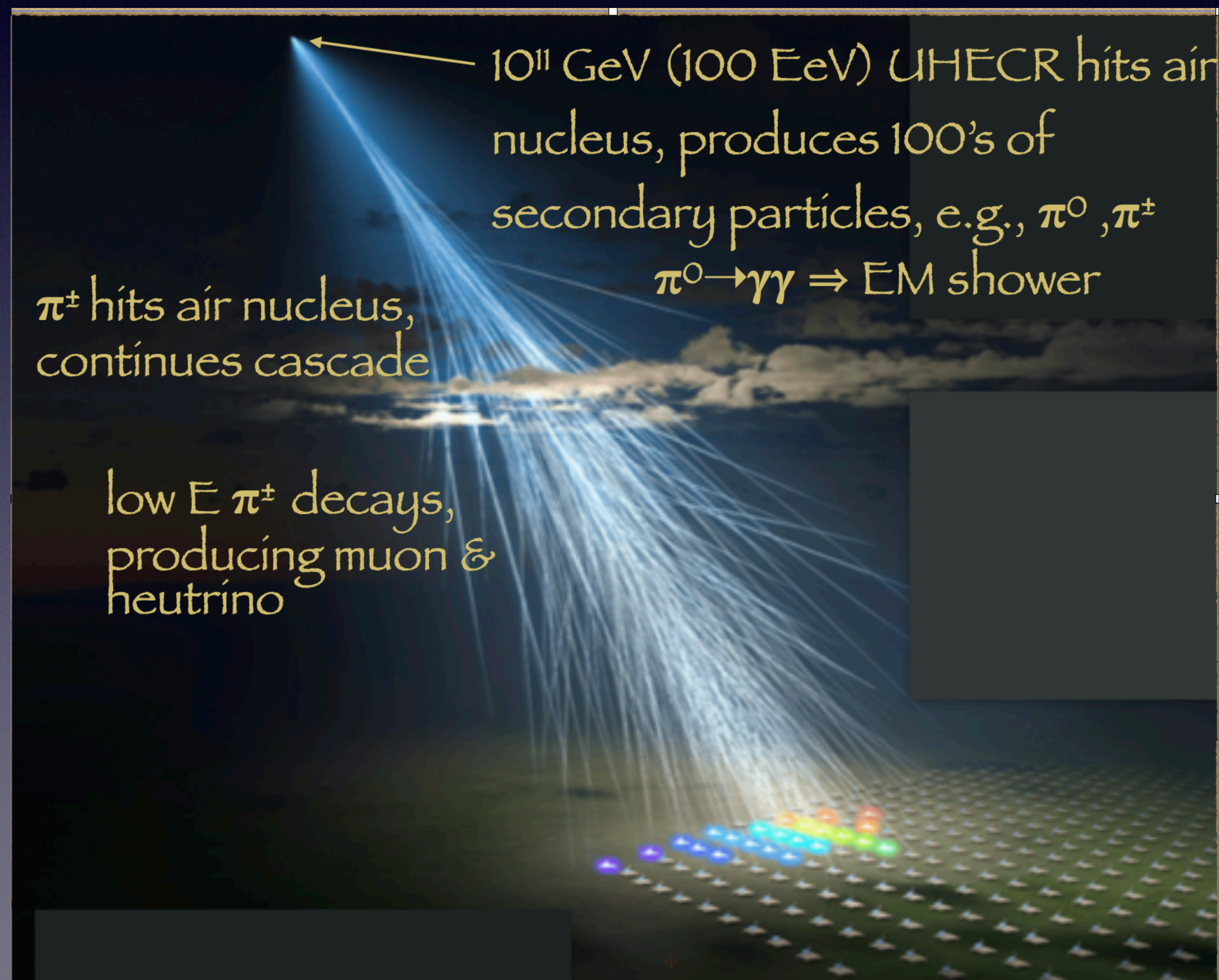


The Origin of Ultra-high Energy* Cosmic Rays



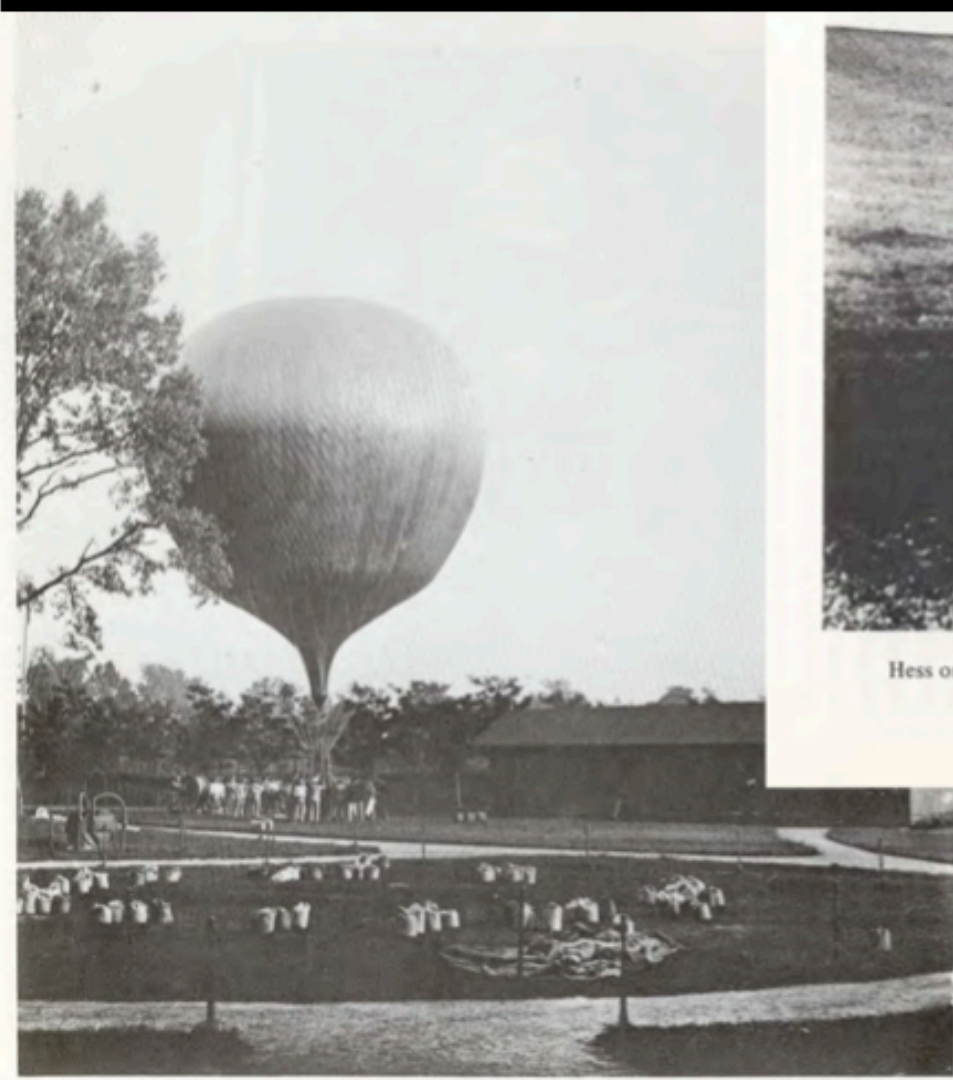
Glennys R. Farrar, New York University
INT workshop, Sept.10, 2025

$$\langle E \text{ per nucleon} \rangle \sim 2 \text{ EeV} \rightarrow \gamma > 10^9$$

Today

- Review of UHECR essentials
- Why BNS mergers are likely to be the site of UHECR production
- UHECR acceleration in the magnetized turbulent outflow of a BNS merger
- Predicting the UHECR spectrum and composition
- Tests
- What we need from YOU!

It's been a long, hard search



Aeronautisches Gelände im Wiener Prater, von dem aus V. F. Hess in den Jahren 1911/12 seine ersten Freiballon-Forschungsfahrten unternommen hatte. (Courtesy of Heeresgeschichtliche Museum, Vienna)
<Ed> Contributed by R. Steinmauer. See p. 17.



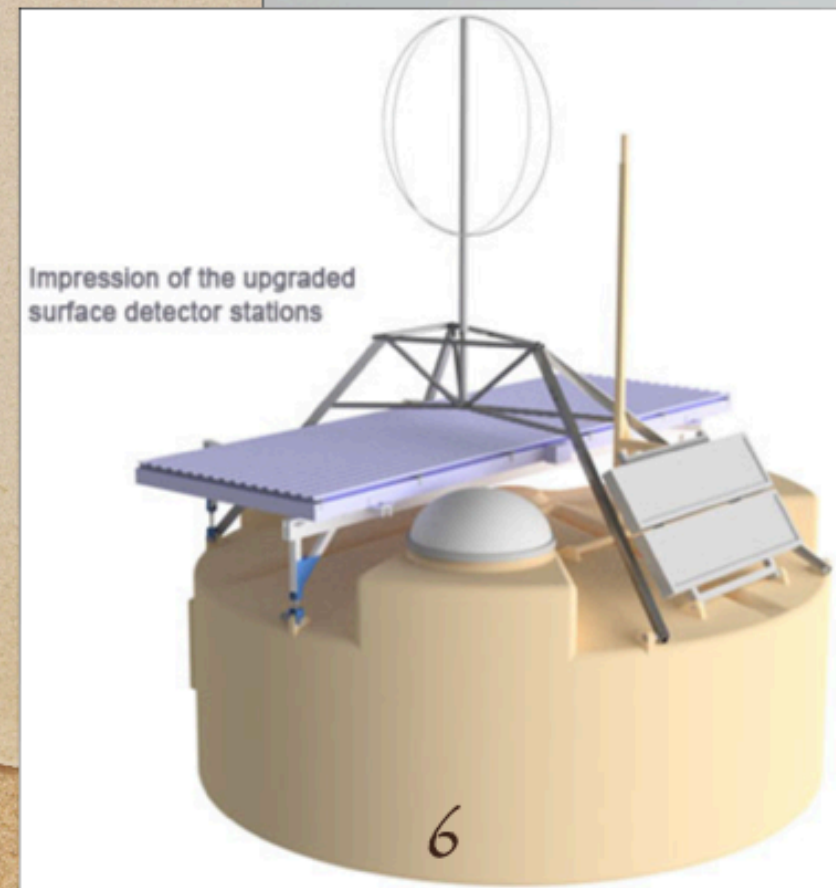
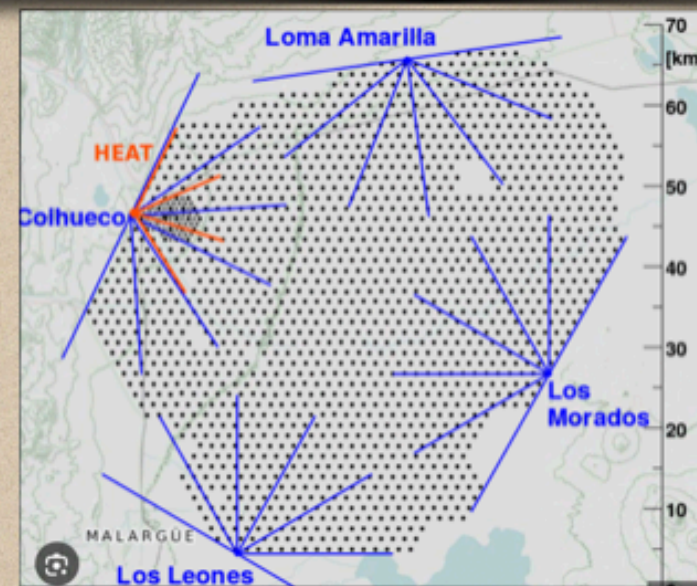
Hess on gondola in 1912 probably in test flight. The date and place is not clear at present.
<Ed> Contributed by R. Steinmauer. See p. 17.

Hess: CRs
1911 or 1912

3

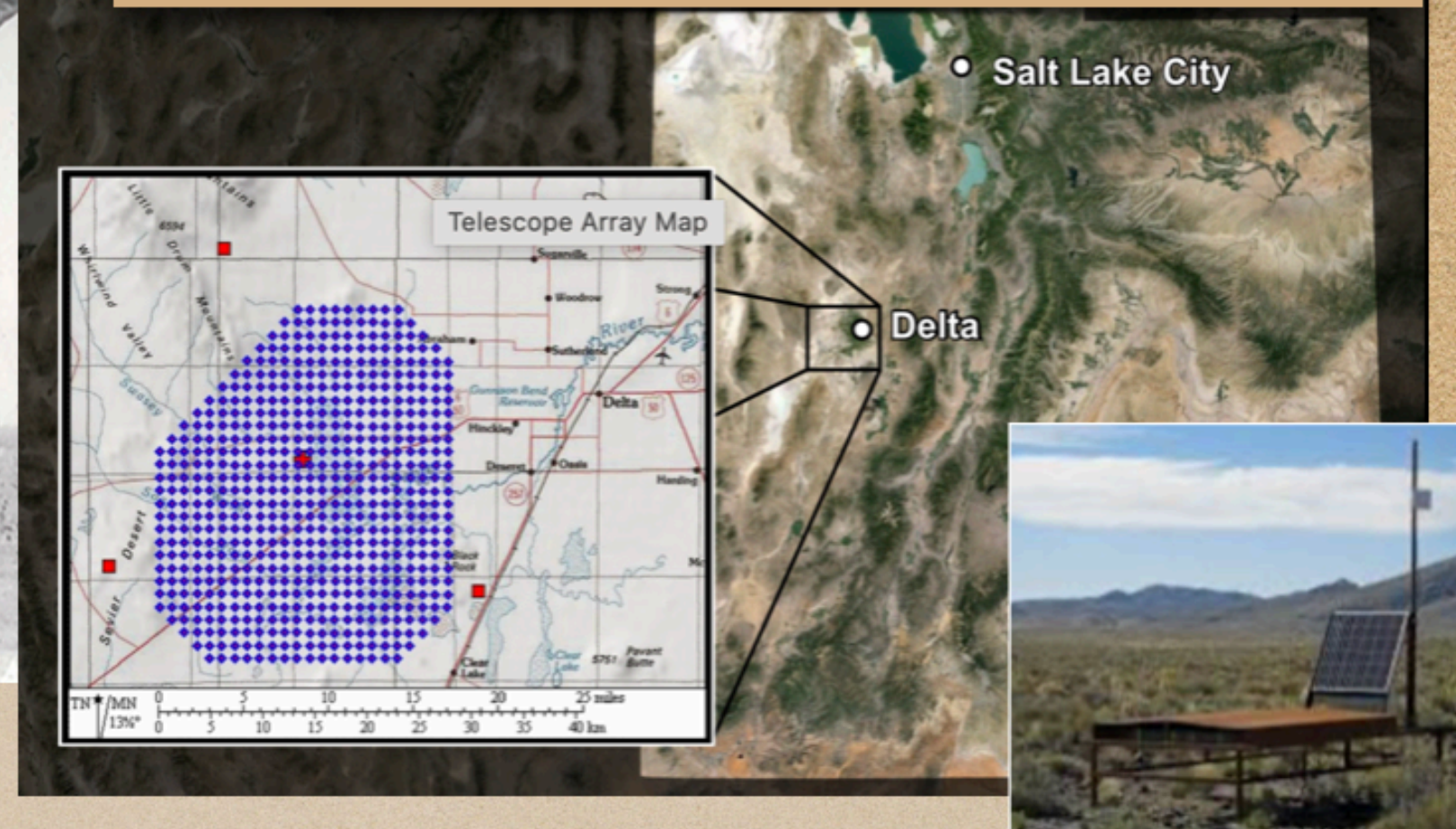


Linsley: 1st evt > 100 EeV
Volcano Ranch, NM ~1962



6

Telescope Array, Utah
Amaterasu ('23): 240 EeV



Pierre Auger Obs., Argentina
~50 evts > 100 EeV

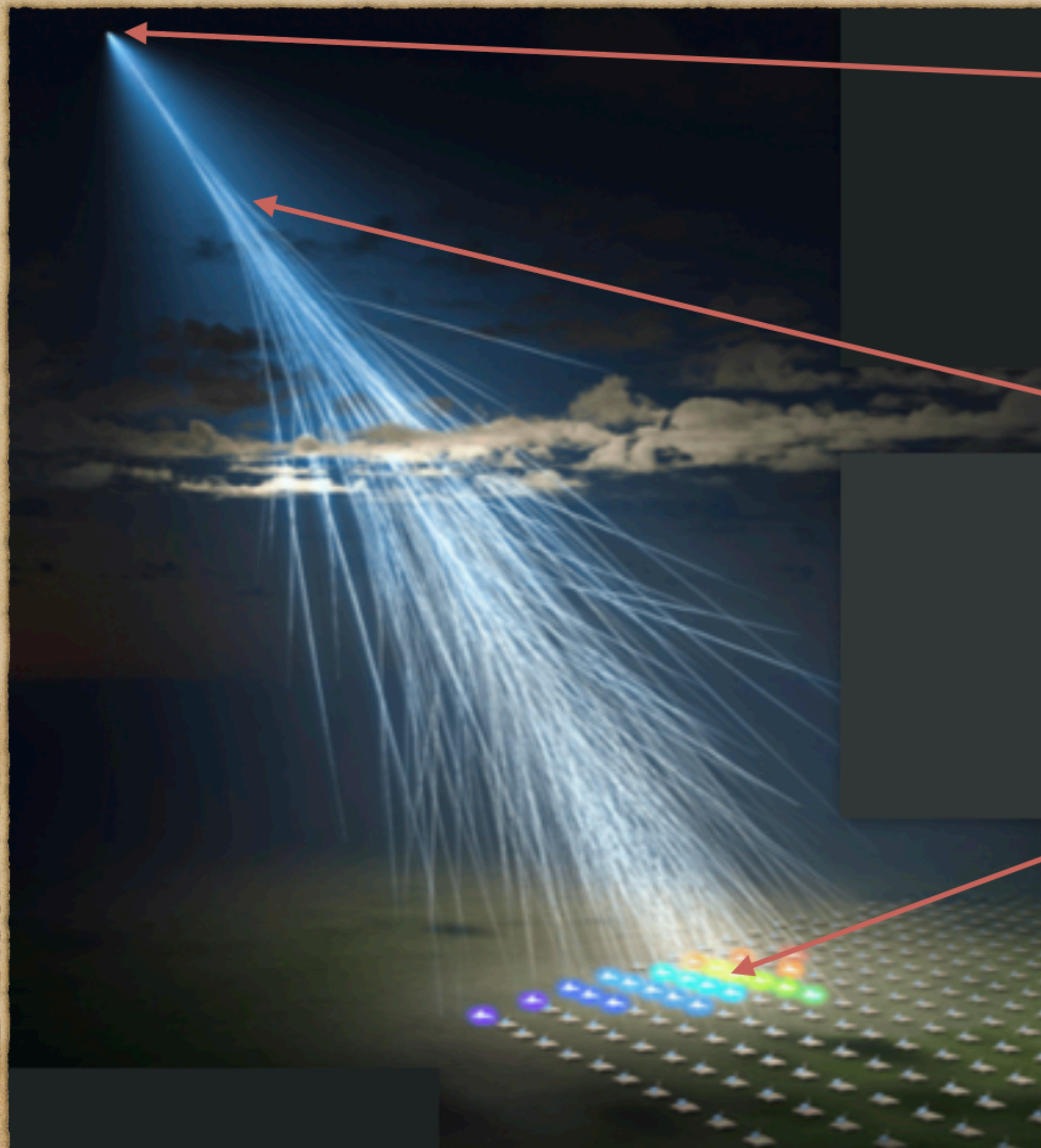


1700 stations, 3000 km^2



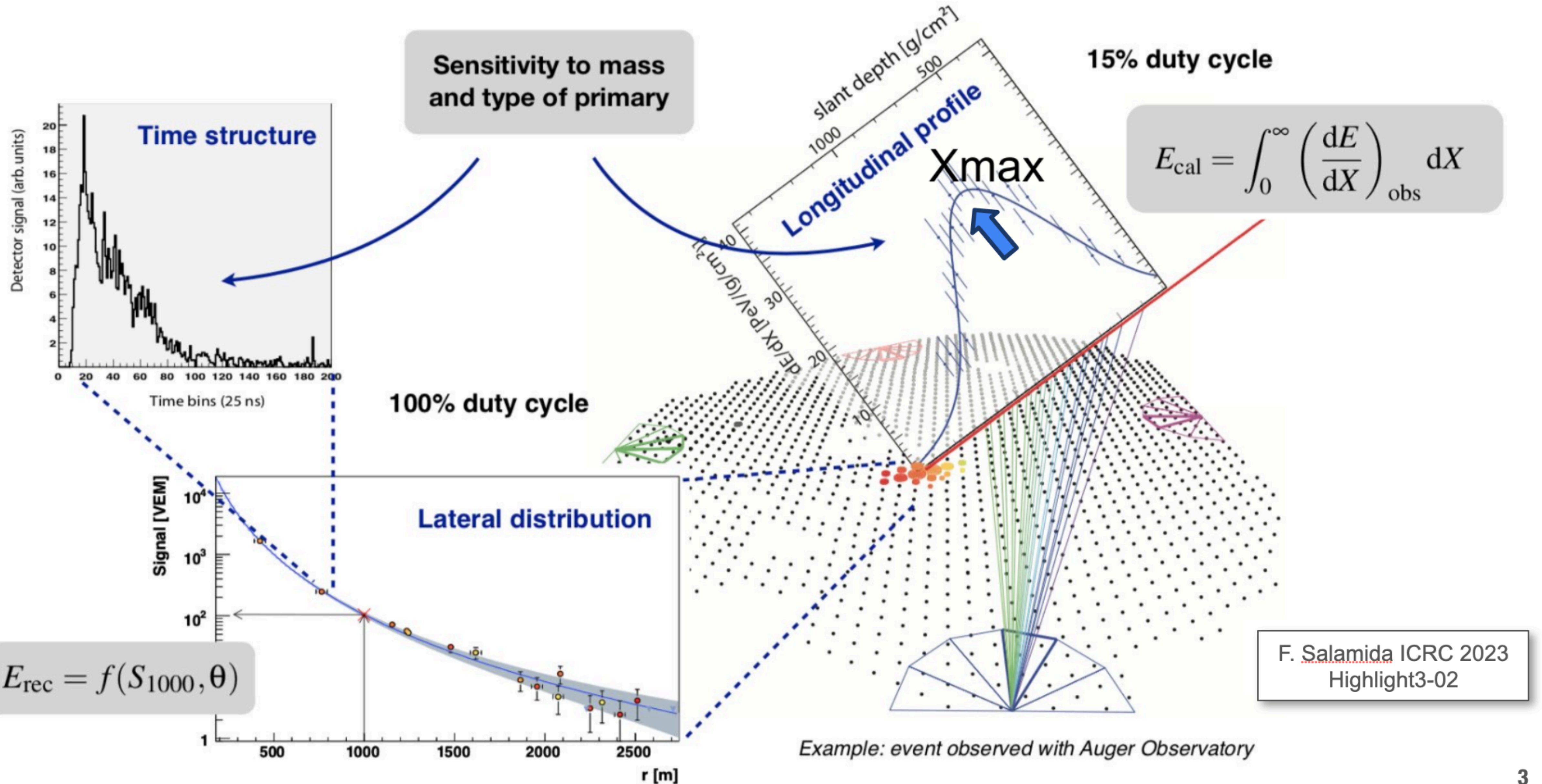
Fly's Eye Utah 1991
OMG: $320(250)$ EeV

How to deduce the mass and energy of a UHECR



- Depth of first interaction:
 - heavy nucleus: interacts quickly (starts high)
 - proton: 1st interaction is deep or shallow
- Shower development:
 - heavy nucleus: shower develops quickly
 - proton: more interactions needed to reach shower max
 - primary energy from integrated fluorescence emission
- Ground signal:
 - EM vs muon components \Rightarrow nuclear mass
 - primary energy from total signal

The Hybrid Observation Method of Auger

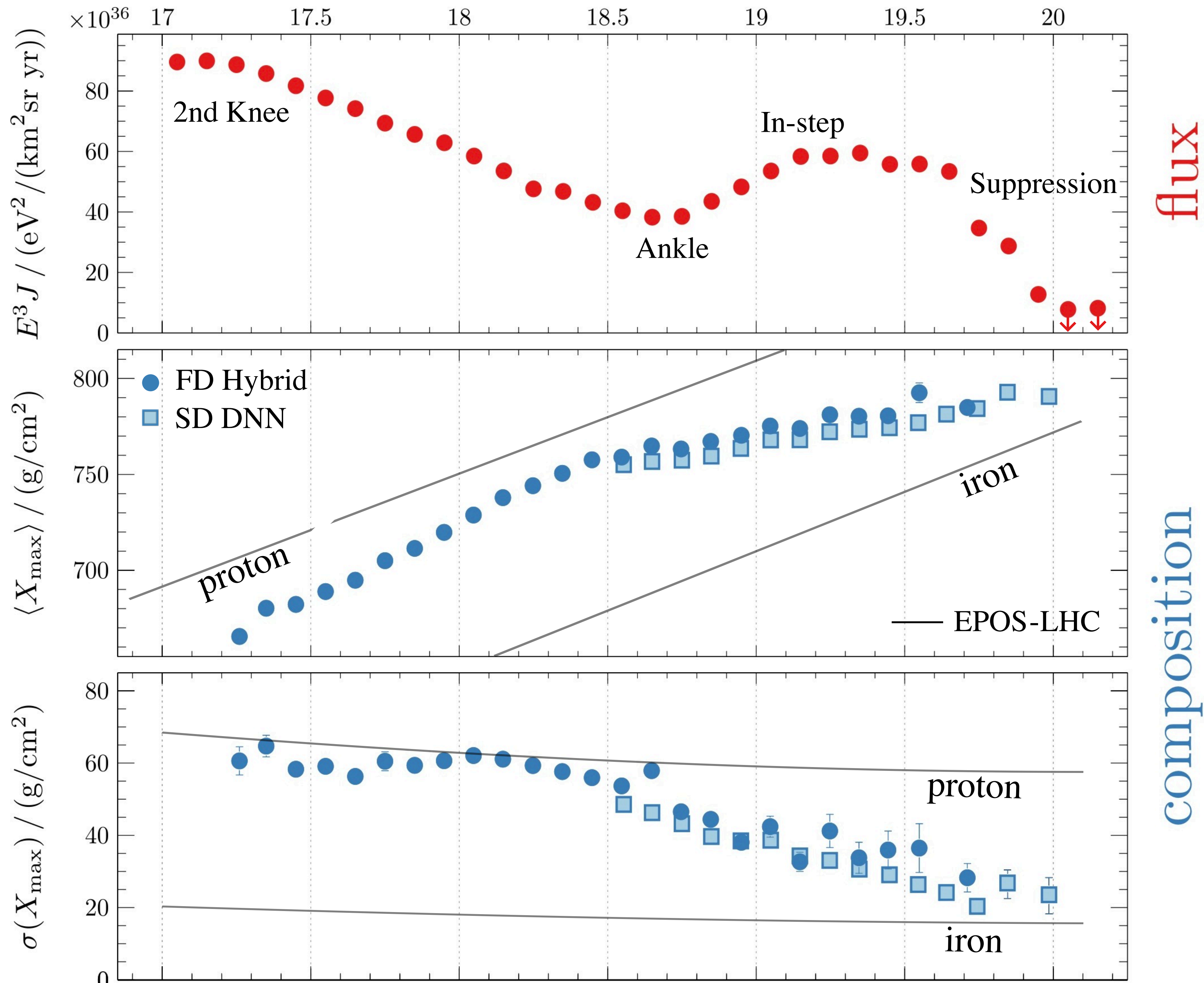
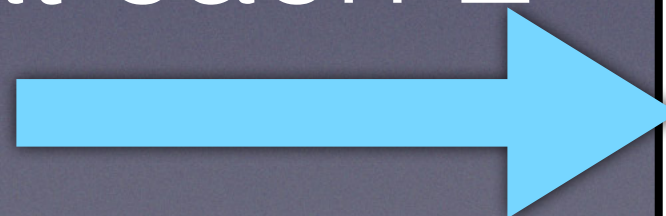


