

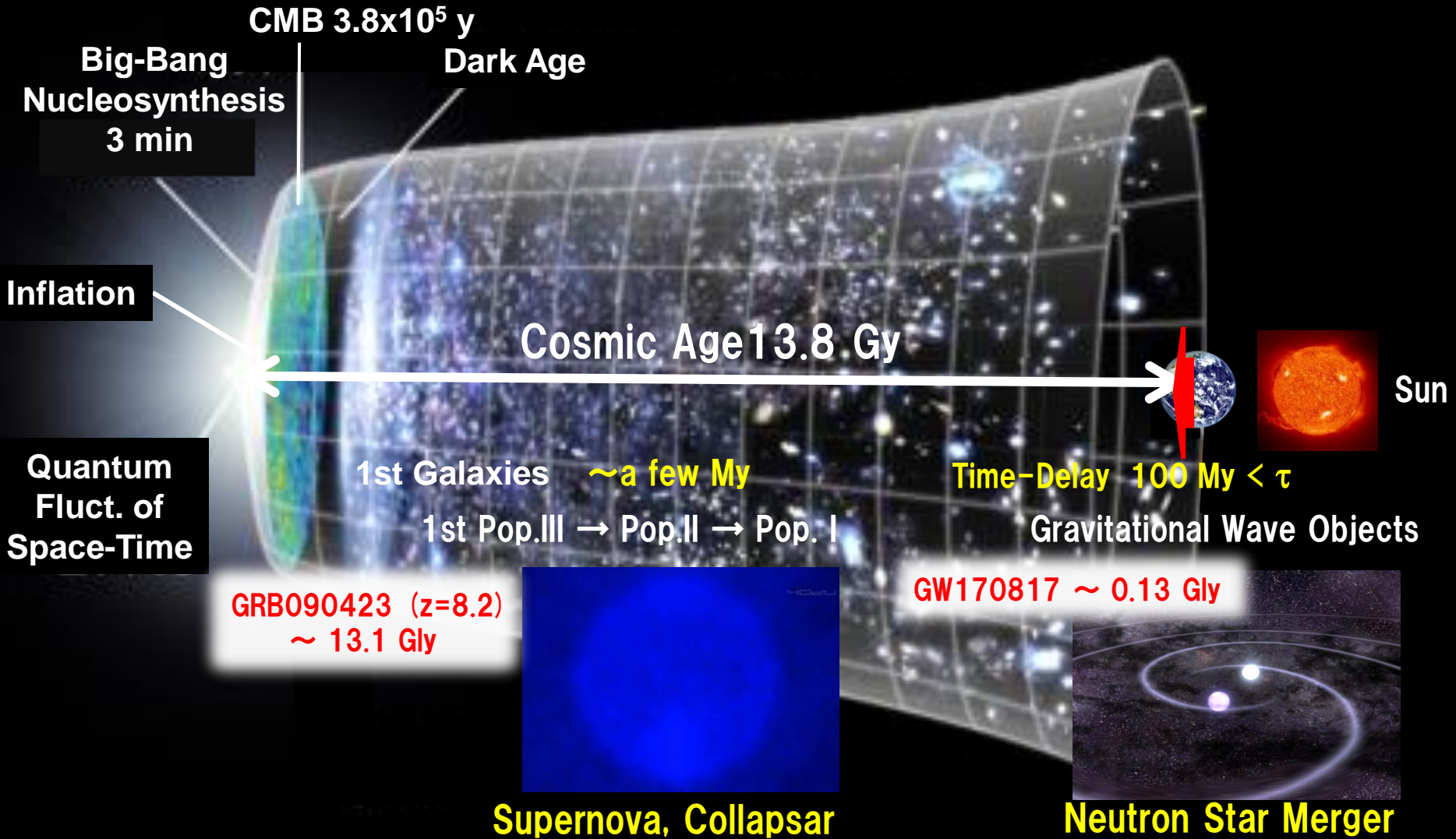
Workshop on “The r–process and the nuclear EOS after LIGO’s third observing run” (INT 20R–1b), May 23–27, 2022

**R–process in GW Objects, Supernovae, Collapsars
& Neutron Star Mergers,
and Galactic Chemo–dynamical Evolution**

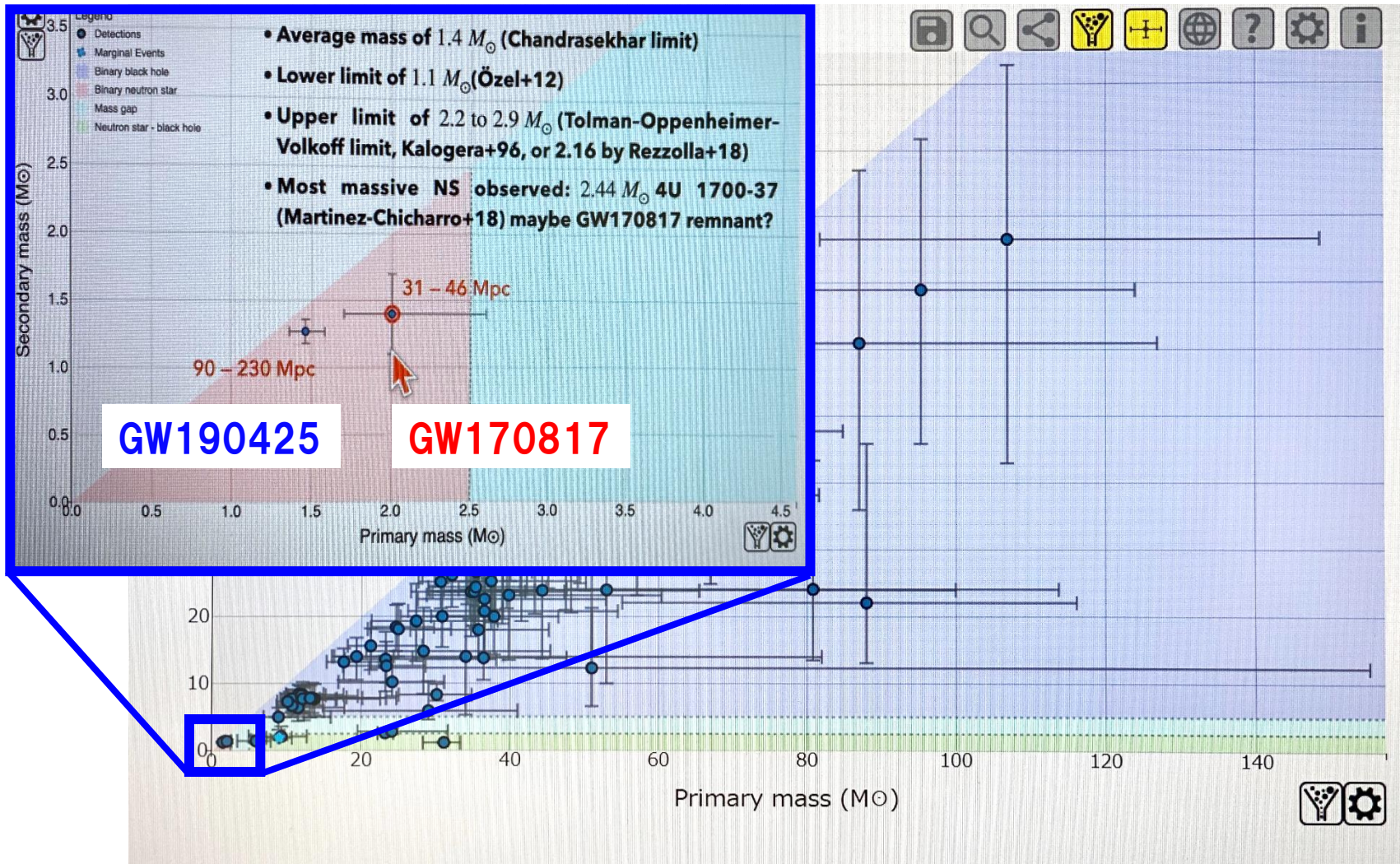
Taka Kajino

Beihang University / NAOJ / University of Tokyo
Kajino@buaa.edu.cn

Origin of Heavy Nuclei in Cosmic & Galactic Evolution



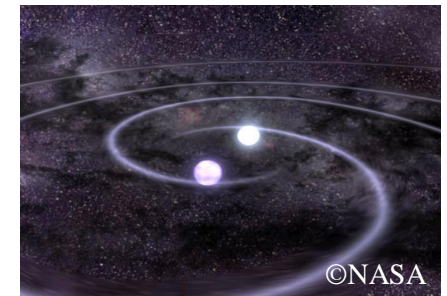
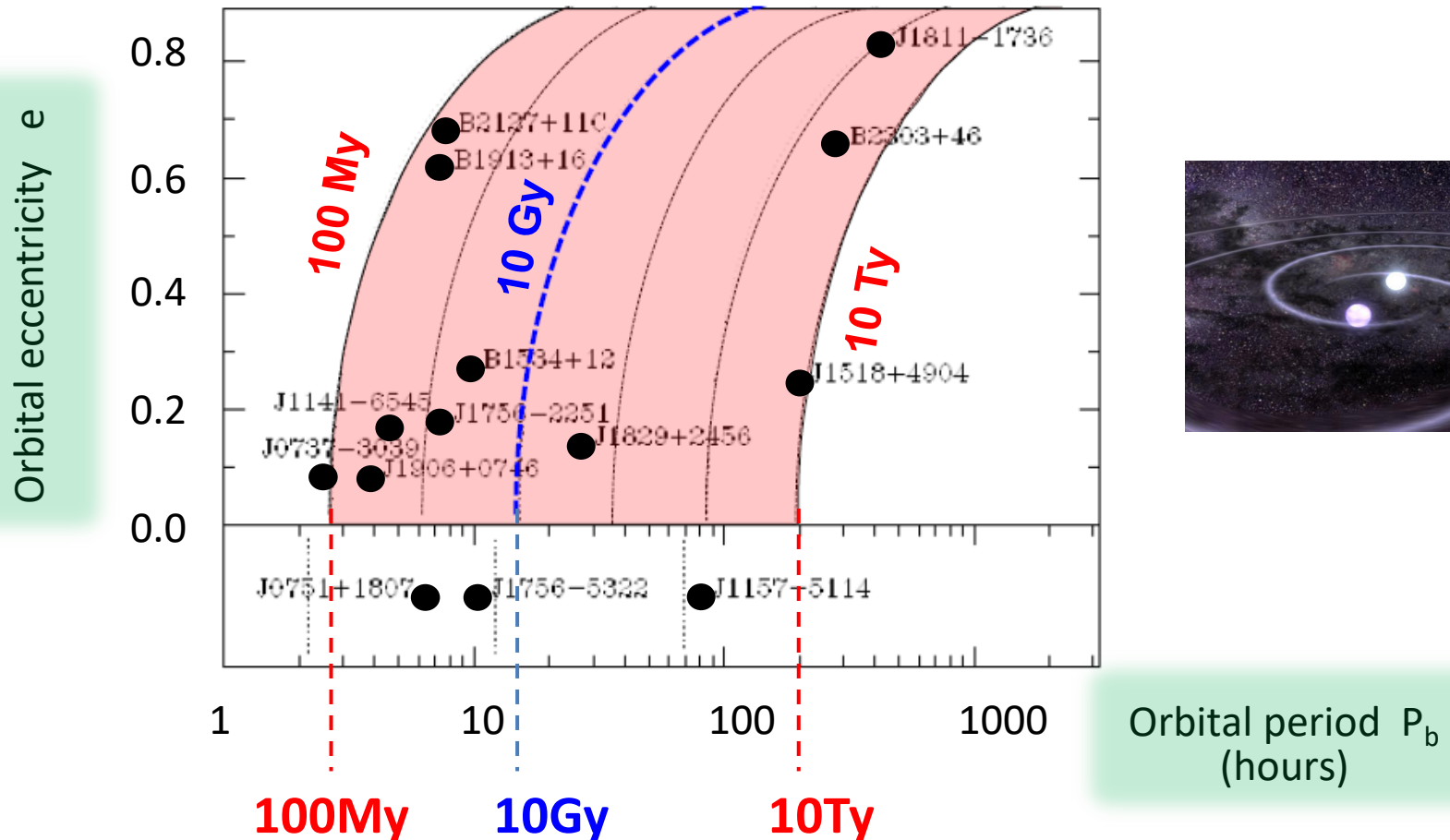
LIGO-Virgo Compact Binary Catalogue



Expected Merger Time-Delay from Binary Pulsars

General Relativity : $\tau_c \simeq 9.83 \times 10^6 \text{ yr} \left(\frac{P_b}{\text{hr}} \right)^{8/3} \times \left(\frac{m_1 + m_2}{M_\odot} \right)^{-2/3} \left(\frac{\mu}{M_\odot} \right)^{-1} (1 - e^2)^{7/2}$

BINARY PULSARS : Lorimer, Living Rev. Rel. 11(2008), 8; Beniamini+ (2019).



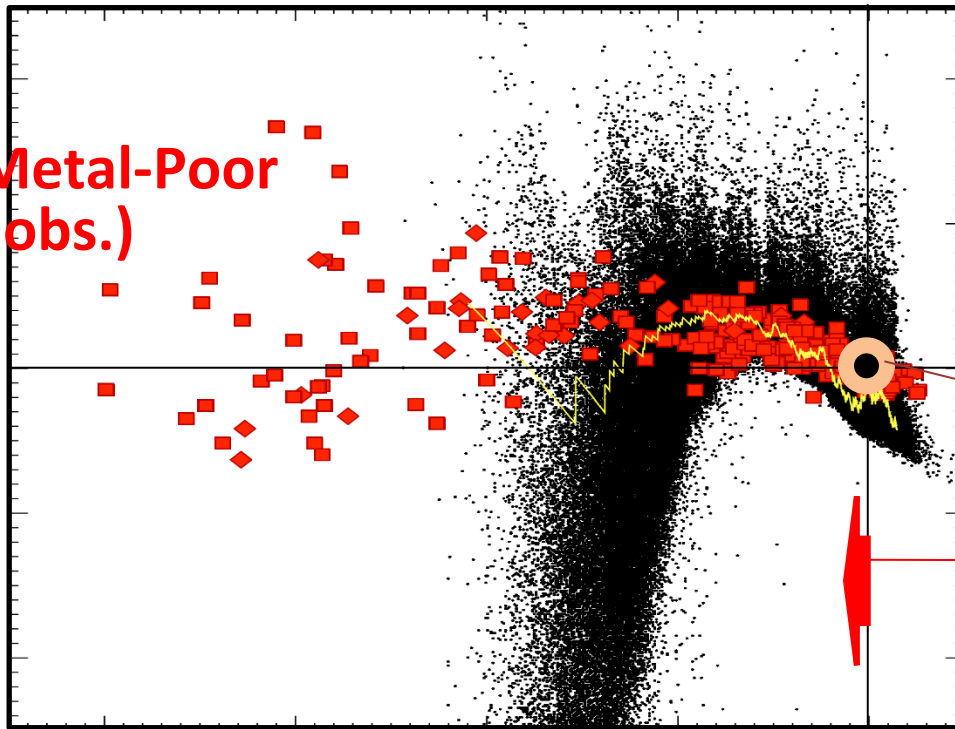
GCE : No Dynamics!

Argast, Samland,
Thielemann, Qian, A&Ap
416 (2004), 997.

$\tau_c = 100 \text{ My}$

Extremely Metal-Poor
Stars (obs.)

$[\text{Ba}_r/\text{Fe}]$

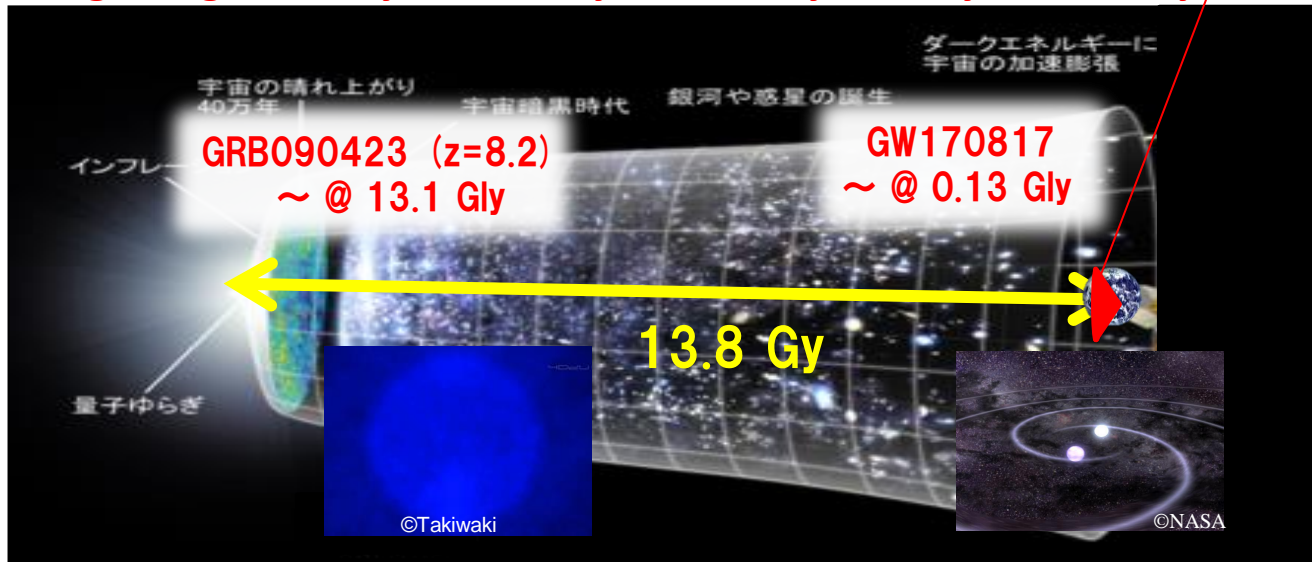


Sun (obs.)

GW170817
@ 0.13 Gly

$t/10\text{Gy}$
 $\approx 10 [\text{Fe}/\text{H}]$

$-\infty$ -4 -3 -2 -1 0 $[\text{Fe}/\text{H}]$
Big-Bang 1My 10My 100My 1Gy 13.8Gy time



Supercomputer Simulation of Binary Neutron Star Mergers

Galactic Chemo-Dynamical Evolution

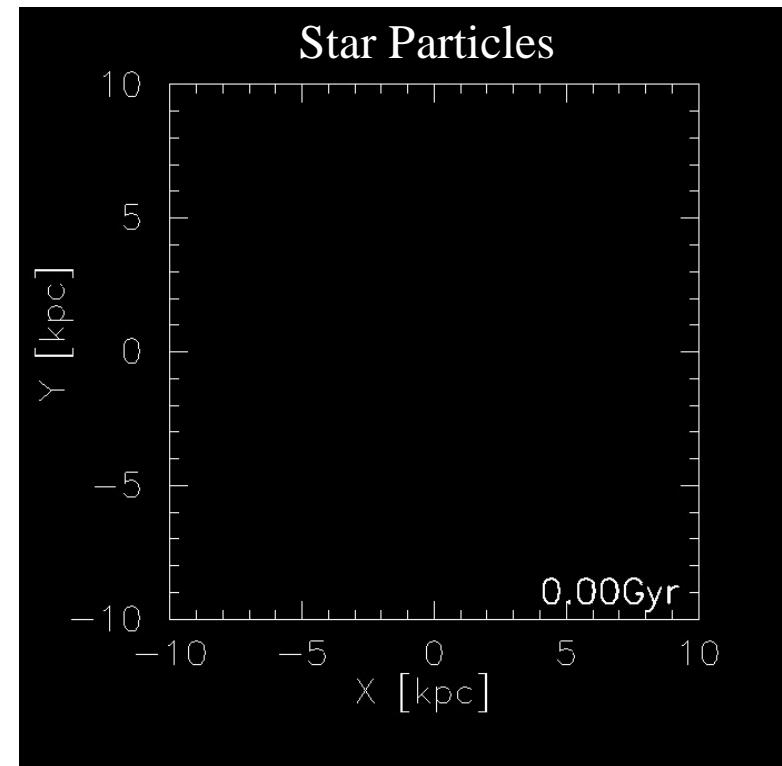
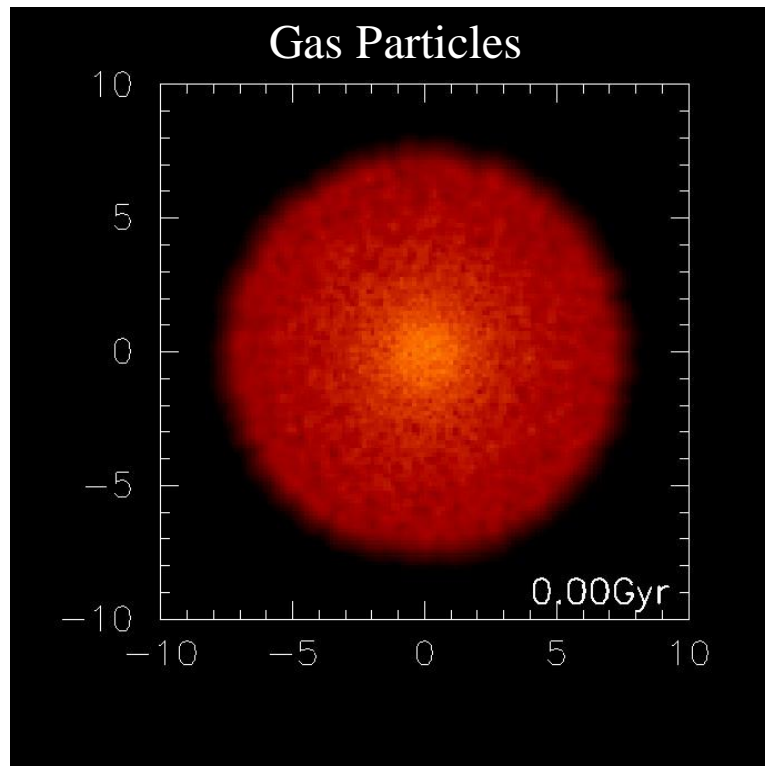
N-Body/SPH Simulation of Dwarf Spheroidal

Gas, DMs & Stars

Star forms in SFR at $T < 10^4\text{K}$, $v < 0$, $n_{\text{H}} > 100\text{ cm}^{-3} \rightarrow 100\text{pc}$, Gas mixing

NSMs($\tau_c=0.1\text{Gyr}$, Ba)+ SNe(1My, Fe) : $M_{\text{tot}} = 7 \times 10^8 M_{\text{sun}}$, $N_i = 5 \times 10^5$ particles, $M_{\star} = 100 M_{\text{sun}}$

Hirai, Kajino, et al., ApJ 814 (2015), 41; MNRAS 466 (2017), 2474



N-Body/SPH Simulation of Chemo-Dynamical Evolution of Dwarf Spheroidal (Building Blocks of MW Halo)

Argast, Samland, Thielemann and Qian,
A&Ap 416 (2004), 997.

