

# Gravitational Laboratories for Nuclear Physics

Prospects for Binary Neutron Star Observations  
and their Impact on the Equation of State

Reed Essick  
[ressick@perimeterinstitute.ca](mailto:ressick@perimeterinstitute.ca)

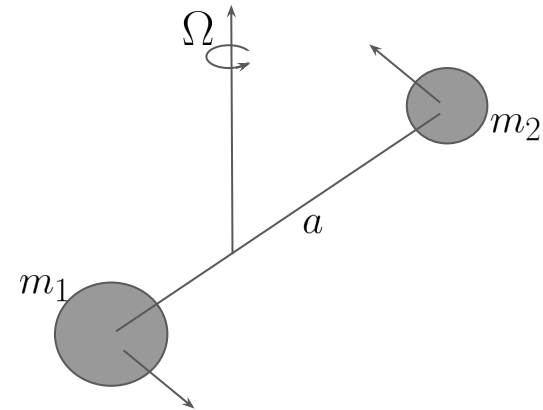
Perimeter Institute for Theoretical Physics

## Basic Physics of Compact Binary Coalescences (CBCs)

$$E = -\frac{Gm_1m_2}{a}$$

$$\Omega^2 = \frac{G(m_1 + m_2)}{a^3}$$

Keplerian Orbit



# Basic Physics of Compact Binary Coalescences (CBCs)

$$\langle P \rangle = \frac{32G^4}{5c^5} \left( \frac{m_1^2 m_2^2 (m_1 + m_2)}{a^5} \right) \left( \frac{1}{(1 - e^2)^{7/2}} \left( 1 + \frac{73}{24}e^2 + \frac{37}{96}e^4 \right) \right)$$

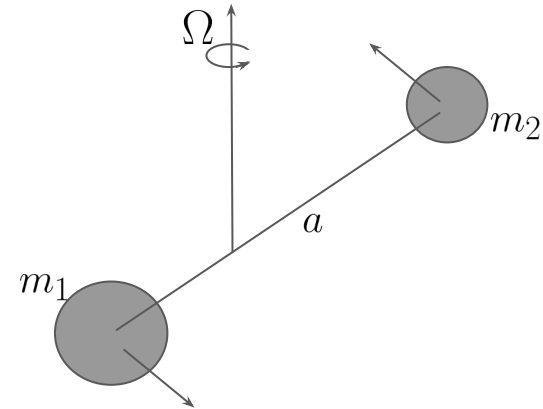
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"Larmor formula"

circular orbit :  $e = 0$



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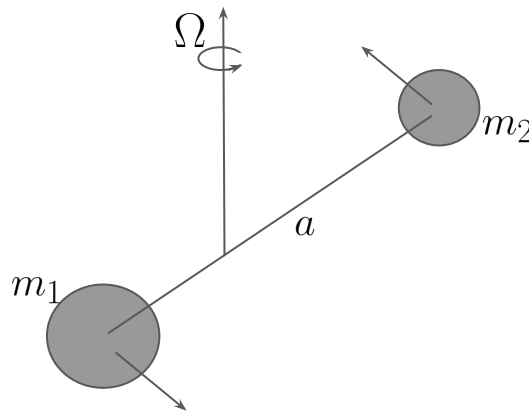
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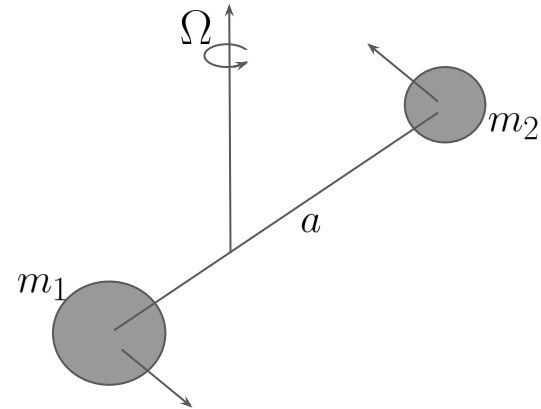
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Post Newtonian Energy Balance

$$f(m_1, m_2, \Omega) = g(m_1, m_2, \Omega) \frac{d\Omega}{dt}$$

evolution of the orbital  
frequency is directly  
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$$h \sim e^{i(2\Omega)t}$$



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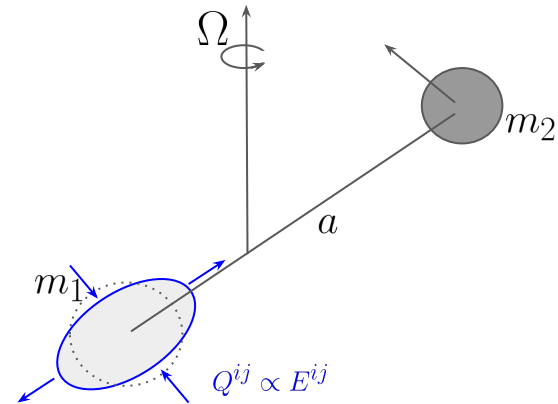
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changes flux and orbit's binding energy



adiabatic tides deform stars

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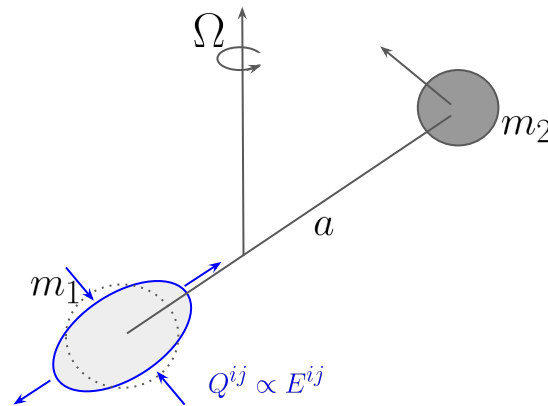
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adiabatic tides  $\Lambda(m) \sim k_2 \left( \frac{R(m)}{m} \right)^5$

$$\tilde{\Lambda} = \left( \frac{16}{13} \right) \frac{(m_1 + 12m_2)m_1^4\Lambda(m_1) + (m_2 + 12m_1)m_2^4\Lambda(m_2)}{(m_1 + m_2)^5}$$



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adiabatic tides

linear tidal resonances

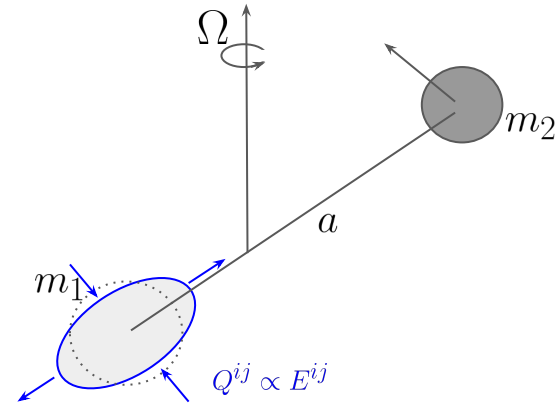
Pratten+ (2021)

nonlinear tidal instabilities

Weinberg (2016)

Abbott+ (2019)

orbital energy transferred to stellar normal modes



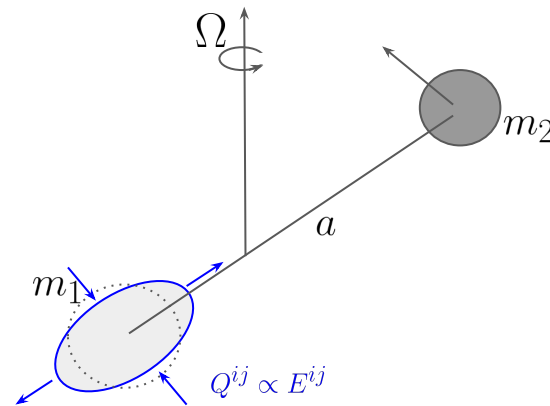
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adiabatic tides

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orbital energy transferred to stellar normal modes

post-merger signals

Most+Raithel (2021)

Weih+ (2020)

## Why are we concerned with orbital phase?

$$\phi(t) = \int_{-\infty}^t d\tau f(\tau)$$

orbital evolution gives time-domain phase

$$\psi(f) = i(2\pi ft - \phi(t)) - \pi/4$$

frequency-domain phase is related to time-domain phase  
(saddle point approximation)

$$\ln \mathcal{L} \sim \int df \frac{\mathcal{R}\{\tilde{d}^* \tilde{h}\}}{S}$$

likelihood of GW data is an integral over a **rapidly oscillating** function of the **difference of freq-domain phases**

$$\sim \int df \frac{|\tilde{d}| |\tilde{h}|}{S} \cos(\Delta\psi)$$

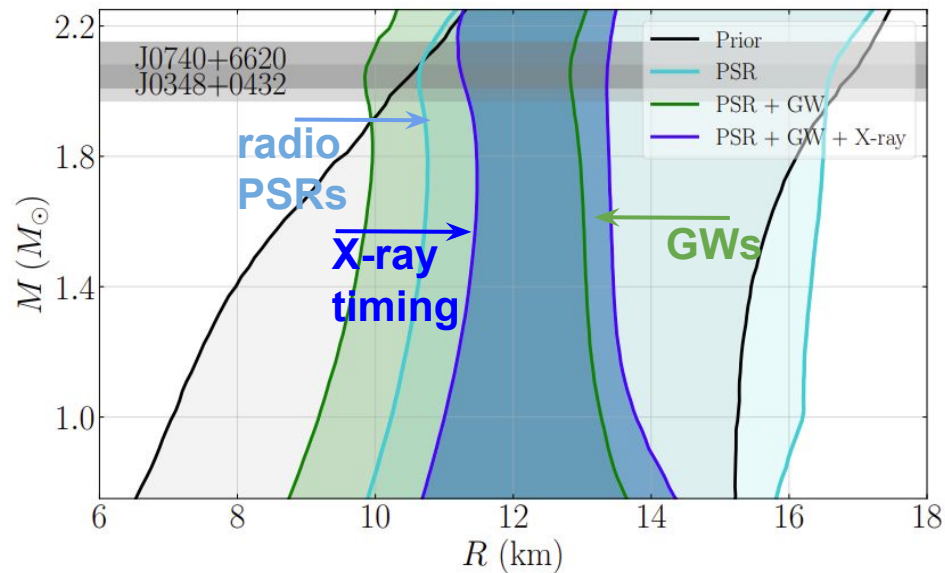
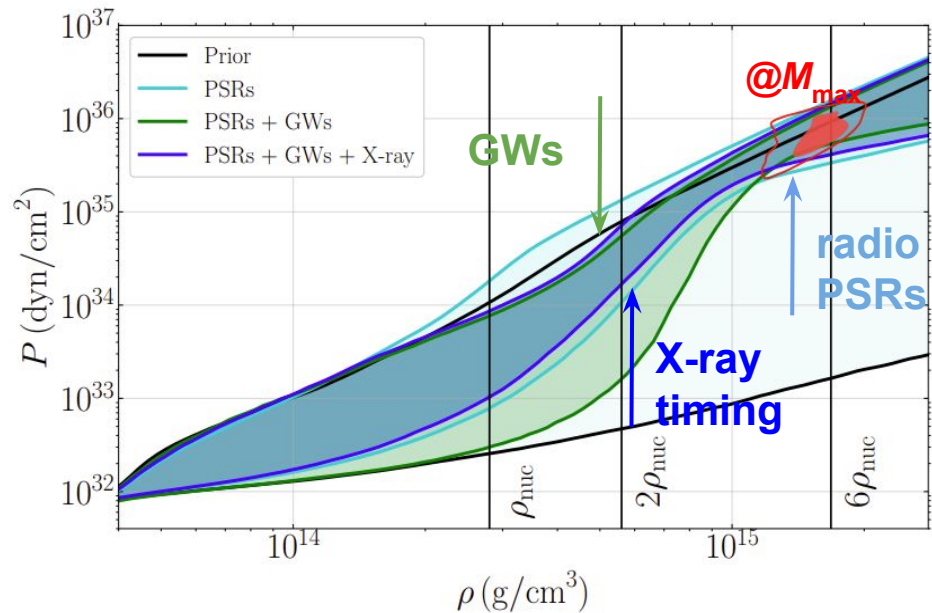
varies slowly

oscillates rapidly

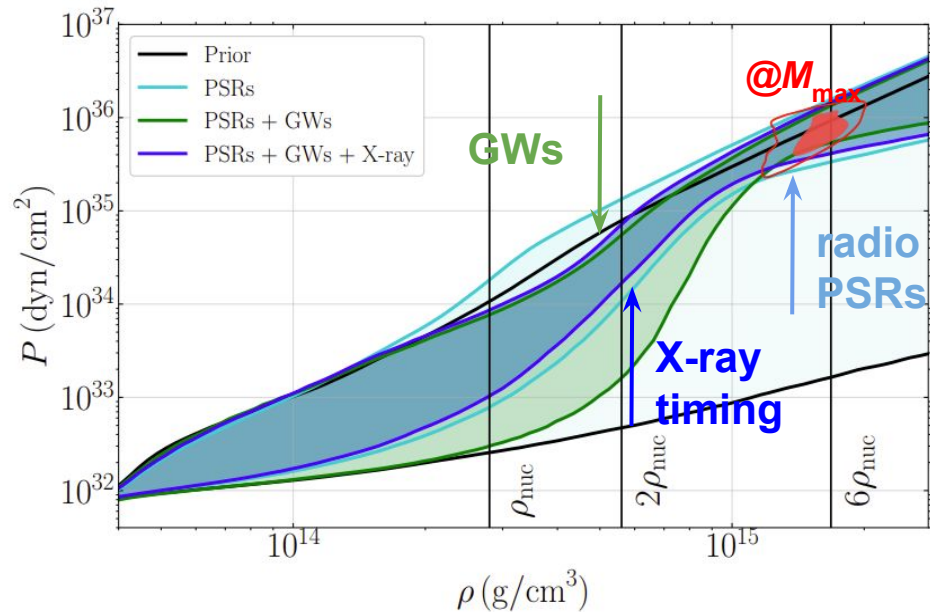
significant likelihood only when  $\Delta\psi$  is small  
and/or varies slowly at all frequencies

