

Terrestrial and Astrophysical Probes of Early-Universe Relics

M.A.F., P. W. Graham, S. Kalia, D. F. Jackson Kimball. Phys. Rev. D **104**, 075023 (2021) [[arXiv:2106.00022](https://arxiv.org/abs/2106.00022)].

M.A.F., P. W. Graham, S. Kalia, D. F. Jackson Kimball. Phys. Rev. D **104**, 095032 (2021) [[arXiv:2108.08852](https://arxiv.org/abs/2108.08852)].

A. Arza, M.A.F., P. W. Graham, S. Kalia, D. F. Jackson Kimball. Phys. Rev. D **105**, 095007 (2022) [[arXiv:2112.09620](https://arxiv.org/abs/2112.09620)].

M.A.F., P.W. Graham, and S. Rajendran. Phys. Rev. D **101**, 115021 (2020) [[arXiv:1911.08883](https://arxiv.org/abs/1911.08883)]

Dark Matter in Compact Objects, Stars, and in Low Energy Experiments

University of Washington Program INT-22-2B

Virtual talk

August 30, 2022

Michael A. Fedderke

mfedderke@jhu.edu
mfedderke.com



Overall plan of my talk

I will try to cover two quite different topics within the theme of the program:

- ❖ Part I: A new signal of ultralight dark-photon and axion-like particle dark matter in terrestrial magnetometer data

M.A.F., P. W. Graham, S. Kalia, D. F. Jackson Kimball. Phys. Rev. D **104**, 075023 (2021) [arXiv:2106.00022].

M.A.F., P. W. Graham, S. Kalia, D. F. Jackson Kimball. Phys. Rev. D **104**, 095032 (2021) [arXiv:2108.08852].

A. Arza, M.A.F., P. W. Graham, S. Kalia, D. F. Jackson Kimball. Phys. Rev. D **105**, 095007 (2022) [arXiv:2112.09620].

- ❖ Part II: Using white dwarfs to constrain a different early-universe relic: charged massive particles (CHAMPs)

M.A.F., P.W. Graham, and S. Rajendran. Phys. Rev. D **101**, 115021 (2020) [arXiv:1911.08883]



Earth as a Transducer for Dark- Matter Detection

Dark Matter in ~~Compact Objects, Stars, and in~~ Low Energy Experiments

Outline - Part I

- ❖ Motivation
 - Ultralight Bosonic Dark Matter (dark photons and ALPs / axions)
- ❖ Phenomenology of the photon—dark-photon or photon—axion system
- ❖ A new signal of DM:
 - Magnetic field oscillating coherently & in-phase across the entire surface of Earth, with a known spatial and vectorial pattern**
- ❖ Results of a search with SuperMAG data
- ❖ Outlook and future directions

Dark Matter

- ❖ ~27% of the energy budget of the universe
- ❖ Want to detect non-gravitationally!
- ❖ A menagerie of possible DM candidates exist
- ❖ Candidates span ~80 orders of magnitude in mass:

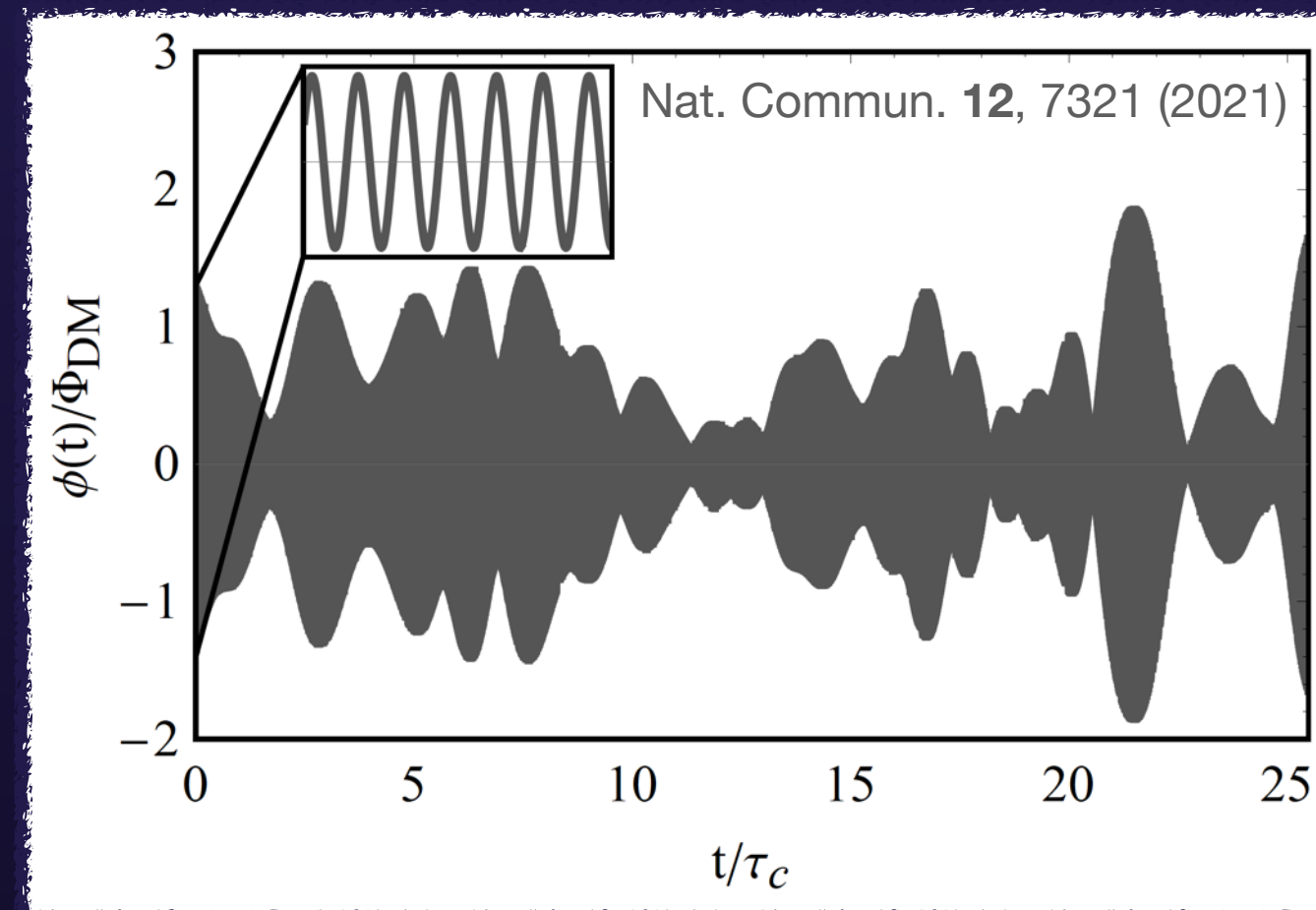
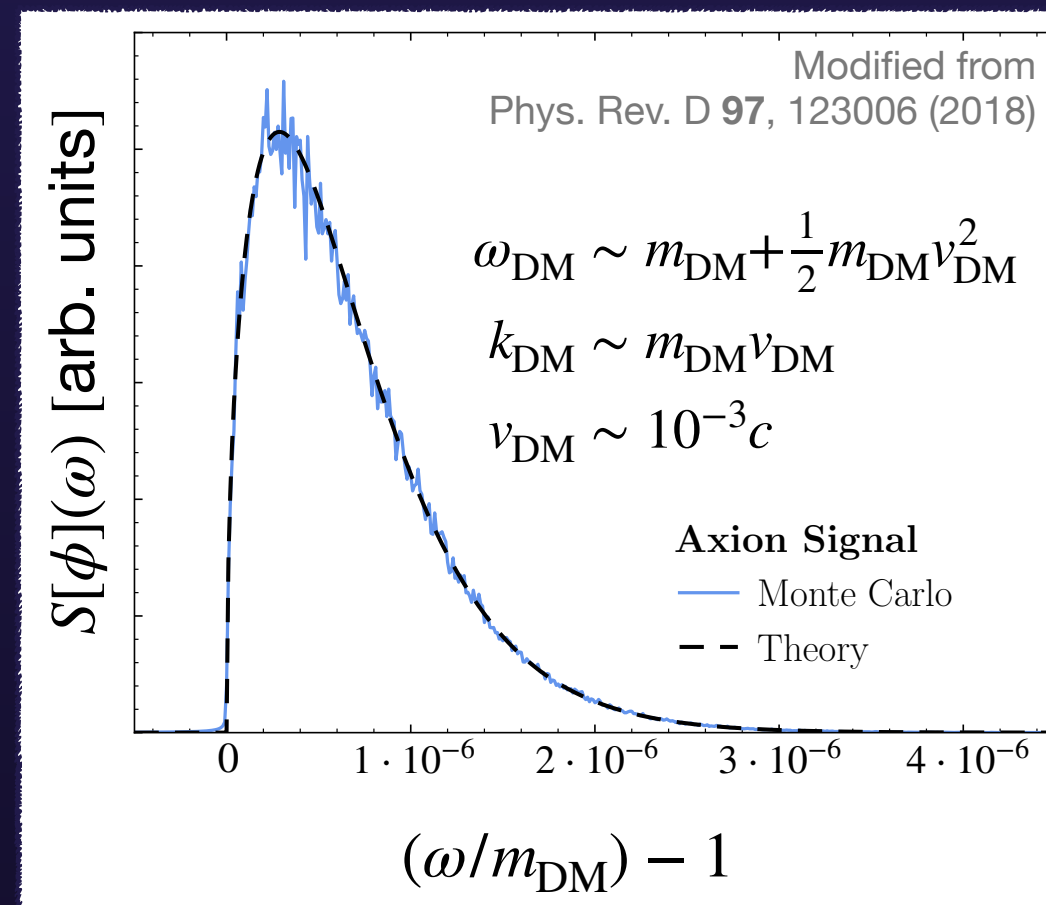


- ❖ No single experimental approach can cover all candidates
- ❖ This talk: ultralight bosonic DM around $m_{\text{DM}} \sim 10^{-17} \text{ eV}$

Ultralight Bosonic Dark Matter I

- ❖ For $m_{\text{DM}} \lesssim 10 \text{ eV}$, overlapping DM particle de Broglie wavelengths
- ❖ Classical field description

$$\phi(\mathbf{x}, t) \sim \frac{\sqrt{\rho_{\text{DM}}}}{m_{\text{DM}}} \cos(\omega_{\text{DM}} t - \mathbf{k}_{\text{DM}} \cdot \mathbf{x}) \sim \frac{\sqrt{\rho_{\text{DM}}}}{m_{\text{DM}}} \cos(m_{\text{DM}} t)$$



$$T_{\text{coh}} \sim v_{\text{DM}}^{-2} T_{\text{osc}} \sim 10^6 T_{\text{osc}}$$

$$\lambda_{\text{coh}} \sim v_{\text{DM}}^{-1} \lambda_{\text{Compton}} \sim 10^3 \lambda_{\text{Compton}}$$

- ❖ Search strategies can **leverage oscillatory phenomenology**

Ultralight Bosonic Dark Matter II

Kinetically Mixed Dark Photon Dark Matter (DPDM)

Phys. Lett. B 166 (1986) 196–198; Phys. Rev. D 84 (2011) 103501 [arXiv:1105.2812]

$$\mathcal{L} = \mathcal{L}_{\text{SM}} - \frac{1}{4}(F')^2 + \frac{1}{2}m_{A'}^2(A')^2 + \frac{\epsilon}{2}FF' \quad (\epsilon \ll 1)$$

$$F_{\mu\nu}^{(\prime)} \equiv \partial_{\mu}A_{\nu}^{(\prime)} - \partial_{\nu}A_{\mu}^{(\prime)}$$

$$\tilde{F}^{\mu\nu} \equiv \frac{1}{2}\epsilon^{\mu\nu\alpha\beta}F_{\alpha\beta}$$

Electromagnetically Coupled Axions

Phys. Rev. Lett. 38 (1977) 1440–1443; Phys. Rev. Lett. 40 (1978) 223–226; Phys. Rev. Lett. 40 (1978) 279–282.
JHEP 06 (2006) 051 [arXiv:hep-th/0605206]; Phys. Rev. D 81 (2010) 123530 [arXiv:0905.4720]

Note: ϵ is dimensionless, $g_{\phi\gamma}$ has dimensions of $(\text{mass})^{-1}$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{2}(\partial_{\mu}\phi)^2 - \frac{1}{2}m_{\phi}^2\phi^2 - g_{\phi\gamma}\phi F\tilde{F}$$

Disclaimer: in this talk “axions” means “axion-like particles (ALPs)”

Mixings with SM EM sector **cause observable EM effects:**

- ▶ oscillatory with period set by fundamental physics parameter m_{DM}
- ▶ narrowband / phase-coherent over long times
- ▶ spatially uniform / phase-coherent over large distances

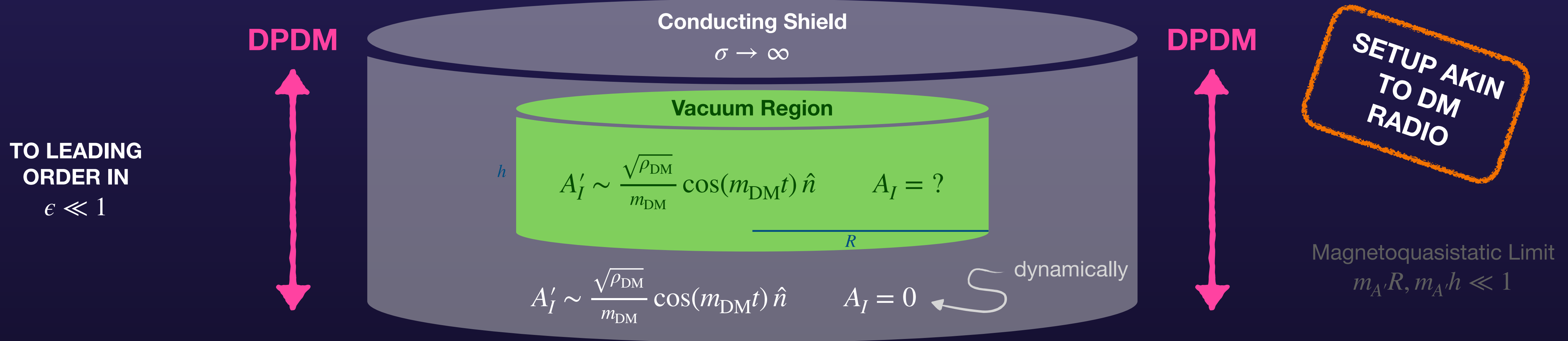
DPDM Pheno I

❖ Focus on $\mathcal{L} = \mathcal{L}_{\text{SM}} - \frac{1}{4}(F')^2 + \frac{1}{2}m_A^2(A')^2 + \frac{\epsilon}{2}FF'$ ($\epsilon \ll 1$)

❖ Same physics, different basis:

Interaction basis $\mathcal{L} \supset -\frac{1}{4}(F_I)^2 - \frac{1}{4}(F'_I)^2 + \frac{1}{2}m_A^2(A'_I + \epsilon A_I)^2 - J_{\text{EM}}A_I$

A_I couples to charges.
 A'_I does not (sterile state).