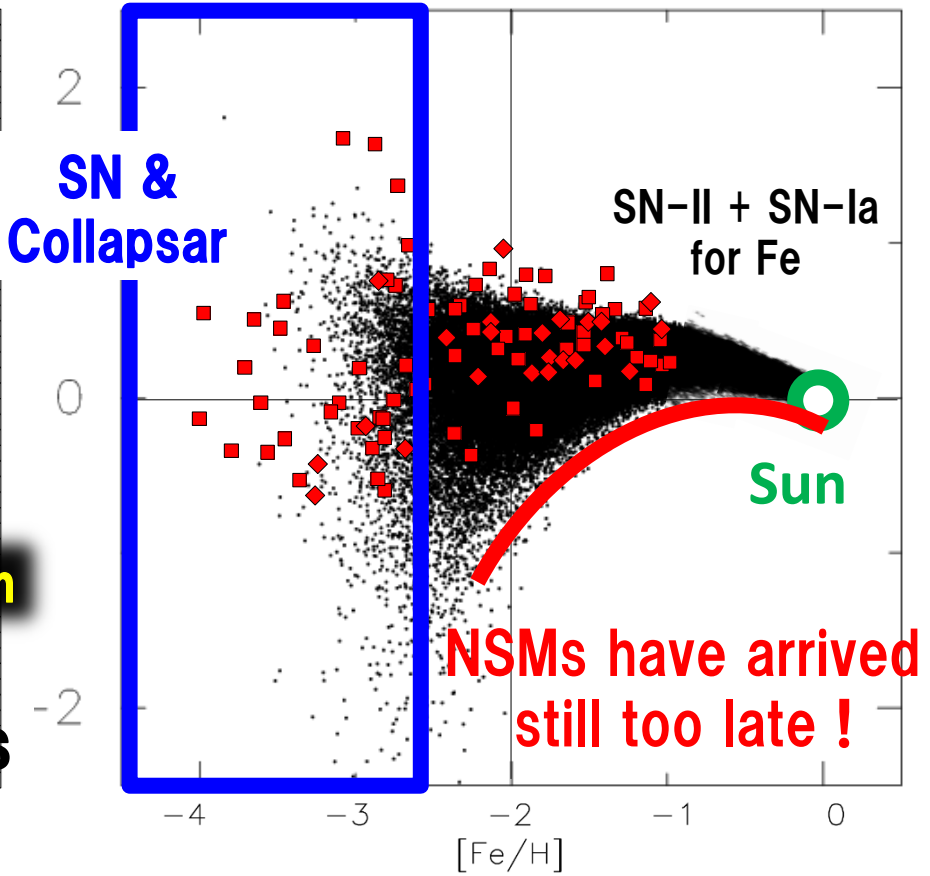
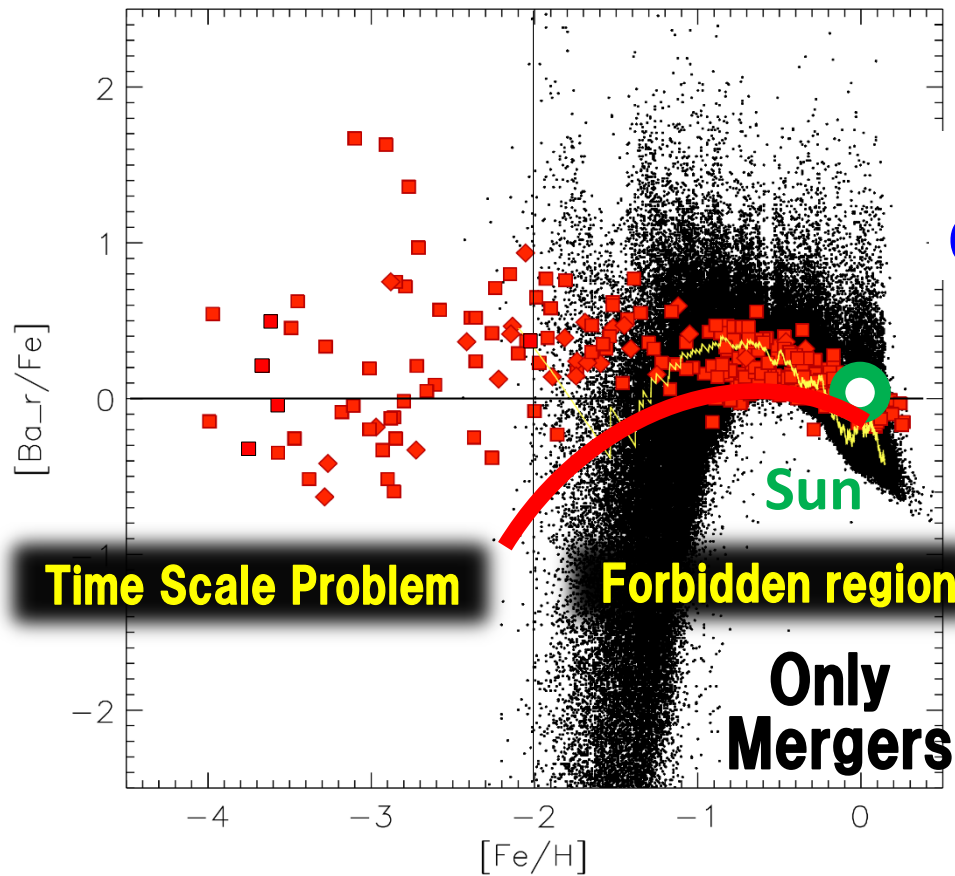
Hirai, Ishimaru, Saitoh, Fujii, Hidaka and Kajino,
ApJ 814 (2015), 41; MNRAS 466 (2017), 2474.

$$\tau_c = 100\text{My}$$

$$\tau_c = 100\text{My}$$

W/O N-body Dynamics & Gas mixing

With N-body Dynamics, Gas mixing



1My **10Gy**

100My **10Gy**

1My **10Gy**

100My **10Gy**

Purpose

- :- to elucidate “when” and “how” the neutron star mergers, CCSNe and Collapsars have contributed differently to the r-process over the entire history of Galactic evolution.**
- :- to elucidate coupling among nuclear physics, neutrino physics and astronomy in the studies of r-process.**

Contents

1. Galactic Chemical Evolution (GCE)

2. R-Process Nucleosynthesis in Various Sites

- Neutron Star Mergers
- Core-Collapse Supernovae (ν -wind & MHD Jet)
- Collapsars

3. ν -Oscillations, Collective & MSW, and Mass Hierarchy

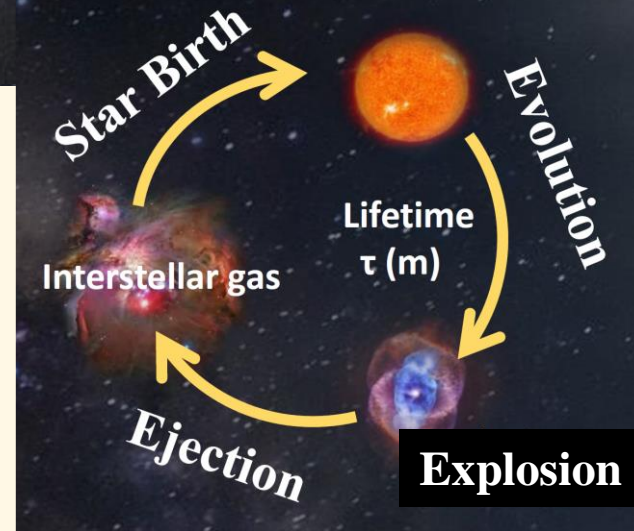
4. Results of GCE

Cosmic & Galactic Evolution

Cosmic Gas- and Nuclear-Evolution

$$\dot{\sigma}_X = \text{Inflow} \cdot \delta_{X,gas} - \frac{\sigma_X}{\sigma_{gas}} \cdot \underbrace{B(\xi_{gas})}_{\text{Stellar Birth Rate}} + \int \underbrace{B(t - \tau(m)) \phi(m) E_X(m)}_{\text{Ejected Nucleus from SNe or NSM}} dm$$

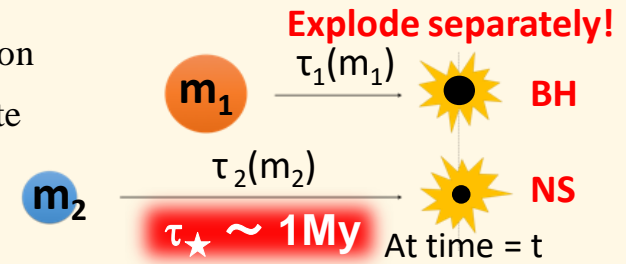
$X = \text{Ejected Nucleus from SNe or NSM}$



CCSN & Hypernova/Collapsar Rate :

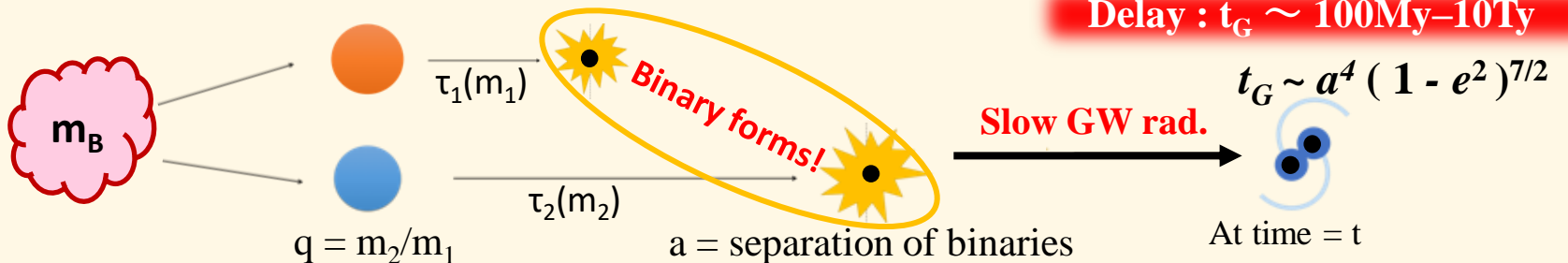
$$R_{SNII} = \int_{m_l}^{m_h} \phi(m) B(t - \tau(m)) dm$$

$\phi(m)$: Initial mass function
 $B(t)$: Star Formation Rate



Binary Neutron Star Merger Rate :

$$R_{NSM} = \epsilon_{NSM} \int_{m_l}^{m_h} dM_B \phi(M_B) \int_{q_l}^1 dq f(q) \int_{a_l}^{a_h} da P(a) B(t - \tau(m_2) - t_G)$$



Observed EVENT RATES

Contribution = Ejected Mass [M_{\odot}] x Event Rate [/Galaxy/Century]

ν SN (Weak r) = 7.4×10^{-4} x $(1.9 \pm 1.1)^a$

MHD Jet SNe = 0.6×10^{-2} x $((0.03 \pm 0.02) \times (1.9 \pm 1.1))^b$

* Binary NSMs (Short-GRB) = $(2 \pm 1) \times 10^{-2}$ x $(1-28) \times 10^{-3}^c$

* Collapsars (Failed SN) = Assuming the same for MHD Jet SNe

Observations a $1.9 \pm 1.1^*$ Diehl, et al., Nature 439, 45 (2006). * 1.3 ± 0.6 (2018)

b 0.03 ± 0.02 Winteler, et al., ApJ 750, L22 (2012).

Obs. Estimate c $(1-28) \times 10^{-3}$ Kalogera, et al., ApJ 614, L137 (2004).

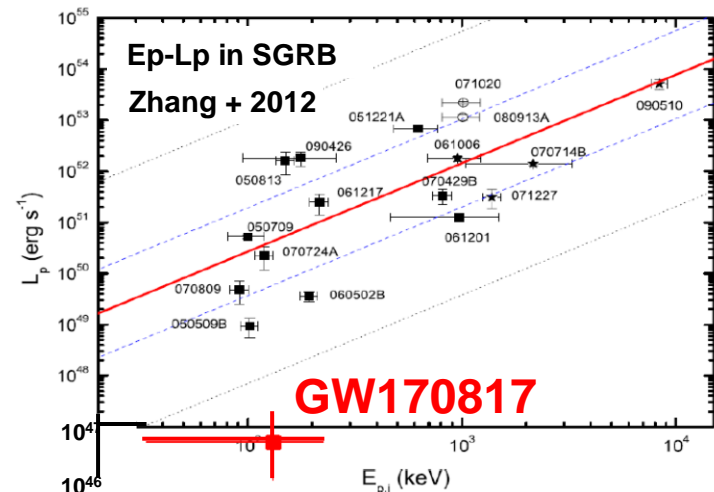
* Binary NSM ← Central engine of short GRB

- GW170817: Why faint ?
- Jet inclination and beaming $< 5^\circ$?

* Collapsar (BH) ← Failed Supernovae, Long GRB

Yamazaki et al. (2022); Harikae et al. (2009, 2010); Nakamura et al. (2015),

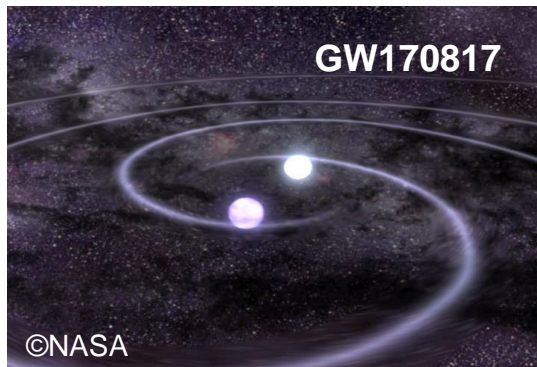
c.f. Siegel et al. (2019) assumed: Super-Luminous SN
Hypernova (Long GRB)



Astrophysical Sites for R-Process

BINARY Stars

Neutron Star Merger



Time Delay : $100 \text{ My} < \tau < 10 \text{ Ty}$

Lorimer, Living Rev. Rel. 11(2008), 8.
Beniamini+ (2019), Timmes+ (1995)

Failed SN \rightarrow Collapsar

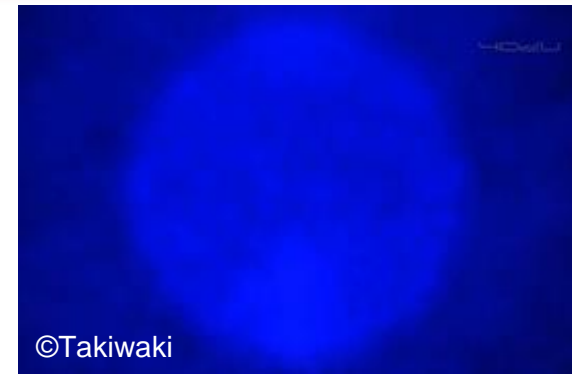
MacFadyen, Woosley, ApJ 524 (1999), 262;
Nakamura, Kajino, Mathews, Sato & Harikae,
A&Ap 582 (2015), A34; Yamazaki, et al. (2022).

Super-Luminous SN/Hypernova

Siegel, Barnes & Metzger, Nature 569 (2019), 243.

