Continuous Gravitational Waves from Neutron Stars

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Purpose

Neutron stars are also possible sources of detectable continuous gravitational waves.

The LIGO Scientific Collaboration and Virgo Collaboration have searched for these waves.

The purpose of this talk is give an overview of these searches, and discuss what can be learned from them compared to searches for gravitational-waves from binary neutron star coalescence. Recent results will be presented.
Periodic Continuous Gravitational waves

• The GW signal from a triaxial pulsar can be modelled as
  
  \[ h(t) = \frac{1}{2} F_+(t; \psi) h_0 (1 + \cos^2 \iota) \cos 2\Psi(t) + F_\times(t; \psi) h_0 \cos \iota \sin 2\Psi(t) \]

• The unknown parameters are
  • \( h_0 \) - amplitude of the gravitational wave signal
  • \( \psi \) - polarization angle of signal; embedded in \( F_{x,y} \)
  • \( \iota \) - inclination angle of the pulsar
  • \( \phi_0 \) - initial phase of pulsar \( \Phi(0) \)

• In the known pulsar searches we usually look for signals at twice the rotation frequency of the pulsars
  • For blind searches for isolated neutron stars the location in the sky and the source’s frequency and its evolution are search parameters.
Generation of Continuous Gravitational Waves

- Radiation generated by quadrupolar mass movements:

\[ h_{\mu\nu} = \frac{2G}{rc^4} \frac{d^2}{dt^2} \left[ I_{\mu\nu} \right] \]

\(I_{\mu\nu} = \) quadrupole tensor, \(r = \) source distance

No GW from axisymmetric object rotating about symmetry axis

- Spinning neutron star with equatorial ellipticity \(\varepsilon_{\text{equat}}\):

\[ \varepsilon_{\text{equat}} = \frac{|I_{xx} - I_{yy}|}{I_{zz}} \]

gives a strain amplitude \(h\) \((f_{GW} = 2f_{\text{Rot}})\):

\[ h = 1.1 \times 10^{-24} \left[ \frac{kpc}{r} \right] \left[ \frac{f_{GW}}{kHz} \right]^2 \left[ \frac{\varepsilon}{10^{-6}} \right] \left[ \frac{I_{zz}}{10^{38} \text{ kg} \cdot \text{m}^2} \right] \]

Courtesy: U. Liverpool
Gravitational CW mechanisms

- **Equatorial ellipticity** (e.g., – mm-high “bulge”):
  \[ h \propto \varepsilon_{\text{equat}} \quad \text{with} \quad f_{GW} = 2f_{\text{rot}} \]

- **Poloidal ellipticity** (natural) + wobble angle (precessing star):
  \[ h \propto \varepsilon_{\text{poloidal}} \times \theta_{\text{wobble}} \quad \text{with} \quad f_{GW} = f_{\text{rot}} \pm f_{\text{precess}} \]
  (precession due to different L and \( \Omega \) axes)

- **Two-component (crust+superfluid) \( \rightarrow \)**
  \[ f_{GW} = f_{\text{rot}} \quad \text{and} \quad 2f_{\text{rot}} \]

- **r modes** (rotational oscillations – CFS-driven instability):
  S. Chandrasekhar PRL 24 (1970) 611

  \[ h \propto \alpha_{\text{r-mode}} \quad \text{with} \quad f_{GW} \approx \frac{4}{3} f_{\text{rot}} \]
Gravitational CW mechanisms

Assumption we (LSC, Virgo) have usually made to date:

- Bulge is best bet for detection
- Look for GW emission at twice the EM frequency

- e.g., look for Crab Pulsar (29.7 Hz) at 59.5 Hz
- (troublesome frequency in North America!)

What is allowed for $\varepsilon_{\text{equat}}$?

