FACULTY OF SCIENCES

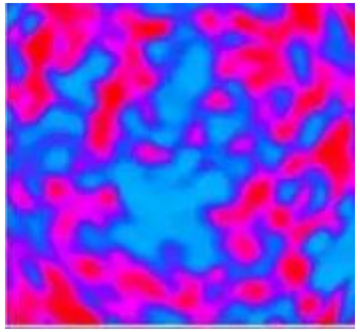
• DEPARTMENT OF PHYSICS AND ASTRONOMY

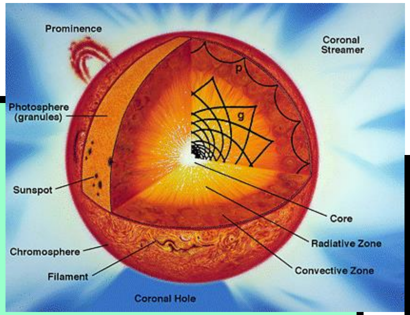
Elastic and Inelastic neutrino-nucleus scattering in the 10s of MeV energy range

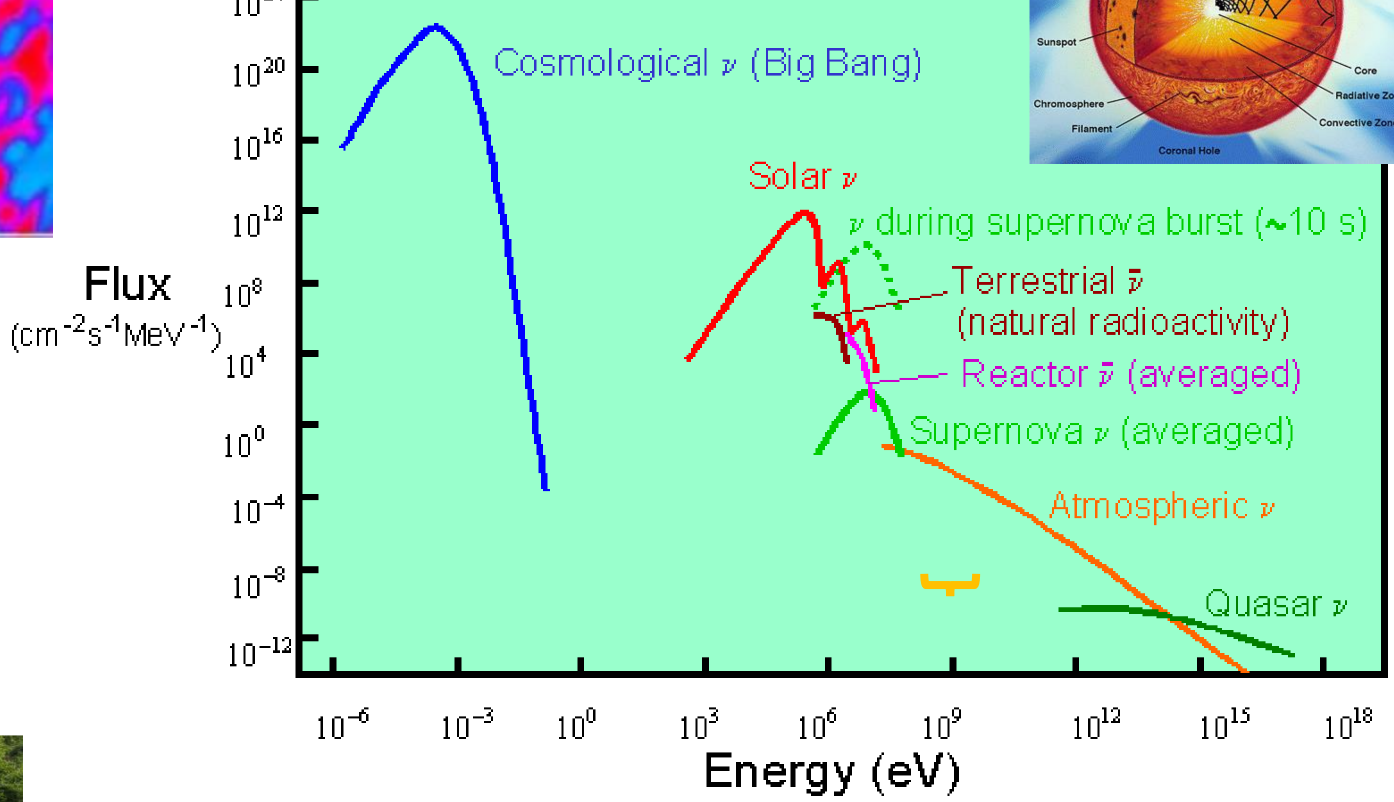
Natalie Jachowicz, K. Niewczas, A. Nikolakopoulos, V. Pandey, P. Vancraeyveld, N. Van Dessel, K. Vantournhout

GHENT UNIVERSITY

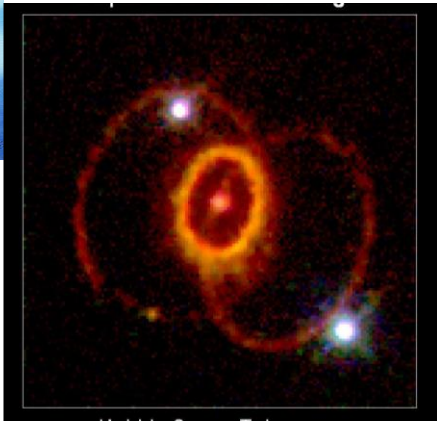
Neutrinos .....

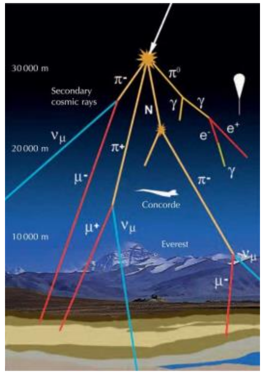


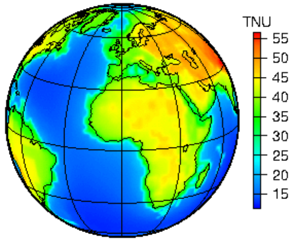


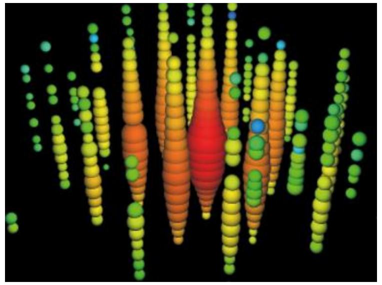


Flux on earth of neutrinos from various sources, in function of energy

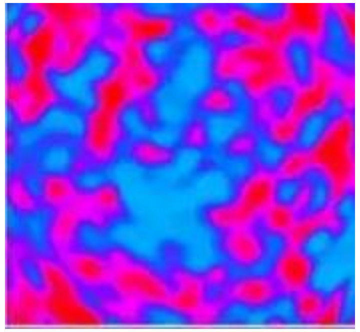


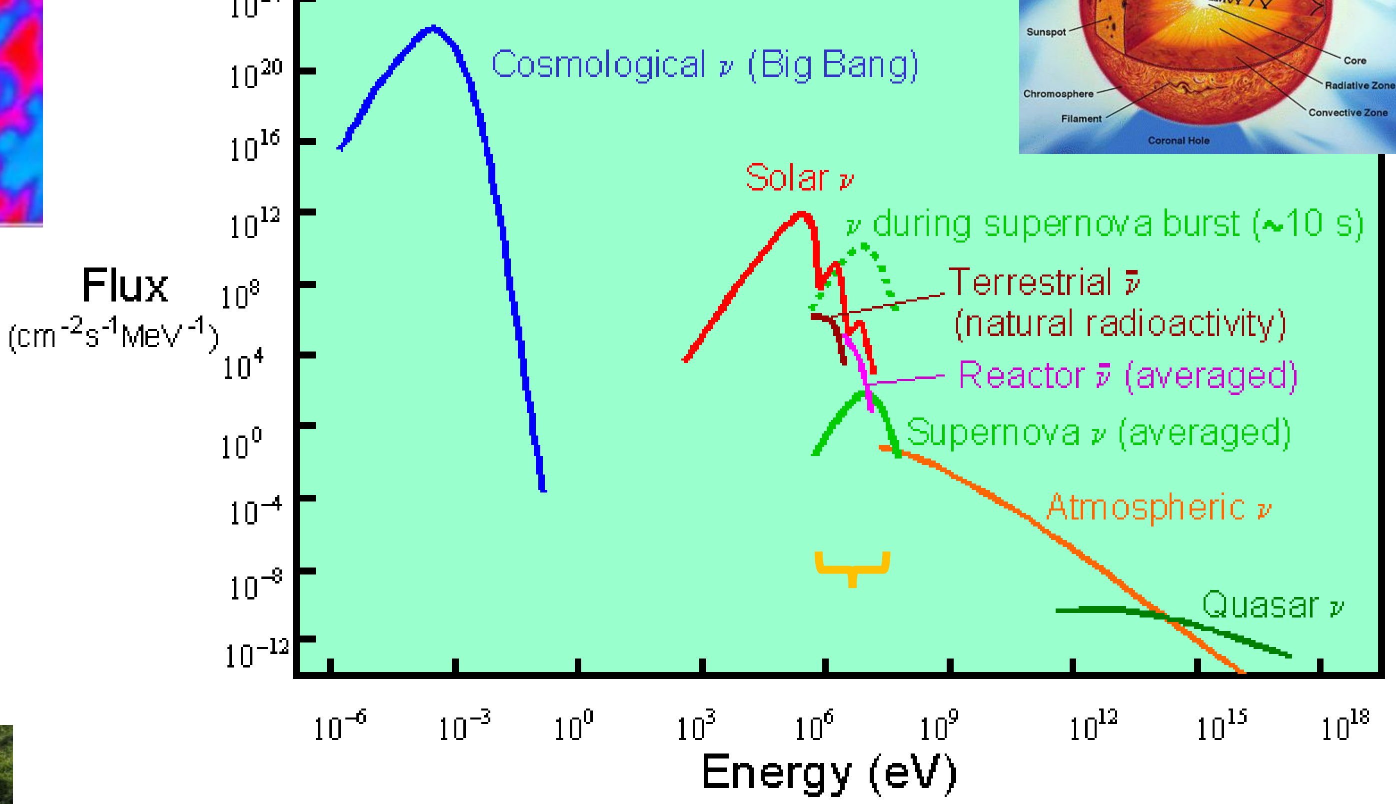




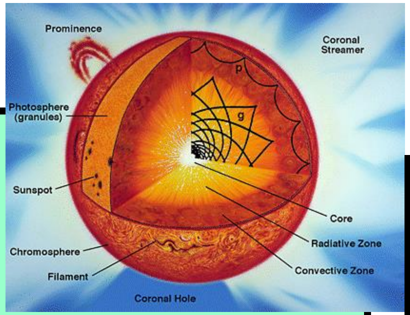


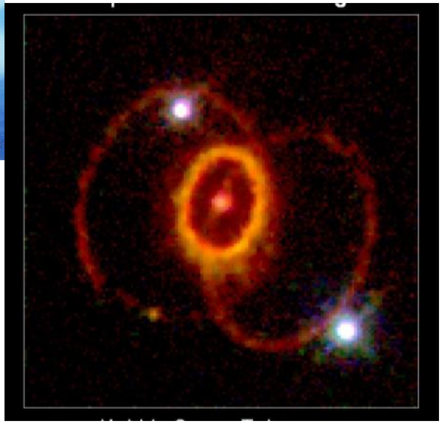
Neutrinos .....

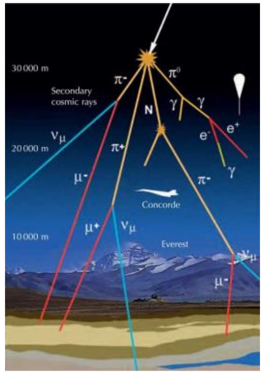


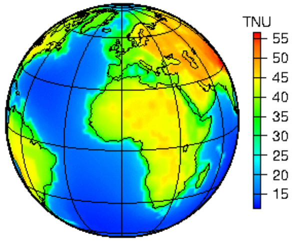


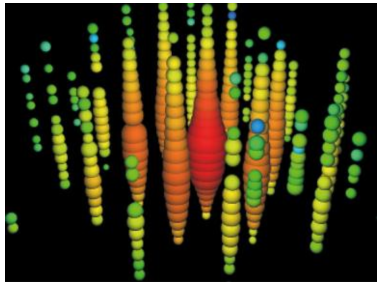
Flux on earth of neutrinos from various sources, in function of energy

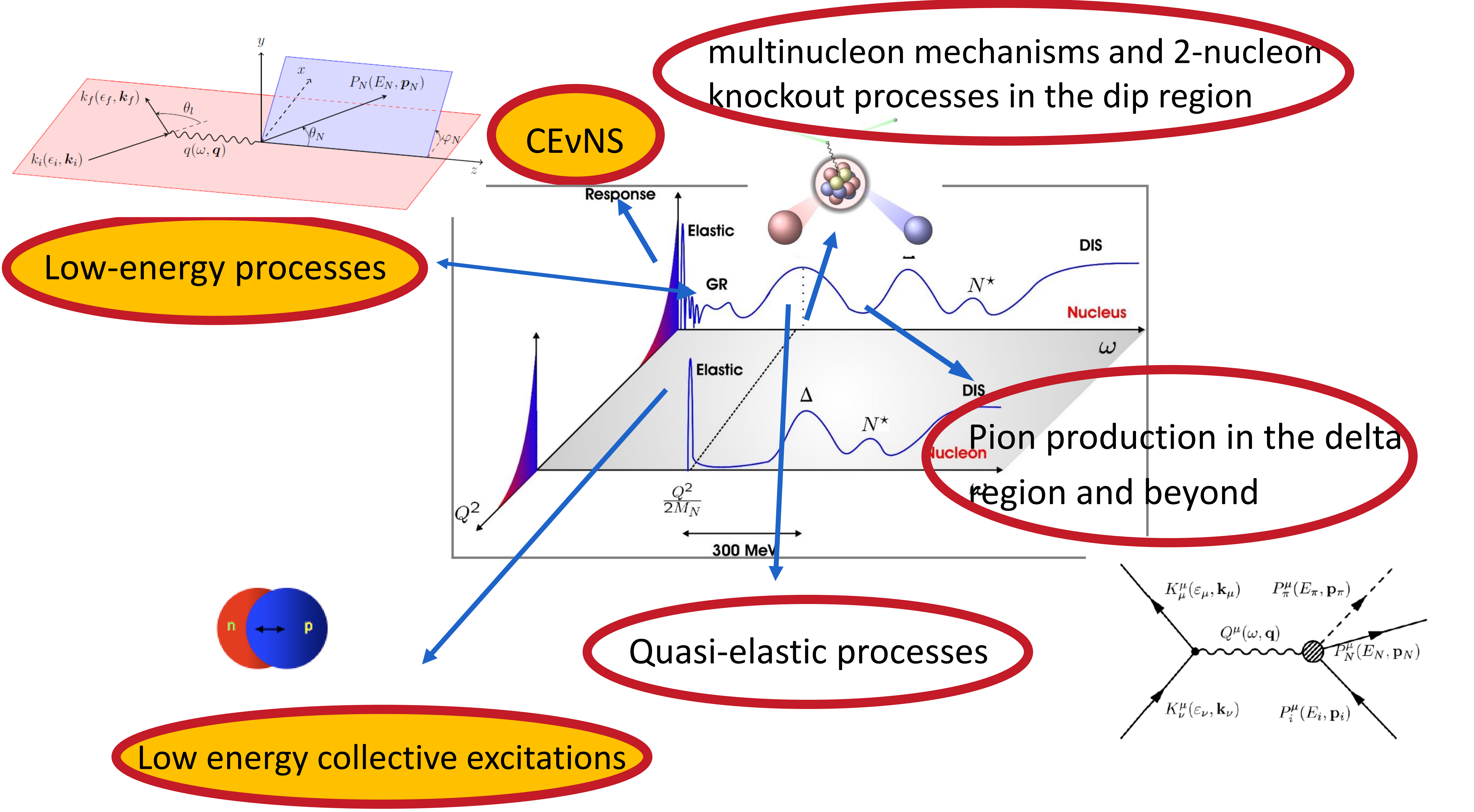











multinucleon mechanisms and 2-nucleon knockout processes in the dip region
CEvNS
Low-energy processes
Pion production in the delta region and beyond
Quasi-elastic processes
Low energy collective excitations
Response
Elastic
GR
DIS
$N^\star$
Nucleus
ω
Elastic
Δ
$N^\star$
DIS
Nucleon
Q^2
$\frac{Q^2}{2M_N}$
300 MeV
n
p
$k_f(\epsilon_f, \boldsymbol{k}_f)$
θ_l
$k_i(\epsilon_i, \boldsymbol{k}_i)$
$q(\omega, \boldsymbol{q})$
$P_N(E_N, \boldsymbol{p}_N)$
θ_N
φ_N
$K^\mu_\mu(\varepsilon_\mu, \mathbf{k}_\mu)$
$P^\mu_\pi(E_\pi, \mathbf{p}_\pi)$
$Q^\mu(\omega, \mathbf{q})$
$P^\mu_N(E_N, \mathbf{p}_N)$
$K^\mu_\nu(\varepsilon_\nu, \mathbf{k}_\nu)$
$P^\mu_i(E_i, \mathbf{p}_i)$

Neutrinos .....

What can we learn from low energy neutrinos ?

- Nuclear structure information
- Electroweak tests
- Neutrino oscillations
- Astrophysical neutrinos : a.o. core-collapse supernovae
- Neutrinonucleosynthesis
- BSM physics

How can we learn from these neutrinos ?

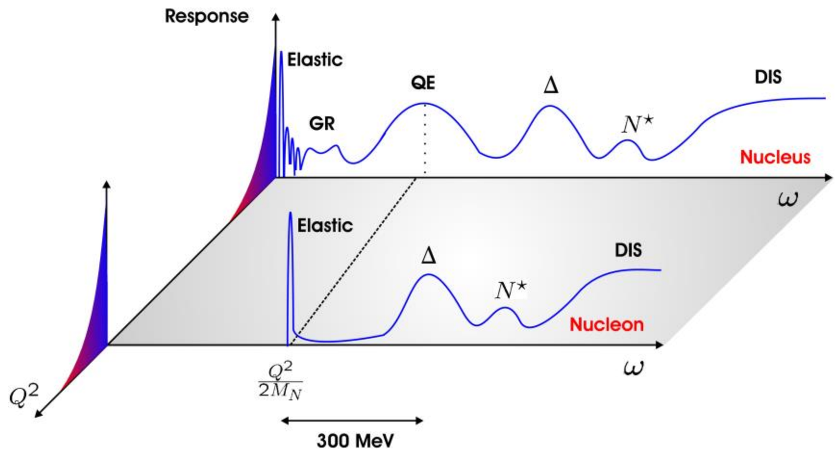
- Study their interactions : theory + experiment
- Detect them
  - Neutrino-electron scattering
  - Neutrino-hadron scattering

Neutrinos .....

Neutrino-hadron scattering ?

•little experimental data is available

- small cross sections
- (almost) no monochromatic neutrino beams



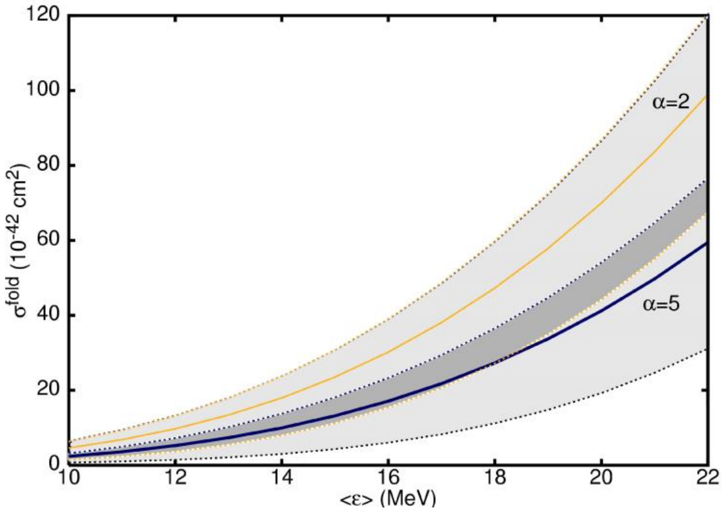
Uncertainties :

- one has to rely on theoretical predictions,
- uncertainties induced by model dependence
- and more fundamental uncertainties …

	B(GT$^-$)
HF+CRPA	8.8
GXPF1J	9.5
DD–ME2	11.3
SGII	12.3
SLy5	14.0
Exp.	9.9 ± 2.5

TABLE II. The total ^{56}Fe B(GT$^-$) strength, tabulated for various models from Ref. [36].

^{208}Pb



N.J. et al, PRC66, 065501 (2002) ;
E. Kolbe et al, PRC63, 025802 (2001) ;
J. Engel et al, PRD67, 013005 (2001)

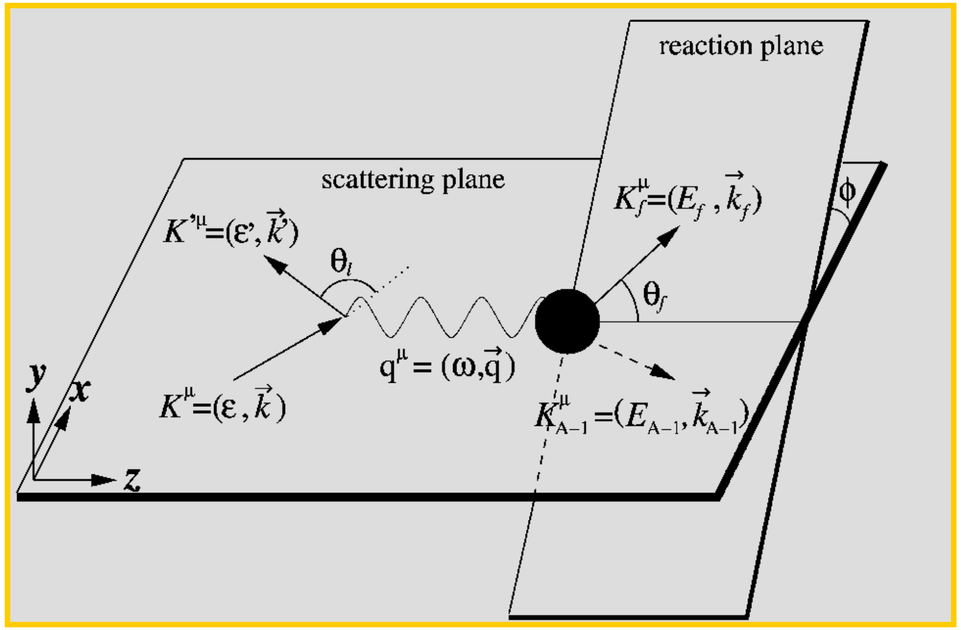
	$\langle\sigma_{DAR}\rangle$ (10^{-42}cm^2)
HF+CRPA	212.9
G–Matrix QRPA [35]	173.5
Phenomenological [47]	214
Hybrid [29]	240
Hybrid [36]	259
RHB+RQRPA [36]	263
LFG+RPA [38]	277
QRPA [64]	352
Exp. (KARMEN) [30]	256 ± 108 ± 43

TABLE I. The total charged–current (ν_e,^{56}Fe) cross section value, folded with a DAR electron neutrino spectrum, tabulated for various models.