Spectrum, $\langle X_{\max} \rangle$
 $\sigma(X_{\max})$

Composition gets
heavier with energy

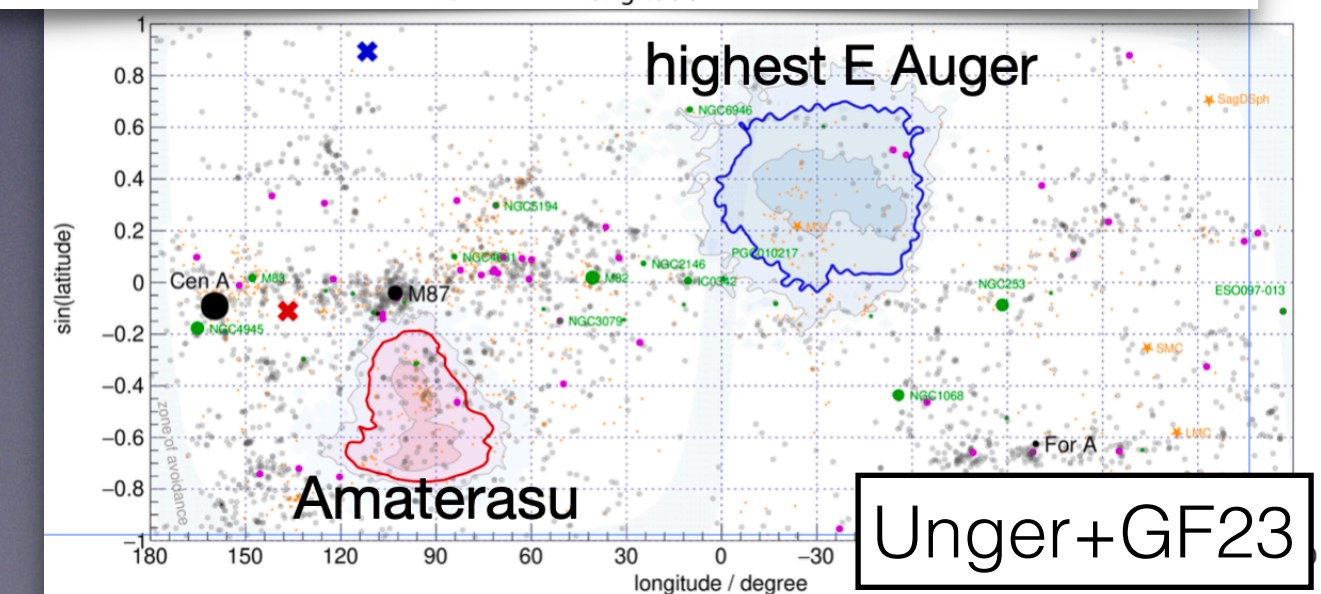
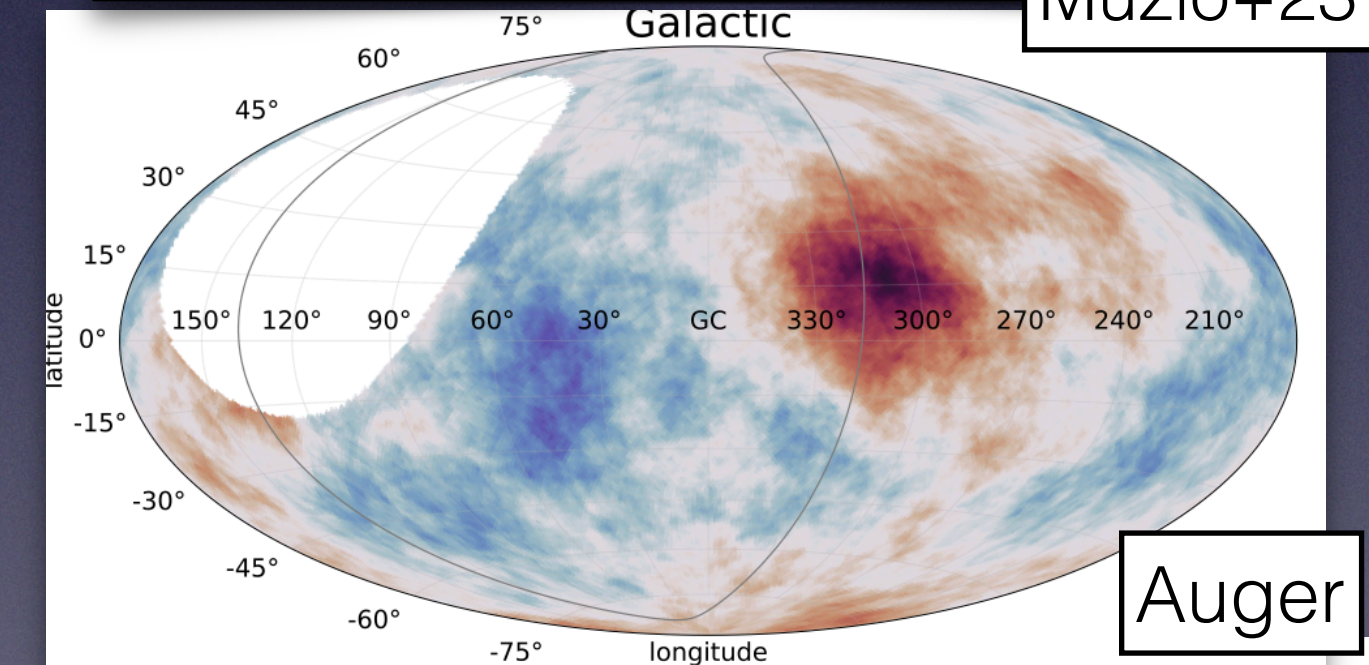
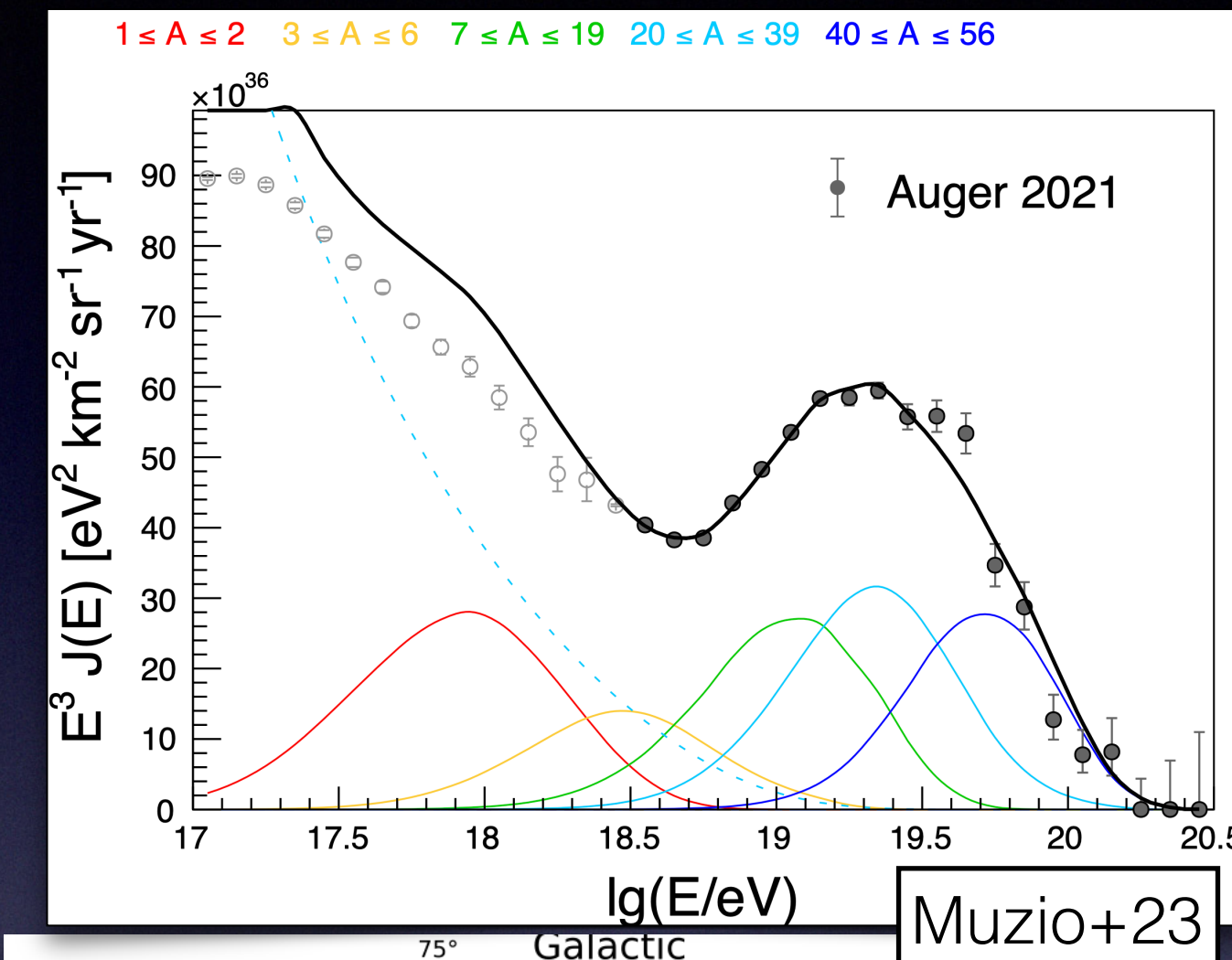


A single mass group
dominates at each E



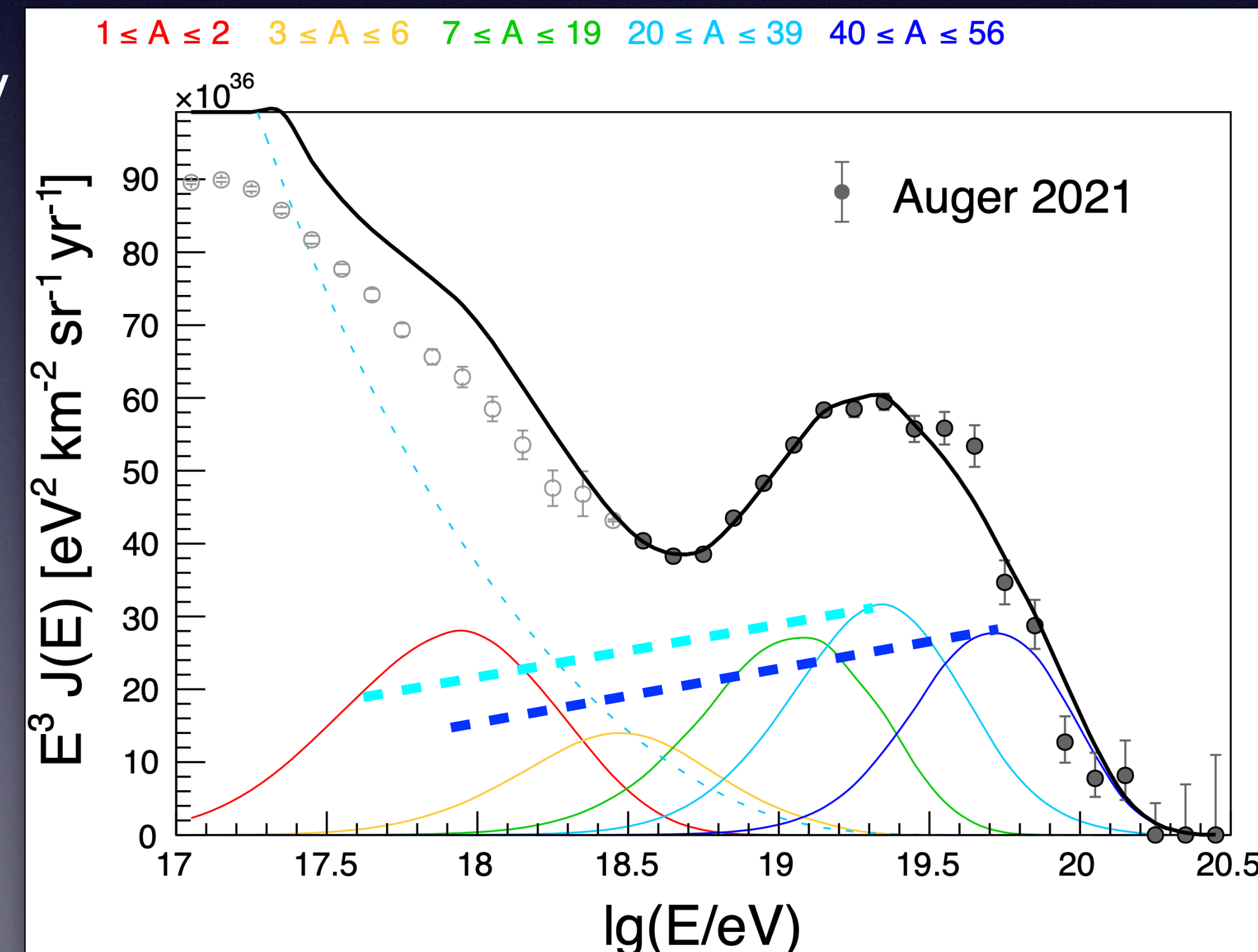
Essential results from observations

- Mixed composition
- *Narrow range of mass at each E :*
 - Rigidity spectrum is narrow. $\langle R \equiv E/Z_e \rangle \approx 4 \text{ EV}$
- Only significant multipole is dipole \rightarrow **sources are not rare.**
- No prominent sources \rightarrow **sources are transients.**



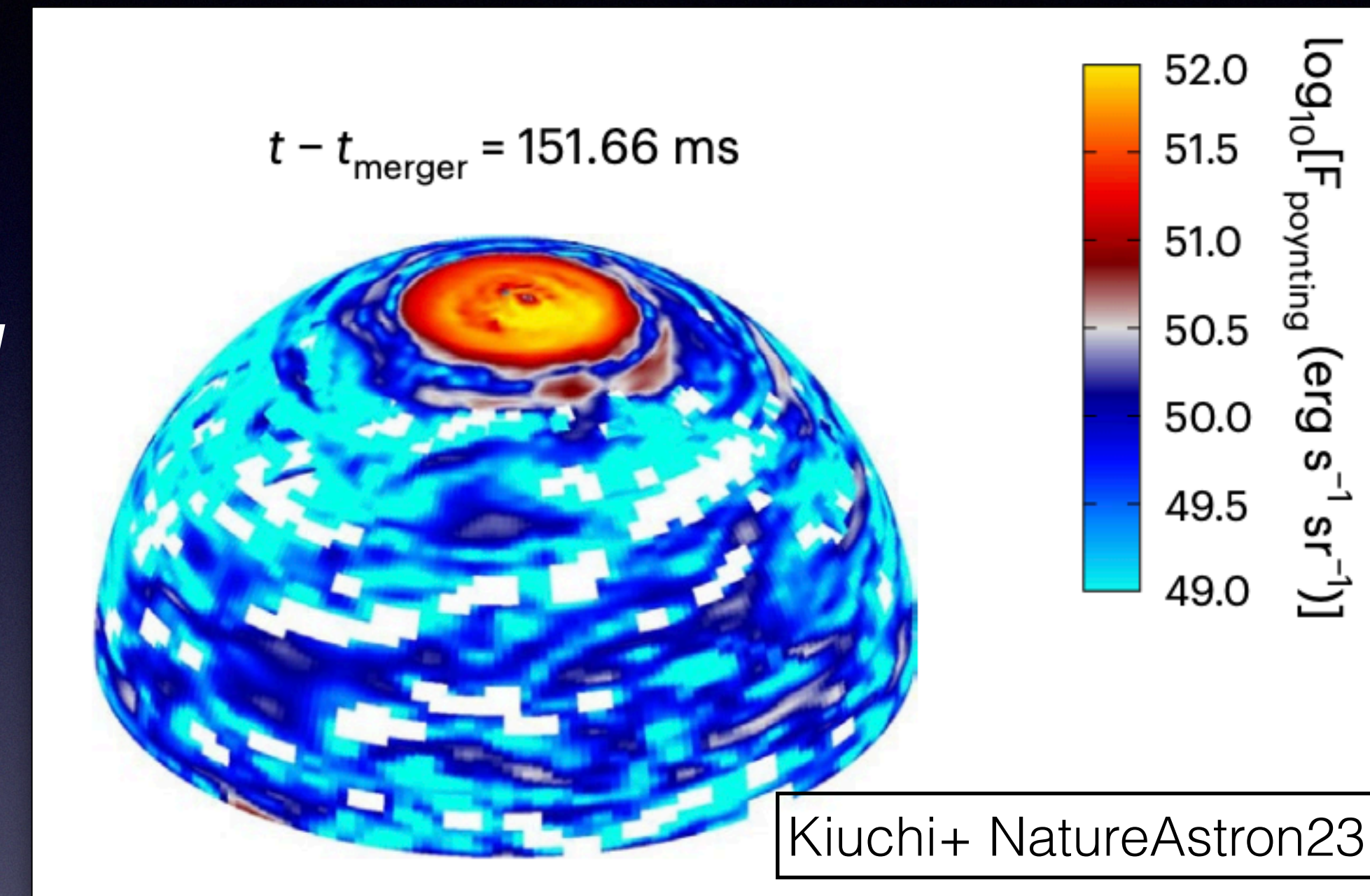
Correct Source model must explain

- Tight mass-energy relation **requires minimal population variance.** Excludes AGNs, long GRBs & TDEs. (Ehlert, Oikonomou, Unger 2023)
- Energy injection rate; source number density
- Maximum rigidity (Hillas criterion)
- Highest energy events: 150-250 EeV (Amaterasu, Fly's Eye, >10 Auger above 150 EeV)
- The highest energy events are from transients.



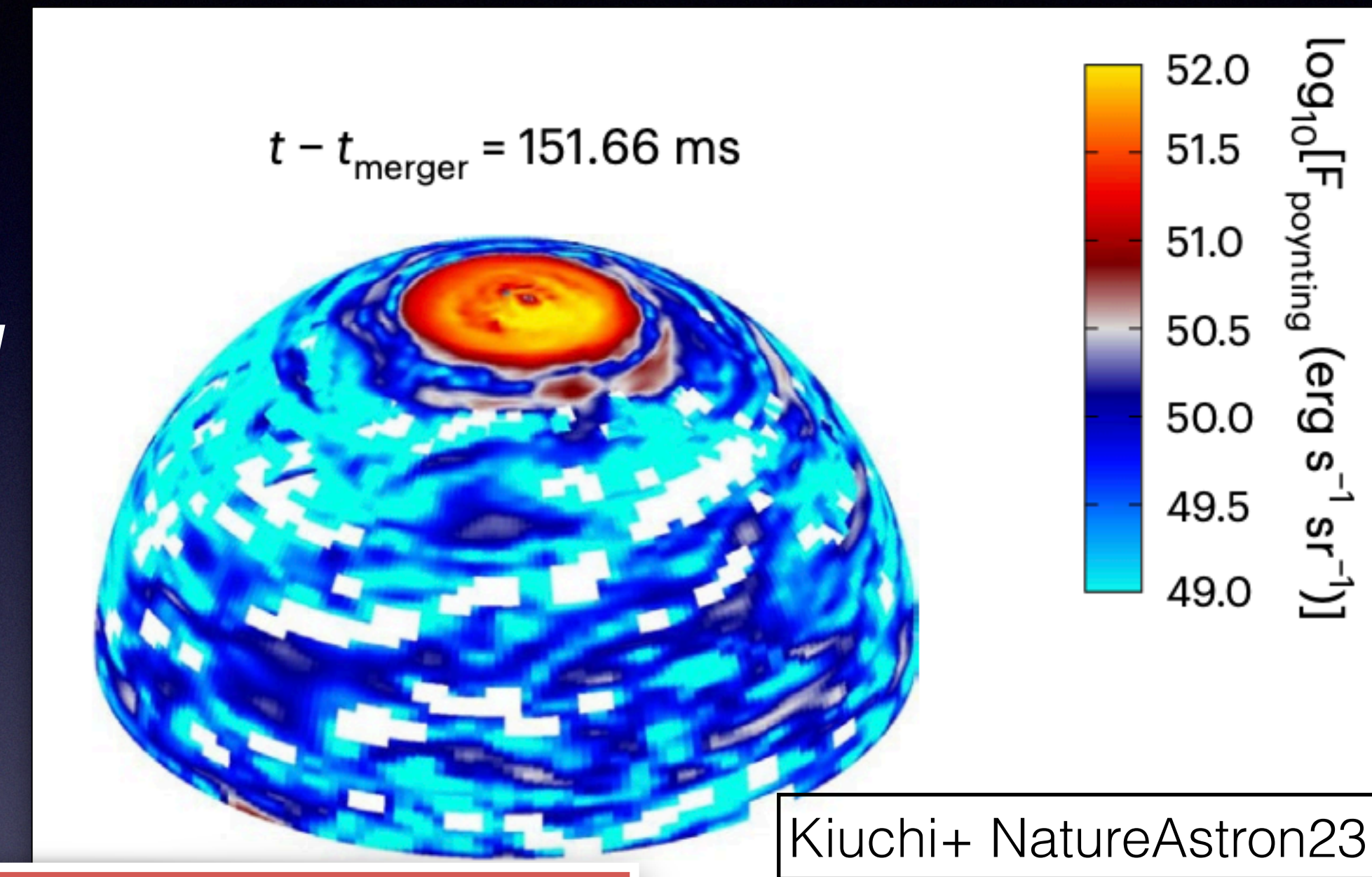
Binary Neutron Star Mergers: the only (currently known) source satisfying all criteria

- Universal rigidity spectrum explained because:
 1. Magnetic field is generated by *gravitational* dynamo.
 2. Mass range of BNS is narrow:
 - ♥♥ known BNS's: $M = 2.64 \pm 0.14$ (5%) M_{\odot}
 - $3.2 M_{\odot} \rightarrow 10\%$ (negligible) spread
- UHECR energy injection rate promising:
 - $\Gamma_{\text{merger}} = (0.3-1.7)10^{-6} \text{ Mpc}^{-3} \text{ yr}^{-1}$; need UHECR energy per merger $\approx 10^{50} \text{ erg}$ (total energy emitted $> 10^{52.5} \text{ erg}$).



Binary Neutron Star Mergers: the only (currently known) source satisfying all criteria

- Universal peak rigidity explained because:
 1. Magnetic field is generated by *gravitational* dynamo.
 2. Mass range of BNS is narrow:
known BNS's: $M = 2.64 \pm 0.14$ (5%) M_{\odot}



- UHECR energy
- $\Gamma_{\text{merger}} = (0.3 \dots 1) \times 10^{50}$
erg (total energy)

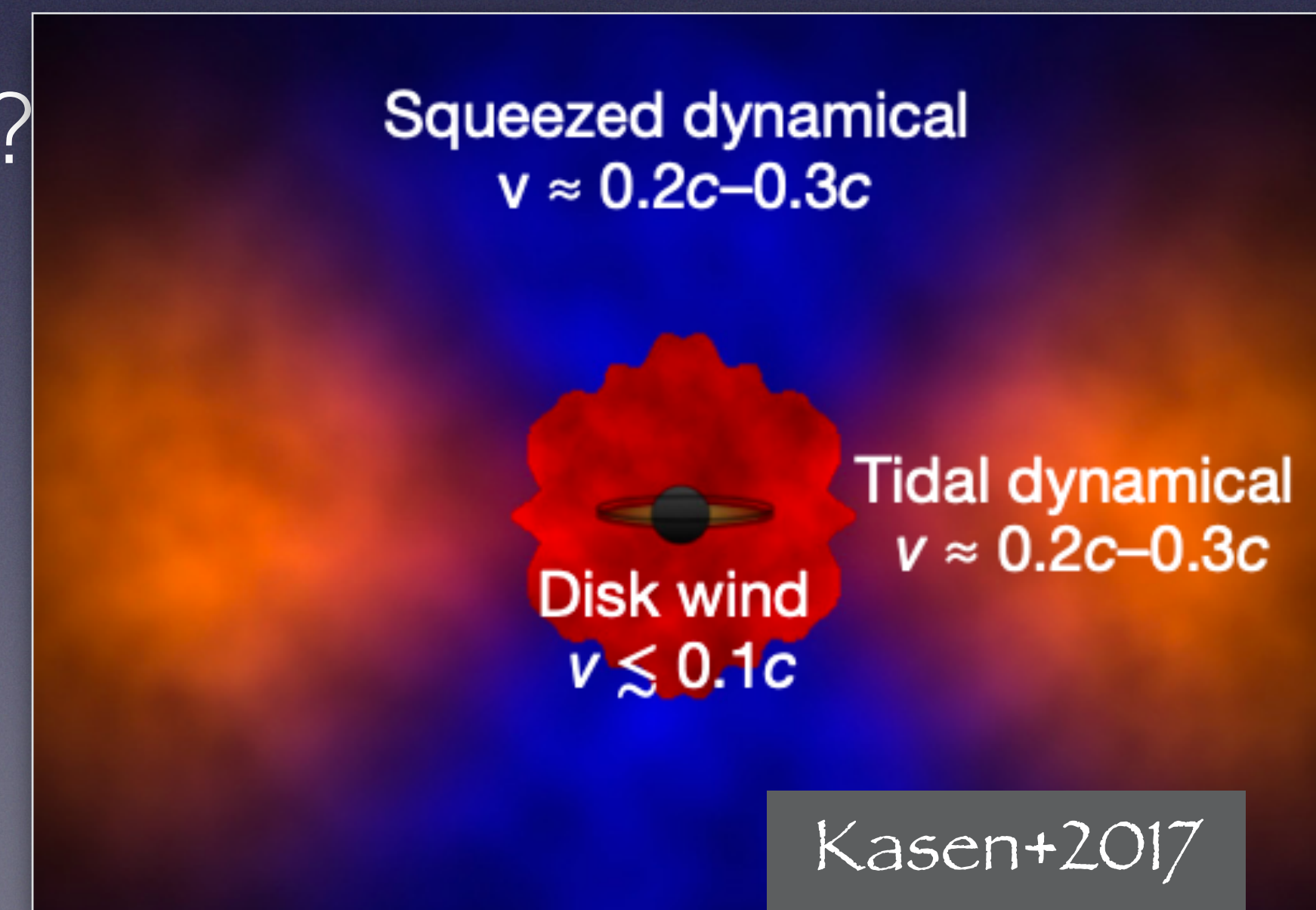
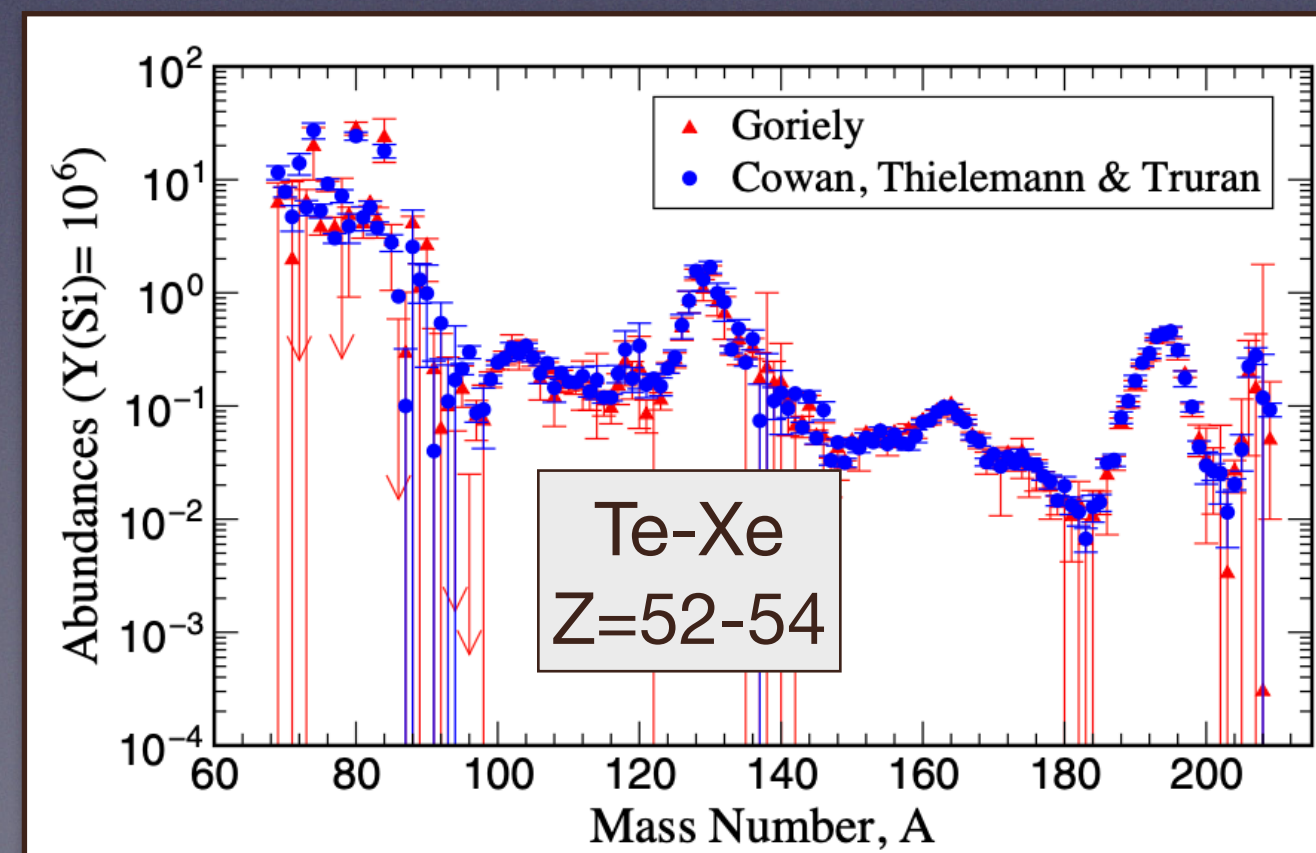
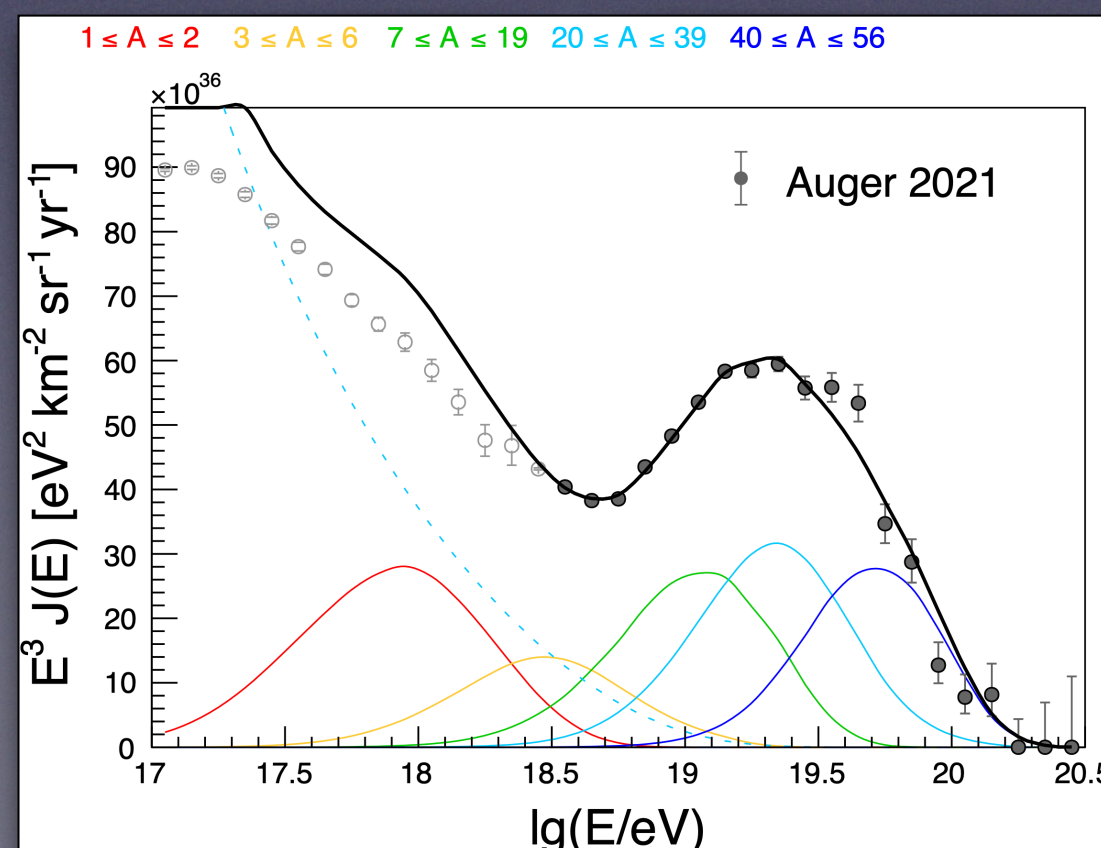
Unique predictions enable definitive tests:

- EHE ν 's \leftrightarrow gravitational waves
- Highest energy UHECRs: $Z > 26$
- BNS merger \rightarrow initial B \rightarrow predict spectrum

merger $\approx 10^{50}$

Topics for today

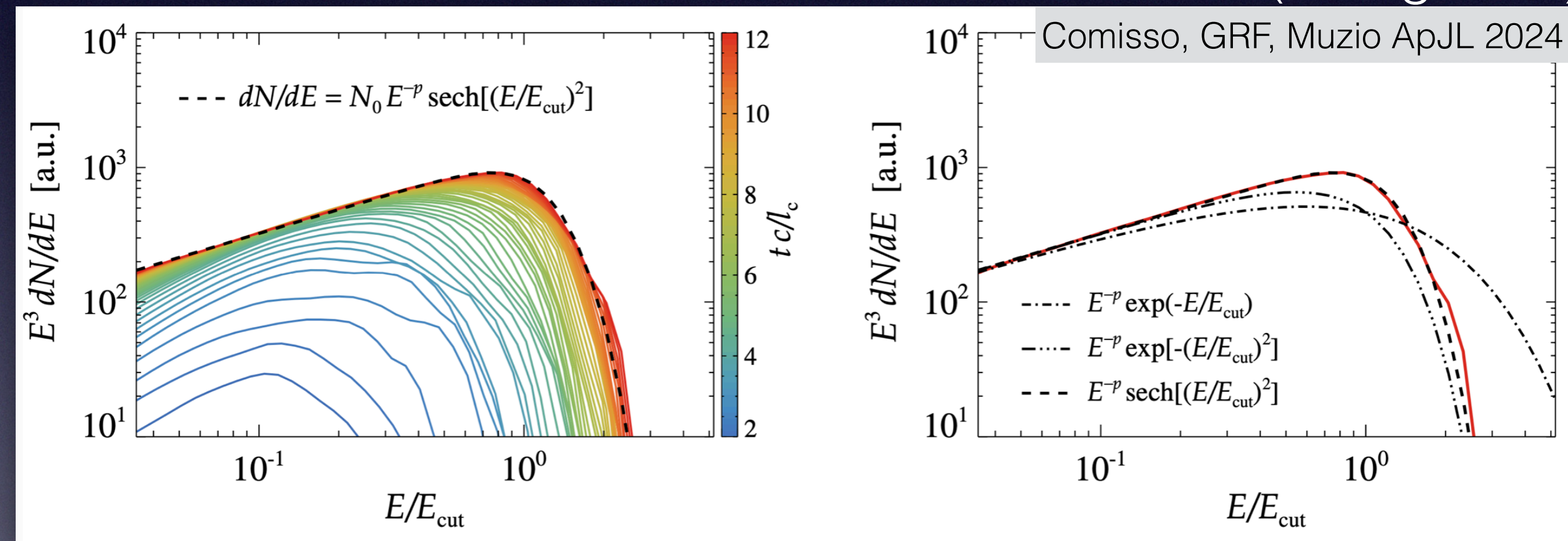
- Where in the merger ejecta does UHECR acceleration occur? → Predict rigidity cutoff R_{cut} and spectral shape
- What is the time profile for UHECR production? (coincidences between GW and EHE ν 's...)
- What are the relative abundances of different nuclei?



