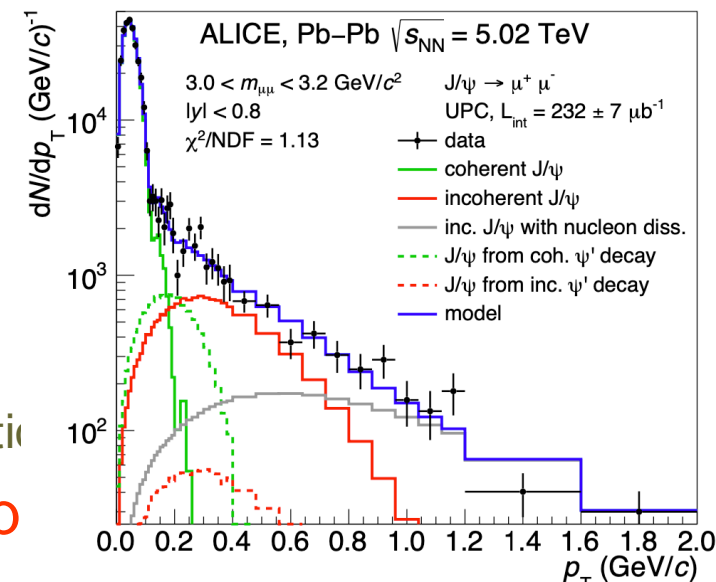
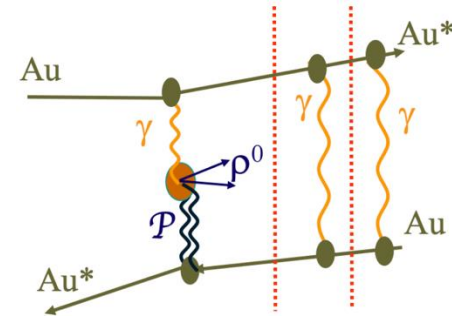
# Incoherent production on protons

- H1 at HERA data on  $J/\psi$  production on protons
- Fluctuations from coherent & incoherent  $J/\psi$  photoproduction.
  - Proton excitations ( $\Delta^+$ )  $\rightarrow$  incoherent
- Two models/calculations of  $d\sigma/dt$  compared
  - Data prefers a fluctuating proton over a smooth proton
- EIC can make precision measurements like this



# Separating coherent & incoherent production on ions

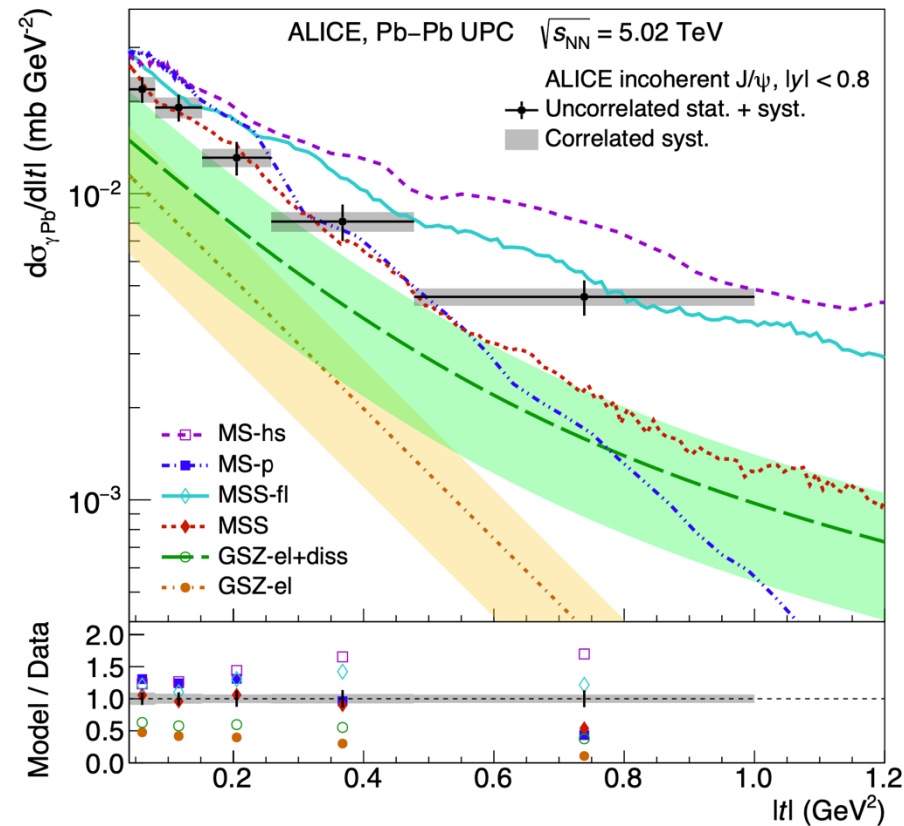
- In UPCs,  $Z\alpha$  is large enough so that the nuclei may exchange additional photons
  - Nuclear breakup complicates separation
  - Photon exchange factorizes
- Coherent dominates at low  $p_T$ 
  - Incoherent dominates at high  $p_T$
  - Subtract one component to get the other
    - Need assumptions re. shape of  $d\sigma/dp_T$ 
      - Shape is based on paradigm
        - $\sigma_{\text{coherent}} = |\sum_i A_i k \exp(ikb)|^2$
      - Somewhat inconsistent to use this paradigm + Good-Walker to find fluctuations
- Presence/absence of neutrons could help separation



