

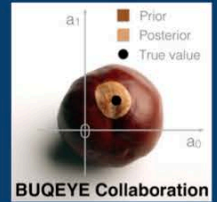
Theoretical advances and uncertainty quantification of neutron star properties

Christian Drischler

INT-22-2a: Neutron Rich Matter on Heaven and Earth

July 21, 2022

MICHIGAN STATE
UNIVERSITY



Ribbon-cutting ceremony
May 2, 2022



 Facility for Rare Isotope Beams
at Michigan State University

Samuel L. Stanley
President of MSU

Jennifer M. Granholm
Secretary of Energy

Recent neutron star observations

relativistic
Shapiro delay

What is the maximum
neutron star mass?



NICER
soft X-ray telescope

GW170817
GRB170817A
AT2017gfo



+ Virgo (Italy)
+ (KAGRA)

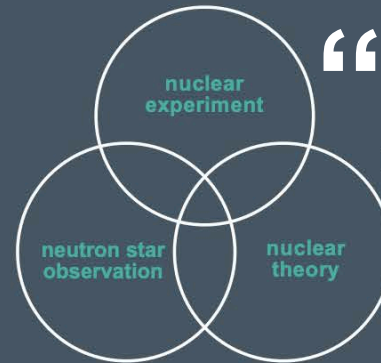
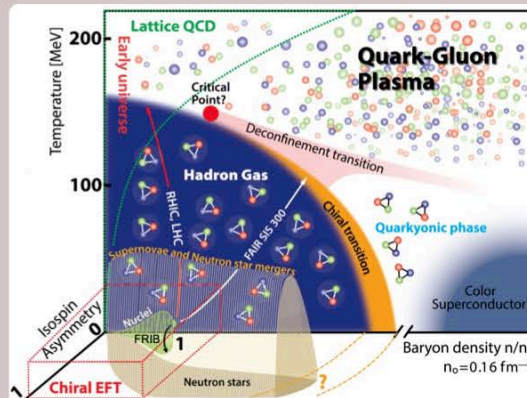
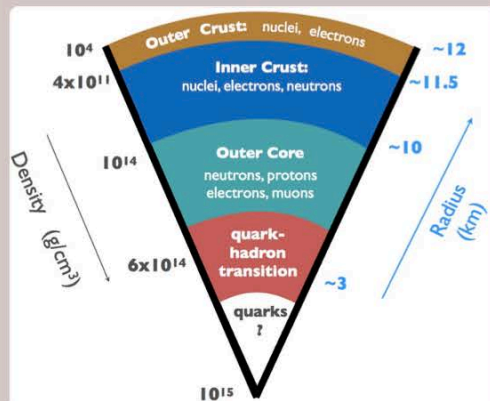
multi-messenger
astronomy

Nuclear theory in the precision era

How can neutron star observations help **improve** nuclear **effective field theories** ?

What are the **phases** of neutron star matter below two times normal densities ?

At what density scale does nuclear effective field theory **break down** ?



“ Enormous progress in theory, experiment, and observation make this new era particularly fruitful for the determination of the equation of state of neutron-rich matter.

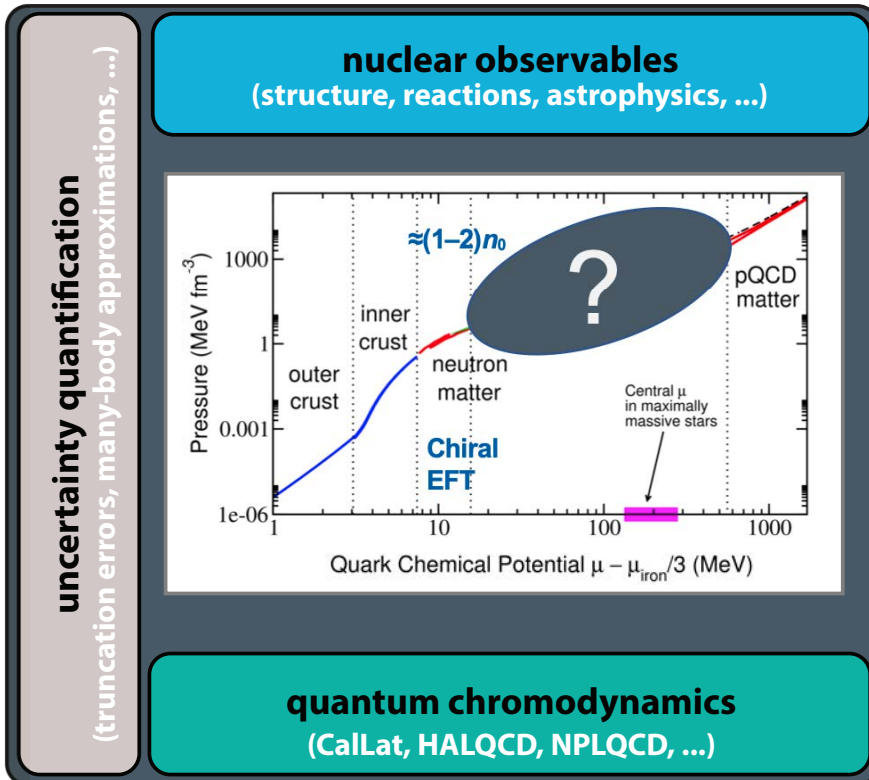
Required: statistically meaningful comparisons ”

Papers presenting the results of theoretical calculations are **expected to include uncertainty estimates...**

- If the authors claim **high accuracy**, or improvements on the accuracy of previous work.
- If the primary motivation for the paper is to make **comparisons with present or future high precision experimental measurements**.
- If the primary motivation is to provide interpolations or extrapolations of known experimental measurements.

Phys. Rev. A: Editorial (April 2011)

Ab initio workflow (idealized)

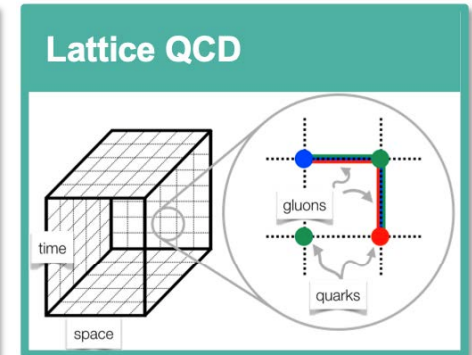
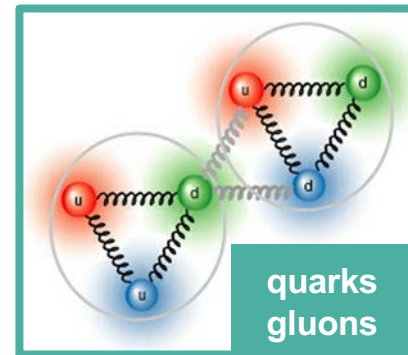


CD & Bogner, *Few Body Syst.* **62**, 109

Here: nuclear equation of state (EOS)
 energy per particle (and derived quantities)

$$\frac{E}{A}(n, \delta, T)$$

baryon density n
 neutron excess δ
 temperature $T (= 0)$

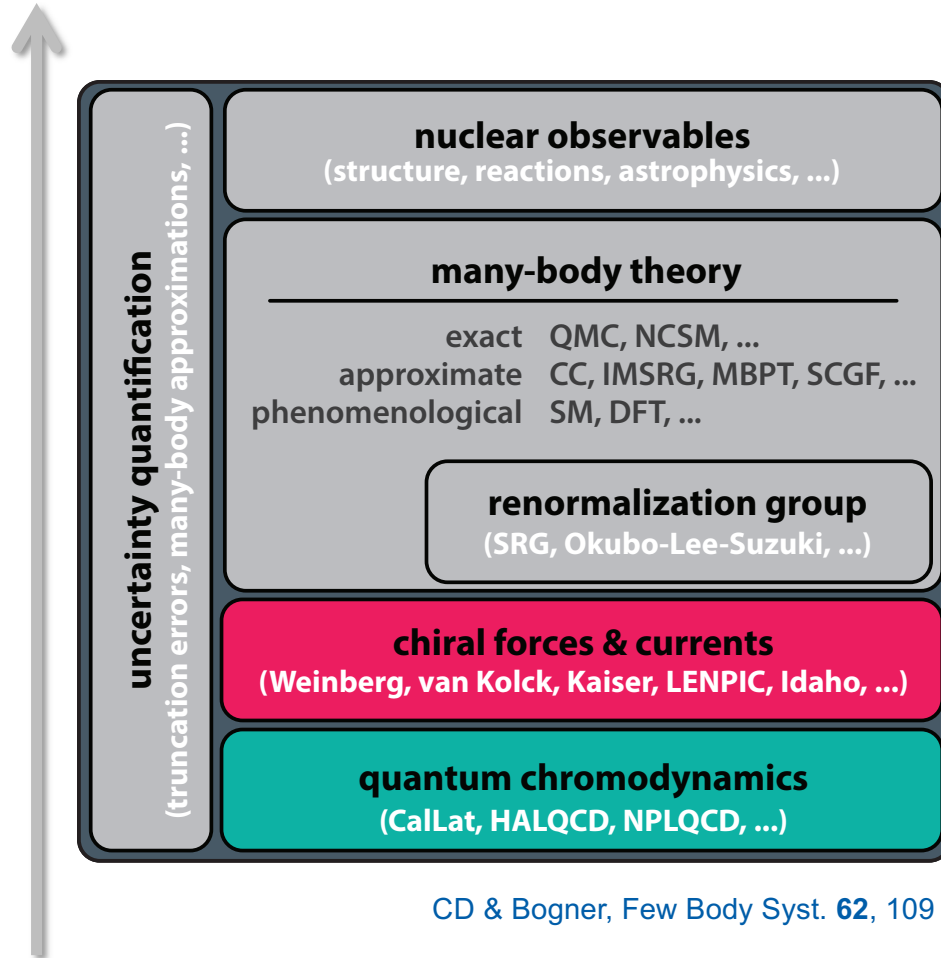


theory of strong interactions

QCD is nonperturbative at the low energies
 relevant for nuclear physics (cf. pQCD & LQCD)

CD, Haxton, McElvain, Mereghetti *et al.*, *PPNP* **121**, 103888

Ab initio workflow (idealized)



CD & Bogner, Few Body Syst. 62, 109

Here: nuclear equation of state (EOS)
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$$\frac{E}{A}(n, \delta, T)$$

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computational framework

solves the (many-body) Schrödinger equation
requires a nuclear potential as input

chiral effective field theory

provides microscopic interactions consistent with
the symmetries of *low-energy* QCD

theory of strong interactions

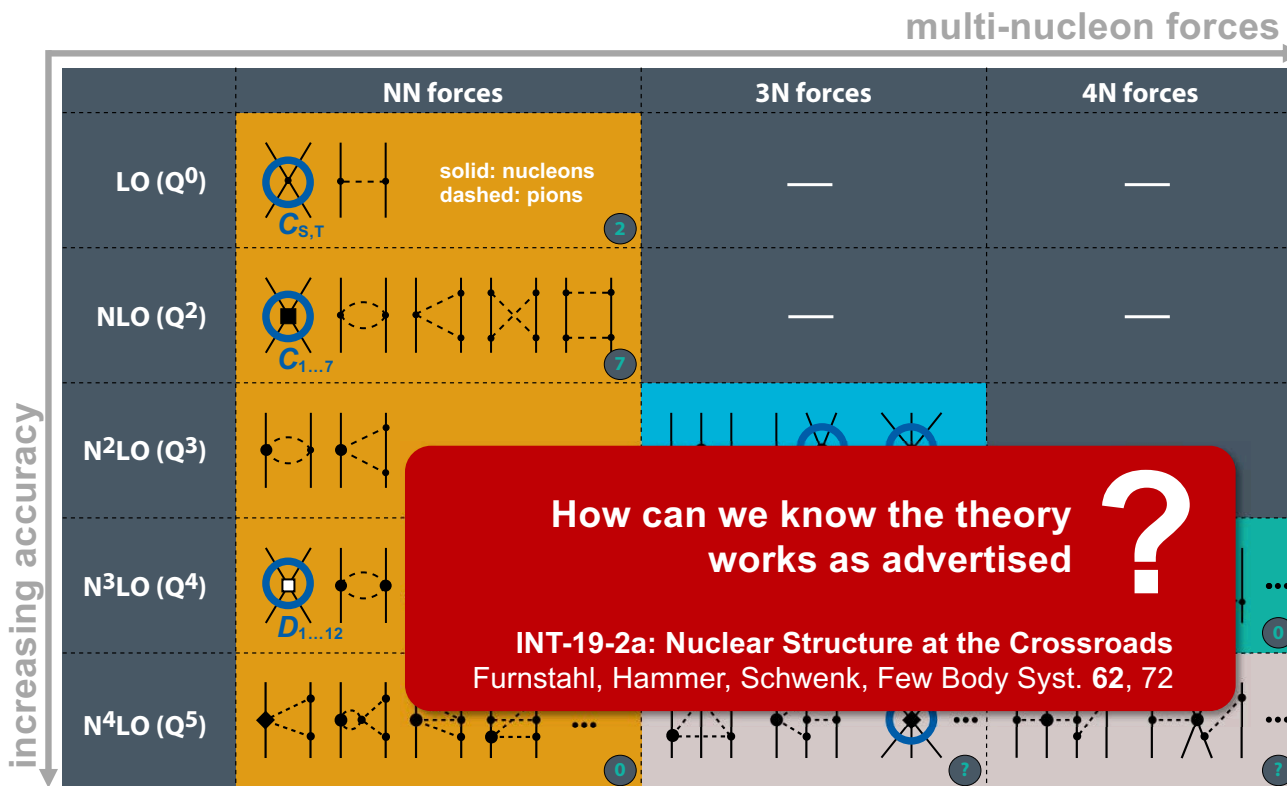
QCD is nonperturbative at the low energies
relevant for nuclear physics (cf. pQCD & LQCD)

Modern theory of nuclear forces



Hierarchy of chiral nuclear forces up to N⁴LO

$$Q = \max\left(\frac{p}{\Lambda_b}, \frac{m_\pi}{\Lambda_b}\right) \lesssim \frac{1}{3}$$



Chiral Effective Field Theory

dominant approach to deriving microscopic interactions consistent with the symmetries of low-energy QCD

degrees of freedom: **nucleons & pions**

fit the **unknown couplings** θ to experimental (or lattice) data

- NN: phase shifts & deuteron
- 3N/4N: binding energies, charge radii

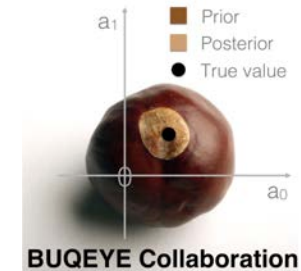
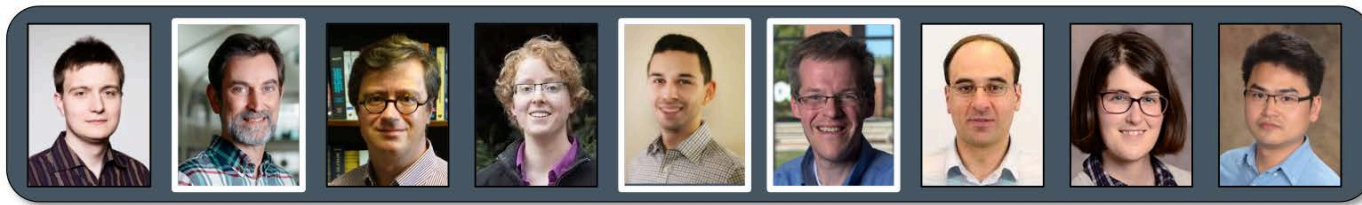
EFT expansion enables **uncertainty quantification** (truncation errors)

Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Kaiser, Krebs, Machleidt, Meißner, ...

see also: Low Energy Nuclear Physics International Collaboration (LENPIC)

Bayesian statistics is ideal for quantifying and propagating theoretical uncertainties

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Bayesian
Uncertainty
Quantification:
Errors for
Your
EFT

Open-source software & tutorials (Jupyter): <https://buqeye.github.io>

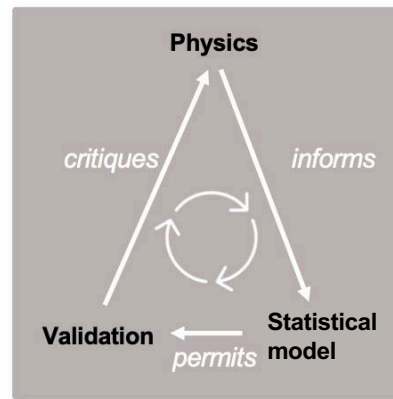
Select papers for UQ in EFT-based calculations:

How well do we know the neutron-matter EOS at the densities inside neutron stars? A Bayesian approach with correlated uncertainties, CD, Furnstahl, Melendez, Phillips, PRL **125**, 202702.