- Old maximum (?) $\approx 5 \times 10^{-7}$ [$\sigma/10^{-2}$] (“ordinary” neutron star)
  with $\sigma =$ breaking strain of crust

- More recent finding: $\sigma \approx 10^{-1}$ supported by detailed numerical simulation

- Recent re-evaluation: $\varepsilon_{\text{equat}} < 10^{-5}$
Strange quark stars could support much higher ellipticities

Maximum $\varepsilon_{\text{equat}} \approx 10^{-1}$ (!)

But what $\varepsilon_{\text{equat}}$ is realistic?

What could drive $\varepsilon_{\text{equat}}$ to a high value (besides accretion)?

Millisecound pulsars have spindown-implied values lower than $10^{-9} – 10^{-6}$

New papers revisiting possible GW emission mechanisms (e.g., buried magnetic fields, accretion-driven r-modes) are also intriguing
What is the “indirect spindown limit”? 

It is useful to define the “indirect spindown limit” for a known pulsar, under the assumption that it is a “gravitar”, i.e., a star spinning down due to gravitational wave energy loss.

Unrealistic for known stars, but serves as a useful benchmark.

Equating “measured” rotational energy loss (from measured period increase and reasonable moment of inertia) to GW emission gives:

$$h_{SD} = 2.5 \times 10^{-25} \left[ \frac{kpc}{d} \right] \sqrt{\left[ \frac{1 kHz}{f_{GW}} \right] \left[ \frac{-df_{GW}/dt}{10^{-10} Hz/s} \right] \left[ \frac{I}{10^{45} g \cdot cm^2} \right]}$$

Example:

Crab $\rightarrow$ $h_{SD} = 1.4 \times 10^{-24}$

($d=2$ kpc, $f_{GW} = 59.5$ Hz, $df_{GW}/dt = -7.4 \times 10^{-10}$ Hz/s)
Finding a completely unknown CW Source

Serious technical difficulty: Doppler frequency shifts
- Frequency modulation from earth’s rotation ($v/c \sim 10^{-6}$)
- Frequency modulation from earth’s orbital motion ($v/c \sim 10^{-4}$)
  → Coherent integration of 1 year gives frequency resolution of 30 nHz
  → 1 kHz source spread over 6 million bins in ordinary FFT!

Additional, related complications:
- Daily amplitude modulation of antenna pattern
- Spin-down of source
- Orbital motion of sources in binary systems
Finding a completely unknown CW Source

Modulations / drifts complicate analysis enormously:

- Simple Fourier transform inadequate
- Every sky direction requires different demodulation

Computational scaling:

Single coherence time \( T_{\text{coherence}} \) – Sensitivity improves as \( (T_{\text{coherence}})^{1/2} \)
but cost scales with \( (T_{\text{coherence}})^{6+} \)

\[ \rightarrow \text{Restricts } T_{\text{coherence}} < 1\text{-}2 \text{ days for all-sky search} \]

\[ \rightarrow \text{Exploit } \text{coincidence} \text{ among different spans} \]

Alternative:

Semi-coherent stacking of spectra (e.g., \( T_{\text{coherence}} = 30 \text{ min} \))

\[ \rightarrow \text{Sensitivity improves only as } (N_{\text{stack}})^{1/4} \]

\[ \rightarrow \text{All-sky survey at full sensitivity} = \text{Formidable challenge} \]

Impossible?
Frequency Modulation

\[ f(t) \equiv \left( 1 + \frac{\ddot{v}(t)}{c} \cdot \hat{n} \right) \left[ f_0 + f_1(t - t_0) + ... \right] \]

Relativistic corrections can be included in the actual code.

\[ \dot{f}(t) \equiv \left( \frac{\ddot{a}(t)}{c} \cdot \hat{n} \right) \left[ f_0 + f_1(t - t_0) \right] + \left( 1 + \frac{\ddot{v}(t)}{c} \cdot \hat{n} \right) f_1 + ... \]

\[ S = \left( \frac{\ddot{a}_{\text{orb}}(t)}{c} \cdot \hat{n} \right) f_0 + f_1 \]