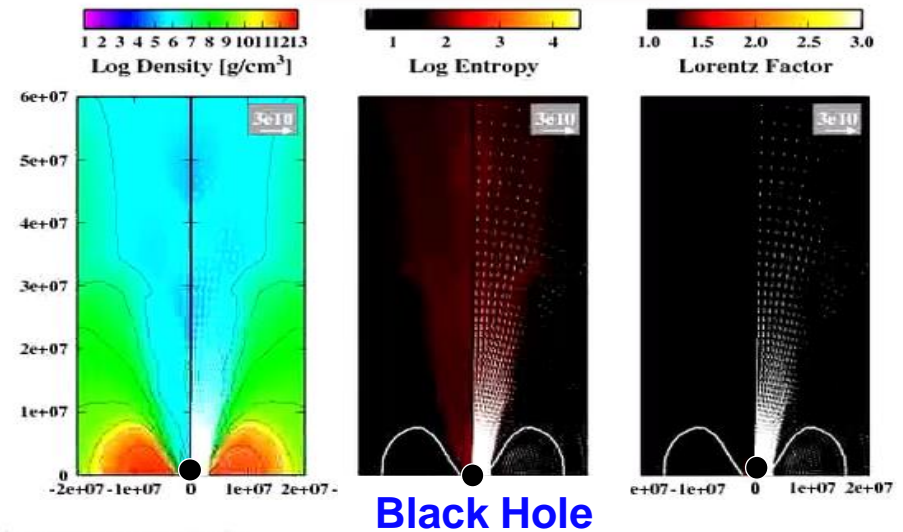
SINGLE Star

CCSN II : ν -DW & MHD Jet



Neutron
Star

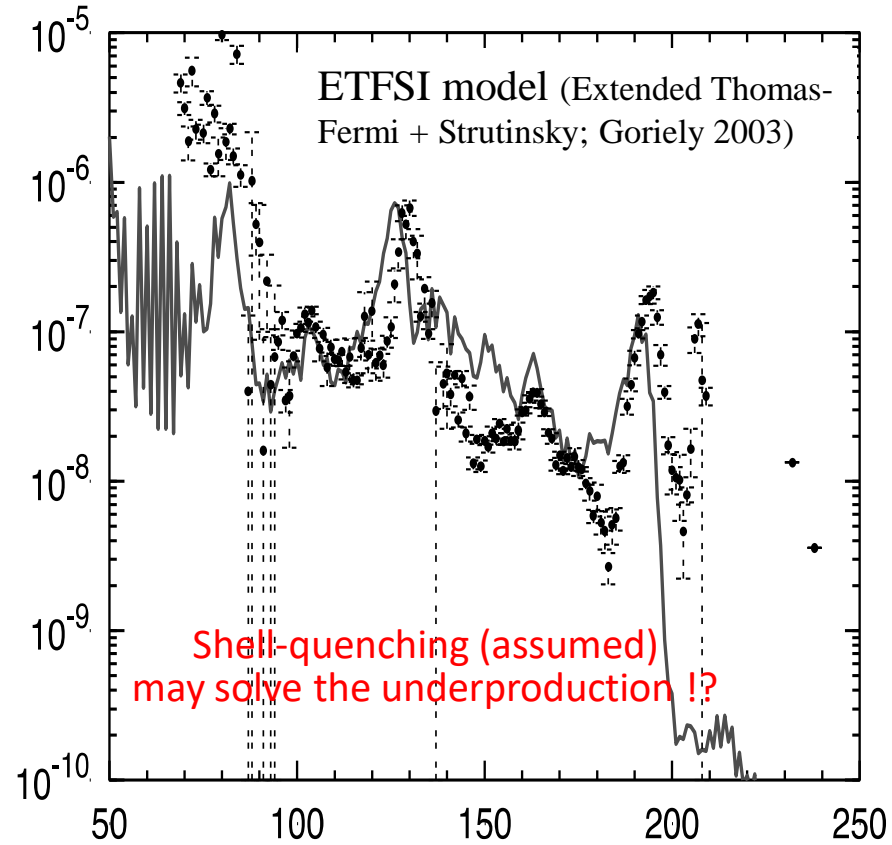
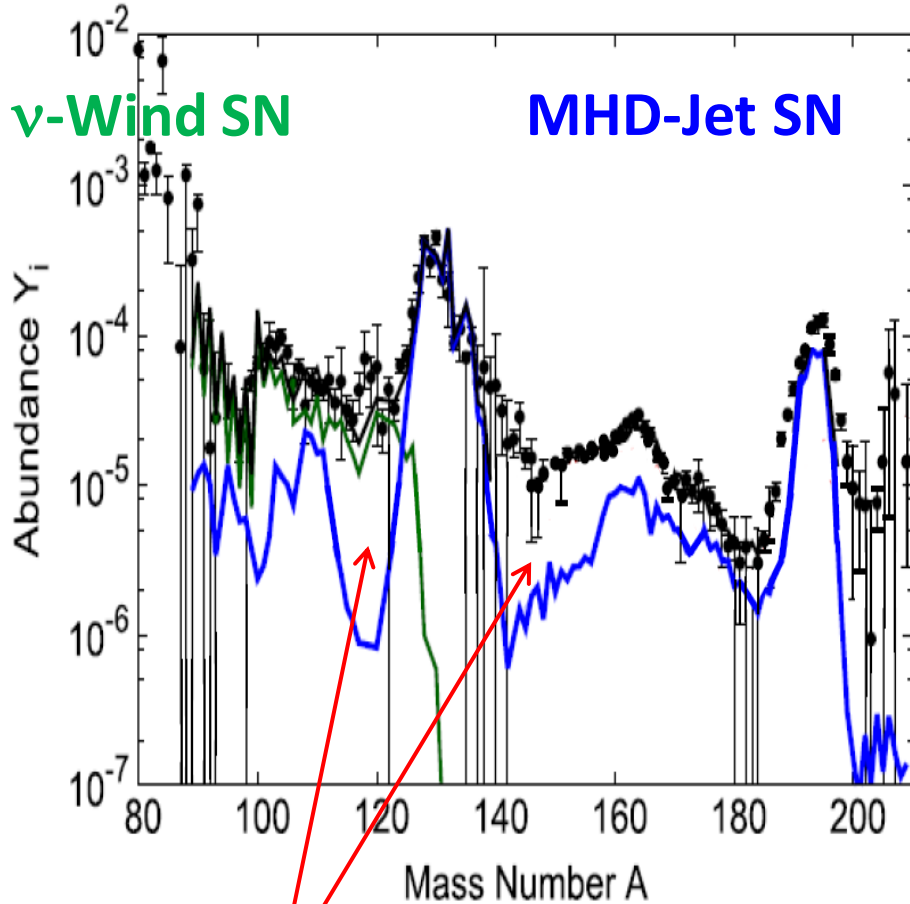
Collapsar Jet



1. Supernovae (ν -driv. Wind & MHDJet)

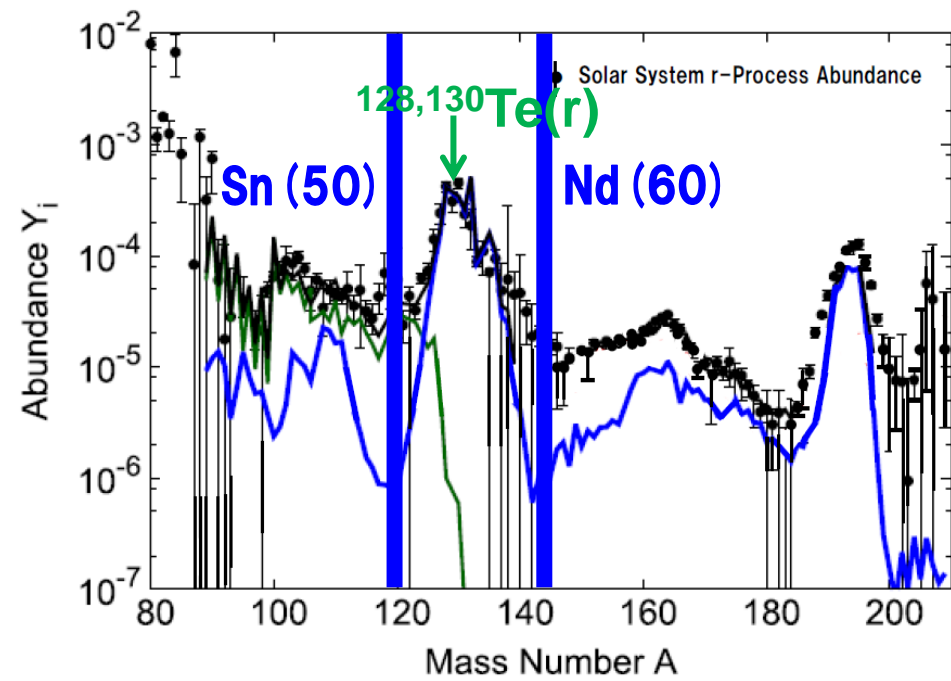
Shibagaki, Kajino, Mathews, Chiba, Nishimura, Lorusso, ApJ 816, 79; Kajino & Mathews (2017), ROPP 80, 084901; Kajino, Aoki, Balantekin, Dihel, Famiano, Mathews (2019), PPNP 107, 109.

©Takiwaki



Underproduction \rightarrow Possible Solution

1. Nucl. Phys. - Shell Quenching ?
2. Astro. - Another Site ?
(Merger or Collapsar ?)



55	Cs122	Cs123	Cs124	Cs125	Cs126	Cs127	Cs128	Cs129	Cs130	Cs131	Cs132	Cs133	Cs134	Cs135	Cs136	Cs137	Cs138	Cs139	Cs140	Cs141	Cs142	Cs143
54	Xe121	Xe122	Xe123	Xe124	Xe125	Xe126	Xe127	Xe128	Xe129	Xe130	Xe131	Xe132	Xe134	Xe135	Xe136	Xe137	Xe138	Xe139	Xe140	Xe141	Xe142	
53	I120	I121	I122	I123	I124	I125	I126	I127	I128	I129	I130	I131	I132	I133	I134	I135	I136	I137	I138	I139	I140	I141
52	Te119	Te120	Te121	Te122	Te123	Te124	Te125	Te126	Te127	Te128	Te129	Te130	Te131	Te132	Te133	Te134	Te135	Te136	Te137	Te138	Te139	Te140
51	Sb118	Sb119	Sb120	Sb121	Sb122	Sb123	Sb124	Sb125	Sb126	Sb127	Sb128	Sb129	Sb130	Sb131	Sb132	Sb133	Sb134	Sb135	Sb136	Sb137	Sb138	Sb139
50	Sn117	Sn118	Sn119	Sn120	Sn121	Sn122	Sn123	Sn124	Sn125	Sn126	Sn127	Sn128	Sn129	Sn130	Sn131	Sn132	Sn133	Sn134	Sn135	Sn136	Sn137	
49	In116	In117	In118	In119	In120	In121	In122	In123	In124	In125	In126	In127	In128	In129	In130	In131	In132	In133	134In	In135		
48	Cd115	Cd116	Cd117	Cd118	Cd119	Cd120	Cd121	Cd122	Cd123	Cd124	Cd125	Cd126	Cd127	Cd128	Cd129	Cd130	Cd131					
47	Ag114	Ag115	Ag116	Ag117	Ag118	Ag119	Ag120	Ag121	Ag122	Ag123	Ag124	Ag125	Ag126	Ag127	Ag128	Ag129						
46																						

N=82

R-process path

128Pd

~~Shell Quenching ?~~

RIKEN-RIBF : Decay Spectroscopy around $A = 100-145$

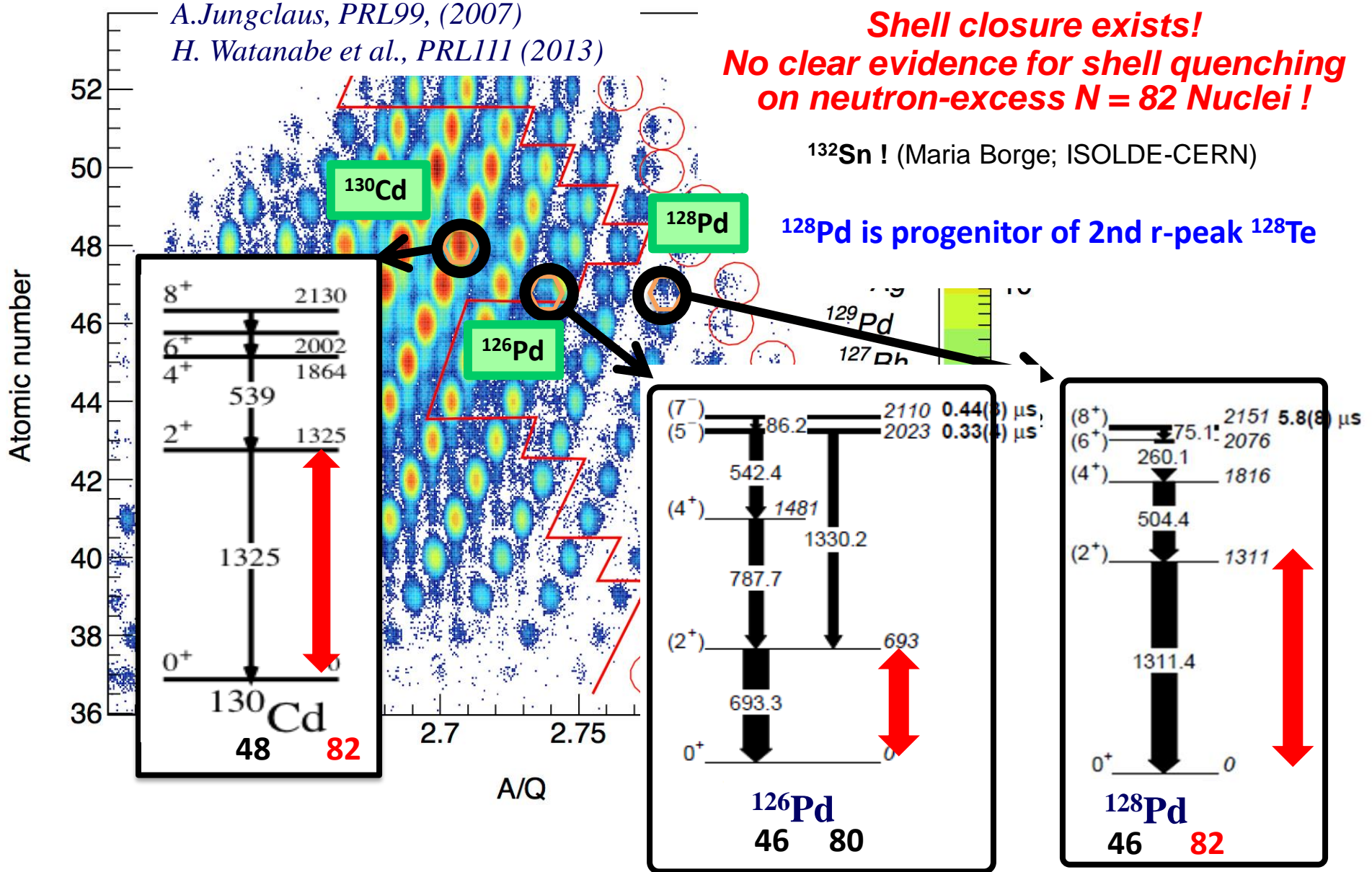
G. Lorusso et al., PRL 114 (2015), 192501.

A. Jungclauss, PRL99, (2007)
H. Watanabe et al., PRL111 (2013)

Shell closure exists!
No clear evidence for shell quenching
on neutron-excess $N = 82$ Nuclei !

^{132}Sn ! (Maria Borge; ISOLDE-CERN)

^{128}Pd is progenitor of 2nd r-peak ^{128}Te

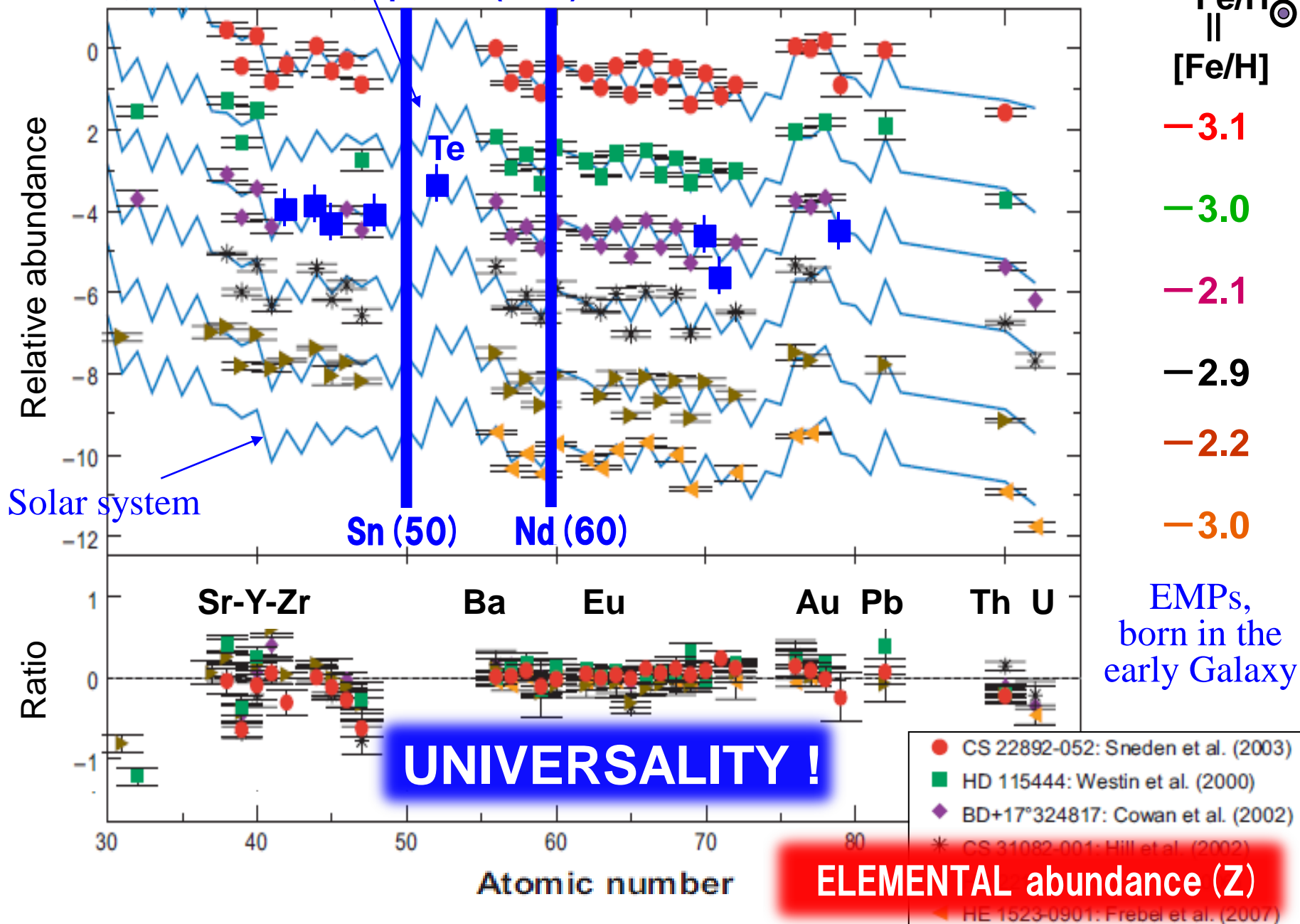


Sneden, Cowan, Gallino, ARAA 46 (2008) 241.

HST-obs., Roederer et al., ApJ 747 (2012) L8.

$$\frac{t}{10^{10}y} \doteq 10^{[Fe/H]}$$

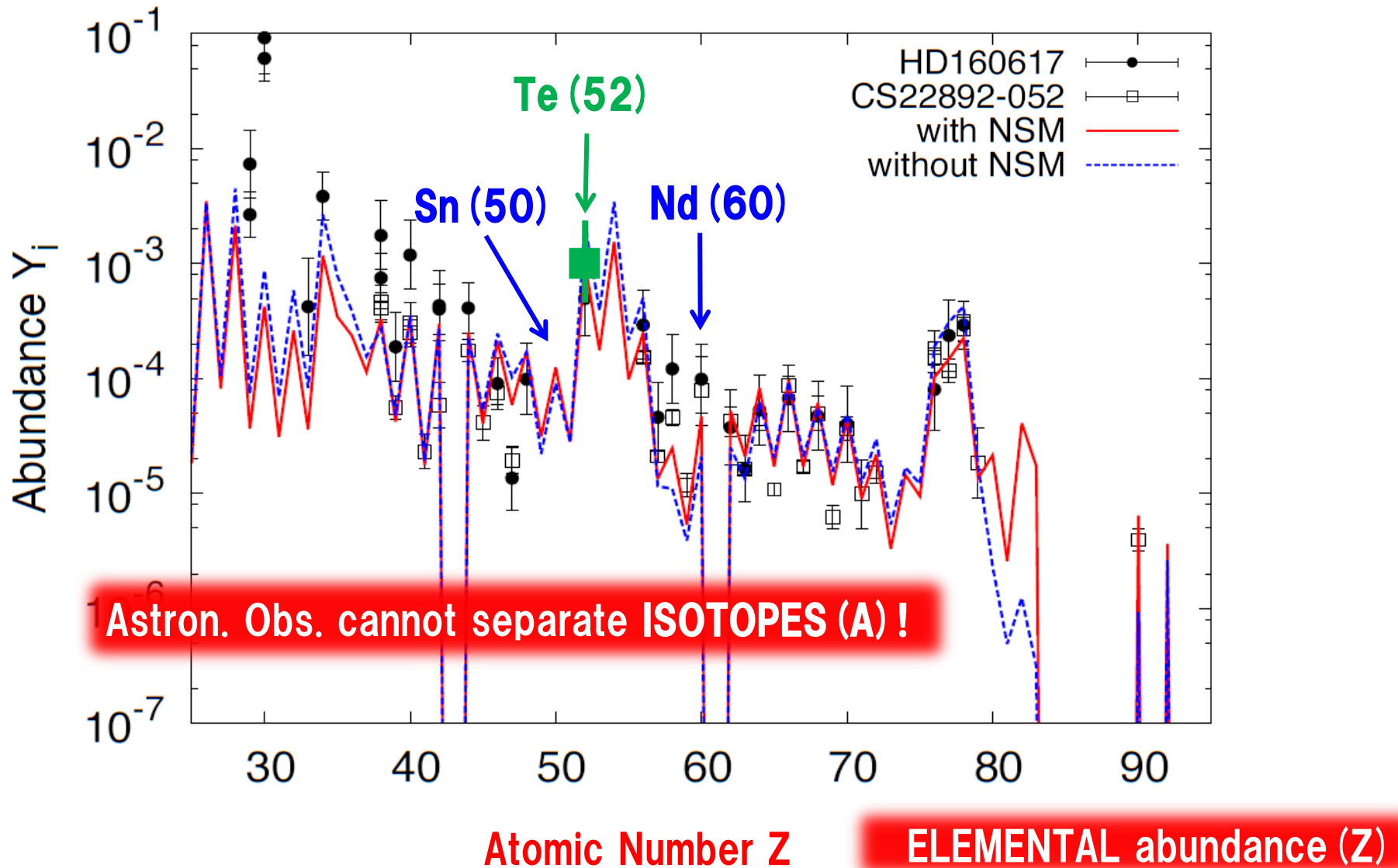
$$\text{Log} \frac{\text{Fe}/H_{\star}}{\text{Fe}/H_{\odot}}$$