ALICE, PRL **132**, 162302 (2024)

# Incoherent $J/\psi$ photoproduction on Pb

- $p_T > 200 \text{ MeV}/c$
- Better agreement with models that include subnucleonic fluctuations
- Large high  $|t|$  tail above expectations from proton form factor



ALICE, PRL **132**, 162302 (2024)

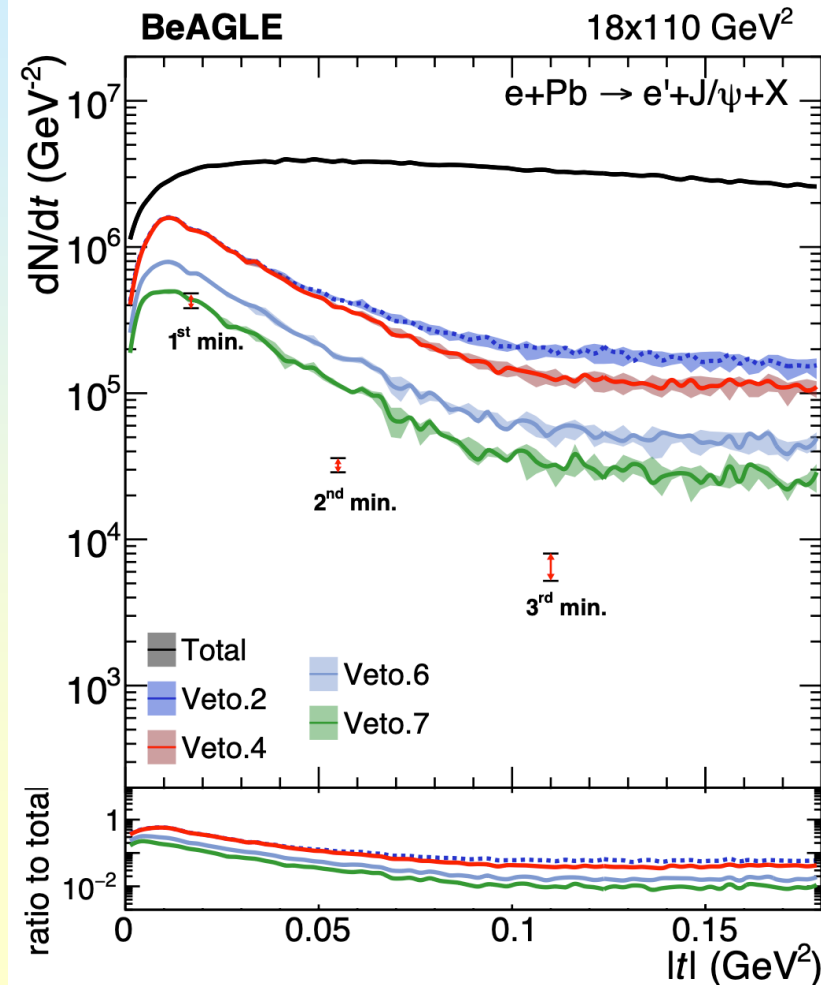
# Theoretical uncertainties in ion breakup