For analysis < 1 yr sky points with small S have small doppler variation making them harder to distinguish GWs from instrument lines at these points.
But three substantial benefits from modulations:

- Reality of signal confirmed by need for corrections
- Corrections give precise direction of source
- Single interferometer can make definitive discovery

Can “zoom in” further with follow-up algorithms once we lock on to source

[V. Dergachev, PRD 85 (2012) 062003
M. Shaltev & R. Prix, PRD 87 (2013) 084057]
Recent results

Targeted (matched-filter) algorithm applied to 195 known pulsars over LIGO S5/S6 and Virgo VSR2/VSR4 data

Lowest (best) upper limit on strain:

\[ h_0 < 2.1 \times 10^{-26} \]

Lowest (best) upper limit on ellipticity:

\[ \varepsilon < 6.7 \times 10^{-8} \]

Crab limit at 1% of total energy loss

Vela limit at 10% of total energy loss

Recent results

Directed-search algorithm applied to the galactic center using LIGO S5 data (knowing direction improves sensitivity)

Uses semi-coherent sums of 630 11.5-hr F-Statistic* powers

Einstein@Home now carrying out similar searches for SNRs

*Jaranowski, Krolak & Schutz, PRD 58 (1998) 063001
Recent results

First all-sky search for unknown binary CW sources

Uses TwoSpect* algorithm:

Sample spectrogram (30-minute FFTs) for simulated strong signal (Earth’s motion already demodulated)

Result of Fourier transforming each row of spectrogram

→ Concentrates power in orbital harmonics

*E. Goetz & K. Riles, CQG 28 (2011) 215006
Recent results

First all-sky search for unknown binary CW sources

TwoSpect results for unknown spinning neutron stars in binary systems:

The blue dots in this plot show the upper limits on the circularly polarized gravitational wave strain amplitude.

The red dots show the upper limit on the randomly polarized gravitational wave strain amplitude.

arXiv:1405.7904
Summary

No discoveries yet, but…

• Still examining data we have taken
  (computationally bound – E@H: 1 Petaflop, 100K volunteers)

• Major upgrade of LIGO & Virgo under way now
  • Advanced LIGO & Virgo
  • Improves range more than an order of magnitude
  • Moore’s Law will help too…

Electromagnetic observations (radio, x-ray, γ-ray) of nearby neutron stars helpful now – and later
Extra Slides
Not all known sources have measured timing

Compact central object in the Cassiopeia A supernova remnant

Birth observed in 1681 – One of the youngest neutron stars known

Star is observed in X-rays, but no pulsations observed

Requires a broad band search over accessible band
Other results

Search for Cassiopeia A – Young age (~300 years) requires search over 2\textsuperscript{nd} derivative

Other results

S5 all-sky results:

- worst case (linear)
- best case (circular)
- non-Gaussian
- 60 Hz

Semi-coherent stacks of 30-minute, demodulated power spectra ("PowerFlux")

Astrophysical reach:

- $\epsilon=1e^{-4}$
- $\epsilon=1e^{-5}$
- $1$ kpc
- $100$ pc
- $\epsilon=1e^{-6}$
- $10$ pc
- $\epsilon=1e^{-7}$
- $1$ pc
- $\epsilon=1e^{-8}$

Frequency derivative (Hz/s)

Frequency (Hz)
Other results

S5 all-sky results:

Einstein@Home semi-coherent sums of 121 25-hour F-Statistic powers (2 interferometers)

Astrophysical reach:

PRD 87 (2013) 042001
Other results

S5 all-sky results:

Hough-transform search based on ~68K 30-minute demodulated spectra (3 interferometers)

arXiv:1311.2409 (Nov 2013)
The Global Interferometer Network

The three (two) LIGO, Virgo and GEO interferometers are part of a Global Network.