❖ Detectable field A_I in the vacuum region? Expand mass term: $\mathcal{L} \supset - (J_{\text{EM}} - \epsilon m_A^2 A'_I) \cdot A_I$

❖ DPDM field A'_I acts as oscillating background current sourcing EM: $J_{\text{eff}}^\mu \sim -\epsilon m_A^2 (A'_I)^\mu$.

❖ Non-relativistic limit: $\mathbf{J}_{\text{eff}} \sim -\epsilon m_A^2 \mathbf{A}'_I \sim -\epsilon m_A \sqrt{\rho_{\text{DM}}} \cos(m_A t) \hat{\mathbf{A}}'_I$.

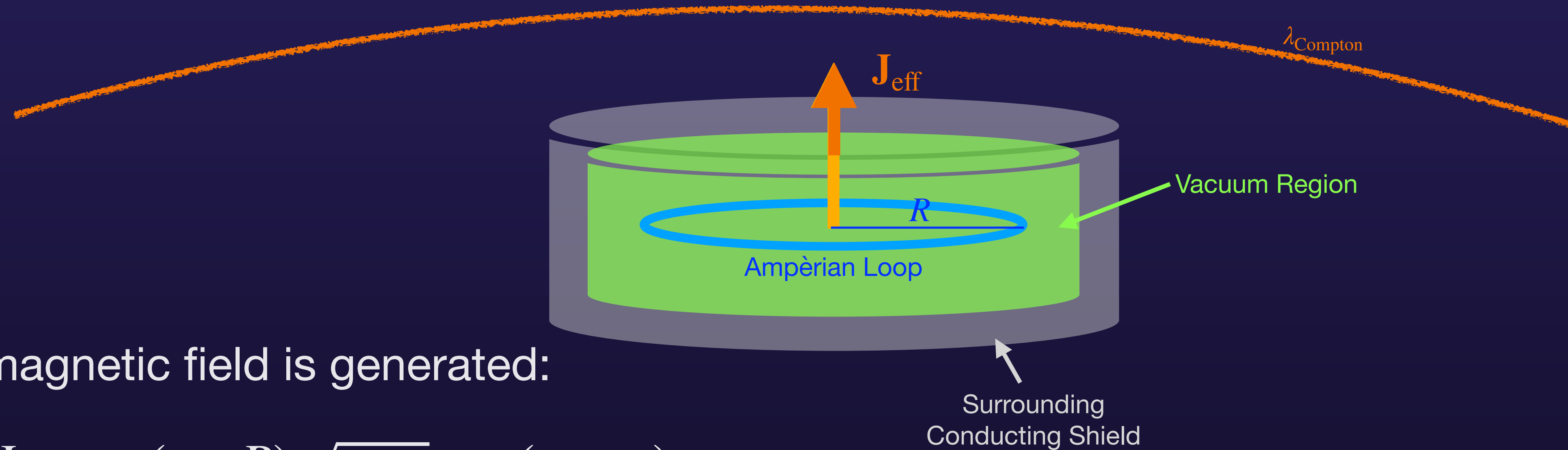
* or region of high plasma frequency

DPDM Pheno II

- ❖ Apply the Ampère-Maxwell Law to this simple geometry

$$\iint \mathbf{J} \cdot d\mathbf{A} = \oint \mathbf{B} \cdot d\mathbf{l} \Rightarrow R^2 J_{\text{eff}} \sim BR$$

Displacement
current term is
higher-order



- ❖ A circumferential magnetic field is generated:

$$B_{\phi}(r \sim R) \sim RJ_{\text{eff}} \sim \varepsilon(m_A, R) \sqrt{\rho_{\text{DM}}} \cos(m_{\text{DM}} t)$$

- ❖ Suppressed by a ratio of length scales: $m_A R \sim R/\lambda_{\text{Compton}}$

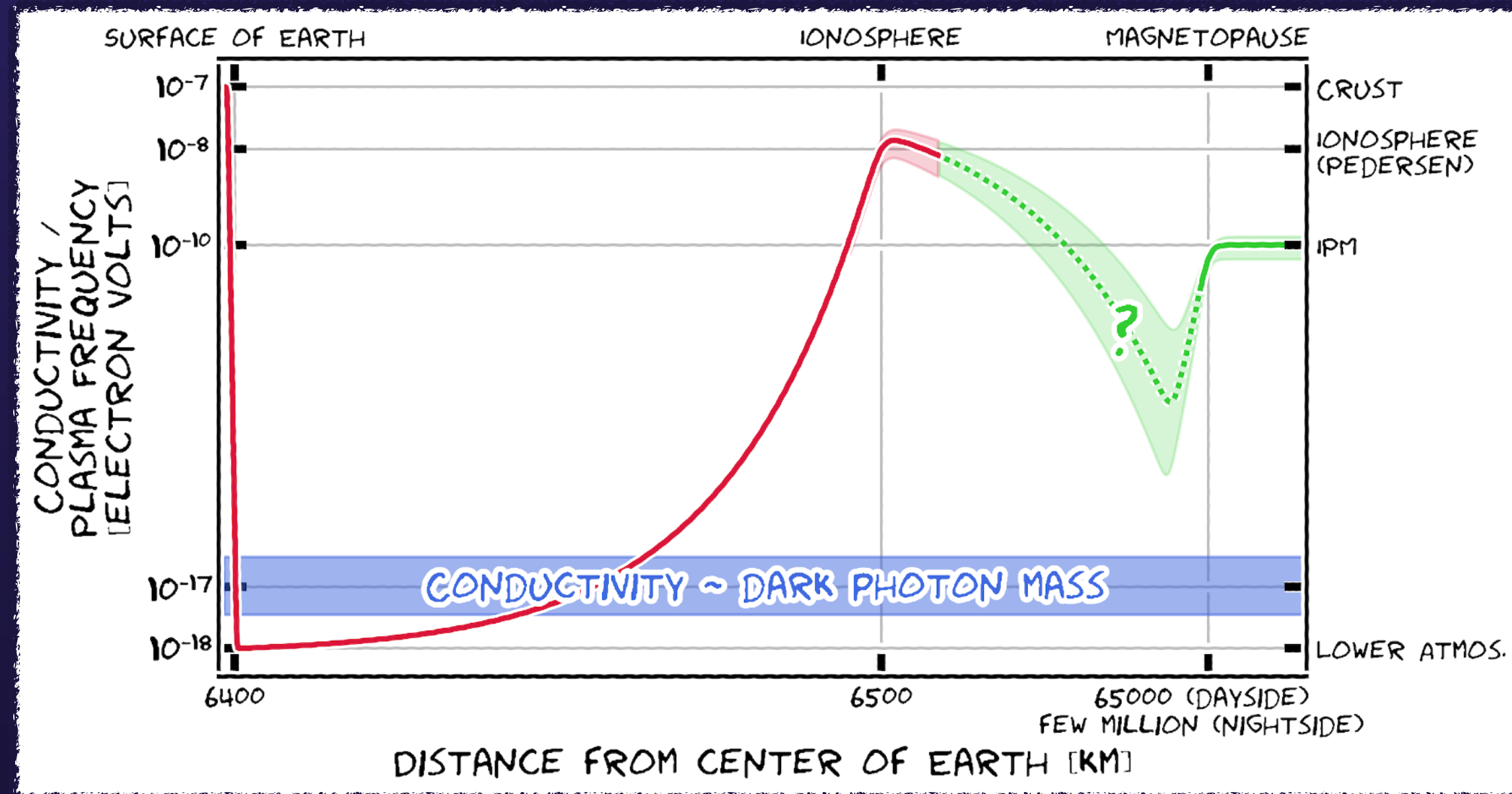
- ❖ Electric field further suppressed by $(m_A R)^2$, $v_{\text{DM}}(m_A R)$

“You’re gonna need a bigger boat”

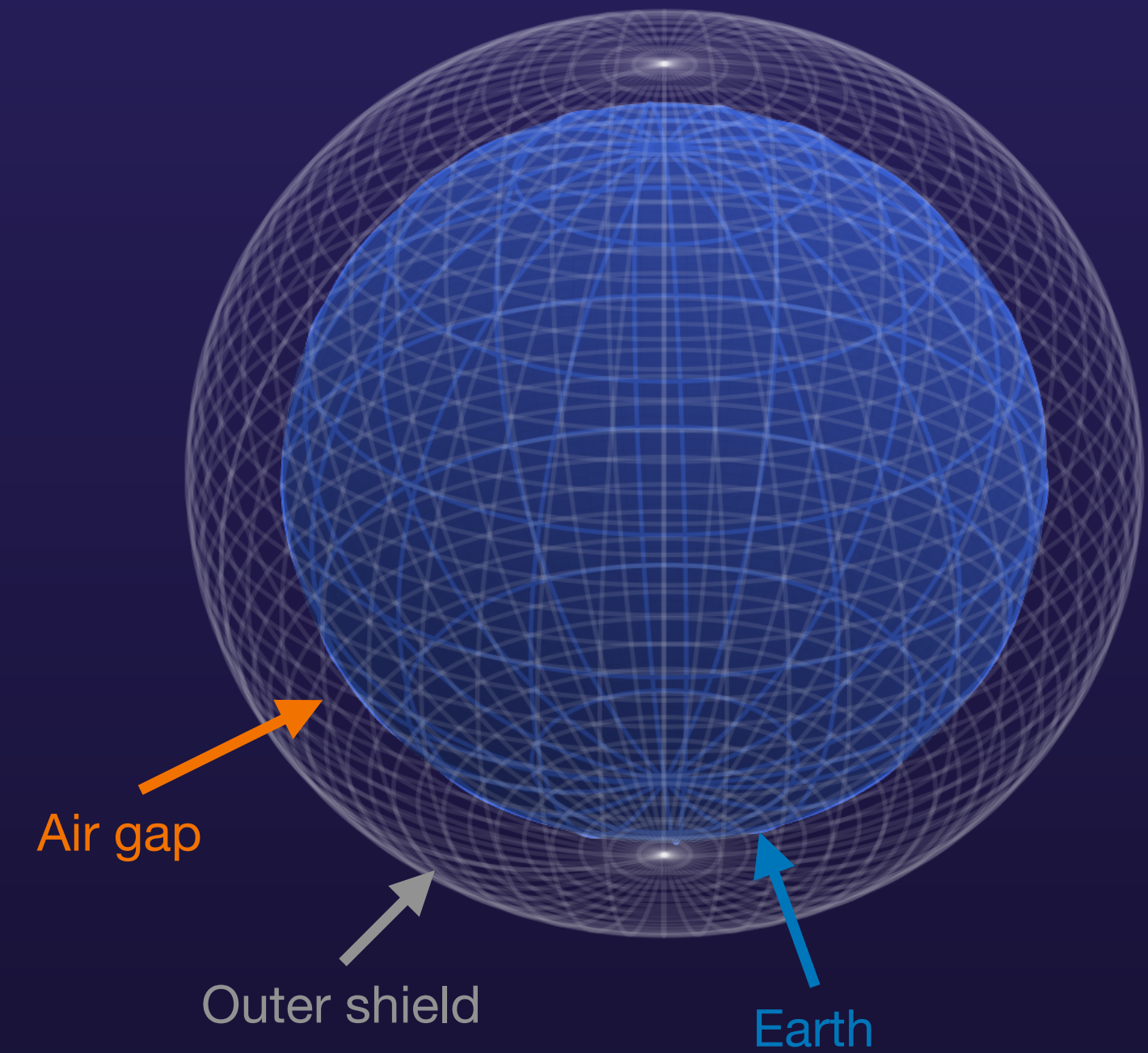
- ❖ DPDM fields source magnetic fields in a vacuum gap in an electromagnetically shielded region.
- ❖ Effects decouple as DPDM mass drops below the inverse size of the box.
 - What if we want to look for DPDM with $m_{A'} \lesssim 1/(1 \text{ meter}) \sim 2 \times 10^{-7} \text{ eV}$?
- ❖ Options? **OUR WORK PICKS UP HERE**
 - ~~Build a more sensitive magnetometer?~~ OK, but noise.
 - ~~Build Death Star 4.0?~~ Impractical, at best.
 - Find a pre-existing natural shield?
- ❖ What about the Earth?
 - It’s big enough: $1/R_{\oplus} \sim 3 \times 10^{-14} \text{ eV}$
 - But is it a shield?

Near-Earth Conductivity Environment

- ▶ Is the Earth a shield? It depends; for some masses, yes!



Simplified model



- ▶ The ground / interior is a region of high conductivity (“inner shield”)
- ▶ The lower atmosphere is a region of low conductivity (“vacuum gap”)
- ▶ The ionosphere / interplanetary medium beyond is a region of high conductivity / plasma frequency (“outer shield”)
- ▶ A nested conductor-insulator-conductor sandwich!

