# PDF implications and potential of PVDIS at JLab

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27 June – 1 July 2022

**PVDIS and EW Physics at JLab12** 

#### **INT-Seattle**

### i **PVDIS** and **PDF uncertainties**, flavor dependence

- 1: precision BSM searches limited by (incomplete) proton structure info
  - → many standard-candle HEP measurements PDF-limited ... including very high x ...
  - → taming PDF dependence: knowledge of hard-to-access phase-space regions
  - $\rightarrow$  PDF studies central to NP QCD at JLab12, maps of hadron structure
- 2: closely related: flavor separation generally needs numerous expts
  - $\rightarrow$  controlling PDF uncertainties requires knowledge of flavor dependence
  - $\rightarrow$  complementary measurements/analyses for  $d/u, \ \bar{d}/\bar{u}, \ s, \ \cdots$

highlight through: i general status of HEP PDFs; ii flavor dependence issues

### iii conclusion(s): implications for EW physics and PVDIS

from NNLO analyses, state-of-the-art predictions for fundamental LHC observables → *e.g.*, total cross sections at 14 TeV

CT18 NNLO, PRD 103 (2021) 1



 significant PDF-driven uncertainties; also, systematic effects: W cross sections sensitive to inclusion of 2016 7 TeV ATLAS inclusive W/Z data

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 $\rightarrow$  these include  $\sigma_H$ ,  $\sin^2 \theta_W$ ,  $m_W$ , ...

| ATLAS, 1701.07240   |                              |               |              |               | for example:   |               |             |            |             |               |
|---------------------|------------------------------|---------------|--------------|---------------|----------------|---------------|-------------|------------|-------------|---------------|
| Channel             | $m_{W^+} - m_{W^-}$<br>[MeV] | Stat.<br>Unc. | Muon<br>Unc. | Elec.<br>Unc. | Recoil<br>Unc. | Bckg.<br>Unc. | QCD<br>Unc. | EW<br>Unc. | PDF<br>Unc. | Total<br>Unc. |
| $W \rightarrow e v$ | -29.7                        | 17.5          | 0.0          | 4.9           | 0.9            | 5.4           | 0.5         | 0.0        | 24.1        | 30.7          |
| $W \to \mu \nu$     | -28.6                        | 16.3          | 11.7         | 0.0           | 1.1            | 5.0           | 0.4         | 0.0        | 26.0        | 33.2          |
| Combined            | -29.2                        | 12.8          | 3.3          | 4.1           | 1.0            | 4.5           | 0.4         | 0.0        | 23.9        | 28.0          |
|                     |                              |               |              |               |                |               |             |            |             |               |

 $\rightarrow$  the PDF uncertainty can be a/the dominant uncertainty!

 $\rightarrow$  frontier efforts at the HL-LHC aim for (sub)percent precision

 $\rightarrow$  large cross-cutting effort spanning theory/expt to improve

- heightened theory accuracy (HO, power corrections)
- novel measurements (JLab12, EIC, LHC, vA)
- generator development Snowmass21, Campbell et al.: 2203.11110

Snowmass21, Amoroso et al.: 2203.13923

 $\rightarrow$  driven by marriage of latest theory, high-energy hadronic data

$$\sigma(AB \to W/Z + X) = \sum_{n} \alpha_{s}^{n} \sum_{a,b} \int dx_{a} dx_{b} f_{a/A}(x_{a}, \mu^{2}) \hat{\sigma}_{ab \to W/Z + X}^{(n)}(\hat{s}, \mu^{2}) f_{b/B}(x_{b}, \mu^{2})$$

$$T.-J. \text{ Hou et al., PRD 103 (2021) 1. } Contemporary NNLO QCD fits$$

$$I.1_{0} \int_{0.9}^{0.9} \int_{0.8}^{0.6} \frac{10^{4}}{10^{4}} \frac{10^{3}}{10^{2}} \frac{10^{-1}}{10^{-1}} \frac{10^{-1}}{0.2} \frac{10^{-1}}{0.5} \frac{10^{-1}}{0.9} \frac{10^{-1}}{0.5} \frac{10^{-1}}{0.9} \frac{10^{-1}}{0.5} \frac{10^{-1}}{0.9} \frac{10^{-1}}{10^{-1}} \frac{10^{-1}}{0.5} \frac{10^{-1}}{0.5} \frac{10^{-1}}{0.9} \frac{10^{-1}}{10^{-1}} \frac{10^{-1}}{0.5} \frac{1$$

periodic benchmarking (PDF4LHC21) valuable to cross-check treatment of data

 $\rightarrow$  seek methodological independence in identifying data-driven PDF features

