

Ab initio calculations of muon capture in light nuclei, and connections to neutrinoless double-beta decay matrix elements

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Canadian Astroparticle Physics Research Institute



Introduction to double-beta decay

Corrections to $0\nu\beta\beta$ -decay nuclear matrix elements

Muon capture as a probe of $0\nu\beta\beta$ decay

Summary and Outlook

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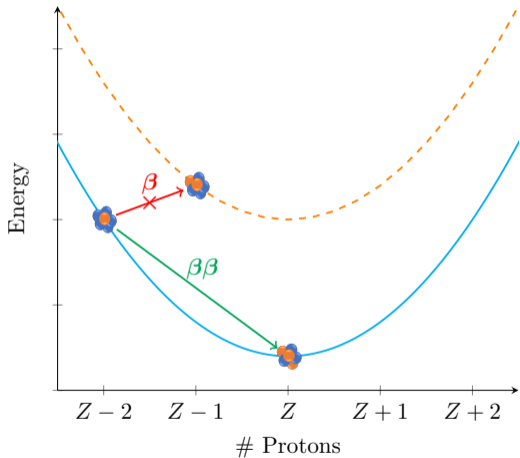
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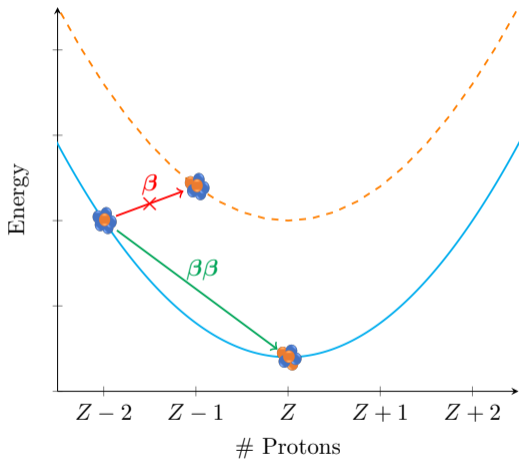


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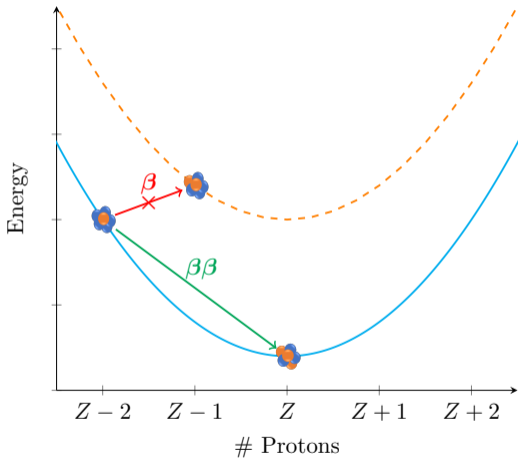


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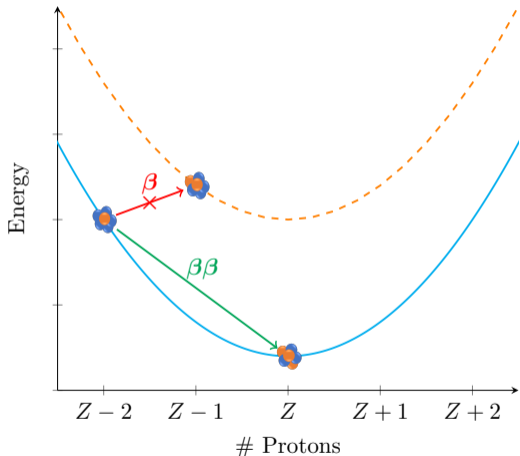


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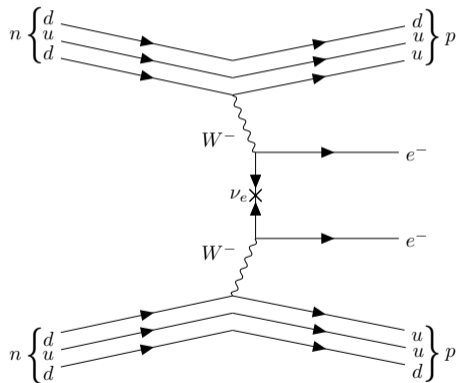
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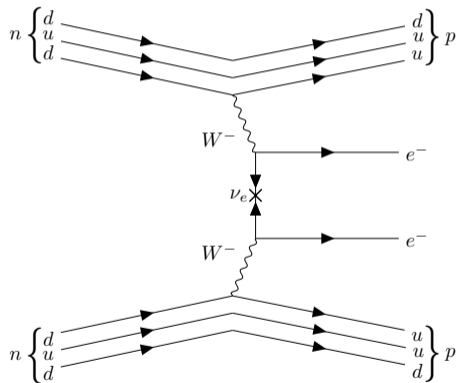
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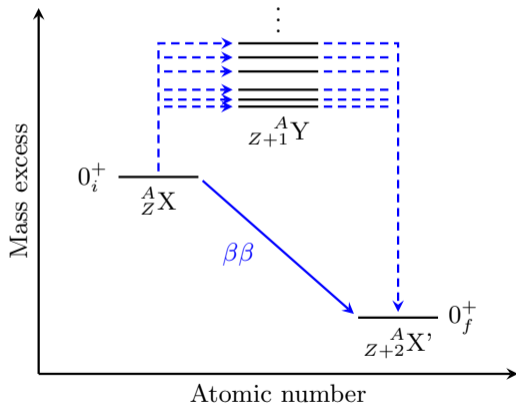
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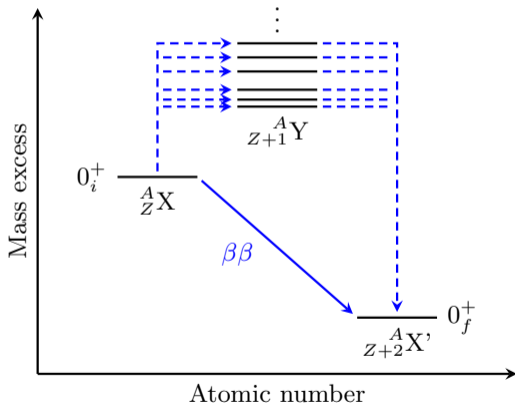
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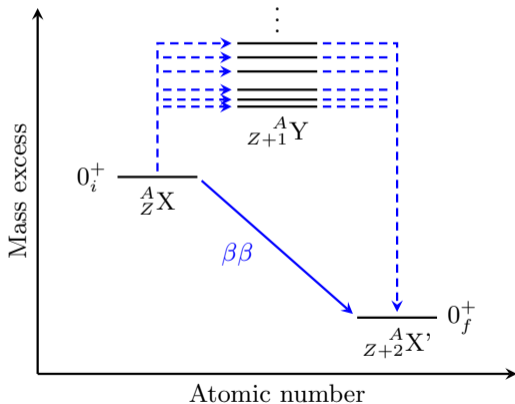
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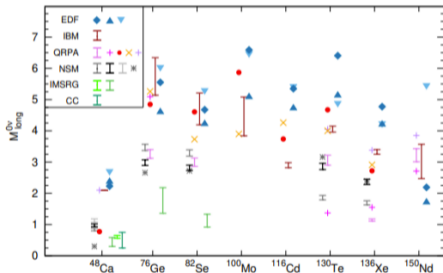
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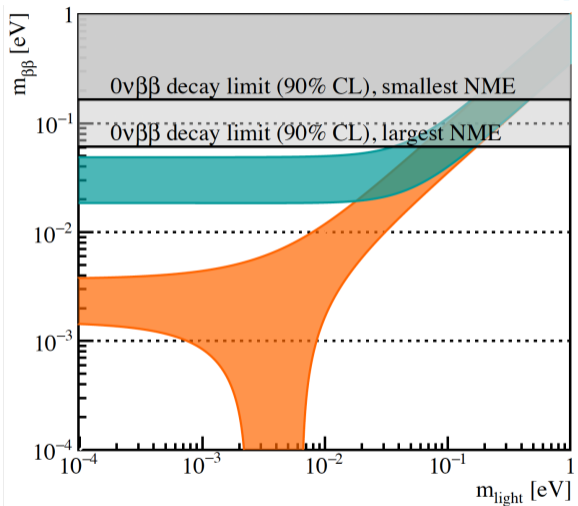
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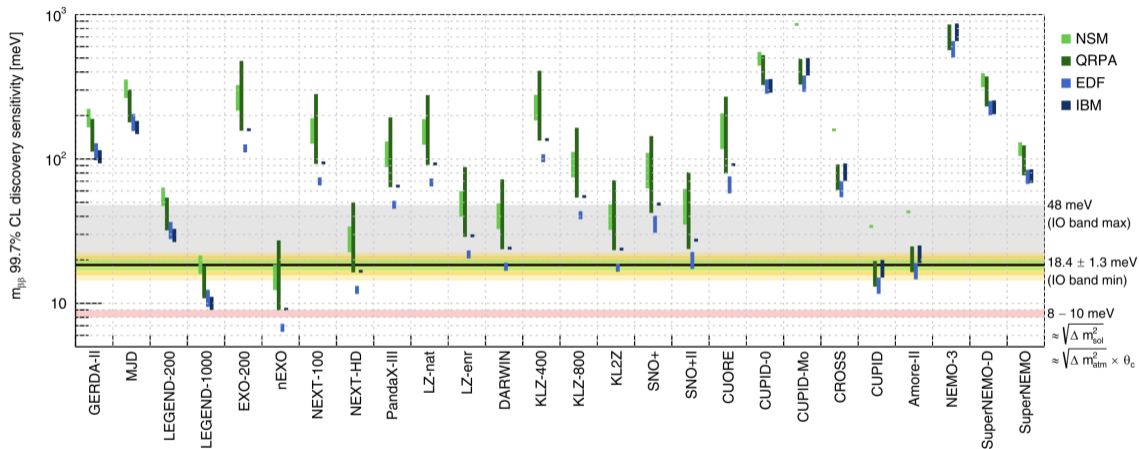
M. Agostini et al., Rev. Mod. Phys. **95**, 025002 (2023)

