

Direct Detection of Sub-GeV Dark Matter with Doped Semiconductors

Peizhi Du

C.N.Yang Institute for Theoretical Physics



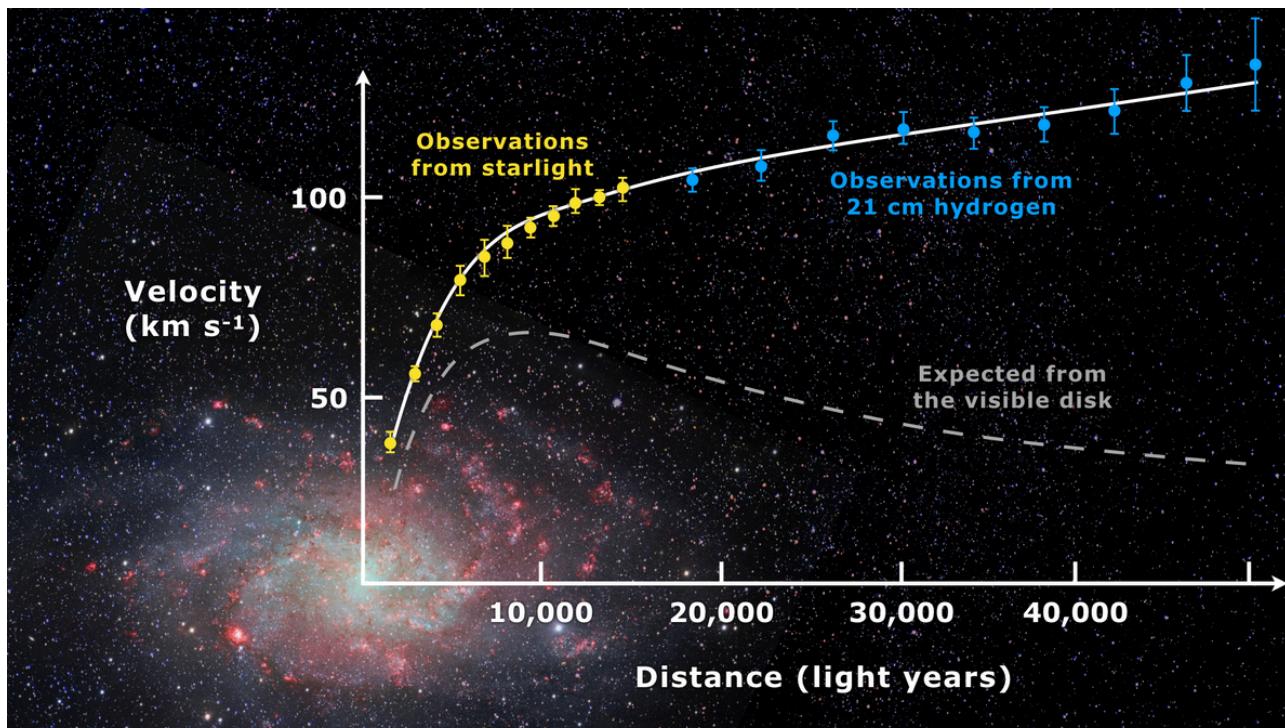
Dark Matter in Compact Objects, Stars, and in Low Energy Experiments

Institute for Nuclear Theory, University of Washington

August 15, 2022

in collaboration with Daniel Egana-Ugrinovic, Rouven Essig and Mukul Sholapurkar (to appear)

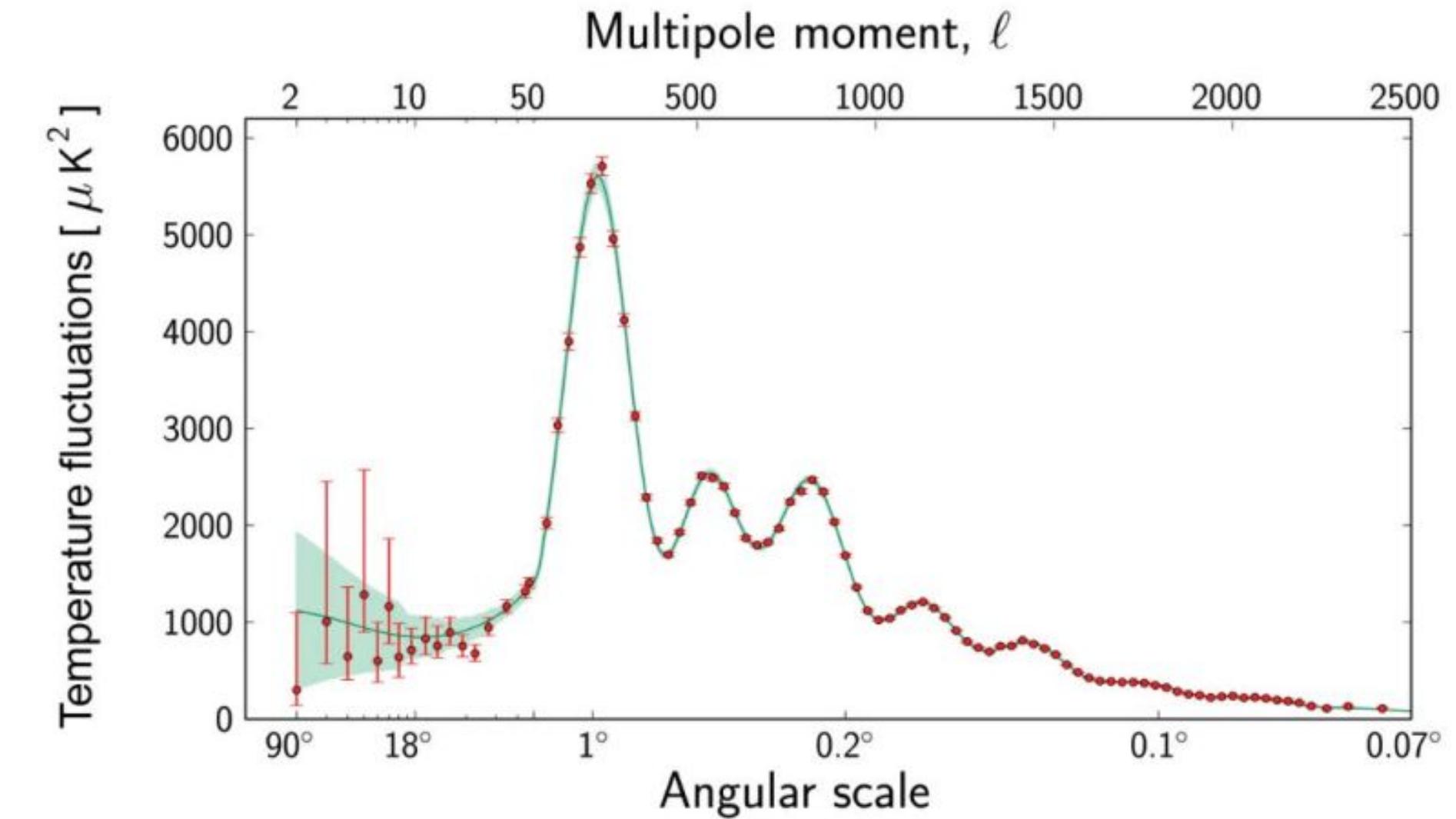
Dark matter



Galaxy



Galaxy Cluster

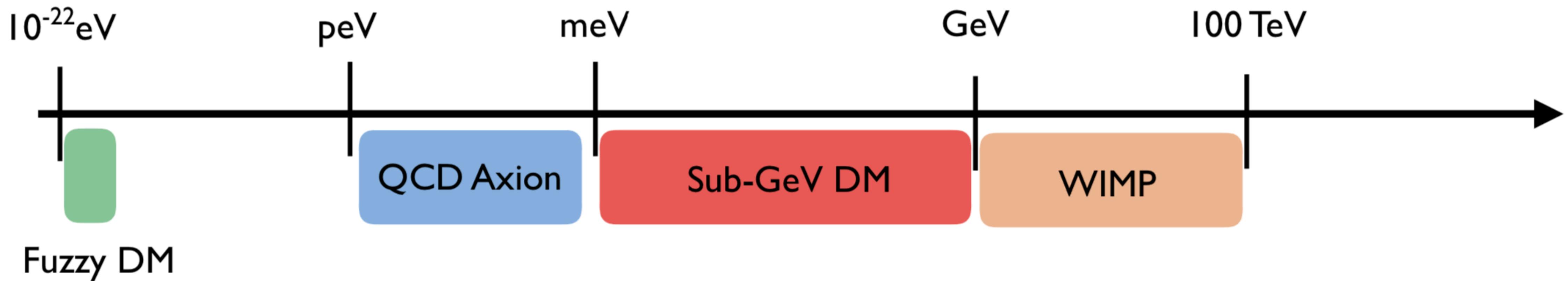


CMB

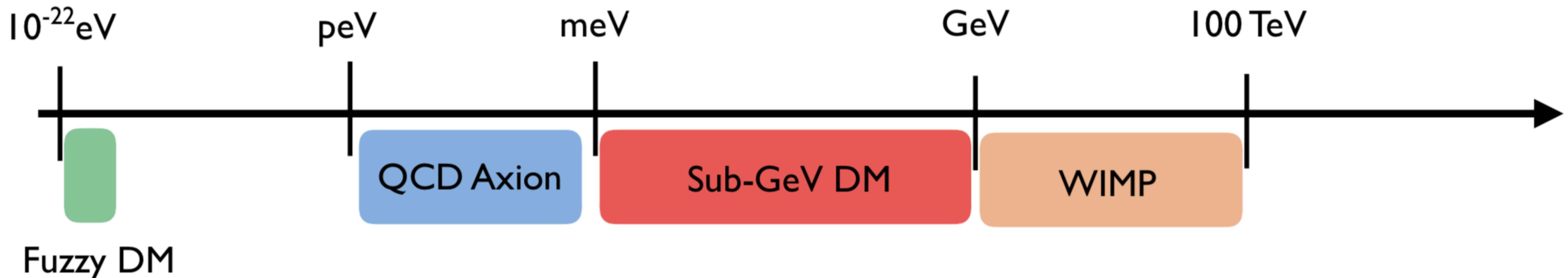
- 85% of matter, 27% total energy density in the Universe
- Evidence for dark matter is currently only gravitational

Particle nature is unknown, a wide range of DM masses are allowed

Dark matter

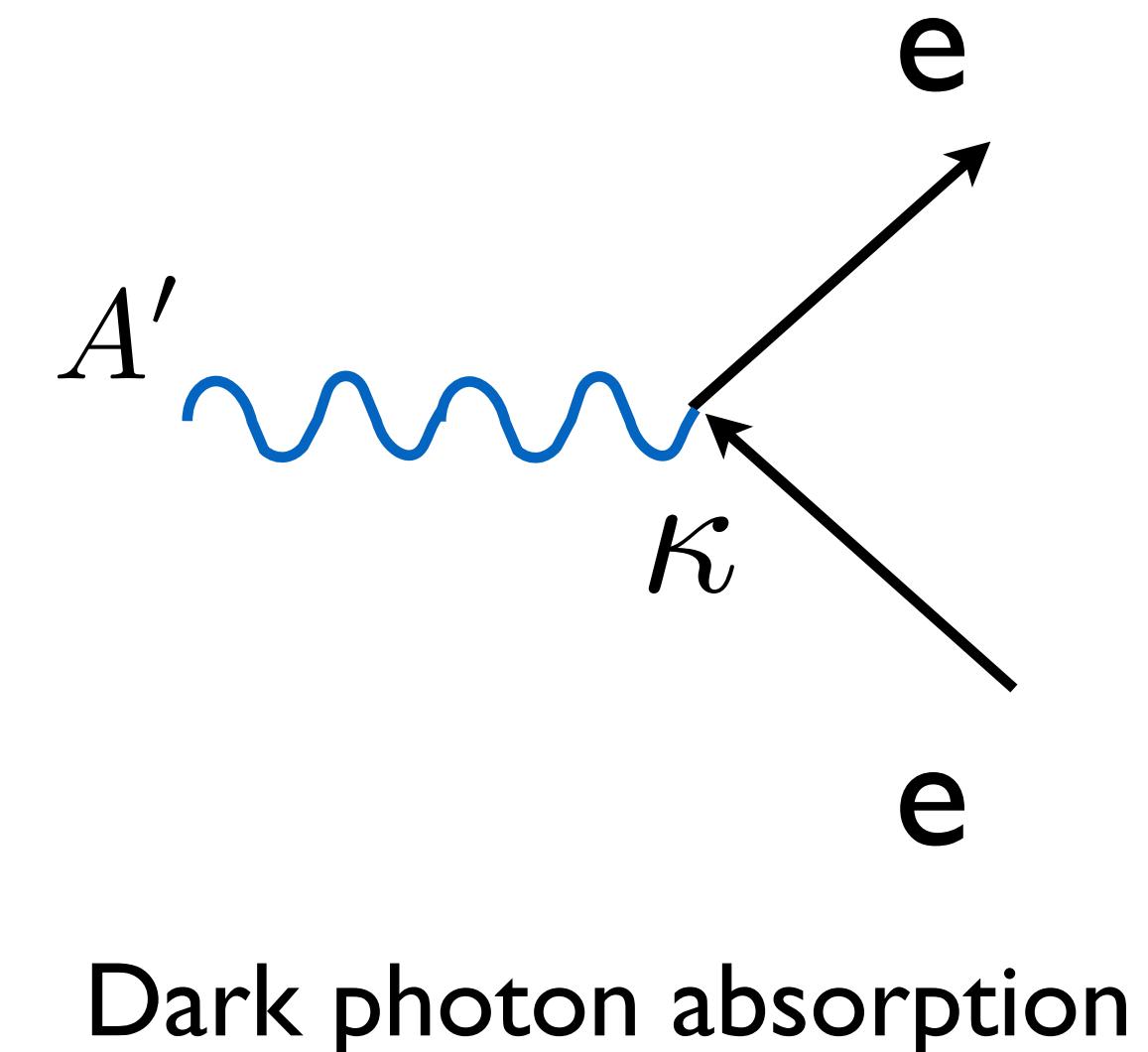
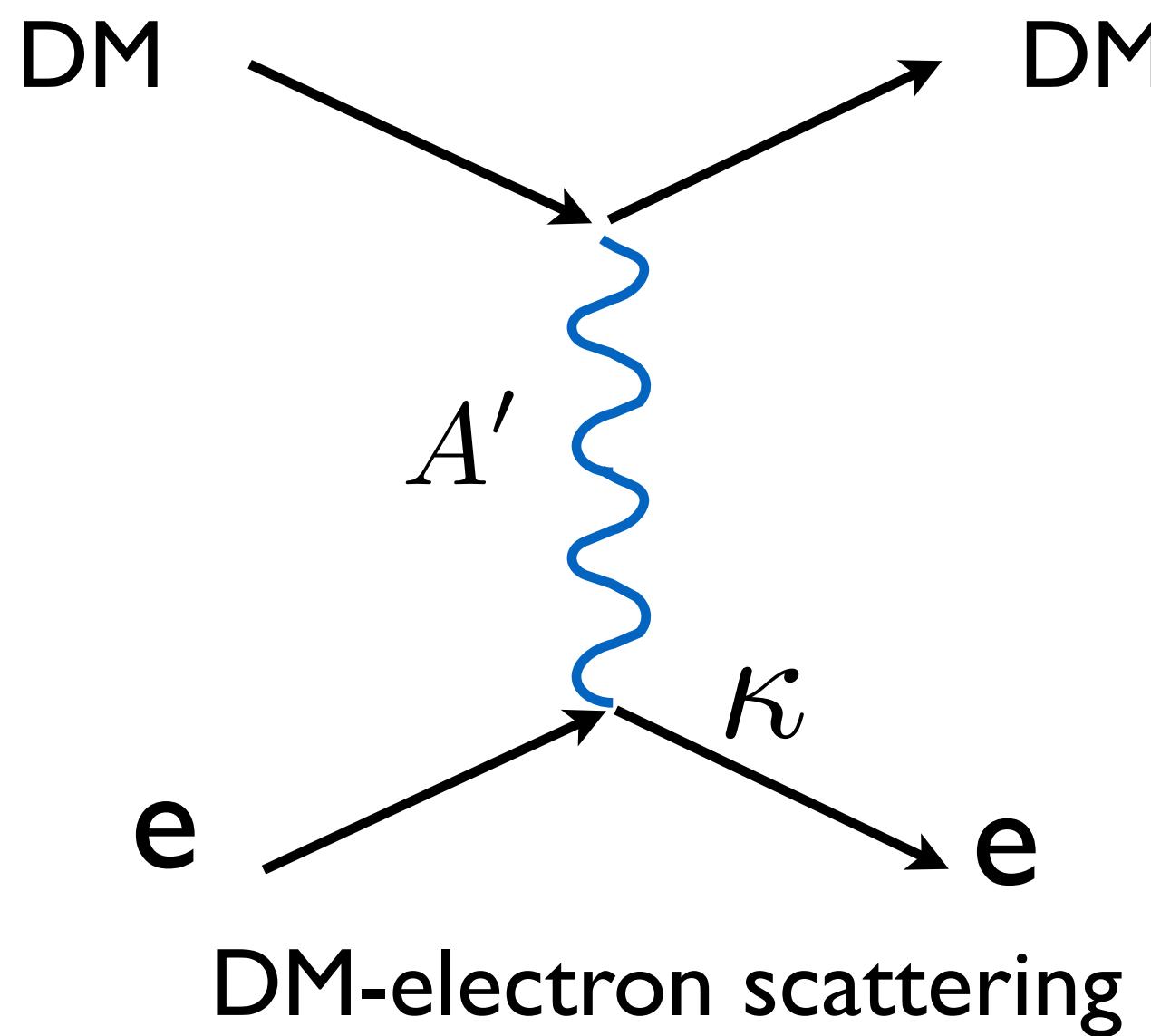


Sub-GeV dark matter

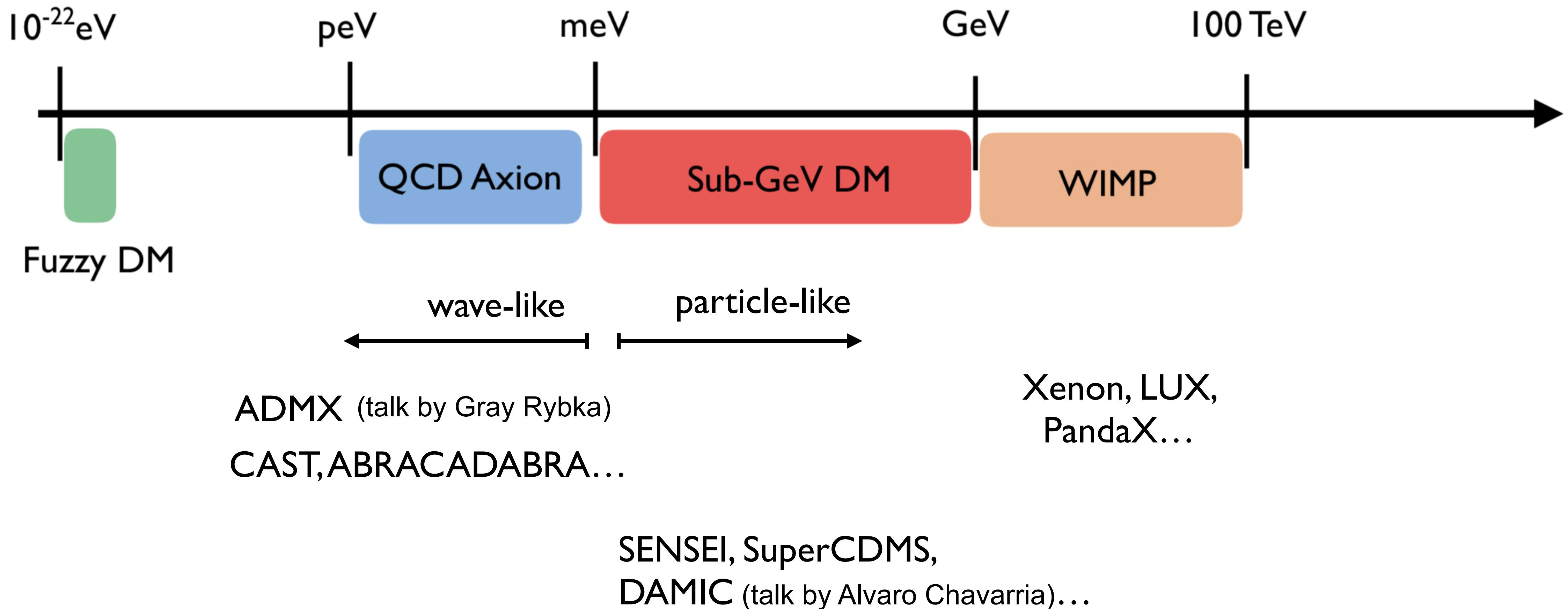


Dark photon model: $\mathcal{L} \supset -\frac{1}{4}F'^{\mu\nu}F'_{\mu\nu} - \frac{\kappa}{2}F^{\mu\nu}F'_{\mu\nu} + \frac{1}{2}m_A^2 A'^{\mu} A'_{\mu}$

Other models, see talk by Robert McGehee

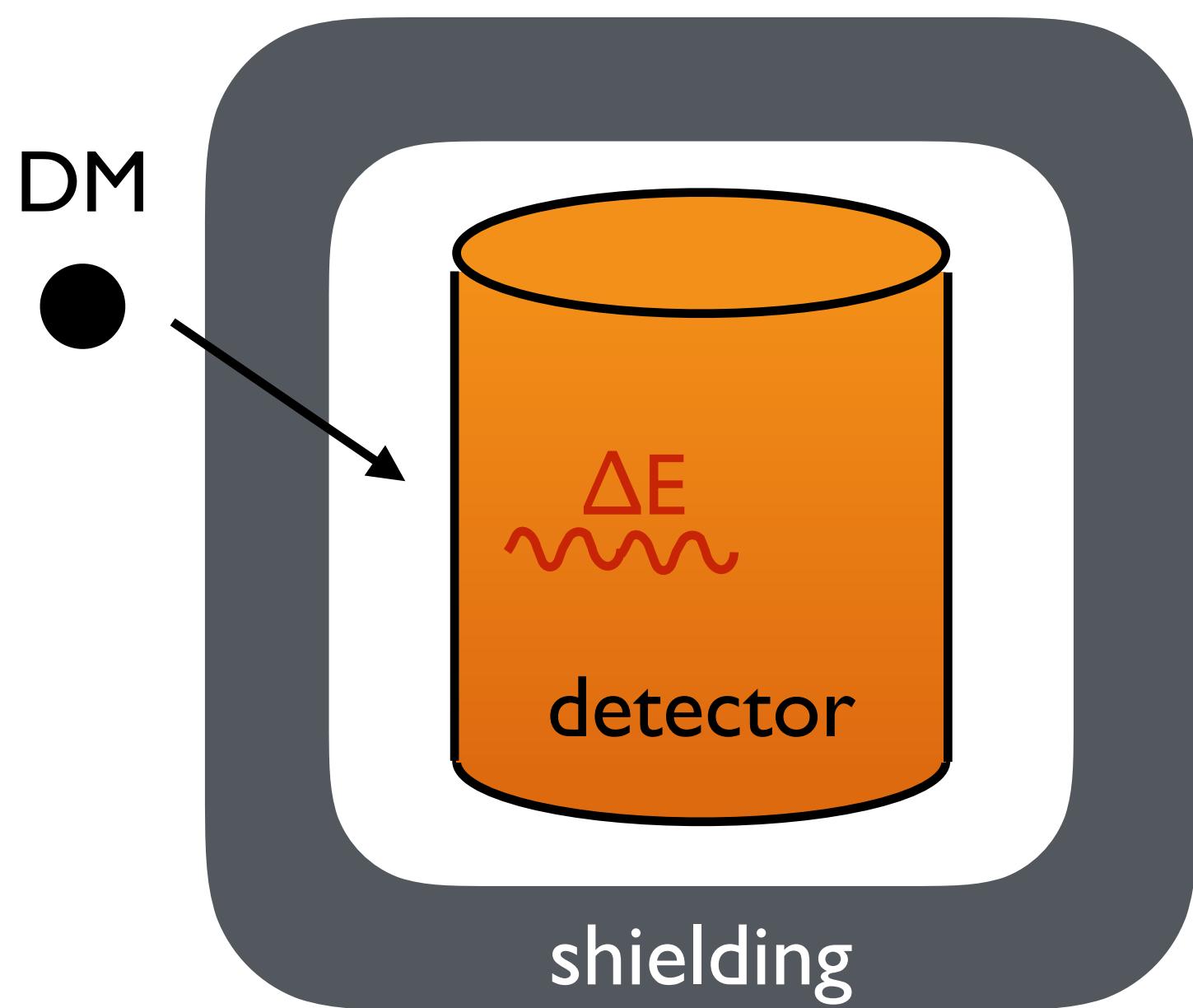


Direct Detection of DM

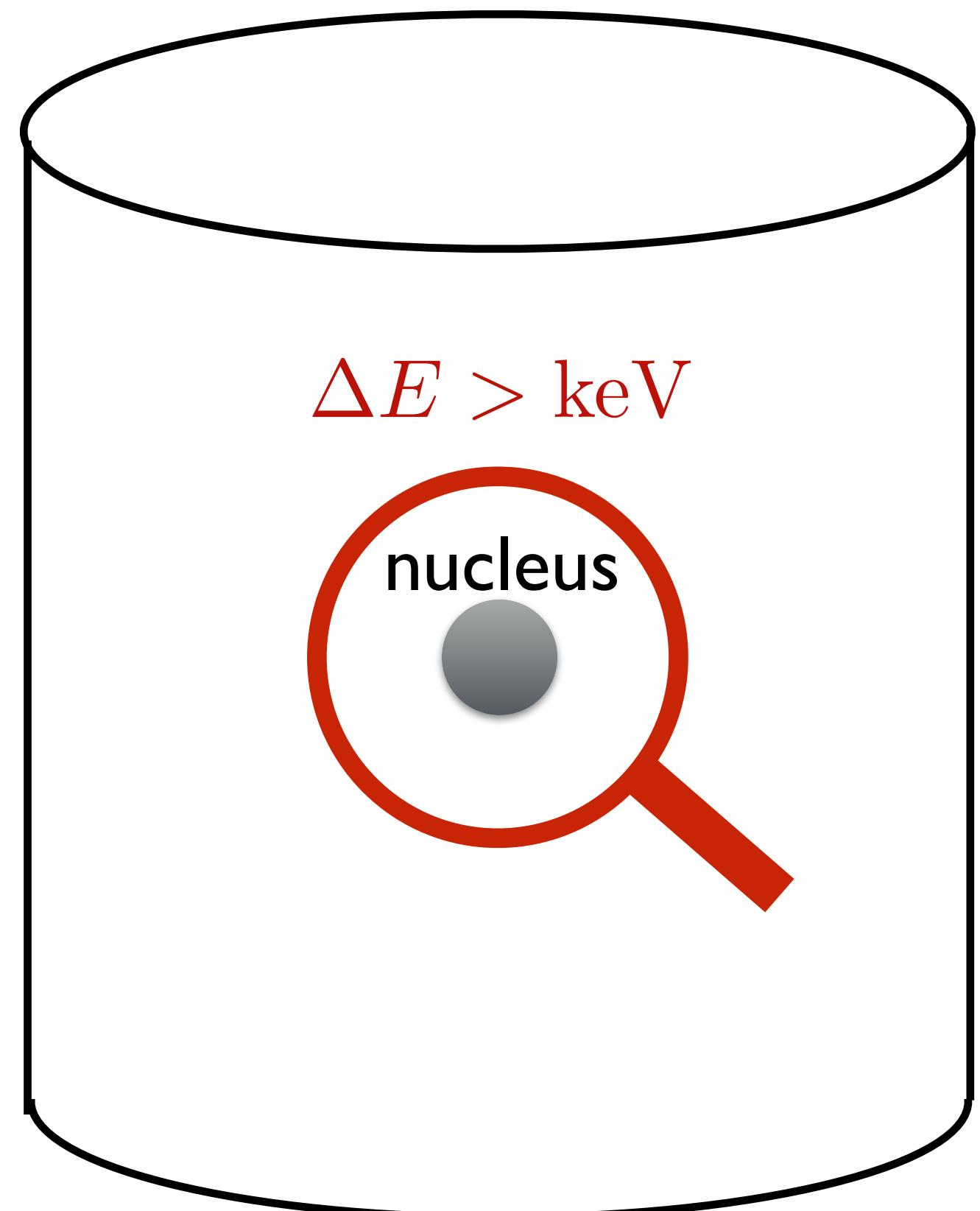


Direct Detection of DM

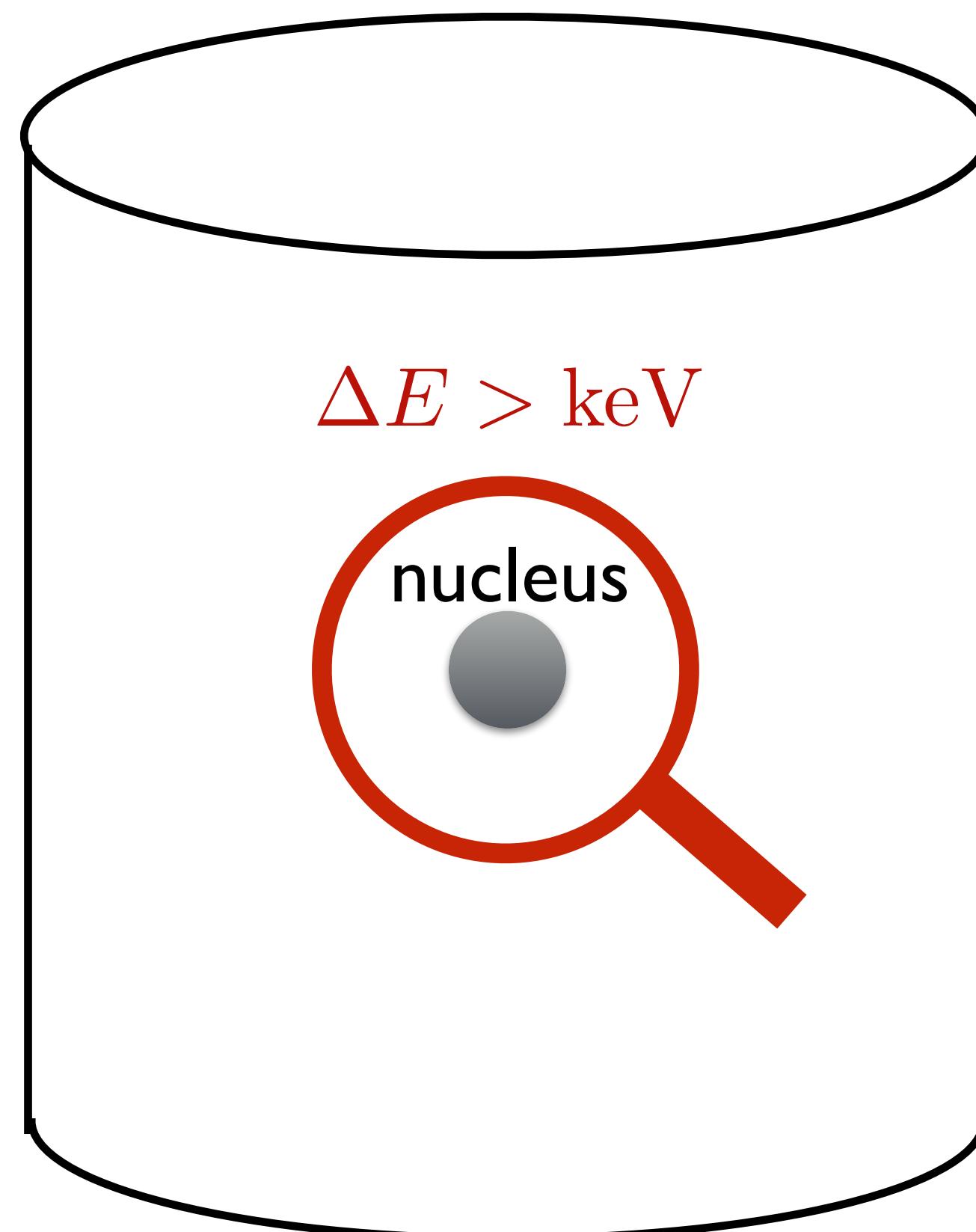
- Assuming DM has more than gravitational interactions with SM
- Clean environment, sensitive detector
- Wait for DM to come!