Rigorous constraints on 3N forces [...] calculations of few-body observables, Wesolowski, Svensson, Ekström *et al.*, PRC **104**, 064001.

Fast & accurate emulation of two-body scattering observables without wave functions, Melendez, CD, Garcia, Furnstahl, Zhang, PLB **821**, 136608.

Designing Optimal Experiments: An Application to Proton Compton Scattering, Melendez, Furnstahl, Grieshammer, McGovern, Phillips, Pratola, EPJ **57**, 81.



BUQEYE aims to use (Bayesian) statistical tools to **answer fundamental problems** in the **construction and application of EFTs** in nuclear physics. The tools include:

- **parameter estimation**
- **model checking & selection**
- **fast & accurate emulators**
- **experimental design**

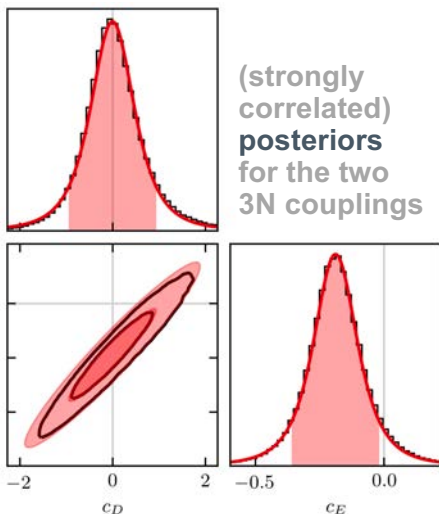


An NSF Cyberinfrastructure Framework designed to facilitate:

- **Bayesian Model Mixing** to quantify model uncertainty;
- **full UQ** for experimentally inaccessible environments such as neutron stars;
- Bayesian **experimental design** to assess the impact of proposed experiments.

BAND Manifesto,
Phillips, Furnstahl, Heinz *et al.*,
JGP: NP **48** 072001

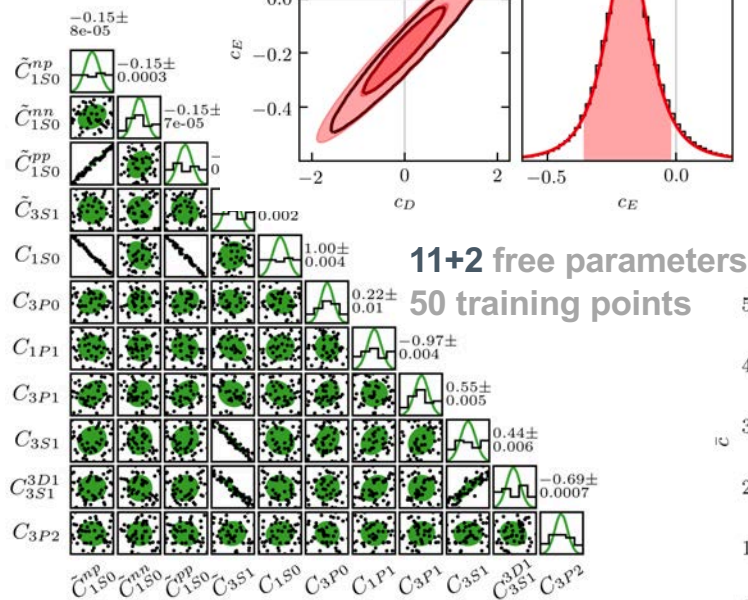
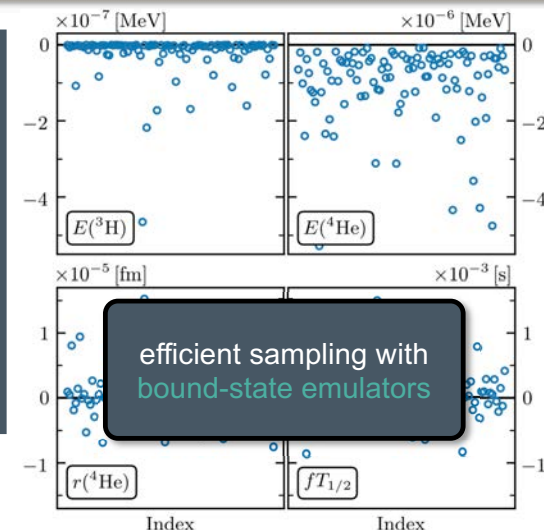
Spotlight: rigorous constraints on chiral forces



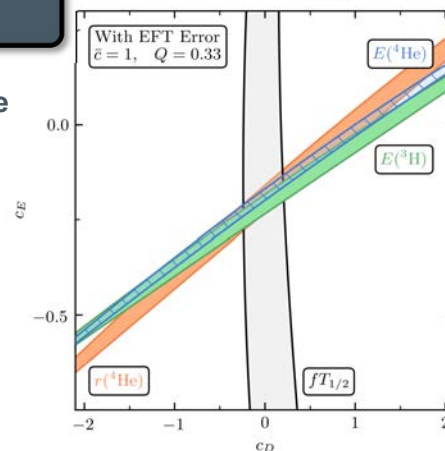
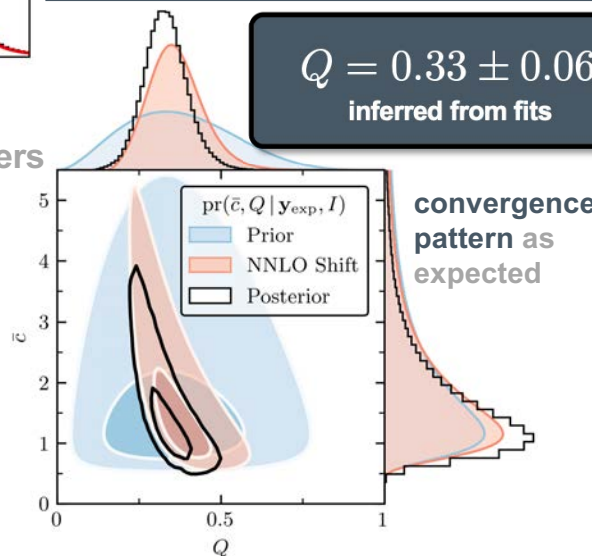
First chiral potentials (up to N²LO) with uncertainties fully quantified

- ✓ EFT truncation errors
- ✓ data errors (LEC fits)
- ✓ method uncertainties

Applications to $A = 6$ nuclei:
Djäv, Ekström *et al.*, PRC 105, 014005

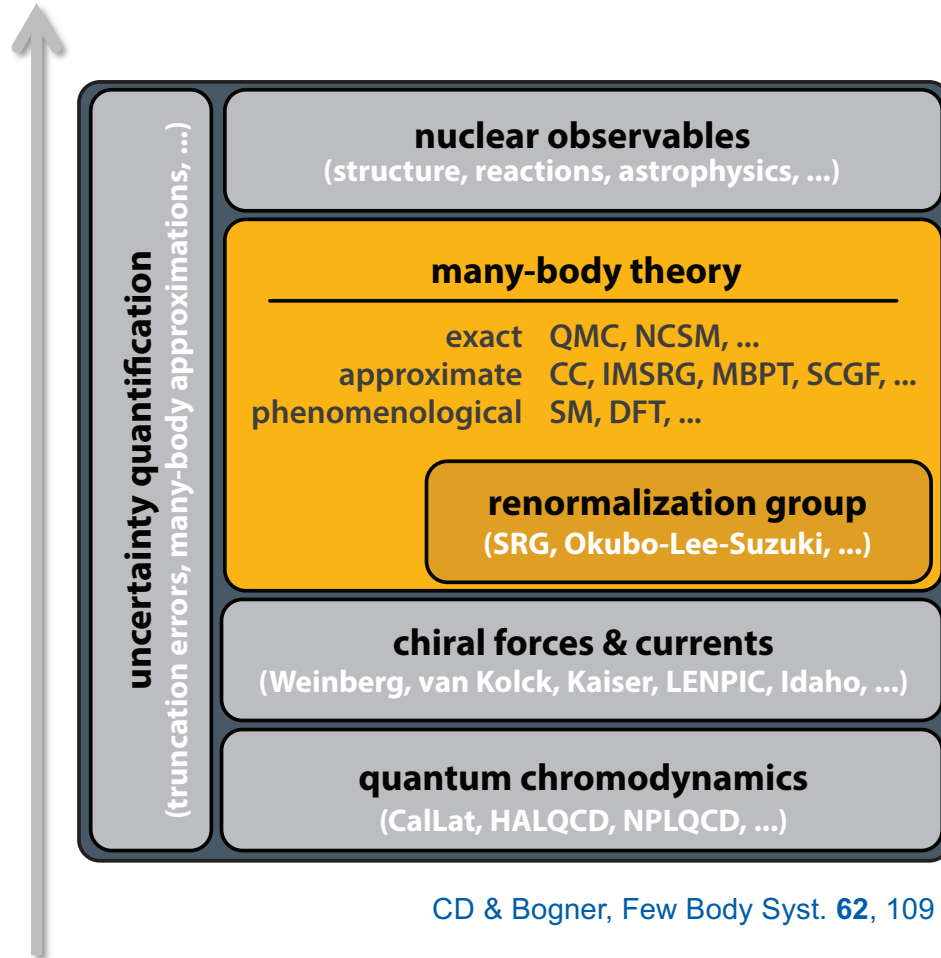


$Q = 0.33 \pm 0.06$
inferred from fits



Simultaneous fits can be problematic without EFT truncation errors

Ab initio workflow (idealized)



CD & Bogner, Few Body Syst. 62, 109

Here: nuclear equation of state (EOS)
energy per particle (and derived quantities)

$$\frac{E}{A}(n, \delta, T)$$

baryon density n
neutron excess δ
temperature $T (= 0)$

Here: many-body perturbation theory (MBPT)


computationally efficient method (HPC-friendly)
allows to estimate many-body uncertainties

Widely applicable:

- ✓ arbitrary proton fractions
- ✓ finite temperature
- ✓ optical potentials, linear response, nuclei, ...

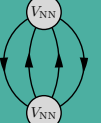
Other frameworks include **quantum Monte Carlo**,
coupled cluster, and self-consistent Green's functions

Many-body perturbation theory (MBPT) in a nutshell



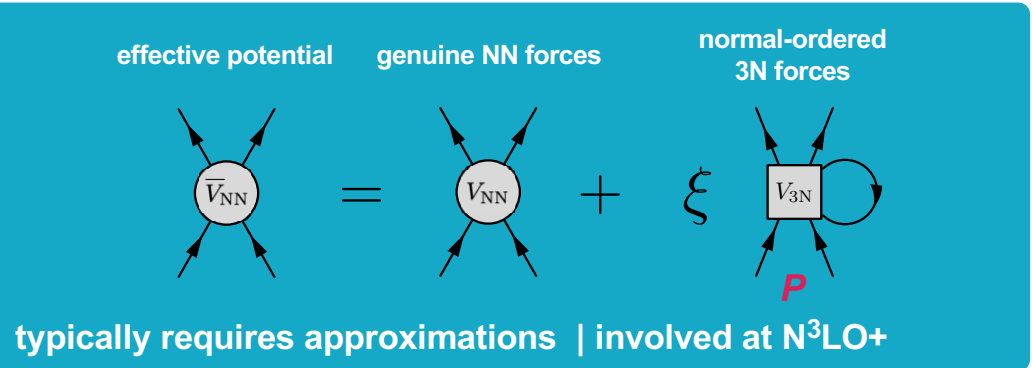
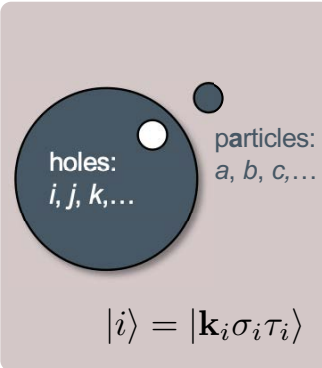
$$\frac{E^{(0)}}{V} = +\frac{1}{2} \sum_{ij} \langle ij | \bar{V}_{NN} | ij \rangle$$


Hartree-Fock



$$\frac{E^{(2)}}{V} = \frac{1}{4} \sum_{ij} \frac{|\langle ij | \bar{V}_{NN} | ab \rangle|^2}{\varepsilon_i + \varepsilon_j - \varepsilon_a - \varepsilon_b}$$

second order





$$\frac{E_{hh}^{(3)}}{V} = +\frac{1}{8} \sum_{ab} \frac{\langle ij | \bar{V}_{NN} | ab \rangle \langle kl | \bar{V}_{NN} | ij \rangle \langle ab | \bar{V}_{NN} | kl \rangle}{D_{ijab} D_{klab}}$$