Current reach of the experiments



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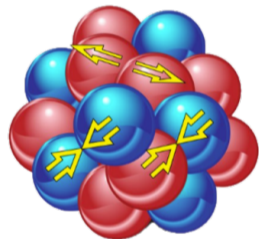
Next generation experiments



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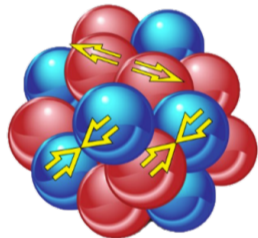
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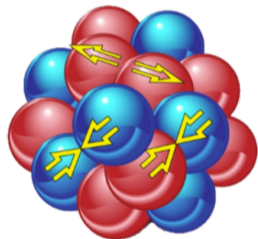
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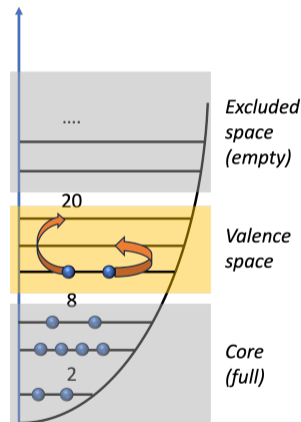
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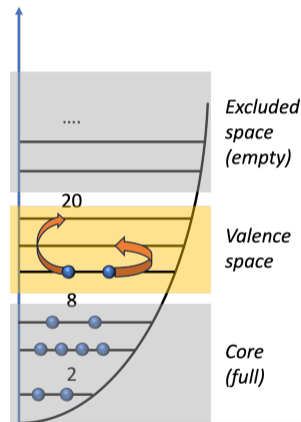
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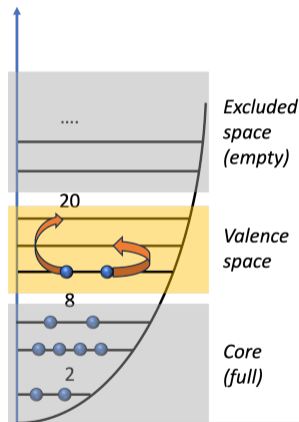
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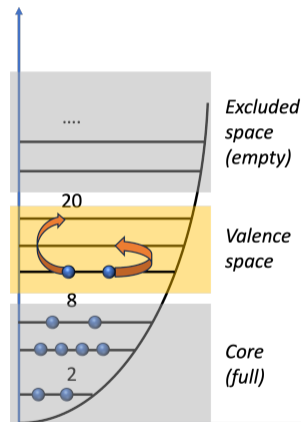
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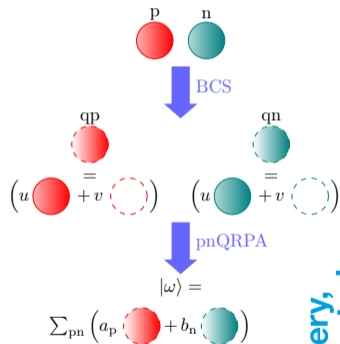
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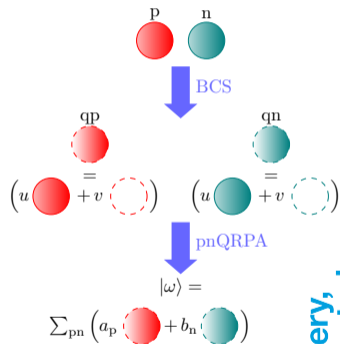
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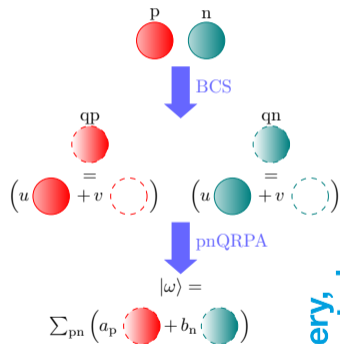
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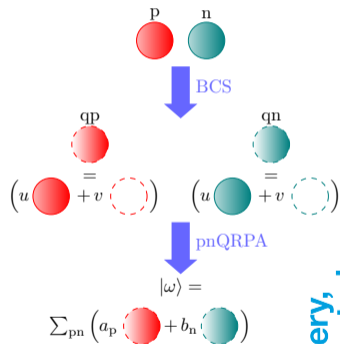
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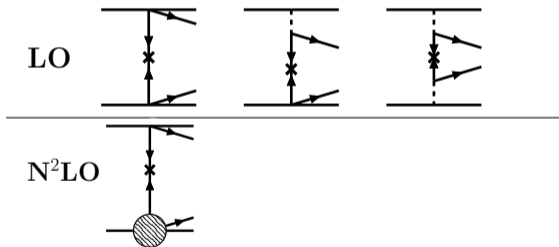
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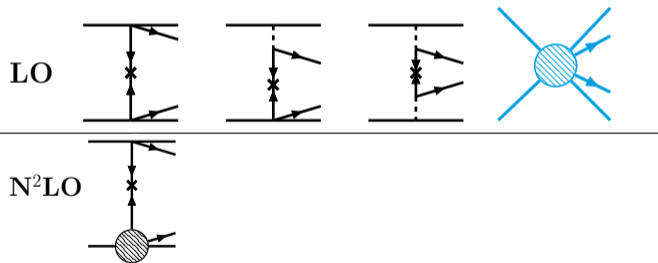
V. Cirigliano et al., *Phys. Rev. C* 97, 065501 (2018), *Phys. Rev. Lett.* 120, 202001 (2018), *Phys. Rev. C* 100, 055504 (2019)



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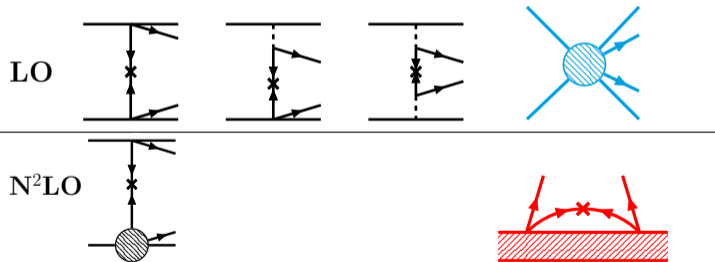
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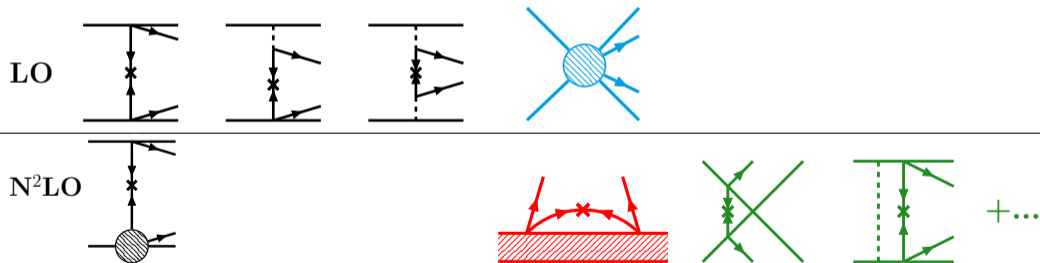
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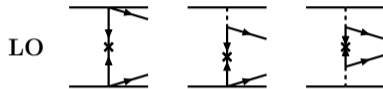
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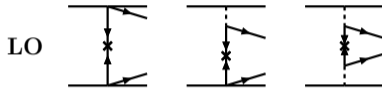