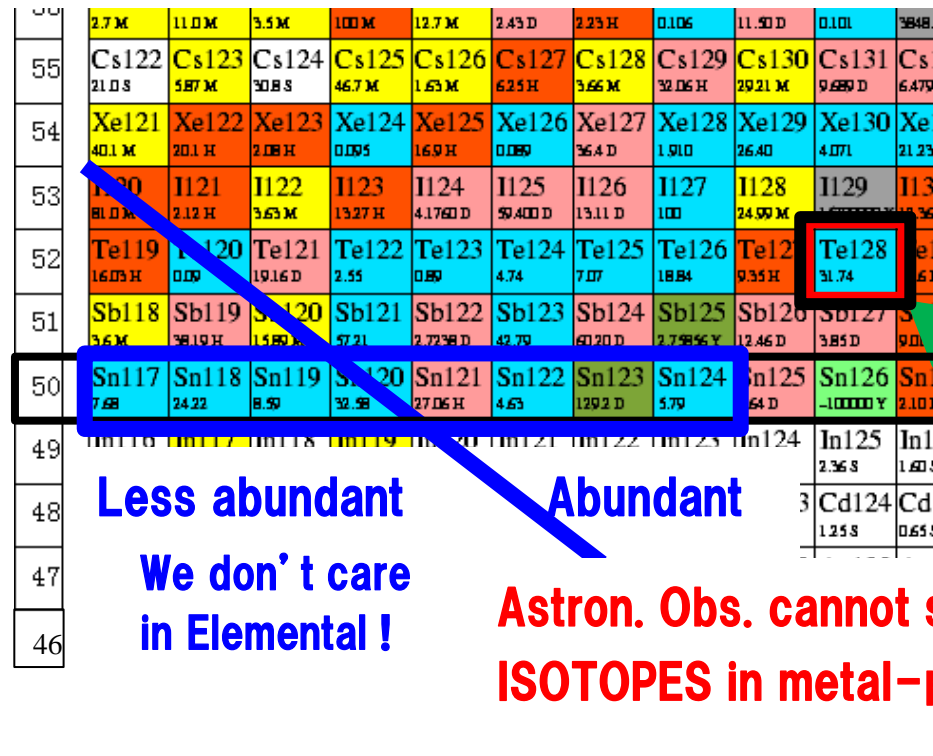
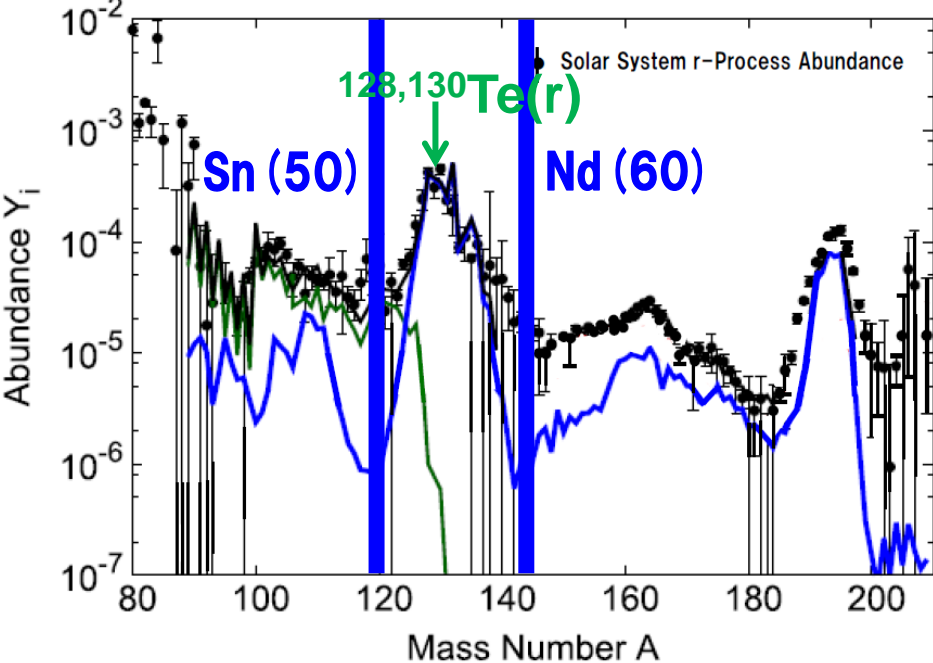


UNIVERSALITY !

Early Galaxy \longleftrightarrow TODAY

Shibagaki et al., ApJ. 816 (2016),79; Kajino & Mathews, ROPP **80** (2017) 08490.





Less abundant
We don't care
in Elemental!

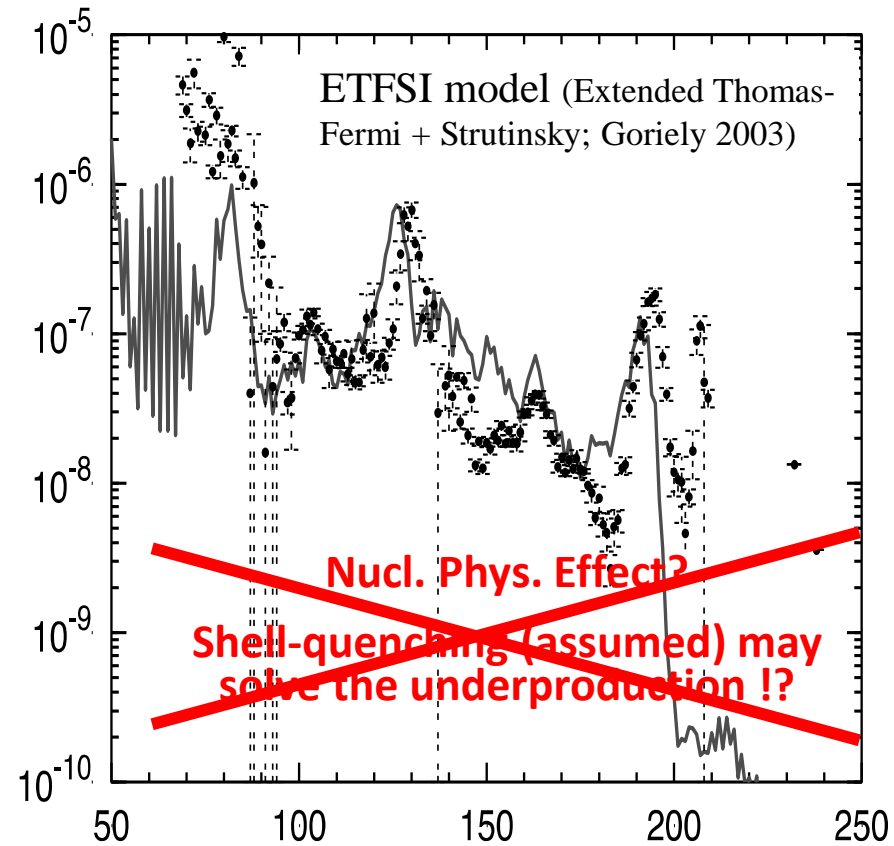
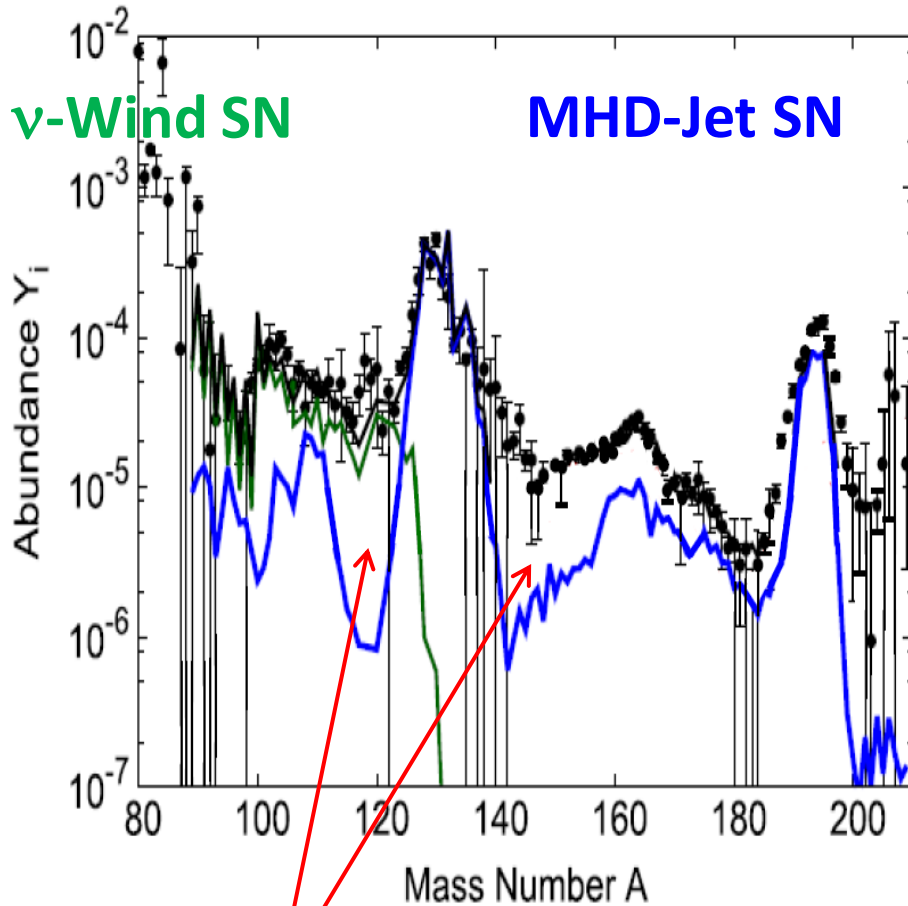
Abundant
Astron. Obs. cannot separate
ISOTOPES in metal-poor stars!

R-process
path

N=82

1. Supernovae (ν -driven & MHD Jet)

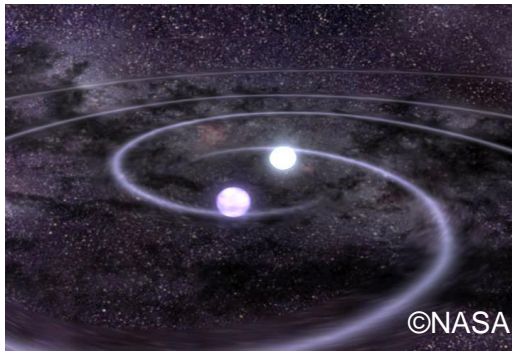
Shibagaki, Kajino, Mathews, Chiba, Nishimura, Lorusso, ApJ 816, 79; Kajino & Mathews (2017), ROPP 80, 084901; Kajino, Aoki, Balantekin, Dihel, Famiano, Mathews (2019), PPNP 107, 109.



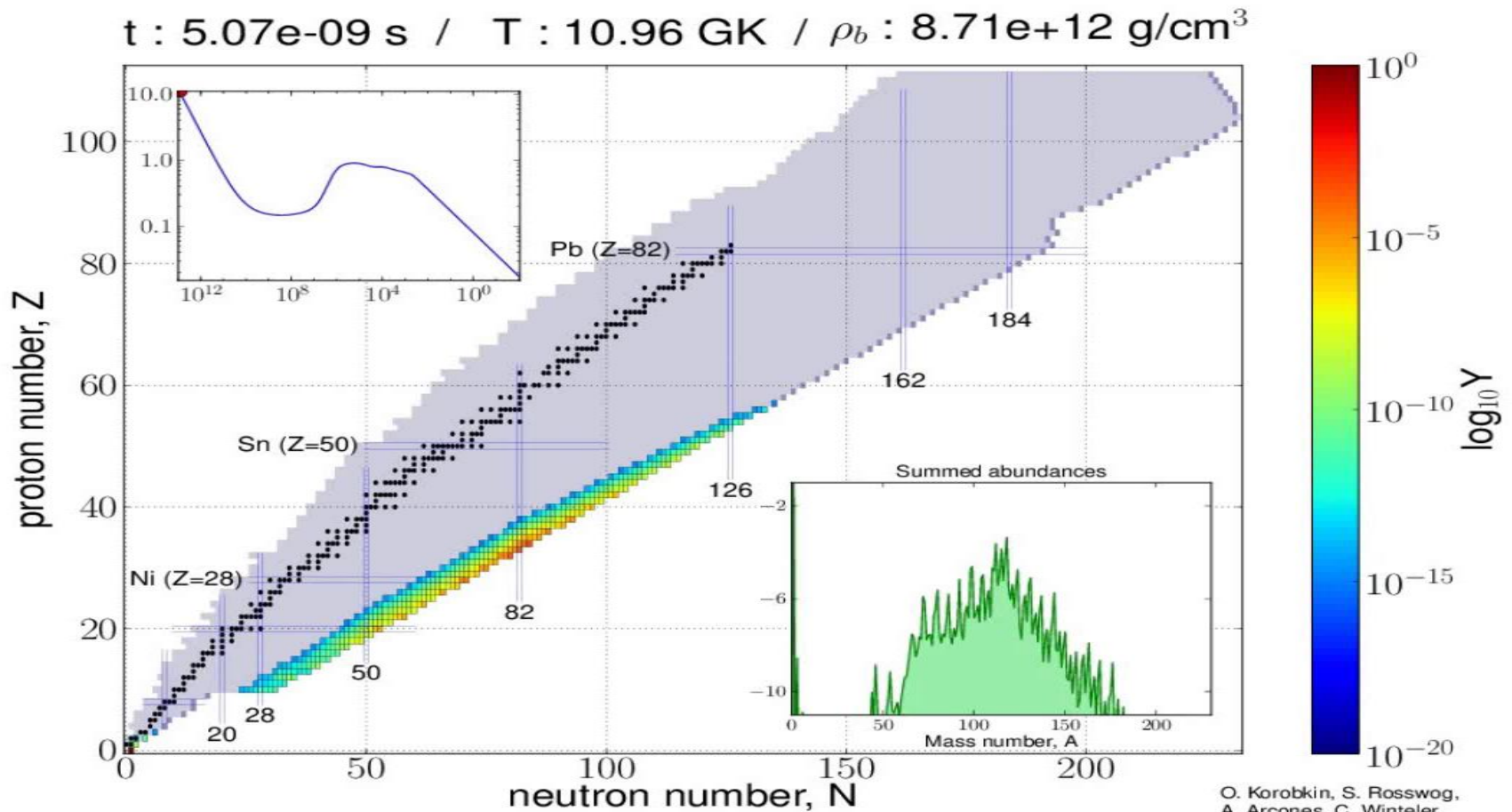
Underproduction \rightarrow Possible Solution

- ~~1. Nucl. Phys. - Shell Quenching ?~~
2. Astro. - Another Site ?
(Merger or Collapsar ?)

2. Neutron Star Merger



Shibagaki, Kajino, Mathews, Chiba, Nishimura, Lorusso, ApJ 816, 79;
Kajino & Mathews (2017), ROPP 80, 084901;
Kajino, Aoki, Balantekin, Dihel, Famiano, Mathews (2019), PPNP 107, 109.



Shape Parametrization

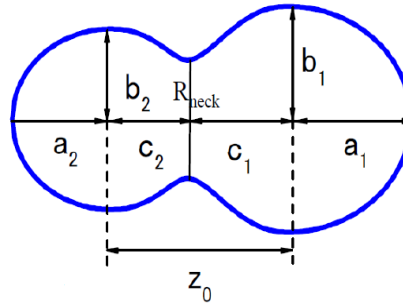
Quantum Tunneling → Dynamical Fission → 3D & 4D Lengevin Eq.

Ishizuka et al., Phys. Rev. C96 (2017), 064616; Ivanyuk et al., Phys. Rev. C97, 054331 (2018);
Okumura et al., J. Nucl. Sci. Tech. 55, 1009 (2018); Usang et al., Sci. Reports, 9, 1525(2019)

Shape parametrization

Two-center model

(Maruhn and Greiner, Z. Phys. 251(1972) 431)



Quantum Tunneling
WKB

Free energy surface F

$$F = E - TS = E \text{ (at } T = 0)$$

Collective coordinates (3 or 4 dynamical variables)

$$\{q\}_{3D} = \{ZZ_0, \delta, \alpha\} \quad \{q\}_{4D} = \{ZZ_0, \delta_1, \delta_2, \alpha\}$$

FRDM+Strutinsky+BCS

★ $ZZ_0 = \frac{z_0}{R}$ Elongation

R : Radius of compound nucleus $= 1.2 A_{CN}^{1/3}$

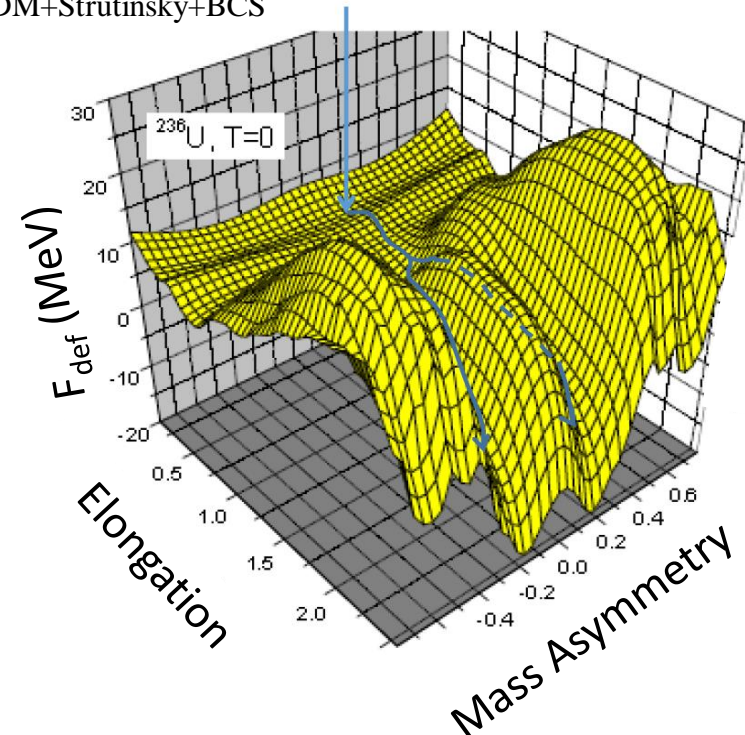
★ $\delta_i = \frac{3(a_i - b_i)}{2a_i + b_i}$ Deformation ($\delta_1 \neq \delta_2 \rightarrow$ 4D Lengevin Eq.)

3D : $\delta_1 = \delta_2 = \delta$ 4D : δ_1, δ_2 are independent

★ $\alpha = \frac{A_1 - A_2}{A_1 + A_2}$ Mass asymmetry A_1 : mass of the right fragment
 A_2 : mass of the left fragment

$\varepsilon = 0.35$ neck parameter : fixed

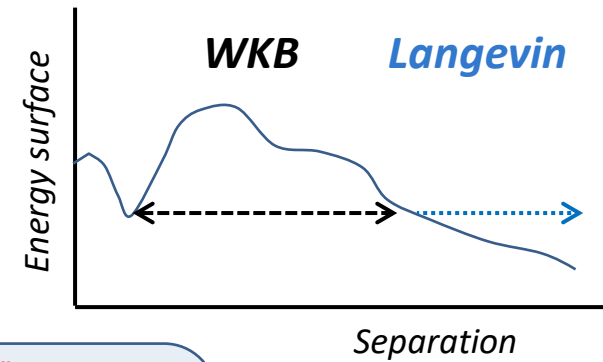
volume conservation condition is applied



Dynamical Fission → 3D & 4D Langevin Eq.

Ishizuka, Chiba et al., Phys. Rev. C96 (2017), 064616.