## Current Astrophysical Constraints on the High-Density, Cold Equation of State

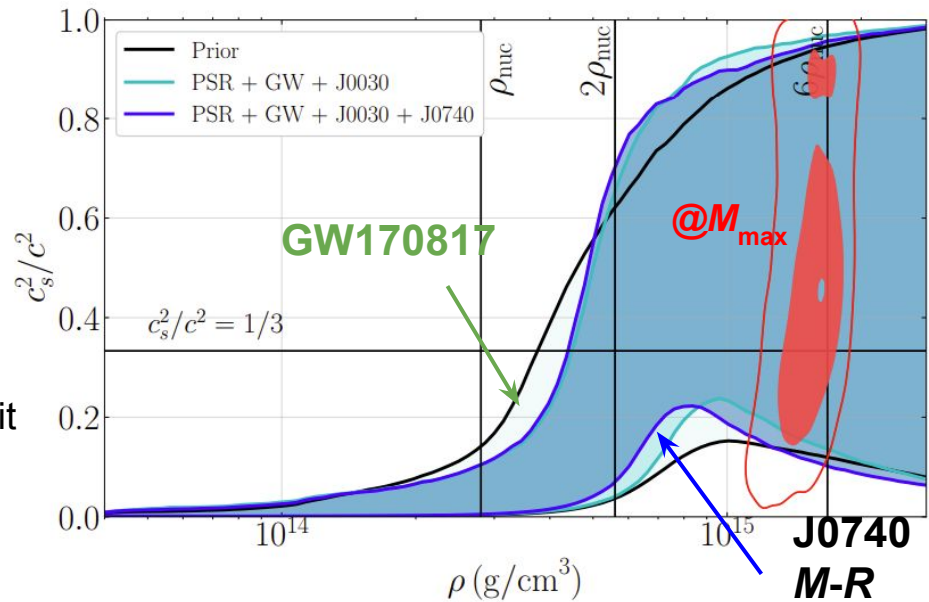
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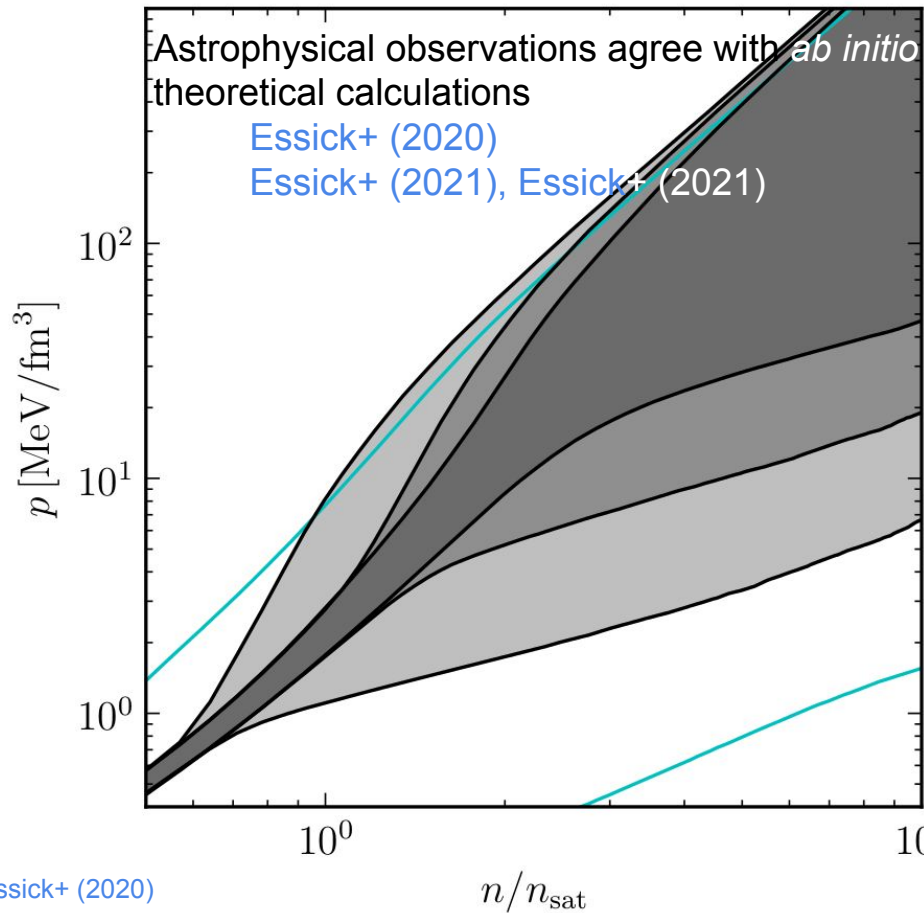
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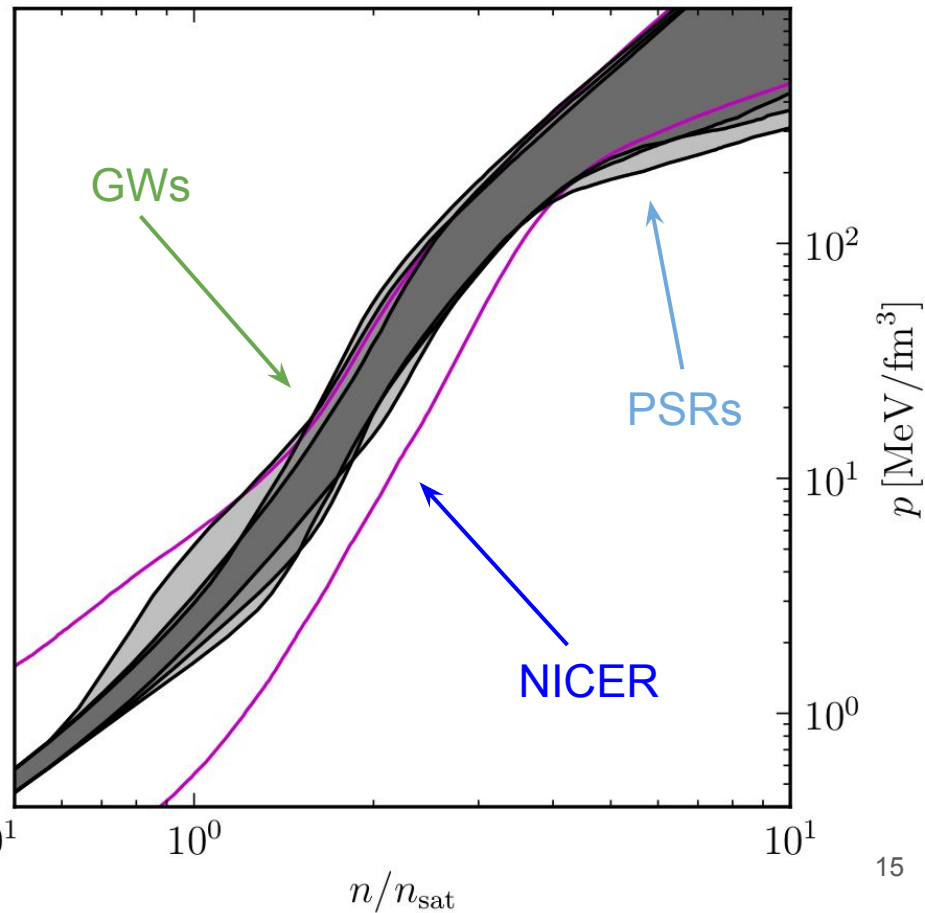
sound speed almost certainly exceeds the conformal limit  
 (suggests strongly-coupled interactions)



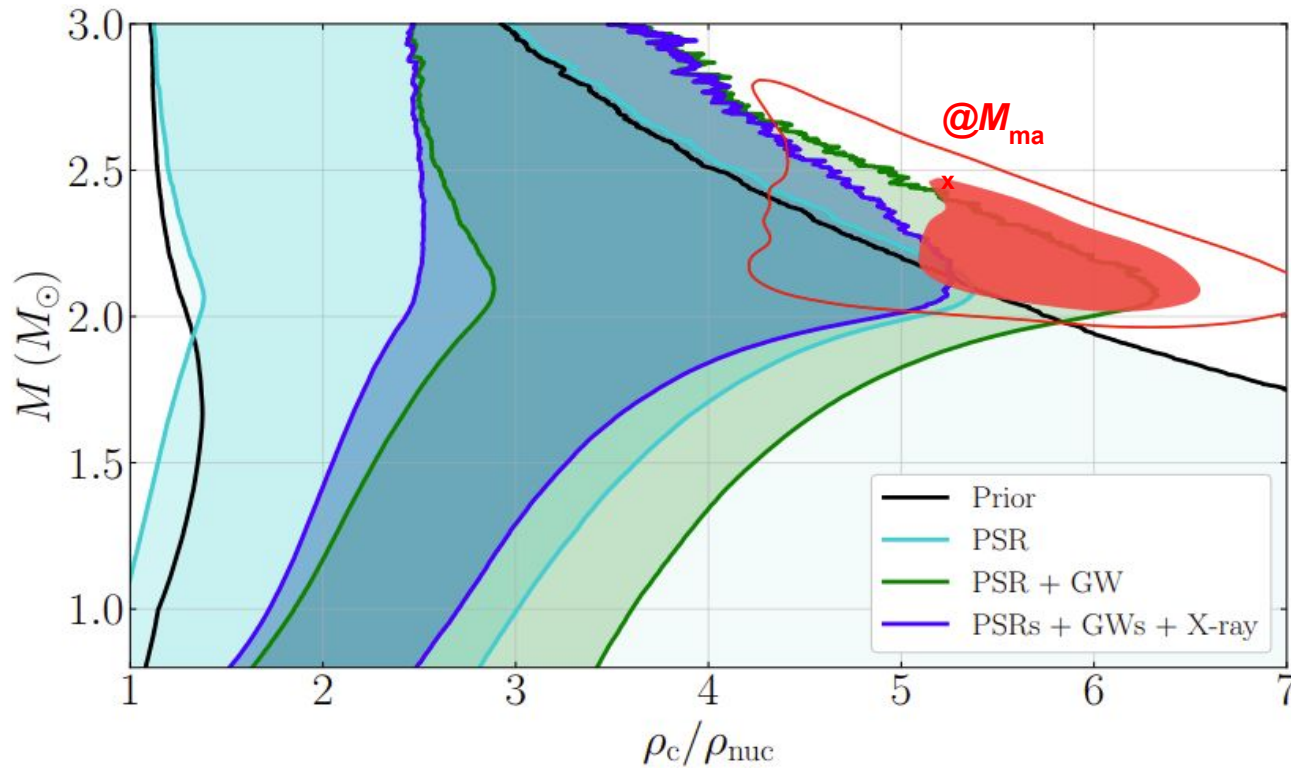
Priors



Posteriors with PSRs+GWs+NICER

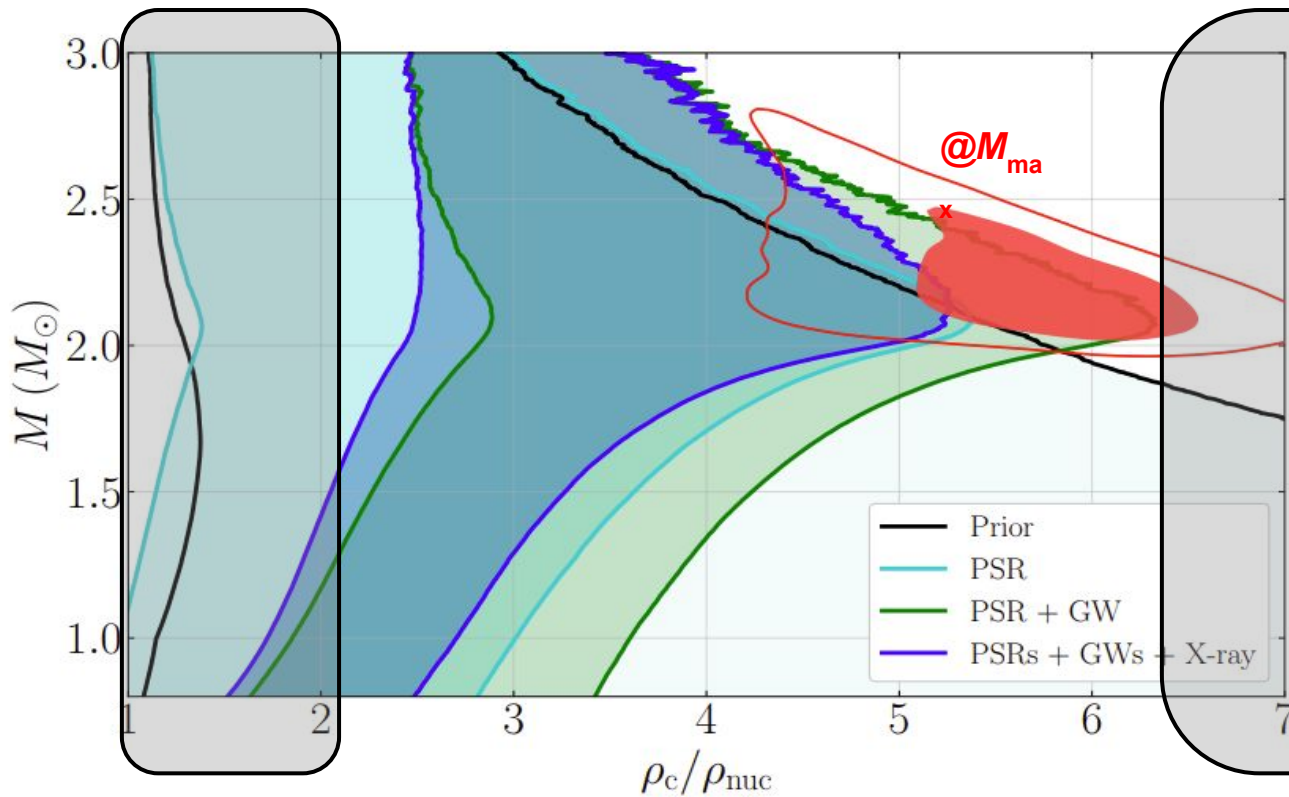


There is a limit to the densities we can probe within NSs





# Current Astrophysical Constraints on the High-Density, Cold Equation of State



**$\chi$ EFT**

Lynn+ (2016)

Drischler+ (2020), Drischler+ (2020)

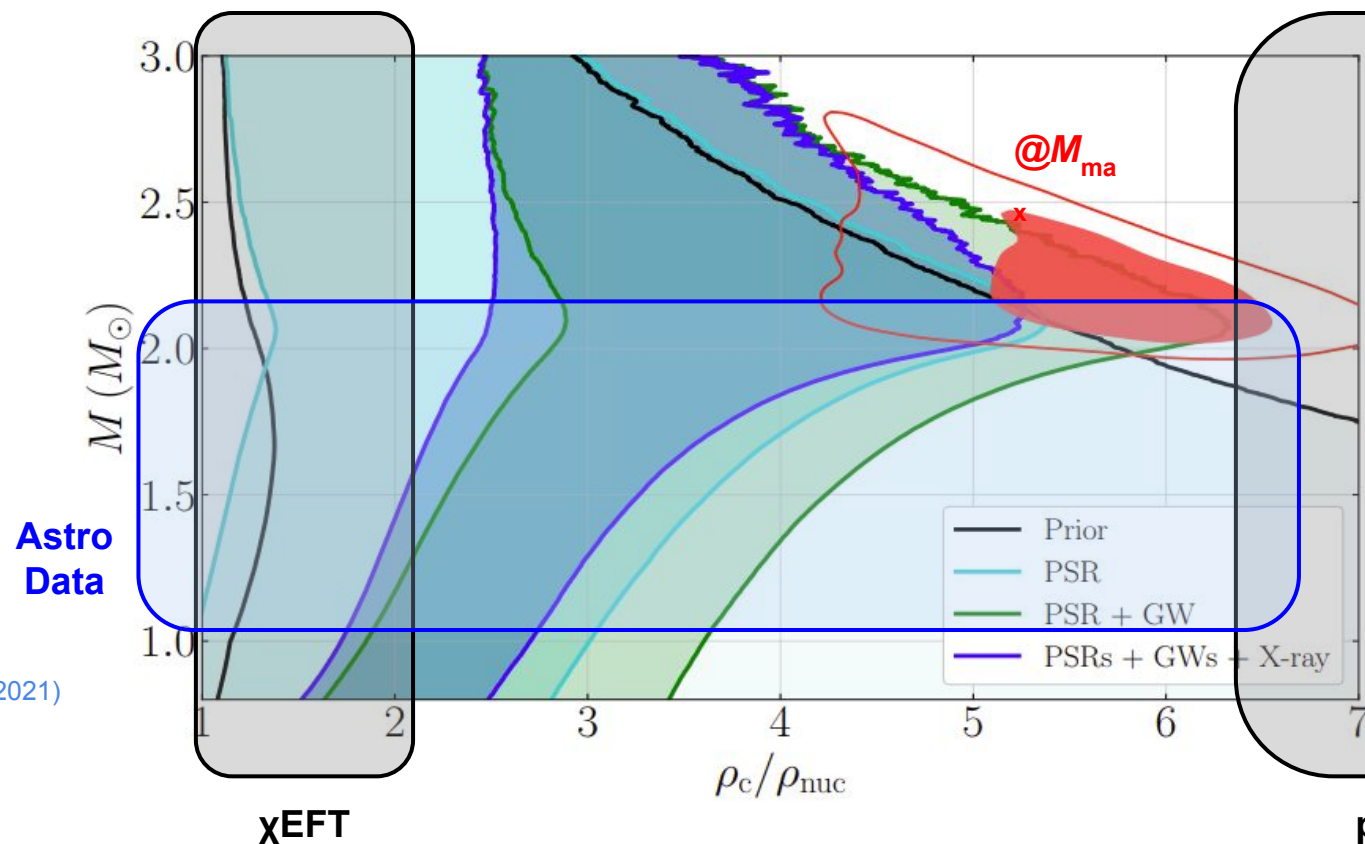
**pQCD**

Komoltsev+Kurkela (2022)