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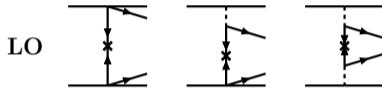
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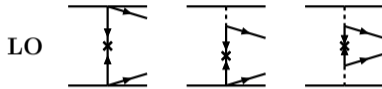
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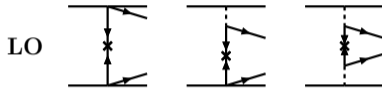
$$\mathbf{J} = \tau \left[g_A(p^2)\boldsymbol{\sigma} - g_P(p^2)\mathbf{p}(\mathbf{p} \cdot \boldsymbol{\sigma}) + ig_M(p^2)\frac{\boldsymbol{\sigma} \times \mathbf{p}}{2m_N} \right]$$

Traditional $0\nu\beta\beta$ -decay operators

- Traditionally, the nuclear current includes the leading-order (LO) transition operators

$$\mathcal{J}^0 = \tau[g_V(0)]$$

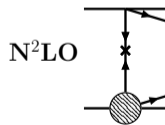
$$\mathbf{J} = \tau[g_A(0)\boldsymbol{\sigma} - g_P(0)\mathbf{p}(\mathbf{p} \cdot \boldsymbol{\sigma})]$$



- and next-to-next-to-leading-order (N²LO) corrections absorbed into **form factors** and **induced weak-magnetism terms**

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Traditional nuclear matrix elements of neutrinoless double-beta decay

$$M^{0\nu} = \frac{R}{g_A^2} \int \frac{d\mathbf{k}}{2\pi^2} \frac{e^{i\mathbf{k}(\mathbf{x}-\mathbf{y})}}{E_\nu} \sum_n \frac{\langle f | J_\mu(\mathbf{x}) | n \rangle \langle n | J^\mu(\mathbf{y}) | i \rangle}{E_\nu + E_n - \frac{1}{2}(E_i + E_f) - \frac{1}{2}(E_1 - E_2)}$$

- Energy of the virtual neutrino typically $E_\nu = \sqrt{m_\nu^2 + \mathbf{k}^2} \sim |\mathbf{k}| \sim k_F \sim 100 \text{ MeV}$ (“soft neutrinos”)

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Closure approximation

Without closure approximation:

$$M^{0\nu} \propto \sum_n \frac{\langle f | J_\mu(\mathbf{x}) | n \rangle \langle n | J^\mu(\mathbf{y}) | i \rangle}{|\mathbf{k}| + E_n - \frac{1}{2}(E_i + E_f)}$$

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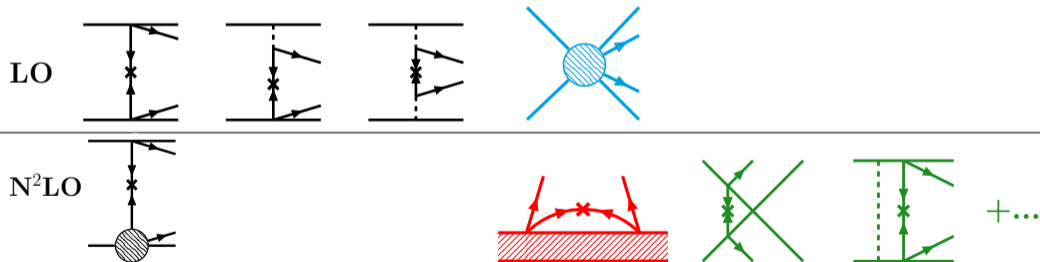
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- ▶ Typically used **with other nuclear methods**

Leading-order short-range contribution to $0\nu\beta\beta$ decay

$$\frac{1}{t_{1/2}^{0\nu}} = g_A^4 G^{0\nu} |M_L^{0\nu} + M_S^{0\nu} + M_{\text{usoft}}^{0\nu} + M_{\text{N}^2\text{LO}}^{0\nu}|^2 \left(\frac{m_{\beta\beta}}{m_e}\right)^2$$

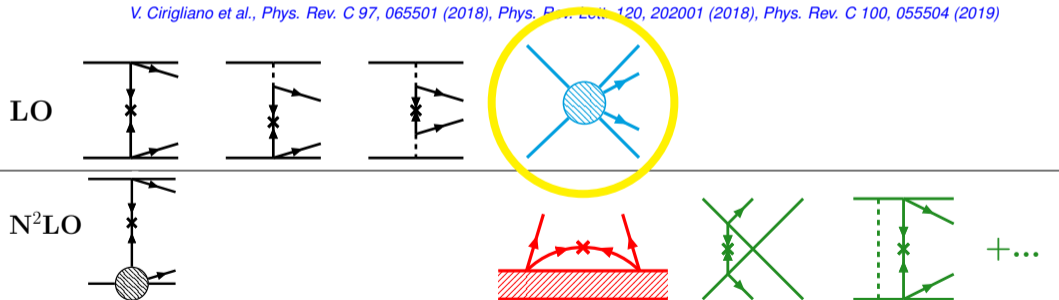
V. Cirigliano et al., *Phys. Rev. C* 97, 065501 (2018), *Phys. Rev. Lett.* 120, 202001 (2018), *Phys. Rev. C* 100, 055504 (2019)



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$$M_S^{0\nu} = \frac{2R}{\pi g_A^2} \langle 0_f^+ | \sum_{m,n} \tau_m^- \tau_n^- \int j_0(qr) h_S(q^2) q^2 dq | 0_i^+ \rangle$$

with

$$h_S(q^2) = 2\mathbf{g}_\nu^{\text{NN}} e^{-q^2/(2\Lambda^2)} .$$

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Contact Term in pnQRPA and NSM

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Couplings (g_ν^{NN}) and scales (Λ) of the Gaussian regulator¹.

g_ν^{NN} (fm ²)	Λ (MeV)
-0.67	450
-1.01	550
-1.44	465
-0.91	465
-1.44	349
-1.03	349

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Contact Term in pnQRPA and NSM

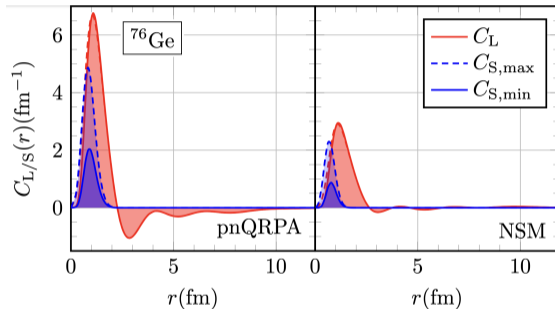
$$\int C_{L/S}(r)dr = M_{L/S}^{0\nu}$$

In pnQRPA:

$$M_S/M_L \approx 30\% - 80\%$$

In NSM:

$$M_S/M_L \approx 15\% - 50\%$$



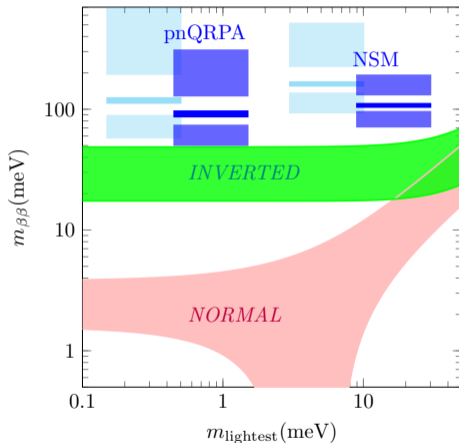
LJ, P. Soriano and J. Menéndez, Phys. Lett. B 823, 136720 (2021)

Effective Neutrino Masses

- Effective neutrino masses combining the likelihood functions of GERDA (^{76}Ge), CUORE (^{130}Te), EXO-200 (^{136}Xe) and KamLAND-Zen (^{136}Xe)

S. D. Biller, Phys. Rev. D **104**, 012002 (2021)

- Middle bands: $M_L^{(0\nu)}$
 Lower bands: $M_L^{(0\nu)} + M_S^{(0\nu)}$
 Upper bands: $M_L^{(0\nu)} - M_S^{(0\nu)}$