A new signal of DPDM I

- ❖ Spherical conductor, surrounded by a (vacuum-like) air gap, surrounded by a shield (ionosphere / IPM)

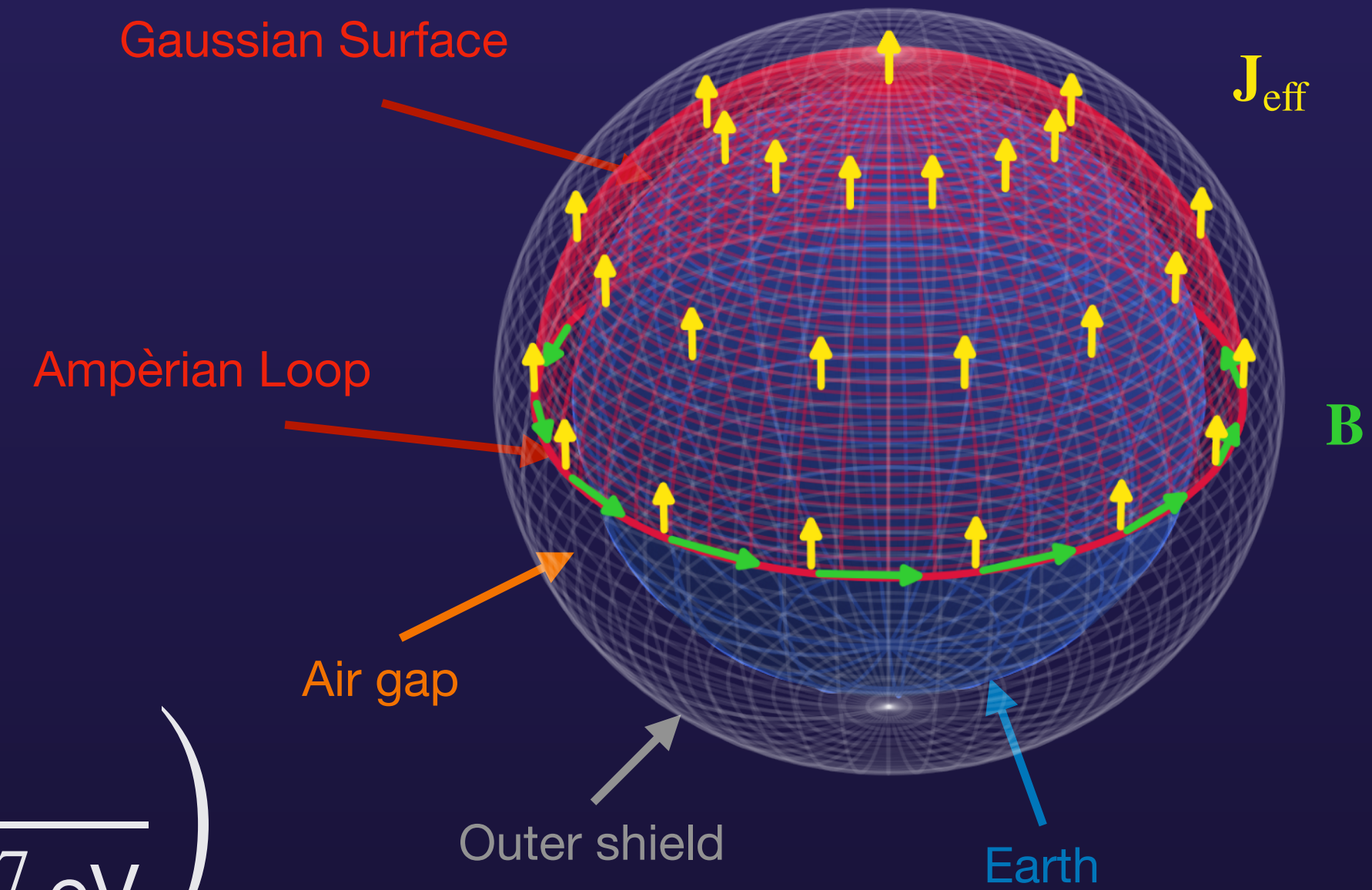
- ❖ What is the field on the ground?

- ❖ Run a similar Ampèrian loop argument:

$$\iint \mathbf{J} \cdot d\mathbf{A} = \oint \mathbf{B} \cdot d\mathbf{l} \Rightarrow R_{\oplus}^2 J_{\text{eff}} \sim BR_{\oplus}$$

Displacement current term is higher-order

$$|B| \sim \varepsilon(m_{A'} R_{\oplus}) \sqrt{\rho_{\text{DM}}} \sim 0.7 \text{ nG} \times \left(\frac{\varepsilon}{10^{-5}} \right) \times \left(\frac{m_{A'}}{4 \times 10^{-17} \text{ eV}} \right)$$



- ❖ **A PERSISTENT, GLOBAL, NARROWBAND MAGNETIC FIELD OSCILLATION WITH A KNOWN PATTERN, IN-PHASE OVER THE ENTIRE SURFACE OF EARTH**

- ❖ The field is **very small**, but **spectral and spatial features differ from noise sources**

- ❖ **Suppressed by $(m_{A'} R_{\oplus})$, not $(m_{A'} h_{\text{atmos.}}) \ll (m_{A'} R_{\oplus})!$**

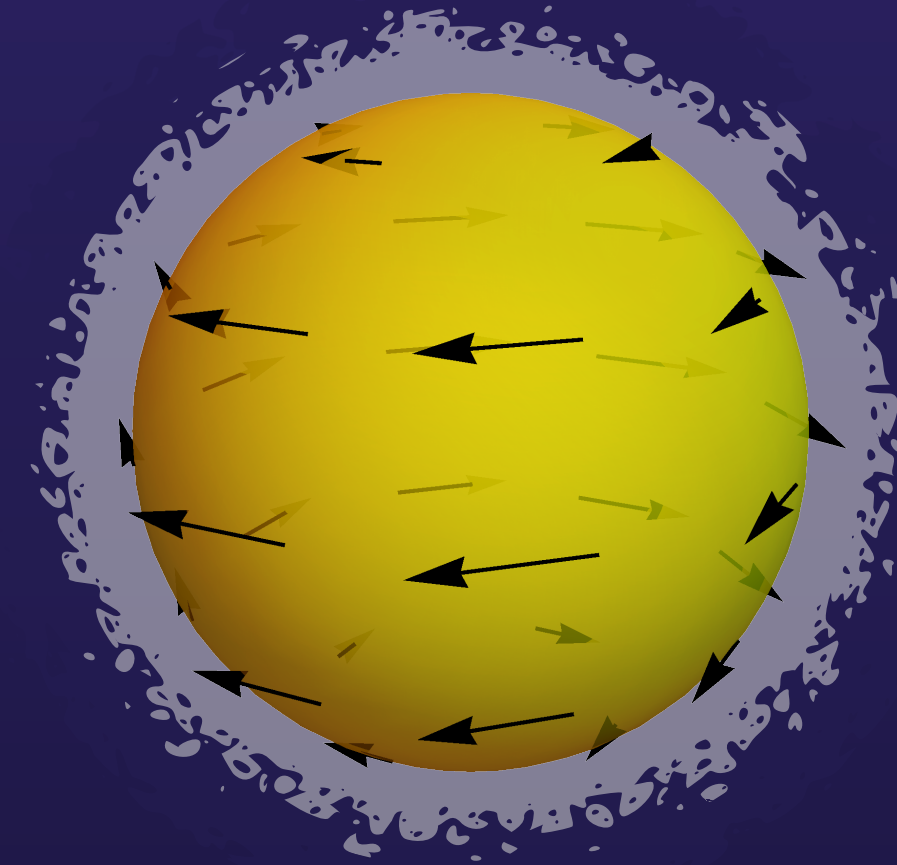
A new signal of DPDM II

- Performing the computation rigorously, we find

$$\mathbf{B}(\Omega, t) = \sqrt{\frac{\pi}{3}} \varepsilon m_{A'}^2 R_{\oplus} \sum_{m=-1}^1 A'_m \Phi_{1m}(\Omega) e^{-i(m_{A'} - 2\pi f_d m)t}$$

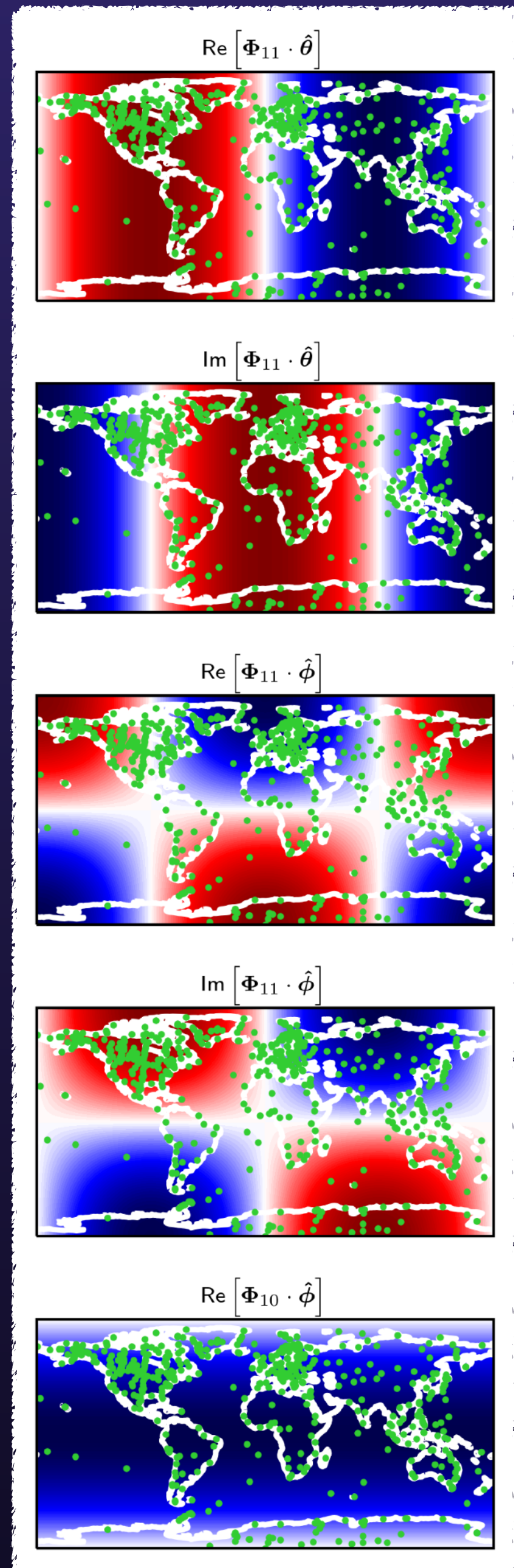
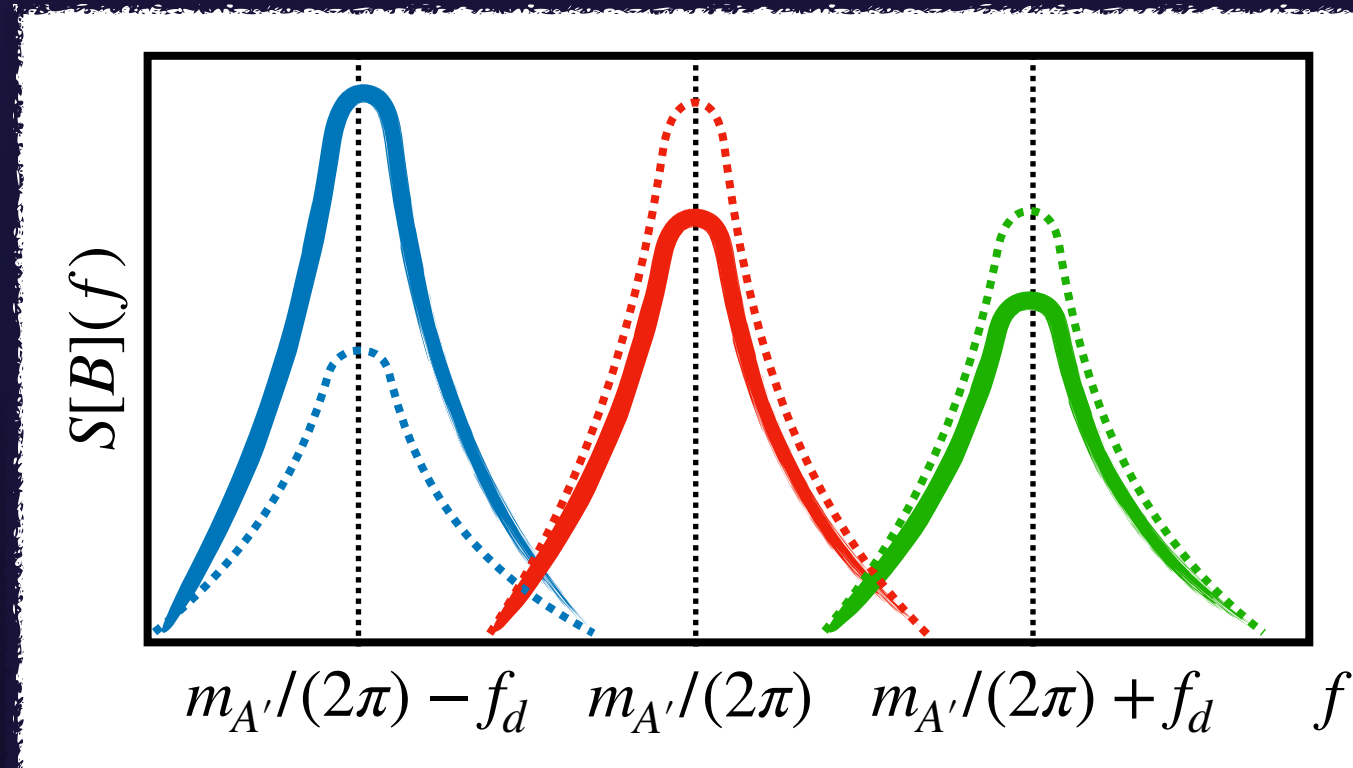
+ higher-order terms + TE fields

where $\frac{1}{2} m_{A'}^2 \langle |A'|^2 \rangle_{T \gg T_{\text{coh}}} = \rho_{\text{DM}}$



- Signal found in one type of Vector Spherical Harmonic (VSH): $\Phi_{1m}(\Omega)$
- Field evaluated at ground magnetometer station at $\Omega = (\theta, \phi)$. Earth frame (rotating).
- DPDM polarization state \mathbf{A}' fixed in inertial space (within coherence time / patch)
- Non-trivial frequency structure
Sensitive to DPDM polarization state

Sidebands offset by $f_d = 1/(\text{sidereal day})$



- Signal derivation holds for $10^{-21} \text{ eV} \lesssim m_{A'} \lesssim 1/R_{\oplus} \sim 3 \times 10^{-14} \text{ eV}$
- Complications and caveats arise in real world: this signal remains (leading-order TM field in an $m_{A'} L$ expansion)

What about the axion?