UHECRs mainly produced in the turbulent outflow (maybe some UHECRs also from jets)

- *Composition:* UHECRs have masses up to (at least) Fe, but jets only have p, He (Perego+22)

- *Spectrum:* $E^{-2.5}$ \times function of (E/E_{cut})



- Outflow: Magnetized turbulence cutoff: $\text{sech}[(E/E_{\text{cut}})^2]$ (Comisso, GRF, Muzio ApJL 2024)
- Jets: Diffusive Shock Acceleration cutoff: $\exp(-E/E_{\text{cut}})$??? (Protheroe+Stanev 1997)

UHECR data strongly favor $\text{sech}[(E/E_{\text{cut}})^2]$ cutoff

Big Picture

time

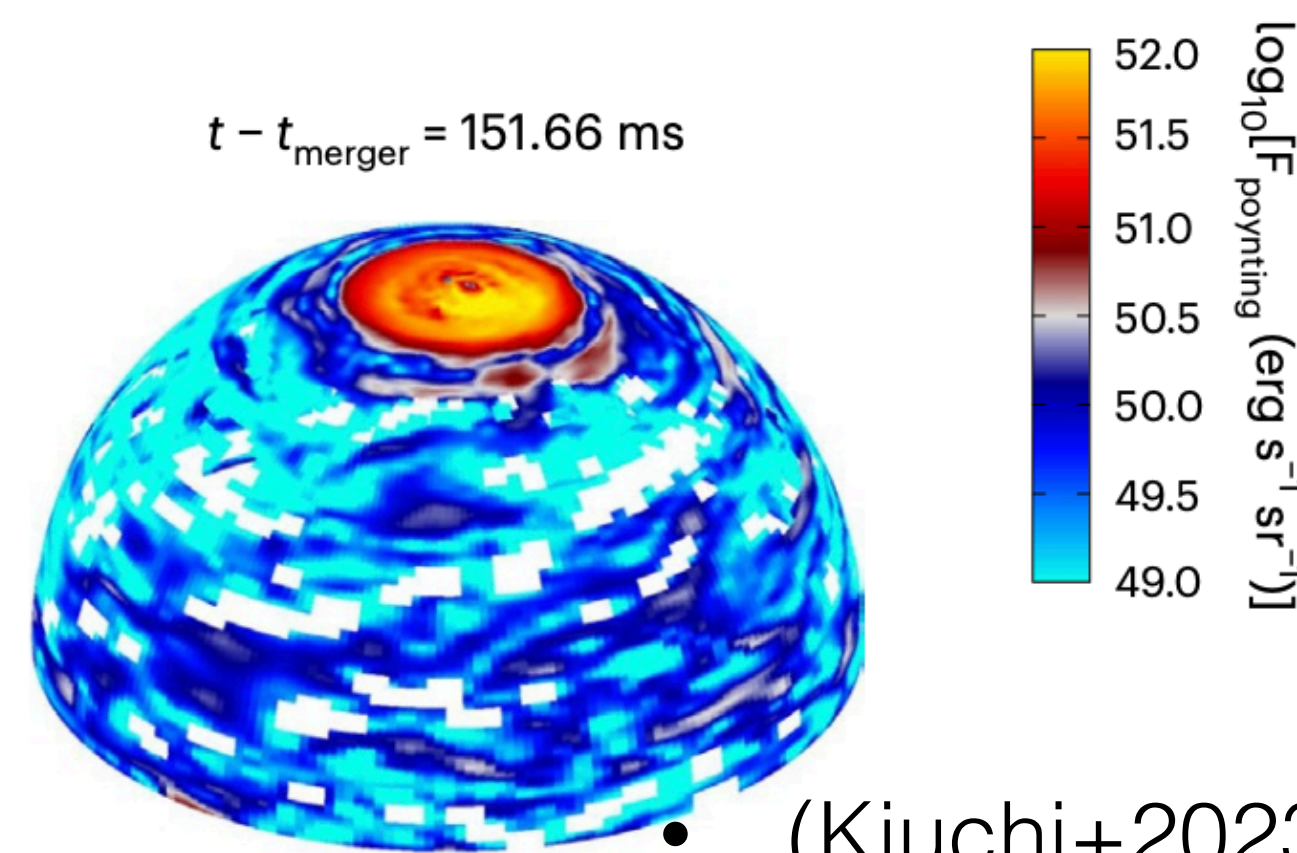
inspinal
GRMHD dynamics
→ huge B

**Gravitational Wave
Emitted**

"Hypervnova"
0-few sec
jet
V ~ 0.1 c
outflow
10°

initialization of B
150ms, 500 km

Kiuchi+23 → B_{ini}



• (Kiuchi+2023)

$r \sim 10^{10} \text{ cm}$

cools to ~ 1 MeV
nucleosynthesis
~ 1 s: $r \sim 10^{10} \text{ cm}$

r-process
nucleosynthesis

$r \sim 10^{14} \text{ cm}$
 $B \approx 400 \text{ G}$

UHECR's
accelerated

B drops till synch losses
are subdominant

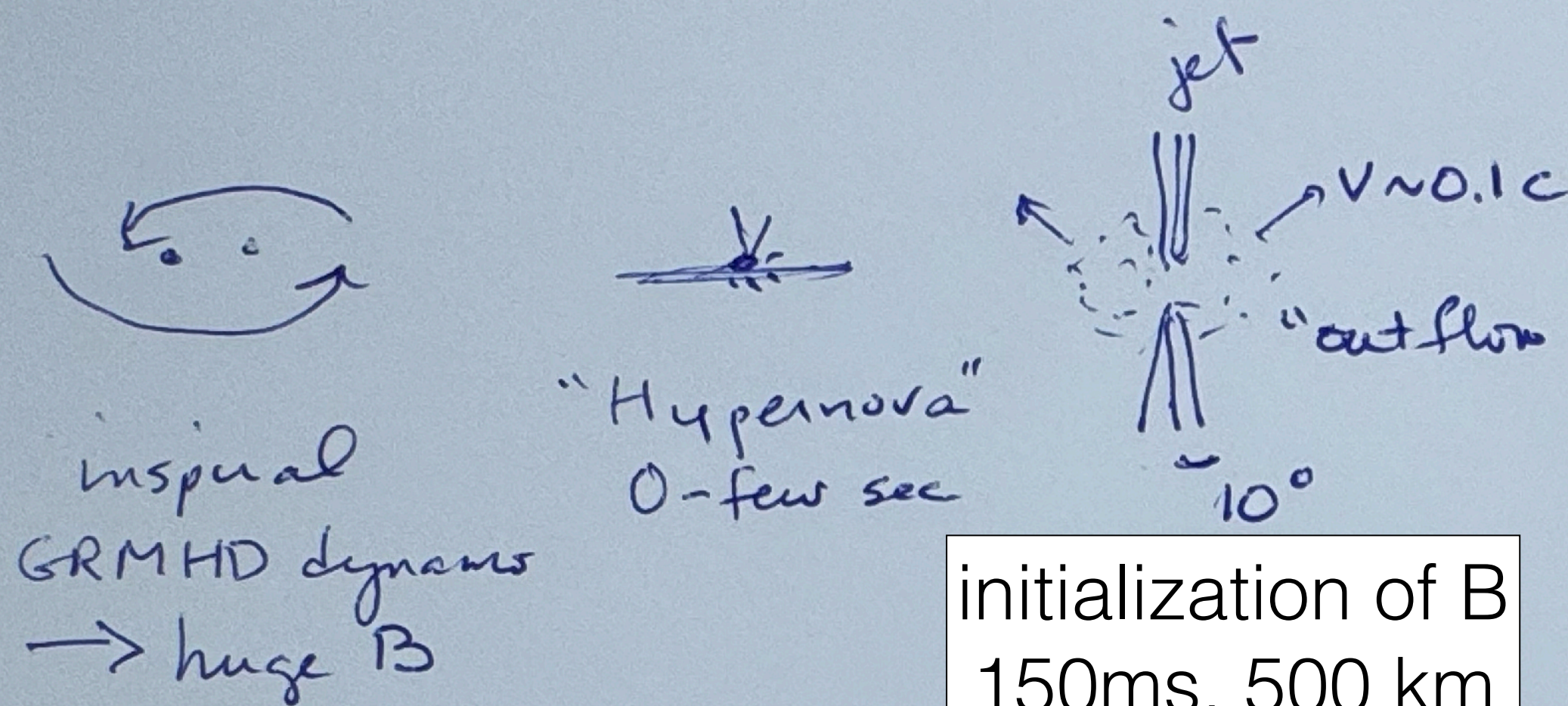
UHECR accel

~ 1 day: 10^{14} cm

$\tau_{\text{accel}} / \tau_{\text{synch-loss}} \sim 1$

Big Picture

time



initialization of B
150ms, 500 km

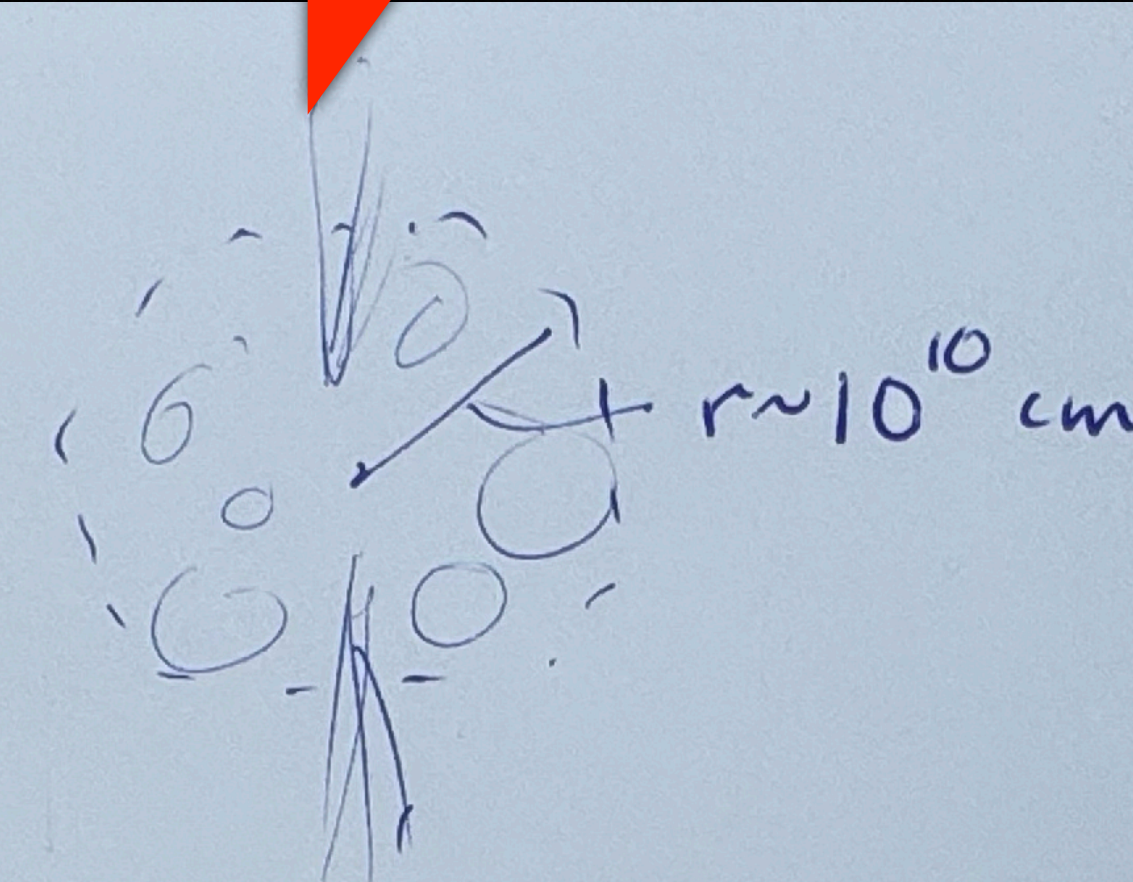
Gravitational Wave
Emitted

Kiuchi+23 → B_{ini}

$t - t_{\text{merger}} = 151.66 \text{ ms}$

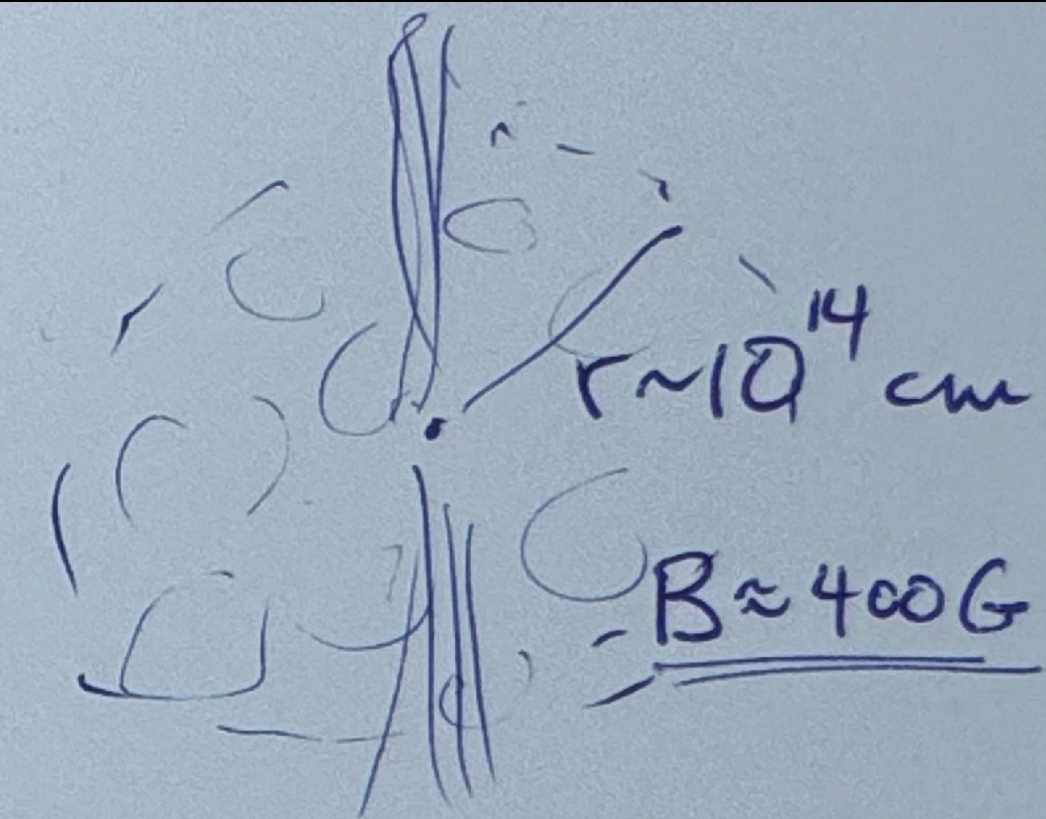


(Kiuchi+2023)



cools to ~ 1 MeV
nucleosynthesis
~ 1 s: $r \sim 10^{10} \text{ cm}$

r-process
nucleosynthesis



UHECR's
accelerated

B drops till

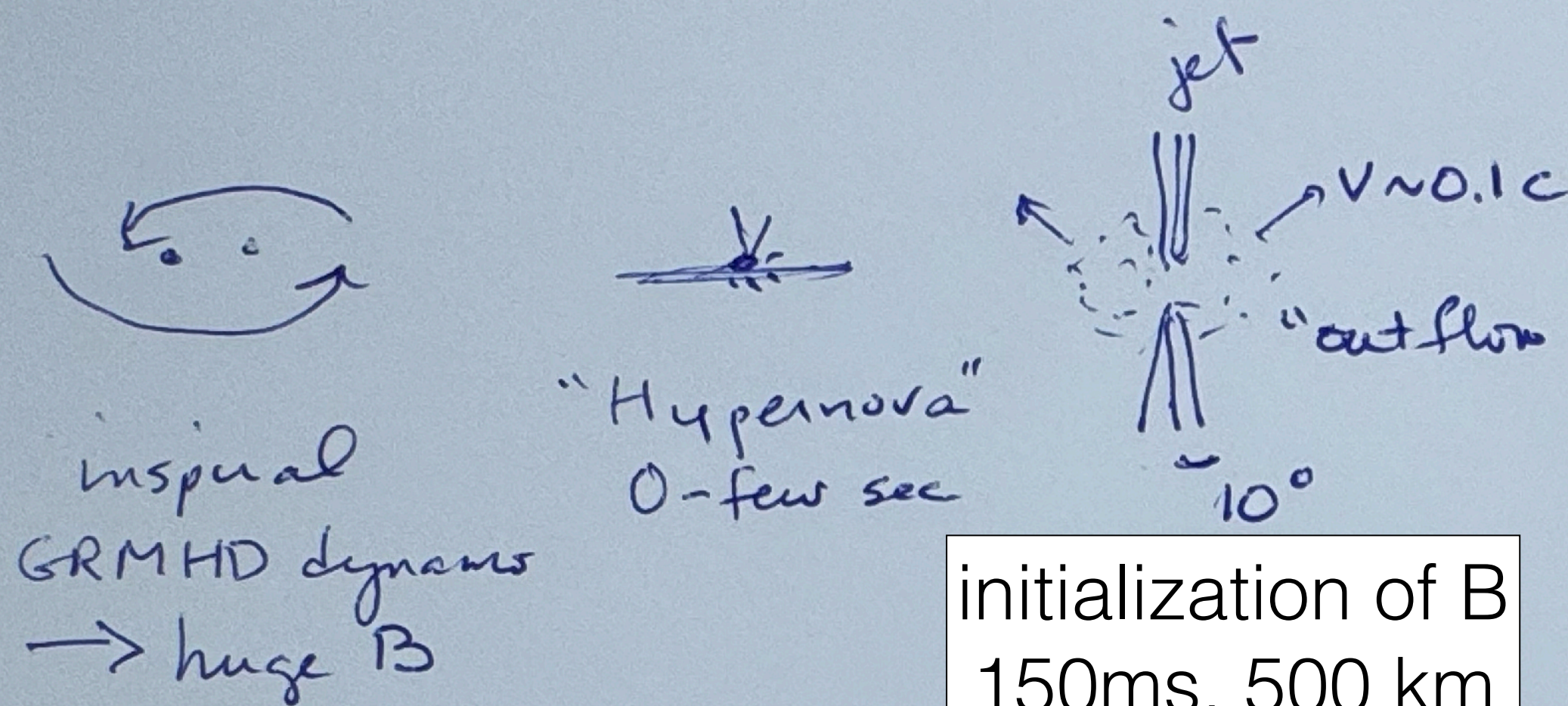
UHECR accel
~ 1 day: 10^{14} cm

Duration of "hypernova"

- depends on NS EoS
- p, He balance in jet
- total energy output

Big Picture

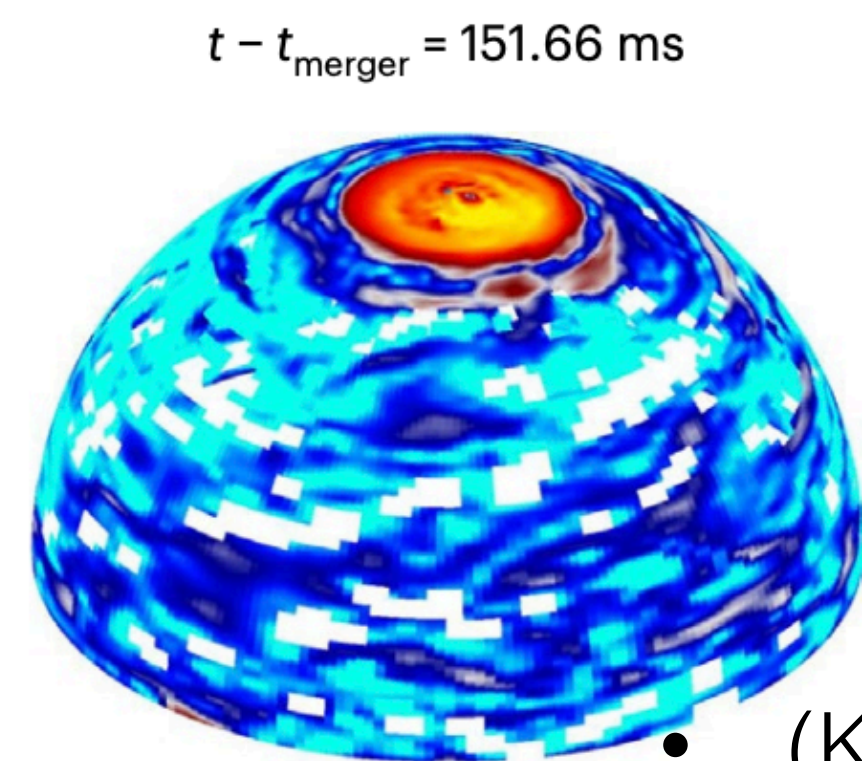
time



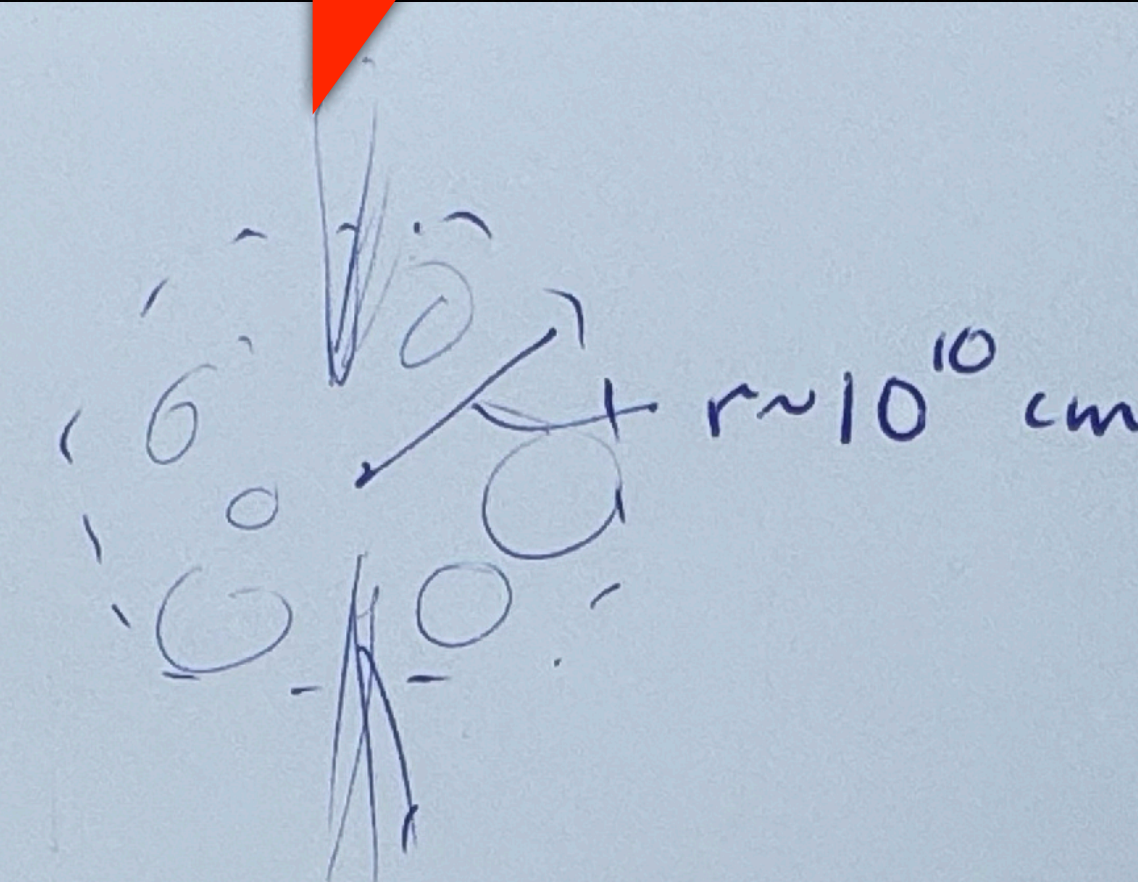
Gravitational Wave
Emitted

initialization of B
150ms, 500 km

Kiuchi+23 → B_{ini}



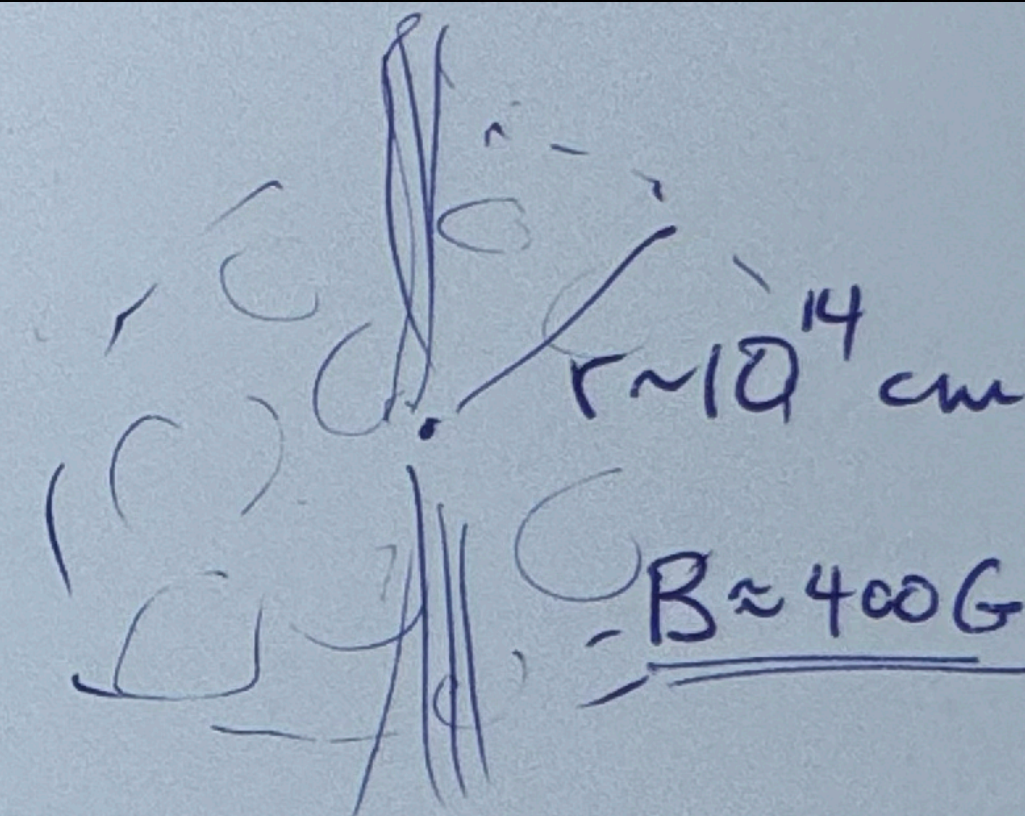
(Kiuchi+2023)



cools to ~ 1 MeV
nucleosynthesis
~ 1 s: $r \sim 10^{10}$ cm

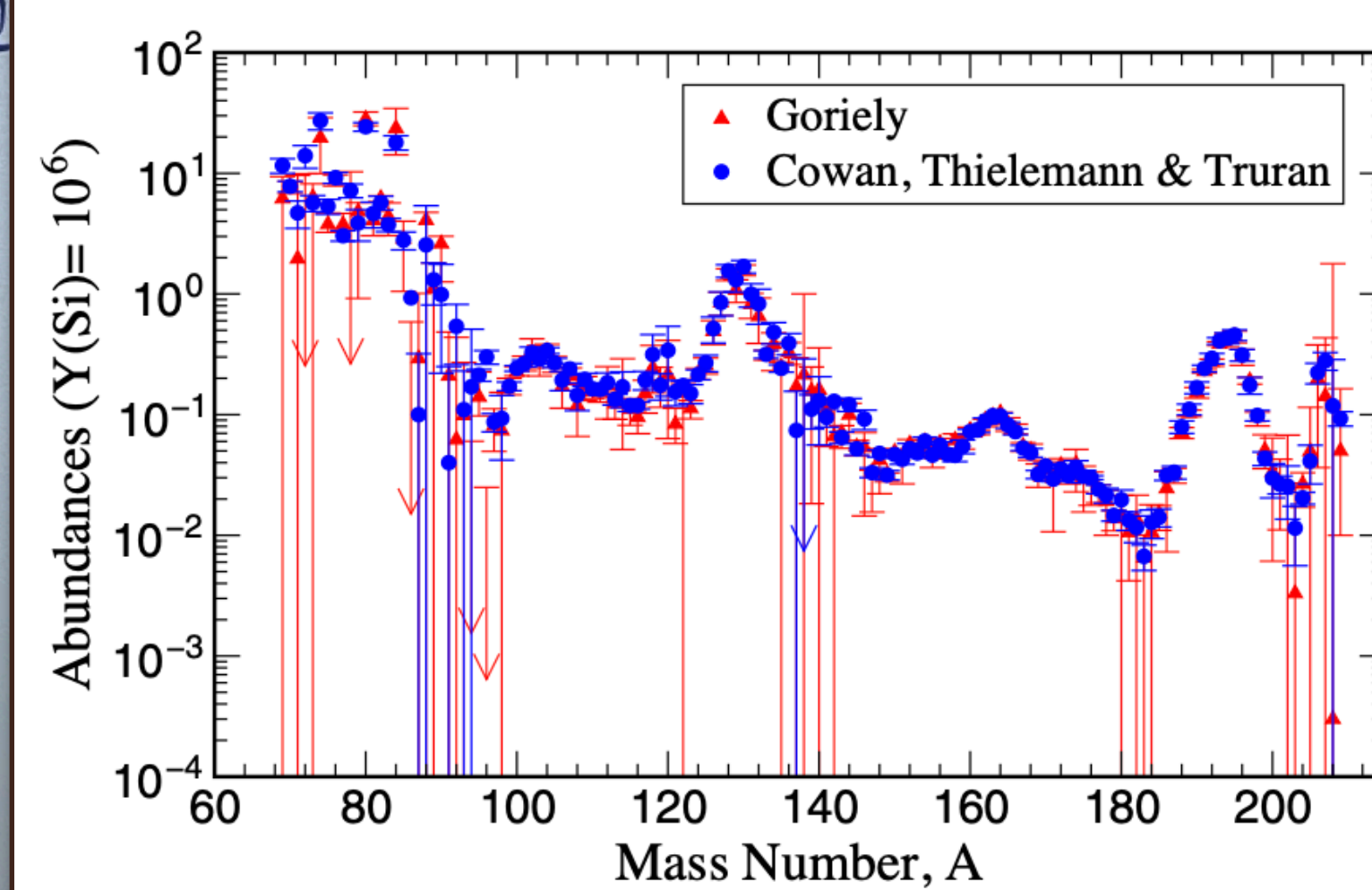
Nucleosynthesis:

- ▶ p, He in jet
- ▶ $Z, N \geq 20$ elsewhere



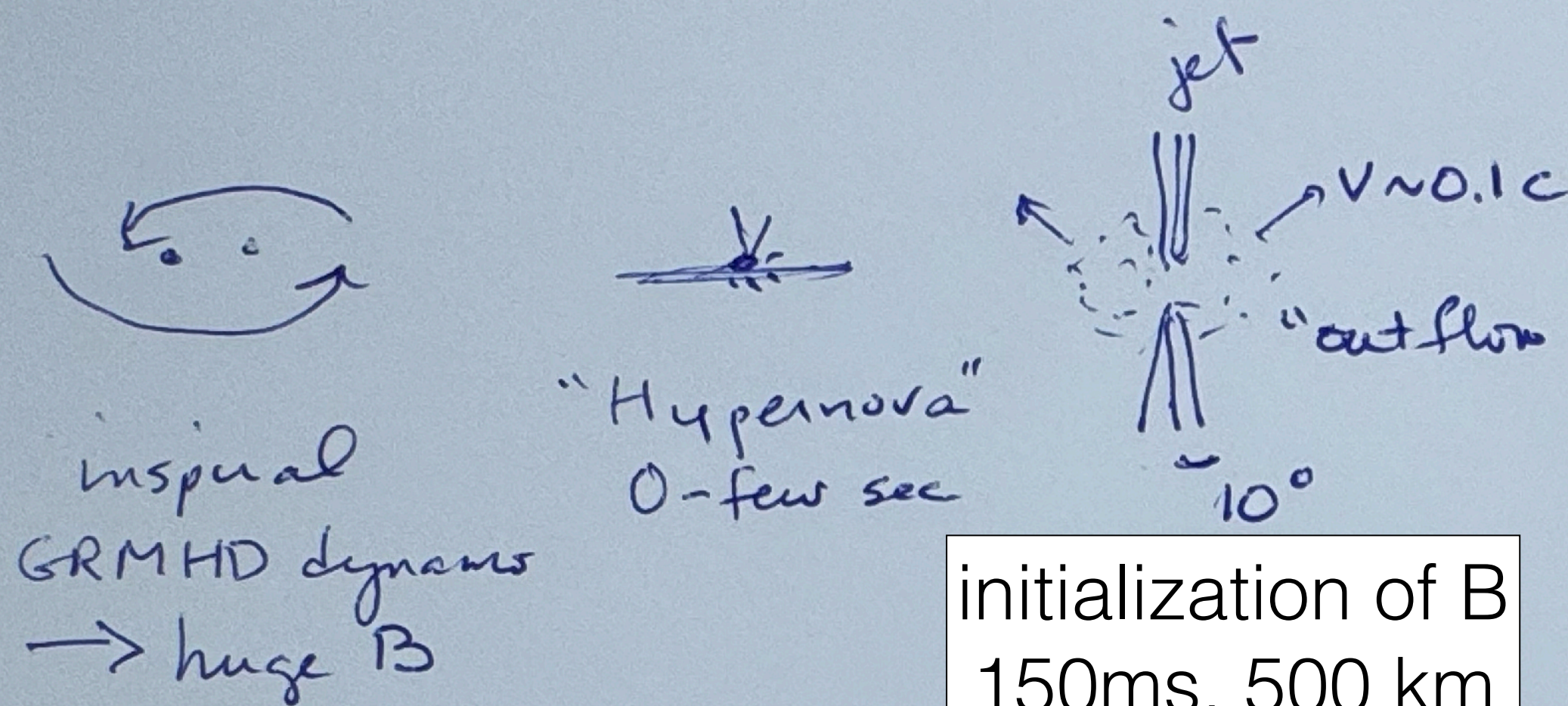
UHECR's
accelerated

~ 1 day



Big Picture

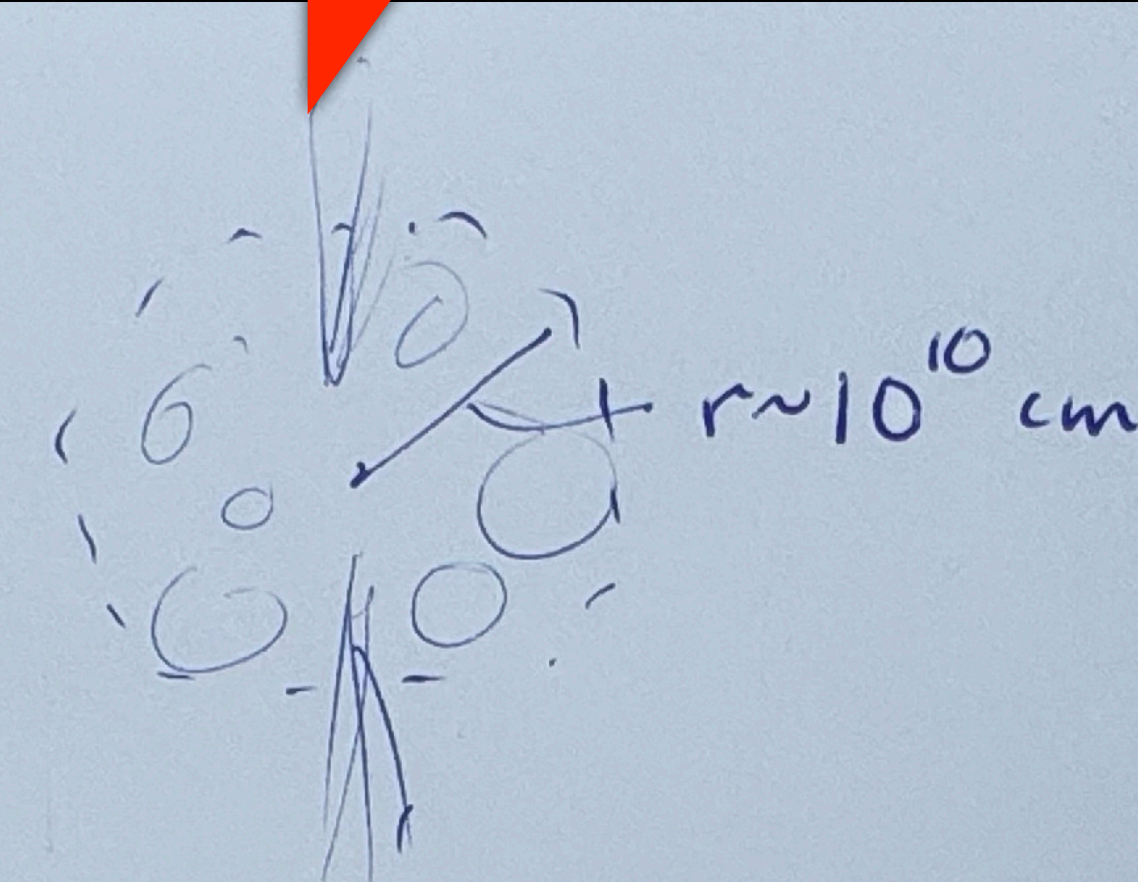
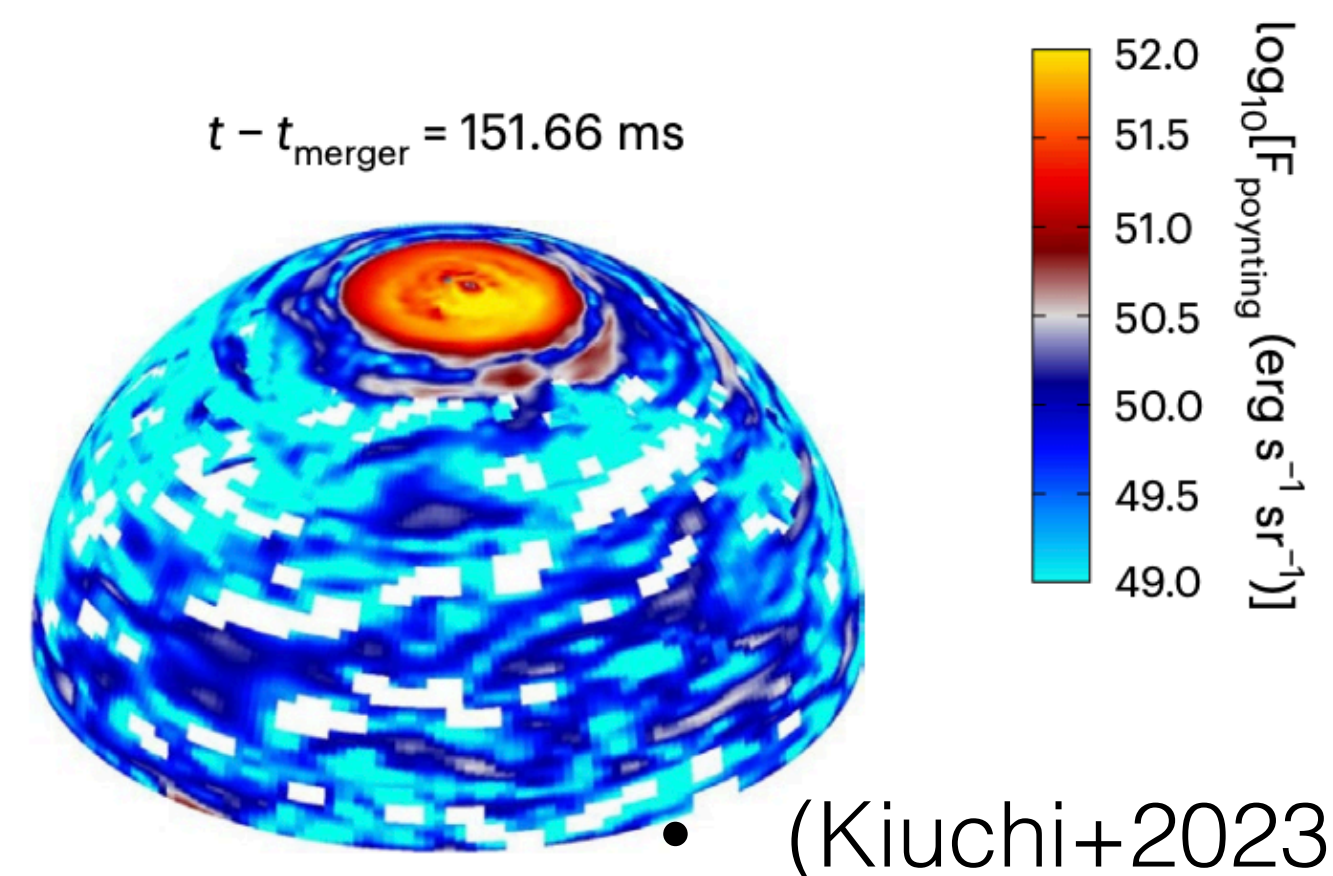
time



initialization of B
150ms, 500 km

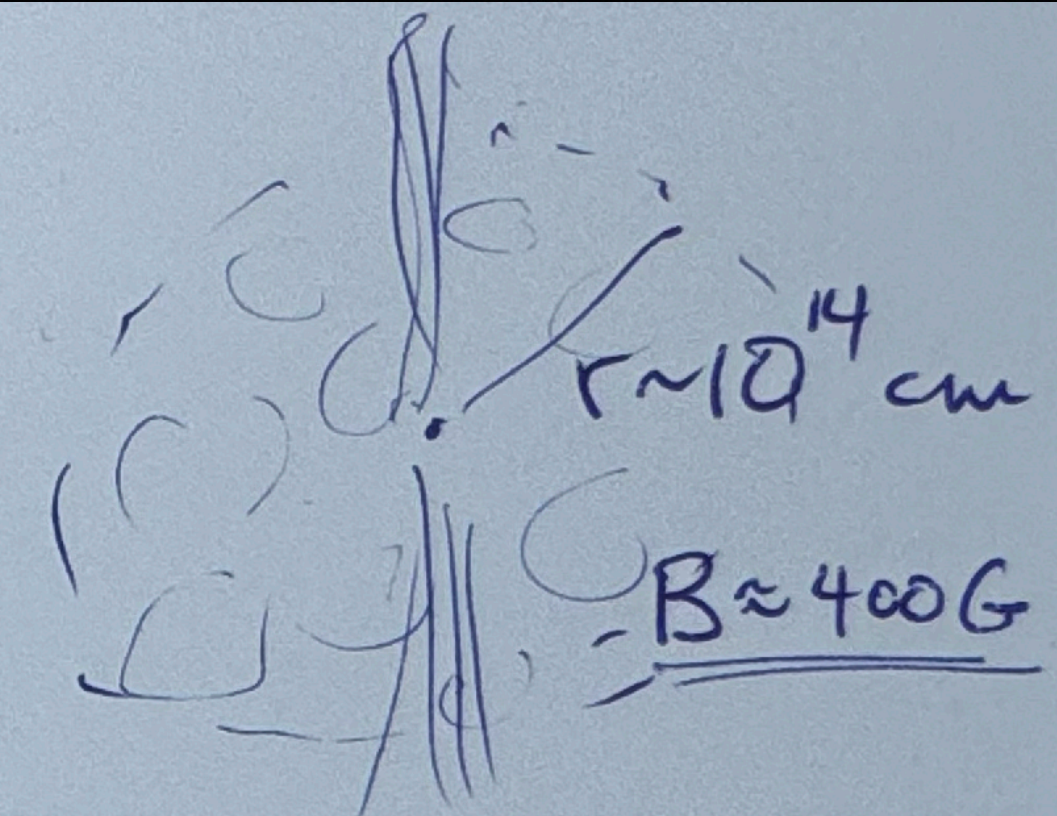
Gravitational Wave
Emitted

Kiuchi+23 → B-in



cools to $\sim 1 \text{ MeV}$
nucleosynthesis
 $\sim 1 \text{ s}$: $r \sim 10^{10} \text{ cm}$

r-process
nucleosynthesis



UHECR's
accelerated

B drops till

$\tau_{\text{accel}} / \tau_{\text{synch-loss}} \sim 1$

- some nuclei accelerated to $\text{KE} > 100 \text{ MeV}$
- collisions → breakup
 - ▶ range of A's
 - ▶ fixes composition

Big Picture

time

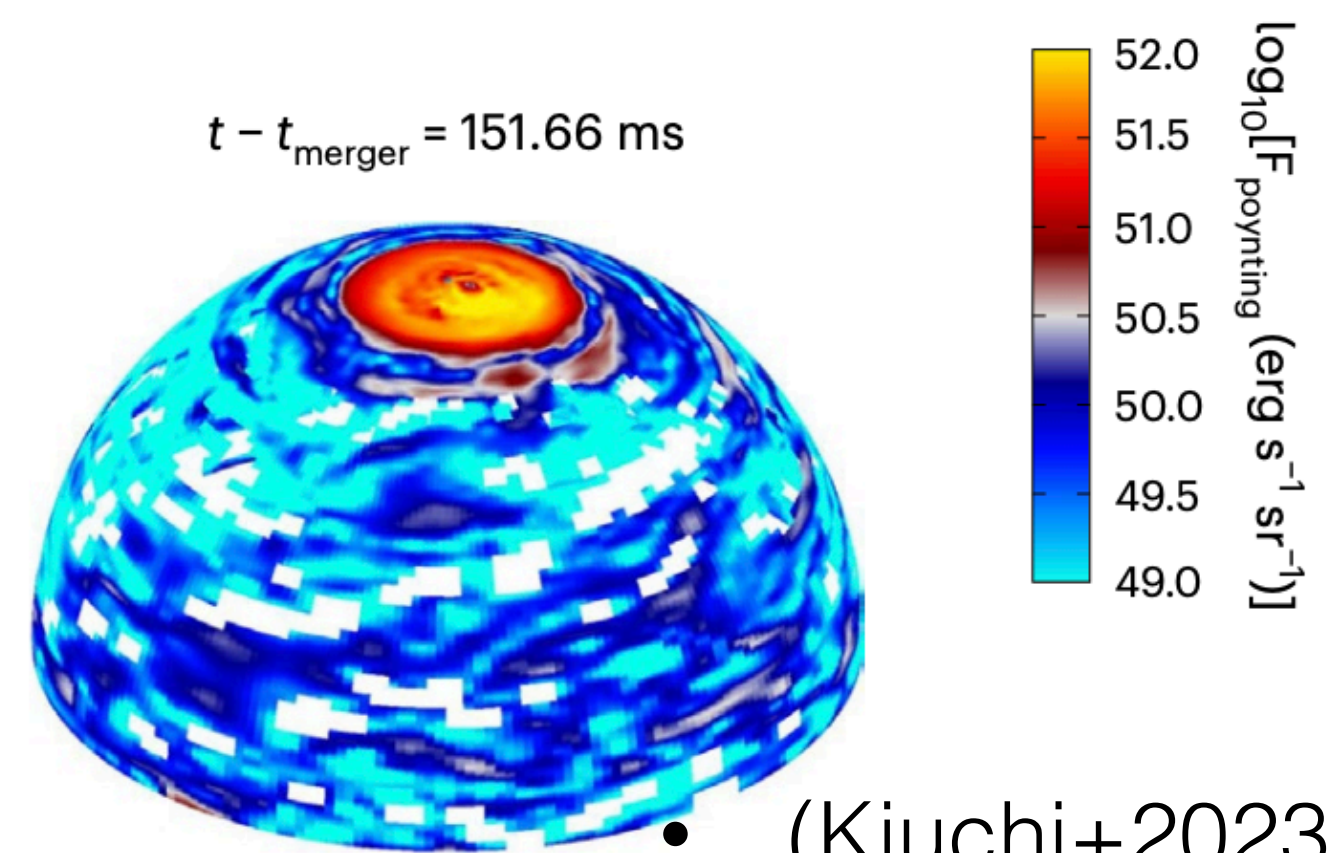
inspinal
GRMHD dynamics
→ huge B

**Gravitational Wave
Emitted**

"Hypervnova"
0-few sec
jet
V ~ 0.1 c
outflow
10°

initialization of B
150ms, 500 km

Kiuchi+23 → B_{ini}



$r \sim 10^{10} \text{ cm}$

cools to ~ 1 MeV
nucleosynthesis
~ 1 s: $r \sim 10^{10} \text{ cm}$

r-process
nucleosynthesis

$r \sim 10^{14} \text{ cm}$
 $B \approx 400 \text{ G}$

UHECR's
accelerated

B drops till

$\tau_{\text{accel}} / \tau_{\text{synch-loss}} \sim 1$

UHECR accel
~ 1 day: 10^{14} cm

Predict R_{cut}
for different $\{Z, A\}$

Predicting the UHECR spectral cutoff, R_{cut} , in the BNS magnetized turbulent outflow (1)

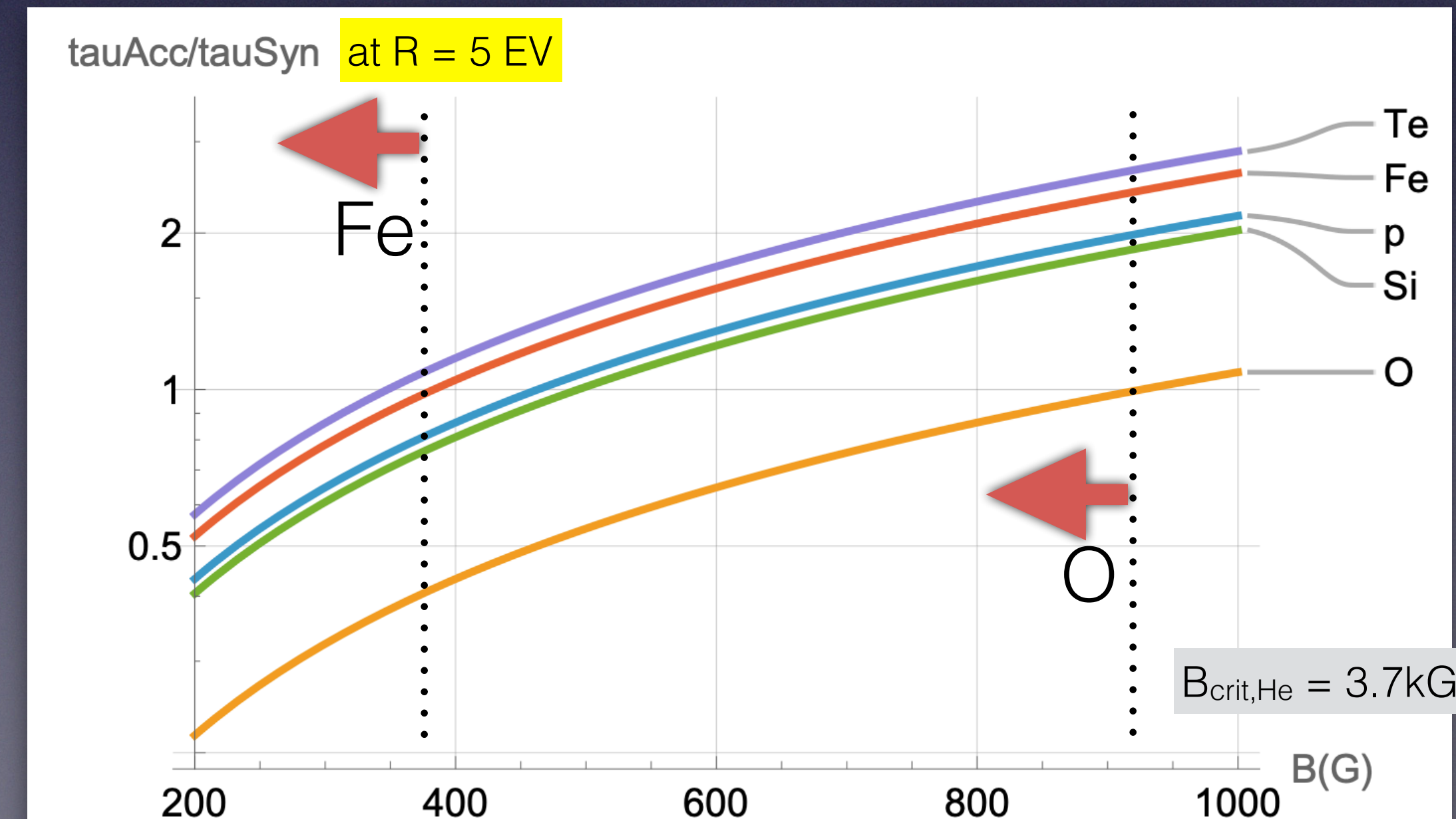
For a given B field, maximum energy of CR is set by $\tau_{\text{synch-loss}} \approx \tau_{\text{accel}}$

- $\tau_{\text{synch-loss}}$ depends on A, Z, B, & **R**igidity ($= \gamma A/Z m_p$):

$$\dot{E}_{\text{synch}} = \frac{4}{9} c \left(\frac{m_e r_0}{m_p} \right)^2 (\gamma \beta B)^2 \frac{Z^4}{A^2}$$

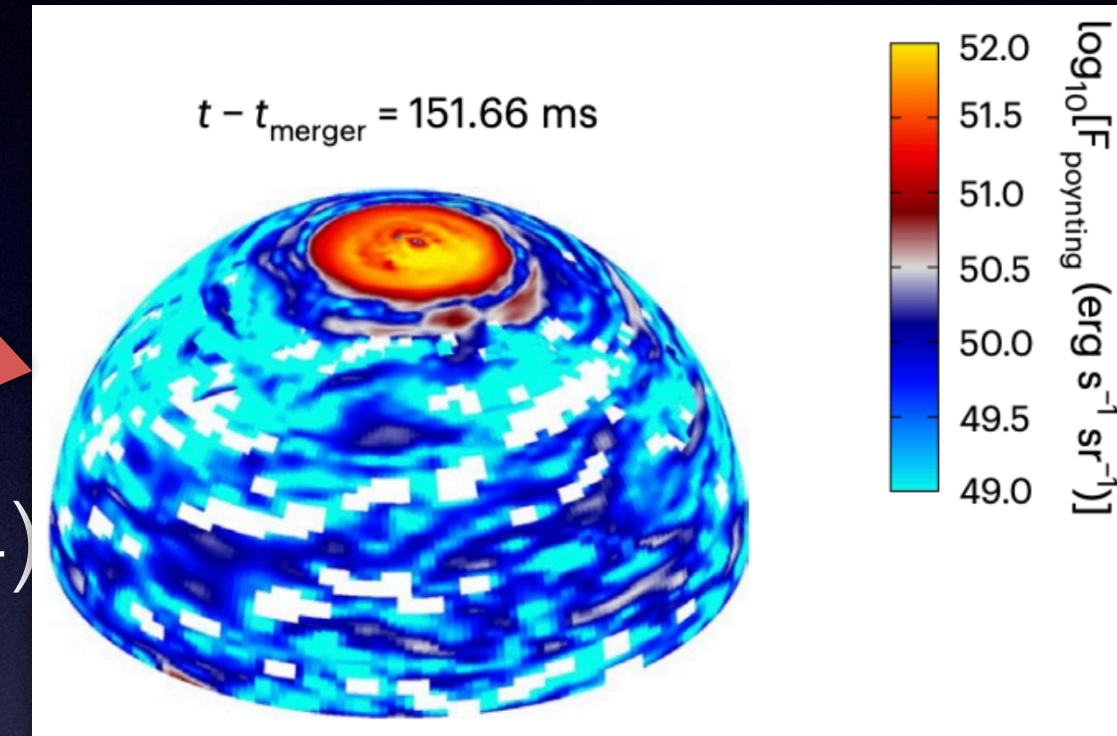
- $\tau_{\text{accel}} \approx r_{\text{Lar}}/c \sim (E/Z = \mathbf{R}) / B$

$\tau_{\text{accel}} / \tau_{\text{synch-loss}} \sim 1 \Rightarrow \text{max } B \text{ for given } R$

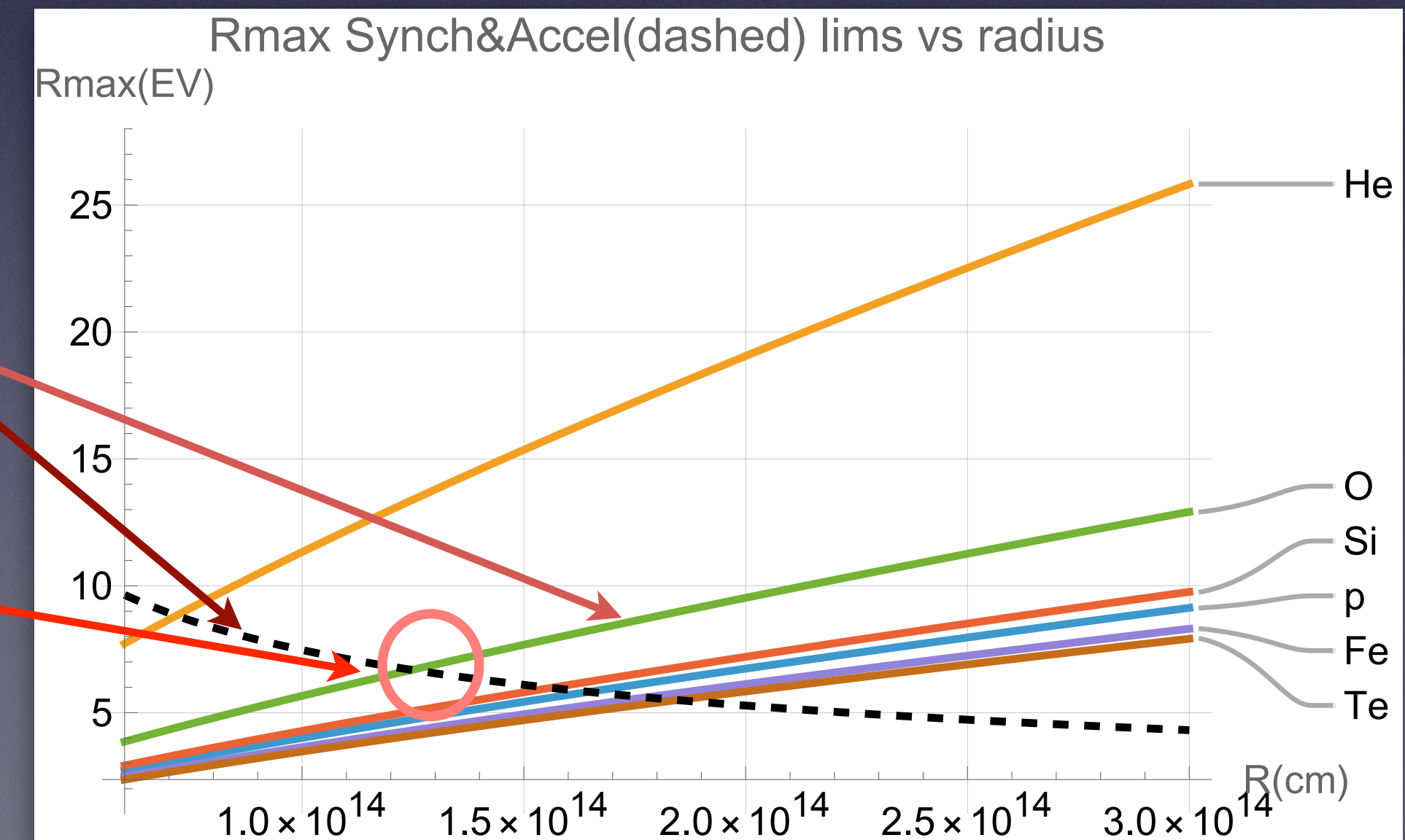


Predicting R_{cut} in the BNS magnetized turbulent outflow (2)

- Initialize $B(r=500 \text{ km}) = 3.3 \cdot 10^{12} \text{ G}$; $L_{\text{coh}} \approx r/3$ (Kiuchi+2023)
- Assume that $B(r) = B_0 (r/r_0)^{-3/2}$; homologous expn (Rosswog+2014)
- Maximum rigidity of CR when radius is r :
 - $R_{\text{max}}(r) = 0.65 B(r) L_{\text{coh}}(r)$ (turb. accel. CFM24)
 - but require $\tau_{\text{synch-loss}}(A, Z, R, r) \gtrsim \tau_{\text{accel}}(R, B)$
 - $\rightarrow R_{\text{cut}}(Z, A, r)$ is intersection



$R_{\text{cut}}\{p, \text{He}, \text{O}, \text{Si}, \text{Fe}, \text{Te}\} \approx \{6.2, 9.4, 7.1, 6.3, 6.0, 5.9\} \text{ EV}$
 fit to Auger data assuming common $\rightarrow R_{\text{cut}} = 6.8 \text{ EV}$



Essential UHECR facts explained!