Gorda+ (2022)

Legred+ (2021)

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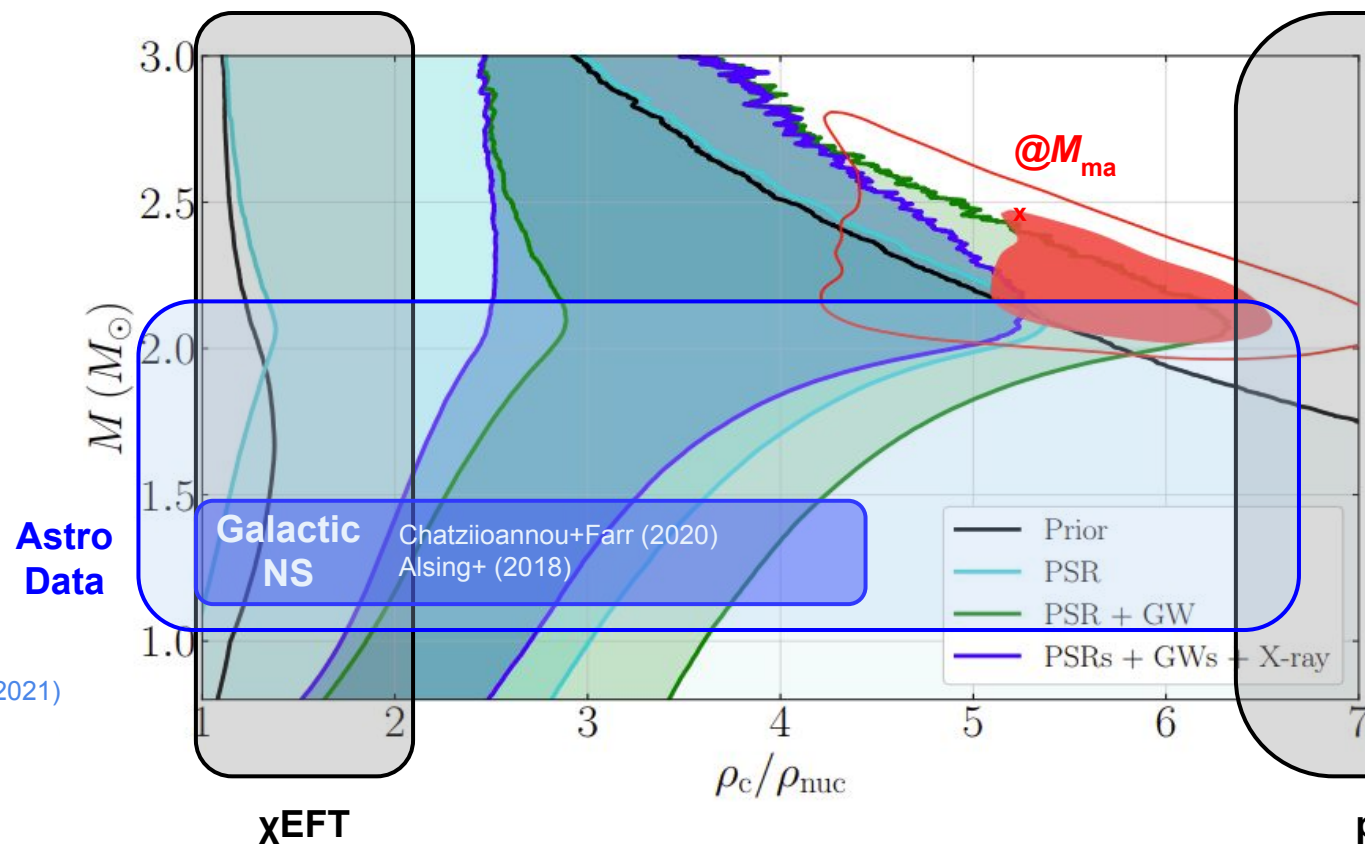
Abbott+ (2021)  
Farah+ (2022)  
Landry+Read (2021)

Lynn+ (2016)  
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# Current Astrophysical Constraints on the High-Density, Cold Equation of State



**Astro Data**

Abbott+ (2021)  
Farah+ (2022)  
Landry+Read (2021)

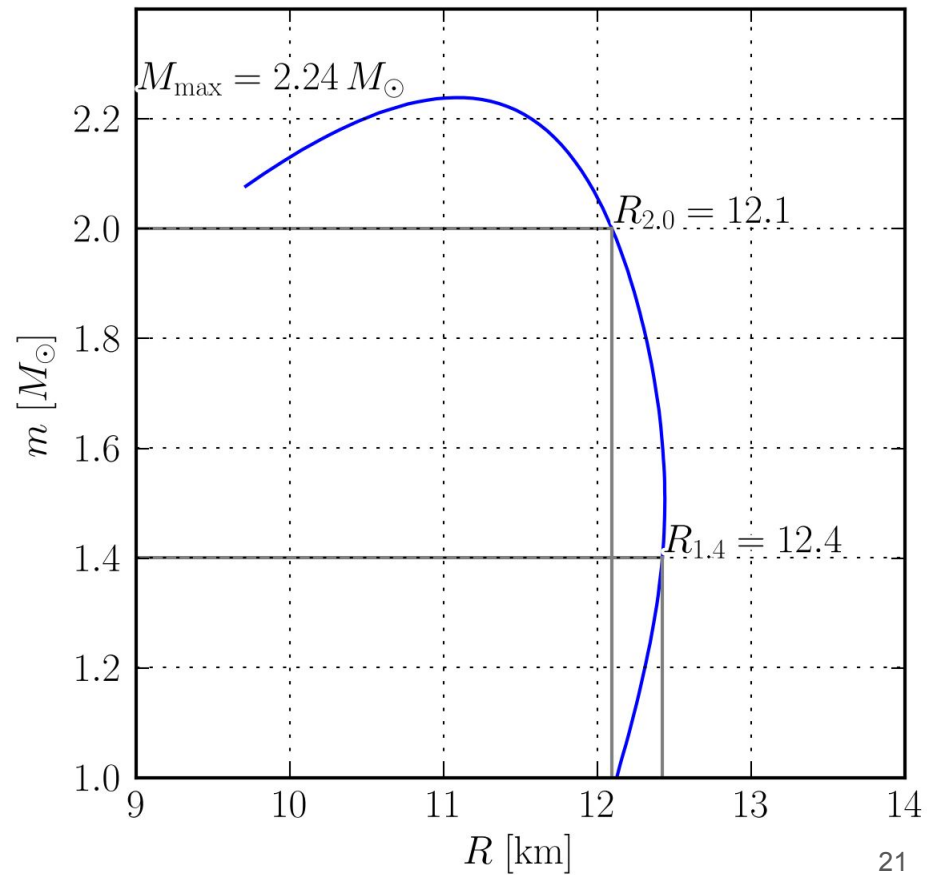
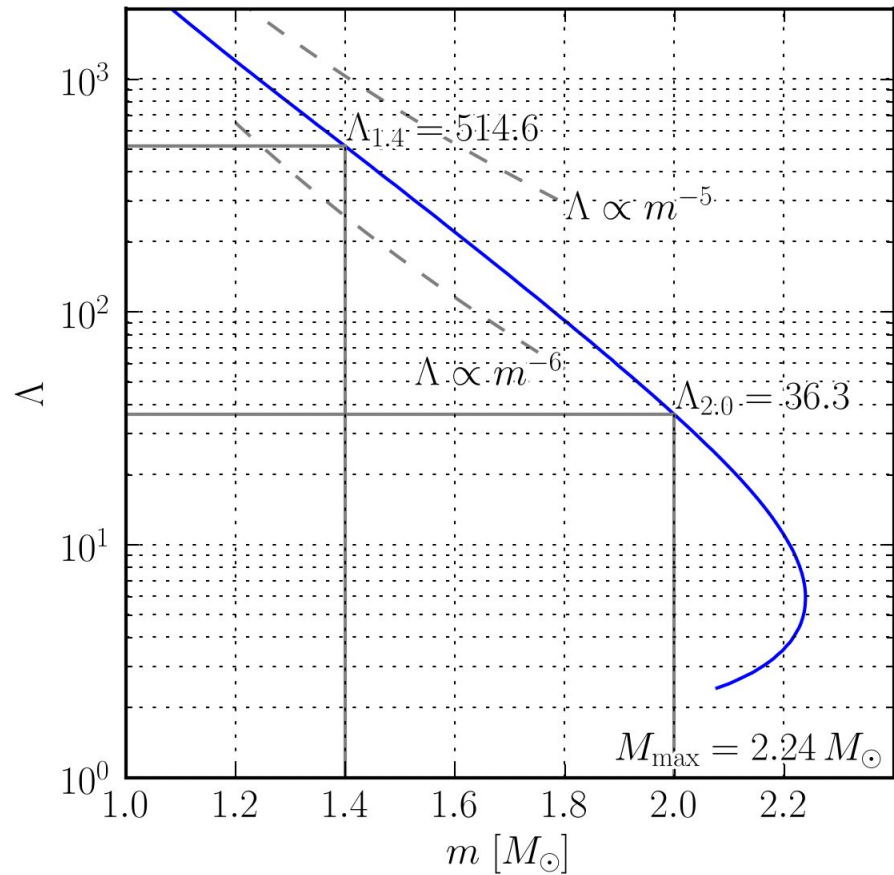
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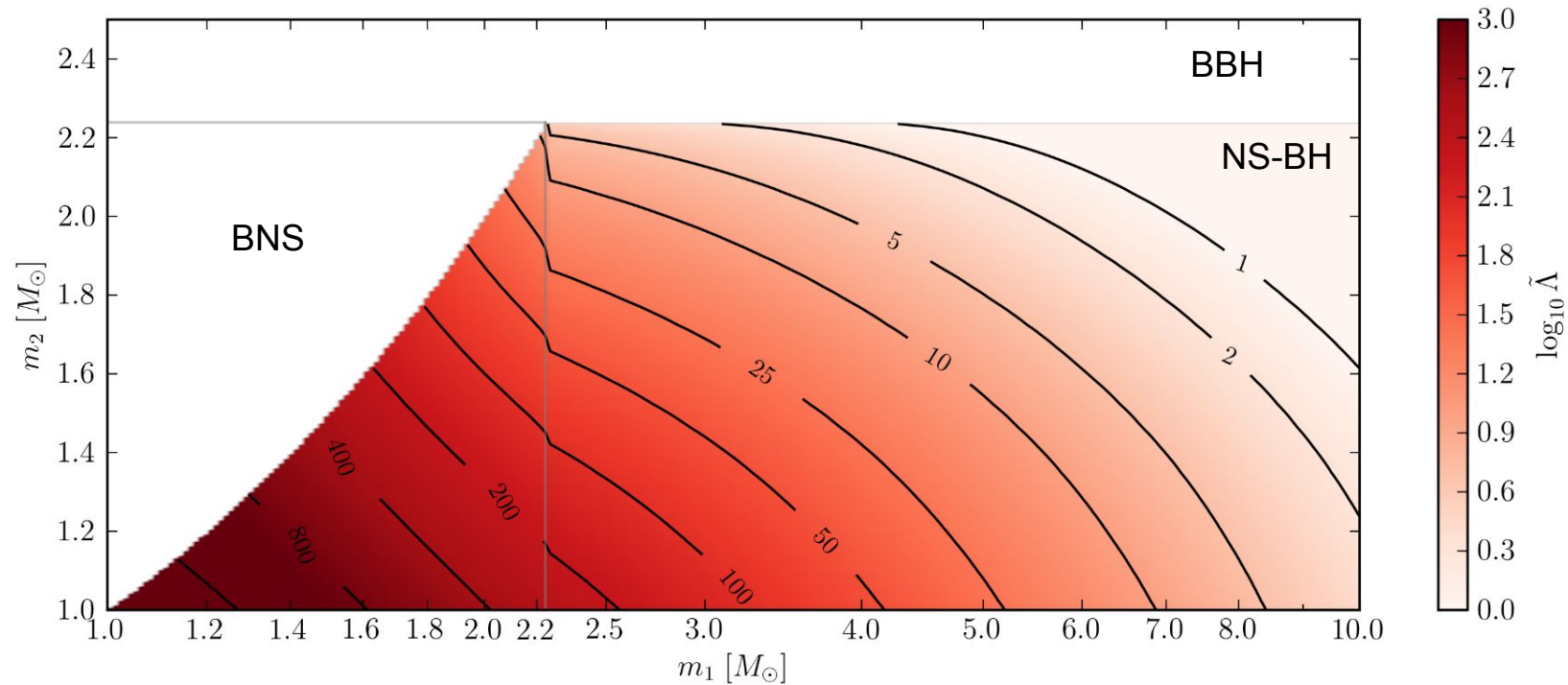
## Prospects for Measuring $\Lambda(m)$