- ❖ A similar signal arises, because of the Earth's static geomagnetic field \mathbf{B}_\oplus !
- ❖ $\mathcal{L} \supset \frac{g_{\phi\gamma}}{4} \phi F \tilde{F} = g_{\phi\gamma} \phi \mathbf{E} \cdot \mathbf{B} \longrightarrow g_{\phi\gamma} (\partial_t \phi) \mathbf{B}_\oplus \cdot \mathbf{A} \sim -i g_{\phi\gamma} m_\phi \phi \mathbf{B}_\oplus \cdot \mathbf{A}$
- ❖ Once again, looks like an effective current: $\mathbf{J}_{\text{eff}} = i g_{\phi\gamma} m_\phi \phi \mathbf{B}_\oplus$
- ❖ $\hat{\mathbf{B}}_\oplus$ plays the role of $\hat{\mathbf{A}}'_I$

[Earth Planets Space 73 (2021) 49]

- ❖ Use the International Geomagnetic Reference Field (IGRF-13) model for \mathbf{B}_\oplus

$$\mathbf{B}(\Omega, t) = -i(g_{\phi\gamma} \phi_0)(m_\phi R_\oplus) \sum_{l,m} \frac{c_{lm}}{l} \Phi_{lm}(\Omega) e^{-im_\phi t}$$

Fixed by IGRF-13

$$B \sim 1 \text{ nG} \times \left(\frac{g_{\phi\gamma}}{10^{-10} \text{ GeV}^{-1}} \right)$$

$$\text{where } \frac{1}{2} m_\phi^2 \langle \phi_0^2 \rangle_{T \gg T_{\text{coh}}} = \rho_{\text{DM}}$$

- ❖ Signal still found in one type of VSH, but get higher multipoles l from Earth's field
- ❖ Sourcing magnetic field is co-rotating with the ground-based detectors: frequency content now only at $f = m_A l (2\pi)$

Rough translation for amplitude

$$\varepsilon m_{A'} \leftrightarrow g_{\phi\gamma} B_\oplus$$

Assuming DM field-amplitude normalisations
[BUT: different pattern]

Signal Properties

- ❖ In both cases, the **Earth acts as a transducer** to convert **dark-matter oscillations** to an **oscillating magnetic field**:
 - ▶ narrowband frequency peaked near $\omega \sim m_{\text{DM}}$
 - ▶ long coherence time and length
 - ▶ present globally & in-phase across the whole surface of the Earth
 - ▶ known pattern
- ❖ We would like to implement a search for these signals. How?
- ❖ This is a perfect signal to search for using a network of geographically distributed vector magnetometers that record time-series field data
- ❖ **They must be unshielded!** Unfortunately, GNOME doesn't work for this [Ann. Phys. (Berl.) **525** (2013) 659]

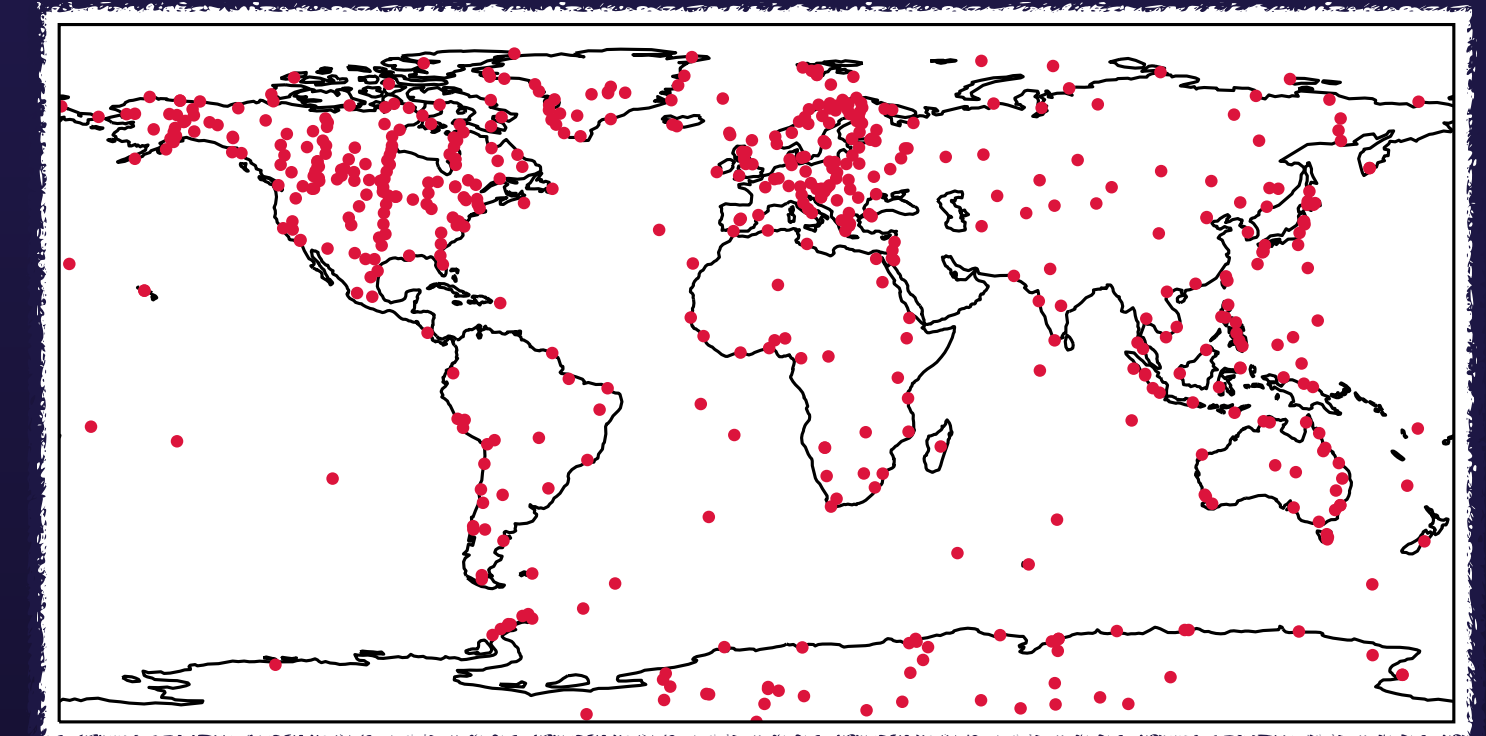
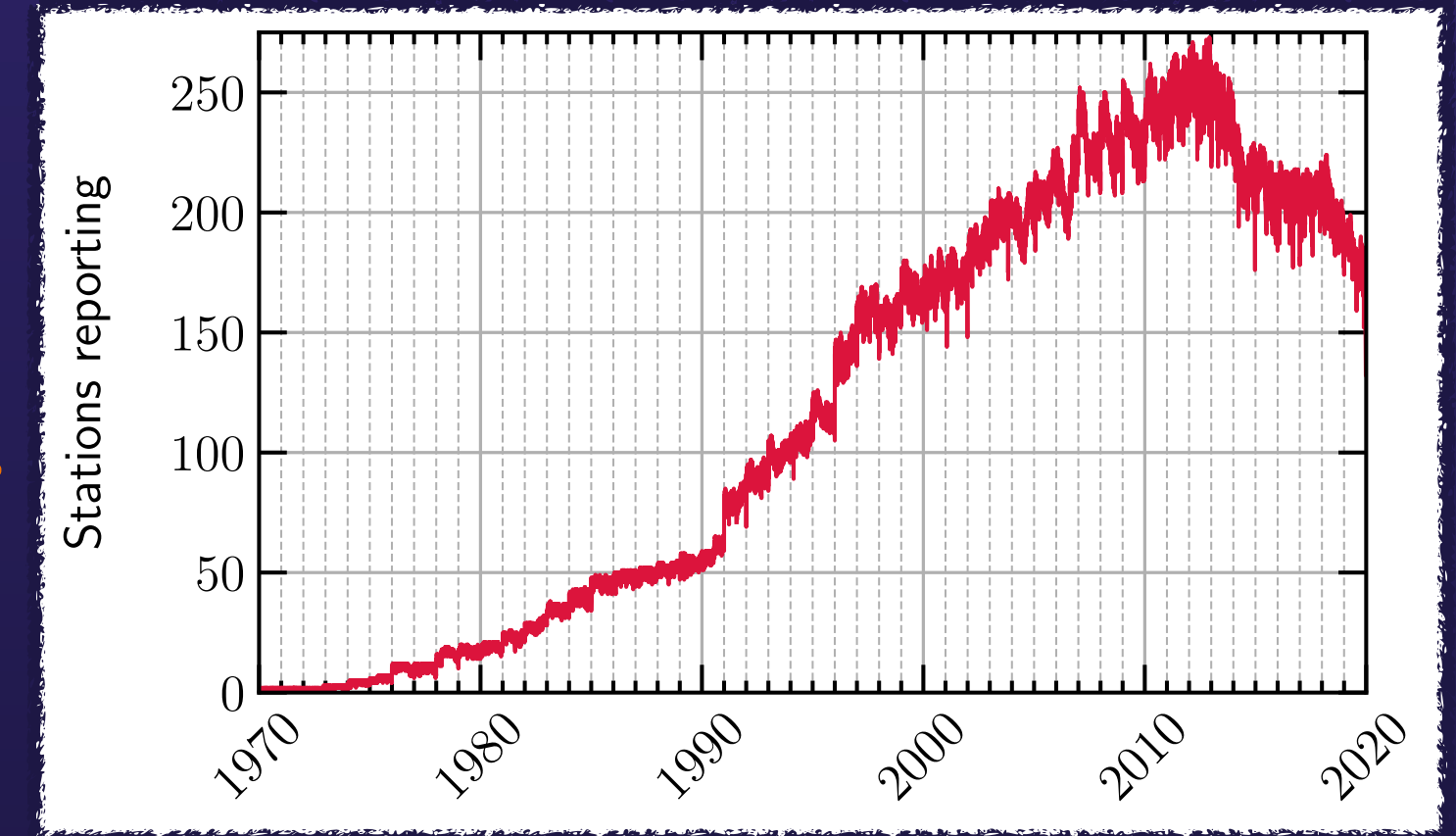
SuperMAG

<https://supermag.jhuapl.edu>

Eos **90** (2009) 230-231

J. Geophys. Res. Space Phys. **117** (2012) A09213

- ❖ Geophysics to the rescue!
- ❖ Collaboration based at JHU-APL (J. W. Gjerloev, PI). Many contributors.
See acknowledgments slide at end of talk
- ❖ Make publicly available a database of:
 - ▶ Unshielded magnetic field measurements
 - ▶ at $\mathcal{O}(500)$ unshielded magnetometer stations; maximally, $\mathcal{O}(250)$ reporting at any single time
 - ▶ widely geographically distributed over the Earth's surface
 - ▶ in time-series
 - ▶ with 60s resolution*
 - ▶ since ~1970 (more data at later times)
 - ▶ in a common data format
 - ▶ in a data-driven co-ordinate reference frame (average magnetic E-W nulled; reference to IGRF)
- ❖ **A near-ideal dataset to analyse for our dark-matter signals**



*a smaller dataset with 1s resolution is also available... see later in the talk

Search for a signal - Analysis Overview (DPDM)

Search for:

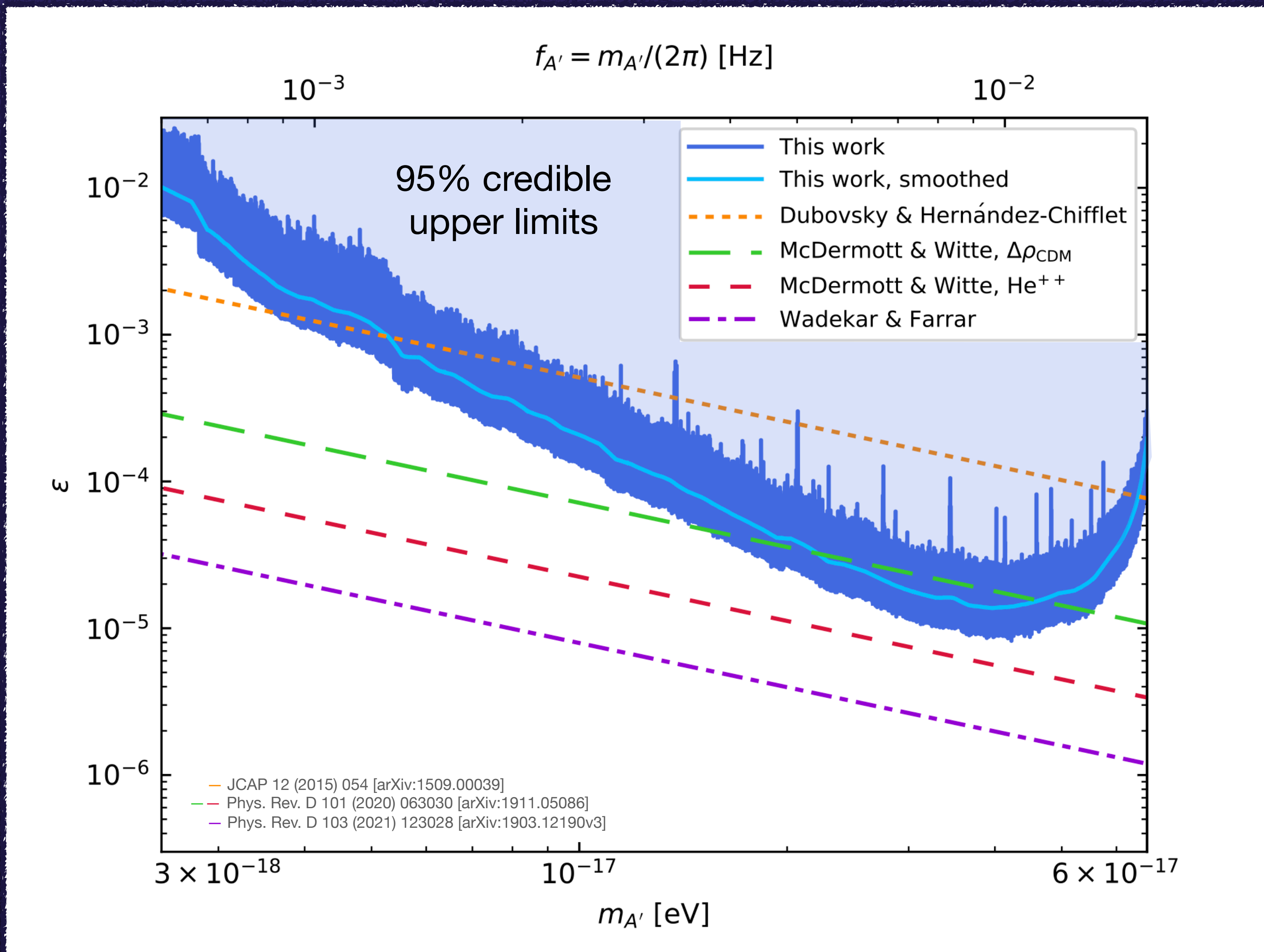
narrowband excess magnetic field power
with the correct spatial configuration