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theory ingredients  $\rightarrow$  <u>higher pQCD accuracy</u>

- future analyses will witness an interplay between pQCD & other dynamics
- NNLO+ necessary to stabilize scale uncertainties; especially over wide scales



 $\rightarrow$  essential to assess PDF dependence in parallel with HO corrections (beyond LO QPM)

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### electroweak (EW) corrections also vital

• at  $\mathcal{O}(\alpha_s^2)$  accuracy, EW corrections and explicit  $\gamma(x, \mu^2)$  needed

Xie, TJH, Hou, Schmidt, Yan, Yuan: PRD105 (2022) 5, 054006

• important for high-energy LHC processes: *e.g.*, 13 TeV W+H production



TeV-scale NLO EW corrections dominated (60%) by single-photon (PDF) contributions

→ requires **delicate** treatment along with QCD perturbative effects

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### i necessary for electroweak precision: photon PDF

- at  $\mathcal{O}(\alpha_s^2)$  accuracy, EW corrections and explicit  $\gamma(x, \mu^2)$  needed

Xie, TJH, Hou, Schmidt, Yan, Yuan: PRD105 (2022) 5, 054006

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following CT14QED, CT18QED now interfaces LUX formalism

 $\mathbf{X}$ 

$$x\gamma(x,\mu^{2}) = \frac{1}{2\pi\alpha(\mu^{2})} \int_{x}^{1} \frac{z}{z} \left\{ \int_{\frac{x^{2}m_{p}^{2}}{1-z}}^{\frac{\mu^{2}}{1-z}} \frac{Q^{2}}{Q^{2}} \alpha_{ph}^{2}(-Q^{2}) \left[ \left( zp_{\gamma q}(z) + \frac{2x^{2}m_{p}^{2}}{Q^{2}} \right) F_{2}(x/z,Q^{2}) - z^{2}F_{L}(x/z,Q^{2}) \right] - \alpha^{2}(\mu^{2})z^{2}F_{2}(x/z,\mu^{2}) \right\} + \mathcal{O}(\alpha^{2},\alpha\alpha_{s})$$

<u>depends on nonperturbative inputs</u> [kinematical cuts alone can't avoid this]



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## i parametrization uncertainty: nonperturbative fitting forms

- initial PDFs still not generally calculable through rigorous QCD at  $Q = Q_0 = m_c$  (to the needed precision...)
  - $\rightarrow$  subject to complex nonperturbative dynamics
  - $\rightarrow$  practice agnosticism w.r.t. initial parametrization

(some guidance from QCD, QCD-inspired models)

 $\rightarrow$  explore model uncertainty with many forms





### ii high-x PDFs remain dominated by large uncertainties

• much of the JLab12 (PVDIS) sensitivity probes high-x (low-W, Q) PDFs

 $\rightarrow$  data tensions

- PDF (Hessian) uncertainties enlarge dramatically in high-*x* limit
  - $\rightarrow$  limited data  $\rightarrow$  extrapolation



flavor separation also challenging; parsing quark sea: various precision data

• PVDIS  $\gamma$ -Z interference further complements standard data sets

### perturbative order ( $\alpha_s$ ), interplay of data sets relevant for high-x d/u



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0.35 0.36 0.37 0.38 0.39 0.4

d/u(x = 0.3, Q = 100 GeV)

۲ ۲ ۳ 10 - 160 HERA I+II

---- 101 BCDMS  $F_2^{pro}$ ---- 281 D02W  $A_{ch}^{e}$ 

---- 250 LHCb8ZW

- 249 CMS8W  $A^{\mu}_{ch}$ 

----- 245 LHCb7ZW

0.41

- influences extrapolation region behavior
- competing pulls of fitted data at high-*x* also restrict precision; *e.g.*,
  - $\rightarrow$  BCDMS,  $F_2^d$
  - $\rightarrow$  LHCb, W/Z 7 TeV