- ePIC can detect nuclear breakup by the presence of neutrons or protons (with near-beam momentum), or of photons in the ZDC from nuclear excitations
  - Typical photon energies are ~ few MeV in emitter frame
    - Lorentz boosted in lab frame
- $^{197}\text{Au}$  has a first excited state at 77 keV (with  $\beta \sim 0.60$ )
  - Not visible in ZDC
  - 2<sup>nd</sup> excited state at 269 keV; boosted to 63 MeV max
- Other low-energy photon lines may be missed, or detected with low efficiency
  - ZDC threshold matters, but background from synch. radiation
- To determine the excitation efficiency accurately, we need a good model of the products of nuclear breakup.
  - Currently use BeAGLE
    - DPMJET + FLUKA for intranuclear cascade
      - Uncertainties are acknowledged to be large



# ePIC veto projected performance

- How well can ePIC veto incoherent  $J/\psi$  production to study coherent?
  - Requires  $\sim 500:1$  to  $1,000:1$  to study coherent production

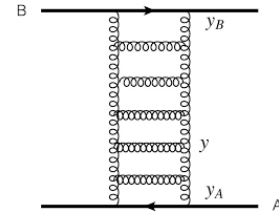


- Veto.1: no activity other than  $e^-$  and  $J/\psi$  in the main detector ( $|\eta| < 4.0$  and  $p_T > 100 \text{ MeV}/c$ );
- Veto.2: Veto.1 and no neutron in ZDC;
- Veto.3: Veto.2 and no proton in RP;
- Veto.4: Veto.3 and no proton in OMDs;
- Veto.5: Veto.4 and no proton in B0;
- Veto.6: Veto.5 and no photon in B0;
- Veto.7: Veto.6 and no photon with  $E > 50 \text{ MeV}$  in ZDC.

Does not reach 500:1  
Modelling will be critical!

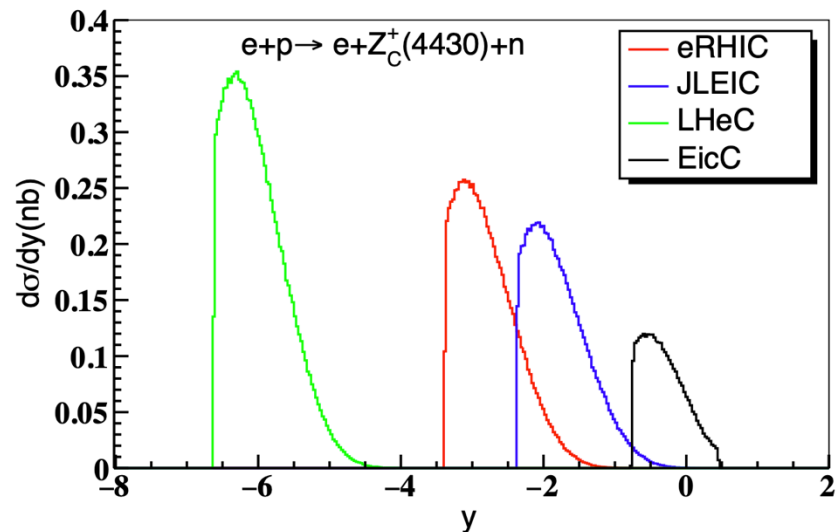
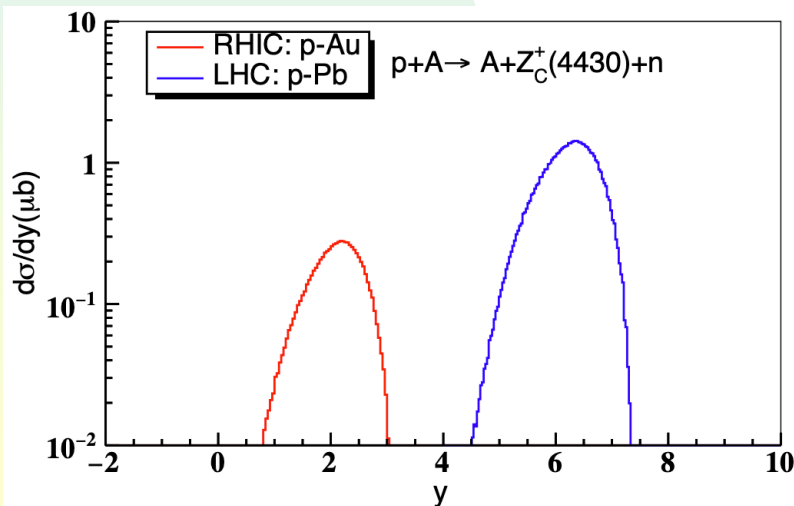
# Beyond Pomerons: Reggeons