Multiple signal detections will increase detection confidence and provide better precision on source locations and wave polarizations.
LIGO Observatories

Hanford

Observation of nearly simultaneous signals 3000 km apart rules out terrestrial artifacts

Livingston
Virgo

Have begun collaborating with Virgo colleagues (Italy/France)

Took data in coincidence for last ~4 months of latest science run

Data exchange and joint analysis underway

Will coordinate closely on detector upgrades and future data taking

3-km Michelson Interferometer just outside Pisa, Italy
GEO600

Work closely with the GEO600 Experiment (Germany / UK / Spain)

• Arrange coincidence data runs when commissioning schedules permit
• GEO members are full members of the LIGO Scientific Collaboration
• Data exchange and strong collaboration in analysis now routine
• Major partners in proposed Advanced LIGO upgrade

600-meter Michelson Interferometer just outside Hannover, Germany
LIGO Detector Facilities

Vacuum System

- Stainless-steel tubes (1.24 m diameter, ~$10^{-8}$ torr)
- Gate valves for optics isolation
- Protected by concrete enclosure
LIGO Detector Facilities

LASER
- Infrared (1064 nm, 10-W) Nd-YAG laser from Lightwave (now commercial product!)
- Elaborate intensity & frequency stabilization system, including feedback from main interferometer

Optics
- Fused silica (high-Q, low-absorption, 1 nm surface rms, 25-cm diameter)
- Suspended by single steel wire
- Actuation of alignment / position via magnets & coils
LIGO Detector Facilities

Seismic Isolation

- Multi-stage (mass & springs) optical table support gives $10^6$ suppression
- Pendulum suspension gives additional $1/f^2$ suppression above $\sim 1$ Hz
Gravitational Wave Detection

- Suspended Interferometers (IFO’s)

  - Suspended mirrors in “free-fall”
  
  - Michelson IFO is “natural” GW detector
  
  - Broad-band response (~20 Hz to few kHz)
  
  - Waveform information (e.g., chirp reconstruction)
LIGO Interferometer Optical Scheme

Michelson interferometer

With Fabry-Perot arm cavities

• Recycling mirror matches losses, enhances effective power by ~ 50x

4 km Fabry-Perot cavity

end test mass
What Limits the Sensitivity of the Interferometers?

- Seismic noise & vibration limit at low frequencies
- Atomic vibrations (Thermal Noise) inside components limit at mid frequencies
- Quantum nature of light (Shot Noise) limits at high frequencies
- Myriad details of the lasers, electronics, etc., can make problems above these levels

Best design sensitivity:
\[ \sim 3 \times 10^{-23} \text{ Hz}^{1/2} @ 150 \text{ Hz} \]
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Best design sensitivity:

- $\approx 3 \times 10^{-23} \text{ Hz}^{-1/2}$ @ 150 Hz
- $< 2 \times 10^{-23}$ (enhanced LIGO)
**LIGO S1 → S5 Sensitivities (“Initial LIGO”)**

2002-2007

Best Strain Sensitivities for the LIGO Interferometers

Comparisons among S1 - S5 Runs  LIGO-G060009-01-Z

\[ h_{\text{rms}} = 3 \times 10^{-22} \]
“Enhanced LIGO” (July 2009 – Oct 2010)

Displacement spectral noise density

Factor of 2 improvement above 300 Hz
Virgo sensitivity in VSR2 (part of LIGO S6)

\[ \geq 10^2 \times \text{better than LIGO below 30 Hz!} \]

Enabled search for Vela at 22 Hz

Comparable to LIGO in sweet spot
“Locking” the Interferometer

Sensing gravitational waves requires sustained resonance in the Fabry-Perot arms and in the recycling cavity

→ Need to maintain half-integer # of laser wavelengths between mirrors
→ Feedback control servo uses error signals from imposed RF sidebands
→ Four primary coupled degrees of freedom to control
→ Highly non-linear system with 5-6 orders of magnitude in light intensity

Also need to control mirror rotation ("pitch" & "yaw")

→ Ten more DOF’s (but less coupled)

And need to stabilize laser (intensity & frequency), keep the beam pointed, damp out seismic noise, correct for tides, etc.,…
Advanced LIGO