Modeling low-energy inelastic neutrino-nucleus scattering

$$\frac{d^2\sigma}{d\Omega\, d\omega} = (2\pi)^4\ k_f \varepsilon_f \sum_{s_f, s_i} \frac{1}{2J_i + 1} \sum_{M_f, M_i} \left| \left\langle f \left| \widehat{H}_W \right| i \right\rangle \right|^2$$

$$\widehat{H}_W = \frac{G}{\sqrt{2}} \int d\vec{x}\, \hat{J}_{\mu,lepton}(\vec{x})\, \hat{J}^{\mu,hadron}(\vec{x})$$



Hadron current

$$J^{\mu} = F_1(Q^2)\gamma^{\mu} + i\frac{\kappa}{2M_N}F_2(Q^2)\sigma^{\mu\nu}q_{\nu} + G_A(Q^2)\gamma^{\mu}\gamma_5 + \frac{1}{2M_N}G_P(Q^2)q^{\mu}\gamma_5$$

Lepton tensor

$$l_{\alpha\beta} \equiv \overline{\sum_{s,s'}}[\overline{u}_l\gamma_{\alpha}(1-\gamma_5)u_l]^{\dagger}[\overline{u}_{\nu}\gamma_{\beta}(1-\gamma_5)u_{\nu}]$$

Neutrino-nucleus cross sections

$$\vec{J}^{\alpha}_{V}(\vec{x}) = \vec{J}^{\alpha}_{convection}(\vec{x}) + \vec{J}^{\alpha}_{magnetization}(\vec{x})$$

with

$$\vec{J}^{\alpha}_{c}(\vec{x}) = \frac{1}{2Mi}\sum_{i=1}^{A} G_E^{i,\alpha}\left[\delta(\vec{x}-\vec{x}_i)\overrightarrow{\nabla}_i - \overleftarrow{\nabla}_i\,\delta(\vec{x}-\vec{x}_i)\right],$$

$$\vec{J}^{\alpha}_{m}(\vec{x}) = \frac{1}{2M}\sum_{i=1}^{A} G_M^{i,\alpha}\,\overrightarrow{\nabla}\times\vec{\sigma}_i\,\delta(\vec{x}-\vec{x}_i),$$

$$\vec{J}^{\alpha}_{A}(\vec{x}) = \sum_{i=1}^{A} G_A^{i,\alpha}\,\vec{\sigma}_i\,\delta(\vec{x}-\vec{x}_i),$$

$$J_V^{0,\alpha}(\vec{x}) = \rho_V^{\alpha}(\vec{x}) = \sum_{i=1}^{A} G_E^{i,\alpha}\,\delta(\vec{x}-\vec{x}_i),$$

$$J_A^{0,\alpha}(\vec{x}) = \rho_A^{\alpha}(\vec{x}) = \frac{1}{2Mi}\sum_{i=1}^{A} G_A^{i,\alpha}\,\vec{\sigma}_i\cdot\left[\delta(\vec{x}-\vec{x}_i)\overrightarrow{\nabla}_i - \overleftarrow{\nabla}_i\,\delta(\vec{x}-\vec{x}_i)\right]$$

$$J_P^{0,\alpha}(\vec{x}) = \rho_P^{\alpha}(\vec{x}) = \frac{m_\mu}{2M}\sum_{i=1}^{A} G_P^{i,\alpha}\,\overrightarrow{\nabla}\cdot\vec{\sigma}_i\,\delta(\vec{x}-\vec{x}_i)$$

for NC reactions

$$G_E^{V,o} = \left(\frac{1}{2}-\sin^2\theta_W\right)\tau_3 - \sin^2\theta_W,$$

$$G_M^{V,o} = \left(\frac{1}{2}-\sin^2\theta_W\right)(\mu_p-\mu_n)\,\tau_3 - \sin^2\theta_W\,(\mu_p+\mu_n)$$

$$G^{A,0} = g_a\,\frac{\tau_3}{2} = -\frac{1.262}{2}\,\tau_3$$

for CC reactions

$$G_E^{V,\pm} = \tau_\pm$$

$$G_M^{V,\pm} = (\mu_p-\mu_n)\,\tau_\pm$$

$$G^{A,\pm} = g_a\,\tau_\pm = -1.262\,\tau_\pm$$

$G = (1+Q^2/M^2)^{-2}$ Q^2 dependence : dipole or BBBA parametrization :

Neutrino-nucleus cross sections

$$\left(\frac{d^2\sigma}{d\omega d\Omega}\right)_\nu = \frac{G_F^2 \cos^2\theta_c}{(4\pi)^2} \left(\frac{2}{2J_i+1}\right) \varepsilon_f \kappa_f \;\; \zeta^2\left(Z', \varepsilon_f, |q|\right) \left[\sum_{J=0}^{\infty} \sigma_{CL,\nu}^J + \sum_{J=1}^{\infty} \sigma_{T,\nu}^J\right]$$

$$\sigma_{CL,\nu}^J = \left[v_\nu^{\mathcal{M}} R_\nu^{\mathcal{M}} + v_\nu^{\mathcal{L}} R_\nu^{\mathcal{L}} + 2\, v_\nu^{\mathcal{M}\mathcal{L}} R_\nu^{\mathcal{M}\mathcal{L}}\right],$$

$$\sigma_{T,\nu}^J = \left[v_\nu^T R_\nu^T \pm 2\, v_\nu^{TT} R_\nu^{TT}\right],$$

Neutrino-nucleus cross sections

$$\left(\frac{d^2\sigma}{d\omega d\Omega}\right)_\nu = \frac{G_F^2\ \cos^2\theta_c}{(4\pi)^2}\left(\frac{2}{2J_i+1}\right)\varepsilon_f\kappa\ \zeta^2\left(Z',\varepsilon_f,|q|\right)\left[\sum_{J=0}^{\infty}\sigma_{CL,\nu}^J+\sum_{J=1}^{\infty}\sigma_{T,\nu}^J\right]$$

$$v_\nu^{\mathcal{M}} = \left[1+\frac{\kappa_f}{\varepsilon_f}\cos\theta\right],$$

$$v_\nu^{\mathcal{L}} = \left[1+\frac{\kappa_f}{\varepsilon_f}\cos\theta-\frac{2\varepsilon_i\varepsilon_f}{|\vec{q}|^2}\left(\frac{\kappa_f}{\varepsilon_f}\right)^2\sin^2\theta\right],$$

$$v_\nu^{\mathcal{ML}} = \left[\frac{\omega}{|\vec{q}|}\left(1+\frac{\kappa_f}{\varepsilon_f}\cos\theta\right)+\frac{m_l^2}{\varepsilon_f|\vec{q}|}\right],$$

$$v_\nu^{T} = \left[1-\frac{\kappa_f}{\varepsilon_f}\cos\theta+\frac{\varepsilon_i\varepsilon_f}{|\vec{q}|^2}\left(\frac{\kappa_f}{\varepsilon_f}\right)^2\sin^2\theta\right],$$

$$v_\nu^{TT} = \left[\frac{\varepsilon_i+\varepsilon_f}{|\vec{q}|}\left(1-\frac{\kappa_f}{\varepsilon_f}\cos\theta\right)-\frac{m_l^2}{\varepsilon_f|\vec{q}|}\right],$$

$$\left[v_\nu^{\mathcal{M}}R_\nu^{\mathcal{M}}+v_\nu^{\mathcal{L}}R_\nu^{\mathcal{L}}\ldots\right.$$

$$, \quad = \quad \left[v_\nu^{T}R_\nu^{T}\pm 2\ldots\right.$$

$$R_\nu^{\mathcal{M}} = |\langle J_f||\widehat{\mathcal{M}}_J^\nu(|\vec{q}|)||J_i\rangle|^2,$$

$$R_\nu^{\mathcal{L}} = |\langle J_f||\widehat{\mathcal{L}}_J^\nu(|\vec{q}|)||J_i\rangle|^2,$$

$$R_\nu^{\mathcal{ML}} = \mathcal{R}\left[\langle J_f||\widehat{\mathcal{L}}_J^\nu(|\vec{q}|)||J_i\rangle\langle J_f||\widehat{\mathcal{M}}_J^\nu(|\vec{q}|)||J_i\rangle^*\right],$$

$$R_\nu^{T} = \left[|\langle J_f||\widehat{\mathcal{J}}_J^{mag,\nu}(|\vec{q}|)||J_i\rangle|^2+|\langle J_f||\widehat{\mathcal{J}}_J^{el,\nu}(|\vec{q}|)||J_i\rangle|^2\right],$$

$$R_\nu^{TT} = \mathcal{R}\left[\langle J_f||\widehat{\mathcal{J}}_J^{mag,\nu}(|\vec{q}|)||J_i\rangle\langle J_f||\widehat{\mathcal{J}}_J^{el,\nu}(|\vec{q}|)||J_i\rangle^*\right].$$

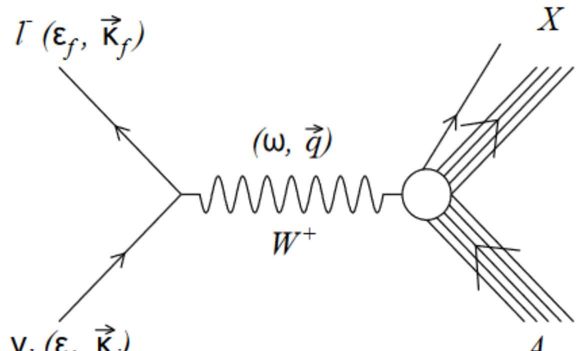
Neutrino-nucleus cross sections

$$\widehat{\mathcal{M}}_{JM}(\kappa) = \int d\vec{x} \; \left[j_J\,(\kappa r)\; Y_J^M\;(\Omega_x) \right] \; \hat{J}_0(\vec{x}) \quad ,$$

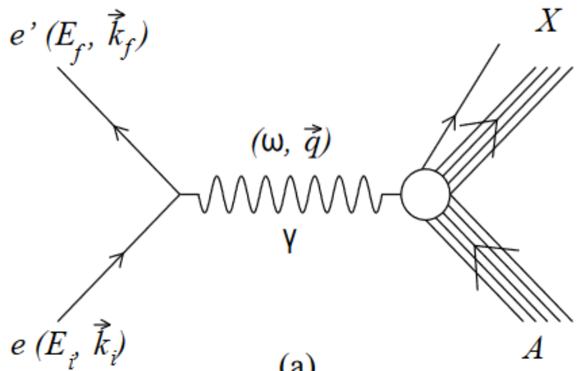
$$\widehat{\mathcal{L}}_{JM}(\kappa) = \frac{i}{\kappa} \int d\vec{x} \; \left[\vec{\nabla} \left(j_J\,(\kappa r)\; Y_J^M\;(\Omega_x) \right) \right] \cdot \widehat{\vec{J}}(\vec{x}) \quad ,$$

$$\widehat{\mathcal{J}}^{el}_{JM}(\kappa) = \frac{1}{\kappa} \int d\vec{x} \; \left[\vec{\nabla} \times \left(j_J\,(\kappa r)\; \vec{\mathcal{Y}}^M_{J,J}\;(\Omega_x) \right) \right] \cdot \widehat{\vec{J}}(\vec{x}) \quad ,$$

$$\widehat{\mathcal{J}}^{mag}_{JM}(\kappa) = \int d\vec{x} \; \left[j_J\,(\kappa r)\; \vec{\mathcal{Y}}^M_{J,J}\;(\Omega_x) \right] \cdot \widehat{\vec{J}}(\vec{x}) \quad .$$



(b)



(a)

Neutrino-nucleus cross sections

$$\widehat{\mathcal{M}}_{JM}(\kappa) = \int d\vec{x}\ [j_J(\kappa r)\ Y_J^M(\Omega_x)]\ \hat{J}_0(\vec{x})\ ,$$

$$\widehat{\mathcal{L}}_{JM}(\kappa) = \frac{i}{\kappa} \int d\vec{x}\ \left[\vec{\nabla}\left(j_J(\kappa r)\ Y_J^M(\Omega_x)\right)\right] \cdot \widehat{\vec{J}}(\vec{x})\ ,$$

$$\widehat{\mathcal{J}}^{el}_{JM}(\kappa) = \frac{1}{\kappa} \int d\vec{x}\ \left[\vec{\nabla} \times \left(j_J(\kappa r)\ \vec{\mathcal{Y}}^M_{J,J}(\Omega_x)\right)\right] \cdot \widehat{\vec{J}}(\vec{x})\ ,$$

$$\widehat{\mathcal{J}}^{mag}_{JM}(\kappa) = \int d\vec{x}\ \left[j_J(\kappa r)\ \vec{\mathcal{Y}}^M_{J,J}(\Omega_x)\right] \cdot \widehat{\vec{J}}(\vec{x})\ .$$

→

$$\langle a||\widehat{\mathcal{M}}_J^{Coul}[\widehat{\rho}_V]||b\rangle = G_E(Q^2) \int \mathrm{d}r\ \langle a||\tau_\pm j_J(qr) Y_J(\Omega_1)||b\rangle_r$$

$$\langle a||\widehat{\mathcal{M}}_J^{Coul}[\widehat{\rho}_A]||b\rangle = \frac{G_A(Q^2)}{2m_N i} \int \mathrm{d}r\ \langle a||\tau_\pm j_J(qr) Y_J(\Omega_1)\,\boldsymbol{\sigma}_1 \cdot \left(\overrightarrow{\nabla}_1 - \overleftarrow{\nabla}_1\right)||b\rangle_r$$

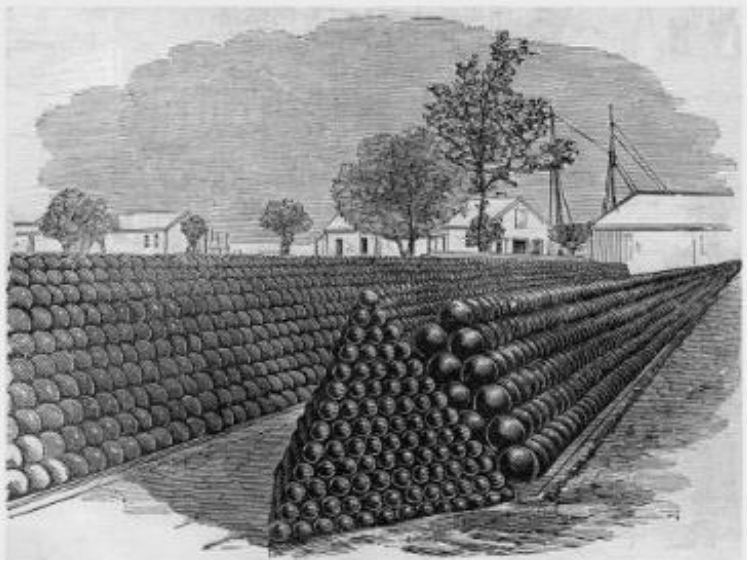
$$\langle a||\widehat{\mathcal{O}}_J^\lambda\left[\widehat{J}_{conv}\right]||b\rangle = \frac{G_E(Q^2)}{2m_N i} \int \mathrm{d}r\ \langle a||\tau_\pm j_{J+\lambda}(qr) \times \left[Y_{J+\lambda}(\Omega_1) \otimes \left(\overrightarrow{\nabla}_1 - \overleftarrow{\nabla}_1\right)\right]_J ||b\rangle_r$$

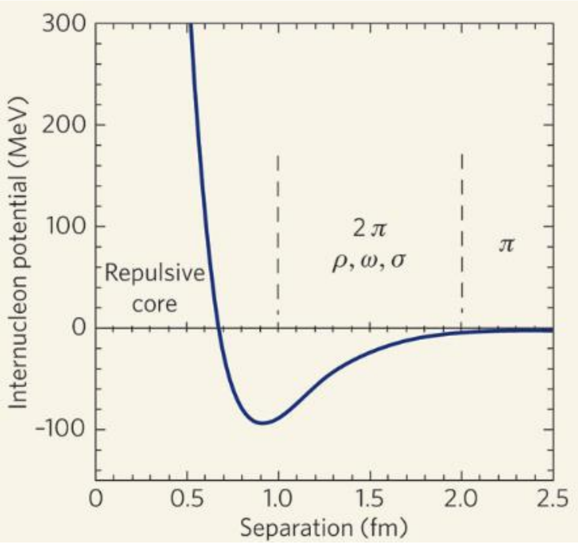
$$\langle a||\widehat{\mathcal{O}}_J^\lambda\left[\widehat{J}_{magn}\right]||b\rangle = i\sqrt{6}q\frac{G_M(Q^2)}{2m_N} \int \mathrm{d}r \sum_{\eta=\pm1} \sqrt{J+\lambda+\delta_{\eta,+1}} \times \begin{Bmatrix} J & J+\lambda & 1 \\ 1 & 1 & J+\lambda+\eta \end{Bmatrix} \times \langle a||\tau_\pm j_{J+\lambda+\eta}(qr)\left[Y_{J+\lambda+\eta}(\Omega_1)\otimes \boldsymbol{\sigma}_1\right]_J||b\rangle_r$$

$$\langle a||\widehat{\mathcal{O}}_J^\lambda\left[\widehat{J}_A\right]||b\rangle = G_A(Q^2) \int \mathrm{d}r\ \langle a||\tau_\pm j_{J+\lambda}(qr)\left[Y_{J+\lambda}(\Omega_1)\otimes \boldsymbol{\sigma}_1\right]_J||b\rangle_r$$

A model for the nucleus

- Nuclear radius $\approx 1.2A^{\frac{1}{3}}$ fm
- Nucleon is a diffuse system
  - Hard core (repulsion) ≈ 0.5 fm
  - RMS charge radius from (e,e') $= 0.897(18)$ fm
- $0.07 \lesssim \text{NPF} \lesssim 0.42$
  - closest packing fraction of spheres ≈ 0.74
  - packing fraction of Argon liquid ≈ 0.032
  - packing fraction of Argon gas $\approx 3.75 \cdot 10^{-5}$
- The nuclear medium is a rather dense quantum liquid

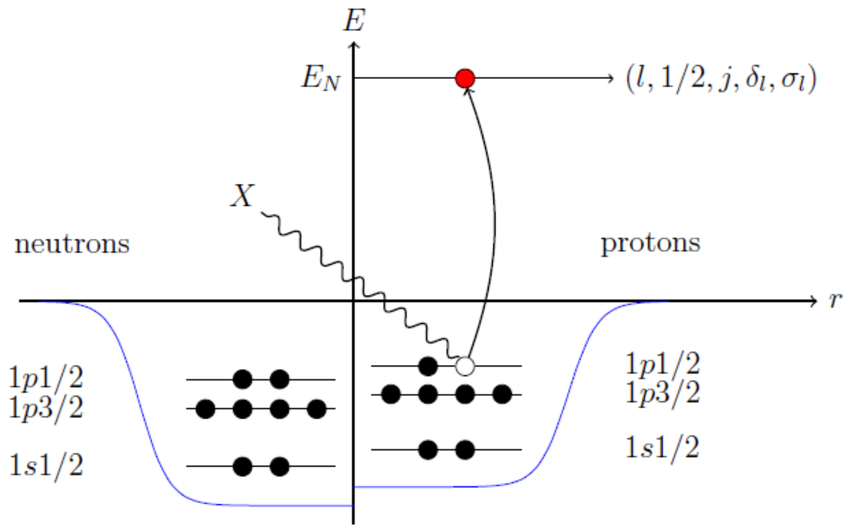




C. Colle, PhD, UGent 2017

→ Identify the right degrees of freedom and main effects for each kinematic region
→ Identify the relevant corrections, correlations to be taken into account

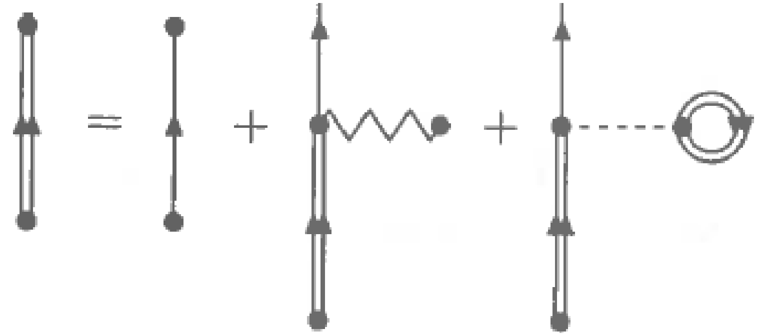
A model for the nucleus



- Starting point : mean-field nucleus with Hartree-Fock single-particle wave functions
- Skyrme SkE2 force used to build the potential
- Pauli blocking
- binding

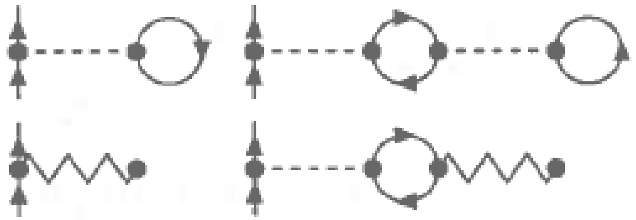
Hartree-Fock mean field

$$G^{HF}(\alpha,\beta;E) = G^{(0)}(\alpha,\beta;E) + \sum_{\gamma,\delta} G^{(0)}(\alpha,\gamma;E)\Sigma^{HF}(\gamma,\delta)G^{HF}(\delta,\beta;E)$$



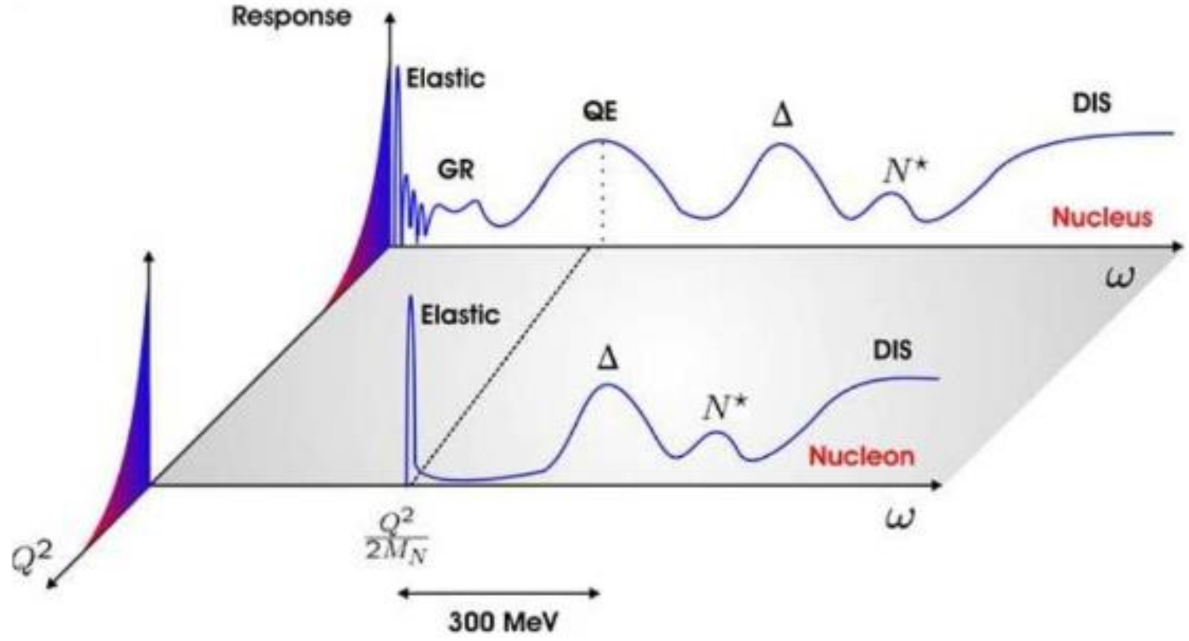
- Mean field already contains correlations !
- Nucleons feel the presence of the others through the averaged field

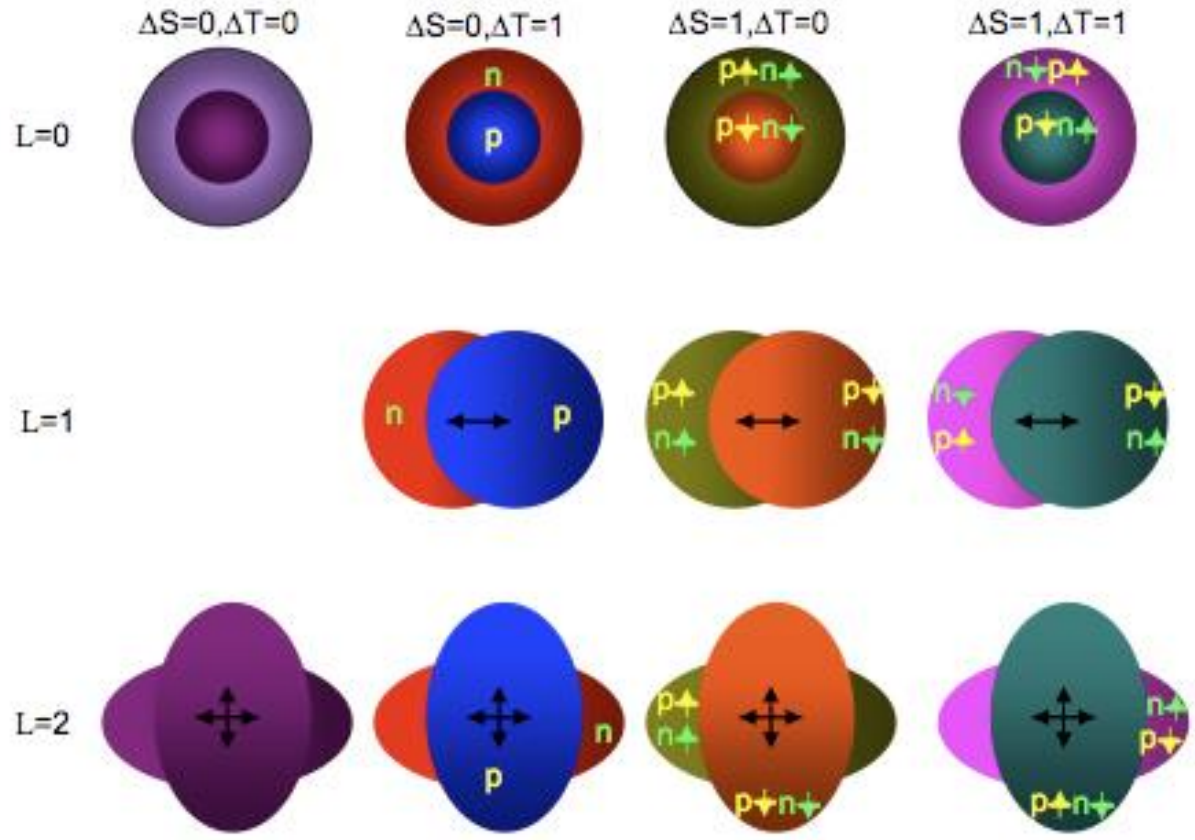
$$\Sigma^{HF}(\gamma,\delta;E) = -\langle\gamma|\,U\,|\delta\rangle - i\int\frac{dE'}{2\pi}\sum_{\mu\nu}\langle\gamma\mu|\,V\,|\delta\nu\rangle\,G^{HF}(\nu\mu;E')$$



Long-range RPA correlations

- Correlations over the whole size of the nucleus
- Redistribute the incoming energy transfer to the nucleus over all the nuclear constituents.
- They manifest themselves in collective excitations such as giant resonances

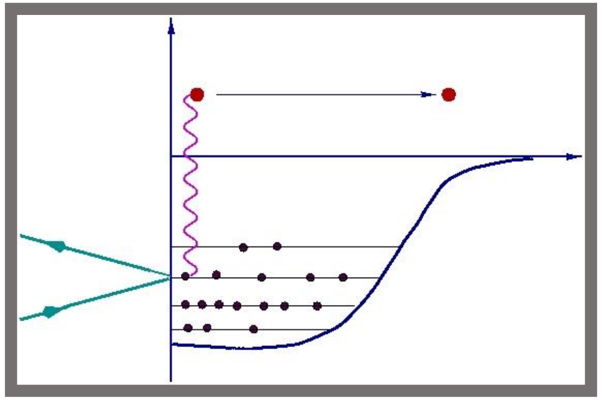


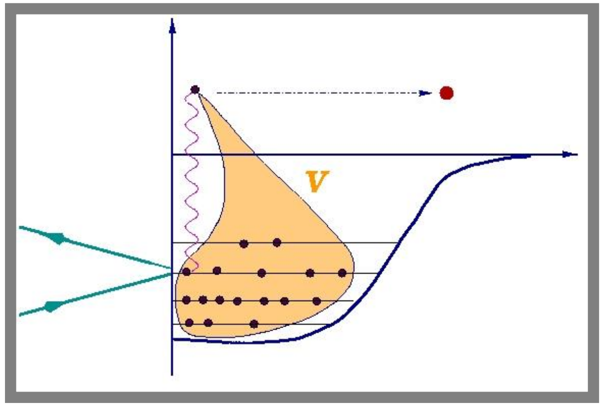


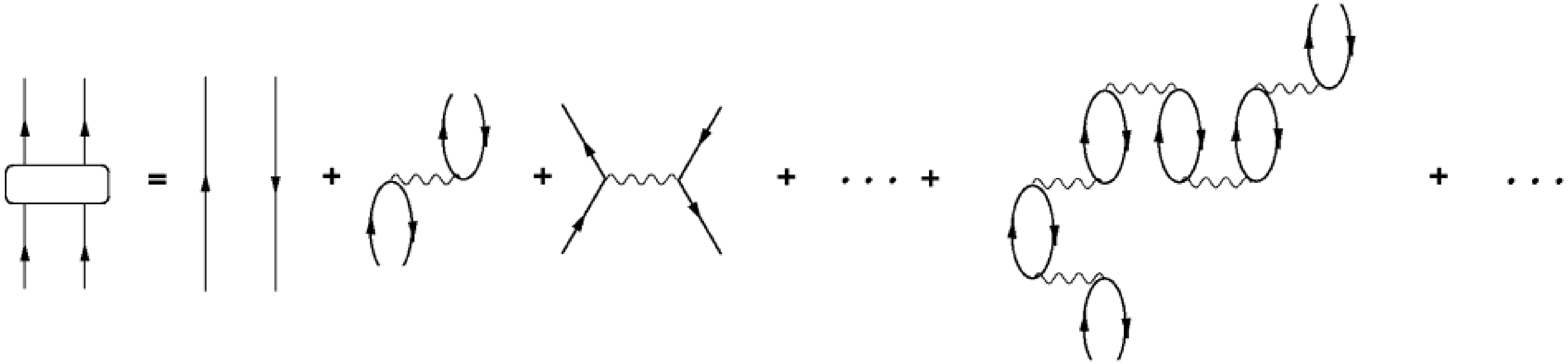
https://cyclotron.tamu.edu/research/nuclear-structure/

Long-range RPA correlations

- Green's function approach
- Skyrme SkE2 residual interaction
- self-consistent calculations







$$|\Psi_{RPA}\rangle = \sum_c \left\{ X_{(\Psi,C)} \left|ph^{-1}\right\rangle - Y_{(\Psi,C)} \left|hp^{-1}\right\rangle \right\}$$

$$\Pi^{(RPA)}(x_1,x_2;\omega) = \Pi^{(0)}(x_1,x_2;\omega) + \frac{1}{\hbar}\int dx \int dx' \; \Pi^{(0)}(x_1,x;\omega)\, \widetilde{V}(x,x')\, \Pi^{(RPA)}(x',x_2;\omega)$$