- **Pomerons carry the quantum numbers of the vacuum**
  - s-channel helicity conservation means that photon + Pomeron interactions lead to  $J^{PC}=1^-$  states
    - Experimentally well tested
  - Mostly gluons
  - Cross-section rises with energy ( $\sigma \sim W_{\gamma p}^{0.22}$ )
- **Reggeons are summed meson Regge trajectories**
  - Mostly quark-antiquark pairs+
  - Can accommodate a wider range of quantum numbers
    - Broad range of physics!
  - Cross-section drops with energy, ( $\sigma \sim W_{\gamma p}^{-1}$ ) so Reggeon interactions are close-ish to threshold
    - Optimum EIC data collection may occur below maximum energy



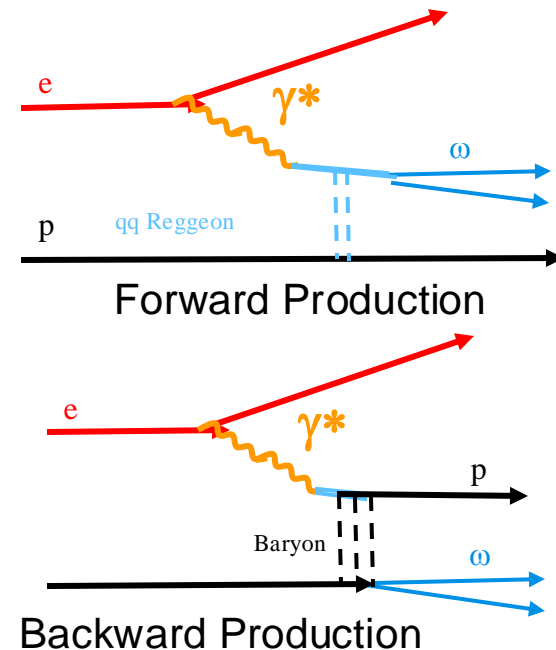
# Production of exotica in UPCs and the EIC

- Exotica with  $J^{PC} = 1^{--}$  can come from  $\gamma$ -Pomeron interaction.
- Other  $J^{PC}$  can only (if at all) come from  $\gamma$ -Reggeon interactions
- In UPCs,  $\gamma$ -Reggeon fusion products are forward, mostly beyond the reach of current detectors.
- $\gamma$ -Reggeon final states are visible at the EIC.
- Predicted rates at the EIC are high enough for characterization
  - $\gamma$ -exotica coupling sensitive to internal structure