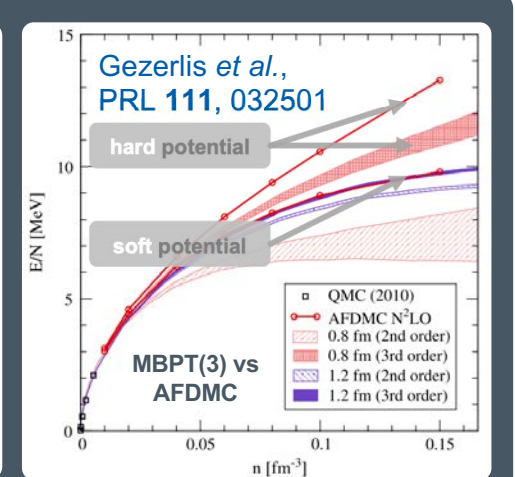
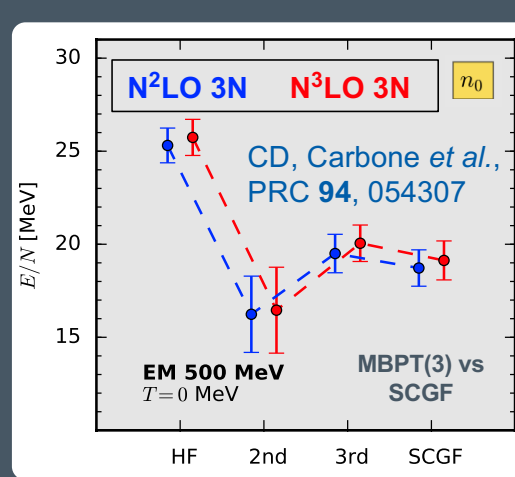
involved partial-wave decomposition

$$\frac{E_{ph}^{(3)}}{V} = +\sum_{abc} \frac{\langle ij | \bar{V}_{NN} | ab \rangle \langle ak | \bar{V}_{NN} | ic \rangle \langle bc | \bar{V}_{NN} | jk \rangle}{D_{ijab} D_{jkc}}$$

see Coraggio, Holt *et al.*, PRC 89, 044321

$$\frac{E_{pp}^{(3)}}{V} = +\frac{1}{8} \sum_{abcd} \frac{\langle ij | \bar{V}_{NN} | ab \rangle \langle ab | \bar{V}_{NN} | cd \rangle \langle cd | \bar{V}_{NN} | ij \rangle}{D_{ijab} D_{ijcd}}$$

third order



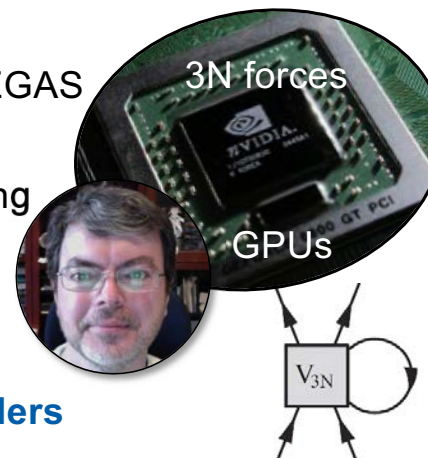
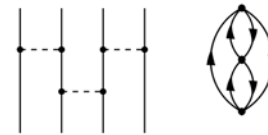
nonperturbative benchmarks in neutron matter

Monte Carlo framework for MBPT (second generation)



efficient evaluation of **MBPT diagrams**
with **NN**, **3N**, and **4N forces** (in a single-particle basis)

- **implementation of arbitrary diagrams** has become **straightforward** (numerically exact)
- multi-dimensional momentum integrals: improved VEGAS with CUDA, OpenMP, and MPI support
- hybrid, GPU-accelerated approach for normal ordering
- propagation of importance sampling distributions (e.g., for mapping the EOS efficiently in density)
- **fast EOS calculations at low MBPT orders; controlled MBPT calculations at higher MBPT orders**



GPU-accelerated
diagram evaluation

automated code
generation

analytic expressions
interaction & MBPT diagrams

CD, McElvain *et al.*, in prep.
CD, Hebeler, Schwenk, PRL **122**, 042501

For applications to the dilute Fermi Gas:
Wellenhofer, CD, Schwenk,
PRC **104**, 014003 & PLB **802**, 135247

The number of diagrams increases rapidly!

	1	3	39	840	27 300	1 232 280	...
$n =$	2	3	4	5	6	7	

Integer sequence A064732:

Number of labeled Hugenholtz diagrams with n nodes.



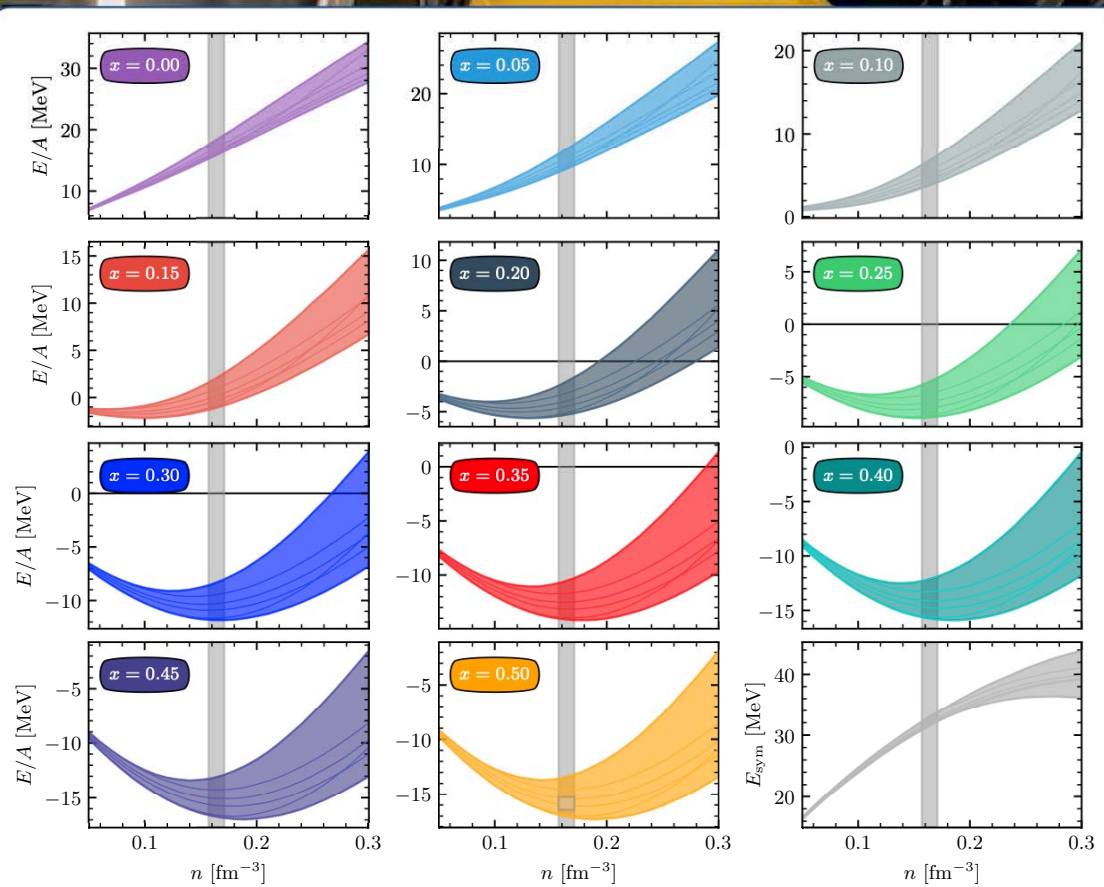
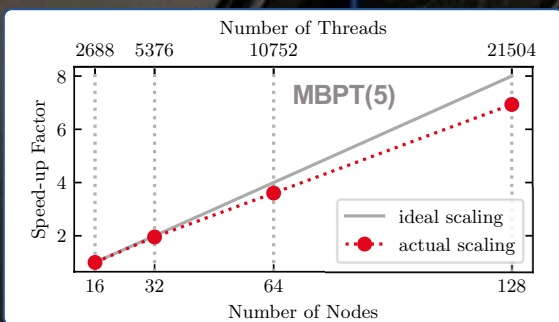
with automated diagram generation



fully automated approach
to MBPT for nuclear matter

Stevenson, *Int. J. Mod. Phys. C* **14**, 1135
Arhuis *et al.*, *Comput. Phys.* **240**, 202

MBPT: an HPC application

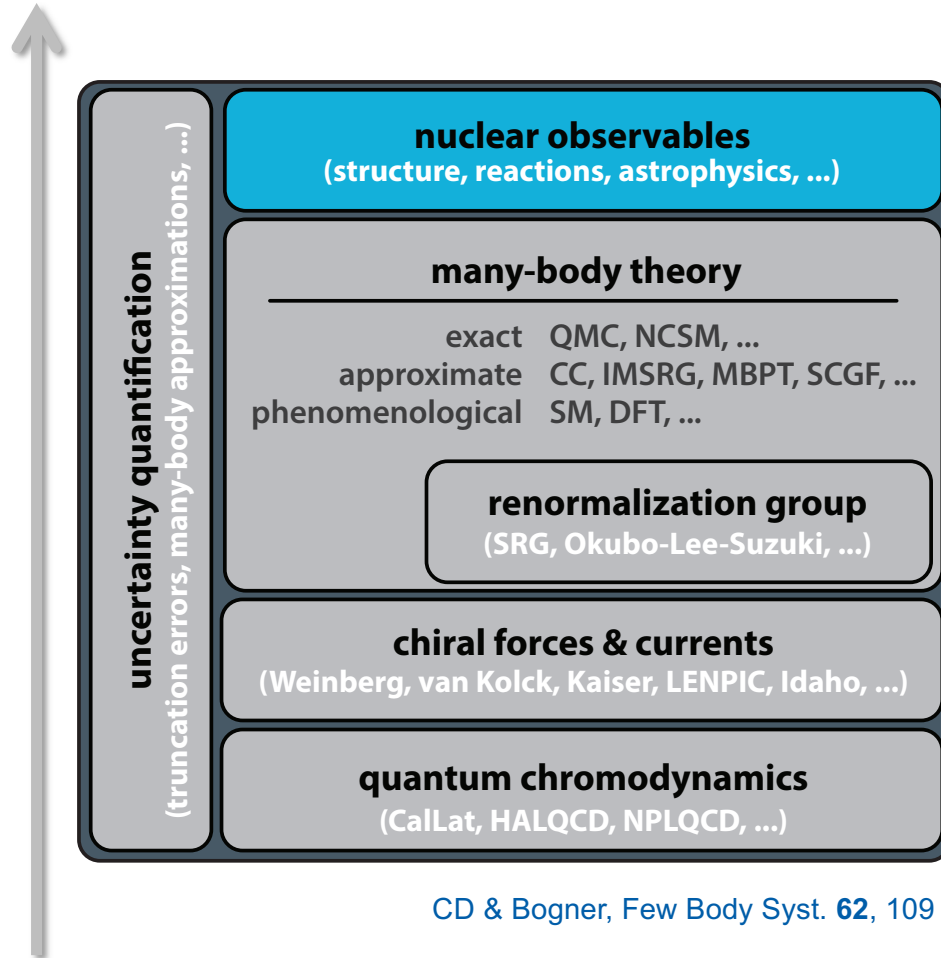


#2 (U.S.)

Summit @ Oak Ridge Leadership Computing Facility

202 752 CPU Cores
27 648 Nvidia GPUs

122.3 peta flops



nuclear EOS with quantified uncertainties
energy per particle (and derived quantities)

$$\frac{E}{A}(n, \delta, T)$$

baryon density n
neutron excess δ
temperature $T (= 0)$

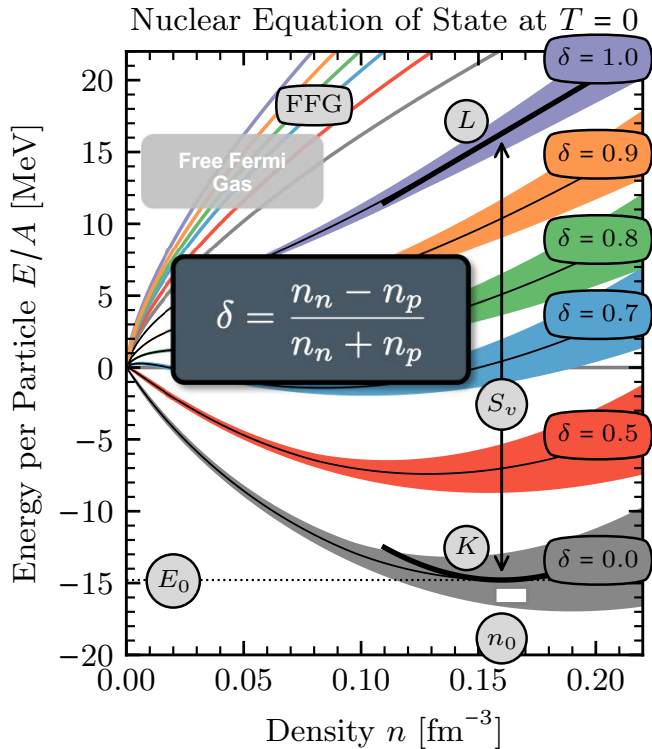
Uncertainty quantification

robust estimates of theoretical uncertainties using Bayesian machine learning via Gaussian Processes
uncertainties in EFT-based calculations due to:

- truncating the EFT expansion
- applying many-body (and other) approximations
- fitting LECs to experimental data

First chiral potentials with uncertainties fully quantified and their application:
Wesolowski, Svensson *et al.*, *PRC* **104**, 064001
Djärv, Ekstöm *et al.*, *PRC* **105**, 014005

Microscopic EOS calculations

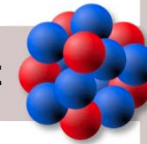


low-density EOS parameters

- symmetry energy S_v
- slope parameter L
- saturation point (n_0, E_0)
- incompressibility K