Sampling of source strengths vis a vis Initial LIGO and Advanced LIGO

Lower $h_{\text{rms}}$ and wider bandwidth both important

“Signal recycling” offers potential for tuning shape of noise curve to improve sensitivity in target band (e.g., known pulsar cluster)
Advanced LIGO

Increased laser power:

10 W → 180 W

Improved shot noise (high freq)

Higher-Q test mass:

Fused silica with better optical coatings

Lower internal thermal noise in bandwidth

Increased test mass:

10 kg → 40 kg

Compensates increased radiation pressure noise
Advanced LIGO

Detector Improvements:

New suspensions:

Single $\rightarrow$ Quadruple pendulum

Lower suspensions thermal noise in bandwidth

Improved seismic isolation:

Passive $\rightarrow$ Active

Lowers seismic “wall” to $\sim 10$ Hz
Neutron Star Binaries:
Average range ~ 200 Mpc
Most likely rate ~ 40/year

The science from the first 3 hours of Advanced LIGO should be comparable to 1 year of initial LIGO

(Range x ~10 \rightarrow Volume x ~1000)

But that sensitivity will not be achieved instantly...

Advanced LIGO

arXiv: 1304.0670
http://www.einsteinathome.org/

- GEO-600 Hannover
- LIGO Hanford
- LIGO Livingston
- Current search point
- Current search coordinates
- Known pulsars
- Known supernovae remnants

Your computer can help too!
Sources

Search methods can detect any type of periodic source.

Upper limits are set on gravitational-wave amplitude, $h_0$, of rotating triaxial ellipsoid.

Credits:
A. image by Jolien Creighton; LIGO Lab Document G030163-03-Z.
B. image by M. Kramer; Press Release PR0003, University of Manchester - Jodrell Bank Observatory, 2 August 2000.
C. image by Dana Berry/NASA; NASA News Release posted July 2, 2003 on Spaceflight Now.
D. image from a simulation by Chad Hanna and Benjamin Owen; B. J. Owen’s research page, Penn State University.
Searching for continuous waves

Several approaches tried or in development:

- Summed powers from many short (30-minute) FFTs with sky-dependent corrections for Doppler frequency shifts → “Semi-coherent” (StackSlide, Hough transform, PowerFlux)

- Push up close to longest coherence time allowed by computing resources (~1 day) and look for coincidences among outliers in different data stretches (Einstein@Home)
Methods

- **Semicoheren Methods**
  - StackSlide: add the power
  - Hough: add weighted 1 or 0
  - PowerFlux: add weighted power
  - Track Doppler shift and $\frac{df}{dt}$

- **Coherent Methods**
  - Bayesian Param. Estimation
  - Maximum Likelihood & Matched Filtering

\[
P(x \mid h) = \frac{1}{\sqrt{2\pi \sigma_1}} e^{-\frac{(x_1-h_1)^2}{2\sigma_1^2}} \frac{1}{\sqrt{2\pi \sigma_2}} e^{-\frac{(x_2-h_2)^2}{2\sigma_2^2}} 
\]

\[
P(h \mid x) = P(h)P(x \mid h) / P(x) \Rightarrow \text{Time Domain}
\]

\[
\chi^2 = \sum_j \frac{(x_j - h_j)^2}{\sigma_j^2} \Rightarrow \left( \sum_j \frac{x_j h_j}{\sigma_j^2} - \frac{1}{2} \sum_j \frac{h_j h_j}{\sigma_j^2} \right)
\]

\[
\Rightarrow \text{Frequency Domain} \Rightarrow (\log \Lambda)_{\text{max}} \Rightarrow F
\]

- Weights depend on both noise and antenna patterns:
- Methods can include multi-detector data and coincidence steps:
- Hierarchical Methods: combine the above to maximize sensitivity.
What is the “indirect spindown limit”?

