

Towards a more precise description of the neutrinoless double beta decay

Mihai Horoi

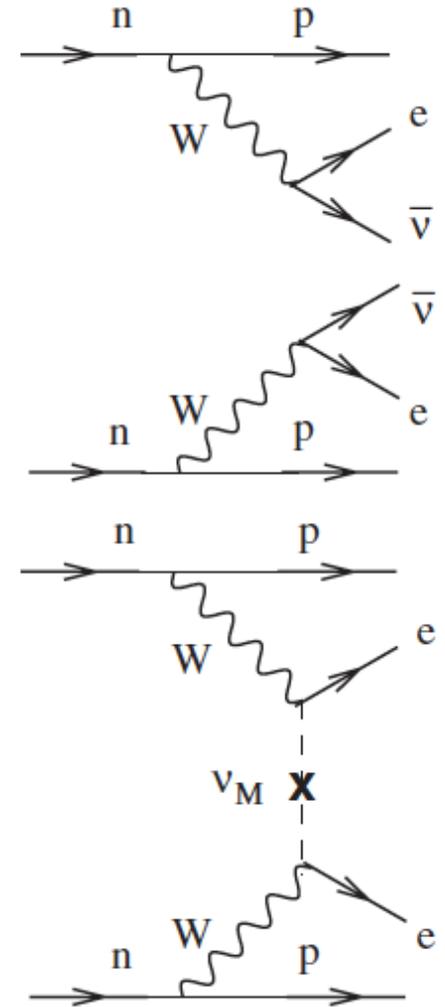
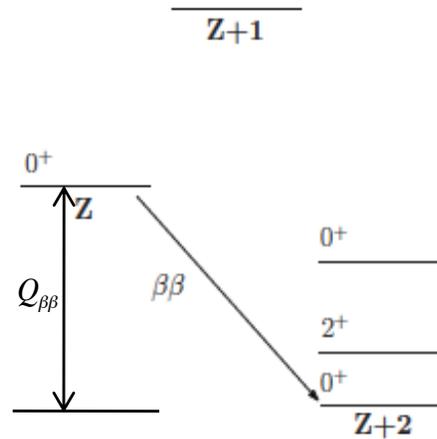
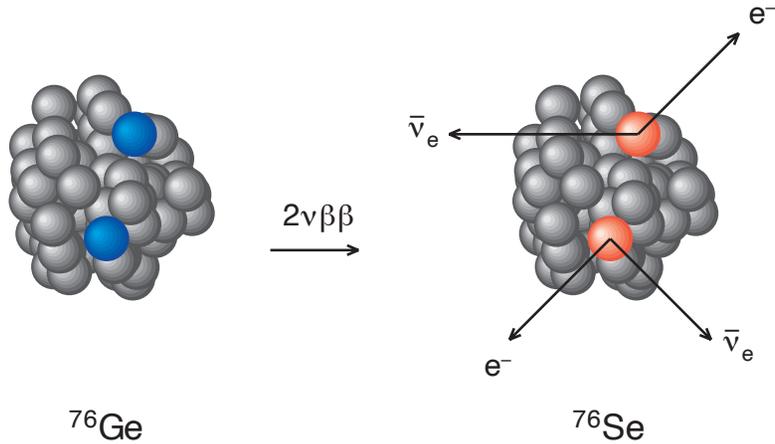
Department of Physics, Central Michigan University, Mount Pleasant, Michigan 48859, USA

Support from DOE grant DE-SC0022538 is acknowledged

Plan of talk

- Neutrinoless double beta decay: Brief introduction
- Nuclear matrix elements
 - Shell model approach
 - Statistical analysis for ^{136}Xe
- Future work: Interference effects
- Conclusions and outlook

Classical Double Beta Decay Problem



$$\langle m_{\beta\beta} \rangle = \left| \sum_k m_k U_{ek}^2 \right|$$

$$T_{1/2}^{-1}(0\nu) = G^{0\nu} (Q_{\beta\beta}) \left[M^{0\nu}(0^+) \right]^2 \left(\frac{\langle m_{\beta\beta} \rangle}{m_e} \right)^2$$

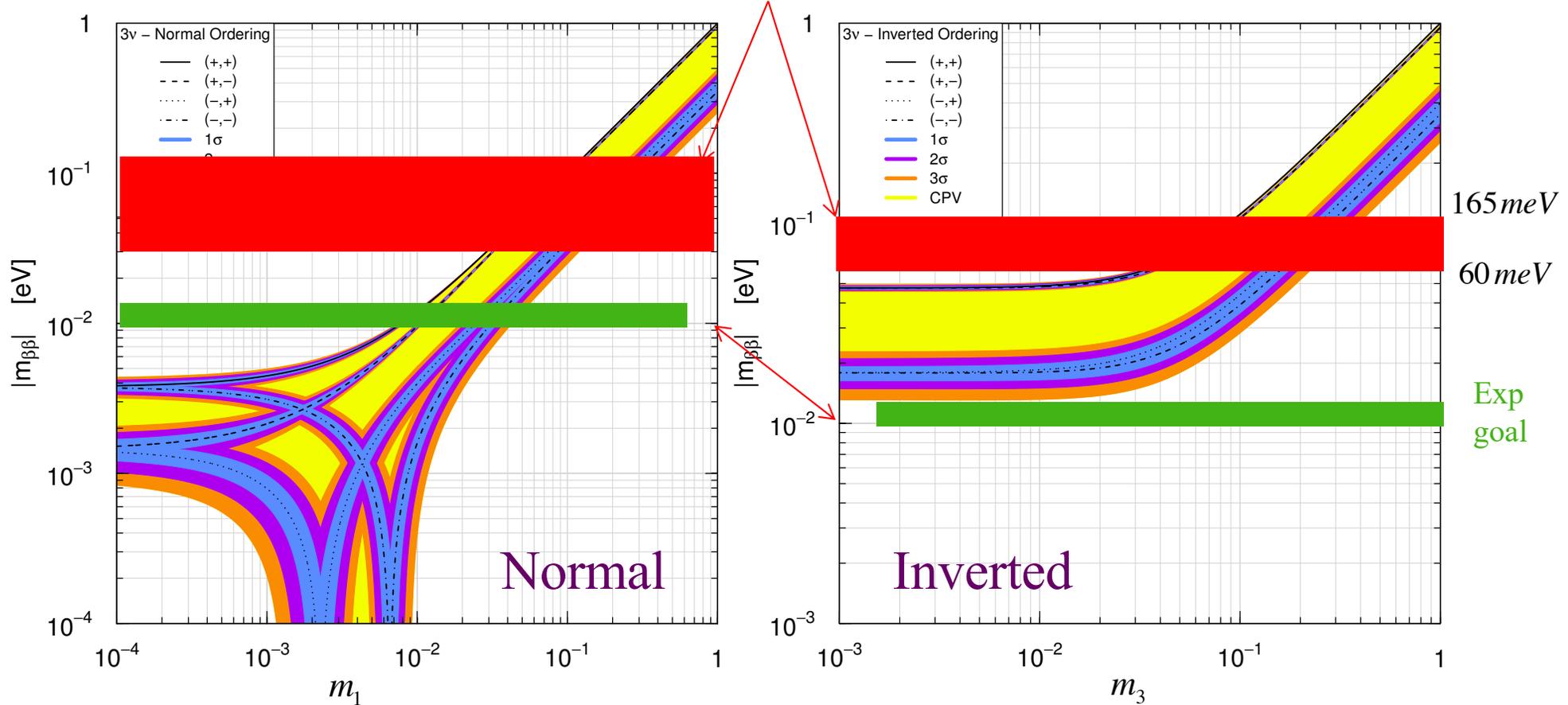
$$\nu_e(x) = \sum_i U_{ei} \nu_i(x)$$

$$|\nu_e\rangle = \sum_i U_{ei}^* |\nu_i\rangle$$

Neutrino $\beta\beta$ effective mass

arxiv:1507.08204

KamLAND – Zen, PRL 117, 082503 (2016): ^{136}Xe



$$|m_{\beta\beta}| = \left| \sum_{k=1}^3 m_k U_{ek}^2 \right| = \left| c_{12}^2 c_{13}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3} \right|$$

$$\phi_2 = \alpha_2 - \alpha_1 \quad \phi_3 = -\alpha_1 - 2\delta$$

$$\Leftrightarrow T_{1/2}^{-1}(0\nu) = G^{0\nu} (Q_{\beta\beta}) \left[M^{0\nu}(0^+) \right]^2 (\eta_{0\nu})^2$$

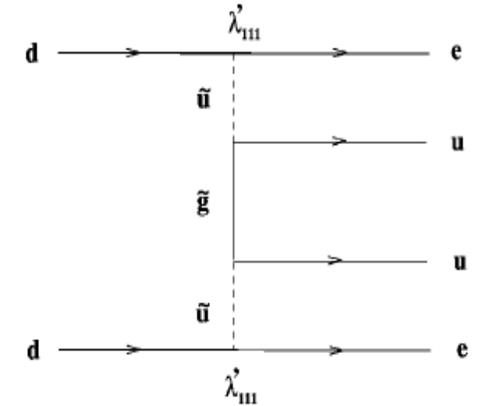
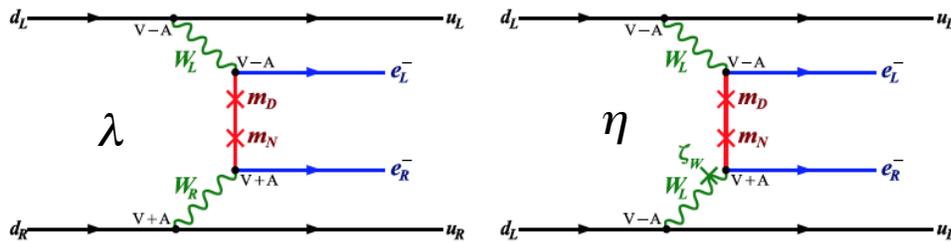
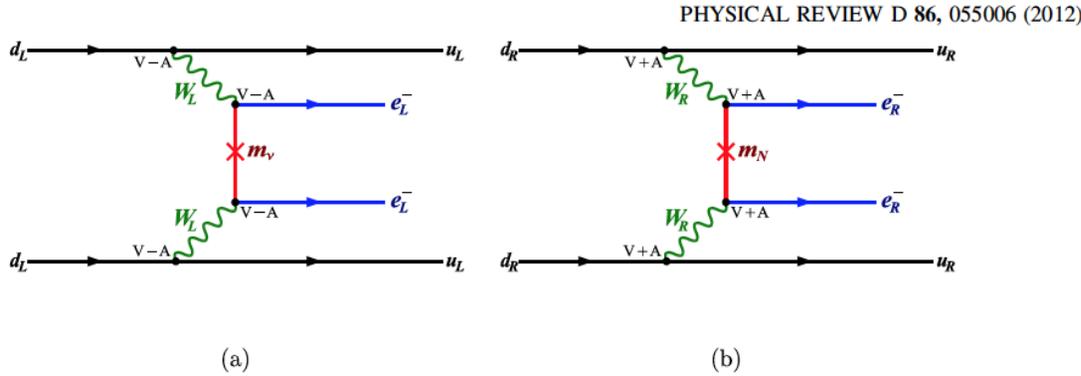
$$\eta_{0\nu} = \frac{|m_{\beta\beta}|}{m_e}$$

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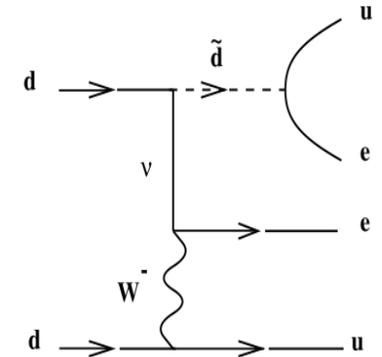
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Other models: Left-Right Symmetric Model (LRSM) and SUSY R-parity violation

DAS *et al.*



Gluino exchange



Squark exchange

$$\left[T_{1/2}^{0\nu} \right]^{-1} = G_{01} g_A^4 \left| \eta_{0\nu} M_{0\nu} + (\eta_{N_R}^L + \eta_{N_R}^R) M_{0N} + \eta_{\tilde{q}} M_{\tilde{q}} + \eta_{\lambda'} M_{\lambda'} + \eta_{\lambda} X_{\lambda} + \eta_{\eta} X_{\eta} \right|^2.$$

(e)