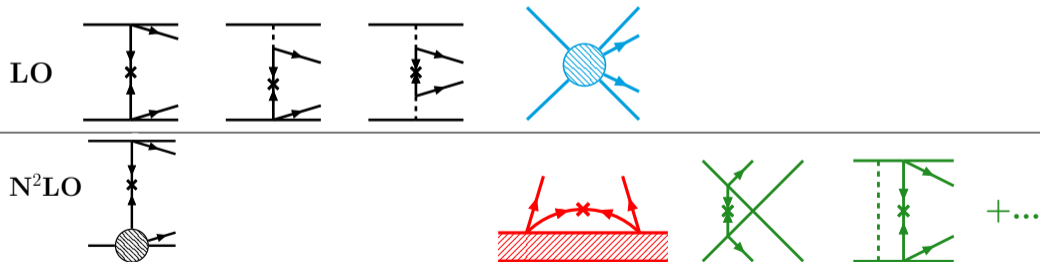


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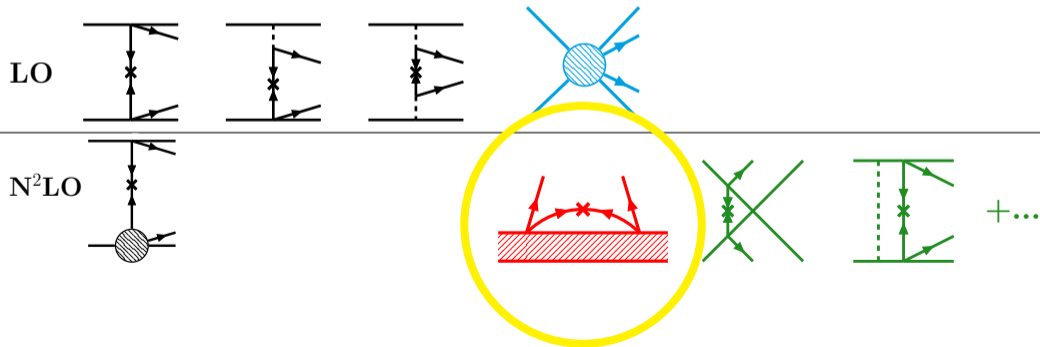
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Contribution of ultrasoft neutrinos

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V. Cirigliano et al., Phys. Rev. C 97, 065501 (2018)

$$M_{\text{usoft}}^{0\nu} = \frac{\pi R}{g_A^2} \sum_n \frac{d^{d-1}k}{(2\pi)^{d-1}} \frac{1}{|\mathbf{k}|} \left[\frac{\langle f | J_\mu | n \rangle \langle n | J^\mu | i \rangle}{|\mathbf{k}| + E_2 + E_n - E_i - i\eta} + \frac{\langle f | J_\mu | n \rangle \langle n | J^\mu | i \rangle}{|\mathbf{k}| + E_1 + E_n - E_i - i\eta} \right]$$

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- Keeping only $\mathbf{k} = \mathbf{0}$ term in the current:

$$\begin{aligned} M_{\text{usoft}}^{0\nu}(\mu_{\text{us}}) = & -\frac{R}{2\pi} \sum_n \langle f | \sum_a \sigma_a \tau_a^+ | n \rangle \langle n | \sum_b \sigma_b \tau_b^+ | i \rangle \\ & \times \left[(E_1 + E_n - E_i) \left(\ln \frac{\mu_{\text{us}}}{2(E_1 + E_n - E_i)} + 1 \right) \right. \\ & \left. + (E_2 + E_n - E_i) \left(\ln \frac{\mu_{\text{us}}}{2(E_2 + E_n - E_i)} + 1 \right) \right] \end{aligned}$$

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Are we missing a factor of 2?

$$\times \left[(E_1 + E_n - E_i) \left(\ln \frac{\mu_{\text{us}}}{2(E_1 + E_n - E_i)} + 1 \right) + (E_2 + E_n - E_i) \left(\ln \frac{\mu_{\text{us}}}{2(E_2 + E_n - E_i)} + 1 \right) \right]$$

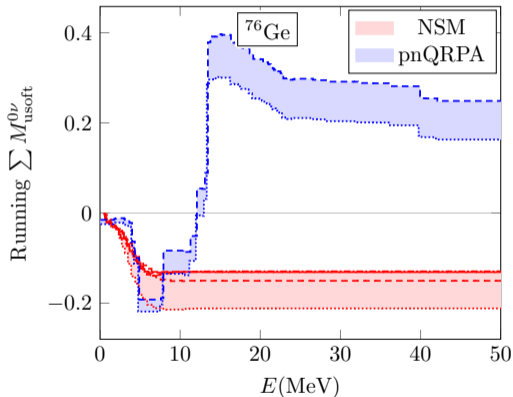
PRELIMINARY Ultrasoft neutrinos in pnQRPA and nuclear shell model

In pnQRPA:

$$|M_{\text{usoft}}^{0\nu} / M_{\text{L}}^{0\nu}| \leq 15\%$$

In NSM:

$$|M_{\text{usoft}}^{0\nu} / M_{\text{L}}^{0\nu}| \leq 5\%$$



LJ, D. Castillo, P. Soriano, J. Menéndez, work in progress

Ultrasoft neutrinos as correction of the closure approximation

$$\frac{1}{t_{1/2}^{0\nu}} = g_A^4 G^{0\nu} |M_L^{0\nu} + M_S^{0\nu} + M_{\text{usoft}}^{0\nu} + M_{\text{N2LO}}^{0\nu}|^2 \left(\frac{m_{\beta\beta}}{m_e}\right)^2$$

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In EFT:

$$M_L^{0\nu} \propto \frac{\langle f | J_\mu(\mathbf{x}) J^\mu(\mathbf{y}) | i \rangle}{|\mathbf{k}|}$$

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PRELIMINARY Ultrasoft neutrinos vs closure approximation in NSM

Nucleus	Interaction	$M^{0\nu}$	$M_{\text{cl}}^{0\nu 2}$	$M^{0\nu} - M_{\text{cl}}^{0\nu}$	$M_{\text{usoft}}^{0\nu}$
^{48}Ca	KB3G	0.92	0.96	-0.04	-0.01
	GXPF1.a42	0.78	0.78	0.00	0.02
^{76}Ge	JUN45	3.37	3.61	-0.24	-0.13
^{82}Se	JUN45	3.16	3.39	-0.23	-0.11

LJ, D. Castillo, P. Soriano, J Menéndez, work in progress

²R. A. Sen'kov, M. Horoi, , PRC **90**, 051301(R) (2014)

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^{82}Se	JUN45	3.16	3.39	-0.23	-0.11

LJ, D. Castillo, P. Soriano, J Menéndez, work in progress

²R. A. Sen'kov, M. Horoi, , PRC **90**, 051301(R) (2014)

PRELIMINARY Ultrasoft neutrinos vs closure approximation in pnQRPA

Nucleus	$M^{0\nu}$	$M_{\text{cl}}^{0\nu}$	$M^{0\nu} - M_{\text{cl}}^{0\nu}$	$M_{\text{usoft}}^{0\nu}$
^{76}Ge	4.83	4.68	0.15	0.25
^{82}Se	4.30	4.20	0.10	0.18
^{96}Zr	4.29	4.04	0.25	0.25
^{100}Mo	3.52	2.71	0.81	0.65
^{116}Cd	4.31	4.47	-0.16	-0.03
^{124}Sn	5.12	4.88	0.24	0.29
^{128}Te	3.99	3.76	0.23	0.27
^{130}Te	3.52	3.36	0.16	0.22
^{136}Xe	2.60	2.71	-0.11	0.06

LJ, D. Castillo, P. Soriano, J. Menéndez, work in progress

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LJ, D. Castillo, P. Soriano, J. Menéndez, work in progress

PRELIMINARY Ultrasoft neutrinos vs closure approximation in pnQRPA

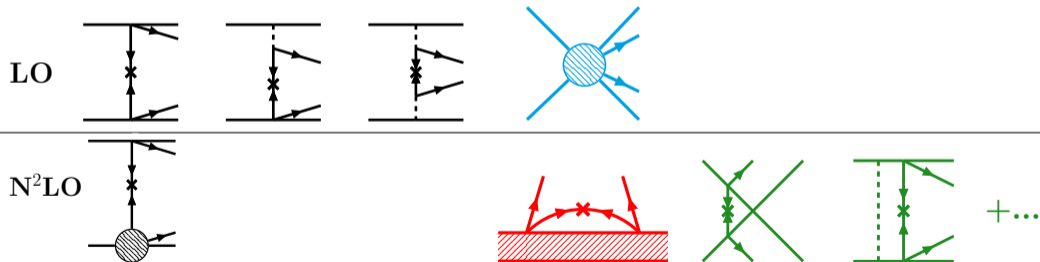
Nucleus	$M^{0\nu}$	$M_{cl}^{0\nu}$	$M^{0\nu} - M_{cl}^{0\nu}$	$M^{0\nu}(1^+) - M_{cl}^{0\nu}(1^+)$	$M_{usoft}^{0\nu}$
⁷⁶ Ge	4.83	4.68	0.15	0.26	0.25
⁸² Se	4.30	4.20	0.10	0.18	0.18
⁹⁶ Zr	4.29	4.04	0.25	0.26	0.25
¹⁰⁰ Mo	3.52	2.71	0.81	0.75	0.65
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¹²⁴ Sn	5.12	4.88	0.24	0.31	0.29
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L.J. D. Castillo, P. Soriano, J. Menéndez, work in progress

Genuine N²LO corrections to $0\nu\beta\beta$ decay

$$\frac{1}{t_{1/2}^{0\nu}} = g_A^4 G^{0\nu} |M_L^{0\nu} + M_S^{0\nu} + M_{\text{usoft}}^{0\nu} + M_{\text{N}^2\text{LO}}^{0\nu}|^2 \left(\frac{m_{\beta\beta}}{m_e}\right)^2$$