Direct Detection: $\Delta E > \text{keV}$



Direct Detection: $\Delta E > \text{keV}$



Elastic DM-nuclear scattering

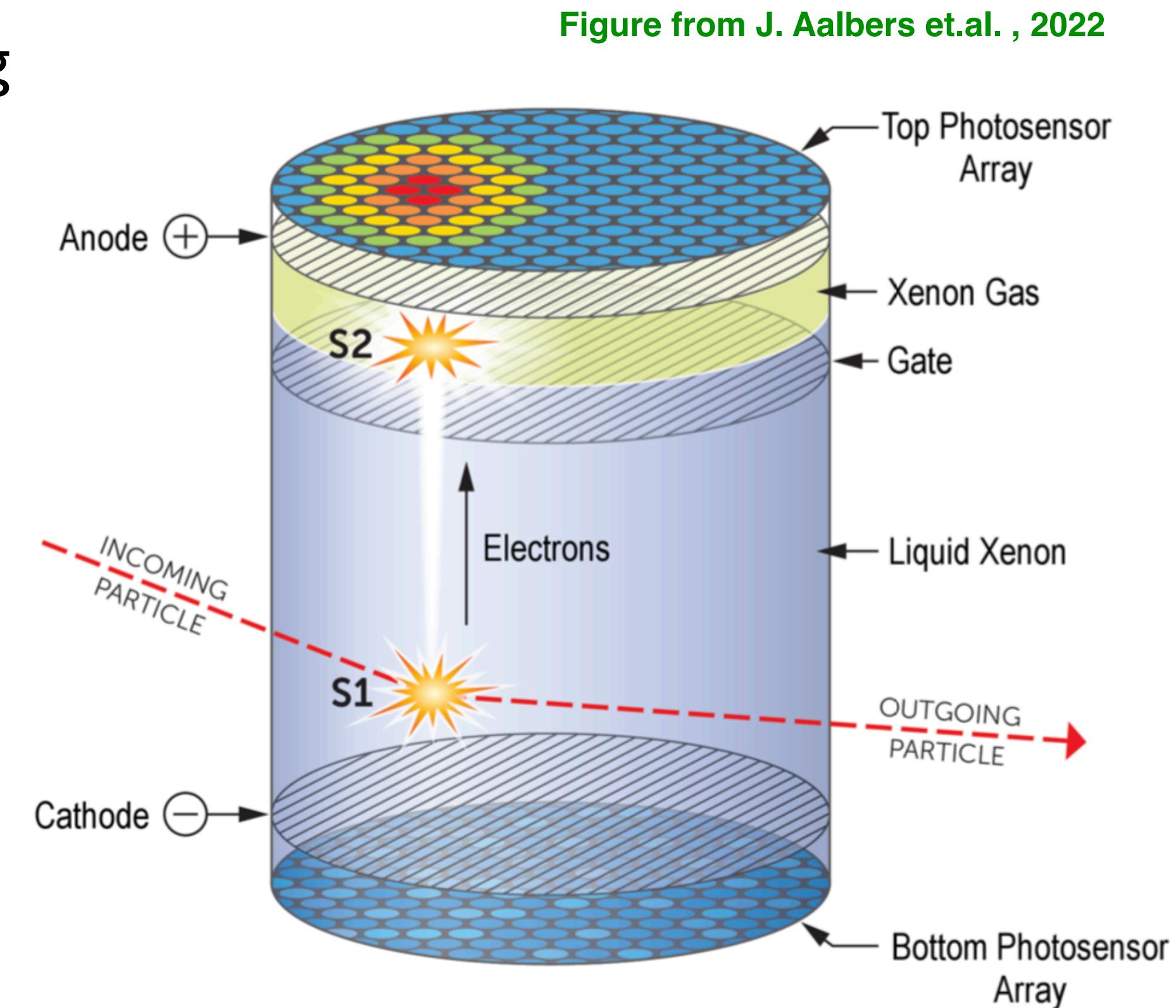
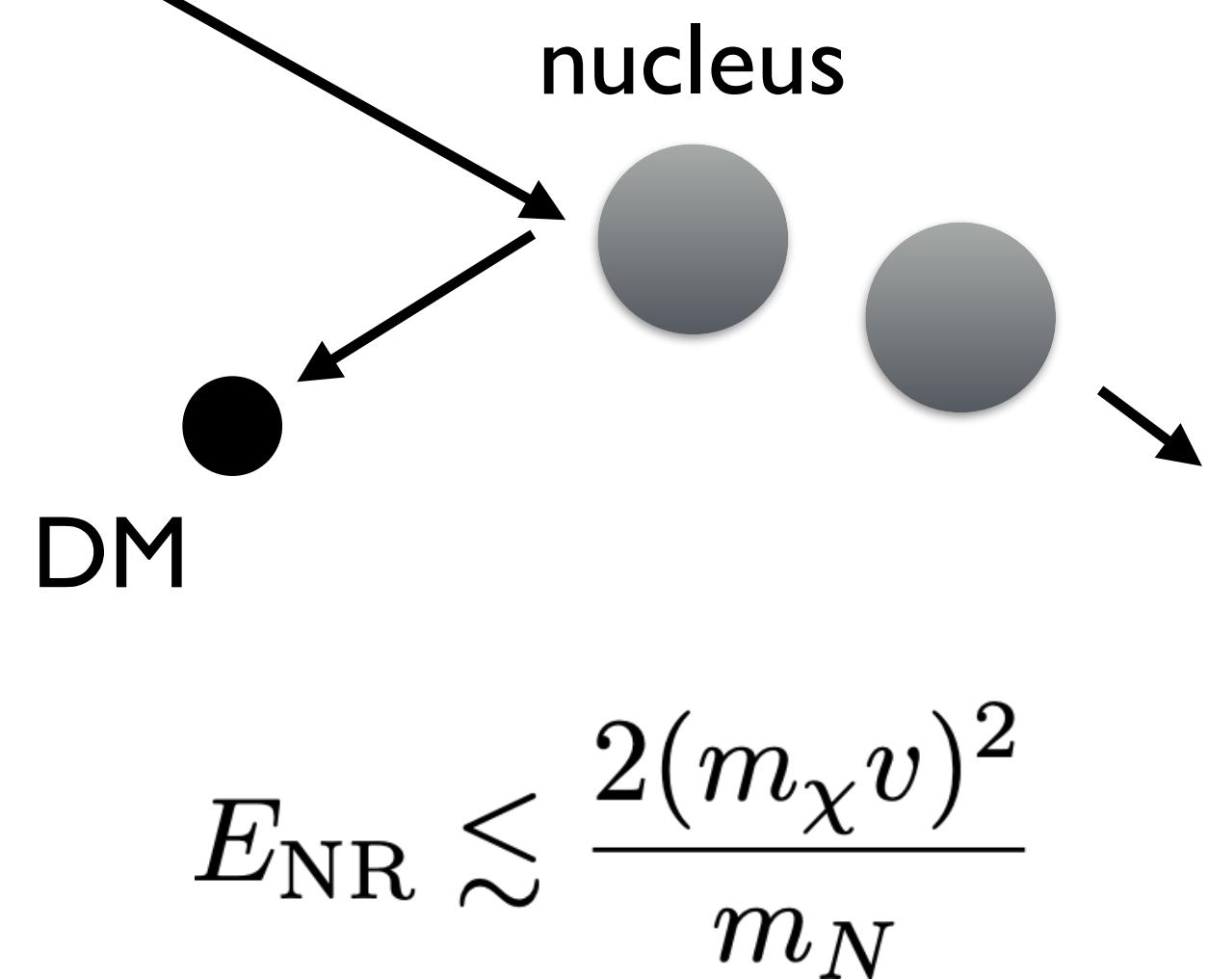
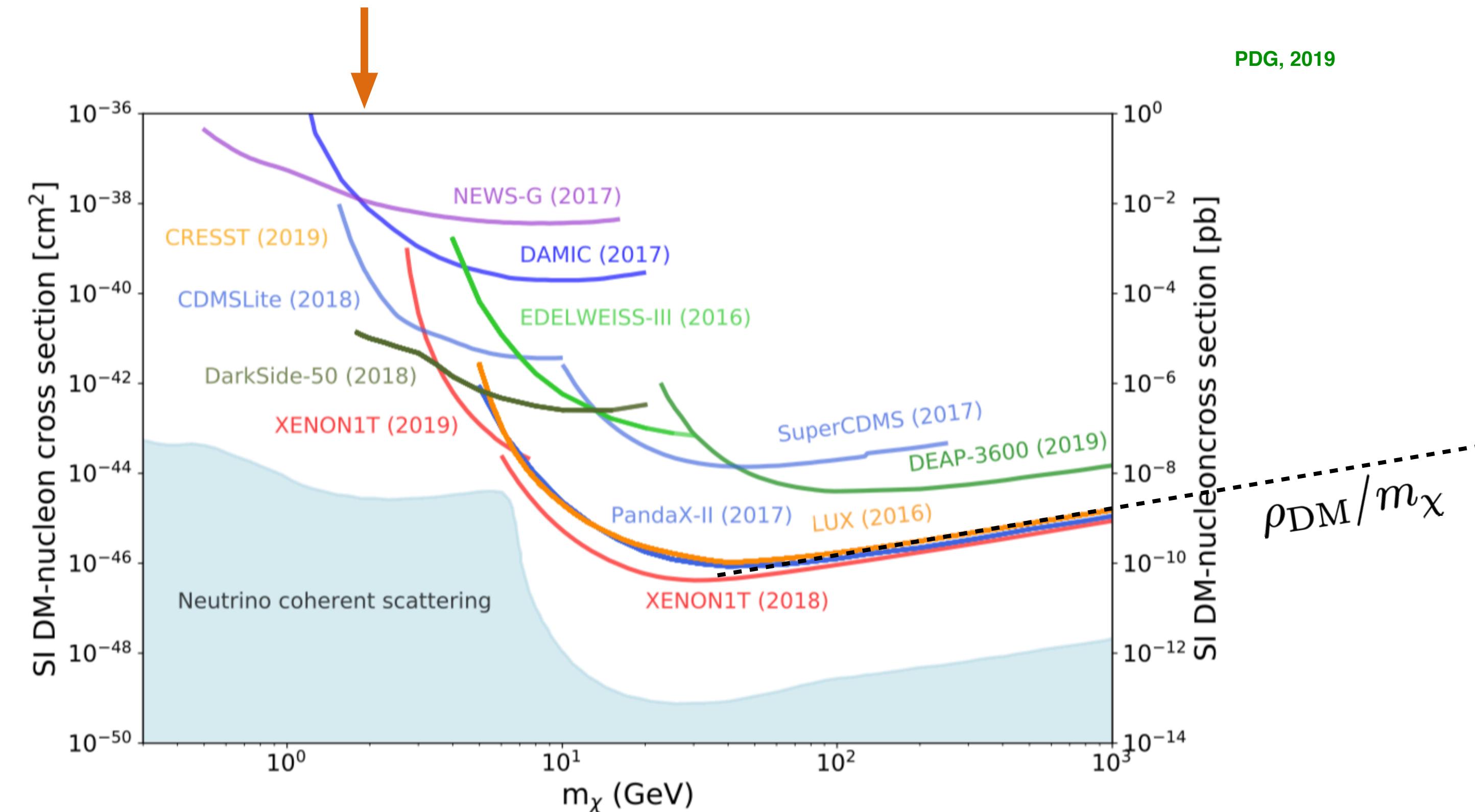


Figure from J. Aalbers et.al. , 2022

Signals: S1+S2
Threshold: $\sim \text{keV}$

Nuclear recoil constraints

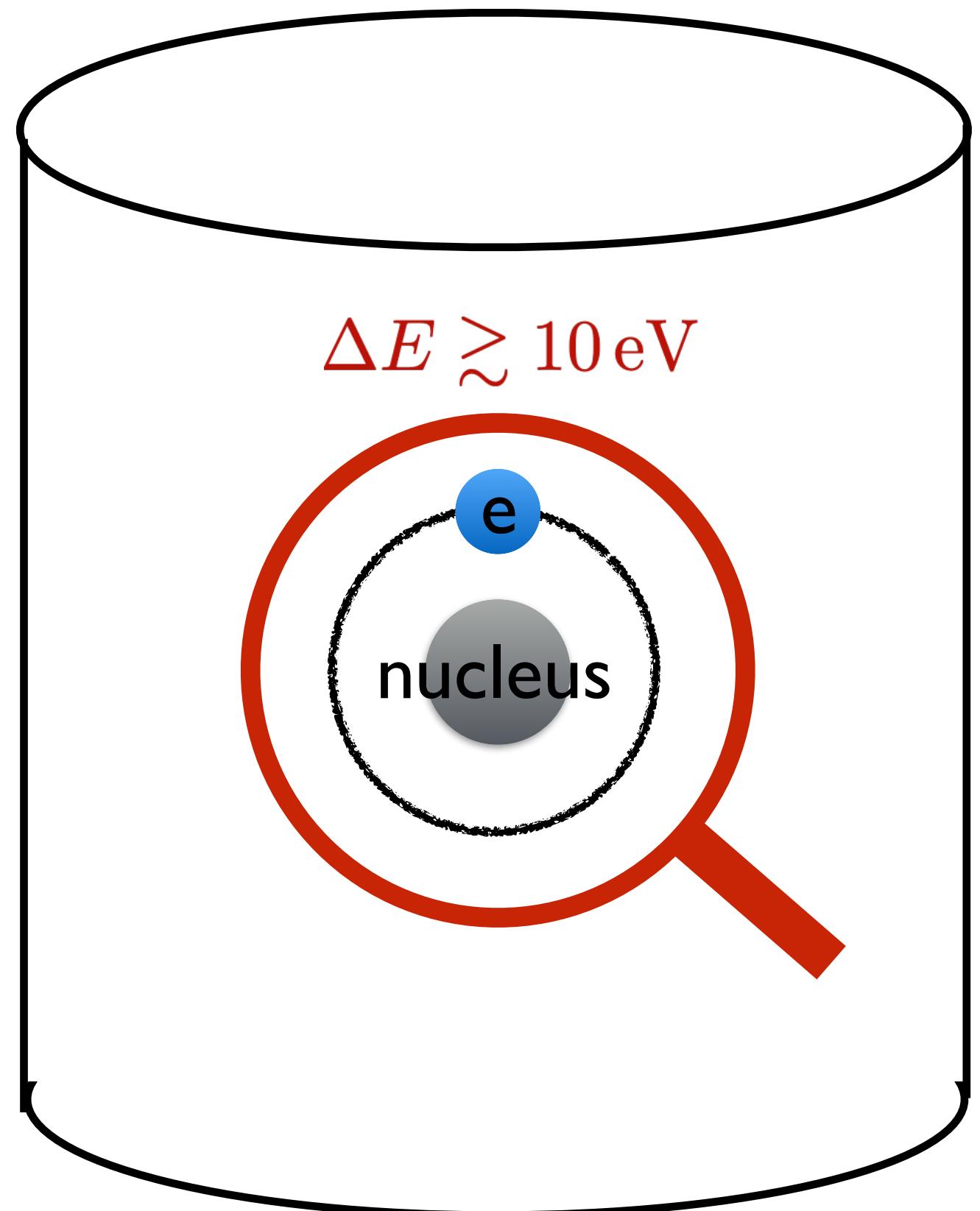
Limited by \sim keV
threshold



Insufficient energy transfer

$$E_{\text{NR}} \lesssim 1 \text{ keV} \left[\frac{m_\chi}{4 \text{ GeV}} \right]^2 \left[\frac{100 \text{ GeV}}{M_N} \right]$$

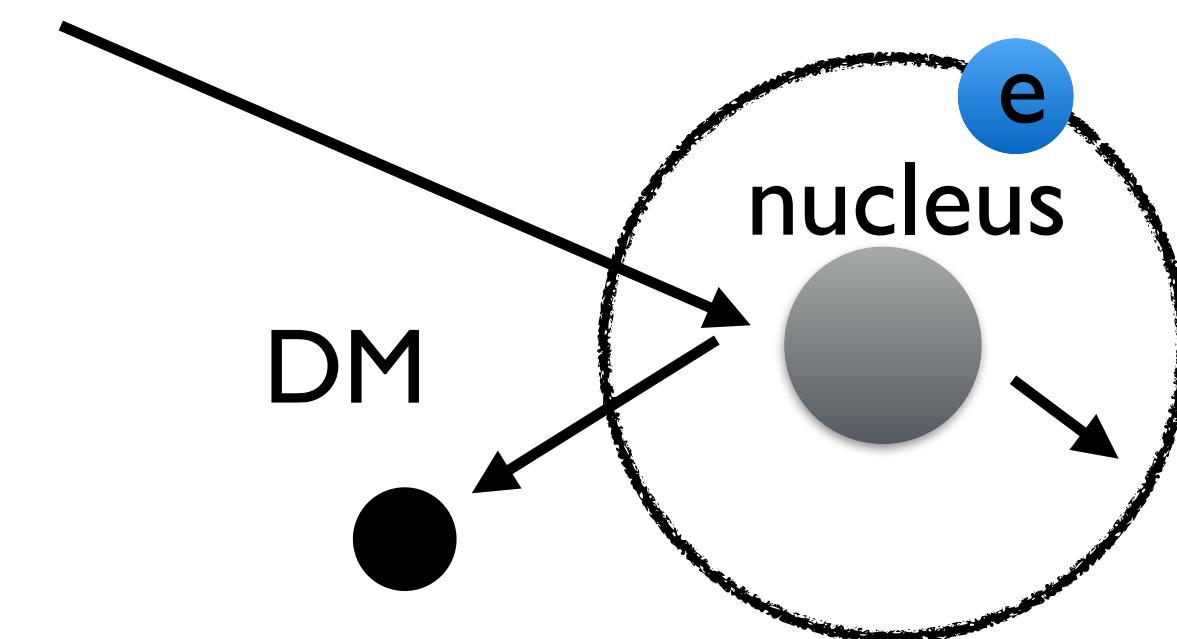
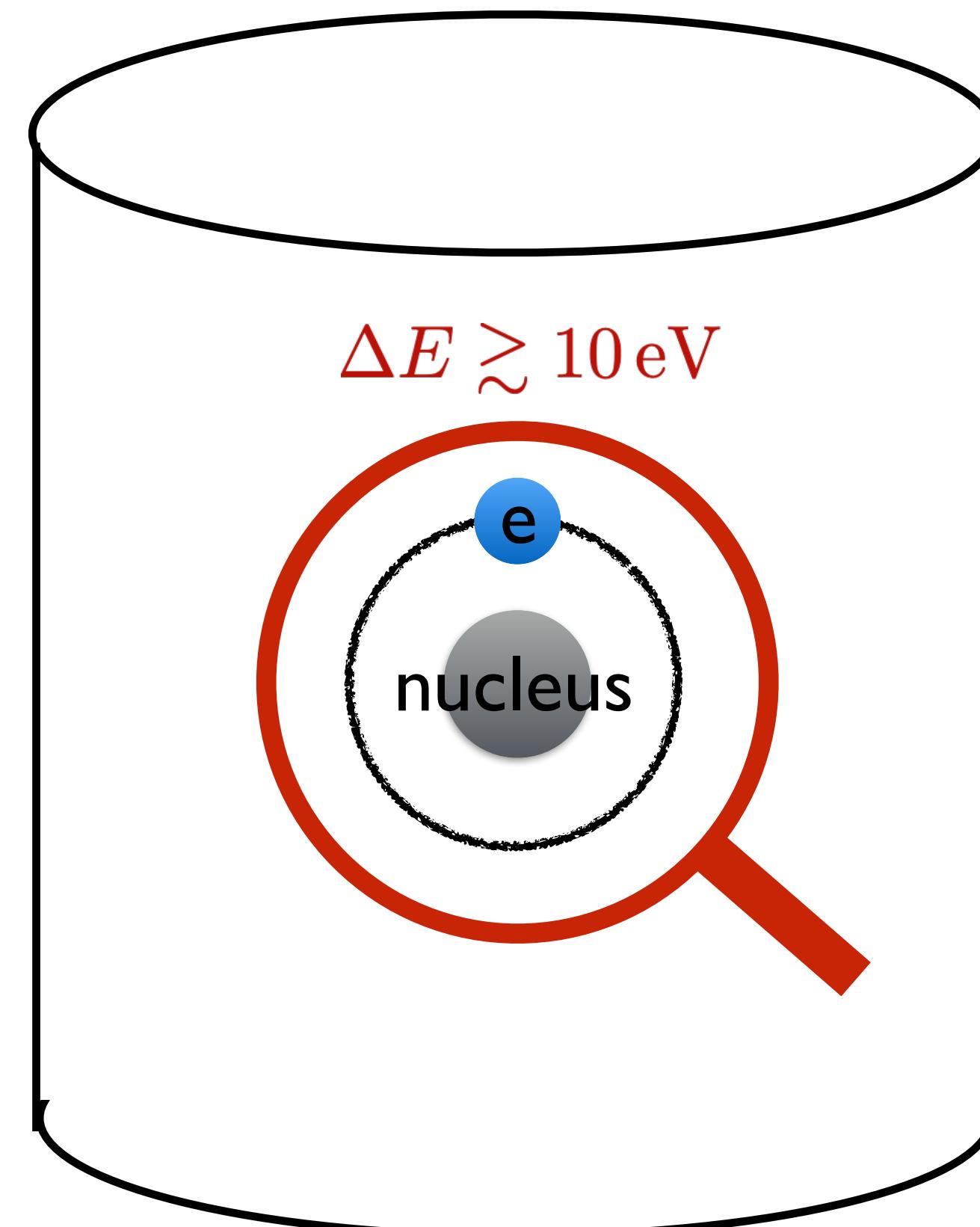
Direct Detection: $\Delta E > \mathcal{O}(10)\text{eV}$



Direct Detection: $\Delta E > O(10)$ eV

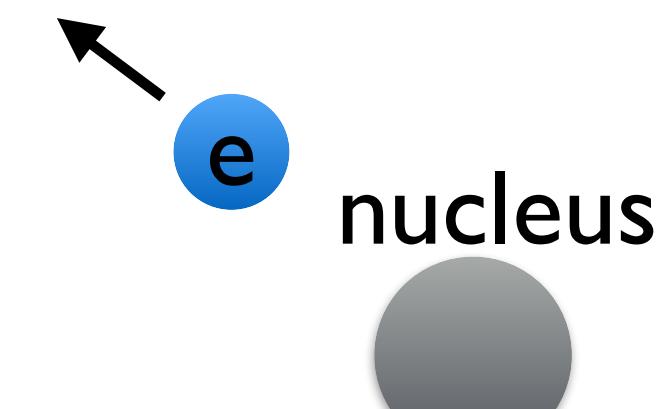
Vergados, Ejiri , 2004

Ive, Nakano, Shoji, Suzuki, 2017



Migdal effect

(See also talks by Nicole Bell and Kim Berghaus)



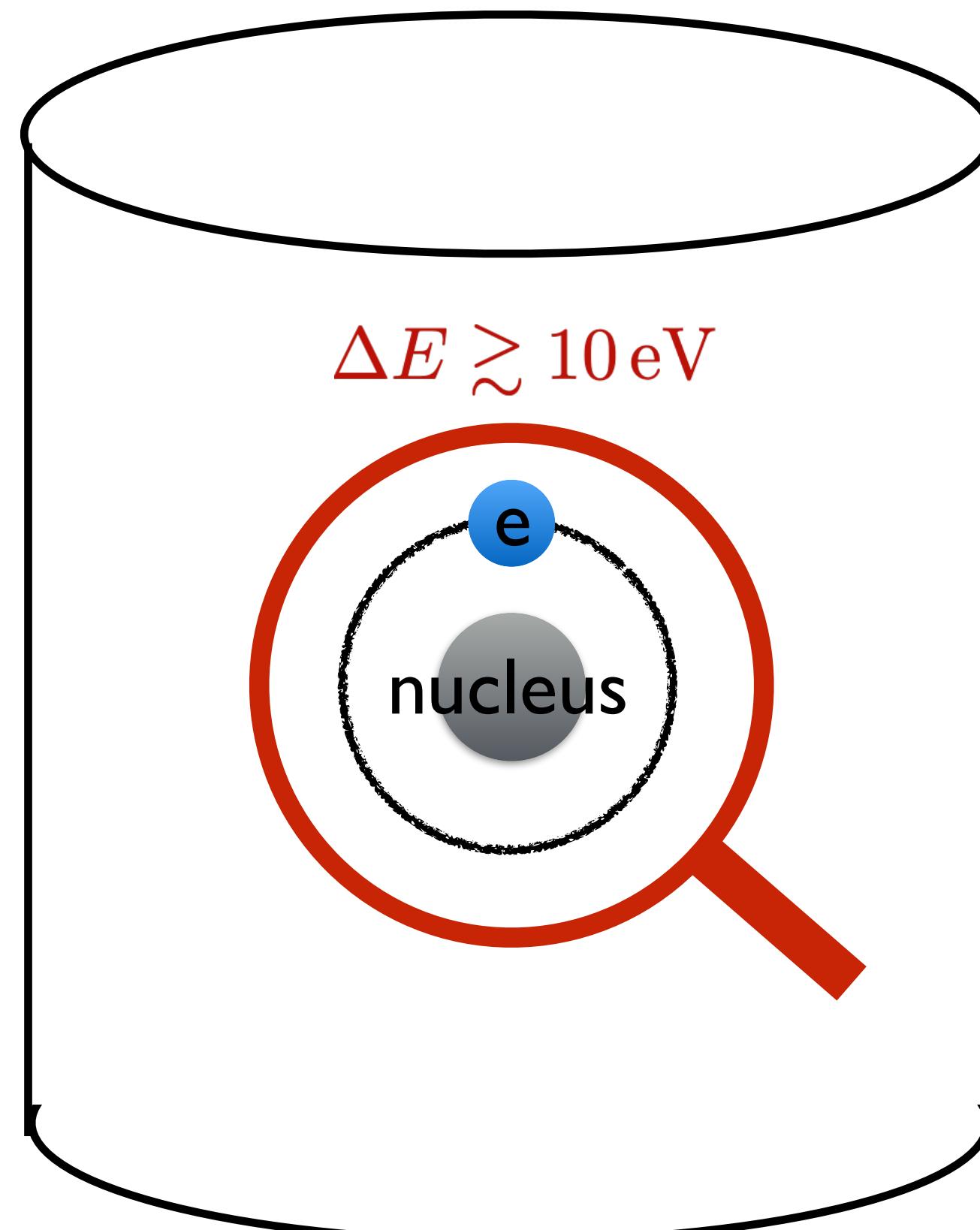
Electron ionization

Signal: electron ionization (S2)

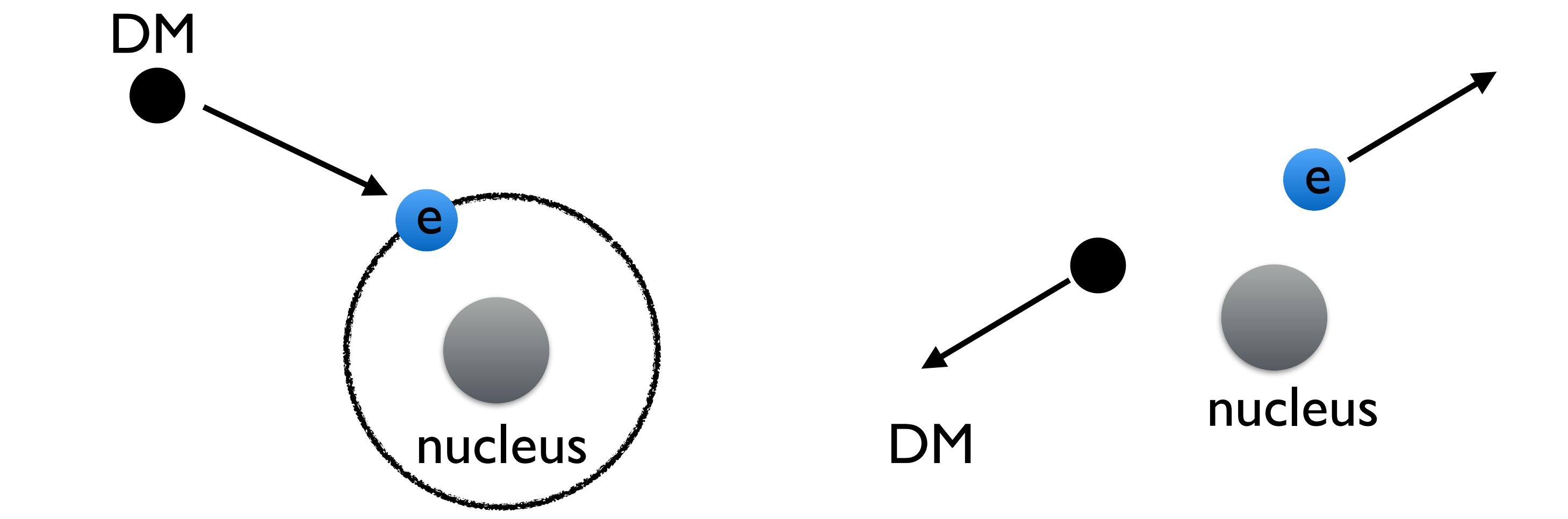
Threshold: ~ 10 eV

Direct Detection: $\Delta E > O(10)$ eV

Essig, Mardon, Volansky, 2011



DM-electron scattering

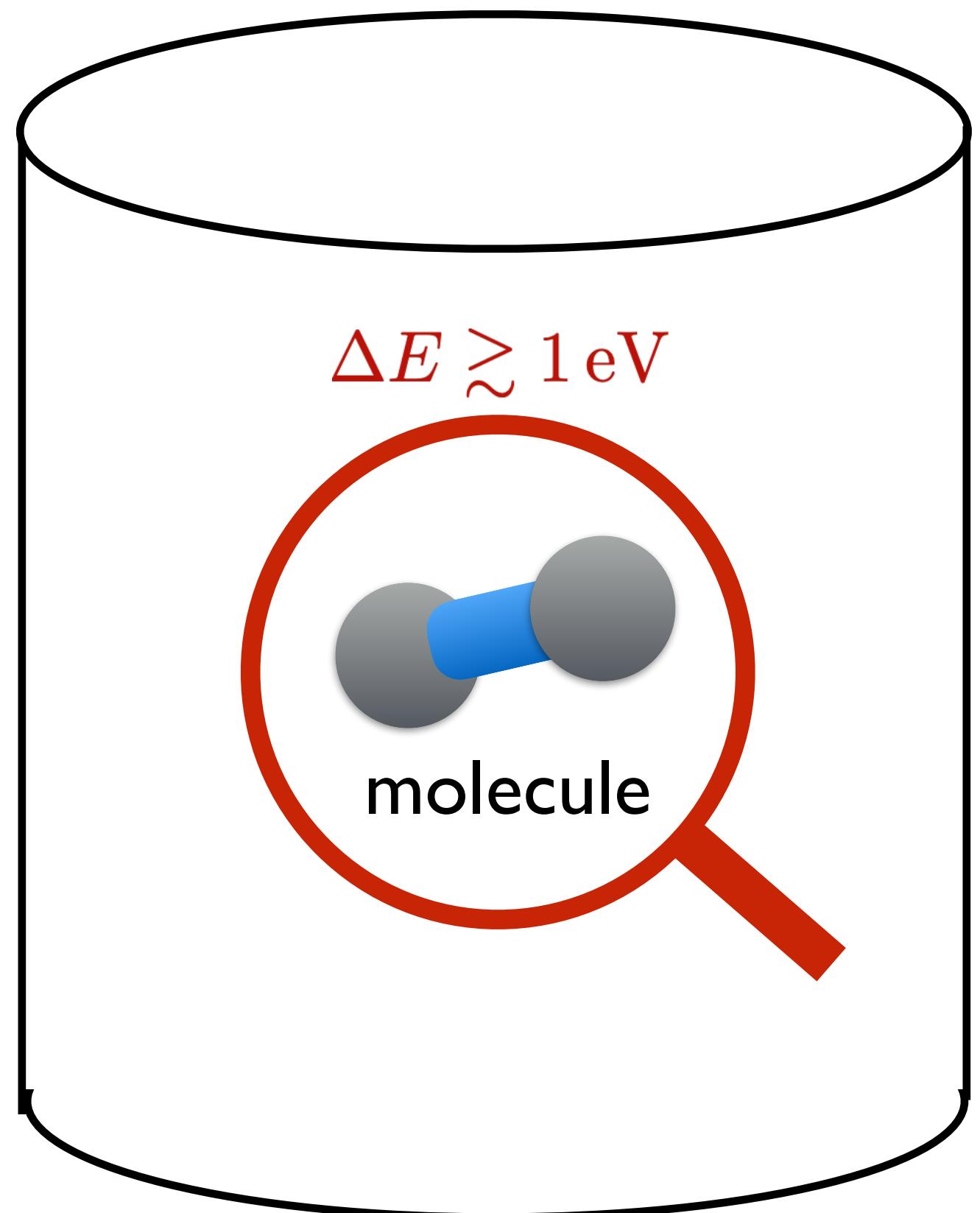


Efficient energy transfer
for light DM

$$E_{\text{ER}} \lesssim \frac{1}{2} m_\chi v^2 \gg E_{\text{NR}} \lesssim \frac{2(m_\chi v)^2}{m_N}$$