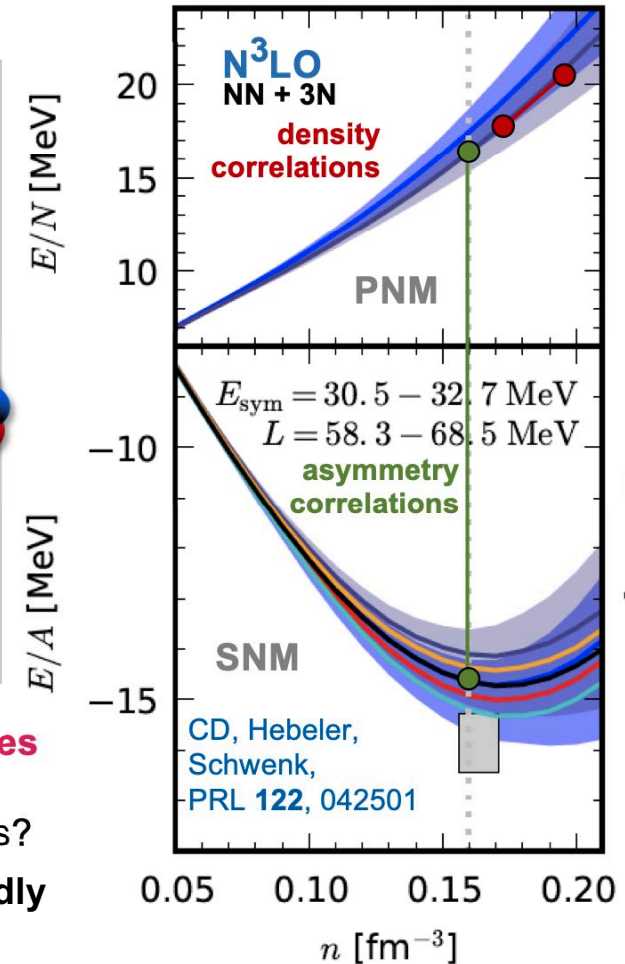
connection to experiment

- binding energies
- neutron skin thicknesses
- giant resonances
- heavy-ion collisions



But: the existing **uncertainty estimates** had only *limited* statistical meaning; *clean* propagation to derived quantities?
EFT truncation errors increase rapidly with increasing density

Great progress in microscopic EOS calculations at densities $\lesssim 2n_0$ and predictions for the neutron star structure



Rigorous UQ for nuclear matter



CD, Furnstahl, Melendez,
Phillips, PRL 125, 202702

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	NN forces	3N forces	4N forces
LO (Q ⁰)			
NLO (Q ²)			
N ² LO (Q ³)			—
N ³ LO (Q ⁴)			
N ⁴ LO (Q ⁵)			

$$Q = \max\left(\frac{p}{\Lambda_b}, \frac{m_\pi}{\Lambda_b}\right)$$

$$\{y_0, y_2, y_3, \dots, y_k\}$$

predict observable y order by order
in the chiral expansion

$$y_k = y_{\text{ref}} \sum_{n=0}^k c_n Q^n$$

make a *falsifiable* model assumption
for the convergence pattern

$$\mathcal{GP}[0, \bar{c}^2 r(x, x'; l)]$$

treat all c_n as independent draws
from a single Gaussian Process

$$\delta y_k = y_{\text{ref}} \sum_{n=k+1}^{\infty} c_n Q^n$$

learn hyperparameters of that GP &
compute to-all-orders truncation error

Recent applications of BUQEYE's GP model for correlated EFT truncation errors:

Gaussian process error modeling for chiral EFT calculations of $np \leftrightarrow dy$ at low energies

Acharya & Bacca, PLB 827, 137011

Ab initio predictions link the neutron skin of ²⁰⁸Pb to nuclear forces

Hu, Jiang, Miyagi *et al.*, arXiv:2112.01125

Ab initio nucleon-nucleus elastic scattering with chiral EFT uncertainties

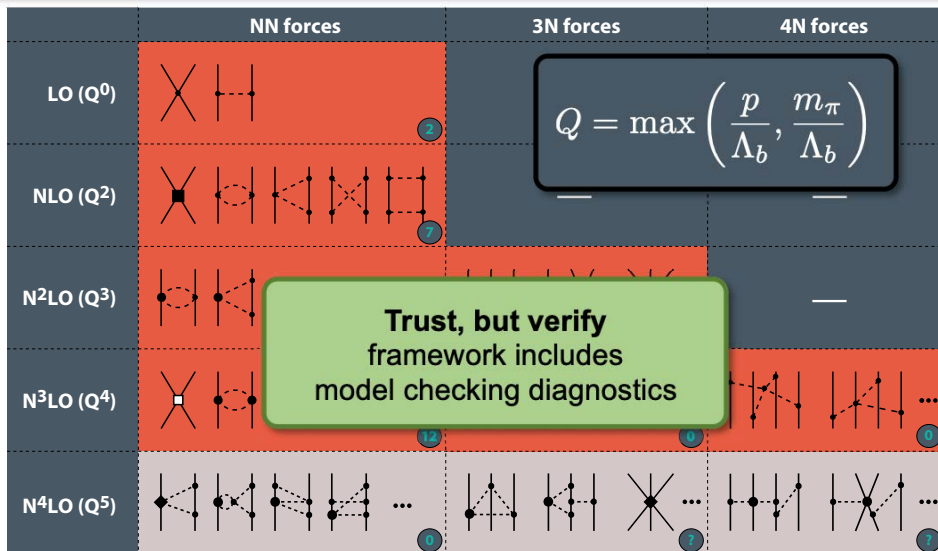
Baker, McClung, Elster *et al.*, arXiv:2112.02442

Rigorous UQ for nuclear matter



CD, Furnstahl, Melendez,
Phillips, PRL 125, 202702

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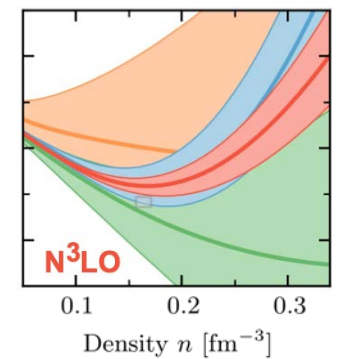
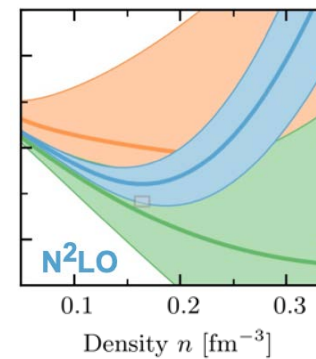
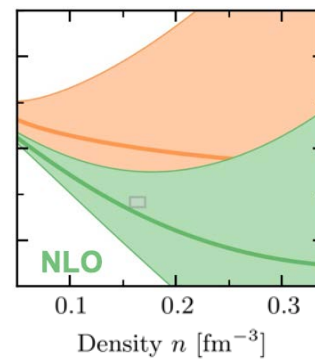
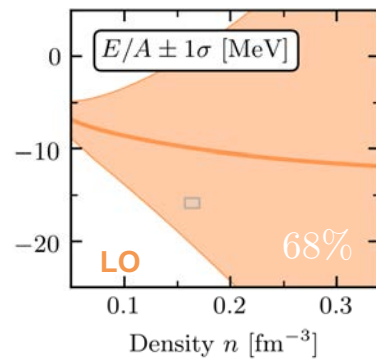
learn hyperparameters of that GP &
compute to-all-orders truncation error

An example: symmetric matter

$$y = \frac{E}{A}, \quad k = 4 \quad (\text{N}^3\text{LO})$$

Uncertainty bands depict
68% credibility regions

$$y = y_k + \delta y_k$$

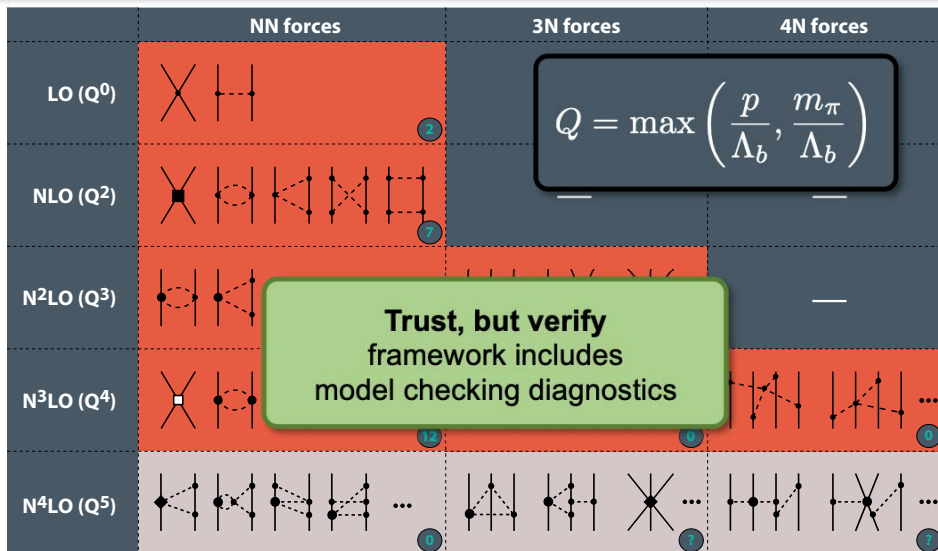


Rigorous UQ for nuclear matter



CD, Furnstahl, Melendez,
Phillips, PRL 125, 202702

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$$\{y_0, y_2, y_3, \dots, y_k\}$$

predict observable **y** order by order
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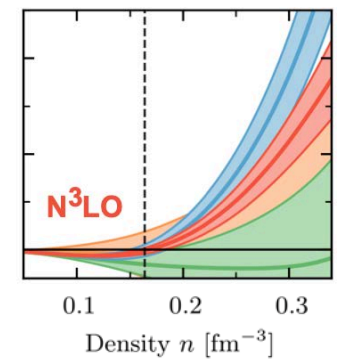
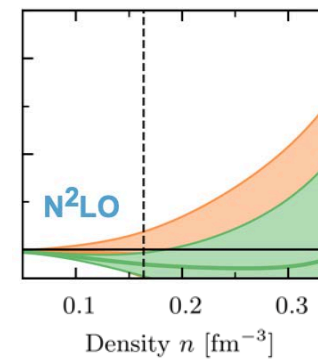
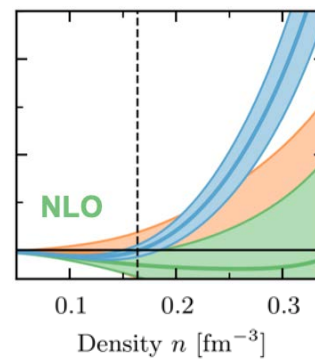
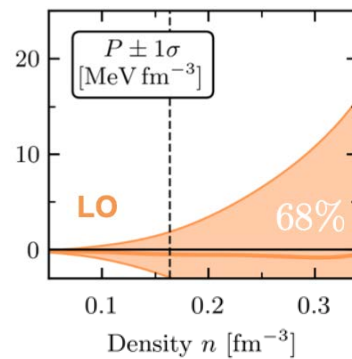
learn hyperparameters of that GP &
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An example: symmetric matter

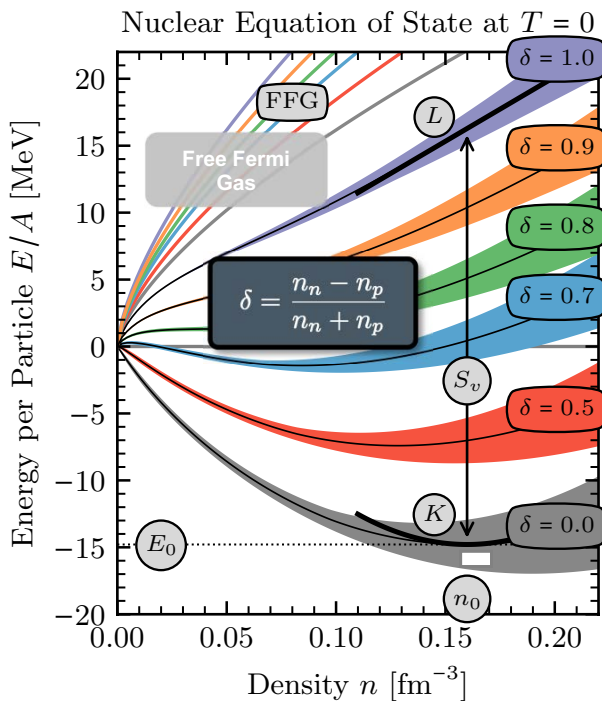
$$y = P \equiv n^2 \frac{dE}{dnA}, \quad k = 4$$

Uncertainty bands depict
68% credibility regions

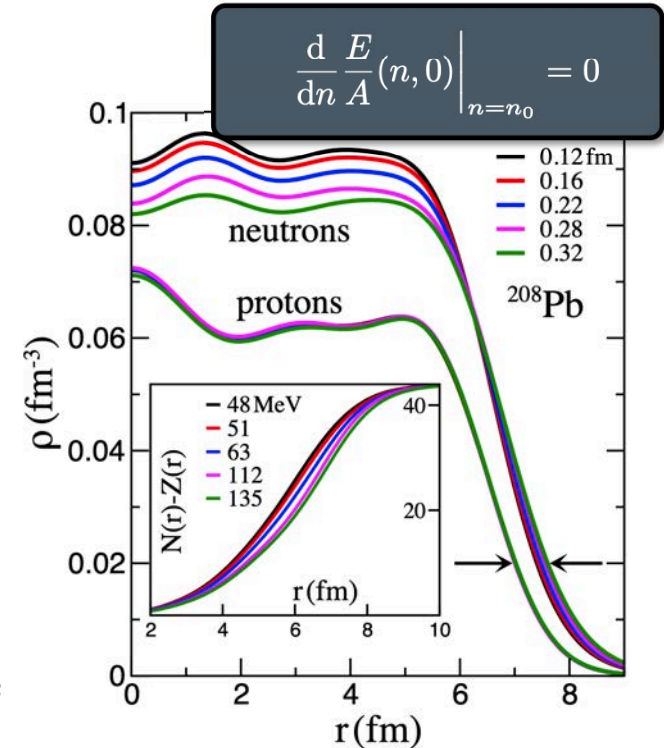
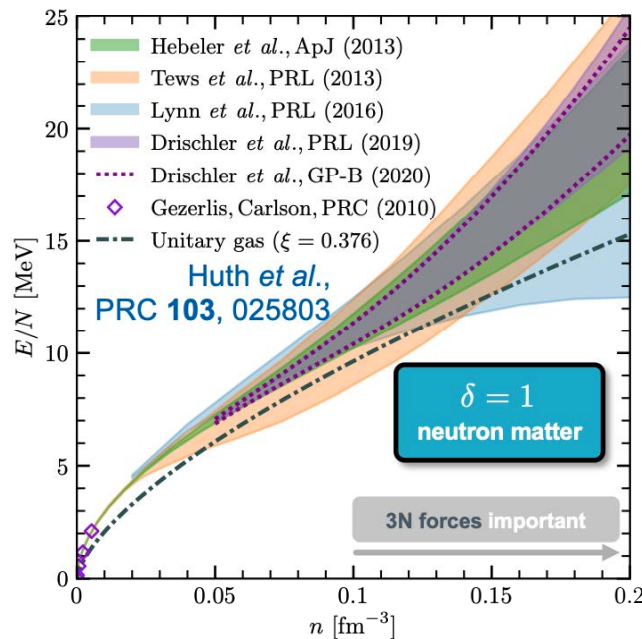
$$y = y_k + \delta y_k$$