M. Horoi, A. Neacsu, PRD 93, 113014 (2016)

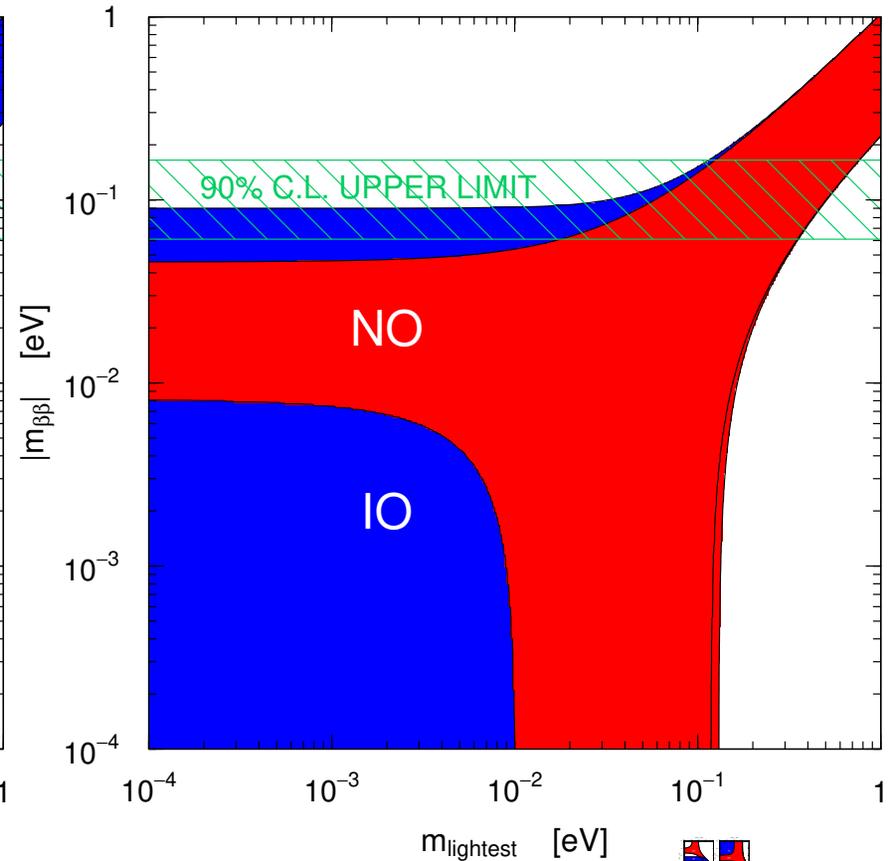
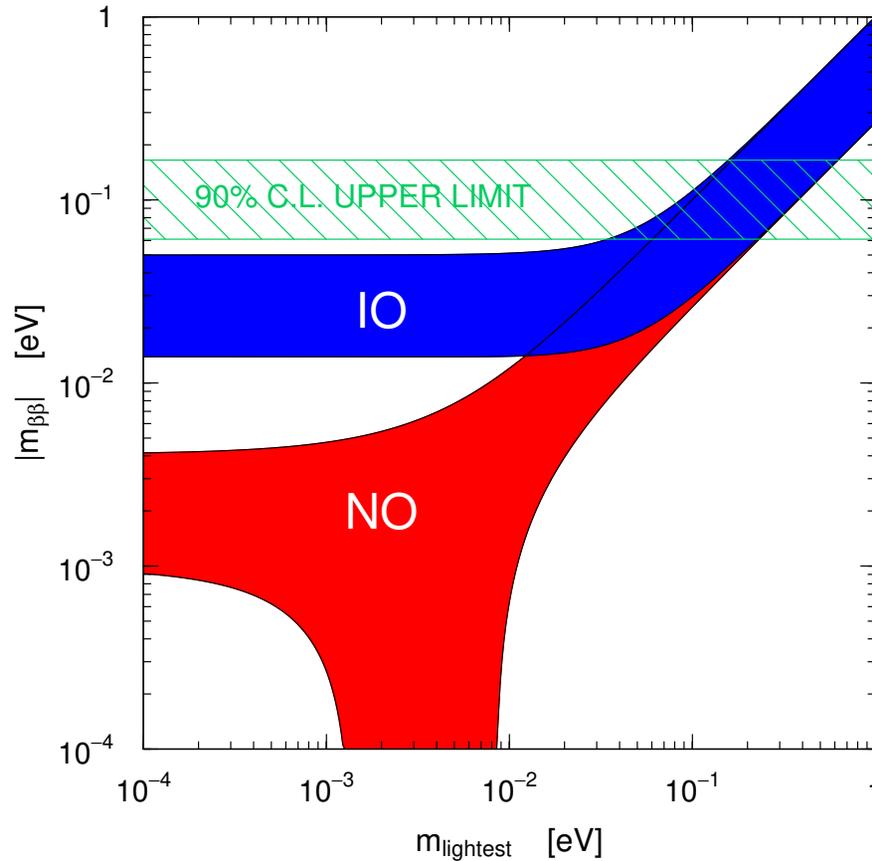
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$0\nu\beta\beta$ decay mass mechanism

3 neutrino flavors

3+1(sterile) neutrino flavors



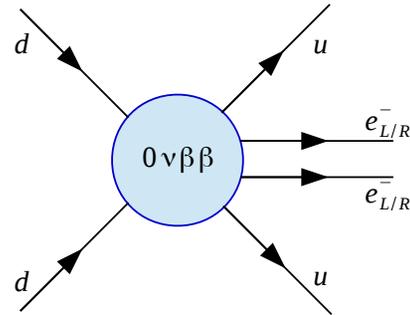
$$|m_{\beta\beta}| = \left| \sum_{k=1}^3 m_k U_{ek}^2 \right| = \left| c_{12}^2 c_{13}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3} \right|$$

$$\phi_2 = \alpha_2 - \alpha_1 \quad \phi_3 = -\alpha_1 - 2\delta$$

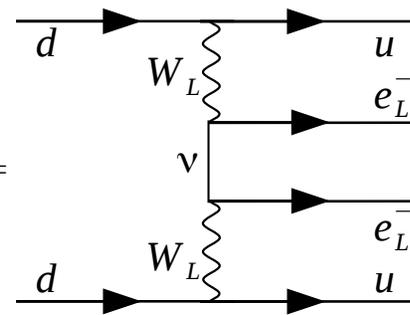
$$\Leftrightarrow T_{1/2}^{-1}(0\nu) = G^{0\nu} (Q_{\beta\beta}) [M^{0\nu}(0^+)]^2 (\eta_{0\nu})^2$$

$$\eta_{0\nu} = \frac{|m_{\beta\beta}|}{m_e}$$

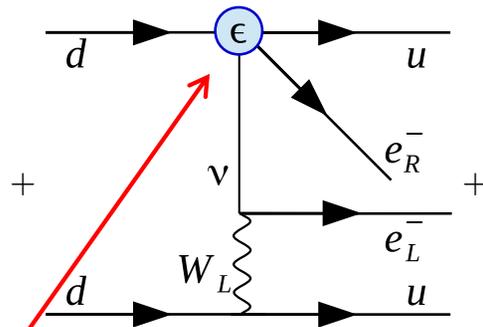
Effective field theory approach



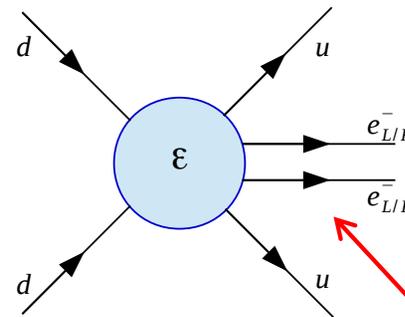
(a) The generic $0\nu\beta\beta$ decay diagram at the quark-level.



(b) Light left-handed neutrino exchange diagram.



(c) The long-range part of the $0\nu\beta\beta$ diagram.



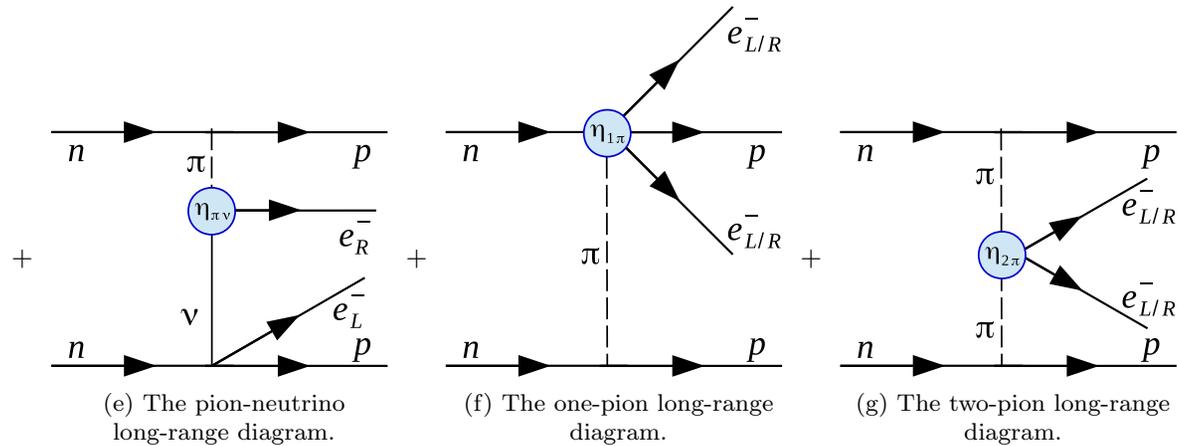
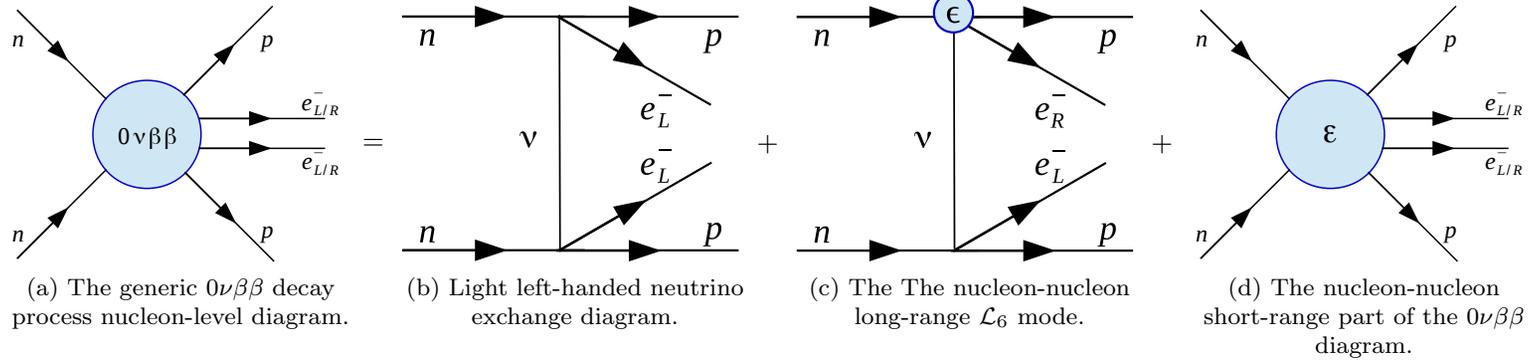
(d) The short-range part of the $0\nu\beta\beta$ diagram.

arxiv:1706.05391

$$\mathcal{L}_6 = \frac{G_F}{\sqrt{2}} \left[j_{V-A}^\mu J_{V-A,\mu}^\dagger + \sum_{\alpha,\beta}^* \epsilon_\alpha^\beta j_\beta J_\alpha^\dagger \right]$$

$$\mathcal{L}_9 = \frac{G_F^2}{2m_p} \left[\epsilon_1 J J j + \epsilon_2 J^{\mu\nu} J_{\mu\nu} j + \epsilon_3 J^\mu J_\mu j \right. \\ \left. + \epsilon_4 J^\mu J_{\mu\nu} j^\nu + \epsilon_5 J^\mu J j_\mu \right],$$

Effective field theory after hadronization



$$\left[T_{1/2}^{0\nu} \right]^{-1} = g_A^4 \left[\sum_i |\mathcal{E}_i|^2 \mathcal{M}_i^2 + \text{Re} \left[\sum_{i \neq j} \mathcal{E}_i \mathcal{E}_j \mathcal{M}_{ij} \right] \right]$$

$$\mathcal{E}_{2-7} = \{ \epsilon_{V-A}^{V+A}, \epsilon_{V+A}^{V+A}, \epsilon_{S \pm P}^{S+P}, \epsilon_{TR}^{TR}, \eta_{\pi\nu} \}$$

$$\mathcal{E}_{8-15} = \{ \epsilon_1, \epsilon_2, \epsilon_3^{LLz(RRz)}, \epsilon_3^{LRz(RLz)}, \epsilon_4, \epsilon_6, \eta_{1\pi}, \eta_{2\pi} \}$$

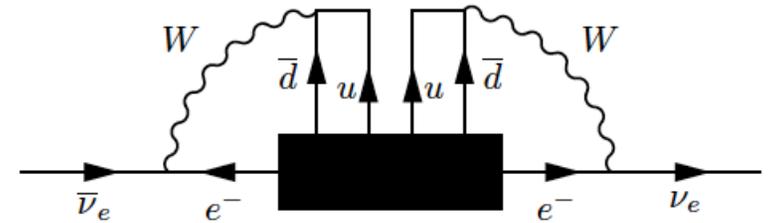
The Black Box Theorems

Black box I (electron neutrino)

J. Schechter and J.W.F Valle, PRD 25, 2951 (1982)

E. Takasugi, PLB 149, 372 (1984)

J.F. Nieves, PLB 145, 375 (1984)



- $0\nu\beta\beta$ observed \Leftrightarrow at some level
- (i) Lepton number conservation is violated by 2 units.
 - (ii) Electron neutrinos are Majorana fermions (with $m > 0$).

However:

M. Duerr et al, JHEP 06 (2011) 91

$$(\delta m_{\nu_e})_{BB} \sim 10^{-24} eV \ll \sqrt{|\Delta m_{32}^2|} \approx 0.05 eV$$

Black box II (all flavors + oscillations)

M. Hirsch, S. Kovalenko, I. Schmidt, PLB 646, 106 (2006)

- $0\nu\beta\beta$ observed \Leftrightarrow at some level
- (i) Lepton number conservation is violated by 2 units.
 - (ii) Neutrinos are Majorana fermions.