Signal: electron ionization (S2)
Threshold: ~ 10 eV

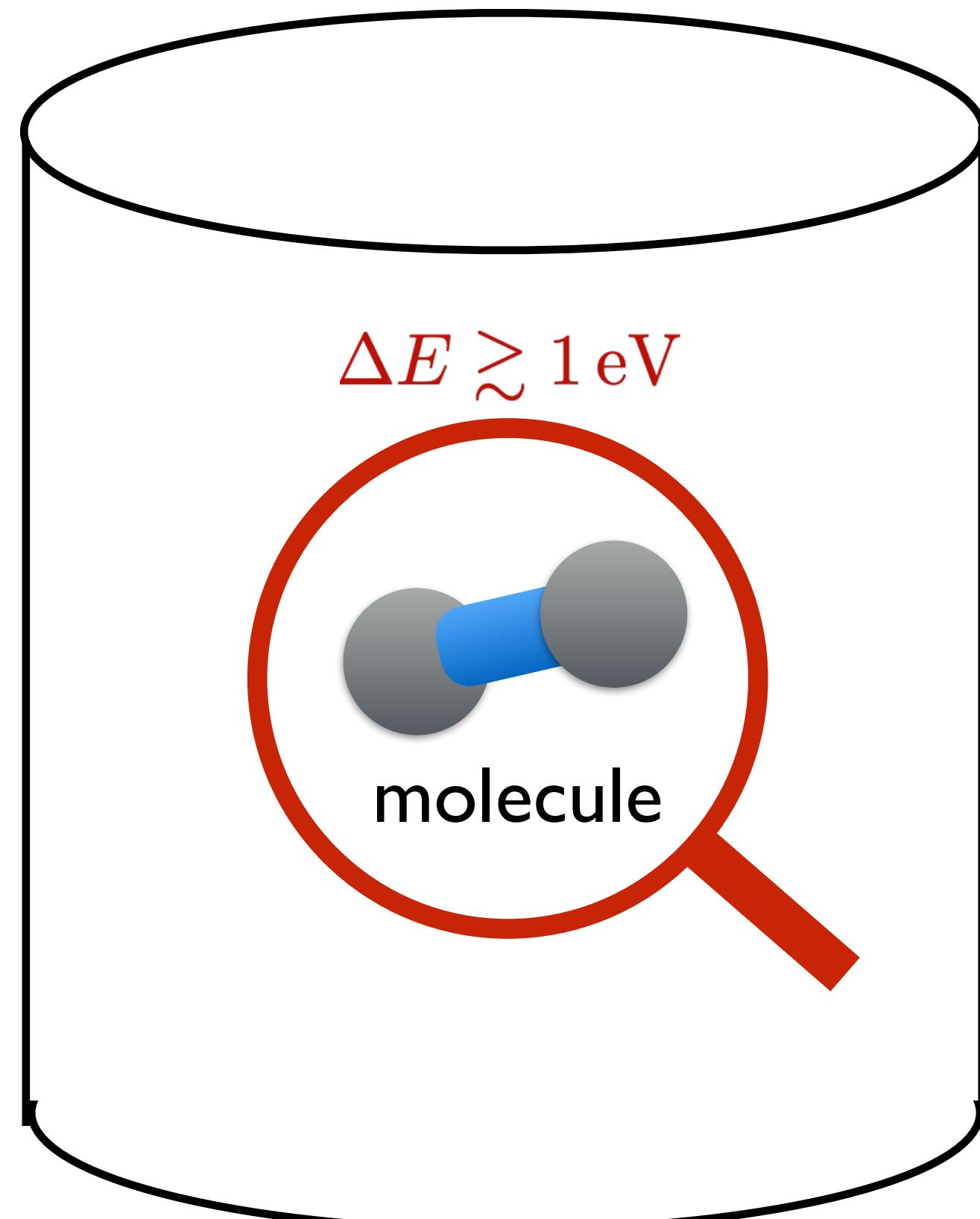
Direct Detection: $\Delta E > O(1)$ eV



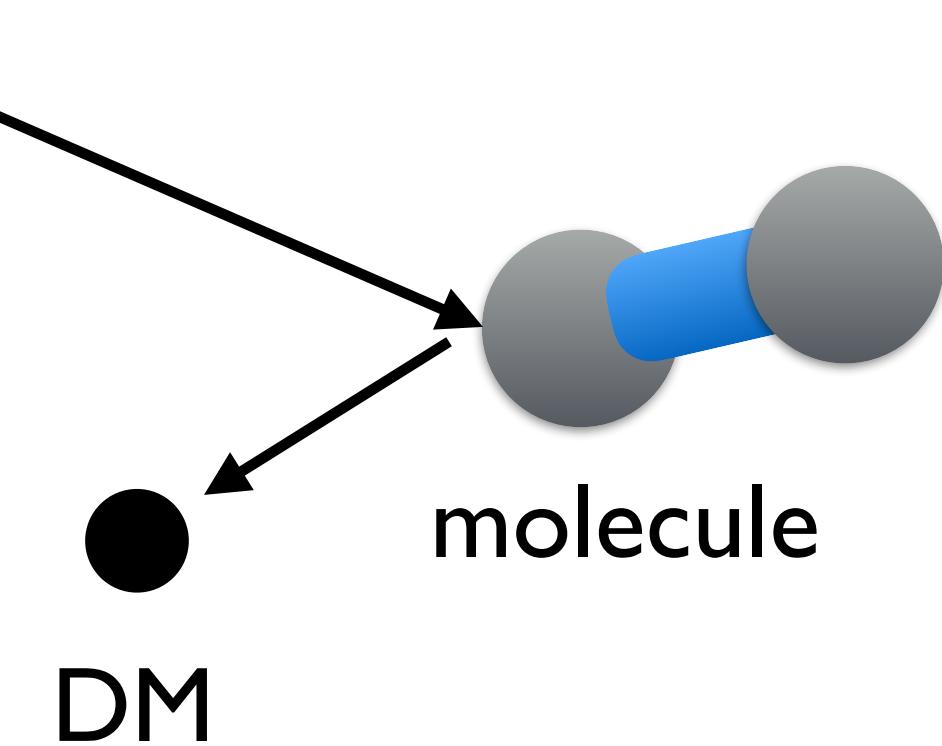
Direct Detection: $\Delta E > O(1)$ eV

Blanco, Collar, Kahn, Lillard, 2019

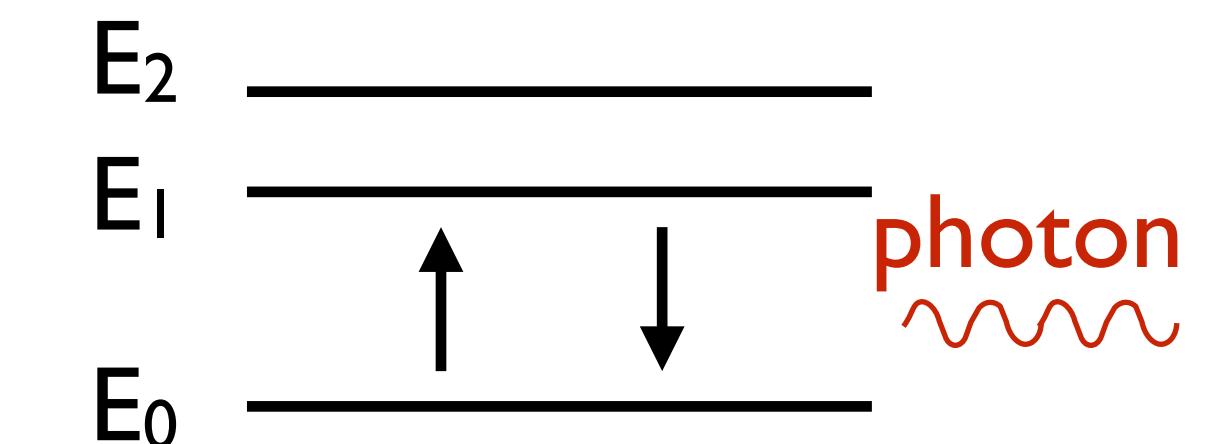
Essig, Perez-Rios, Ramani, Slone, 2019



Excitation in molecules

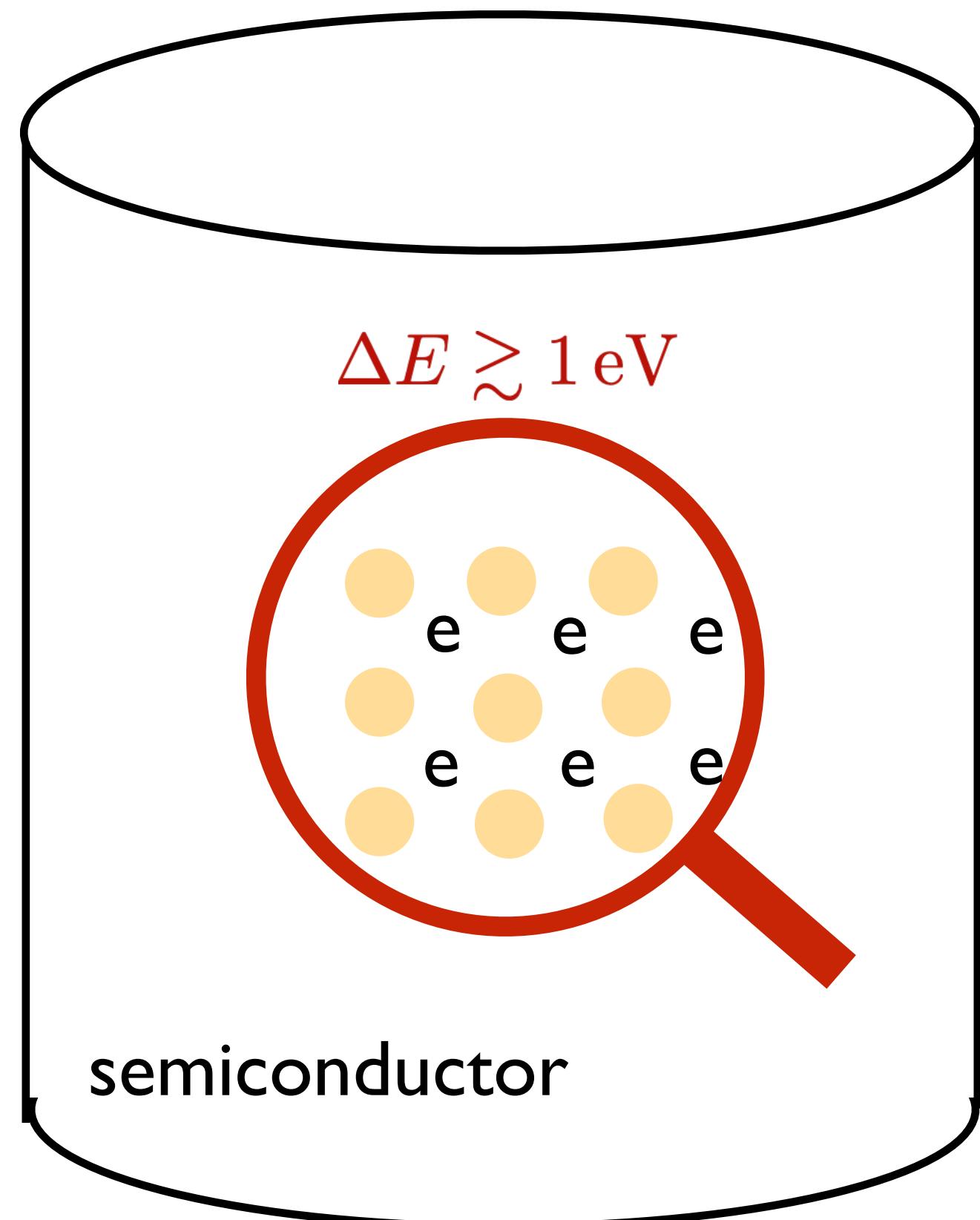


Signal: photons
Threshold: $O(1)$ eV

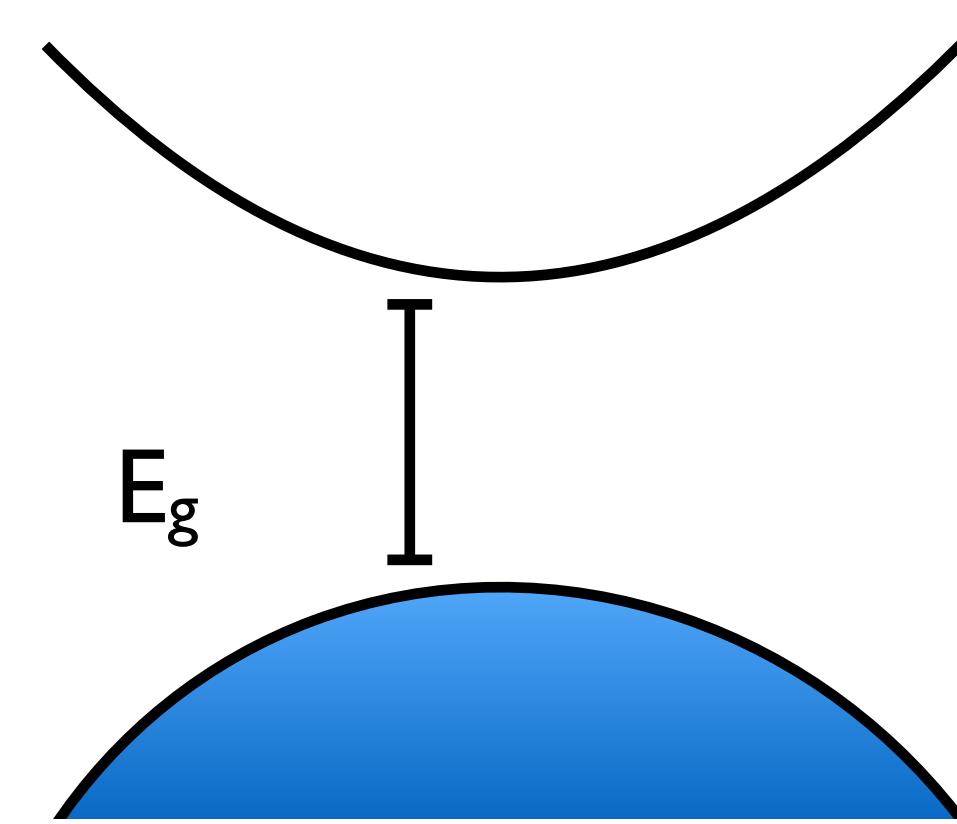


Direct Detection: $\Delta E > O(1) \text{ eV}$

Essig, Mardon, Volansky, 2011

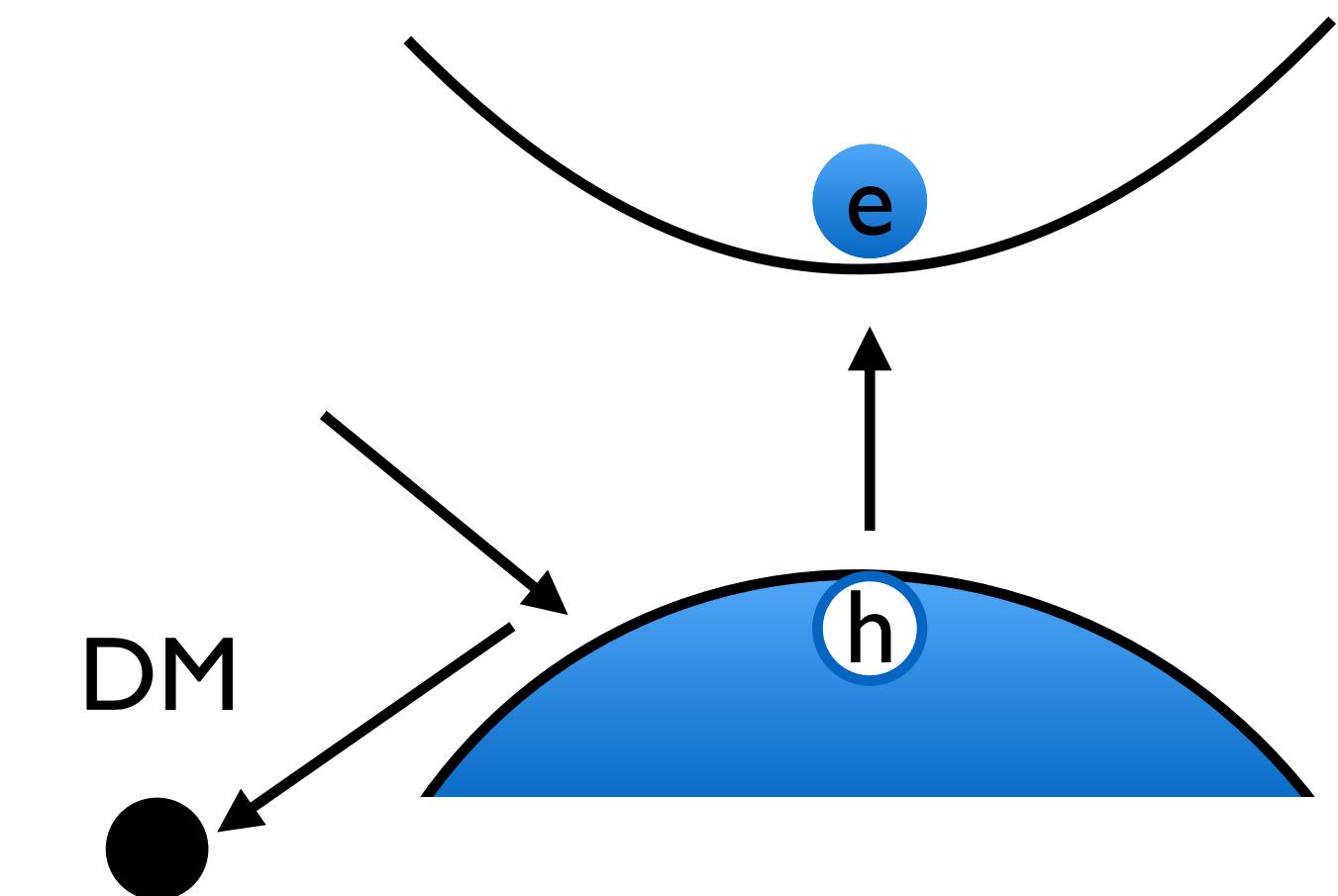


Electron ionization in semiconductors



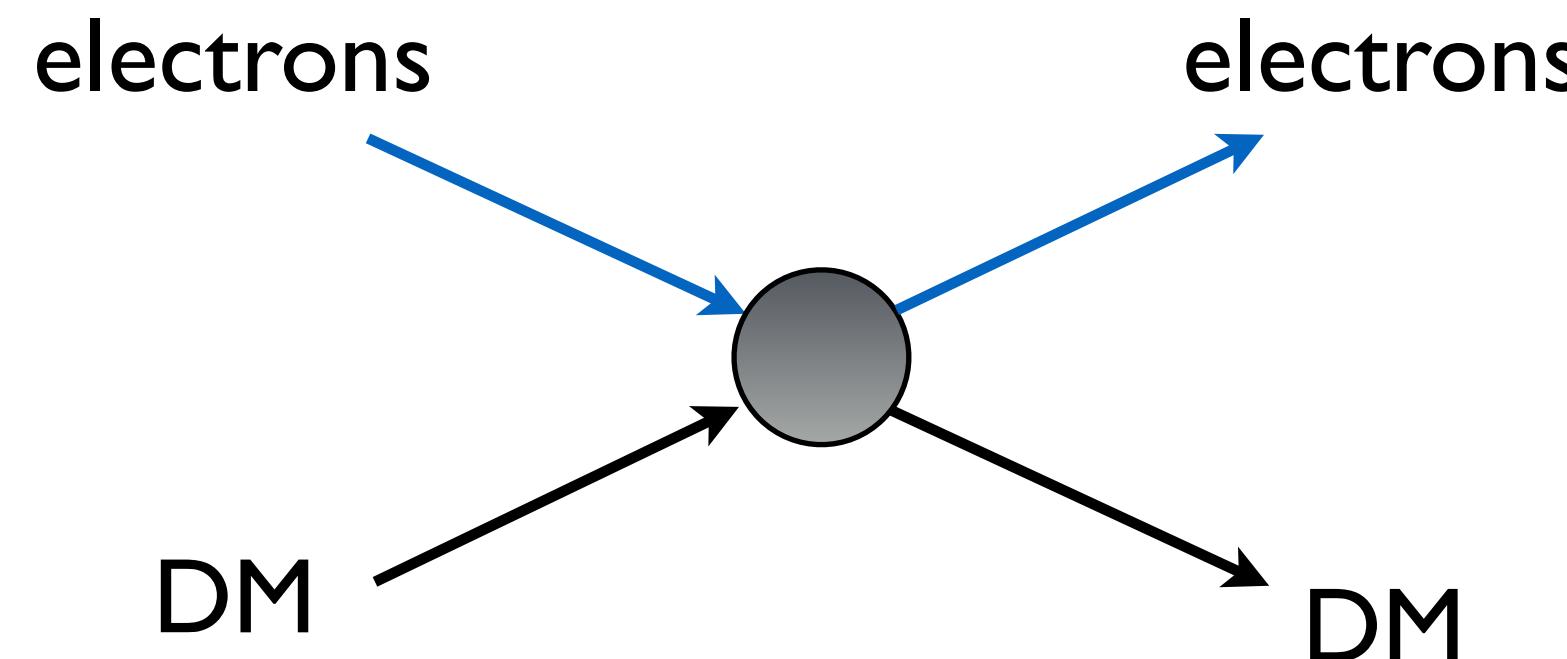
Signals: eh pairs

Threshold: $E_g \sim 1 \text{ eV}$



Direct detection of sub-GeV DM

Electron recoils



Access to whole kinetic energy:

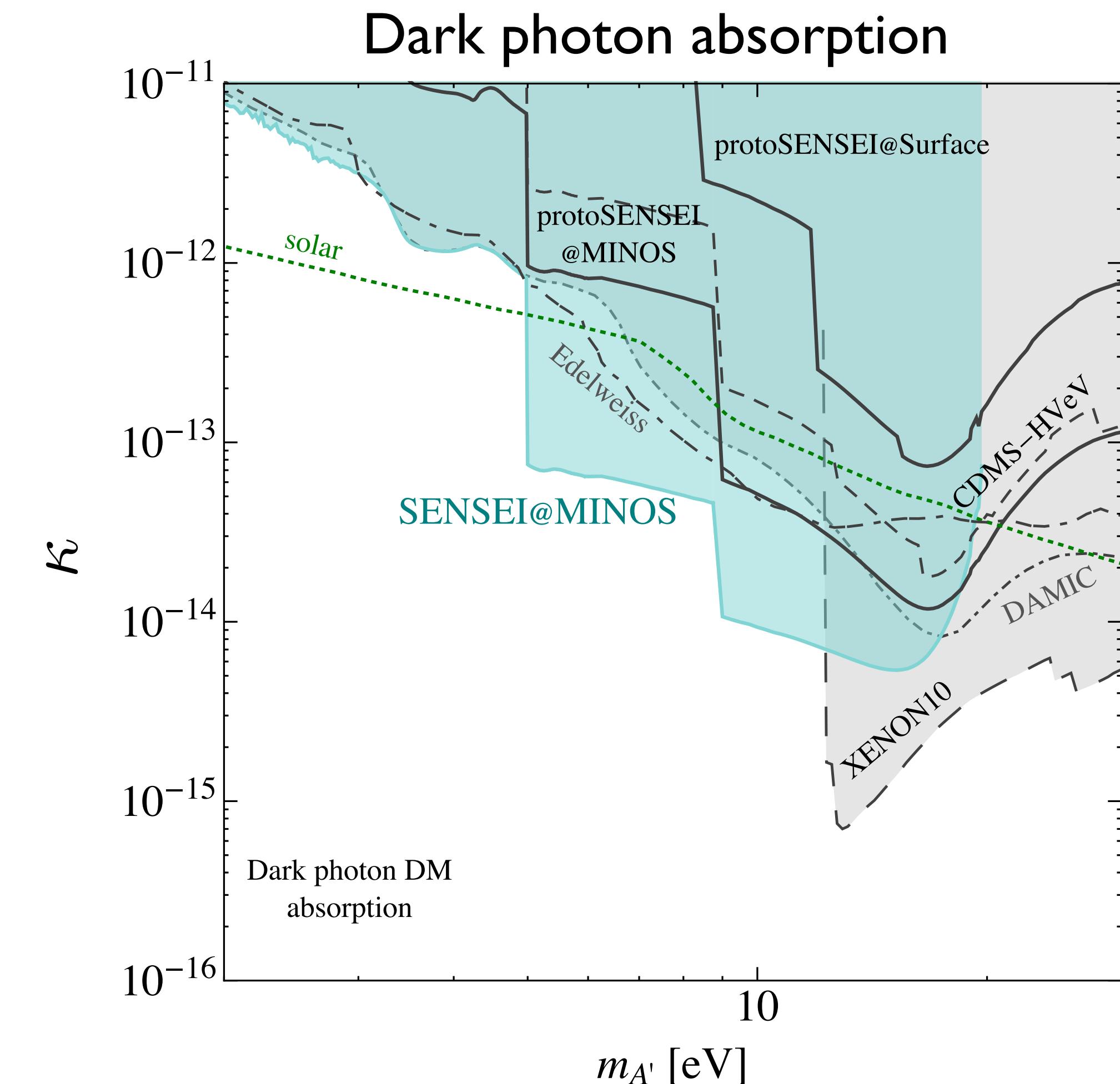
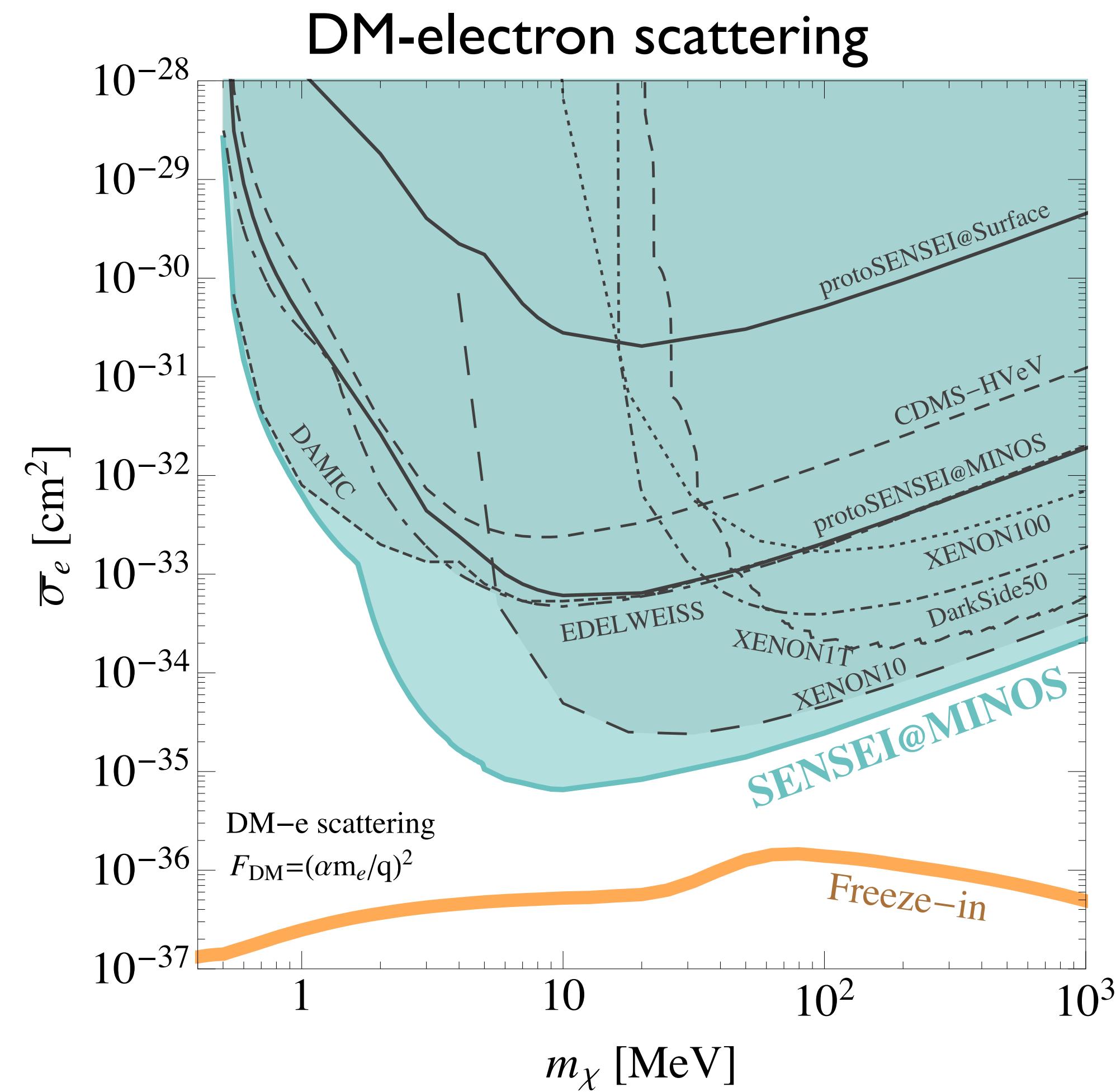
$$E_{\text{ER}} \lesssim \frac{1}{2} m_\chi v^2 \approx 1 \text{ eV} \left[\frac{m_\chi}{0.5 \text{ MeV}} \right]$$

Current targets

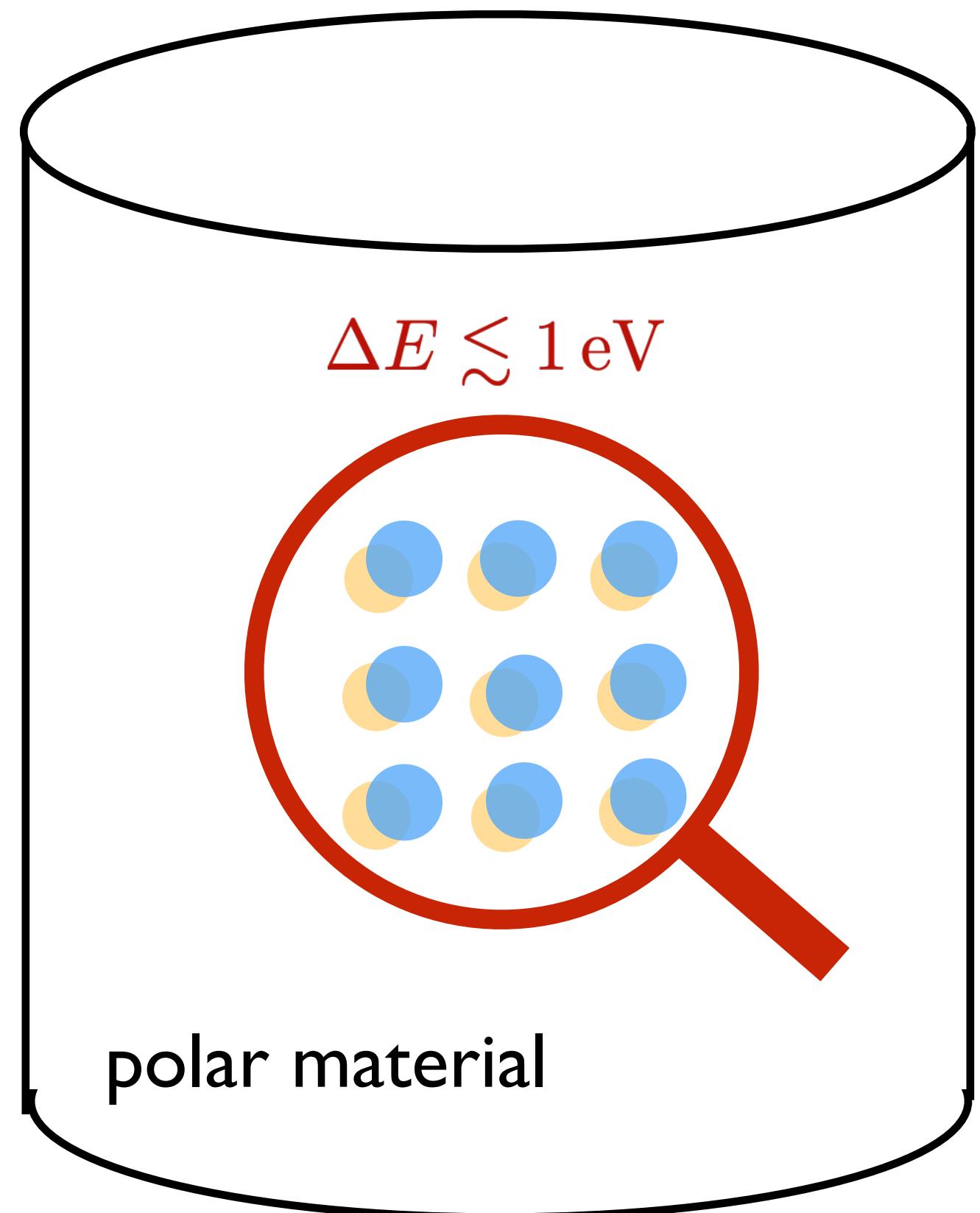
Target	Signal	Threshold	DM Mass range
Noble Liquid	electron ionization	$\sim 10 \text{ eV}$ (atom ionization)	$> 10 \text{ MeV}$
Semiconductors	eh pairs	$\sim 1 \text{ eV}$ (bandgap)	$> \text{MeV}$

Direct detection of sub-GeV DM

Figure from SENSEI, 2020

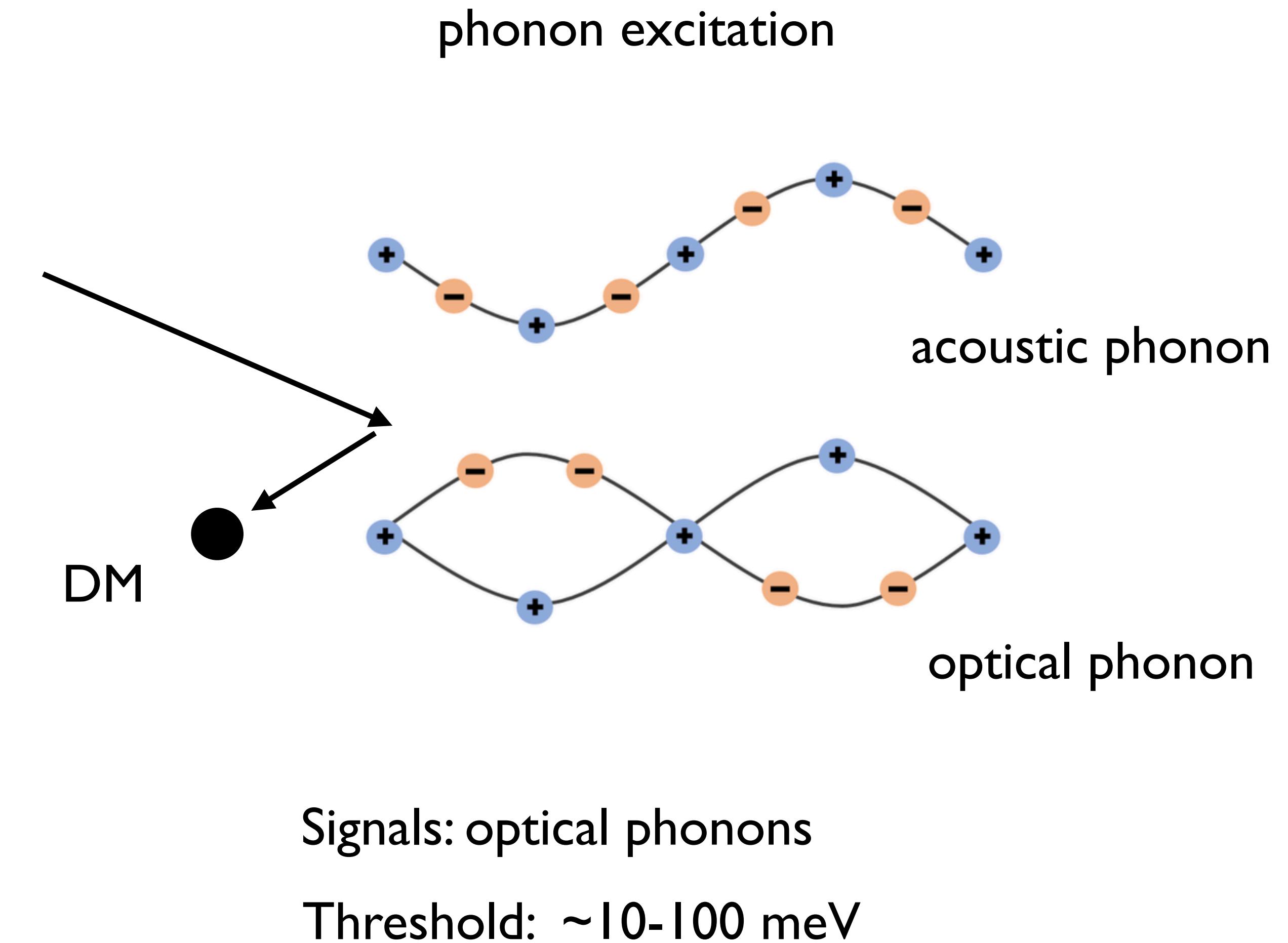
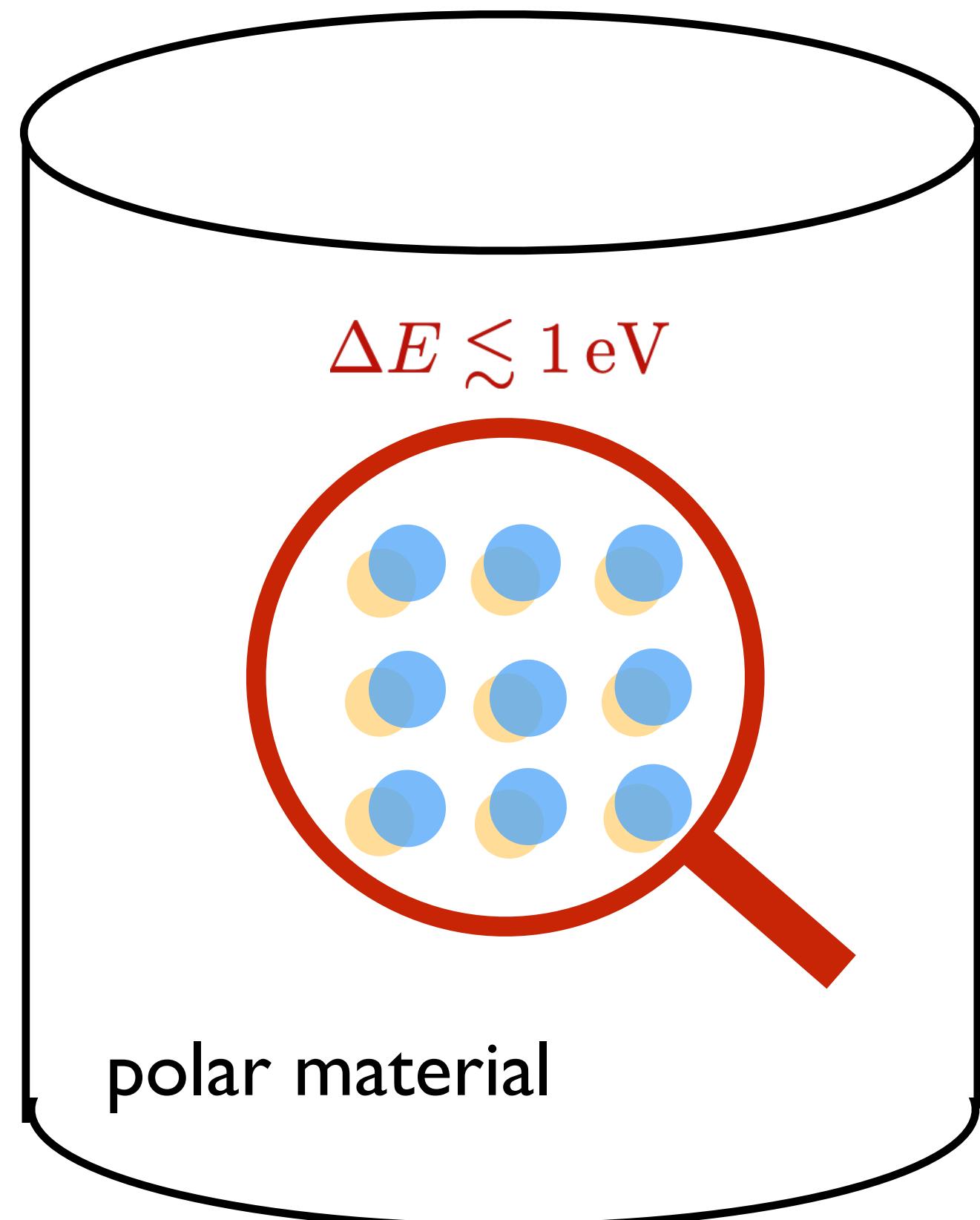