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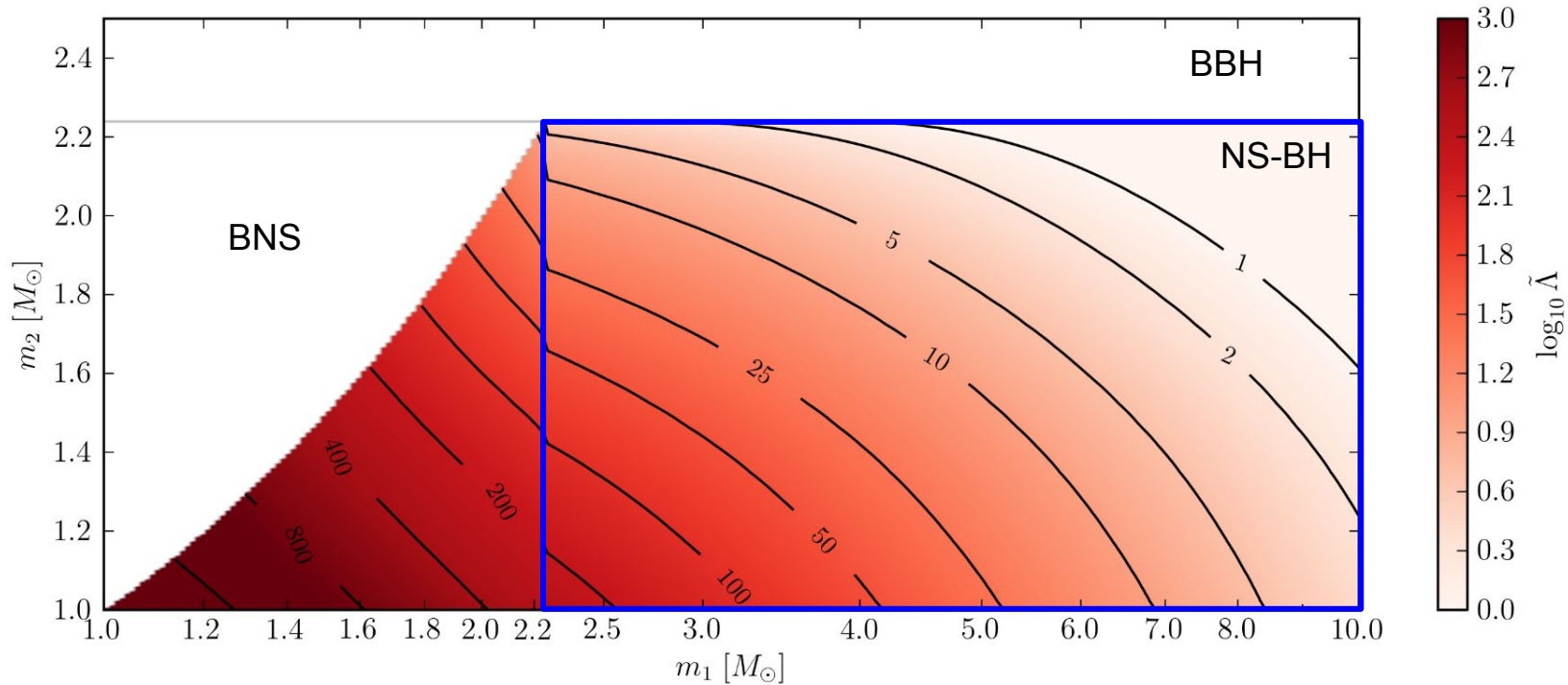
Leading-order adiabatic tidal term



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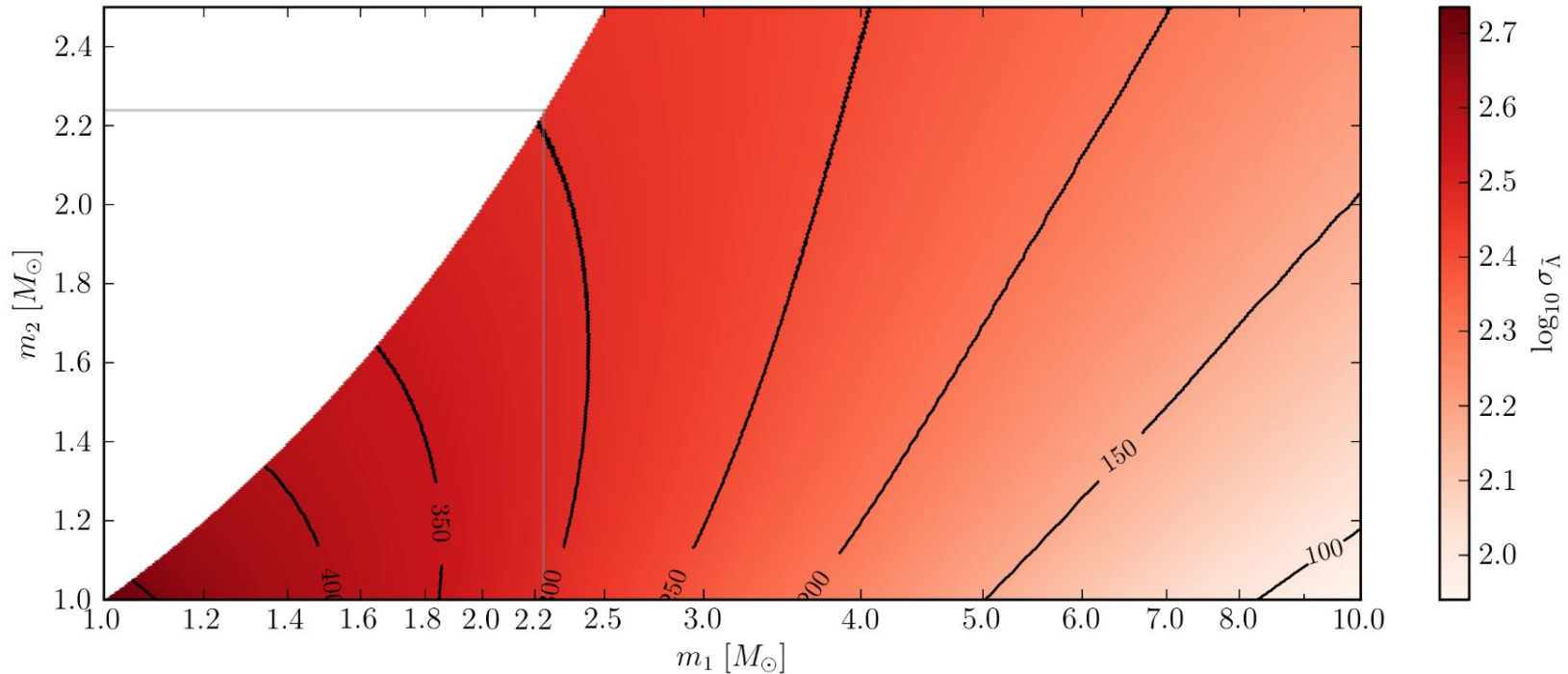
Leading-order adiabatic tidal term

$$\tilde{\Lambda} = \frac{16(q+12)q^4\Lambda(m_2)}{13(1+q)^5} \quad q = \frac{m_2}{m_1}$$



## Prospects for Measuring $\Lambda(m)$

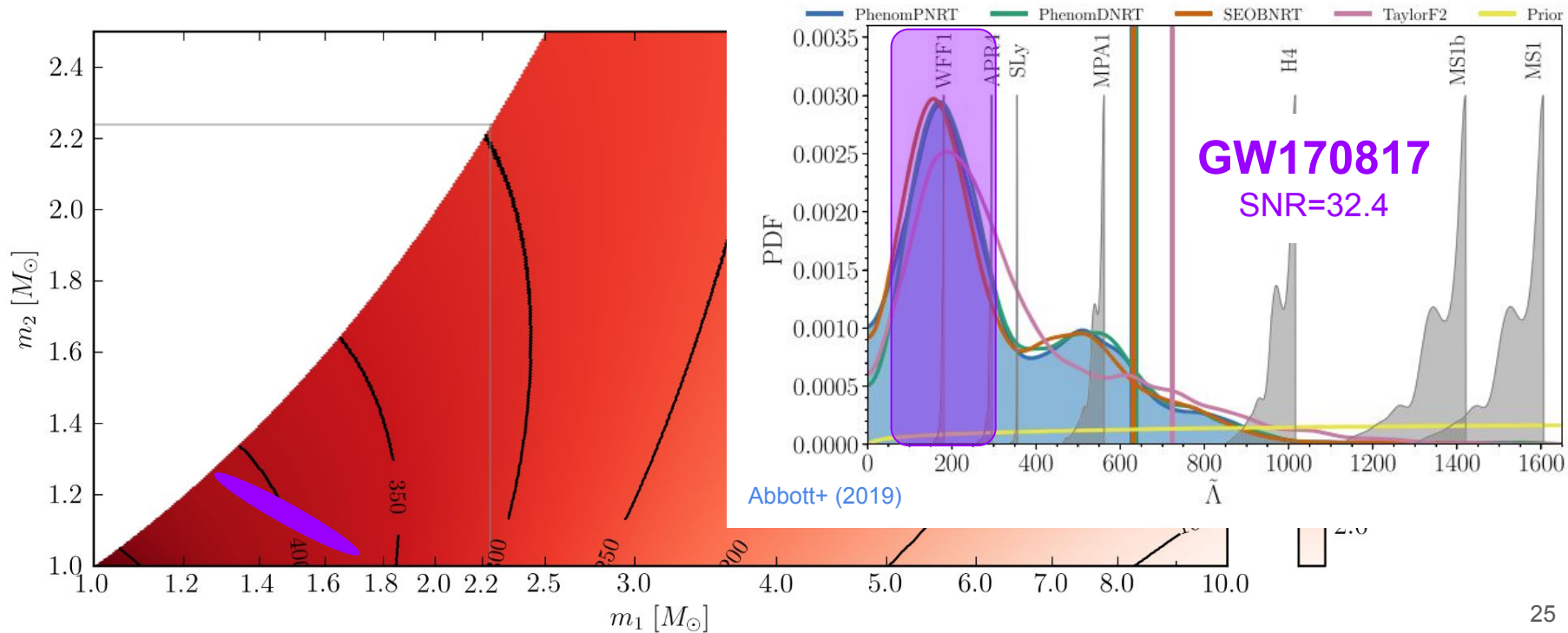
Fisher Matrix : simplistic PN phasing, aLIGO design sensitivity, **SNR=10**, wide priors on spins, mass ratio  
→ “best case” scenario (Cramer-Rao bound) that strictly holds only in the high-SNR limit





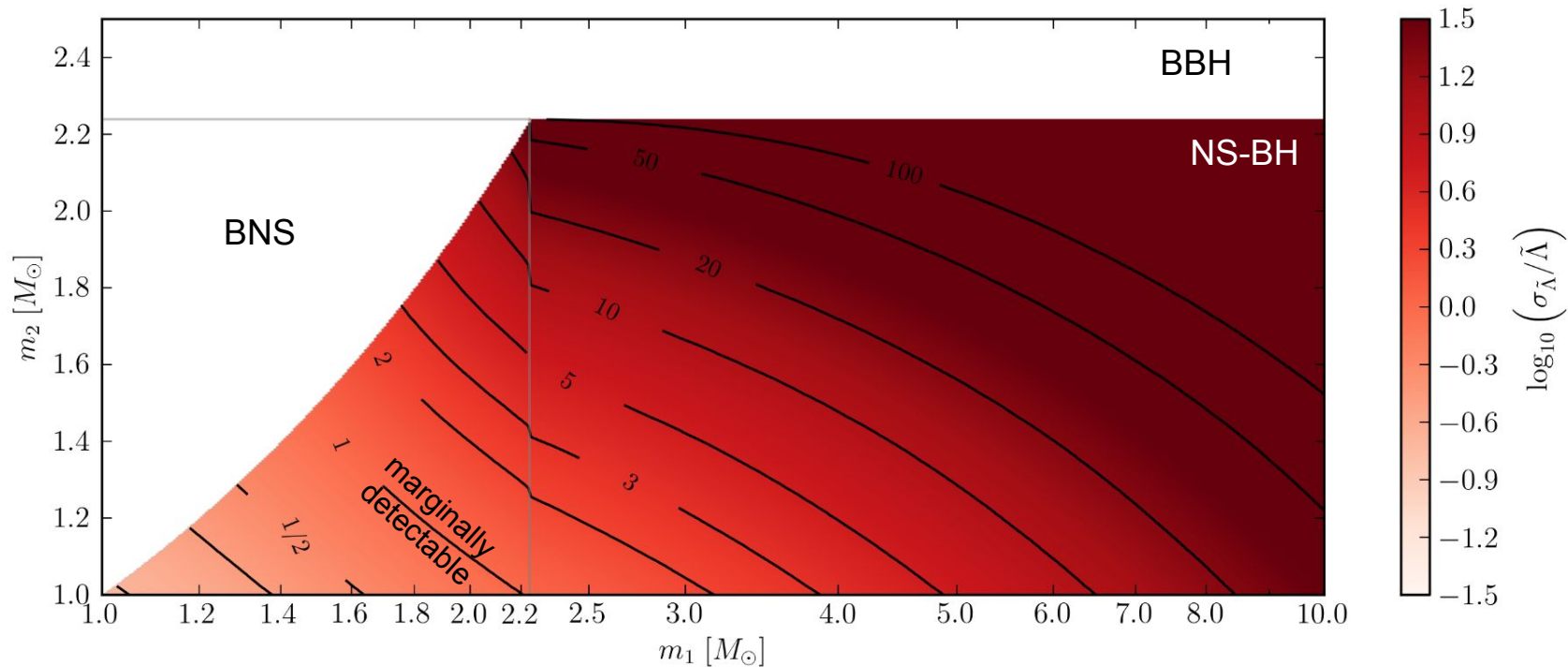
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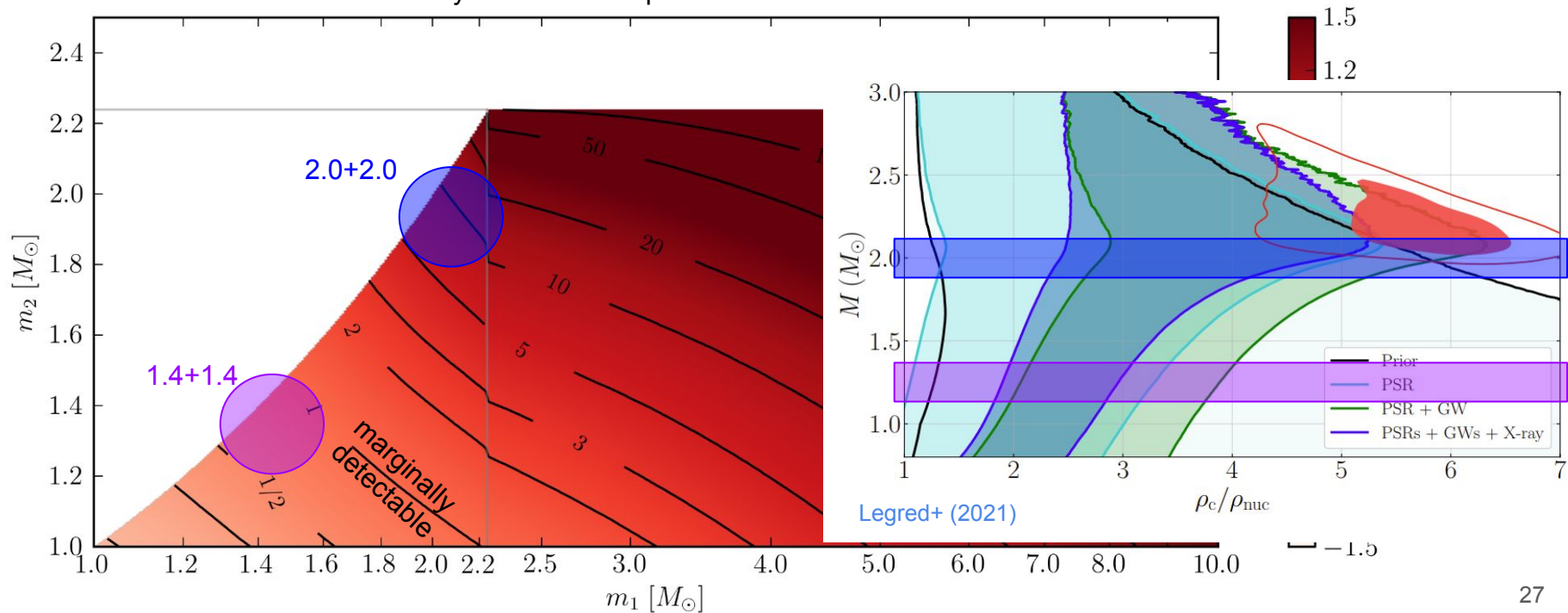
Uncertainty in leading order tidal term is still broad :  $\Lambda_{1.4} \sim 500 \pm 250$  at 90% credibility ([Legred+ 2021](#))  
Fisher Matrix may **underpredict uncertainty** by a factor of O(few)