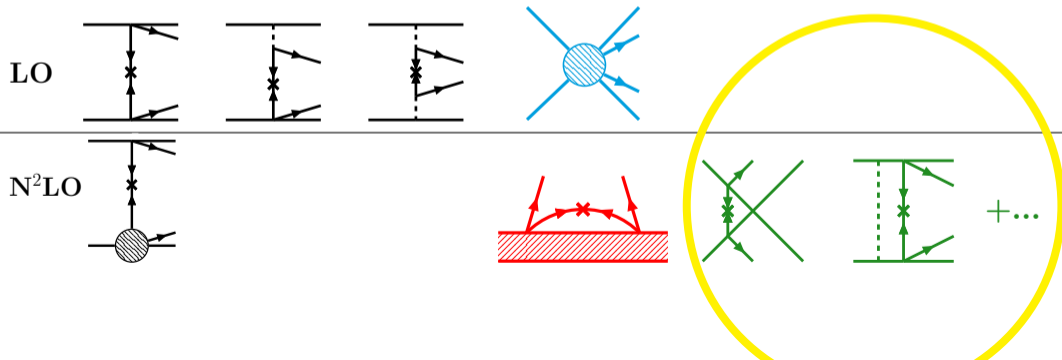
V. Cirigliano et al., *Phys. Rev. C* 97, 065501 (2018), *Phys. Rev. Lett.* 120, 202001 (2018), *Phys. Rev. C* 100, 055504 (2019)



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PRELIMINARY Genuine N²LO Corrections in pnQRPA

Nucleus	$M_L^{0\nu}$	$M_{N^2LO}^{0\nu}$	$ M_{N^2LO}^{0\nu}/M_L^{0\nu} $
⁷⁶ Ge	4.83	-0.04–0.53	$\lesssim 10\%$
⁸² Se	4.30	0.28–0.44	6% – 10%
⁹⁶ Zr	4.29	-0.04–0.42	$\lesssim 10\%$
¹⁰⁰ Mo	3.52	-0.05–0.62	$\lesssim 18\%$
¹¹⁶ Cd	4.31	-0.02–0.29	$\lesssim 7\%$
¹²⁴ Sn	5.12	-0.04–0.66	$\lesssim 13\%$
¹²⁸ Te	3.99	-0.04–0.55	$\lesssim 14\%$
¹³⁰ Te	3.52	-0.03–0.52	$\lesssim 15\%$
¹³⁶ Xe	2.60	-0.02–0.07	$\lesssim 3\%$

LJ, D. Castillo, P. Soriano, J Menéndez, work in progress

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Caveats:

LJ, D. Castillo, P. Soriano, J Menéndez, work in progress

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LJ, D. Castillo, P. Soriano, J Menéndez, work in progress

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Caveats:

- Unknown parameters
- Scale dependence

LJ, D. Castillo, P. Soriano, J Menéndez, work in progress

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Caveats:

- Unknown parameters
- Scale dependence
- Regulator dependence

LJ, D. Castillo, P. Soriano, J Menéndez, work in progress

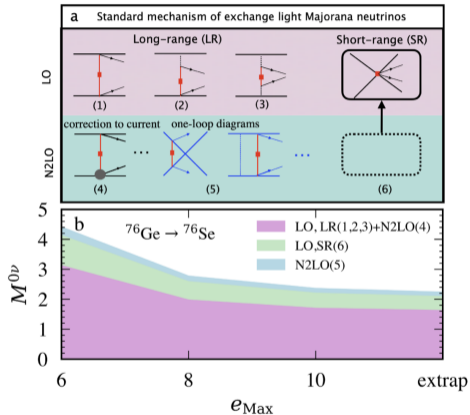
Similar effects found in *ab initio* studies

● In ^{76}Ge :

$$M_S^{0\nu} / M_L^{0\nu} \sim 40\% ,$$

$$M_{N^2LO}^{0\nu} / M_L^{0\nu} \sim 5\%$$

A. Belley et al. arXiv:2308.15634 (2023)



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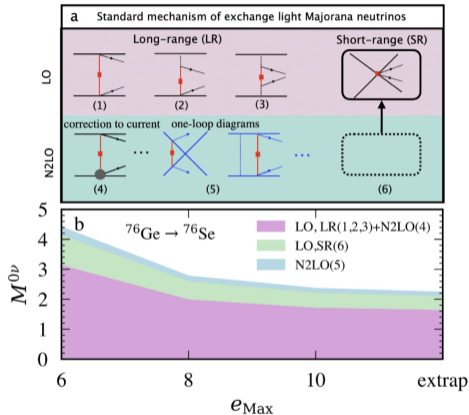
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A. Belley et al. arXiv:2308.15634 (2023)

- In ^{130}Te and ^{136}Xe :

$$M_S^{0\nu} / M_L^{0\nu} \sim 20\% - 120\%$$

A. Belley et al. arXiv:2307.15156 (2023)



A. Belley et al. arXiv:2308.15634 (2023)

Introduction to double-beta decay

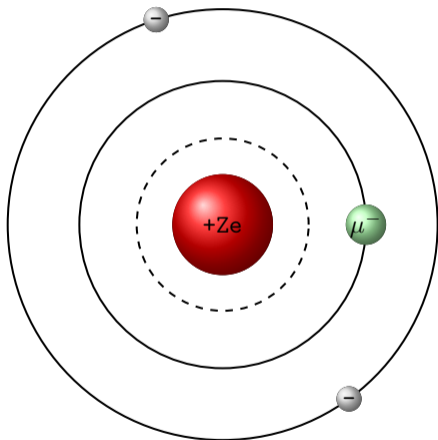
Corrections to $0\nu\beta\beta$ -decay nuclear matrix elements

Muon capture as a probe of $0\nu\beta\beta$ decay

Summary and Outlook

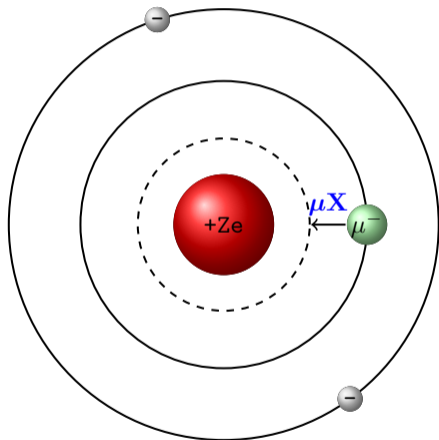
Ordinary Muon Capture (OMC)

- A muon can replace an electron in an atom, forming a *muonic atom*



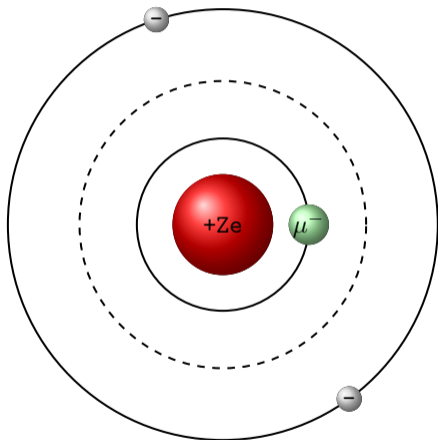
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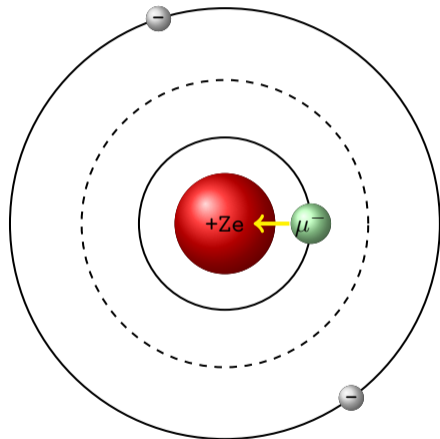
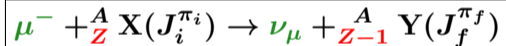
Ordinary Muon Capture (OMC)

- A muon can replace an electron in an atom, forming a *muonic atom*
 - ▶ Eventually bound on **the $1s_{1/2}$ orbit**



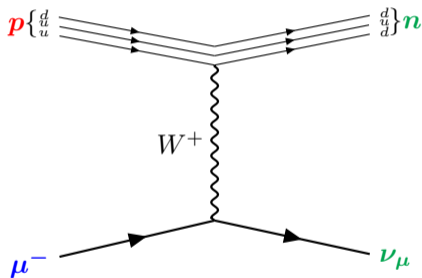
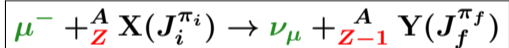
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- The *muon* can then be captured by the nucleus