- Spectrum from magnetized turbulence: $\sim E^{-2.1} \text{sech}[(E/E_{\text{cut}})^2]$
- At each E , the range of masses is narrow:
 - Minimal source-to-source variation
 - Rigidity spectrum of individual sources is narrow (sech)
 - $R_{\text{cut}} \approx 6\text{-}9 \text{ EV}$ (observed: $R_{\text{cut}} = 6.3^{+6.3}_{-2.3} \text{ EV}$)
- First explanation for highest energy events: EXCELLENT interpretation as originating in $A \sim 130$ Tellurium/Xenon peak.

$$E = R Z_{\text{Te-Xe}} \approx 4.5 \text{ EV} \times (52\text{-}54) = 240 \text{ EeV}$$

$$E_{\text{OMG}} \approx 250 \pm 70 \text{ EeV}, \quad E_{\text{Amaterasu}} \approx 212 \pm 25 \text{ EeV}$$

Tests of BNS merger scenario

- ✓ Eroded r-process nuclei among UCRs ($A > 56$)
- ✓ Successful prediction of R_{cut}
- Characteristic pattern of R_{cut} slightly varying with A reflecting synchrotron-limited acceleration
- flux of EHE neutrinos
- GW- ν coincidences for EHE ν 's ($E > \text{PeV}$) - time delay ≈ 1 day

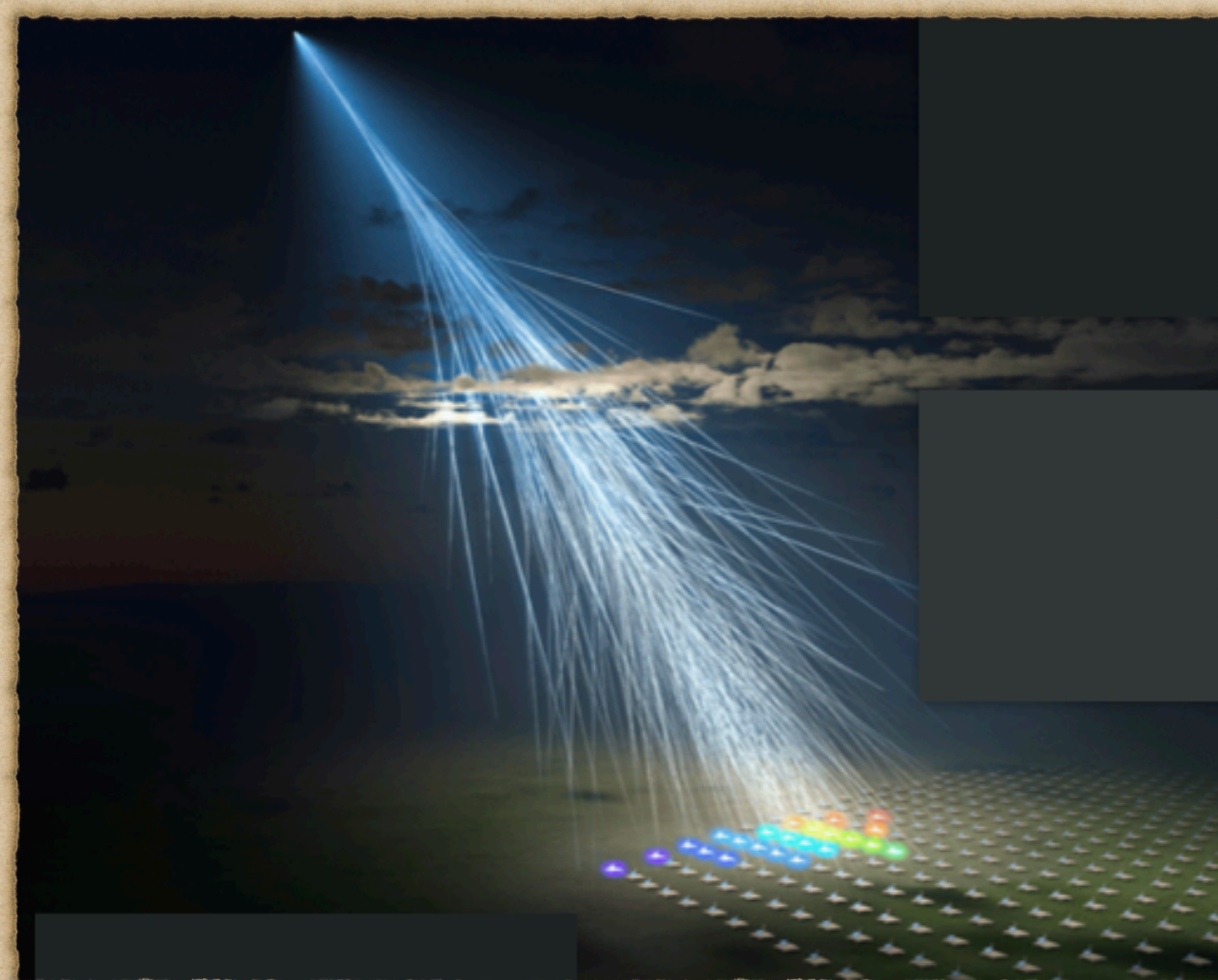
EHE Neutrino production

- Neutrinos come from UHECRs:
 - spallation neutrons beta decay: $E_\nu \approx 10^{-3} (E_n \approx \mathbf{R}/2) \sim 2 \text{ PeV}$
 - photo-pion production: $E_\nu \approx (E/A)/20 \sim \mathbf{R}/40 \sim 100 \text{ PeV}$
- Time delay for ν is $>$ time for ejecta to expand to 10^{14} cm where UHECRs form:
 - $\langle v_{ej} \rangle = 0.1 c \rightarrow$ delay of $O(\text{day})$

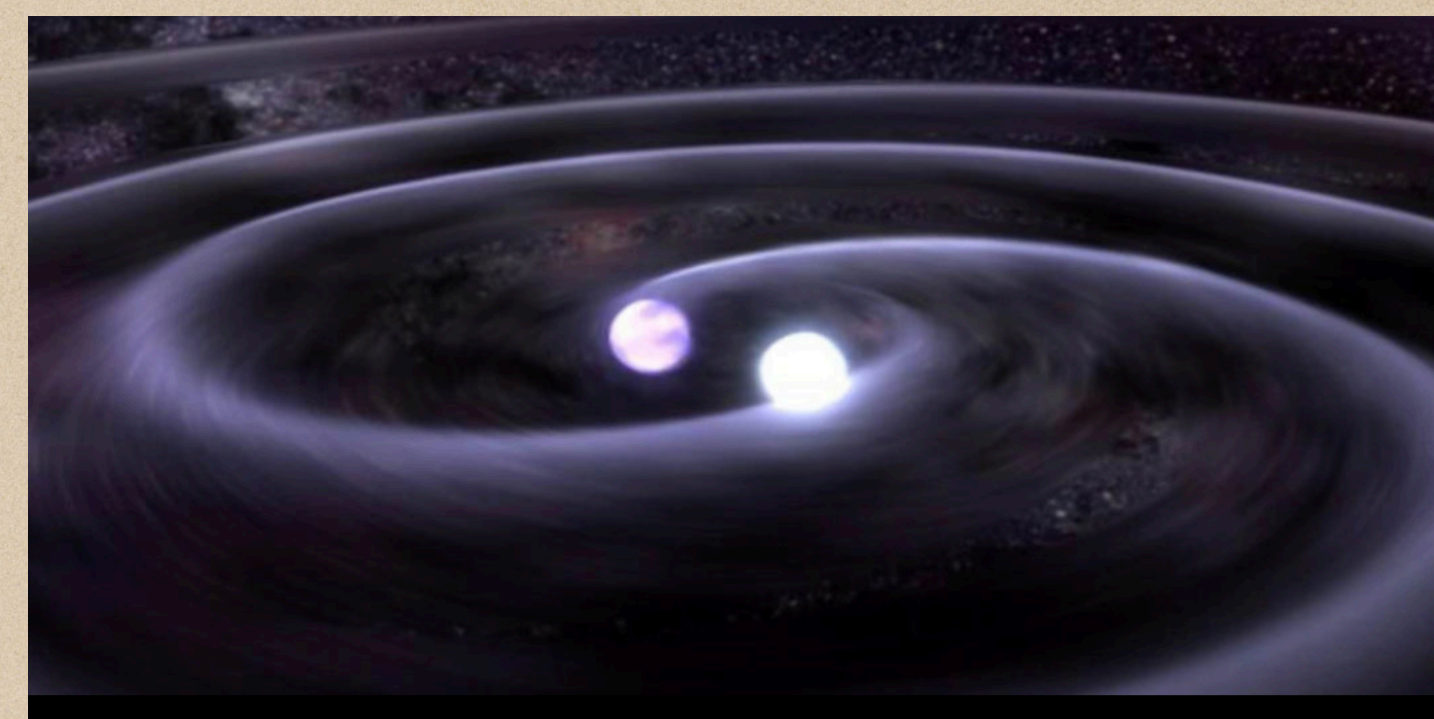
To do: make better estimate of rate of observing GW-nu coincidences
very crude: 1/yr with Cosmic Explorer, Einstein telescope & ICGen2

Path to predict UCR composition

- Initialize from r-process simulations
 - Depends on NS EoS & ... but will improve
 - Beyond $\sim 30^\circ$, most nuclei are heavier than $\sim \text{Fe}$
- Follow expansion as outflow radius increases
 - CRs accelerated to $\gtrsim 100$ MeV collide, producing spallation
 - lighter nuclei are formed
- Need PIC simulations
 - Expanding magnetized turbulence
 - “Uptake probability” into acceleration chain, to calculate absolute magnitude of CR component.



Summary



New suggestion: UHECRs are produced in binary NS mergers.

- ◆ Uniquely, potentially satisfies all requirements:
 - * Universal Maximum Rigidity explained — predicted value agrees with data.
 - * TBD: Total UHECR power (depends on BNS merger rate & power in CRs)
- ◆ Highest energy events are r-process nuclei — testable by Auger
- ◆ EHE neutrinos, and coincidences with GWs or sGRBs from BNS-merger

Needed from YOU!!!

- BNS merger (simulation) community:
 - photon field at large radii, to calculate spallation
 - nucleosynthesis abundances
- Nuclear physics community:
 - breakup cross sections for reaction networks to predict UHECR composition
- Particle Acceleration/Plasma Physics community:
 - PIC simulation for spherically expanding system
 - PIC simulations to understand uptake

Backup slides

Source candidates vs key constraints

	$n_s \gtrsim 10^{-3.5}$ Mpc ⁻³	energy injection	ordinary galaxy	Universal R_{\max}	Highest energy events
Powerful AGN	[X]	✓	X	X	X
Long GRBs	[X]	X	X	X	X
Tidal Disruption Events	?	?	✓	X	X
Accretion Shocks	?	?	[X]	X	X
BNS mergers	✓	[✓]	✓	✓	✓

(All can satisfy Hillas size > Larmor radius)

Very highest energy UHECRs explained!

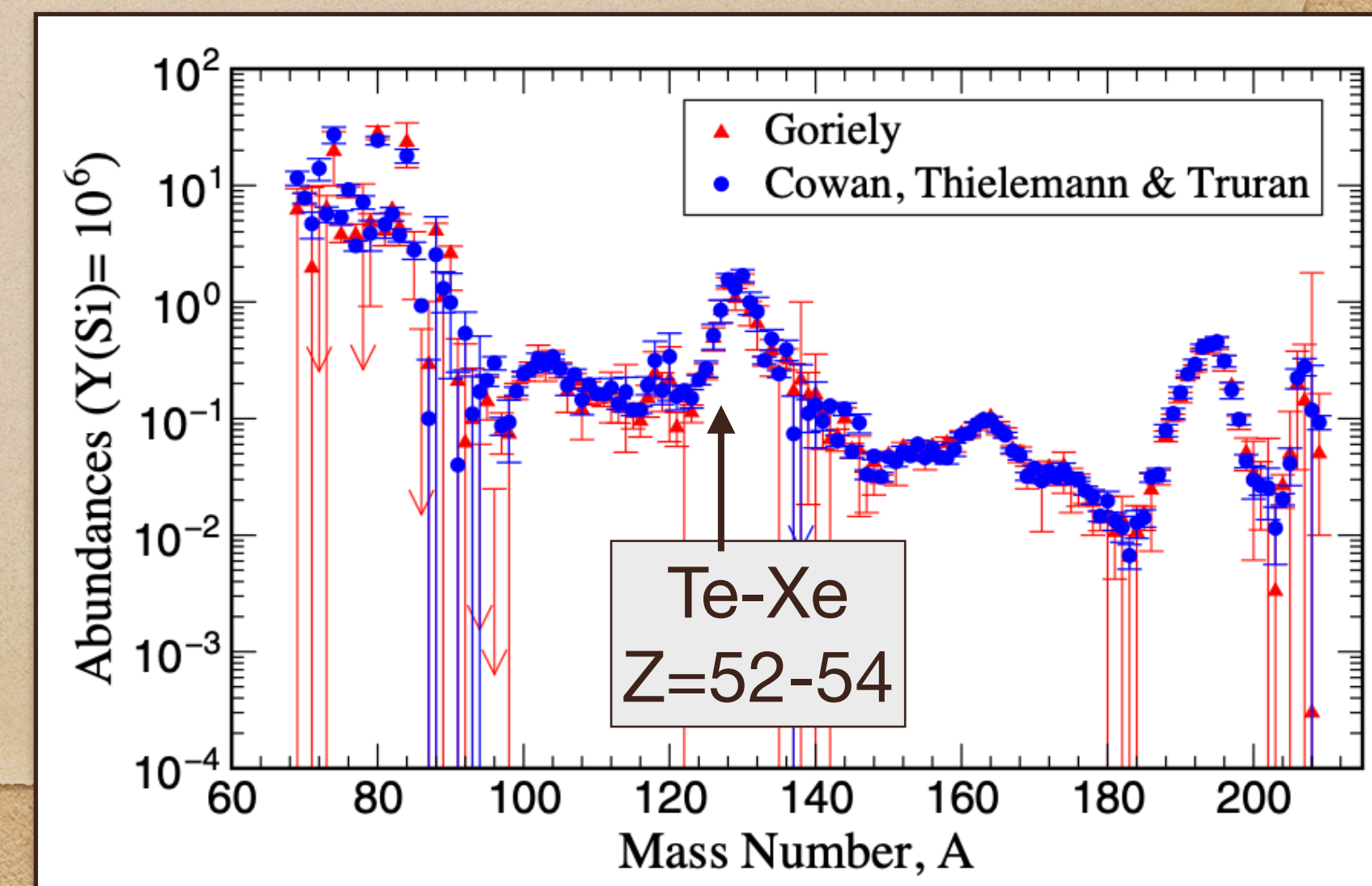
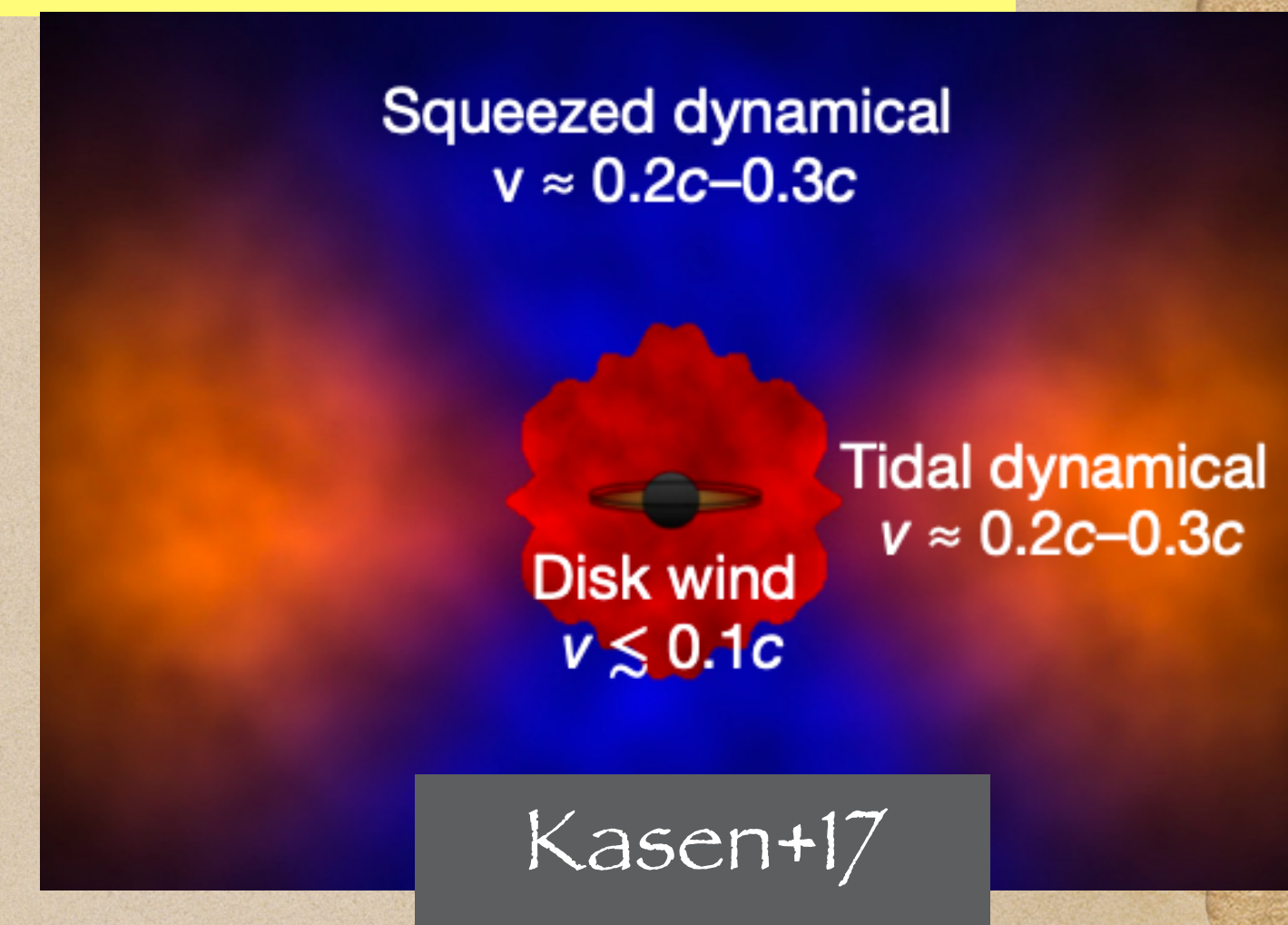
(with narrow Rigidity ≈ 4.5 EV, where do UCR's above 100 EeV come from?)

- ◆ r-process nucleosynthesis takes place in BNS mergers
- ◆ sometimes an r-nucleus is swept up and accelerated

$$\rightarrow E \approx R Z_{\text{Te-Xe}} \approx 4.5 \text{ EV} \times (52-54) \approx 240 \text{ EeV}$$

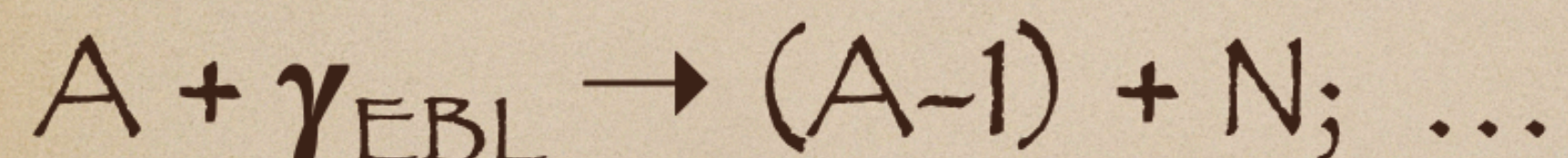
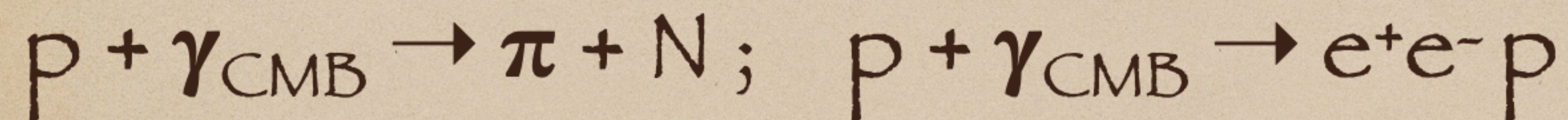
- Excellent agreement with OMG and Amaterasu!
- $E_{\text{OMG}} \approx 250 \pm 70 \text{ EeV}^*$, $E_{\text{Amaterasu}} \approx 212 \pm 25 \text{ EeV}^{**}$

*with modern air fluorescence yield **higher if a proton



ENERGY LOSS IN PROPAGATION \Rightarrow GZK Horizon

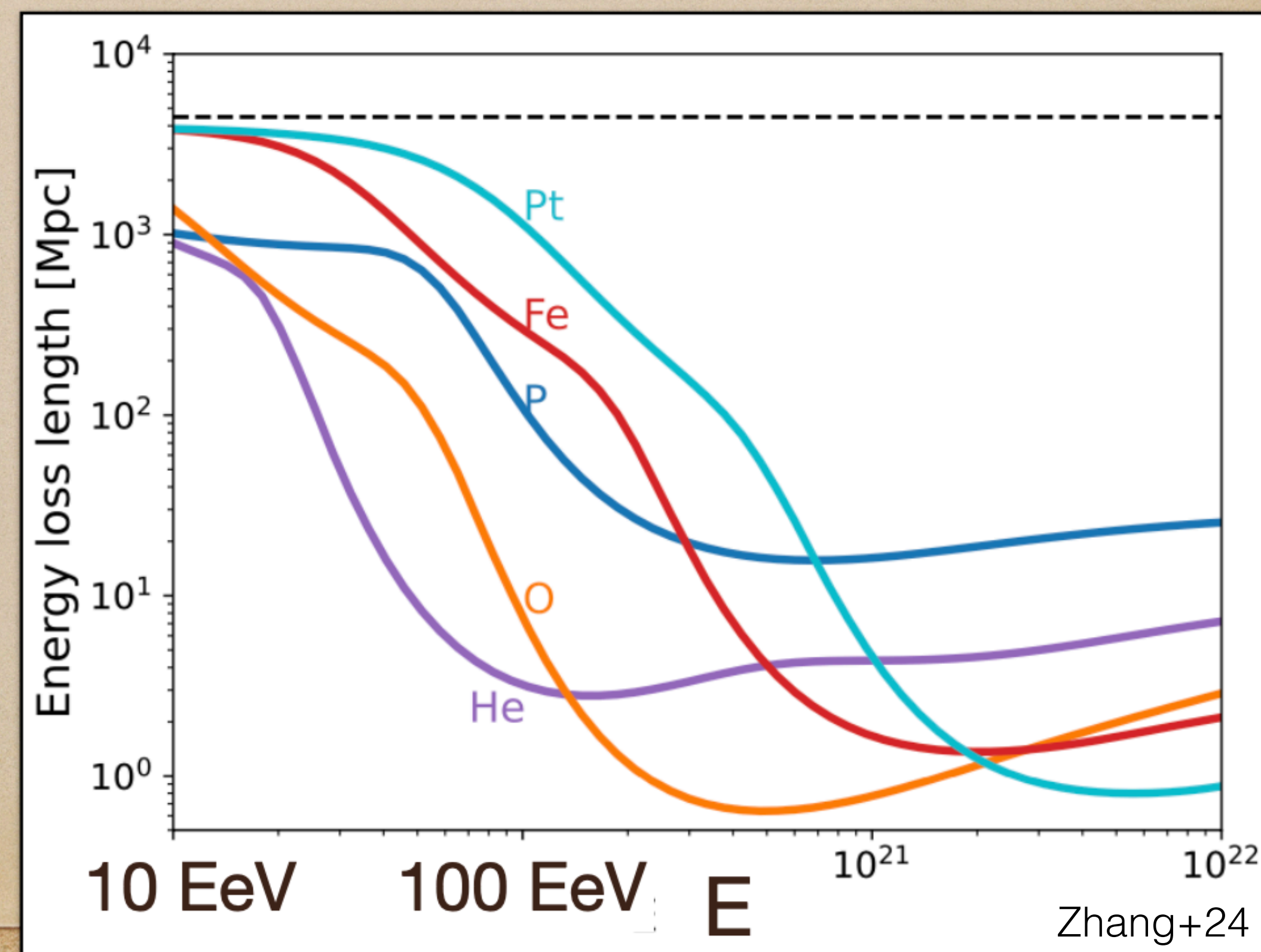
Greisen, Zatsepin, Kuzmin (GZK) bound



$$\langle \text{dist} \rangle_{E > 8 \text{ EeV}} = 210 \text{ Mpc}$$

$$\langle \text{dist} \rangle_{E > 32 \text{ EeV}} = 70 \text{ Mpc}$$

(for ref: M87 is 16 Mpc away)



Spectrum for Magnetized Turbulence

$$\sigma > 1$$

Ultra-High-Energy Cosmic Rays Accelerated by **Magnetically Dominated Turbulence**

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³Columbia Astrophysics Laboratory, Columbia University, New York, NY 10027, USA

⁴Center for Cosmology and Particle Physics, Department of Physics, New York University, New York, NY 10003, USA

⁵Department of Physics, Pennsylvania State University, University Park, PA 16802, USA

⁶Department of Astronomy and Astrophysics, Pennsylvania State University, University Park, PA 16802, USA

⁷Institute of Gravitation and the Cosmos, Center for Multi-Messenger Astrophysics, Pennsylvania State University, University Park, 16802, USA

ABSTRACT

Ultra-High-Energy Cosmic Rays (UHECRs), particles characterized by energies exceeding 10^{18} eV, are generally believed to be accelerated electromagnetically in high-energy astrophysical sources. One promising mechanism of UHECR acceleration is magnetized turbulence. We demonstrate from first principles, using fully kinetic particle-in-cell simulations, that magnetically dominated turbulence accelerates particles on a short timescale, producing a power-law energy distribution with a rigidity-dependent, sharply defined cutoff well approximated by the form $f_{\text{cut}}(E, E_{\text{cut}}) = \text{sech}[(E/E_{\text{cut}})^2]$. Particle escape from the turbulent accelerating region is energy-dependent, with $t_{\text{esc}} \propto E^{-\delta}$ and $\delta \sim 1/3$. The resulting particle flux from the accelerator follows $dN/dEdt \propto E^{-s} \text{sech}[(E/E_{\text{cut}})^2]$, with $s \sim 2.1$. We fit the Pierre Auger Observatory's spectrum and composition measurements, taking into account particle interactions between acceleration and detection, and show that the turbulence-associated energy cutoff is well supported by the data, with the best-fitting spectral index being $s = 2.1^{+0.06}_{-0.13}$. Our first-principles results indicate that particle acceleration by magnetically dominated turbulence may constitute the physical mechanism responsible for UHECR acceleration.

Fully kinetic treatment of the plasma

- ▶ The evolution of the particle density $f_s(\mathbf{x}, \mathbf{p}, t)$ of species s in a collisionless plasma is described by the Vlasov equation

$$\frac{\partial f_s}{\partial t} + \frac{\mathbf{p}}{m_s \gamma_s} \cdot \nabla_{\mathbf{x}} f_s + \mathbf{F} \cdot \nabla_{\mathbf{p}} f_s = 0$$

$$\text{where } \gamma_s^2 = 1 + \frac{|\mathbf{p}|^2}{m_s^2 c^2} \text{ and } \mathbf{F} = q_s \left(\mathbf{E} + \frac{\mathbf{p}}{\gamma_s m_s c} \times \mathbf{B} \right).$$

- ▶ $\mathbf{E}(\mathbf{x}, t)$ and $\mathbf{B}(\mathbf{x}, t)$ are determined from Maxwell's equations

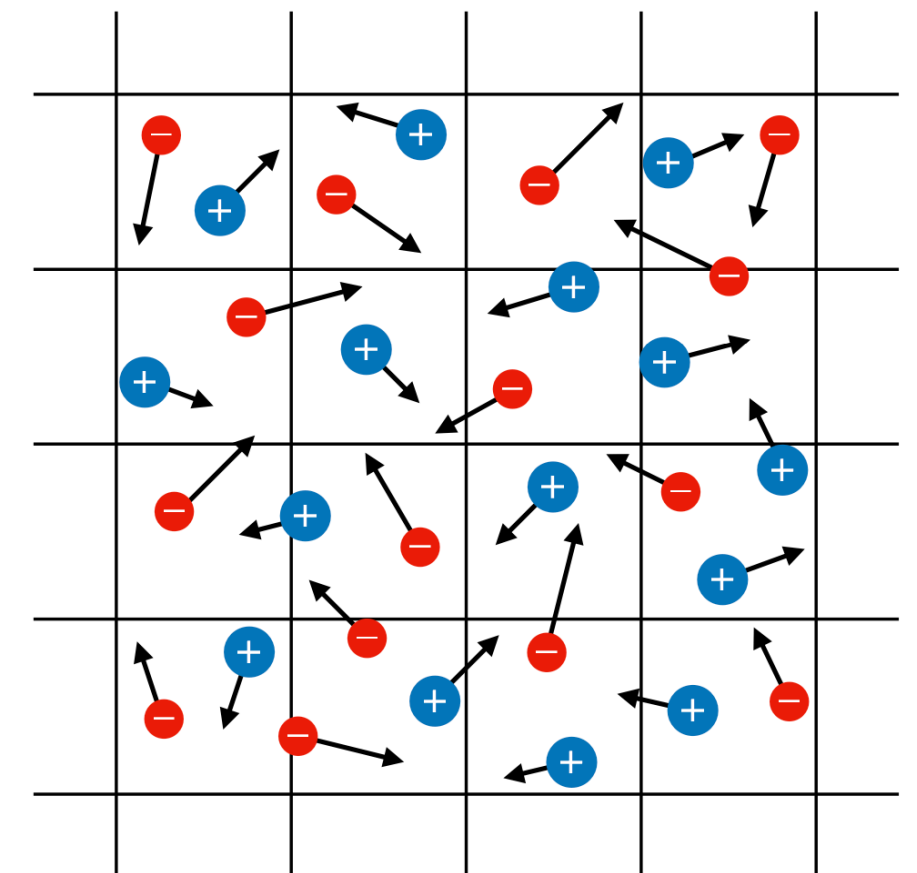
$$\frac{\partial \mathbf{E}}{\partial t} - c \text{curl} \mathbf{B} = -4\pi \mathbf{J}, \quad \text{div} \mathbf{E} = 4\pi \rho,$$

$$\frac{\partial \mathbf{B}}{\partial t} + c \text{curl} \mathbf{E} = 0, \quad \text{div} \mathbf{B} = 0,$$

where the source terms are computed by

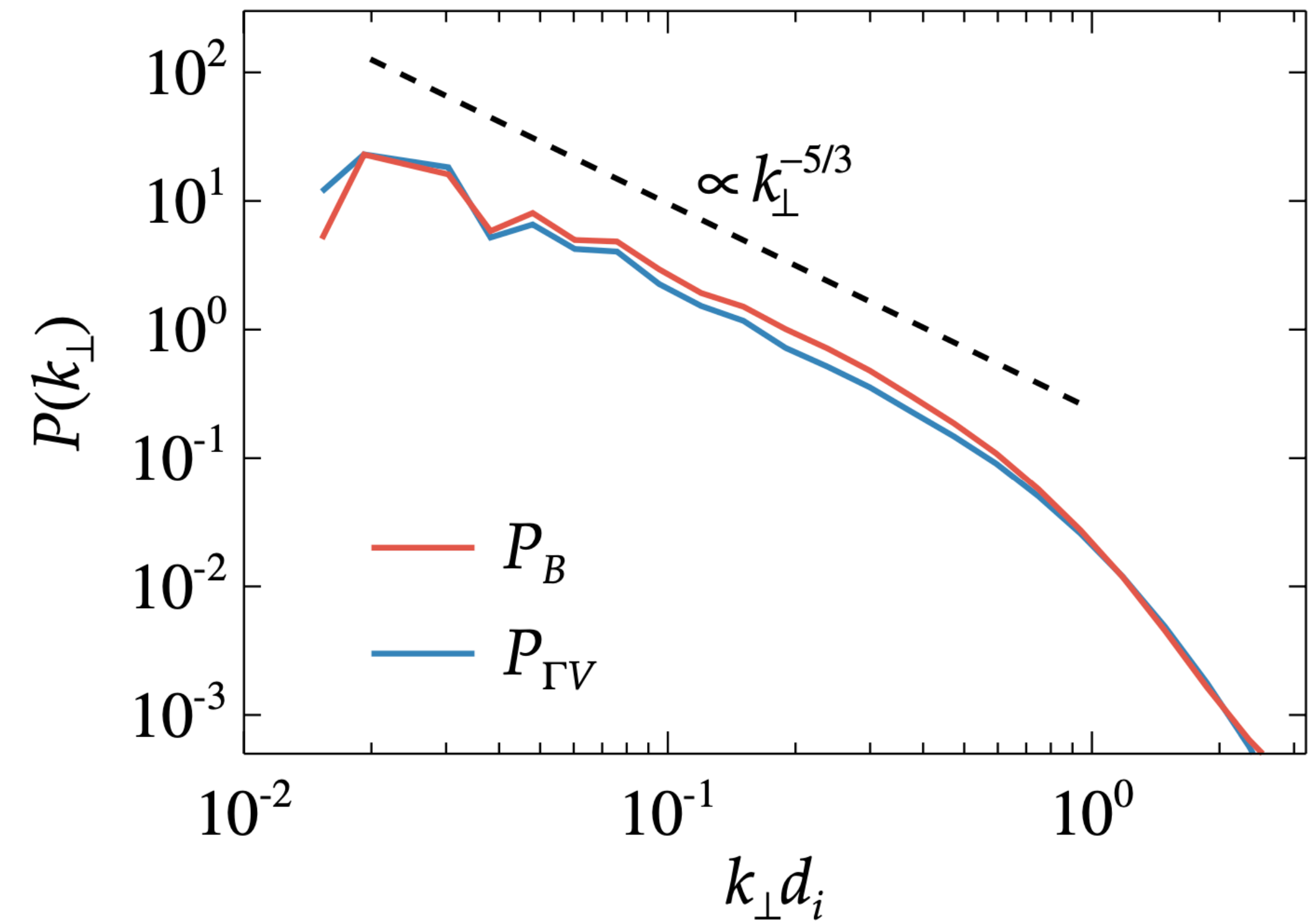
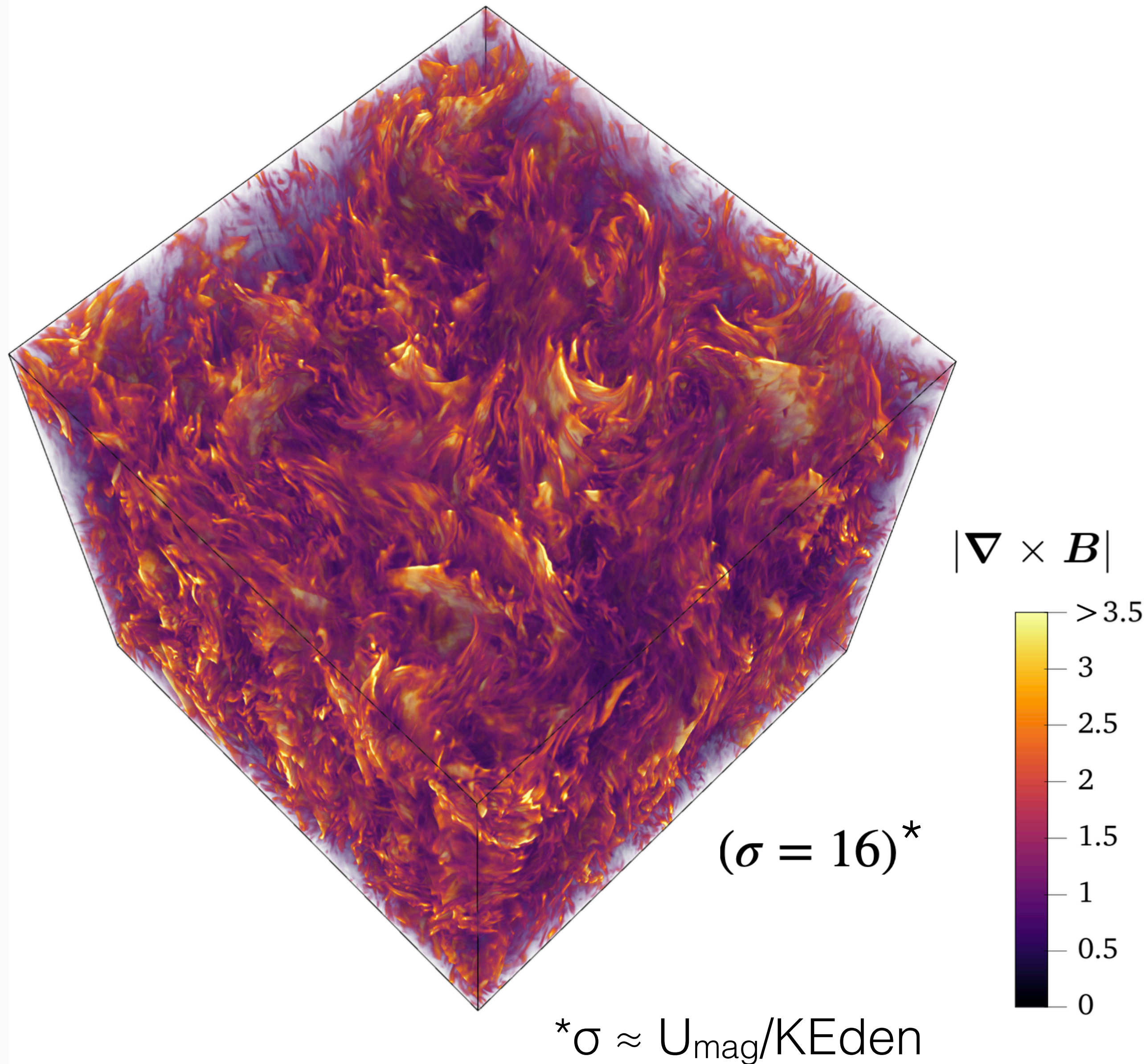
$$\rho = \sum_s q_s \int_{\mathbb{R}^3} f_s d\mathbf{p}, \quad \mathbf{J} = \sum_s \frac{q_s}{m_s} \int_{\mathbb{R}^3} f_s \frac{\mathbf{p}}{\gamma_s} d\mathbf{p}.$$

