

Three-Body Valence Interactions with the IMSRG

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IMSRG

In-Medium Similarity Renormalization Group

- IMSRG

Perform a unitary transformation on the Hamiltonian

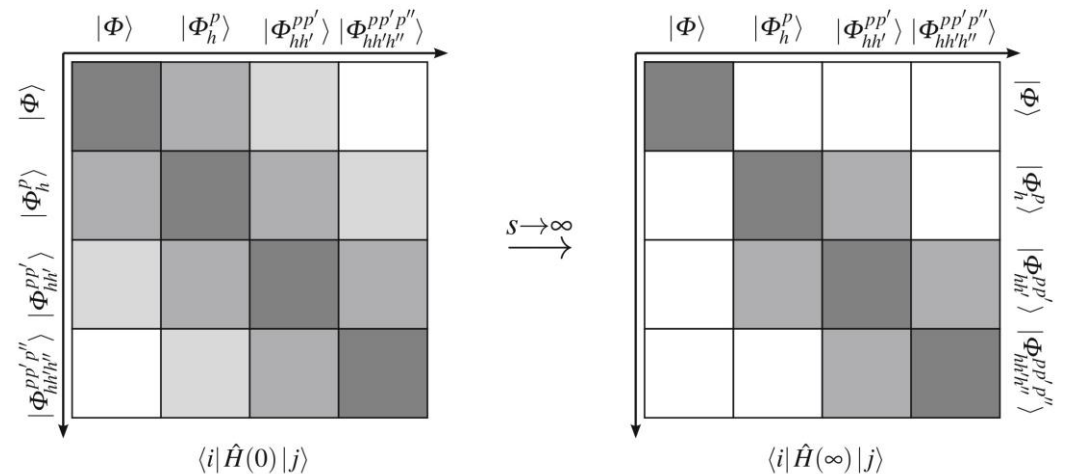