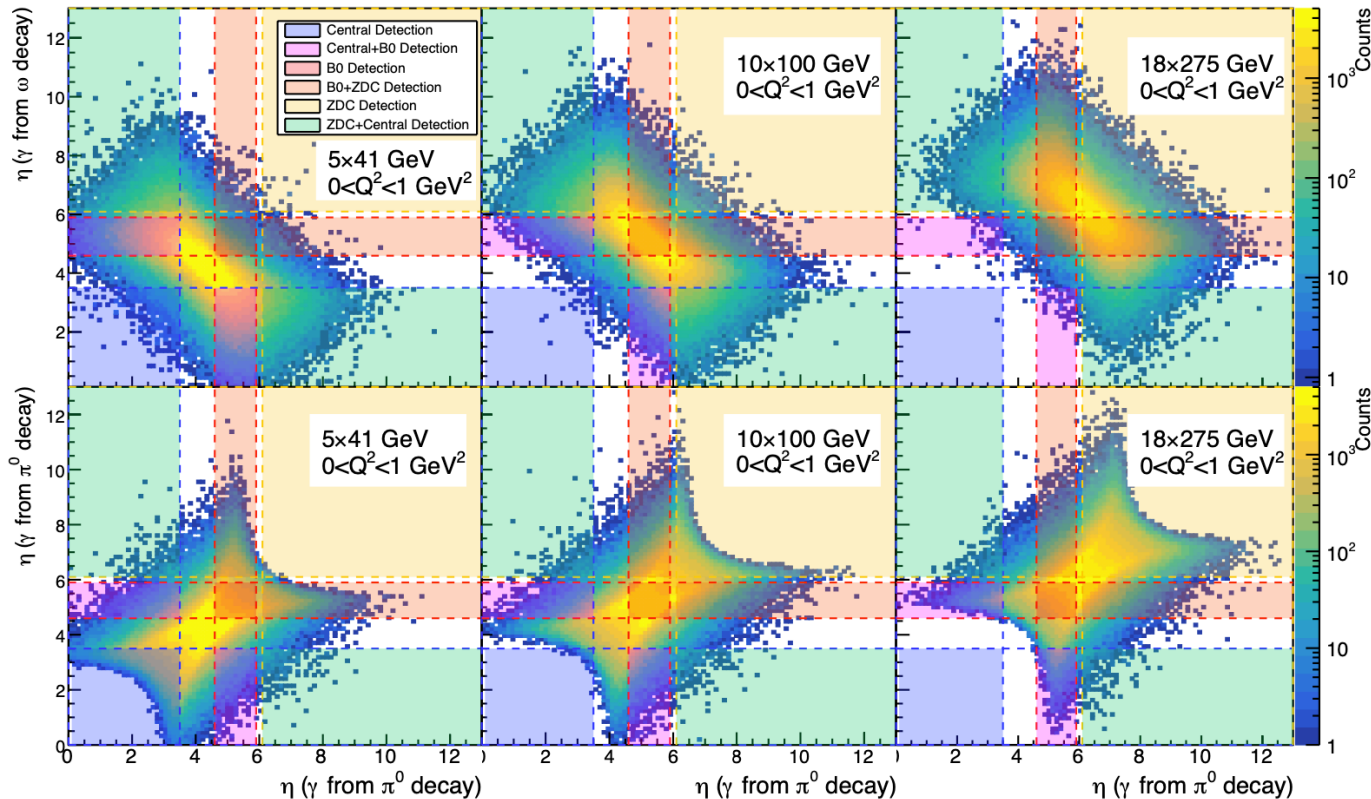
# Backward (u-channel) production

- Reggeons reactions that carry baryon number like  $\gamma p \rightarrow \rho/\omega/\pi^0/\gamma p$ 
  - $d\sigma/dt$  is large  $d\sigma/du$  is small
- Seen by many fixed-target experiments, including at JLab
  - Parameterize using Regge trajectories
  - Rate  $\sim 1/1000$  of forward production
  - Similar to baryon stopping in heavy-ion collisions baryon junction models
- The  $\gamma$ /meson takes most of the proton momentum (so is far forward), while  $t$ 
  - The proton  $\sim$  stops  $\rightarrow$  at mid-rapidity
- $\gamma$ /meson rapidity depends on its mass
  - $\rho \rightarrow \pi\pi$ ,  $\omega \rightarrow \gamma\pi^0$  in B0 detector at lower beam energies
  - $\pi^0$  and  $\gamma$  in ZDC - best at higher beam energies



# $\omega \rightarrow \gamma \pi^0$ backward production kinematics

- $\rho \rightarrow \pi\pi$  kinematics are similar
- B0 is the key detector
- Best detection for 10 GeV e on 100 GeV p
  - $Q^2$  doesn't change kinematics a lot