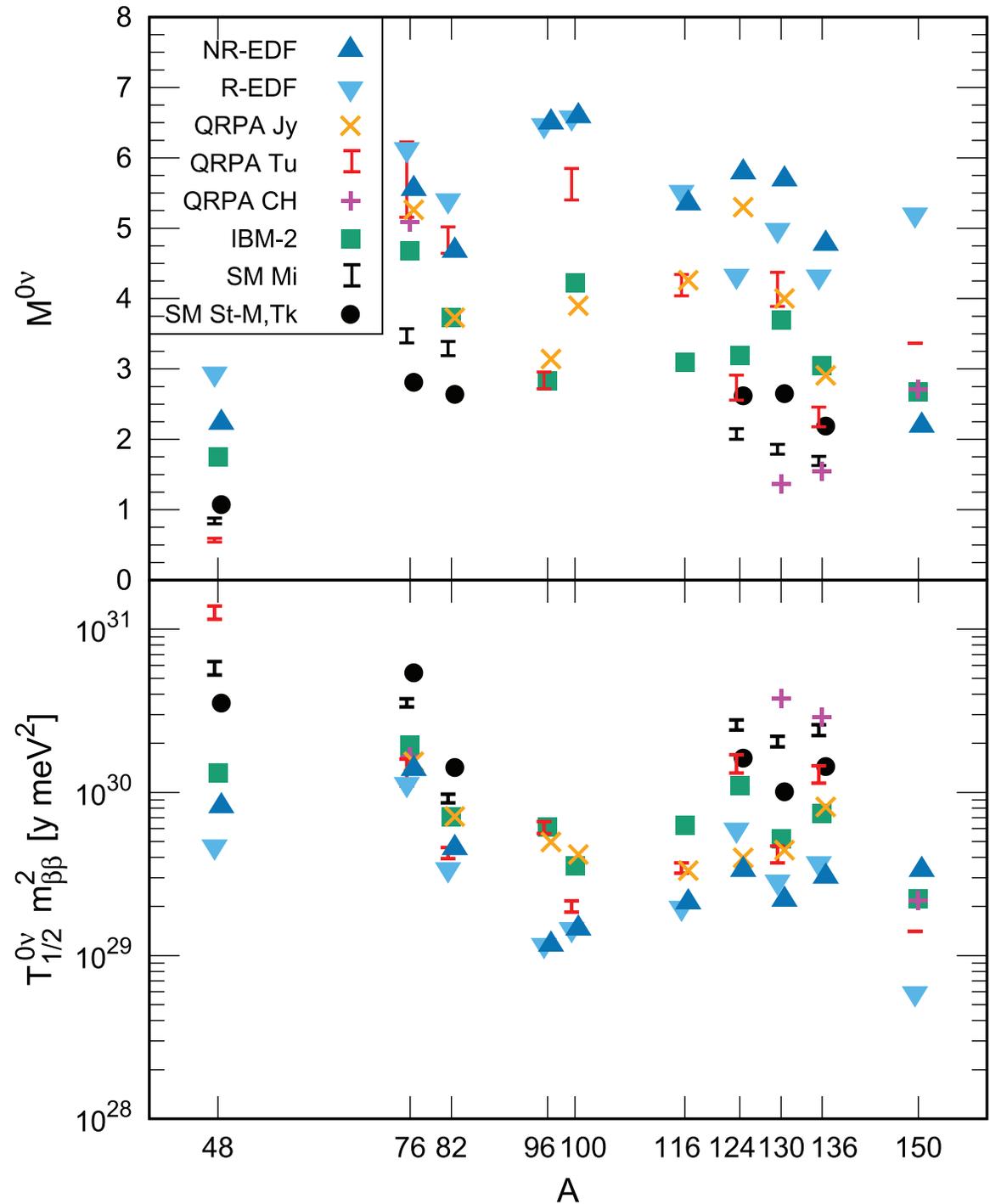
Regardless of the dominant $0\nu\beta\beta$ mechanism!

$$(iii) \quad \langle m_{\beta\beta} \rangle = \left| \sum_{k=1}^3 m_k U_{ek}^2 \right| = \left| c_{12}^2 c_{13}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3} \right| > 0$$

Nuclear matrix element: mass mechanism

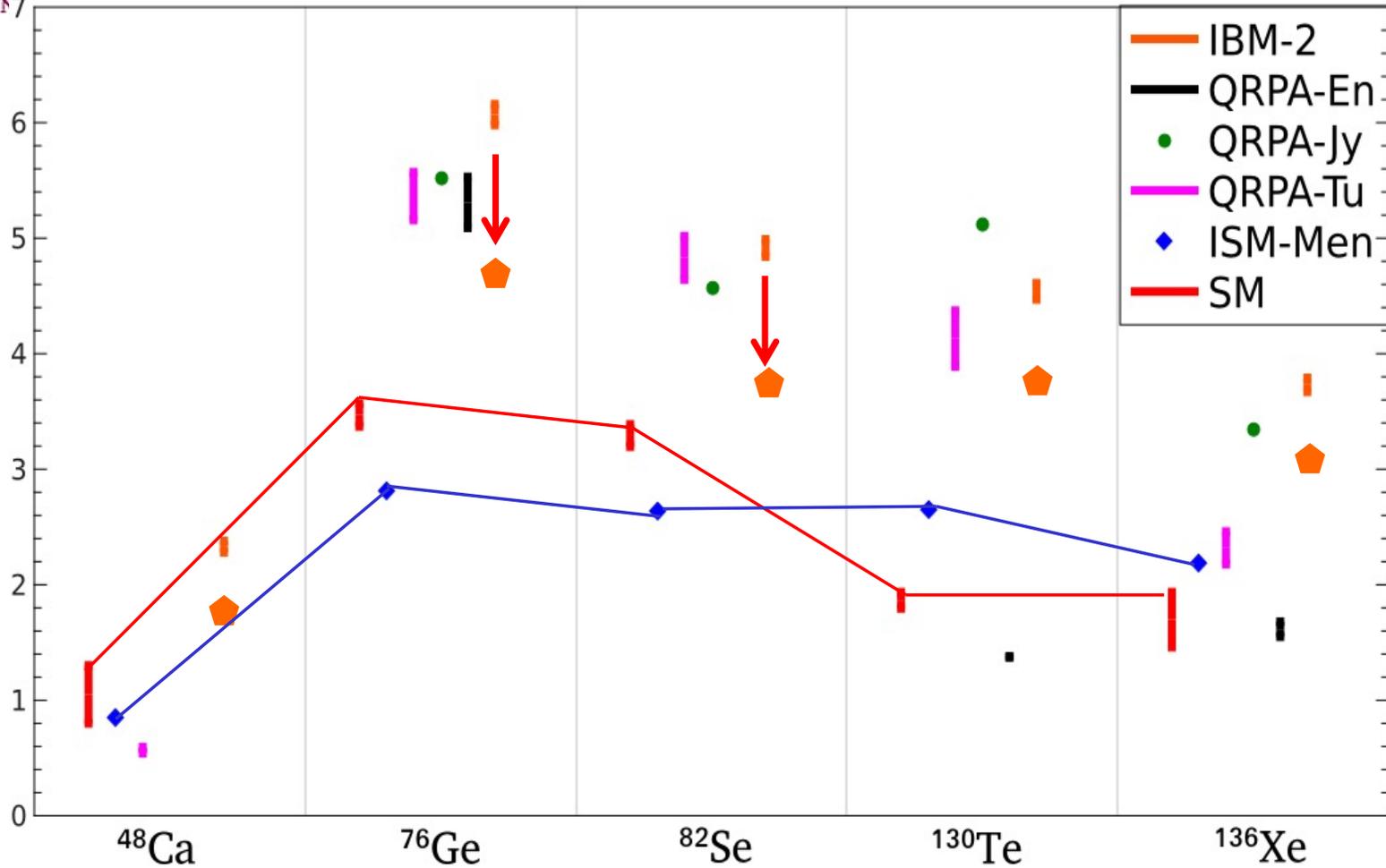
$$M_{0\nu} = M_{GT} + M_F + M_T + \dots$$

$$(T_{1/2}^{0\nu})^{-1} = G_{01} g_A^4 |M_{0\nu}|^2 |m_{\beta\beta}|^2$$



NME for the light-neutrino exchange mechanism

$M_{0\nu}$



IBA-2 J. Barea, J. Kotila, and F. Iachello, Phys. Rev. C **87**, 014315 (2013). **→** **IBM-2** PRC **91**, 034304 (2015)

QRPA-En M. T. Mustonen and J. Engel, Phys. Rev. C **87**, 064302 (2013).

QRPA-Jy J. Suhonen, O. Civitarese, Phys. NPA **847** 207–232 (2010).

QRPA-Tu A. Faessler, M. Gonzalez, S. Kovalenko, and F. Simkovic, arXiv:1408.6077

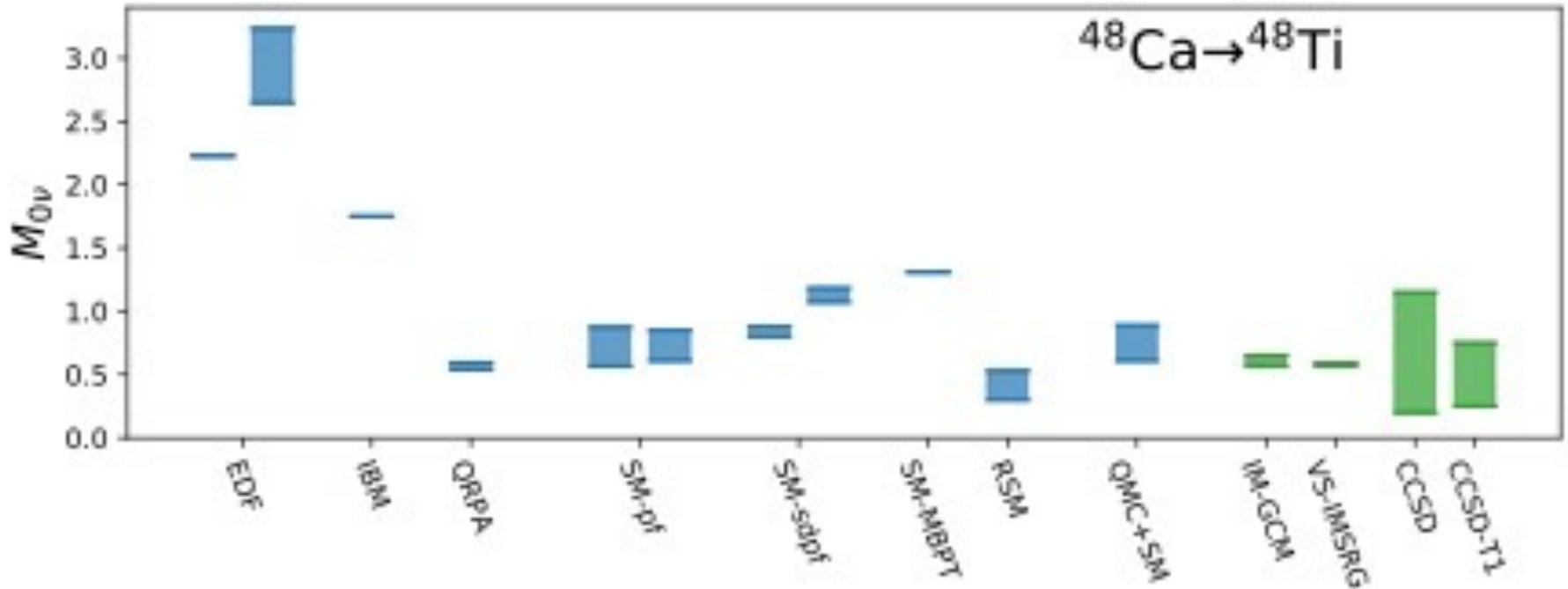
ISM-Men J. Menéndez, A. Poves, E. Caurier, F. Nowacki, NPA **818** 139–151 (2009).

SM M. Horoi et. al. PRC **88**, 064312 (2013), PRC **89**, 045502 (2014), PRC **89**, 054304 (2014), PRC **90**, 051301(R) (2014), PRC **91**, 024309 (2015), PRL **110**, 222502 (2013), PRL **113**, 262501(2014).