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**ii** 

### ii light-nuclear corrections: **flavor dependence**, high-*x* PDFs

• *d*-PDF information from deuteron scattering; nuclear corrections relevant

$$f^d(x,Q^2) = \int \frac{dz}{z} \int dp_N^2 \,\mathcal{S}^{N/d}(z,p_N^2) \,\widetilde{f}^N(x/z,p_N^2,Q^2)$$

CJ-CT: Accardi, TJH, Jing, Nadolsky: EPJC81 (2021) 7, 603

• corrections are generally ~percent-level, but can become larger, especially at <u>high x</u>





 impacts LHC observables; necessary for high precision

### ii nuclear modifications of increasing importance with greater PDF precision

• example: nuclear effects in deuteron influence *d*-quark PDF ( $x \ge 0.1$ )

CJ-CT: Accardi, TJH, Jing, Nadolsky: EPJC81 (2021) 7, 603

• comparative study: additional effects at lower-*x*; gluon impacts (from flavor correlations)



subtle interplay with nuclear corrections in large-A targets

 $\rightarrow$  demands more attention in future (free nucleon) PDF analyses

<sup>[</sup>will revisit]

pulls of fitted data on d/u influenced by deuteron corrections

- $L_2$  sensitivity: quickly assess  $\Delta \chi^2$  from upward shift of PDFs by  $1\sigma$  uncertainty
  - $\rightarrow$  useful for apples-apples comparisons, especially among Hessian fits

https://ct.hepforge.org/PDFs/ct18/figures/L2Sensitivity/



• in CJ fits *without* deuteron correction (left), large tensions at high x:  $\gamma$ -jet [1] vs. DIS-deuterium [4] data; especially at x > 0.5

 $\rightarrow$  with fixed correction (right), tensions relieved significantly

CJ-CT: Accardi, TJH, Jing, Nadolsky: EPJC81 (2021) 7, 603

### ii extracting high-x dependence in PDF fits

- high-x PDFs, ratios [e.g., d/u] connected to details of proton WF
- behavior at  $x \to 1$  an important nonpert. discriminator
- CT18, parametrize  $f_{a/A}(x, Q_0^2) = x^{A_{1,a}}(1-x)^{A_{2a}} \times \Phi_a(x)$



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Courtoy and Nadolsky, PRD103, 054029 (2021)

ii high-*x* sea-quark asymmetries: SeaQuest (E906) and STAR Drell-Yan data

see talk: A. Tadepalli

- potential overlap between PVDIS and PDF sensitivities of other high-*x* expts  $\bar{d} > \bar{u}$
- recently-released SeaQuest data have moderate sensitivity to high-x



SeaQuest prefers larger dbar/ubar at x > 0.2; may *somewhat* reduce strange-PDF uncertainty

- interplays: treatment of nucl. data (left); inclusion of STAR *W*-prod., E866 expts. (right)
  - → E906, STAR constructively enhance dbar/ubar; removing E866 augments this effect

### ii high-*x* sea-quark asymmetries: SeaQuest (E906) and STAR Drell-Yan data

- potential overlap between PVDIS and PDF sensitivities of other high-x expts  $d > \bar{u}$
- $L_2$  sensitivities: excluding E866, minimal tensions between E906, other sets



- forthcoming study: exploration of PDF pulls in post-E906 fits...
- $\rightarrow$  E906, STAR constructively enhance dbar/ubar; removing E866 augments this effect

#### ii nucleon strangeness from neutrino and Drell-Yan data

- strangeness PDF, breaking of flavor SU(3) remain highly uncertain
- CT18: indications of tension between neutrino (need nuclear corrections), precise LHC W/Z data

CT18Z NNLO, s(x, 2 GeV)



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#### explored in PDF4LHC21 benchmarking studies, especially for reduced PDF fits



[again, see talk: T. Cridge]

dominant experiments exert qualitatively similar pulls on strangeness in CT, MSHT

NuTeV neutrino scattering; 2016 ATLAS W/Z; HERA combined DIS data; ...