Hartree-Fock-CRPA

Solving the RPA equations in coordinate space

$$|\Psi_C(E)\rangle = \left|ph^{-1}(E)\right\rangle + \int dx_1 \int dx_2 \; \widetilde{V}(x_1, x_2)$$

$$\sum_{c'} \mathcal{P} \int d\varepsilon_{p'} \left[\frac{\psi_{h'}(x_1)\psi^{\dagger}_{p'}(x_1, \varepsilon_{p'})}{E - \varepsilon_{p'h'}} \left|p'h'^{-1}(\varepsilon_{p'h'})\right\rangle \right.$$

$$\left. - \frac{\psi^{\dagger}_{h'}(x_1)\psi_{p'}(x_1, \varepsilon_{p'})}{E + \varepsilon_{p'h'}} \left|h'p'^{-1}(-\varepsilon_{p'h'})\right\rangle \right] \left\langle \Psi_0 \left| \hat{\psi}^{\dagger}(x_2)\hat{\psi}(x_2) \right| \Psi_C(E) \right\rangle$$

What we really need is transition densities :

$$\langle \Psi_0 || X_{\eta J} || \Psi_C(J; E) \rangle_r = - \langle h || X_{\eta J} || p(\varepsilon_{ph}) \rangle_r$$

$$+ \sum_{\mu,\nu} \int dr_1 \int dr_2 \; U^J_{\mu\nu}(r_1, r_2) \; \mathcal{R} \left(R^{(0)}_{\eta\mu;J}(r, r_1; E) \right) \; \langle \Psi_0 || X_{\nu J} || \Psi_C(J; E) \rangle_{r_2}$$

$$\int dr \int dr' \; R^{(0)}_{\eta\mu;JM}(r, r'; E) = \frac{1}{\hbar} \int dx \int dx' \;\; X_{\eta JM}(x) \;\; \Pi^{(0)}(x, x'; \omega) \;\; X^{\dagger}_{\eta' JM}(x')$$

So in the end we have to solve a set of coupled equations, that after discretizing on a mesh in coordinate space, translates into a matrix inversion :

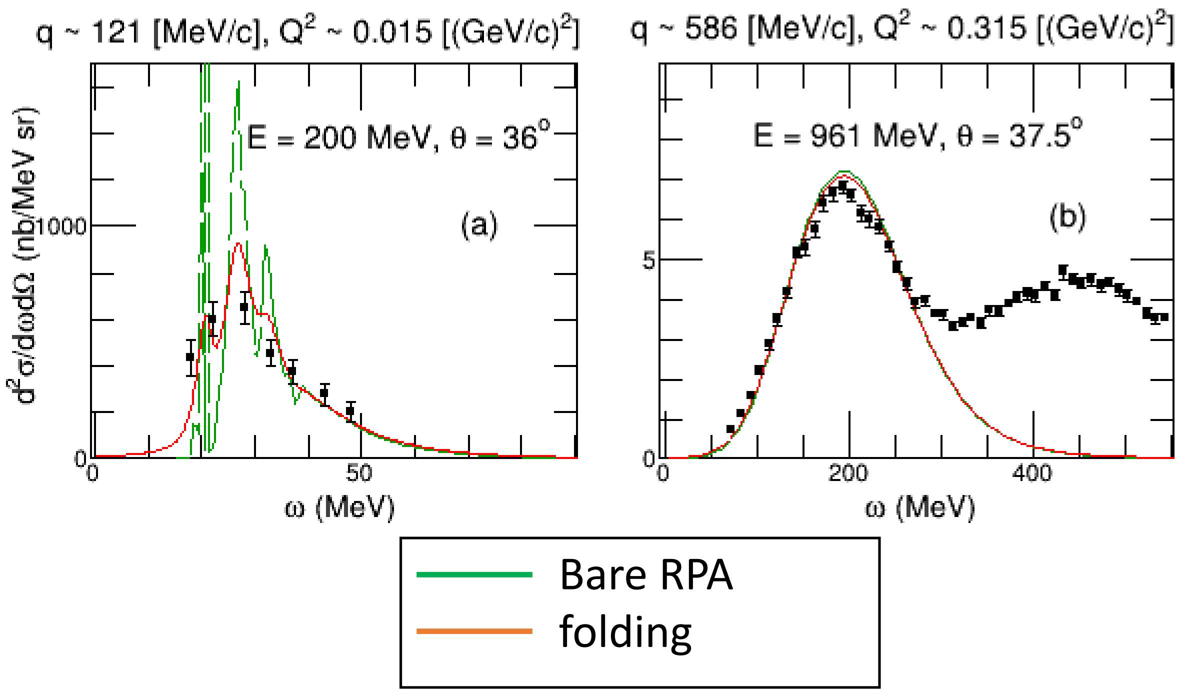
$$\rho^{RPA}_C = - \frac{1}{1 - R\,U} \;\; \rho^{HF}_C$$

Hartree-Fock-CRPA

- Final state interactions :

-taken into account through the calculations of the wave function of the outgoing nucleon in the (real) nuclear potential generated using the Skyrme force

-influence of the spreading width of the particle states is implemented through a folding procedure



$$R'(q,\omega') = \int_{-\infty}^{\infty} d\omega \; R(q,\omega) \; L(\omega,\omega'),$$

$$L(\omega,\omega') = \frac{1}{2\pi}\left[\frac{\Gamma}{(\omega-\omega')^2+(\Gamma/2)^2}\right].$$

Hartree-Fock-CRPA

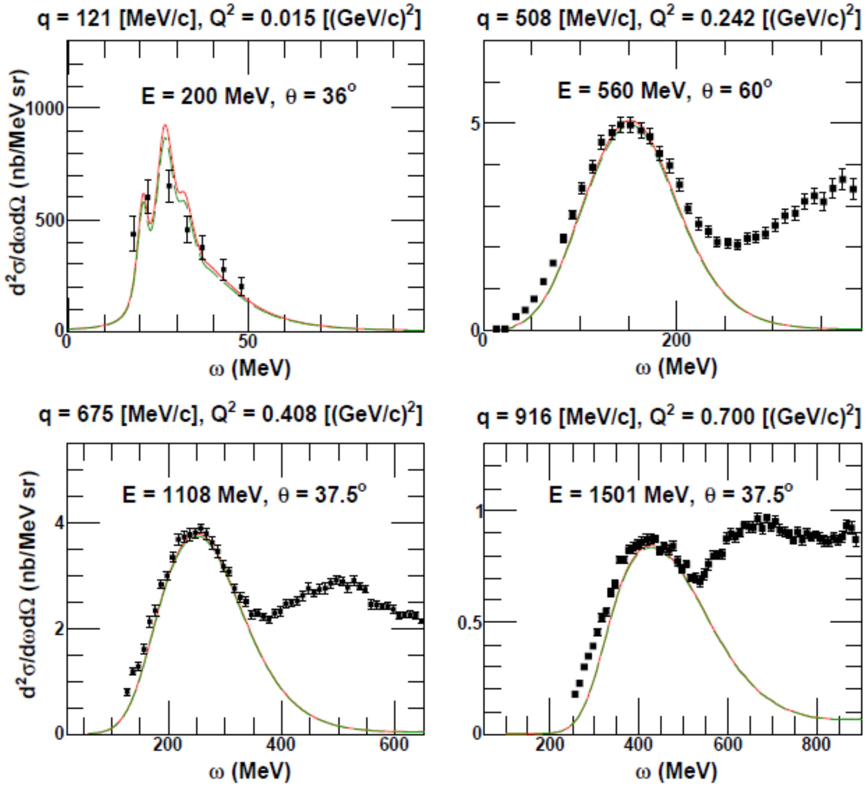
•Coulomb correction for the outgoing lepton in charged-current interactions :

- ✓ Low energies : Fermi function

$$F(Z', E) = \frac{2\pi\eta}{1 - e^{-2\pi\eta}} \qquad \eta \sim \mp Z'\alpha$$

- ✓ High energies : modified effective momentum approximation (J. Engel, PRC57,2004

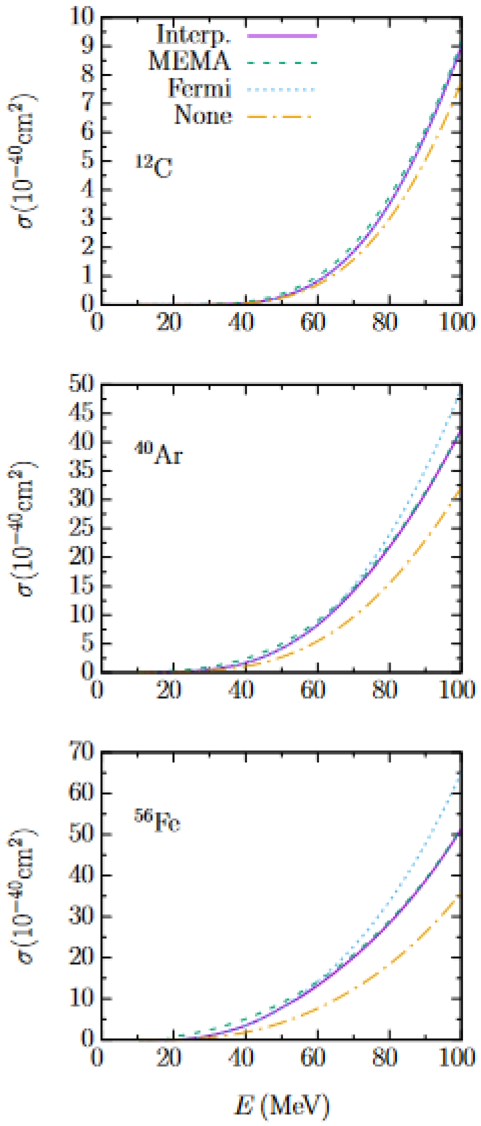
$$q_{eff} \;=\; q + 1.5 \left(\frac{Z'\alpha\hbar c}{R}\right),$$



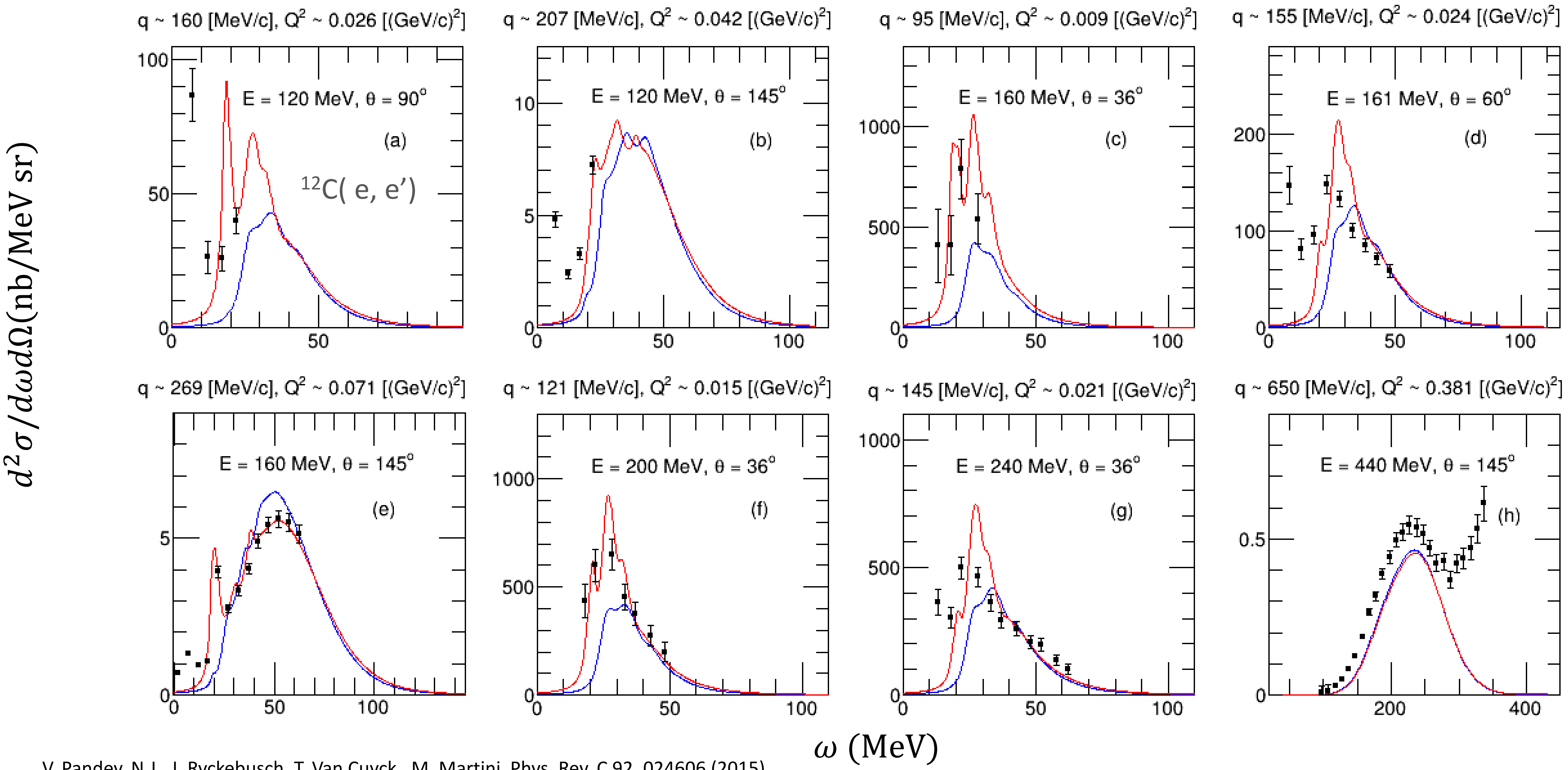
$$\Psi_l^{eff} \;=\; \zeta(Z', E, q)\; \Psi_l\,,$$

$$\zeta(Z', E, q) \;=\; \sqrt{\frac{q_{eff}E_{eff}}{qE}}$$

Uncorrected	MEMA



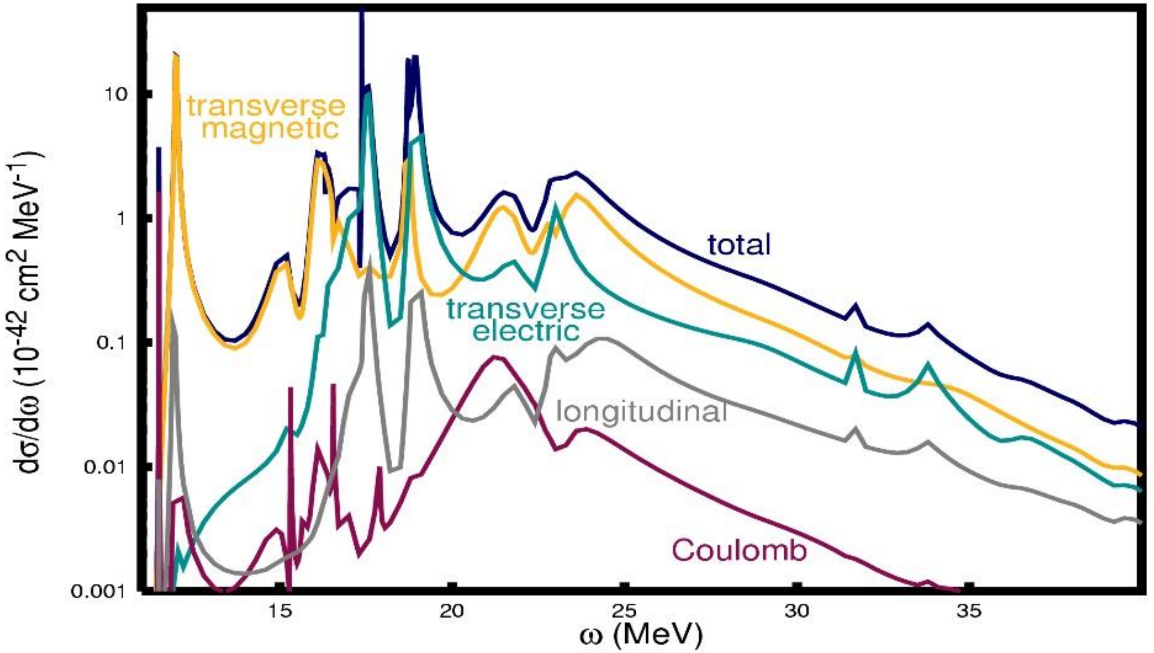
Comparison with electron scattering data

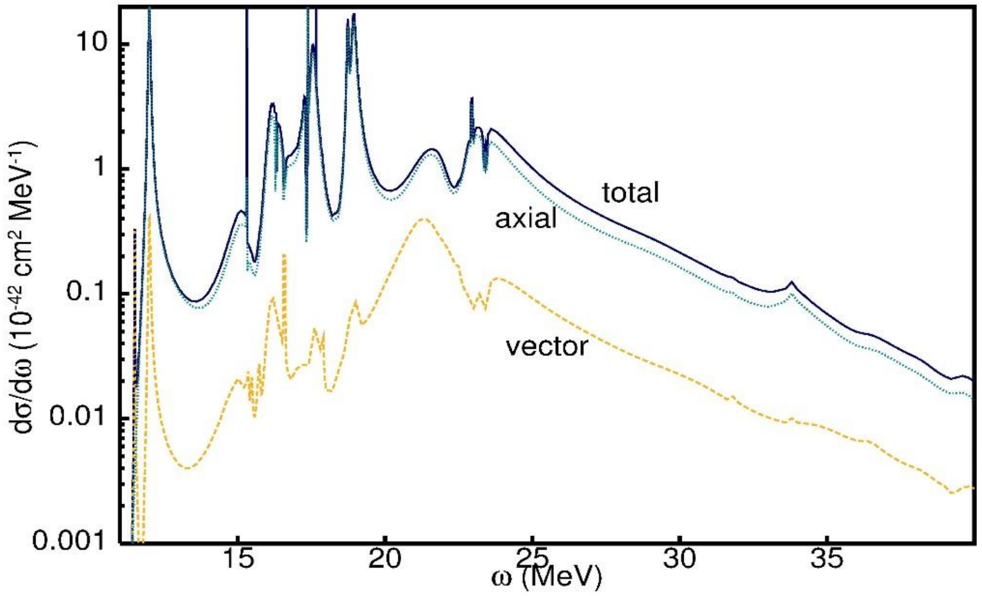


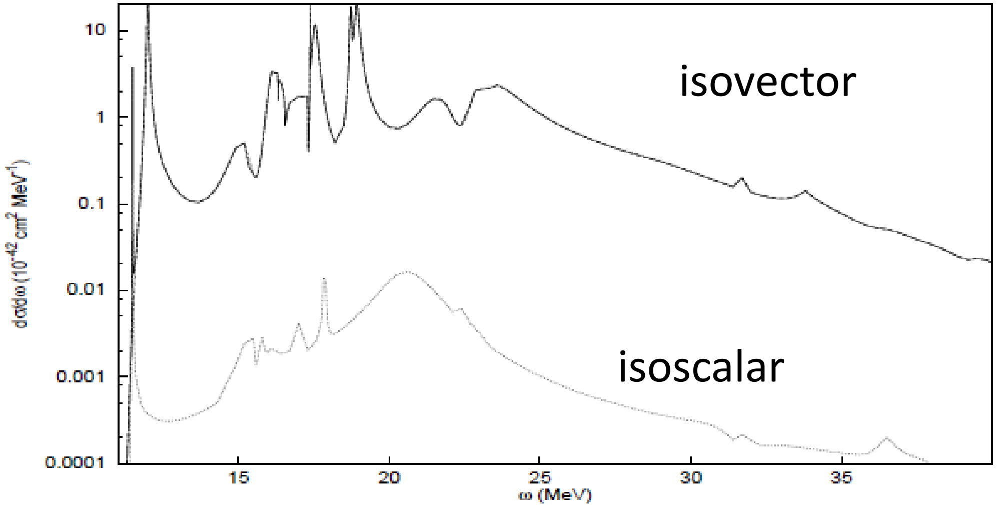
V. Pandey, N.J., J. Ryckebusch, T. Van Cuyck , M. Martini, Phys. Rev. C 92, 024606 (2015)