- ▶ Solution via particle-in-cell method



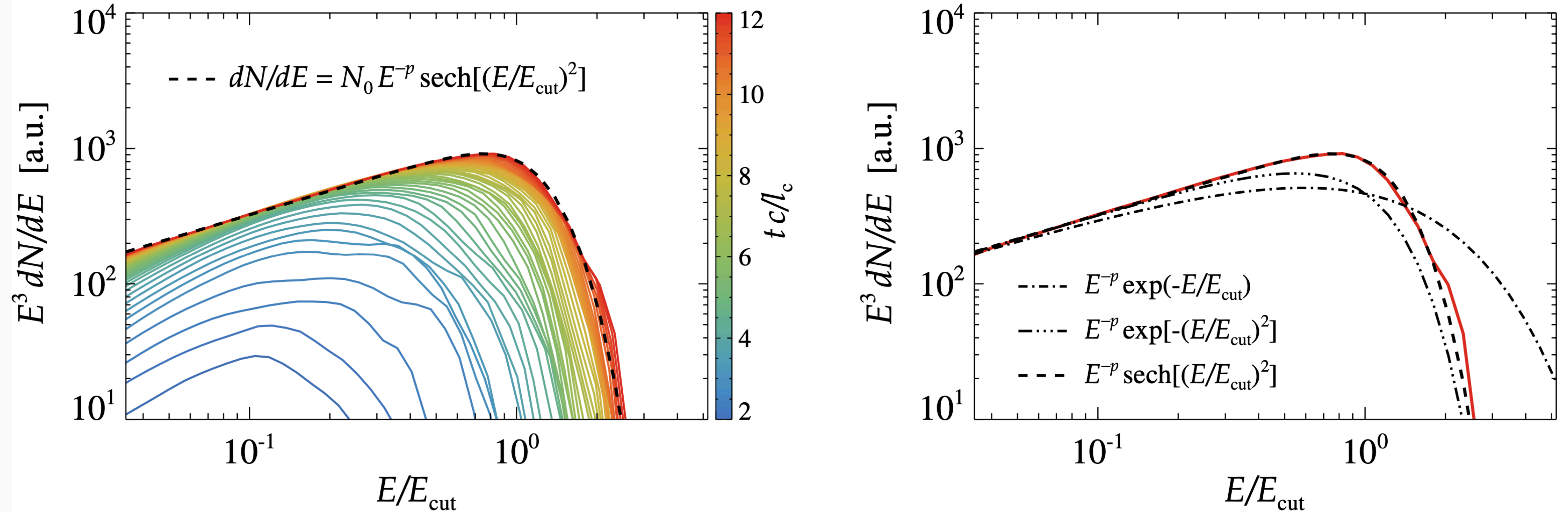
PIC code: TRISTAN-MP
(Spitkovsky 2005)

Magnetically dominated turbulence from first-principles PIC simulations



- The large computational domain allow us to capture both the MHD cascade at large scales and the kinetic cascade at small scales

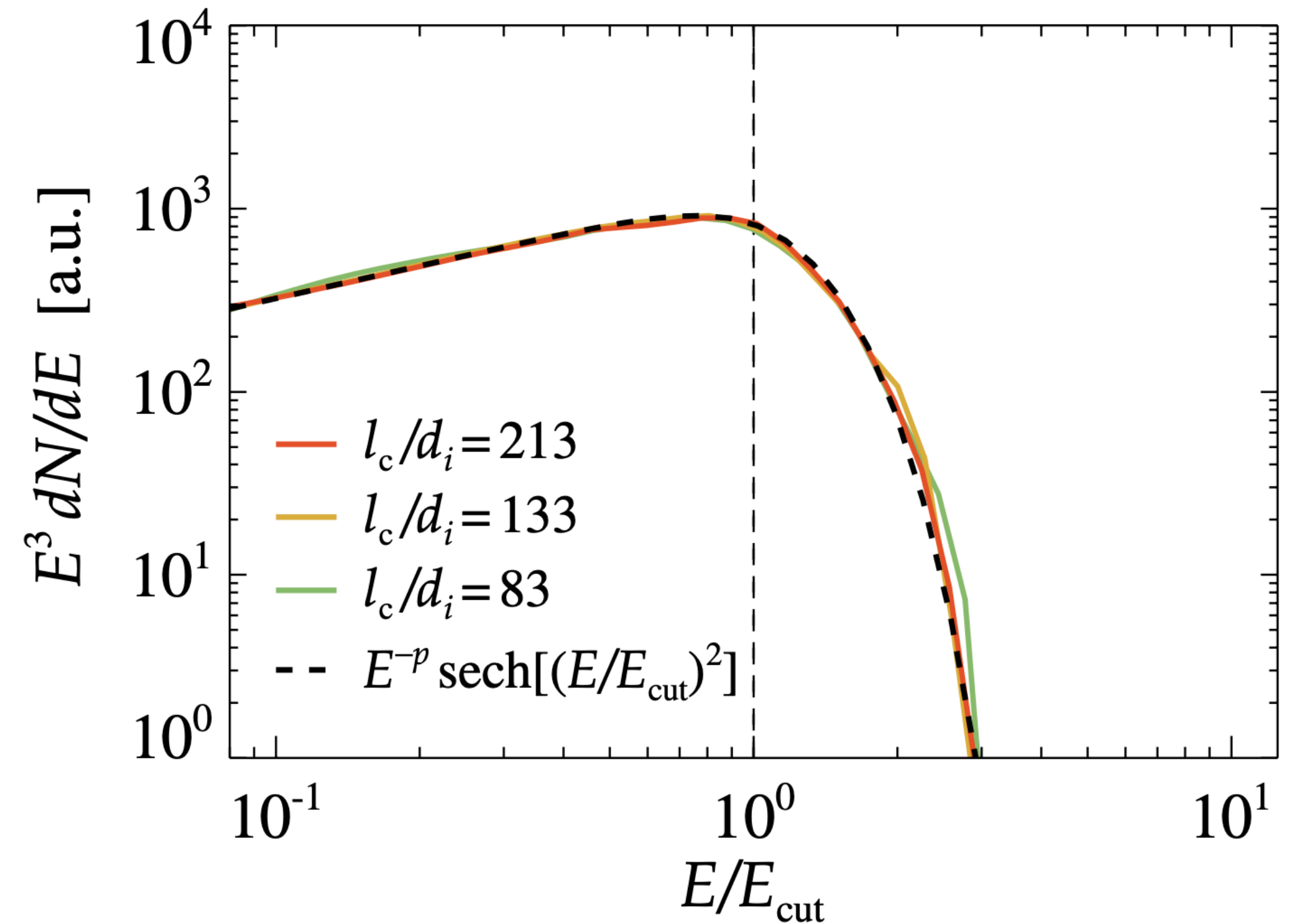
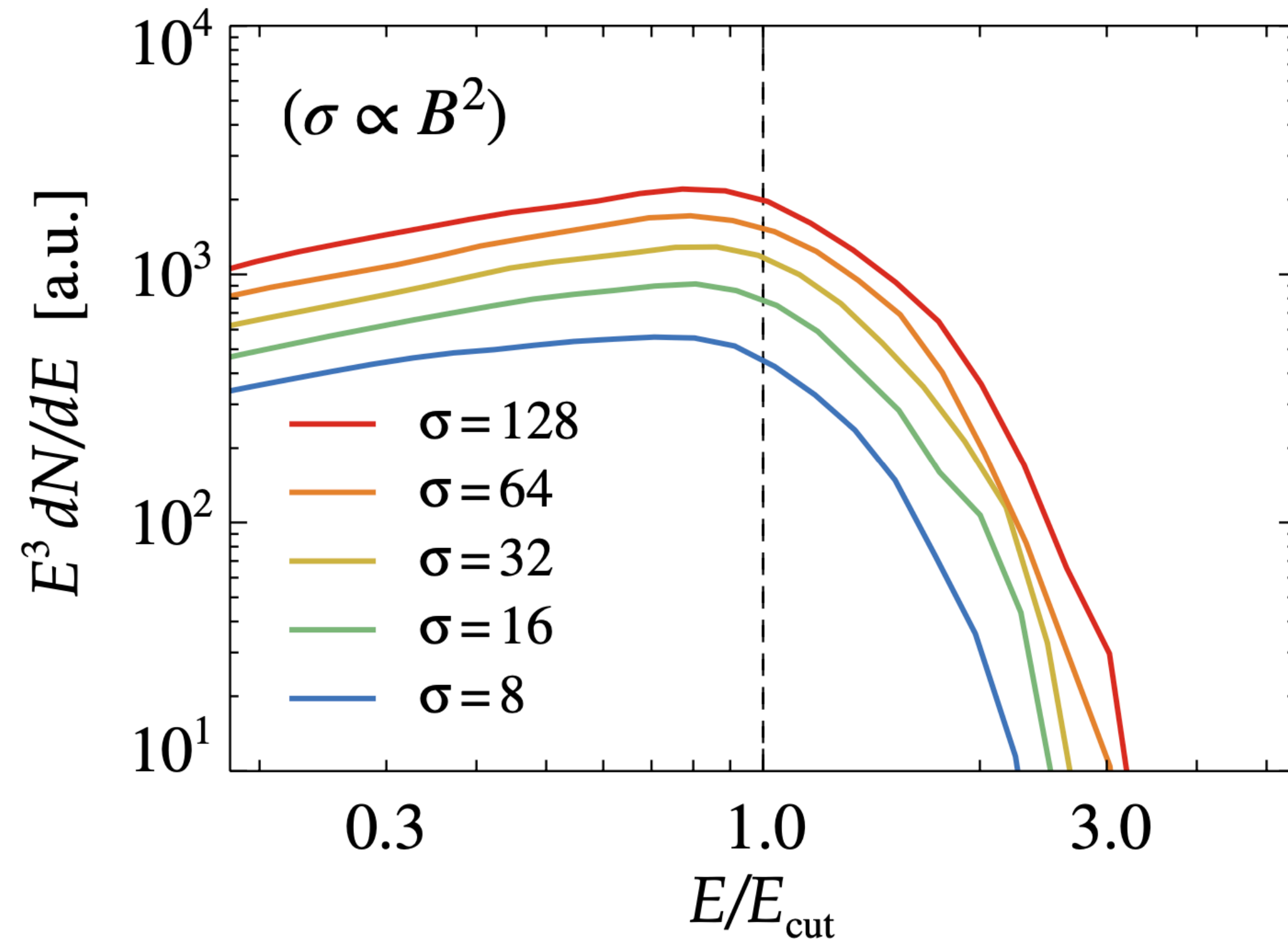
Particle acceleration via magnetized turbulence: nonthermal particle spectrum



► magnetized turbulence accelerates particles into a spectrum of the form:

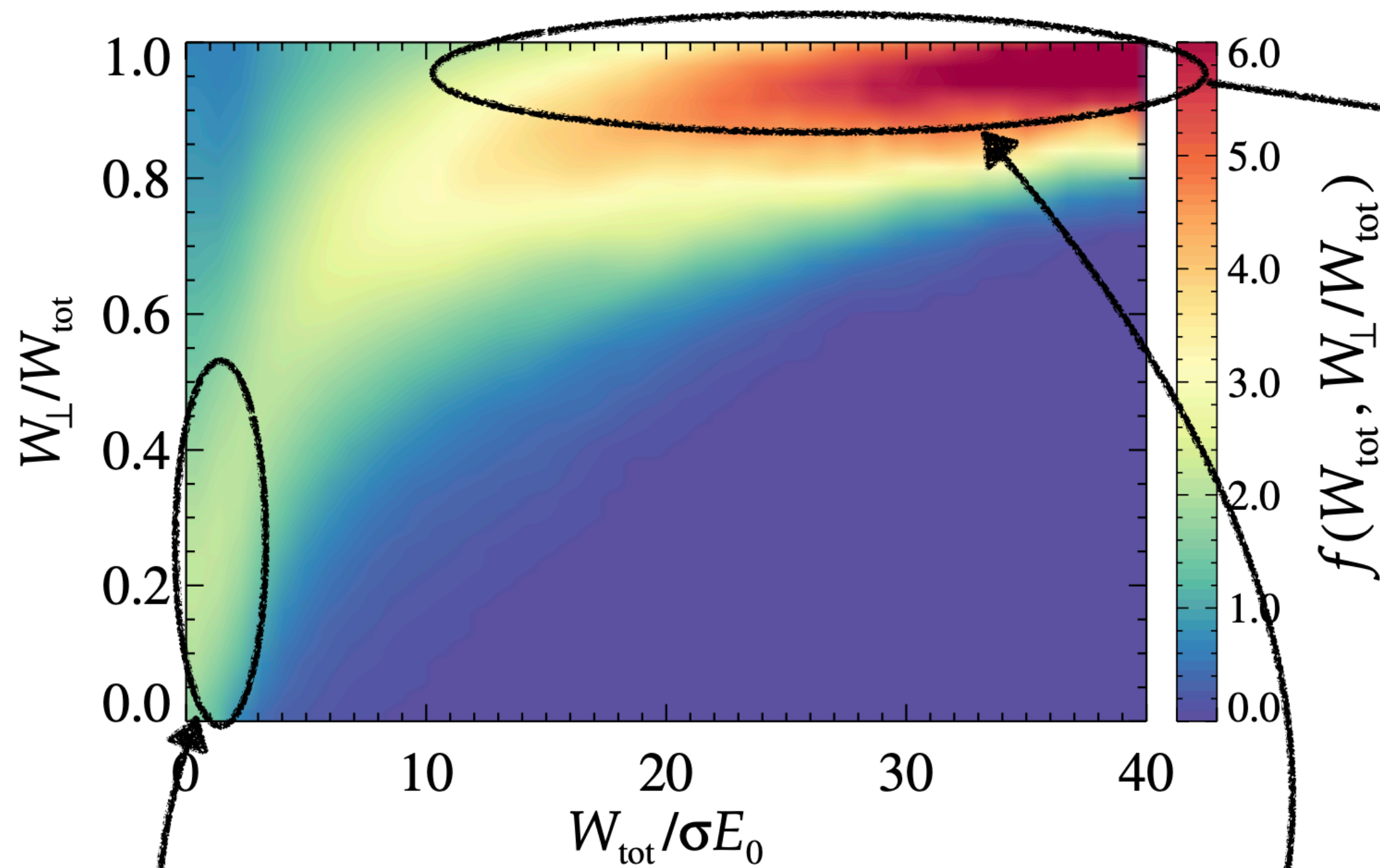
$$\frac{dN}{dE} = N_0 E^{-p} f_{\text{cut}}(E, E_{\text{cut}}) \quad \text{with} \quad f_{\text{cut}}(E, E_{\text{cut}}) = \text{sech}\left[\left(E/E_{\text{cut}}\right)^2\right]$$

Particle acceleration via magnetized turbulence: cutoff energy



- ▶ cutoff $\text{sech}[(E/E_{\text{cut}})^2]$ scales with $E_{\text{cut}} = ZeR_{\text{cut}} = Ze(B_{\text{rms}}\kappa l_c)$, where $\kappa = 0.65$ from the fits
 - ▶ magnetized turbulence does accelerate particles to the “Hillas limit” if one assumes $R_{\text{size}} = l_c$
- $R_{\text{cut}} = 0.65 B L_{\text{coh}}$

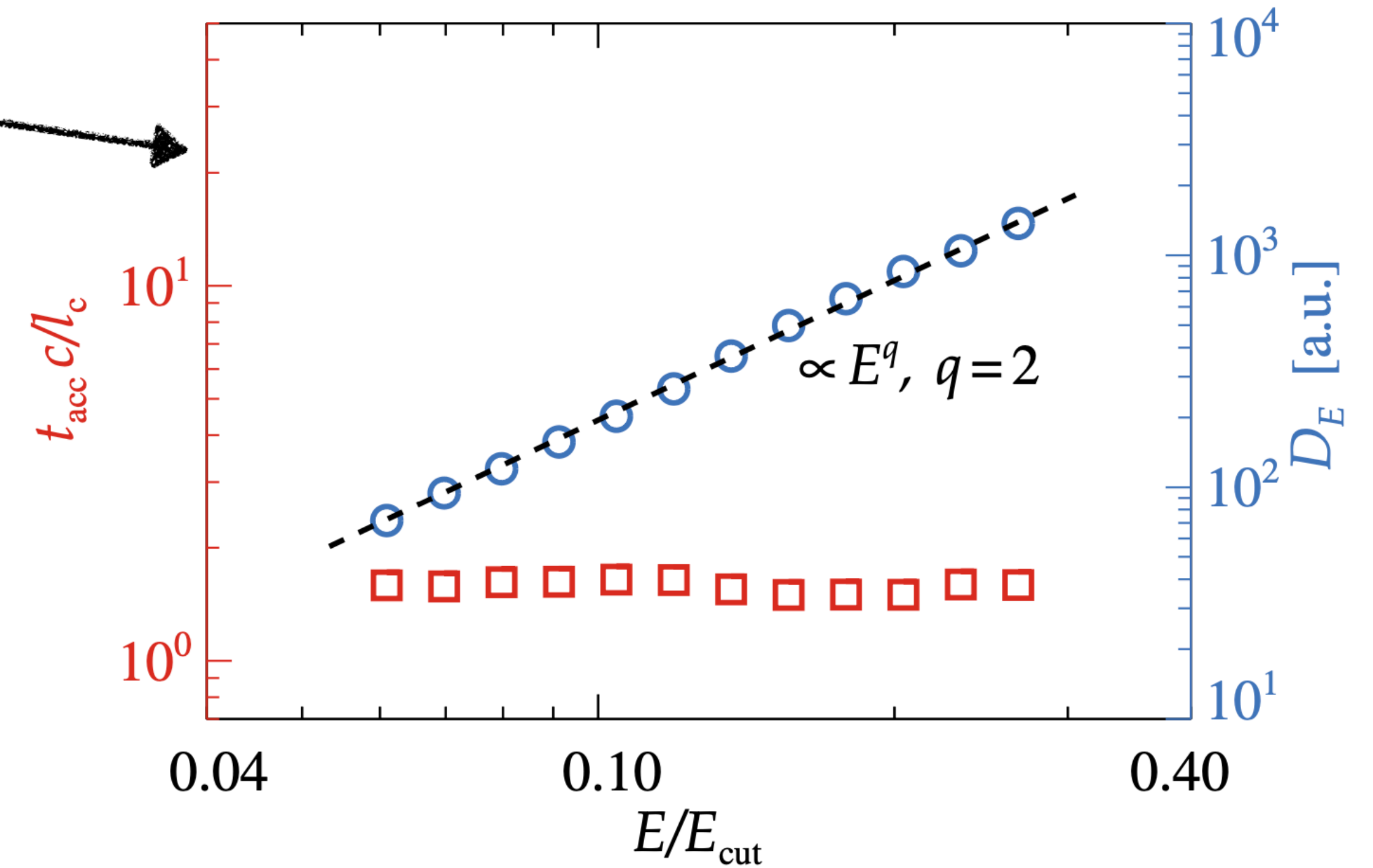
Particle acceleration via magnetized turbulence: particle acceleration elements



$$W_{\parallel,\perp}(t) = q \int_0^t \mathbf{E}_{\parallel,\perp}(t') \cdot \mathbf{v}(t') dt'$$

$$\mathbf{E}_{\parallel} = (\mathbf{E} \cdot \mathbf{B})\mathbf{B}/B^2$$

$$\mathbf{E}_{\perp} = \mathbf{E} - \mathbf{E}_{\parallel} \simeq -(\mathbf{V}/c) \times \mathbf{B}$$



$$t_{\text{acc}} = \frac{E^2}{4D_E} \simeq \frac{1}{4\kappa_{\text{acc}}\delta u^2} \frac{B_{\text{rms}}^2}{\delta B_{\text{rms}}^2} \frac{l_c}{c}$$

$\kappa_{\text{acc}} \simeq 0.1$ from PIC simulations
(see Comisso & Sironi 2019)

Particle acceleration via magnetized turbulence: spectrum out o

- residence time within the accelerator:

$$t_{\text{esc}} \simeq \frac{L^2}{\lambda_s c} \simeq \frac{L^2}{l_c c} \left(\frac{E_{\text{cut}}}{E} \right)^\delta \propto E^{-\delta}$$

- flux of particles escaping the accelerator is given by

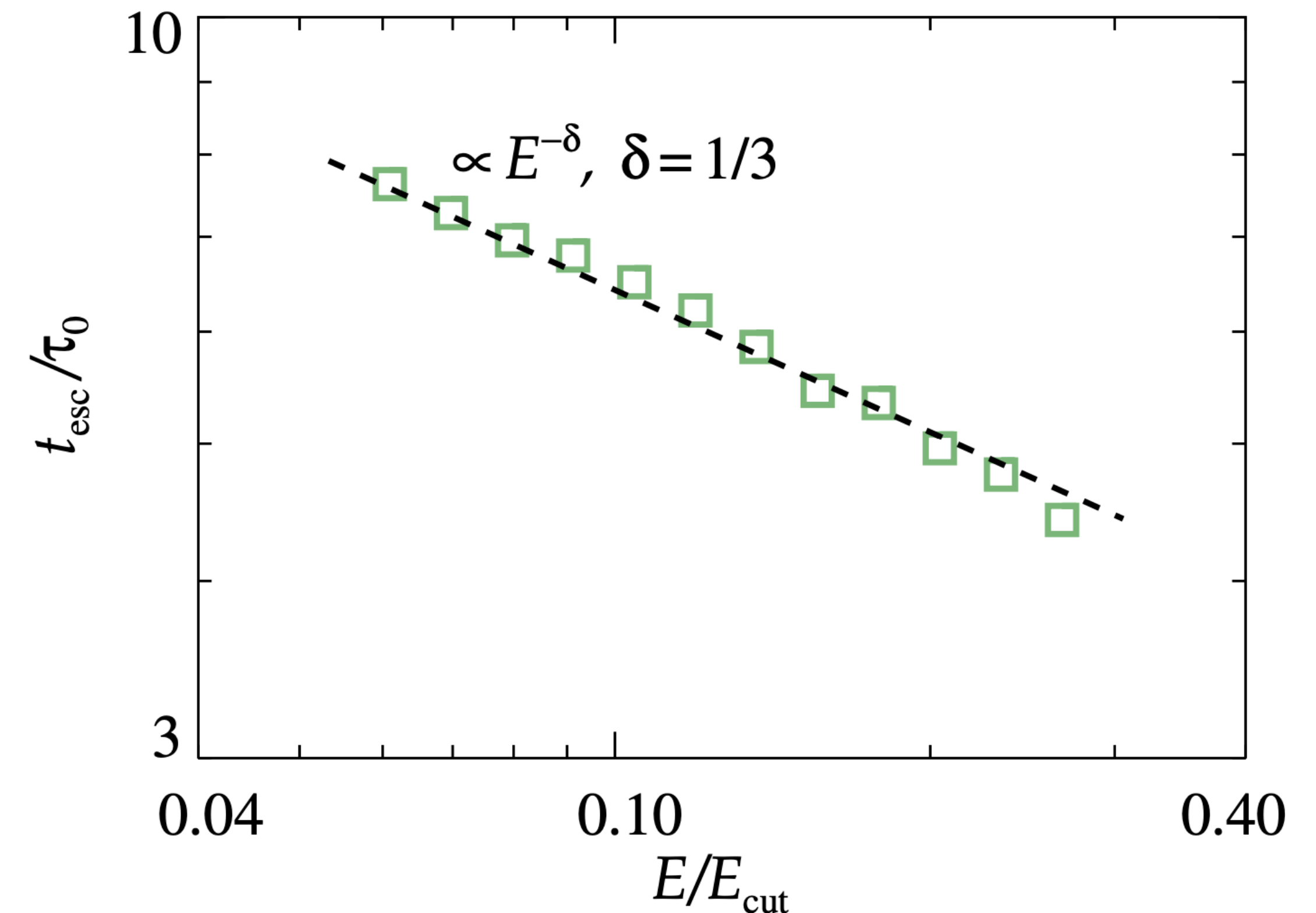
$$\phi(E) = \frac{dN}{dE dt} = \frac{1}{t_{\text{esc}}} \frac{dN}{dE} \propto E^{-s} \text{sech} \left[\left(E/E_{\text{cut}} \right)^2 \right]$$

with $s = \underbrace{p}_{\sim 2.4} + \underbrace{\delta}_{\sim 1/3} \sim 2.1$

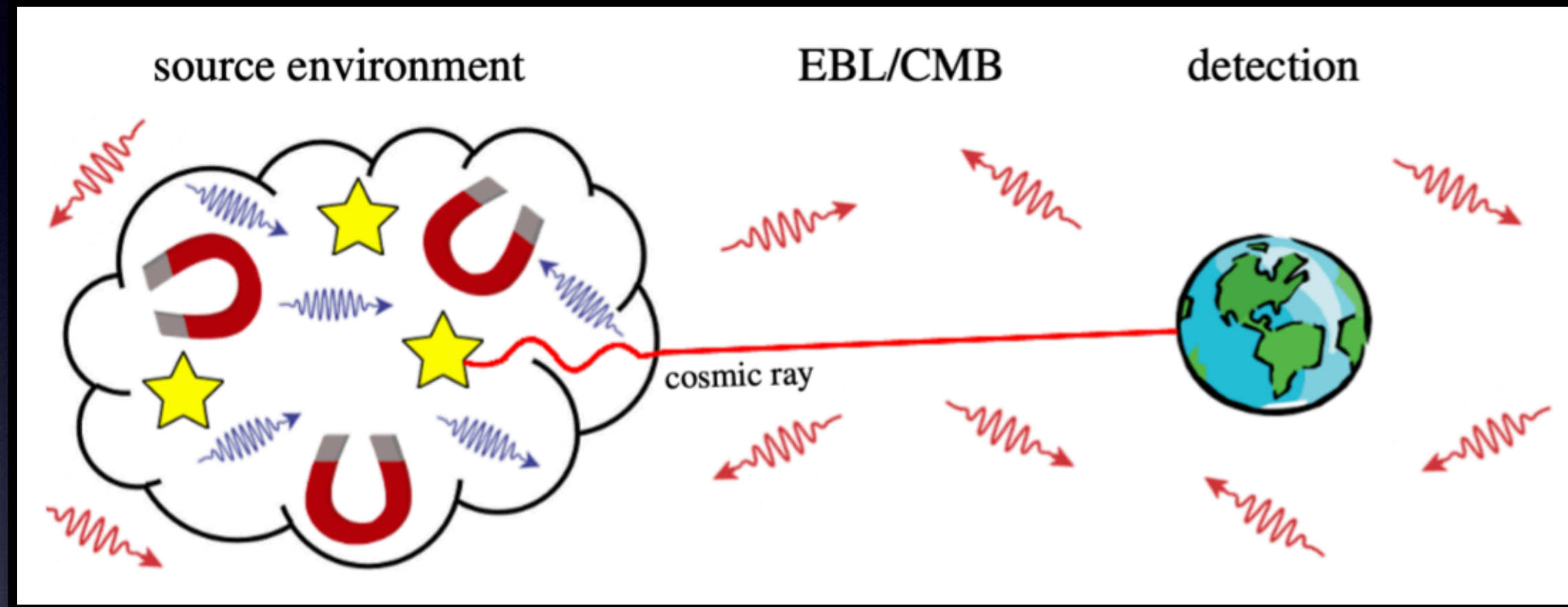
$p \sim 2.4$

$\delta \sim 1/3$

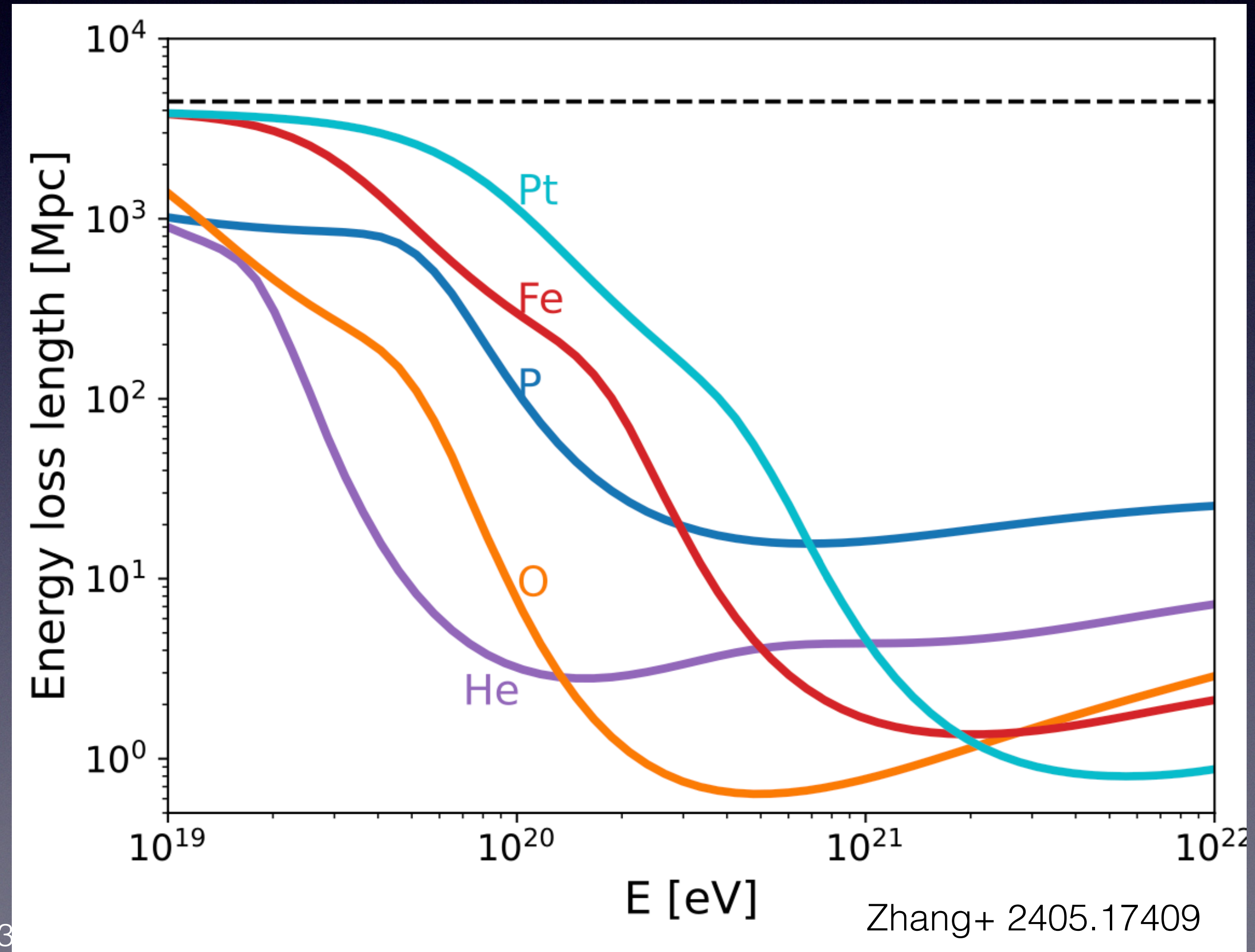
from PIC simulations of highly magnetized ($\sigma \gg 1$) turbulence