- $H(s) = U(s)HU(s)^\dagger$

Split Hamiltonian into diagonal and off-diagonal pieces

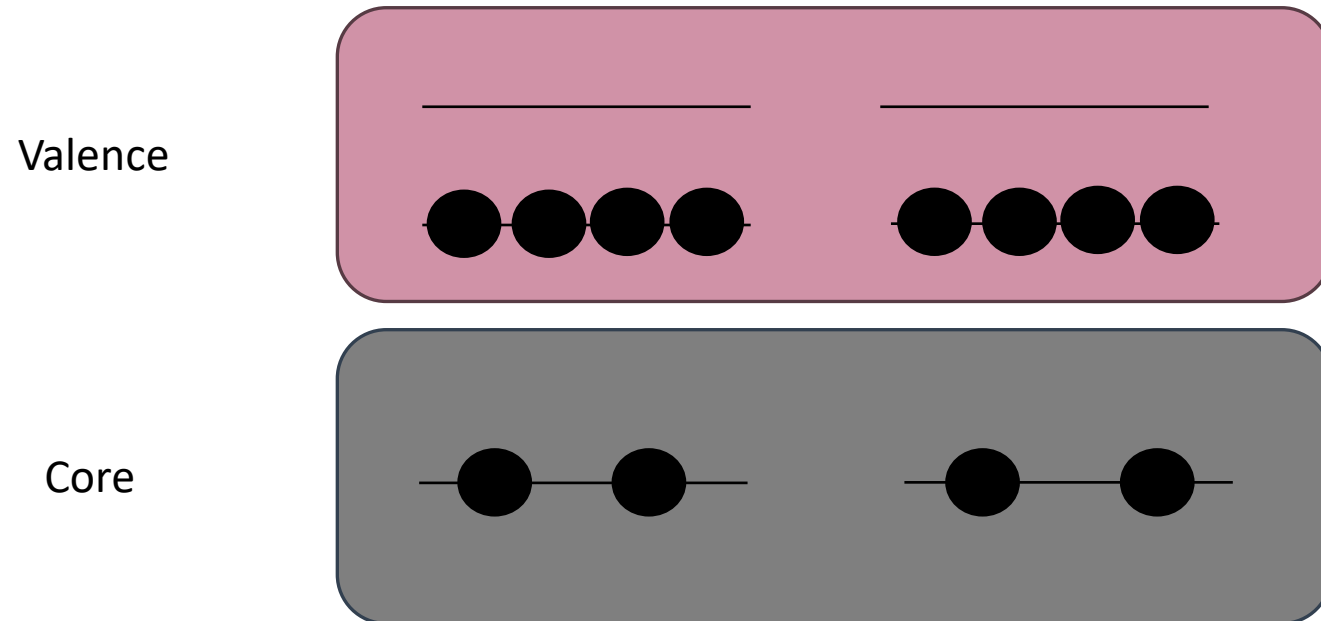
- $H = H^d + H^{od}$

Key truncations for the theory

- IMSRG(2), IMSRG(3f2), IMSRG(3n7), IMSRG(3)



Valence Space