Inelastic neutrino cross sections

^{16}O, 50 MeV, NC

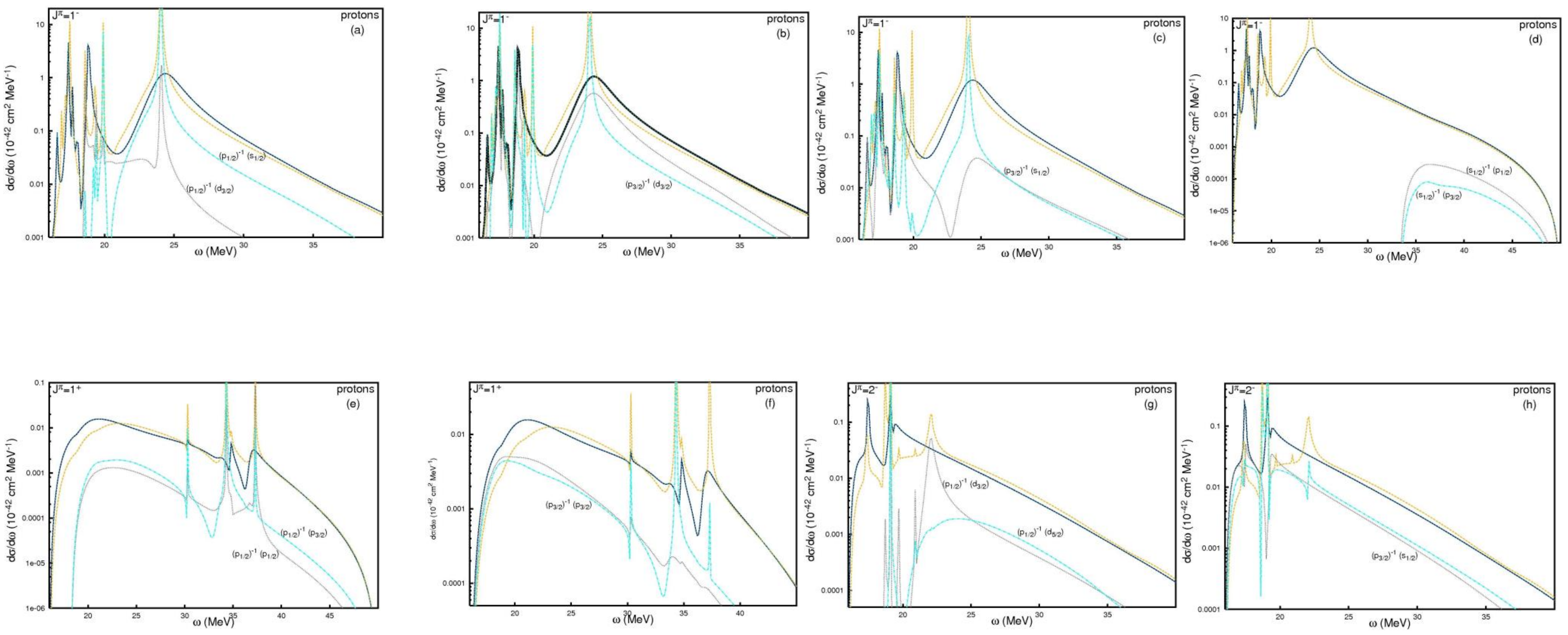




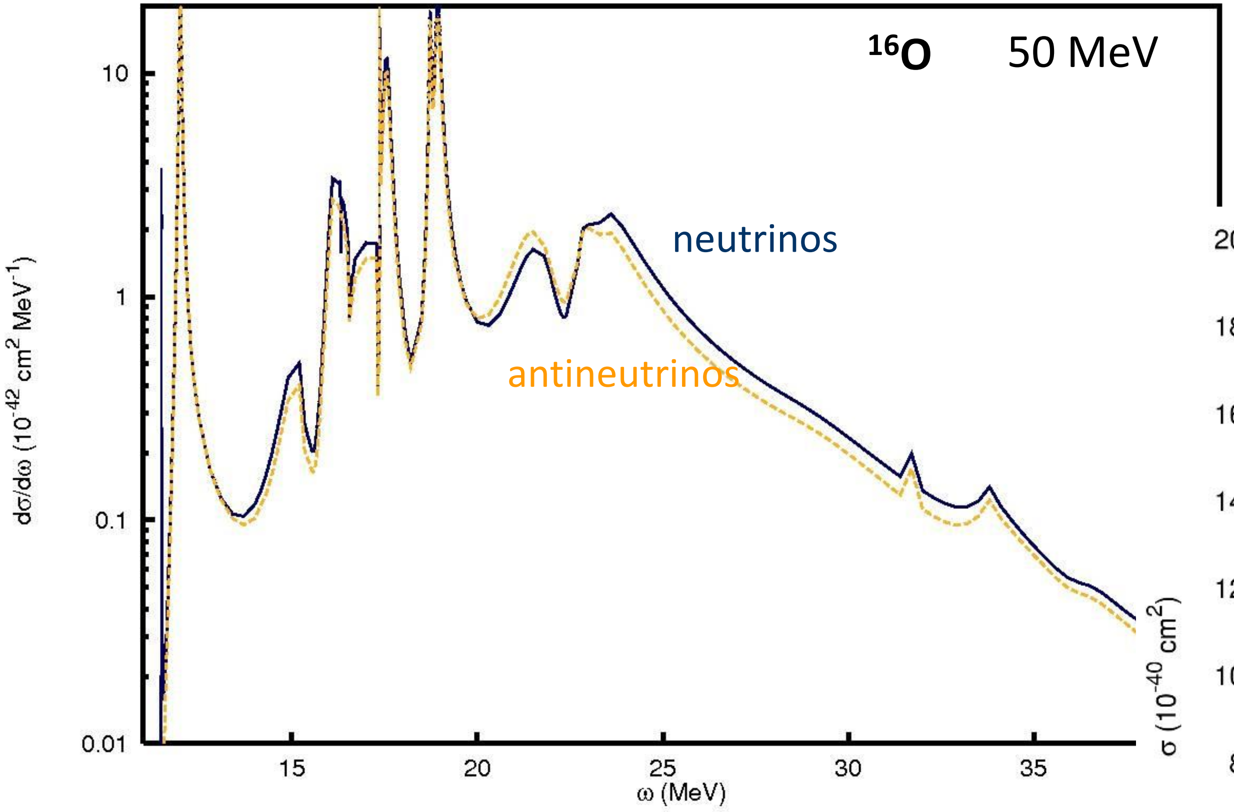


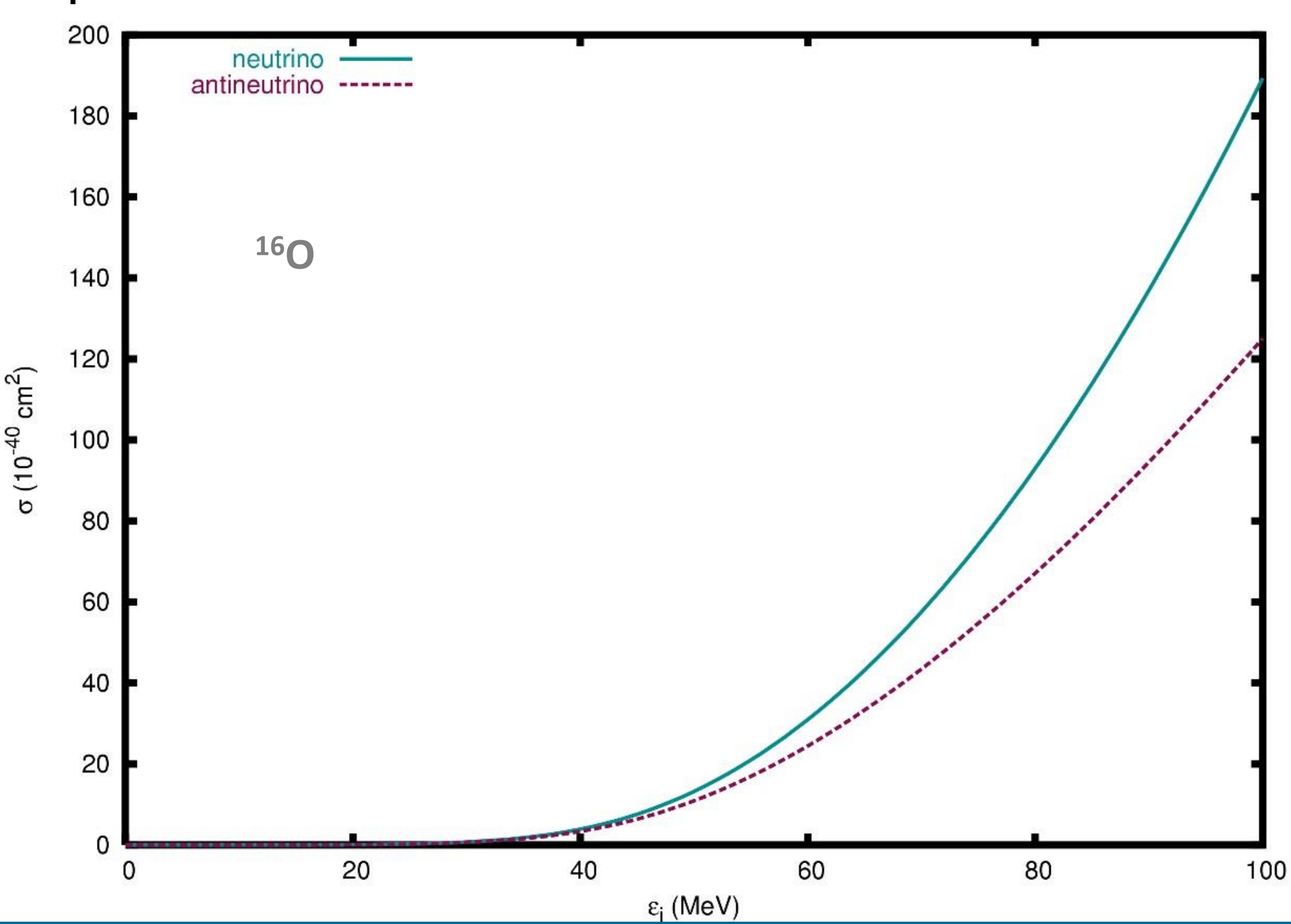
Inelastic neutrino cross sections

Contribution of different single-particle channels in ^{12}C

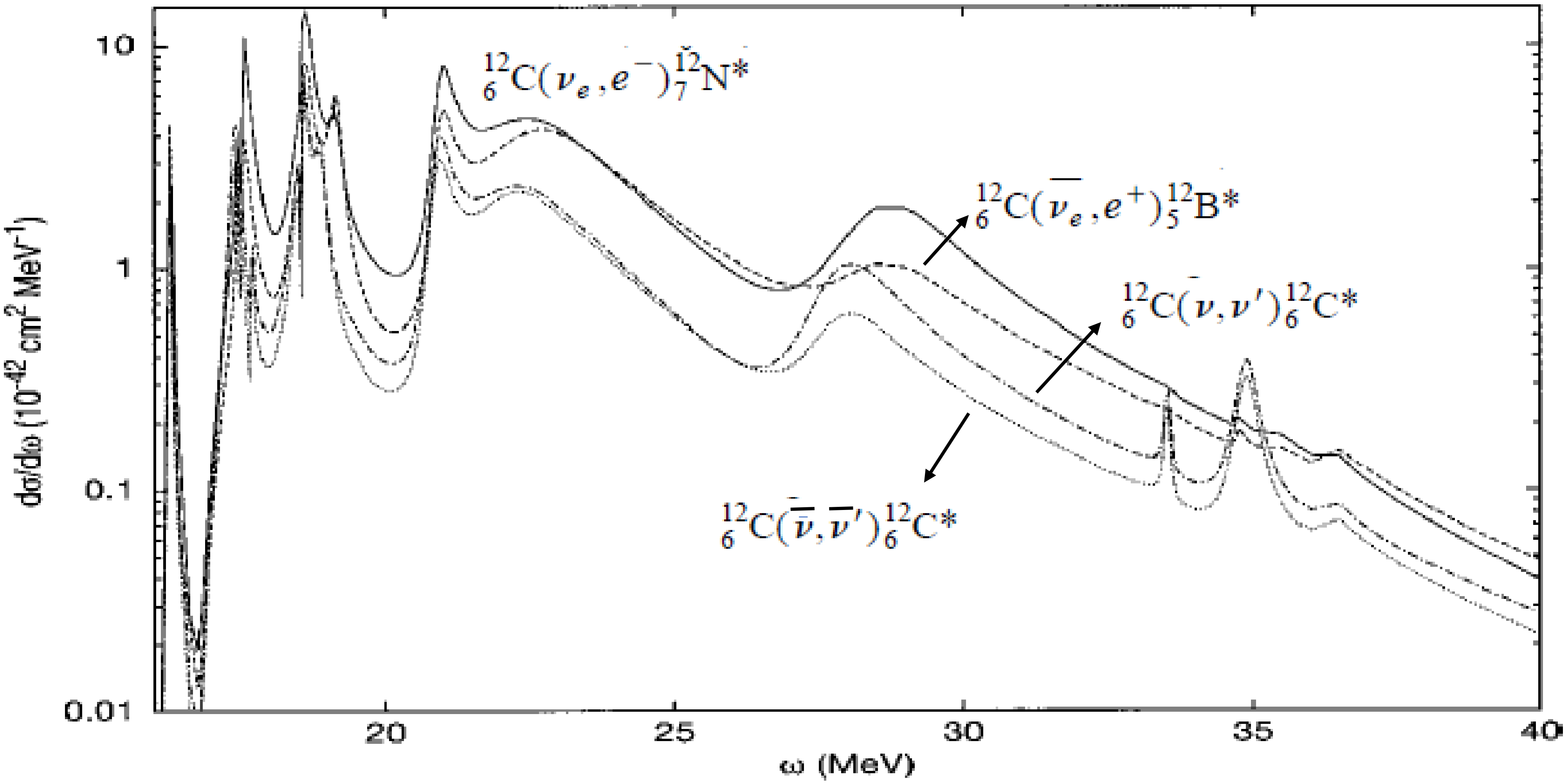


Inelastic neutrino cross sections



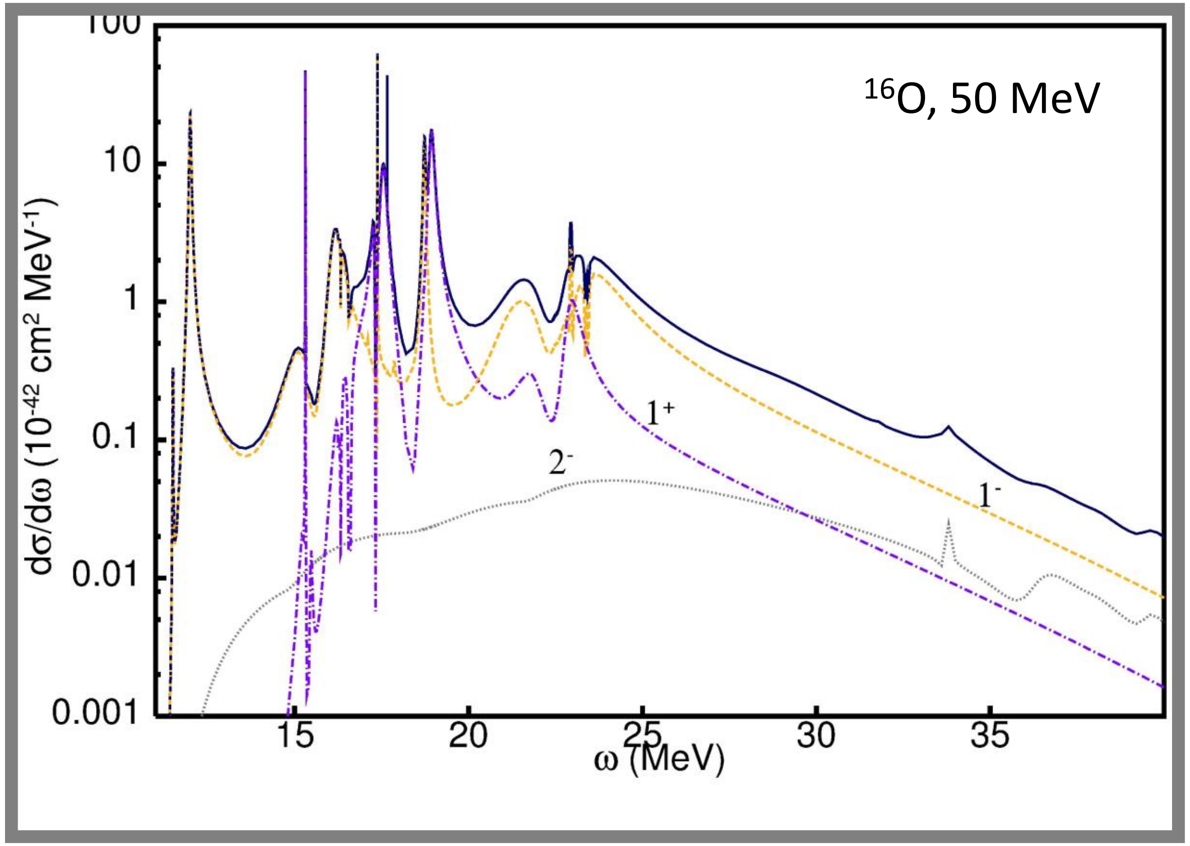


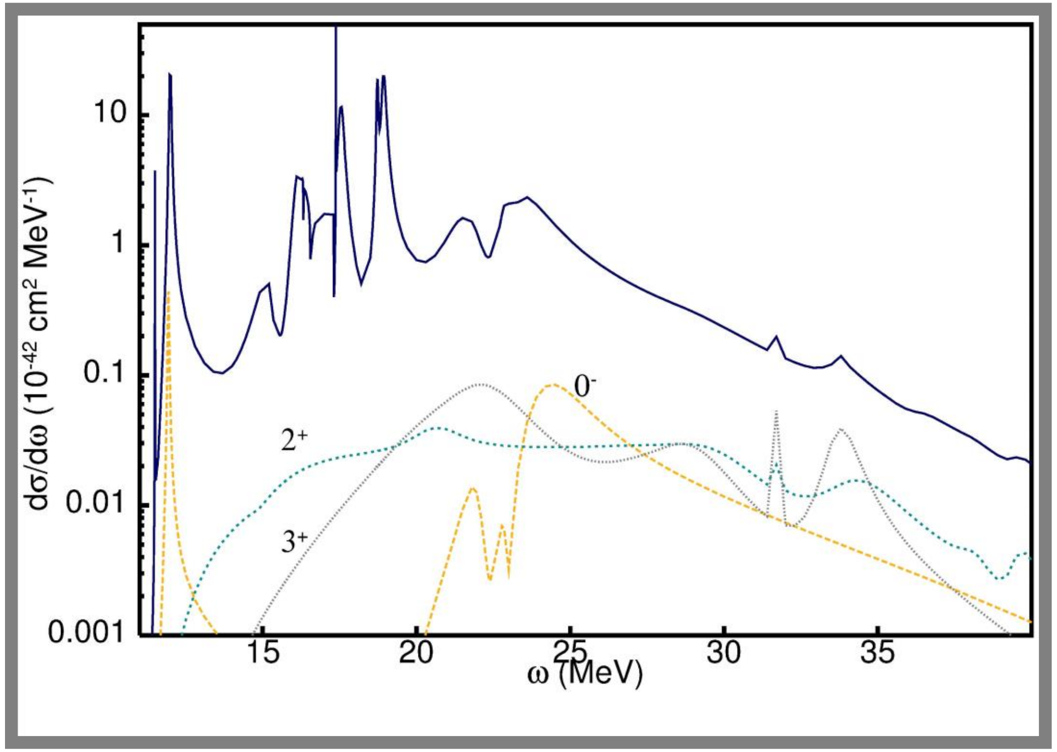
Inelastic neutrino cross sections



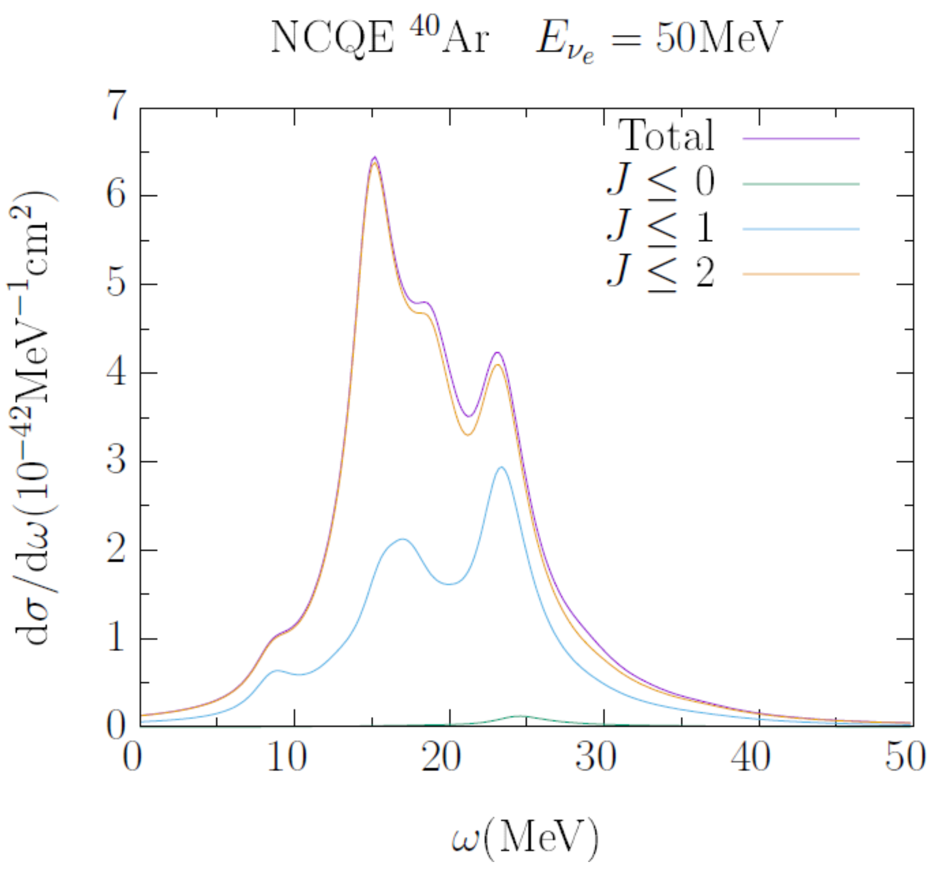
Multipole decomposition

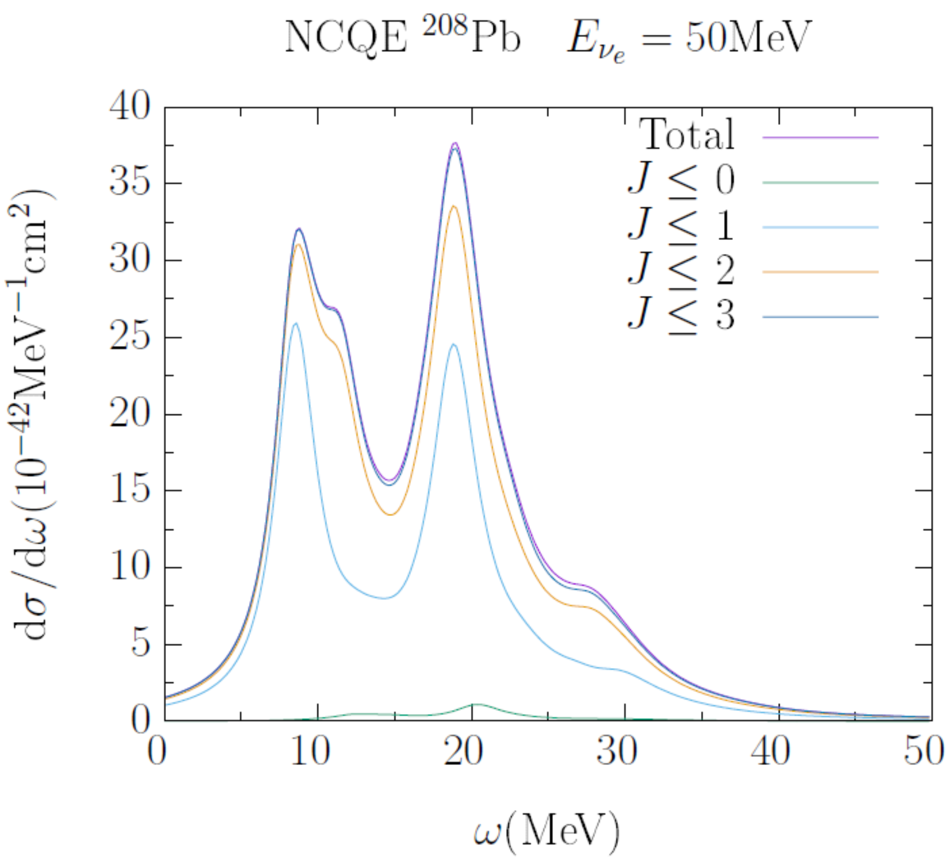
Higher order multipoles important :



