

Comment on "Jahn-Teller effect for the negatively charged C₆₀ molecule: Analogy with the silicon vacancy"

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(Received 27 December 1991)

We point out that the effective U due to vibrational couplings is a factor of 2 larger than calculated by Lannoo *et al.*, due to off-diagonal interactions neglected in their work.

Lannoo *et al.*¹ found an effective interaction between valence electrons on C₆₀ due to Jahn-Teller distortions of magnitude $U = -0.05$ eV. This is twice the energy shift of a single electron in the harmonic approximation, and thus agrees with the single-electron energy shift $\Delta E = -0.024$ eV reported in Ref. 2. We wish to point out here that these calculations neglect significant off-diagonal matrix elements of the electron-vibration coupling. Although Lannoo *et al.* write down a Hamiltonian with off-diagonal matrix elements, they minimize in a space that has a static valence electron and static deformation. In the harmonic approximation, it is equivalent to the perturbation expression

$$\Delta E_{\text{static}} = - \sum_{\alpha} \frac{\langle i, 0 | V | i, \alpha \rangle^2}{\hbar \omega_{\alpha}}$$

where the notation $|i, \alpha\rangle$ is a state with an electron in orbital i and the vibrational wave function with a quantum of excitation in the α mode. The valence electron is in the t_{1u} multiplet of states, and in Ref. 1 the state i was chosen to minimize the energy. However, the full perturbation expression should have an additional sum over final electron states in the t_{1u} multiplet,

$$\Delta E_{\text{full}} = - \sum_{\alpha, j} \frac{\langle i, 0 | V | j, \alpha \rangle^2}{\hbar \omega_{\alpha}}. \quad (1)$$

We have recalculated these quantities using the electron Hamiltonian of Refs. 3 and 4 and the vibrational Hamiltonian of Ref. 5. We find values $\Delta E_{\text{static}} = -0.023$ eV and $\Delta E_{\text{full}} = -0.049$ eV, i.e., a doubling when the off-diagonal terms are included.

These relative magnitudes are easy to understand if one replaces the finite-symmetry group by the full rotation group. The t_{1u} state behaves like a state of orbital symmetry $l = 1$, and the dominant H_g vibration behaves like an $l = 2$ state in a spherical basis. Then the ratio of perturbations would be given by a Clebsch-Gordan coefficient, $\Delta E_{\text{full}}/\Delta E_{\text{static}} = (1020|10)^{-2} = 2.5$.

The effective U should be calculated using the full perturbation formula in the two-particle state with symmetry $t_{1u} \times t_{1u} = A$. This is the state $|A\rangle = \sum_j |jj^*\rangle / \sqrt{3}$. The resulting U is twice ΔE_{full} , while the corresponding U in the static state $|ii^*\rangle$ is twice ΔE_{static} .

Thus we agree with Refs. 1 and 2 for the calculations they reported, but the full vibrational contribution to U should be -0.1 eV. This is still too small to overcome the Coulomb repulsion and make an attractive interaction that would explain the superconductivity.

¹M. Lannoo *et al.*, Phys. Rev. B **44**, 12 106 (1991).

²V. de Coulon, J.L. Martins, and F. Reuse, Phys. Rev. B **45**, 13 671 (1992).

³D. Tomanek and M. Schlueter, Phys. Rev. Lett. **67**, 2331

(1991).

⁴G. Bertsch *et al.*, Phys. Rev. Lett. **67**, 2690 (1991).

⁵D. Weeks and W. Harter, Chem. Phys. Lett. **144**, 366 (1988).