Direct Detection: $\Delta E < 1 \text{ eV}$



Direct Detection: $\Delta E < 1 \text{ eV}$

Knapen, Lin, Pyle, Zurek, 2017

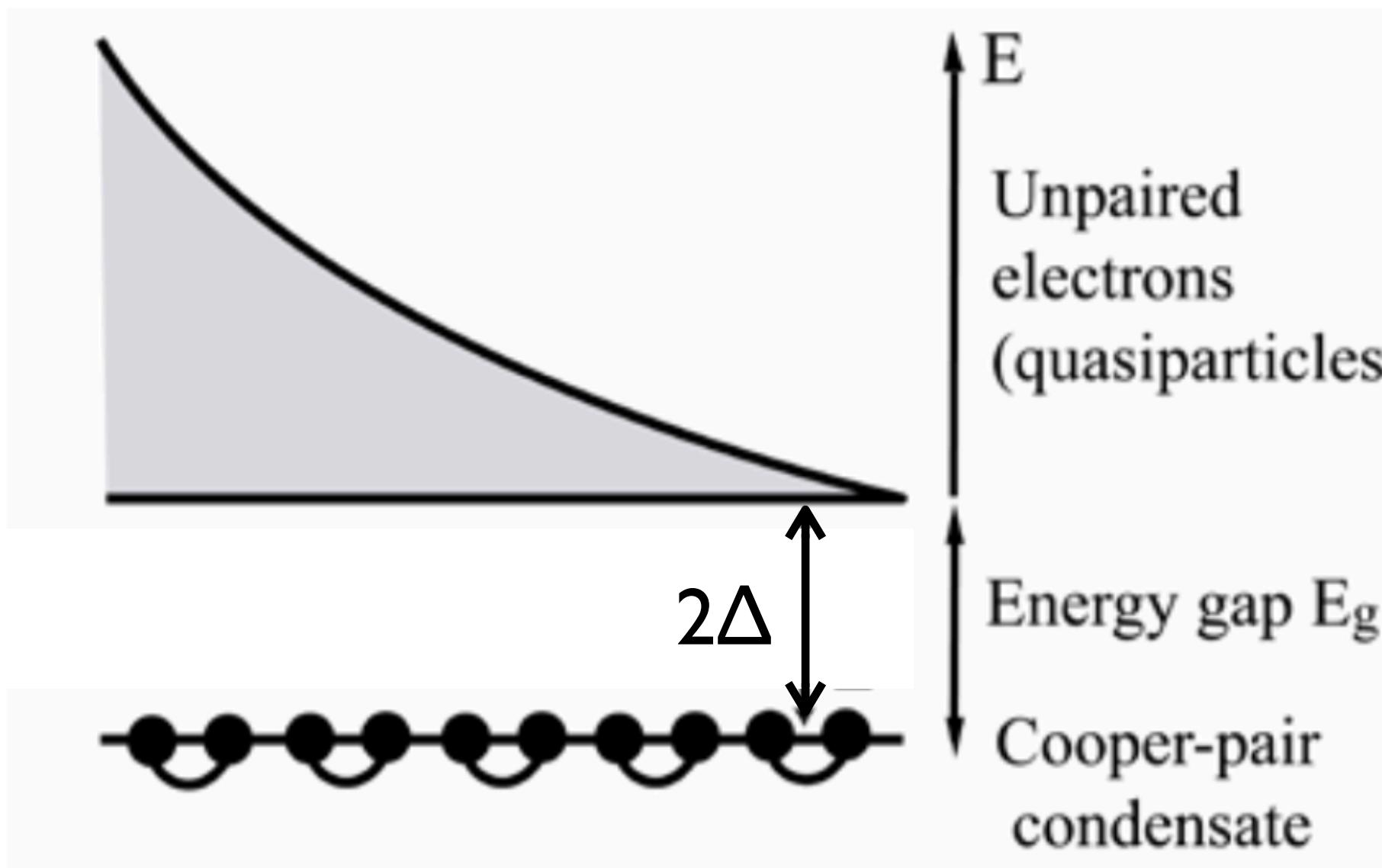


Direct Detection: $\Delta E < 1 \text{ eV}$

Hochberg, Zhao, Zurek, 2015

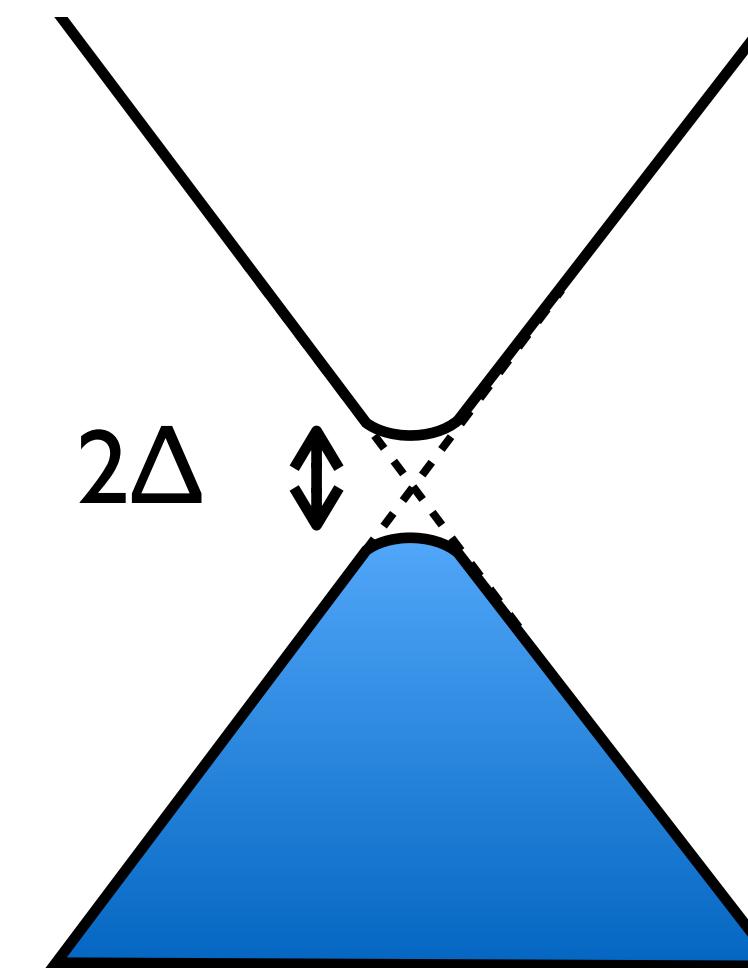
Hochberg, Kahn, Lisanti, Zurek, et.al, 2017

Superconductor



$$\Delta = \mathcal{O}(1) \text{ meV}$$

Dirac material



$$\Delta = \mathcal{O}(1) \text{ meV}$$

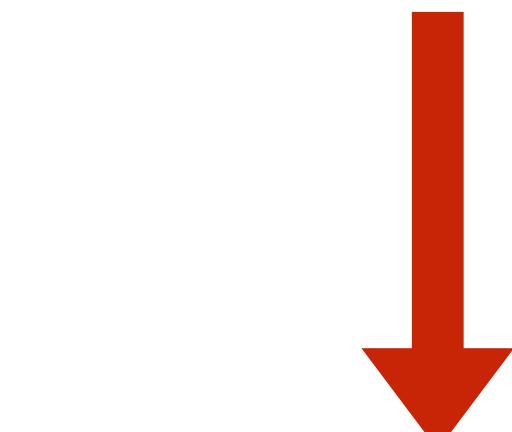
Signals: quasiparticles/phonons

Threshold: $\sim 1 \text{ meV}$

Probing sub-MeV (sub-eV) DM

Hochberg, Zhao, Zurek, 2015
 Schutz, Zurek, 2016
 Knapen, Lin, Pyle, Zurek, 2017
 Hochberg, Kahn, Lisanti, Zurek, et.al, 2017
 Bunting, Gratta, Melia, Rajendran, 2017
 D. M. Mei, et.al. 2017
 ...

Target	Signal	Threshold	DM Mass range
Nobel Liquid	electron ionization	~10 eV (atom ionization)	>10 MeV
Semiconductors	eh pairs	~1 eV (bandgap)	>MeV
Polar materials	phonon	10-100meV	>10-100 keV
Superconductor	phonon/ quasiparticle	~1 meV	>1 keV


**Low threshold
can probe
low DM masses**

Dirac materials, superfluid helium, magnetic bubble chamber, Ge detector with charge amplification ...

Probing sub-MeV (sub-eV) DM

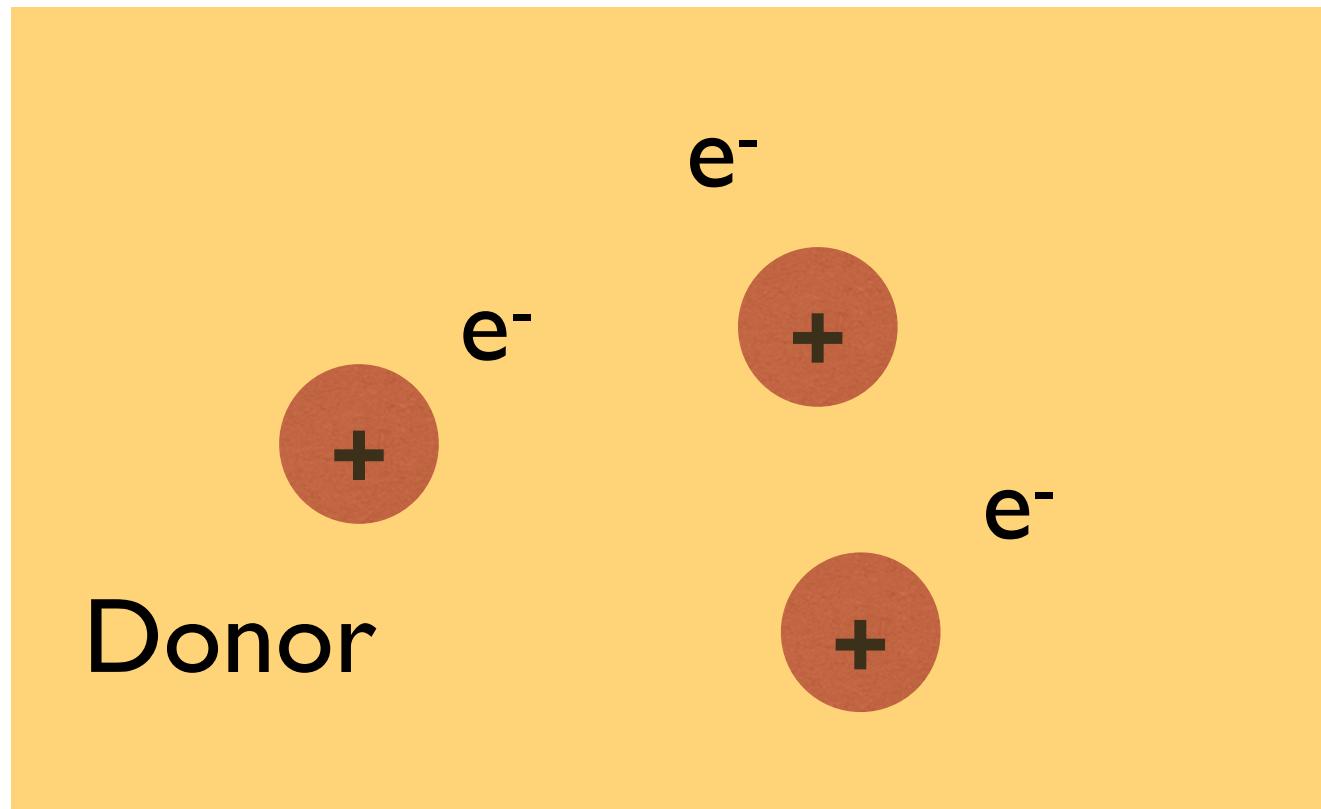
Target	Signal	Threshold	DM Mass range
Nobel Liquid	electron ionization	~10 eV (atom ionization)	>10 MeV
Semiconductors	eh pairs	~1eV (bandgap)	>MeV
Polar materials	phonon	10-100meV	>10-100 keV
Doped Semiconductors	phonon/ electron ionization/ eh pairs	10-100meV	>10-100 keV
Superconductor	phonon/ quasiparticle	~1meV	>1keV

Doped semiconductors

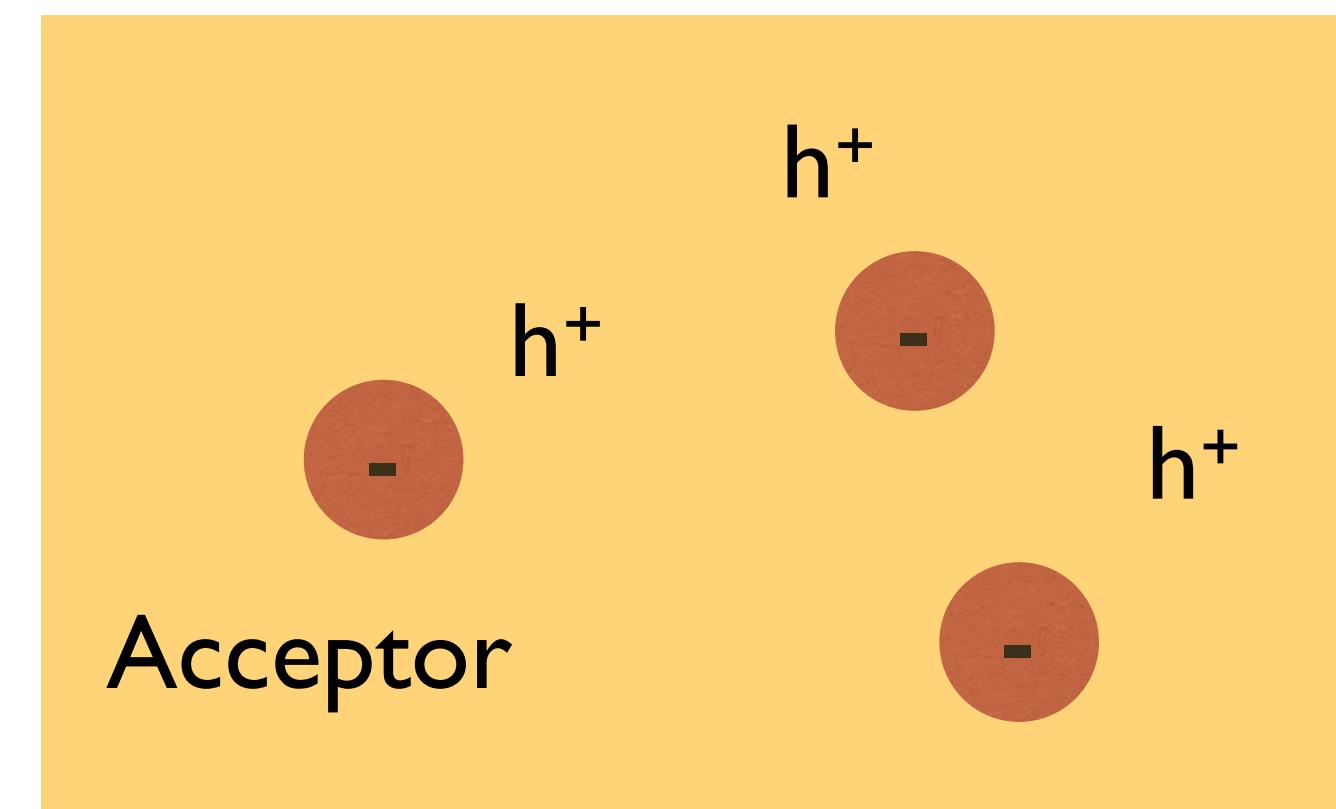


Doped semiconductors

n-type semiconductor



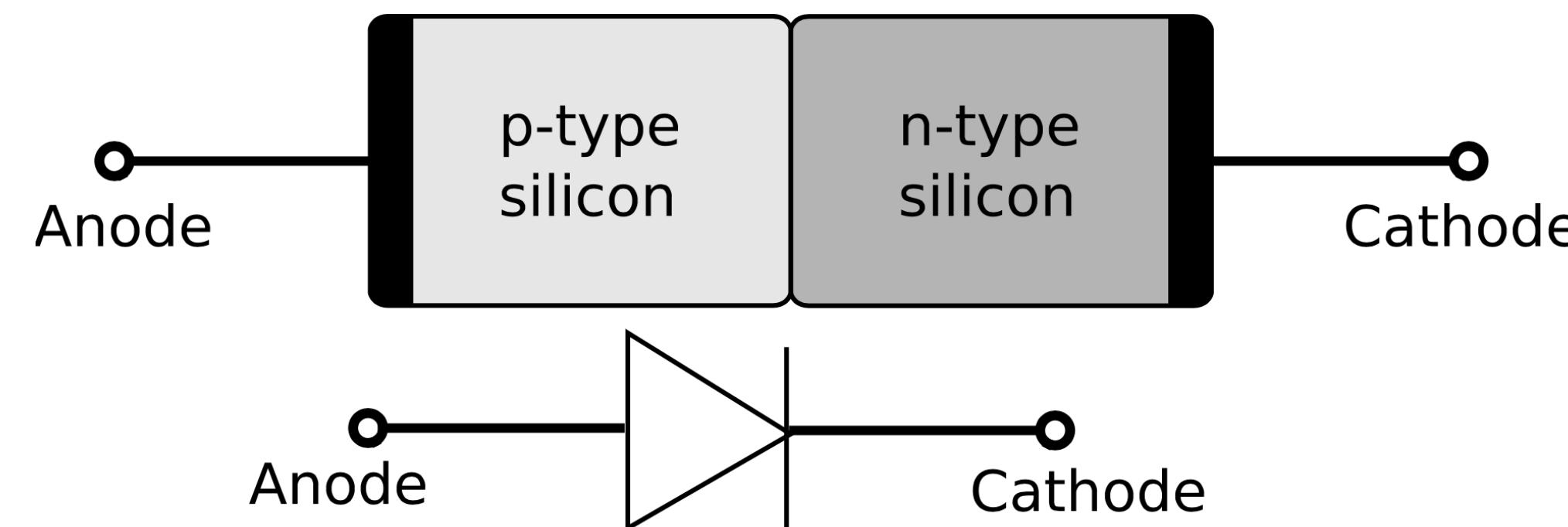
p-type semiconductor



Donors in Silicon: P ,As ... (group V elements)

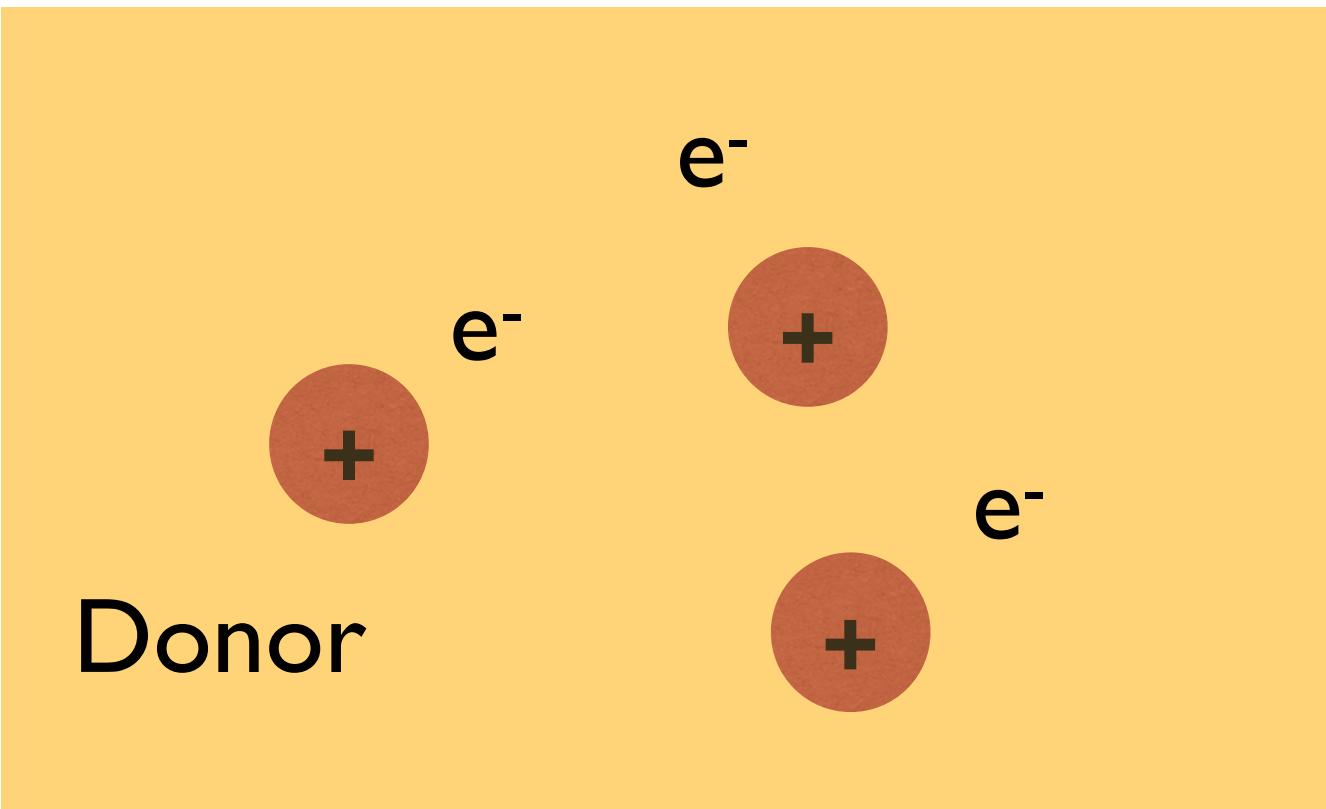
Commonly used: p-n junction, diodes

Acceptors in Silicon: B ,Al ... (group III elements)

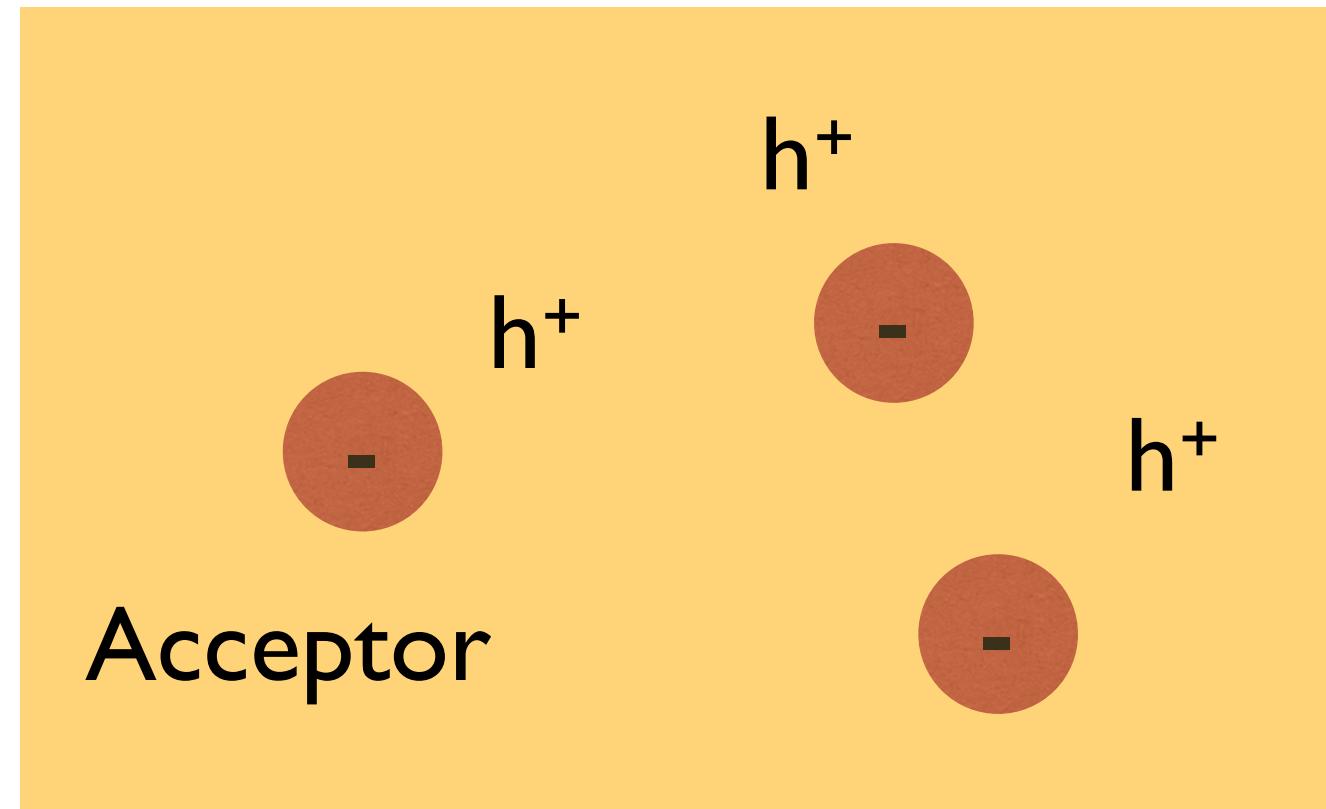


Doped semiconductors

n-type semiconductor

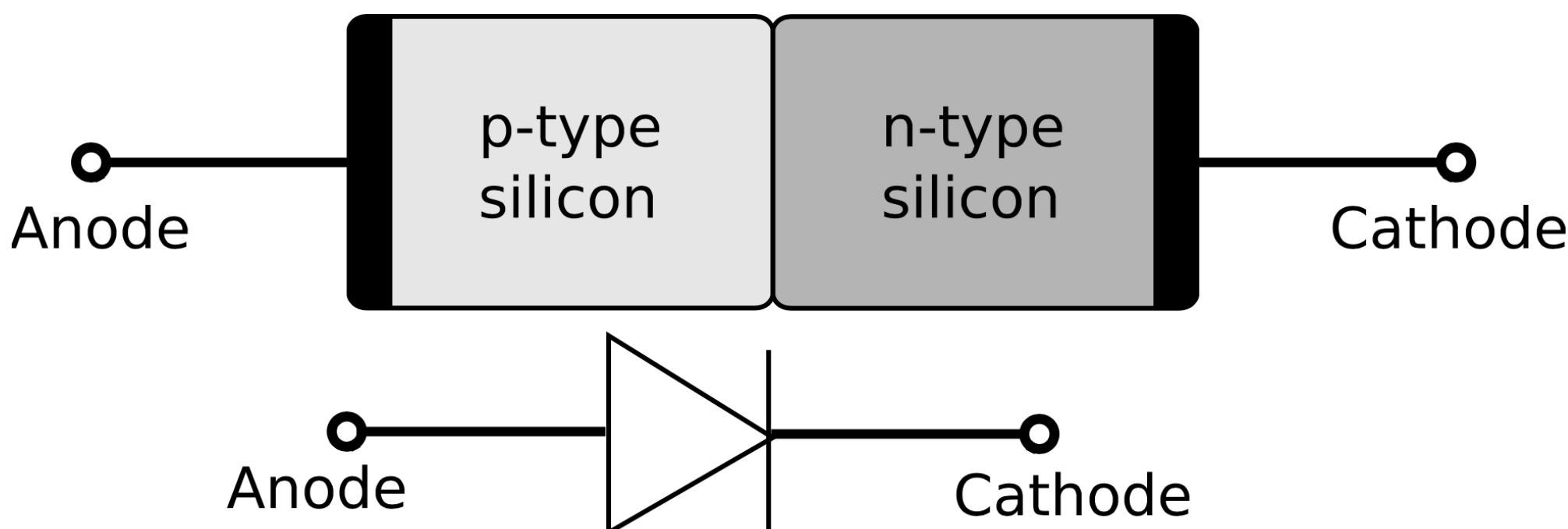


p-type semiconductor



Donors in Silicon: P ,As ... (group V elements)

Commonly used: p-n junction, diodes

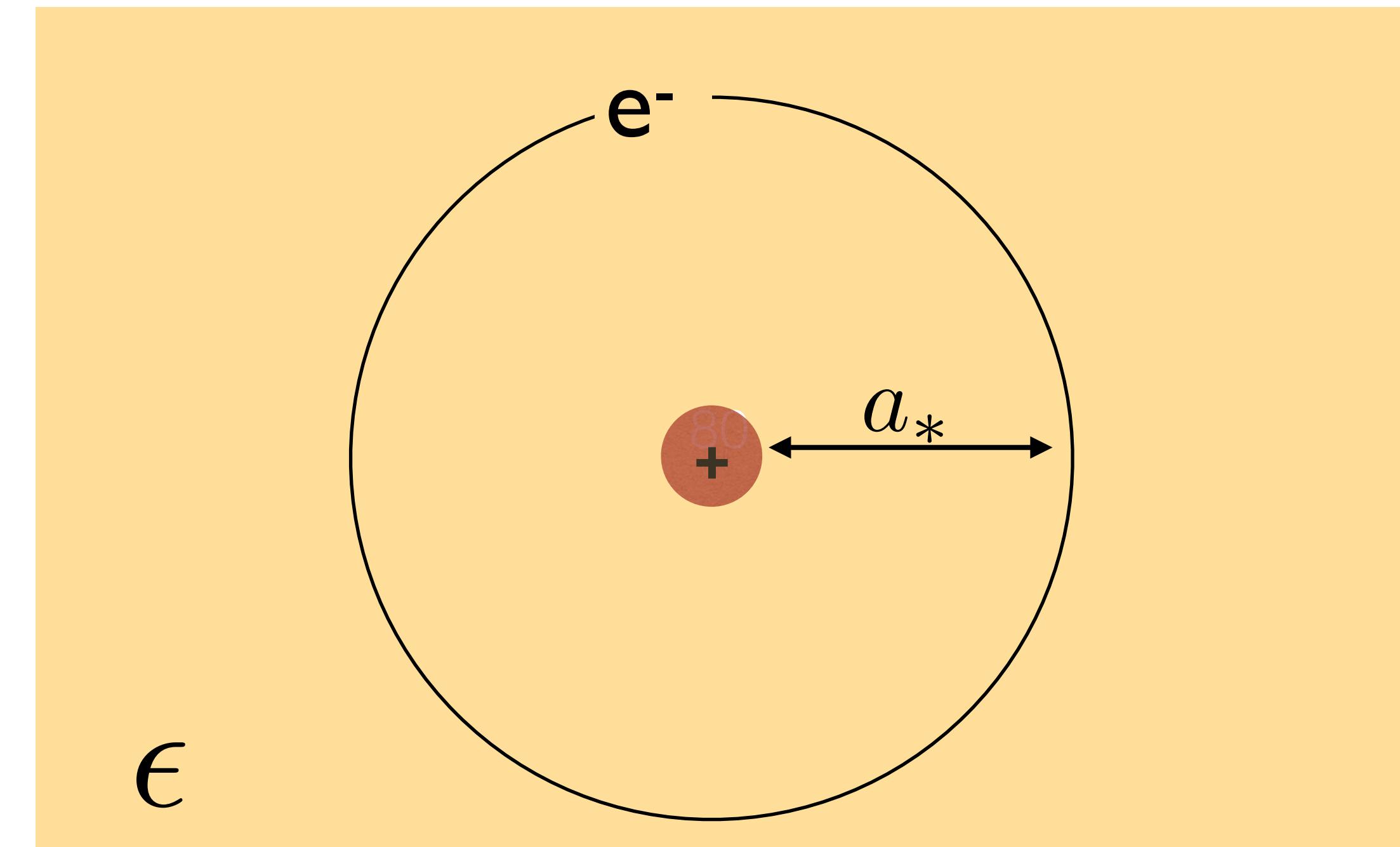
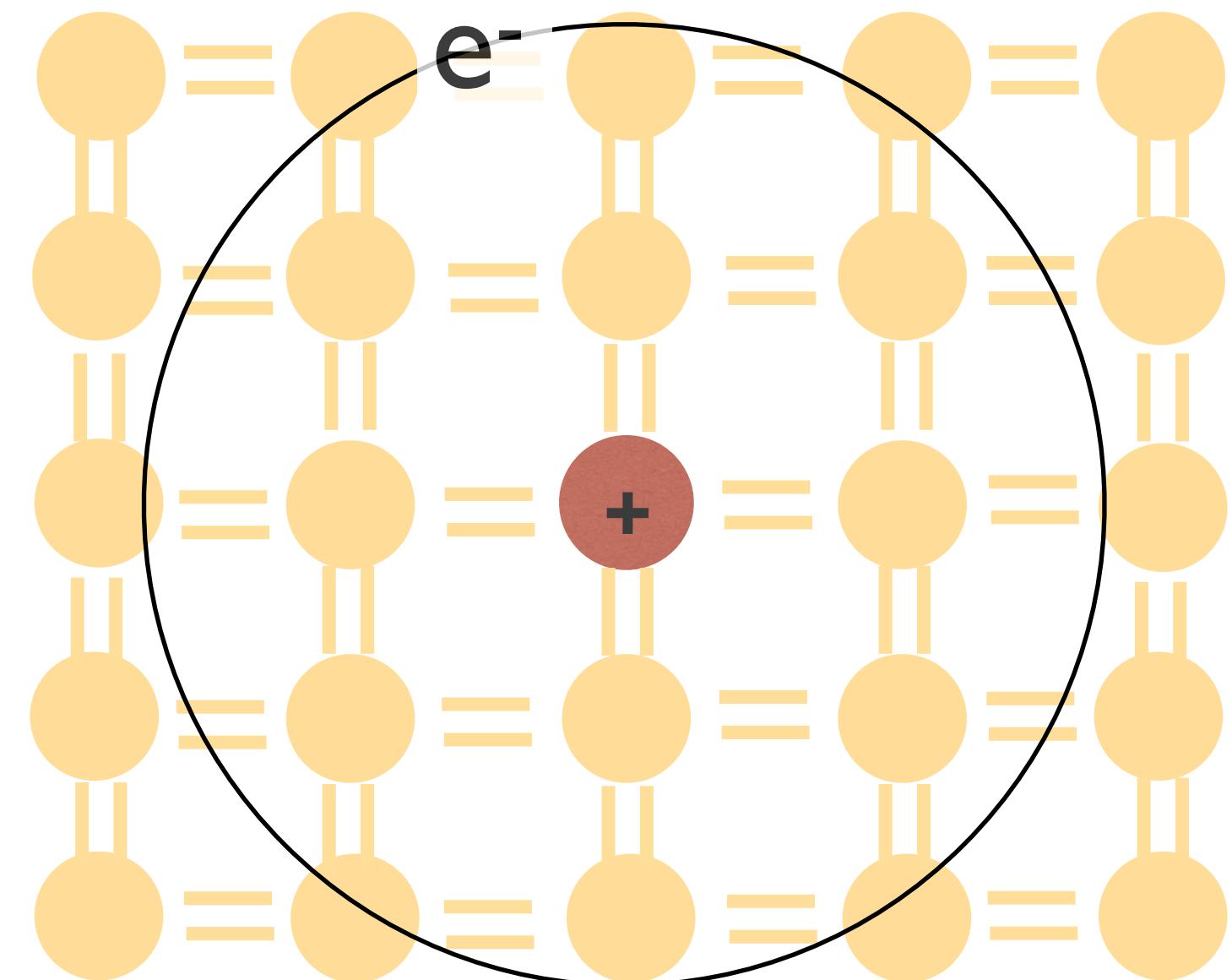


Acceptors in Silicon: B ,Al ... (group III elements)