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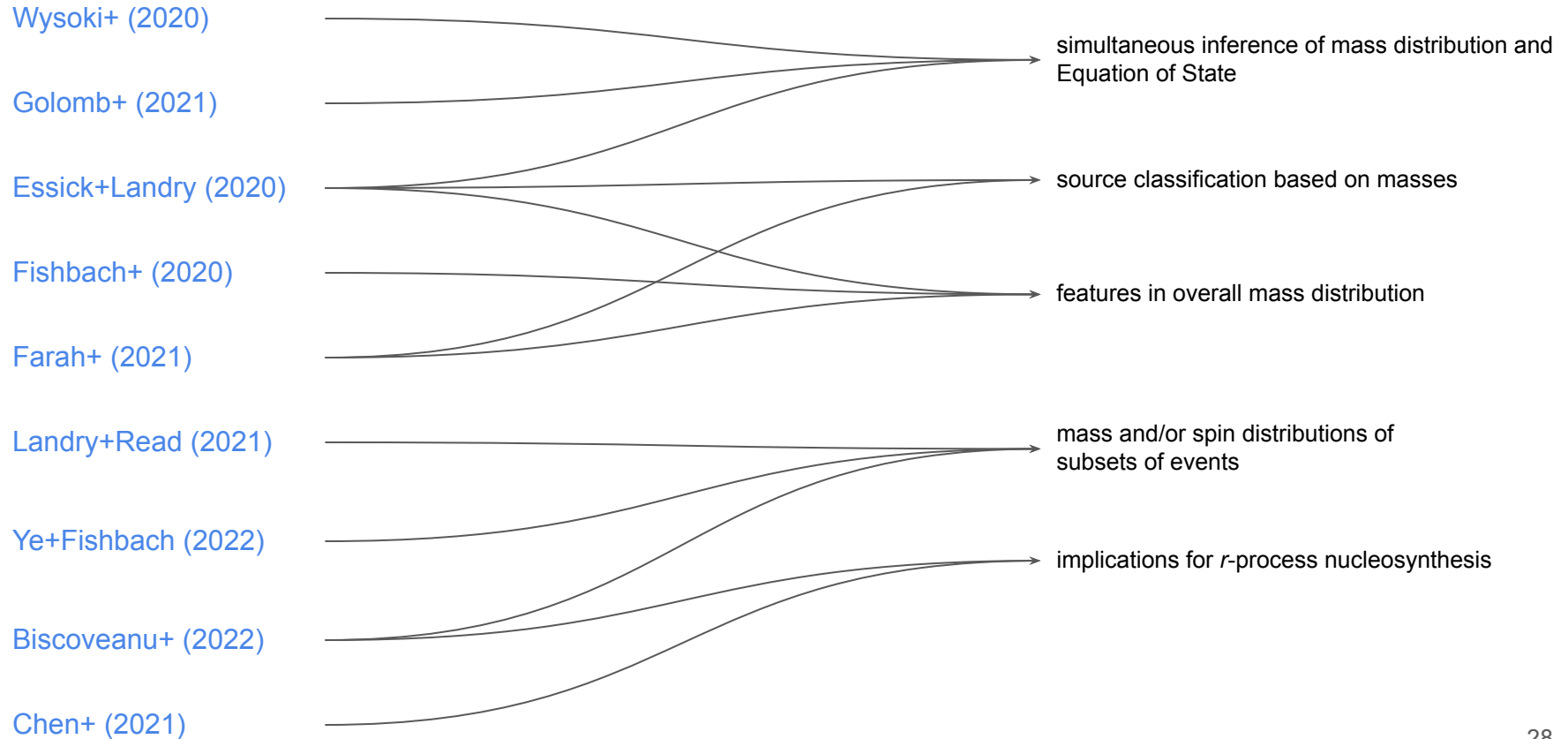
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relative precision could be 10x worse for 2.0+2.0 compared to 1.4+1.4  
→ need ~100x as many events to compensate



# Prospects for Measuring $\Lambda(m)$

## Connections to astrophysical mass distribution



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Wysocki+ (2020)

Golomb+ (2021)

Essick+Landry (2020)

Fishbach+ (2020)

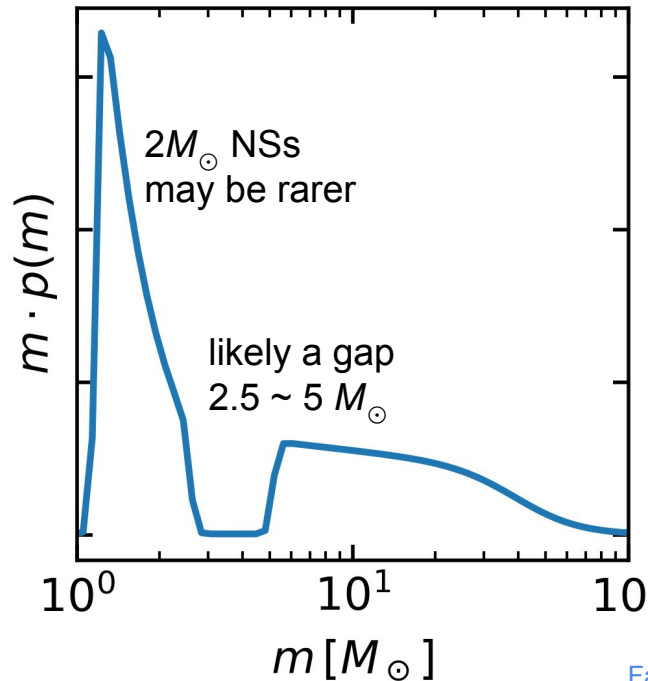
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Landry+Read (2021)

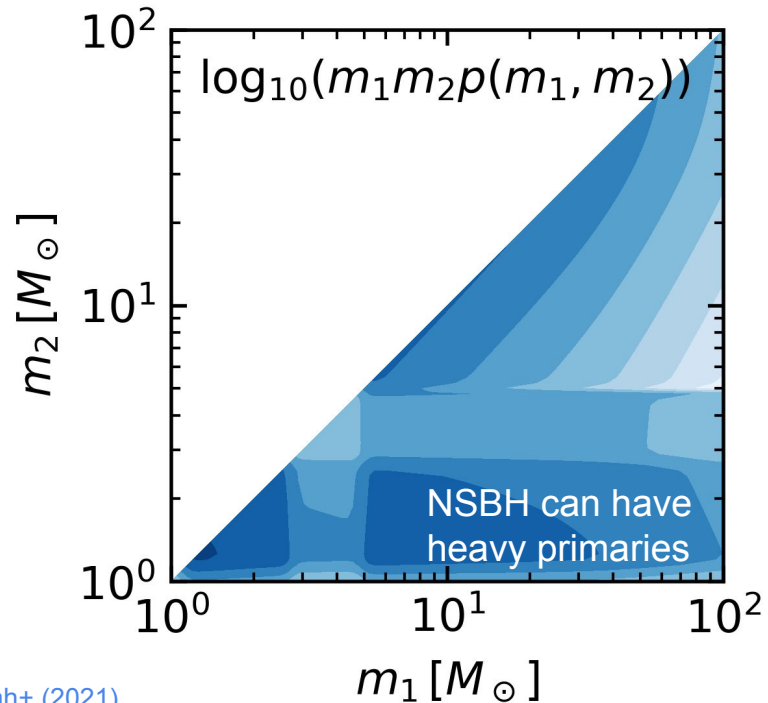
Ye+Fishbach (2022)

Biscoveanu+ (2022)

Chen+ (2021)



Farah+ (2021)



## Prospects for O4 starts March 2023

We will hold a one-hour Zoom meeting on **21 July 2022 at 14:00 UTC** (09:00 Central Time). Please register in advance for the Zoom meeting at

<https://bit.ly/3PbWa48>

The agenda is being developed at

<https://wiki.gw-astronomy.org/OpenLVEM/Telecon20220721>

Similar to previous calls, the first ~30 minutes will be devoted to **updates from LIGO-Virgo-KAGRA relating plans for O4** including the run schedule and some planned changes to the public alert infrastructure.

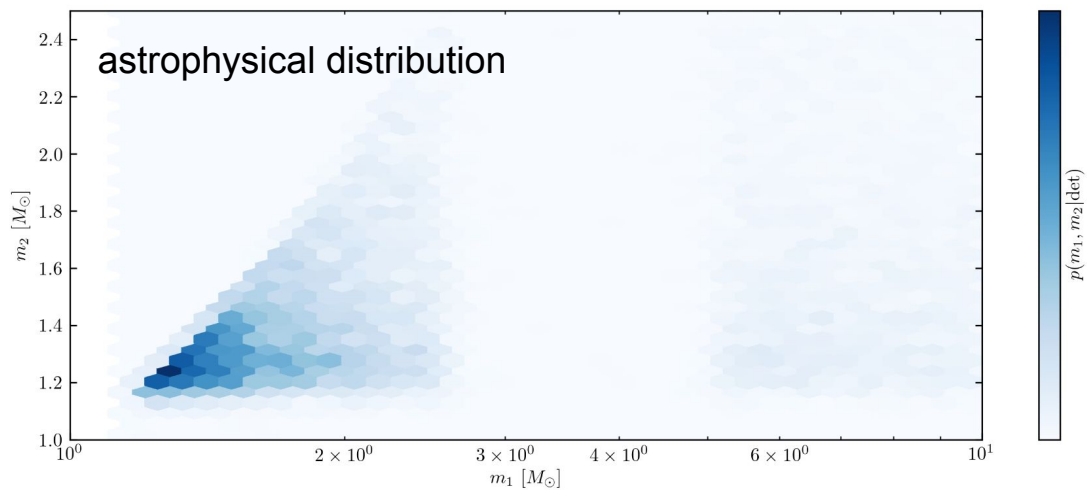
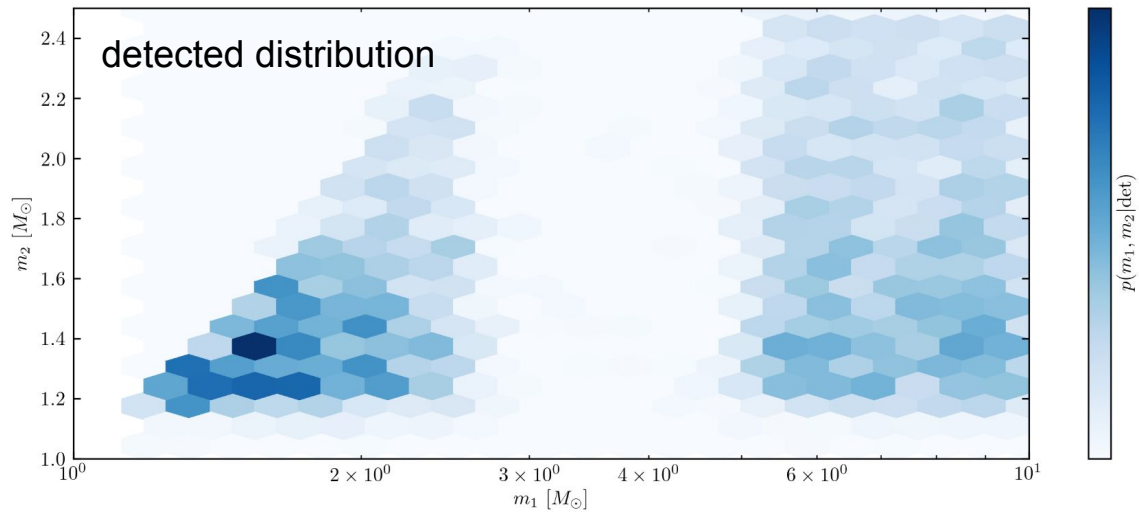
We encourage the OpenLVEM community to use these town hall meetings to share plans and progress with us and each other. In this first call of the O4 series, we have assigned 25 minutes for short presentations from the community. If you would like to make a short presentation at this, or a future call, please submit a request to

<https://forms.gle/tieqoa2xsnm7Sqpr7>

We look forward to your participation; feel free to forward this announcement to others who may be interested.

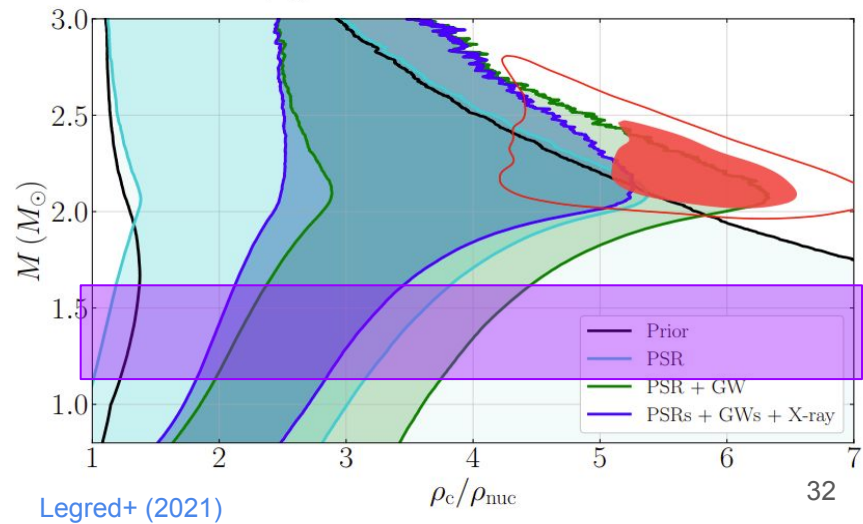
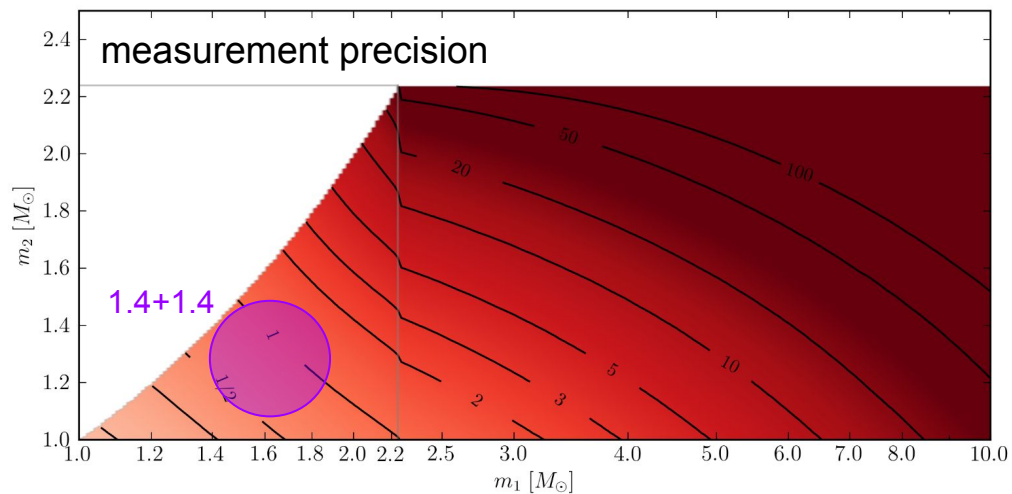
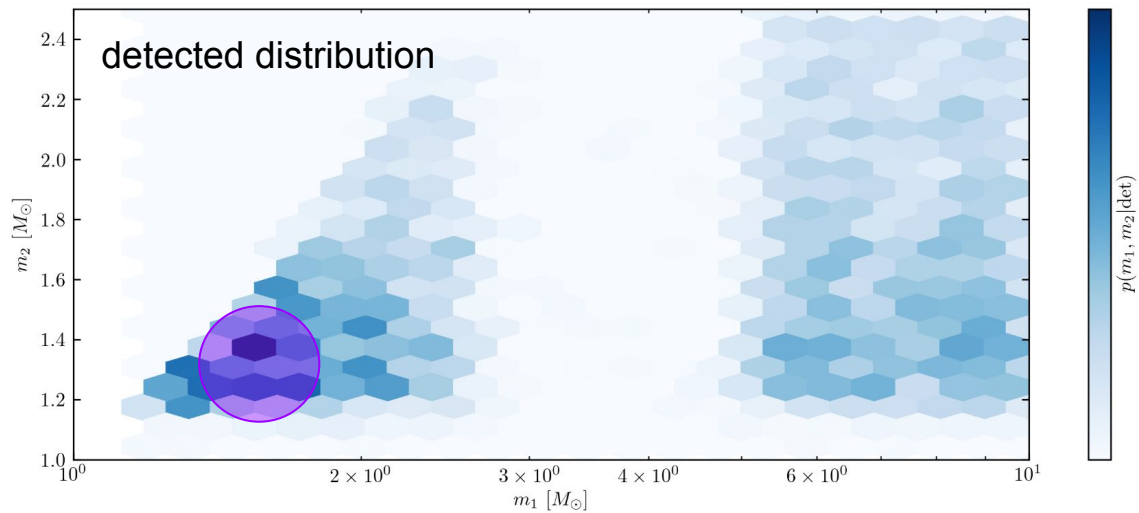
## Prospects for O4