over a

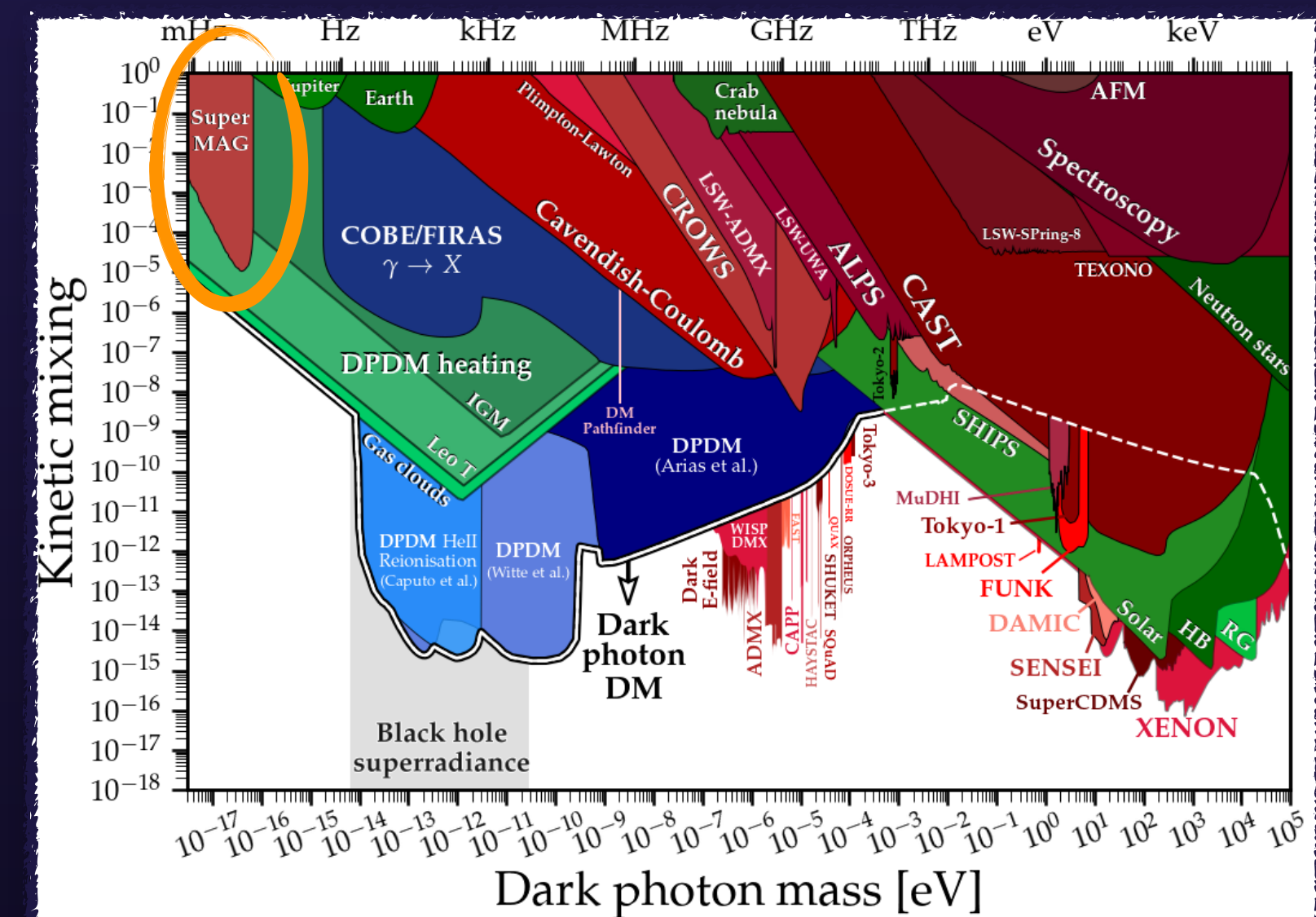
broadband magnetic field background
estimated from data

PRD 104, 095032 (2021)

DPDM Results I - Limits

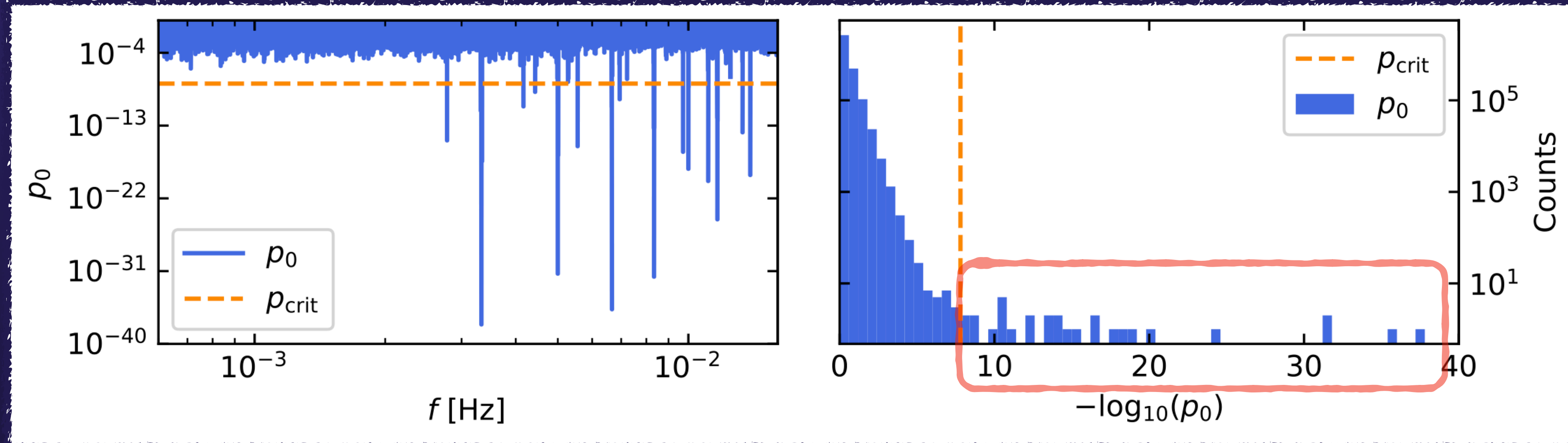


- ❖ Complementary to astrophysical limits
- ❖ Very different systematics!
- ❖ First “direct” constraints on this piece of parameter space.



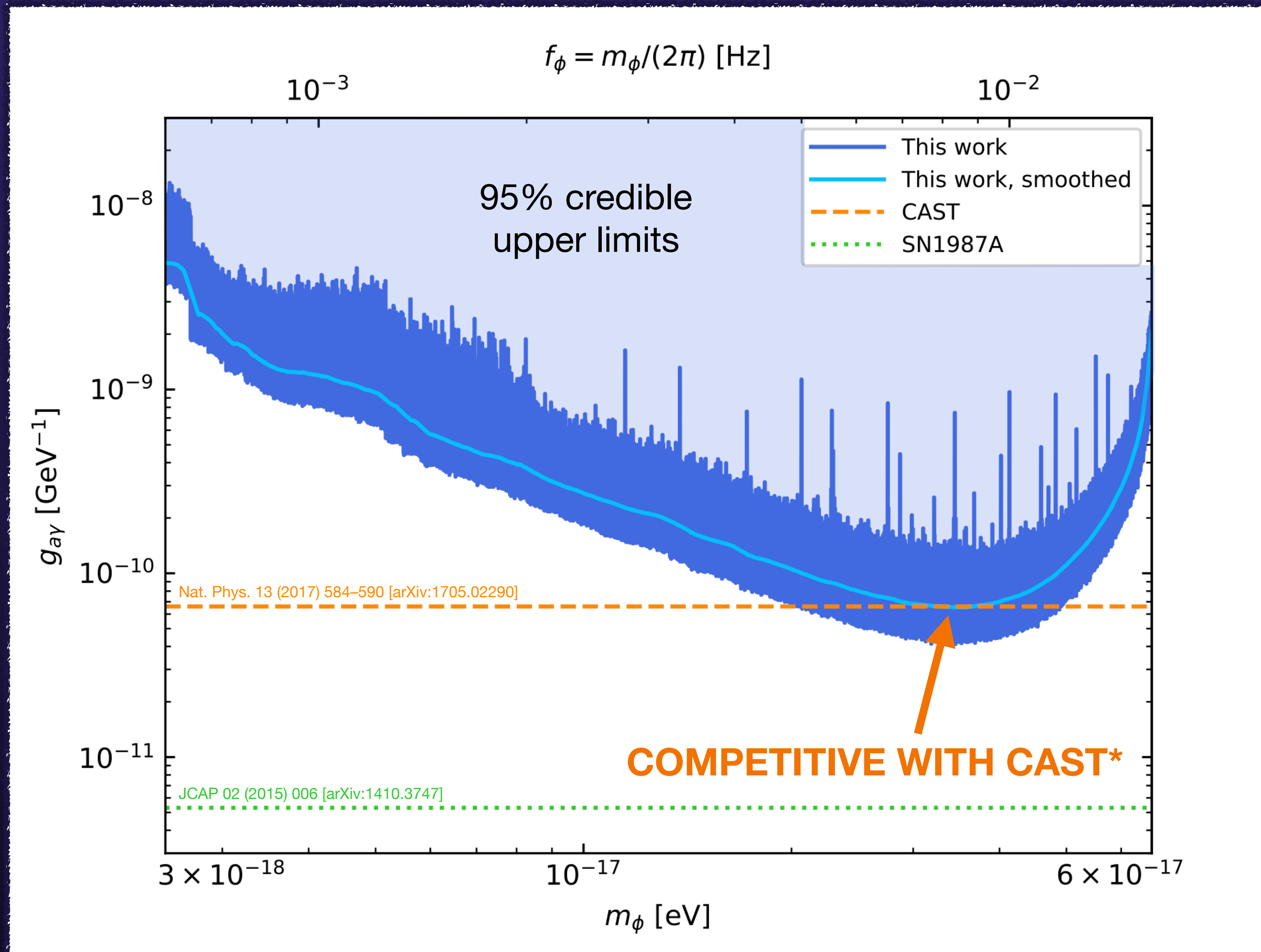
<https://cajohare.github.io/AxionLimits/>
PRD 104, 095029 (2021) [arXiv:2105.04565]

DPDM Results II - Signal candidates?



- ❖ 30 naive signal candidates identified
- ❖ Reject 24 completely on the basis of **robustness tests** for spatial consistency, temporal constancy.
- ❖ 6 remaining candidates are in tension with at least one of the checks (or are otherwise problematic); further work might be required to definitively exclude.
- ❖ E.g.: candidate at 3.344939mHz has a 6.1σ global significance, but $p = 0.034$ on the temporal constancy robustness check ($p = 0.97$ on the spatial check).
- ❖ **We do not claim any robust evidence for a DPDM signal**