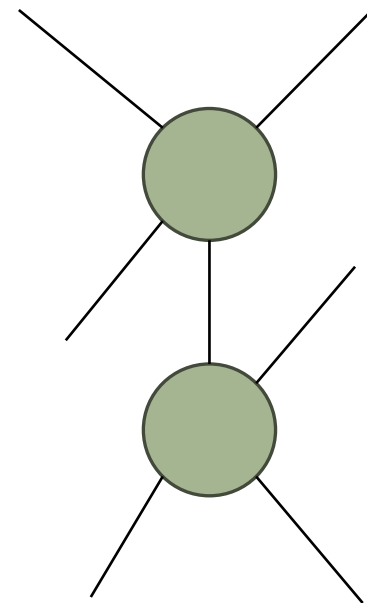
Using Three-Body Operators

Where Three-Body Operators come from

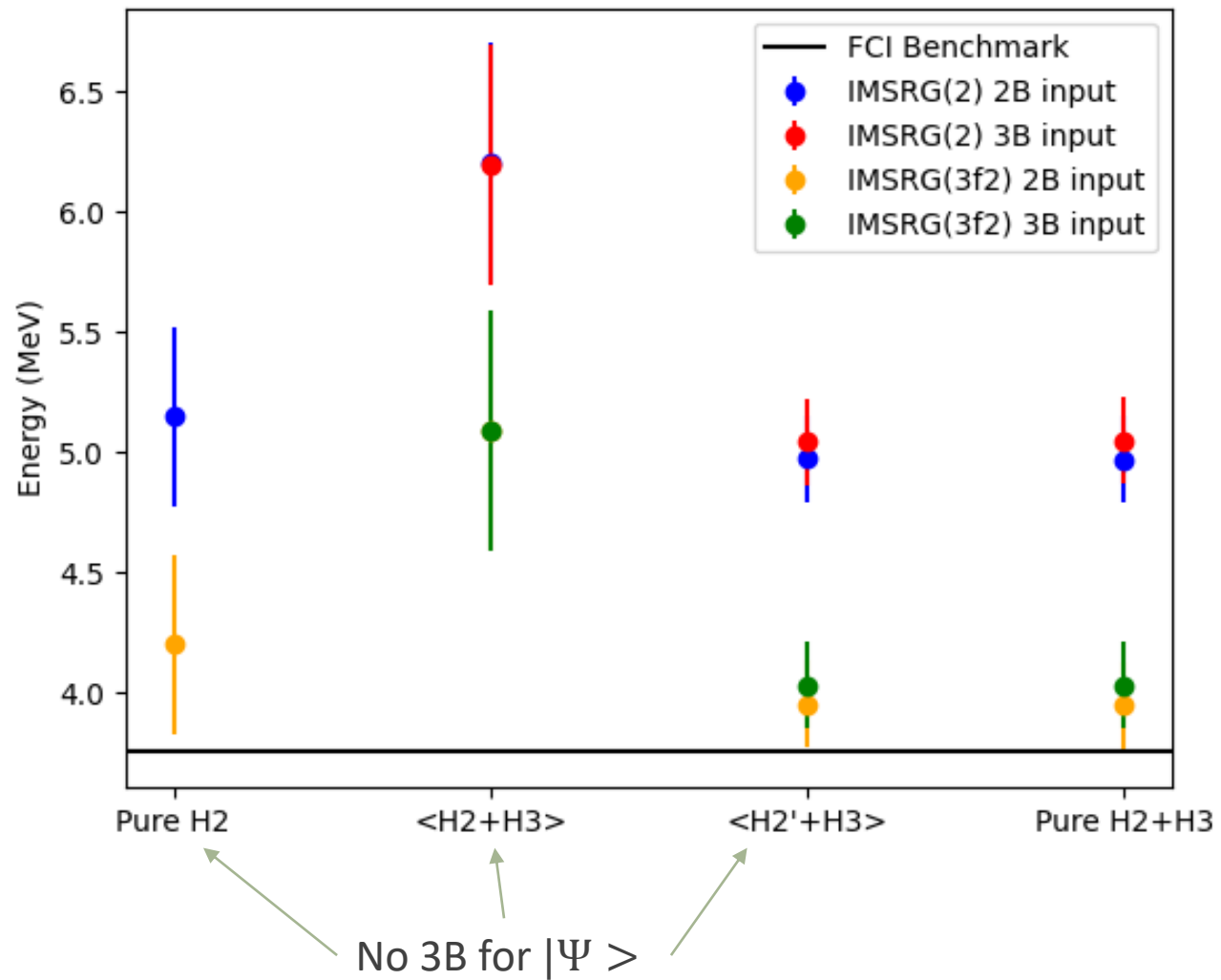
- Perturbative Three-Body
 - $[\Omega_2, H_2]_3 = H_3$
- Input Interaction Three-Body Piece