100% duty cycle at design sensitivity  
uniformly distributed in  
co-moving-volume+source-frame time



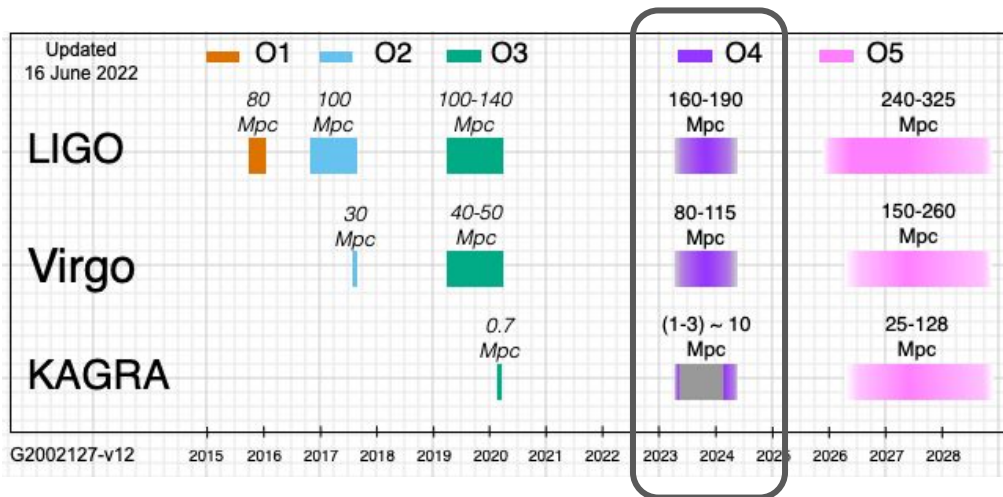
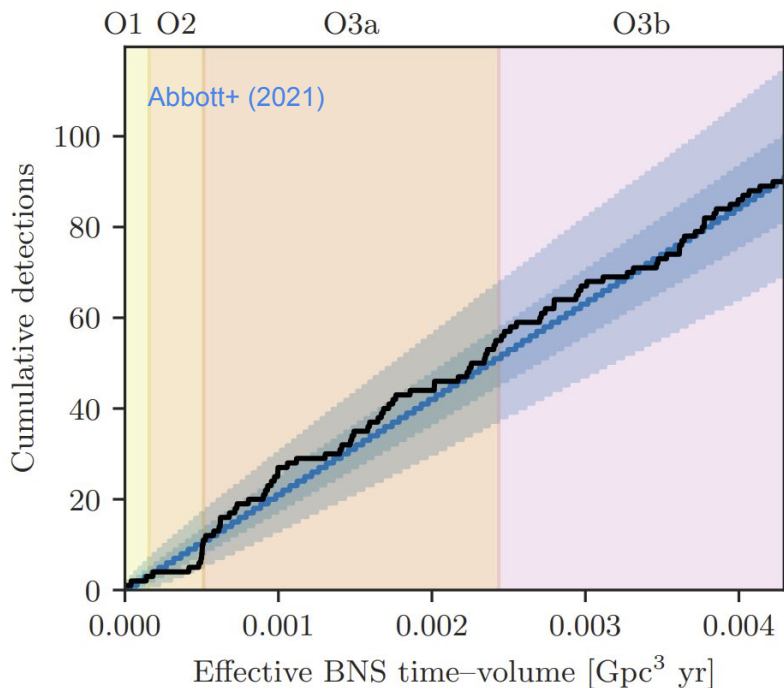
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# Prospects for O4



$$\langle VT \rangle_{O4} \sim (0.012 - 0.020) \text{ Gpc}^3 \text{ yr}$$

Expect 3 - 5 times more VT from O4  
 → 6 - 10 expected BNS detections in O4

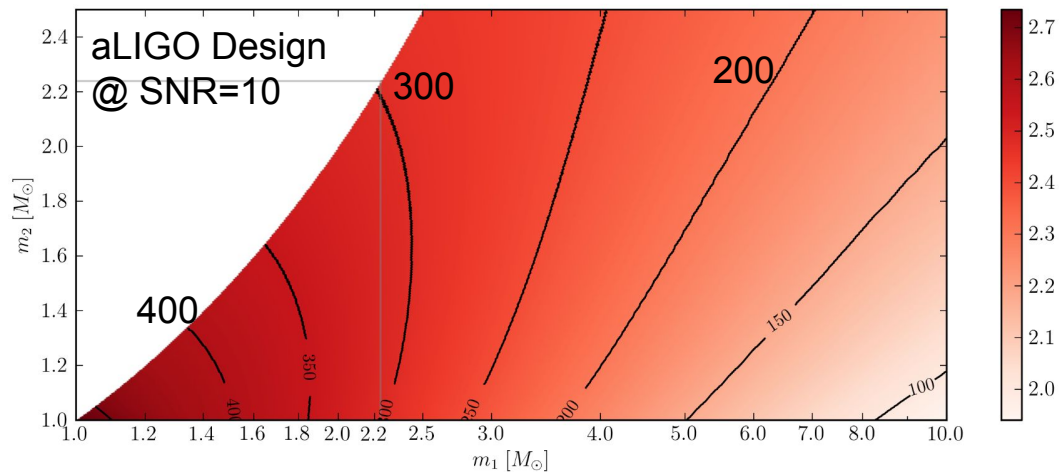
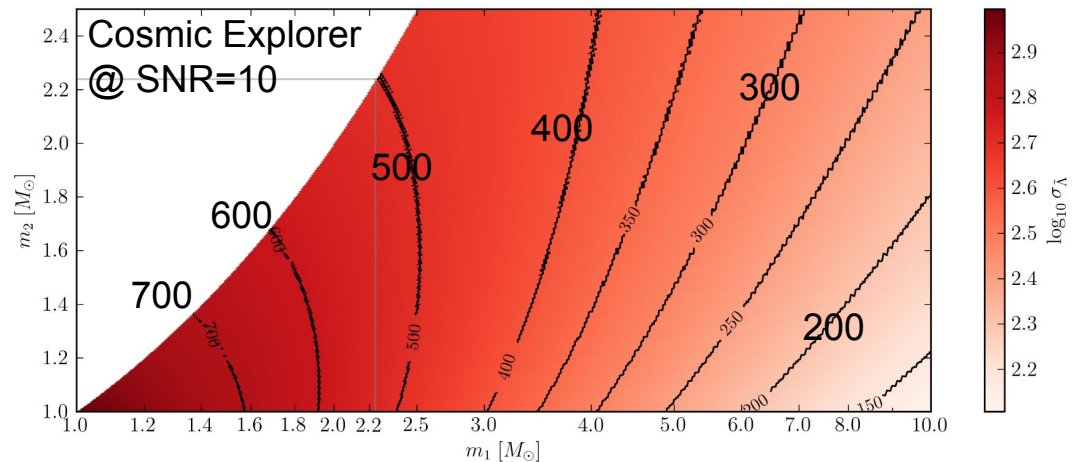
Estimate of rates span 10 - 1700 BNS/Gpc<sup>3</sup>yr (Abbott+ 2021)  
 → 0.12 - 32 expected BNS detections in O4

What about 3G Detectors?

At a **fixed SNR**, **measurement of tides is worse**

low-freq sensitivity increases more than high-freq sensitivity for “nominal” CE (e.g., [Essick 2022](#))

detectors may be tuned to target tidal effects  
[Srivastava+ \(2022\)](#)

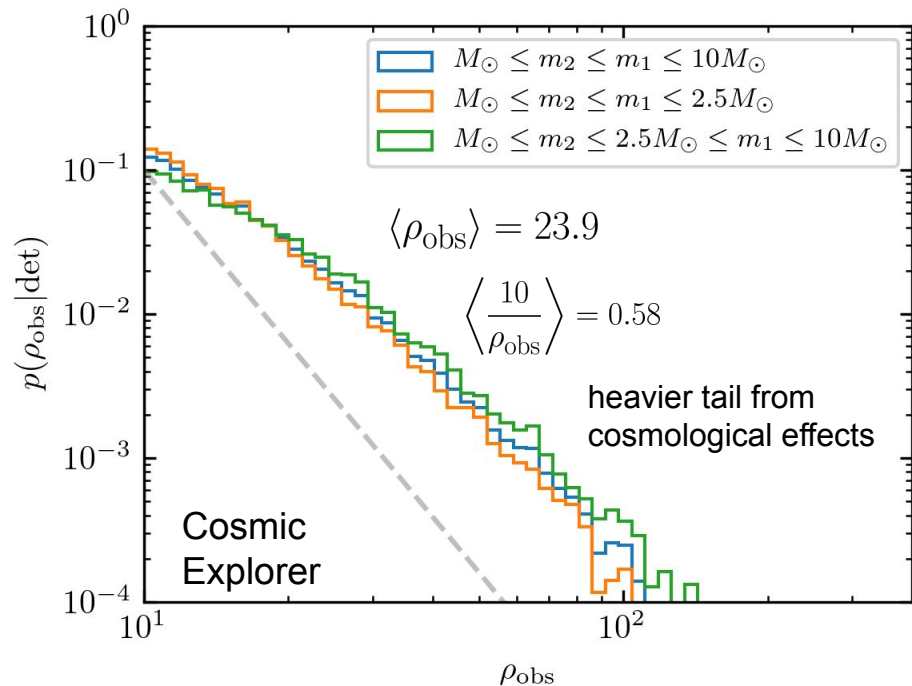
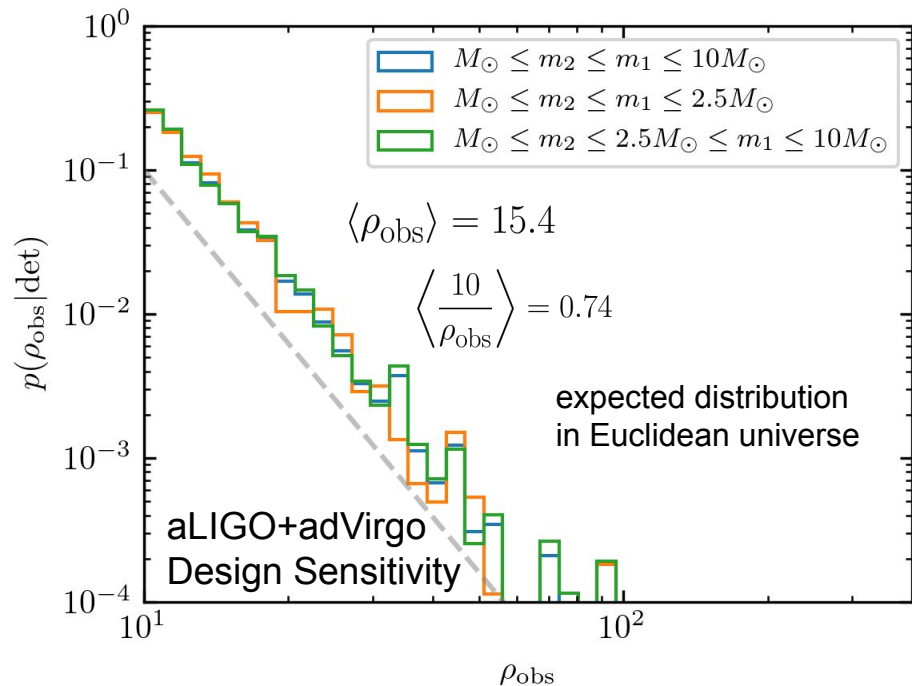


Each individual source will have a higher SNR in 3G than in aLIGO.

→ will the proportion of high-SNR signals be larger in 3G detectors?