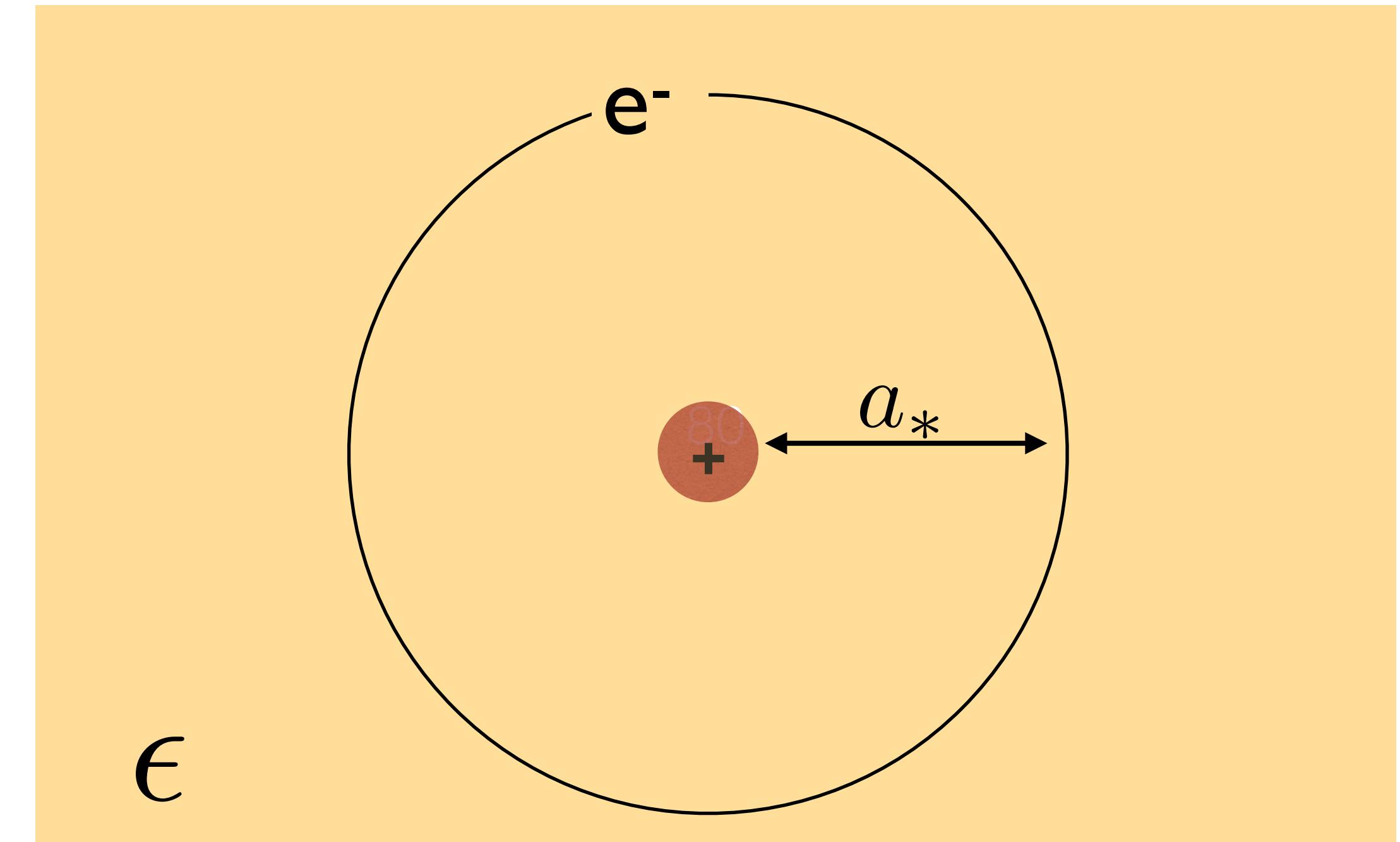
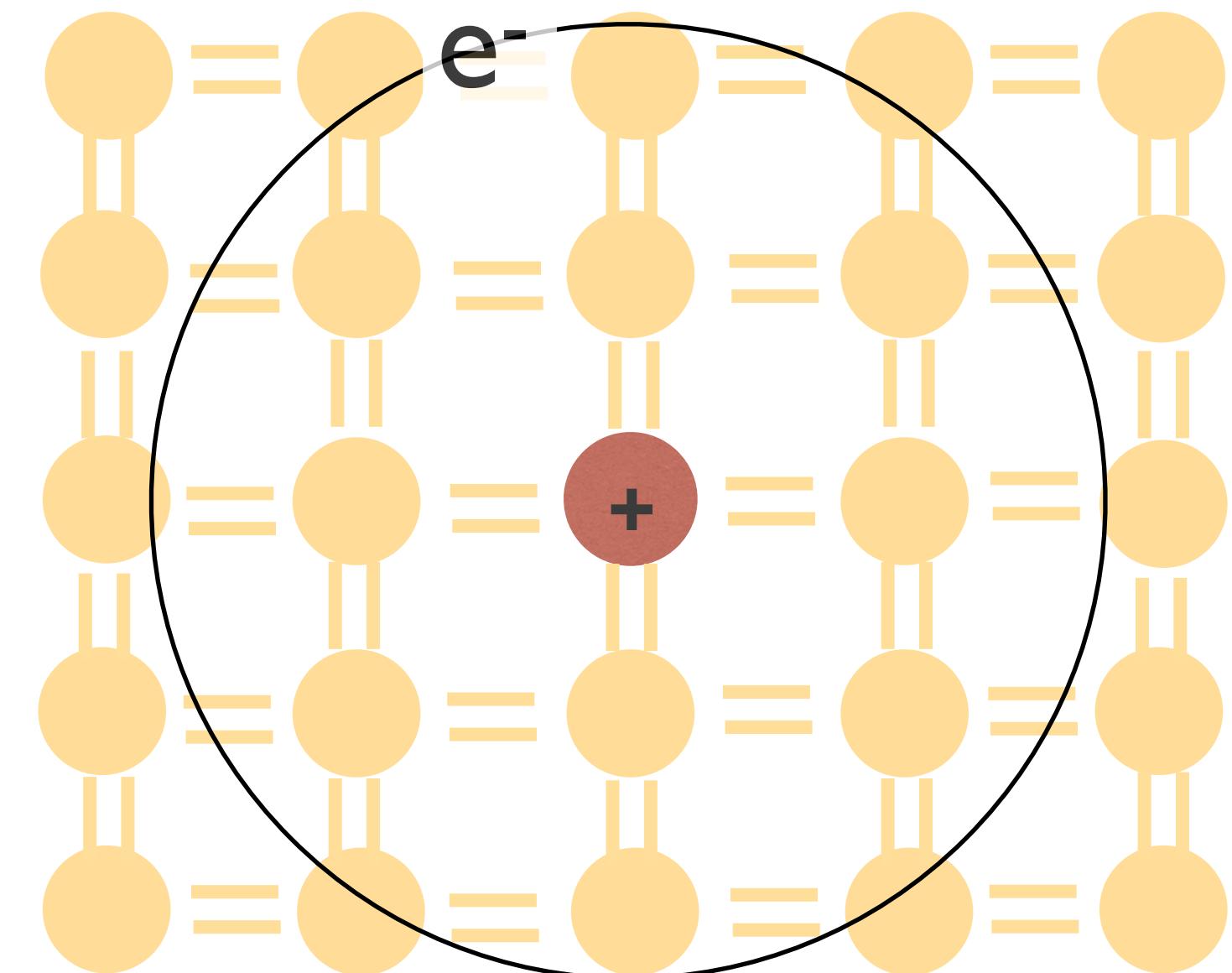
Dopants in semiconductors

Dopants: “Hydrogen atoms” in a background with a large dielectric constant



Dopants in semiconductors

Dopants: “Hydrogen atoms” in a background with a large dielectric constant

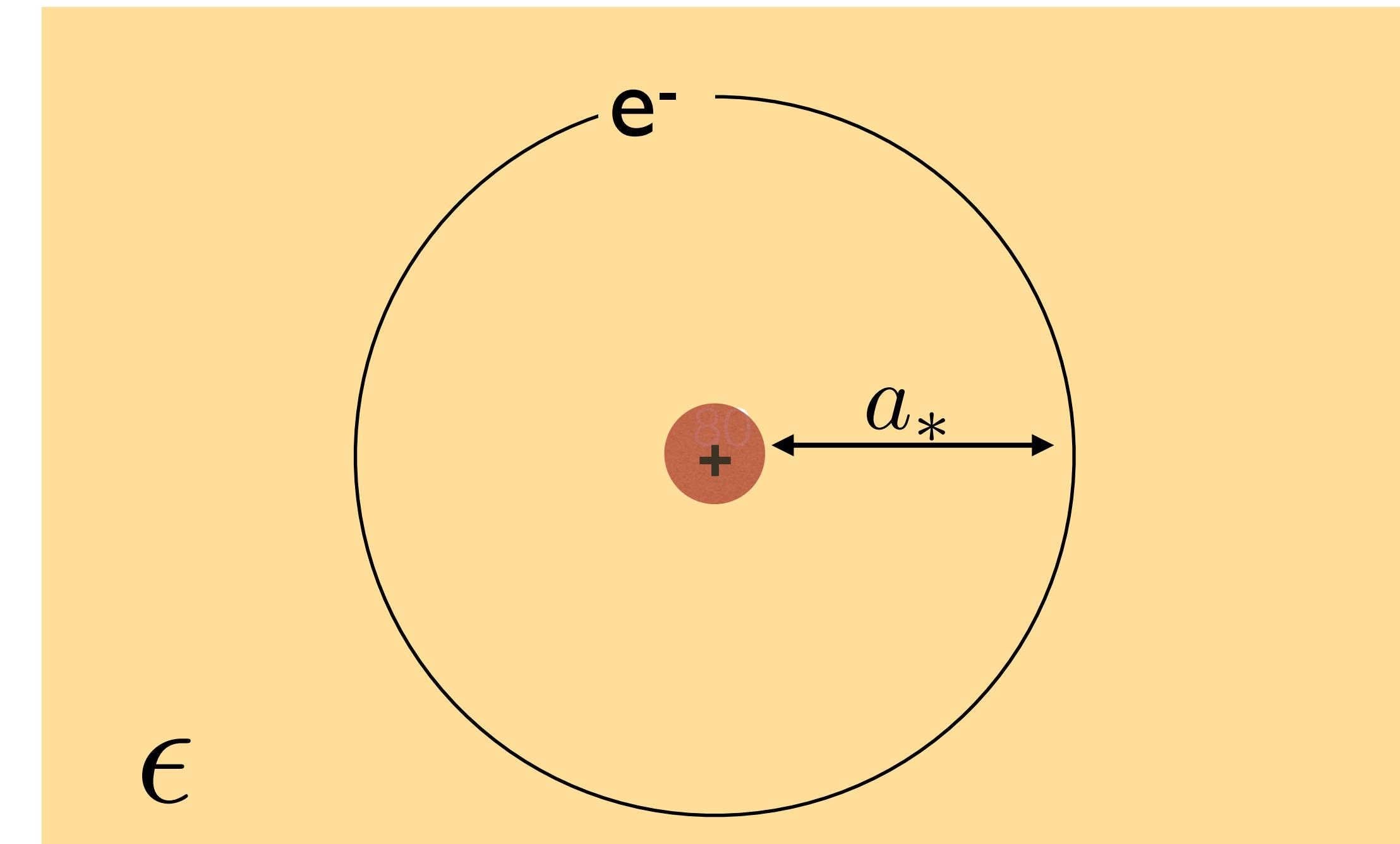
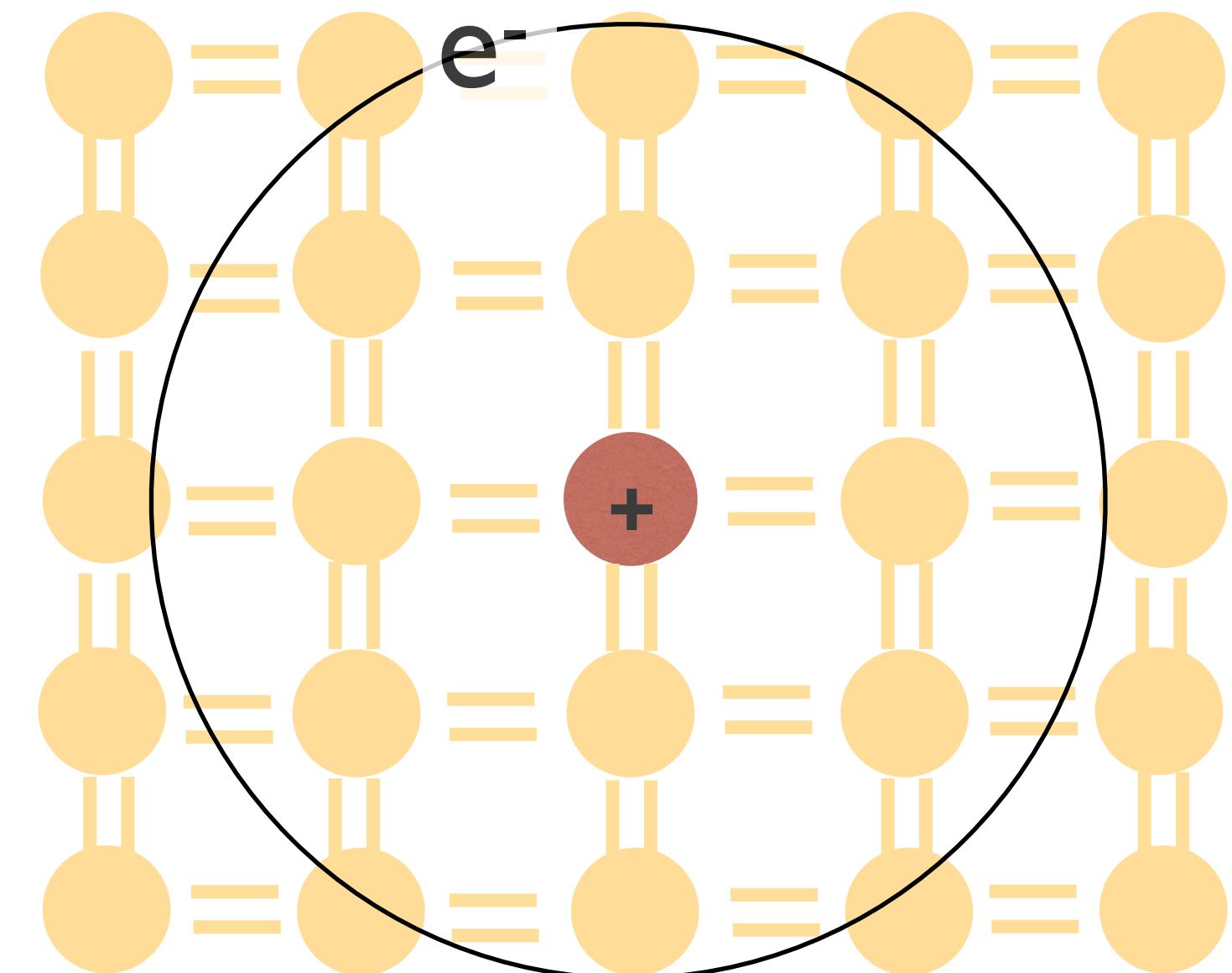


For $\epsilon \sim 10$ electron effective mass $a_* \sim (\frac{\alpha}{\epsilon} m_*)^{-1} \sim O(10) a_0$ Bohr radius $q_* \sim a_*^{-1} \sim O(100) \text{ eV}$ $v_* = \frac{q_*}{m_*} \sim 10^{-3}$

$$E_{\text{ionization}} \sim \frac{1}{2} \left(\frac{\alpha}{\epsilon} \right)^2 m_* \sim 10 - 100 \text{ meV}$$

Dopants in semiconductors

Dopants: “Hydrogen atoms” in a background with a large dielectric constant



For $\epsilon \sim 10$ electron effective mass $a_* \sim \left(\frac{\alpha}{\epsilon} m_*\right)^{-1} \sim O(10) a_0$ Bohr radius $q_* \sim a_*^{-1} \sim O(100) \text{ eV}$

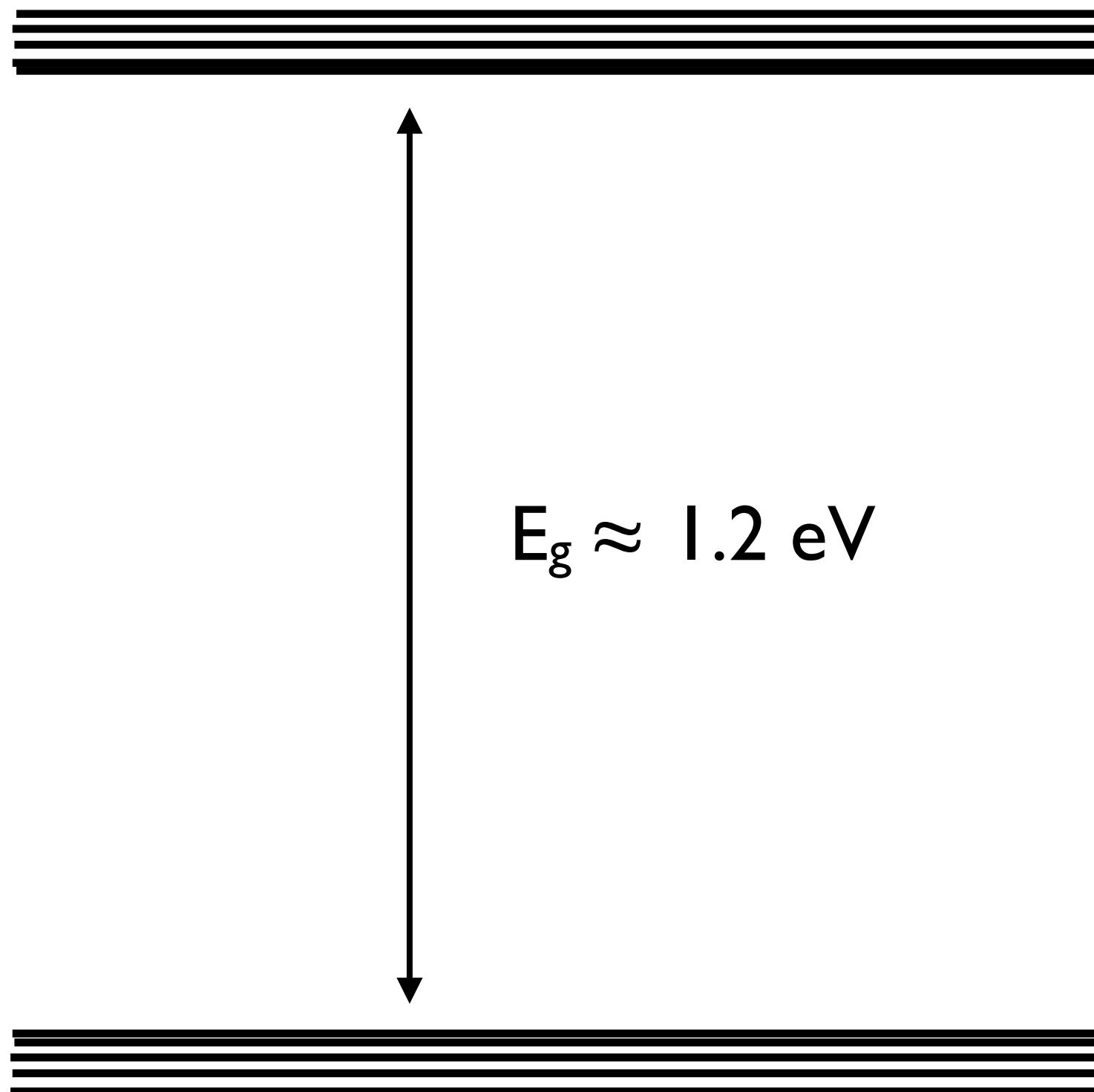
$$E_{\text{ionization}} \sim \frac{1}{2} \left(\frac{\alpha}{\epsilon}\right)^2 m_* \sim 10 - 100 \text{ meV}$$



Dopant energy levels in silicon

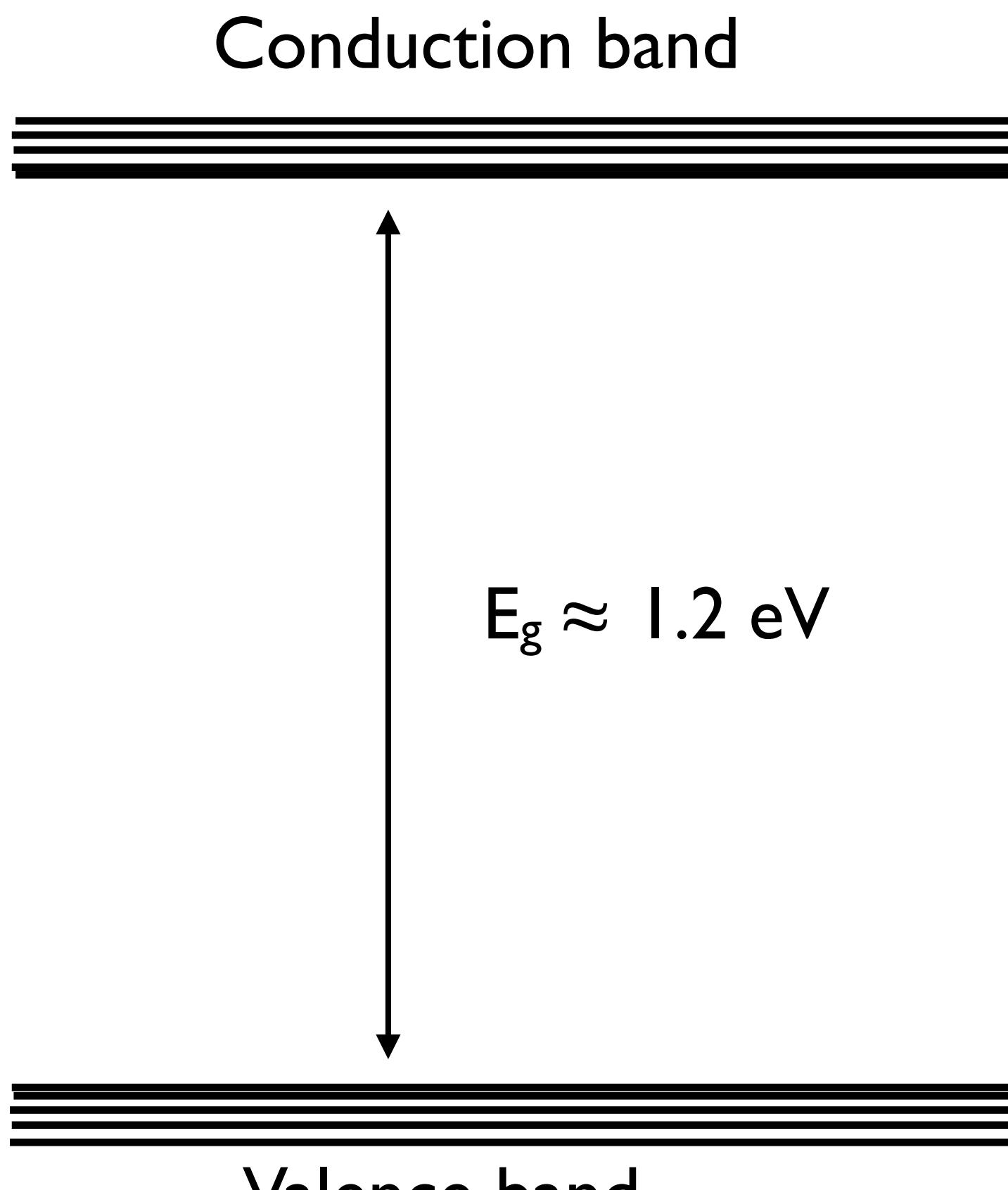
Undoped Si

Conduction band

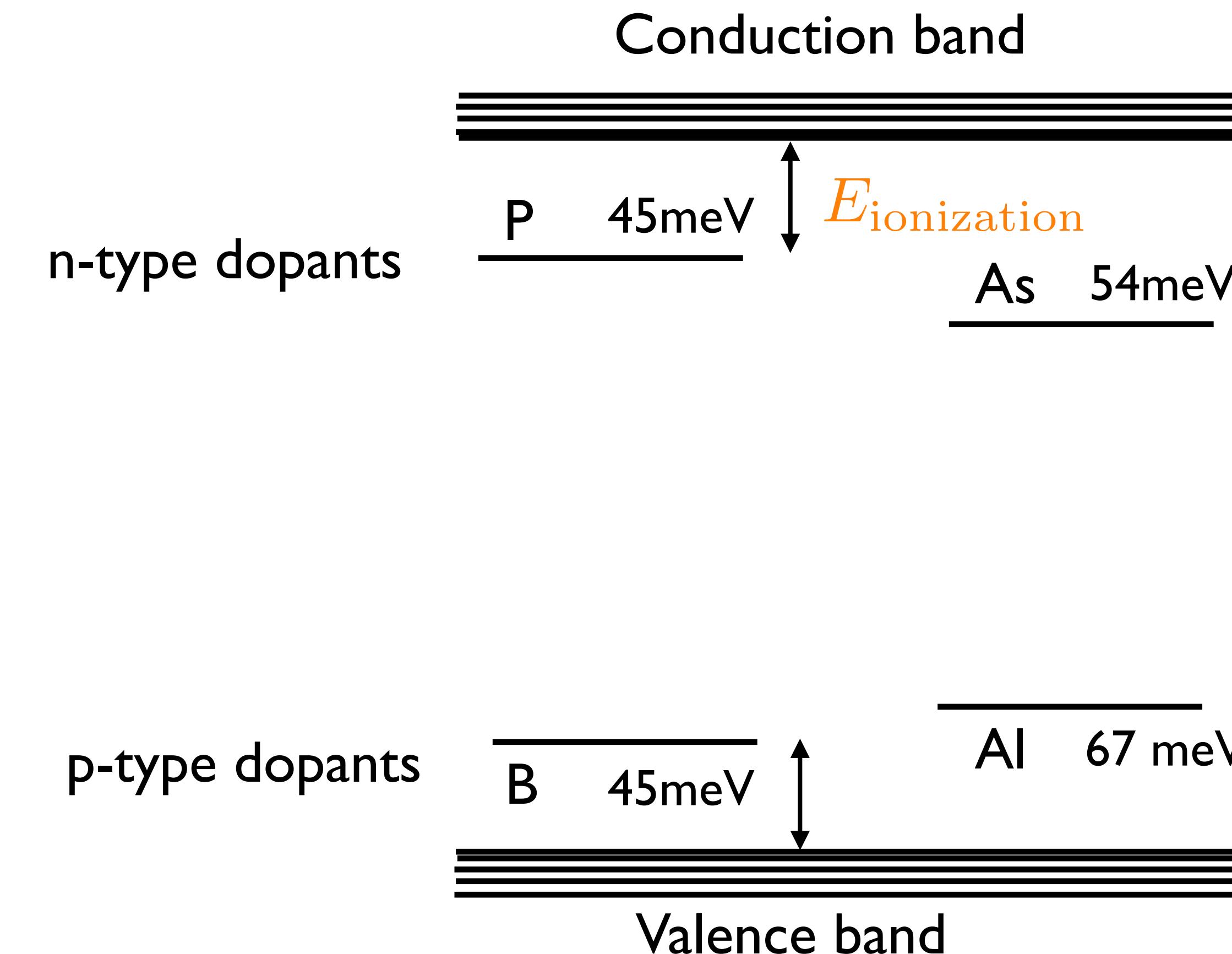


Dopant energy levels in silicon

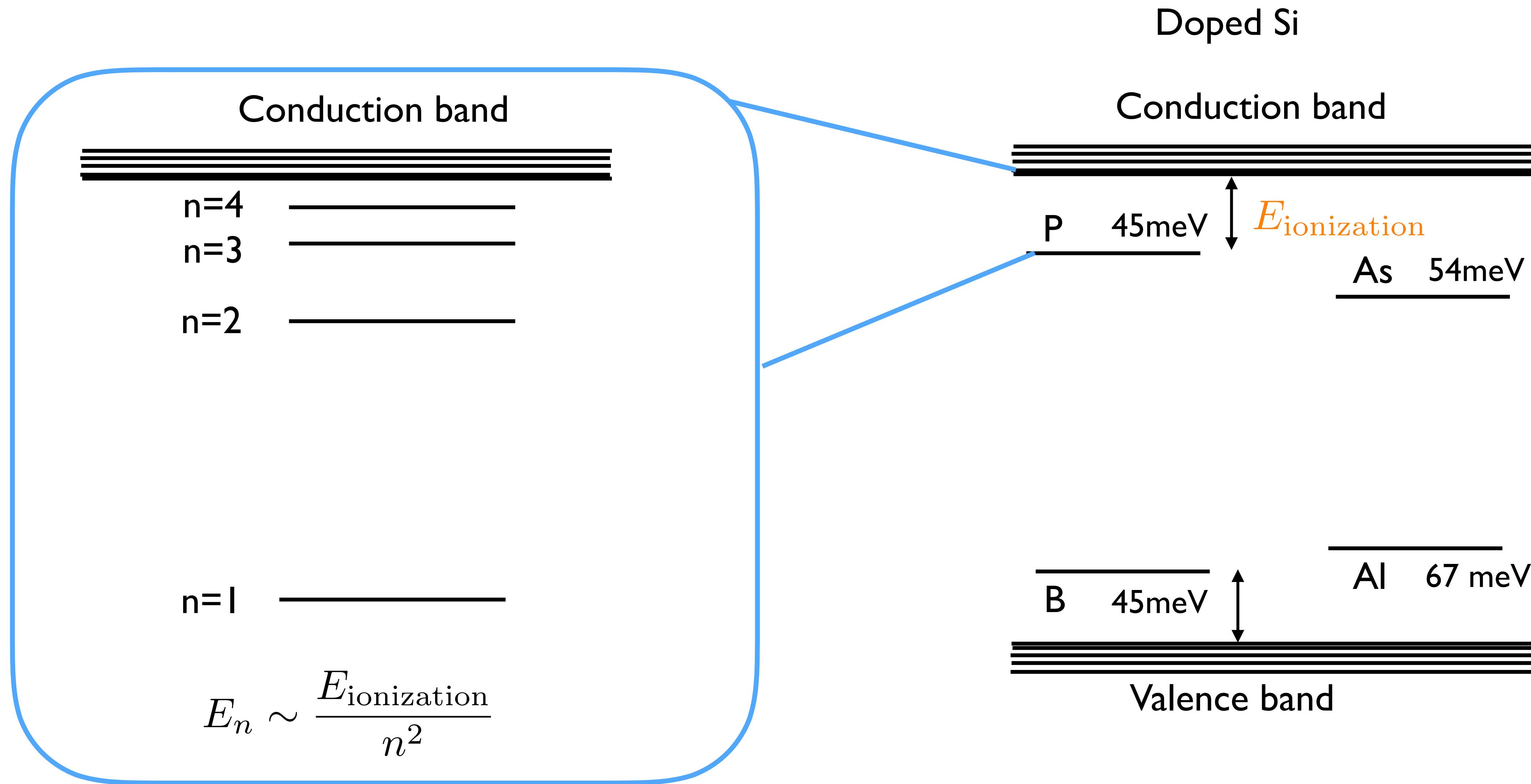
Undoped Si



Doped Si

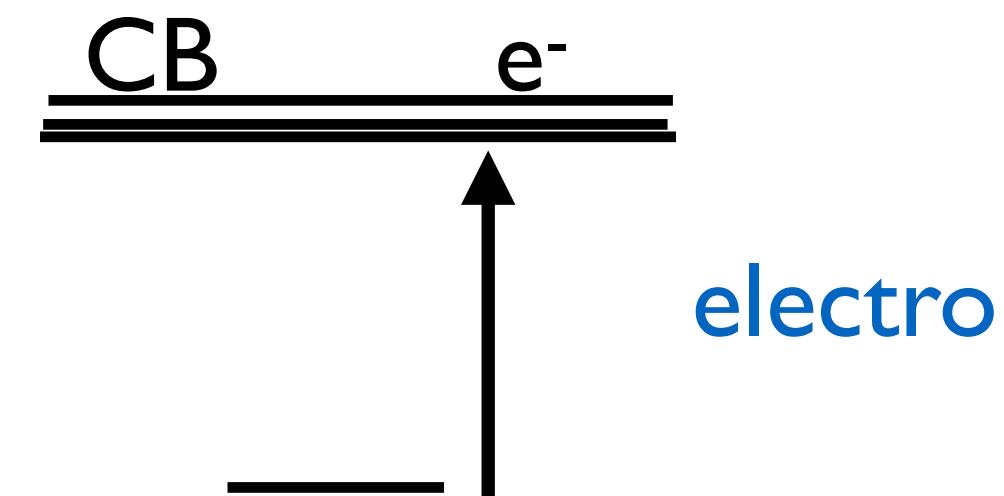
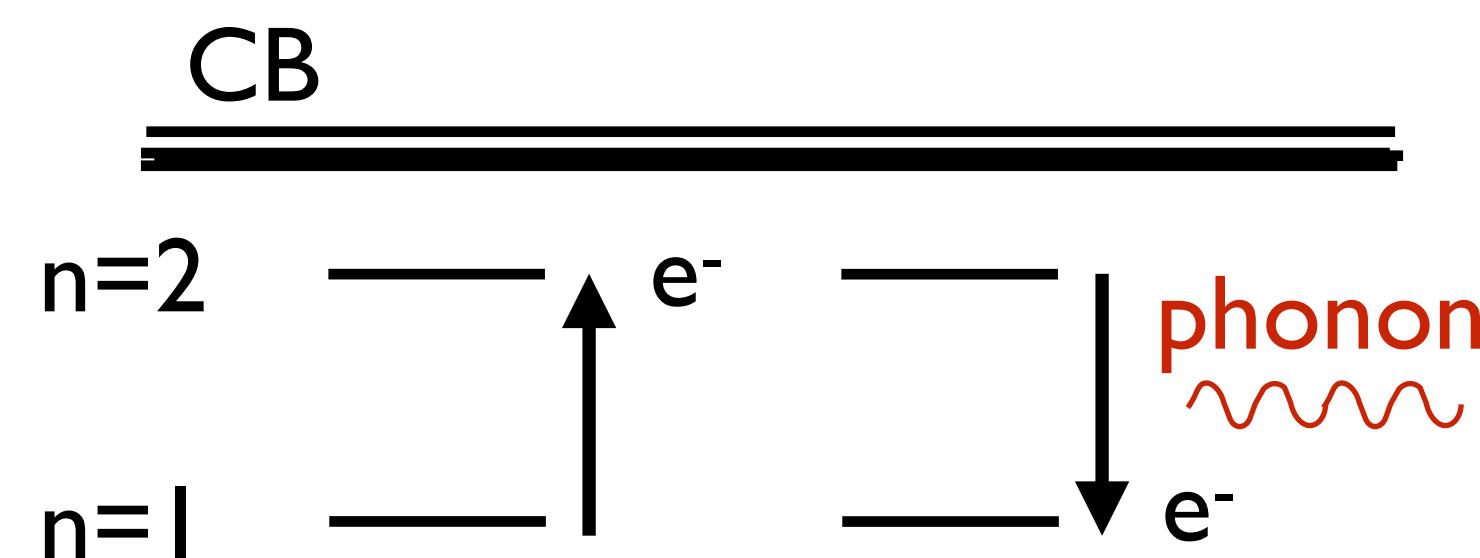
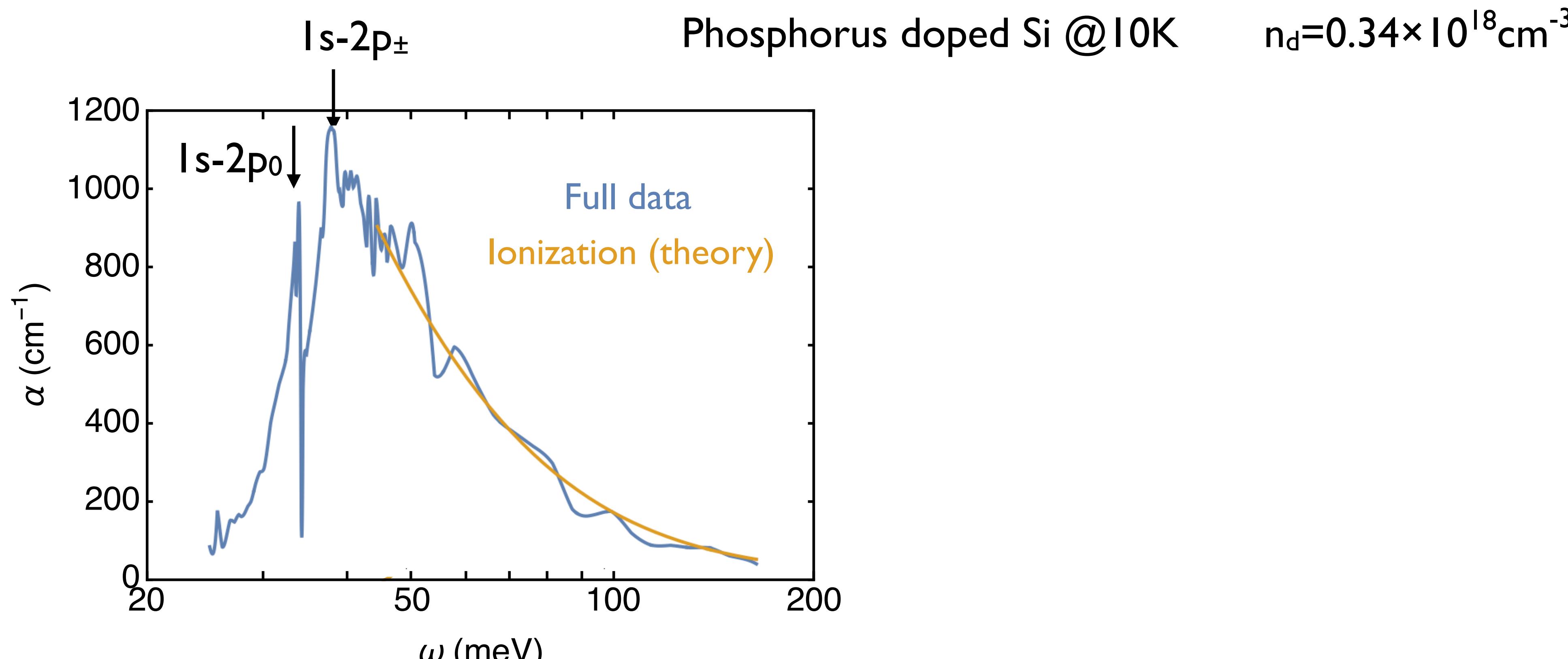