4 Collective Coordinates $\{q_i; i = 1..4\} = \{ZZ_n, \alpha, \delta_l, \delta_p\}$



$$\left\{ \begin{aligned} \frac{dp_i}{dt} &= - \frac{\partial F}{\partial q_i} - \frac{1}{2} \frac{\partial}{\partial q_i} (m^{-1})_{jk} p_j p_k - \gamma_{ij} (m^{-1})_{jk} p_k + g_{ij} R_j(t) \\ \frac{dq_i}{dt} &= (m^{-1})_{ij} p_j \end{aligned} \right.$$

Drift term Friction term Wiener term

F : Helmholtz' free energy, $F = E - TS$
 2-center Woods-Saxon model

q_i : Nuclear shape motion

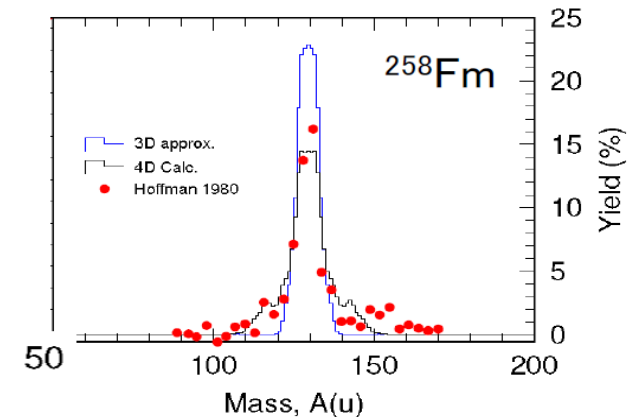
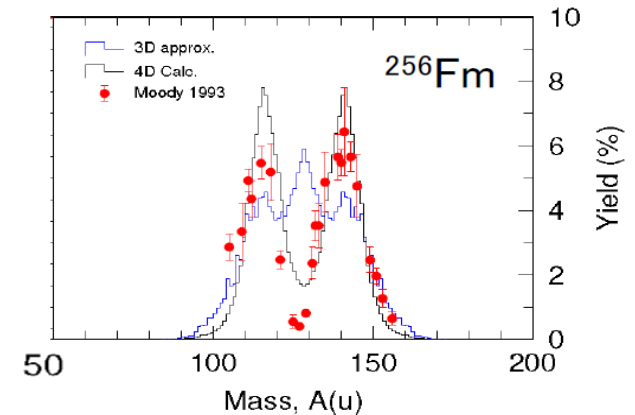
p_i : Momentum conjugate to q_i

m_{ij} : Inertia tensor } Hydrodynamical model
 γ_{ij} : friction tensor }

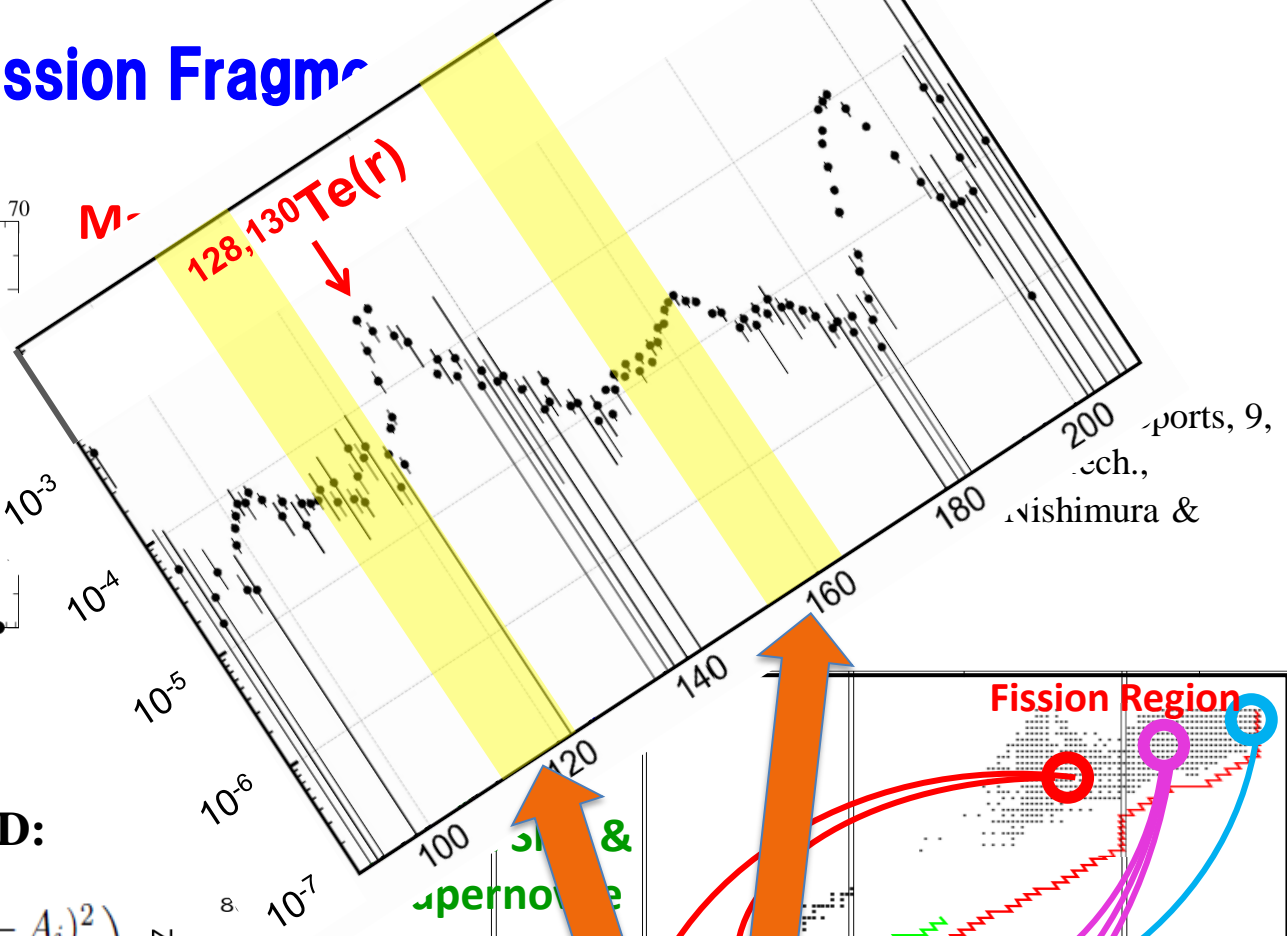
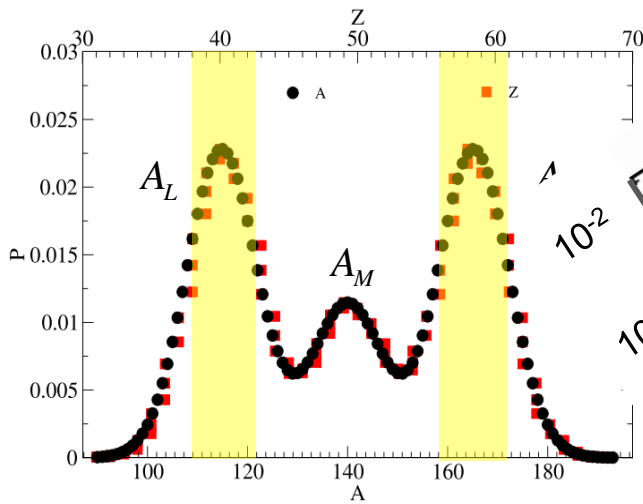
$g_{ij}g_{ij} = \gamma_{ij}T$: Fluctuation dissipation theorem
 (+Einstein relation)

$$T = \sqrt{\frac{E^* - \frac{1}{2} m_{ij} p_i p_j - E_{rot}}{a}}$$

E^* : Total excitation energy of the system



Fission Fragments



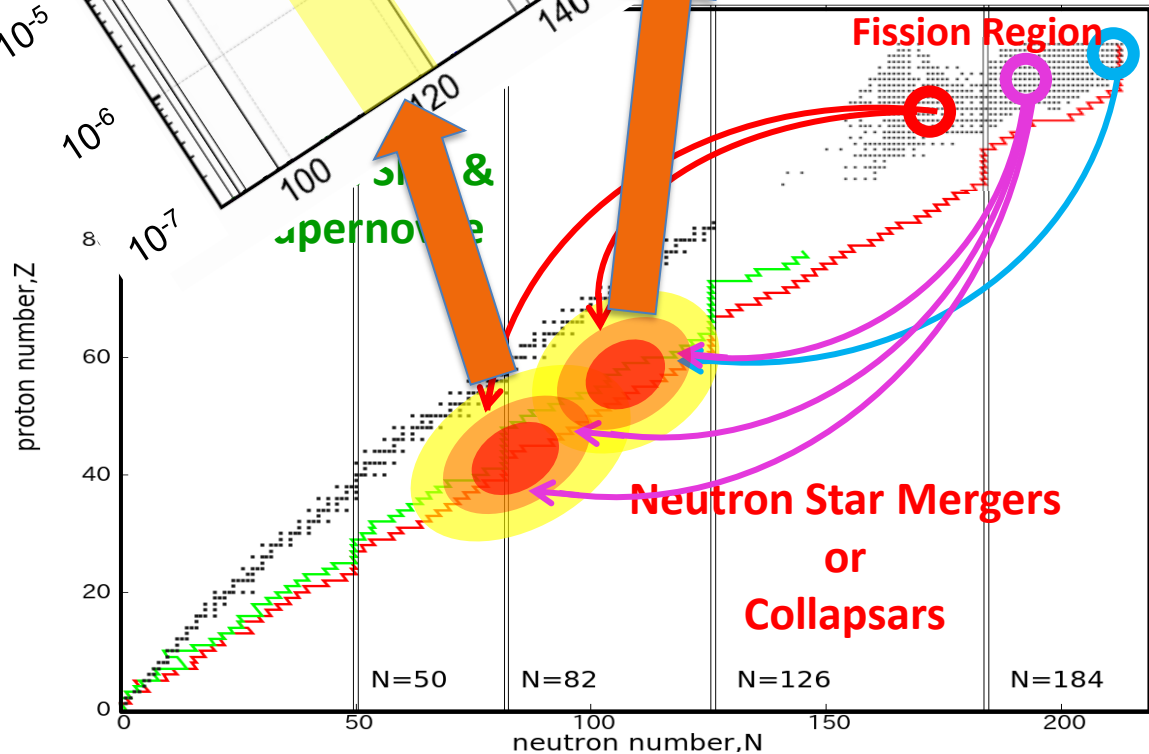
Bimodal or Trimodal FFD:

$$f(A, A_p) = \sum_{A_i} \frac{1}{\sqrt{2\pi}\sigma} W_i \exp\left(\frac{-(A - A_i)^2}{2\sigma^2}\right)$$

$$A_H = (1 + \alpha)(A_p - N_{loss})/2$$

$$A_L = (1 - \alpha)(A_p - N_{loss})/2$$

$$A_M = (A_H + A_L)/2.$$



Fission Region

Neutron Star Mergers
or
Collapsars

M-
128, 130Te(r)

ports, 9,
ech.,
ishimura &

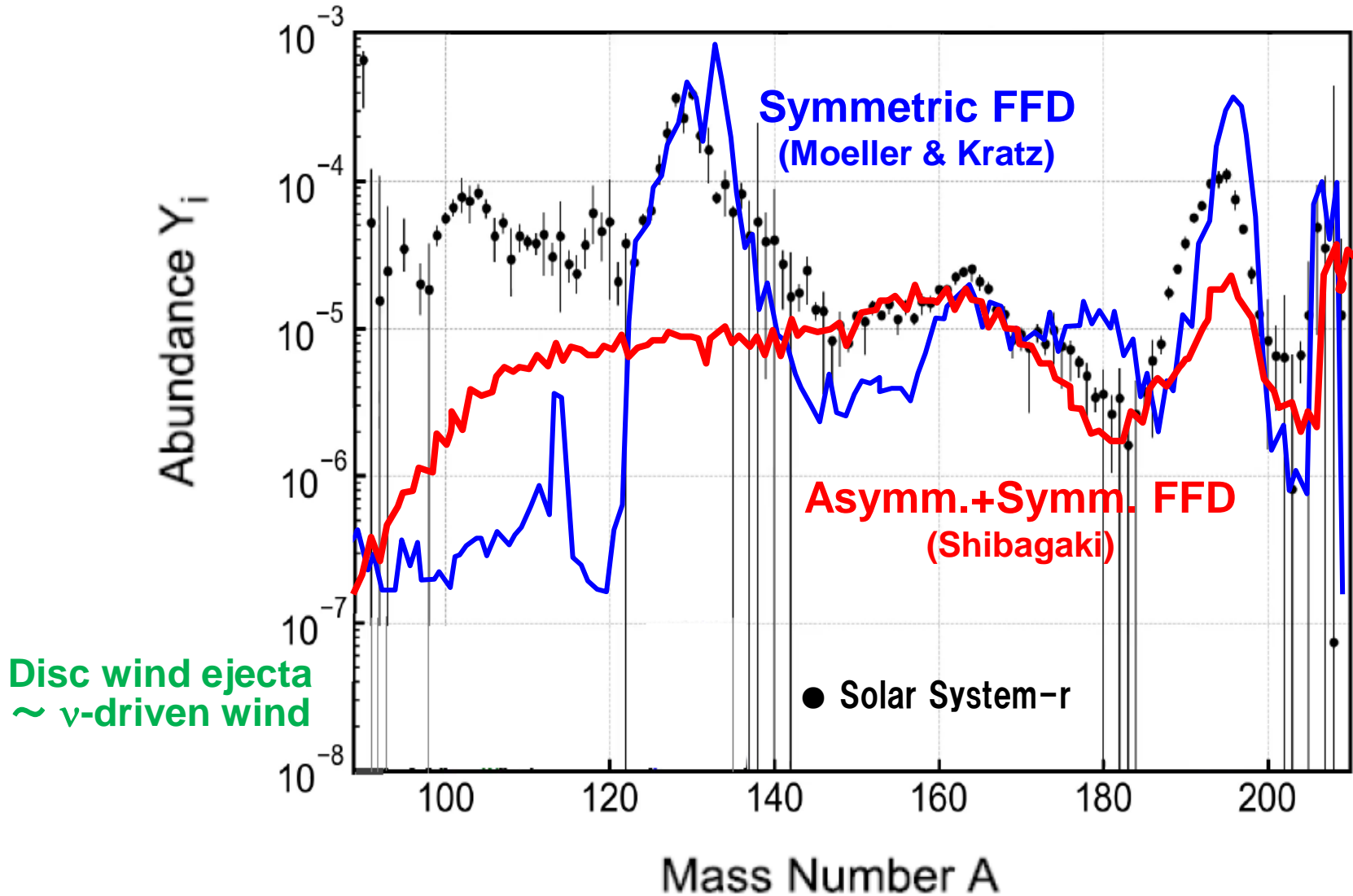
N=50 N=82 N=126 N=184

neutron number, N

proton number, Z

Neutron Star Merger r-Process

Shibagaki, Kajino, Mathews, Chiba, Nishimura, Lorusso, ApJ 816, 79; Suzuki, et al., ApJ 859 (2018), 133; Kajino & Mathews (2017), ROPP 80, 084901; Kajino, Aoki, Balantekin, Dihel, Famiano, Mathews (2019), PPNP 107, 109.



3. Collapsar Jet

Numerical Hydrodynamic Model

MacFadyen, Woosley, ApJ 524 (1999), 262;
 Nagataki et al., ApJ 659 (2007), 512.

2D Collapsar Hydrodynamic Model

Takiwaki et al., ApJ 691 (2009), 1360;
 Harikae et al., ApJ 704 (2009), 354; ApJ 713 (2010) 304;

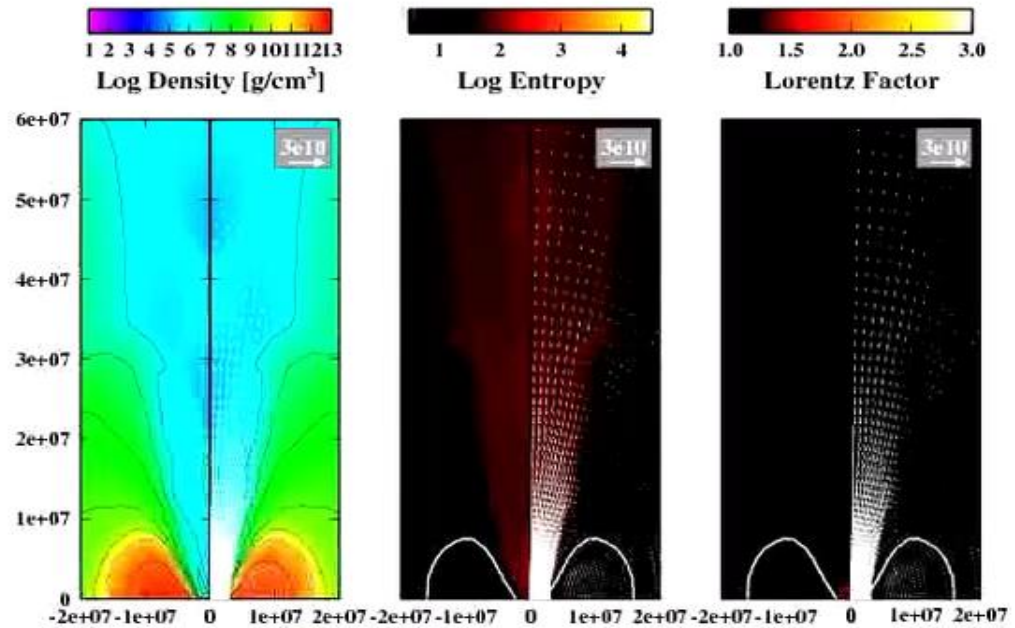
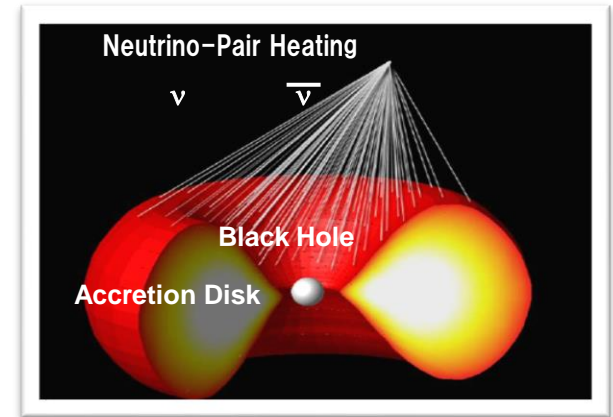
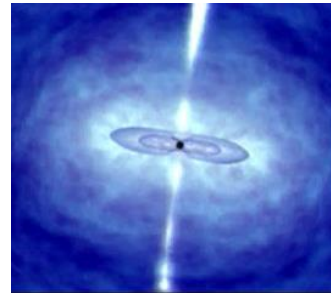
Nakamura, Kajino, Mathews, Sato, Harikae, A&Ap 582 (2015), A34.

- **Collapsar = Failed Supernova**
- Produces a **Black Hole** and a high temperature **Accretion Disk**.
- **MHD + ν -heating triggers Energetic Jet.**
- **A Model of central engine for Long-GRB**



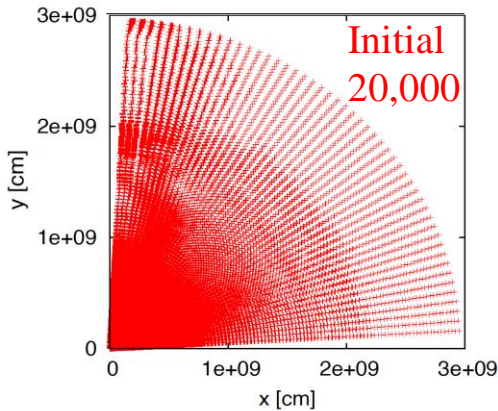
Hypernova = Super-Luminous SN

Siegel, Barnes & Metzger,
 Nature 569 (2019), 243.