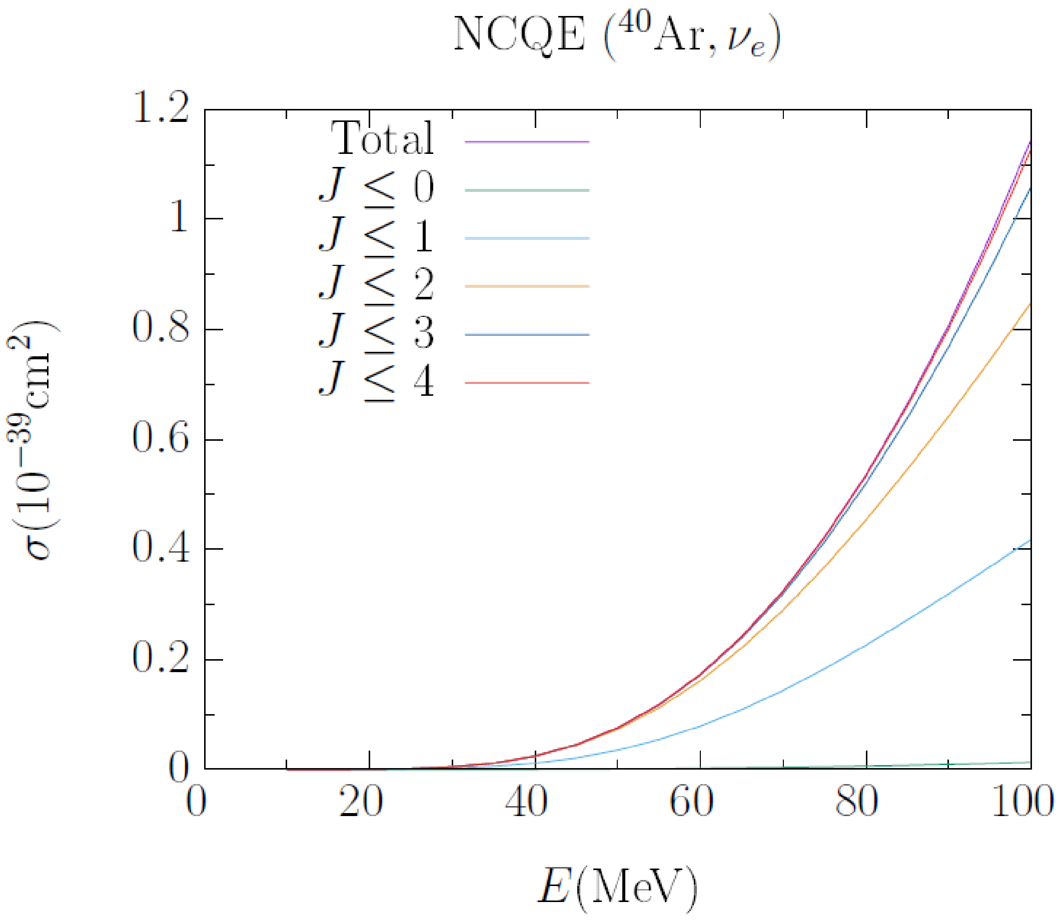


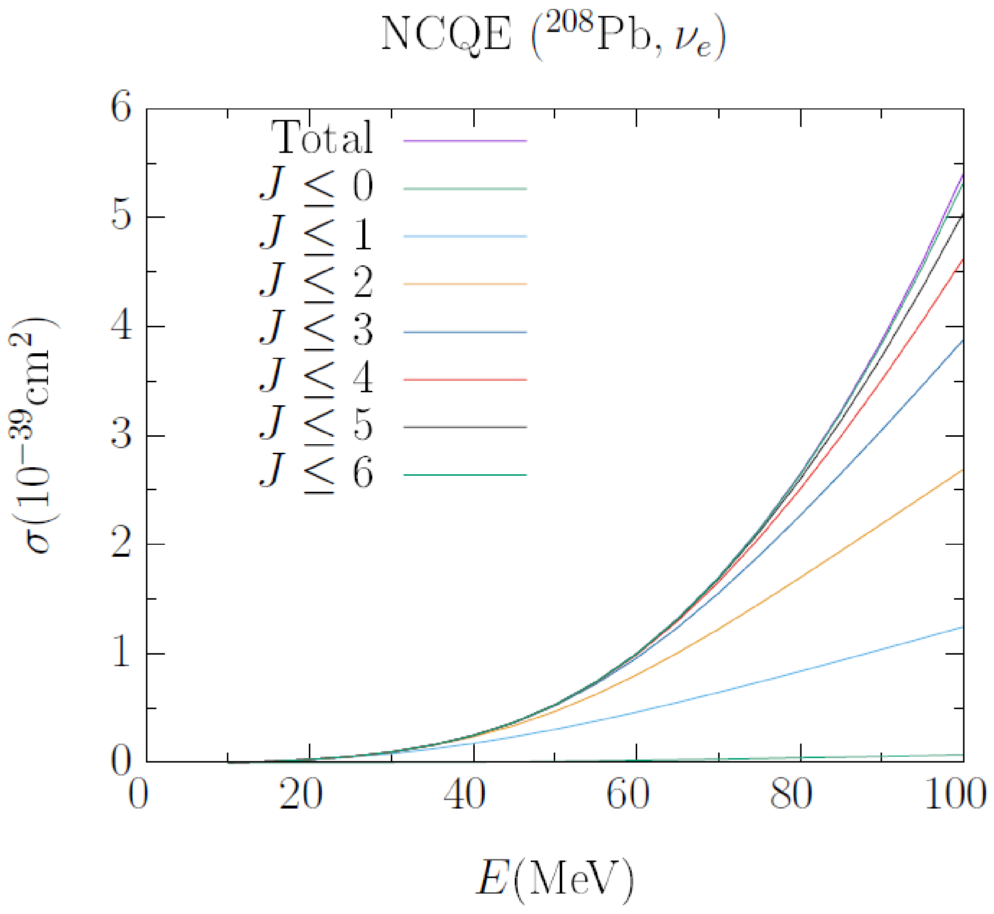
Multipole decomposition



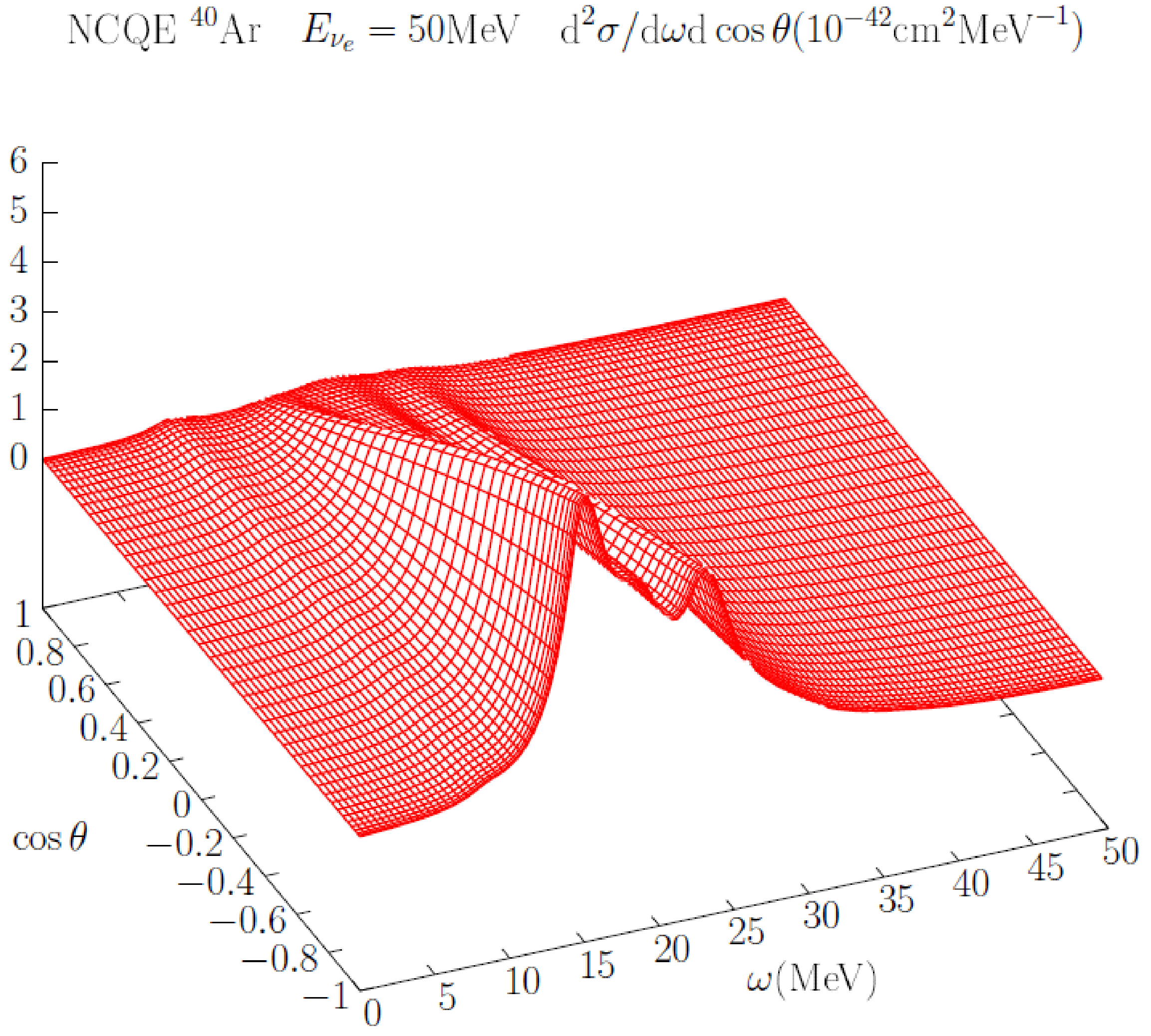


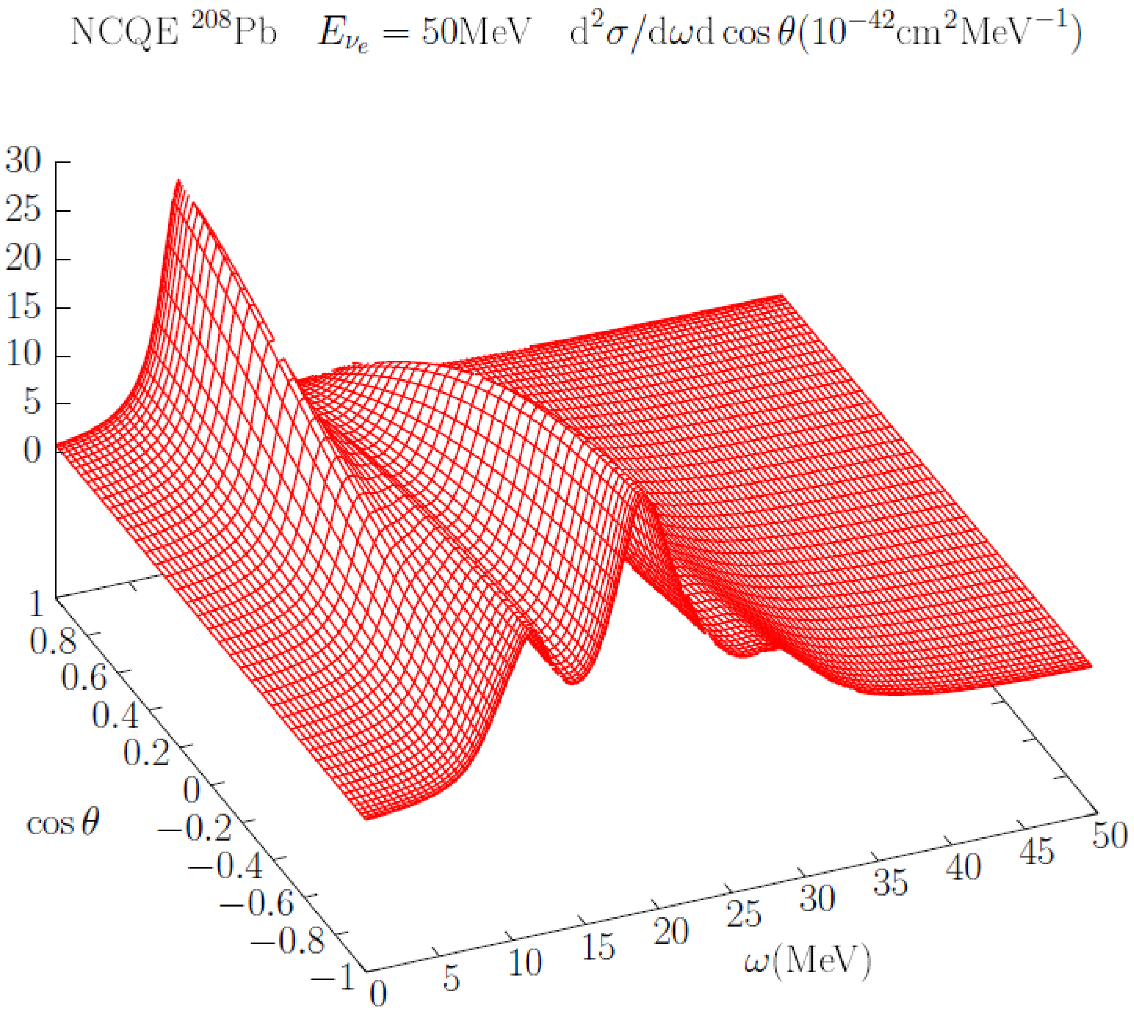
Multipole decomposition



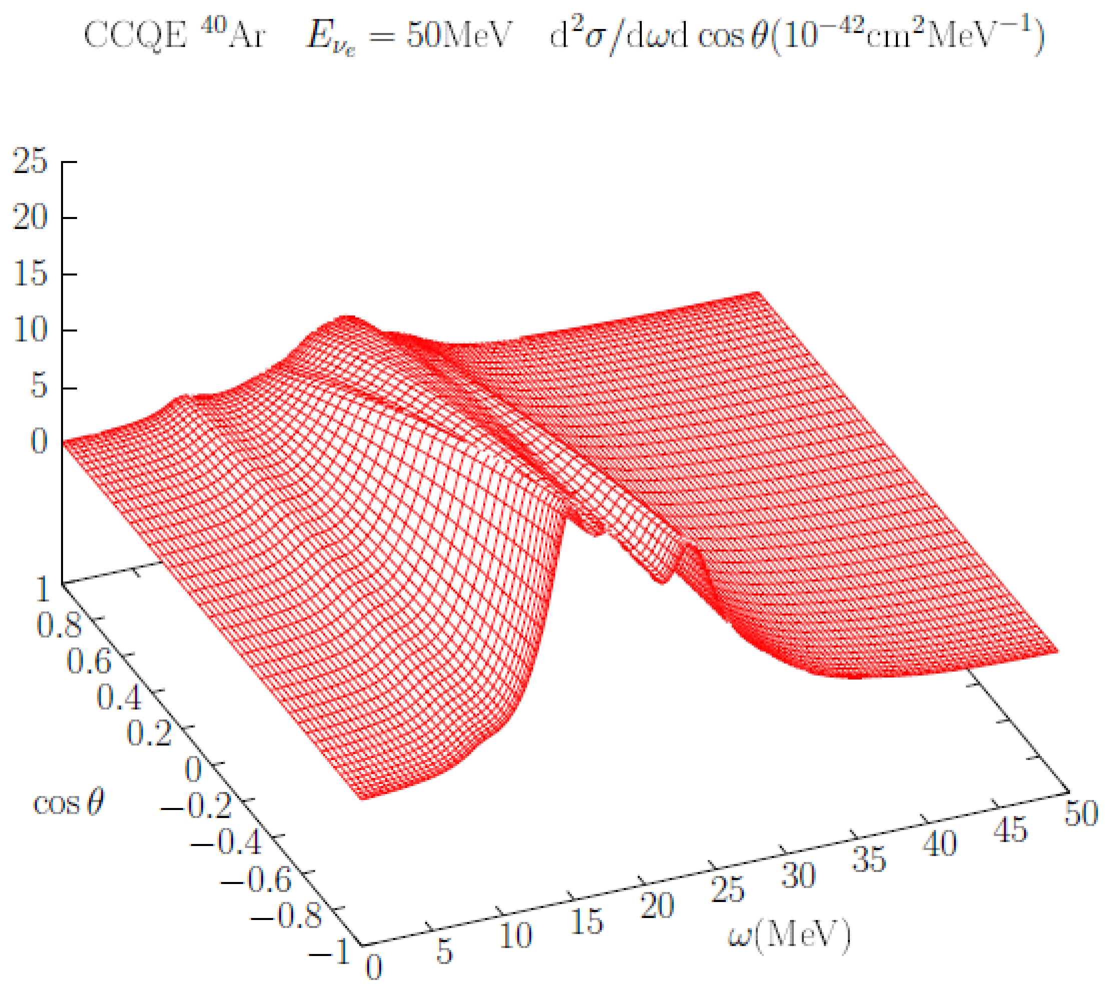


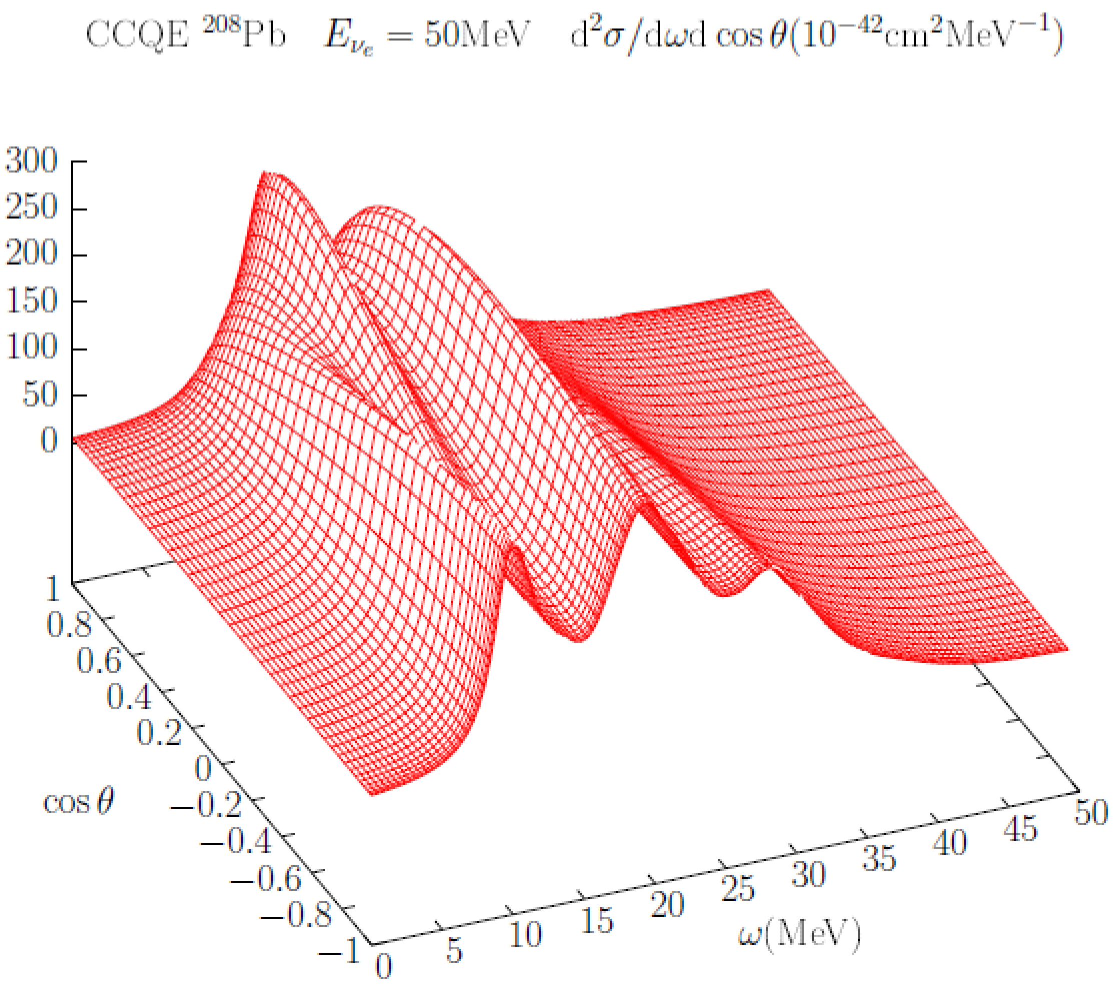
Angular dependence



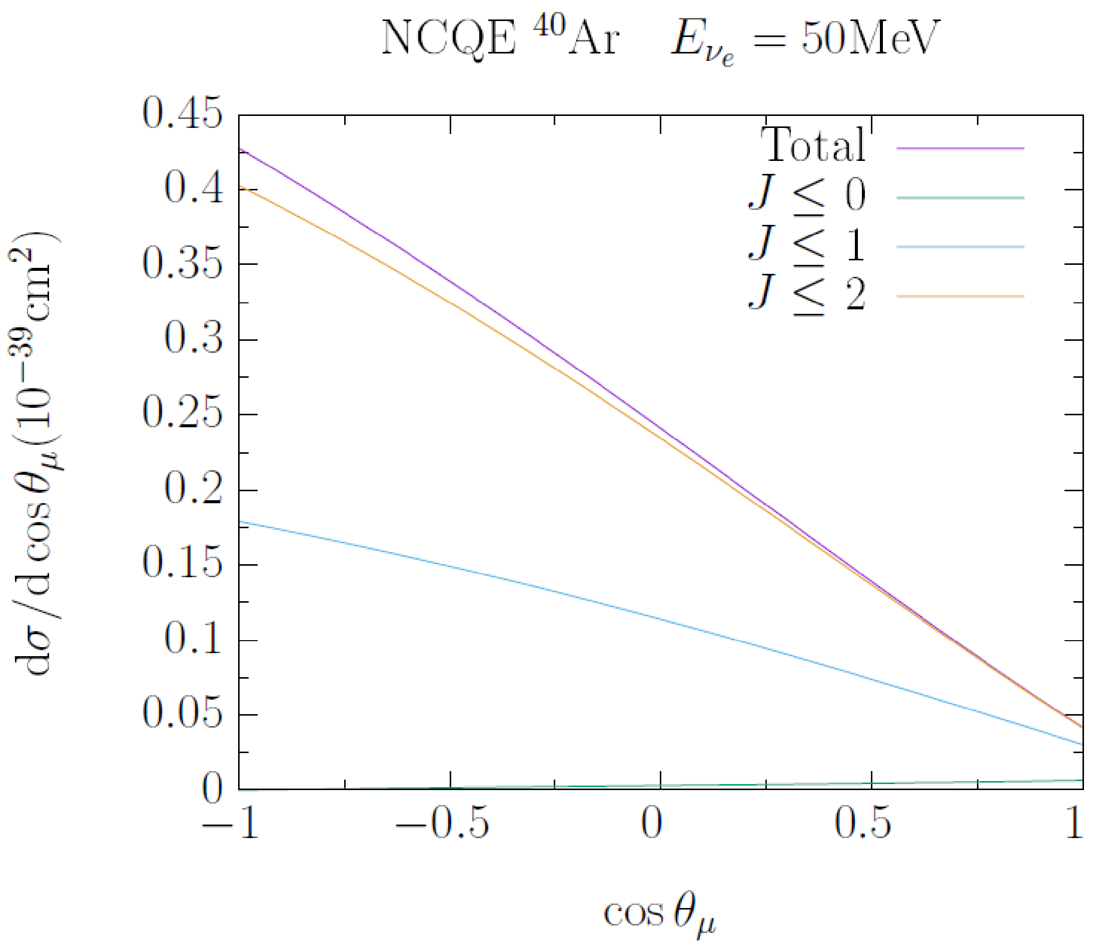


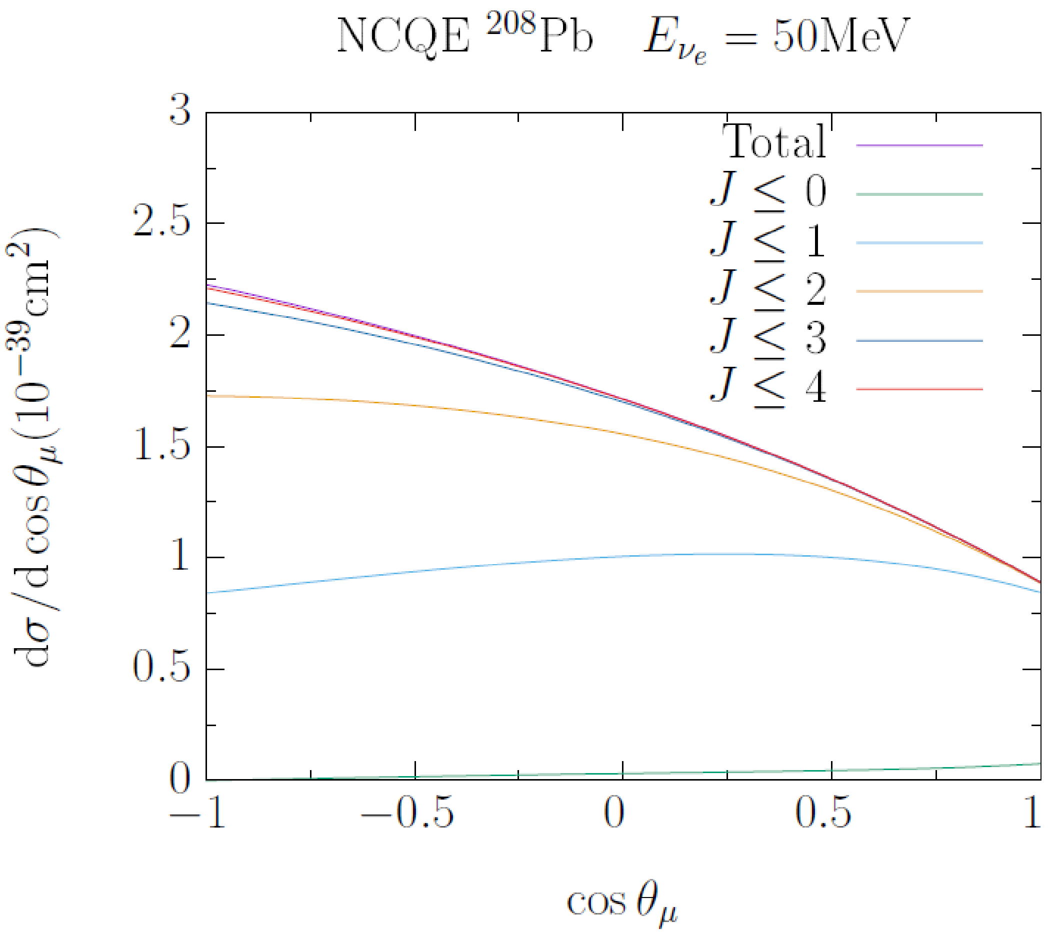
Angular dependence





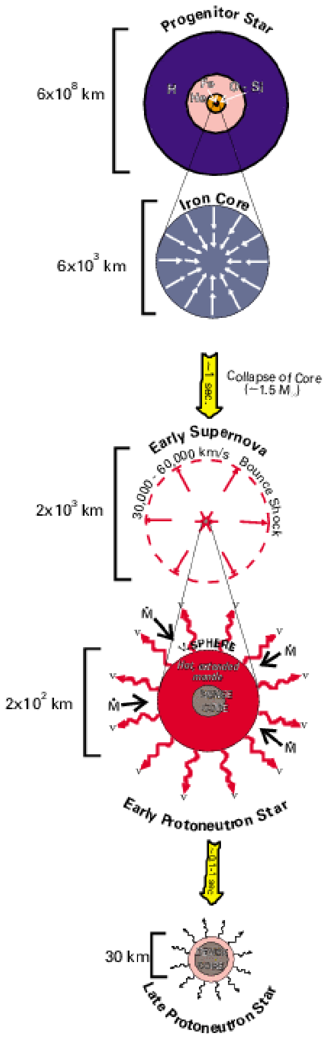
Angular dependence



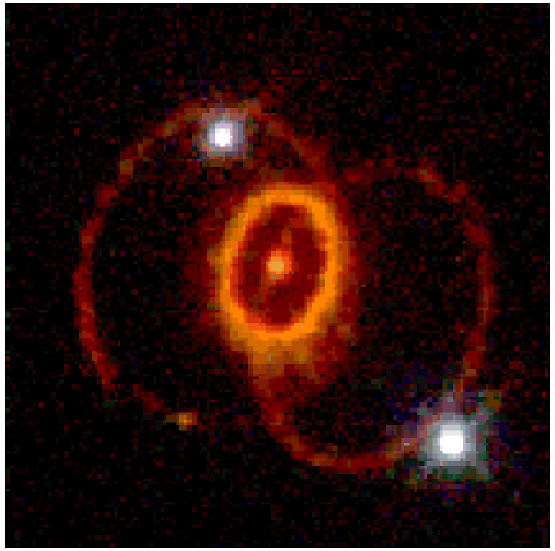


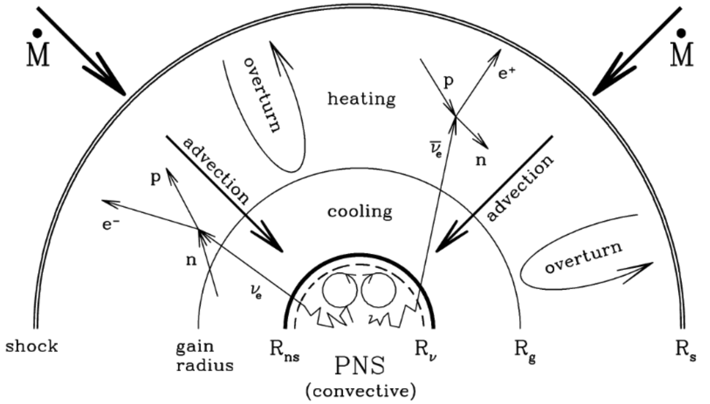
Supernovaneutrinos

Core-collapse supernova



- weak interactions are important
- neutrinos are produced in the neutronization processes characterizing the gravitational collapse
- neutrinos are responsible for the cooling of the proto-neutron star
- neutrinonucleosynthesis
- energy deposition by neutrinos might reheat the stalled shock wave and cause a delayed explosion
- terrestrial detection of supernova neutrinos



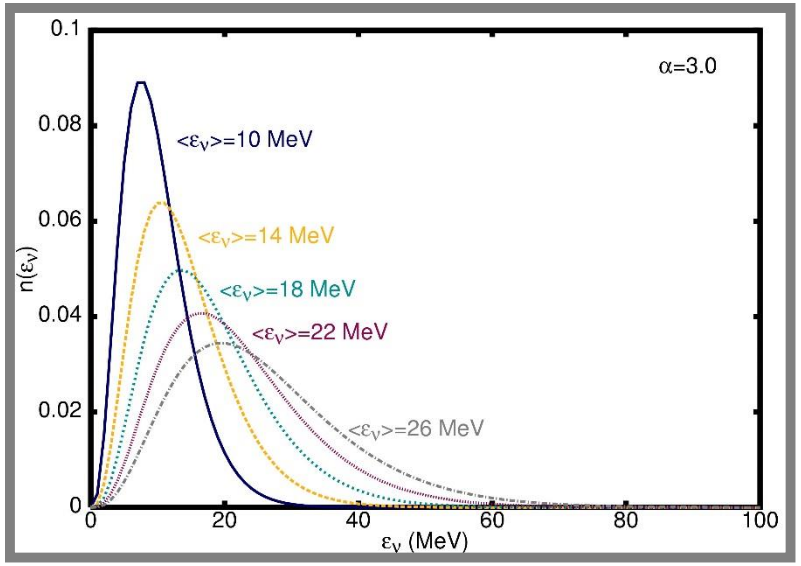


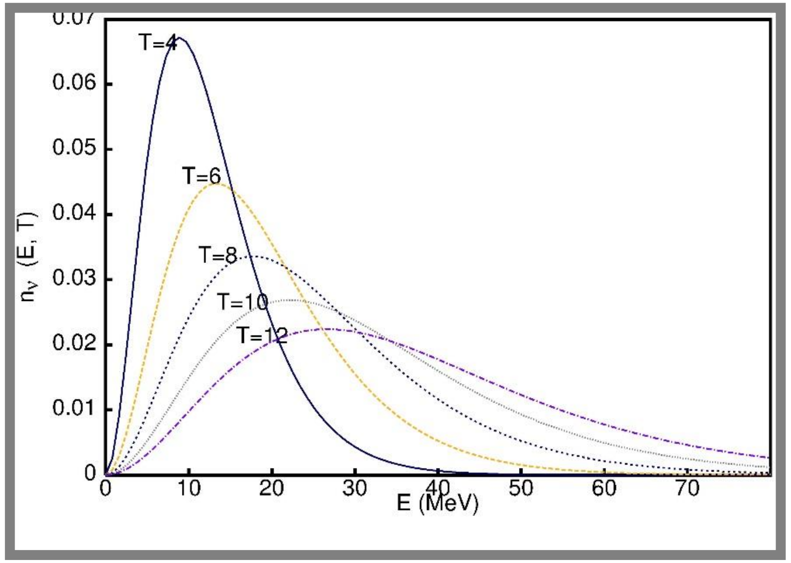
H.-T. Janka astro-ph/0008432

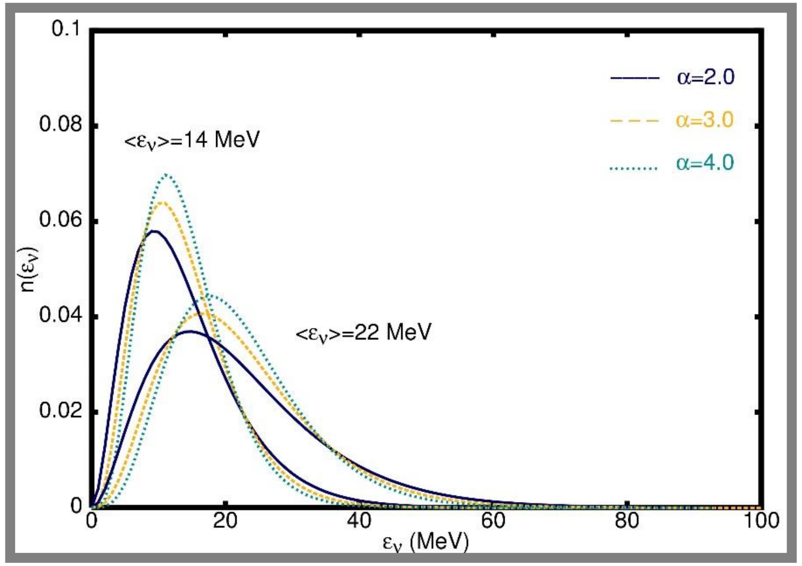
Supernovaneutrino : Energy spectra

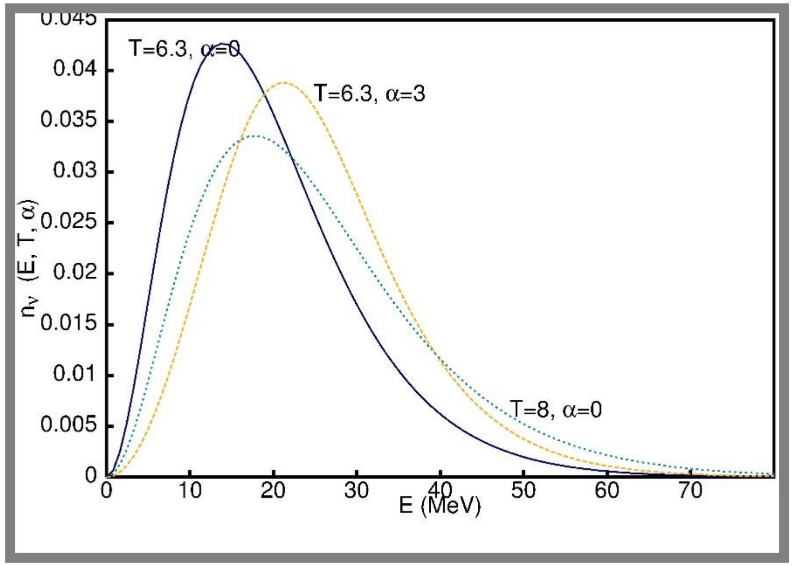
$$n_{SN[\langle\varepsilon\rangle,\alpha]}(\varepsilon) = \left(\frac{\varepsilon}{\langle\varepsilon\rangle}\right)^{\alpha} e^{-(\alpha+1)\frac{\varepsilon}{\langle\varepsilon\rangle}}$$