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Light left-handed neutrino-exchange NME update arxiv:2207.01085



Predicting the neutrinoless double- β -decay matrix element of ^{136}Xe using a statistical approach

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²International Center for Advanced Training and Research in Physics (CIFRA), Magurele 077125, Romania

³Horia Hulubei National Institute of Physics and Nuclear Engineering (IFIN-HH), Magurele 077125, Romania

PHYSICAL REVIEW C **106**, 054302 (2022)

Statistical analysis for the neutrinoless double- β -decay matrix element of ^{48}Ca

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IOP Publishing

Journal of Physics G: Nuclear and Particle Physics

J. Phys. G: Nucl. Part. Phys. **49** (2022) 120502 (30pp)

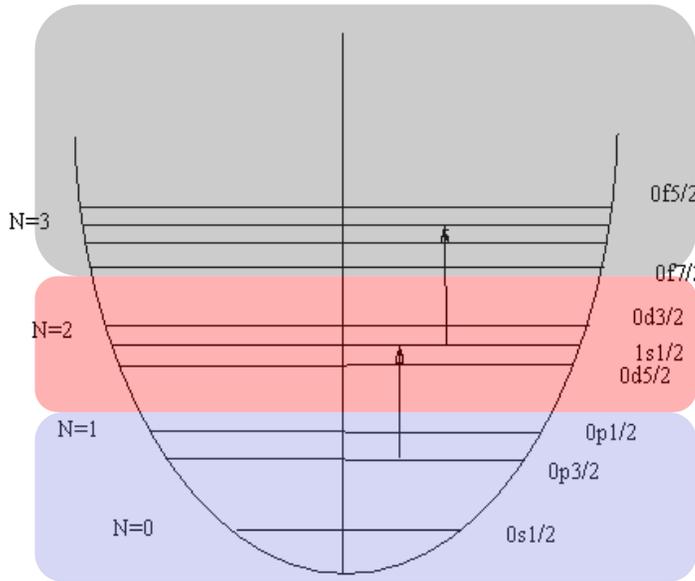
<https://doi.org/10.1088/1361-6471/aca03e>

Towards precise and accurate calculations of neutrinoless double-beta decay

V Cirigliano ¹, Z Davoudi ², J Engel ^{3,*}, R J Furnstahl ⁴,
G Hagen ^{5,6}, U Heinz ⁴, H Hergert ⁷, M Horoi ⁸,
C W Johnson ⁹, A Lovato ^{10,11,12}, E Mereghetti ¹³,
W Nazarewicz ^{14,*}, A Nicholson ^{3,15}, T Papenbrock ^{5,6},
S Pastore ¹⁶, M Plumlee ^{17,18}, D R Phillips ^{19,*},
P E Shanahan ²⁰, S R Stroberg ¹⁰, F Viens ²¹,
A Walker-Loud ¹⁵, K A Wendt ²² and S M Wild ^{18,23}

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Shell Model Effective Hamiltonians



$$g_A \sigma \tau \xrightarrow{\text{quenched}} g_A 0.77 \sigma \tau$$

empty

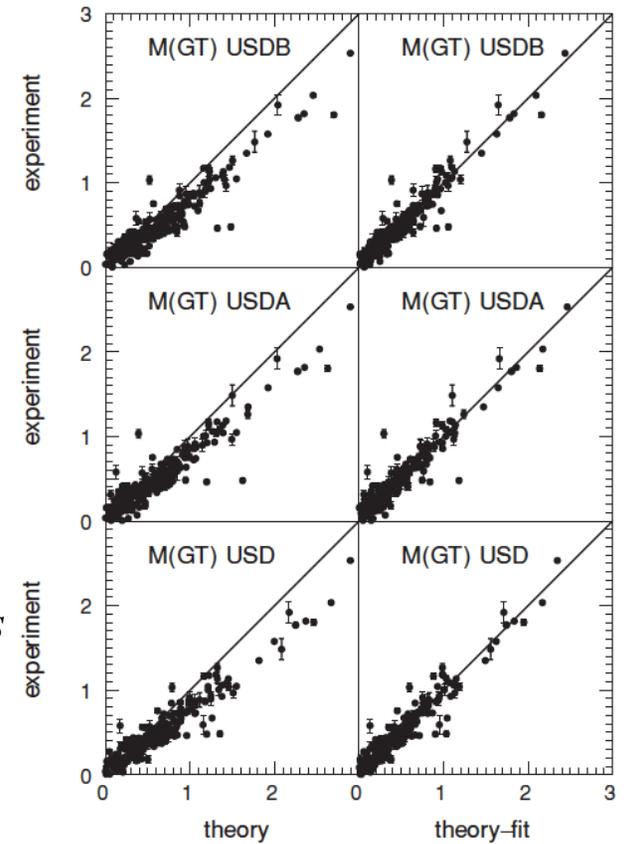
valence

frozen core

$$H_{\text{valence}} = H_{2\text{-body}}$$

can describe most correlations around the Fermi surface!

$$H_{\text{valence}} \Psi = E_n \Psi$$

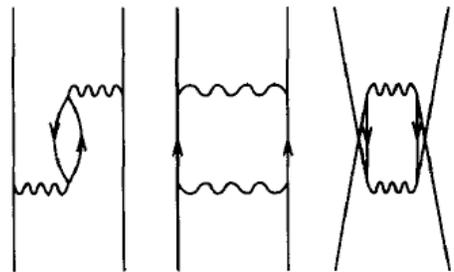


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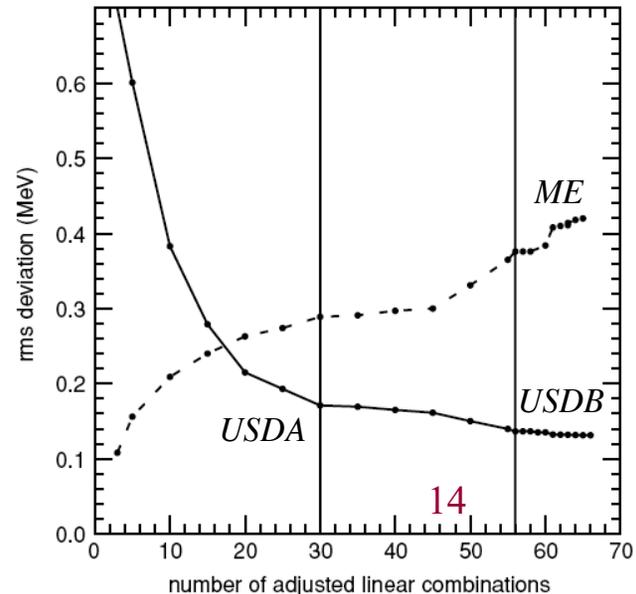
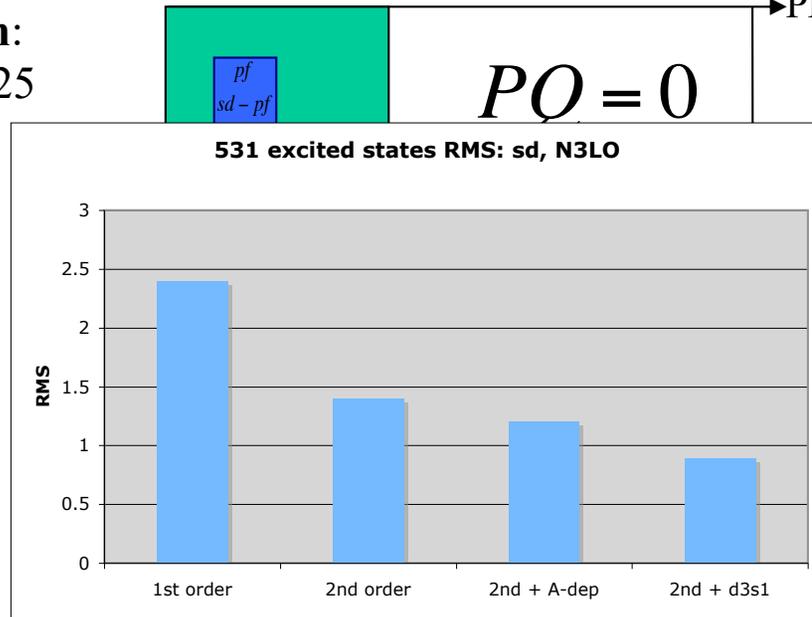
→ PRC 74, 34315 (2006), 78, 064302 (2008)



core polarization:
Phys.Rep. 261, 125
(1995)



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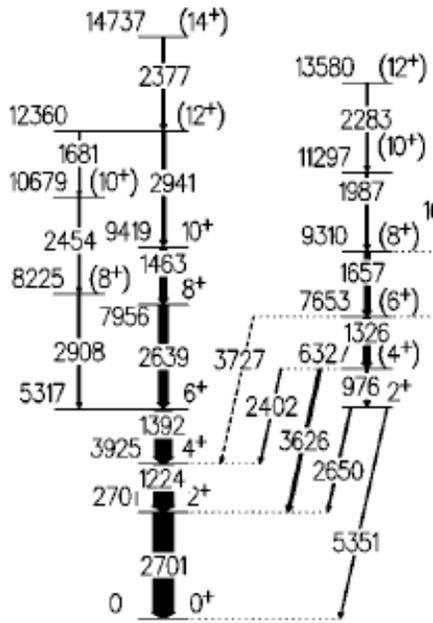


Koeln Low Spin States in ^{56}Ni : complete spectroscopy

$^{56}_{28}\text{Ni}_{28}$

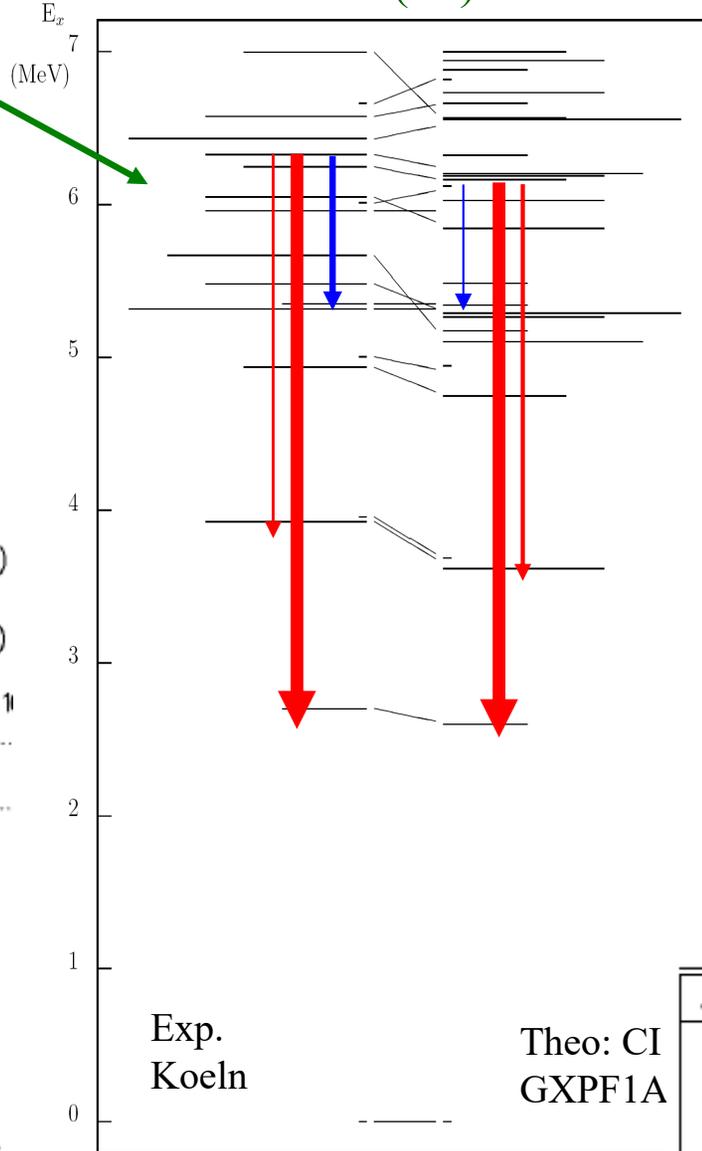
(sph)

①

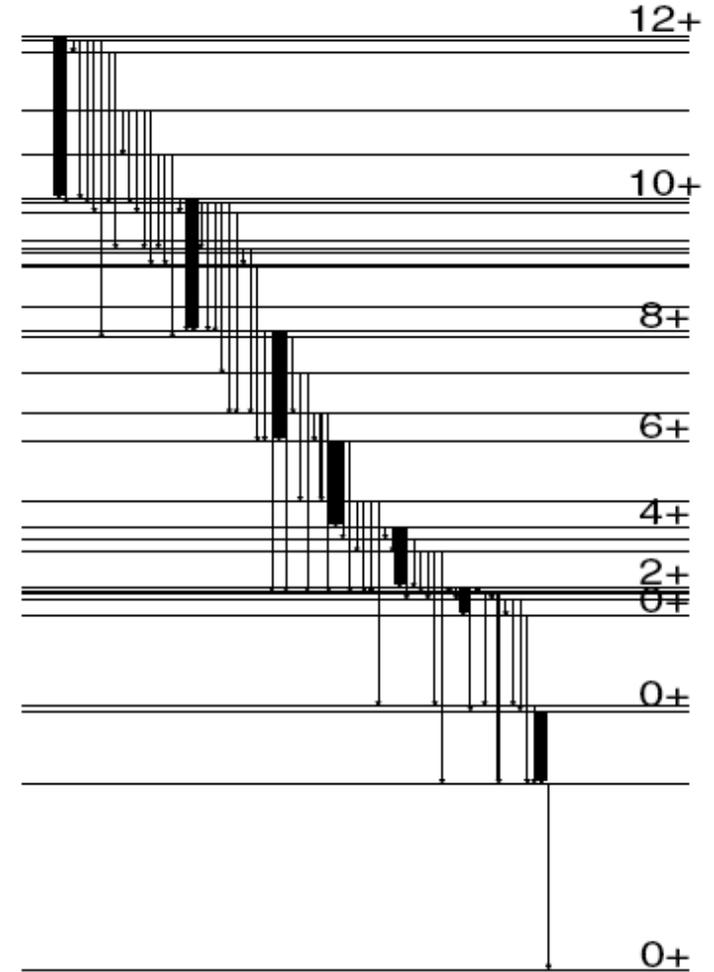


LBNL high spin experiment
D. Rudolf et al., PRL 88, 1999

Configuration Interaction (CI)