D. Cebra *et al.*,  
PRC **106**, 015204 (2022);  
Z. Sweger *et al.*,  
PRC **108**, 055205 (2023)

# How do UPCs and the EIC compare?

- UPCs reach higher energy, so lower Bjorken- $x$ 
  - Photons are nearly real, but  $Q^2$  comes from the hard scale of the final state
- The EIC photons cover a wide range of  $Q^2$ , and ePIC can detect the scattered electron, and so tag the photon
- Between the scattered electron and the nearly-hermetic detector, ePIC has very strong power to completely reconstruct exclusive interactions with low backgrounds.
  - The proposed ALICE 3 has coverage out to  $|y| < 4$ , so will partially compete.

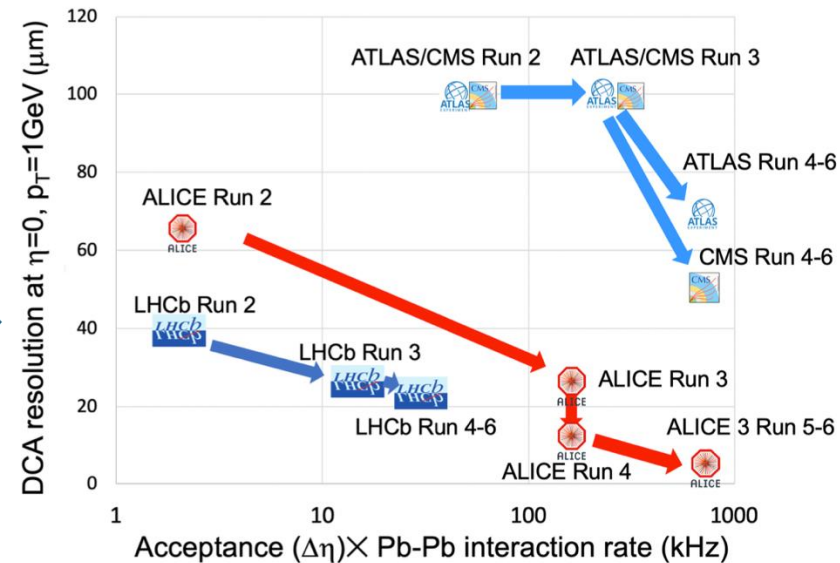
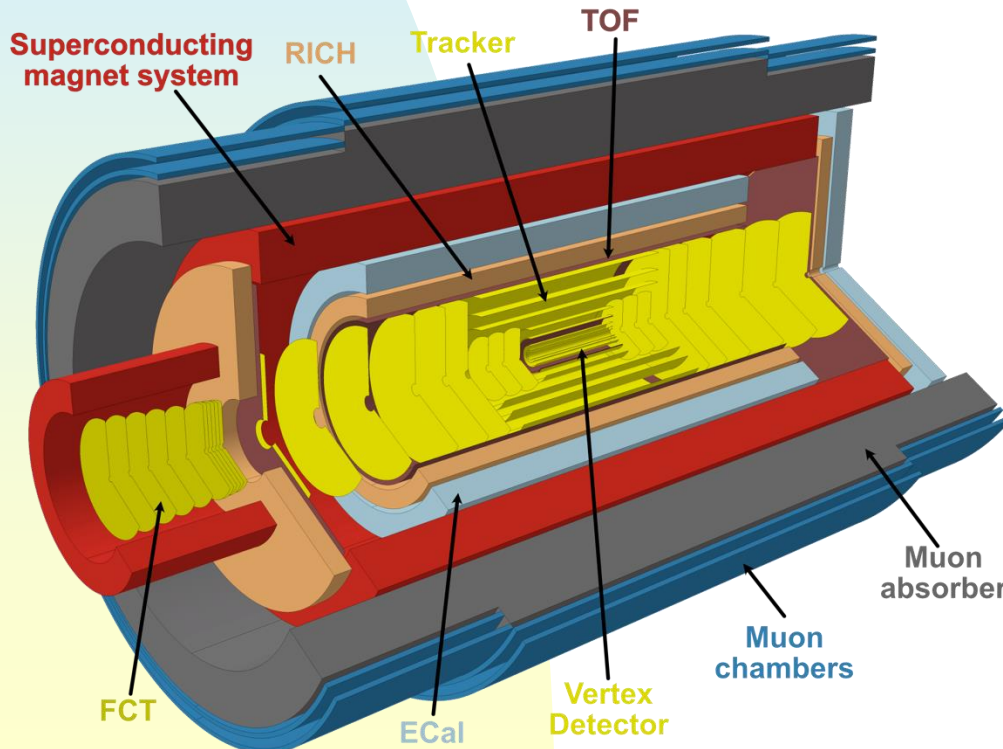