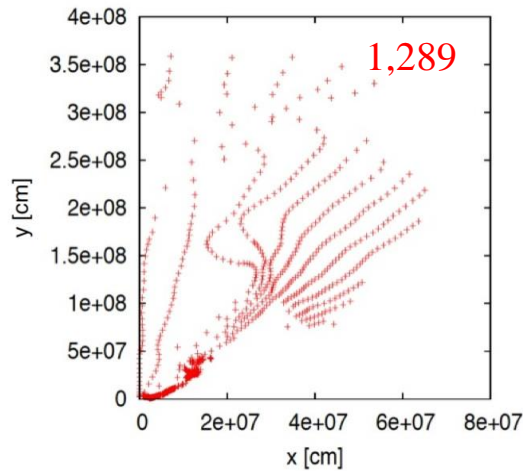


Modeling the r-process

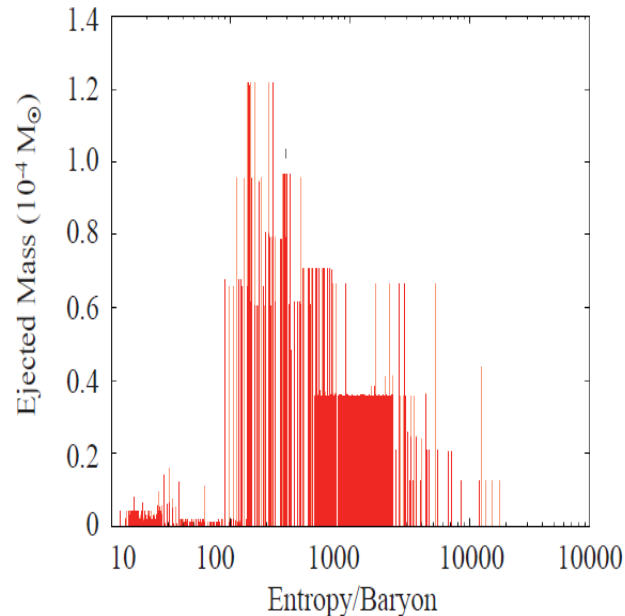
Surman et al. 2008; Fujimoto et al. 2008; Ono et al. 2012;
Nakamura, Kajino, Mathews, Sato, Harikae, A&Ap 582 (2015), A34.



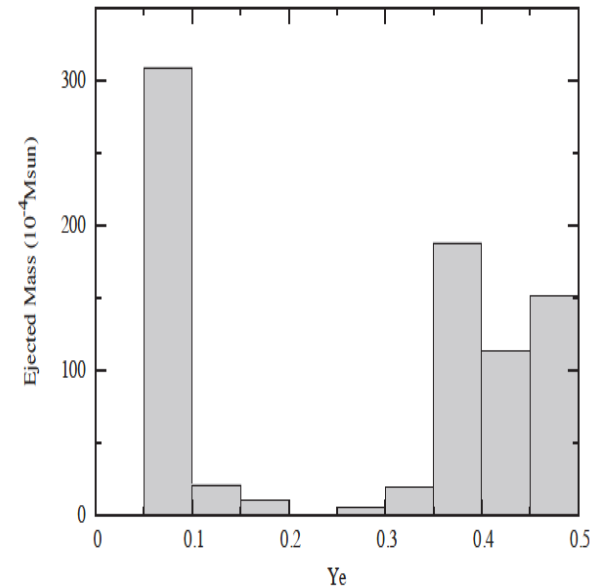
- **Extend the jet beyond the MHD+neutrino pair heating using 2D hydo** (Takiwaki et al., 2009; Harikae et al., 2009).
- **Attach 20,000 tracer particles** to evolve the flow of material into the accretion disk and out into the jet.
- **1,289 trajectories** are ejected with positive energy.



Entropy distribution



Y_e distribution

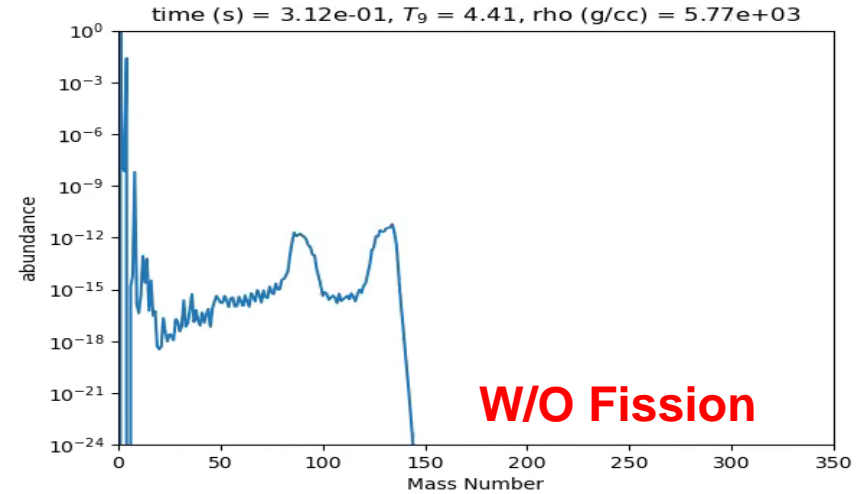
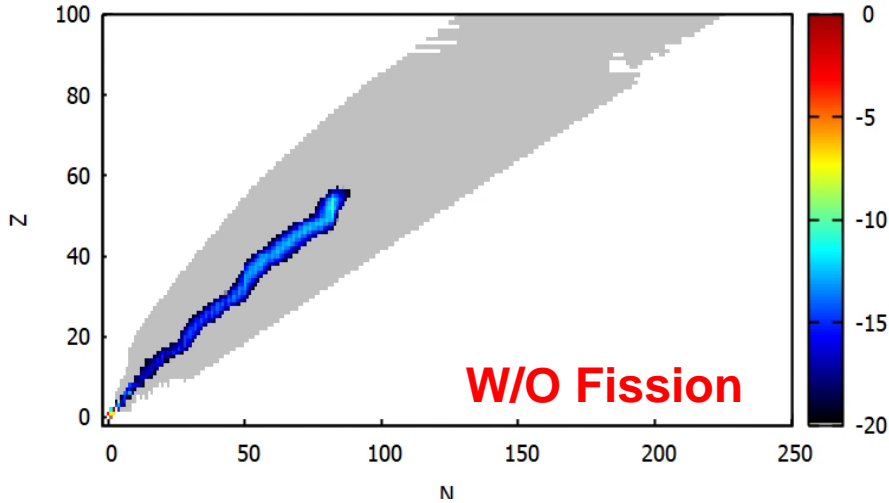


Collapsar Jet r-process

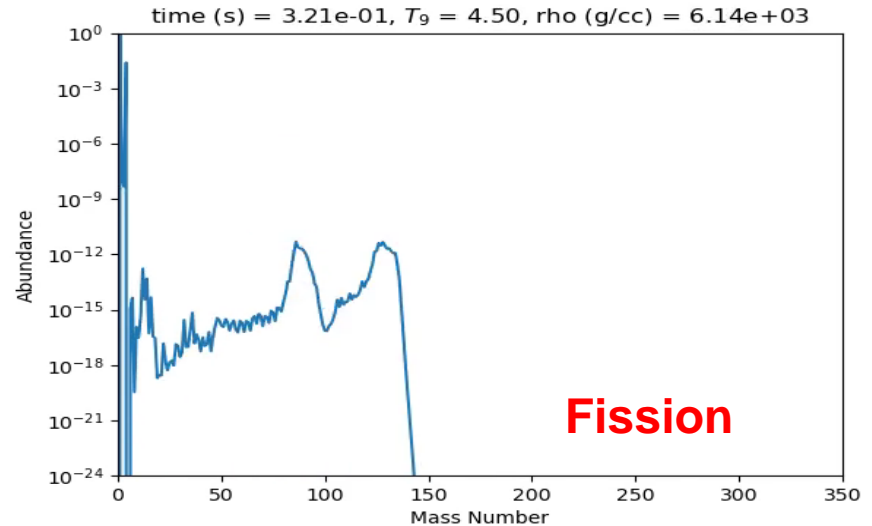
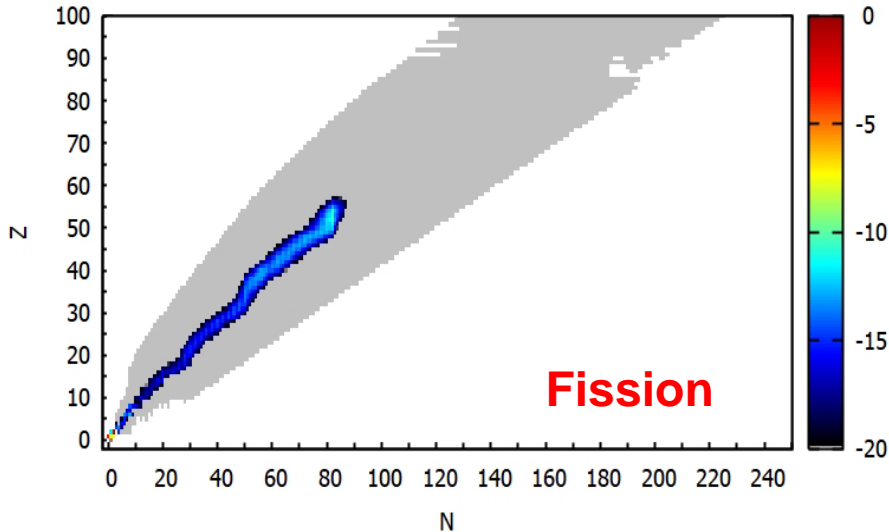
Fission Effect

Zhenyu He, M. Kusakabe, T. Kajino, G. Mathews, Y. Yamazaki, et al., (2022).

time(s) = 3.8641E-01; $t_9 = 3.8903$; $\rho(\text{gcc}) = 4.3707\text{E}+03$; $Y_e = 0.0504$; $Y_n = 0.8994$



time(s) = 3.7343E-01; $t_9 = 4.1059$; $\rho(\text{gcc}) = 4.8898\text{E}+03$; $Y_e = 0.0504$; $Y_n = 0.8995$



Strong Magnetic Field : $B \neq 0$

Relativistic Fermi-Dirac Screening of β -decays
Landau Quantization

Relativistic Fermi-Dirac Screening Effect on e's ($B = 0$)

$$n = -\frac{e}{4\pi} \int_0^\infty d^3p \left[\frac{1}{e^{(E-\mu-e\phi)/T} + 1} - \frac{1}{e^{(E+\mu+e\phi)/T} + 1} \right]$$

Relativistic Screening ($B \neq 0$) with Landau Quantization

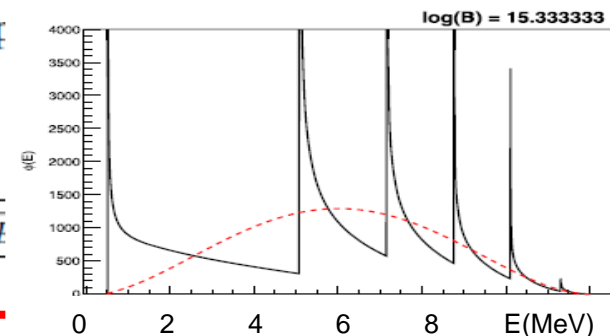
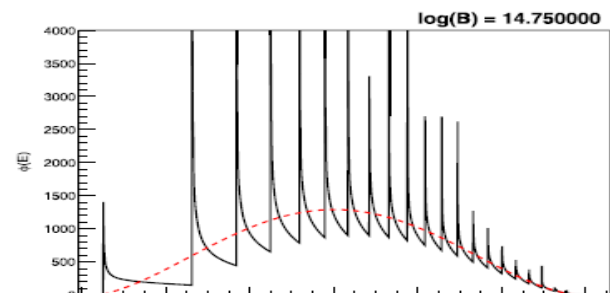
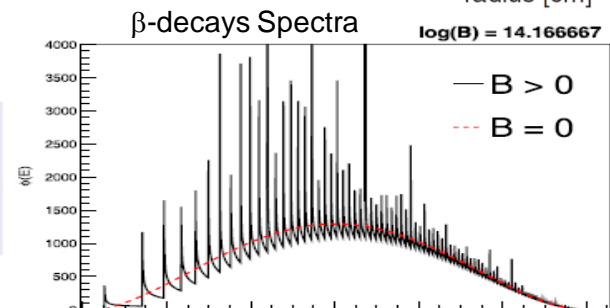
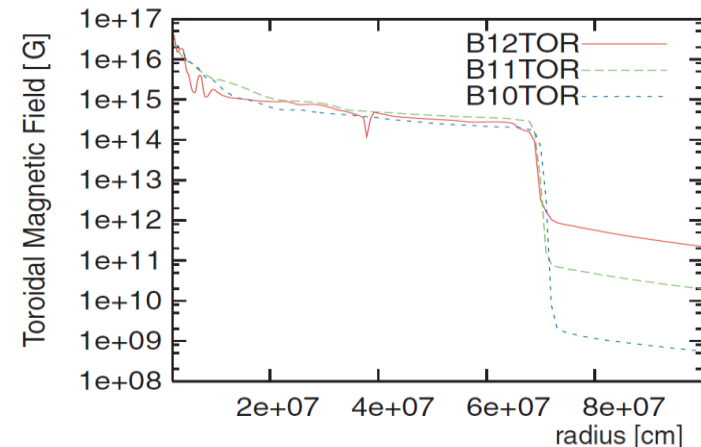
$$n = -\frac{eB}{2\pi^2} \sum_{\nu=0}^{\infty} g_\nu \int_0^\infty dp_z \left(\left[\exp \left(\frac{\sqrt{p_z^2 + m_e^2 + 2\nu eB} - \mu - e\phi}{T} \right) \right] \right)$$

Poisson-Boltzmann Equation ($B \neq 0$)

$$\nabla^2 \phi(r) = -4\pi Z e \delta^3(\mathbf{r}) - 4\pi \sum_{z>0} z e n_z \exp \left[-\frac{z e \phi_r}{T} \right] + 4\pi$$

$$+ \frac{eB}{\pi} \sum_{n=0}^{\infty} g_n \int_0^\infty dp \left[\frac{1}{\exp \left[\frac{\sqrt{E^2 + 2neB} - \mu - e\phi_r}{T} \right] + 1} - \frac{1}{\exp \left[\frac{\sqrt{E^2 + 2neB} - \mu + e\phi_r}{T} \right] + 1} \right]$$

$$- \frac{eB}{\pi} \sum_{n=0}^{\infty} g_n \int_0^\infty dp \left[\frac{1}{\exp \left[\frac{\sqrt{E^2 + 2neB} - \mu}{T} \right] + 1} - \frac{1}{\exp \left[\frac{\sqrt{E^2 + 2neB} - \mu}{T} \right] + 1} \right]$$

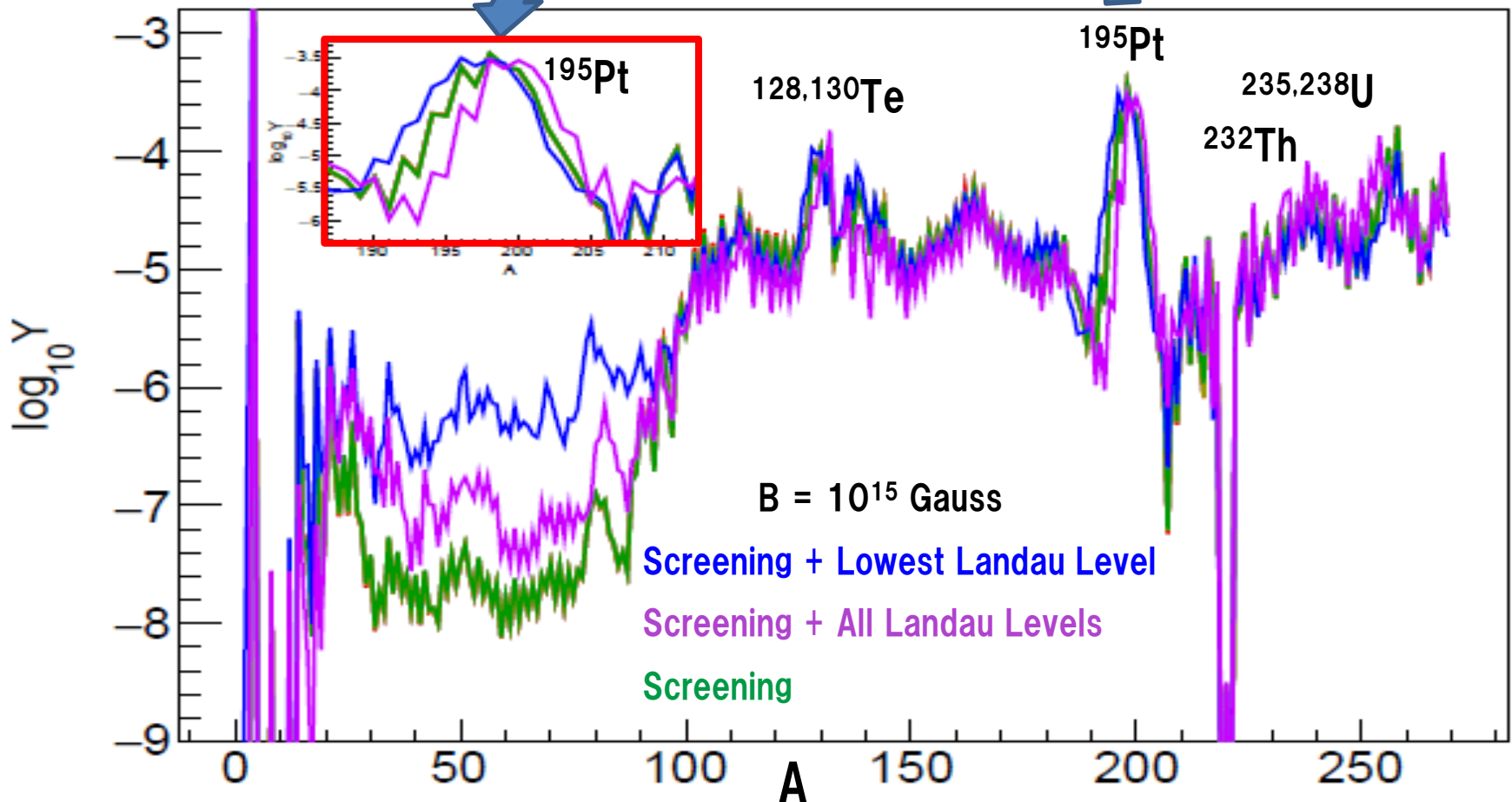


Strong Magnetic Field Effect on COLLAPSAR r-process

Famiano, Balantekin, Kajino, Kusakabe, Mori, Luo, ApJ **898** (2020), 163.

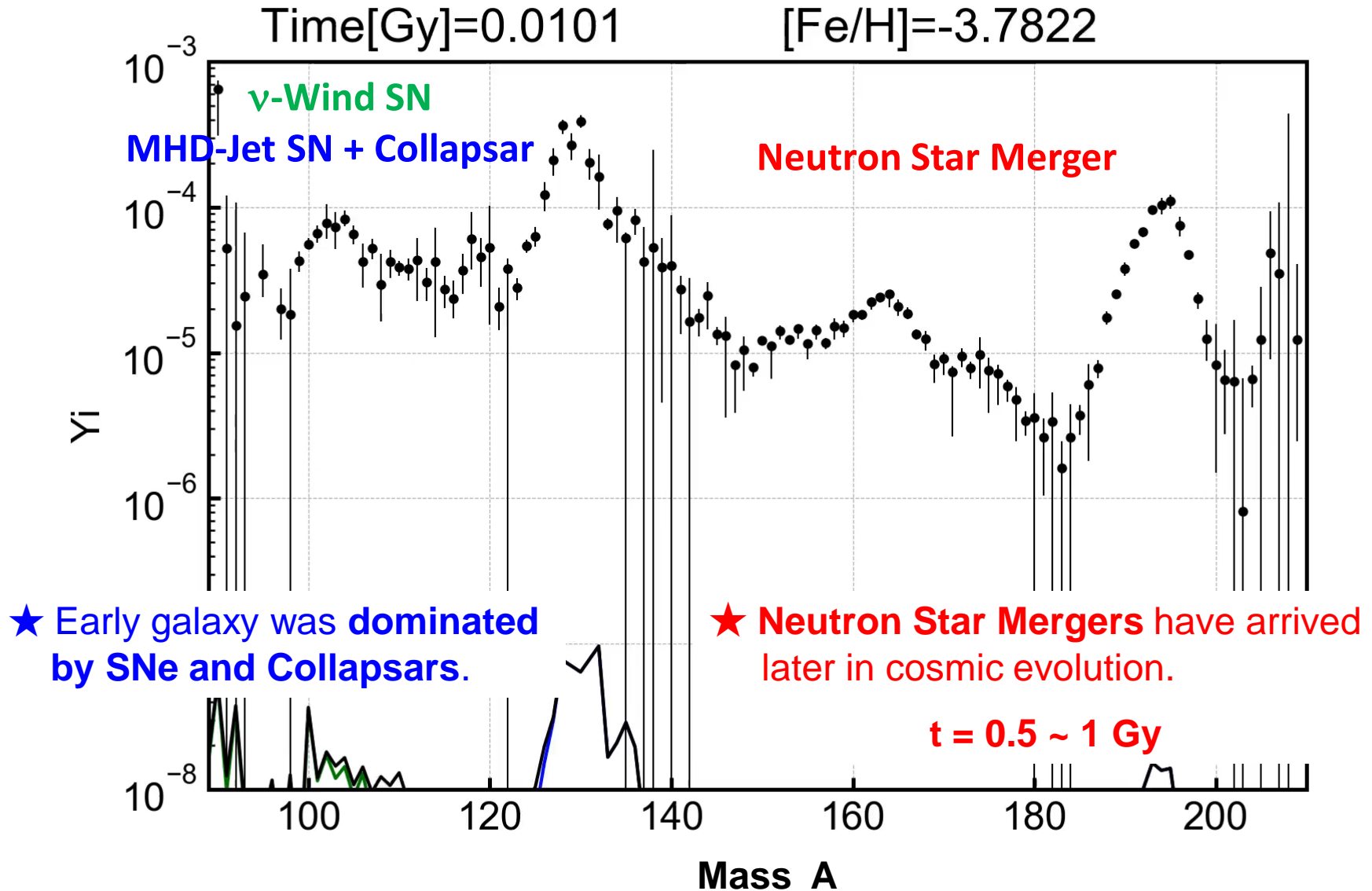
Yamazaki, He, Kajino, Mathews, Famiano, Tang, Shi, ApJ (2022), in press.