Neutron matter | saturation in symmetric matter



CD, Holt, and Wellenhofer, Annu. Rev. Nucl. Part. Sci. **71**, 403



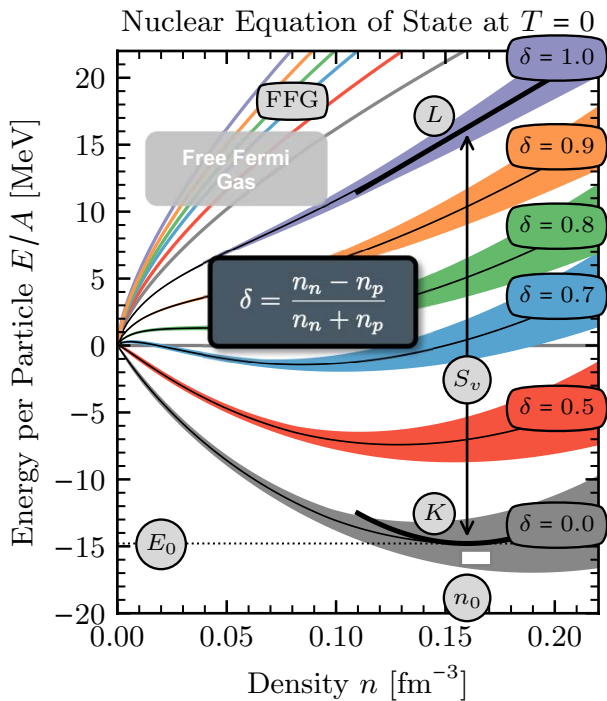
Piekarewicz & Fattoyev, Phys. Today **72**, 7

saturation point: **fine-tuned cancellation** between the kinetic and interaction contributions (ideal testbed for chiral EFT)

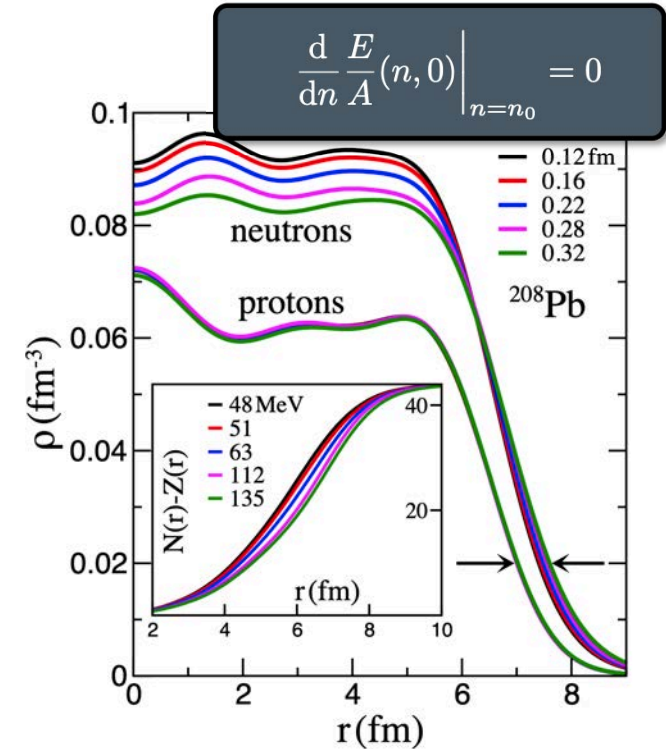
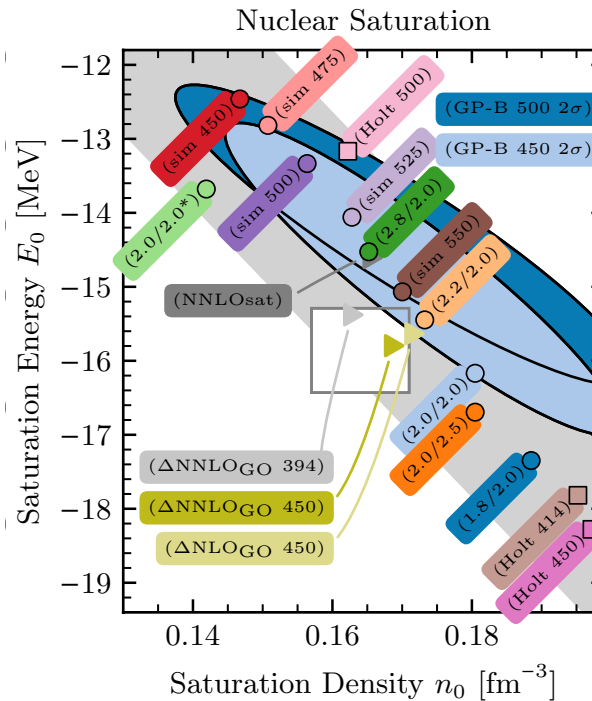
Coester band overlaps with the empirical box (but limited meaning without errors)
 Annotations: (λ / Λ_{3N}) in fm^{-1} or (Λ) in MeV

needed: improved predictions with novel NN+3N interactions and robust uncertainty quantification

Neutron matter | saturation in symmetric matter



CD, Holt, and Wellenhofer, Annu. Rev. Nucl. Part. Sci. **71**, 403



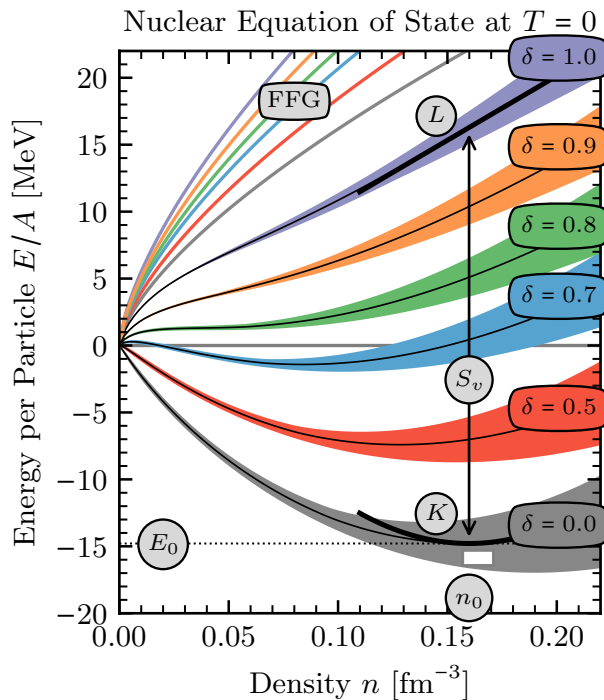
Piekarewicz & Fattoyev, Phys. Today **72**, 7

saturation point: **fine-tuned cancellation** between the kinetic and interaction contributions (ideal testbed for chiral EFT)

Coester band overlaps with the empirical box (but limited meaning without errors)
Annotations: (λ / Λ_{3N}) in fm^{-1} or (Λ) in MeV

needed: improved predictions with novel NN+3N interactions and robust uncertainty quantification

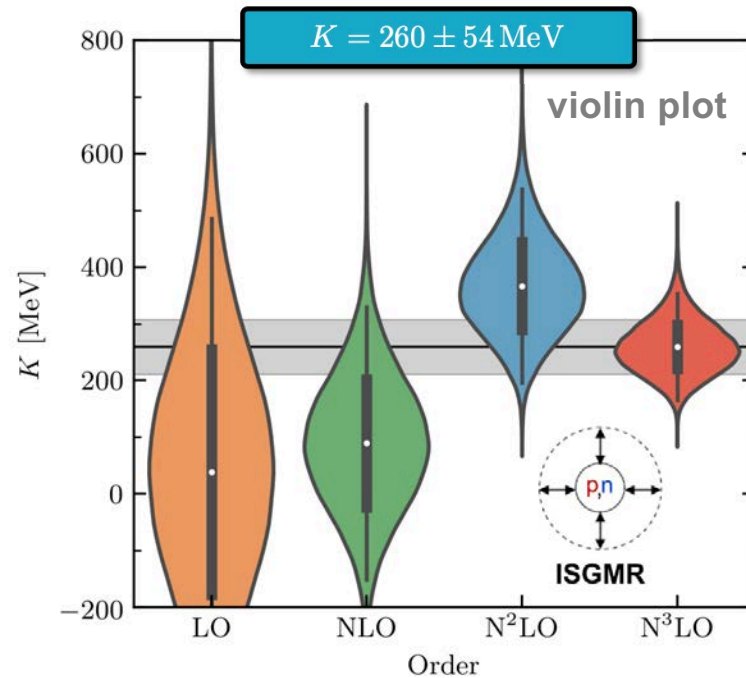
Incompressibility (in symmetric matter)



CD, Holt *et al.*, ARNPS **71**, 403

$$\text{pr}(K | \mathcal{D}) = \int \text{pr}(K | \mathcal{D}, n_0) \text{pr}(n_0 | \mathcal{D}) dn_0$$

$$\text{pr}(n_0 | \mathcal{D}) \approx 0.17 \pm 0.01 \text{ fm}^{-3}$$

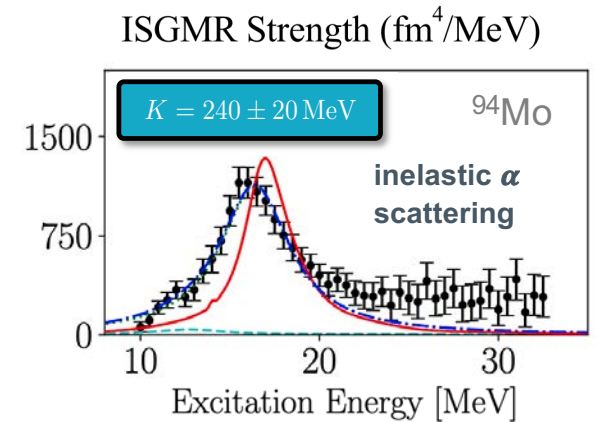


CD, Melendez *et al.*, PRC **102**, 054315

order-by-order convergence pattern

uncertainties due to the predicted saturation density included via marginalization

$$K = 9n_0^2 \left. \frac{d^2 E}{dn^2} \frac{E}{A}(n) \right|_{n=n_0}$$



Howard, Garg *et al.*, PLB **807** 135608
Roca-Maza & Paar, PPNP **101**, 96

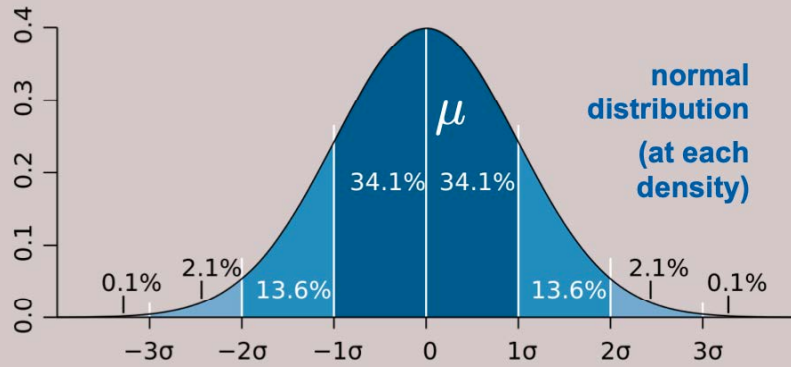
Approved FRIB experiment:

"The ISGM R in ^{132}Sn : Implications on the Nuclear Incompressibility"

Randhawa *et al.* (experiment: 21056)

Why correlations are important: symmetry energy

Reminder: Statistics 101



$$S_2 \sim \mathcal{N}(\mu_{S_2}, \sigma_{S_2}^2)$$

$$\mu_{S_2} = \mu_{\text{PNM}} - \mu_{\text{SNM}}$$

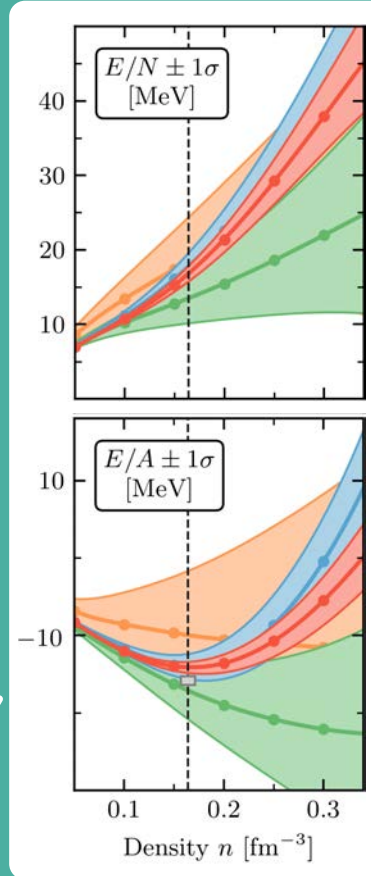
$$\sigma_{S_2}^2 = \sigma_{\text{PNM}}^2 + \sigma_{\text{SNM}}^2$$