Dopant energy levels in silicon



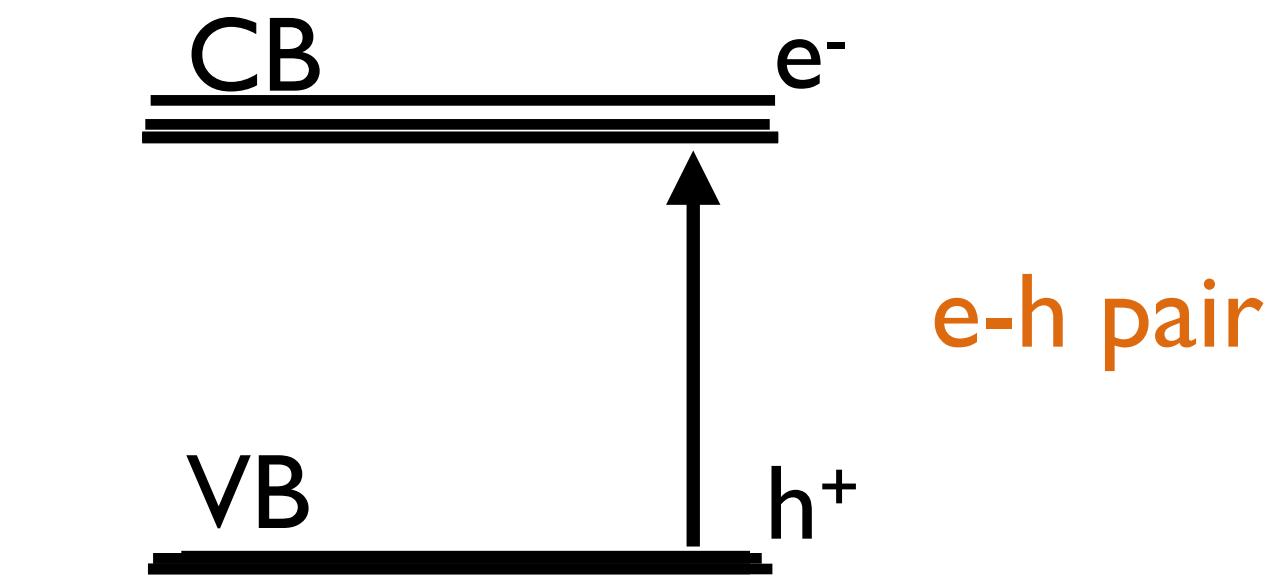
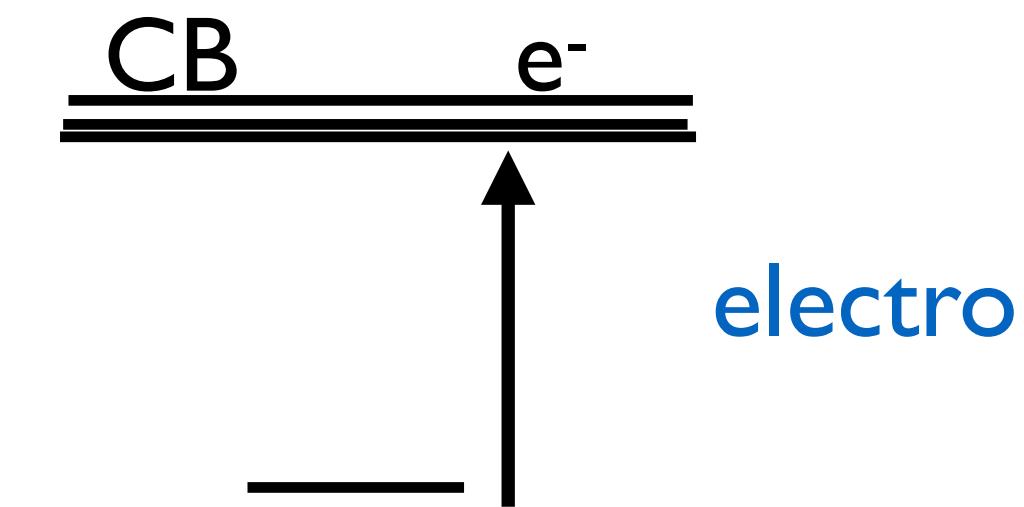
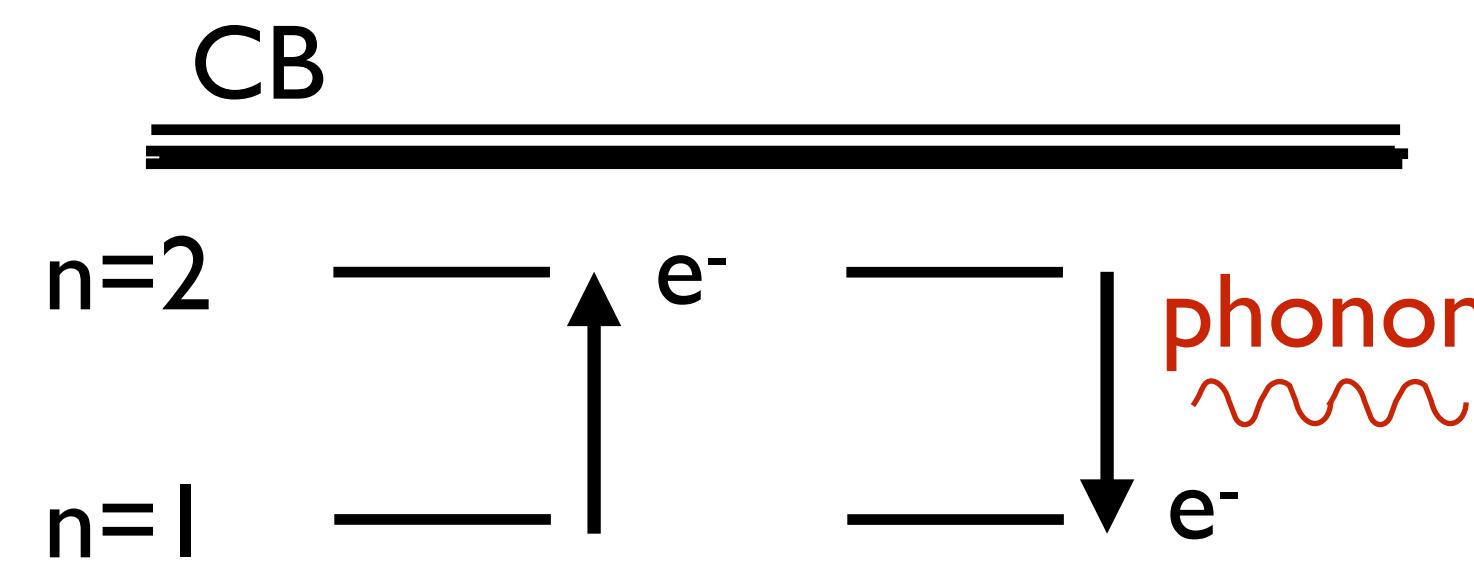
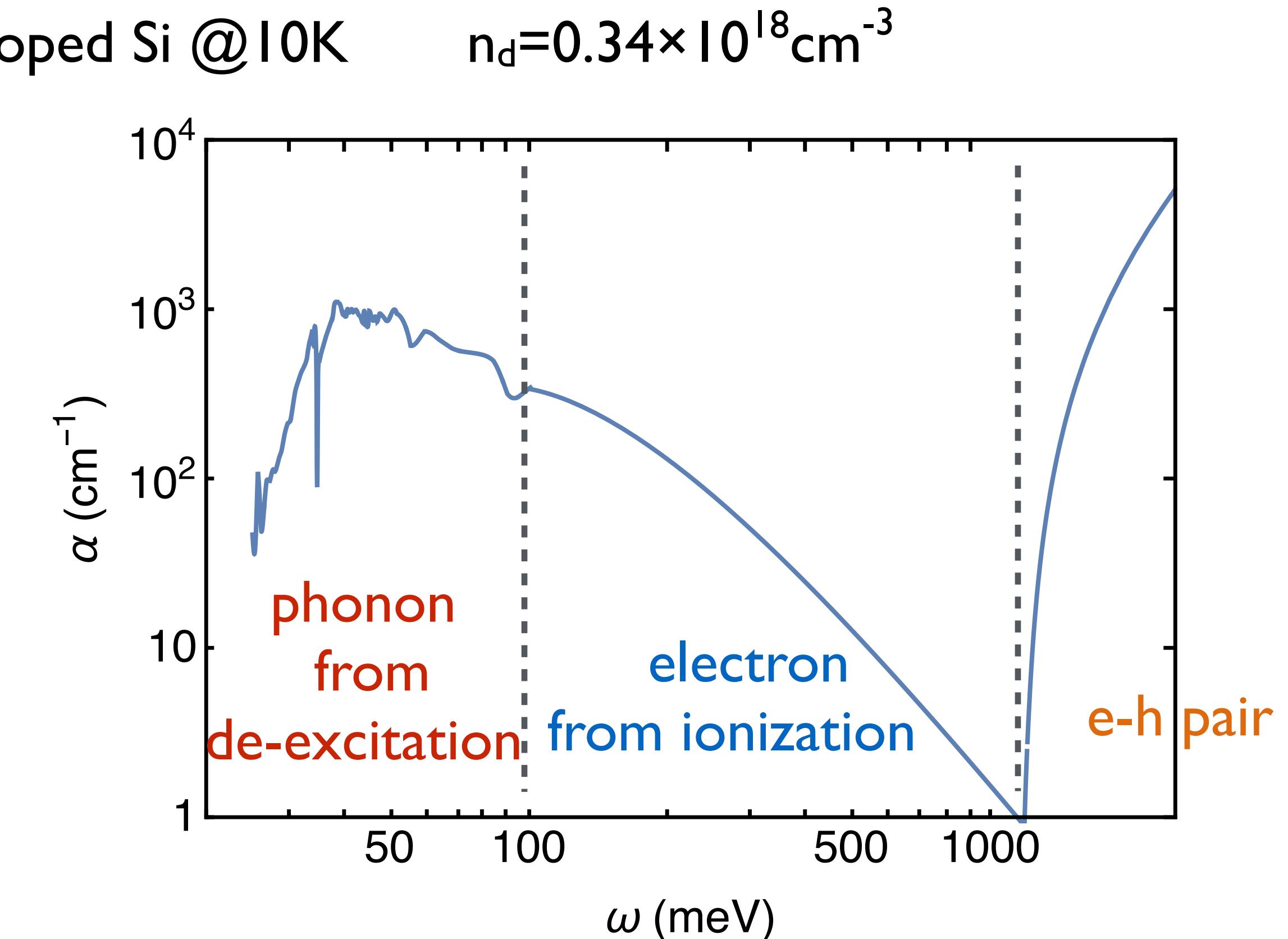
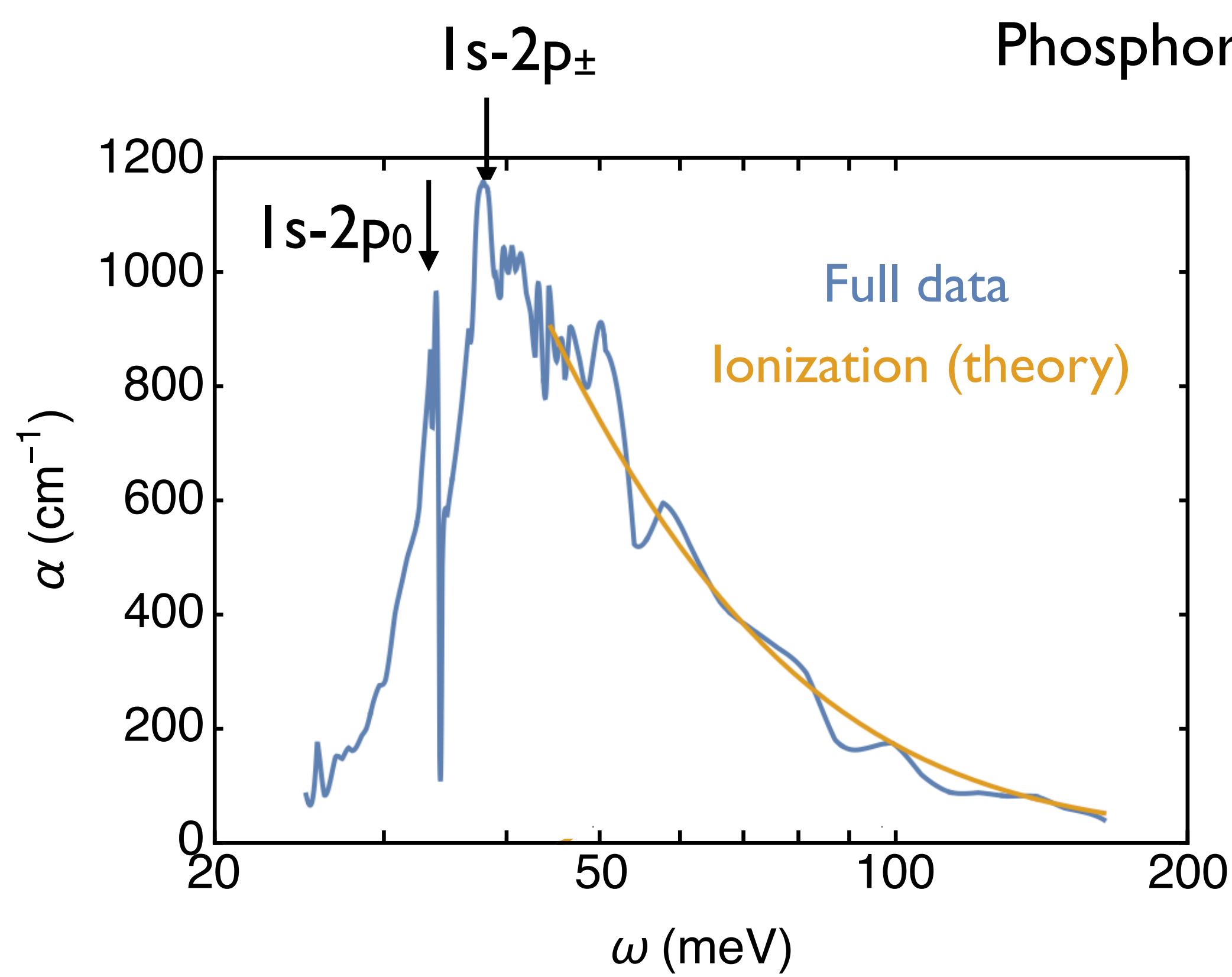
Signals in doped silicon

Gaymann, Geserich, Lohneysen, 95



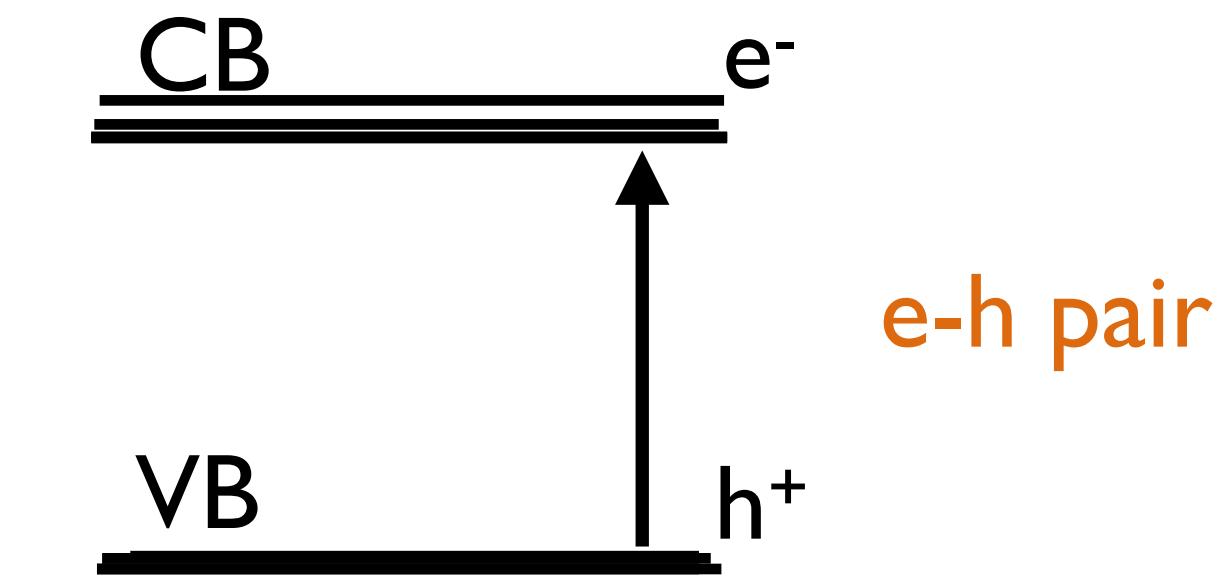
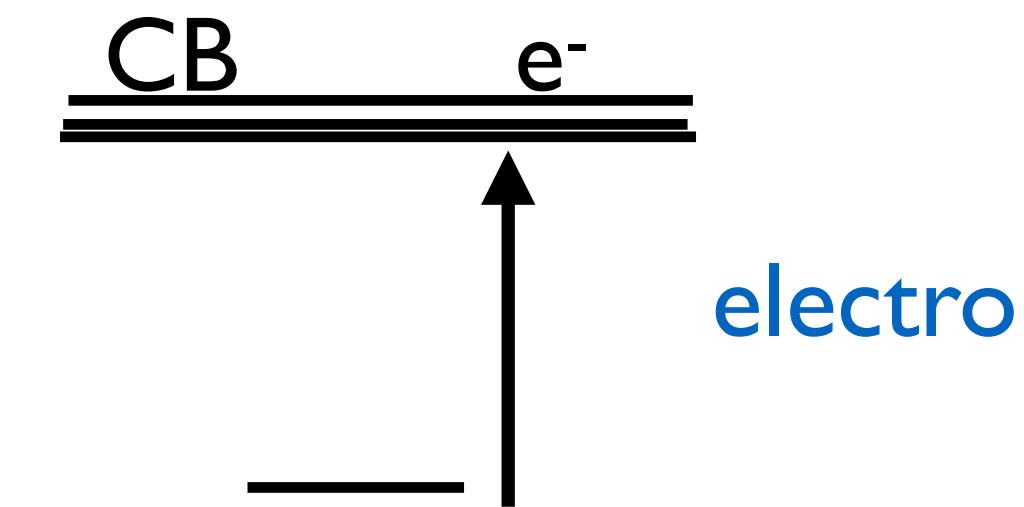
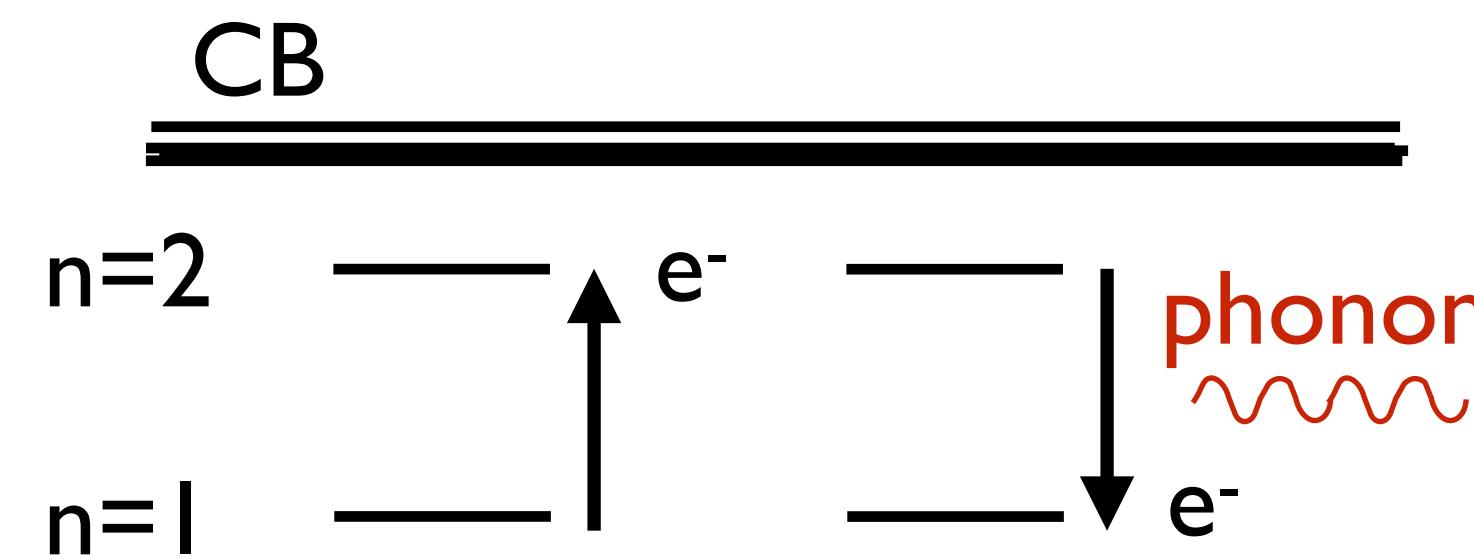
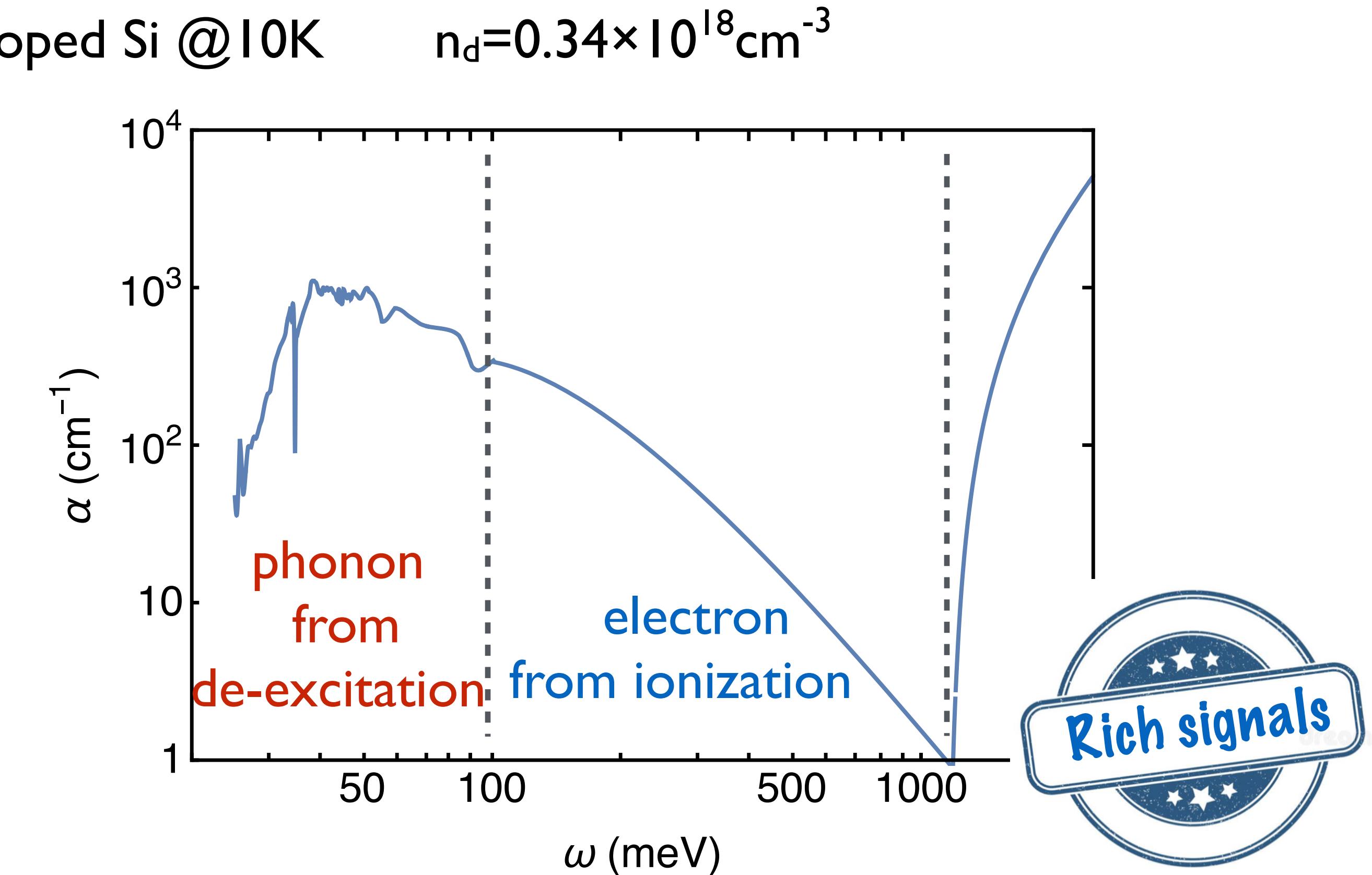
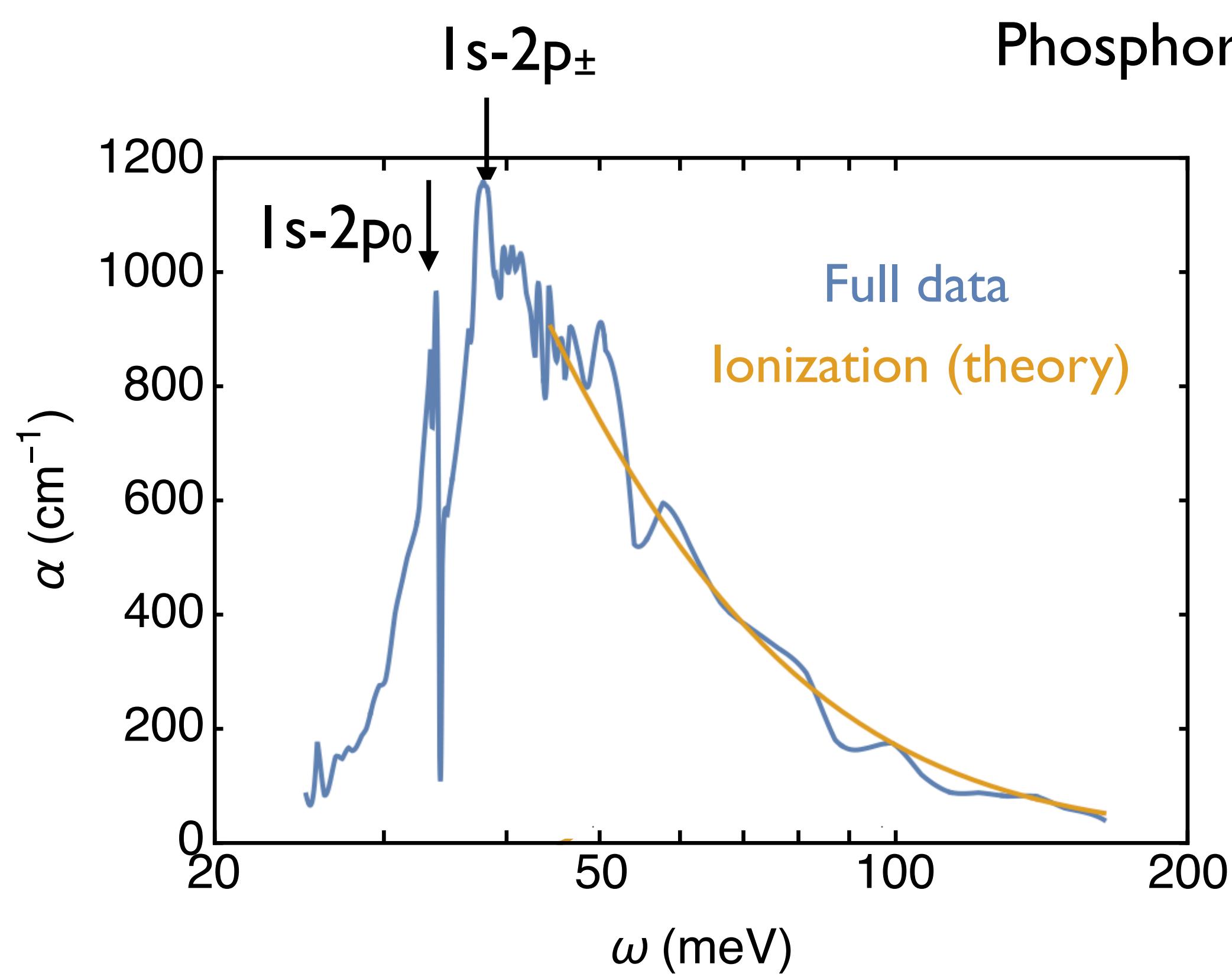
Signals in doped silicon

Gaymann, Geserich, Lohneysen, 95



Signals in doped silicon

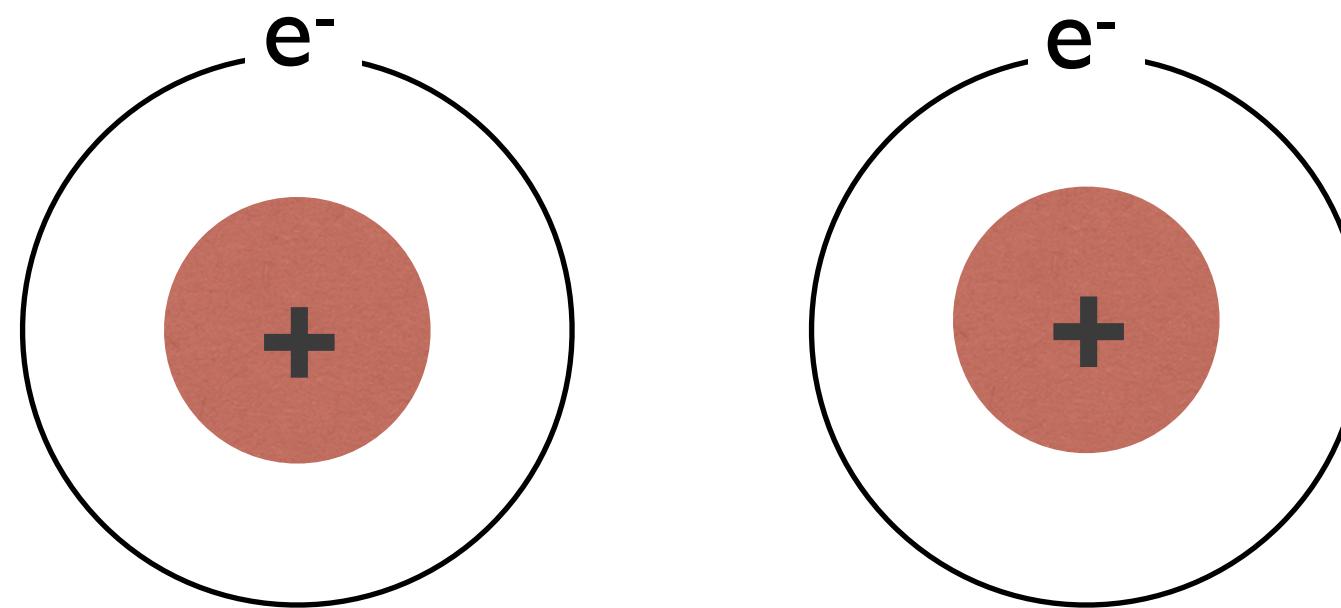
Gaymann, Geserich, Lohneysen, 95



What is the optimal n_d for DM searches?

Metal-insulator transition

Electrons are localized on dopants

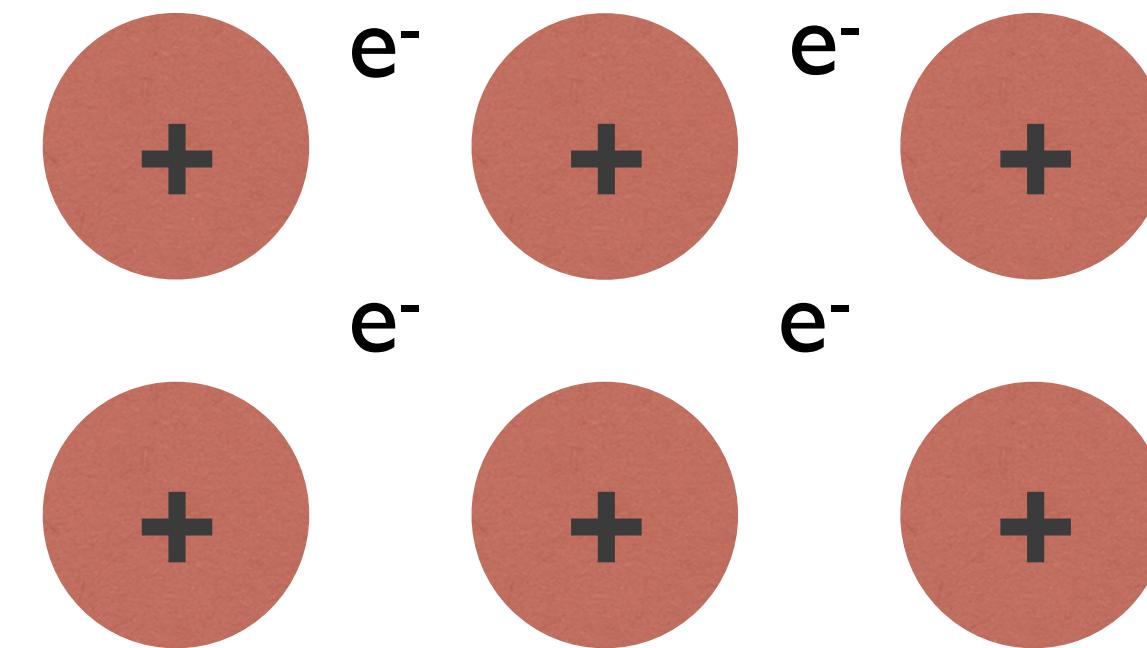


Insulating

Good for DM searches

$$n_d < n_c$$

Electrons are delocalized



Metallic

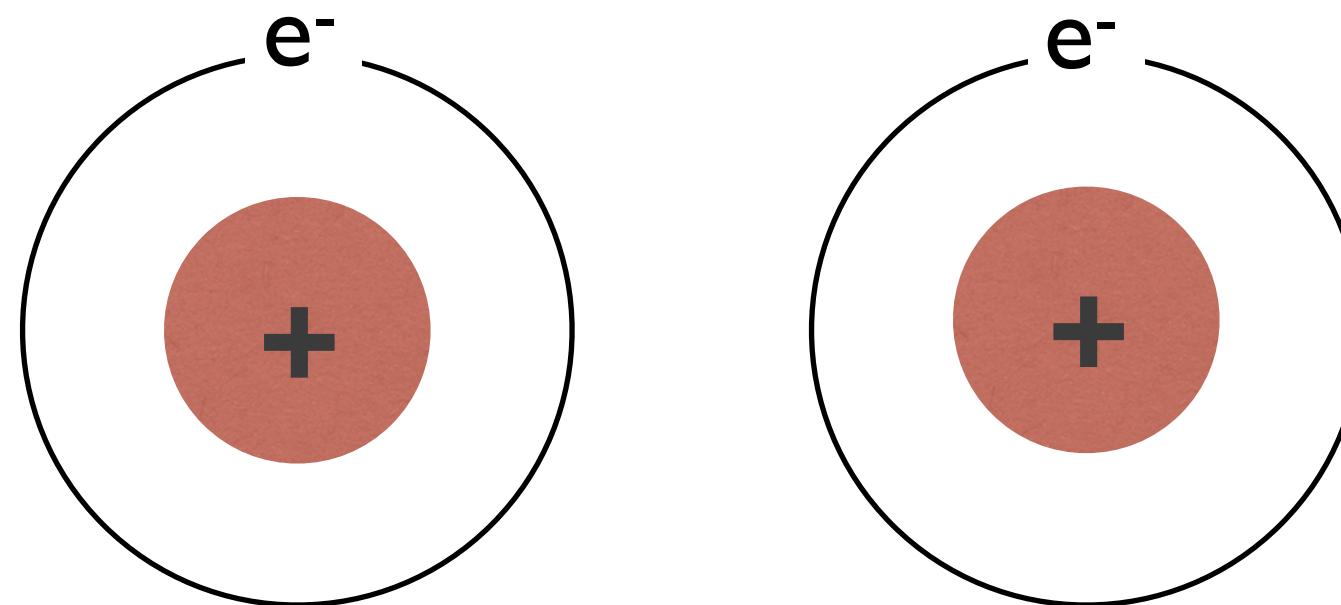
Metallic targets have no gap, hard to control noise

$$n_d > n_c$$

What is the optimal n_d for DM searches?