- beyond few-body systems, CT, other analyses use heavy-nuclei for flavor separation
- requires knowledge of nuclear corrections; these directly fitted by nPDF analyses
  - $\rightarrow$  better control over *x*, *A* dependence can benefit nucleon PDF extractions



Muzakka, Duwentäster, TJH et al., 2204.13157

• more systematic studies of interdependence of nucleon, nuclear analyses needed

 $\rightarrow$  high-*x* PVDIS from proton, nuclei helpful to building this picture

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### ii nonperturbative theory developments: lattice QCD inputs

recent years: progress in *ab initio* hadron-structure calculations from LQCD
 → quasi-PDFs, pseudo-PDFs, quasi-TMDs, ...

there are be important synergies between PDF fitting and lattice QCD

[overlaps with PVDIS]

 lattice data can potentially inform high-x behavior of quark sea



TJH, Wang, Nadolsky, Olness, PRD100 (2019) 9, 094040

u-d at  $\mu_F$  = 3GeV

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## iii nonperturbative the PVDIS and HEP PDF fits CD inputs

#### have illustrated: PDF flavor dependence, including high x, is an HEP frontier



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### achieving highest (PDF) impact of PVDIS: mastery of small Q

TJH and Melnitchouk, PRD77, 114023 (2008)

 $\rightarrow \gamma$ -Z interference accesses unique flavor currents in nucleon

$$A^{\rm PV} = -\left(\frac{G_F Q^2}{4\sqrt{2}\pi\alpha}\right) \left[g_A^e Y_1 \frac{F_1^{\gamma Z}}{F_1^{\gamma}} + \frac{g_V^e}{2} Y_3 \frac{F_3^{\gamma Z}}{F_1^{\gamma}}\right] \quad A^{\rm PV} \text{ potentially subject to finite-}Q^2 \text{ corrections}$$

$$Y_{1} = \frac{1 + (1 - y)^{2} - y^{2}(1 - r^{2}/(1 + R^{\gamma Z})) - xyM/E}{1 + (1 - y)^{2} - y^{2}(1 - r^{2}/(1 + R^{\gamma})) - xyM/E} \left(\frac{1 + R^{\gamma Z}}{1 + R^{\gamma}}\right) \qquad \qquad R = \sigma_{L}/\sigma_{T} \neq 1$$

$$(= 1, \text{ Callan-Gross})$$

 in principle, could complicate PDF sensitivity of PV asymmetry

- effectively, proxy for various low-Q<sup>2</sup> corrections which must be investigated/controlled
- substantial theory, phenomenological progress over intervening years



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### iii PVDIS will have subtle relationship with HTs, TMCs, ...

Brady, Accardi, TJH, and Melnitchouk, PRD84, 074008 (2011)

[see talks: A. Accardi, S. Li, N. Sato, ...]

• at lowest  $Q^2$ , TMCs (and HTs) can represent multi-percent **proton** A<sup>PV</sup> corrections

 $\rightarrow$  as precision effect, ~comparable to deuteron corrections shown earlier



- <u>TMC prescription dependence mild</u>; deuteron A<sup>PV</sup> (mostly) insulated from effects
  - → BSM, SMEFT analyses based on deuteron have much weaker PDF dependence [see talk: K. Simsek]
  - $\rightarrow$  A<sup>PV</sup> deuteron more sensitive to charge-symmetry violation (complement EIC)
  - → TMCs more relevant for CSV; consequential for  $\gamma(x)$ , CT EW phenomenology

PDF4LHC21 benchmarking: 2203.05506

- MC sampling of high-*x* PDFs can sometimes produce irregularities
  - $\rightarrow$  *e.g.*, positive-definiteness not always guaranteed for  $x \rightarrow 1$



• high-x PVDIS data: perhaps explore PDF uncertainties with representative sampling  $\rightarrow$  discriminating power Courtoy, Huston Nadolsky, Xie, Yan, Yuan: 2205.10444

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## conclusions

- HEP analyses increasingly confront problem of taming high-*x* uncertainties
  - $\rightarrow$  limit to EW precision, sensitivity of BSM tests
  - $\rightarrow$  closely related to challenge of flavor separation
- QCD analysis progress: mix of theory corrections, pheno considerations
   [higher-order QCD, EW; light, heavy nuclear effects; power-suppressed corrections; ...]

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- PVDIS, EW measurements at JLab12 hold intriguing possibilities
  - → significant PDF (proton), BSM/EW (deuteron) sensitivity
  - → extending HEP fits past kinematical cuts nontrivial but may be worthwhile meanwhile, could play valuable benchmarking role
    - $\rightarrow$  clear complementarity to EIC (both PV and positron) measurements

EIC YR, 2103.05419

supplementary material -

#### strangeness (and other) PDF pulls have modest tolerance dependence

CT & MSHT groups, in preparation.