# UPCs in the EIC era

- CMS, ATLAS and LHCb will continue to take data with improved vertexing and other smaller upgrades
  - More mass reach than the EIC
- ALICE 3 is a proposed completely-new detector
- A broad UPC program in  $\gamma\gamma$  and  $\gamma p$  interactions is on-going, and will continue
- What can UPCs do that the EIC can't? <- Key question for US
  - Higher collision-energy  $\gamma\gamma$  and  $\gamma p$  interactions
    - Lower Bjorken-x values (but only at large  $|y|$ )
  - Physics in a strong (EM) field environment
    - UPCs act as a 2-source interferometer
      - Interference seen with single mesons
    - The LHC can extend this to interferometers with two or more mesons

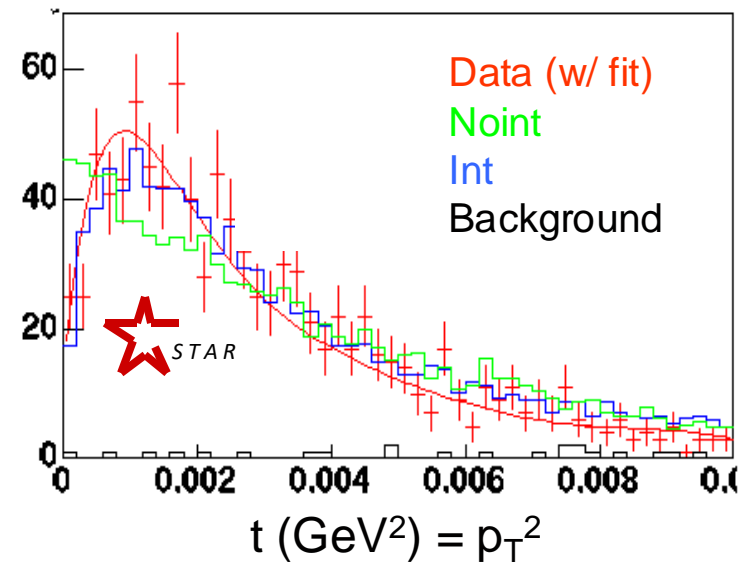
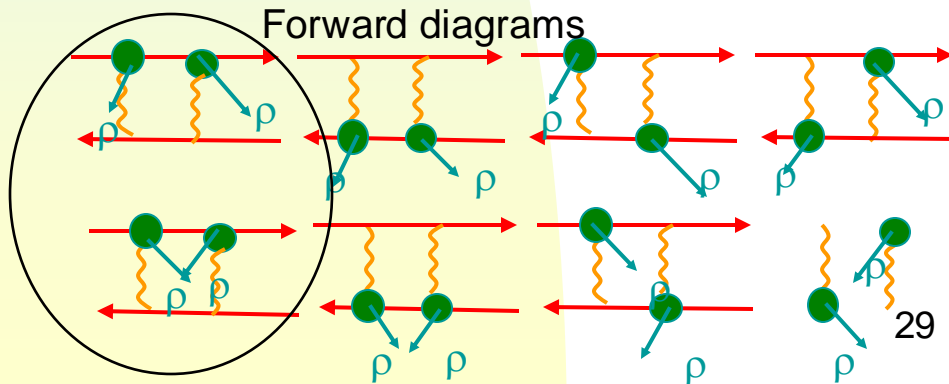
# ALICE 3