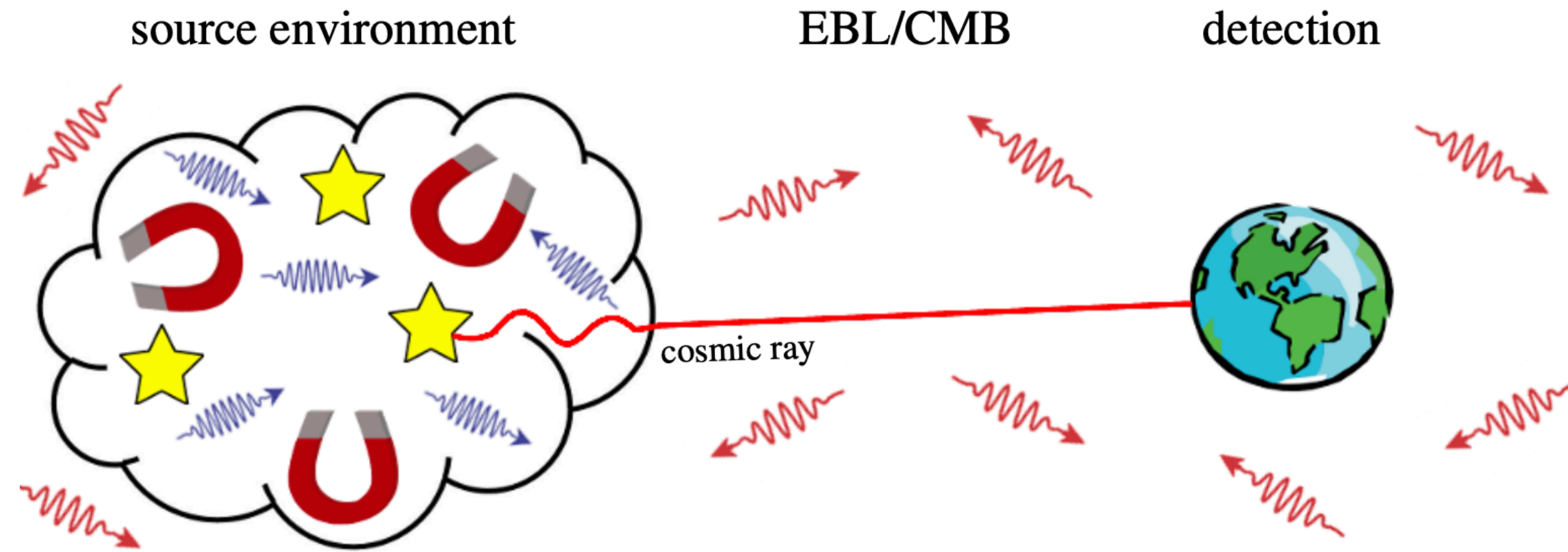
Interactions in surroundings and propagation from source



- Flexible treatment of processing leaving source
Unger, GF, Anchordoqui 2015, Muzio+GF 2023
- Hardens the spectrum, since highest rigidity particles escape more readily
- UCRs spallate in source environment, CMB and EGBL: E/A (thus E/Z) approximately constant



Particle injected by the accelerator: spectral index from fit to data

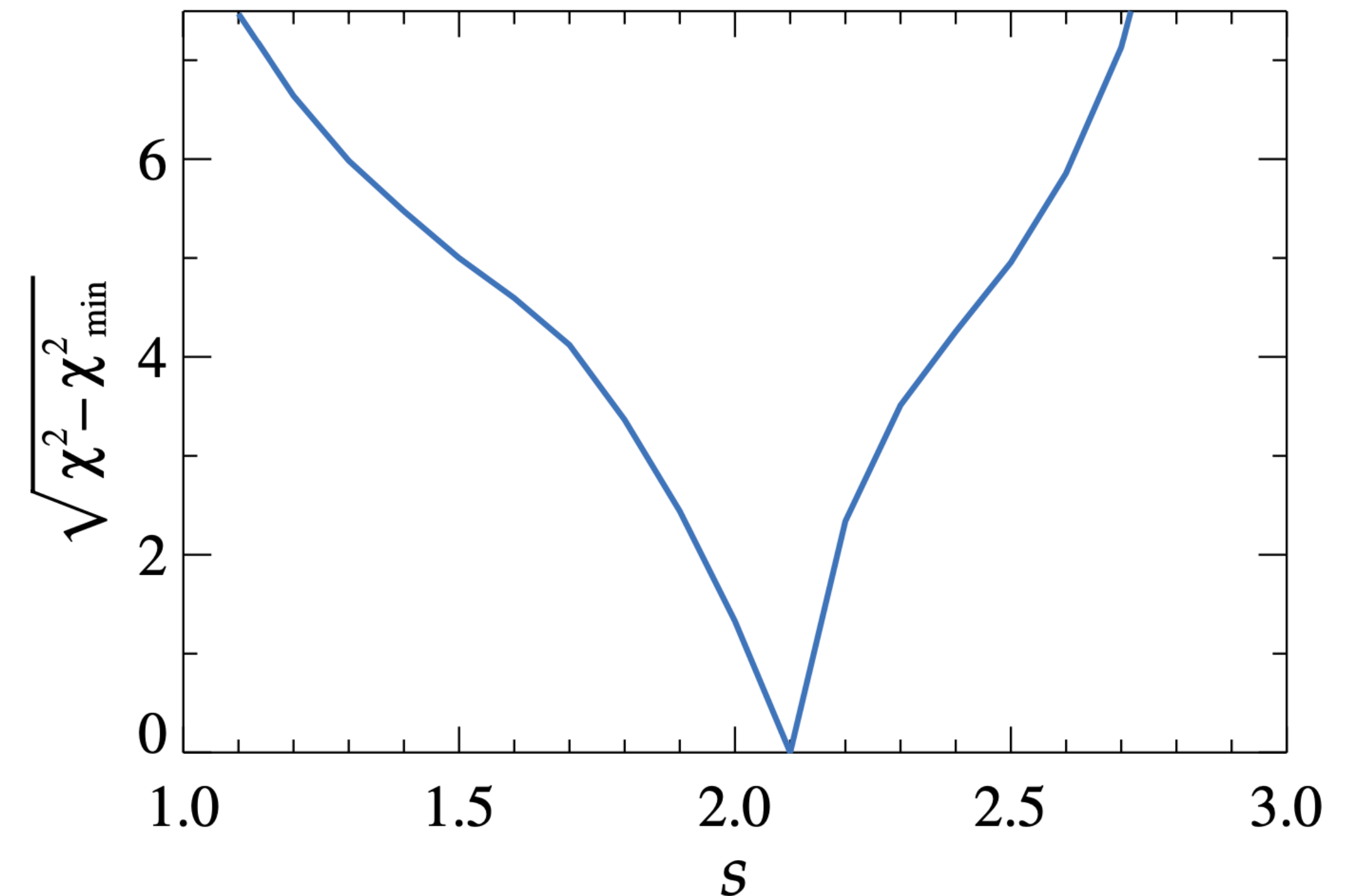


- We computed particle interaction and propagation according to Unger, Farrar, Anchordoqui 2015 (see also Muzio and Farrar 2023)

$$\chi^2 = \sum_i^{N_{\text{spec}}} \frac{(J_{m,i} - J_i)^2}{\sigma_{J,i}^2} + \sum_j^{N_{\text{comp}}} \frac{(\langle \ln A \rangle_{m,j} - \langle \ln A \rangle_j)^2}{\sigma_{\langle \ln A \rangle,j}^2} + \sum_j^{N_{\text{comp}}} \frac{(\text{Var}(\ln A)_{m,j} - \text{Var}(\ln A)_j)^2}{\sigma_{\text{Var}(\ln A),j}^2}$$

- Best fit to data return $s = 2.1^{+0.06}_{-0.13}$

$E^{-s} \exp(-E/E_{\text{cut}})$ require $s \approx 1$ and get much worse fit



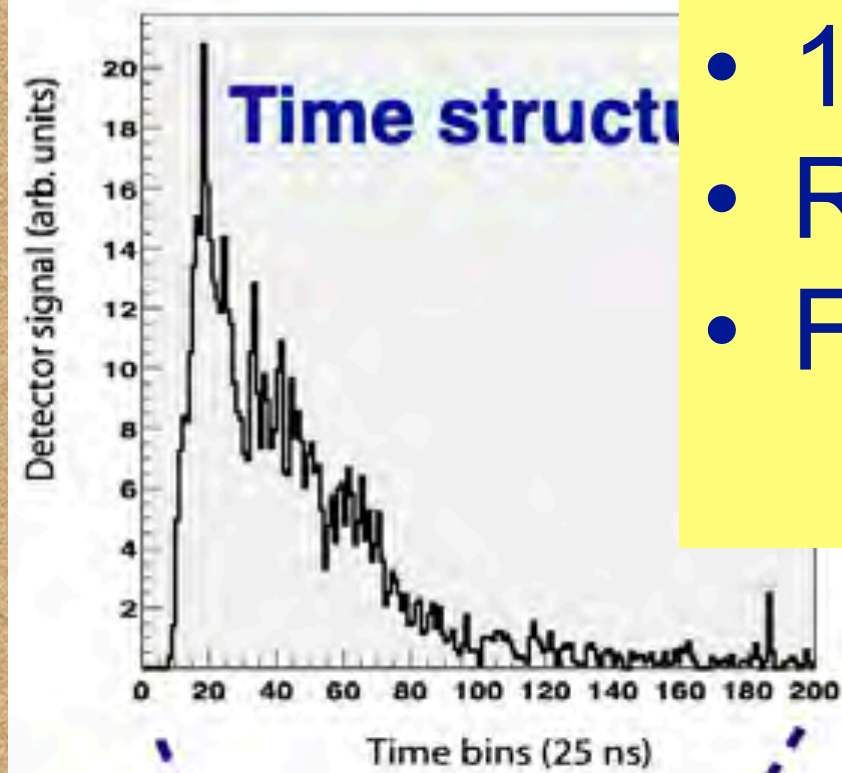
Does spectral cutoff discriminate between acceleration mechanisms?

- The $\text{sech}[(E/E_{\text{cut}})^2]$ spectral cutoff of magnetized turbulence fits well, but is it generic?
- Analytic treatment (Protheroe+Stanev 1999): $\text{DSA} \rightarrow E^{-2} \exp(-E/E_{\text{cut}})$ or softer
 - $\exp(-E/E_{\text{cut}})$ cutoff gives poor fit to UCR data while $\text{sech}[(E/E_{\text{cut}})^2]$ cutoff fits well (Comisso, GRF, Muzio ApJL 2024)
- **Must measure spectral cutoff in PIC simulations, for other acceleration mechanisms!**
- **Should also measure:**
 - “Uptake efficiency” versus Z & A $\sim (Z/A), (Z/A)^2, \dots???$
 - What is the low energy (rigidity) cutoff? What governs it?
 - Evolution of U_B while CRs are accelerated? Does CR acceleration sap U_B and “shut down”?

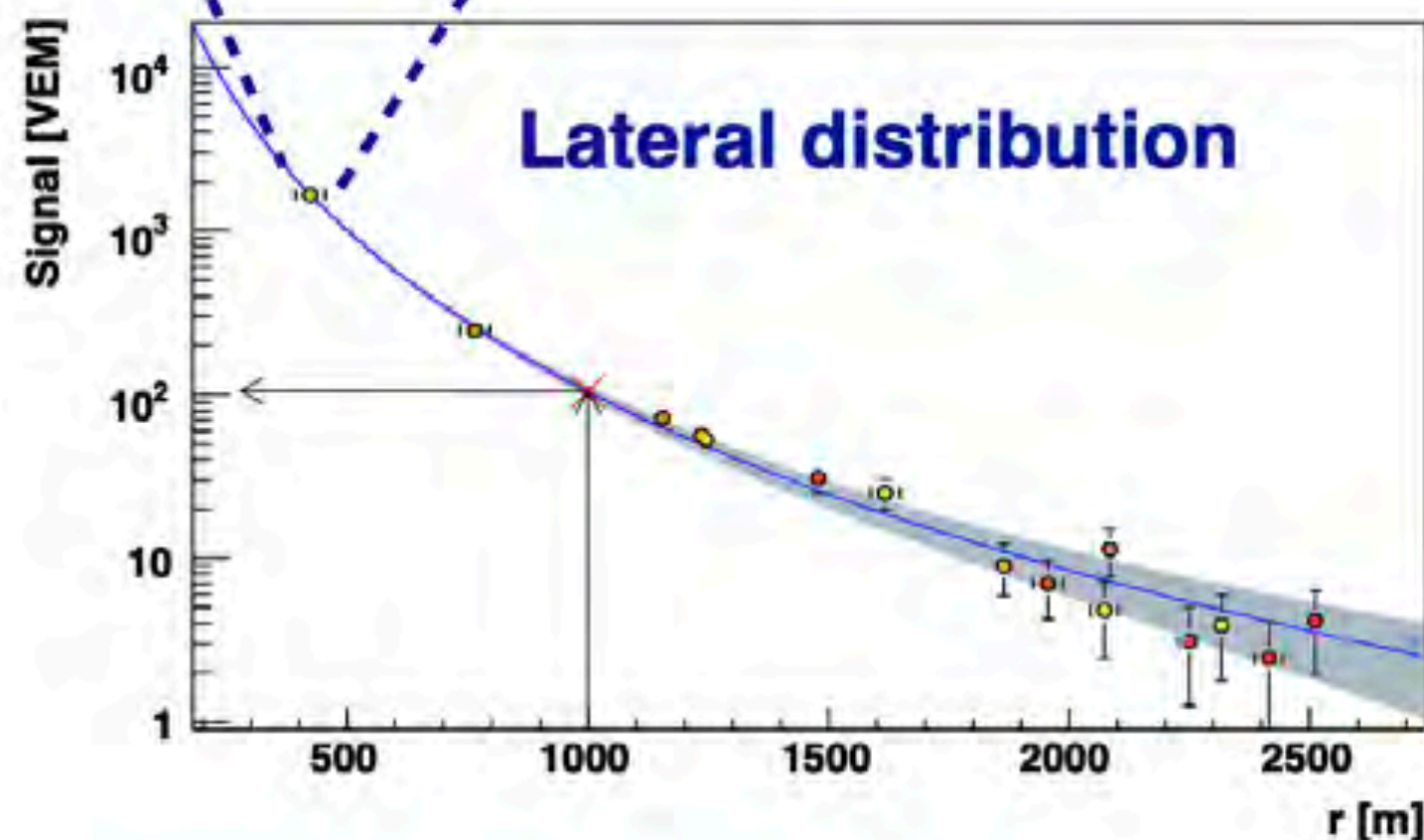
Air shower observables (hybrid observation)

Key components of UHECR observatory:

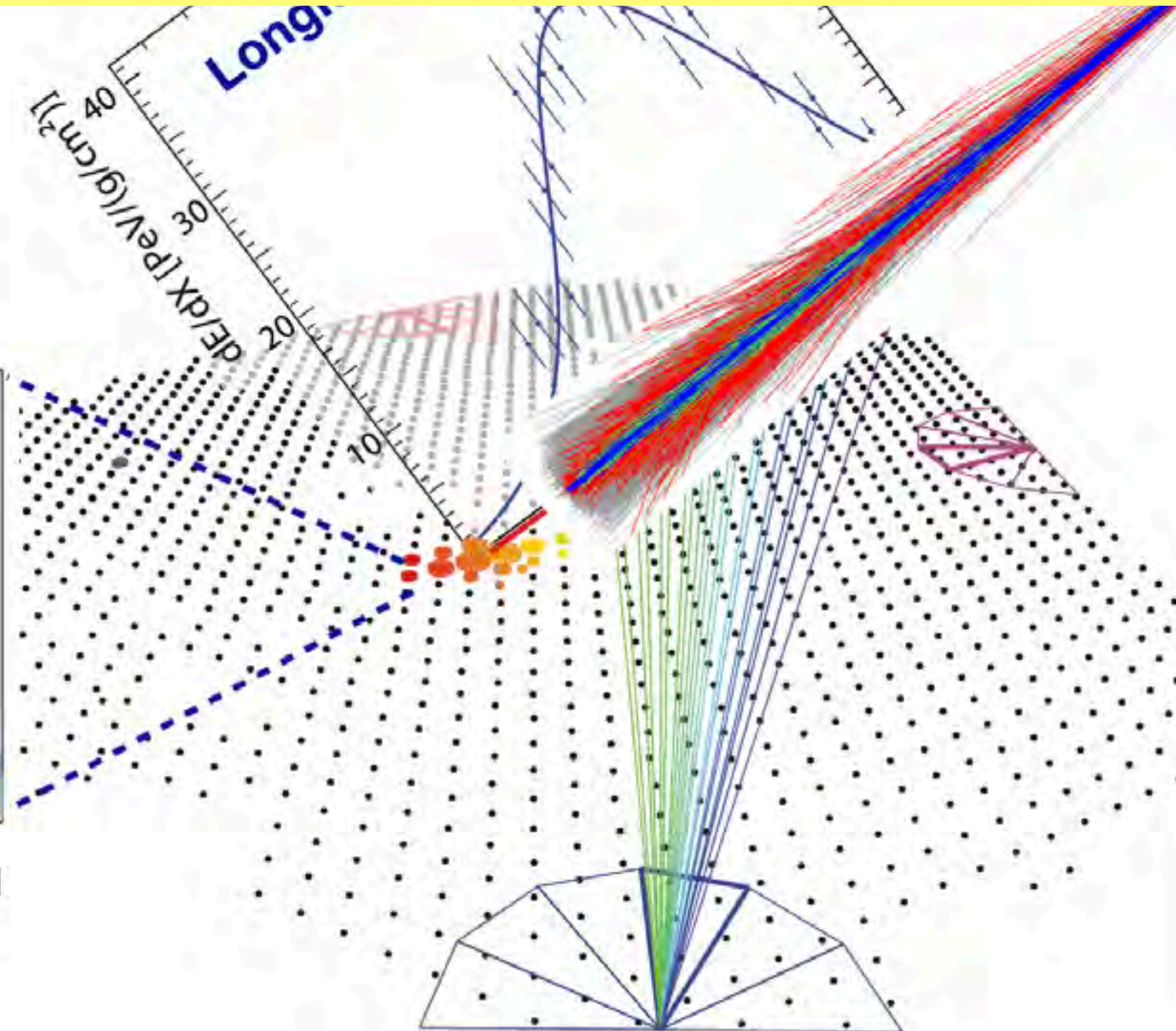
- 1600 (ea) Water Cherenkov & Scintillator Detectors, 1.5 km spacing (100%)
- Radio (100%, best for large zenith angle)
- Fluorescence Detector → Longitudinal profile (15%)



$$E_{\text{rec}} = f(S_{1000}, \theta)$$

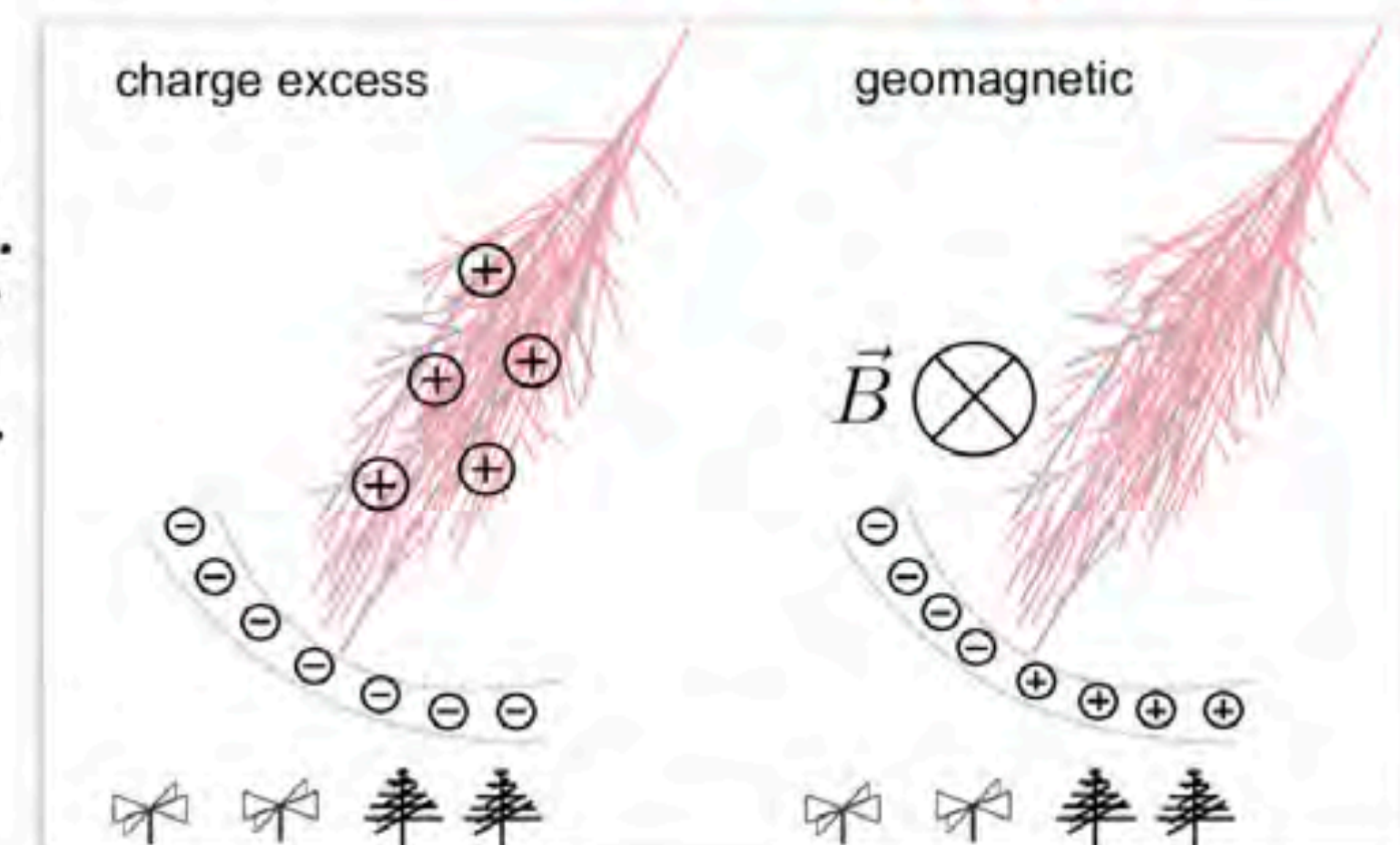


Surface Detector (SD)
100% duty cycle



$$E_{\text{cal}} = \int_0^\infty \left(\frac{dE}{dX} \right)_{\text{obs}} dX$$

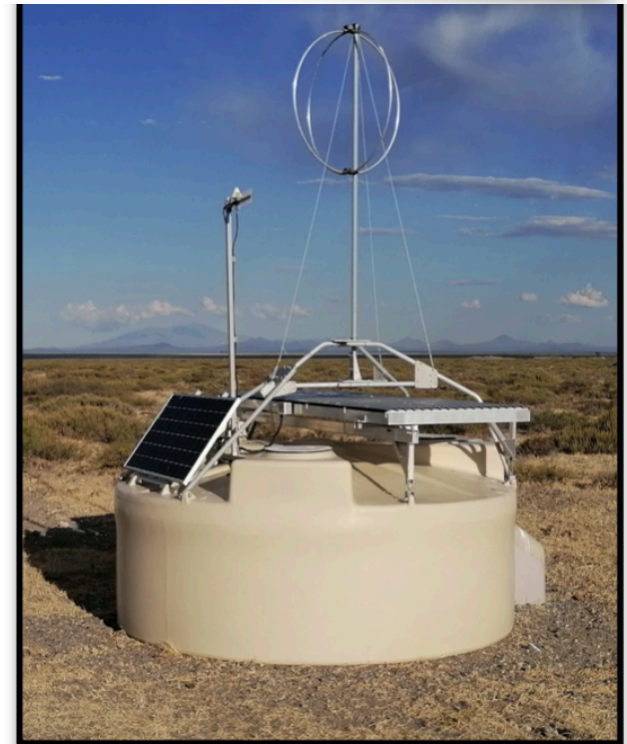
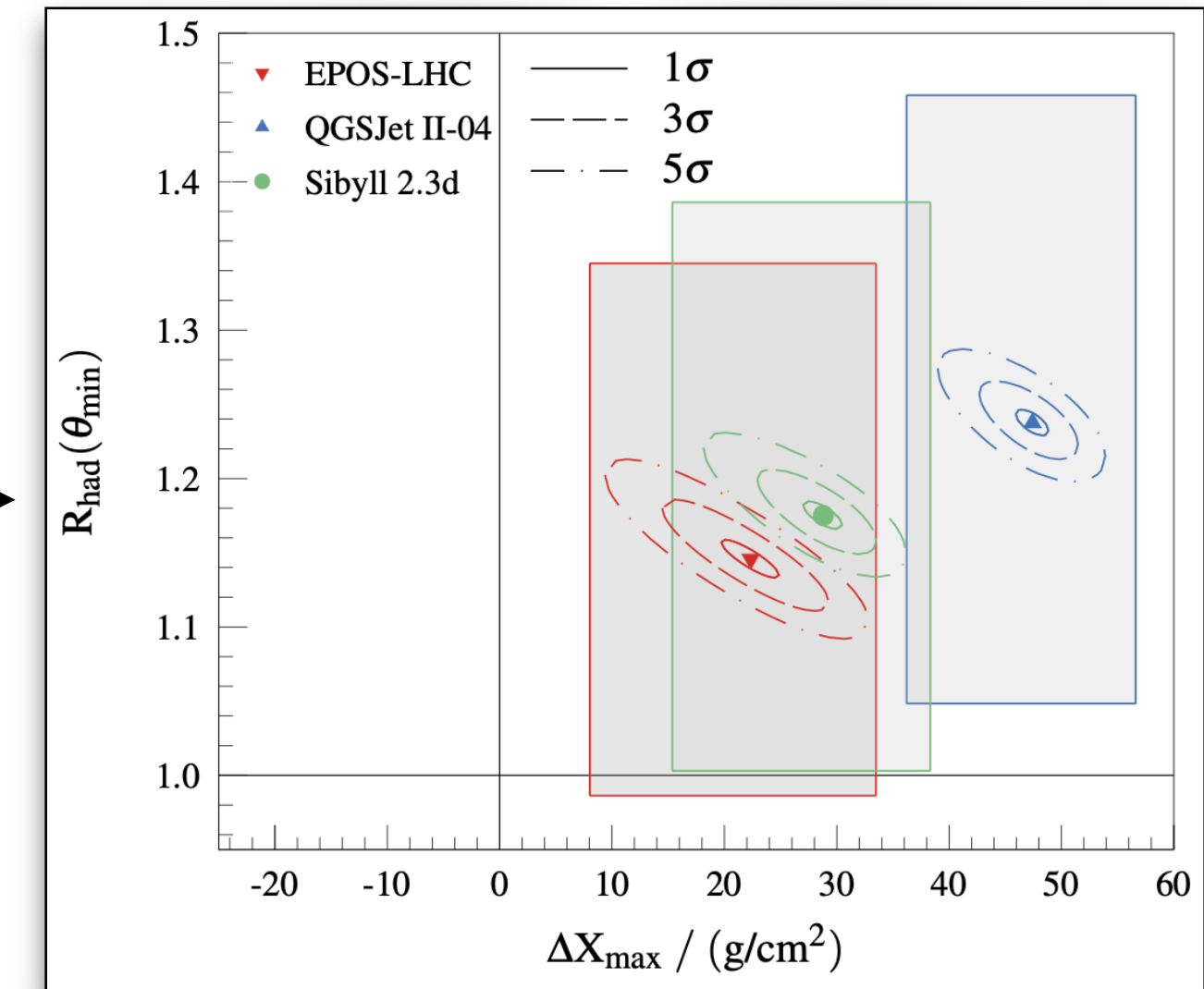
Radio Detector (RD):
100% duty cycle



Achieving correct hadronic interaction models (HIMs)

Auger *Phys.Rev.D* 109 (2024) 10, 102001

- No accelerator-tuned HIM accurately describes the muon content and X_{\max} seen in UHECR shower observations
- Phase II tools should identify source of the problem
 - Underground muons \rightarrow *muon spectral info*
 - SSD/RD \rightarrow *more precise EM/hadronic separation*
- WCD + SSD/RD + FD + UMC \rightarrow **MULTI-HYBRID** composition assignment
 - Phase II + Machine Learning *enables quality composition estimation for all Phase I data* (>60k events above the ankle)
- **Accurate HIMs + multi-hybrid composition \rightarrow robust A, Z inference**



“Muon Problem”