Axion Results



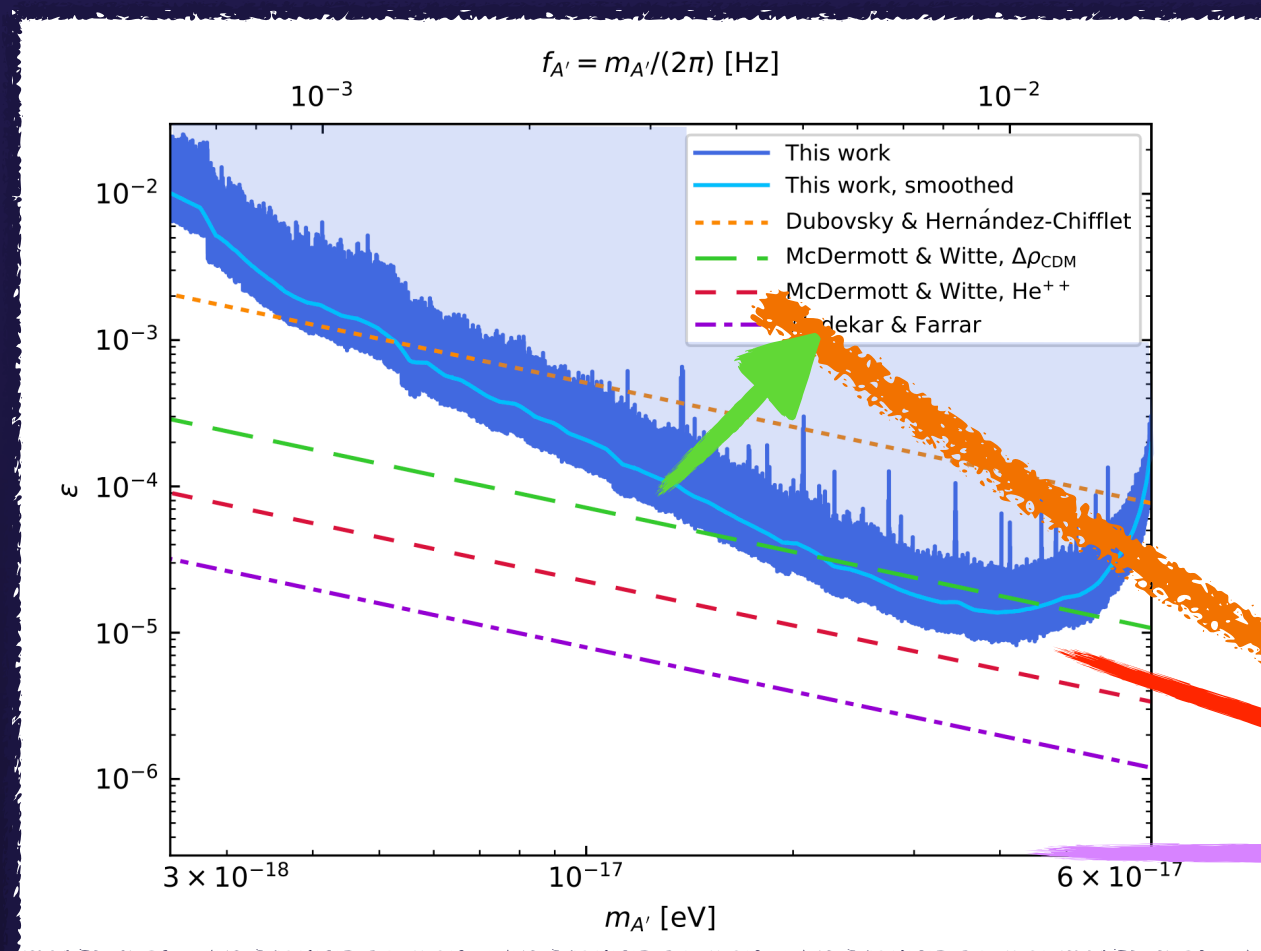
- ❖ Analysis proceed similarly, but simpler.
- ❖ Similar naive signal candidates exist (27).
- ❖ Again reject most (23) outright with robustness checks.
- ❖ Some weak signals remain (4), but have tensions with checks.
- ❖ **We do not claim any robust evidence for an axion signal**

*CAST does not need to assume the axion is all of the dark matter

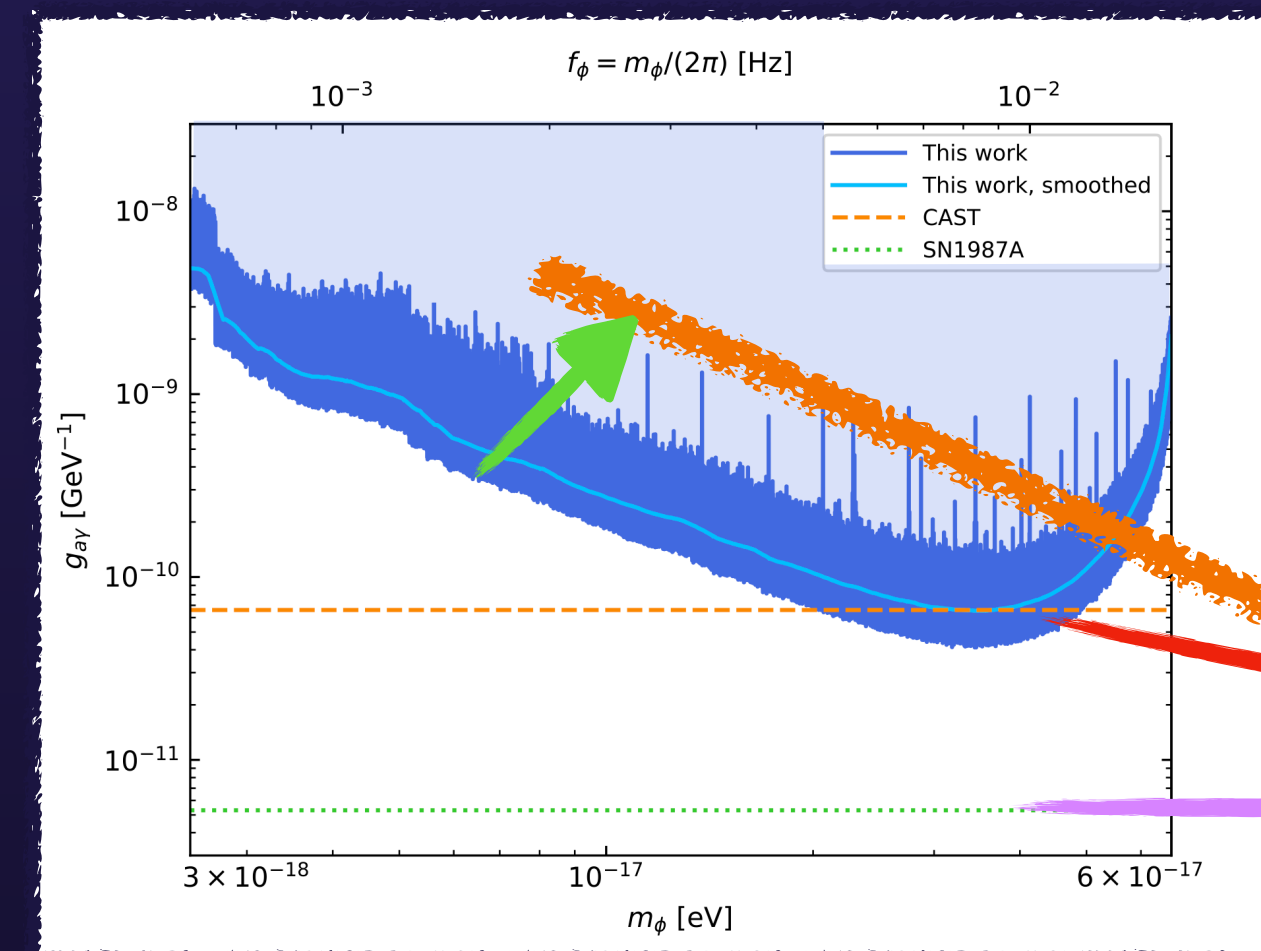
Outlook and Future Directions

- ❖ SuperMAG has 1s resolution data available. But less of it (fewer stations, less time).
- ❖ Limit of our 60s-cadence analysis at high frequencies is the measurement sampling frequency.
- ❖ In 60s data, noise falls with increasing frequency.
- ❖ Potential to **sacrifice some sensitivity*** with data quantity, but **push to higher frequency**.

*In fact, we might not even lose sensitivity. Less total time duration at 1s, but more data per time.



?



?

- ❖ **Win in a relative sense?** More competitive with astro limits? CAST?
- ❖ **Watch this space. Analysis in progress** (in formal collaboration with SuperMAG).
- ❖ Even higher frequencies? A dedicated network of higher-frequency magnetometers: SNIPE Hunt

There was a poster about this at the APS DAMOP Meeting in June

Summary

- ❖ Earth is a transducer to generate new signals of dark-photon and axion dark matter
 - Coherent, narrowband oscillating magnetic fields that have particular vectorial spatial pattern over whole surface of the Earth**
- ❖ Distributed networks of unshielded magnetometers are DM detectors
- ❖ Search in SuperMAG magnetic field data:
 - ▶ **No robust candidates of either DM type**
 - ▶ **New limits, complementary to existing constraints**
- ❖ Search in other existing data is underway. Extend reach to new parameter space?
- ❖ Dedicated search push to even higher frequencies being considered

Questions on this part of the talk?

White Dwarfs as Amplifiers for CHAMP Detection

THE SHORT STORY

(Maybe some very tiny fraction of) Dark Matter in Compact Objects, ~~Stars, and in Low-Energy Experiments~~

... get dragged into galaxies, distributed like DM $\sigma/m_X \ll (\sigma/m)_{SM} \dots$

X^\pm
Charged Massive Particles (CHAMPs)
[$m_X \gtrsim 10^{11}$ GeV]

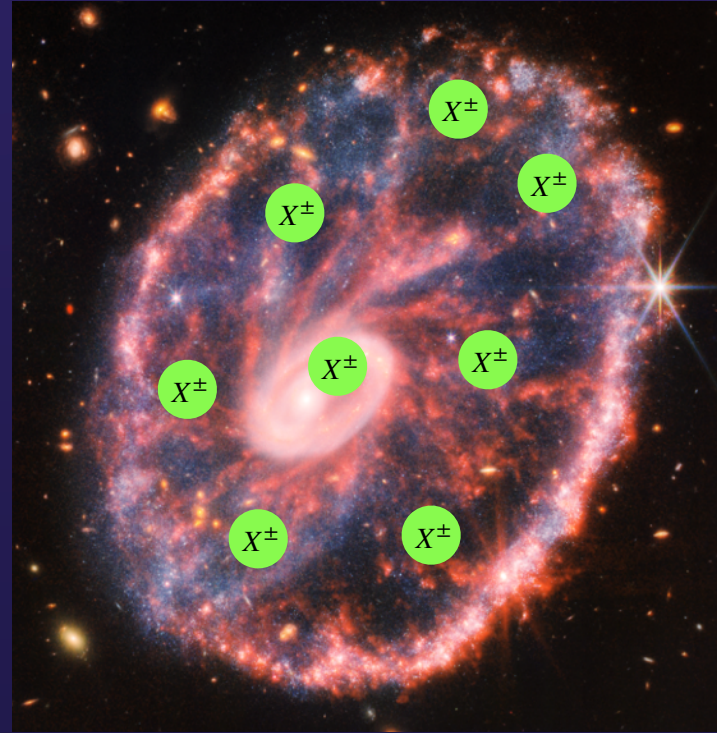
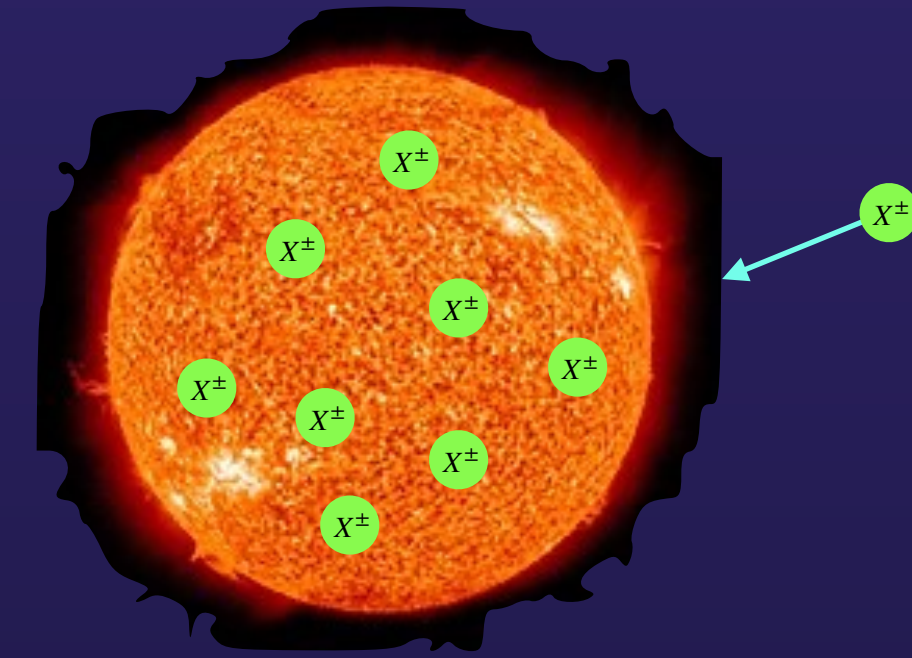
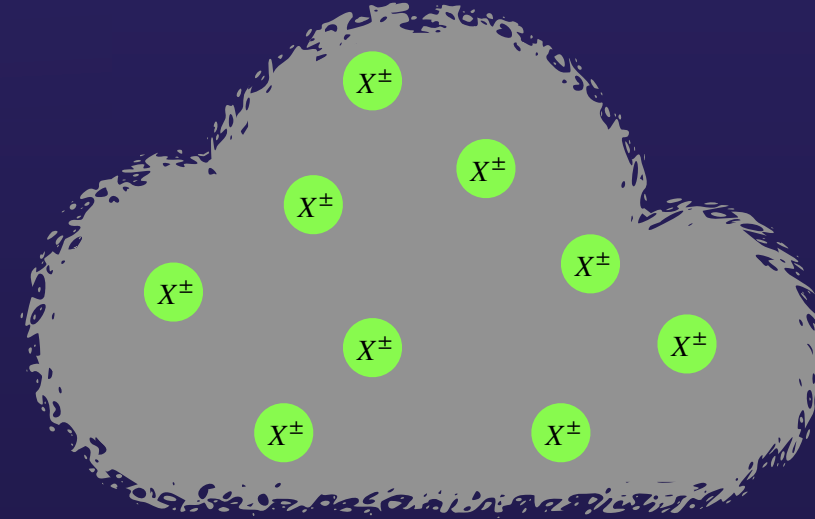


IMAGE: NASA, ESA, CSA, STScI, Webb ERO Production Team

... but they can get trapped in denser protostellar gas clouds...



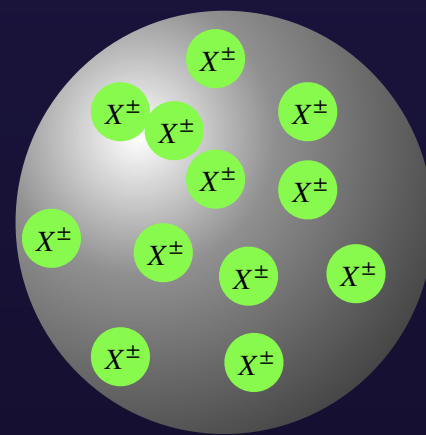
... and contaminate stars from birth (or get captured directly)

$$M \sim 10 - 30 M_\odot$$

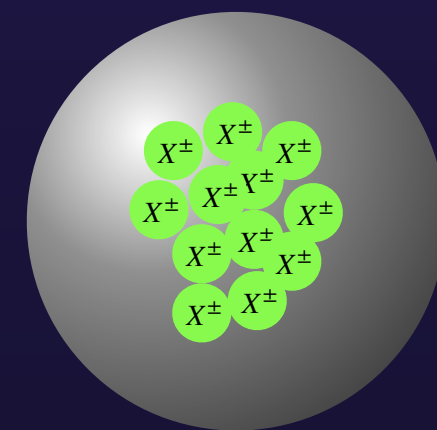
which evolve to...