$$- 2\sigma_{\text{PNM}}\sigma_{\text{SNM}}\rho$$

correlation coefficient $-1 \leq \rho \leq +1$

may result in smaller uncertainties than one might naively expect

difference



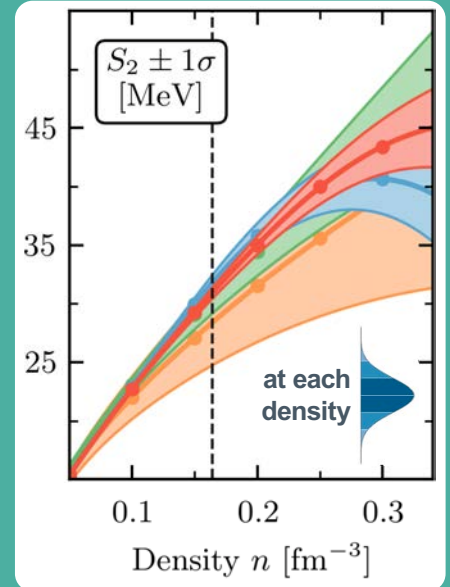
at each density

at each density

$$E/A \sim \mathcal{N}(\mu, \sigma^2)$$

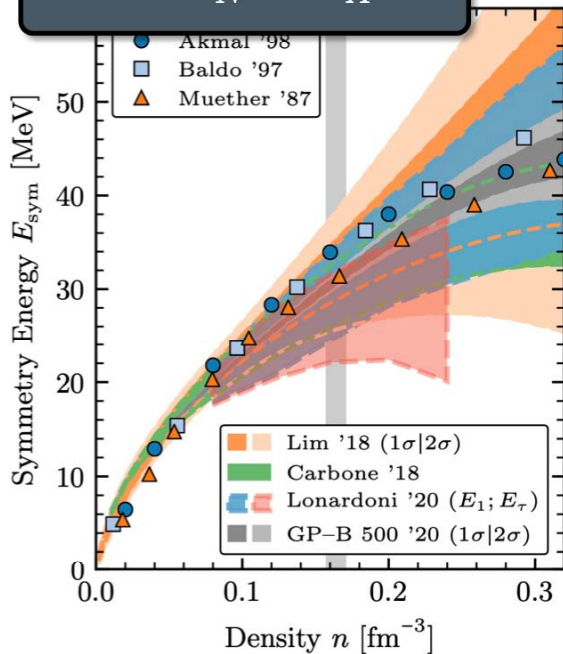
symmetry energy | multi-task GPs

$$S_2(n) \approx \frac{E}{N}(n) - \frac{E}{A}(n)$$



Nuclear symmetry energy

$$S_2(n) \approx \frac{E}{N}(n) - \frac{E}{A}(n)$$

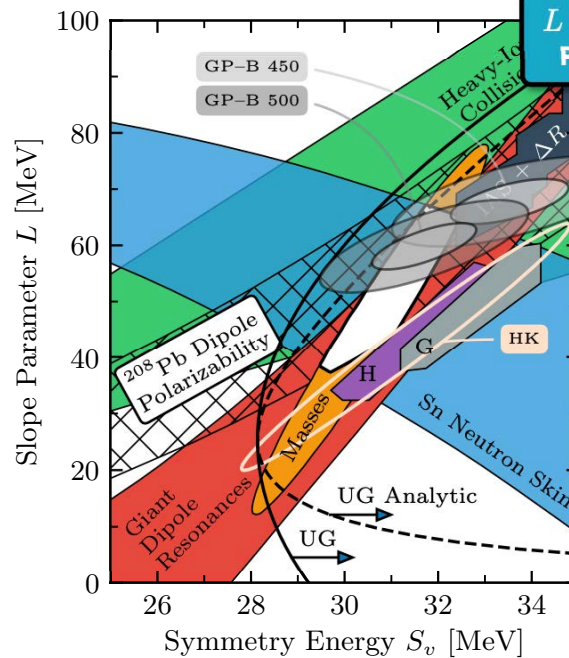


CD, Holt *et al.*, ARNPS **71**, 403

$$\text{pr}(S_v, L | \mathcal{D}) = \int \text{pr}(S_v, L | \mathcal{D}, n_0) \text{pr}(n_0 | \mathcal{D}) dn_0$$

$$\text{pr}(n_0 | \mathcal{D}) \approx 0.17 \pm 0.01 \text{ fm}^{-3}$$

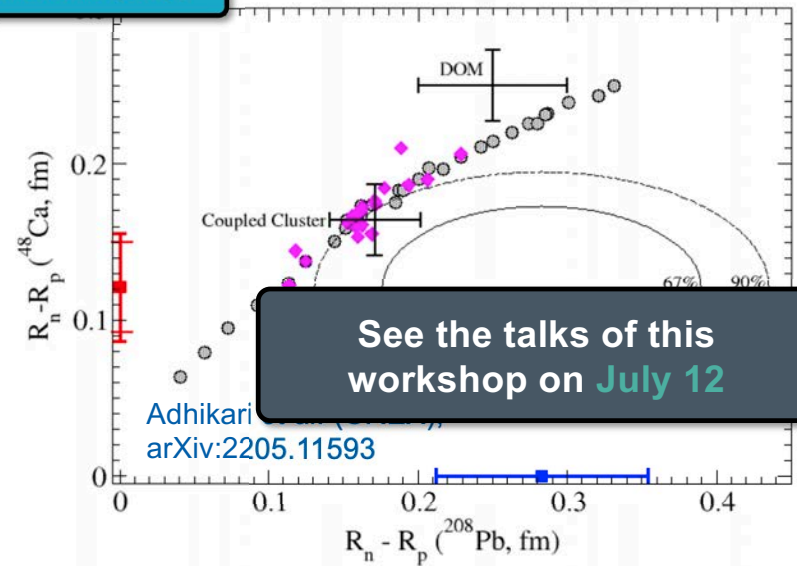
excellent agreement with experiment



CD, Holt *et al.*, ARNPS **71**, 403
Lattimer & Lim, APJ **771**, 51

marginalization over predicted saturation density

$$S_2(n) \equiv S_v + \frac{L}{3} \left(\frac{n - n_0}{n_0} \right) + \dots$$



See the talks of this workshop on July 12

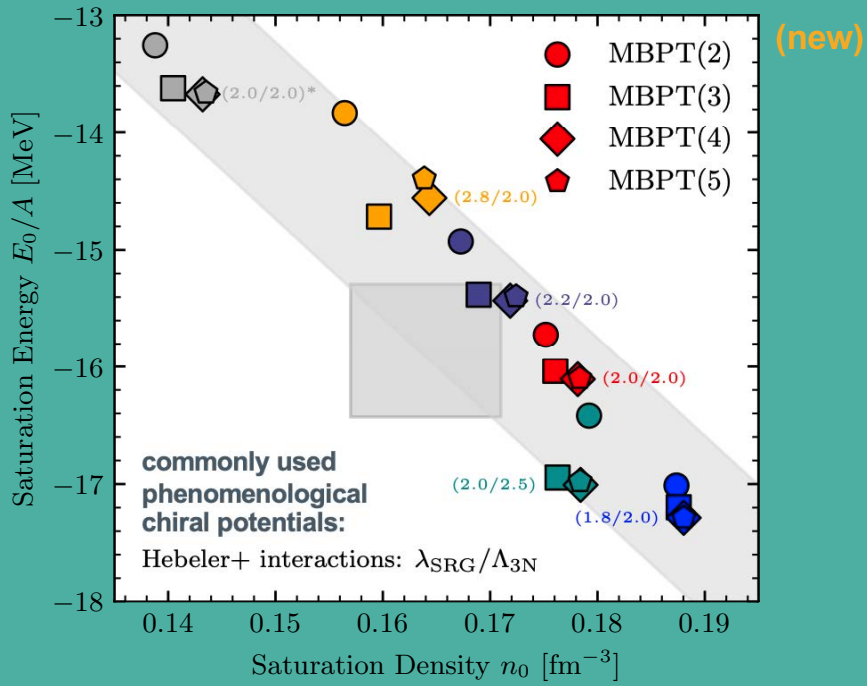
Reinhard *et al.*, PRL **127**, 232501
Reed, Fattoyev *et al.*, PRL **126**, 172503
Piekarewicz, PRC **104**, 024329

“Tension” between PREX-II and different theoretical approaches at the ~68-95% level

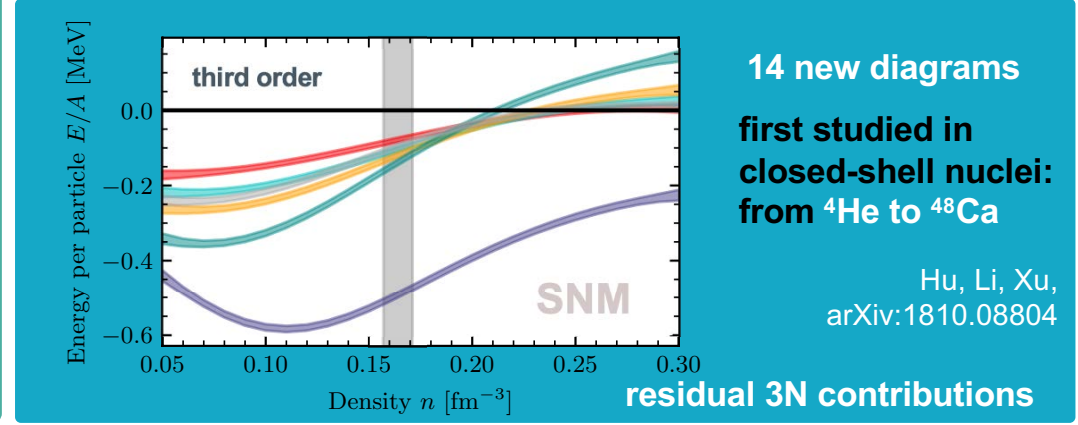
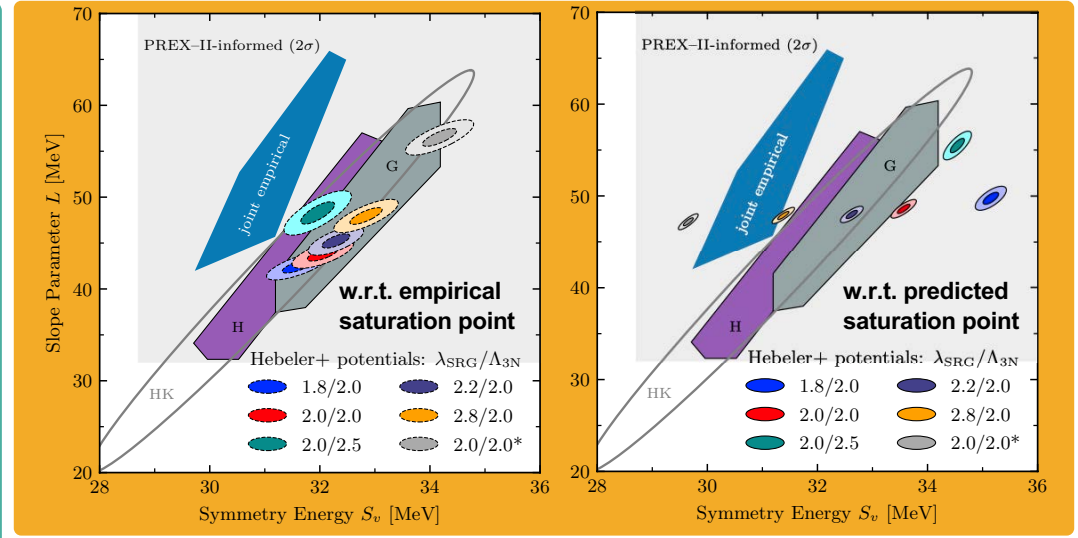
3N forces beyond normal ordering?

Drischler, McElvain *et al.*, in prep.

MBPT(<i>n</i>)	2	3	4	5
NN+3N norm. ord.	✓	✓	(✓, 3N)	✓
residual 3N	✓ (1)	✓ (14)	X	X

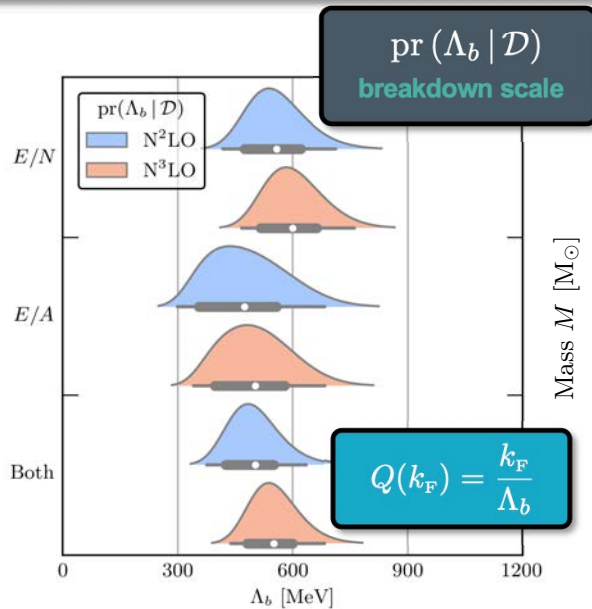


nuclear saturation | symmetric nuclear matter



residual 3N contributions

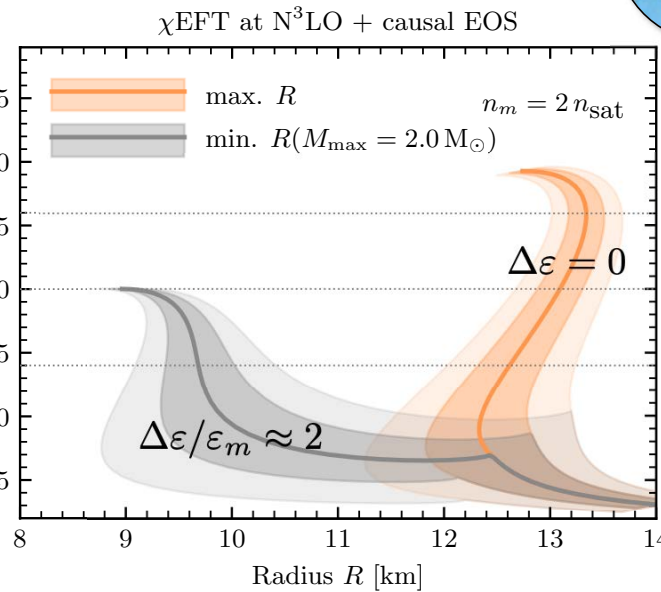
Exploring the limits of chiral EFT



CD, Melendez *et al.*, PRC **102**, 054315

Bayesian inference of the in-medium breakdown scale

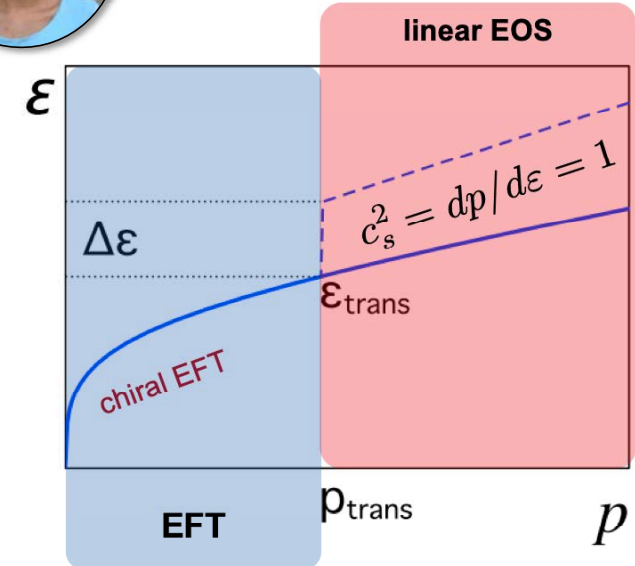
But: at what density does chiral EFT break down?