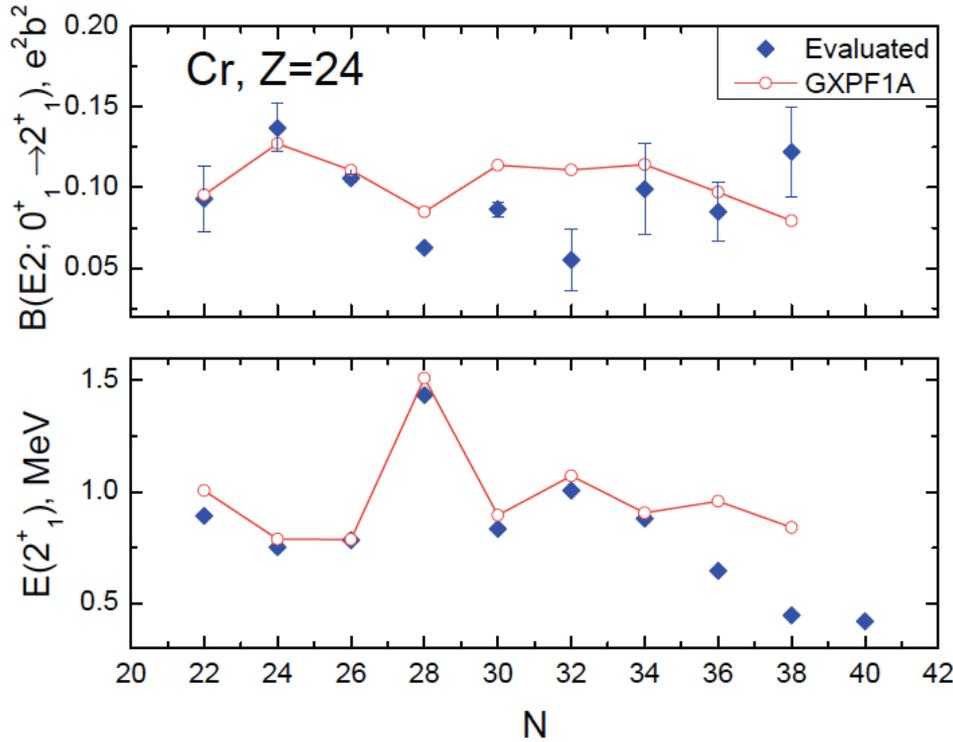
M. Horoi et al., 2006



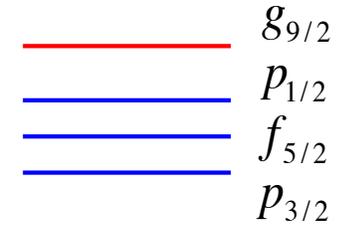
J_n	E_x (MeV)	$\langle Q \rangle_{J0}$	$(Q_0)_{sp}$	$B(E2; J \rightarrow J-2)$	$ Q_0 _{tr}$
2_3	5.342	-41.6	+145.3	413.2	144.1
4_5	6.027	-55.2	+151.9	598.0	143.8
6_3	7.556	-56.2	+140.6	609.3	139.5
8_4	9.300	-47.2	+112.2	558.4	130.5
10_6	10.782	-63.9	+147.0	591.1	132.5
12_5	13.071	-62.7	+141.1	612.3	133.7

$e_{\text{eff}}^p = 1.5 \quad e_{\text{eff}}^n = 0.5$

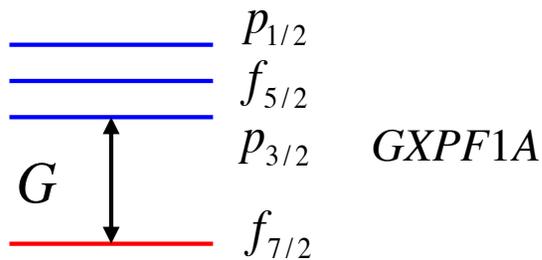
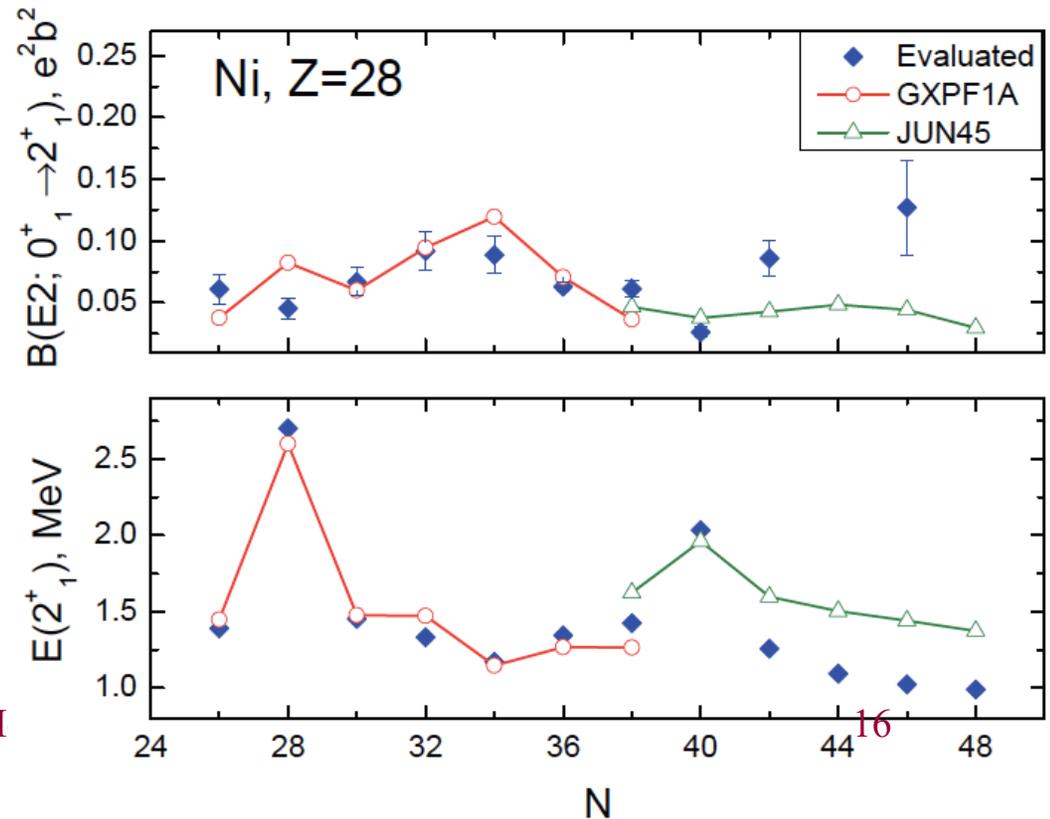
Shell-model Calculations



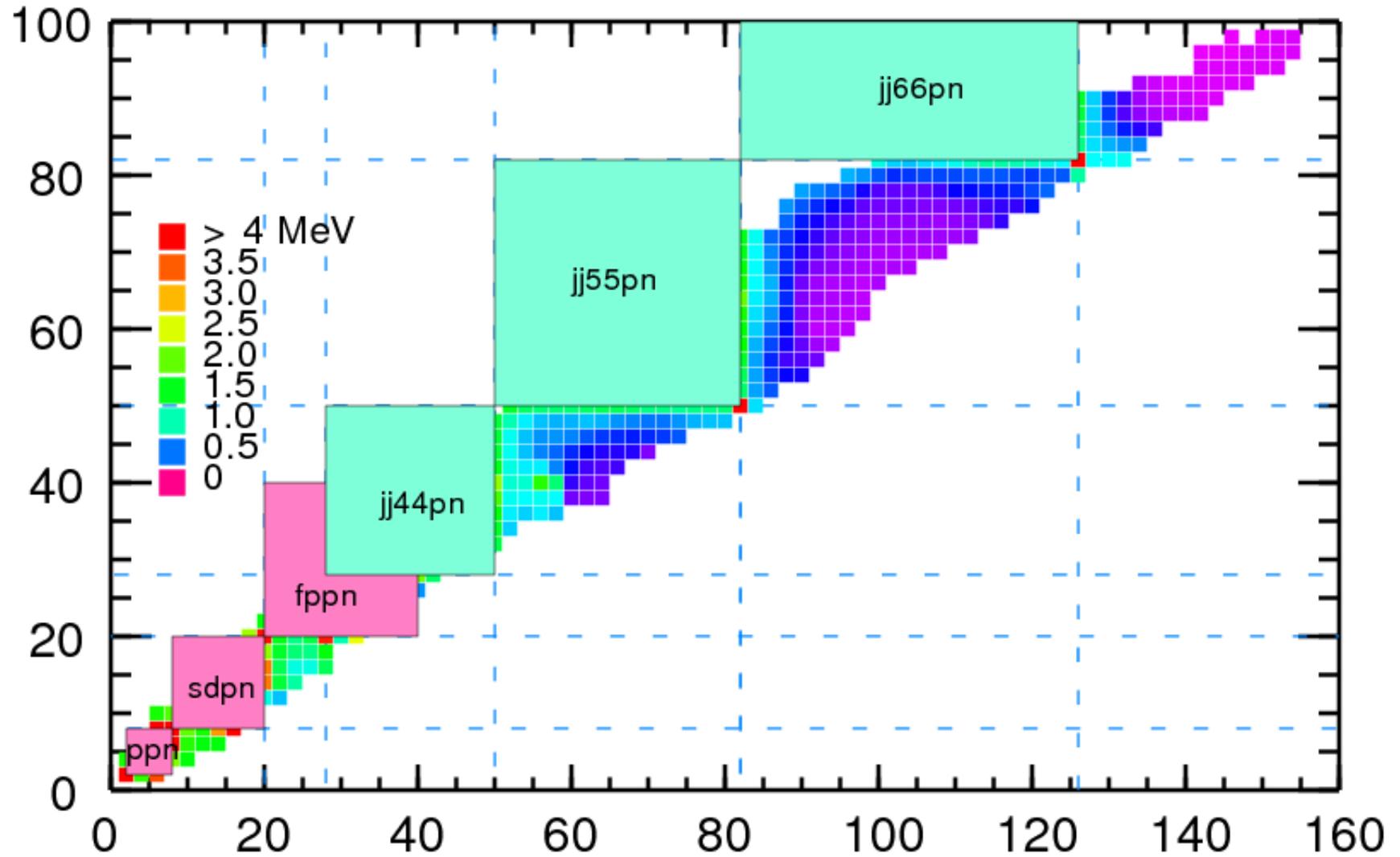
have been
 GXPF1A and
 $e_{eff}^p = 1.5$ $e_{eff}^n = 0.5$



JUN45



Shell Model Spaces



Statistical Model: why ^{136}Xe ?

- One of the longest half-life limit measured, $\sim 2 \times 10^{26}$ years
- One of the potential isotopes chosen for the next “tonne experiment”
- Relatively low shell model dimensions in the $jj55$ ($0g_{7/2}, 2s, 1d, 0h_{11/2}$) model space

Statistical Model: Shell Model Effective Hamiltonians

- Isospin conserving effective Hamiltonians in $jj55$ model space ($0g_{7/5}$, $2s1d$, $0h_{11/2}$)
- 327 TBME of starting Hamiltonians $\pm 10\%$ uniformly random changes
- Three starting effective Hamiltonians:
 - \Rightarrow SVD, PRC 86 044323 (2012)
 - \Rightarrow GCN5082, PRC 82 064304 (2010)
 - \Rightarrow $jj55t$, PRL 110 222502 (2013)

Statistical Model: Observables (24)

- 0nbb NME (1): M_{0v} , short-range correlator CD-Bonn
- 2nbb NME (1): M_{2v} , $q = 0.7$
- Energies 2+, 4+, 6+ (6): PE_{2+} , PE_{4+} , PE_{6+} , DE_{2+} , DE_{4+} , DE_{6+}
- B(GT) to 1+1 state (2): PGT, DGT, $q = 0.7$
- B(E2) $2^+ \rightarrow 0^+$ (2): PBE2, DBE2 $e_{eff}^p = 1.5$ $e_{eff}^n = 0.5$
- Proton occupation (8): POP_{g_7} , POP_{s_1} , POP_d , $POP_{h_{11}}$, DOP_{g_7} , DOP_{s_1} , DOP_d , $DOP_{h_{11}}$
- Neutron vacancies (4): DVN_{g_7} , DVN_{s_1} , DVN_d , $DVN_{h_{11}}$
- Muon Capture rate (0)
- B(M1) (0)