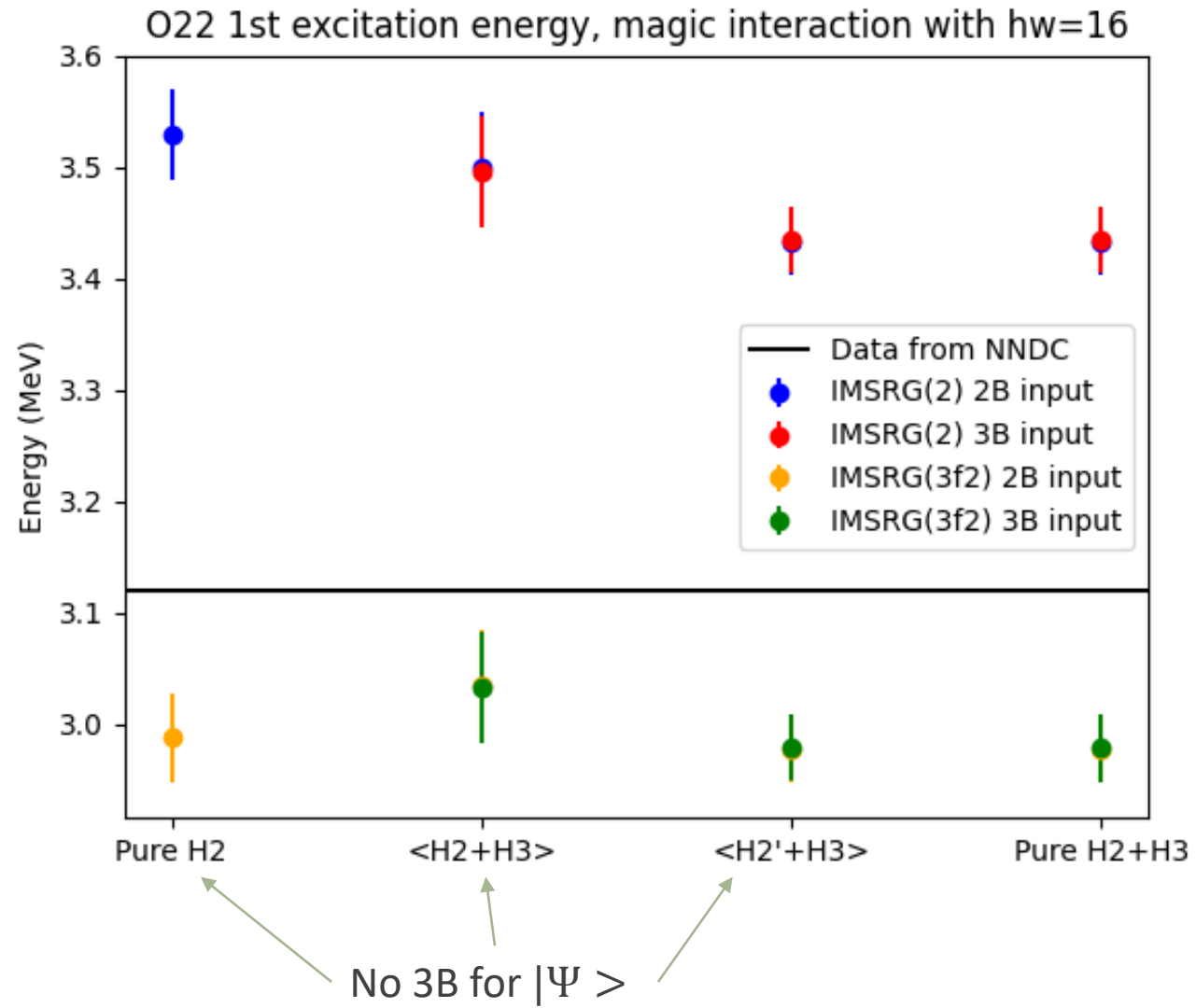
Where we use the Three-Body Operators

- Use in solving for wavefunction $|\Psi\rangle$
- Use in Operator expectation value, $\langle \Psi | O | \Psi \rangle$
 - $O = O_1 + O_2 + O_3$



C12 1st excitation energy, magic interaction with hw=16





Summary

1. Can approximate full three-body diagonalization result by including residual one-body and two-body pieces along with three-body piece at the operator level
2. Three-Body versions of input interactions give a small correction(if any) compared to the two-body versions of the same interactions

Thank you!

Ragnar Stroberg

Zhonghao Sun

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K. Hebeler, S. K. Bogner, R. J. Furnstahl, A. Nogga, and A. Schwenk, Phys. Rev. C 83, 031301(R) (2011).