- Ground signal (S_{38}) & X_{\max} distribution should not depend on zenith

- Muon problem → muon AND X_{\max} problems

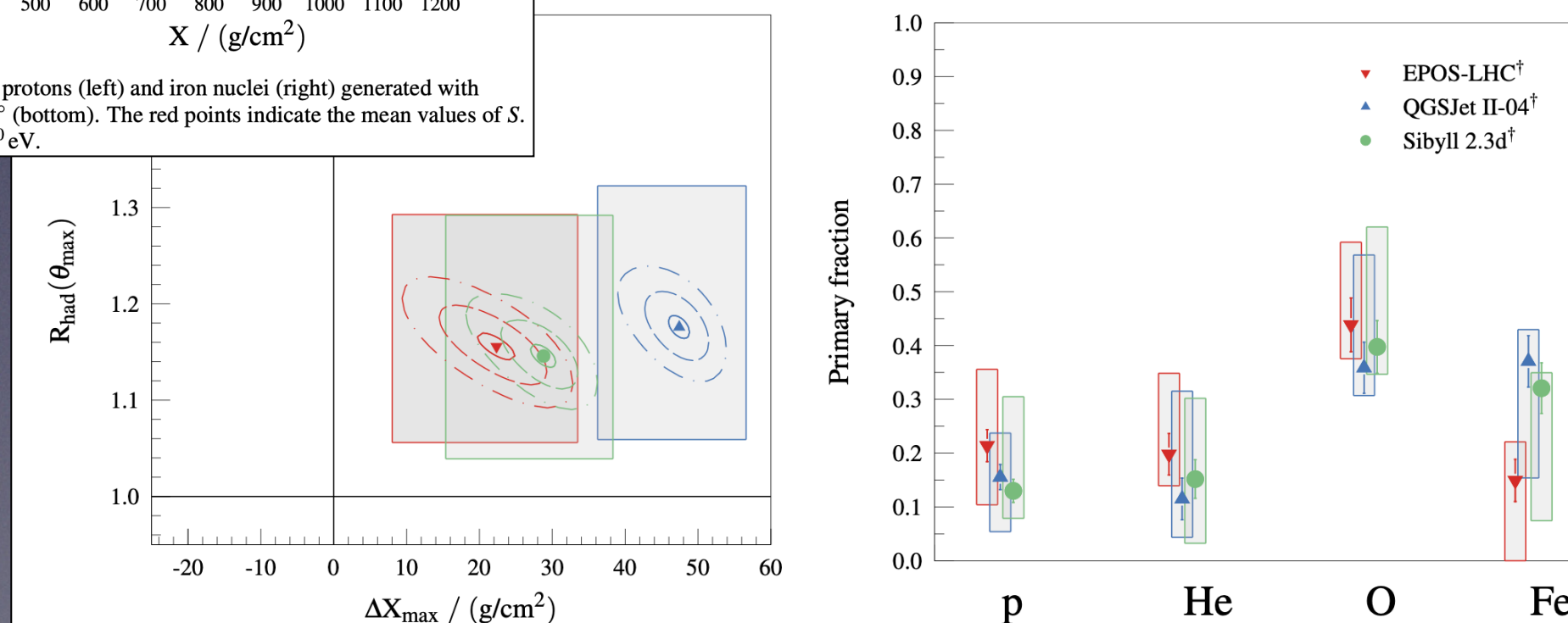
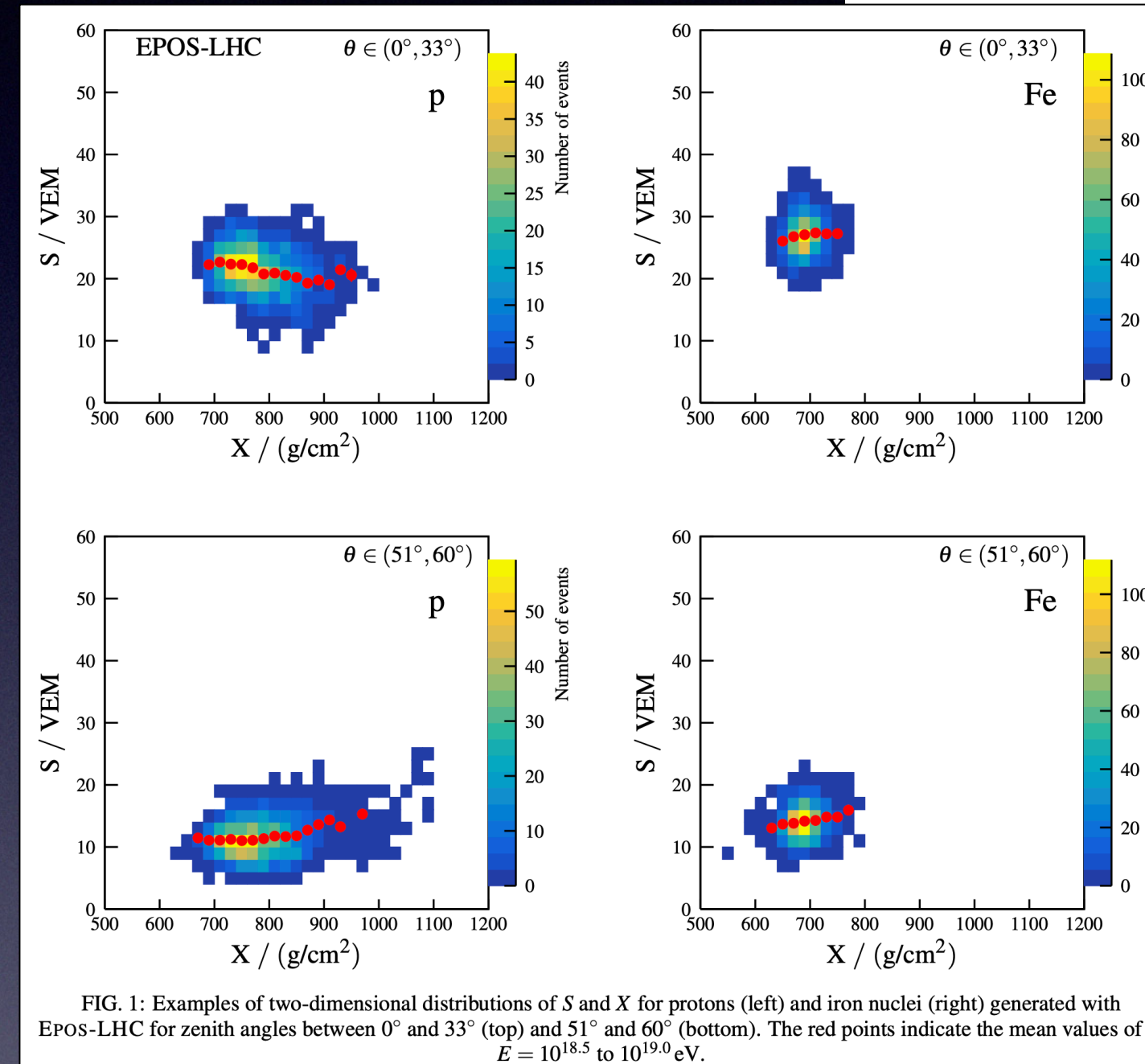
We found that for the best description of the data distributions in the energy range $10^{18.5}$ to $10^{19.0}$ eV for $\theta < 60^\circ$ the MC predictions of X_{\max} should be deeper in the atmosphere by about 20 to 50 g/cm², and the hadronic signal should be increased by about 15 to 25% in all three models. These modifications reduce the differences between the models in X_{\max} and $S(1000)$, and as a consequence, lead to smaller uncertainties on the estimated fractions of the primary nuclei. Due to the deeper MC X_{\max} scale and, correspondingly, a heavier mass composition inferred from the data compared with non-modified models, the scaling factors for the hadronic signal are found to be smaller than in previous estimations not considering any modifications to the MC X_{\max} scales. The

- After shift, composition determination agrees between models, and becomes heavier than before.

Testing Hadronic-Model Predictions of Depth of Maximum of Air-Shower Profiles and Ground-Particle Signals using Hybrid Data of the Pierre Auger Observatory

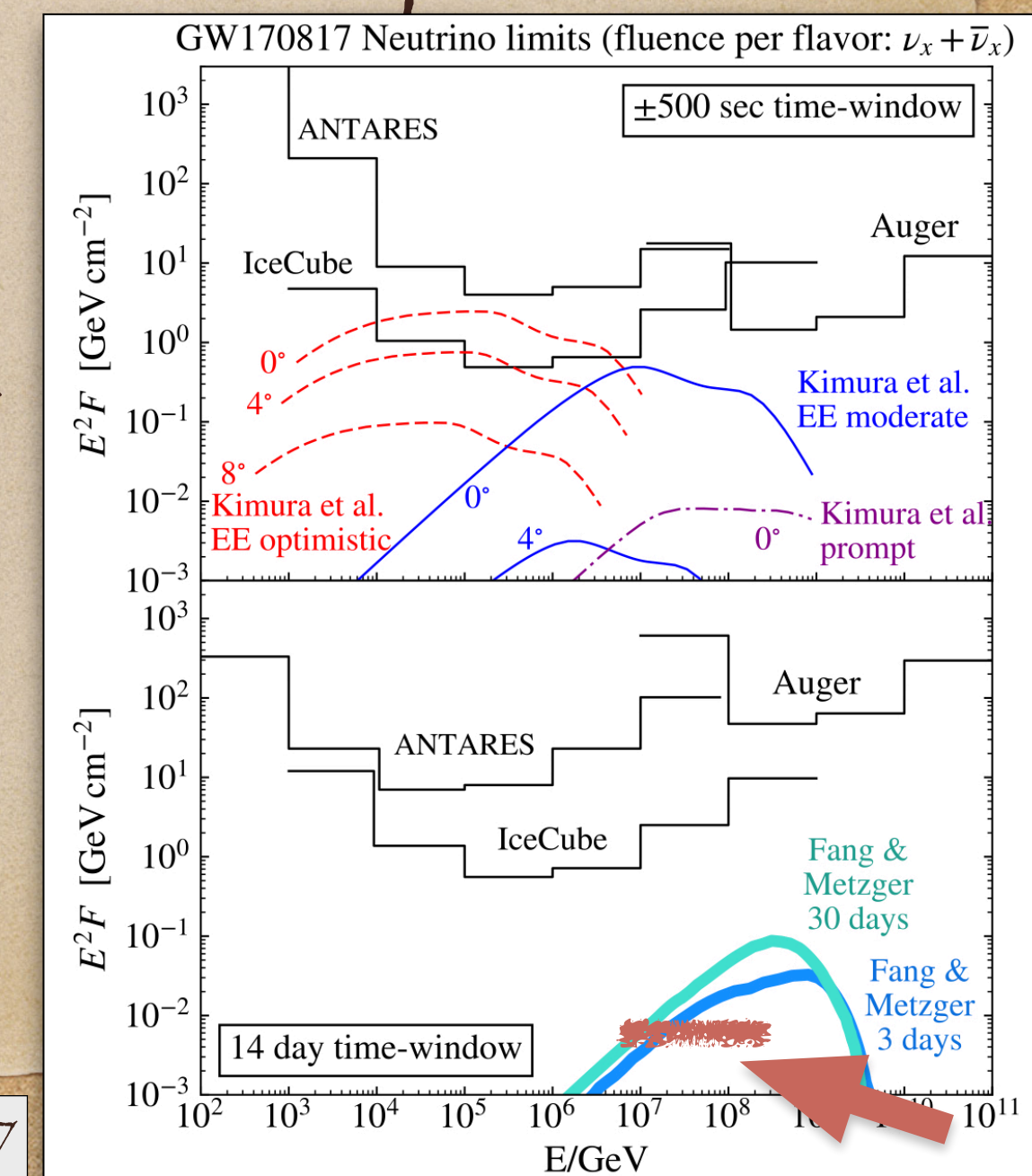
The Pierre Auger Collaboration
The Pierre Auger Observatory, Av. San Martín Norte 306,
5613 Malargüe, Mendoza, Argentina;
<http://www.auger.org>*
(Dated: February 19, 2024)

PRD2024



Future test of BNS-merger origin: EHE neutrino \approx coincident with GW from BNS merger

- EVERY EHE ν should be accompanied by a gravitational wave from the NS merger.
- Cosmic Explorer+Einstein Telescope+IceCube-Gen2 x few yrs: very promising
- GW170817 also accompanied by EHE neutrinos but estimated fluence for favorable case of aligned jet $\ll 0.15 \text{ GeV cm}^{-2}$ per flavor.
Sensitivity not adequate by orders of magnitude



r-process nucleosynthesis B²FH

REVIEWS OF MODERN PHYSICS

VOLUME 29, NUMBER 4

OCTOBER, 1957

Synthesis of the Elements in Stars*

E. MARGARET BURBIDGE, G. R. BURBIDGE, WILLIAM A. FOWLER, AND F. HOYLE

*Kellogg Radiation Laboratory, California Institute of Technology, and
Mount Wilson and Palomar Observatories, Carnegie Institution of Washington,
California Institute of Technology, Pasadena, California*

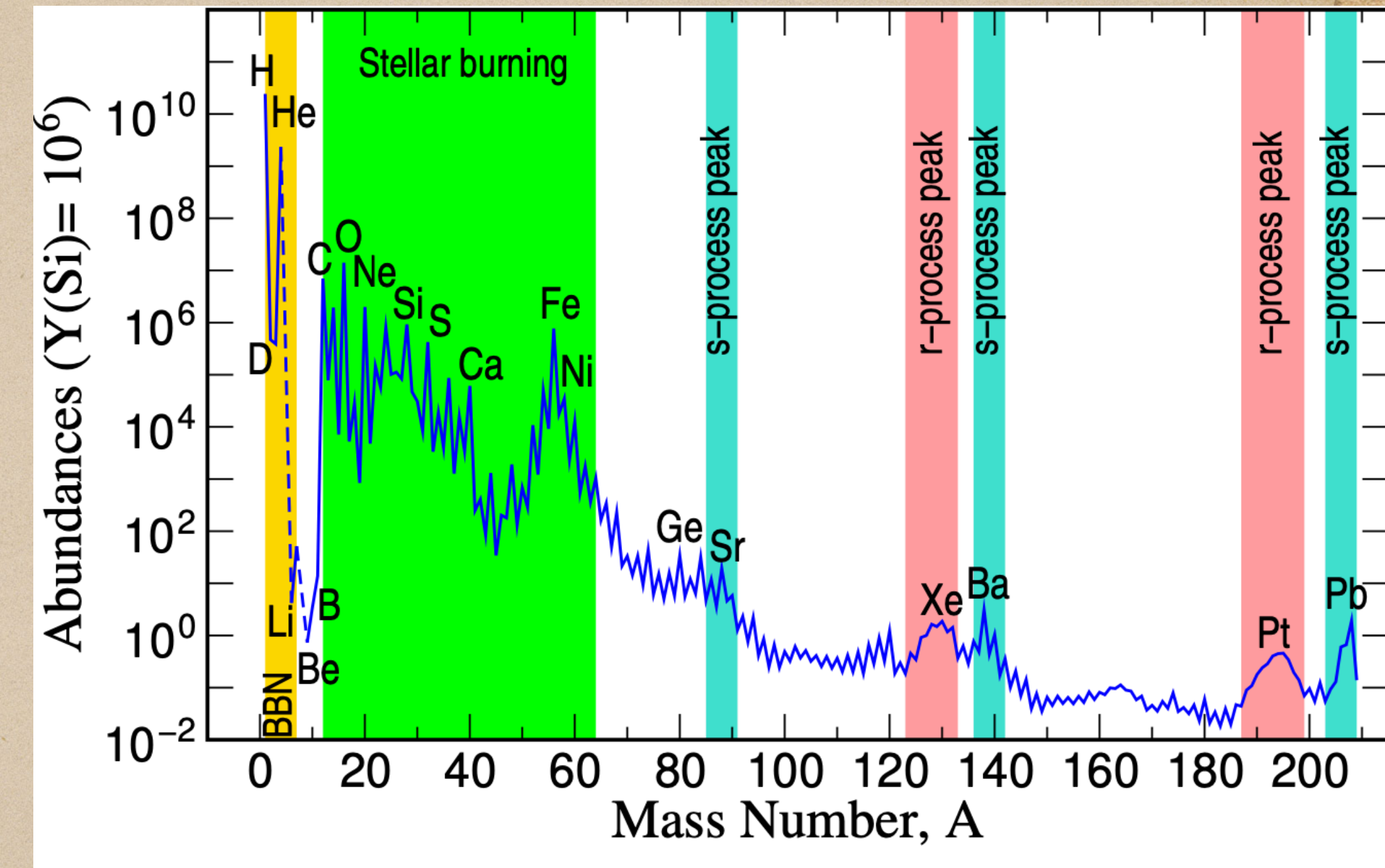
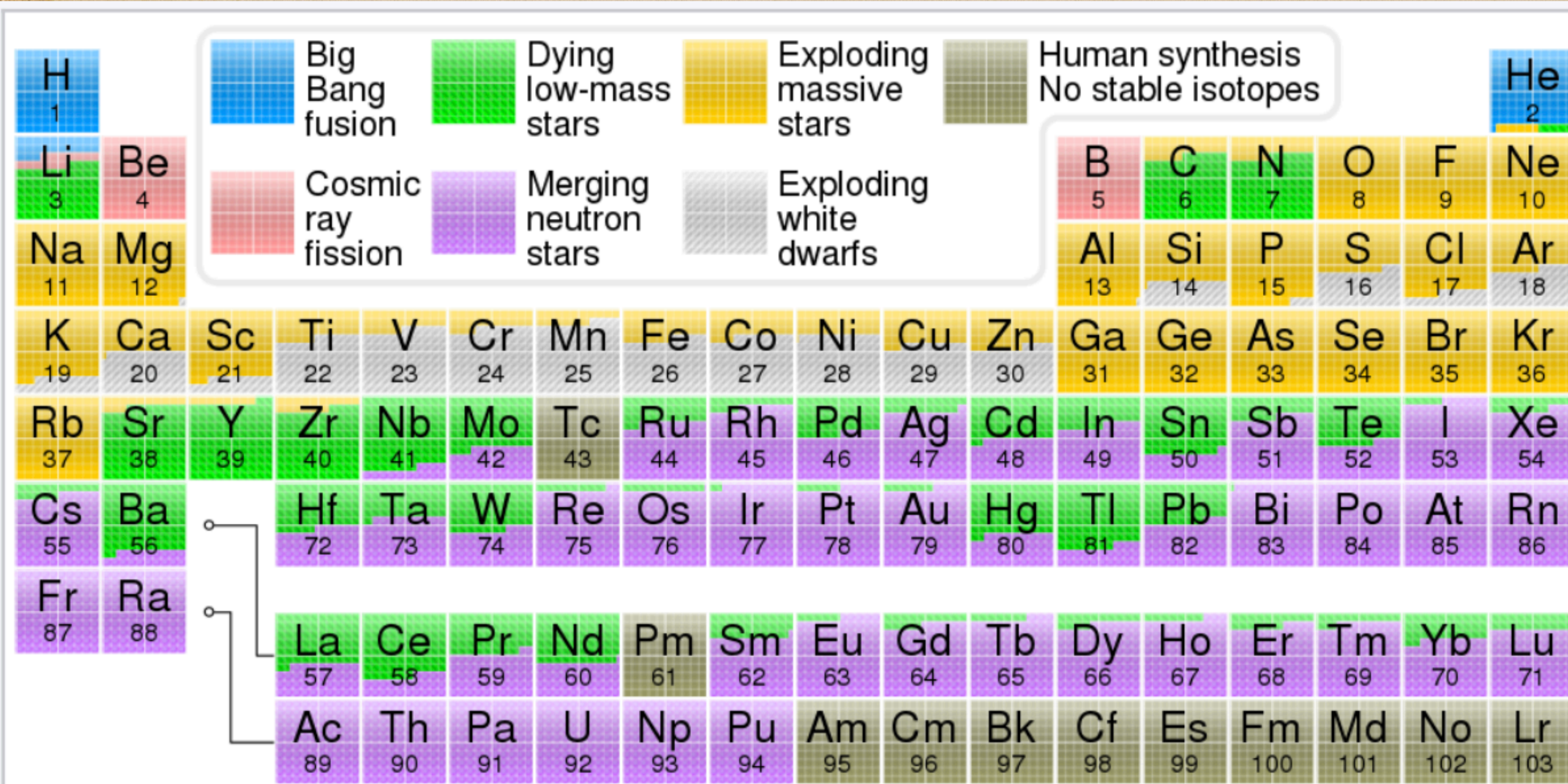
"It is the stars, The stars above us, govern our conditions";
(*King Lear*, Act IV, Scene 3)

but perhaps

"The fault, dear Brutus, is not in our stars, But in ourselves,"
(*Julius Caesar*, Act I, Scene 2)

r process.—The nuclear physics of this process demands that neutrons be added extremely rapidly, so that the total time-scale for the addition of a maximum of about 200 neutrons per iron nucleus is ~ 10 – 100 sec.

Merging NS's produce "r-process" elements



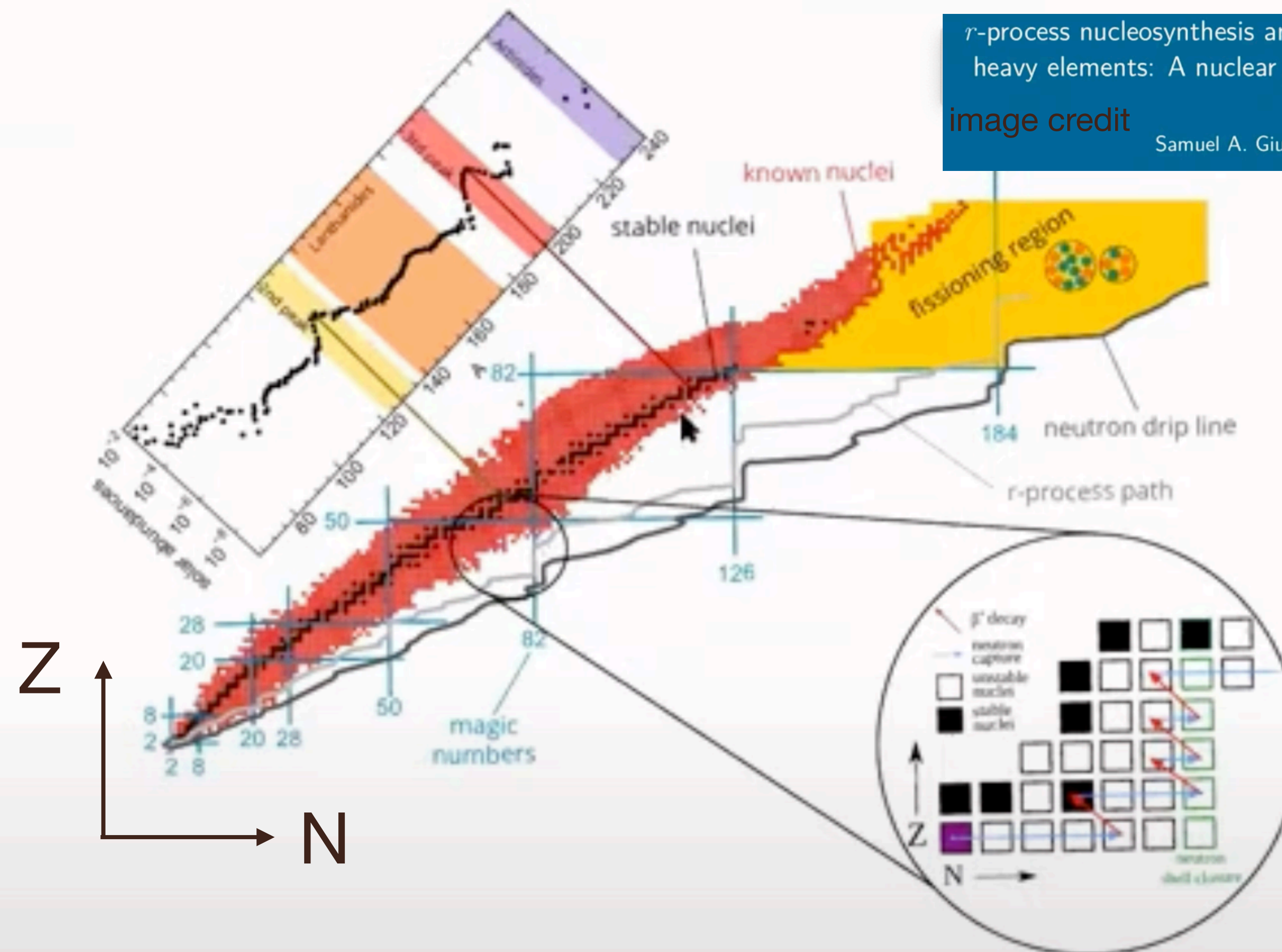
Periodic table showing the cosmogenic origin of each element. The elements heavier than iron with origins in supernovae are typically those produced by the *r*-process, which is powered by supernova neutron bursts

The r process

B²FH, Rev. Mod. Phys. 29, 547 (1957) ; A. Cameron, Report CRL-41 (1957)

r (apid neutron capture) process: $\tau_{(n,\gamma)} \ll \tau_{\beta-}$

And see Nicole Vassh movie <https://www.google.com/search?client=safari&rls=en&q=r-process+nucleosynthesis+movie&ie=UTF-8&oe=UTF-8#fpstate=ive&vld=cid:94d4d99d,vid:P1tHGLdXRTw,st:361>



- The path to heavier nuclei goes through **neutron-rich nuclei**.

Major Upgrades Underway

V.Versi, UHECR22

AugerPrime

**Malargüe Mendoza
(Argentina)**
35° S latitude

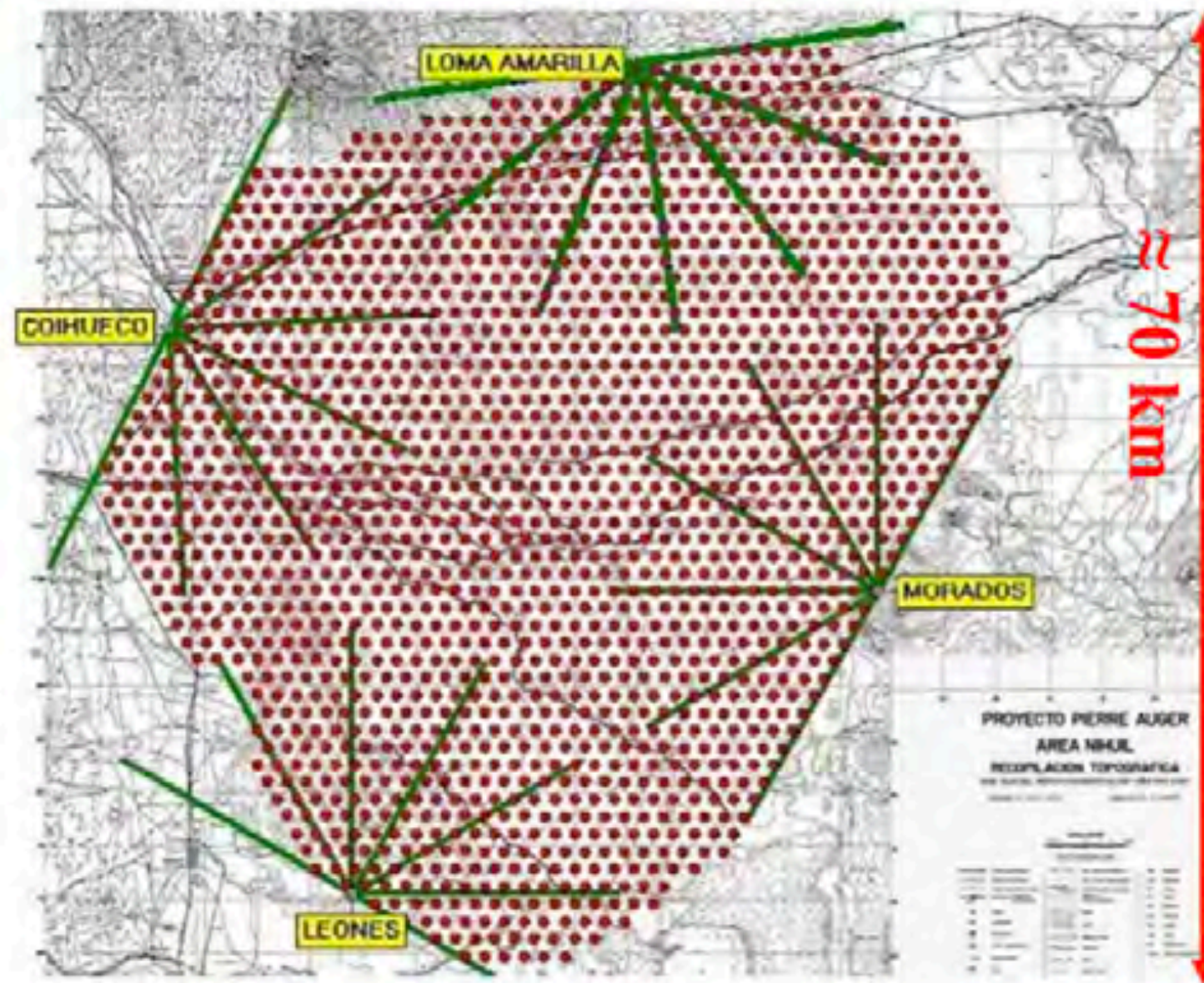
3000 km²

1660 WCDs
1500 m spacing
triangular grid

4 FD sites

WCD = water Cherenkov detector; FD = fluorescence detector

Adding capability for better particle ID:
Scintillators; underground muon detectors...



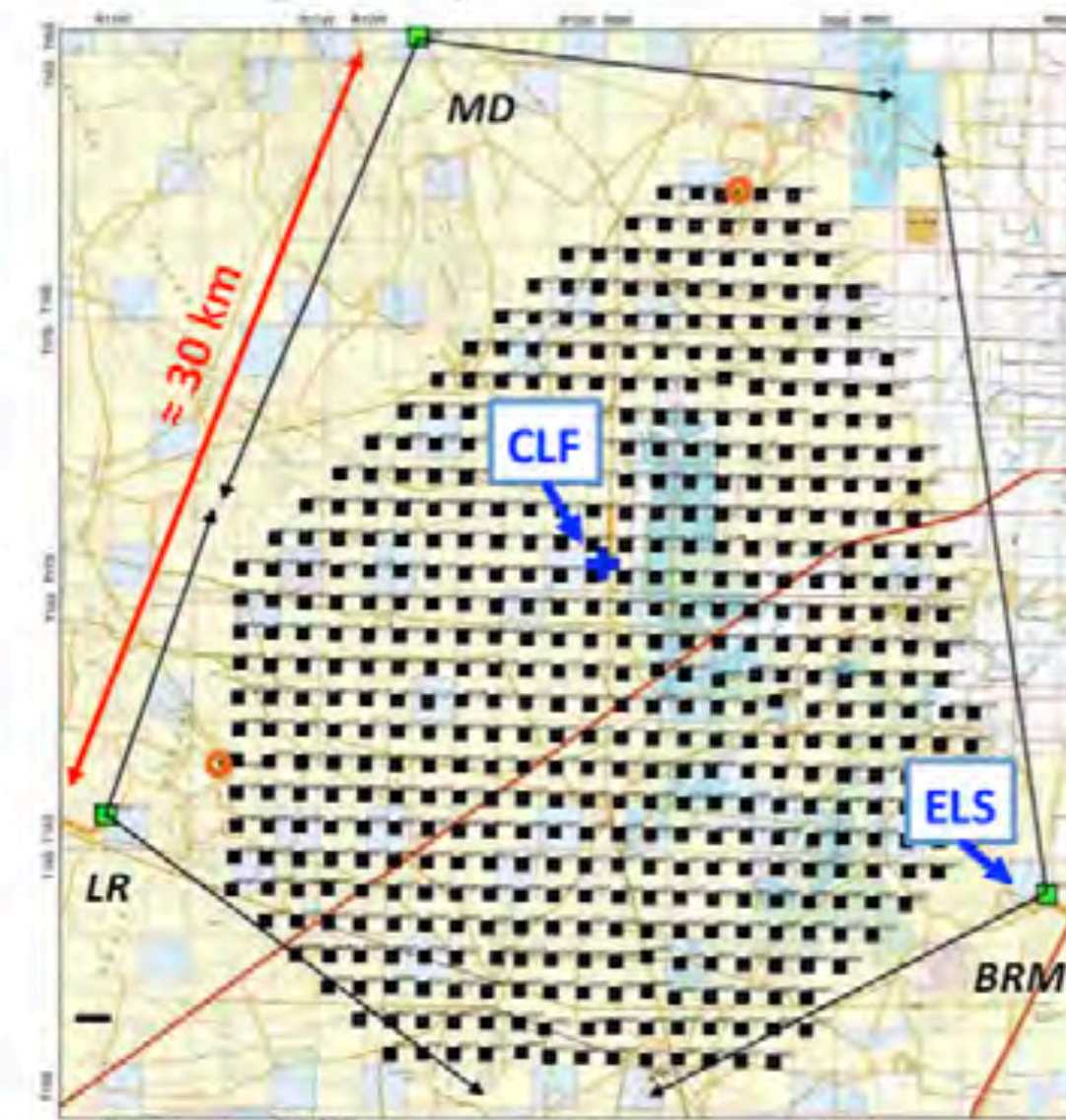
TA x 4

**Millard County
Utah (USA)**
39° N latitude

700 km²

507 scintillators
1200 m spacing
square grid

3 FD sites



Scintillator



Fluorescence Detector