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### quantifying PDF preferences of fitted data: the L, method

 $S_{f,L_2}(E)$ : fast approximation of the Lagrange Multiplier scan of  $\chi^2_E$ along direction of  $f_a(x_i, Q_q)$ .

> $\rightarrow$  estimated  $\Delta \chi_E^2$  for experiment E when  $f_a(x_i, Q_i)$  increases by its +68% c.l. PDF uncertainty

$$Y = \chi_E^2 \quad X = f_a(x_i, Q_i) \qquad S_{f,L_2} \equiv \Delta Y(\vec{z}_{m,X}) = \vec{\nabla} Y \cdot \vec{z}_{m,X}$$
  
CT18 NNLO, BcdF2dCor (102), Q=2 GeV  

$$= \vec{\nabla} Y \cdot \frac{\vec{\nabla} X}{|\vec{\nabla} X|} = \Delta Y \cos \varphi$$
  
extension of L<sub>1</sub> sensitivity (PDFSense)  
method used to explore  
• HEP data pulls for CT18  
Phys.Rev.D 103 (2021) 1, 014013  
• PDF-Lattice sensitivities  
Phys.Rev.D 100 (2019) 9, 094040  
• EIC potential

EIC YR, arXiv: 2103.05419

flavor structure only from inclusive data is challenging!

![](_page_29_Figure_1.jpeg)

note PDFs' different orders-of-mag.!

NC DIS: sensitivity to d-type quarks  $\frac{1}{4}$  that of u-type

$$\sigma \propto \frac{4}{9}(u_+ + c_+) + \frac{1}{9}(d_+ + s_+ + b_+)$$

CC DIS: lower accuracy (1/10 lumi.)

 $\rightarrow u$ -quark dominates

 $\rightarrow d$ -quark  $\frac{1}{2}$  of u, but harder to access in NC DIS (above)

 $\rightarrow \bar{d} + \bar{u} ~\sim~ {\rm few}~ {\rm percent}~ {\rm of}~ u$ 

→ for x~0.1,  

$$s \approx \bar{s} \approx \bar{d} - \bar{u} < 0.1(\bar{d} + \bar{u})$$

 $\rightarrow$  at x>0.5, no separation for  $\bar{u}, \bar{d}, \bar{s}$ 

### the electroweak sector and New Physics searches at EIC

- if measured to sufficient precision, the quark-level electroweak couplings may be sensitive to an extended EW sector, e.g.,  $Z^\prime$ 

$$\mathcal{L}^{\mathrm{PV}} = \frac{G_F}{\sqrt{2}} \left[ \bar{e} \gamma^{\mu} \gamma_5 e \left( C_{1u} \bar{u} \gamma_{\mu} u + C_{1d} \bar{d} \gamma_{\mu} d \right) + \bar{e} \gamma^{\mu} e \left( C_{2u} \bar{u} \gamma_{\mu} \gamma_5 u + C_{2d} \bar{d} \gamma_{\mu} \gamma_5 d \right) \right]$$

$$C_{1u} = -\frac{1}{2} + \frac{4}{3}\sin^2\theta_W$$

![](_page_30_Figure_4.jpeg)

 a unique strength of an EIC is its combination of very high precision and beam polarization, which allows the observation of parity-violating helicity asymmetries:

$$A^{\rm PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \quad (R/L : e^- \text{ beam helicities})$$

selects  $\gamma$ -Z interference diagrams!

TJH and Melnitchouk, PRD**77**, 114023 (2008).

$$A^{\rm PV} = -\left(\frac{G_F Q^2}{4\sqrt{2}\pi\alpha}\right) (Y_1 \ a_1 \ + \ Y_3 \ a_3)$$
$$a_1 = \frac{2\sum_q e_q \ C_{1q} \ (q+\bar{q})}{\sum_q e_q^2 \ (q+\bar{q})} \qquad a_3 = \frac{2\sum_q e_q \ C_{2q} \ (q-\bar{q})}{\sum_q e_q^2 \ (q+\bar{q})}$$

### the electroweak sector and New Physics searches at EIC

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$$C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W$$

- with sufficient precision, an EIC (which will be statistics-limited in these measurements) can extract  $\sin^2 \theta_W$ 
  - this measurement is potentially sensitive to the TeV-scale in a complementary fashion to energy-frontier searches!

![](_page_31_Figure_5.jpeg)