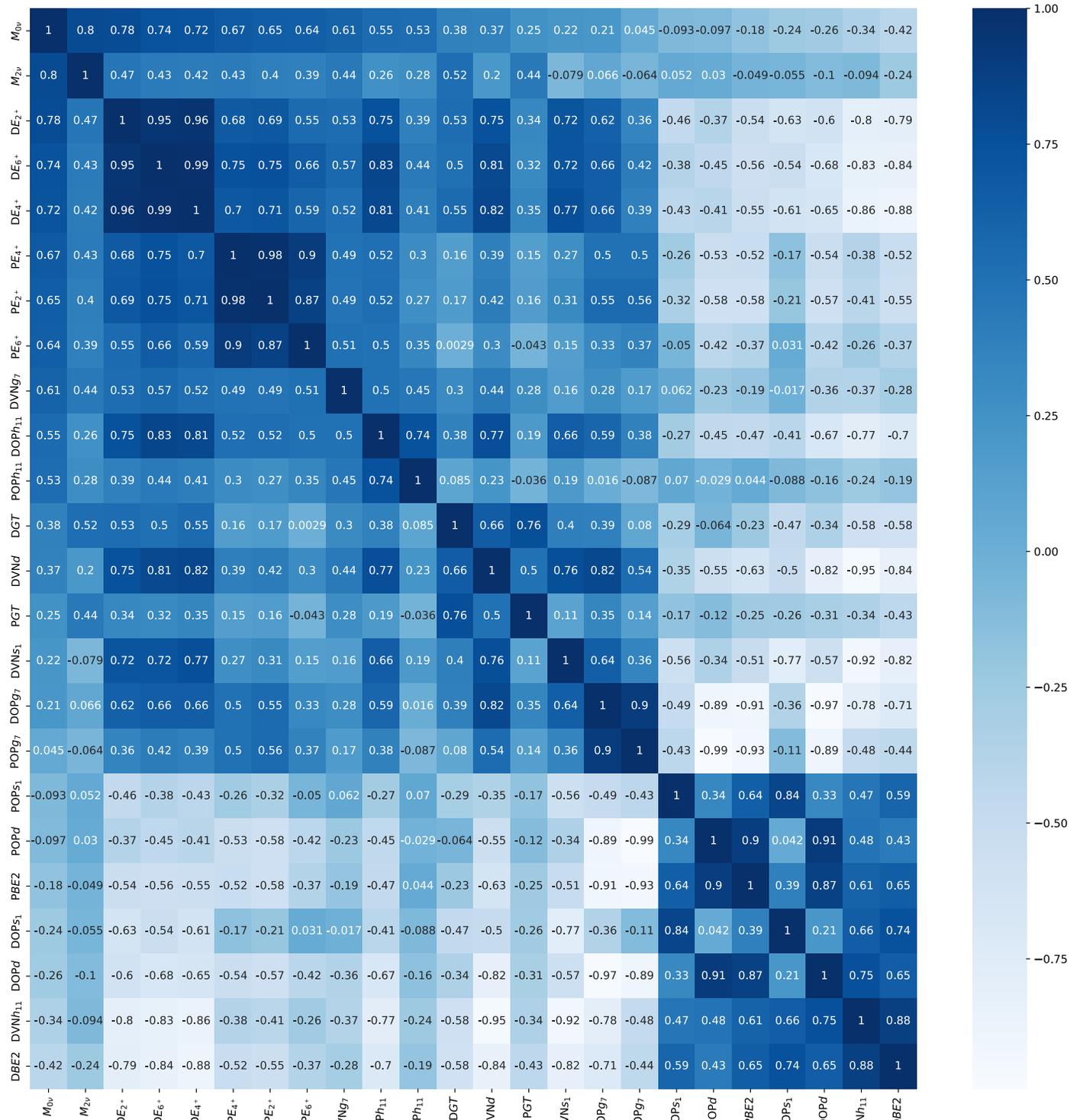
Overview

Observable	Data	Error	svd_s	gcn_s	$j5t_s$	μ_{svd}	σ_{svd}	μ_{gcn}	σ_{gcn}	μ_{j5t}	σ_{j5t}
$M_{0\nu}$	N/A	N/A	1.763	2.645	2.314	1.749	0.111	2.632	0.135	2.306	0.156
$M_{2\nu}$	0.018	0.001	0.025	0.069	0.060	0.022	0.003	0.061	0.007	0.052	0.007
PGT	0.150	0.021	0.163	0.545	0.512	0.141	0.059	0.457	0.105	0.333	0.220
$PBE2$	0.286	0.081	0.154	0.121	0.096	0.153	0.009	0.122	0.013	0.099	0.012
PE_{2+}	1.313	0.150	1.498	1.363	1.513	1.494	0.089	1.352	0.089	1.507	0.098
PE_{4+}	1.694	0.150	2.073	1.747	2.012	2.070	0.089	1.740	0.107	2.011	0.107
PE_{6+}	1.892	0.150	2.178	1.892	2.254	2.192	0.088	1.884	0.125	2.212	0.096
POP_{g7}	2.930	0.100	2.705	2.716	3.143	2.702	0.187	2.705	0.209	3.082	0.267
POP_{s1}	0.057	0.006	0.089	0.025	0.020	0.090	0.018	0.025	0.006	0.021	0.006
POP_{h11}	0.400	0.040	0.190	0.375	0.265	0.189	0.020	0.373	0.050	0.265	0.045
POP_d	0.520	0.030	1.016	0.884	0.572	1.019	0.180	0.896	0.197	0.632	0.250
DGT	0.012	0.005	0.001	0.009	0.004	0.001	0.000	0.008	0.003	0.003	0.003
$DBE2$	0.413	0.011	0.342	0.194	0.158	0.337	0.023	0.195	0.026	0.163	0.028
DE_{2+}	0.819	0.150	0.662	0.842	0.917	0.660	0.067	0.836	0.056	0.919	0.049
DE_{4+}	1.867	0.150	1.389	1.873	2.113	1.403	0.131	1.861	0.116	2.087	0.082
DE_{6+}	2.207	0.150	2.157	2.196	2.502	2.171	0.151	2.197	0.090	2.507	0.117
DVN_{g7}	0.000	0.150	0.102	0.174	0.130	0.100	0.010	0.172	0.014	0.132	0.023
DVN_{s1}	0.080	0.020	0.271	0.251	0.415	0.286	0.117	0.255	0.058	0.407	0.110
DVN_{h11}	1.680	0.130	1.205	0.726	0.347	1.177	0.237	0.724	0.132	0.385	0.162
DVN_d	0.240	0.050	0.423	0.850	1.108	0.437	0.132	0.850	0.118	1.076	0.158
DOP_{g7}	3.860	0.100	3.189	3.475	4.145	3.187	0.209	3.477	0.249	4.078	0.436
DOP_{s1}	0.200	0.020	0.263	0.083	0.049	0.264	0.047	0.084	0.020	0.052	0.017
DOP_{h11}	0.620	0.060	0.264	0.658	0.625	0.269	0.049	0.658	0.093	0.613	0.121
DOP_d	1.290	0.080	2.285	1.785	1.181	2.280	0.227	1.781	0.265	1.258	0.447

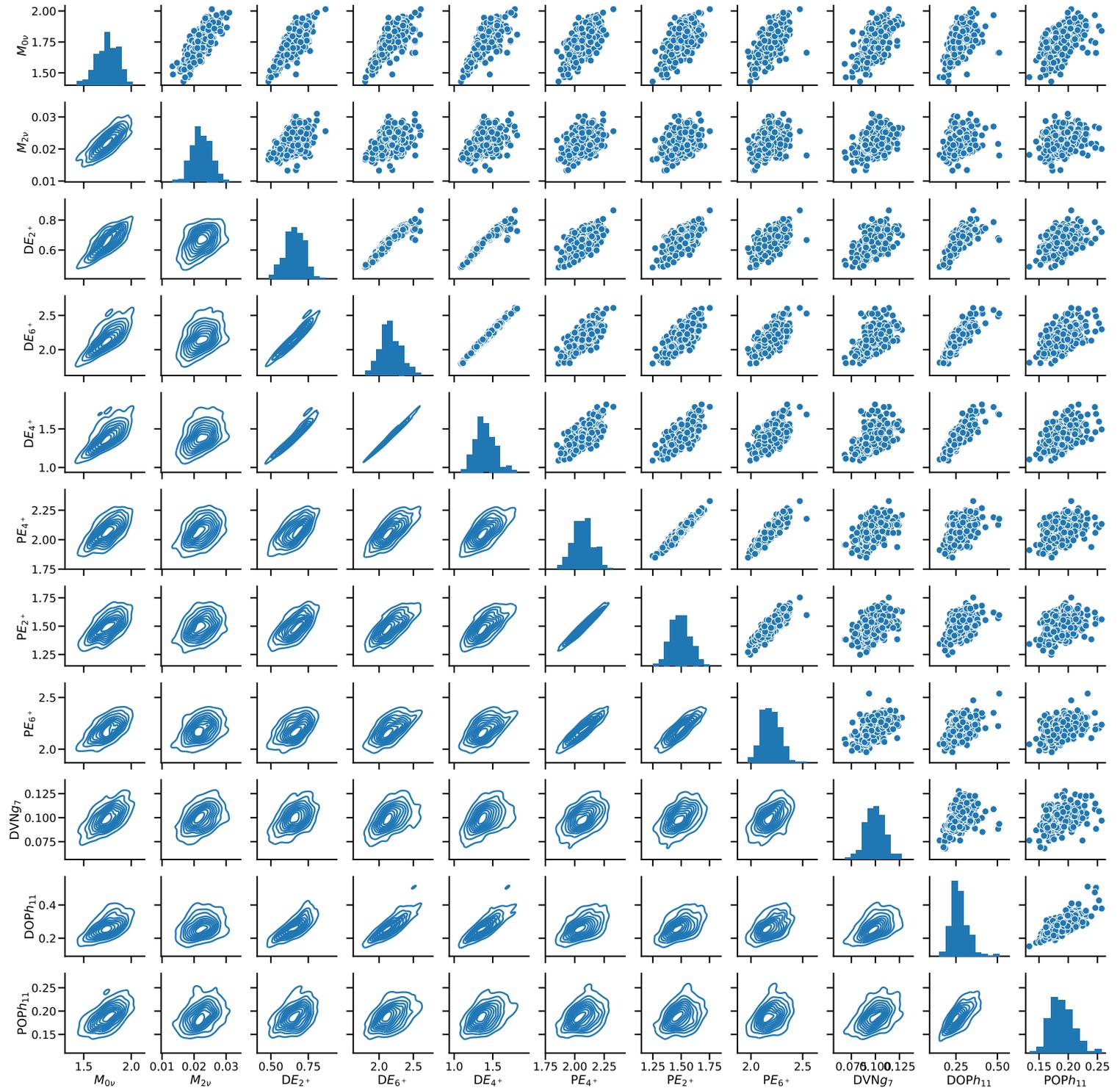
Table I. All relevant data and statistics for all selected observables.