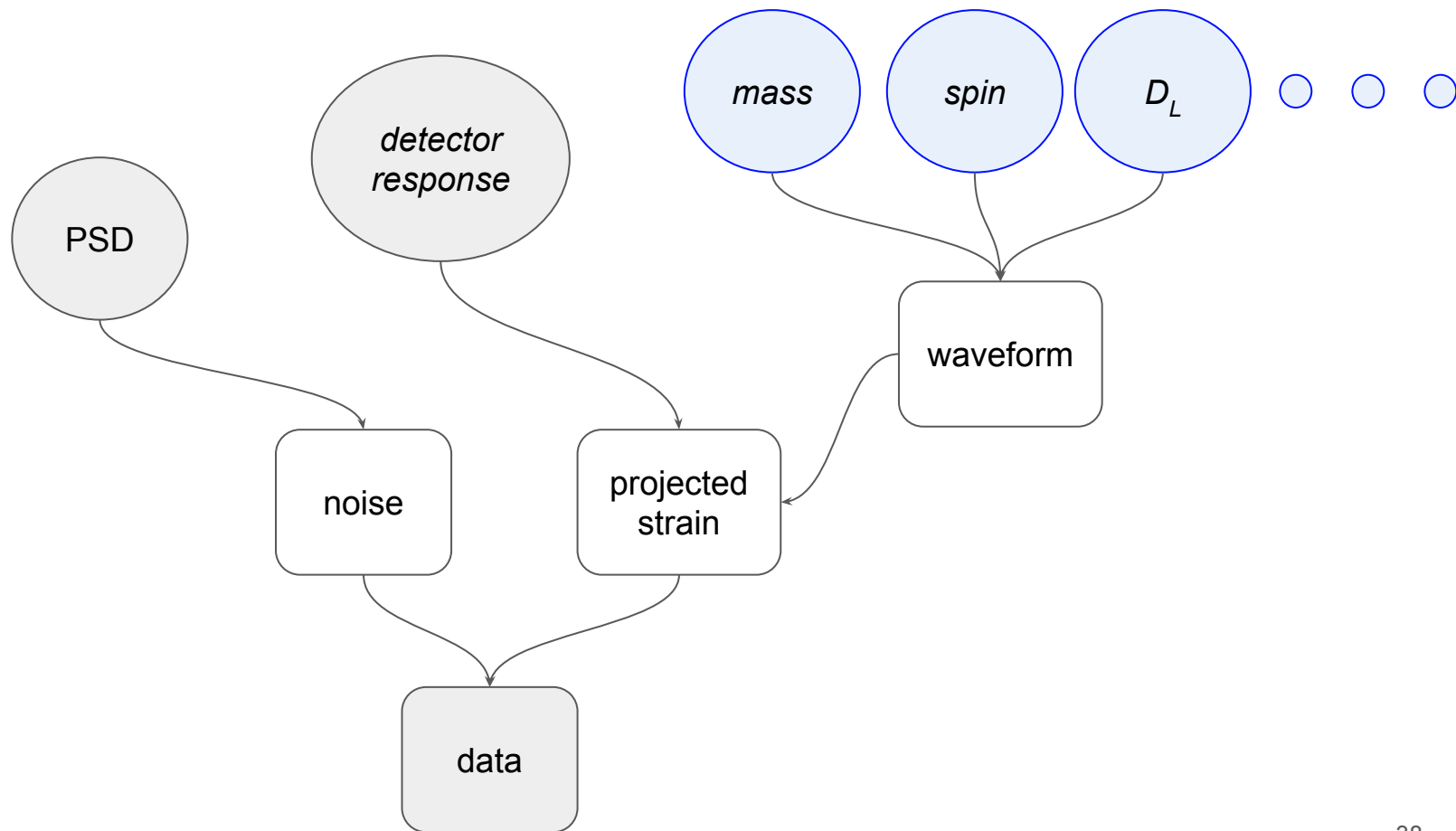
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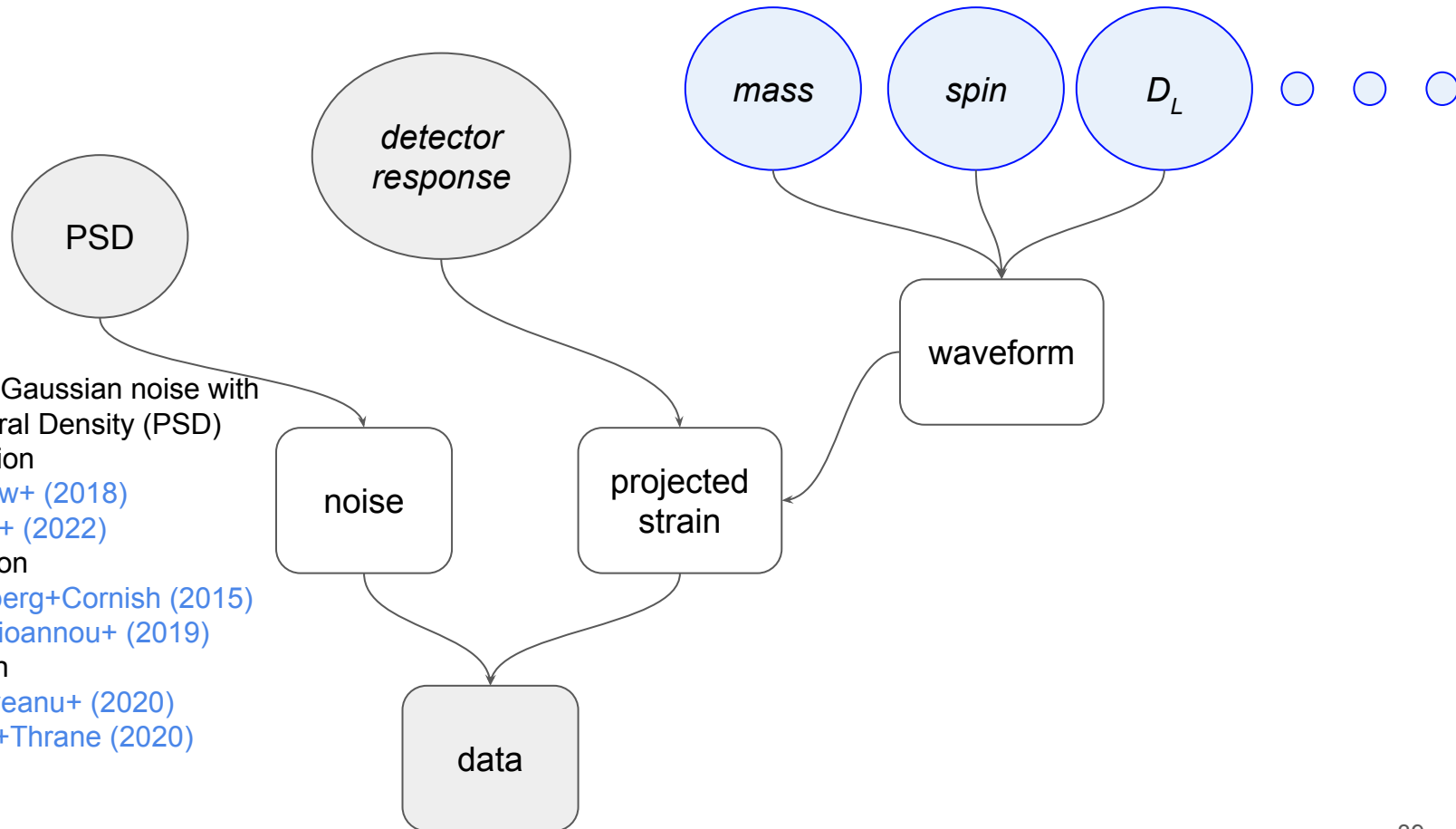
Most low-mass events in 3G will still be near the detection threshold (compare to [Vitale 2016](#))

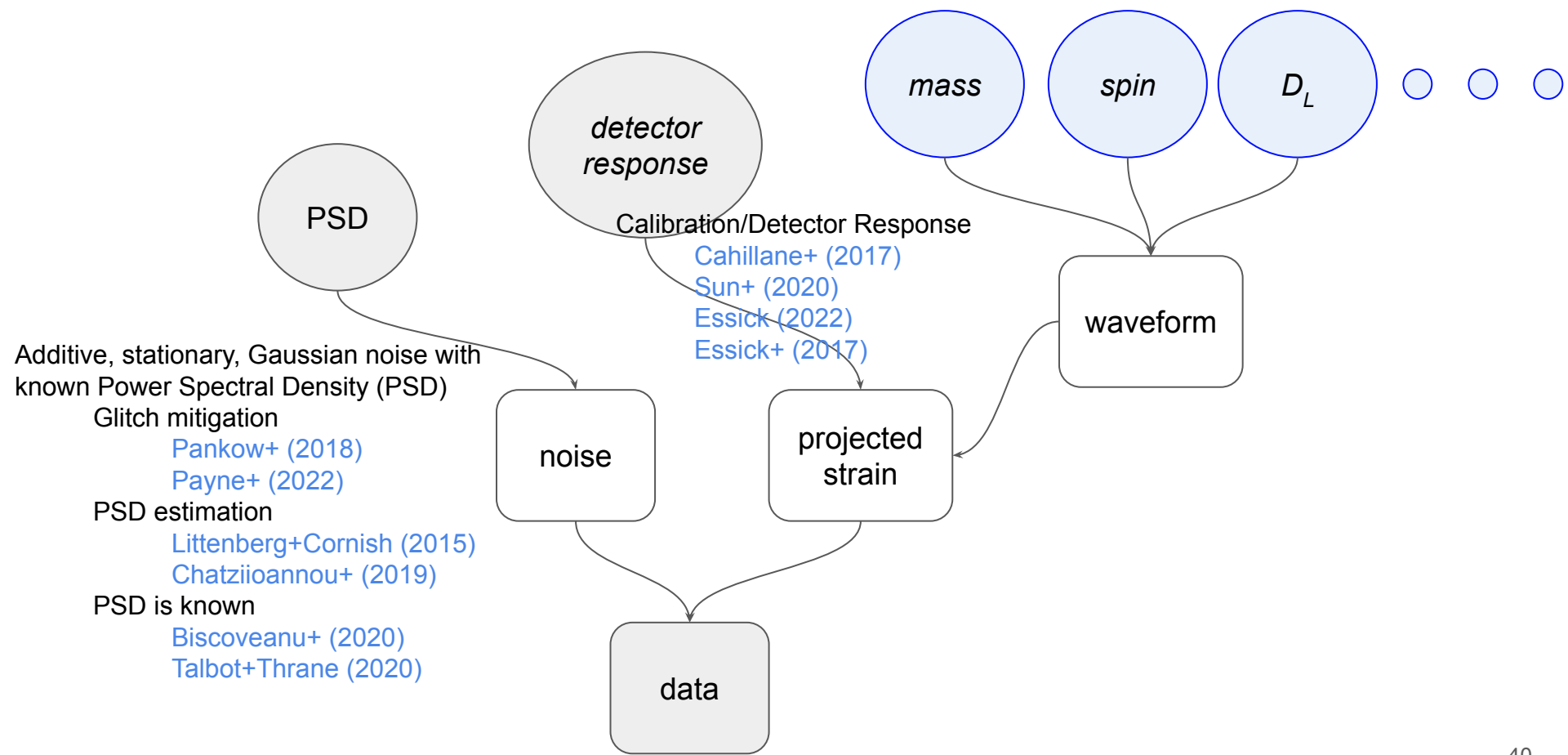


For the average event, increased SNR with CE will **not** overcome the decreased precision in adiabatic tidal measurements