Fermi-Dirac spectra



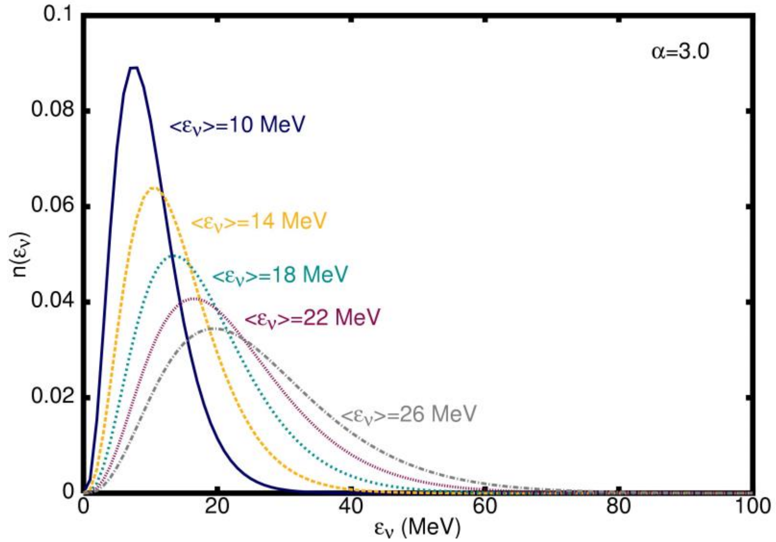


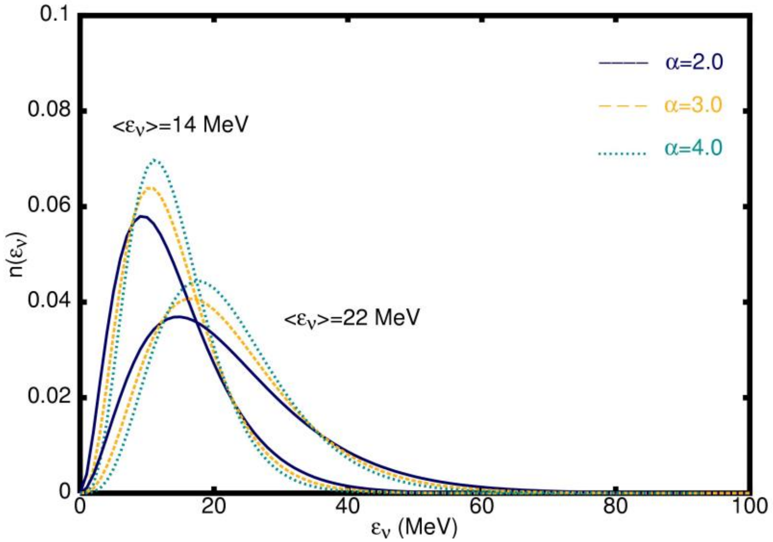


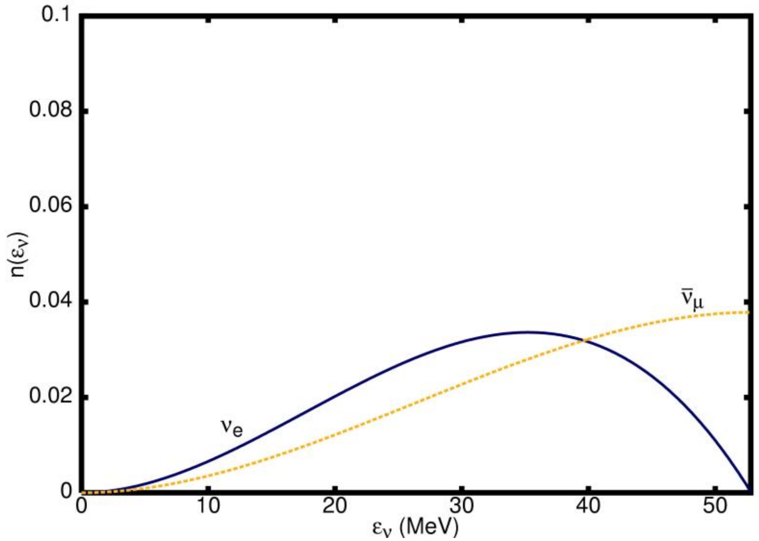


Supernovaneutrino : Energy spectra

Supernova neutrino spectra :







Michel spectra :

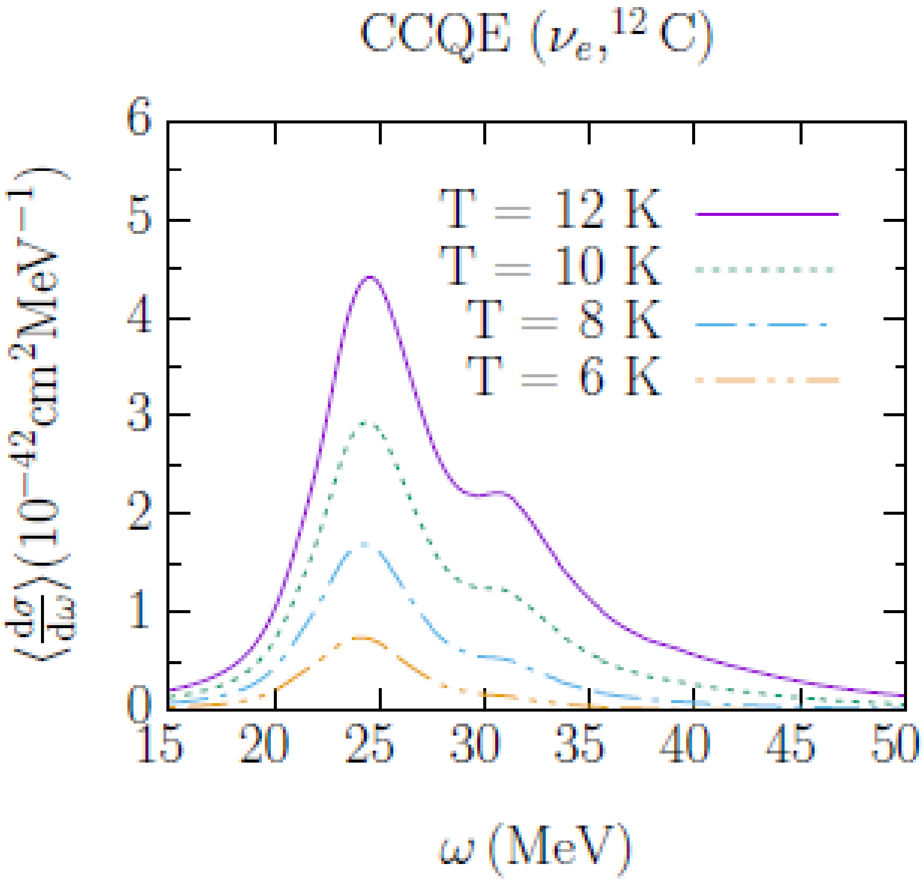
$$n(E_{\nu_e}) = \frac{96E_{\nu_e}^2}{m_\mu^4}(m_\mu - 2E_{\nu_e})$$

$$n(E_{\bar{\nu}_\mu}) = \frac{32E_{\bar{\nu}_\mu}^2}{m_\mu^4}\left(\frac{3}{2}m_\mu - 2E_{\bar{\nu}_\mu}\right)$$

Supernovaneutrino cross sections

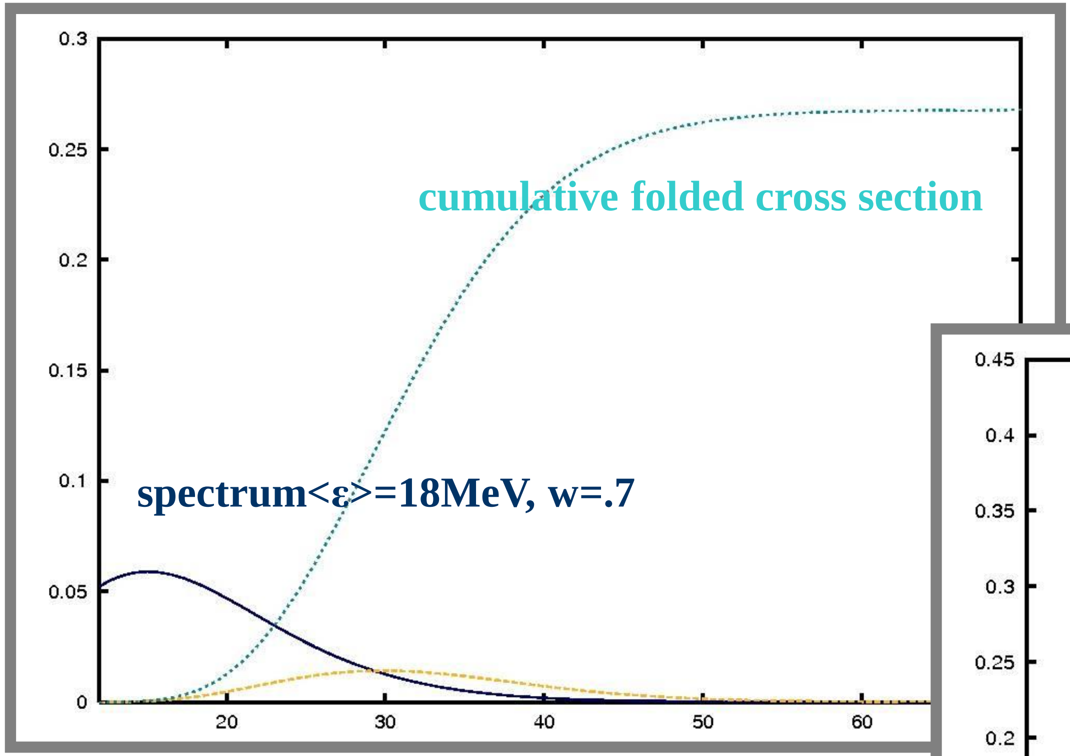
Folded cross sections supernova neutrino spectra :

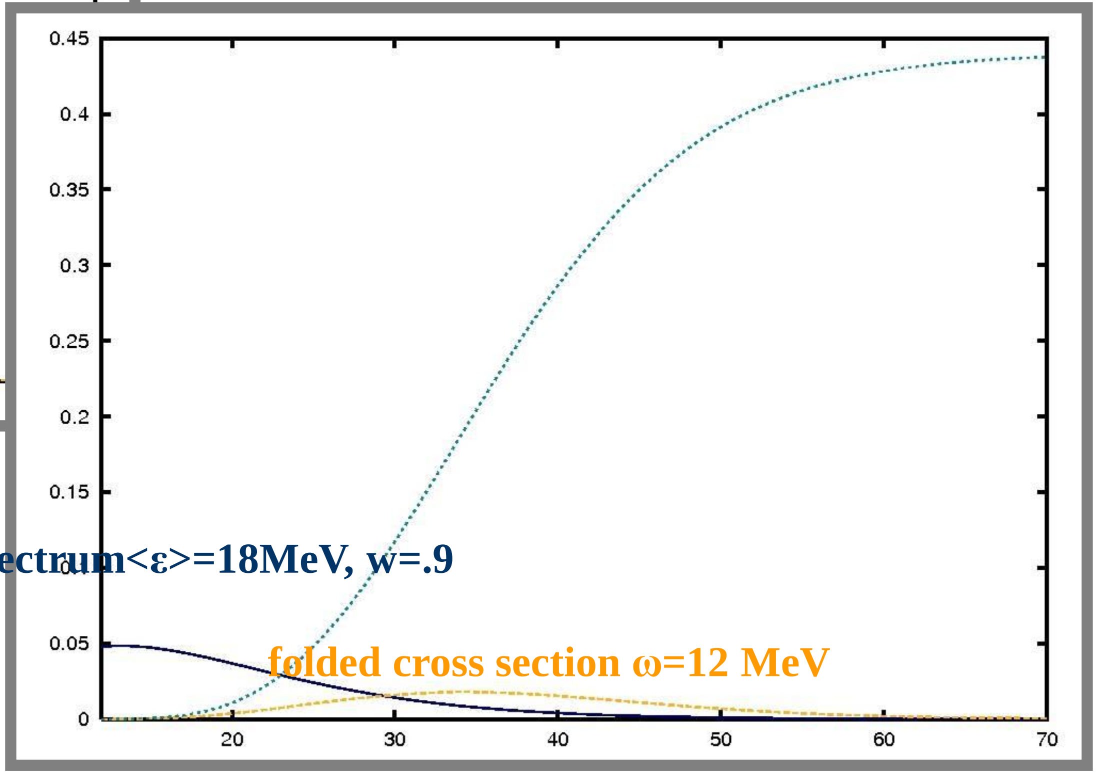
→strong dependence on average energy of the spectrum



Supernovaneutrino cross sections

Cumulative folded cross sections:



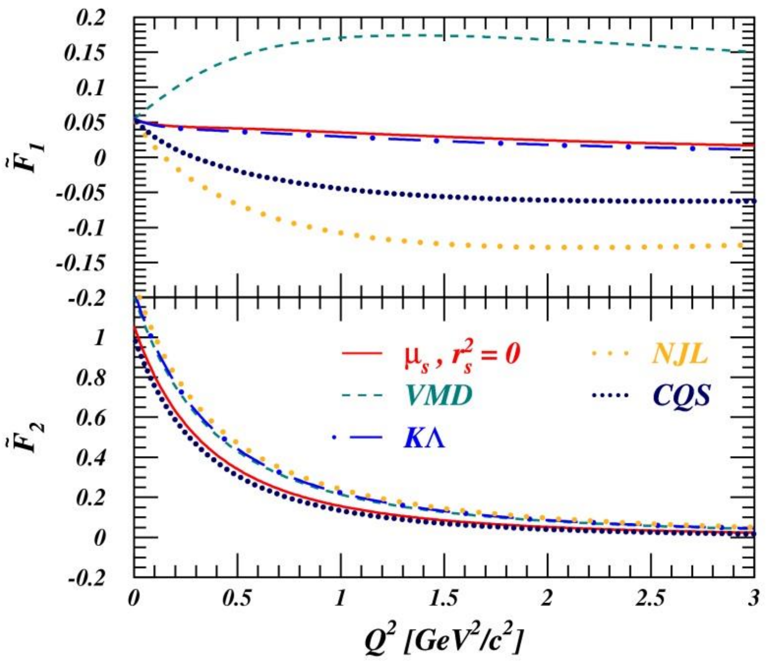


Strangeness

Axial form factor :

$$G_A(Q^2) = -\frac{(\tau_3 g_A - g_A^s)}{2} G(Q^2), \qquad g_A = 1.262$$

$$G(Q^2) = (1 + Q^2/M^2)^{-2}, \qquad M = 1.032$$



Model	$\mu_s(\mu_N)$	$r_s^2(\mathrm{fm}^2)$
VMD	-0.31	0.16
$K\Lambda$	-0.35	-0.007
NJL	-0.45	-0.17
CQS (K)	0.115	-0.095

Weak vector form factors :

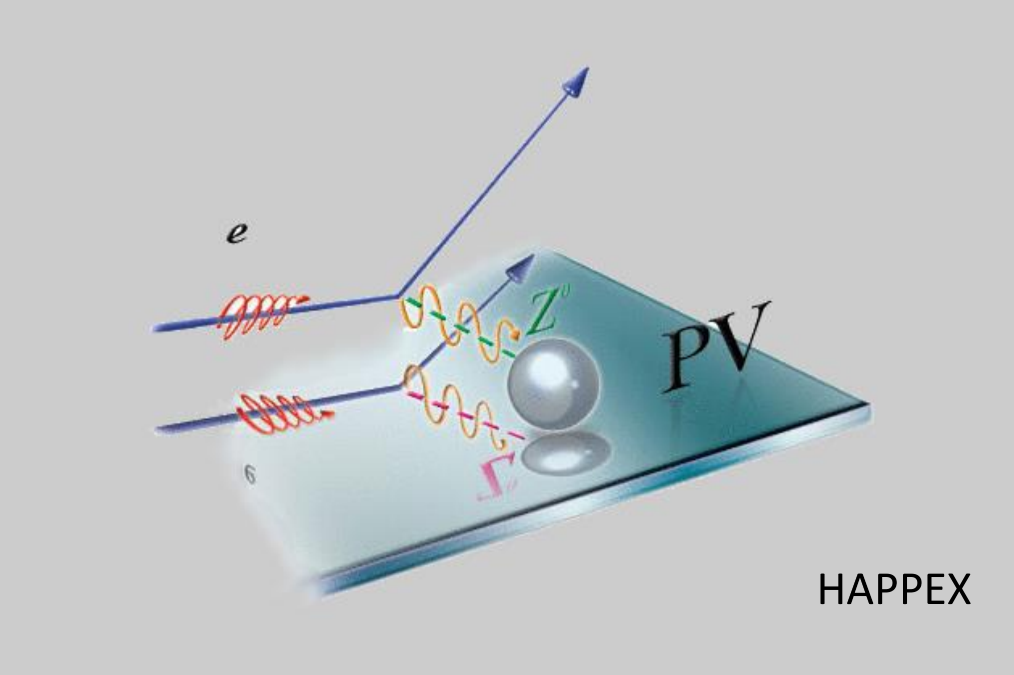
$$F_1^s = \frac{1}{6} \frac{-r_s^2 Q^2}{(1 + Q^2/M_1^2)^2}, \qquad M_1 = 1.3$$

$$F_2^s = \frac{\mu_s}{(1 + Q^2/M_2^2)^2}, \qquad M_2 = 1.26$$

Parity violating electron scattering

Using polarized electrons, one gets access to parity violating electron scattering (HAPPEX, G0, SAMPLE, A4)

- Axial-vector interference terms
- Information about axial vector form factor
- Information about strangeness in the nucleon in the vector as well as axial sector
- Larger cross sections
- Prone to radiative corrections



Parity violating asymmetry :

$$A^{PV} = \frac{d\sigma_R - d\sigma_L}{d\sigma_R + d\sigma_L}$$

Strangeness

Traditionally :

- strangeness contribution to the ***weak vector formfactors*** : Parity Violating Electron Scattering (Sample, HAPPEX, G0, …)
- strangeness contribution to the ***axial current*** : neutrino scattering

-vector current contributions are suppressed

-no radiative corrections

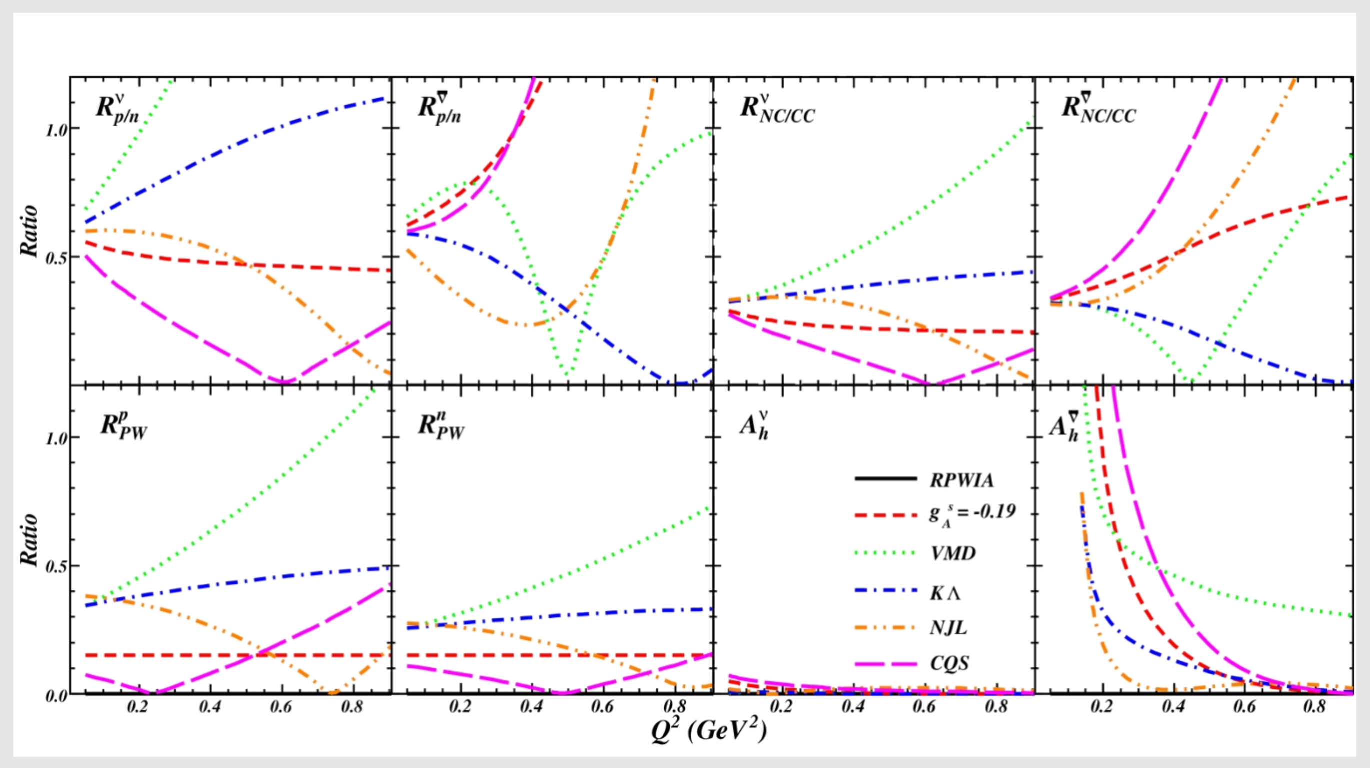
Strangeness

Traditionally :

- strangeness contribution to the ***weak vector formfactors*** : Parity Violating Electron Scattering (Sample, HAPPEX, G0, …)

Correlated !

- strangeness contribution to the ***axial current*** : neutrino scattering



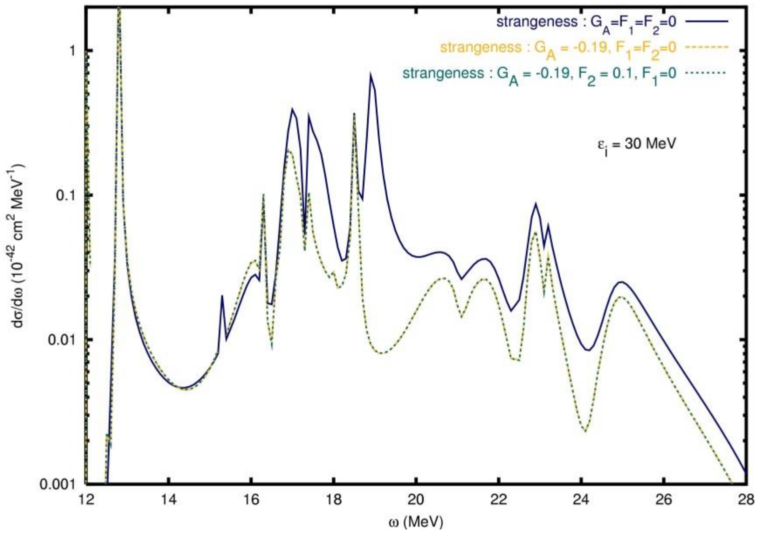
-vector current contributions are suppressed

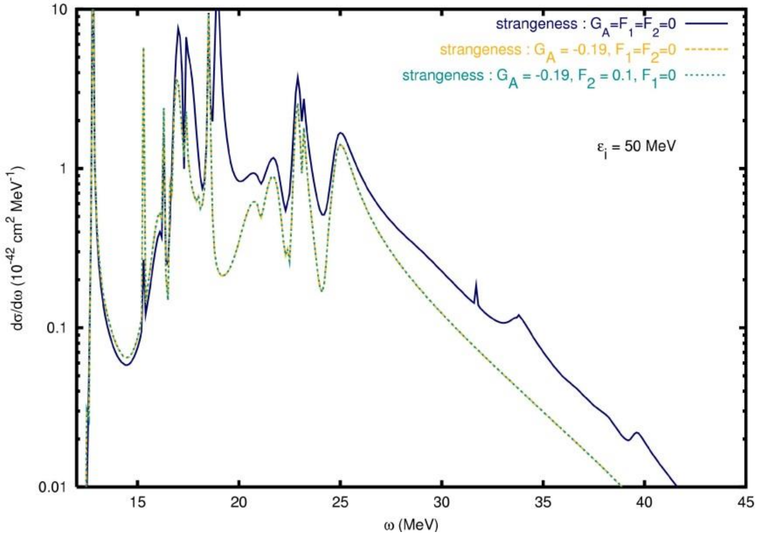
-no radiative corrections

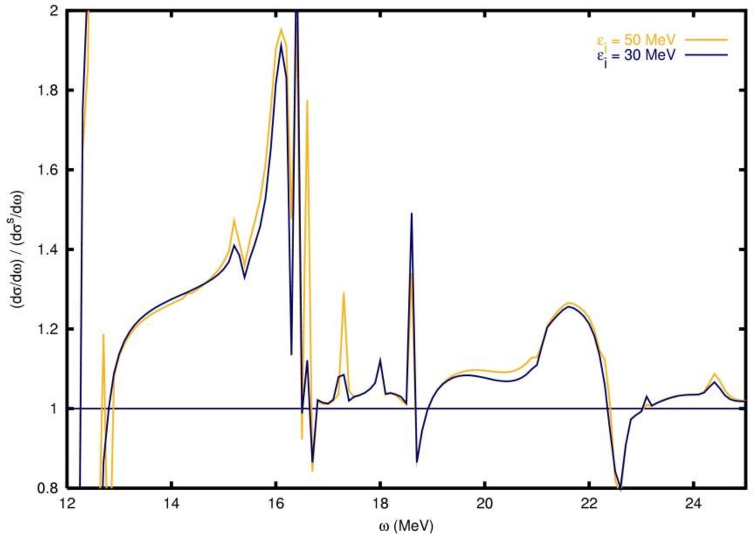
N.J., P. Vancraeyveld, P. Lava, J. Ryckebusch, PRC76, 055501 (2007).