Statistical Model: Heat Map

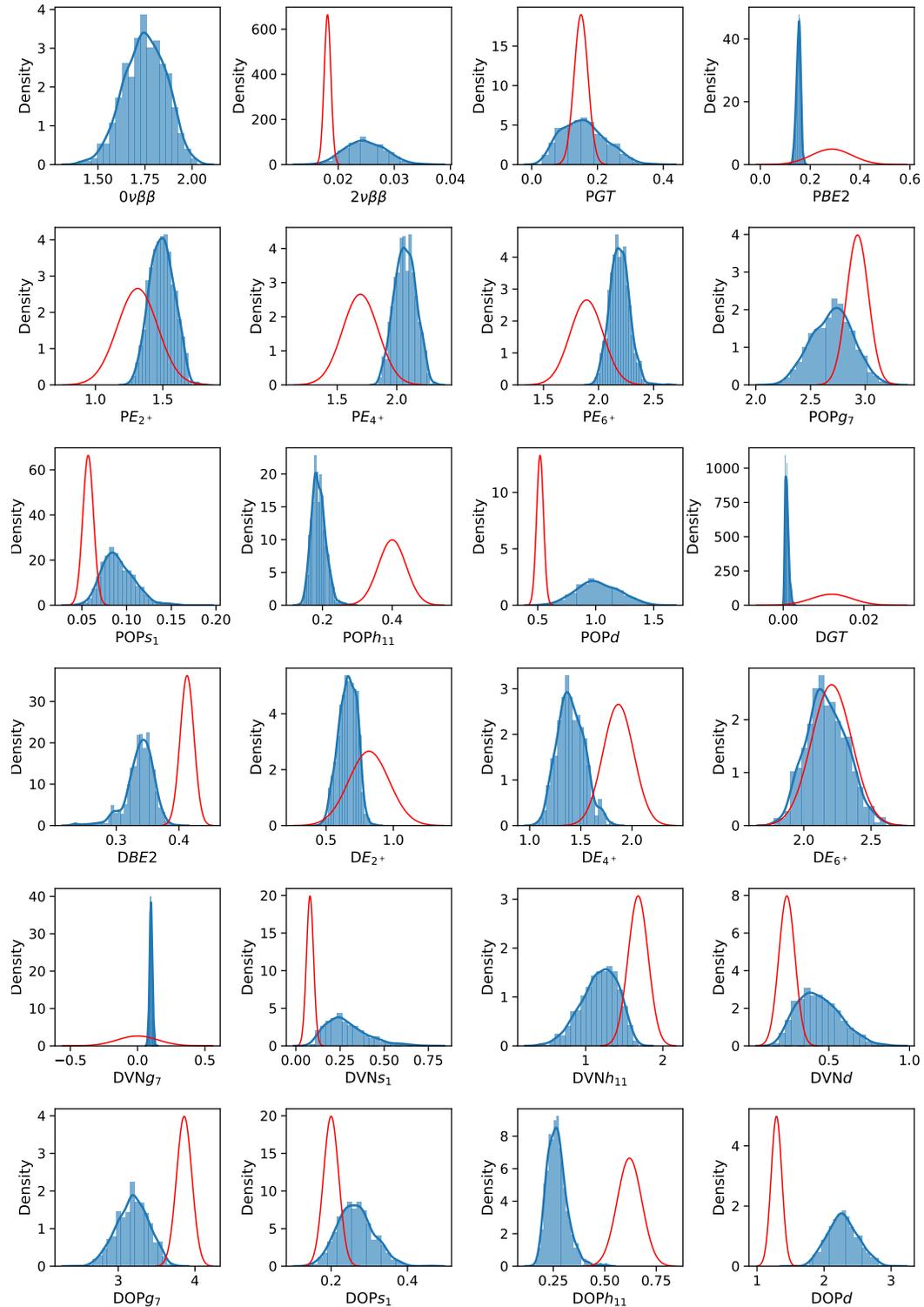
samples=1,000



Statistical Model:
Correlations Map



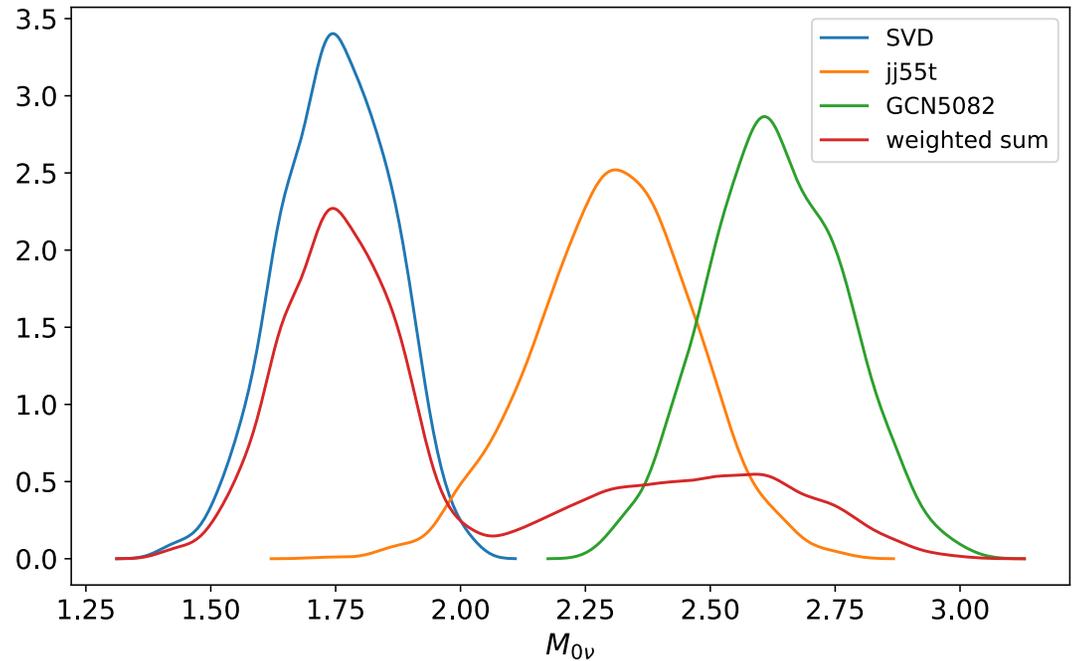
Statistical Model: Density Distributions



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Statistical Model: joint distribution

$$P(x = M_{0\nu}) = W_{svd}P_{svd}(x) + W_{gcn}P_{gcn}(x) + W_{j5t}P_{j5t}(x),$$



Statistical Model: Bayesian Model Averaging

$$p(x = M_{0\nu} | y_e, \sigma_e) = \sum_{k=svd}^{j5t} p(x = M_{0\nu} | y_e, \sigma_e, \mathcal{M}_k) p(\mathcal{M}_k | y_e, \sigma_e),$$

$$M_{0\nu} = 1.99 \pm 0.37$$

$$p(\mathcal{M}_k | y_e, \sigma_e) = \frac{p(y_e, \sigma_e | \mathcal{M}_k) \pi(\mathcal{M}_k)}{\sum_{k=svd}^{j5t} p(y_e, \sigma_e | \mathcal{M}_k) \pi(\mathcal{M}_k)}$$



$$p(y_e, \sigma_e | \mathcal{M}_k) = \int \prod_i^{N_{obs}} dy_i p_{y_e, \sigma_e}(y_i) \left[\int \prod_j^{N_{tbme}} d\theta_j p(y_i | \theta_j, \mathcal{M}_k) \pi(\theta_j | \mathcal{M}_k) \right]$$

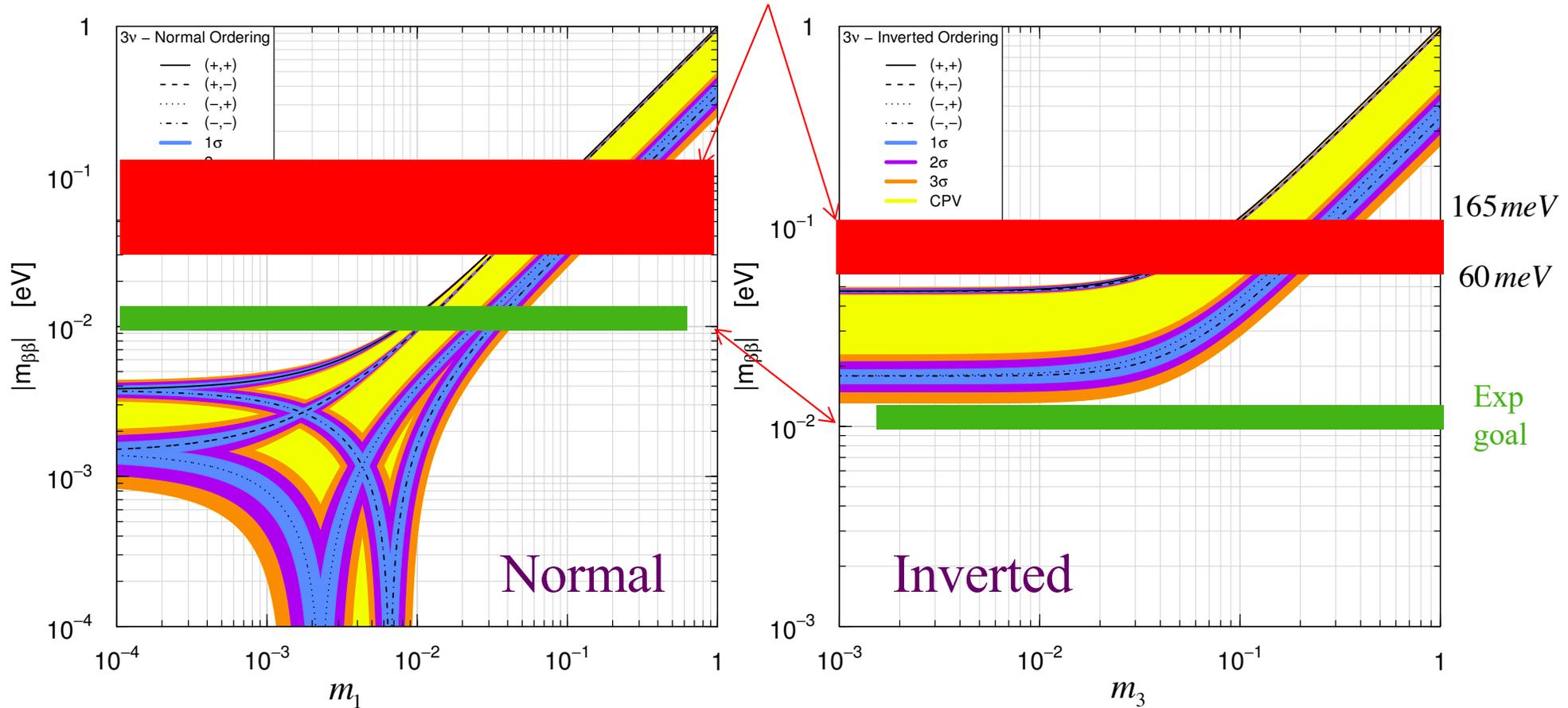
Evidence integrals

$$p_{y_e, \sigma_e}(y_i) \propto \prod_i^{N_{obs}} \exp[-(y_i - y_{e_i})^2 / (2\sigma_{e_i}^2)]$$

Neutrino $\beta\beta$ effective mass

arxiv:1507.08204

KamLAND – Zen, PRL 117, 082503 (2016): ^{136}Xe



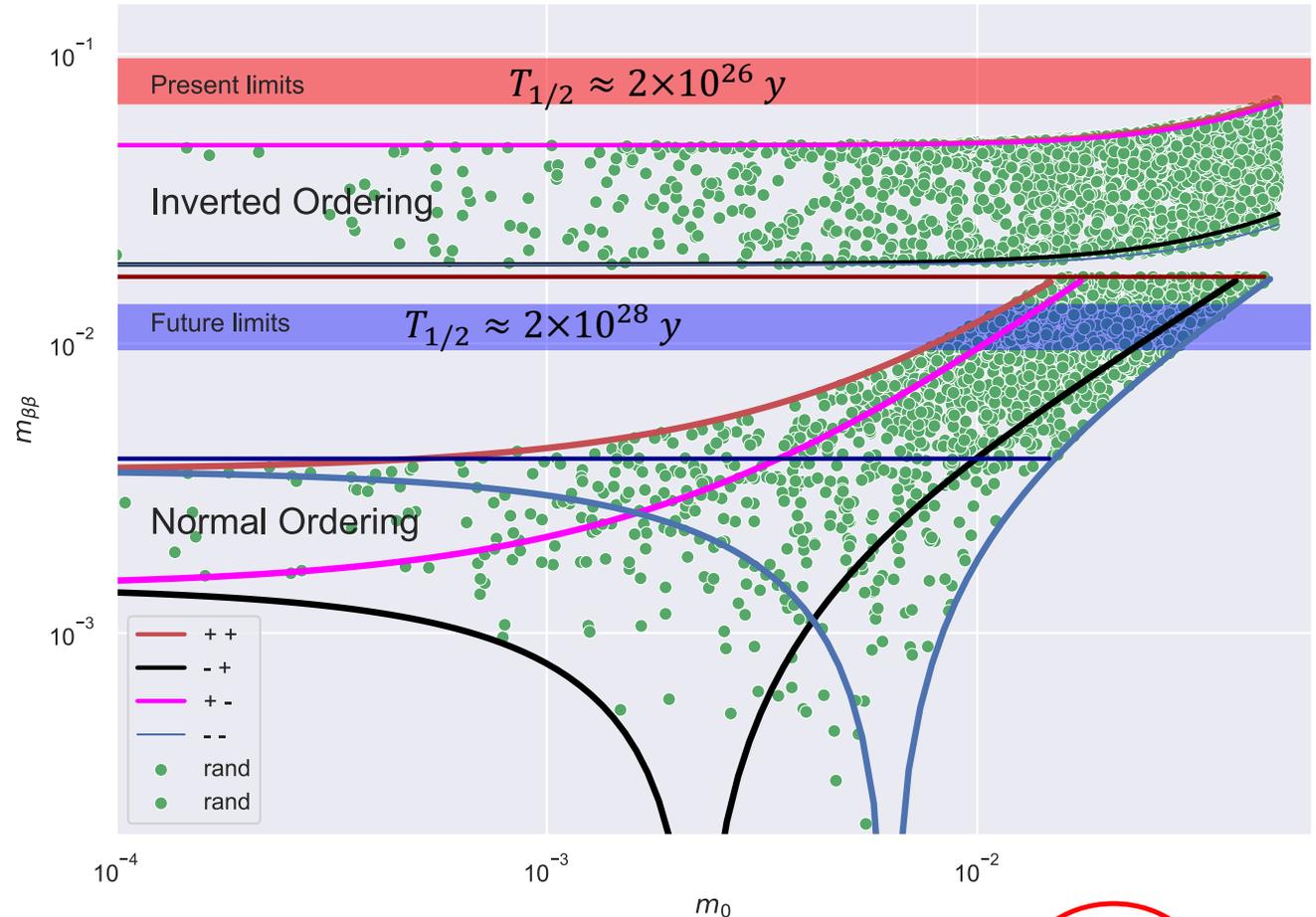
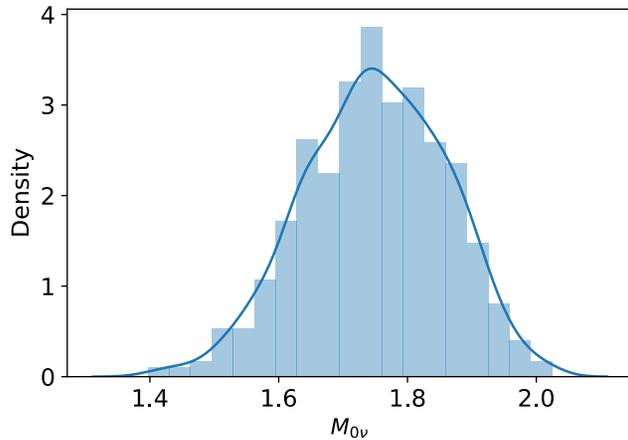
$$|m_{\beta\beta}| = \left| \sum_{k=1}^3 m_k U_{ek}^2 \right| = \left| c_{12}^2 c_{13}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3} \right|$$

$$\phi_2 = \alpha_2 - \alpha_1 \quad \phi_3 = -\alpha_1 - 2\delta$$

$$\Leftrightarrow T_{1/2}^{-1}(0\nu) = G^{0\nu} (Q_{\beta\beta}) \left[M^{0\nu}(0^+) \right]^2 (\eta_{0\nu})^2$$