Zhenyu He, M. Kusakabe, T. Kajino, G. J. Mathews, Y. Yamazaki, et al., (2022).



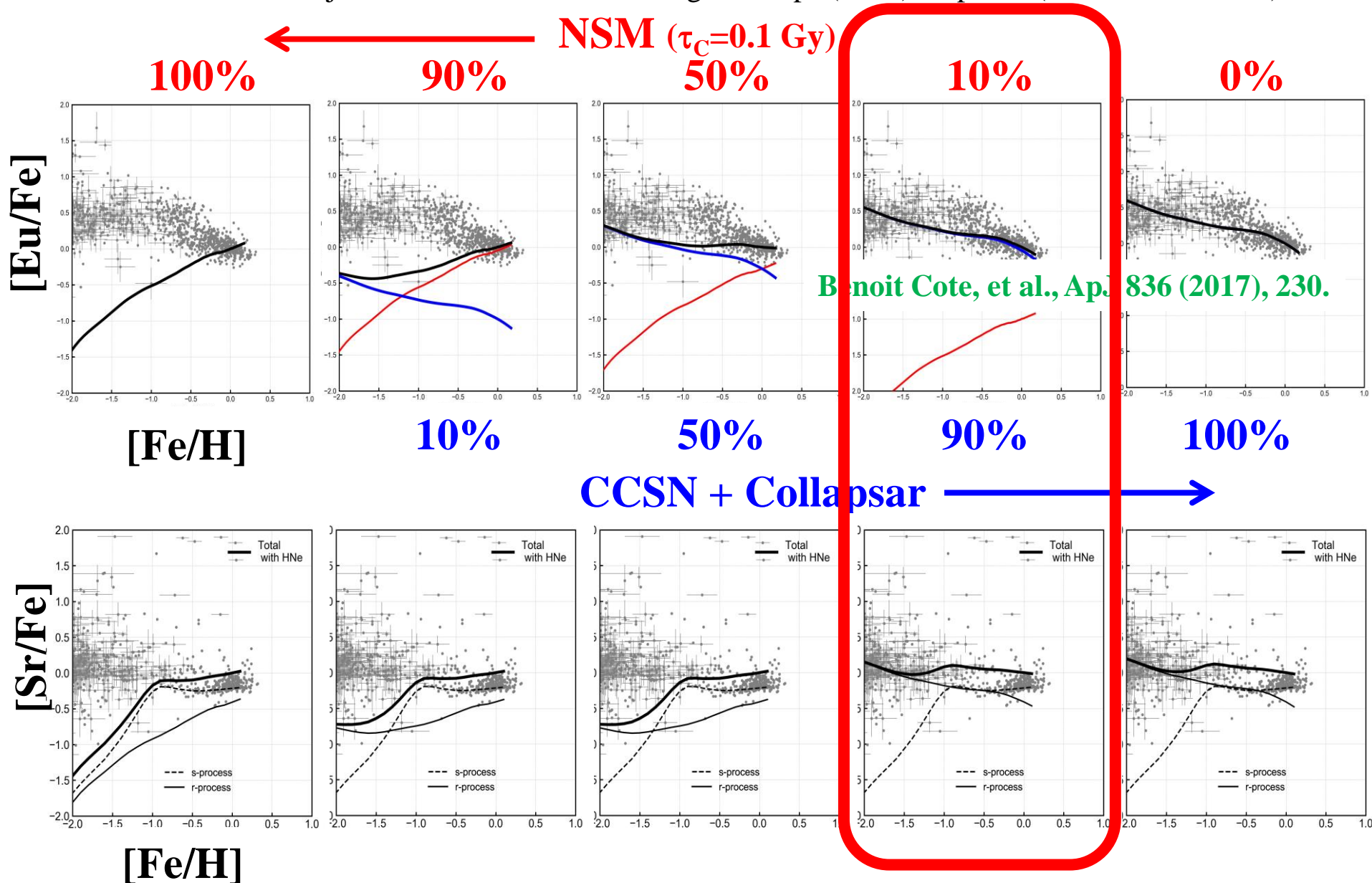
Cosmic & Galactic Evolution of R-Process Nuclei

Yamazaki, He, Kajino, Mathews, Famiano, Tang, Shi, ApJ (2022), in press. (arXiv:2102.05891)



Milky Way : Coalescence Time Delay of **NSM**

Yamazaki, He, Kajino, Mathews, Famiano, Tang, Shi, ApJ (2022), in press. (arXiv:2102.05891)

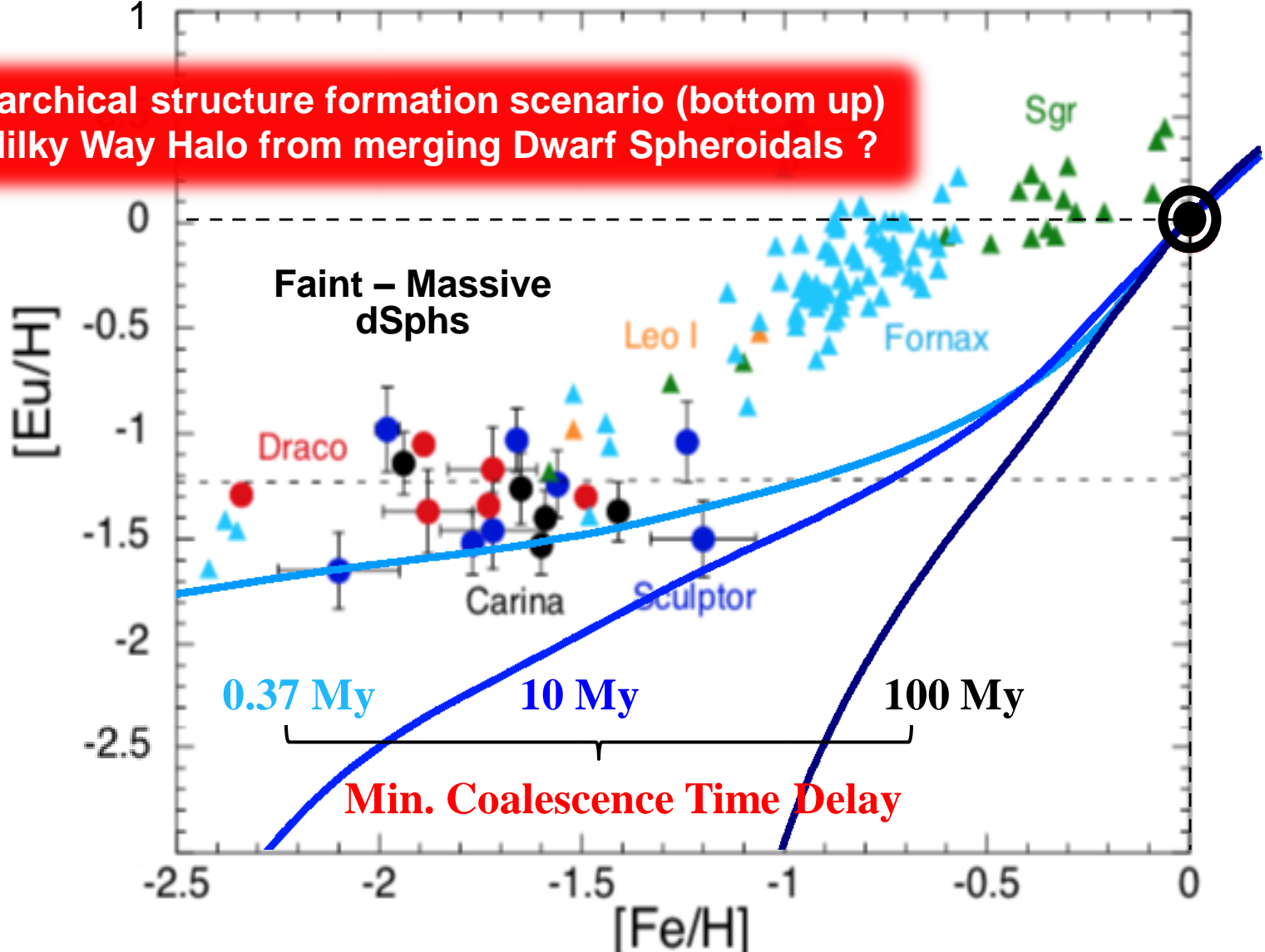


Dwarf Spheroidals

Yamazaki, He, Kajino, Mathews, Tang, Shi, Famiano et al. (2022), ApJ, in press.

1

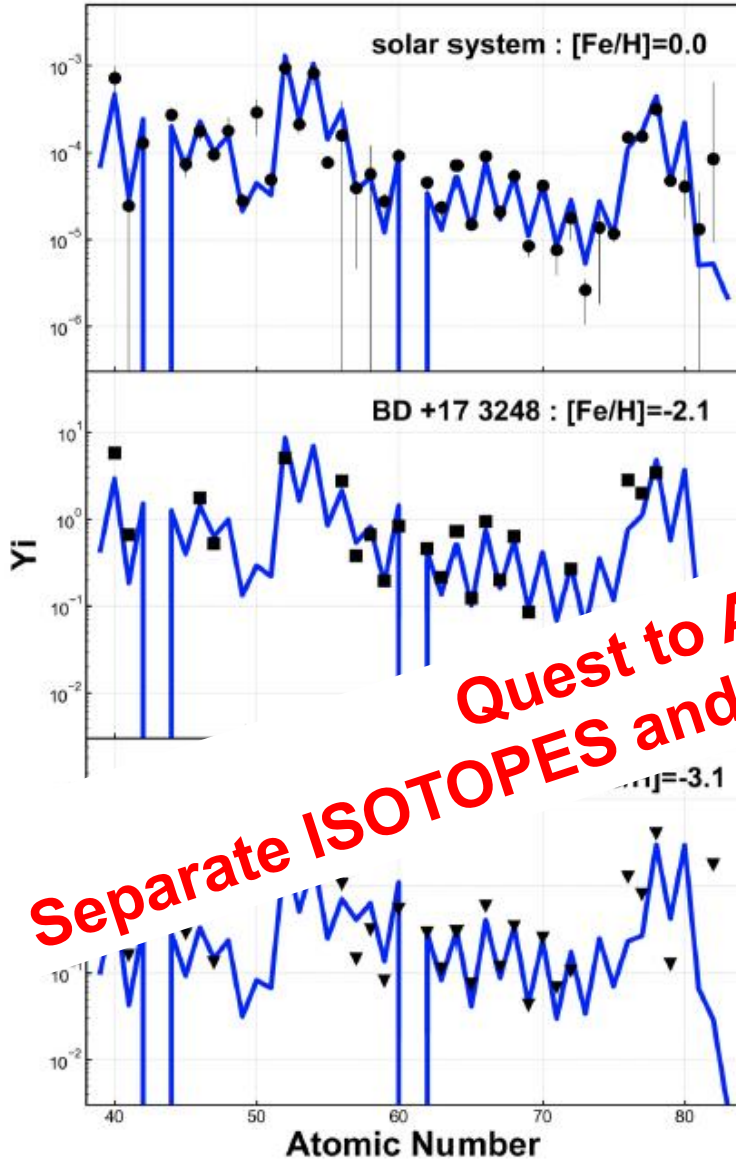
Hierarchical structure formation scenario (bottom up)
of Milky Way Halo from merging Dwarf Spheroidals ?



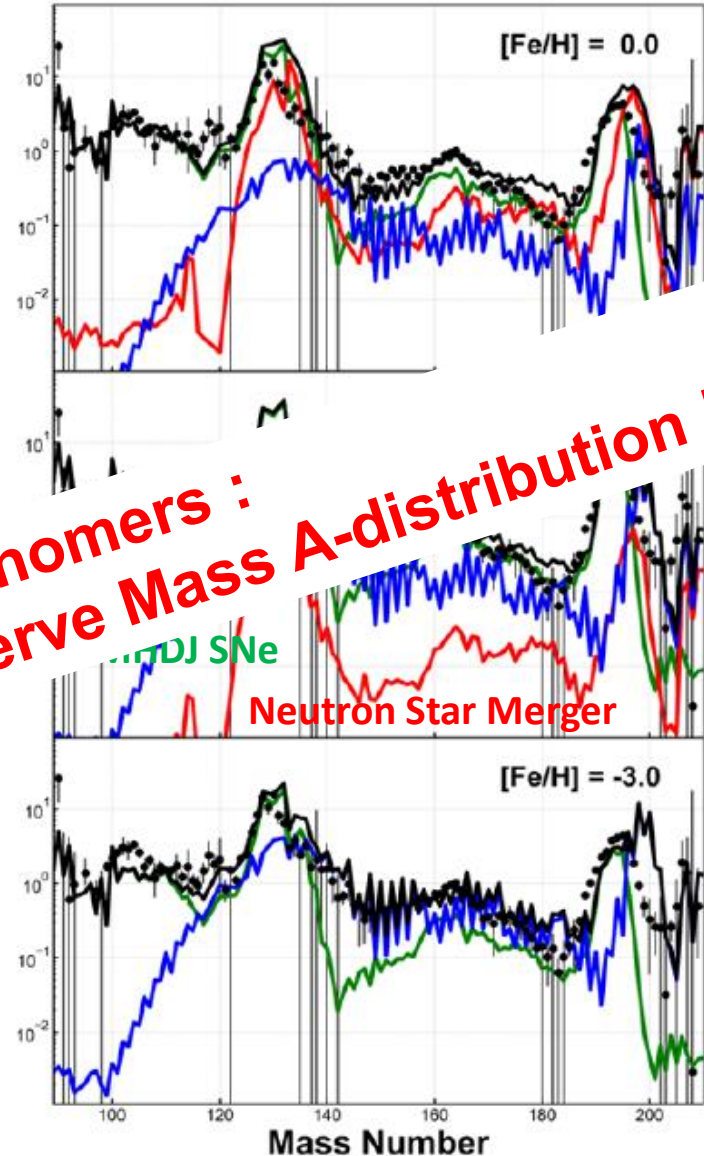
Galactic Chemical Evolution

Yamazaki, He, Kajino, Mathews, Famiano, Tang, Shi, ApJ (2022), in press. (arXiv:2102.05891)

Symm. & Asymm. fission



Symm. fission



Quest to Astronomers :
Separate ISOTOPES and observe Mass A-distribution !



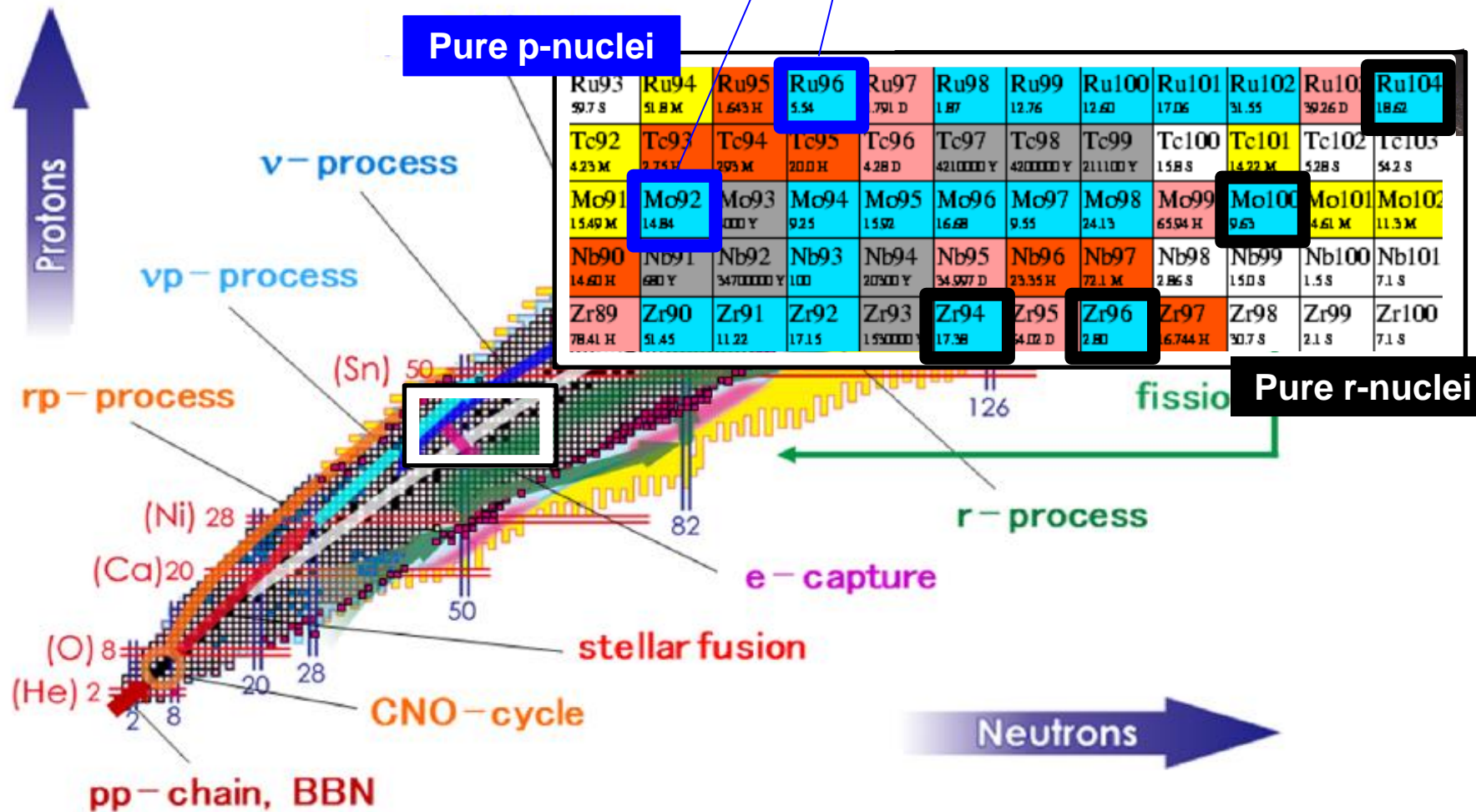
Burbidge², Fowler & Hoyle
RMP29 (1957), 547.

宇宙核物理
Guo, Hahn, Wang, Wu

Origin of p-nuclei and the roles of νp -process

Liu, Liu,

⁹²Mo 14.53% All other p-nuclei ~ 0.1-1%
⁹⁶Ru 5.54%

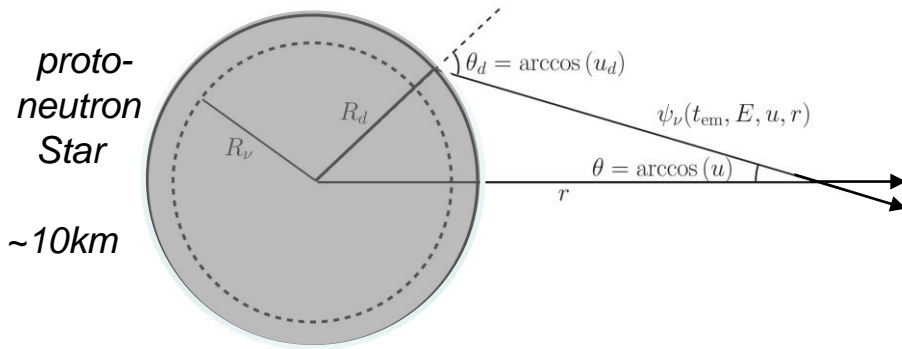


Collective + MSW ν Oscillations — Many Body Quantum Effect

Balantekin, Pehlivan & Kajino, PR D84 (2011), 065008; PR D90 (2014), 065011; PR D98 (2018), 083002
 Duan, Fuller, Carlson & Qian, PRL 97 (2006), 241101; Fogli, Lisi, Marrone & Mirizzi, JCAP 12 (2007) 010.