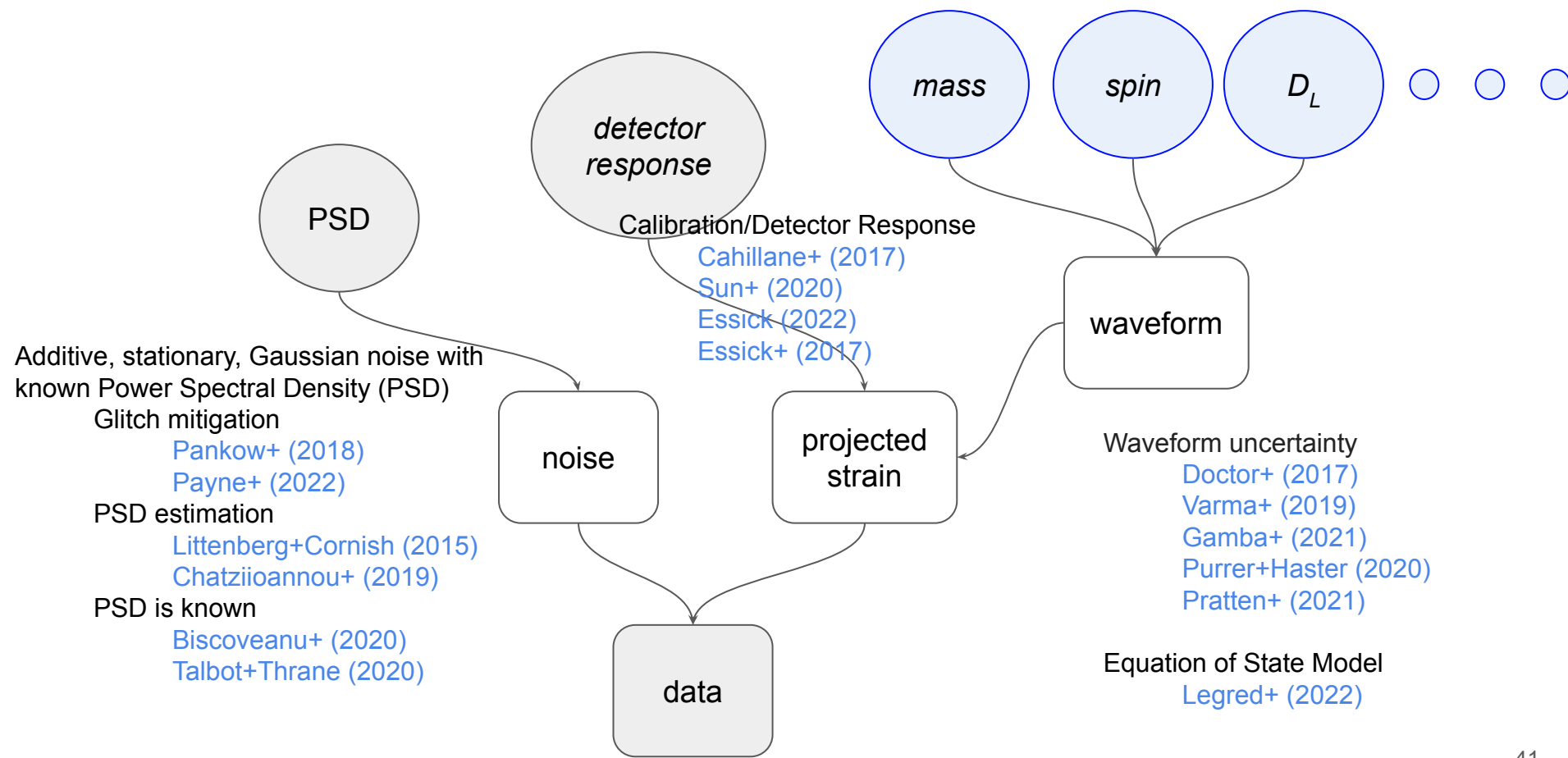
## Pain Points

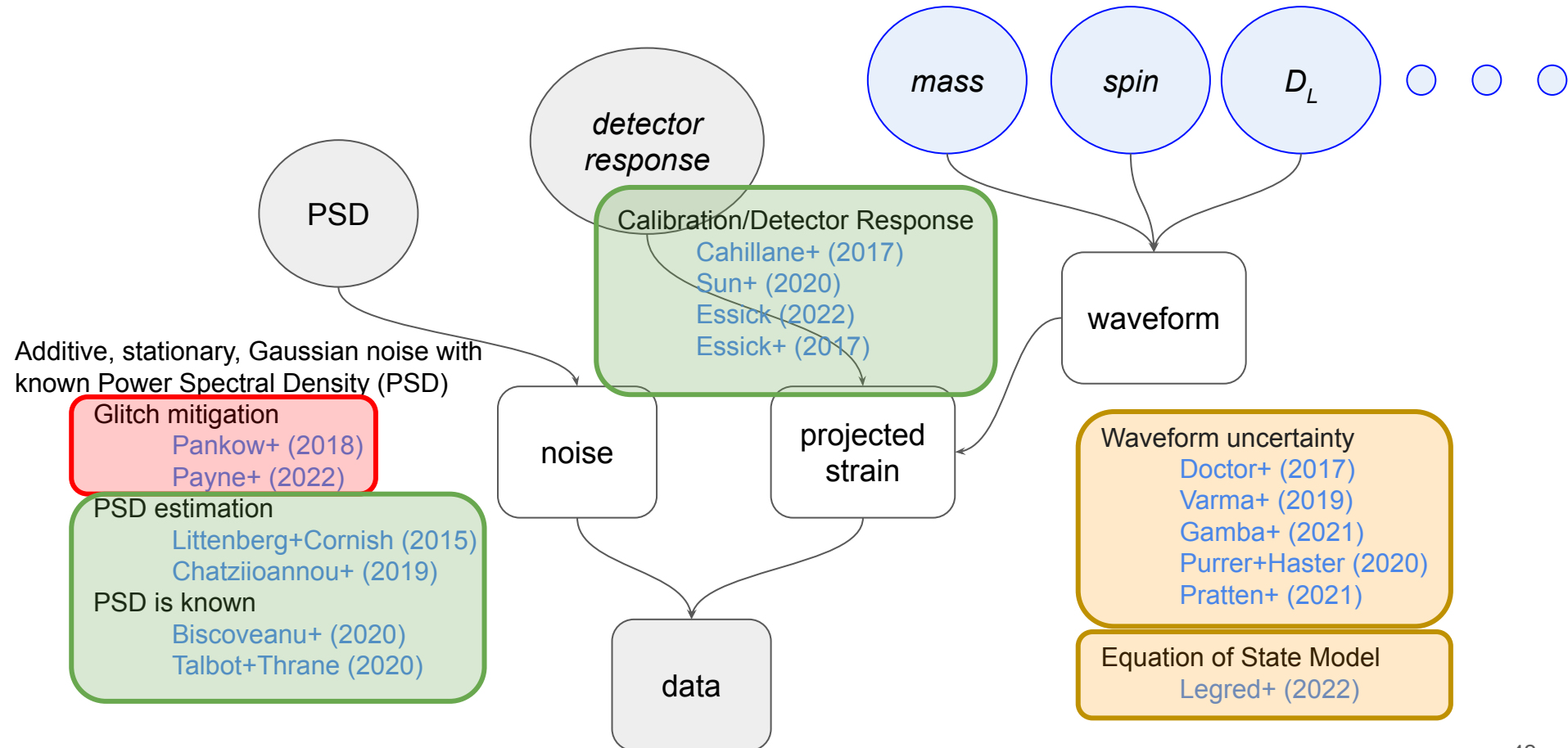




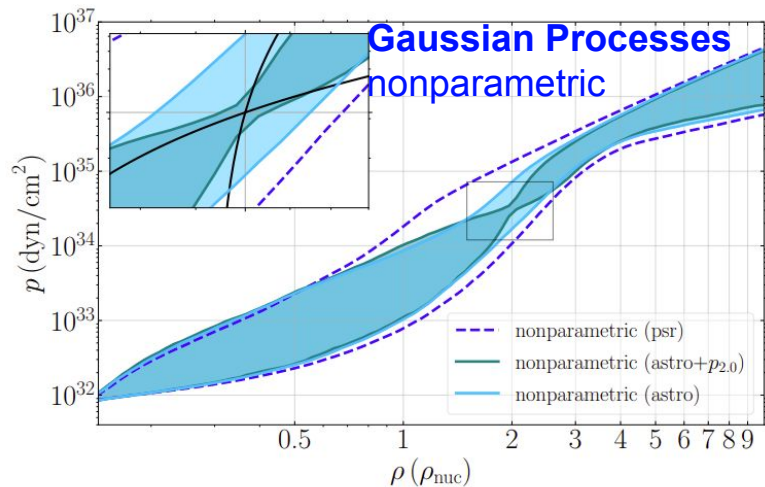






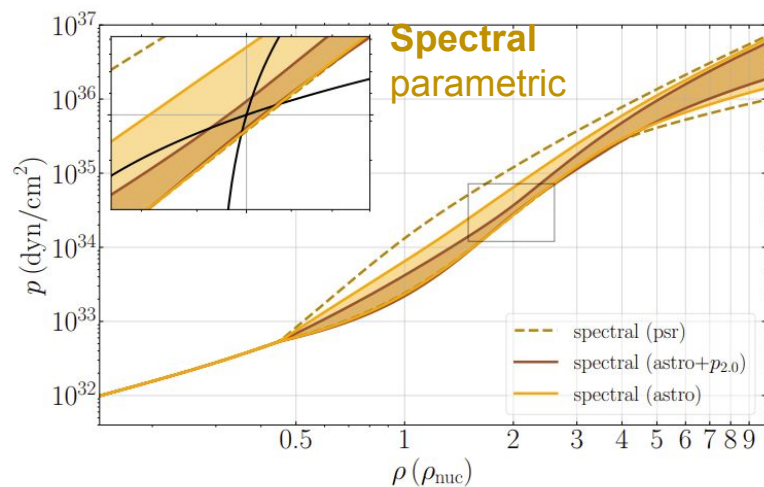


# Pain Points : Phenomenological EoS Models and Implicit Assumptions

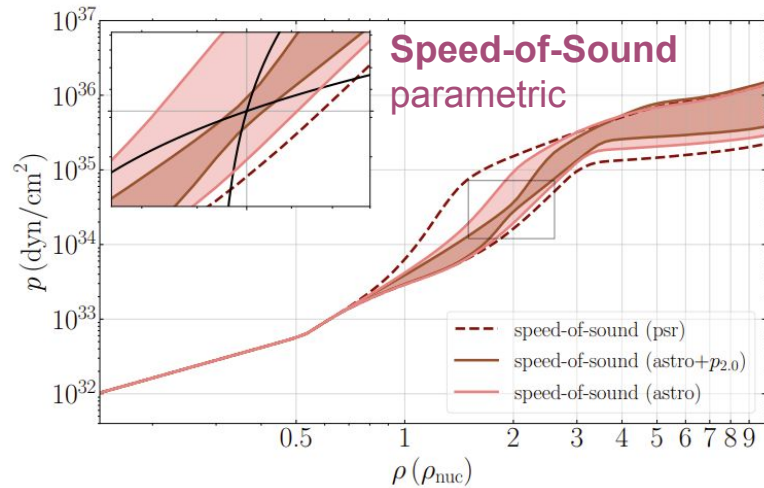
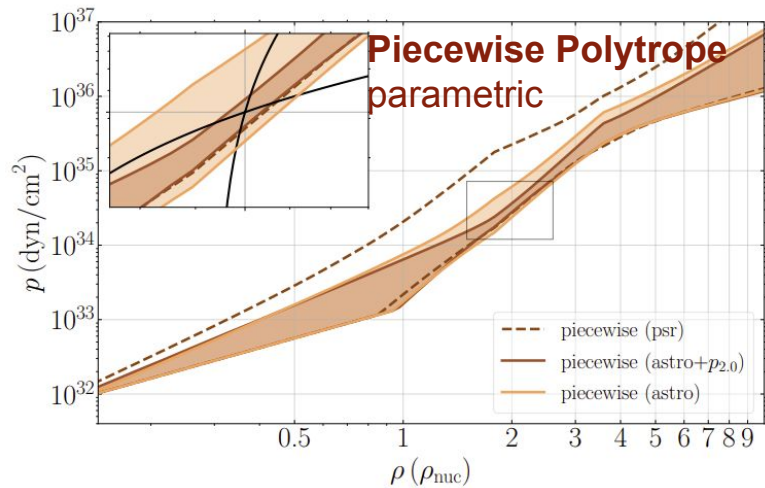


priors – dashed

astro data – light



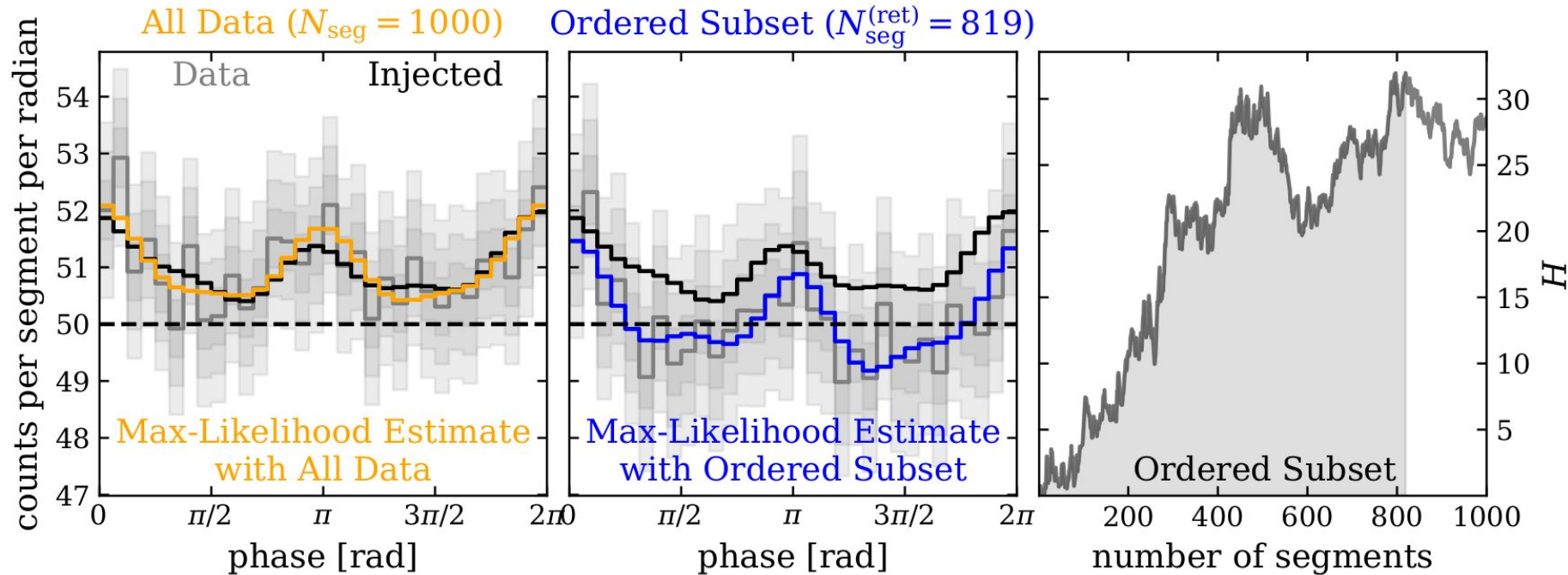
astro +  $p(2n_{\text{sat}})$  – dark



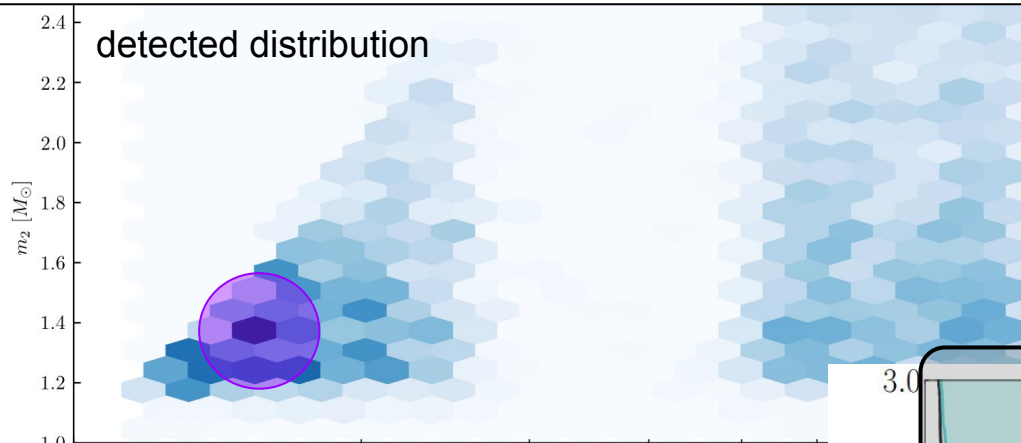
## Pain Points : (X-ray) Data Selection

These analyses have a lot of moving parts, and seemingly small choices can have unexpected consequences.

→ NICER's *data selection procedure* (for J0740) introduces small biases

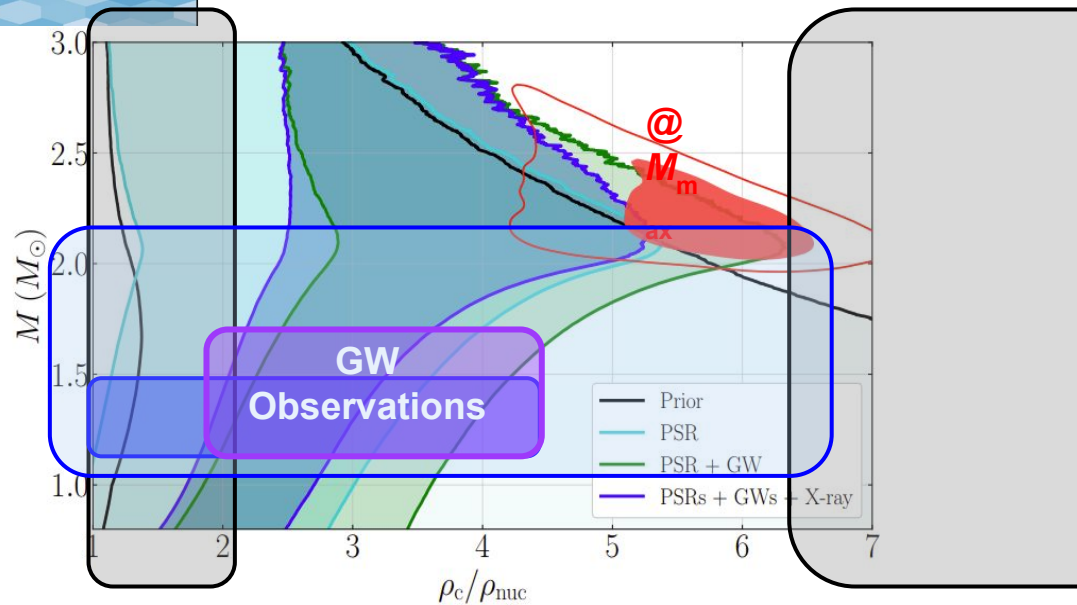
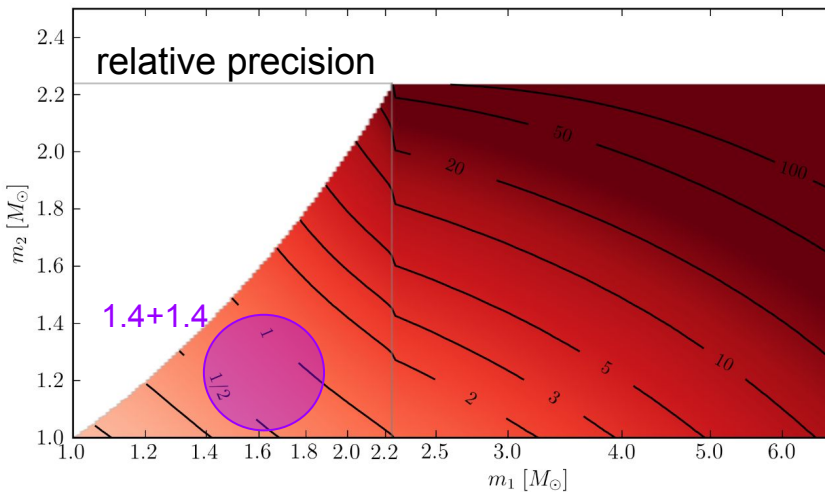


# Summary



Expect 3 - 5 times more VT from O4  
→ 6 - 10 expected BNS detections in O4

Estimate of rates span 10 - 1700 BNS/Gpc<sup>3</sup>yr (Abbott+ 2021)  
→ 0.12 - 32 expected BNS detections in O4



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