X^- bind as $(\text{He}X)^+$

...where the CHAMPs sink...

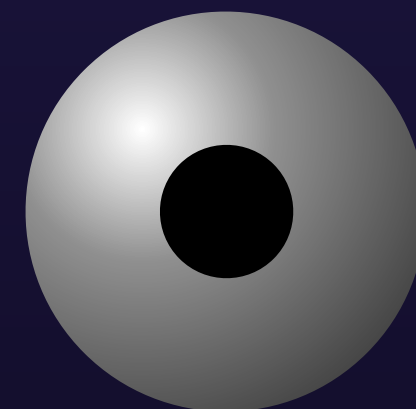


Neutron Stars

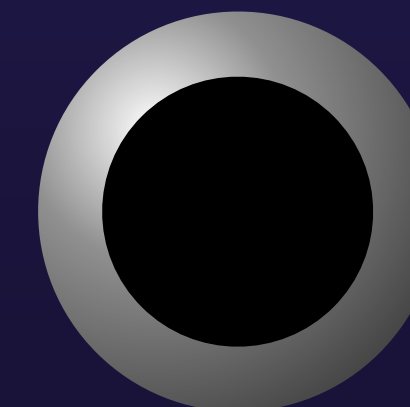


...and self-gravitate...
($M_X^{\text{tot}} > M_{s.g.}$)

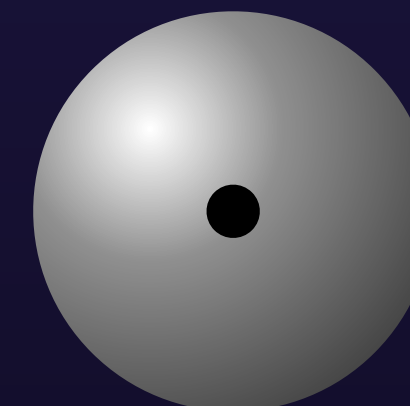
... and can form a mini black hole inside the NS...



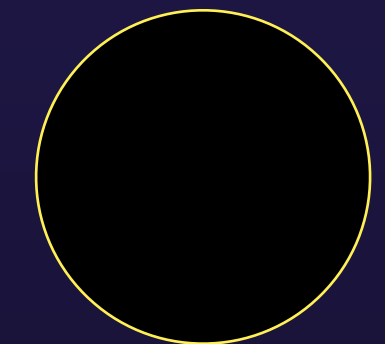
...which can...



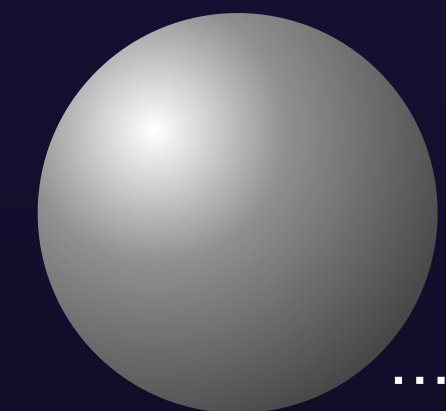
grow



evaporate



... and destroy the NS



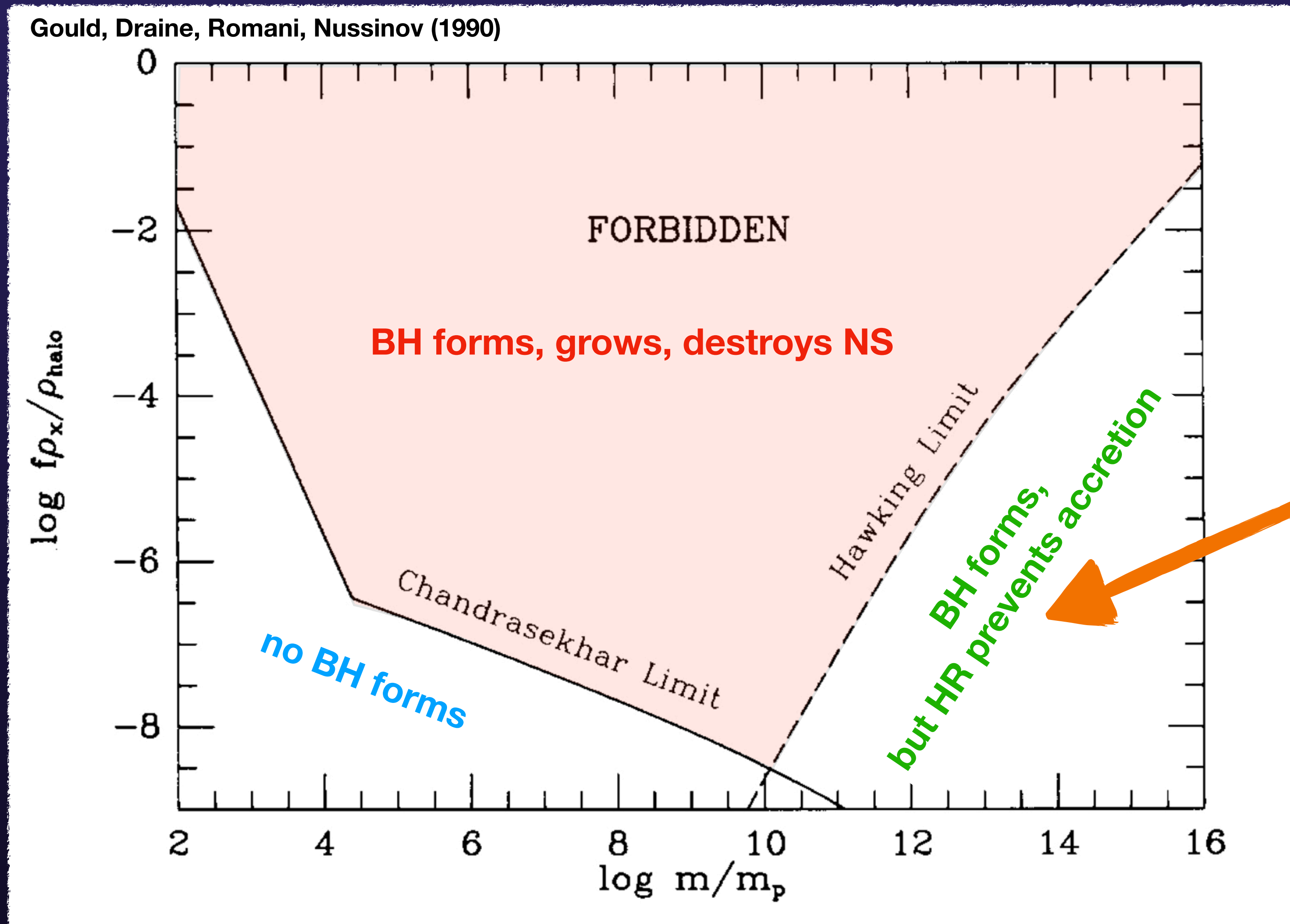
... and leave the NS effectively untouched

$$M_{\text{Chand}}^X \sim 1.4 M_\odot \times \left(\frac{Q_X m_p}{m_X} \right)^2$$

**Old NS existence bounds
CHAMP abundances**

Old Neutron Star Bounds

$$M_{\text{Chand.}}^X \sim 1.4M_{\odot} \times \left(\frac{Q_X m_p}{m_X} \right)^2$$



OUR QUESTION

Can we find a system where this region CAN be bounded?

Possibly orders of magnitude of improvement to be had

... get dragged into galaxies, distributed like DM $\sigma/m_X \ll (\sigma/m)_{SM} \dots$

X^\pm
Charged Massive Particles (CHAMPs)
[$m_X \gtrsim 10^{11} \text{ GeV}$]

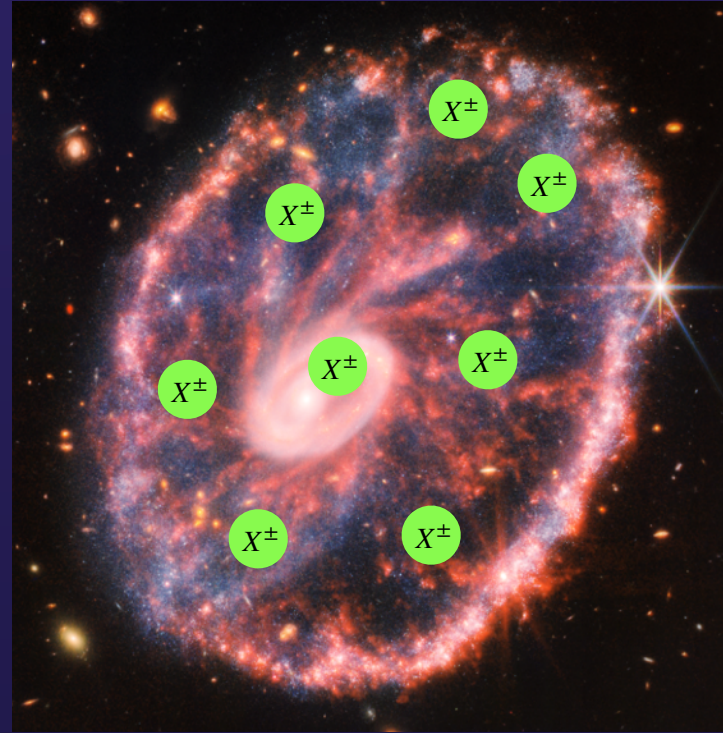
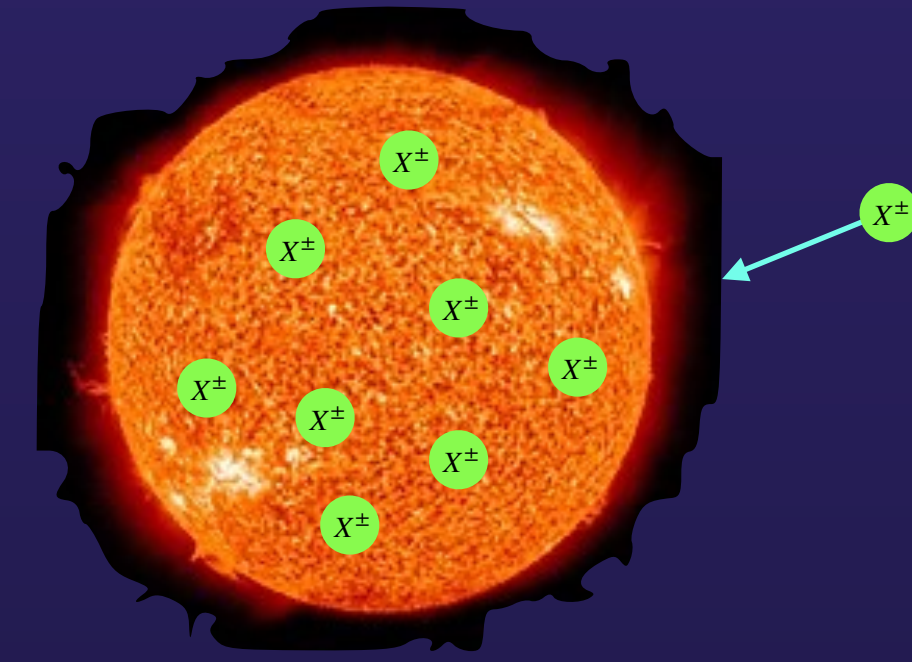
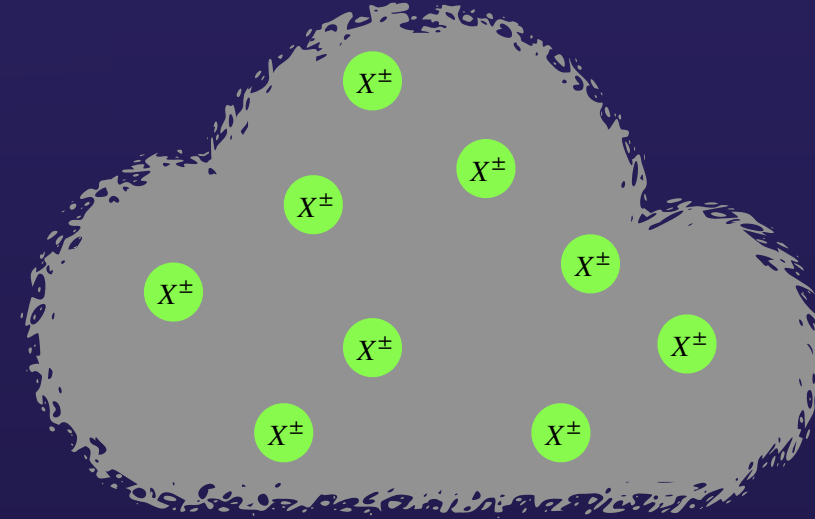


IMAGE: NASA, ESA, CSA, STScI, Webb ERO Production Team

... but they can get trapped in denser protostellar gas clouds...



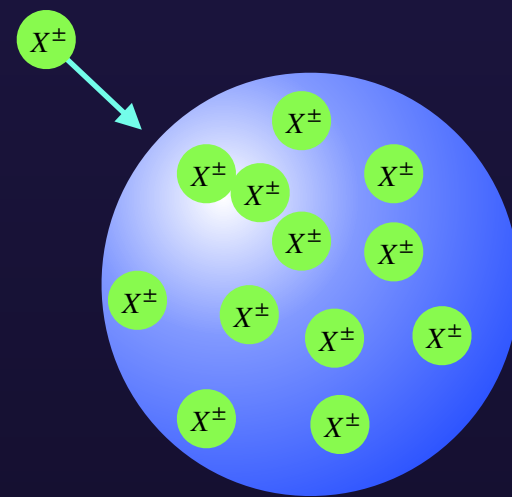
... and contaminate stars from birth (or get captured directly)

$M \sim 2 - 8 M_\odot$

which evolve to...