CD, Han, Lattimer *et al.*, PRC **103**, 045808
CD, Han, and Reddy, PRC **105**, 035808

derived **bounds on the neutron star radius** (and sound speed) assuming chiral EFT is valid up to a given critical density (here: $2n_0$) could already be challenged by NICER

$$R_{2.0} = (11.4 - 16.1) \text{ km} \quad \text{Riley et al., AJL 918, L27} \\ \text{Miller et al., AJL 918, L28}$$



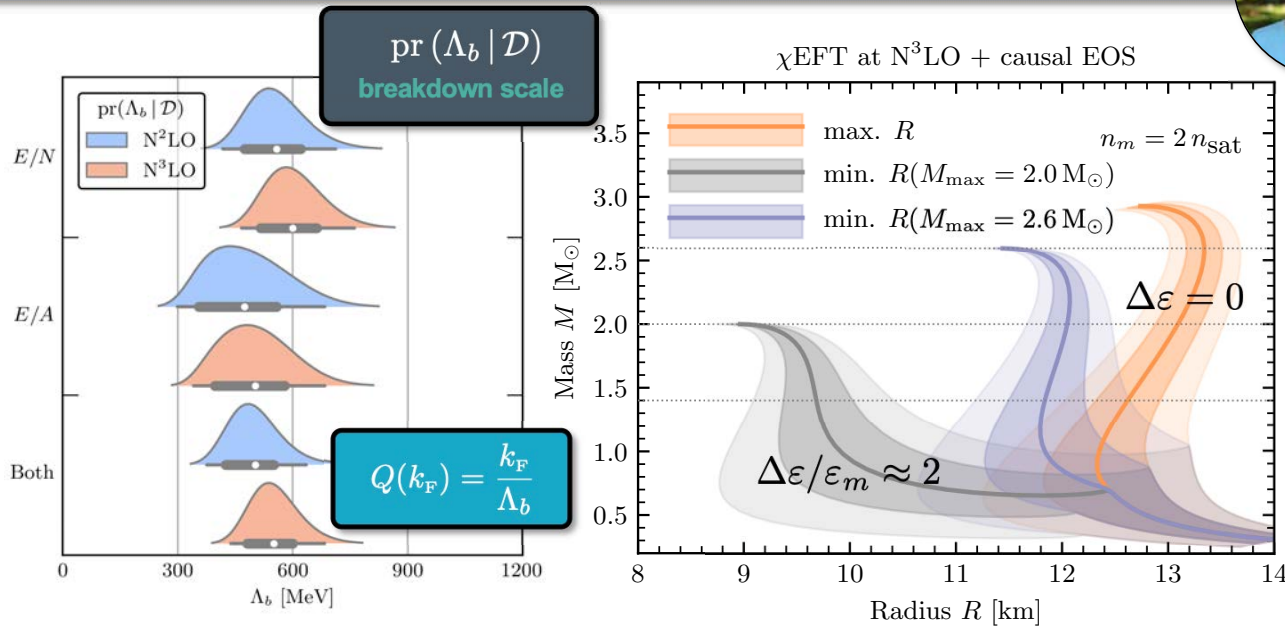
Han & Prakash, APJ **899**, 2
Alford *et al.*, JPG: NPP **46**, 114001

extend EFT EOS at n_m to linear EoS with finite discontinuity (softening)

continuous match sets upper bound

use **lower limit on M_{max}** from observation to adjust $\Delta\epsilon$ and constrain R_{min}

Exploring the limits of chiral EFT



CD, Melendez *et al.*, PRC **102**, 054315

Bayesian inference of the in-medium breakdown scale

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CD, Han, Lattimer *et al.*, PRC **103**, 045808
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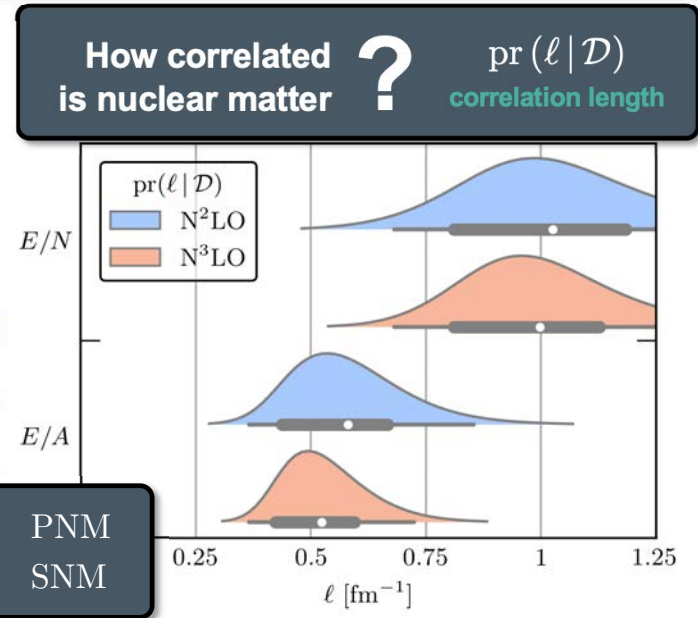
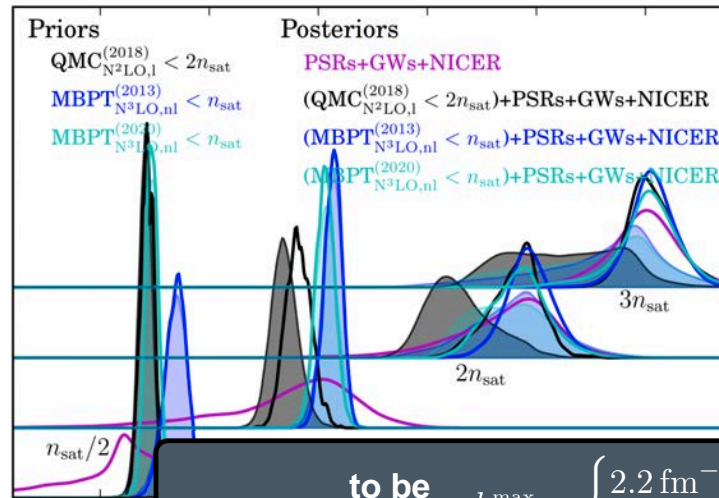
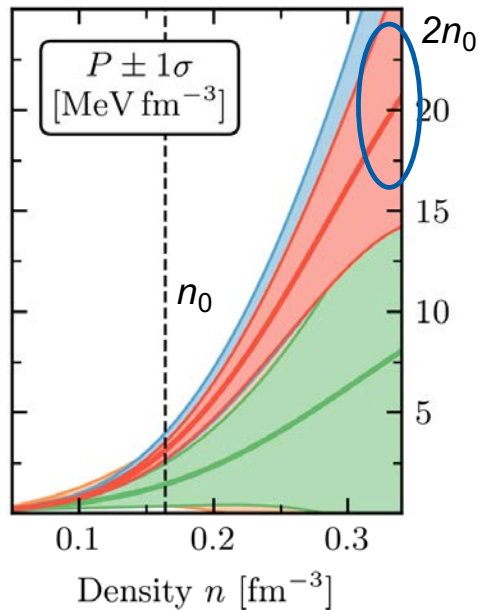
Han & Prakash, APJ **899**, 2
Alford *et al.*, JPG: NPP **46**, 114001

extend EFT EOS at n_m to linear EoS with finite discontinuity (softening)

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use **lower limit on M_{max}** from observation to adjust $\Delta\varepsilon$ and constrain R_{min}

Direct astrophysical tests at supranuclear densities



CD, Furnstahl *et al.*, PRL **125**, 202702

see also: Essick *et al.*, PRC **102**, 055803

CD, Melendez *et al.*, PRC **102**, 054315



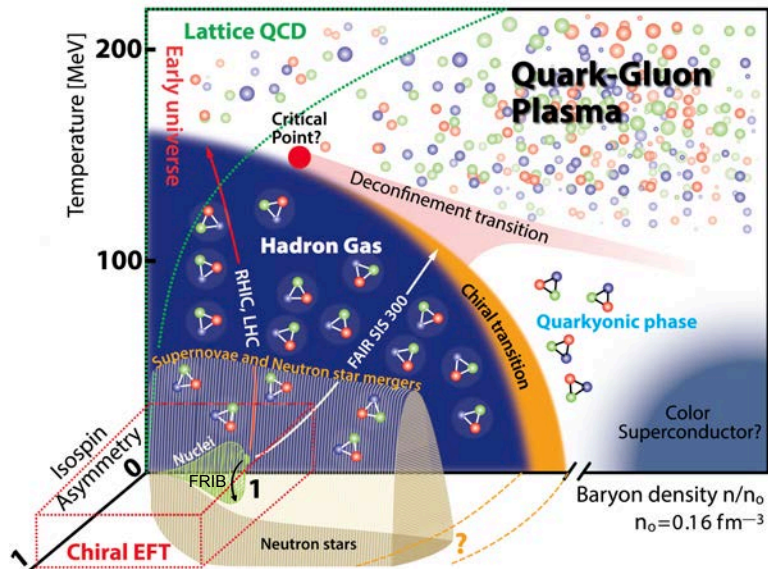
Neutron star observations could be used for:
 Model checking & selection of chiral interactions
 Constraints on coupling constants in nuclear forces

EFT truncation error is highly correlated

$$P(n = 0.32 \text{ fm}^{-3}) \approx \begin{cases} 20 \pm 6 \text{ MeV fm}^{-3} & \text{MBPT: nonlocal} \\ 15 \pm 5 \text{ MeV fm}^{-3} & \text{QMC: local } V_{E,1} \end{cases}$$

More details? Recent review article

MICHIGAN STATE
UNIVERSITY



Chiral Effective Field Theory and the High-Density Nuclear Equation of State

Annual Review of Nuclear and Particle Science

Vol. 71:403-432 (Volume publication date September 2021)

First published as a Review in Advance on July 6, 2021

<https://doi.org/10.1146/annurev-nucl-102419-041903>



C. Drischler,^{1,2,3} J.W. Holt,⁴ and C. Wellenhofer^{5,6}

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²Nuclear Science Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA

³Facility for Rare Isotope Beams, Michigan State University, East Lansing, Michigan 48824, USA; email: drischler@frib.msu.edu

⁴Cyclotron Institute and Department of Physics and Astronomy, Texas A&M University, College Station, Texas 77843, USA

⁵Institut für Kernphysik, Technische Universität Darmstadt, 64289 Darmstadt, Germany

⁶ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung GmbH, 64291 Darmstadt, Germany

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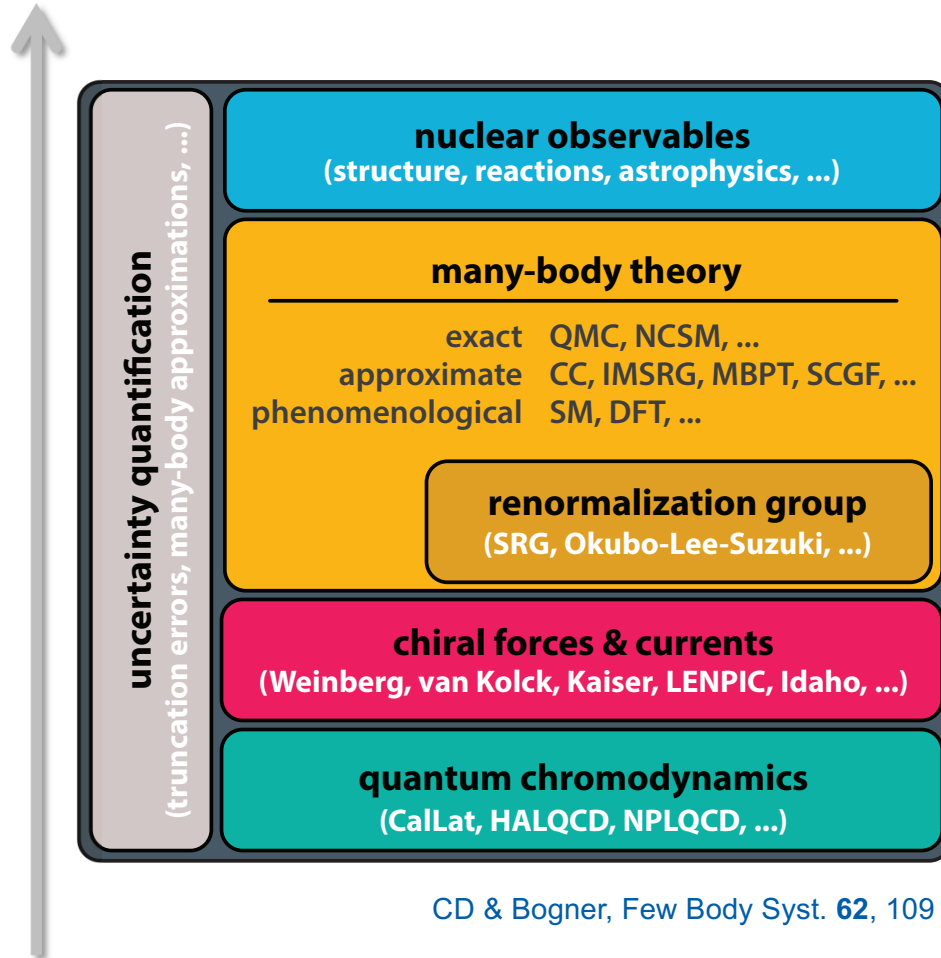
Keywords:

Chiral EFT | neutron stars | MBPT
nuclear matter at zero and finite temperature
Bayesian uncertainty quantification
recent neutron star observations

see also in the same journal:
James Lattimer, Annu. Rev. Nucl. Part. Sci. 71, 433

Open Access

Ab initio workflow (idealized)



Here: nuclear equation of state (EOS)
energy per particle (and derived quantities)

$$\frac{E}{A}(n, \delta, T)$$

baryon density n
neutron excess δ
temperature $T (=0)$

Emulators for nuclear physics

computationally *inexpensive* algorithms
capable of approximating exact model
calculations with *high accuracy*

*surrogate
models*

See also applications in GW astronomy, such as:

Fast Prediction and Evaluation of Gravitational Waveforms Using Surrogate Models

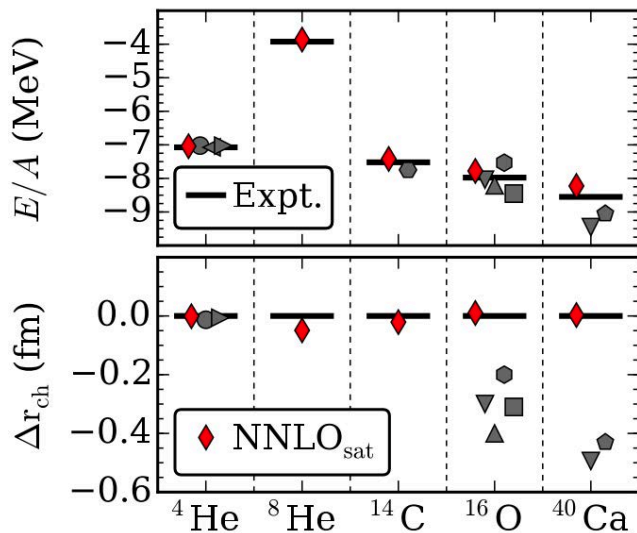
Field, Galley, Hesthaven, Kaye, Tiglio, PRX 4, 031006