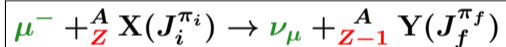
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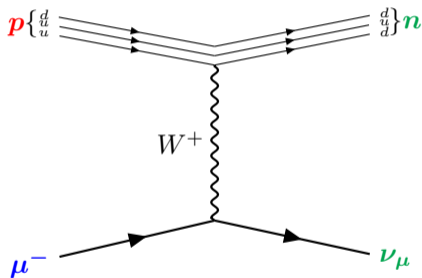
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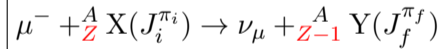
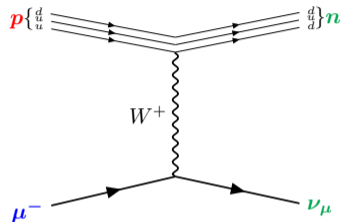
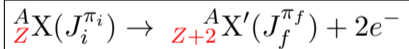
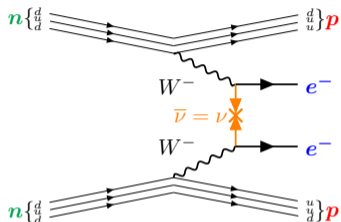
Ordinary = non-radiative

(Radiative muon capture (RMC):)

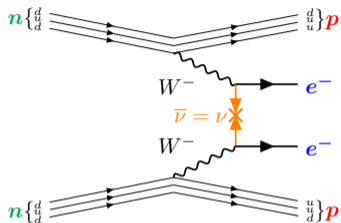
$$\mu^- + {}^A_Z X(J_i^{\pi_i}) \rightarrow \nu_\mu + {}^A_{Z-1} Y(J_f^{\pi_f}) + \gamma$$



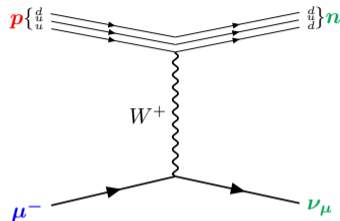
$0\nu\beta\beta$ Decay vs. Muon Capture



$0\nu\beta\beta$ Decay vs. Muon Capture



$$\frac{A}{Z} X(J_i^{\pi_i}) \rightarrow \frac{A}{Z+2} X'(J_f^{\pi_f}) + 2e^-$$

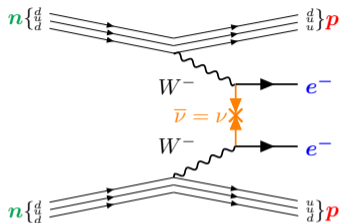


$$\mu^- + \frac{A}{Z} X(J_i^{\pi_i}) \rightarrow \nu_\mu + \frac{A}{Z-1} Y(J_f^{\pi_f})$$

Both involve hadronic current:

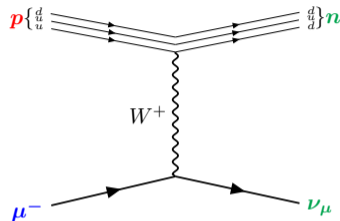
$$\langle \mathbf{p} | j^{\alpha\dagger} | \mathbf{p} \rangle = \bar{\Psi} \left[g_V(q^2) \gamma^\alpha - g_A(q^2) \gamma^\alpha \gamma_5 - g_P(q^2) q^\alpha \gamma_5 + i g_M(q^2) \frac{\sigma^{\alpha\beta} q_\beta}{2m_p} \right] \tau^\pm \Psi$$

$0\nu\beta\beta$ Decay vs. Muon Capture



$$\frac{A}{Z} X(J_i^{\pi_i}) \rightarrow \frac{A}{Z+2} X'(J_f^{\pi_f}) + 2e^-$$

- $q \approx 1/|r_1 - r_2| \approx 100 - 200 \text{ MeV}$

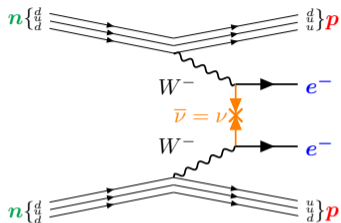


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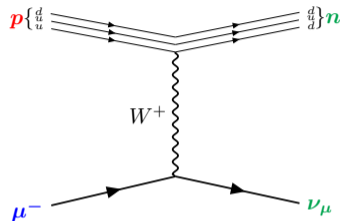
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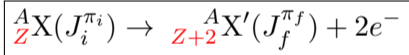
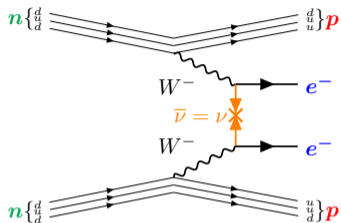
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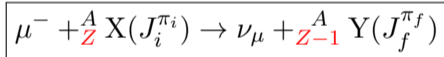
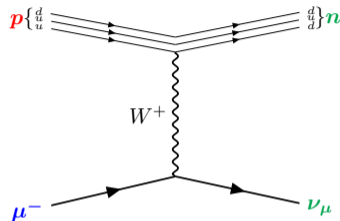
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$0\nu\beta\beta$ Decay vs. Muon Capture



- $q \approx 1/|r_1 - r_2| \approx 100 - 200 \text{ MeV}$
- **Yet hypothetical**

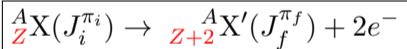
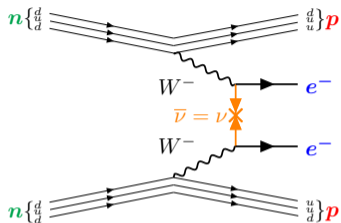


- $q \approx m_\mu + M_i - M_f - m_e - E_X \approx 100 \text{ MeV}$

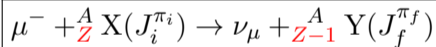
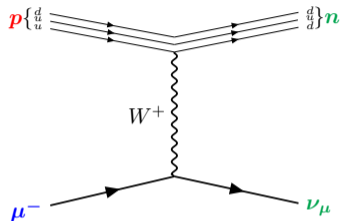
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- **Yet hypothetical**



- $q \approx m_\mu + M_i - M_f - m_e - E_X \approx 100 \text{ MeV}$
- **Has been measured!**

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Ab initio No-Core Shell Model (NCSM)

- Solve nuclear many-body problem

$$H^{(A)}\Psi^{(A)}(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A) = E^{(A)}\Psi^{(A)}(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A)$$

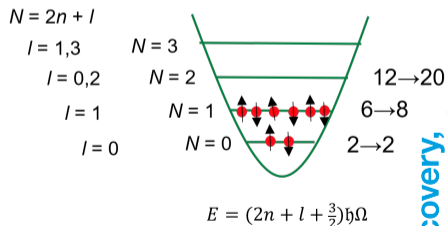
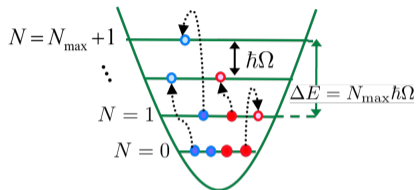


Figure courtesy of P. Navrátil

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- **Two- (NN)** and **three-nucleon (3N)** forces from χ EFT

$$H^{(A)} = \sum_{i=1}^A \frac{p_i^2}{2m} + \sum_{i<j=1}^A V^{NN}(\mathbf{r}_i - \mathbf{r}_j) + \sum_{i<j<k=1}^A V_{ijk}^{3N}$$

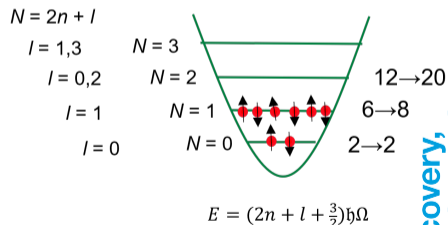
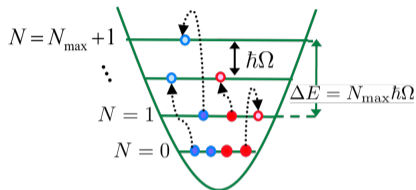


Figure courtesy of P. Navrátil

Ab initio No-Core Shell Model (NCSM)


- Solve nuclear many-body problem

$$H^{(A)}\Psi^{(A)}(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A) = E^{(A)}\Psi^{(A)}(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A)$$

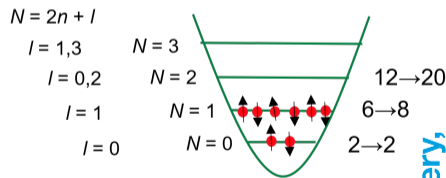
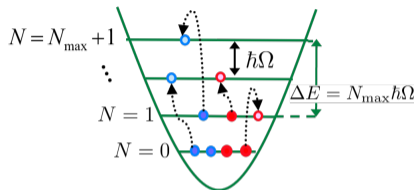
- **Two- (NN)** and **three-nucleon (3N)** forces from χ EFT

$$H^{(A)} = \sum_{i=1}^A \frac{p_i^2}{2m} + \sum_{i<j=1}^A V^{NN}(\mathbf{r}_i - \mathbf{r}_j) + \sum_{i<j<k=1}^A V_{ijk}^{3N}$$