Sasaki, Kajino, Takiwaki, Hayakawa, Balantekin and Pehlivan, PR D96 (2017), 043013

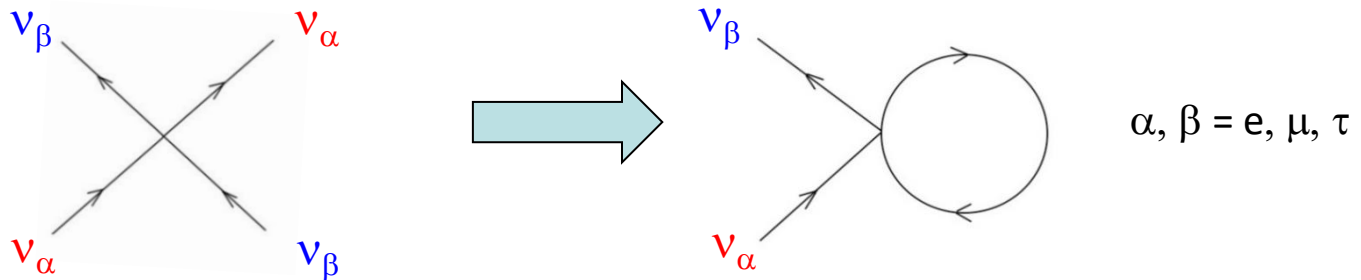
ν -sphere



$$\begin{cases} i \frac{d\psi_\nu}{dt} = (H_\nu + H_e - H_\nu) \psi_\nu(t_{em}, E, u, r), \\ H_\nu = U \frac{M^2}{2E} U^\dagger, & \text{Vacuum} \\ H_e = \sqrt{2} G_F n_e(r) \text{diag}(1, 0, 0), & \text{MSW} \end{cases}$$

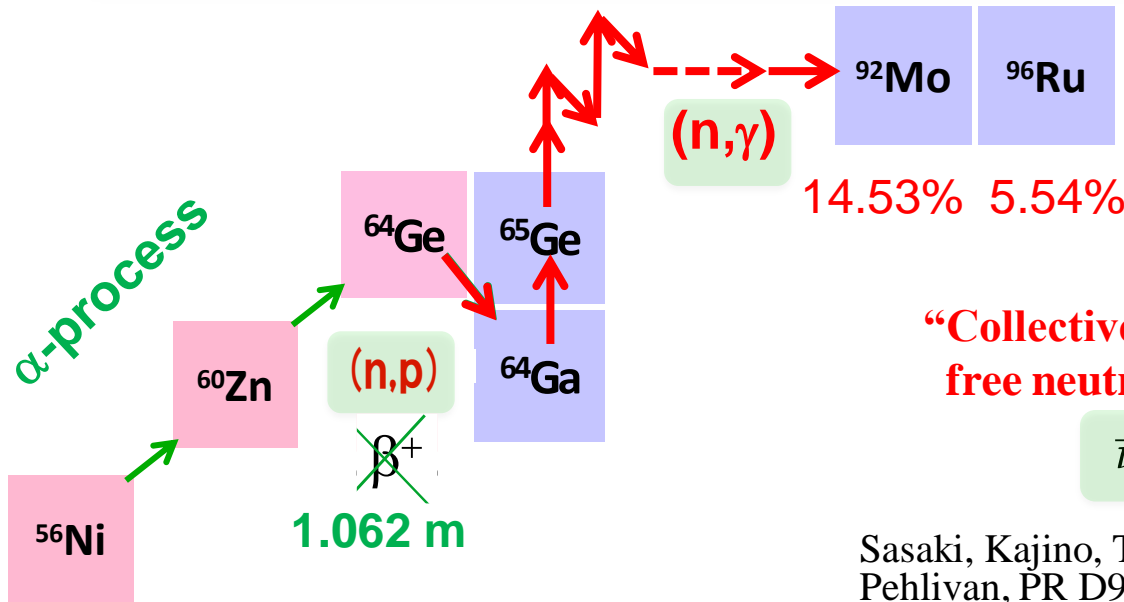
Collective flavor oscillation in coherent ν - ν scattering

$$H_\nu = \sqrt{2} G_F \sum_\alpha \int dE' d\Omega' \frac{(1 - uu')}{\nu \text{ angle dep!}} \left[\frac{d^2 n_{\nu_\alpha}}{dE' d\Omega'} \rho_{\nu_\alpha}(t'_{em}, E', u', r) - \frac{d^2 n_{\bar{\nu}_\alpha}}{dE' d\Omega'} \rho_{\bar{\nu}_\alpha}^*(t'_{em}, E', u', r) \right].$$



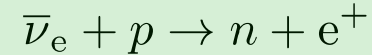
10^{48} ν 's with 3-flavors & multi-angles! \rightarrow Mean Field Approx.

vp-process including Collective Oscillation in Supernova

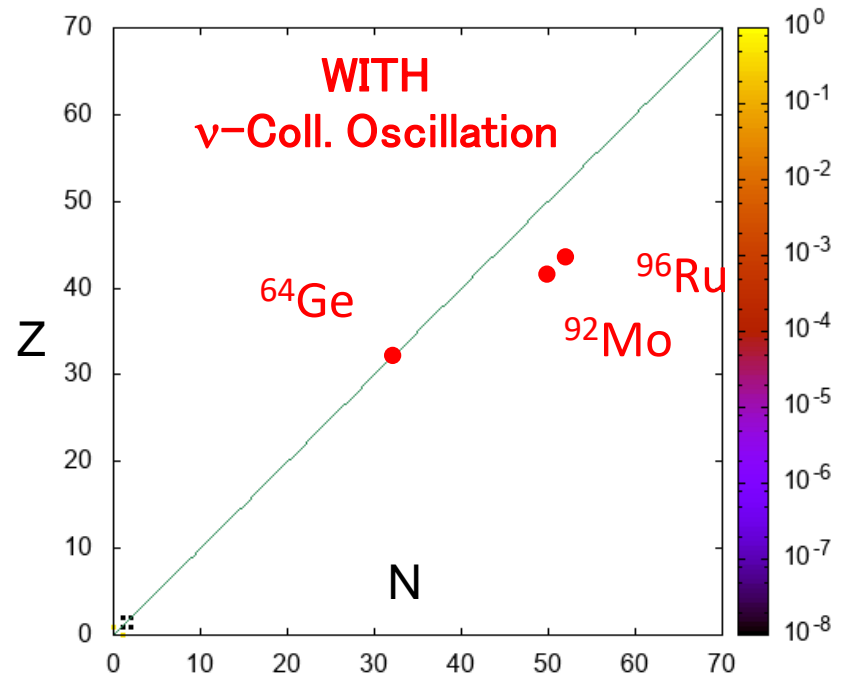
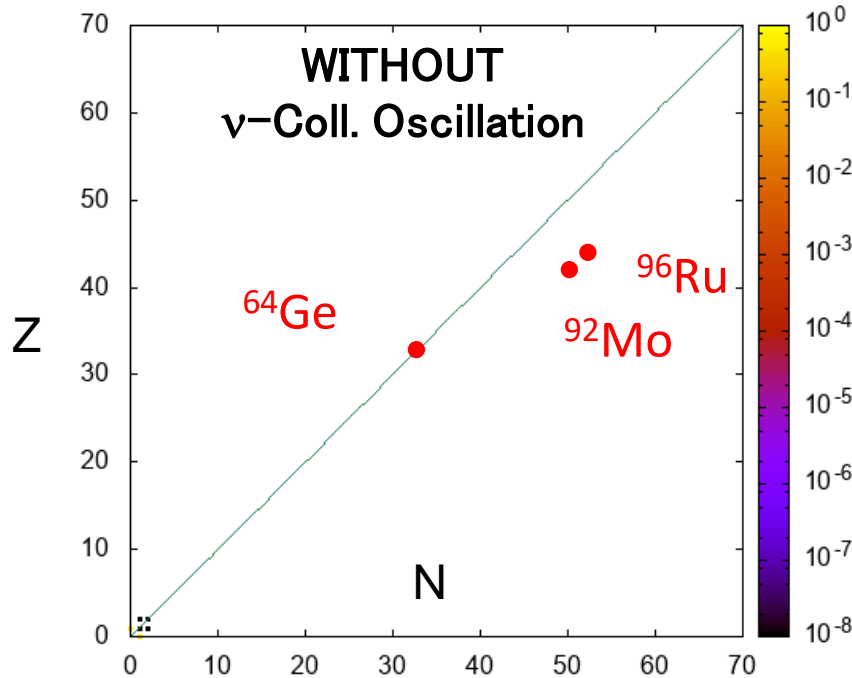


Isotopic ratio of p-nuclei
~ 0.1-1%

“Collective ν -oscillation” produces free neutrons continuously.



Sasaki, Kajino, Takiwaki, Hayakawa, Balantekin, Pehlivan, PR D96 (2017), 043013.



Cosmic & Galactic Chem. Evolution

Mo

⁹²Mo

p

14.53%

⁹⁴Mo

p + s

9.15%

⁹⁵Mo

s

⁹⁶Mo

s + r

⁹⁷Mo

s

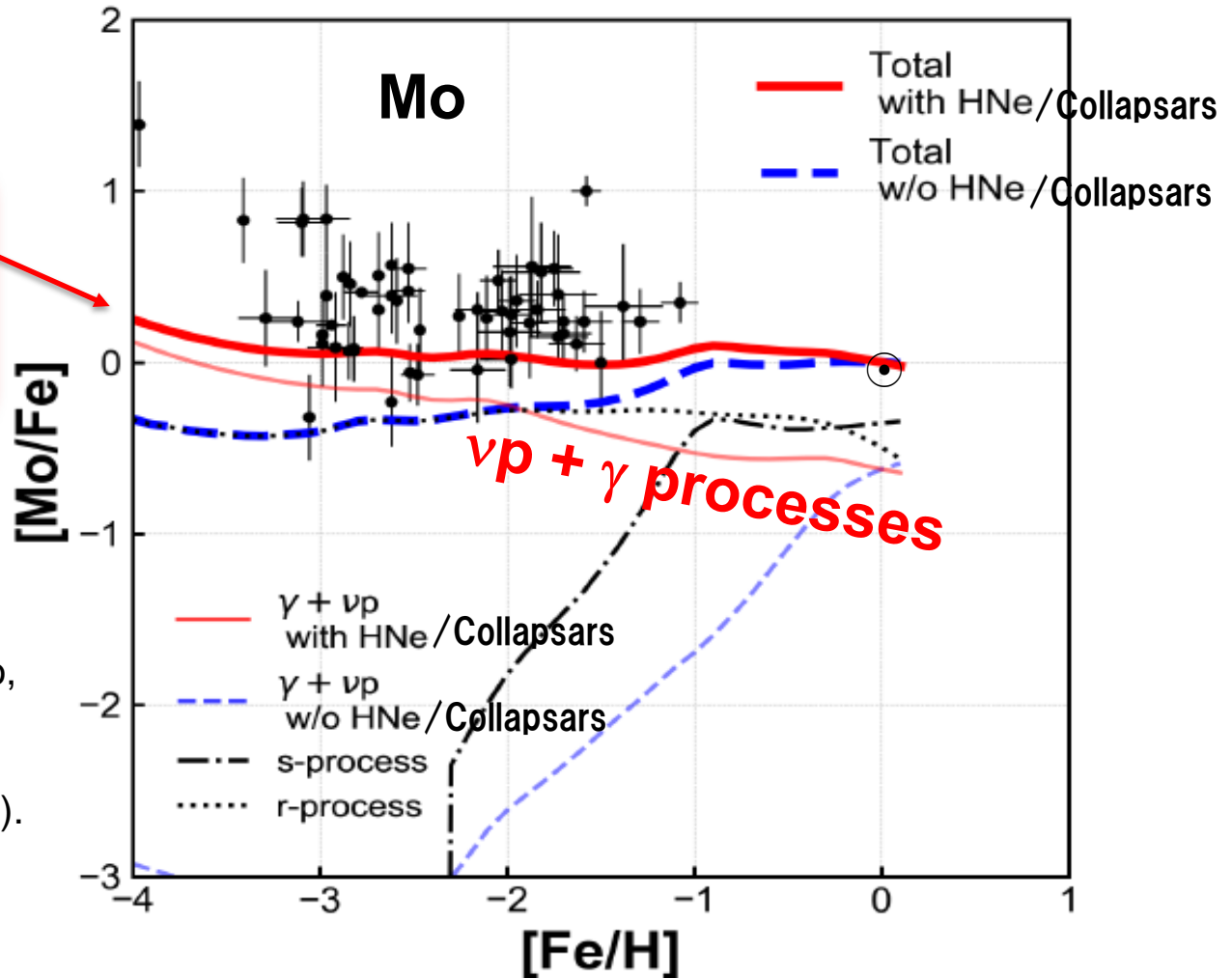
⁹⁸Mo

s + r

¹⁰⁰Mo

r

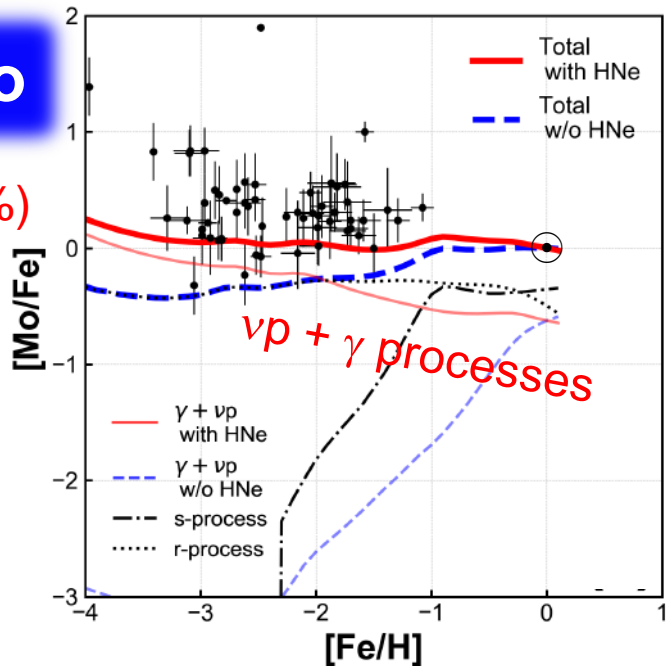
**vp-process
in Collapsar
+ MHDJ-SN**



Sasaki, Yamazaki, Kajino,
Kusakabe, Hayakawa,
Cheoun, Ko, Mathews,
ApJ 924 (2022), 29 (7 pp).

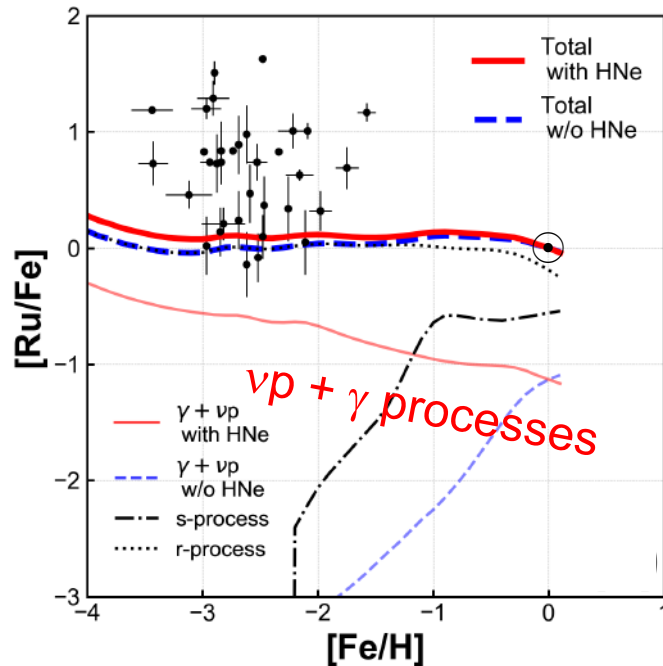
92-100Mo

(p= 23.7%)



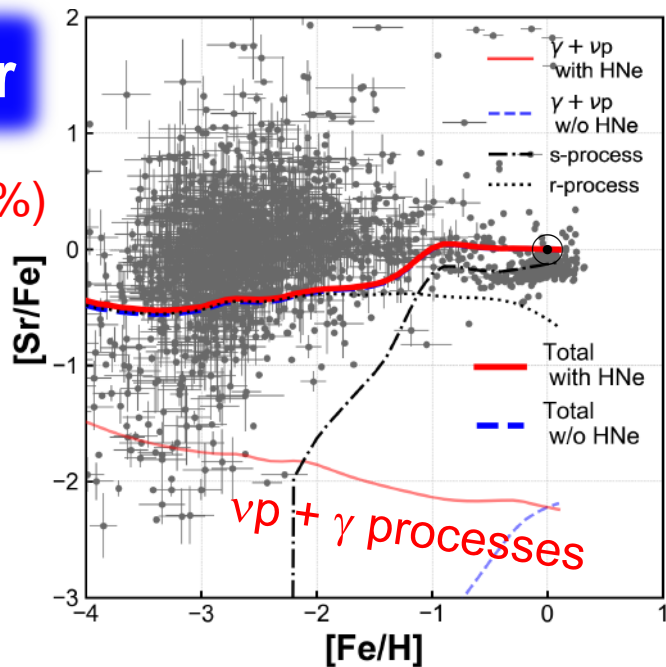
96-104Ru

(p= 7.4%)



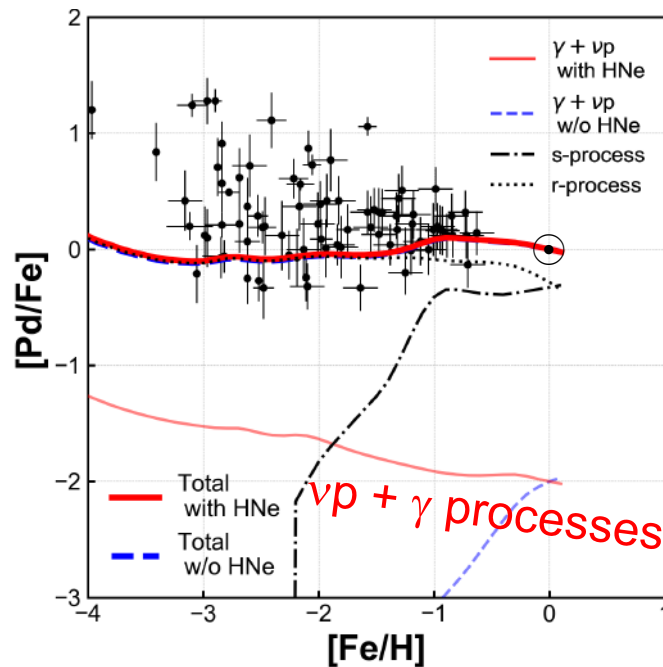
84-88Sr

(p= 0.56%)



102-110Pd

(p= 1.02%)



Summary

1. Galactic Chemical Evolution R-process Elements

Supernovae and **collapsars** are the main sites for the heavy element production over the entire history of galactic evolution. **Neutron star mergers** have arrived later and contribute partially to the solar-system elements.

2. Origin of $^{92,94}\text{Mo}$ and $^{96,98}\text{Ru}$

Origin of abundant p-nuclei $^{92,94}\text{Mo}$ and $^{96,98}\text{Ru}$ in the early Galaxy is dominated by the **vp-process** in **collapsar** nucleosynthesis.

3. Isomers in R-process

New isomer in ^{127}Cd could affect the 2nd r-process peak through the β -decay and β -delayed one & two neutron emissions from ^{128}Ag . This effect is **independent** of astrophysical models of **supernova, collapsar or neutron star merger**.

4. Origin of $^{180}\text{Ta}/^{138}\text{La}$, ^{92}Nb and ^{98}Nb

Solar system ^{180}Ta and ^{138}La are explained consistently by supernova v-process. Intra and inter **isomer-intermediate-ground** transitions of ^{180}Ta is found to be critical.

5. Flavor Oscillation and the ν -mass Hierarchy

Supernova **v-process** of ^7Li and ^{11}B is quantitatively the most sensitive probe of the **v-flavor oscillation** in high-density matter and could constraint the **v-mass hierarchy**.