Influence of strangeness on neutrino cross sections

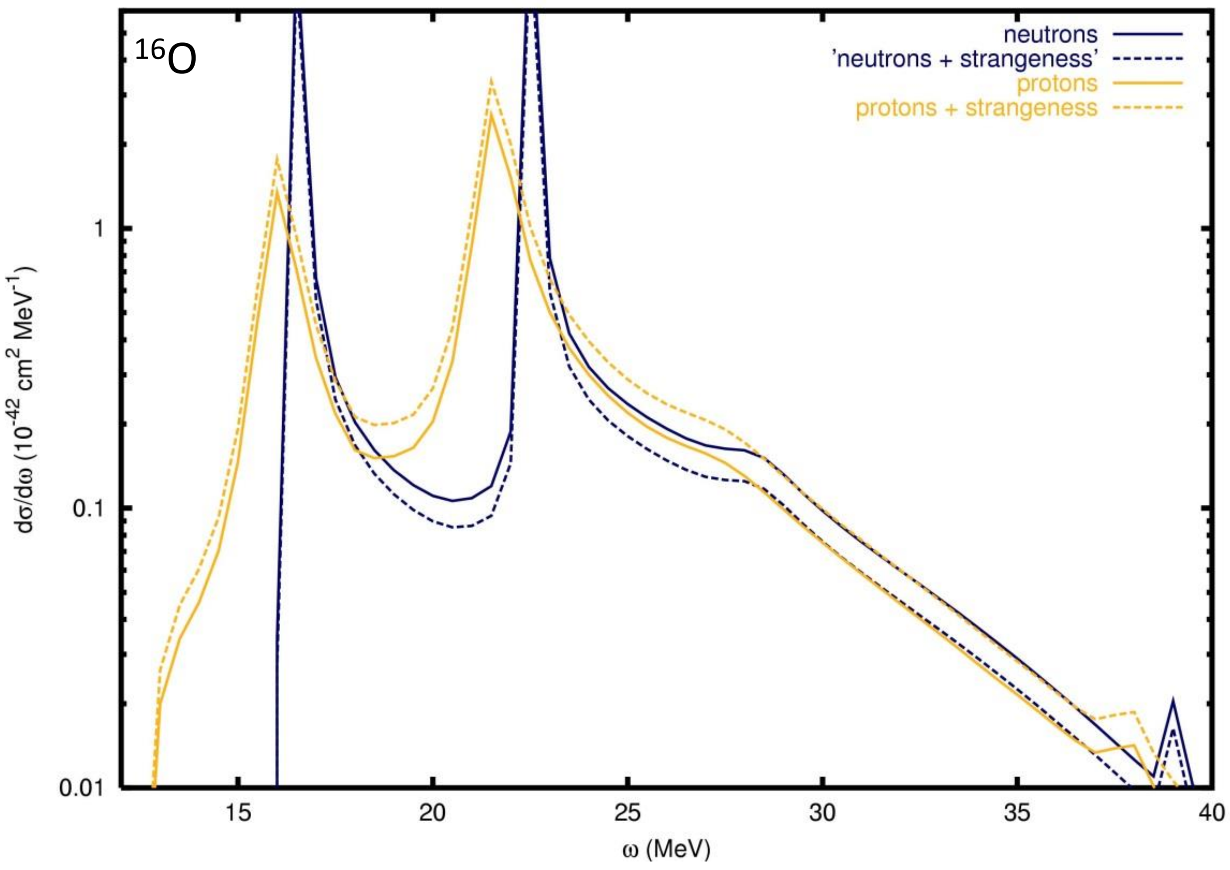
- Generally : net strangeness effect vanishes for isoscalar targets
- close to particle knockout threshold the influence becomes larger due to binding energy differences between protons and neutrons
- differential cross sections differ, energy of reaction products can be very different







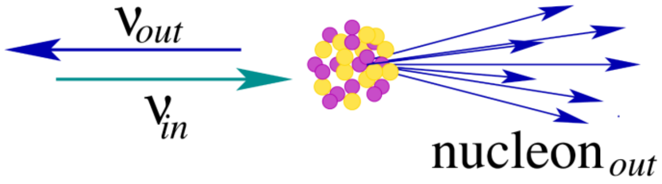
Influence of strangeness on p/n neutrino cross sections



- differences up to 20%
- opposite effect for protons and neutrons

Spin of the outgoing nucleon

Helicity dependence of the cross section:



$$
\begin{aligned}
(l_- l_-^* h_+ h_+^* + l_+ l_+^* h_- h_-^*)_{\frac{\nu}{\bar{\nu}}} \\
&= S(h_+ h_+^* + h_- h_-^*) \mp A(h_+ h_+^* - h_- h_-^*) \\
&= (S \mp A) h_+ h_+^* + (S \pm A) h_- h_-^*,
\end{aligned}
$$

For neutrinos :

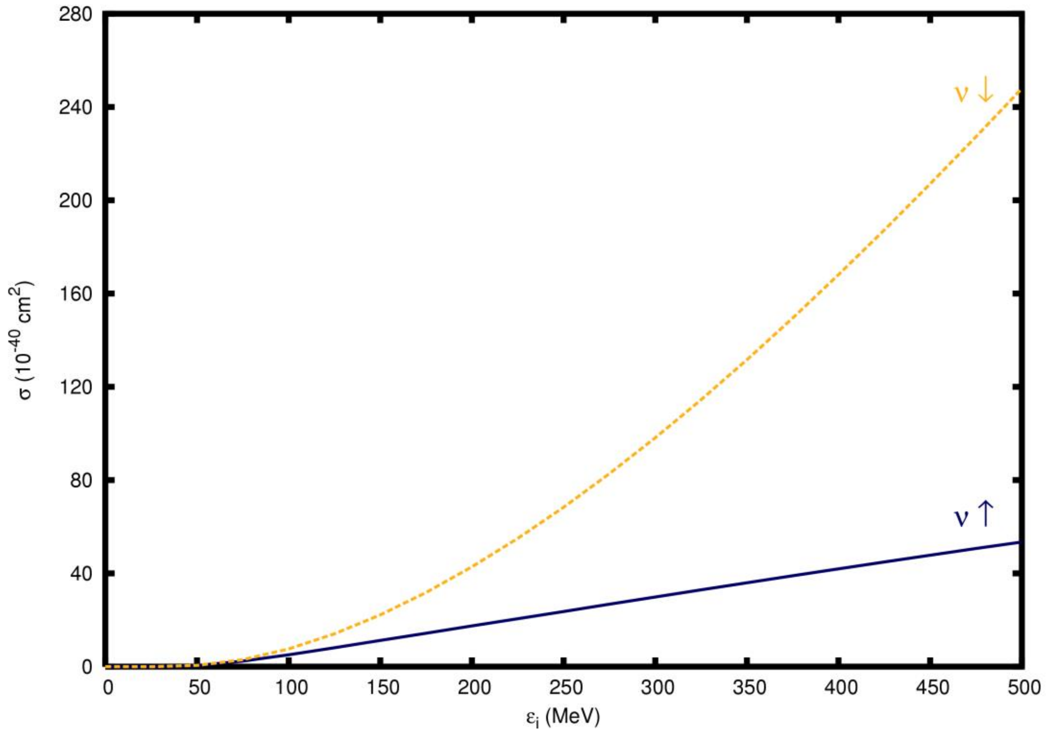
S+hA = S-A, is small

S-hA = S+A, is large

$\Downarrow$

$h_- h_-^*$ dominates

^{16}O



For antineutrinos :

S+hA = S+A, is large

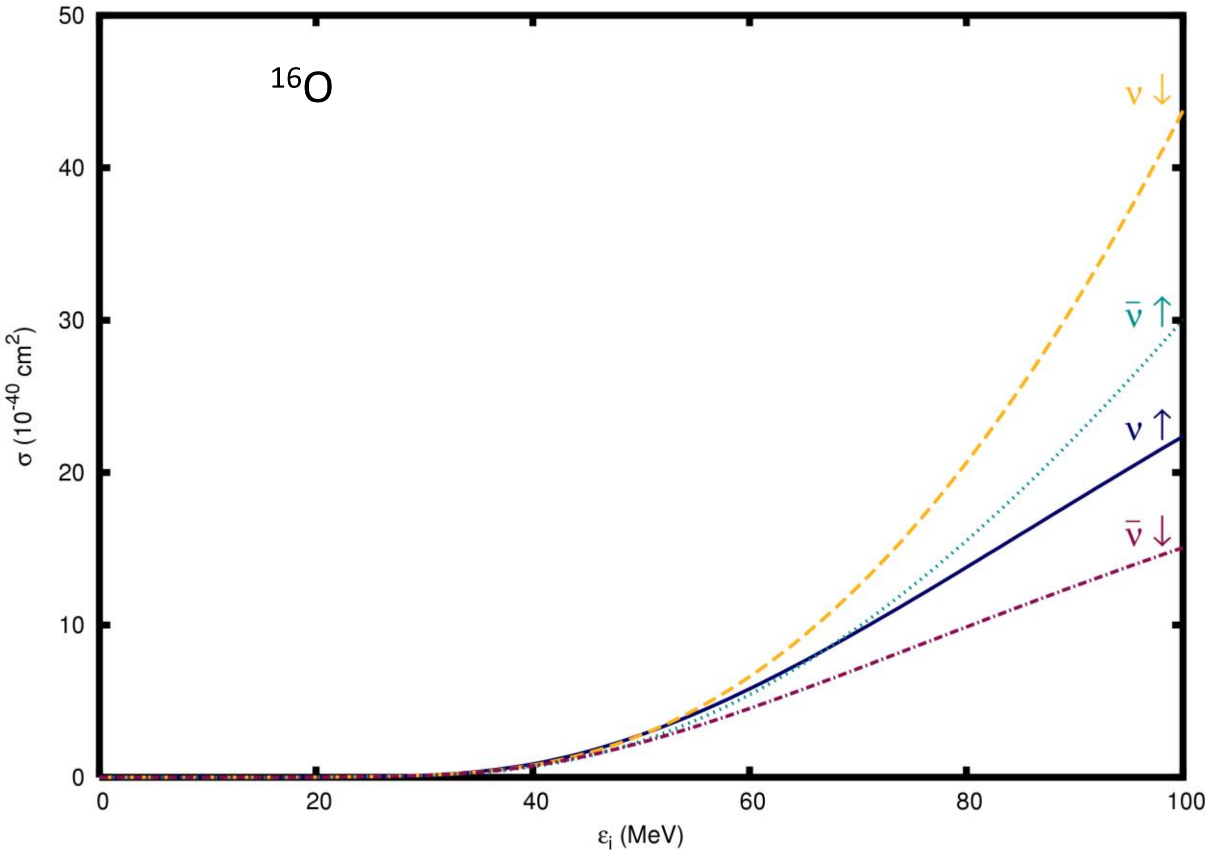
S-hA = S-A, is small

$\Downarrow$

$h_+ h_+^*$ dominates

Spin of the outgoing nucleon

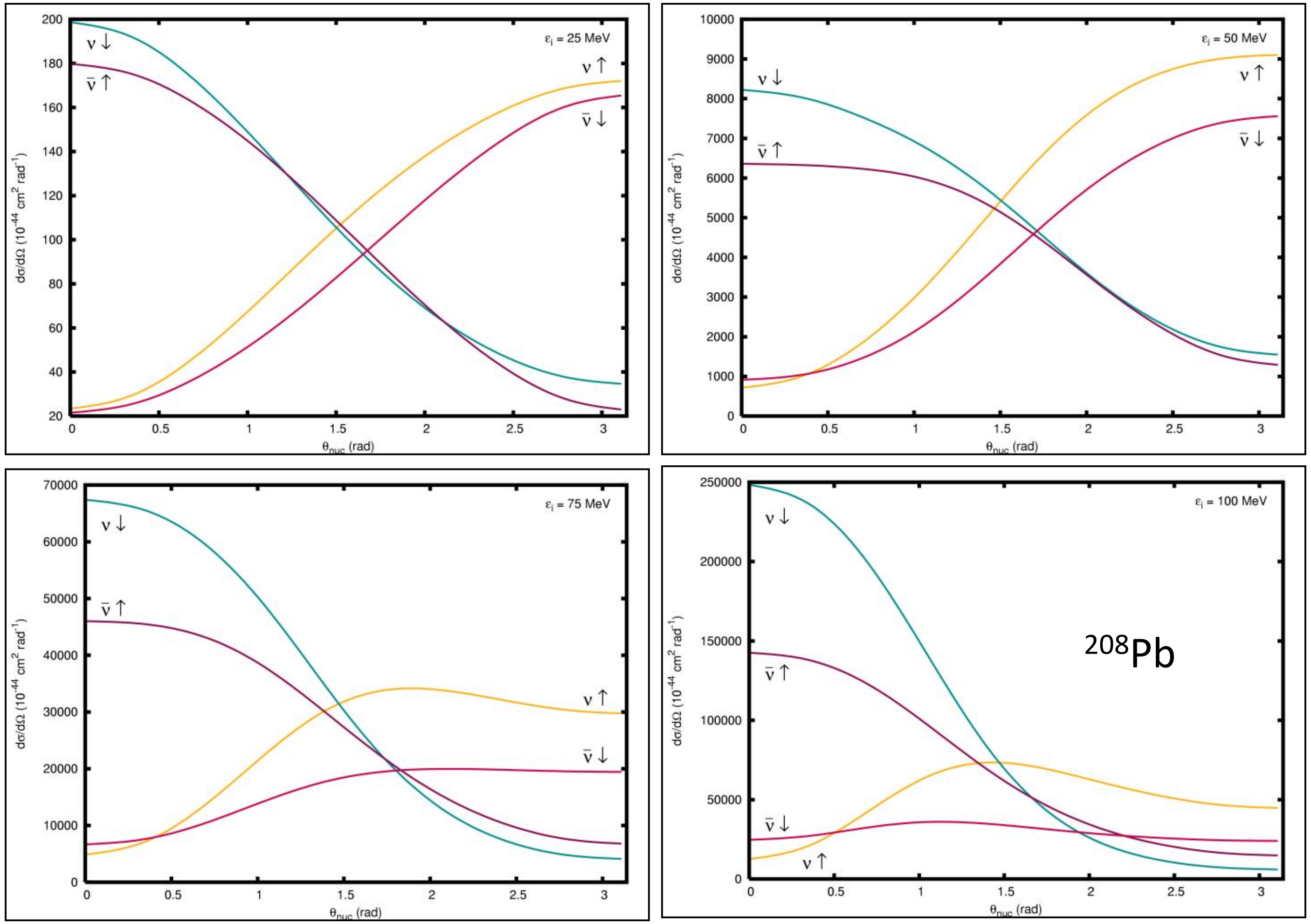
Adding antineutrinos to the picture :



- Neutrinos favor 'spin down' nucleon knockout
- Antineutrinos mainly induce 'spin up' knockout reactions
- Polarization differences increase with incoming neutrino energies

Spin of the outgoing nucleon

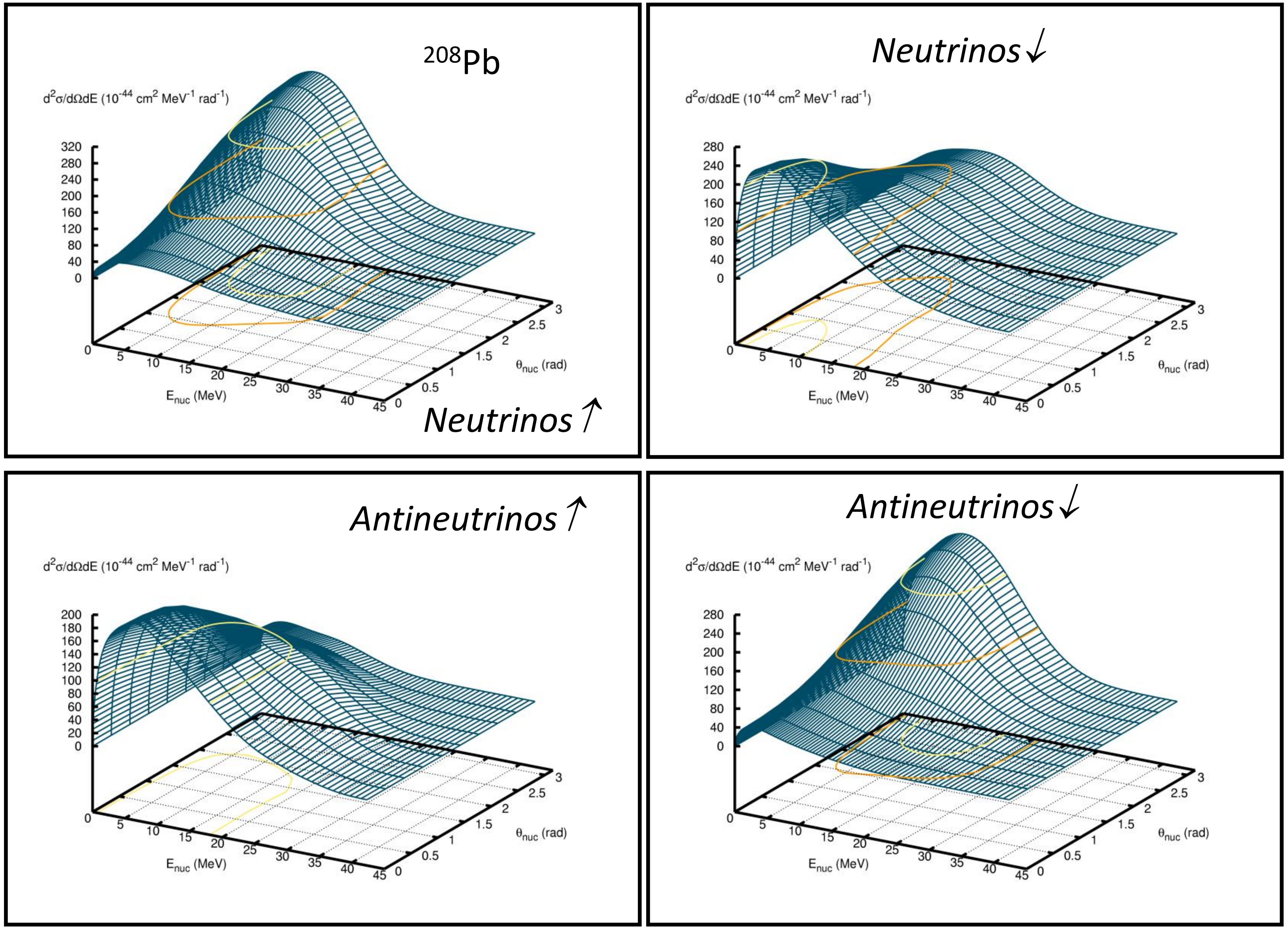
The asymmetry and the dissimilarities between neutrinos and antineutrinos are most clear considering the angular cross section :



The asymmetry is most prominent for forward nucleon knockout, and remains large over a broad angular range.

For the suppressed backward scattering reactions, the asymmetry is completely reversed

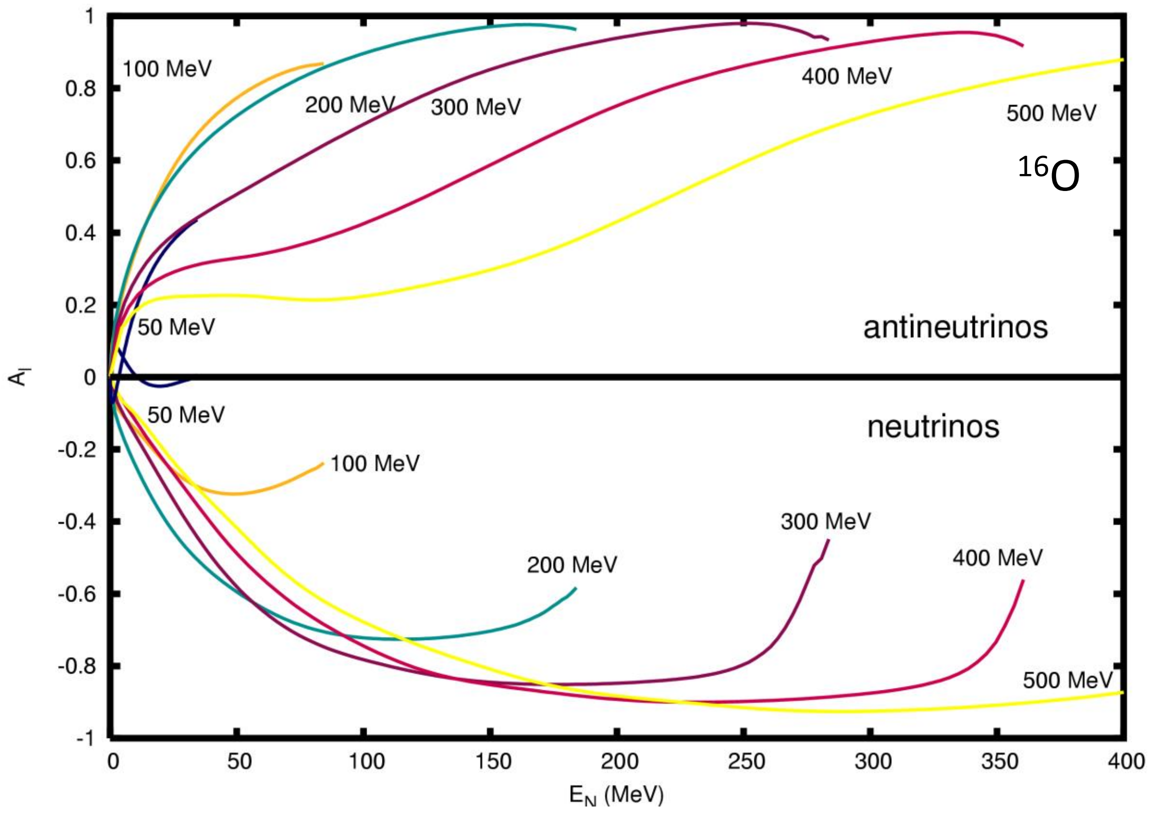
Spin of the outgoing nucleon



Polarization asymmetry

Longitudinal **polarization asymmetry :**

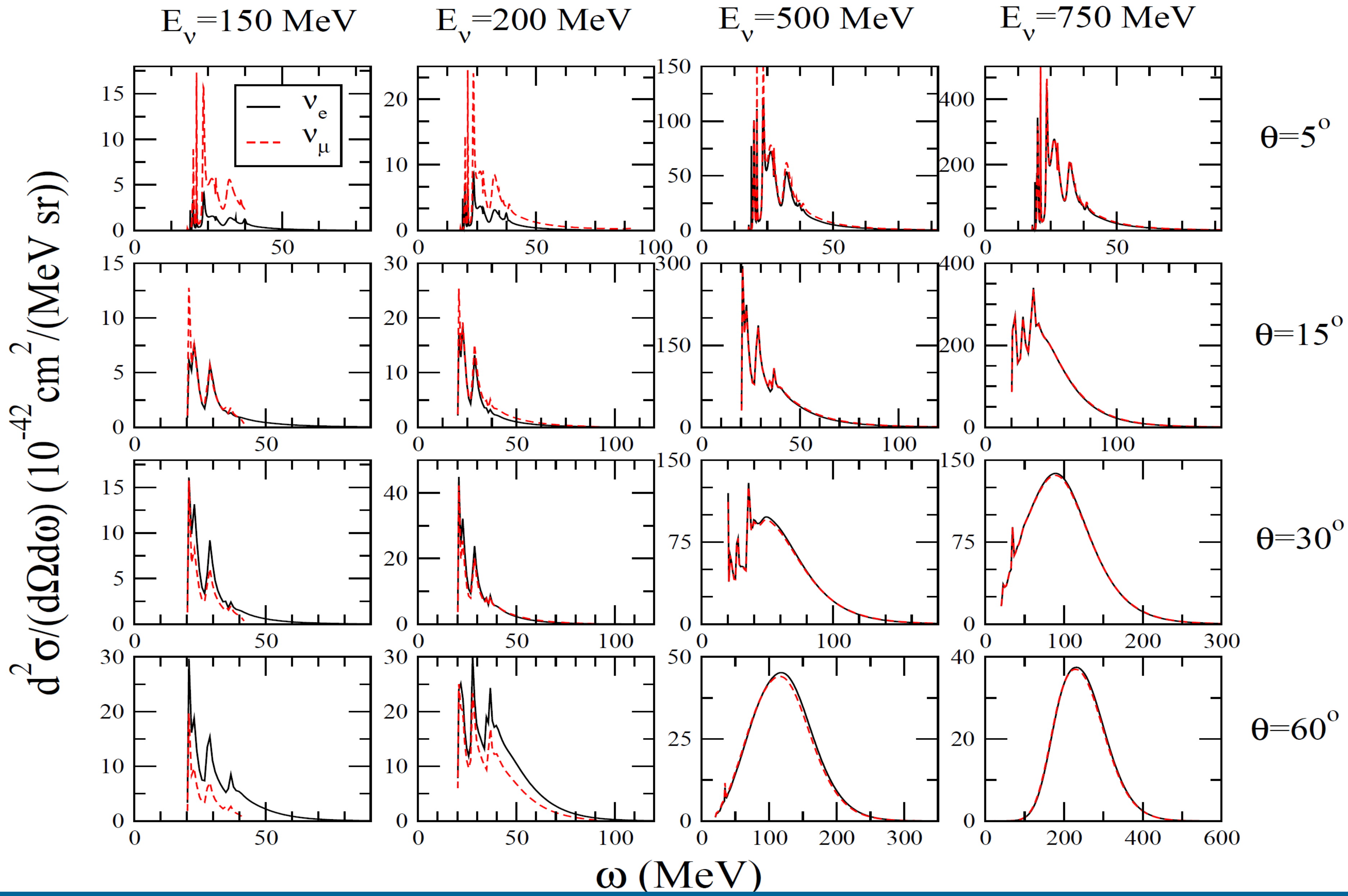
$$A_l = \frac{\sigma(s_N^l = \uparrow) - \sigma(s_N^l = \downarrow)}{\sigma(s_N^l = \uparrow) + \sigma(s_N^l = \downarrow)}$$



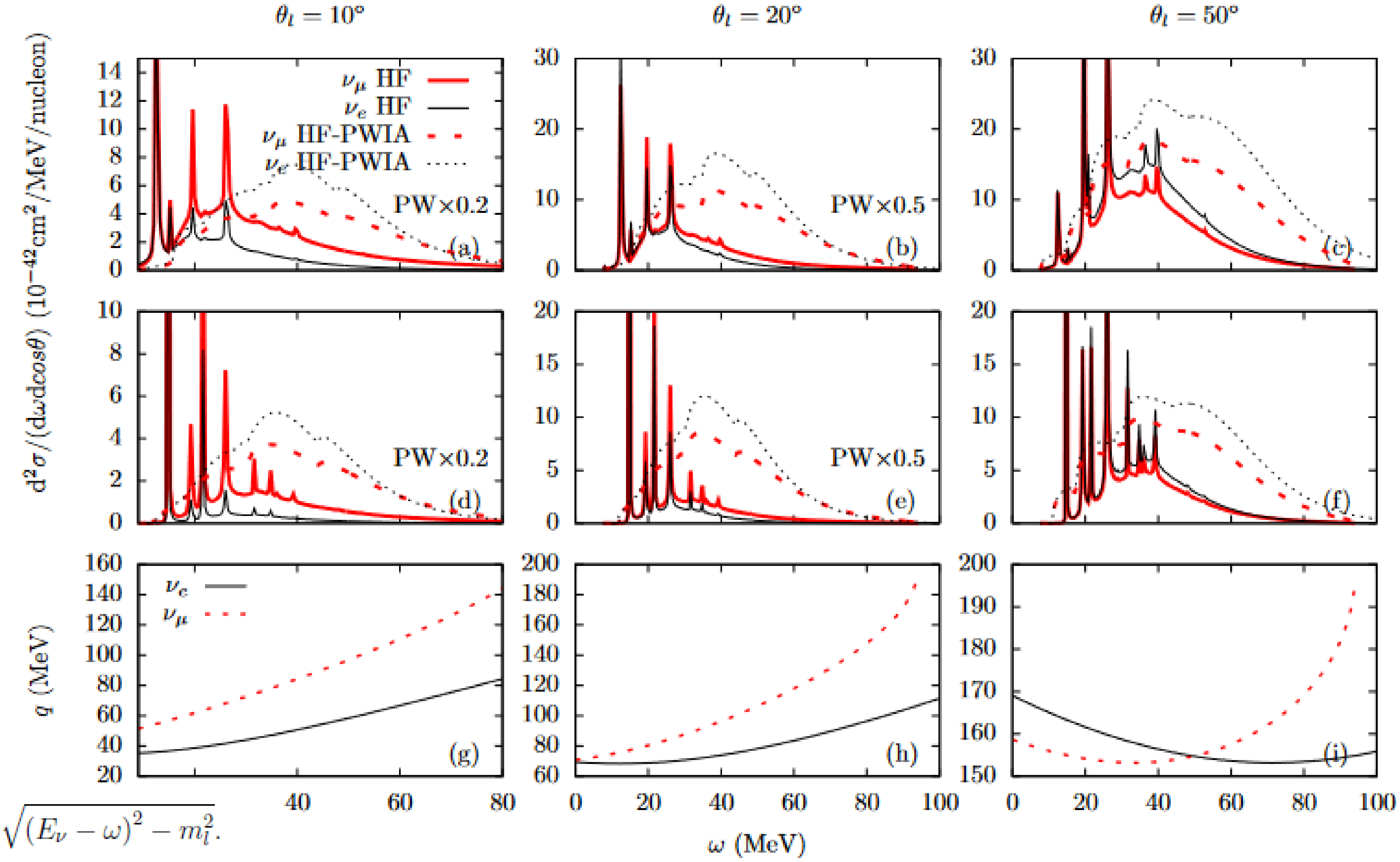
- For antineutrinos, A_l is large and positive
- For neutrinos, A_l is large and negative

N. J., K. Vantournhout, J. Ryckebusch, K. Heyde, PRL 93, 082501 (2004) ; N. J. , K. Vantournhout, J. Ryckebusch, K. Heyde, PRC71, 034604 (2005).