If a star’s age is known (e.g., historical SNR), but its spin is unknown, one can still define an indirect spindown upper limit by assuming gravitar behavior has dominated its lifetime:

\[ \tau = \frac{f}{4 \left( \frac{df}{dt} \right)} \]

And substitute into \( h_{SD} \) to obtain

[K. Wette, B. Owen,… CQG 25 (2008) 235011]

\[
\begin{align*}
\hSD &= 2.2 \times 10^{-24} \left[ \frac{\text{kpc}}{d} \right] \sqrt{\left[ \frac{1000 \text{yr}}{\tau} \right]} \left[ \frac{I}{10^{45} \text{g} \cdot \text{cm}^2} \right] \\
\text{Example:} \\
\text{Cassiopeia A} &\rightarrow \hSD = 1.2 \times 10^{-24} \\
(d=3.4 \text{ kpc, } \tau=328 \text{ yr})
\end{align*}
\]
What is the “X-ray flux limit”?


\[
h_{X-ray} \approx 5 \times 10^{-27} \sqrt \left[ \frac{600 \text{Hz}}{f_{\text{sig}}} \right] \left[ \frac{F_x}{10^{-8} \text{erg} \cdot \text{cm}^{-2} \cdot \text{s}^{-1}} \right]
\]

Example: Scorpius X-1

\[
\Rightarrow h_{X-ray} \approx 3 \times 10^{-26} \left[ \frac{600 \text{ Hz}}{f_{\text{sig}}} \right]^{1/2}
\]

\((F_x = 2.5 \times 10^{-7} \text{ erg} \cdot \text{cm}^{-2} \cdot \text{s}^{-1})\)

Courtesy: McGill U.
CW observational papers to date

S1:
Setting upper limits on the strength of periodic gravitational waves from PSR J1939+2134 using the first science data from the GEO 600 and LIGO detectors - PRD 69 (2004) 082004

S2:
First all-sky upper limits from LIGO on the strength of periodic gravitational waves using the Hough transform - PRD 72 (2005) 102004

Limits on gravitational wave emission from selected pulsars using LIGO data - PRL 94 (2005) 181103 (28 pulsars)

Coherent searches for periodic gravitational waves from unknown isolated sources and Scorpius X-1: results from the second LIGO science run - PRD 76 (2007) 082001
CW observational papers to date

S3-S4:

*Upper Limits on Gravitational Wave Emission from 78 Radio Pulsars* - PRD 76 (2007) 042001

*All-sky search for periodic gravitational waves in LIGO S4 data* – PRD 77 (2008) 022001

*The Einstein@Home search for periodic gravitational waves in LIGO S4 data* – PRD 79 (2009) 022001

*Upper limit map of a background of gravitational waves* – PRD 76 (2007) 082003 (Cross-correlation – Sco X-1)
**CW observational papers to date**

**S5:**

*Beating the spin-down limit on gravitational wave emission from the Crab pulsar* – APJL 683 (2008) 45

*All-sky LIGO Search for Periodic Gravitational Waves in the Early S5 Data* – PRL 102 (2009) 111102

*Einstein@Home search for periodic gravitational waves in early S5 LIGO data* – PRD 80 (2009) 042003

*Searches for gravitational waves from known pulsars with S5 LIGO data* – APJ 713 (2010) 671 (116 pulsars)

*First search for gravitational waves from the youngest known neutron star* – APJ 722 (2010) 1504

*All-sky search for periodic gravitational waves in the full S5 LIGO data* – PRD 85 (2012) 022001
CW observational papers to date

S5:

Einstein@Home all-sky search for periodic gravitational waves in LIGO S5 data – PRD 87 (2013) 042001

A directed search for continuous Gravitational Waves from the Galactic Center – PRD 88 (2013) 102002

CW observational papers to date

S6 / VSR2 / VSR4:

*Beating the spin-down limit on gravitational wave emission from the Vela pulsar* – APJ 737 (2011) 93

*Gravitational-waves from known pulsars: results from the initial detector* – arXiv:1309.4027 (Sept 2013)