$$\eta_{0\nu} = \frac{|m_{\beta\beta}|}{m_e}$$

Mass mechanism interference effects



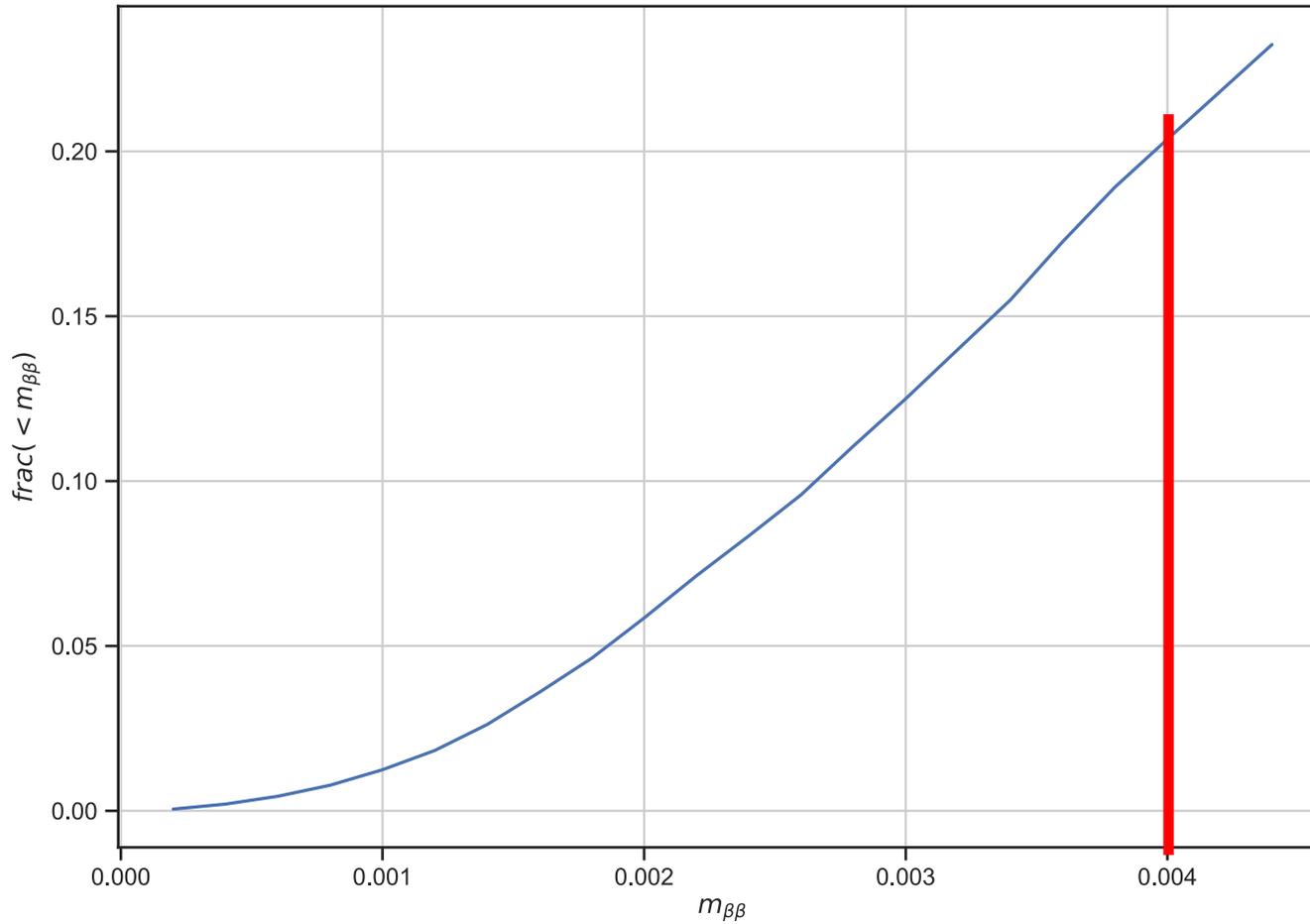
$$|m_{\beta\beta}| = \left| \sum_{k=1}^3 m_k U_{ek}^2 \right| = \left| c_{12}^2 c_{13}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3} \right|$$

$$M^{0\nu} = M_{GT}^{0\nu} - \left(\frac{g_V}{g_A} \right)^2 M_F^{0\nu} + M_T^{0\nu}$$

$$\Leftrightarrow T_{1/2}^{-1}(0\nu) = G^{0\nu} (Q_{\beta\beta}) [M^{0\nu}(0^+)]^2 (\eta_{0\nu})^2$$

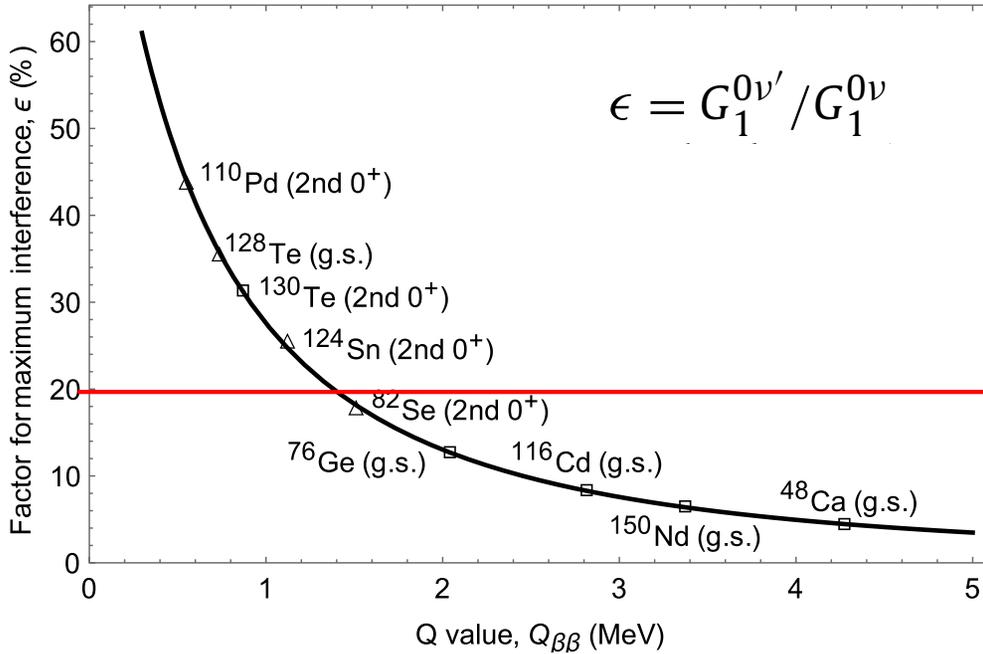
$$\eta_{0\nu} = \frac{|m_{\beta\beta}|}{m_e}$$

Mass mechanism interference effects



LRSM interference effects

$$\left[T_{1/2}^{0\nu} \right]^{-1} \approx 2g_A^4 G_1^{0\nu} |\eta|^2 |M|^2 [1 - \epsilon \cos \phi]$$



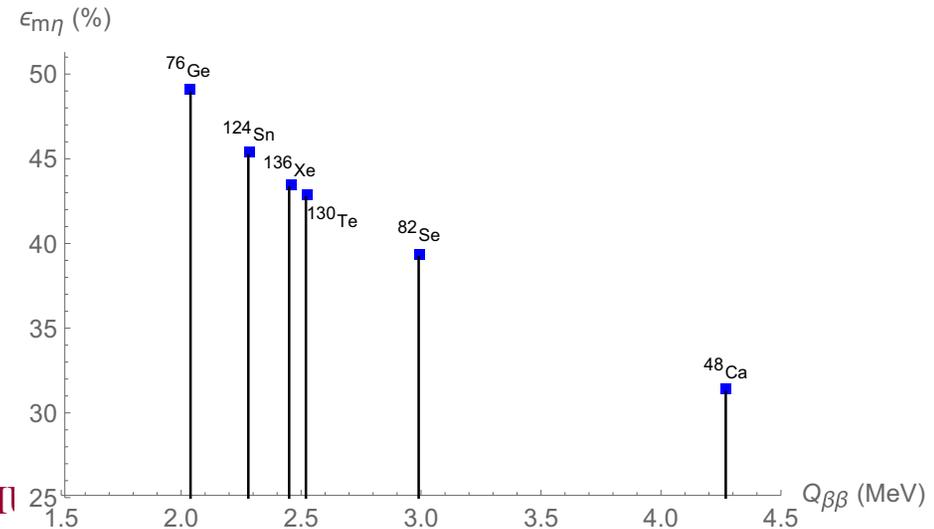
Interference between mass mechanism and heavy neutrino mechanism: F. Ahmed, A. Neacsu, and M. Horoi, Phys. Lett. B 769, 299 (2017).

In LRSM:

$$\begin{aligned} [T_{1/2}^{0\nu}]^{-1} = & |\mathcal{M}_{GT}^{0\nu}|^2 \left[C_m |\eta_m|^2 + C_N |\eta_N|^2 + C_\lambda |\eta_\lambda|^2 + C_\eta |\eta_\eta|^2 \right. \\ & + C_{mN} |\eta_m| |\eta_N| \cos(\phi_m - \phi_N) + C_{m\lambda} |\eta_m| |\eta_\lambda| \cos(\phi_m - \phi_\lambda) \\ & + C_{m\eta} |\eta_m| |\eta_\eta| \cos(\phi_m - \phi_\eta) + C_{N\lambda} |\eta_N| |\eta_\lambda| \cos(\phi_N - \phi_\lambda) \\ & \left. + C_{N\eta} |\eta_N| |\eta_\eta| \cos(\phi_N - \phi_\eta) + C_{\lambda\eta} |\eta_\lambda| |\eta_\eta| \cos(\phi_\lambda - \phi_\eta) \right] \end{aligned}$$

Interference between mass mechanism and eta mechanism: F. Ahmed, and M. Horoi, PRC 101, 035504 (2020).

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M. Horoi CMI

FIG. 5. Coefficient of maximum interference $\epsilon_{m\eta}(1)$ plotted against $Q_{\beta\beta}$ values.

LRSM interference effects

$$[T_{1/2}^{0\nu}]^{-1} \simeq g_A^4 [C_i |\eta_i|^2 + C_j |\eta_j|^2 + C_{ij} |\eta_i| |\eta_j| \cos(\phi_i - \phi_j)],$$

$$C_{m\eta} = \left[M_{GT} - \left(\frac{g_V}{g_A} \right)^2 M_F + M_T \right] \\ \times [G_{03} \mathcal{M}_{2+} - G_{04} \mathcal{M}_{1-} - G_{05} M_P + G_{06} M_R]$$

Phase space factors

$$C_j |\eta_j|^2 = \alpha C_i |\eta_i|^2 \Rightarrow |\eta_j| = \sqrt{\alpha \frac{C_i}{C_j}} |\eta_i|.$$

$$\epsilon_{ij}(\alpha) = \frac{\sqrt{\alpha}}{1 + \alpha} \frac{|C_{ij}|}{\sqrt{|C_i| |C_j|}}$$

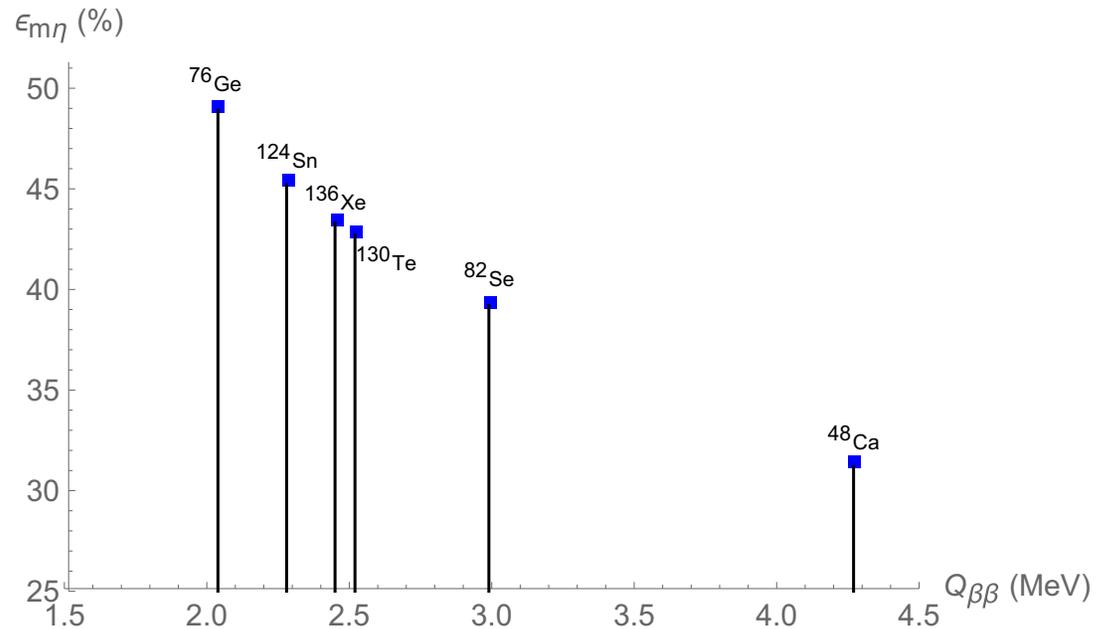


FIG. 5. Coefficient of maximum interference $\epsilon_{m\eta}(1)$ plotted against $Q_{\beta\beta}$ values.

$$\alpha = 1$$

Statistical Model: Results

- Shell model observables show robustness relative to small random changes of the effective Hamiltonian
- $M_{0\nu}$ shows significant correlations with about half of the observables considered
- The outcome of our **Bayesian Averaging model** is a joint distribution of the ^{136}Xe $M_{0\nu}$ with a range of 1.55 – 2.65 at 90% confidence level
- We propose a $M_{0\nu}$ NME of 1.99 ± 0.37 for ^{136}Xe

Future work:

- Statistical investigation of the maximal interference effects
- Investigate other random changes, e.g. 15%, GOA, etc