- Proposed detector for LHC Runs 5 and 6 (starting ~ 2035)
- Tracking and calorimetry for  $|\eta| < 4$
- Particle identification
- Vertex detector inside beampipe ( $\sim 4 \mu\text{m}$  resolution @ 1 GeV/c)

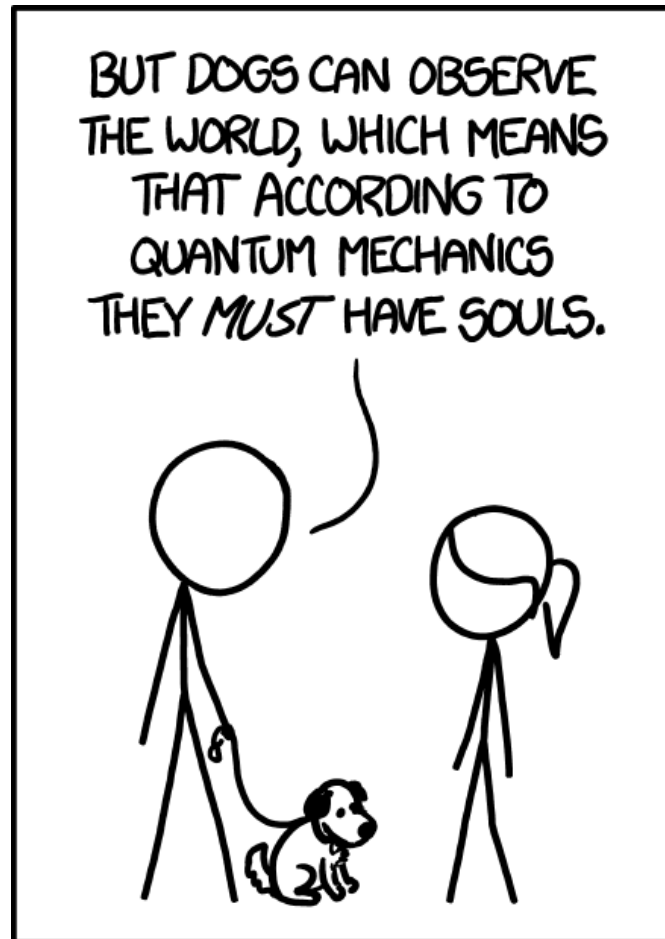


# Two-meson interferometry

- For 1 meson, :  $\sigma \sim |A_1 - A_2 e^{ip \cdot b}|^2$ 
  - At midrapidity  $A_1=A_2$  and ,  $\sigma \rightarrow 0$  as  $p_T \rightarrow 0$
- With 2 identical mesons, the possibilities multiply.
  - Like an interferometer containing two photons.
- For  $|y| \gg 0$ , the two photons are from the same nucleus
  - Superradiant emission: N meson probability is enhanced by N!
    - Like a laser
    - $\langle p_T \rangle \sim \langle p_{T1} \rangle / N$
- Stimulated decay?
  - e. g.  $\pi^+$  from  $\rho$  decay close in phase



# Quantum interferometry – an alternate view



**PROTIP:** YOU CAN SAFELY IGNORE ANY SENTENCE THAT INCLUDES THE PHRASE "ACCORDING TO QUANTUM MECHANICS"

xkcd.com

# Conclusions

- Exclusive interactions can probe many interesting physics topics, including the low- $x$  structure of matter, including its spatial distribution
- The nearly-hermetic ePIC detector at the EIC is well suited to pursue high-statistics measurements with small systematic errors, over a wide range of  $Q^2$ .
  - Precise measurements of gluon saturation.
  - Transverse distribution of gluons in nucleus
  - Event-by-event fluctuations in gluon content (hot spots)
    - Measurements of  $d\sigma/dt$  are limited by the limited  $t$  resolution.
- ePIC will also study backward production, exotica, etc.
- UPCs at the LHC will retain their interest during the EIC era, providing unique data on multi-meson production in high fields, and of nuclear structure at lower Bjorken- $x$  than the EIC can reach.