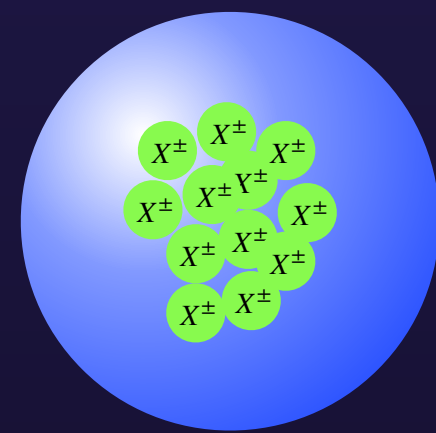
X^- bind as $(\text{He}X)^+$

...where the CHAMPs sink...

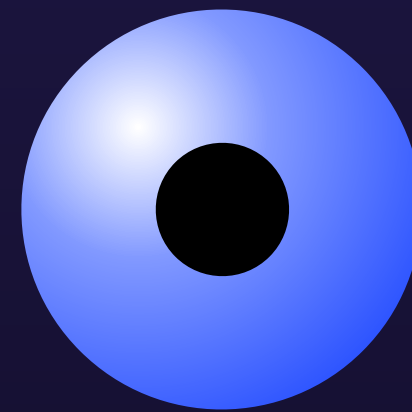


White Dwarfs

which maybe also still accrete some CHAMPs

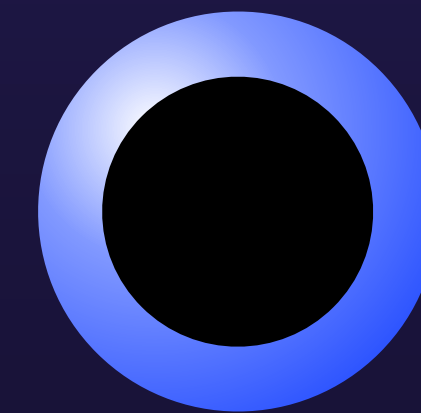


...and self-gravitate...
($M_X^{\text{tot}} > M_{s.g.}$)

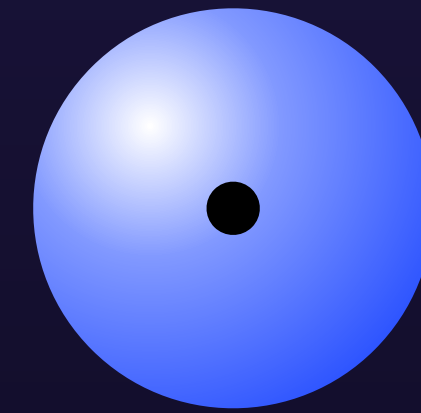


... and can form a mini black hole inside the WD...

...which can...

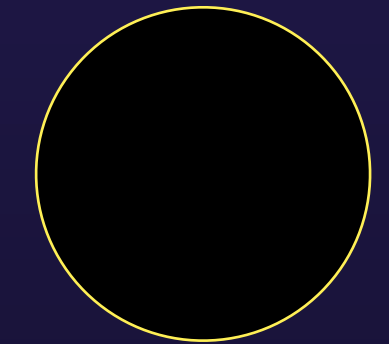


grow



evaporate

... and destroy the WD



... and trigger the supernova instability

Under very general conditions, a BH born inside a WD such that it evaporates will do so within the WD lifetime ($\sim \text{Gyr}$), and will trigger the Type-Ia SN instability (runaway carbon burning)

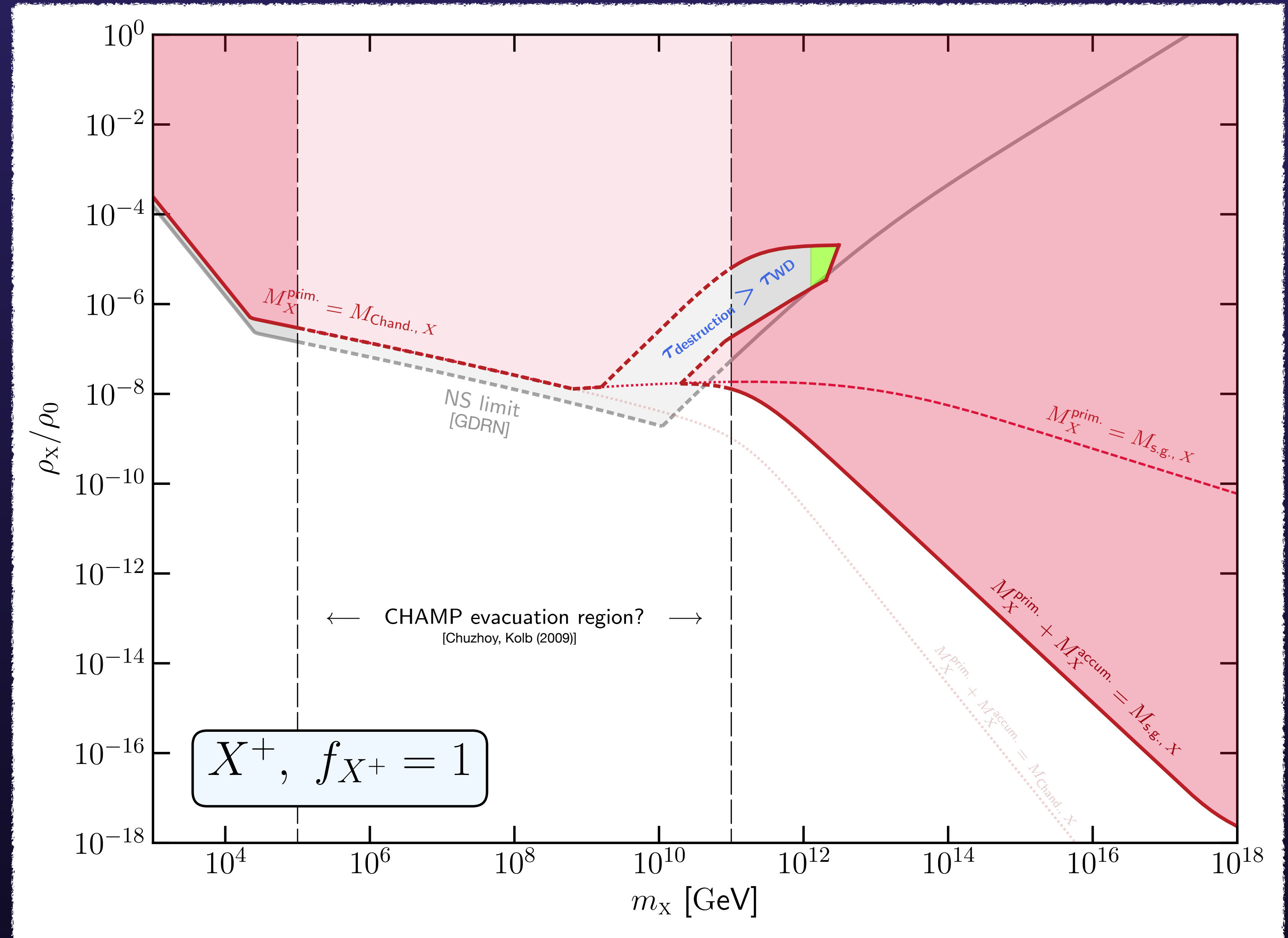
Bramante, Acevedo (2019)

Janish, Narayan, Riggins (2019)

Old WD existence bounds
CHAMP abundances

Old White Dwarf Bounds - X^+

Require
 enough CHAMPs to get into the WD
 via
 primordial contamination
 and
 accretion
 to achieve the larger of a
 self-gravitating mass of CHAMPs
 or
 a Chandrasekhar mass of CHAMPs
 AND
 destroy the WD within its cooling age



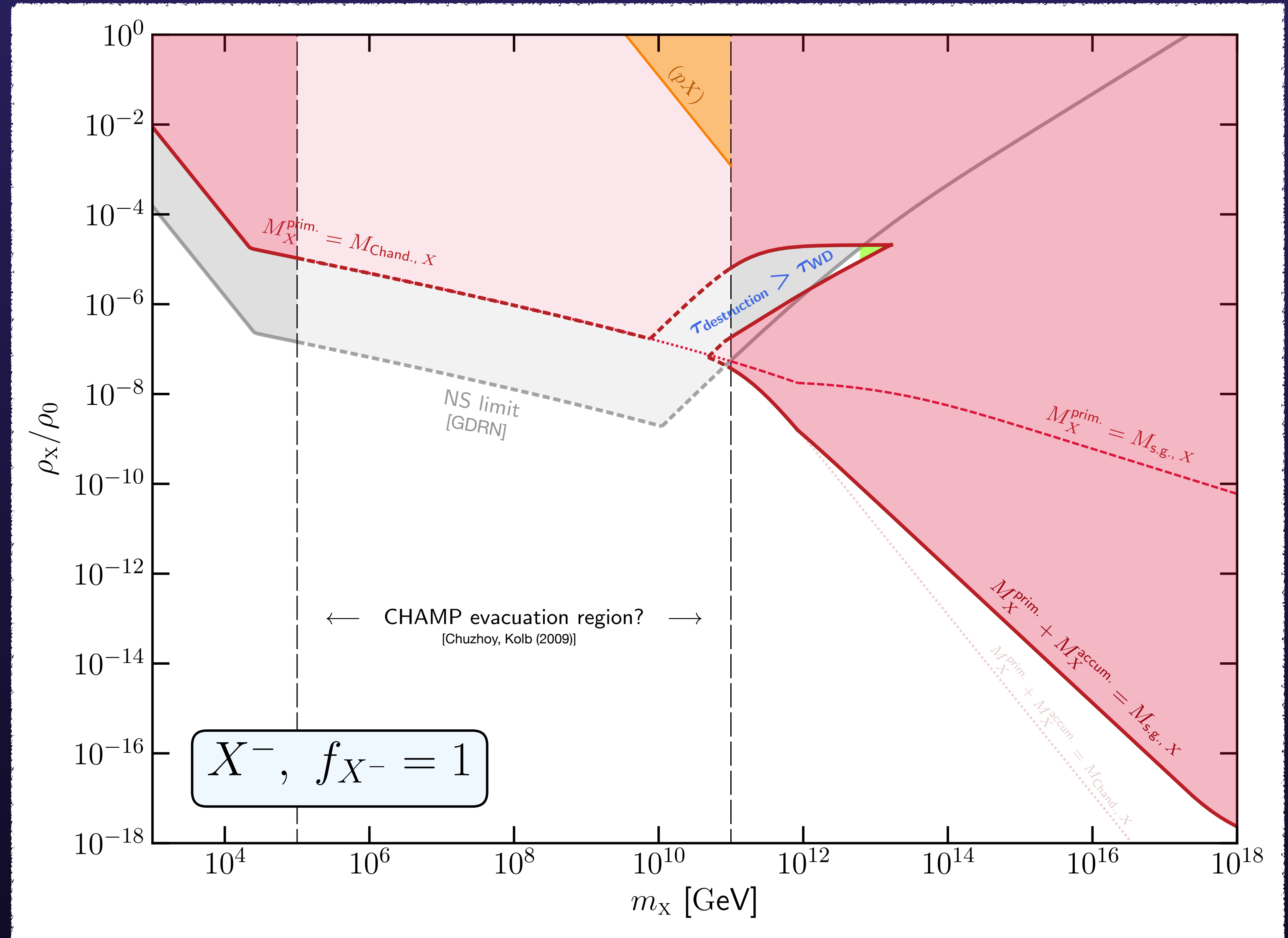
| Name | $M_{WD} [M_{\odot}]$ | $B [MG]$ | $t_{cool.} [Gyr]$ |
|------------------------|----------------------|----------|-------------------|
| WDJ062144.86+753011.67 | 1.18–1.23 | — | 4.1 |
| WDJ013839.12-254233.40 | 1.17–1.22 | — | 4.2 |
| WD 2202-000 | 1.08 | 1.0 | 2.19 |

Old White Dwarf Bounds - X^-

X^- exist in WD as
(CX)⁵⁺ or (OX)⁷⁺
objects.

Binding energies are
3-4 MeV.

$$M_{\text{Chand.}}^X \sim 1.4M_{\odot} \times \left(\frac{Q_{(NX)} m_p}{m_X} \right)^2$$



| Name | $M_{\text{WD}} [M_{\odot}]$ | B [MG] | $t_{\text{cool.}}$ [Gyr] |
|------------------------|-----------------------------|----------|--------------------------|
| WDJ062144.86+753011.67 | 1.18-1.23 | — | 4.1 |
| WDJ013839.12-254233.40 | 1.17-1.22 | — | 4.2 |
| WD 2202-000 | 1.08 | 1.0 | 2.19 |

Summary

- ❖ Charged massive particles (CHAMPs) contaminate white dwarfs
- ❖ If sufficient total CHAMP mass in the WD (larger of Chandrasekhar and self-gravitating masses), form mini black hole inside the WD
- ❖ BH either accretes up in mass, or Hawking radiates down in mass. Either way destroys WD, except if it takes too long.
 - Other supernova trigger mechanisms discussed in the paper
- ❖ Dramatic, orders of magnitude improvement of GDRN galactic abundance bounds on high- m_X CHAMPs
- ❖ Speculation: trigger for Ca-Rich Gap Transients? Correct spatial morphology?

Concluding remarks

I told you about:

- ❖ A new signal of ultralight dark-photon and axion-like particle dark matter in terrestrial magnetometer data, and new limits

M.A.F., P. W. Graham, S. Kalia, D. F. Jackson Kimball. Phys. Rev. D **104**, 075023 (2021) [arXiv:2106.00022].

M.A.F., P. W. Graham, S. Kalia, D. F. Jackson Kimball. Phys. Rev. D **104**, 095032 (2021) [arXiv:2108.08852].

A. Arza, M.A.F., P. W. Graham, S. Kalia, D. F. Jackson Kimball. Phys. Rev. D **105**, 095007 (2022) [arXiv:2112.09620].

- ❖ Using white dwarfs to constrain a different early-universe relic: charged massive particles (CHAMPs), and new limits

M.A.F., P.W. Graham, and S. Rajendran. Phys. Rev. D **101**, 115021 (2020) [arXiv:1911.08883]

Thanks!

Questions?