Metal-insulator transition

Electrons are localized on dopants

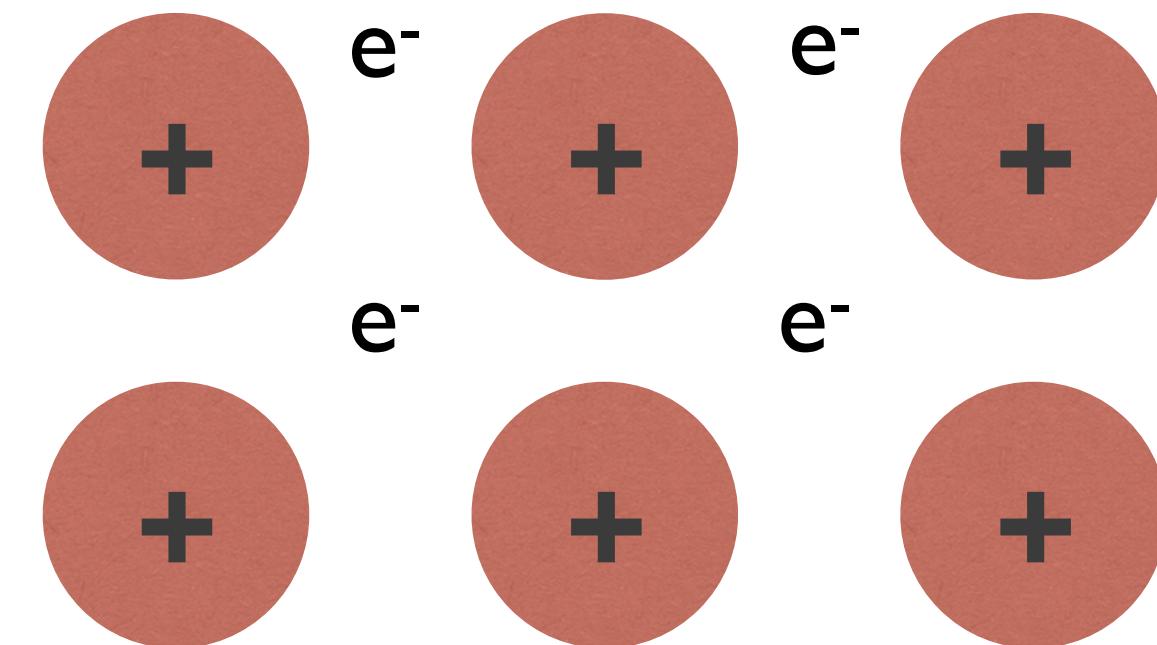


Insulating

Good for DM searches

$$n_d < n_c$$

Electrons are delocalized



Metallic

Metallic targets have no gap, hard to control noise

$$(n_c)^{-1/3} \sim a_*$$

For Phosphorus doped Si: $n_c = 3.5 \times 10^{18} \text{ cm}^{-3}$ We choose $1 \times 10^{18} \text{ cm}^{-3}$ for DM reach projection

DM-electron scattering rate

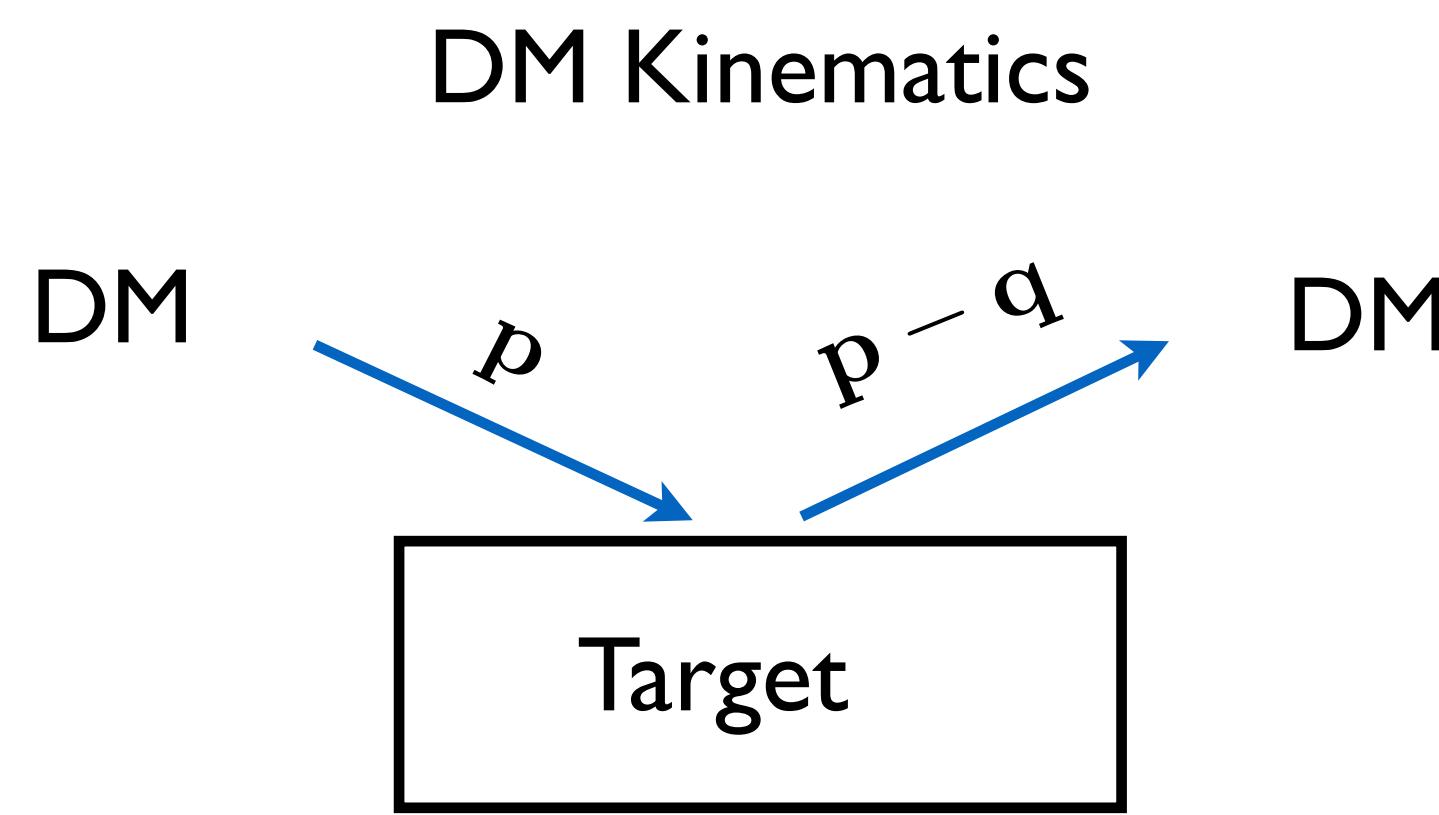
DM velocity distribution Particle interaction Target response

$$R \sim \int d^3\mathbf{v} f(\mathbf{v}) \int d^3\mathbf{q} F^2(\mathbf{q}) S(\mathbf{q}, \omega_{\mathbf{q}})$$

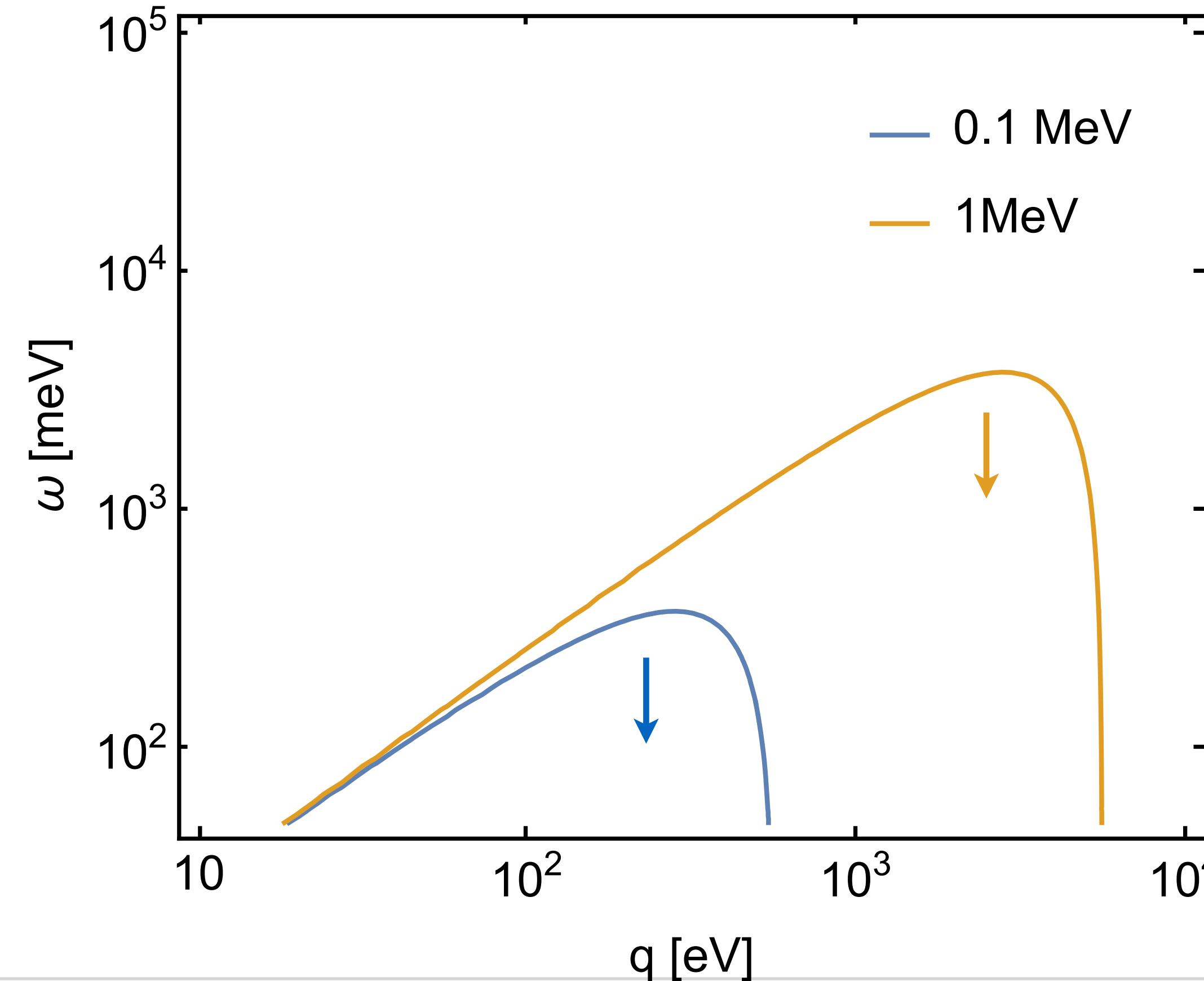
DM-electron scattering rate

DM velocity distribution Particle interaction Target response

$$R \sim \int d^3\mathbf{v} f(\mathbf{v}) \int d^3\mathbf{q} F^2(\mathbf{q}) S(\mathbf{q}, \omega_{\mathbf{q}})$$



$$\omega_{\mathbf{q}} = \frac{\mathbf{p}^2}{2m_\chi} - \frac{(\mathbf{p} - \mathbf{q})^2}{2m_\chi} = \mathbf{q} \cdot \mathbf{v} - \frac{q^2}{2m_\chi}$$



Target response

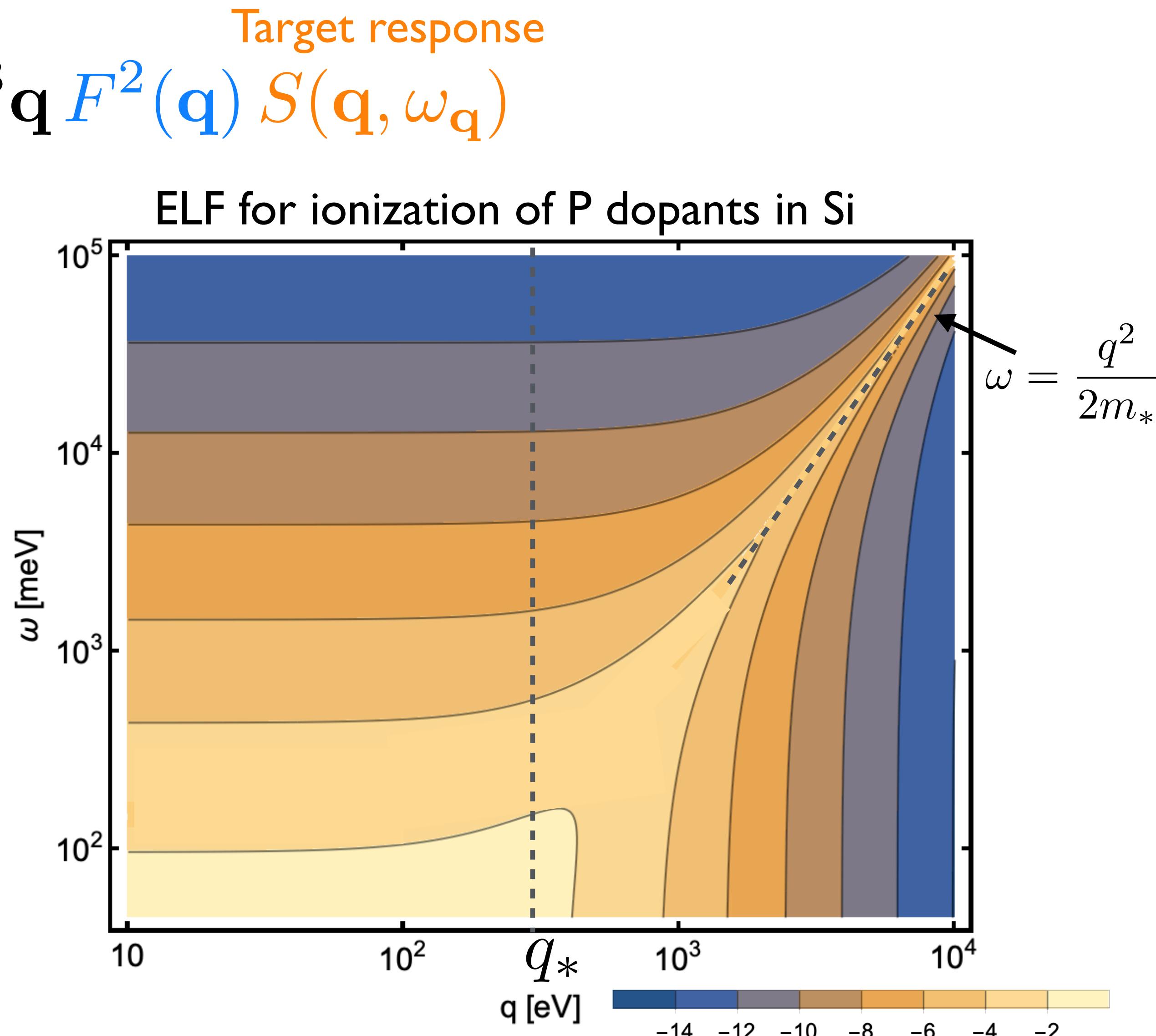
Knapen, Kozaczuk, Lin, 2021

Hochberg, Kahn, Kurinsky, Lehmann, Yu, Berggren, 2021

$$R \sim \int d^3\mathbf{v} f(\mathbf{v}) \int d^3\mathbf{q} F^2(\mathbf{q}) S(\mathbf{q}, \omega_{\mathbf{q}})$$

$$S(\mathbf{q}, \omega_{\mathbf{q}}) = \frac{q^2}{2\pi\alpha} \text{Im} \left[\frac{-1}{\epsilon(\mathbf{q}, \omega_{\mathbf{q}})} \right]$$

Energy loss function (ELF)



Target response

Knapen, Kozaczuk, Lin, 2021

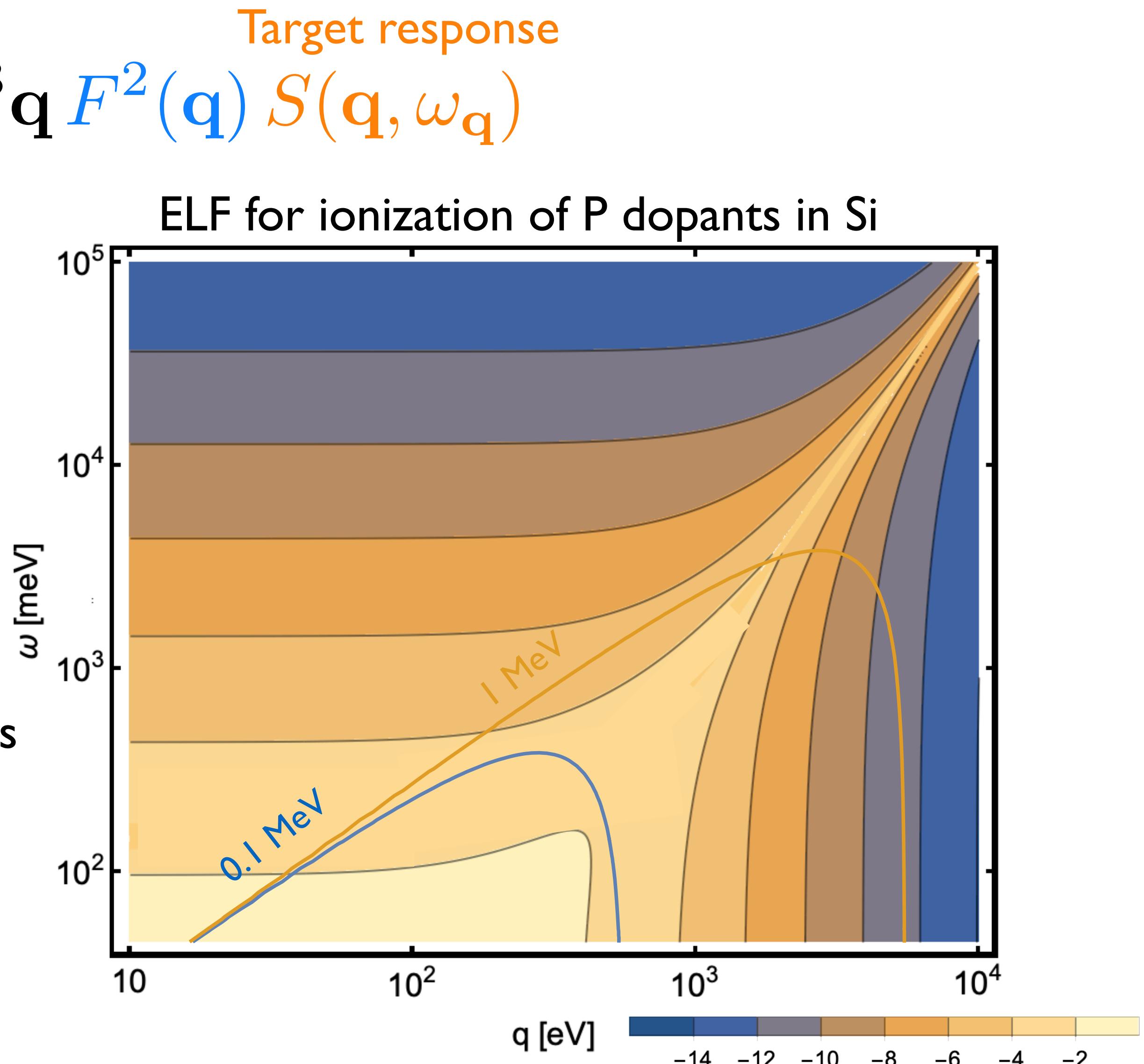
Hochberg, Kahn, Kurinsky, Lehmann, Yu, Berggren, 2021

$$R \sim \int d^3\mathbf{v} f(\mathbf{v}) \int d^3\mathbf{q} F^2(\mathbf{q}) S(\mathbf{q}, \omega_{\mathbf{q}})$$

$$S(\mathbf{q}, \omega_{\mathbf{q}}) = \frac{q^2}{2\pi\alpha} \text{Im} \left[\frac{-1}{\epsilon(\mathbf{q}, \omega_{\mathbf{q}})} \right]$$

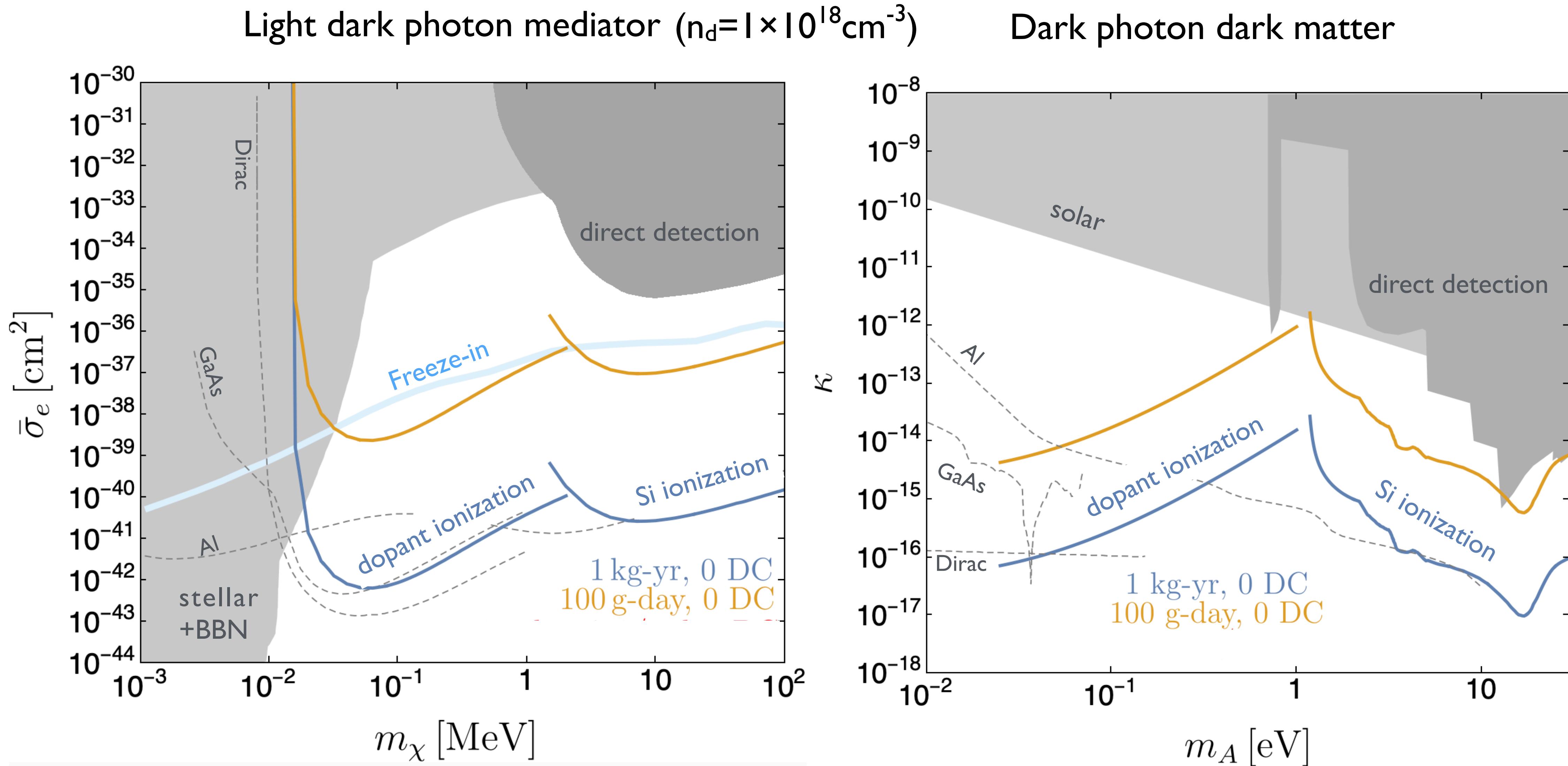
Energy loss function (ELF)

good reach for low mass DM with light mediators



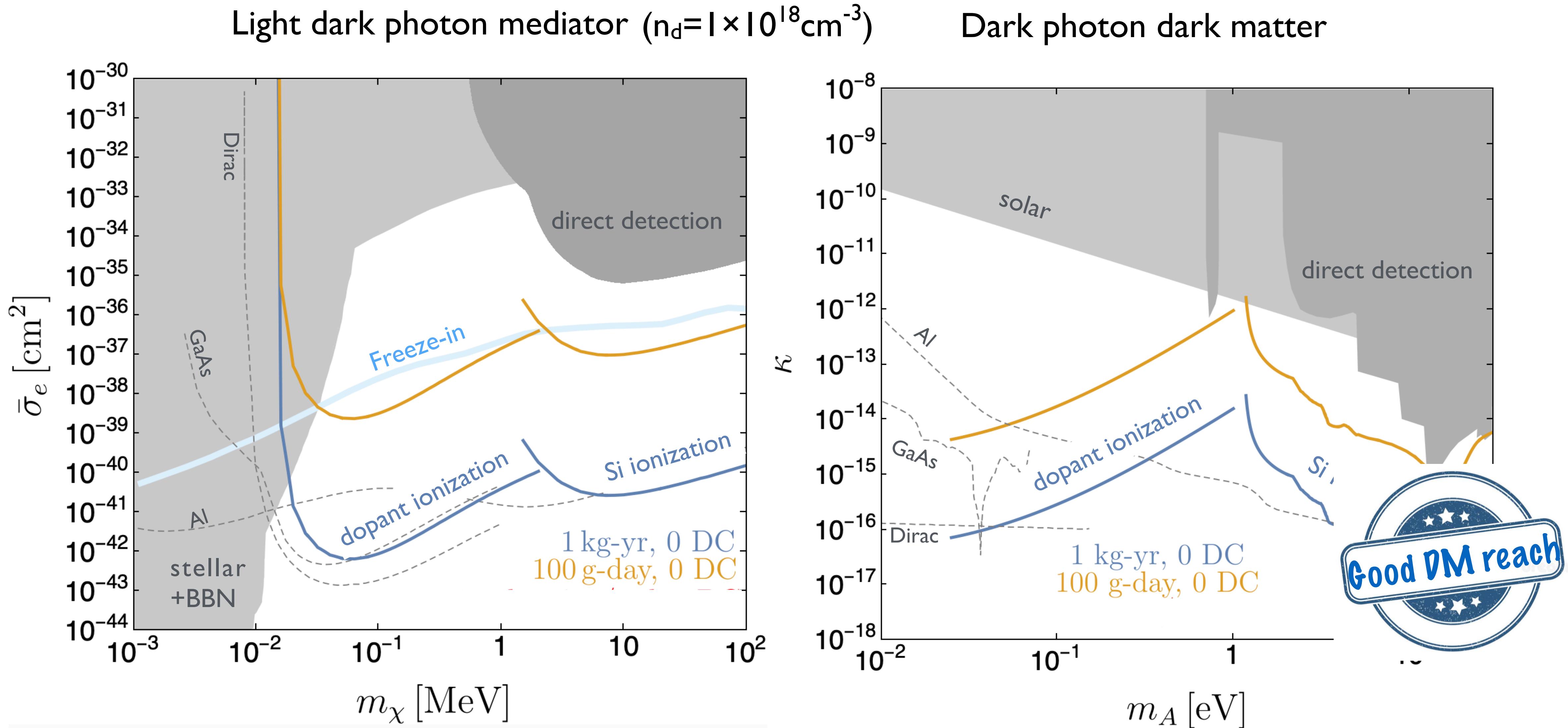
DM-electron scattering rate with doped silicon

PD, Egana-Ugrinovic, Essig, Sholapurkar, (in prep)

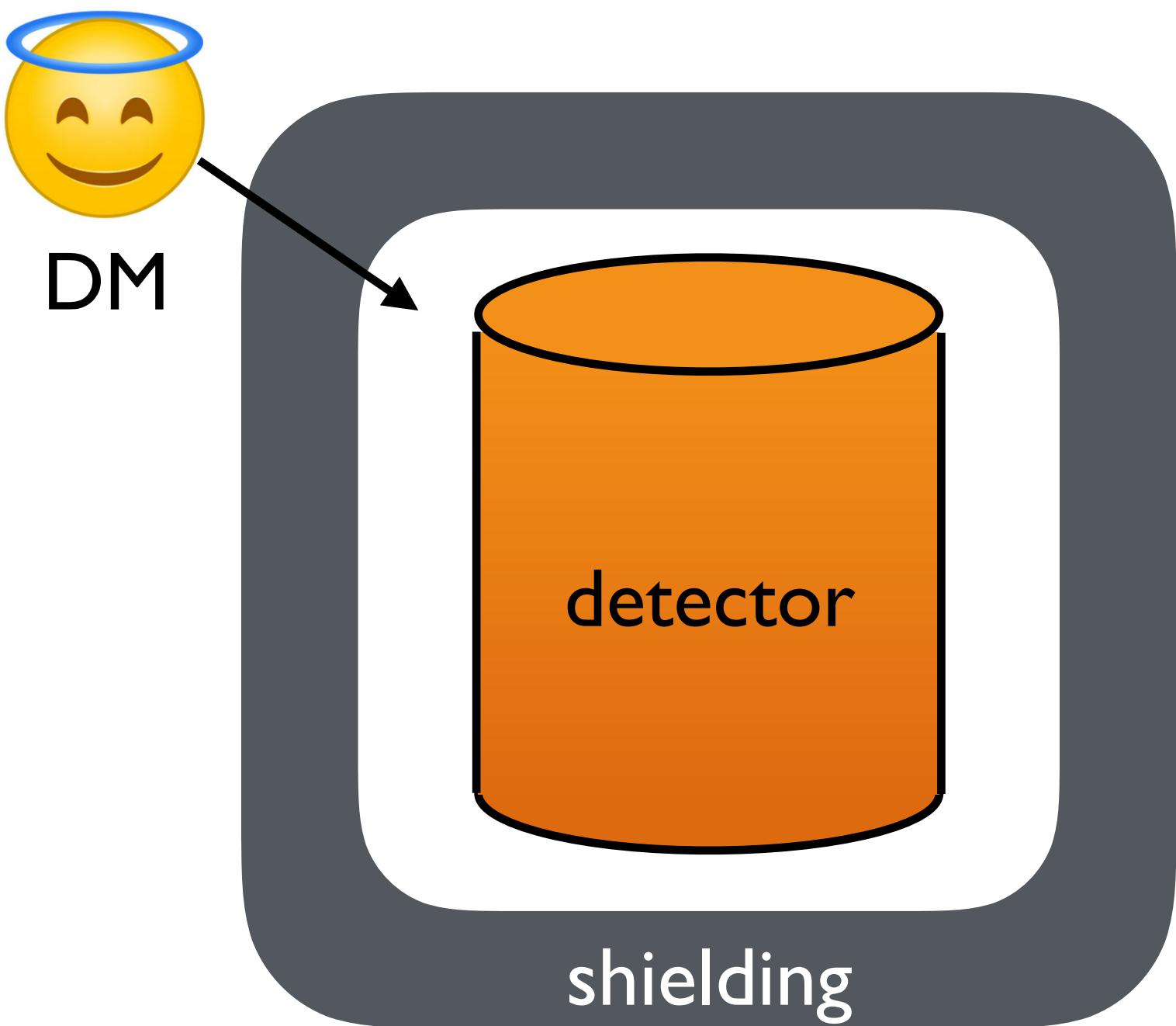


DM-electron scattering rate with doped silicon

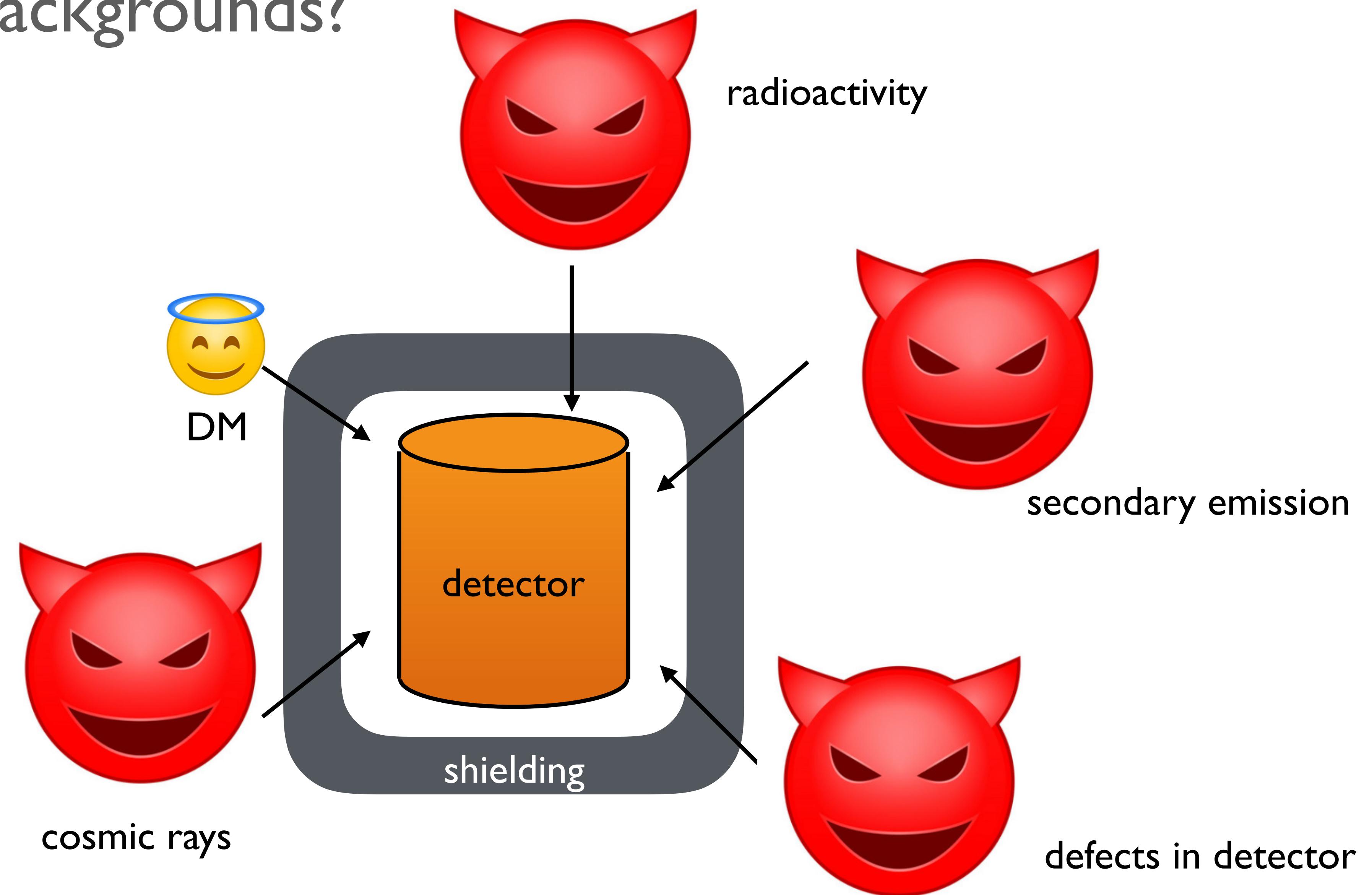
PD, Egana-Ugrinovic, Essig, Sholapurkar, (in prep)



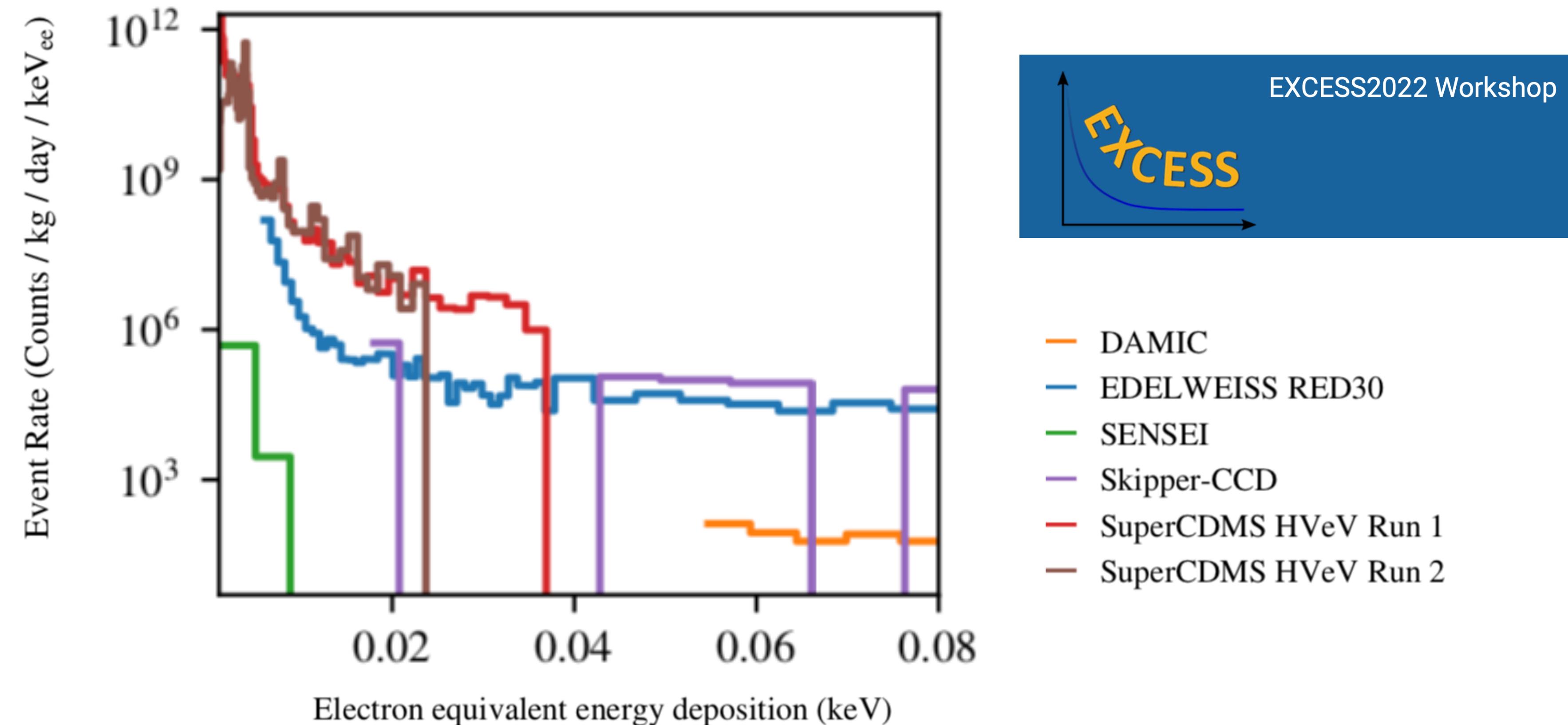
What about backgrounds?