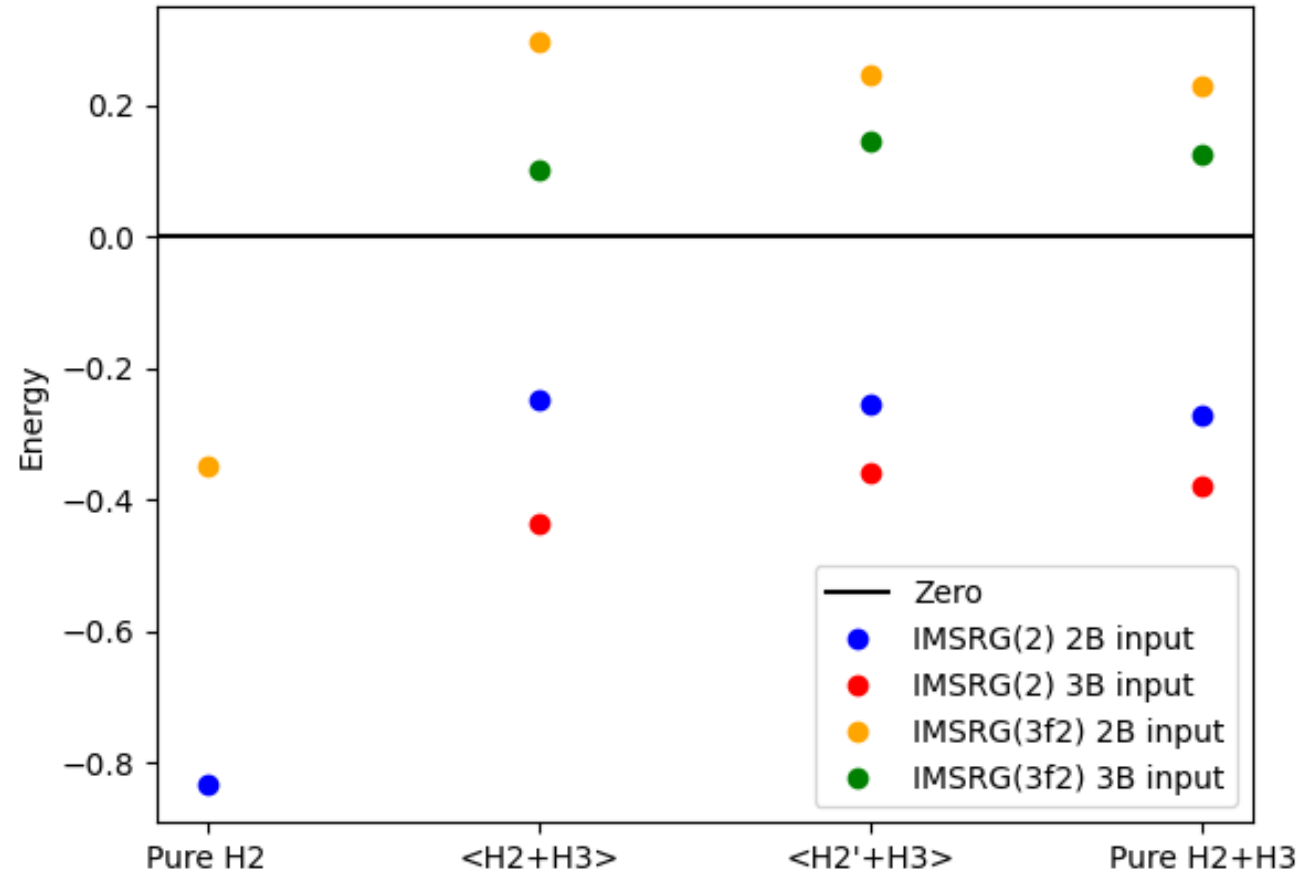
Backup

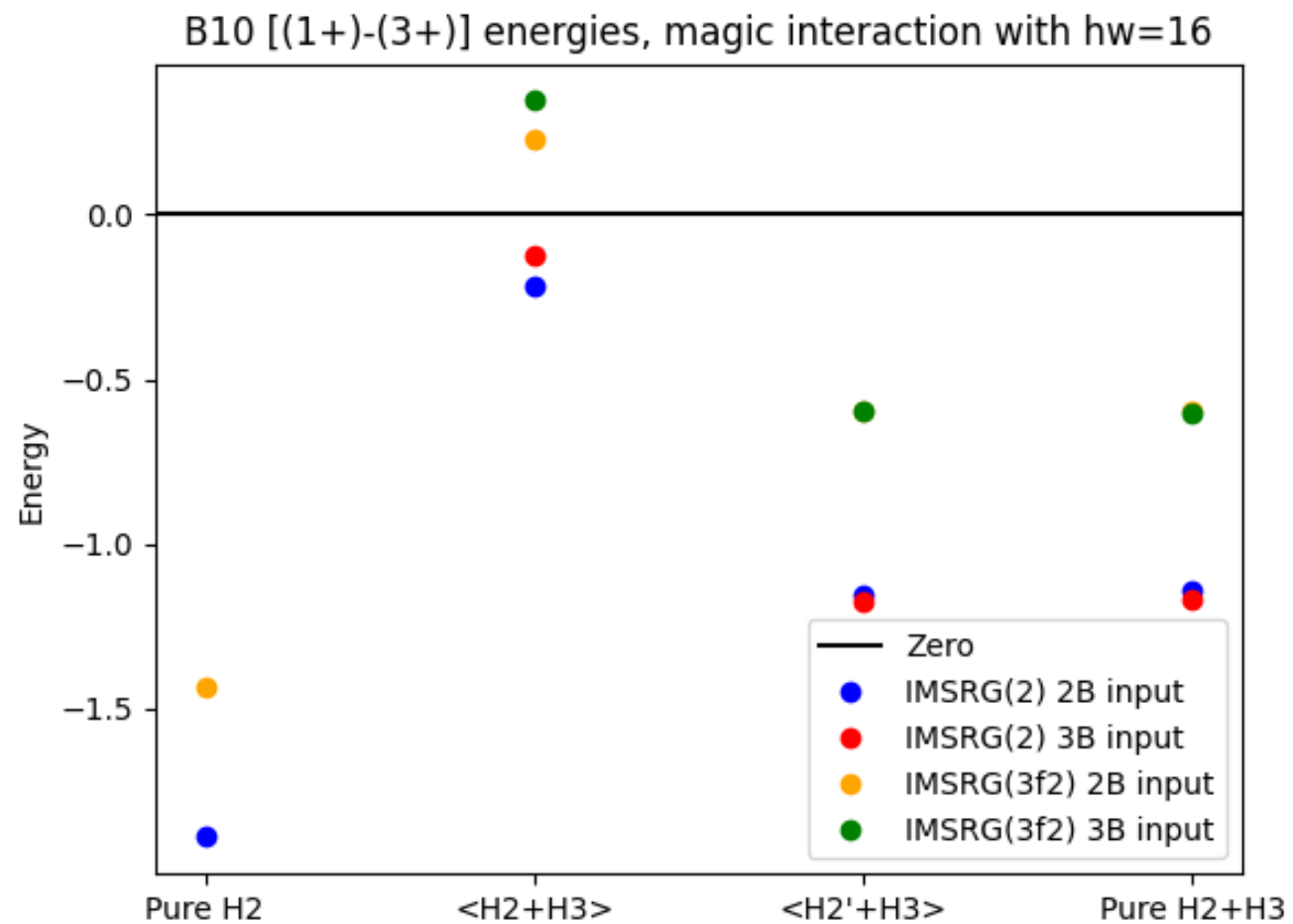
Approximate Three-Body Density

$$\rho_{ijklmn} \approx A\{\lambda_{il}\lambda_{jm}\lambda_{kn}\} + A\{\lambda_{il}\lambda_{jkmn}\}$$

$$\lambda_{ijkl} = \rho_{ijkl} - A\{\lambda_{ik}\lambda_{jl}\}$$

Na22 [(1+)-(3+)] energies, magic interaction with hw=16





IMSRG \rightarrow OSM

IMSRG decouples the interaction

- We extract the zero-body energy

OSM takes decoupled interaction as input

- Along with target nuclei
- Produces the energy spectra for target nuclei

We combine zero-body energy with energy spectra to get binding energy

BCH

Baker-Campbell-Hausdorff (BCH) formula

- $\frac{d}{ds} \Omega(s) = \sum_{k=0}^{\infty} \frac{B_k}{k!} [\Omega(s), \eta(s)]^{(k)}$

What we do

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Magnus Formulation

- $U(s) = e^{\Omega(s)}$, $\Omega(s)$ is anti-Hermitian

Using Baker-Campbell-Hausdorff (BCH) formula

- $H(s) = e^{\Omega(s)}H(0)e^{-\Omega(s)} = \sum_{k=0}^{\infty} \frac{1}{k!} [\Omega(s), H(0)]^{(k)}$

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Dependence of $U(s)$ on flow parameter s

- $\frac{d}{ds}U(s) = \eta(s)U(s)$, $\eta(s)$ is our generator

SRG flow equation

- $\frac{d}{ds}H(s) = [\eta(s), H(s)]$

Reference and modelspace

Reference is what we normal order with respect to

- In this case, Sn132

Modelspace is what we define our allowed transitions

- Nuclei with $Z < 50$ would need a different modelspace compared to nuclei with $Z > 50$
- Ex: 'jj45' for $Z < 50$, and 'jj55' for $Z > 50$