Frequency-domain reduced order models for

gravitational waves from aligned-spin compact binaries

Pürrer, Class. Quantum Grav. 31, 195010

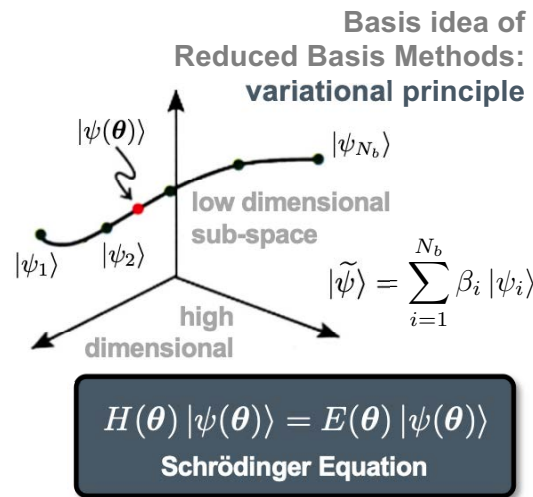
Game changers in nuclear physics: emulators



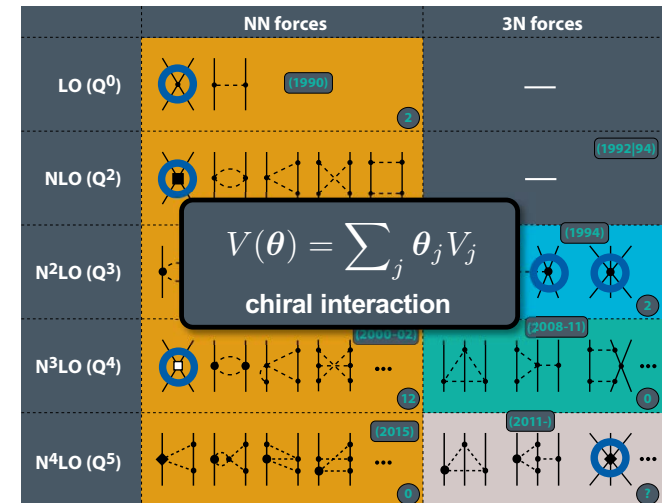
Ekström *et al.*, PRC **91**, 051301(R)

Open questions in chiral EFT: power counting, regulator artifacts, and differing predictions for medium-mass nuclei

statistical analyses of scattering data could provide valuable insights



Frame *et al.*, PRL **121**, 032501
Melendez, CD *et al.*, arXiv:2203.05528



CD, Holt, and Wellenhofer, ARNPS **71**, 403

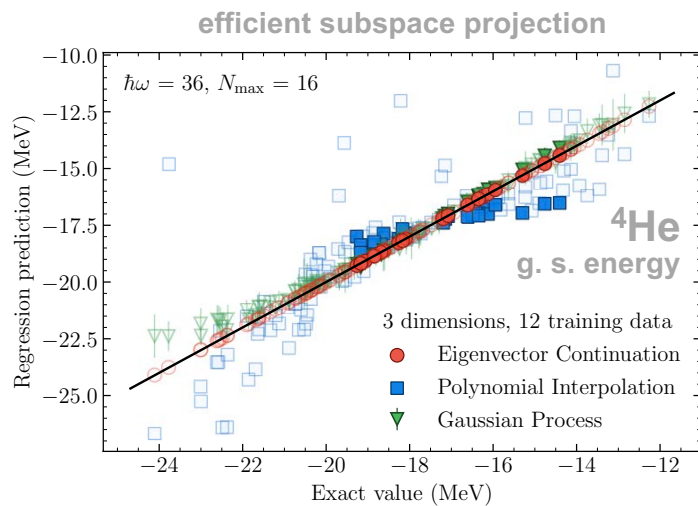
Here: Reduced Basis Methods (Eigenvector Continuation)

Construct **reduced-order models** by removing superfluous information in HiFi models

Emulators enable applications thought to be prohibitively slow (MC sampling)

Example: **Bayesian parameter estimation**

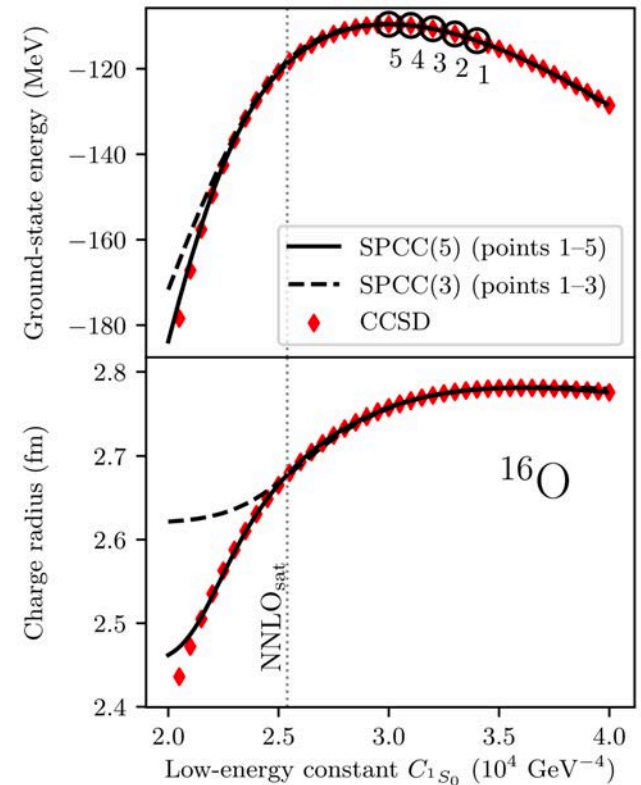
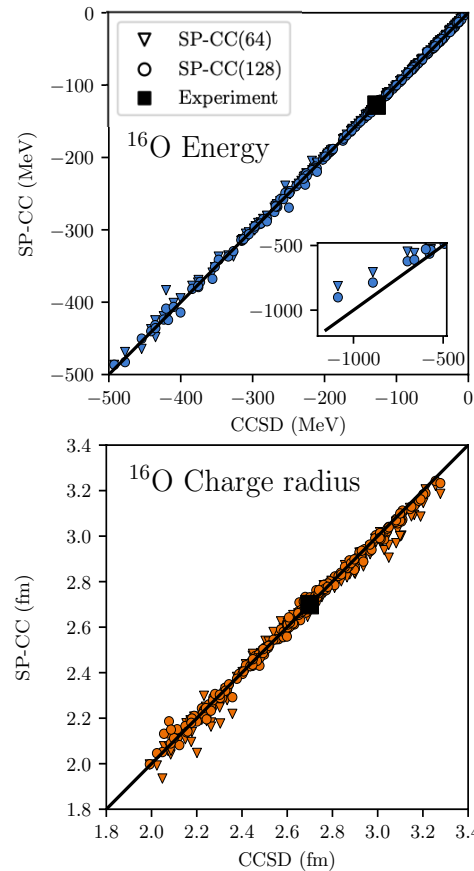
Emulators for bound-state calculations



König, Ekström *et al.*, PLB **810**, 135814

Fast & accurate emulation via subspace projection methods (RBM)

RBM-driven emulators have **accurately approximated g. s. properties** binding energies and charge radii

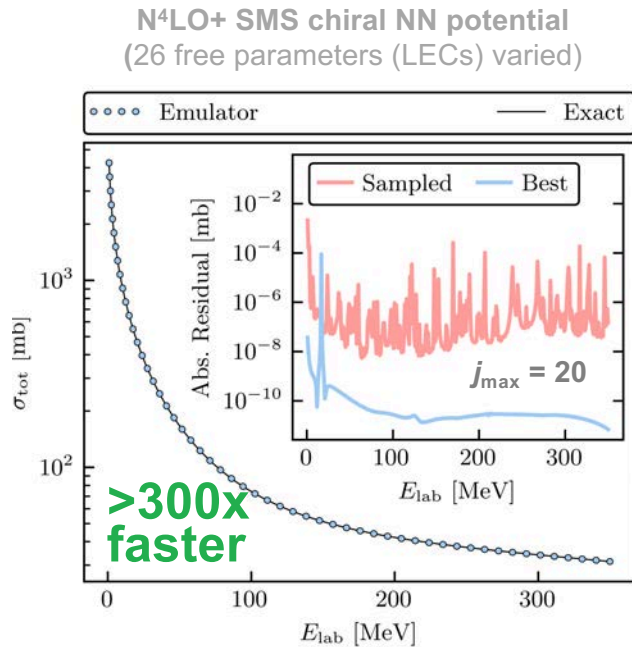


Ekström & Hagen, PRL **123**, 252501

Millions of sampling points computed in one hour on a standard laptop.
An equivalent set of exact CC computations would require 20 years.

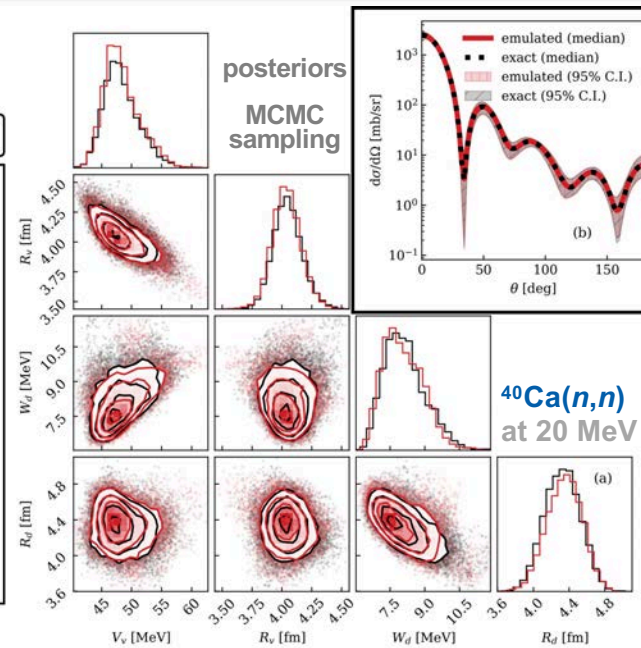
Emulators: mining scattering & reaction data, and more

BAND
Bayesian Analysis of Nuclear Dynamics



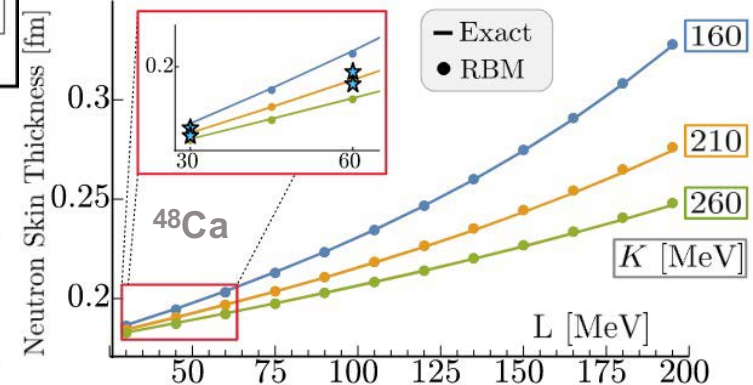
Melendez, CD *et al.*, PLB **821**, 136608

Proof of principle: emulation of two-body scattering observables with and without wavefunctions (KVP vs NVP)
26 free parameters (LECs) varied (N⁴LO+ SMS chiral NN potential)



CD, Quinonez *et al.*, PLB **823**, 136777

both interpolation and extrapolation in the parameter space θ with negligible errors
efficient Bayesian parameter estimation for improving next-generation chiral interactions and optical models



Bonilla, Giuliani, Godbey, Lee, arXiv:2203.05284
Anderson, O'Donnell, Piekarewicz, arXiv:2206.14889
see also: J. Melendez on GitHub

towards **calibrating modern EDFs** using Bayesian optimization and RBM emulators
Proof of principle: emulated the entire single-particle spectrum of a variety of nuclei

Emulator classification

data-driven (non-intrusive)

Gaussian Processes, Artificial Neural Networks, etc.

model-driven (intrusive)

reduced-order equations from high-fidelity equations

J. A. Melendez,^{1,*} C. Drischler,^{2,†} R. J. Furnstahl,^{1,‡} A. J. Garcia,^{1,§} and Xilin Zhang^{2,¶}

¹*Department of Physics, The Ohio State University, Columbus, OH 43210, USA*

²*Facility for Rare Isotope Beams, Michigan State University, MI 48824, USA*

Many pointers to the MOR literature: [arXiv:2203.05528](https://arxiv.org/abs/2203.05528) (accepted Guide in J. Phys. G Nucl. Part.)

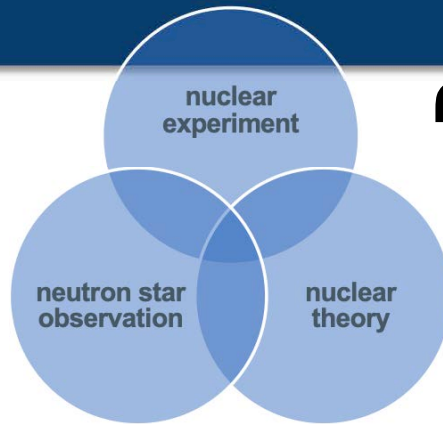
The field of model order reduction (MOR) is growing in importance due to its ability to extract the key insights from complex simulations while discarding computationally burdensome and superfluous information. We provide an overview of MOR methods for the creation of fast & accurate emulators of memory- and compute-intensive nuclear systems. As an example, we describe how “eigenvector continuation” is a special case of a much more general and well-studied MOR formalism for parameterized systems. We continue with an introduction to the Ritz and Galerkin projection methods that underpin many such emulators, while pointing to the relevant MOR theory and its successful applications along the way. We believe that this will open the door to broader applications in nuclear physics and facilitate communication with practitioners in other fields.

see also: CD & Zhang, Chap. 8, pp. 29–36, in [arXiv:2202.01105](https://arxiv.org/abs/2202.01105) (collective pieces edited by Tews, Davoudi, Ekström, Holt)



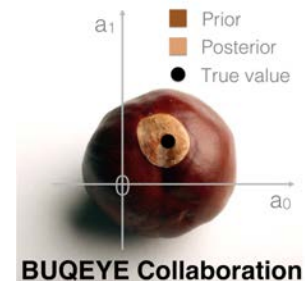
Take-away points

multi-messenger
nuclear precision
FRIB } era



“ Enormous progress in theory, experiment, and observation make this new era particularly fruitful for the determination of the equation of state of neutron-rich matter. ”

- 1 Upcoming observational (and experimental) campaigns will provide **stringent constraints** on the properties of neutron stars
- 2 Chiral EFT enables **microscopic predictions** of nuclear matter (and nuclei) **with quantified uncertainties** to interpret these empirical constraints
- 3 Bayesian methods are powerful tools for quantifying & propagating **correlated uncertainties** in EFT-based calculations (*model checking* is important)
- 4 Emulators have been **game changers in nuclear physics**; and much can be learned from the well-established MOR field in applied mathematics.



Many thanks to my collaborators:

R. Furnstahl J. W. Holt J. Melendez K. McElvain D. Phillips
S. Han J. Lattimer M. Prakash S. Reddy C. Wellenhofer T. Zhao