Electron vs muon neutrino CC cross sections



Electron vs muon neutrino CC cross sections

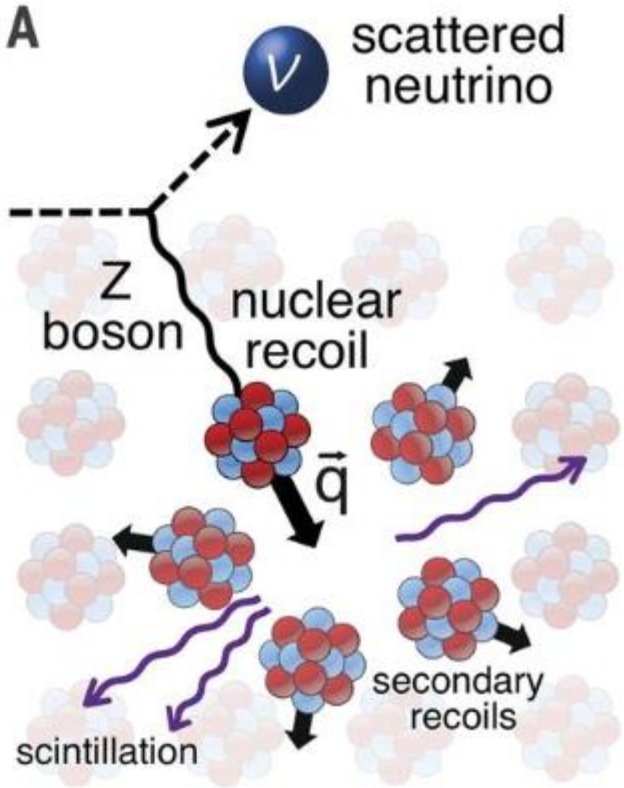


$$q = \sqrt{E_\nu^2 + P_l^2 - 2\cos\theta_l E_\nu P_l} \approx E_\nu - \sqrt{(E_\nu - \omega)^2 - m_l^2}.$$

A. Nikolakopoulos, N.J, N. Van Dessel, K. Niewczas, R. Gonzalez-Jimenez, J.M. Udias, V. Pandey, Phys. Rev. Lett. 123, 052501 (2019)

Coherent Scattering

Science, September 2017 : **The First Observation of Coherent Elastic Neutrino Nucleus Scattering**



Coherent Scattering

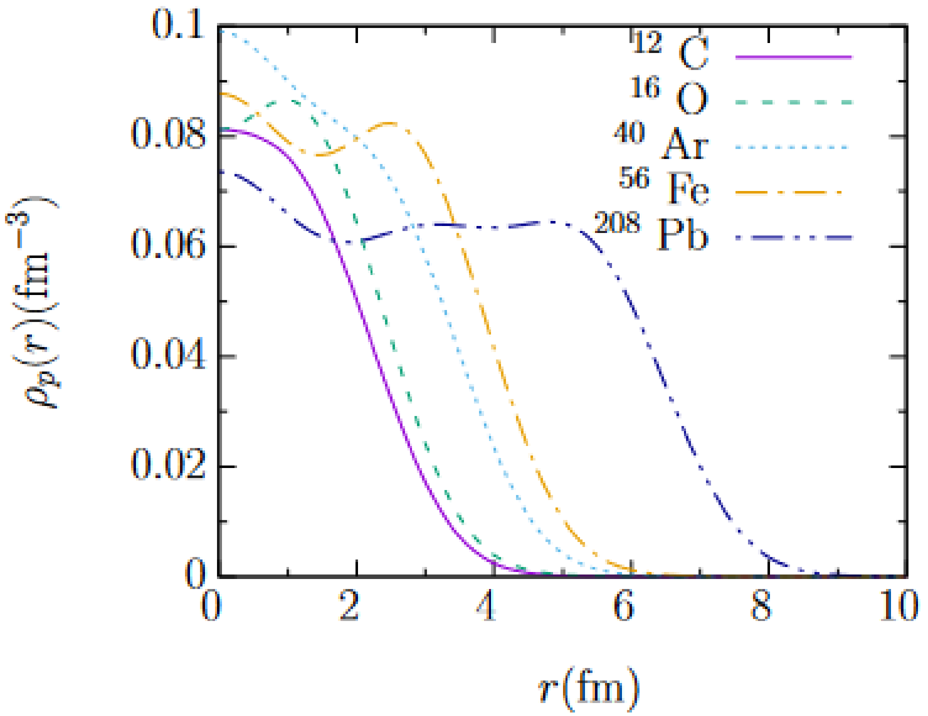
Coherent cross section as a function of nuclear recoil energy :

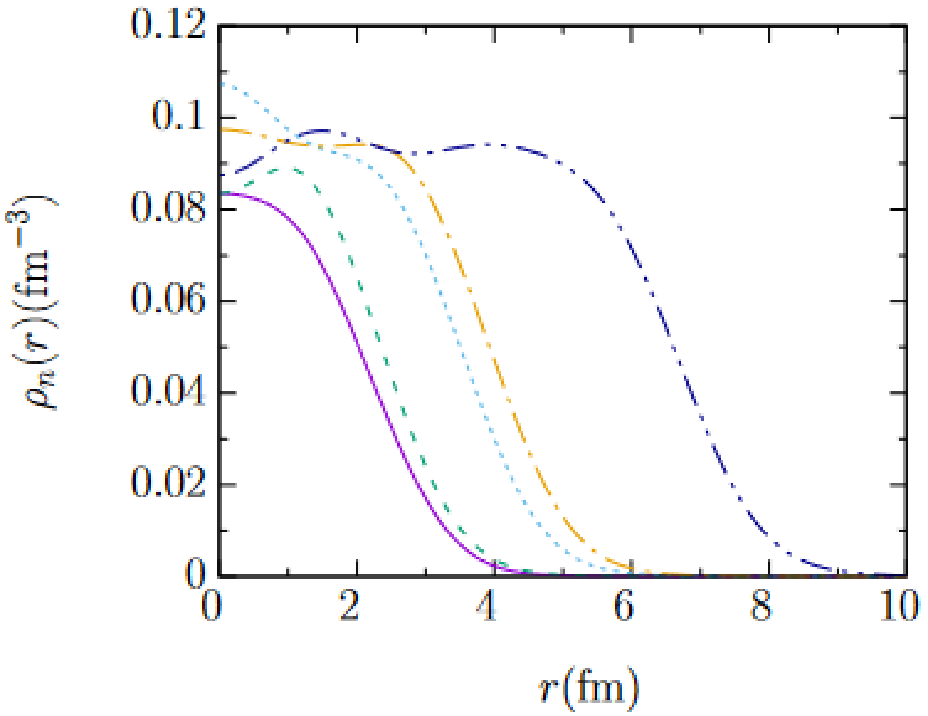
$$\frac{d\sigma}{dT} = \frac{G_F^2}{\pi} M_A \left(1 - \frac{T}{E_i} - \frac{M_A T}{2E_i^2}\right) \frac{Q_W^2}{4} F_W^2(Q^2),$$

Mainly sensitive to neutron distributions

$$F(Q^2) = \frac{4\pi}{Q_W} \int \left((1 - 4\sin^2\theta_W)\rho_p(r) - \rho_n(r)\right) \frac{\sin(qr)}{qr} r^2 dr,$$

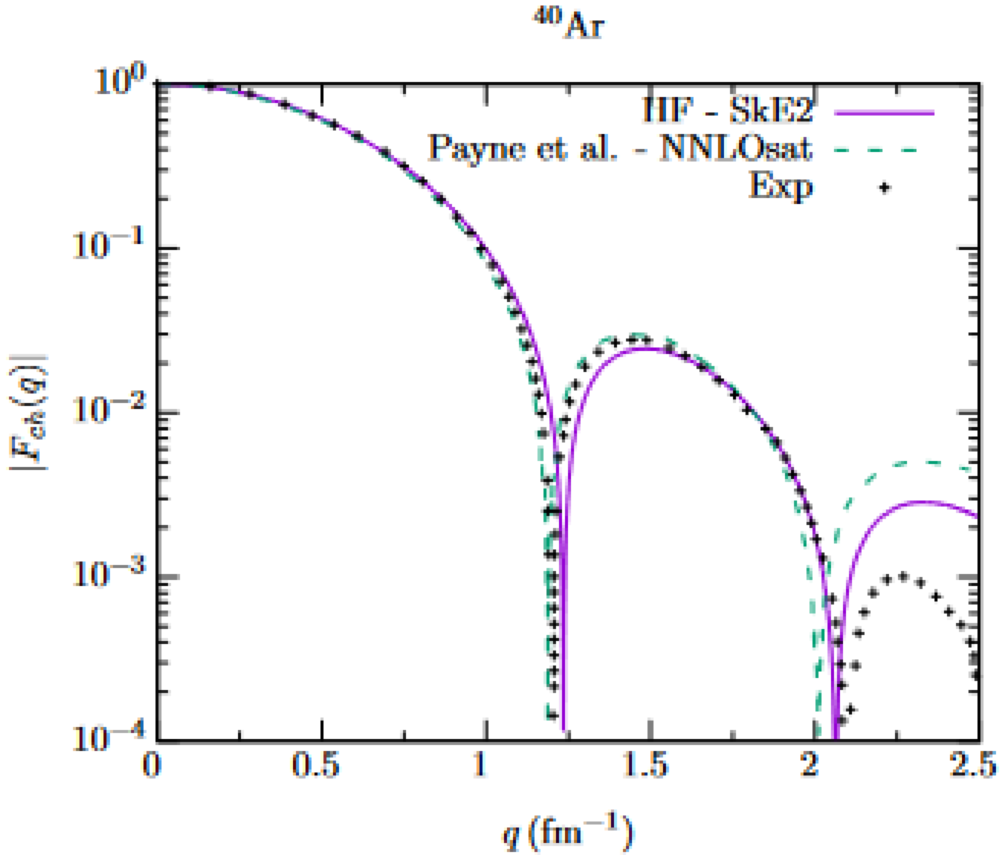
$$F_W(Q^2) = \frac{1}{Q_W} \left[(1 - 4\sin^2\theta_W) f_p(\vec{q}) F_p(Q^2) \; - f_n(\vec{q}) F_n(Q^2)\right]$$

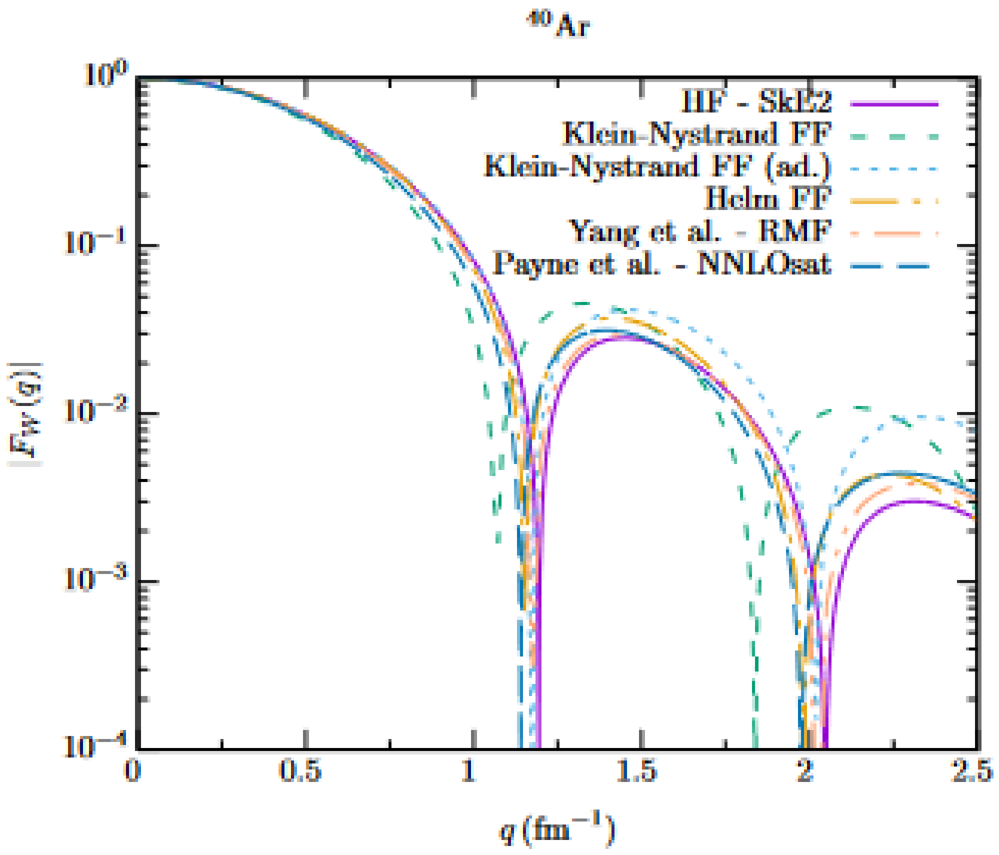




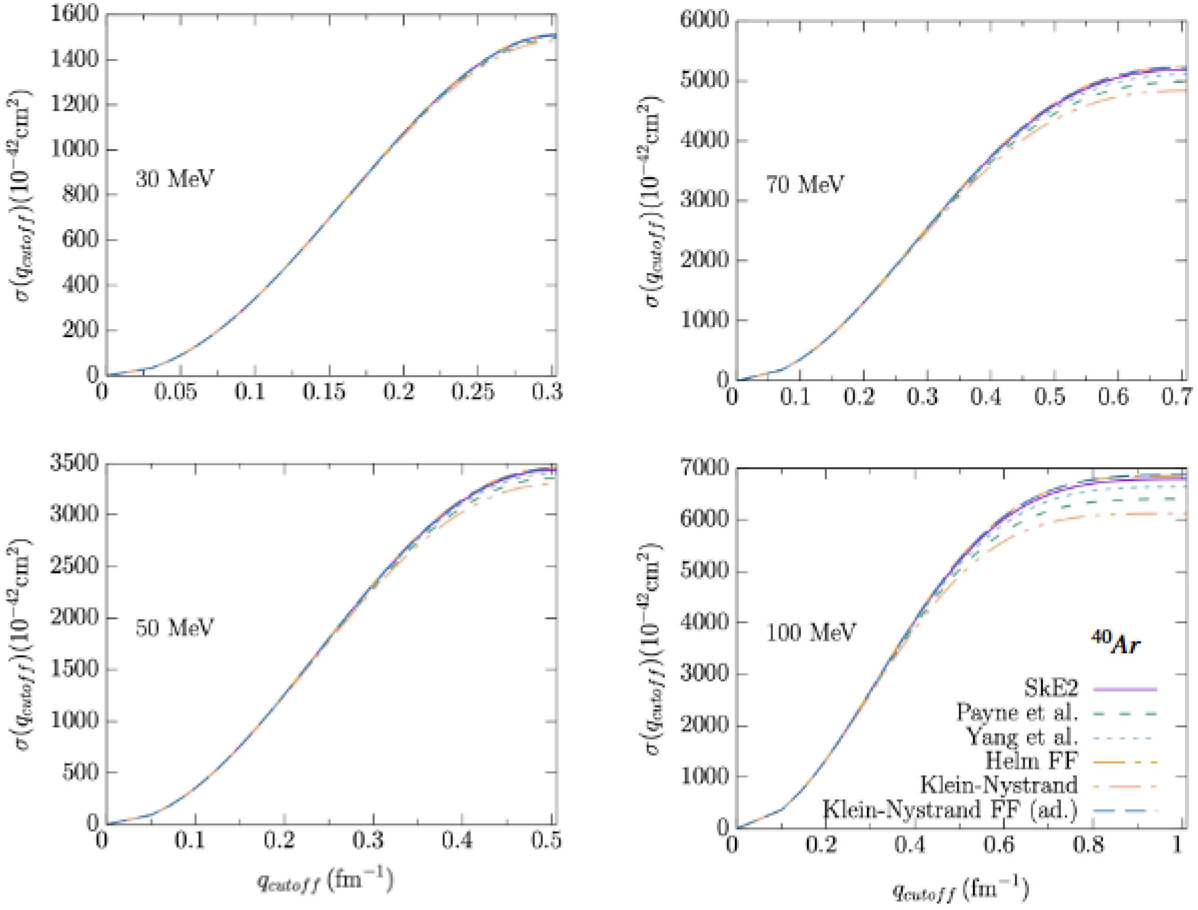
Coherent Scattering

Weak form factor : Model comparison

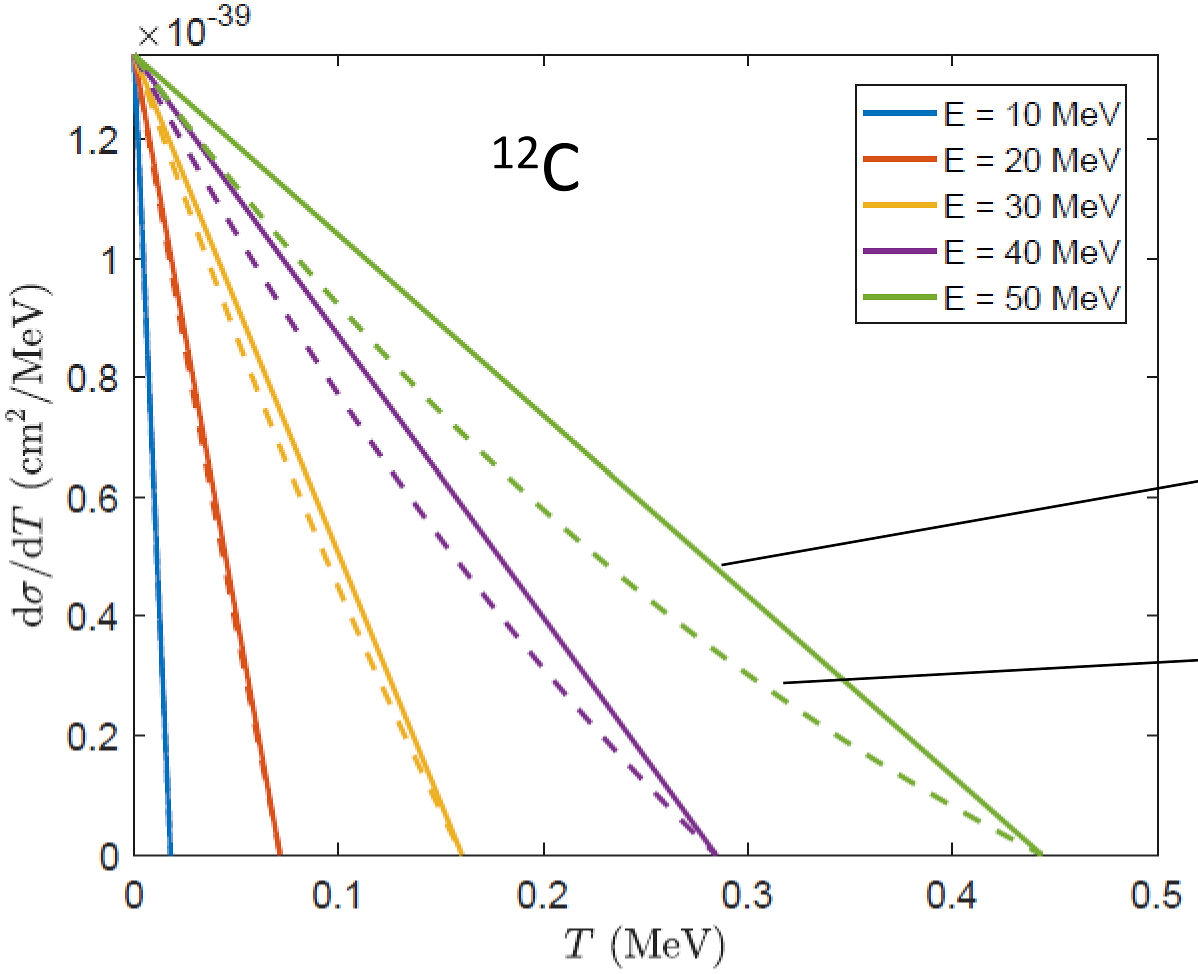




Weak form factor : Model comparison



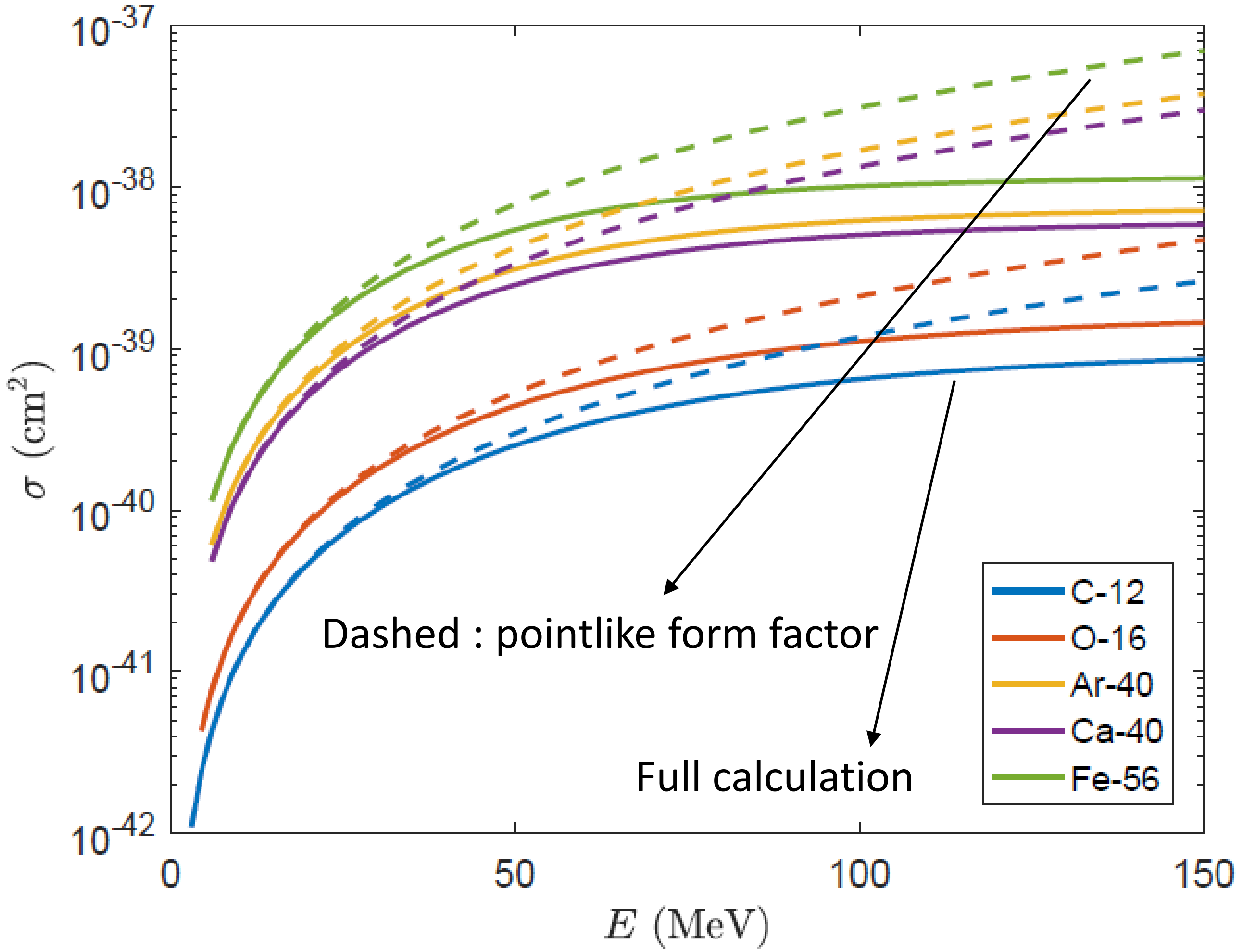
Coherent Scattering



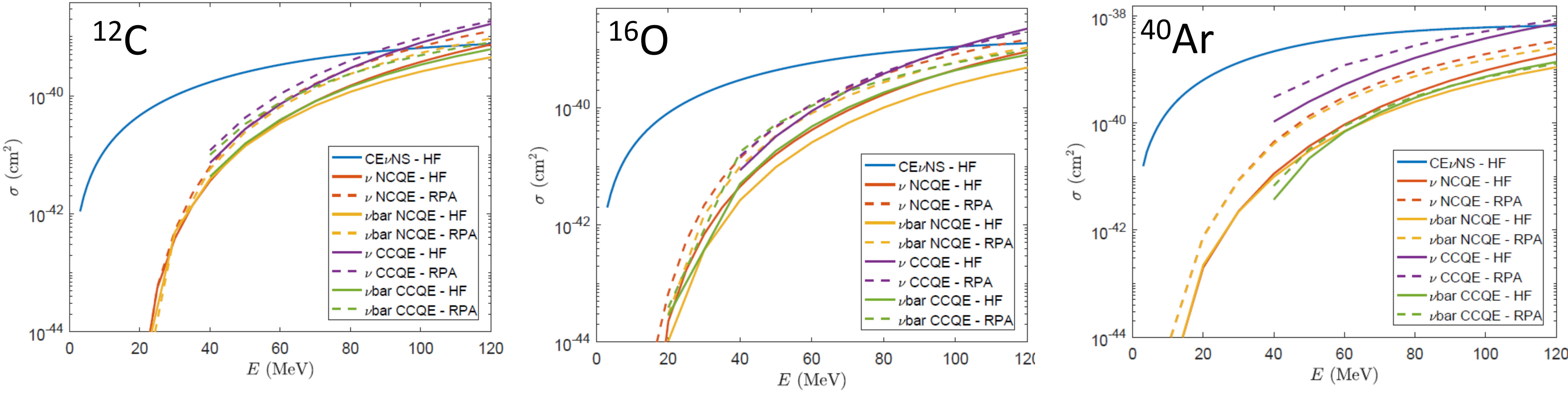
Full lines : pointlike form factor

Dashed : full calculation

Coherent Scattering



Coherent Scattering



- Strong mass dependence of coherent cross section
- Coherent process stronger than inelastic over a large kinematic range

Summary & Outlook

Neutrino-nucleus scattering at low energies provides a very rich source of information about the weak interaction and nuclear structure effects, of interest for weak particle, nuclear as well as astrophysics !