- A -nucleon wave functions expanded in harmonic oscillator (HO) basis



$$\Psi^{(A)} = \sum_{N=0}^{N_{\max}} \sum_j c_{Nj} \Phi_{Nj}^{\text{HO}}(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A)$$



$$E = (2n + l + \frac{3}{2})\hbar\Omega$$

Figure courtesy of P. Navrátil

Dependency on the Harmonic-Oscillator Frequency

$$\Psi^{(A)} = \sum_{N=0}^{N_{\max}} \sum_j c_{Nj} \Phi_{Nj}^{\text{HO}}(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_A)$$

- The expansion depends on the HO frequency because of the N_{\max} truncation

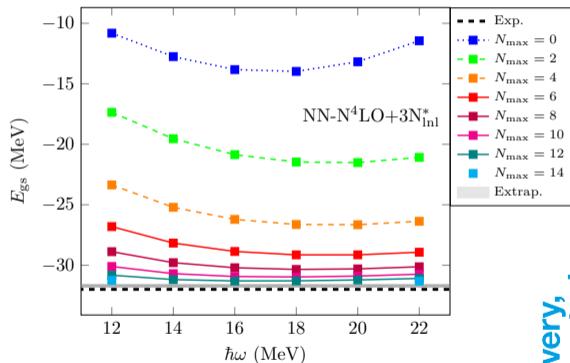
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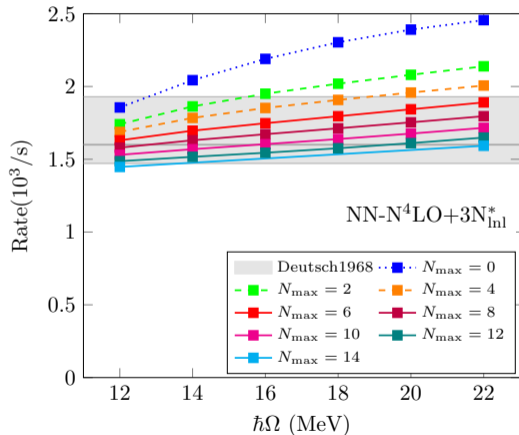
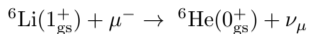
- Increasing N_{\max} leads towards converged results

Ground-state energy of ${}^6\text{Li}$



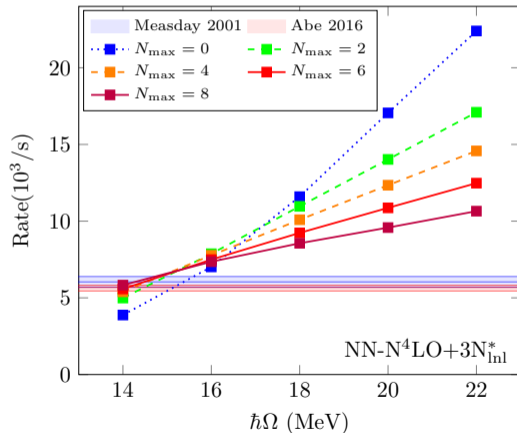
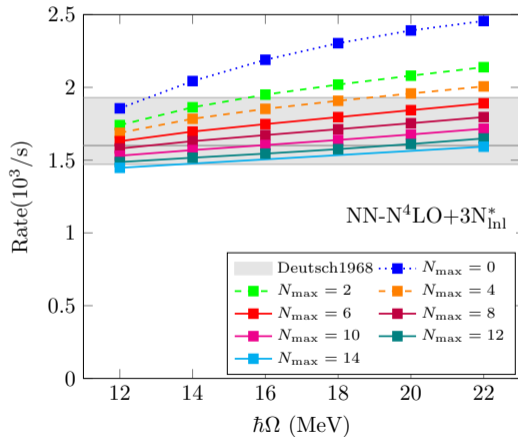
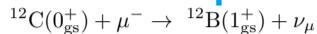
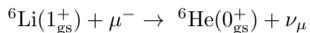
LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

Harmonic-Oscillator Frequency Dependence of Muon Capture



LJ, Navrátil, Kotila and Kravvaris, arXiv:2403.XXXX

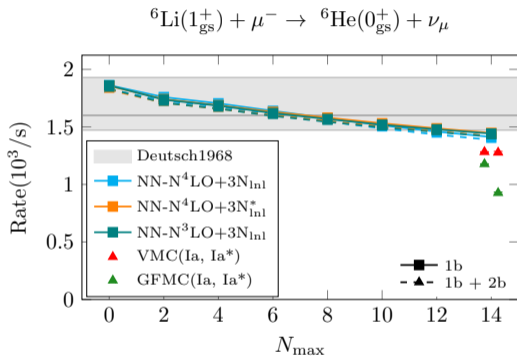
Harmonic-Oscillator Frequency Dependence of Muon Capture



LJ, Navrátil, Kotila and Kravvaris, arXiv:2403.XXXX

Muon Capture on ${}^6\text{Li}$

- NCSM slightly underestimating experiment

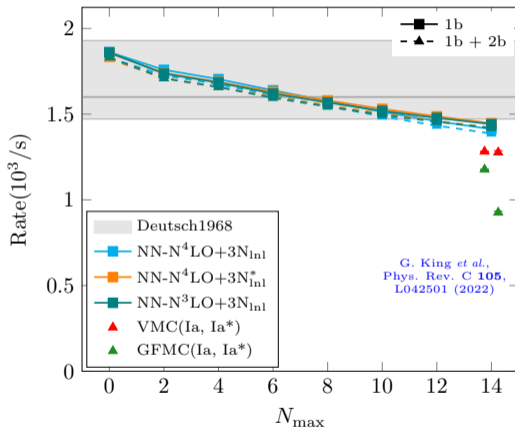
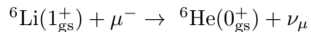


LJ, Navrátil, Kotila, Kravvaris, *arXiv:2403.XXXX*

- NCSM slightly underestimating experiment
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King *et al.*, Phys. Rev. C **105**, L042501 (2022)

Muon Capture on ${}^6\text{Li}$



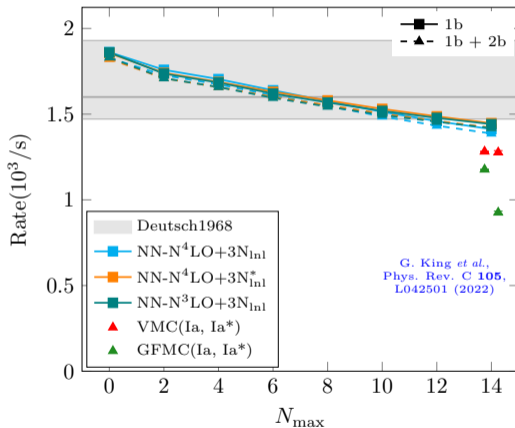
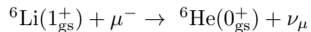
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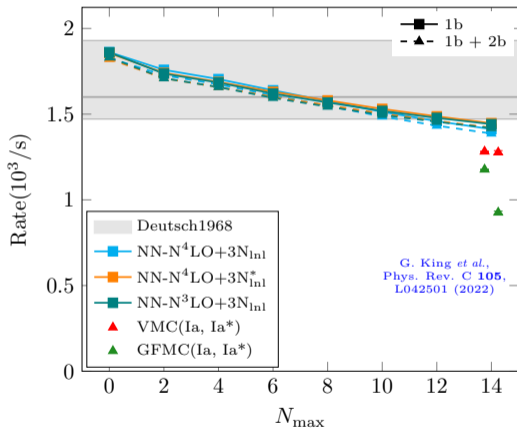
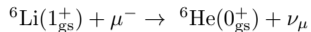
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 - ▶ **NCSM with continuum (NCSMC) might give better results?**

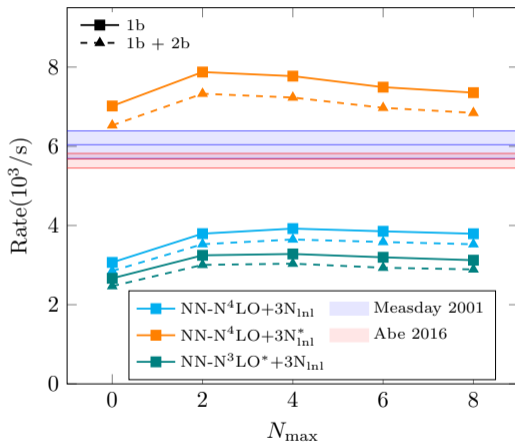
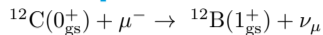
Muon Capture on ${}^6\text{Li}$