What about backgrounds?



Low energy backgrounds at current detectors



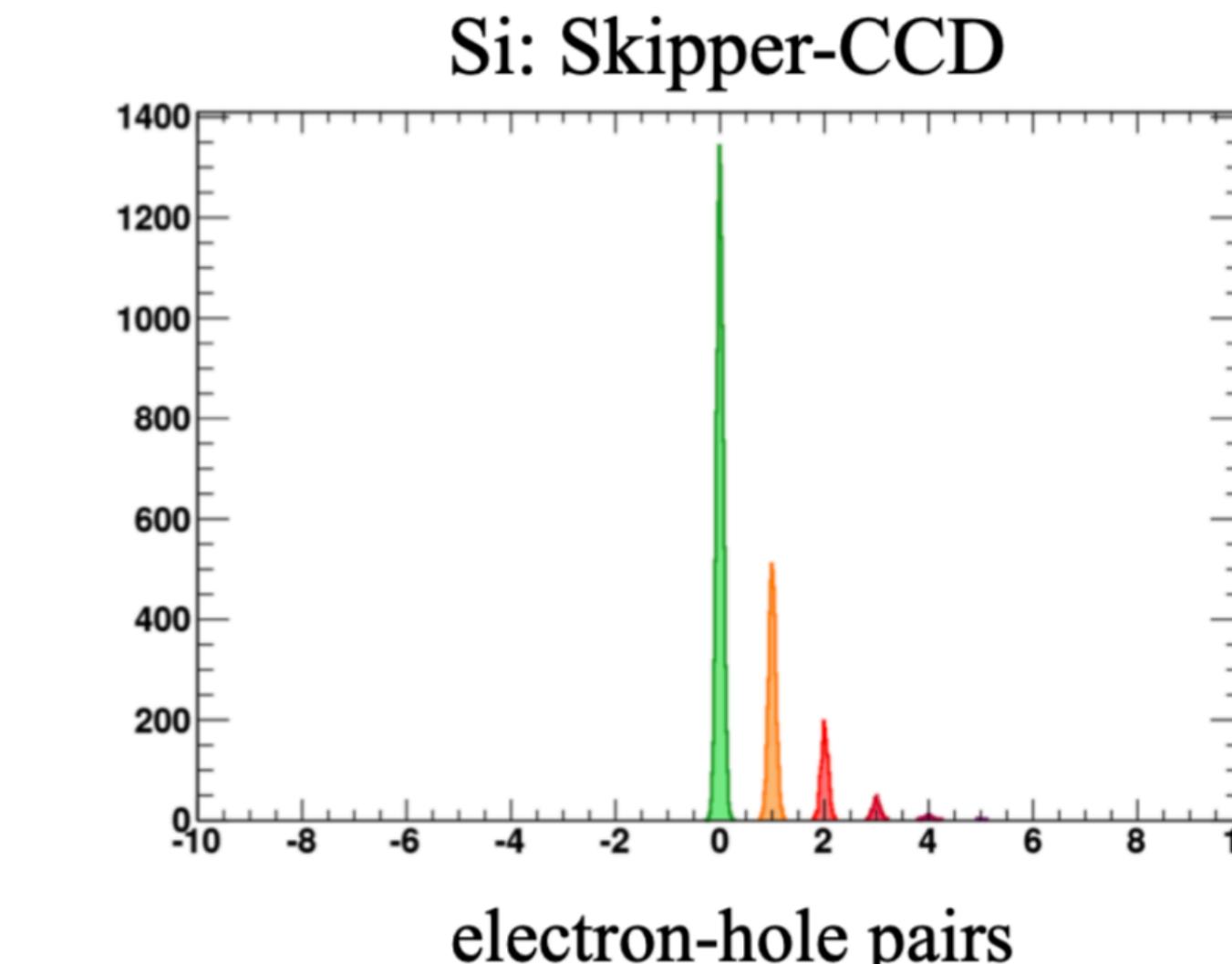
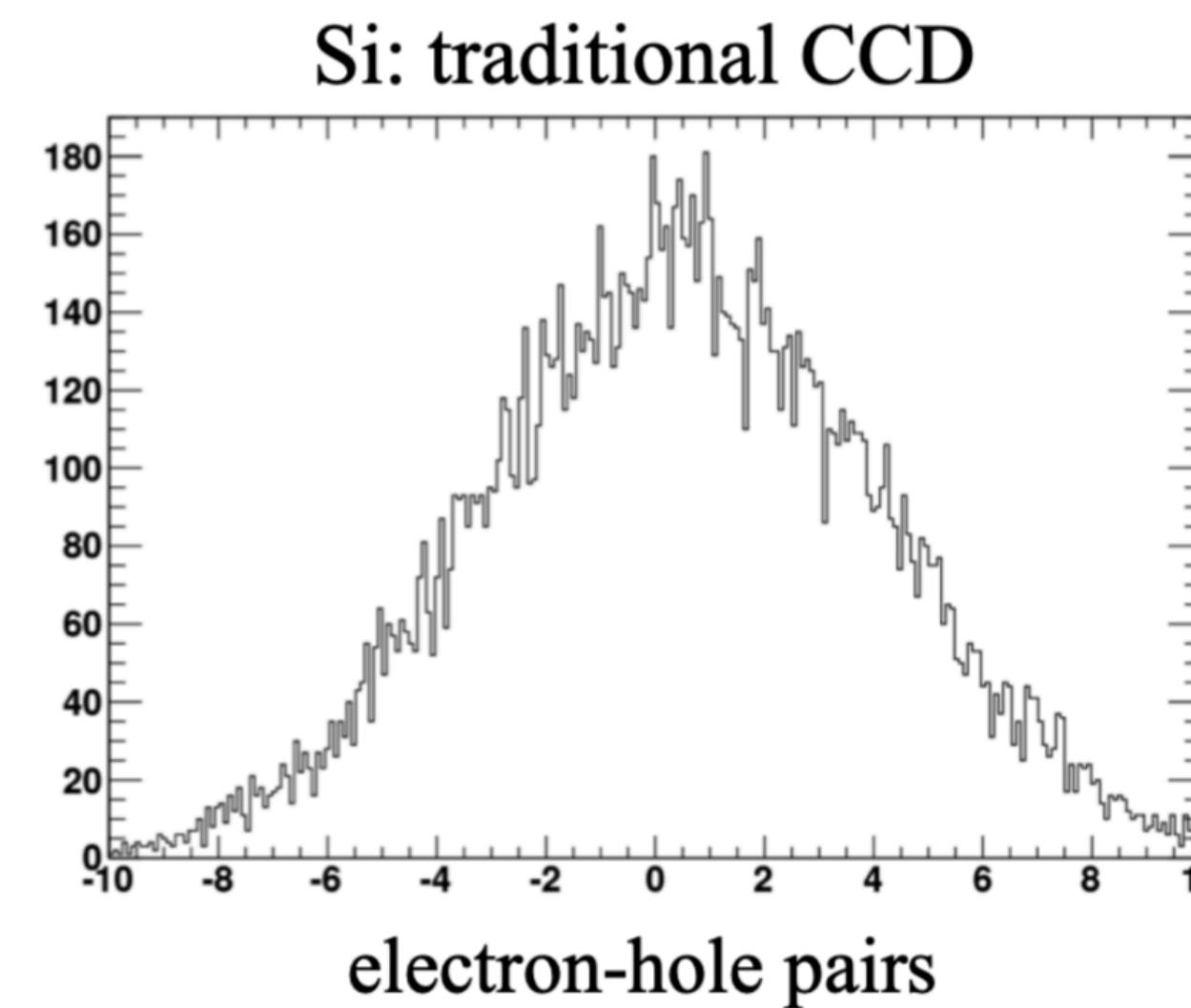
- Cherenkov radiation and radiative recombination may explain SENSEI and superCDMS excess
- There are likely more sources of excess: crystal cracking/microfracture...

PD, Egana-Ugrinovic, Essig, Sholapurkar, 2020

See talk by Mukul Sholapurkar

Skipper CCD

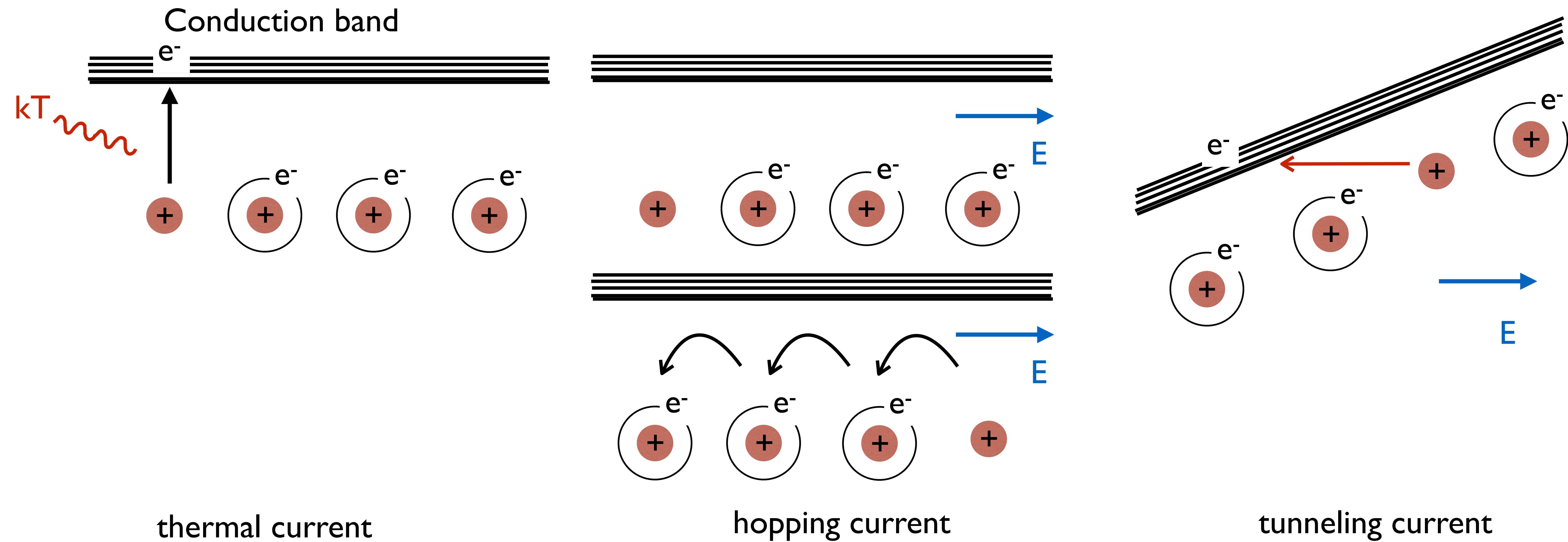
- Fully depleted and excellent spatial resolution
- Skipper readout: noise $\sim 1/\sqrt{N}$



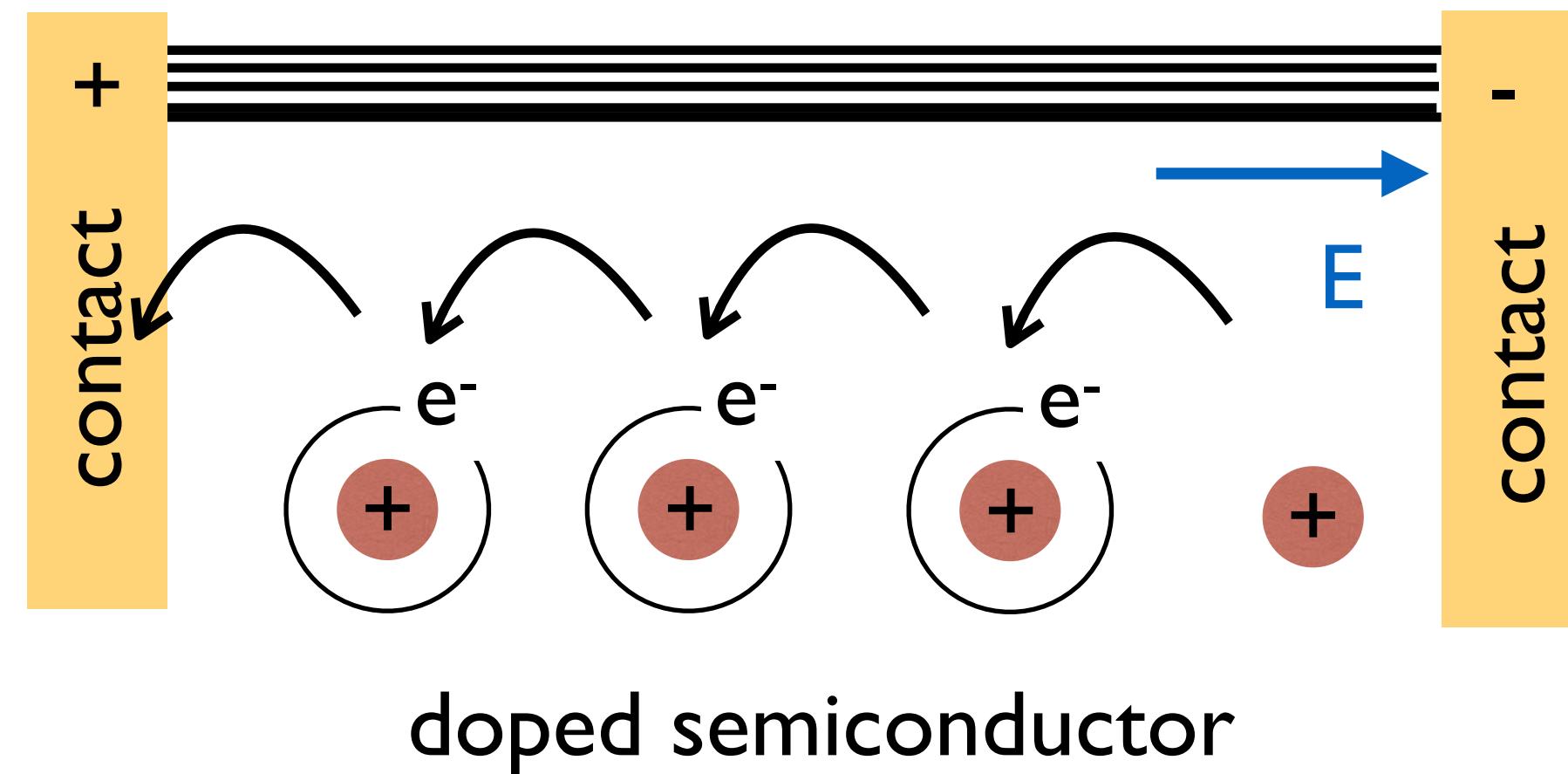
Essig, ICHEP 2020

- Single electron resolution and ultra-low dark current: $O(10^{-4})$ e/pixel/day

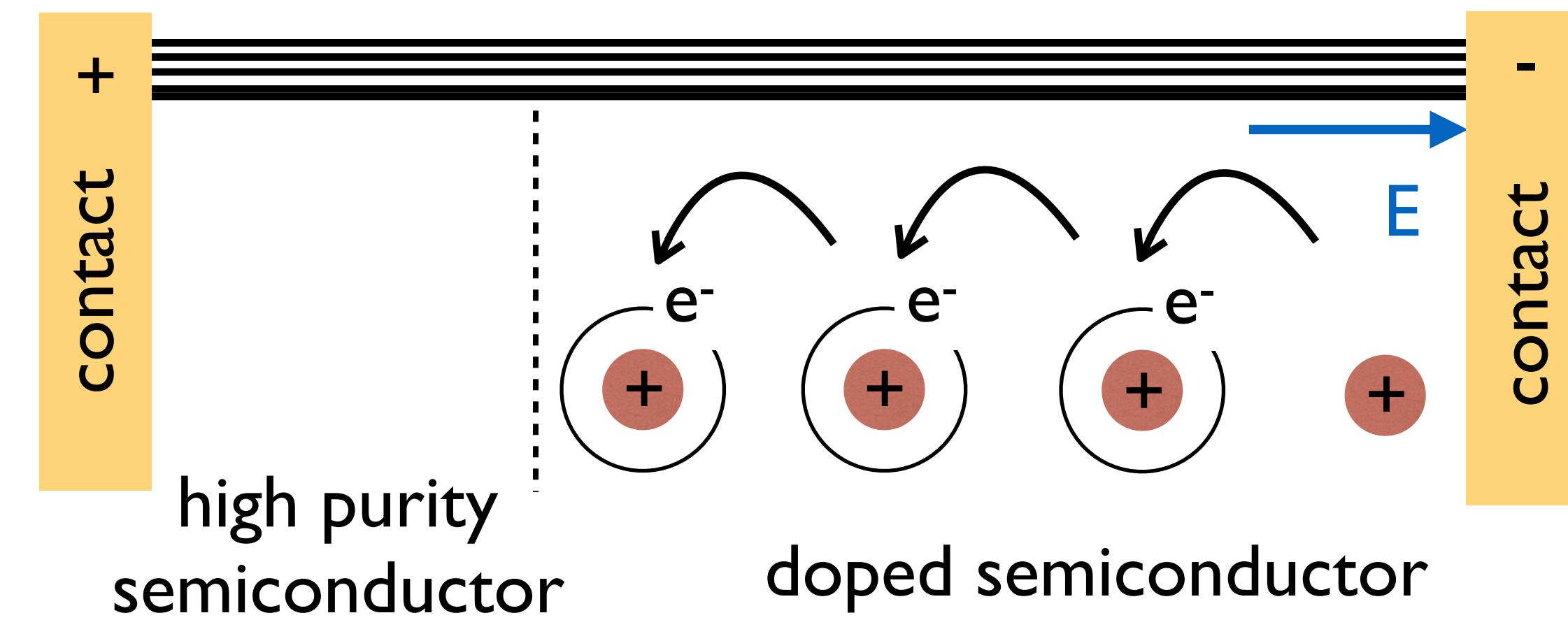
Dark current in doped semiconductor



Dark current in doped semiconductor



blocking layer blocks hopping current

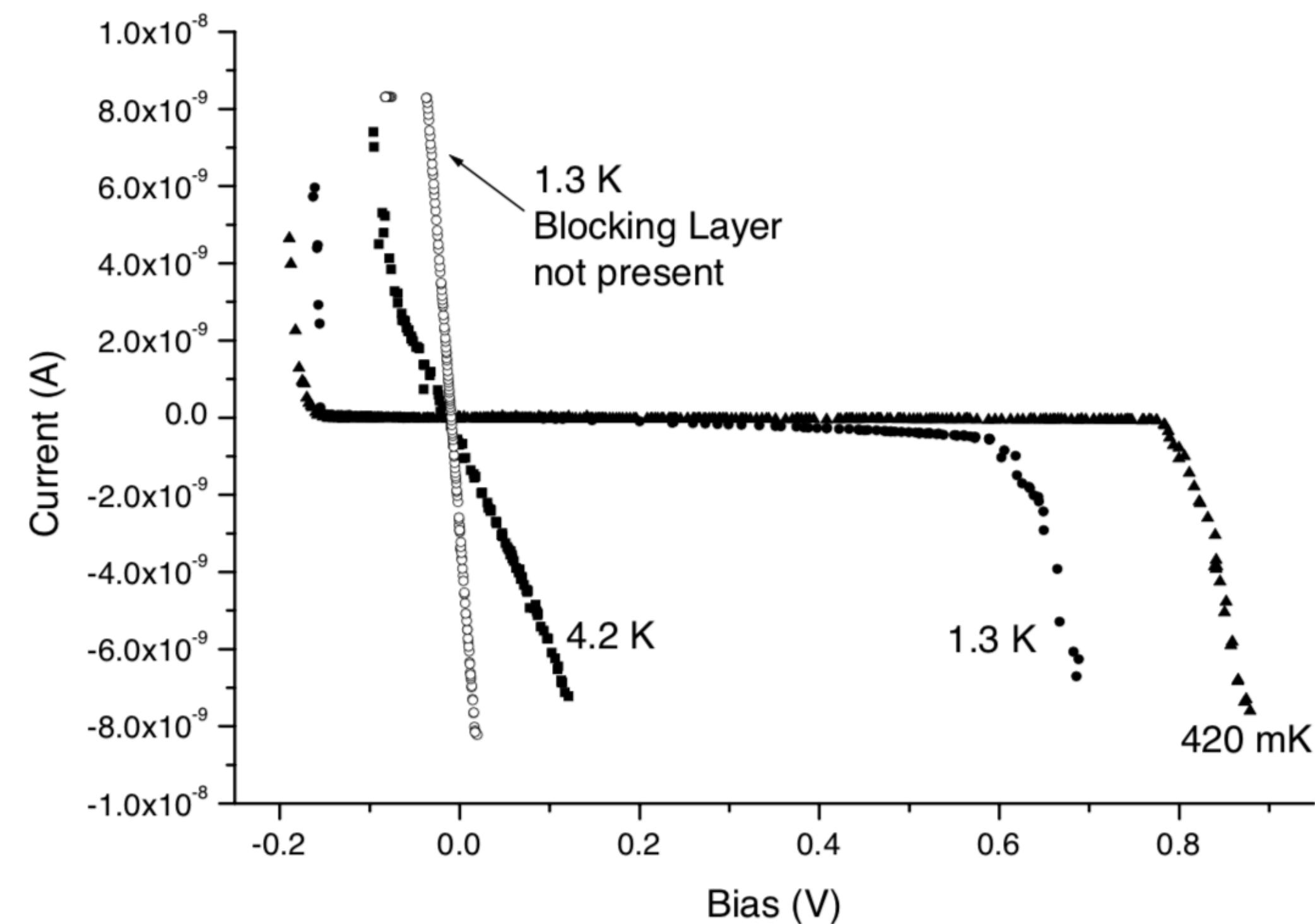


blocked impurity band detector: low dark current at low temperature with a blocking layer

Petroff, Stapelbroek, 1980

Dark current in BIB detector

GaAs:Te BIB detector



Benjamin Lewin Cardozo, 2004

- JWST uses Si:As BIB detector with dark current: $\mathcal{O}(10^{-2})$ e/pixel/s

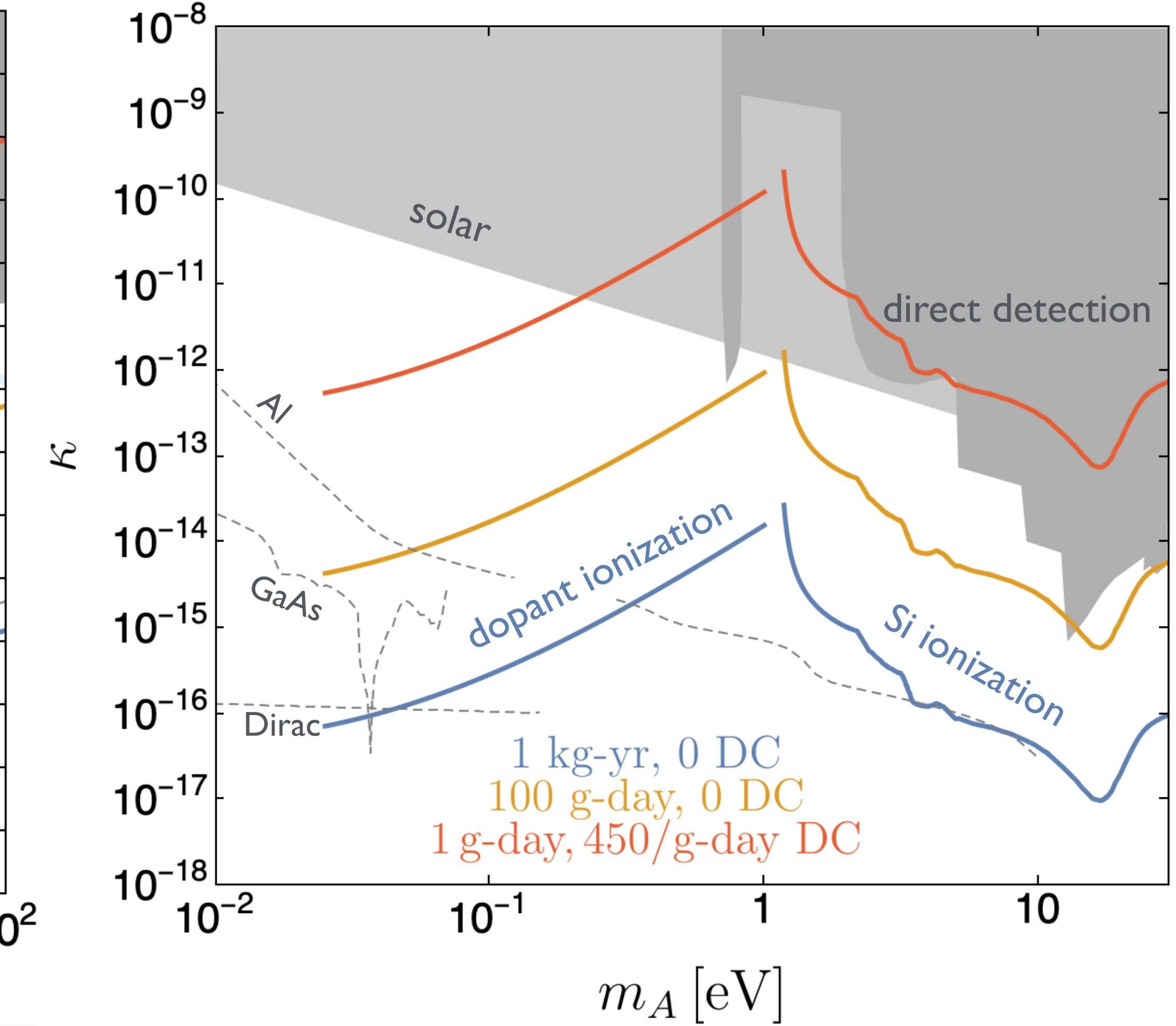
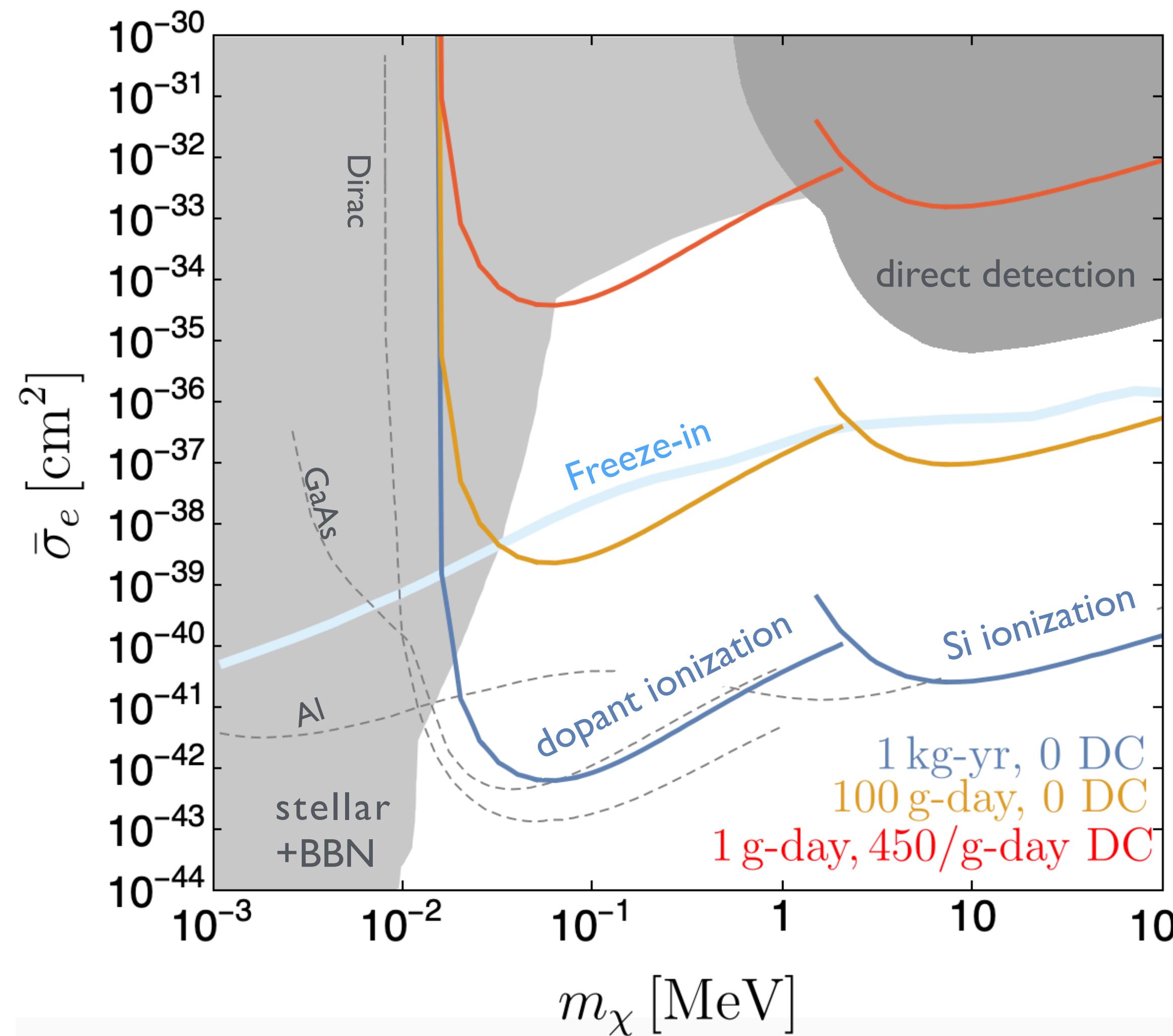
Rieke et.al., The Mid-Infrared Instrument for JWST

SENSEI dark current: $\mathcal{O}(10^{-4})$ e/pixel/day

- Skipper CCD with doped Si and blocking layers may achieve low dark current

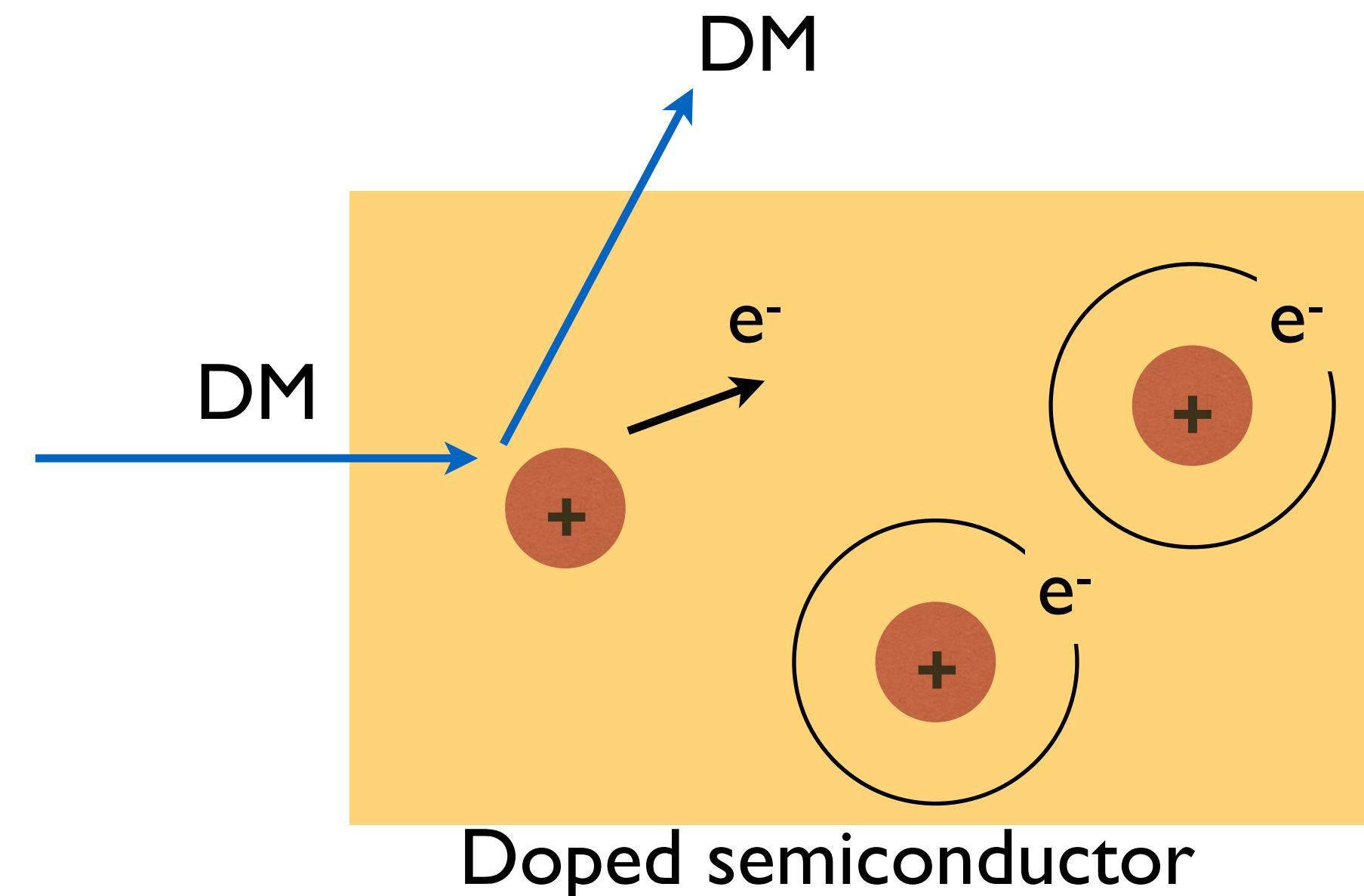
DM reach with backgrounds

PD, Egana-Ugrinovic, Essig, Sholapurkar, (in prep)



Conclusions

- Dopants in semiconductors can be thought as “Hydrogen atom” in a background with a large dielectric constant
- Doped semiconductors can be detector targets with $O(10\text{-}100)$ meV threshold and have sensitivity over a wide range of DM masses: >10 keV for DM scattering and >10 meV for DM absorption
- Skipper CCD with doped Si and blocking layers may achieve low dark current



Thank you

Summary of current experiments

Experiment	Location	Cherenkov contribution	Domiant Source of Cherenkov
SENSEI	~100m underground	likely dominant with radiative recombination	ambient high energy particles hitting detector
SuperCDMS HVeV	surface	likely dominant	ambient high energy particles hitting holders
EDELWEISS	~1800m underground	subdominant	radioactivity from impurities in holders
CRESST	~1400m underground	vetoed everything near the detector is instrumented	-

Good spatial resolution — SENSEI

Good timing resolution — EDELWEISS, CRESST

High ambient backgrounds — SuperCDMS HVeV

Low ambient backgrounds — EDELWEISS, CRESST

EDELWEISS and CRESST excess may dominantly come from crystal cracking/microfracture