LJ, Navrátil, Kotila, Kravvaris, *arXiv:2403.XXXX*

- The **NN-N⁴LO+3N_{inl}^{*}** interaction with the additional spin-orbit 3N-force term most consistent with experiment

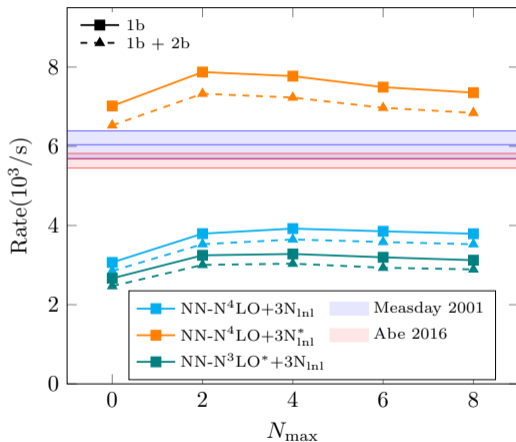
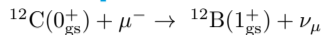
Muon capture on ¹²C



LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

- The $NN-N^4LO+3N_{Inl}^*$ interaction with the additional spin-orbit $3N$ -force term most consistent with experiment
- Capture rates to excited states in ^{12}B also well reproduced

Muon capture on ^{12}C

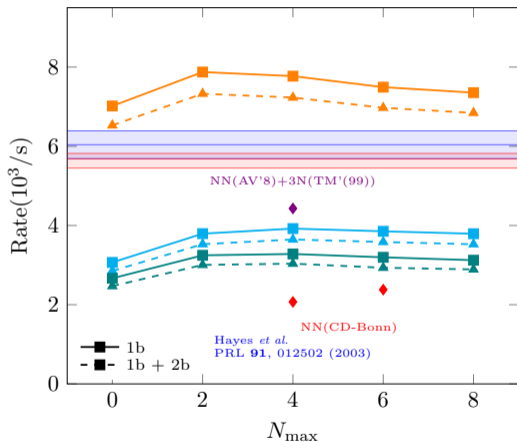
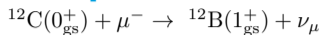


LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

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Hayes *et al.*, *Phys. Rev. Lett.* **91**, 012502 (2003)

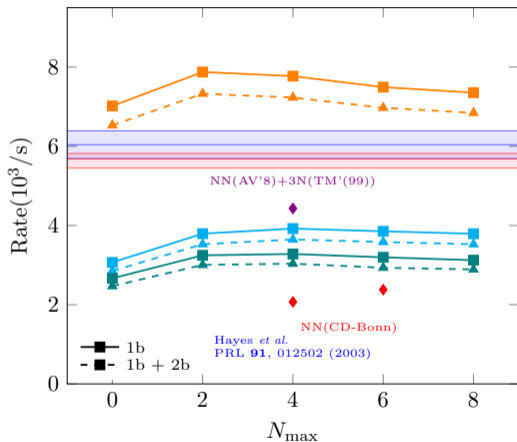
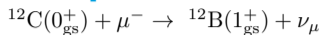
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LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

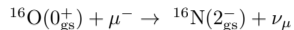
- The **NN-N⁴LO+3N_{int}*** interaction with the additional spin-orbit 3N-force term most consistent with experiment
 - Capture rates to excited states in ¹²B also well reproduced
 - Rates comparable with earlier NCSM results
- Hayes et al., Phys. Rev. Lett. 91, 012502 (2003)*
- 3N-forces essential to reproduce the measured rate

Muon capture on ¹²C

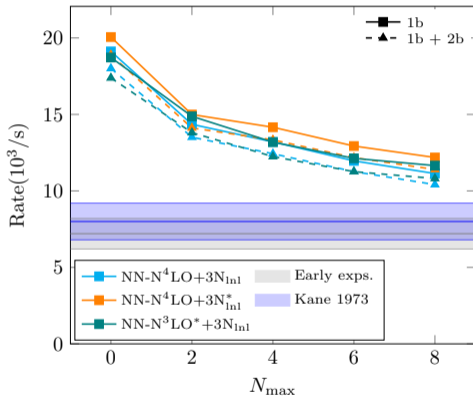


LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

Muon capture on ^{16}O

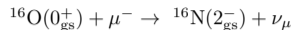


- NCSM describes well the complex systems ^{16}O and ^{16}N

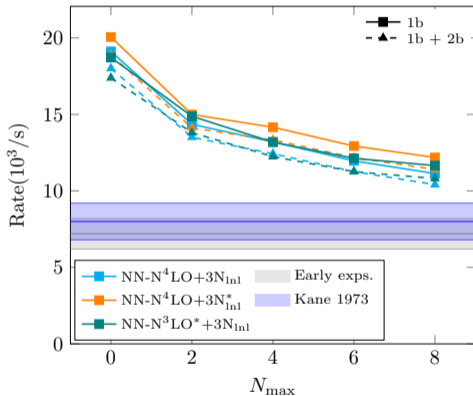


LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

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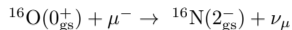


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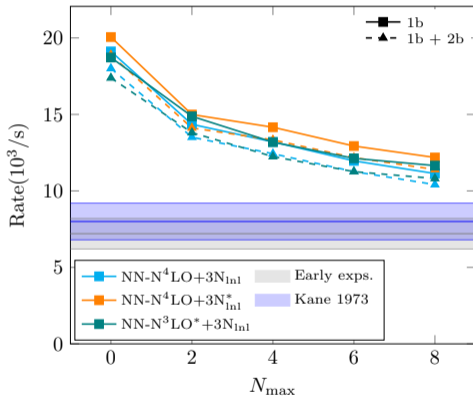


LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

Muon capture on ^{16}O

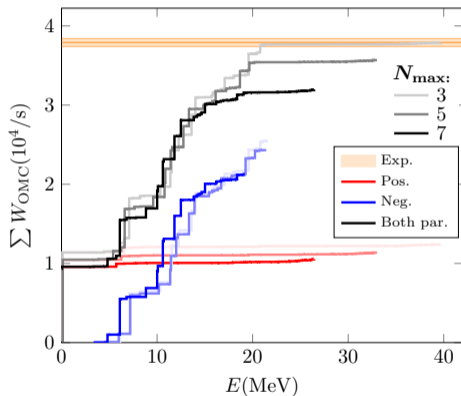
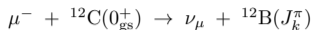


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LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

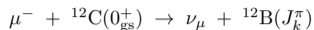
Total Muon-Capture Rates



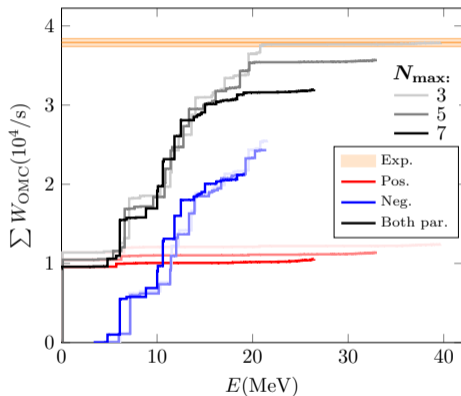
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LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

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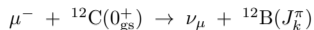


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- Summing up the rates, we capture $\sim 85\%$ of the total rate in both ${}^{12}\text{B}$ and ${}^{16}\text{N}$

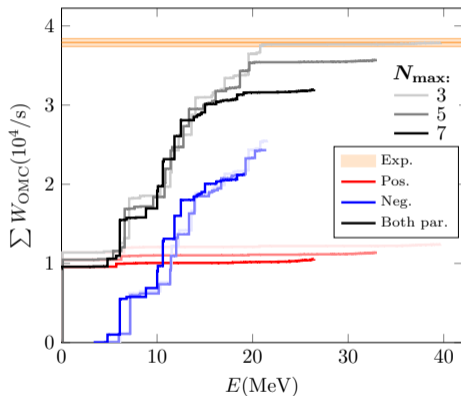


LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

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- Better estimation with the Lanczos strength function method underway



LJ, Navrátil, Kotila, Kravvaris, arXiv:2403.XXXX

Introduction to double-beta decay

Corrections to $0\nu\beta\beta$ -decay nuclear matrix elements

Muon capture as a probe of $0\nu\beta\beta$ decay

Summary and Outlook

- Newly introduced contact term significantly enhances the $0\nu\beta\beta$ -decay NMEs

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- Studying the contribution from ultrasoft neutrinos may help us estimate the closure correction to the $0\nu\beta\beta$ -decay NMEs
- Ab initio muon-capture studies could shed light on g_A quenching at finite momentum exchange regime relevant for $0\nu\beta\beta$ decay

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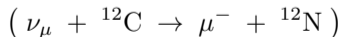
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 - ▶ ^{12}C and ^{16}O are both of interest in **neutrino-scattering experiments**



Thank you
Merci

